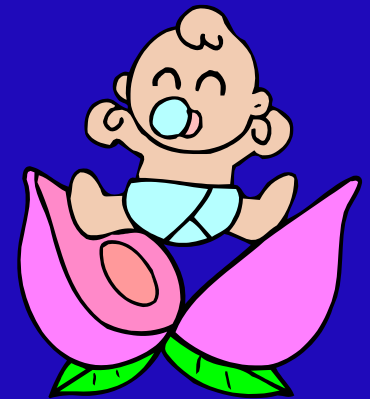
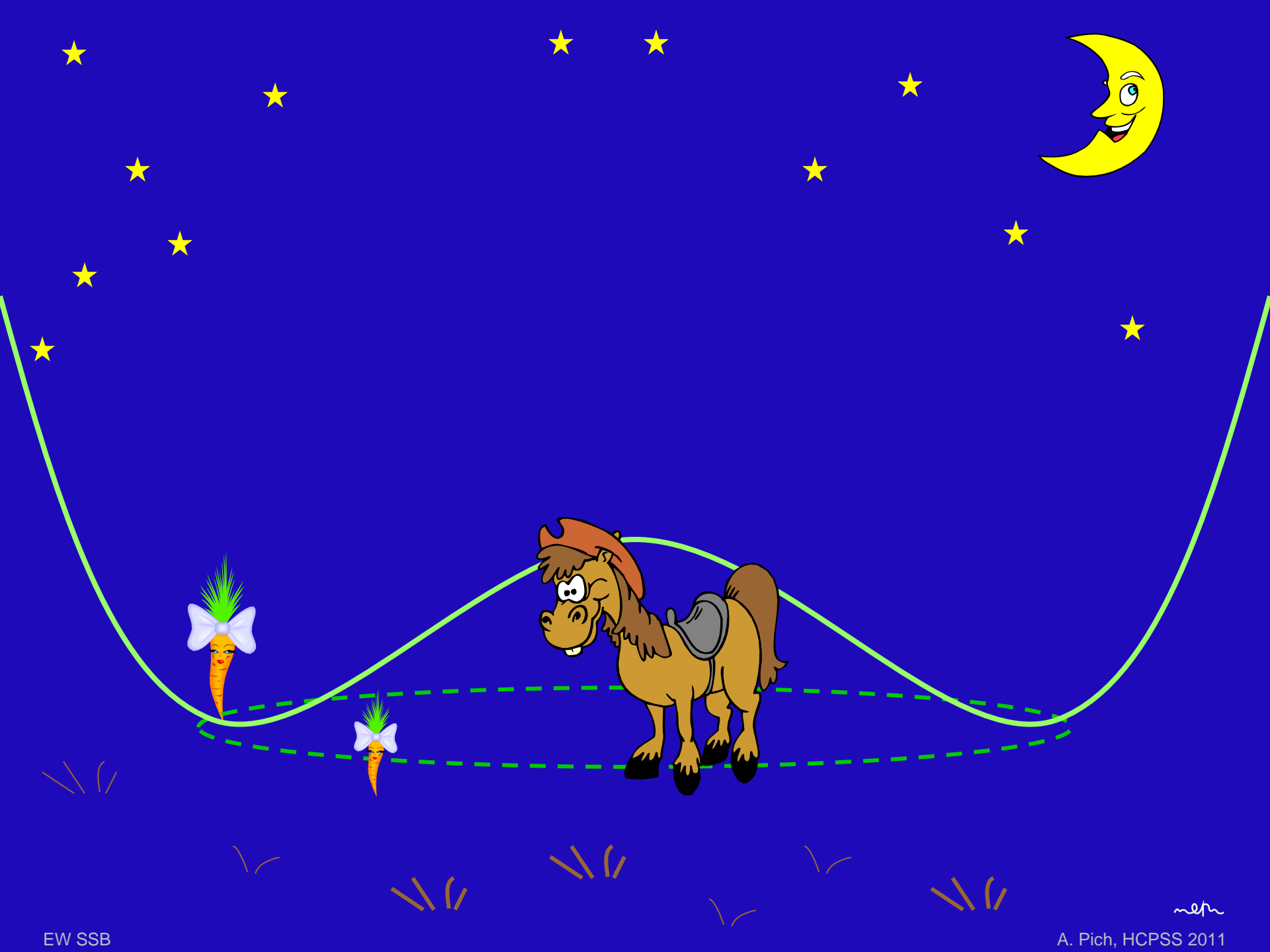


2. The Higgs Boson

- Goldstone Theorem
- Higgs Mechanism
- Higgs Properties
- Higgs Searches





SCALAR POTENCIAL

$$\mathcal{L}(\phi) = \partial_{\mu}\phi^{\dagger}\partial^{\mu}\phi - V(\phi)$$

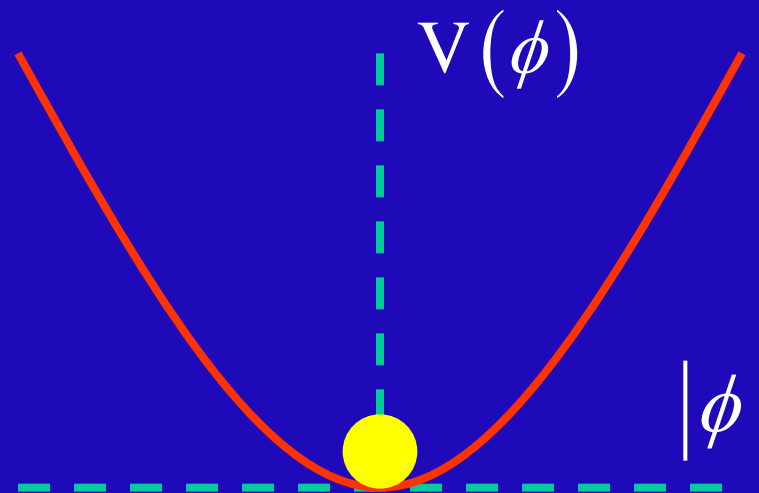
$$V(\phi) = \mu^2\phi^{\dagger}\phi + h(\phi^{\dagger}\phi)^2$$

Phase Symmetry:

$$\phi(x) \rightarrow e^{i\theta}\phi(x)$$

$$h > 0 \quad ; \quad \mu^2 > 0$$

$$M_{\phi} = \mu$$



Trivial Minimum (Ground State / Vacuum):

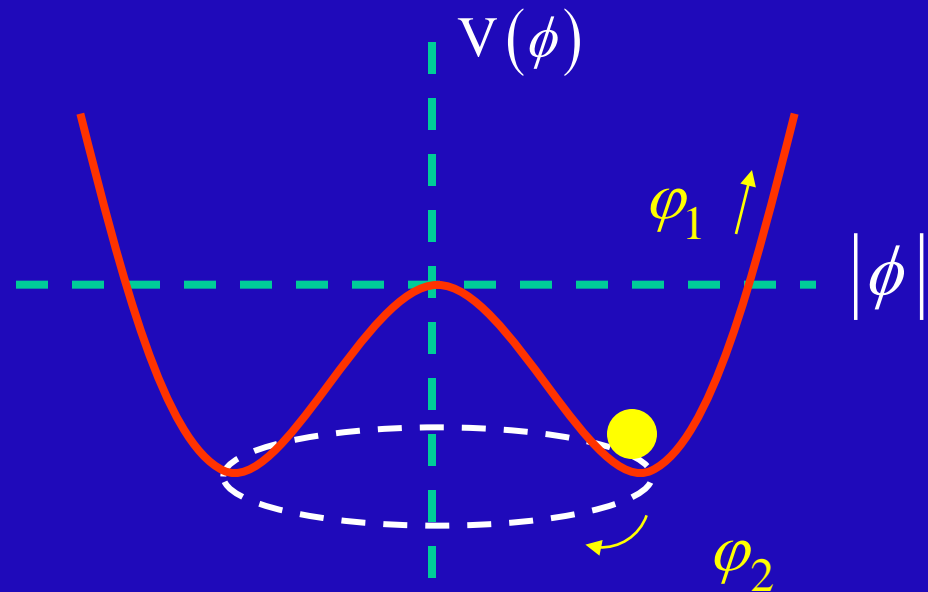
$$\phi = \phi_0 = 0$$

$$V(\phi) = \mu^2 \phi^\dagger \phi + h (\phi^\dagger \phi)^2$$

Phase Symmetry:

$$\phi(x) \rightarrow e^{i\theta} \phi(x)$$

$$\mu^2 < 0$$



Degenerate Minima

(Ground State / Vacuum)

$$|\phi_0| = \sqrt{\frac{-\mu^2}{2h}} \equiv \frac{v}{\sqrt{2}} > 0 \quad ; \quad V(\phi_0) = -\frac{1}{4} h v^4$$

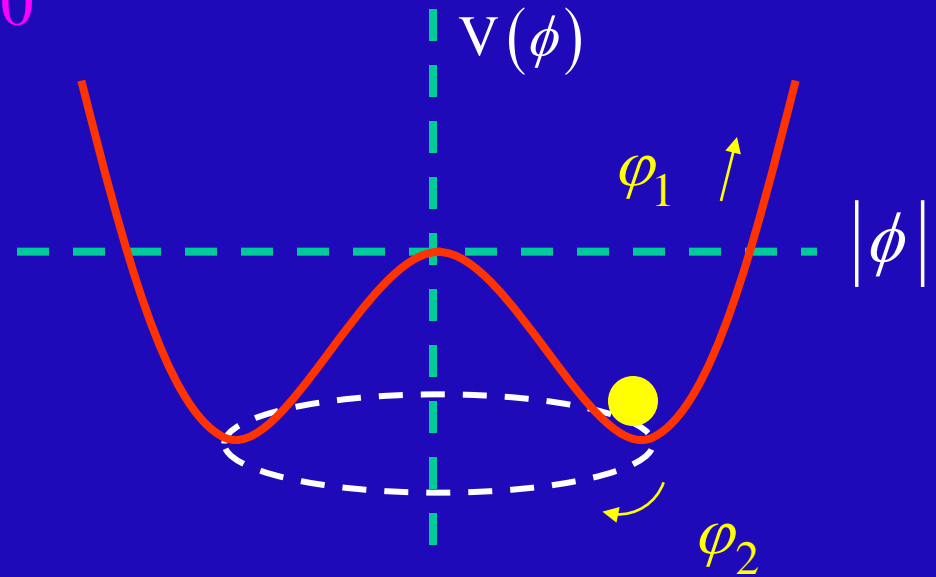
Spontaneous Symmetry Breaking:

$$\phi \equiv \frac{1}{\sqrt{2}} [v + \phi_1(x)] e^{i\phi_2(x)/v}$$

Vacuum Choice

Spontaneous Symmetry Breaking

$$\mu^2 < 0$$



$$\phi \equiv \frac{1}{\sqrt{2}} [v + \varphi_1(x)] e^{i\varphi_2(x)/v}$$



$$\mathcal{L}(\phi) = \frac{1}{2} \partial_\mu \varphi_1 \partial^\mu \varphi_1 + \frac{1}{2} \left(1 + \frac{\varphi_1}{v}\right)^2 \partial_\mu \varphi_2 \partial^\mu \varphi_2 - V(\phi)$$

$$V(\phi) = V(\phi_0) + \frac{1}{2} M_{\varphi_1}^2 \varphi_1^2 + h v \varphi_1^3 + \frac{1}{4} h \varphi_1^4$$

$$M_{\varphi_1}^2 = -2\mu^2 > 0 \quad ; \quad M_{\varphi_2}^2 = 0$$

**1 Massless
Goldstone Boson**

ELECTROWEAK SSB

New Scalar Doublet

$$\phi(x) \equiv \begin{pmatrix} \phi^{(+)}(x) \\ \phi^{(0)}(x) \end{pmatrix} ; \quad y_\phi = Q_\phi - T_3 = \frac{1}{2}$$

$$\mathcal{L}(\phi) = (\mathbf{D}_\mu \phi)^\dagger \mathbf{D}^\mu \phi - \mu^2 \phi^\dagger \phi - h (\phi^\dagger \phi)^2$$

$$\mathbf{D}^\mu \phi = \left[\partial^\mu + i g \mathbf{W}^\mu + i g' y_\phi B^\mu \right] \phi ; \quad \mathbf{W}^\mu = \frac{\vec{\tau}}{2} \cdot \vec{W}^\mu$$

$SU(2)_L \otimes U(1)_Y$ Symmetry

Degenerate Vacuum States:

$$(\mu^2 < 0, h > 0)$$

$$\left| \langle 0 | \phi^{(0)} | 0 \rangle \right| = \sqrt{\frac{-\mu^2}{2h}} \equiv \frac{v}{\sqrt{2}}$$

Spontaneous Symmetry Breaking:

$$\phi(x) = \exp \left\{ i \frac{\vec{\tau}}{2} \cdot \vec{\theta}(x) \right\} \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix}$$

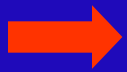
HIGGS MECHANISM

$$\phi(x) = \exp \left\{ i \frac{\vec{\tau}}{2} \cdot \vec{\theta}(x) \right\} \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix}$$

$SU(2)_L$ Invariance \longrightarrow $\vec{\theta}(x)$ can be gauged away

Unitary Gauge: $\vec{\theta}(x) = 0$

$$(\mathbf{D}_\mu \phi)^\dagger \mathbf{D}^\mu \phi \longrightarrow \frac{1}{2} \partial_\mu H \partial^\mu H + \frac{g^2}{4} (v + H)^2 \left\{ W_\mu^\dagger W^\mu + \frac{1}{2 \cos^2 \theta_W} Z_\mu Z^\mu \right\}$$



$$M_Z \cos \theta_W = M_W = \frac{1}{2} v g$$

**Massive
Gauge Bosons**

Bosonic Degrees of Freedom

Massless W^\pm, Z

$$3 \times 2 \text{ polarizations} = 6$$

+

3 Goldstones $\vec{\theta}$

SSB



Massive W^\pm, Z

$$3 \times 3 \text{ polarizations} = 9$$

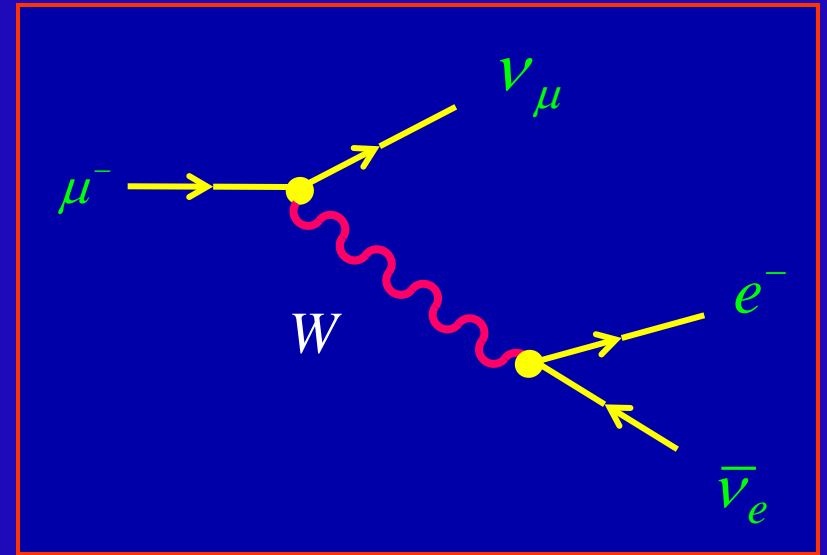
**SAME
PHYSICS**

$$M_Z \cos \theta_W = M_W = \frac{1}{2} v g$$

$$M_Z = 91.1875 \text{ GeV} > M_W = 80.399 \text{ GeV} \quad \longrightarrow \quad \sin^2 \theta_W = 1 - \frac{M_W^2}{M_Z^2} = 0.223$$

$$\frac{g^2}{M_W^2 - q^2} \approx \frac{g^2}{M_W^2} \equiv 4\sqrt{2} G_F$$

$$\frac{1}{\tau_\mu} \equiv \Gamma = \frac{G_F^2 m_\mu^5}{192 \pi^3}$$



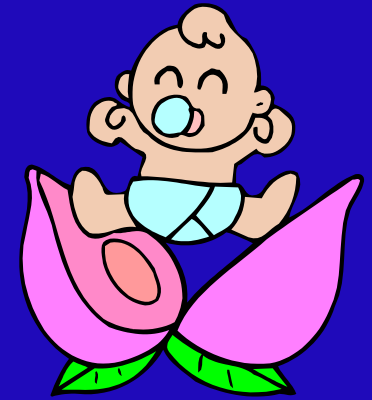
$$G_F = 1.16637 \times 10^{-5} \text{ GeV}^{-2}$$

$$g = \frac{e}{\sin \theta_W}, \quad M_W$$

$$\sin^2 \theta_W = 0.215$$

$$v = \left(\sqrt{2} G_F \right)^{-1/2} = 246 \text{ GeV}$$

THE HIGGS BOSON



$$\mathcal{L}_S = \frac{h v^4}{4} + \mathcal{L}_H + \mathcal{L}_{HG^2}$$

$$\mathcal{L}_H = \frac{1}{2} \partial_\mu H \partial^\mu H - \frac{1}{2} M_H^2 H^2 - \frac{M_H^2}{2v} H^3 - \frac{M_H^2}{8v^2} H^4$$

$$\mathcal{L}_{HG^2} = \left[M_W^2 W_\mu^\dagger W^\mu + \frac{1}{2} M_Z^2 Z_\mu Z^\mu \right] \left\{ 1 + \frac{2}{v} H + \frac{H^2}{v^2} \right\}$$

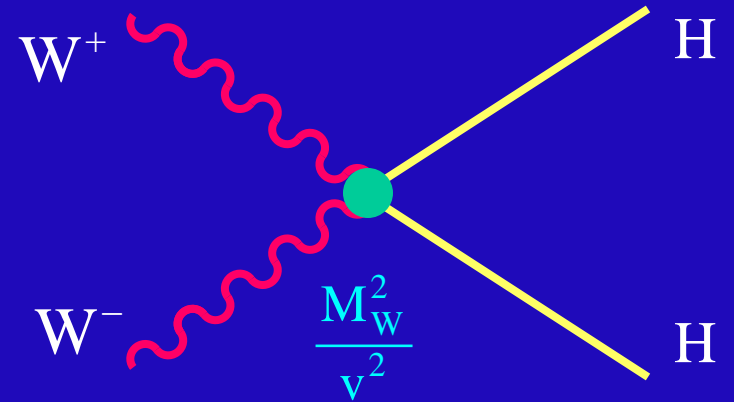
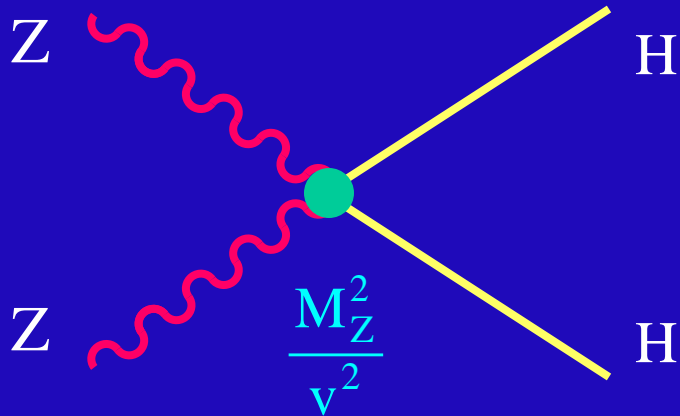
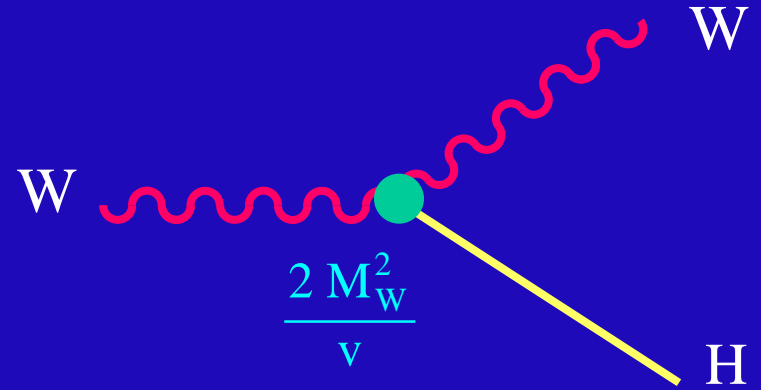
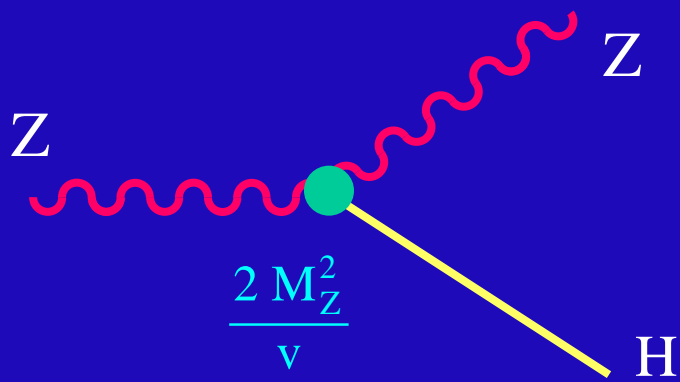
1 Scalar Particle H^0 to be Discovered

$$M_H = \sqrt{-2 \mu^2} = \sqrt{2 h} v$$

Free Parameter

LEP: 114.4 GeV < M_H < 158 GeV (95% CL)
(Direct) (Indirect)

Higgs Couplings \propto Masses



$$v = \left(\sqrt{2} G_F \right)^{-1/2} = 246 \text{ GeV}$$

FERMION MASSES

Scalar – Fermion Couplings allowed by Gauge Symmetry

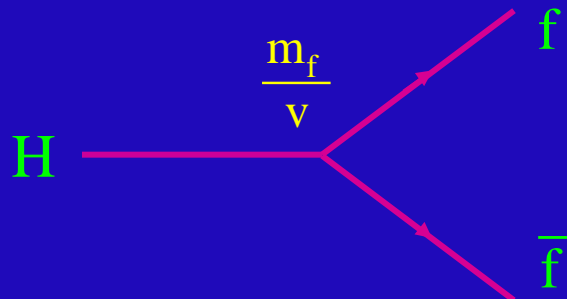
$$\mathcal{L}_Y = (\bar{q}_u, \bar{q}_d)_L \left[c^{(d)} \begin{pmatrix} \phi^{(+)} \\ \phi^{(0)} \end{pmatrix} (q_d)_R + c^{(u)} \begin{pmatrix} \phi^{(0)\dagger} \\ -\phi^{(+)\dagger} \end{pmatrix} (q_u)_R \right] + (\bar{\nu}_l, \bar{l})_L c^{(l)} \begin{pmatrix} \phi^{(+)} \\ \phi^{(0)} \end{pmatrix} l_R + \text{h.c.}$$

SSB

$$\mathcal{L}_Y = - \left(1 + \frac{H}{v} \right) \left\{ m_{q_d} \bar{q}_d q_d + m_{q_u} \bar{q}_u q_u + m_l \bar{l} l \right\}$$

Fermion Masses are
New Free Parameters

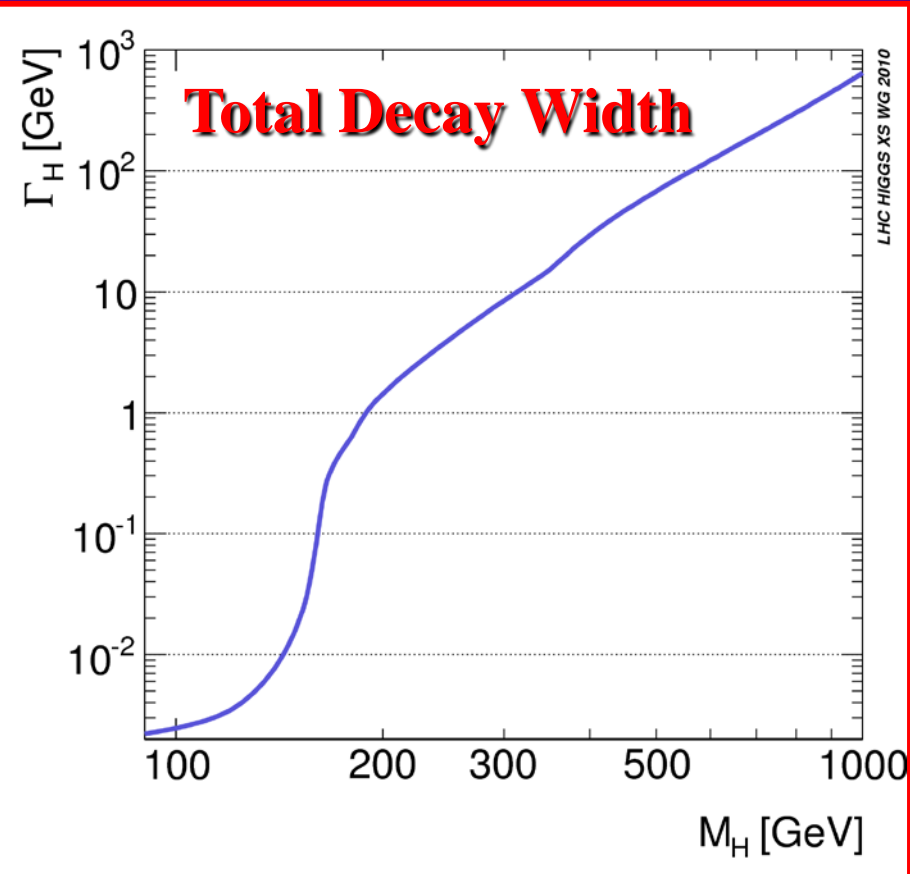
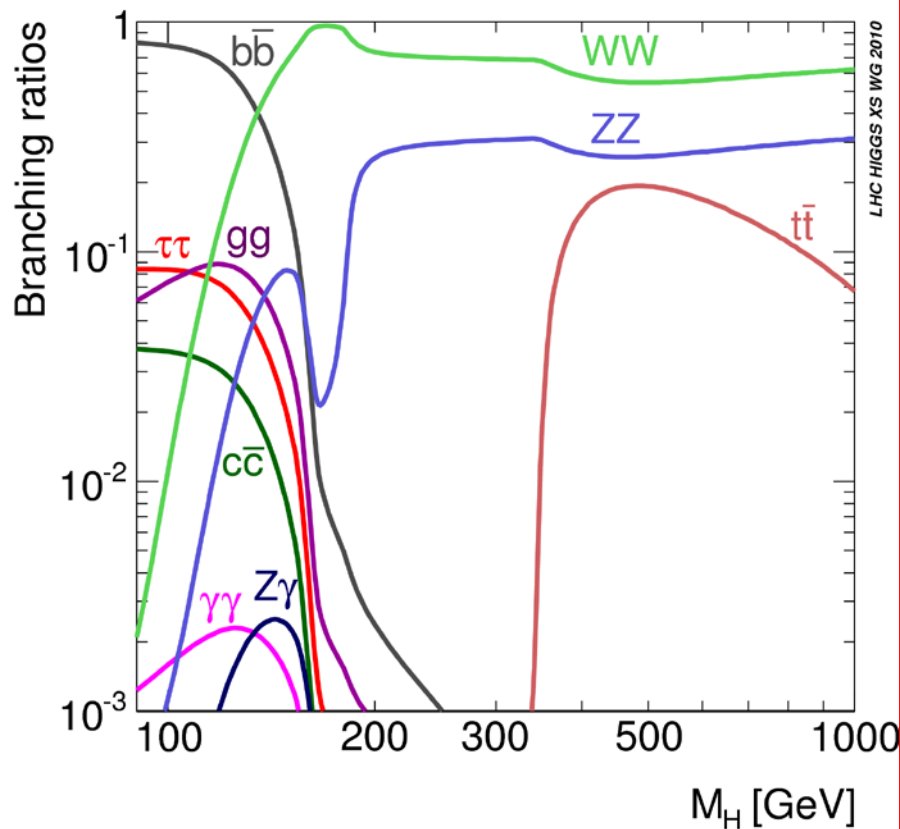
$$\left[m_{q_d}, m_{q_u}, m_l \right] = - \left[c^{(d)}, c^{(u)}, c^{(l)} \right] \frac{v}{\sqrt{2}}$$



Couplings Fixed:

$$g_{Hf\bar{f}} = \frac{m_f}{v}$$

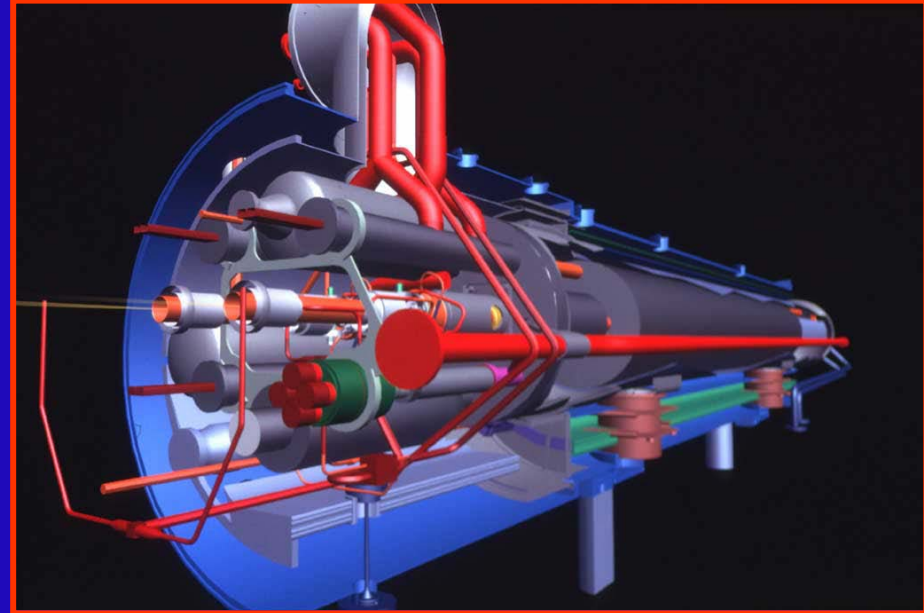
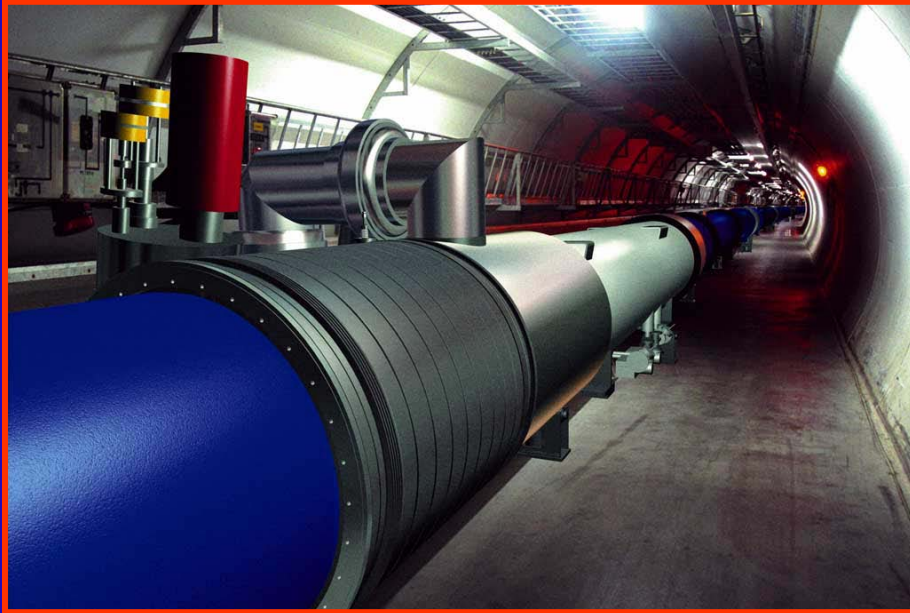
Searching for the HIGGS



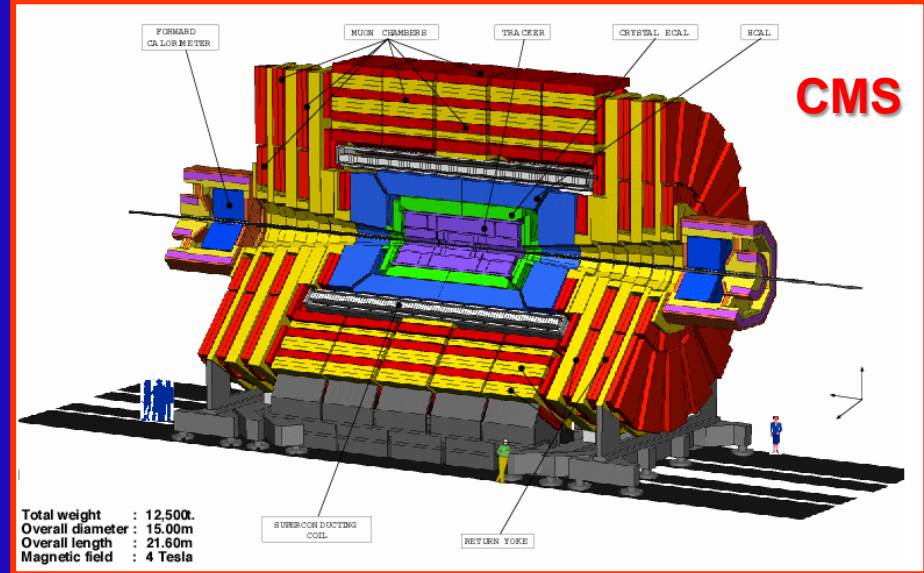
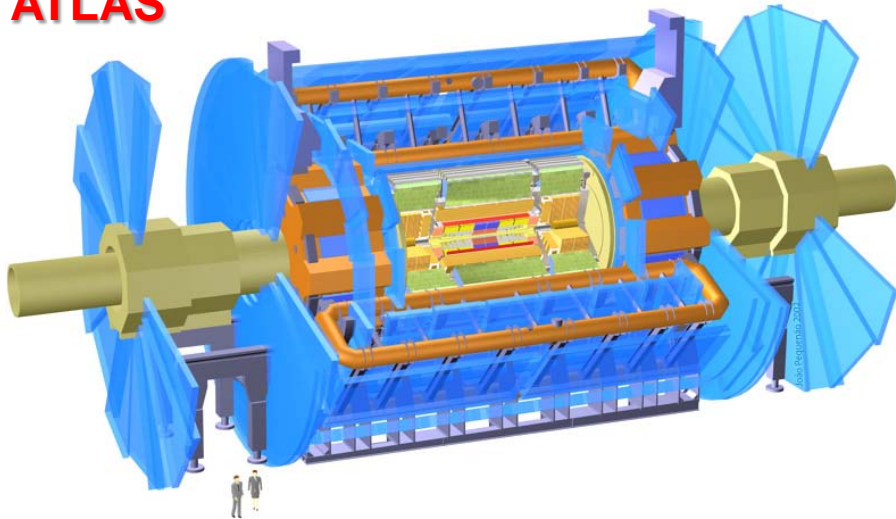
Interaction proportional
to mass (M_W^2, M_Z^2, m_f)

The Higgs decays into the
heaviest possible particles

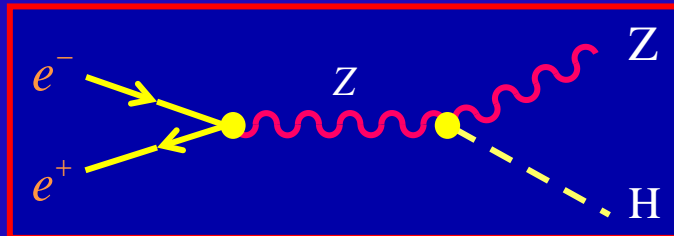
The Large Hadron Collider



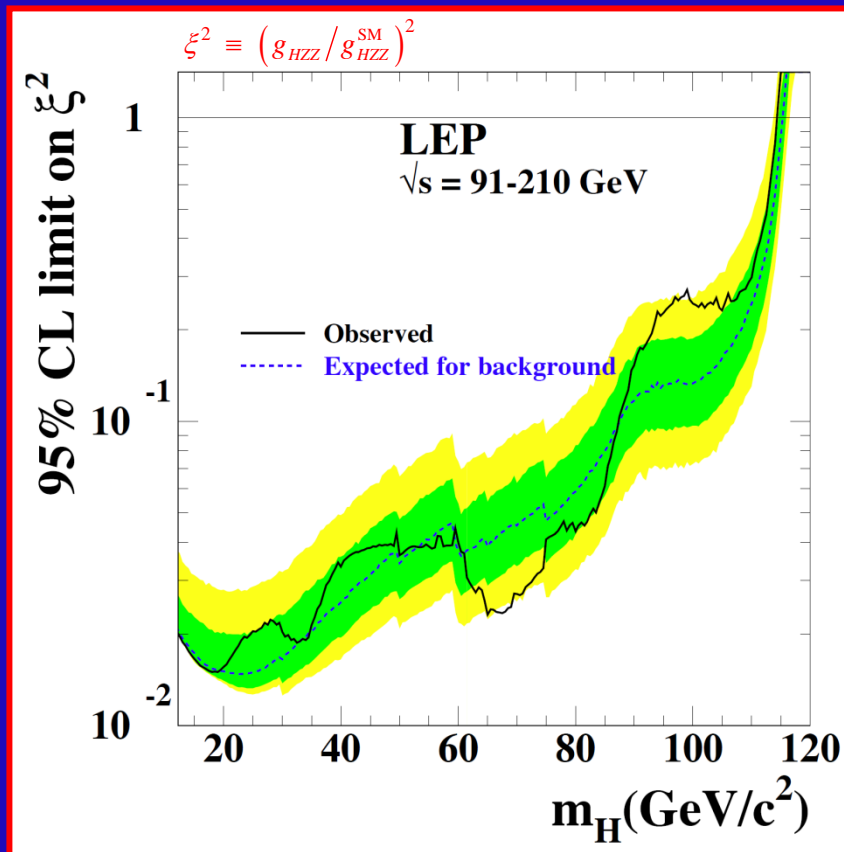
ATLAS



LEP Searches



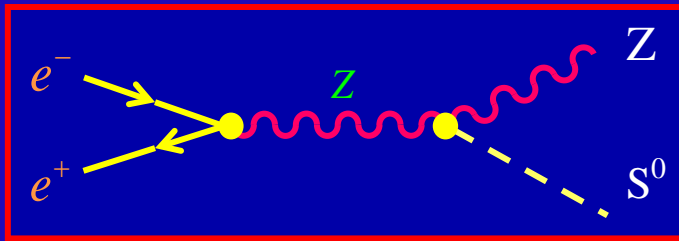
- 1) $H \rightarrow b\bar{b}$, $Z \rightarrow q\bar{q}$
- 2) $H \rightarrow \tau^+\tau^-$, $Z \rightarrow q\bar{q}$
 $H \rightarrow b\bar{b}$, $H \rightarrow \tau^+\tau^-$
- 3) $H \rightarrow b\bar{b}$, $Z \rightarrow \nu\bar{\nu}$
- 4) $H \rightarrow b\bar{b}$, $Z \rightarrow e^+e^-, \mu^+\mu^-$



$M_H > 114.4 \text{ GeV}$ (95% CL)

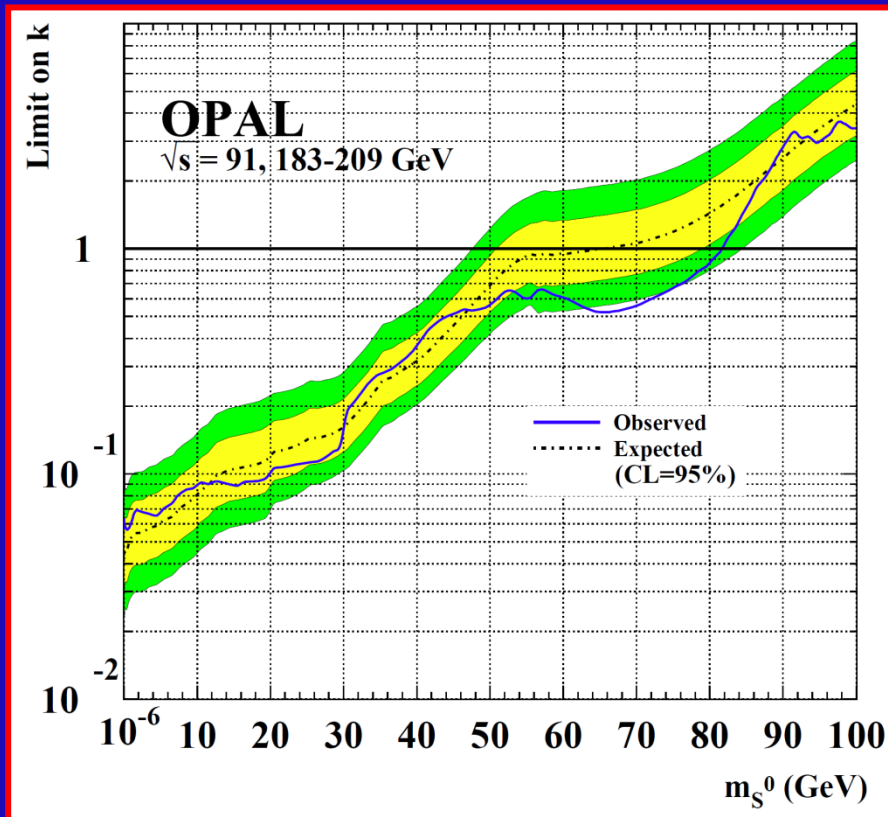
3σ excess @ 115 GeV (ALEPH)

Not seen at DELPHI, L3 and OPAL



- 1) $S^0 \rightarrow \text{all}, Z \rightarrow e^+e^-, \mu^+\mu^-$
- 2) $S^0 \rightarrow e^+e^-, Z \rightarrow \nu\bar{\nu}$

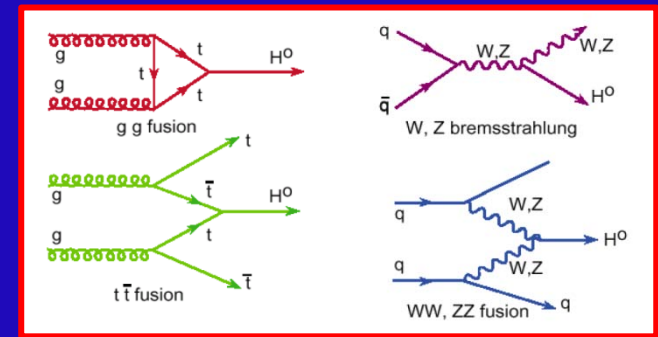
$$k \equiv \sigma(e^+e^- \rightarrow ZS^0) / \sigma(e^+e^- \rightarrow ZH)_{\text{SM}}$$



Low-mass Higgs excluded

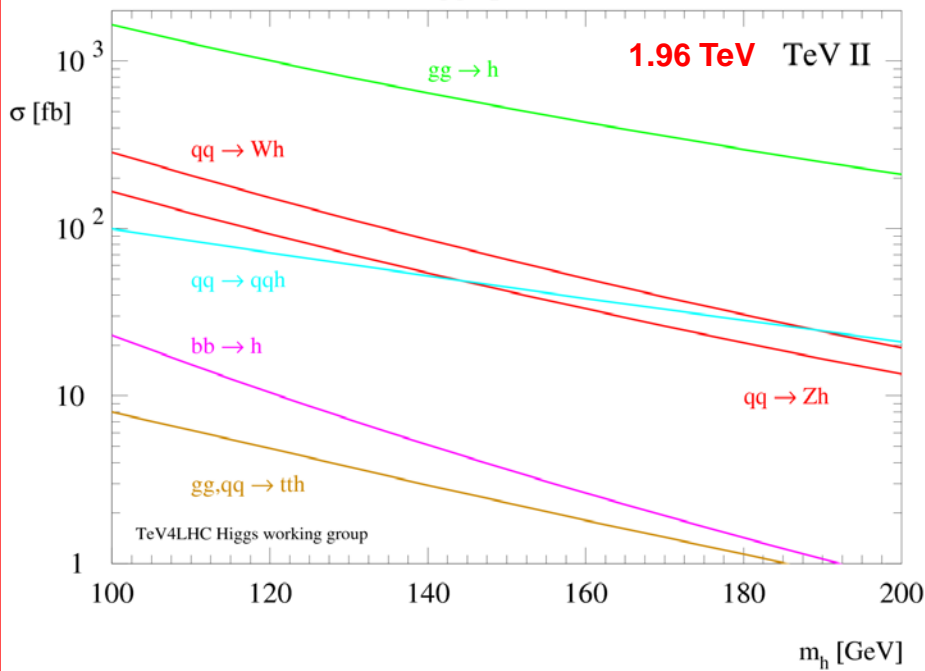
Decay-mode-independent limits

Higgs Production @ Colliders

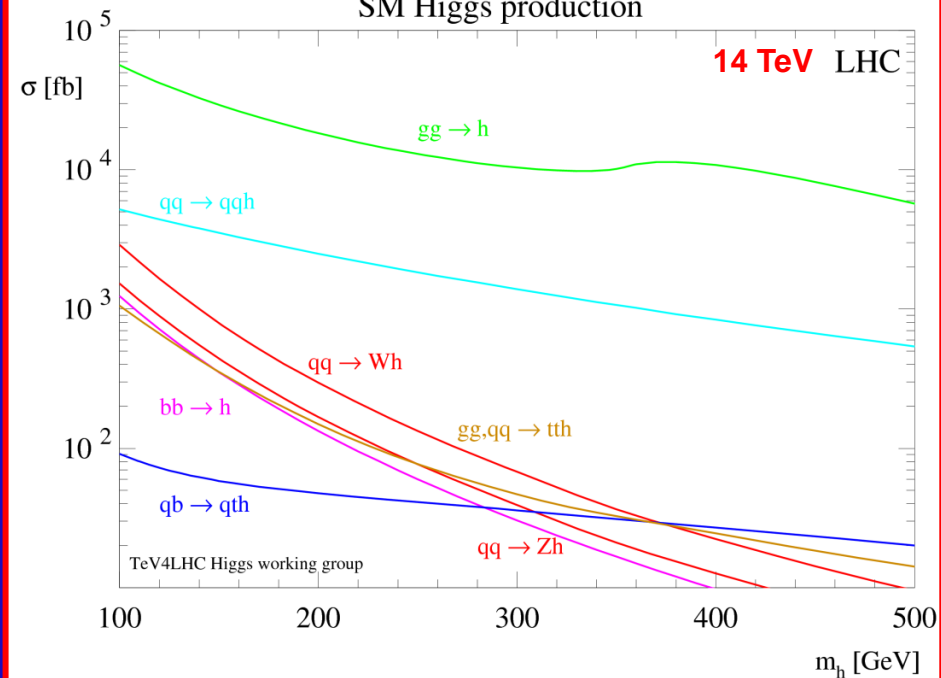


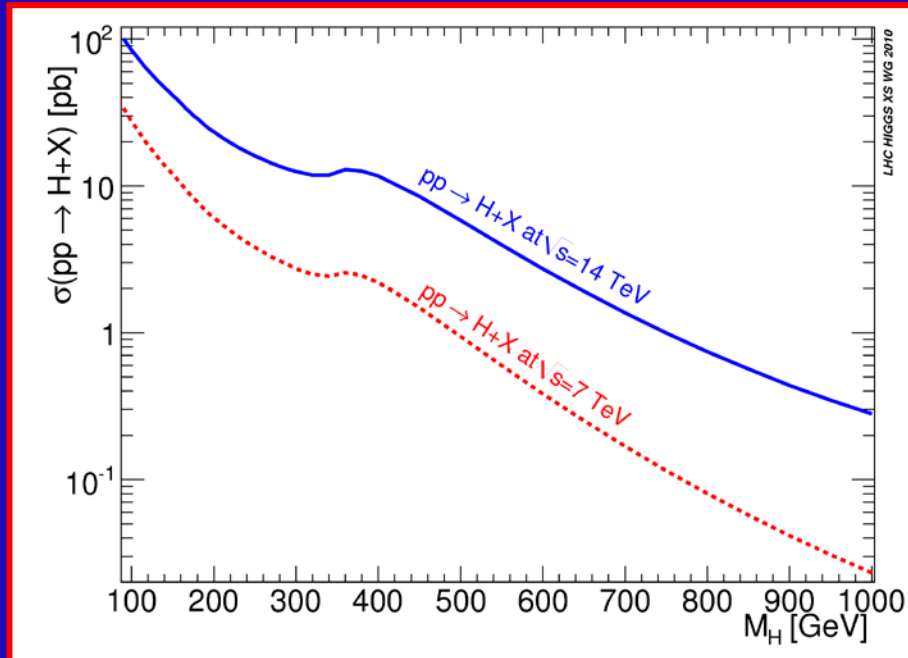
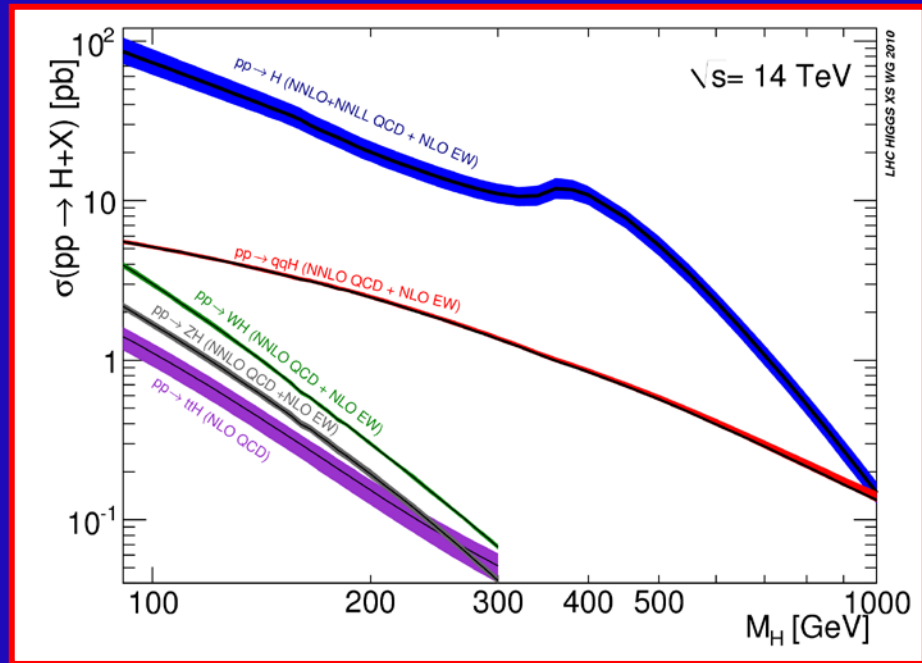
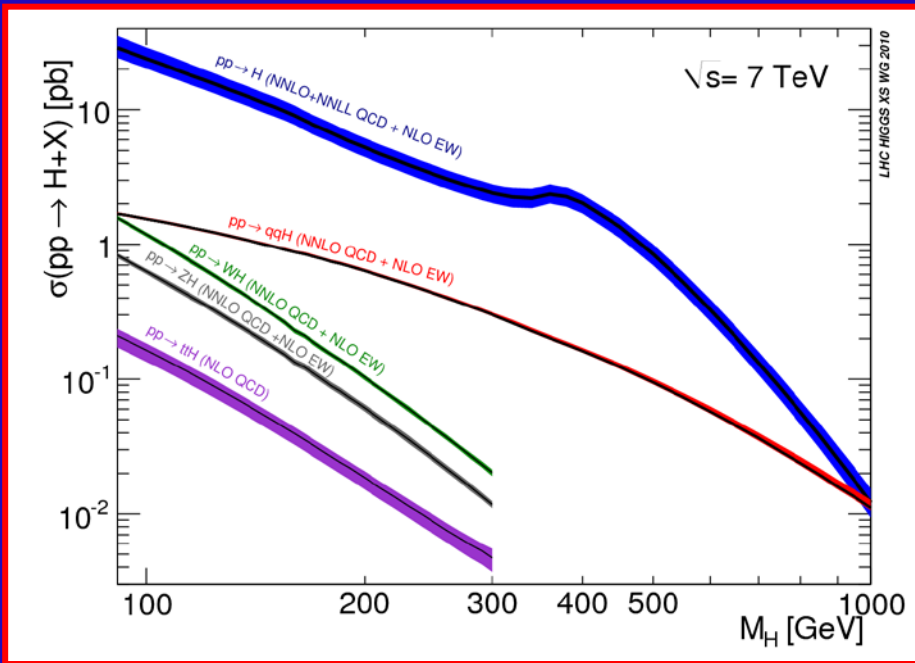
F. Maltoni

SM Higgs production



SM Higgs production



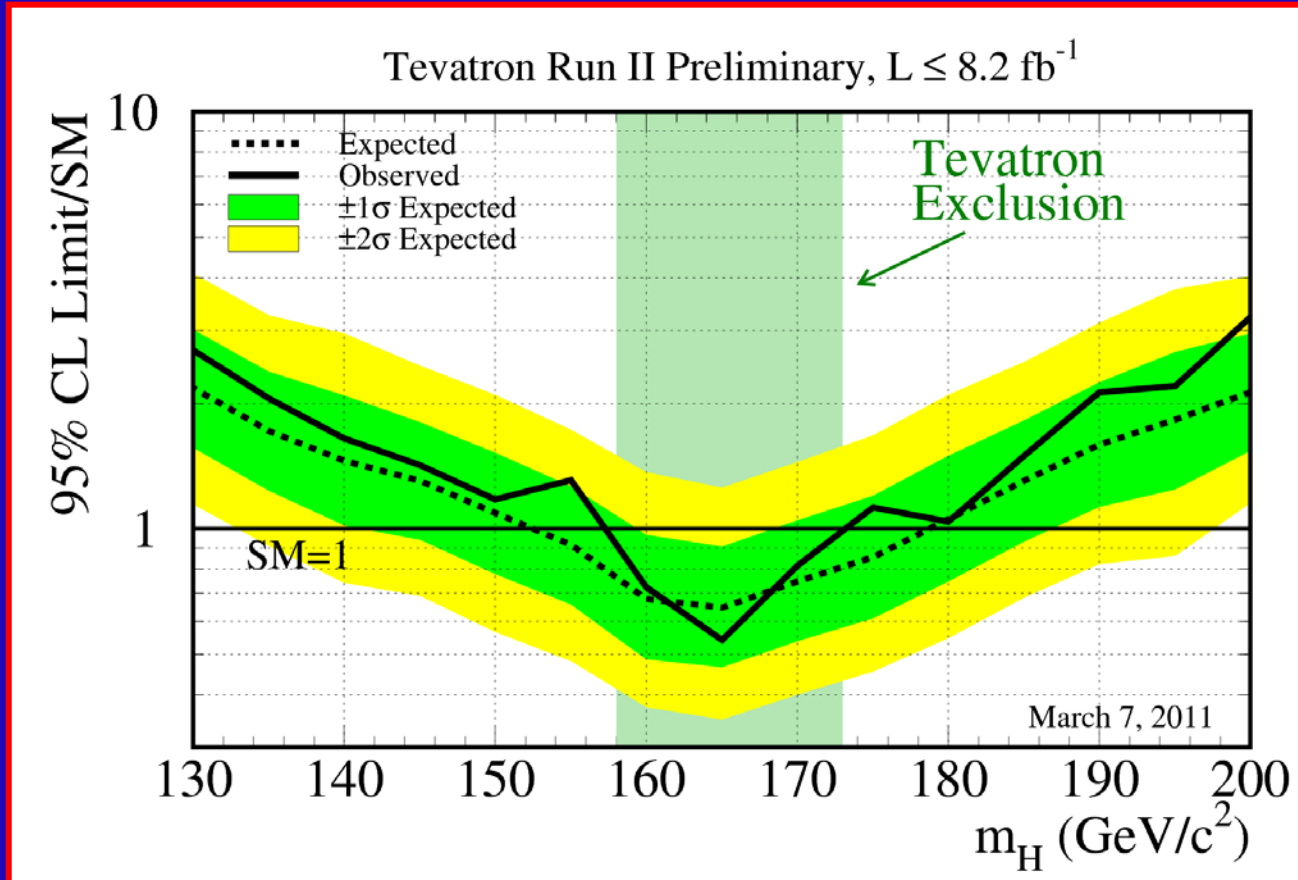


Tevatron Limits

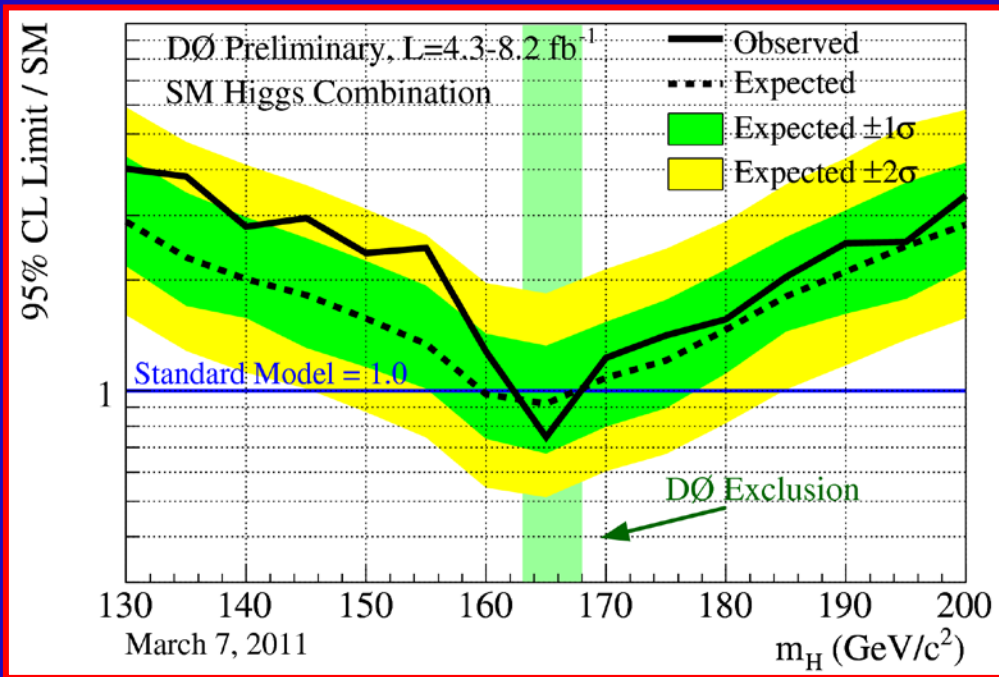
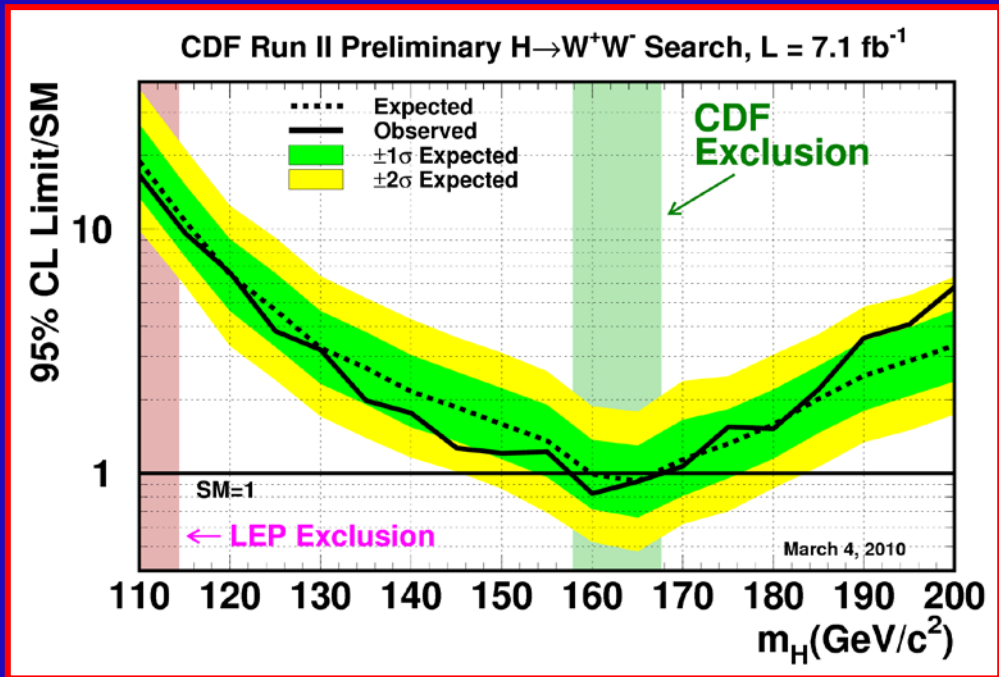
arXiv:1103.3233

$L \leq 7.1 \text{ fb}^{-1}$ (CDF), 8.2 fb^{-1} (D0)

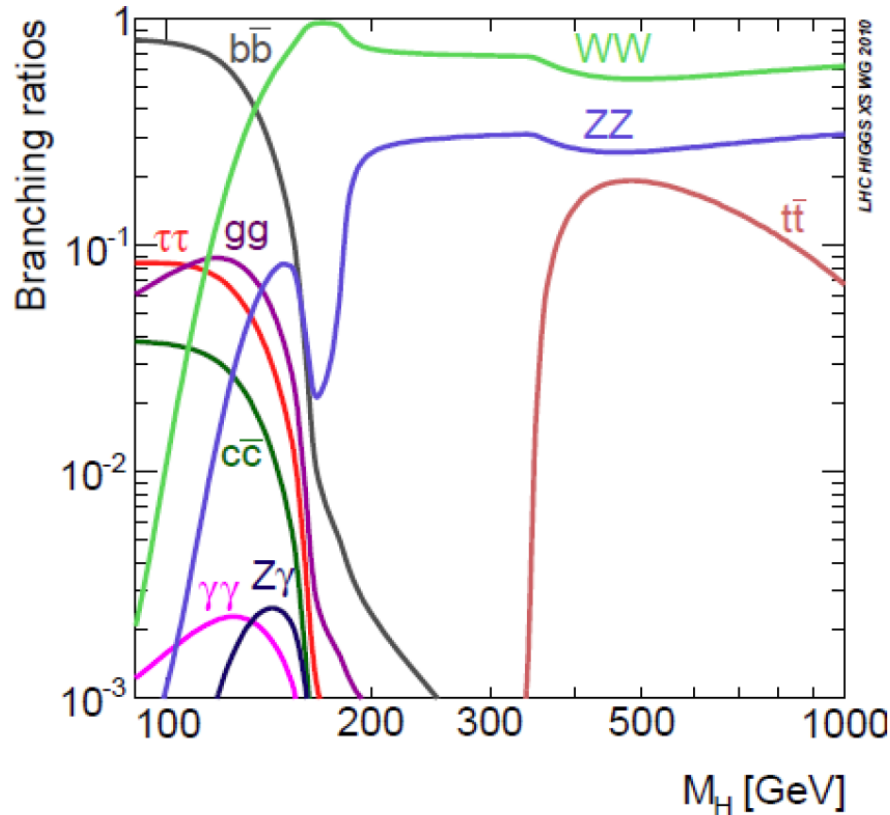
$H \rightarrow W^+W^-$



$M_H \notin [158, 173] \text{ GeV}$ (95% CL)



SM Higgs Search Channels

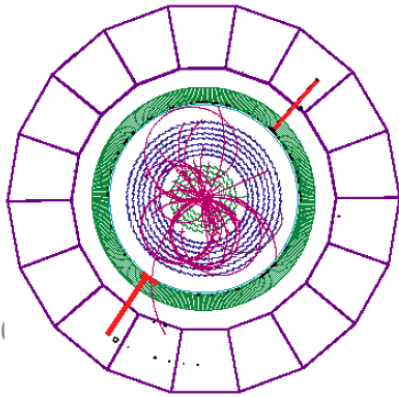
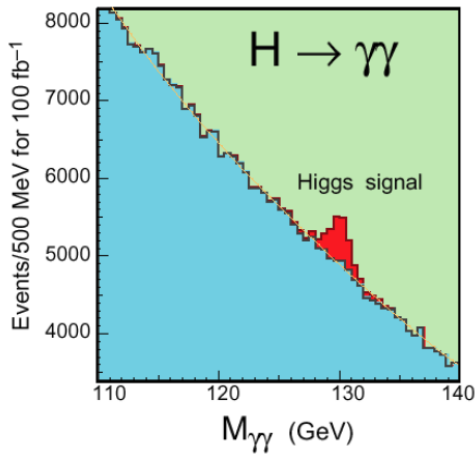
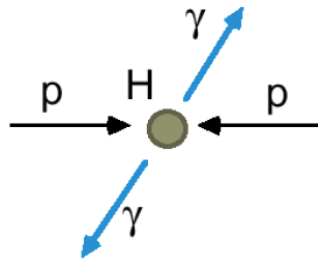


Channels included	Higgs mass range used in analyses (GeV)
$H \rightarrow \gamma\gamma$	115-150
VBF $H \rightarrow \tau\tau$	115-145
VH, $H \rightarrow b\bar{b}$ (highly boosted)	115-125
VH, $H \rightarrow WW \rightarrow l\nu jj$	130-200
$H \rightarrow WW \rightarrow 2l2\nu + 0/1$ jets	120-600
VBF $H \rightarrow WW \rightarrow 2l2\nu$	130-500
$H \rightarrow ZZ \rightarrow 4l$	120-600
$H \rightarrow ZZ \rightarrow 2l2\nu$	200-600
$H \rightarrow ZZ \rightarrow 2l2b$	300-600

Channels with $H \rightarrow \gamma\gamma$, $H \rightarrow \tau\tau$, $H \rightarrow WW^*$, $H \rightarrow ZZ^*$ are all used for the search
 $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ are the channels where mass can be measured with $\sim 1\%$ res.

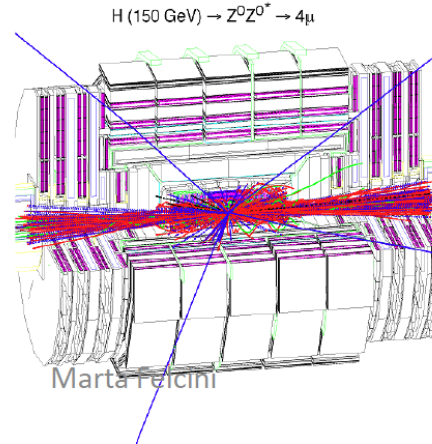
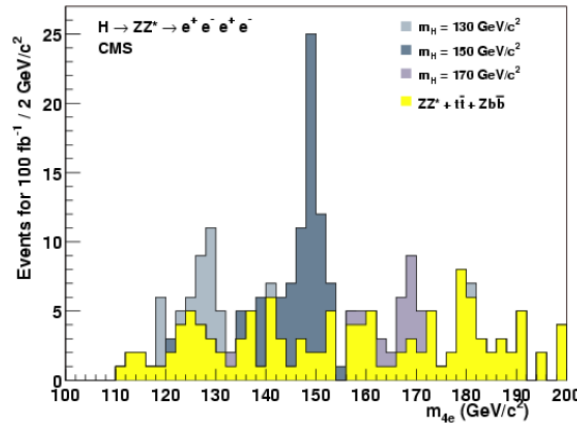
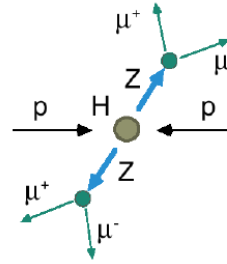
SM Higgs Search Strategies - Examples

Low $M_H < 140 \text{ GeV}$

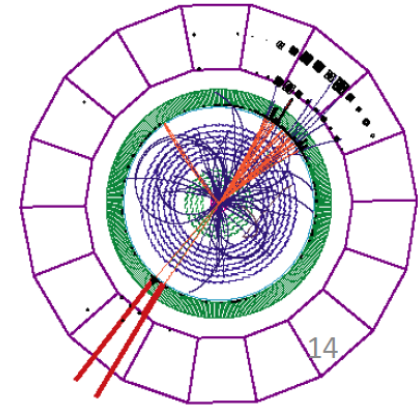
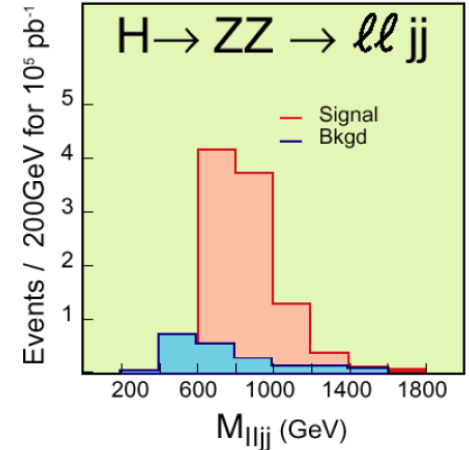
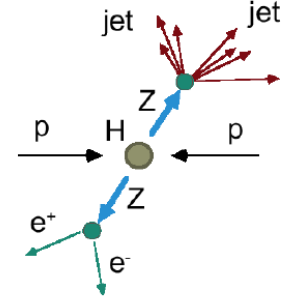


2/2/21

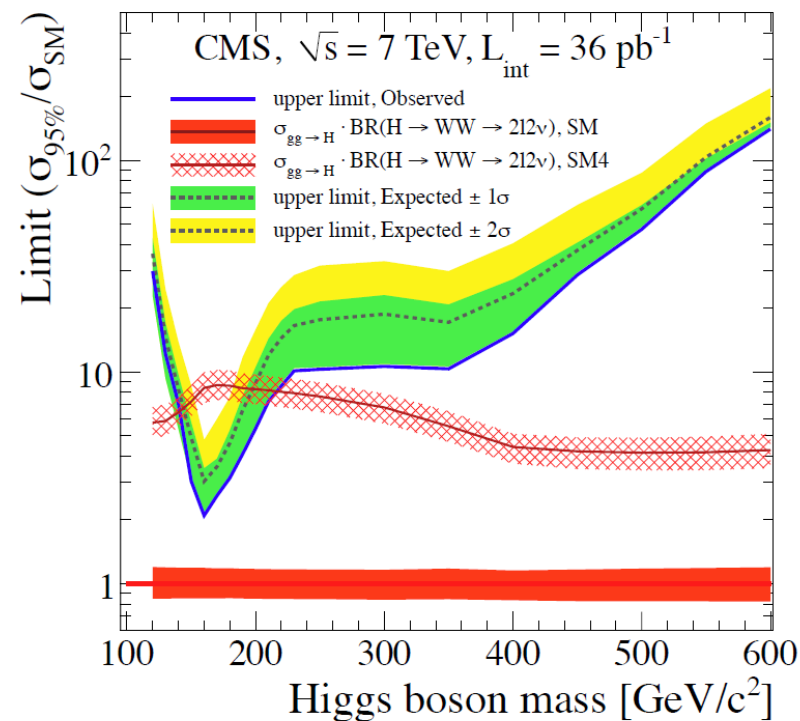
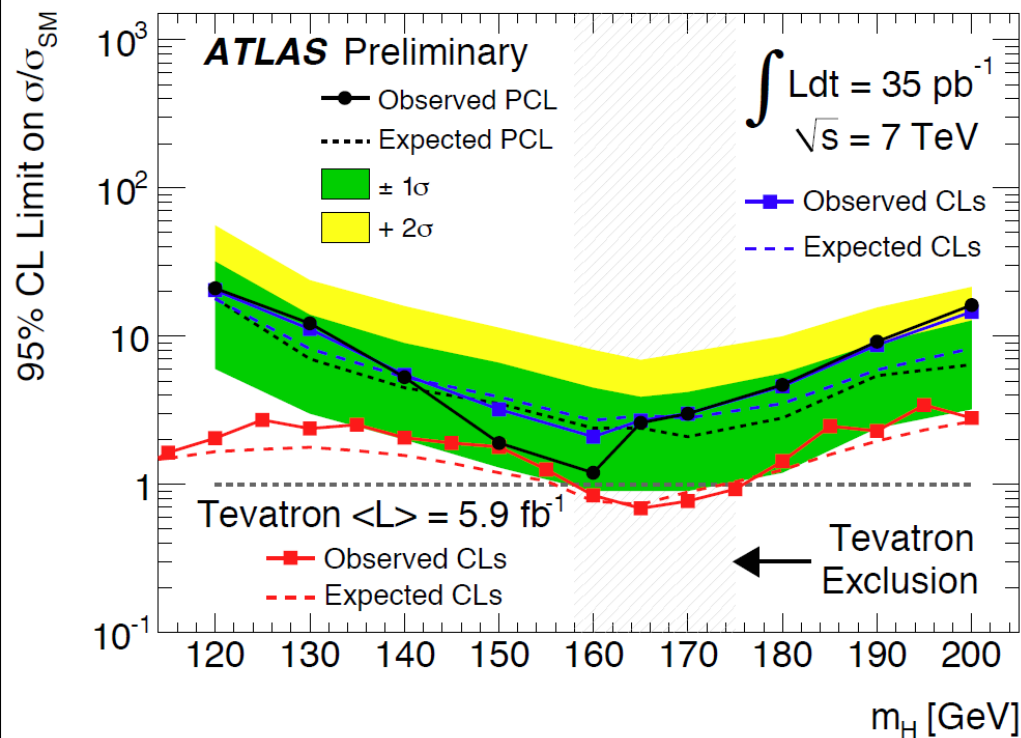
Medium $130 < M_H < 500 \text{ GeV}$



High $M_H > \sim 500 \text{ GeV}/c^2$



H → WW → lνlν : Exclusion Limits



⇒ At $M_H = 160 \text{ GeV}$

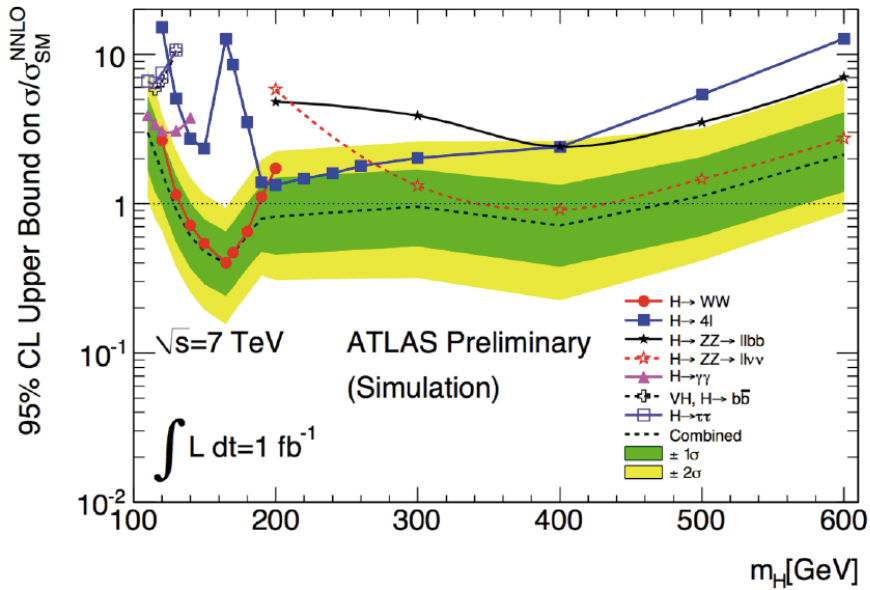
→ ATLAS excludes $2.1 \times \sigma_{\text{SM}}$ (obs) and $2.7 \times \sigma_{\text{SM}}$ (exp) CLs

→ CMS excludes $2.1 \times \sigma_{\text{SM}}$ (obs) and $3.0 \times \sigma_{\text{SM}}$ (exp) bayesian

⇒ Neither experiment excludes any M_H in the Standard Model

⇒ Assuming a heavy 4th generation CMS excludes a Higgs with $144 \text{ GeV} < M_H < 207 \text{ GeV}$

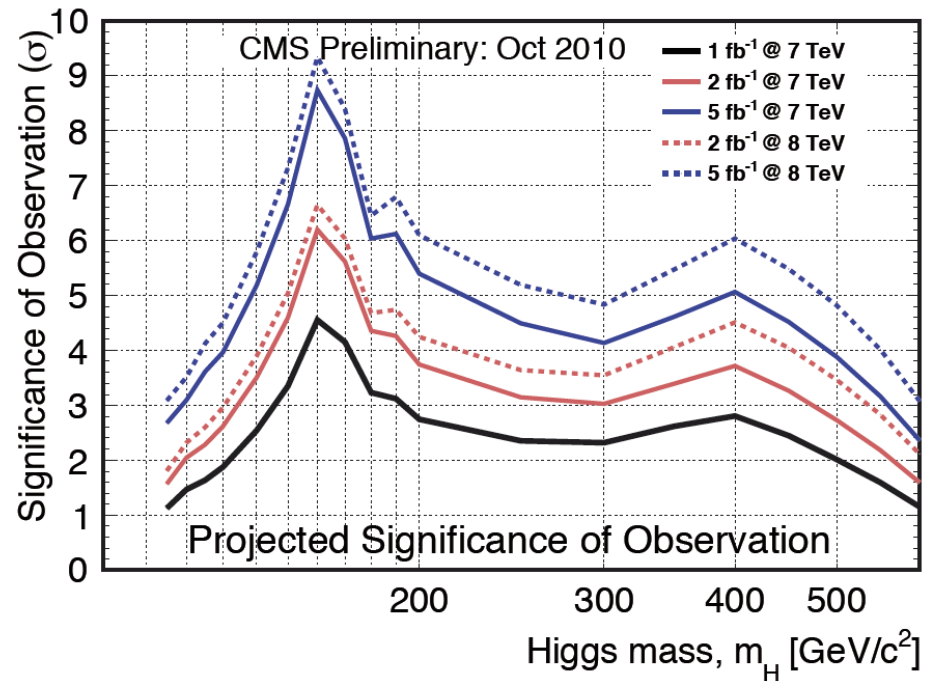
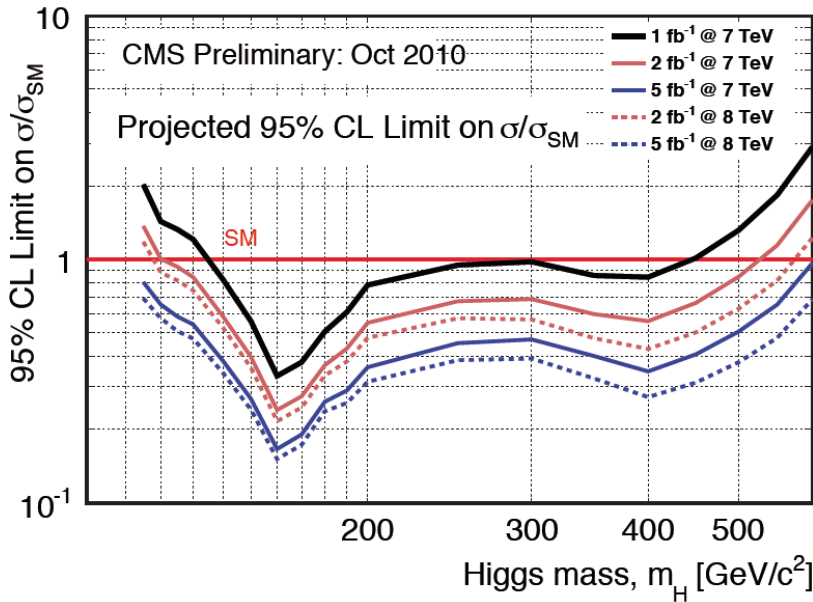
Higgs boson prospects for the (near) future



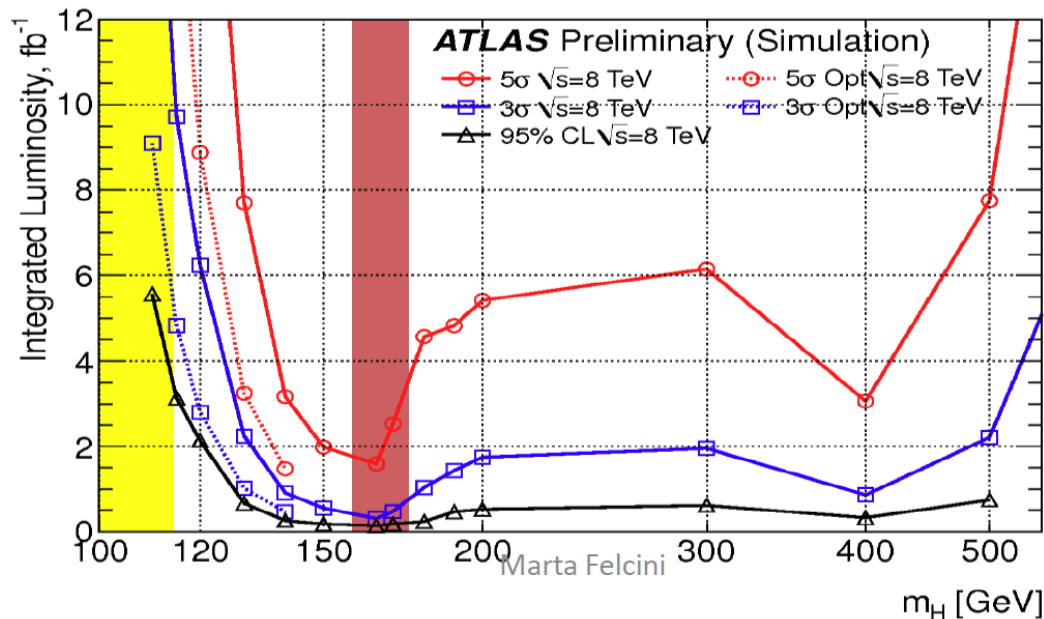
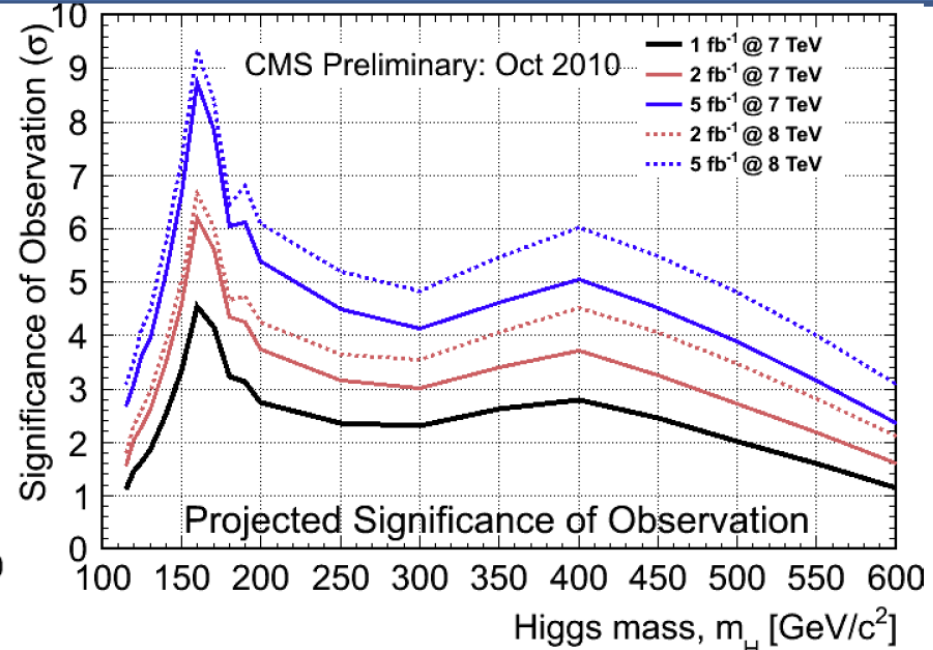
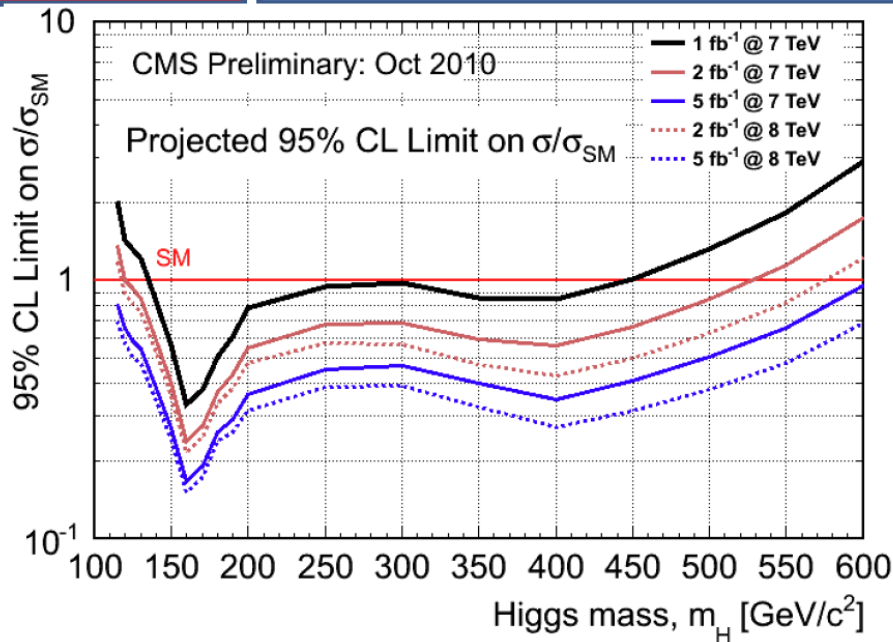
With $\sim 1 \text{ fb}^{-1}$ for EPS and 2 fb^{-1} by end of summer for each experiment
 \Rightarrow Sensitive to SM Higgs boson cross sections this summer!

If SM Higgs is not Nature's choice we'll know by the end of this summer

Discovery more difficult \rightarrow
 BUT upward fluctuations may be observed!



SM Higgs Search Performance at 7 TeV





ner