#### Flavour Physics (II) Sixth CERN-Fermilab Hadron Collider Physics Summer School 8 June - 17 June 2011, Geneva, Switzerland

#### Tatsuya NAKADA Laboratory for High Energy Physics (LPHE) Swiss Federal Institute of Technology Lausanne (EPFL) Lausanne, Switzerland





### Plan of the lecture today

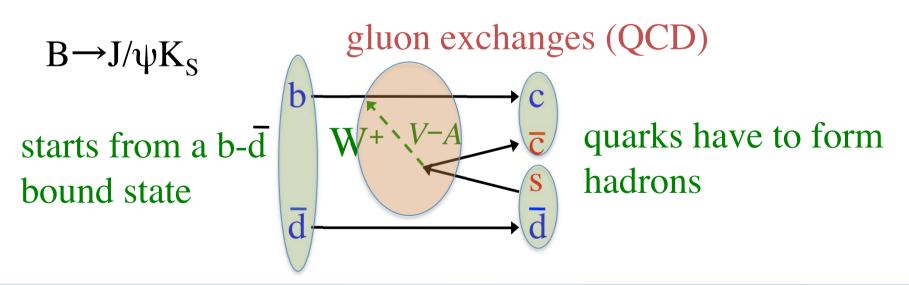
- More on Weak Decays
- Current Status of  $V_{\rm CKM}$

decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) quark decay

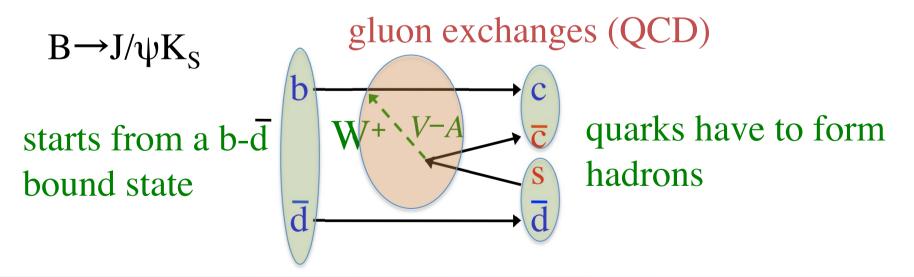
$$b \rightarrow c + W^{-}(\rightarrow \overline{c}s)$$
  $b \xrightarrow{\kappa} V^{-A} \xrightarrow{c}s$ 

decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ )

quark decay to hadron decay

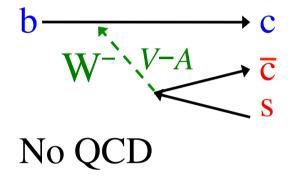


decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) Theoretical tool to describe the decay amplitude for  $M \rightarrow F$  $A(M \rightarrow F) = \langle F | H_{\text{effective}}^{\text{weak decay}} | M \rangle = \frac{G_F}{\sqrt{2}} \sum_i \xi_{\text{CKM}}^i C_i(\mu) \langle F | Q_i(\mu) | M \rangle$  $Q_i$ : quark operators



decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ )

lowest order weak interactions ( $\Delta F = 1$ )

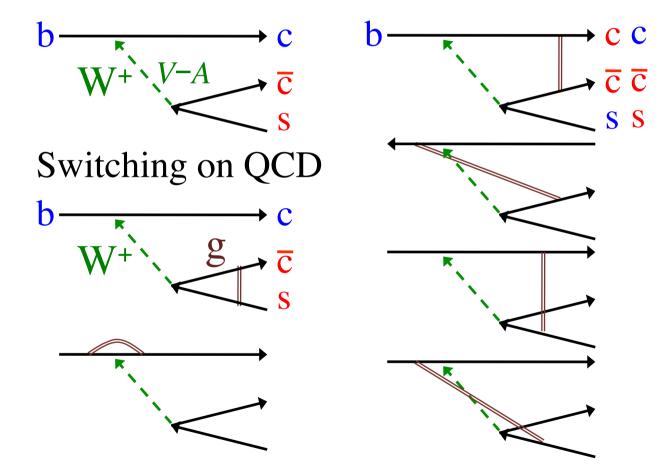


decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ )

lowest order weak interactions ( $\Delta F = 1$ )

 $(\overline{c}_i b_i)_{V-A} (\overline{s}_j c_j)_{V-A}$ 

decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) lowest order weak interactions ( $\Delta F = 1$ )

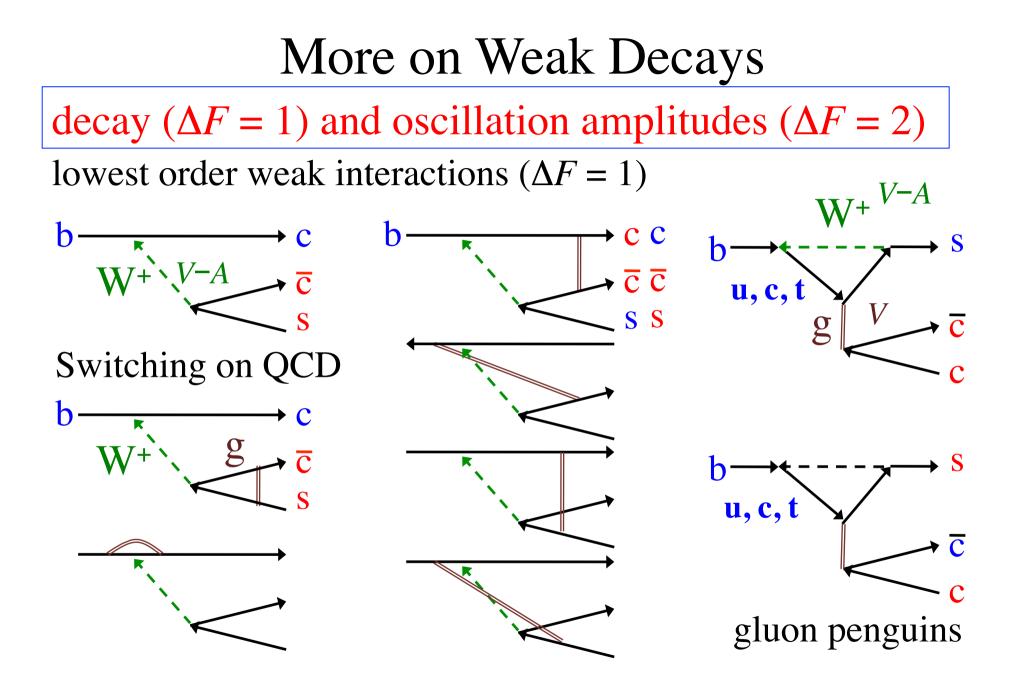


decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) lowest order weak interactions ( $\Delta F = 1$ )

$$(\overline{c}_i b_i)_{V-A} (\overline{s}_j c_j)_{V-A}$$
  
No-QCD tree diagram

+ one gluon tree diagrams with two different colour structures

$$(\overline{c}_{i} b_{i})_{V-A} (\overline{s}_{j} c_{j})_{V-A}$$
$$(\overline{c}_{j} b_{i})_{V-A} (\overline{s}_{i} c_{j})_{V-A}$$



decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) lowest order weak interactions ( $\Delta F = 1$ )

$$(\overline{c}_i b_i)_{V-A} (\overline{s}_j c_j)_{V-A}$$
  
No QCD tree diagram

+ one gluon tree diagrams with two different colour structures  $(\overline{s_i} b_i)_{V-A} (\overline{c_i} c_j)_V$ 

 $(\overline{c}_i b_i)_{V-A} (\overline{s}_i c_j)_{V-A}$  $(\overline{c}_i b_i)_{V-A} (\overline{s}_i c_j)_{V-A}$ 

+ gluon penguins with two different colour structure gluon = V

$$(\overline{s}_{j}b_{i})_{V-A}(\overline{c}_{i}c_{j})_{V}$$

decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) lowest order weak interactions ( $\Delta F = 1$ )

- $(\overline{c}_i b_i)_{V-A} (\overline{s}_j c_j)_{V-A}$ No QCD tree diagram
- + one gluon tree diagrams with two different colour structures

 $(\overline{c}_{i} b_{i})_{V-A} (\overline{s}_{j} c_{j})_{V-A}$  $(\overline{c}_{j} b_{i})_{V-A} (\overline{s}_{i} c_{j})_{V-A}$ 

es 
$$(\overline{s_i} b_i)_{V-A} (\overline{c_j} c_j)_{V-A}$$

+ gluon penguins with two different colour structure gluon = V  $\rightarrow$  split to (V-A) + (V+A)  $(\overline{s_i} b_i)_{V-A} (\overline{c_i} c_j)_{V-A}$ (needed for the Q<sup>2</sup> evolution)  $(\overline{s_i} b_i)_{V-A} (\overline{c_i} c_j)_{V+A}$ 

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decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) lowest order weak interactions ( $\Delta F = 1$ )

operators

 $Q_5$ 

 $Q_4$ 

 $Q_6$ 

#### $(\overline{c}_i b_i)_{V-A} (\overline{s}_j c_j)_{V-A}$ No QCD tree diagram

+ one gluon tree diagrams with two different colour structures

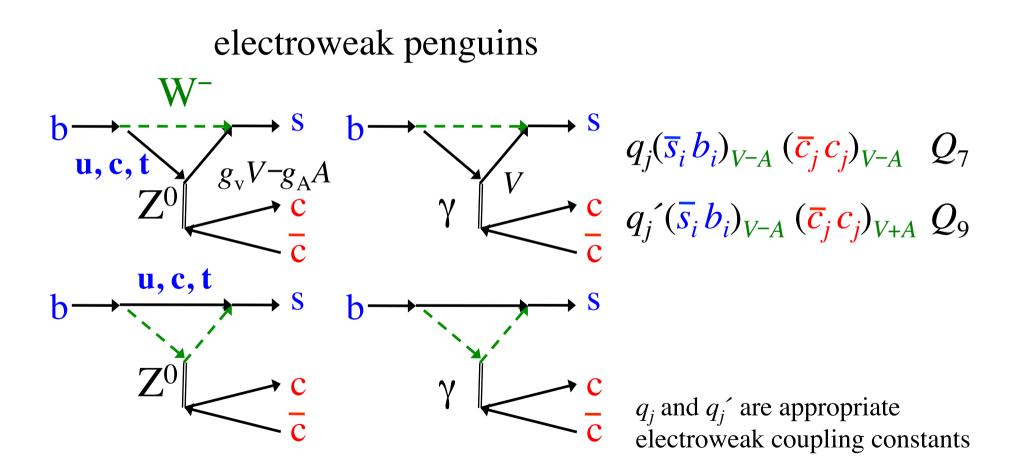
$$(\overline{c}_i b_i)_{V-A} (\overline{s}_j c_j)_{V-A} \qquad Q_2$$

$$(\overline{c}_{j} b_{i})_{V-A} (\overline{s}_{i} c_{j})_{V-A} \qquad Q_{1}$$

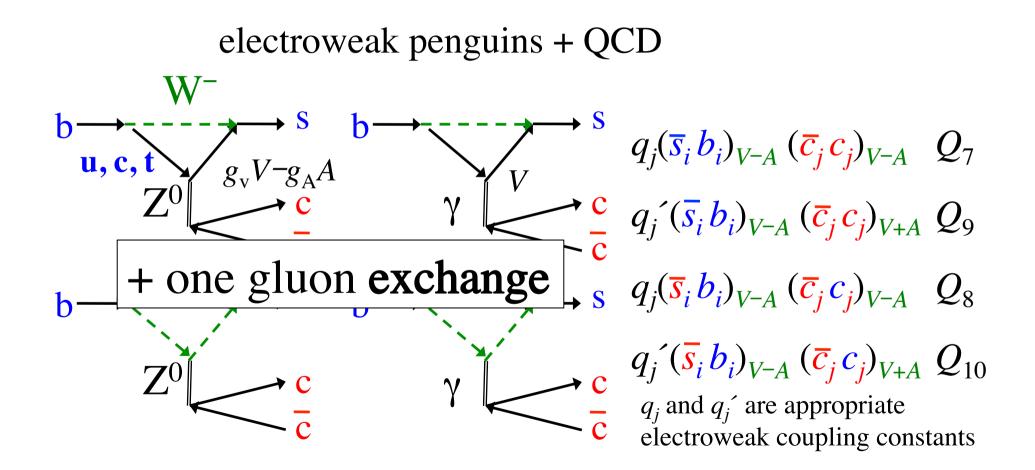
$$\lim_{i \text{ res }} (\overline{s_i} \, b_i)_{V-A} \, (\overline{c_j} \, c_j)_{V-A} \qquad Q_3$$

+ gluon penguins with two different colour structure gluon = V  $\rightarrow$  split to  $(V-A) + (V+A) (\overline{s_j} b_i)_{V-A} (\overline{c_i} c_j)_{V-A}$ (needed for the Q<sup>2</sup> evolution)  $(\overline{s_i} b_i)_{V-A} (\overline{c_i} c_j)_{V+A}$ 

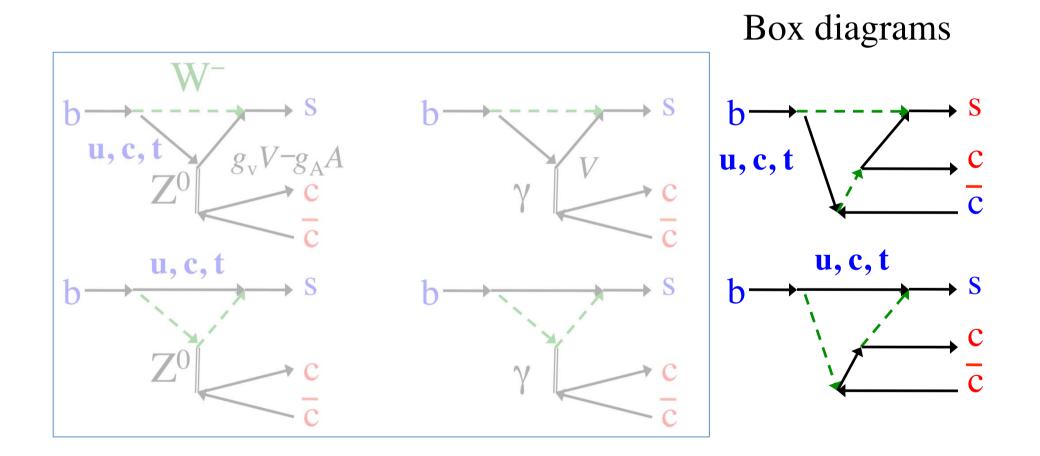
decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) the second order electroweak interactions ( $\Delta F = 1$ )



decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) the second order electroweak interactions ( $\Delta F = 1$ )



decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) also the second order electroweak interactions,  $\Delta F = 2$ 



decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) Theoretical tool to describe the decay amplitude for M $\rightarrow$ F  $A(M \rightarrow F) = \langle F | H_{\text{effective}}^{\text{weak decay}} | M \rangle = \frac{G_F}{\sqrt{2}} \sum_i \xi_{\text{CKM}}^i C_i(\mu) \langle F | Q_i(\mu) | M \rangle$ 

 $G_{\rm F}$ : Fermi constant,

- $Q_i(\mu)$ : Local four-fermion operators evaluated at energy scale  $\mu$  calculable in perturbation
- $C_i(\mu)$ : Coupling constants for  $Q_i(\mu)$  at energy scale  $\mu$  i.e. Wilson coefficient, calculable in perturbation
- <*F*|*Q<sub>i</sub>*(µ)|*M*>: Hadronic matrix element long distance effect
- $\xi_i^{\text{CKM}}$ : Combination of the CKM elements the ultimate interest for Flavour Physics extraction of the CKM matrix, search for new physics

decay ( $\Delta F = 1$ ) and oscillation amplitudes ( $\Delta F = 2$ ) Theoretical tool to describe the decay amplitude for M $\rightarrow$ F  $A(M \rightarrow F) = \langle F | H_{\text{effective}}^{\text{weak decay}} | M \rangle = \frac{G_F}{\sqrt{2}} \sum_i \xi_{\text{CKM}}^i C_i(\mu) \langle F | Q_i(\mu) | M \rangle$ 

- Comparing the full and effective theory at  $\mu = m_W$  $\rightarrow C_i(\mu = m_W)$ 

- Scale  $C_i$  down to  $\mu \approx 1$  GeV (K),  $m_c$  (D),  $m_b$  (B)

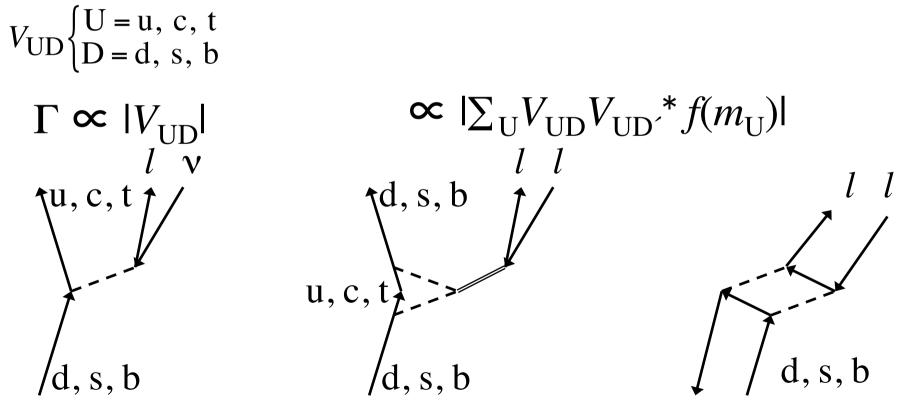
 $C_i(\mu) = U_{ij}(\mu, \mu = m_W)C_j(\mu = m_W)$  $U_{ij}$  not diagonal  $\Rightarrow$  mixing of the operators in the evolution

 Evaluate <*F*|*Q<sub>i</sub>*(μ)|*M*> (hadronic matrix element) with non perturbative methods at μ lattice, HQET, QCD sum rule, etc.
 major source of uncertainties

# Current Status of $V_{\rm CKM}$

Can be extract from decay widths generated by the tree, penguin, and box processes

examples of semileptonic and leptonic decays



Current Status of 
$$V_{\text{CKM}}$$

$$V_{\rm CKM} = \begin{pmatrix} V_{\rm ud} & V_{\rm us} & V_{\rm ub} \\ V_{\rm cd} & V_{\rm cs} & V_{\rm cb} \\ V_{\rm td} & V_{\rm ts} & V_{\rm tb} \end{pmatrix}$$

First 2×2 sub-matrix: four  $|V_{ij}|$  are measured by nucleus, pion, kaon and charm hadron decays It is "almost" unitary with one single parameter  $\lambda (\equiv \sin \theta_{\text{Cabibbo}}) = |V_{\text{us}}| = 0.2246 \pm 0.0012$  (PDG 2010)

$$V_{\rm CKM} \approx \begin{pmatrix} 1 & \lambda & V_{\rm ub} \\ \neg \lambda & 1 & V_{\rm cb} \\ V_{\rm td} & V_{\rm ts} & V_{\rm tb} \end{pmatrix}$$

Current Status of 
$$V_{\text{CKM}}$$
  
 $V_{\text{CKM}} \approx \begin{pmatrix} 1 & \lambda & V_{\text{ub}} \\ -\lambda & 1 & V_{\text{cb}} \\ V_{\text{td}} & V_{\text{ts}} & V_{\text{tb}} \end{pmatrix}$ 

 $|V_{cb}|$  and  $|V_{ub}|$  measured by semileptonic  $B_u$  and  $B_d$  decays

$$|V_{cb}| = \begin{cases} (41.5 \pm 0.7) \times 10^{-3} \text{ inclusive decays} \\ (38.7 \pm 1.1) \times 10^{-3} \text{ exclusive decays} \\ -\text{errors limited theoretically-} \end{cases}$$

$$2.1\sigma \text{ discrepancy}$$

$$(PDG 2010)$$

 $|V_{ub}| = \begin{cases} (4.27 \pm 0.38) \times 10^{-3} \text{ inclusive decays} \\ (3.38 \pm 0.36) \times 10^{-3} \text{ exclusive decays} \\ -\text{errors very limited theoretically-} \end{cases} 1.7\sigma \text{ discrepancy} \\ \text{(PDG 2010)} \end{cases}$ 

exclusives systematically smaller than exclusives...?

Current Status of 
$$V_{\text{CKM}}$$
  
 $V_{\text{CKM}} \approx \begin{pmatrix} 1 & \lambda & V_{\text{ub}} \\ -\lambda & 1 & V_{\text{cb}} \\ V_{\text{td}} & V_{\text{ts}} & V_{\text{tb}} \end{pmatrix}$ 

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 $|V_{cb}|$  and  $|V_{ub}|$  measured by semileptonic  $B_u$  and  $B_d$  decays arg  $V_{cb} = 0$  by a phase convention

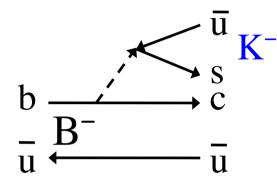
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 $|V_{cb}|$  and  $|V_{ub}|$  measured by semileptonic  $B_u$  and  $B_d$  decays arg  $V_{cb} = 0$  by a phase convention arg  $V_{ub}$  by CP violation in B $\rightarrow$ DK

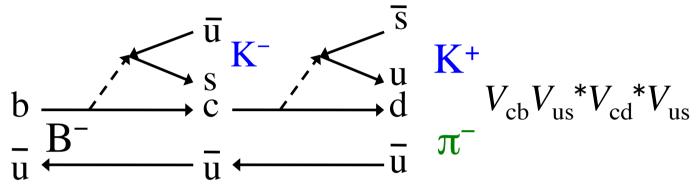
arg  $V_{\rm ub}$  so called angle " $\gamma$ "

two decay diagrams producing identical final states



$$V_{\rm cb}V_{\rm us}^{*}$$

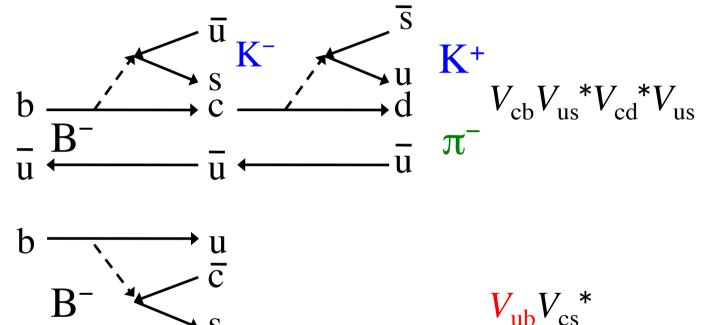
arg  $V_{\rm ub}$  so called angle " $\gamma$ "



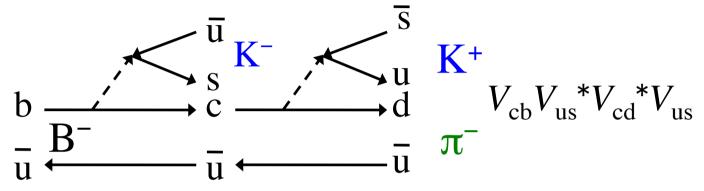
arg  $V_{\rm ub}$  so called angle " $\gamma$ "

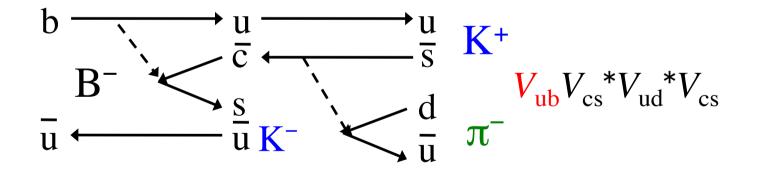
**K**<sup>-</sup>

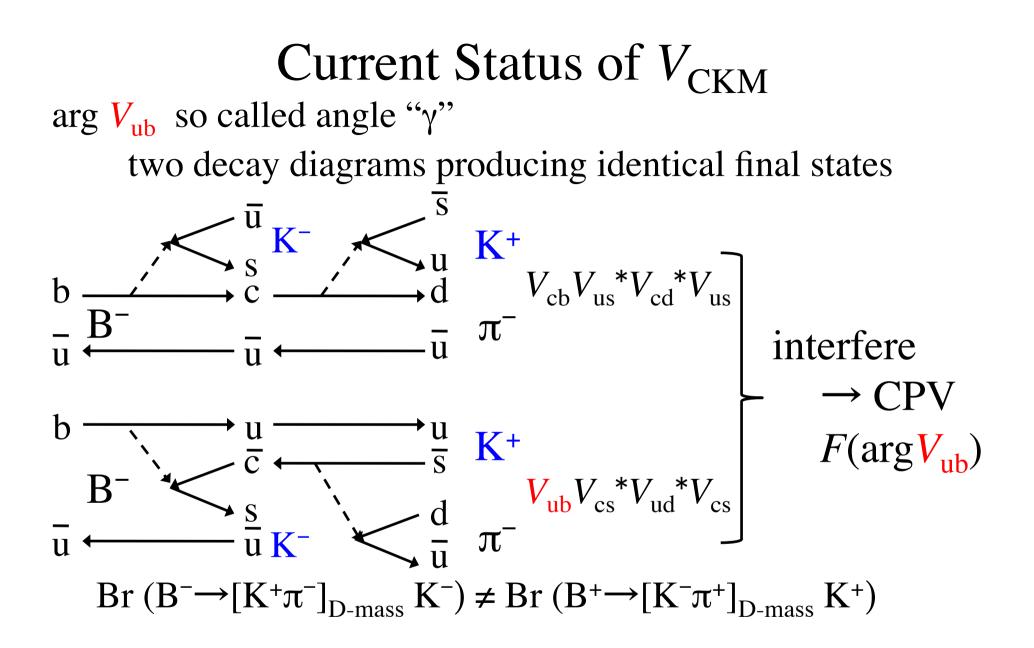
u



arg  $V_{\rm ub}$  so called angle " $\gamma$ "

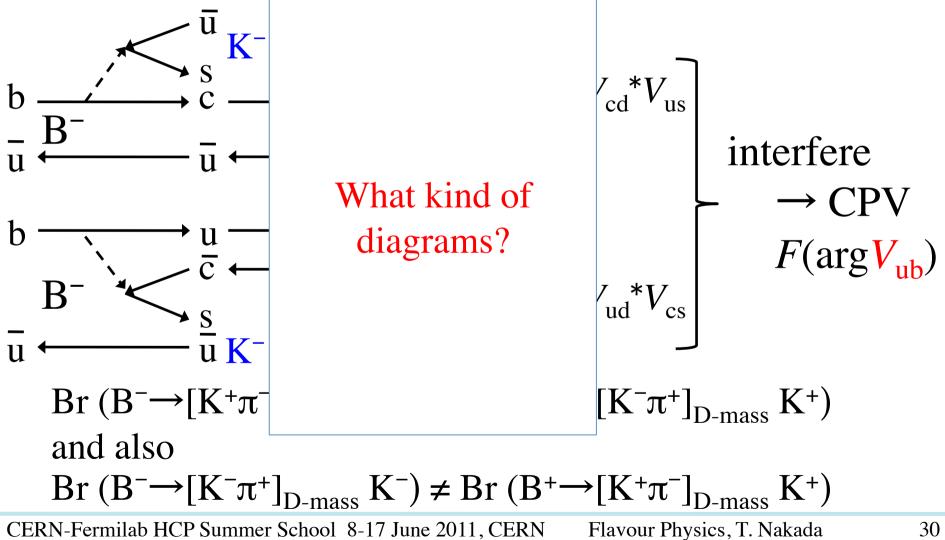


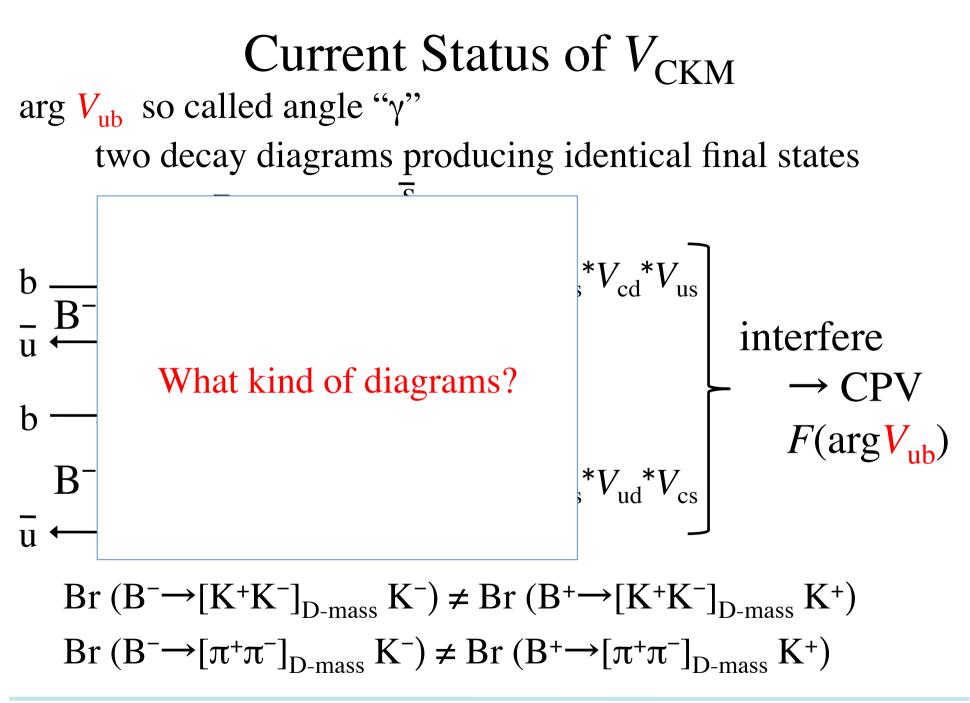


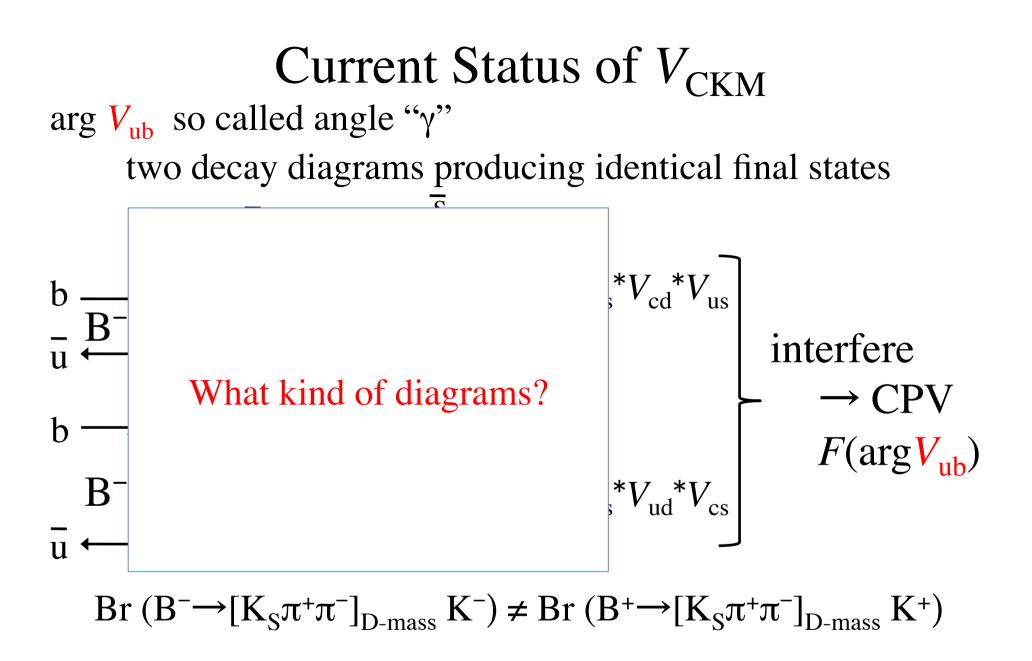


### Current Status of $V_{CKM}$

arg  $V_{\rm ub}$  so called angle " $\gamma$ "







arg  $V_{\rm ub}$  so called angle " $\gamma$ "

two decay diagrams producing identical final states

Current average =  $(73 + 22)^{-25}$  )° (PDG 2010)

arg  $V_{\rm ub}$  so called angle " $\gamma$ "

two decay diagrams producing identical final states

Current average = 
$$(73 + 22)^{\circ}$$
 (PDG 2010)

-Determined by the "tree" level amplitude interference between  $V_{cb}$  and  $V_{ub}$  no "New Physics" effect

#### -So far measured only by the e<sup>+</sup>e<sup>-</sup> B factory experiments: BABAR and BELLE

-In future, hadron machine will over take this...

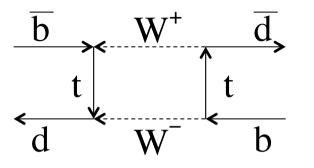
$$V_{\rm CKM} \approx \begin{pmatrix} 1 & \lambda & V_{\rm ub} \\ -\lambda & 1 & V_{\rm cb} \\ V_{\rm td} & V_{\rm ts} & V_{\rm tb} \end{pmatrix}$$

 $|V_{cb}|$  and  $|V_{ub}|$  measured by semileptonic  $B_u$  and  $B_d$  decays arg  $V_{cb} = 0$  by a phase convention arg  $V_{ub}$  by CP violation in B $\rightarrow$ DK  $V_{tb} \approx 1$  if we assume  $V_{CKM}$  to be unitary

$$V_{\rm CKM} \approx \begin{pmatrix} 1 & \lambda & V_{\rm ub} \\ -\lambda & 1 & V_{\rm cb} \\ V_{\rm td} & V_{\rm ts} & V_{\rm tb} \end{pmatrix}$$

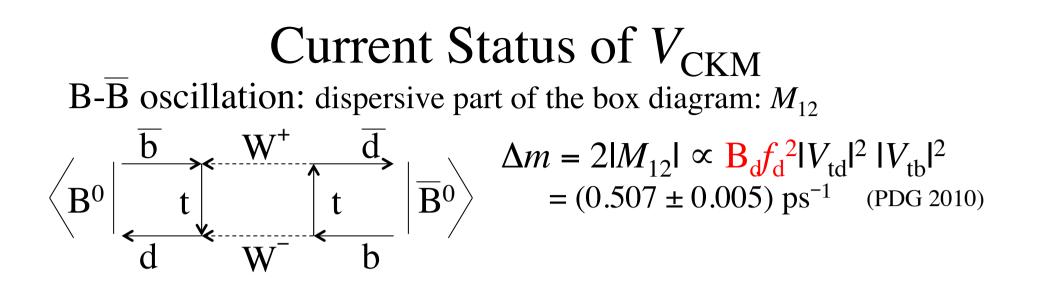
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B- $\overline{B}$  oscillation: dispersive part of the box diagram:  $M_{12}$ 

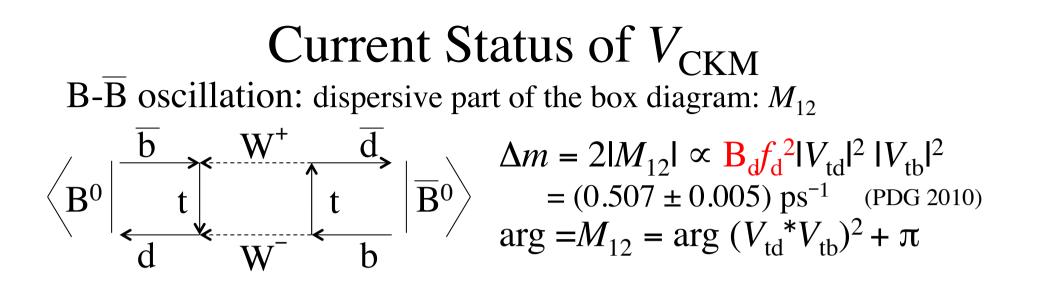


 $\begin{array}{c} \underline{d} \\ \underline{t} \\ \underline{t} \end{array} \qquad \Delta m = 2|M_{12}| \propto |V_{td}|^2 |V_{tb}|^2 \\ = (0.507 \pm 0.005) \text{ ps}^{-1} \quad (\text{PDG 2010}) \end{array}$ 

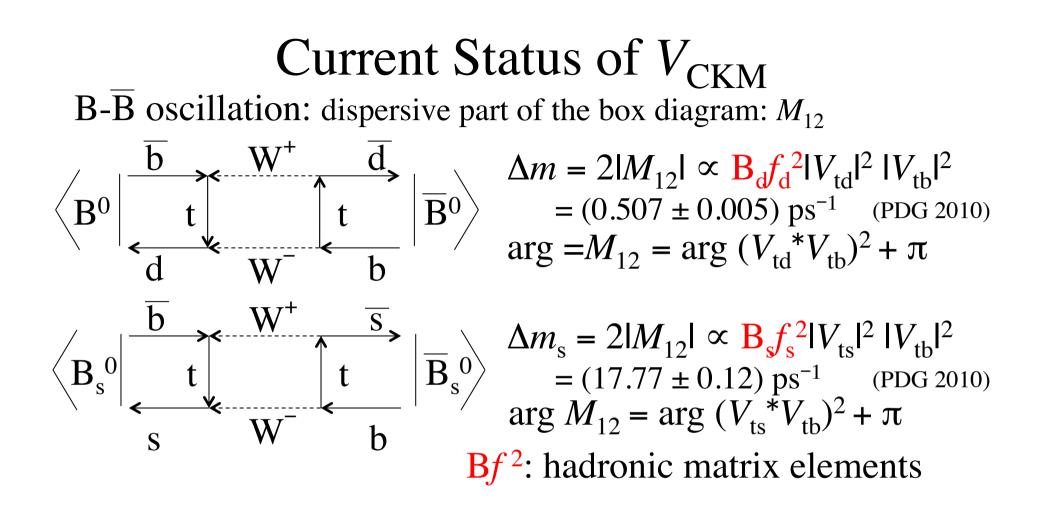
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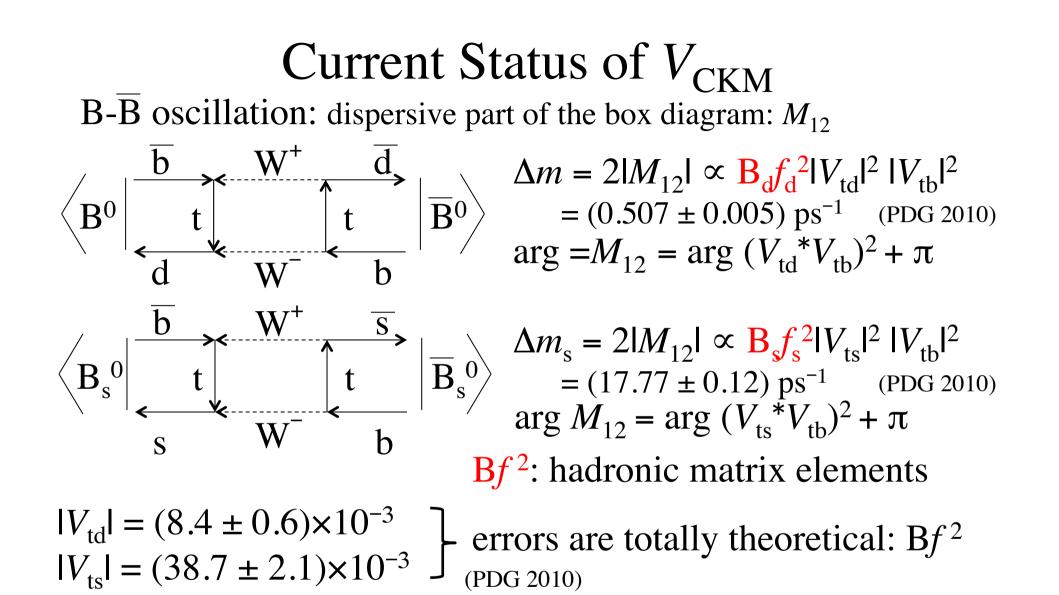


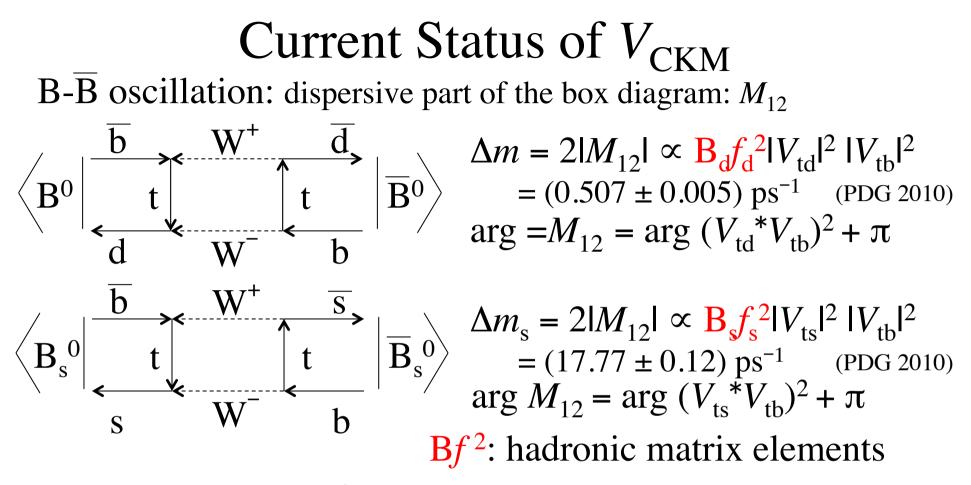
#### $\mathbf{B}f^2$ : hadronic matrix elements



 $\mathbf{B}f^2$ : hadronic matrix elements







 $\frac{|V_{td}| = (8.4 \pm 0.6) \times 10^{-3}}{|V_{ts}| = (38.7 \pm 2.1) \times 10^{-3}} = \text{errors are totally theoretical: } Bf^{2}_{(PDG\ 2010)}$  $\frac{|V_{td}|}{|V_{ts}| = 0.211 \pm 0.001 \pm 0.005} = (B_{d}f_{d}^{-2})/(B_{s}f_{s}^{-2}): \text{ smaller error}$ 

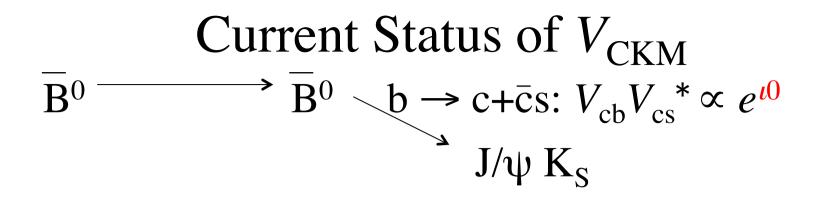
 $^{(\text{PDG 2010})}\Delta m_{\text{s}}$  measured only at the hadron machines

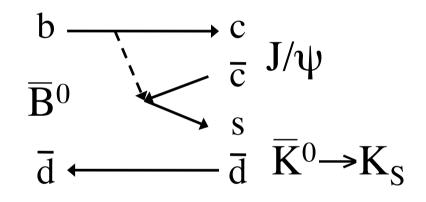
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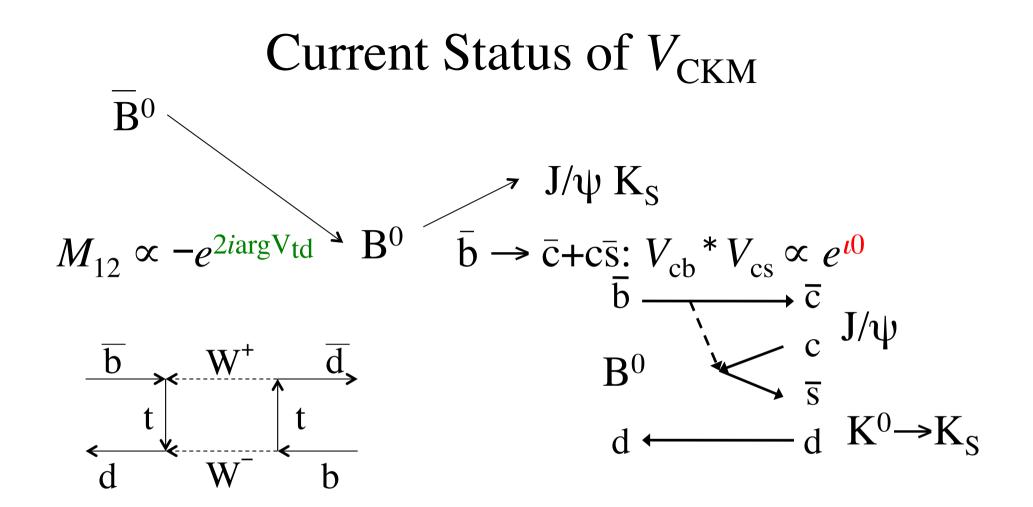
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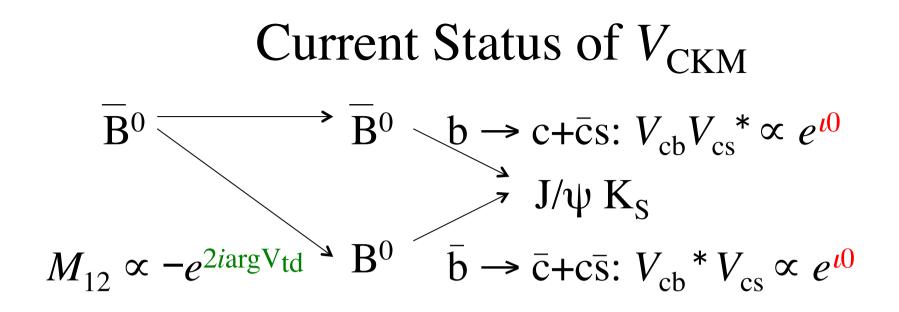
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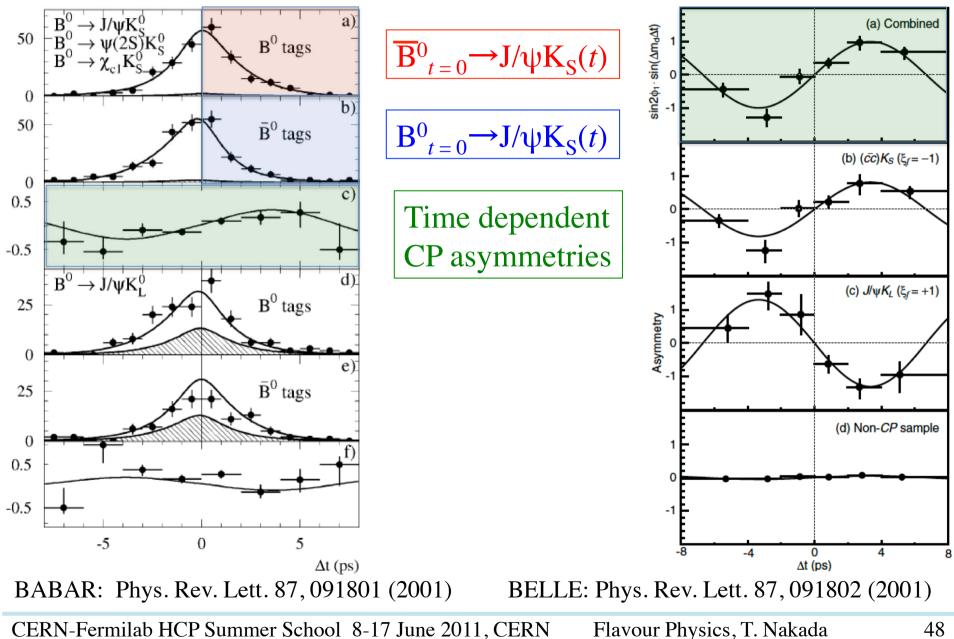


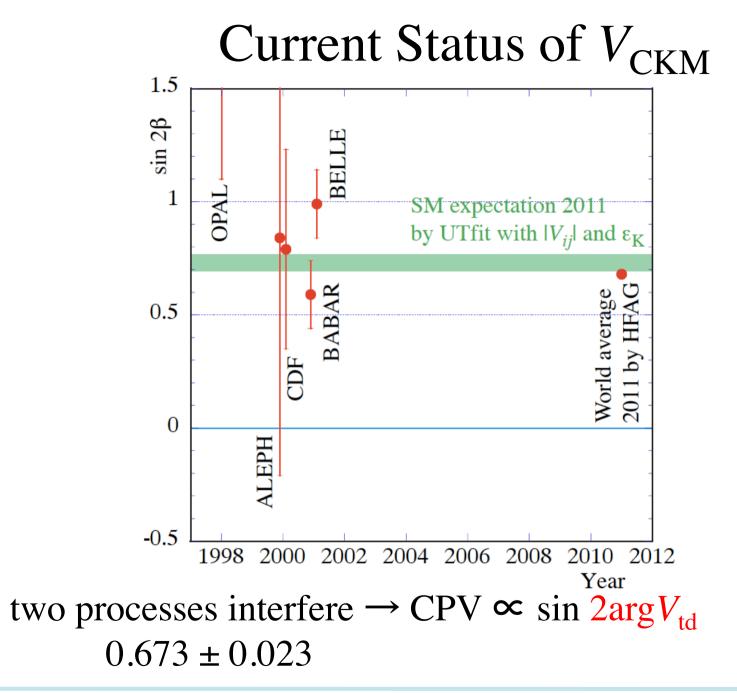


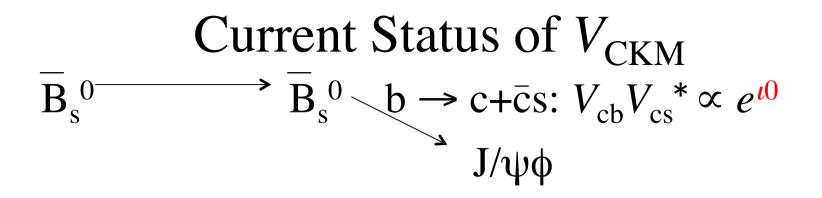


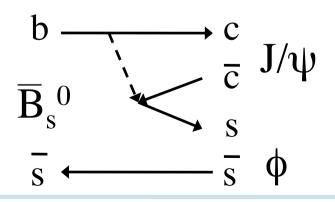


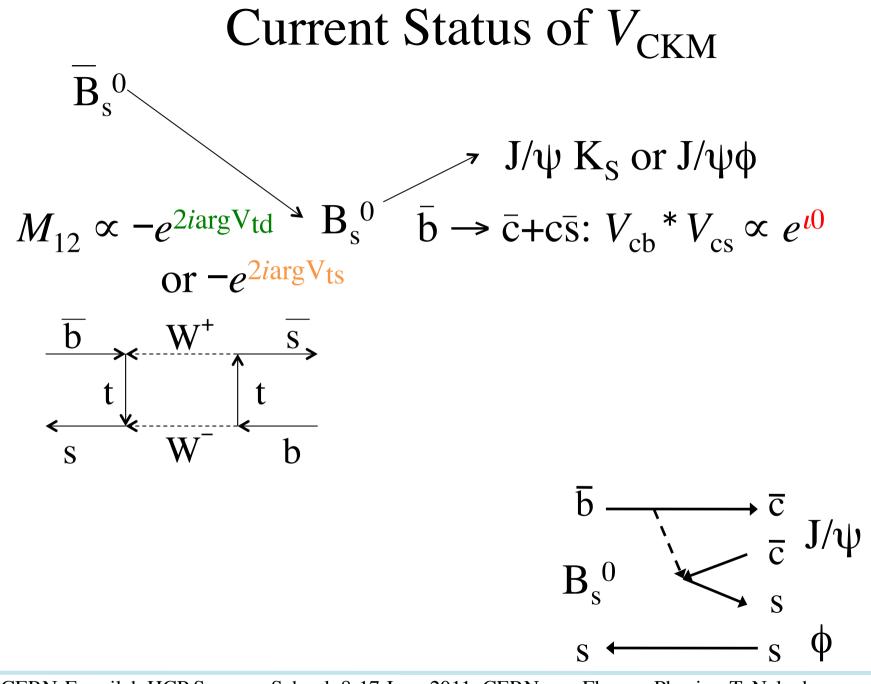
#### two processes interfere $\rightarrow CPV \propto \sin 2 \arg V_{td}$

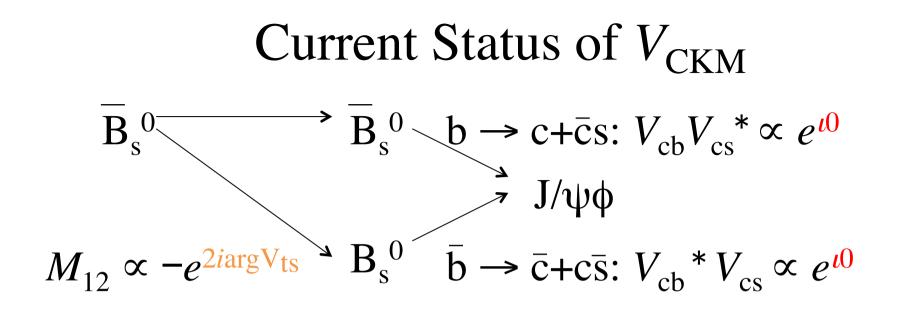




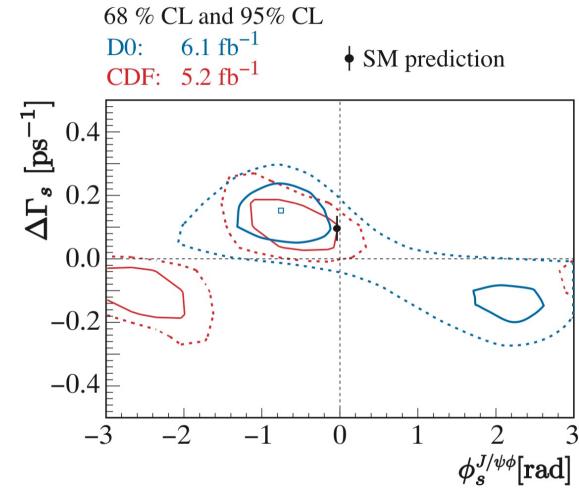




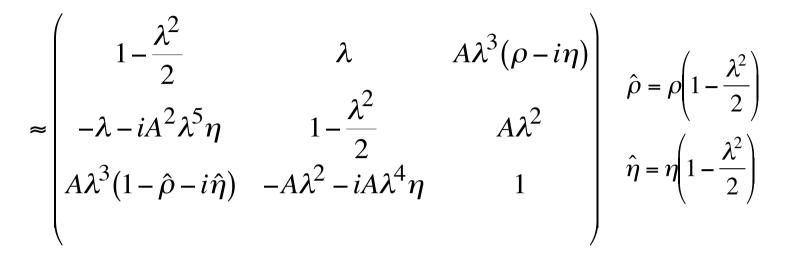




two processes interfere  $\rightarrow$  CPV  $\propto \sin 2 \arg V_{ts}$ not yet well measured

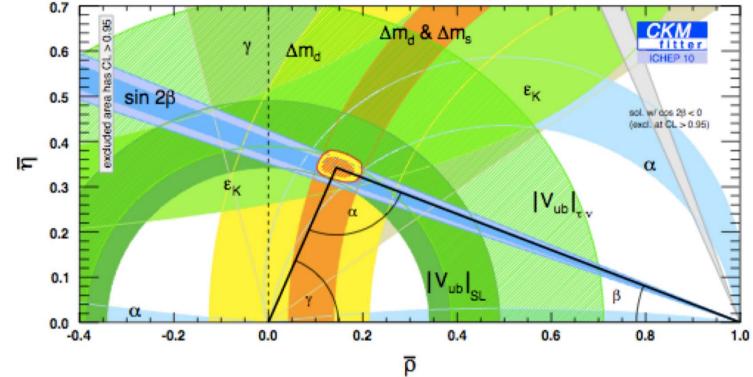


two processes interfere  $\rightarrow \text{CPV} \propto \sin 2 \arg V_{\text{ts}}$ not yet well measured



A from 
$$|V_{cb}|$$
,  $\rho$  and  $\eta$  from 
$$\begin{bmatrix} |V_{ub}| \text{ and arg } V_{ub} \\ |V_{tb}| \text{ and arg } V_{tb} \\ |V_{ub}| \text{ and } |V_{tb}| \\ |V_{ub}| \text{ and } |V_{tb}| \\ |V_{td}| \text{ and arg } V_{ub} \end{bmatrix}$$
many solutions i.e. consistency can be checked

• All input from B factories, except  $\varepsilon_{\rm K}$  and  $\Delta m_{\rm s}$ 



 BABAR and PEP-II completed in 2008 Belle and KEKB completed in 2010 with a total of ~1.2 ab<sup>-1</sup> data, i.e. ~1.3×10<sup>9</sup> BB!
 →Looking forward to seeing many key results with full statistics data in the coming conferences.