

Flavour Physics (III)

Sixth CERN-Fermilab Hadron Collider Physics Summer School

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Tatsuya NAKADA

Laboratory for High Energy Physics (LPHE)

Swiss Federal Institute of Technology Lausanne (EPFL)

Lausanne, Switzerland

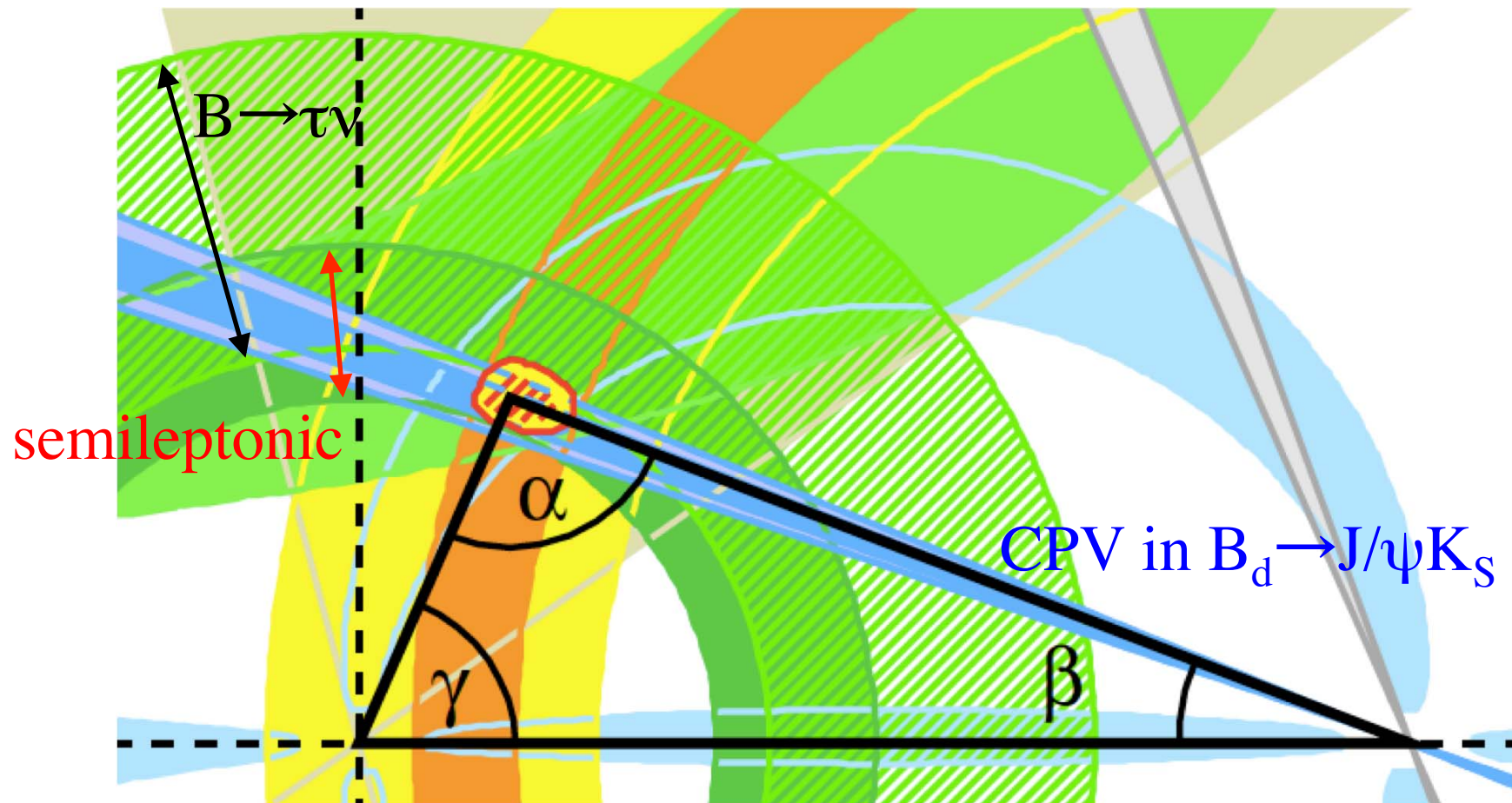


Plan of the lecture today

- Closer Look
- Flavour Physics at Hadron Machines

Closer Look

- If one looks closer, there exists some hint of discrepancies
 - “ $\sin 2\beta$ ” extracted from CPV in $B_d \rightarrow J/\psi K_S$ somewhat small
 - $|V_{ub}|$ extracted from $B \rightarrow \tau \nu$ decays larger than $|V_{ub}|$ extracted from the **semileptonic** decays.



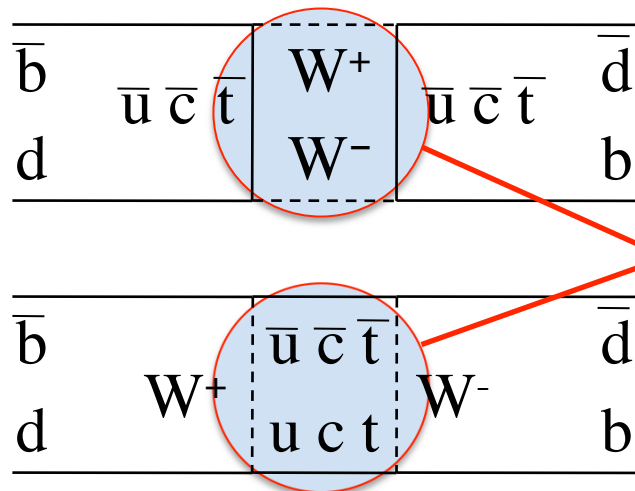
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- This could be due to
 1. Problem with extracting $|V_{ub}/V_{cb}|$ due to the hadronic uncertainties

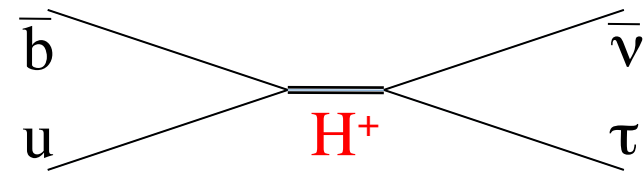
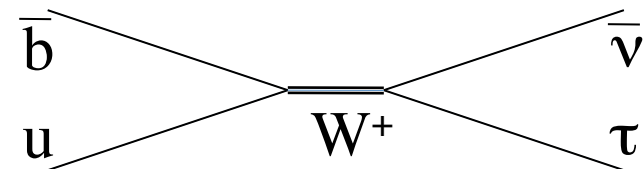
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- This could be due to
 1. Problem with extracting $|V_{ub}/V_{cb}|$ due to the hadronic uncertainties
 - OR
 2. New Physics in B^0 - \bar{B}^0 oscillations and charged Higgs in $B \rightarrow \tau \nu$



+ new particles



Closer Look

$|V_{cb}|$ and $|V_{ub}|$

Errors are **dominated by the theoretical uncertainties** in the strong interaction for the semileptonic decays.

⇒ can be reduced by studying the decay kinematics,
e.g. lepton momentum, hadronic-mass distribution, etc.
with **higher statistics in a clean environment**

Closer Look

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Case for Super B Factories: One B fully reconstructed, all remaining particles belong to other B!

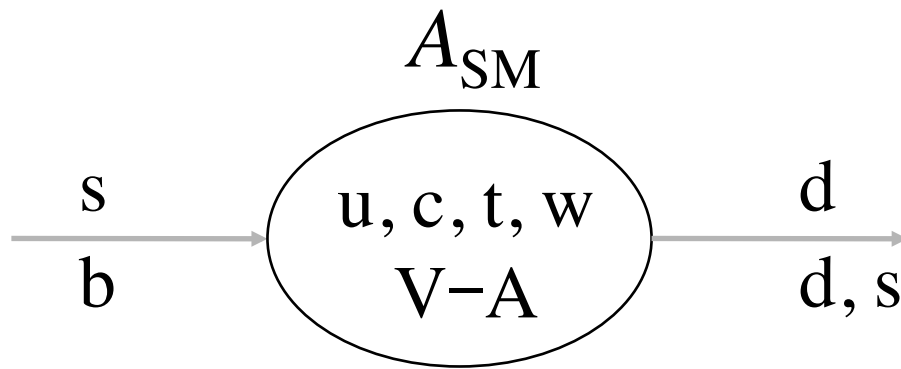
$|V_{ub}|$ from $B^\pm \rightarrow \tau^\pm \nu$ error still **statistics limited**.

Another case Super B Factories

But there are more...

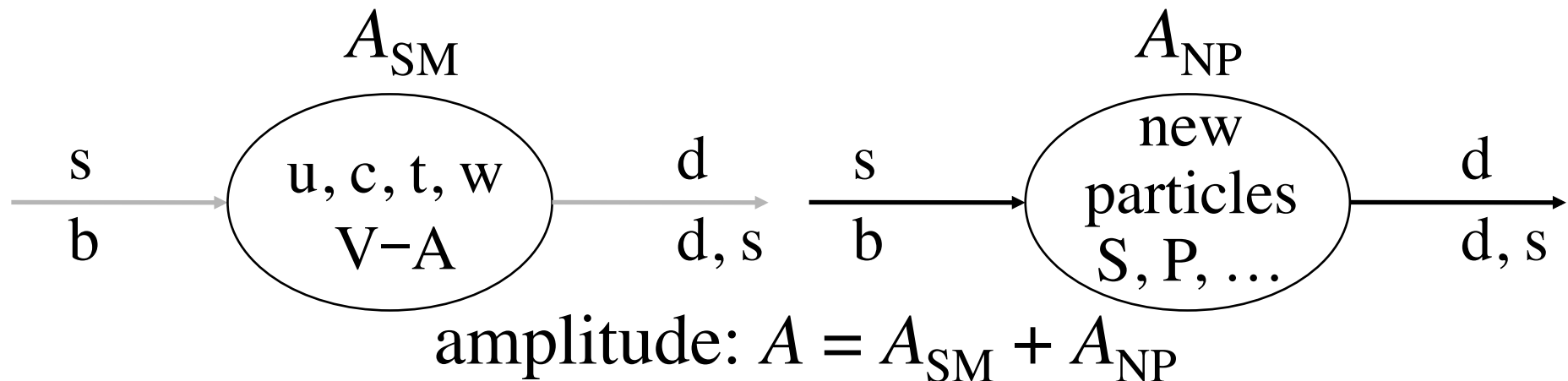
Closer Look

In general, new physics can enter in the loop diagrams as virtual states



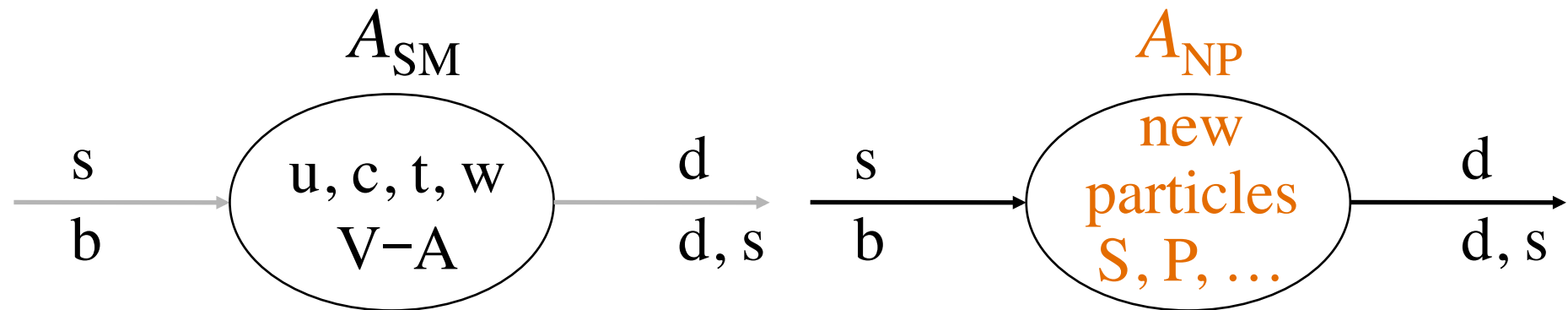
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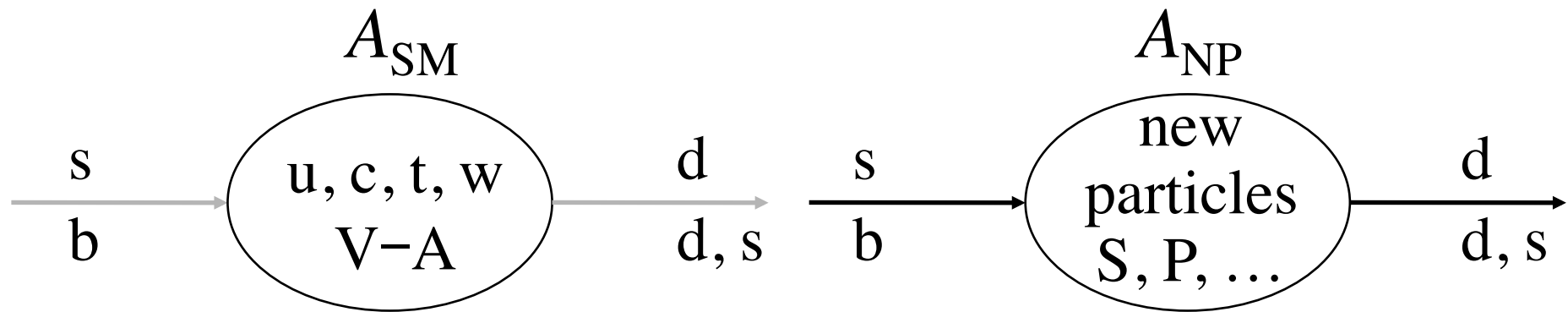
$$\text{amplitude: } A = A_{SM} + A_{NP}$$

$|A|$:

rare decays, Δm

Closer Look

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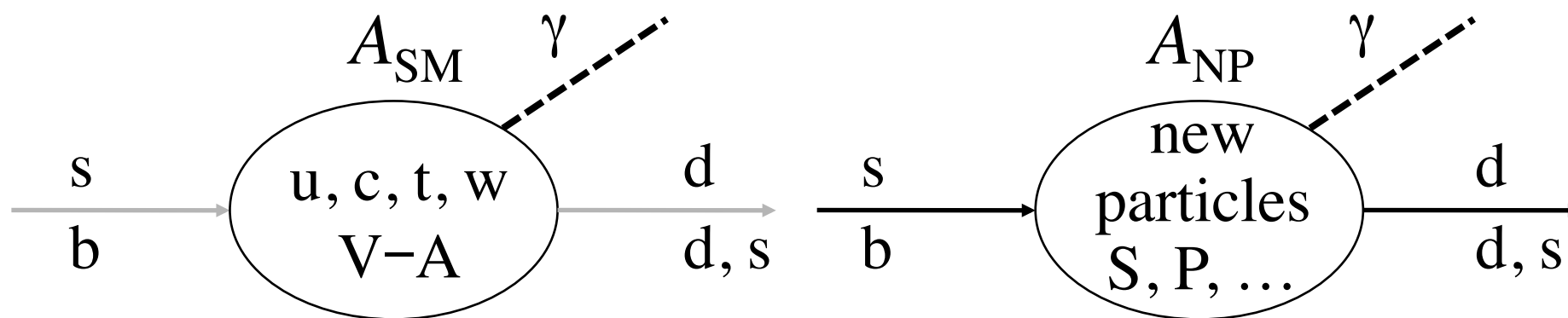
rare decays, Δm

$\arg A$:

CP violation

Closer Look

In general, new physics can enter in the loop diagrams as virtual states



$$\text{amplitude: } A = A_{SM} + A_{NP}$$

|A|: **rare decays, Δm**

arg A: **CP violation**

Lorentz structure of A: “photon” polarization via

final state angular distribution or mixing-decay CP violation

Closer Look

- Current experimental limits on new physics are still very large, up to $O(10)$ or more above the SM values:
 - $B_s \rightarrow \mu^+ \mu^-$
 - CPV in $B_s \rightarrow J/\psi \phi$
 - Lorentz structure in $b \rightarrow s$ radiative decays, $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, CPV in $B \rightarrow \phi \gamma$, etc.
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- Comparison of (ρ, η) determined from the tree processes, i.e. $|V_{ub}|$ and γ ($B \rightarrow DK$), and (ρ, η) from the loop processes, i.e. ε_K , β , and $|V_{td}|$.

Those measurements can be done at LHC, in particular with LHCb!

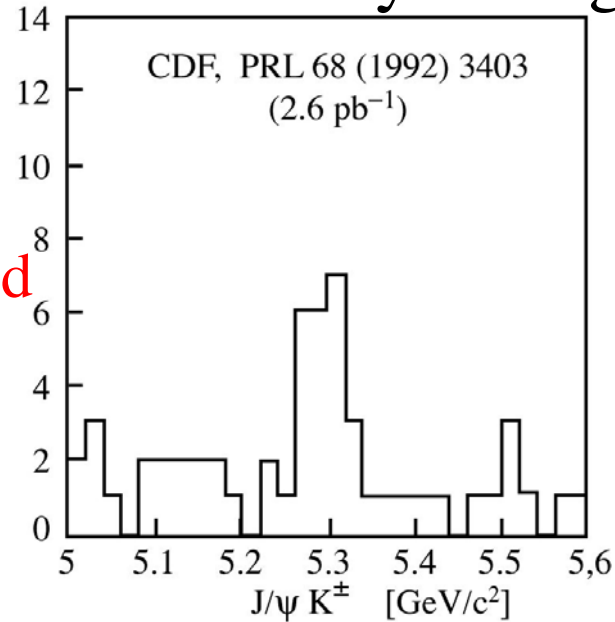
Flavour Physics at Hadron Machines

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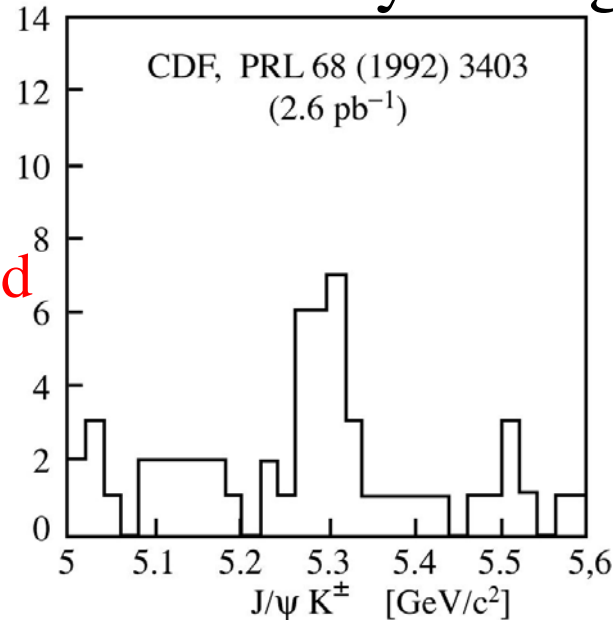
largest number of reconstructed $B^\pm \rightarrow J/\psi K^\pm$ at that time



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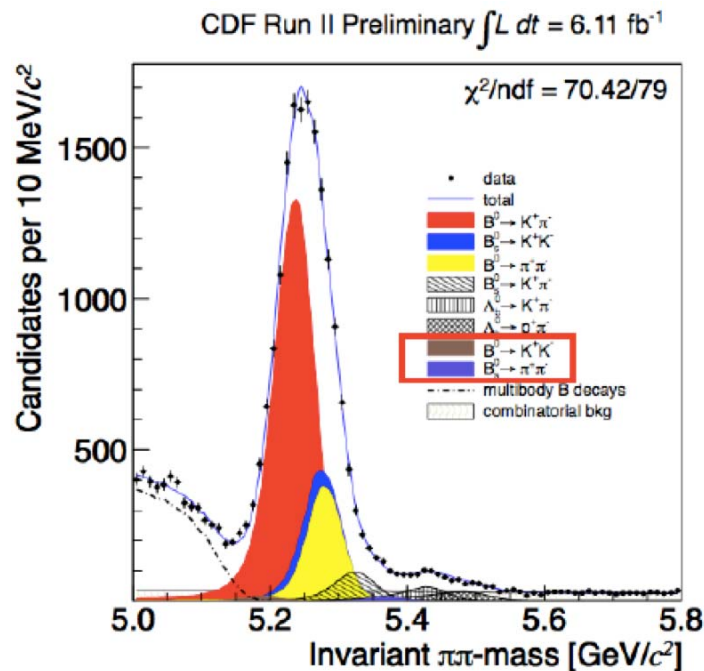
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- But significant contributions are during Run II: due to improved L , good vertex detectors, and trigger (DCF): b-baryon spectroscopy, lifetimes, $B_s - \bar{B}_s$ oscillation (Δm_s), and CP in $B \rightarrow J/\psi \phi$, CP in $B_s - \bar{B}_s$ oscillations, $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, $B_s \rightarrow \mu^+ \mu^-$, and D physics...

Flavour Physics at Hadron Machines

- CDF and D0 demonstrated that
 - Exclusive b-hadron decay modes (with charged particles, including semileptonic decays) can be well reconstructed
 - b-baryon and B_s : very unique and B_s oscillation can be resolved
 - B_d : for some decay modes as good as B factories or even better



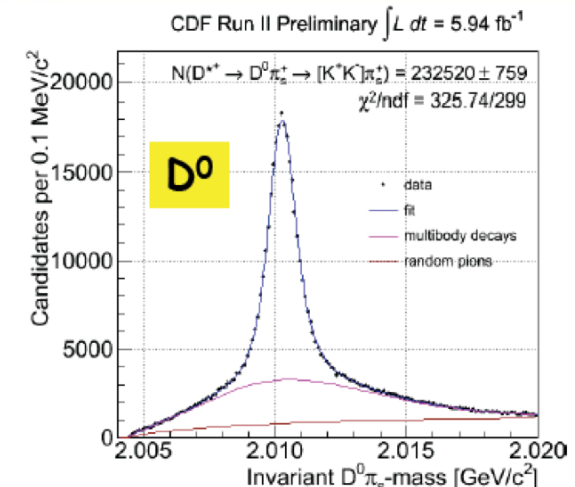
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CDF-Pub-10296

$$A_{CP}(D^0 \rightarrow \pi^+ \pi^-) = [+0.22 \pm 0.24 \pm 0.11]\%$$

$$A_{CP}(D^0 \rightarrow K^+ K^-) = [-0.24 \pm 0.22 \pm 0.10]\%$$



Flavour Physics at Hadron Machines

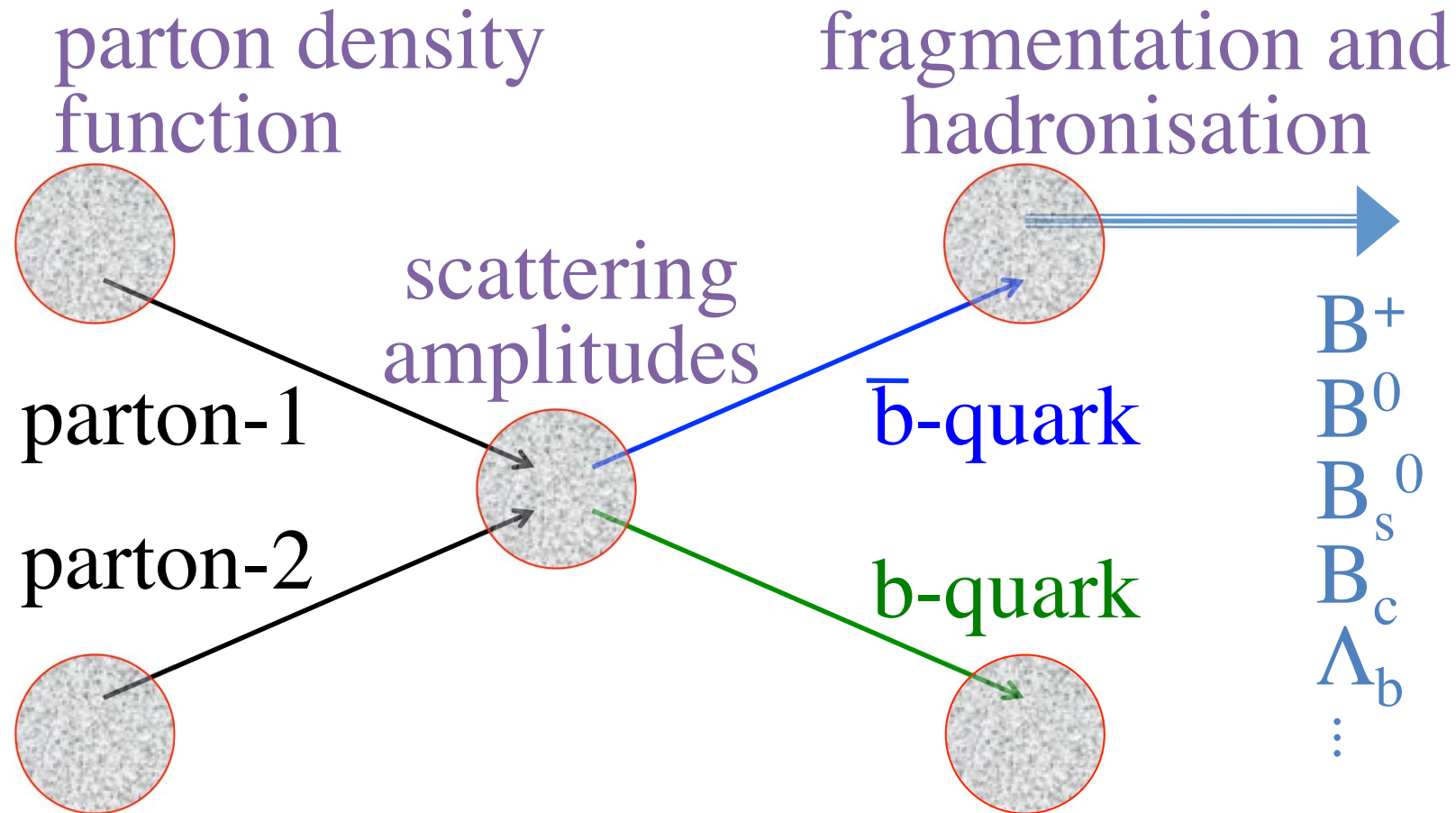
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- **But, LHC has started with an experiment dedicated for flavour physics, LHCb.**

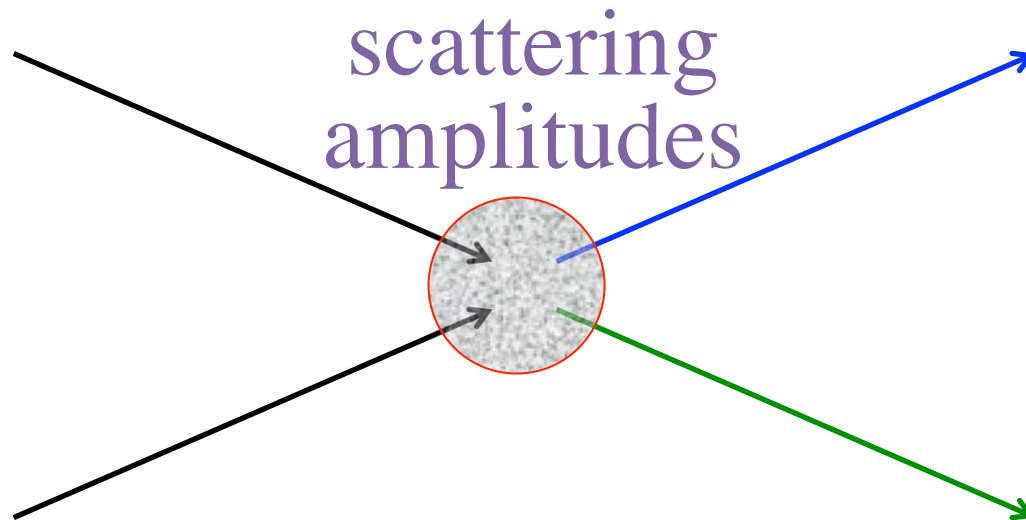
Flavour Physics at Hadron Machines

Production of heavy flavour



Flavour Physics at Hadron Machines

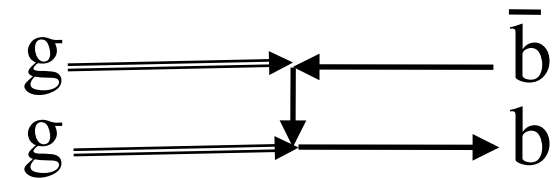
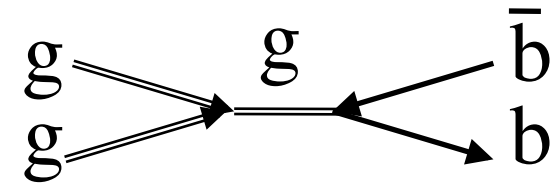
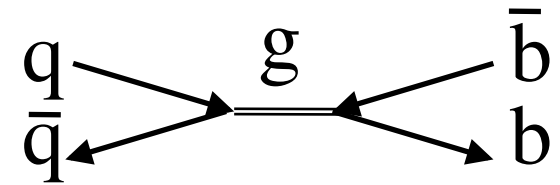
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Leading Order

$$q\bar{q} \rightarrow b\bar{b}$$

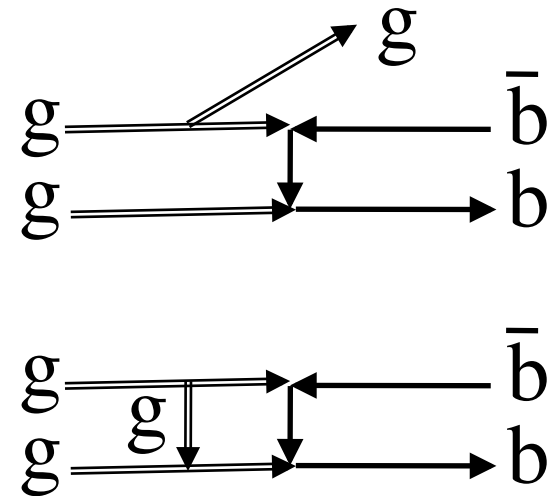
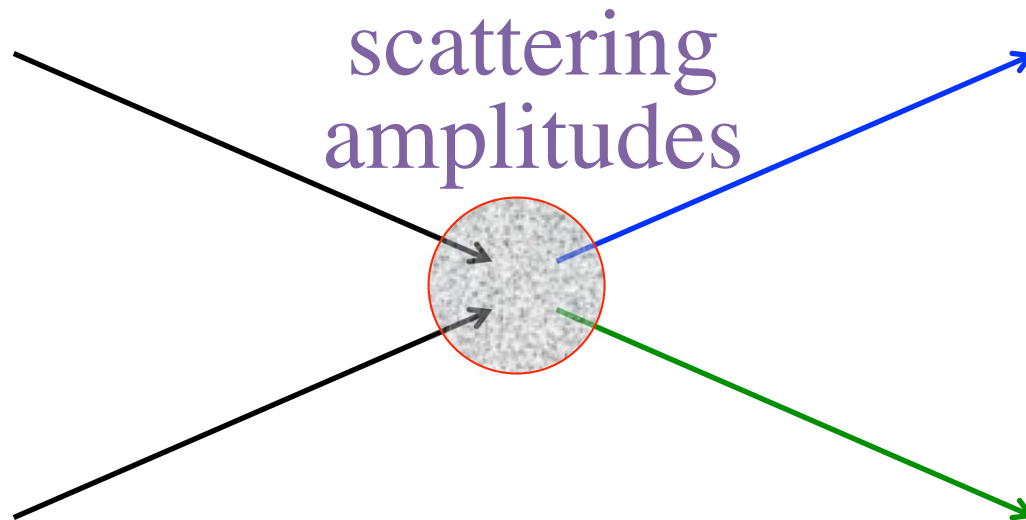
$$gg \rightarrow b\bar{b}$$



Flavour Physics at Hadron Machines

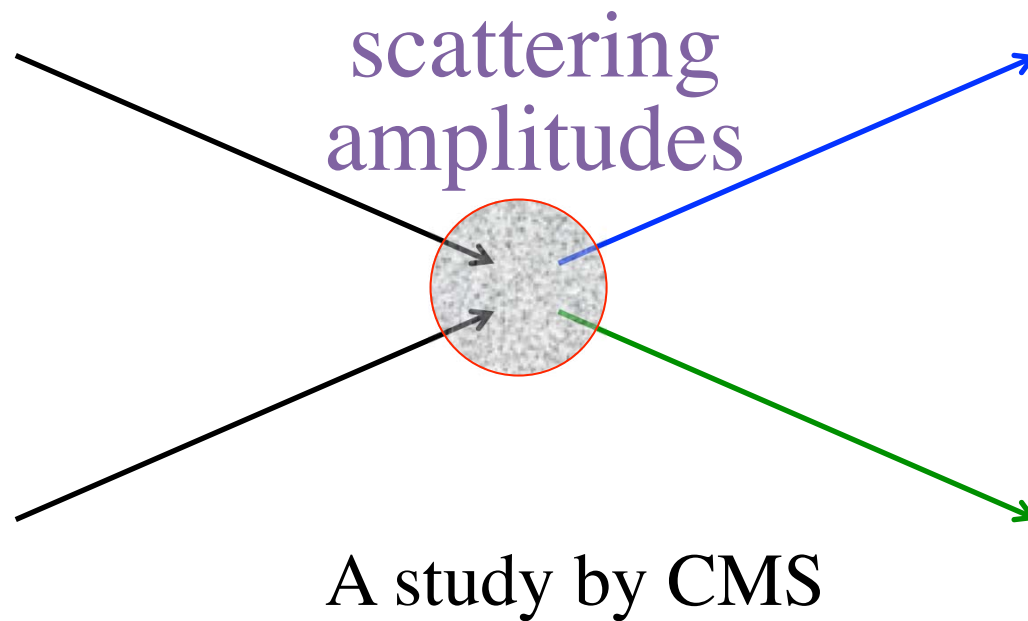
Production of heavy flavour

Leading Order
 $q\bar{q} \rightarrow b\bar{b}$
 $gg \rightarrow b\bar{b}$
+ Next to LO
i.e. adding one g

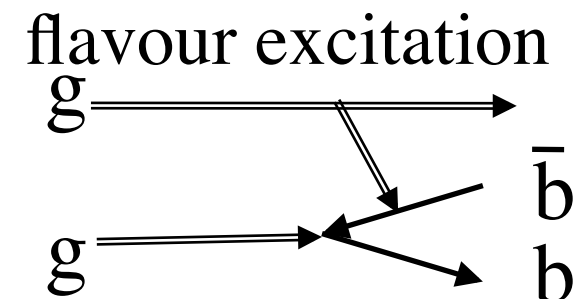
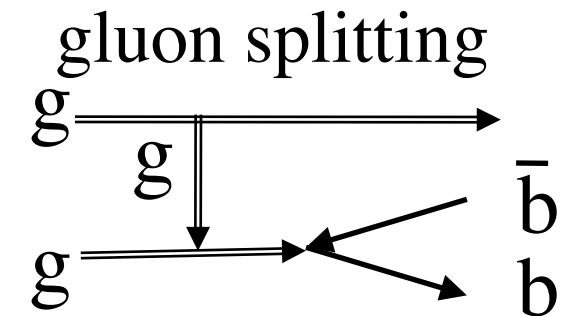


Flavour Physics at Hadron Machines

Production of heavy flavour



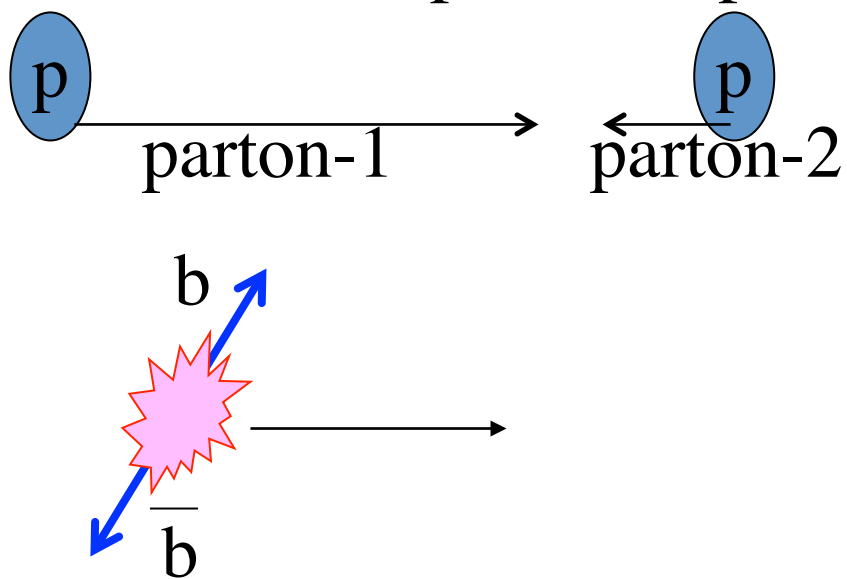
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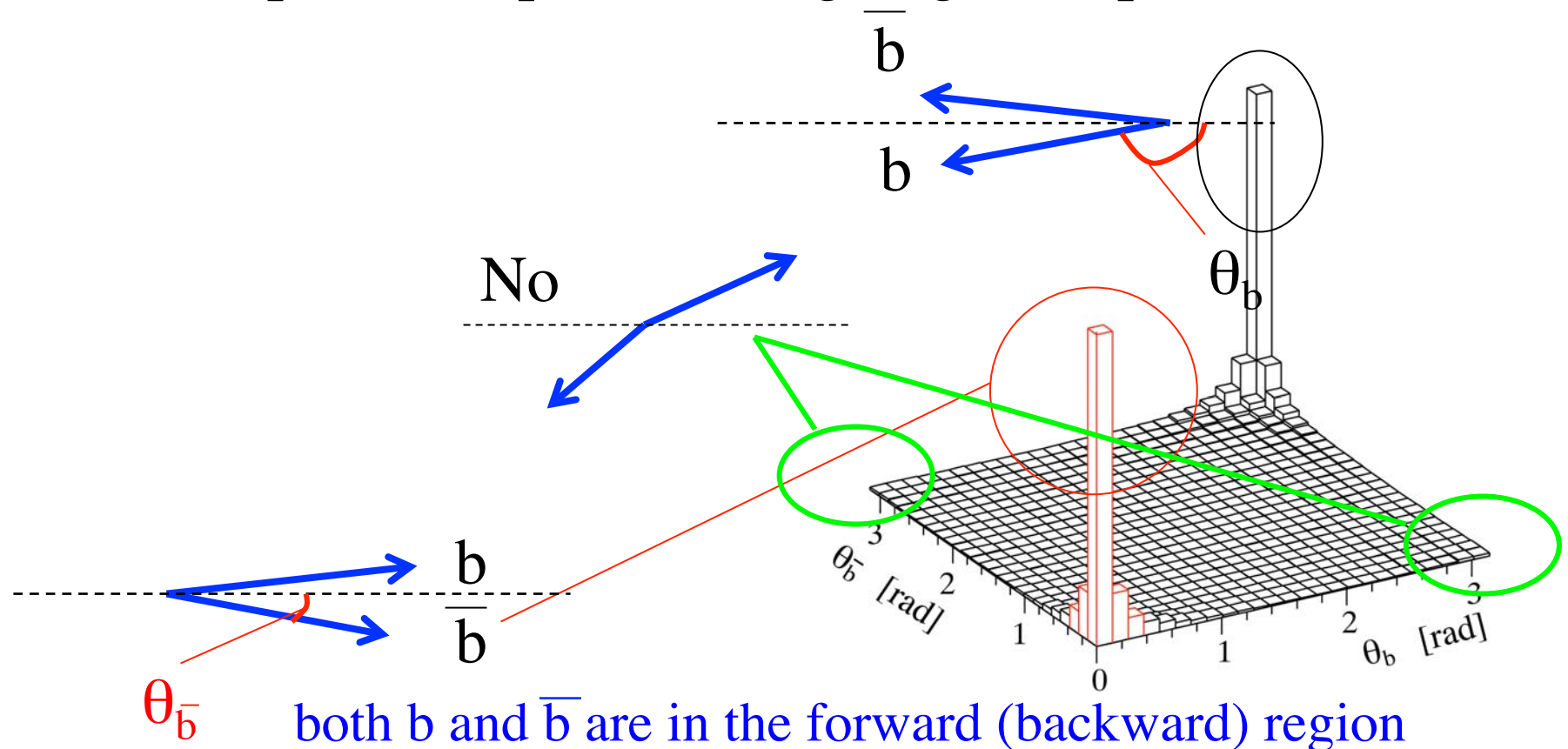
important input for designing an experiment



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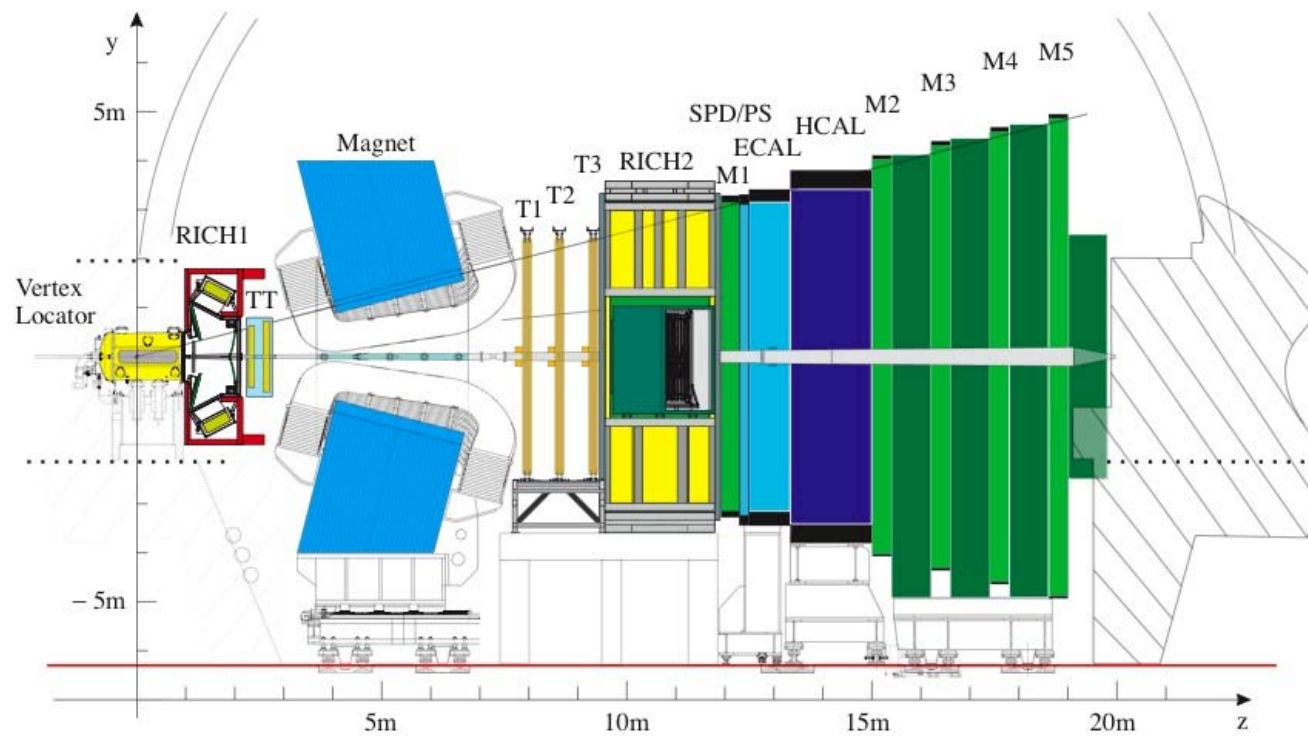
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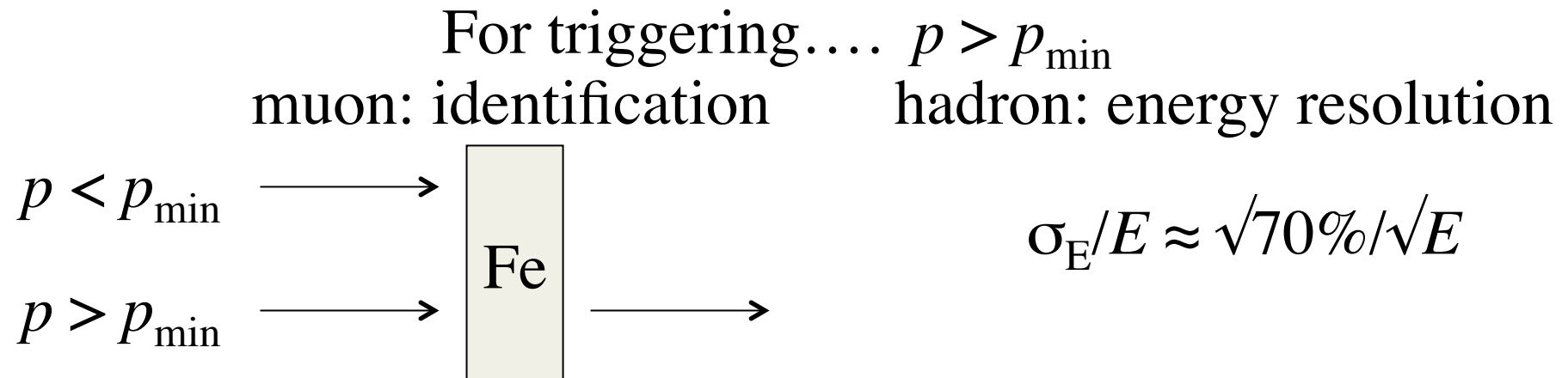
Flavour Physics at Hadron Machines

that is why LHCb is a forward spectrometer



Flavour Physics at Hadron Machines

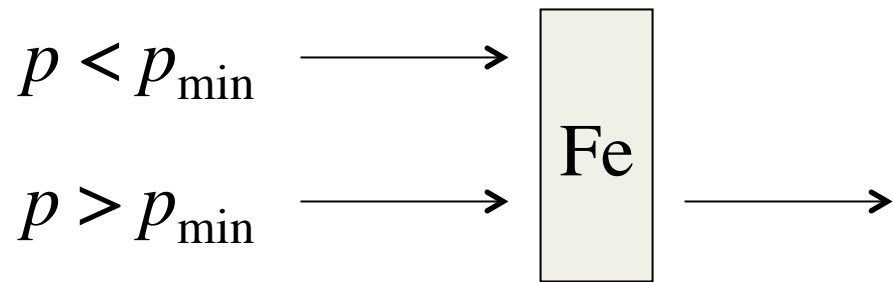
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Flavour Physics at Hadron Machines

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For triggering.... $p > p_{\min}$
 muon: identification hadron: energy resolution

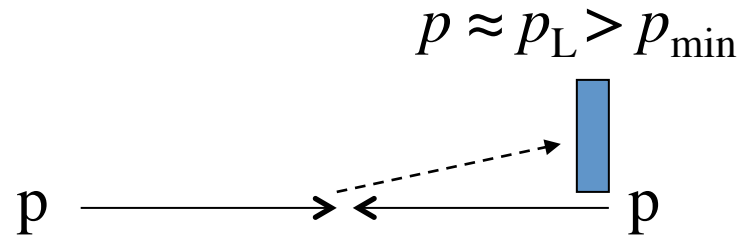
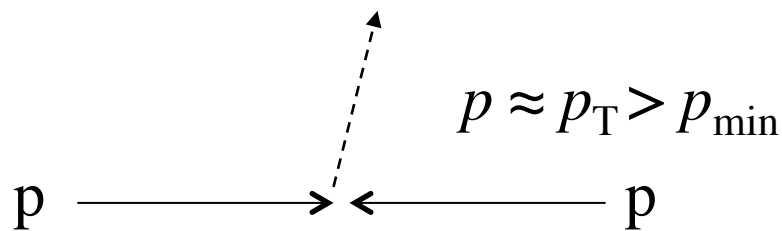


$$\sigma_E/E \approx \sqrt{70\%}/\sqrt{E}$$

central detector



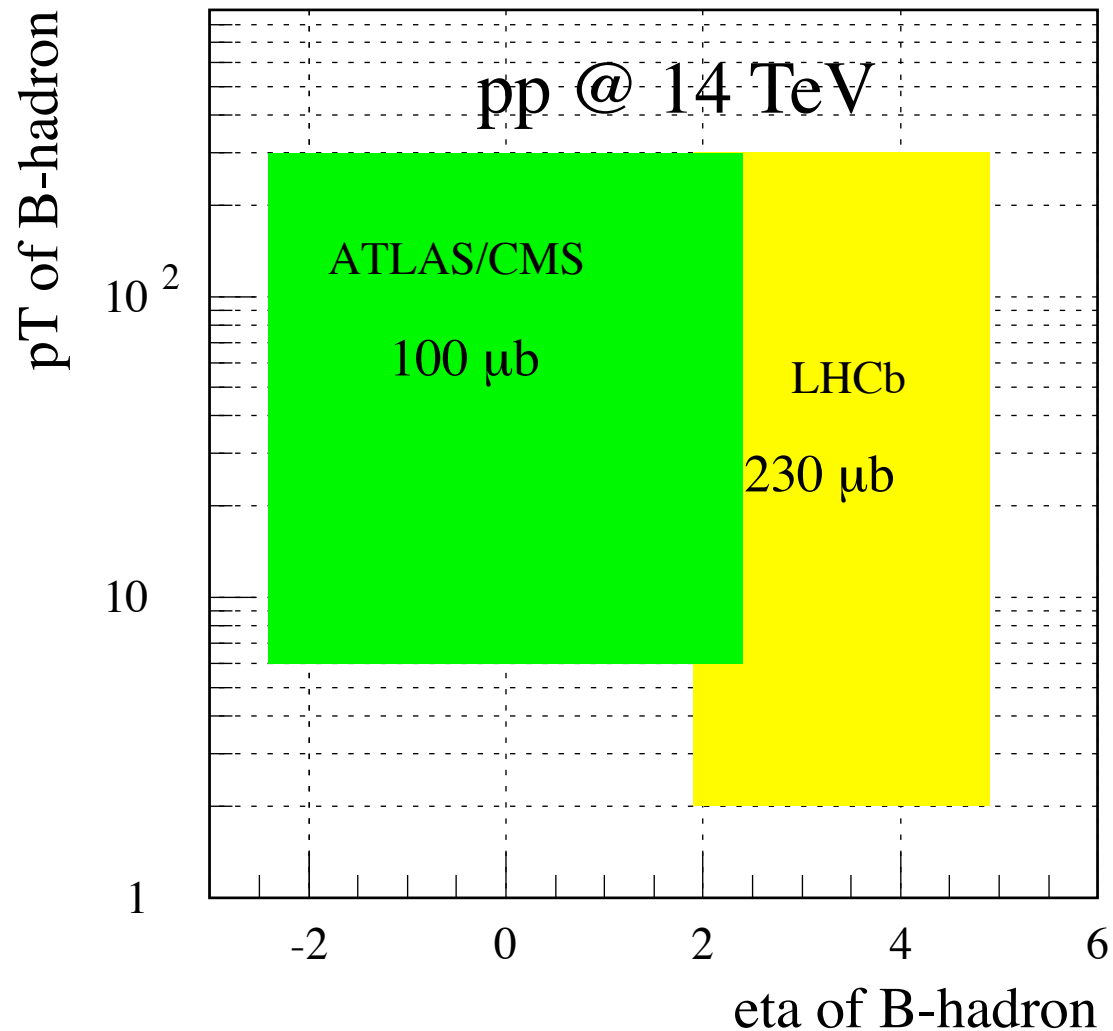
forward detector



forward: p_T threshold can be set low: \rightarrow high b efficiency

Flavour Physics at Hadron Machines

can exploit low p_T particles



Flavour Physics at Hadron Machines

Reconstruction of B **decay vertex with a good resolution**
is essential to **reduce combinatorial background**:

decay vertex: >1 well reconstructed tracks

well reconstructed track =

- charged particle seen by vertex detector
- reconstructed particle from tracks measured
by vertex detector

$D^0(K^-\pi^+)$, $D_s(K^+K^-\pi^+)$, etc., also K_S

examples are

$B_{(s)}^0 \rightarrow l^+l^-$, h^+h^- , $B_s^0 \rightarrow D_s(K^+K^-\pi^-) \pi^+$, $B^+ \rightarrow D(K_S\pi^+\pi^-) K^+$

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K_S not seen by the vertex detector, π^0 and γ can be **associated**
to a reconstructed vertex (if not too many)

$B^0 \rightarrow J/\psi K_S$, $K^{*0}(K^+\pi^-)\gamma$, $\rho^0(\pi^+\pi^-)\pi^0$, etc. are possible

but not

$B^0 \rightarrow K_S\pi^0$, $\rho^+(\pi^+\pi^0)\pi^0$, $\pi^0\nu\nu$, etc.

$B^+ \rightarrow \mu^+\nu$, $K^+\nu\nu$, $\tau^+\nu$

Flavour Physics at Hadron Machines

Advantage of LHCb over CDF and D0 are:

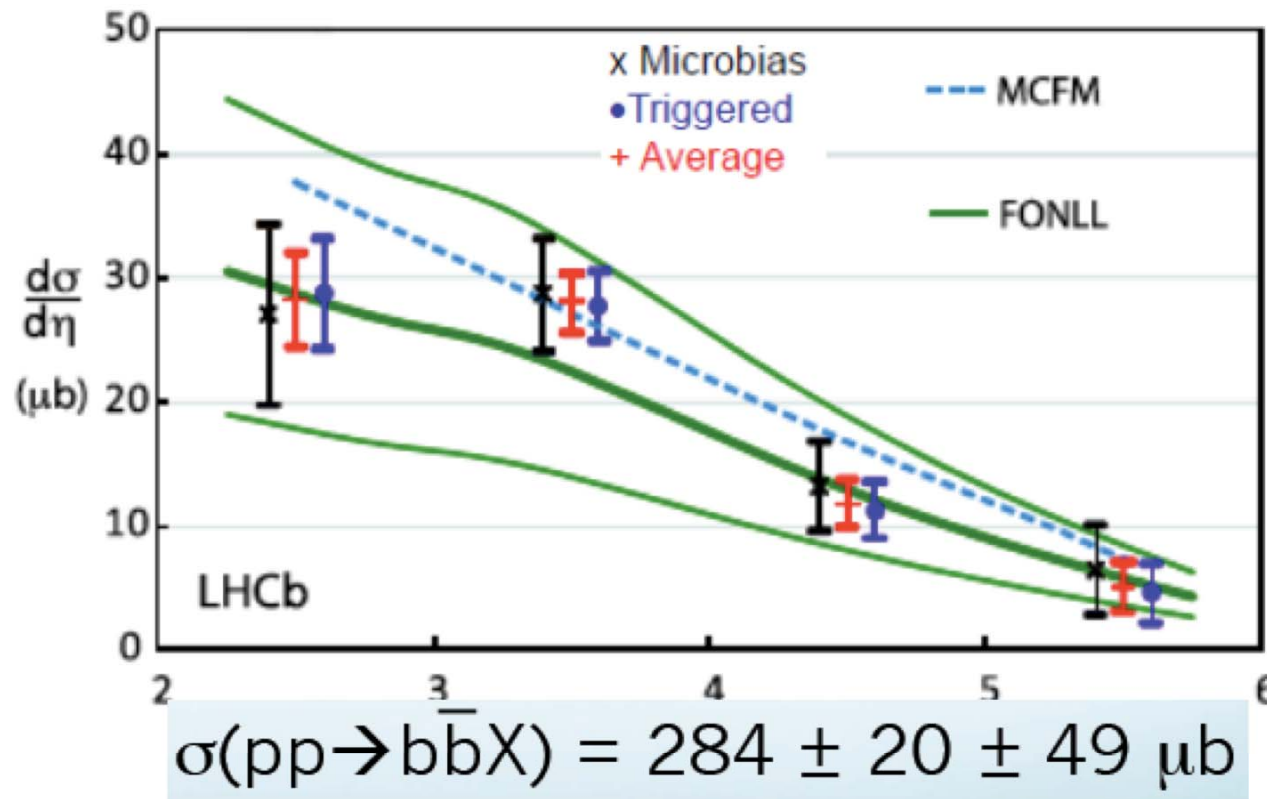
-larger cross section (measured also by ATLAS and CMS)

Flavour Physics at Hadron Machines

Advantage of LHCb over CDF and D0 are:

- larger cross section (measured also by ATLAS and CMS)
- one of the very early measurements at LHCb*

Phys Lett B 694 (2010)

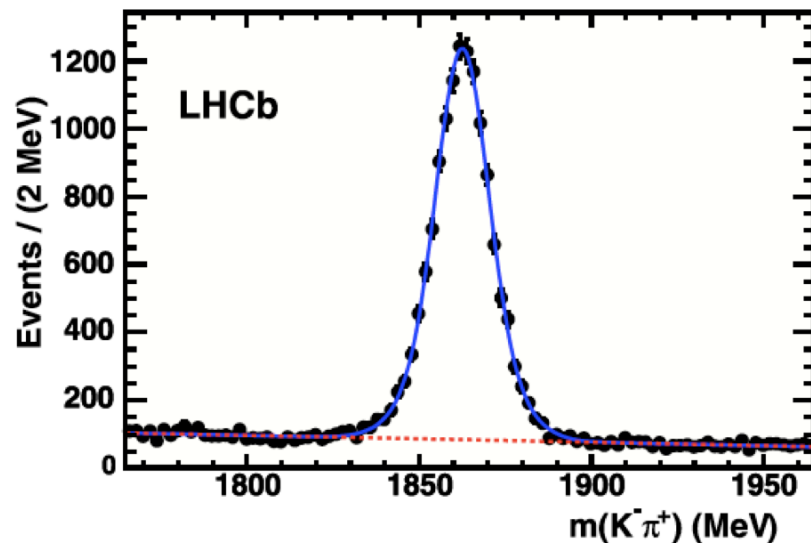


Flavour Physics at Hadron Machines

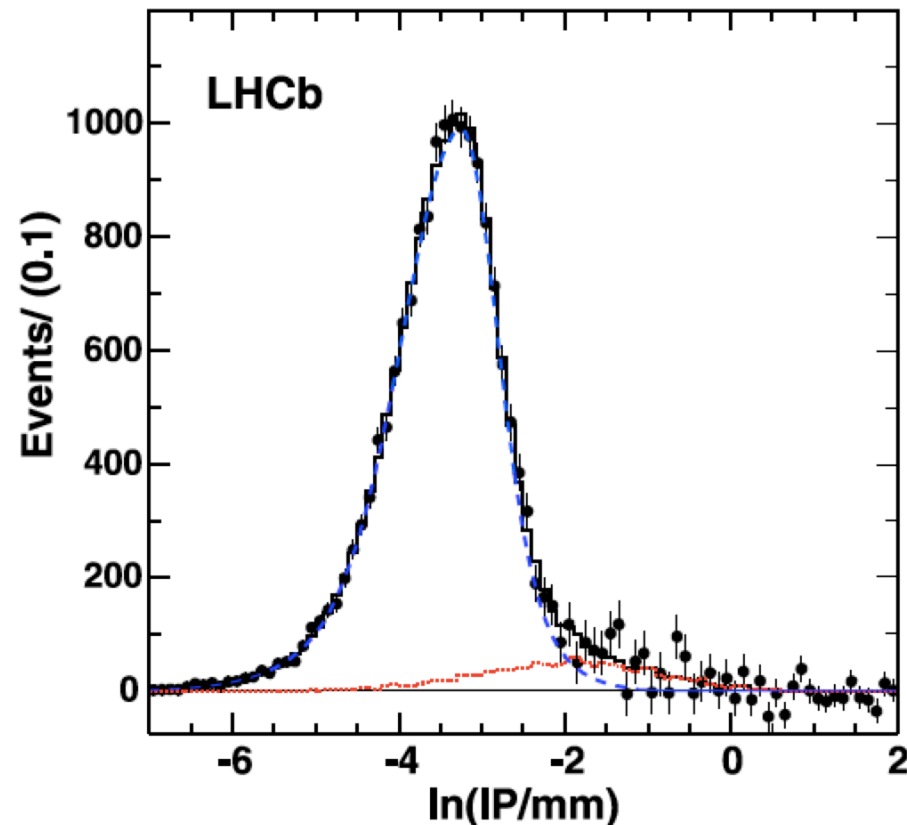
An interesting example of semi inclusive reconstruction

$$B^+ \rightarrow \bar{D}^0 (\rightarrow K^- \pi^+) X \mu^+ \nu$$

Phys Lett B 694 (2010)



$$\bar{D}^0 (\rightarrow K^- \pi^+) \\ + D^0 (\rightarrow K^- \pi^+)$$



\bar{D}^0 and D^0

created at the primary vertex: pointing to PV, small IP

decay from b: not pointing to PV, large IP

Flavour Physics at Hadron Machines

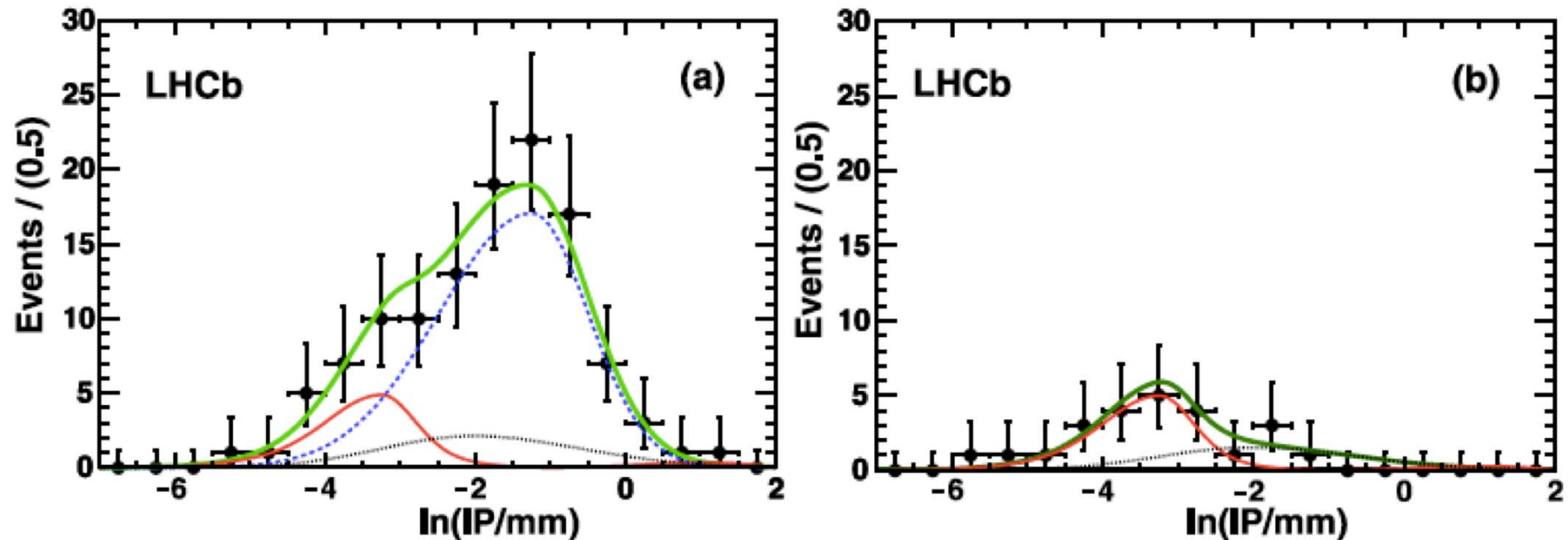
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$$B^+ \rightarrow \bar{D}^0 (\rightarrow K^- \pi^+) X \mu^+ \nu$$

B component can be increased by combining with μ^+ with appropriate $D(K^- \pi^+) \mu^+$ invariant mass

Non-B component form $D(K^- \pi^+) \mu^-$

Phys Lett B 694 (2010)



Flavour Physics at Hadron Machines

Advantage of LHCb over CDF and D0 are:

-larger cross section

-better invariant mass and proper time resolutions

B_s^0 - \bar{B}_s^0 oscillations: measure time dependent rates for

$$B_s^0 \text{ initial} \Rightarrow \bar{B}_s^0 \text{ at } t$$

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flavour tagging of the initial state

B_s^0 initial

Flavour Physics at Hadron Machines

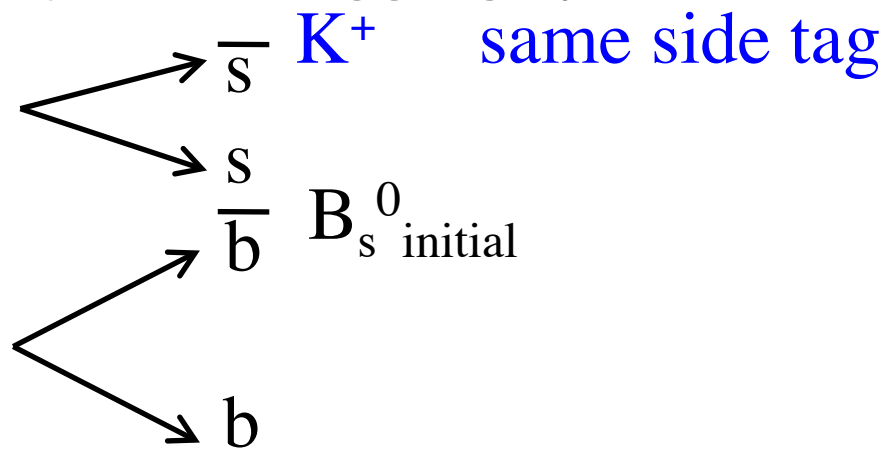
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↳ l^-X , $c \rightarrow s \rightarrow K^-$, secondary vertex with - charge

oposit side tag

Flavour Physics at Hadron Machines

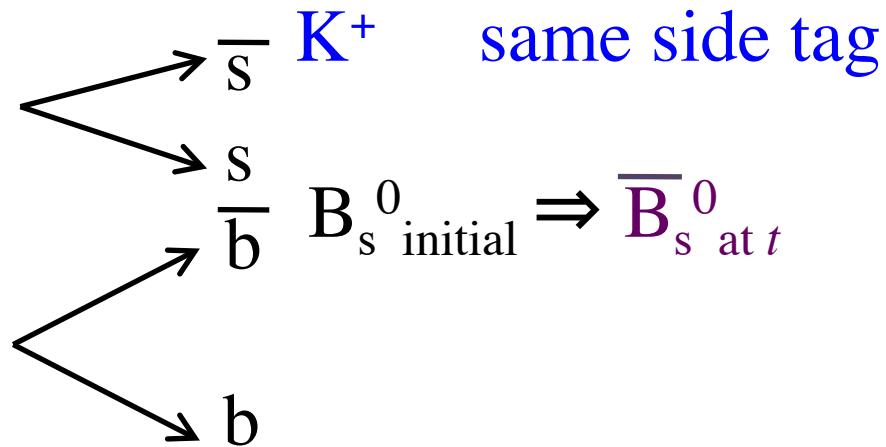
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flavour tagging of the final state

↳ $l^- X, c \rightarrow s \rightarrow K^-$, secondary vertex with **- charge**

oposit side tag

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\bar{s} K^+ same side tag

\bar{B}_s^0 initial \Rightarrow \bar{B}_s^0 at t ($\rightarrow D_s^+ \pi^-$) *flavour tagging of the final state*

$\hookrightarrow K^+ K^- \pi^+$

b

$\hookrightarrow l^- X, c \rightarrow s \rightarrow K^-$, secondary vertex with - charge

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Flavour Physics at Hadron Machines

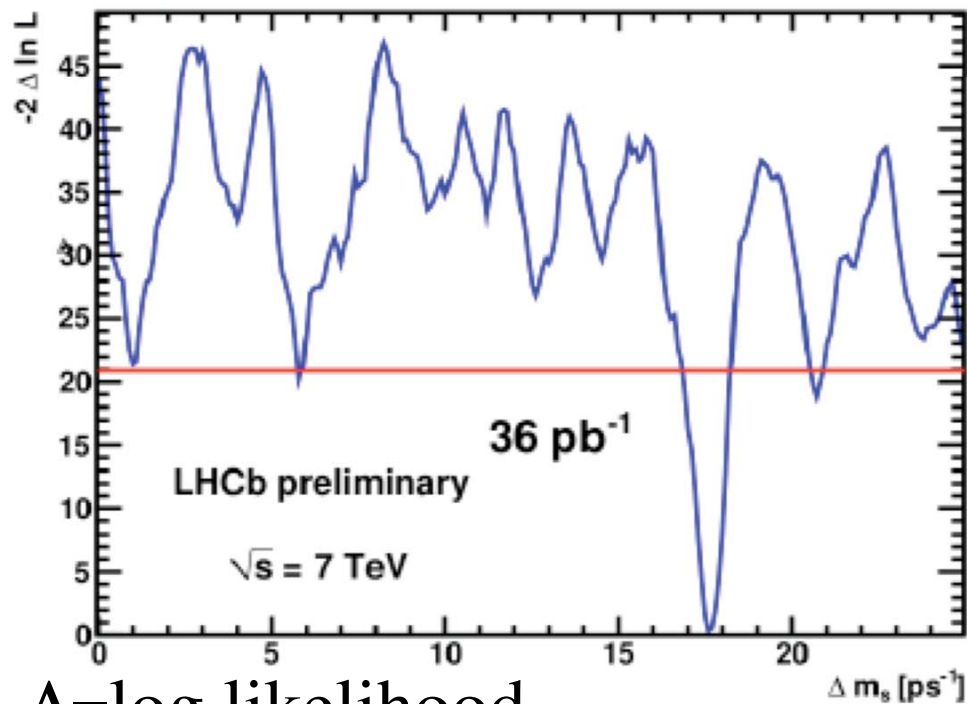
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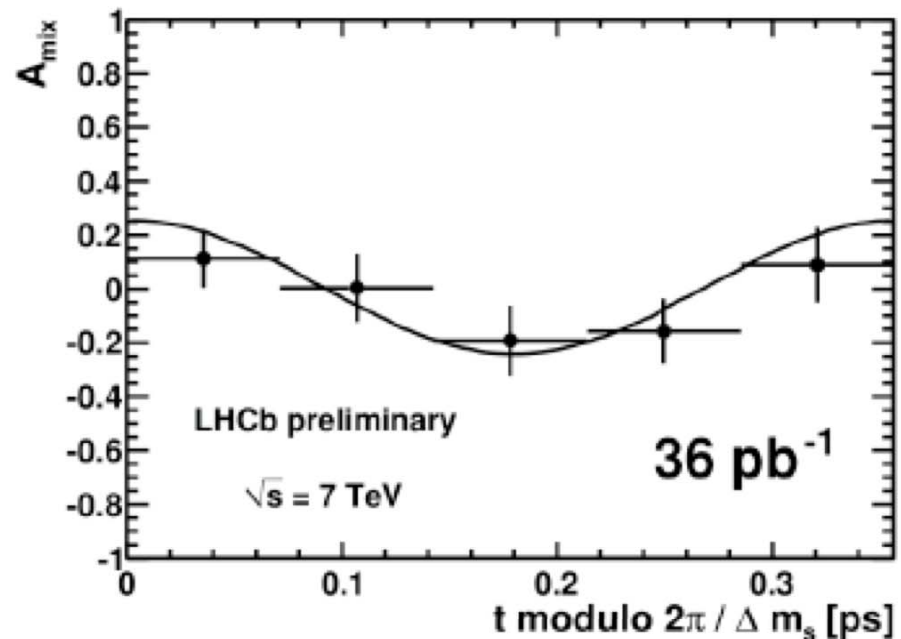
$$B_s^0 - \bar{B}_s^0 \text{ oscillations: } \frac{1-2\omega}{2} \cos \Delta m_s \times t \otimes \sigma_t\text{-effect}$$

LHCb-CONF-2011-005



Δ -log likelihood

$\omega = \text{wrong-tag} / \text{all-tag}$



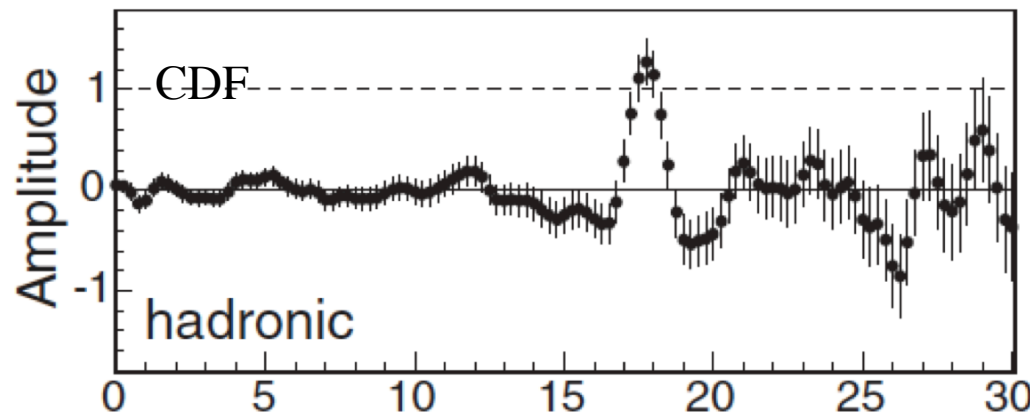
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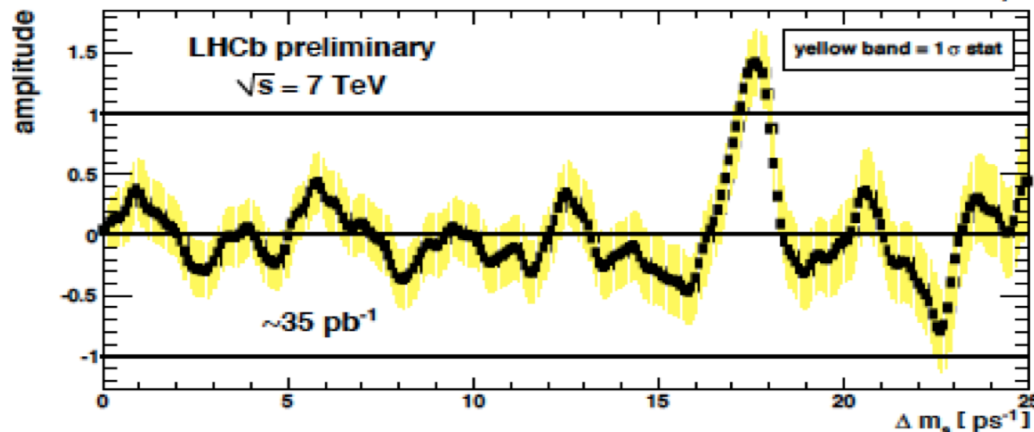
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$B_s^0 - \bar{B}_s^0$ oscillations: $A \times \cos \Delta m_s \times t \otimes \sigma_t$ -effect



CDF:
1 fb⁻¹

PhysRevLett.97.242003



LHCb:
36 pb⁻¹

LHCb-CONF-2011-005

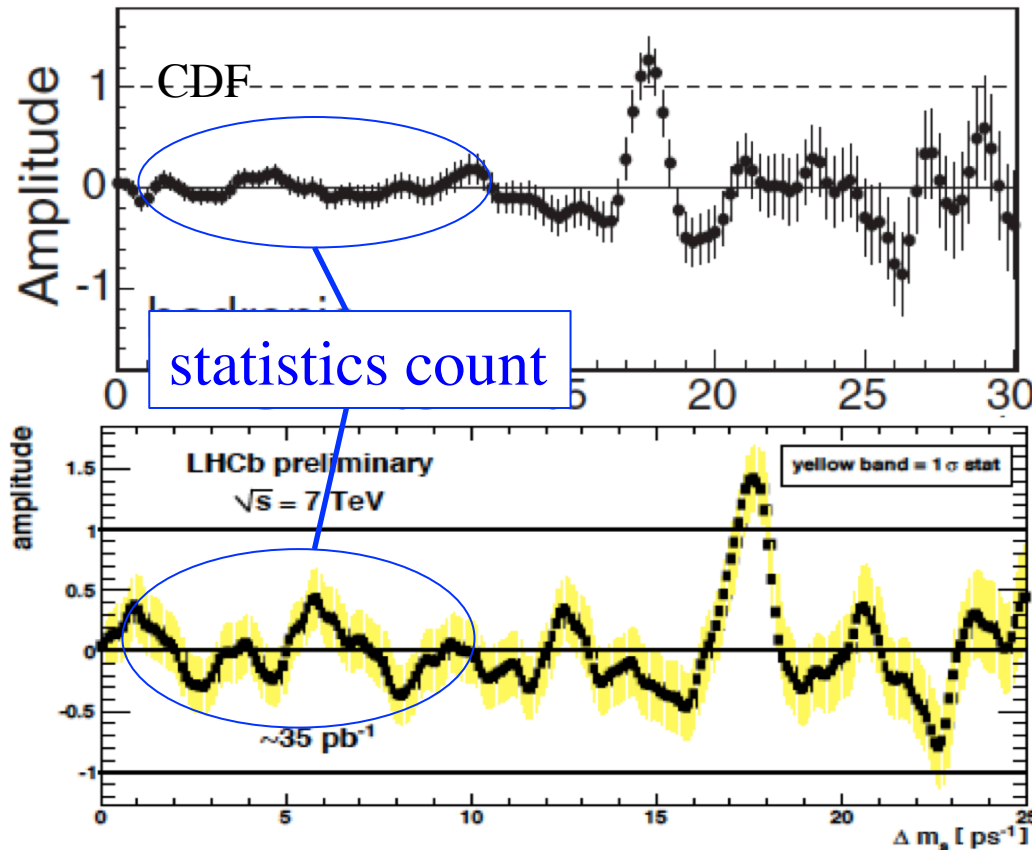
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CDF:

1 fb⁻¹

5600 signal

LHCb:

36 pb⁻¹

1350 signal

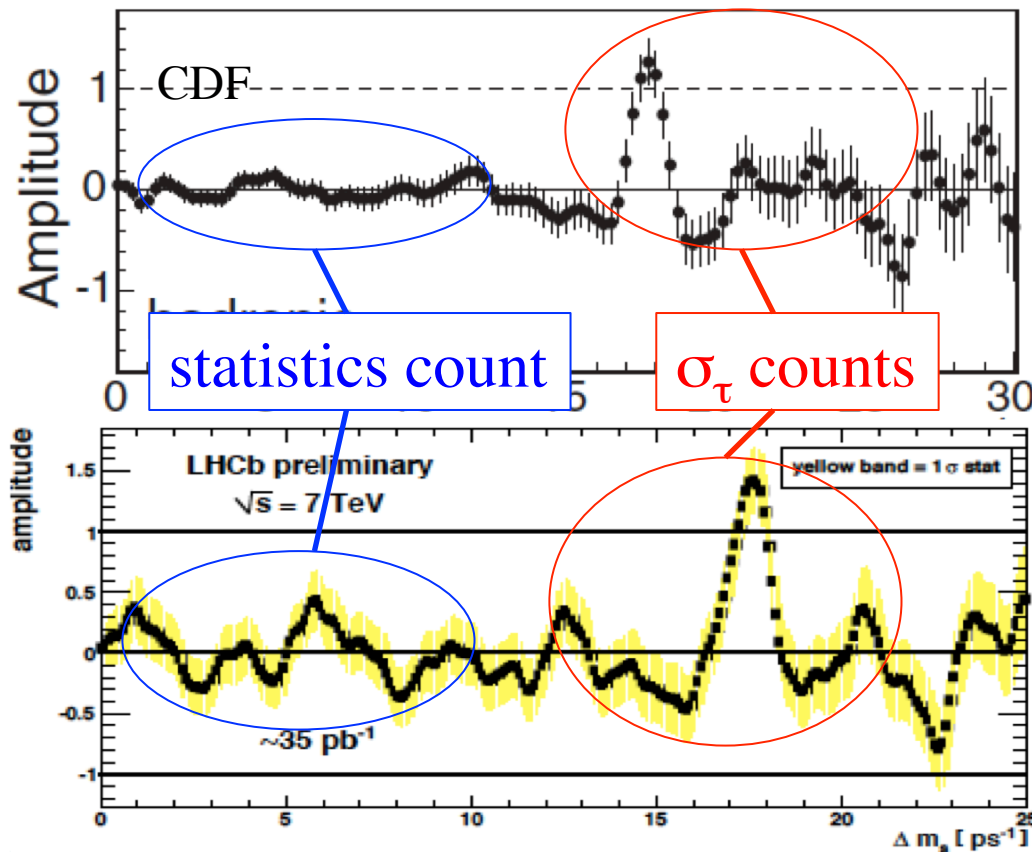
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CDF:

1 fb⁻¹

5600 signal, $\sigma_\tau=87$ fs

LHCb:

36 pb⁻¹

1350 signal, $\sigma_\tau=36$ or 44 fs

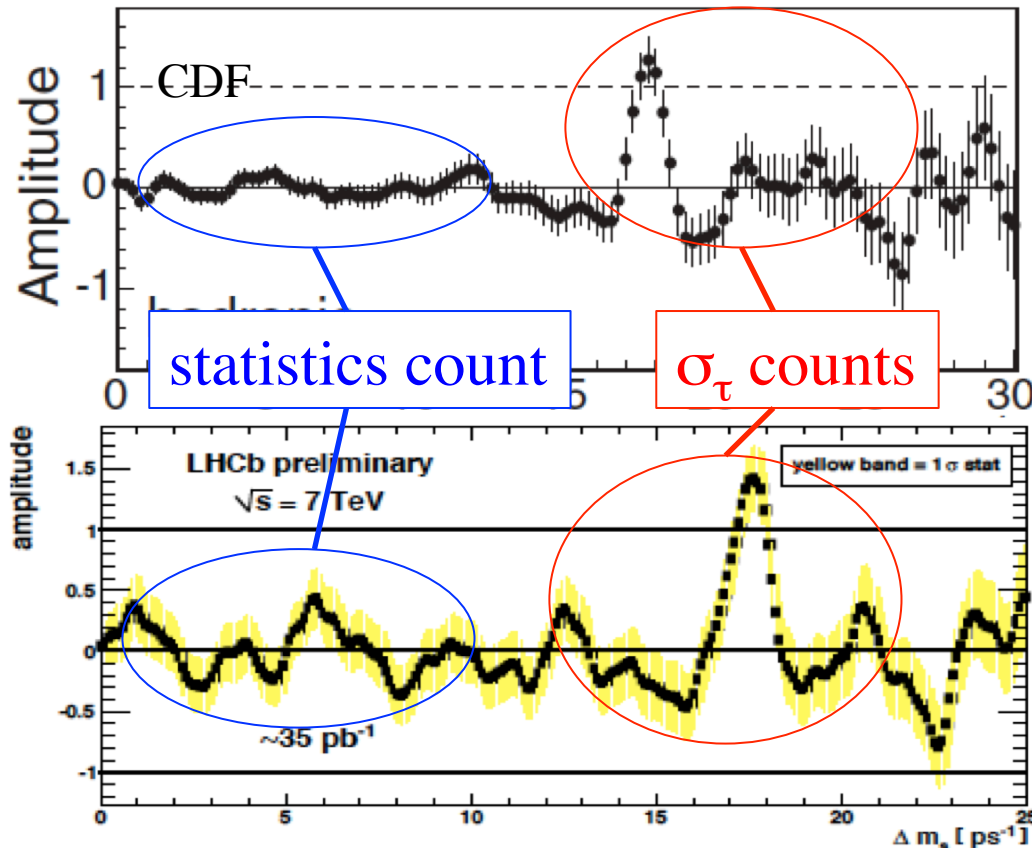
Flavour Physics at Hadron Machines

Advantage of LHCb over CDF and D0 are:

-larger cross section

-better invariant mass and proper time resolutions

$B_s^0-\bar{B}_s^0$ oscillations: $A \times \cos \Delta m_s \times t \otimes \sigma_t$ -effect



CDF:

1 fb⁻¹

5600 signal, $\sigma_\tau=87$ fs

$\Delta m_s = 17.77 \pm 0.10 \pm 0.07$ ps⁻¹

LHCb:

36 pb⁻¹

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$\Delta m_s = 17.63 \pm 0.11 \pm 0.04$ ps⁻¹

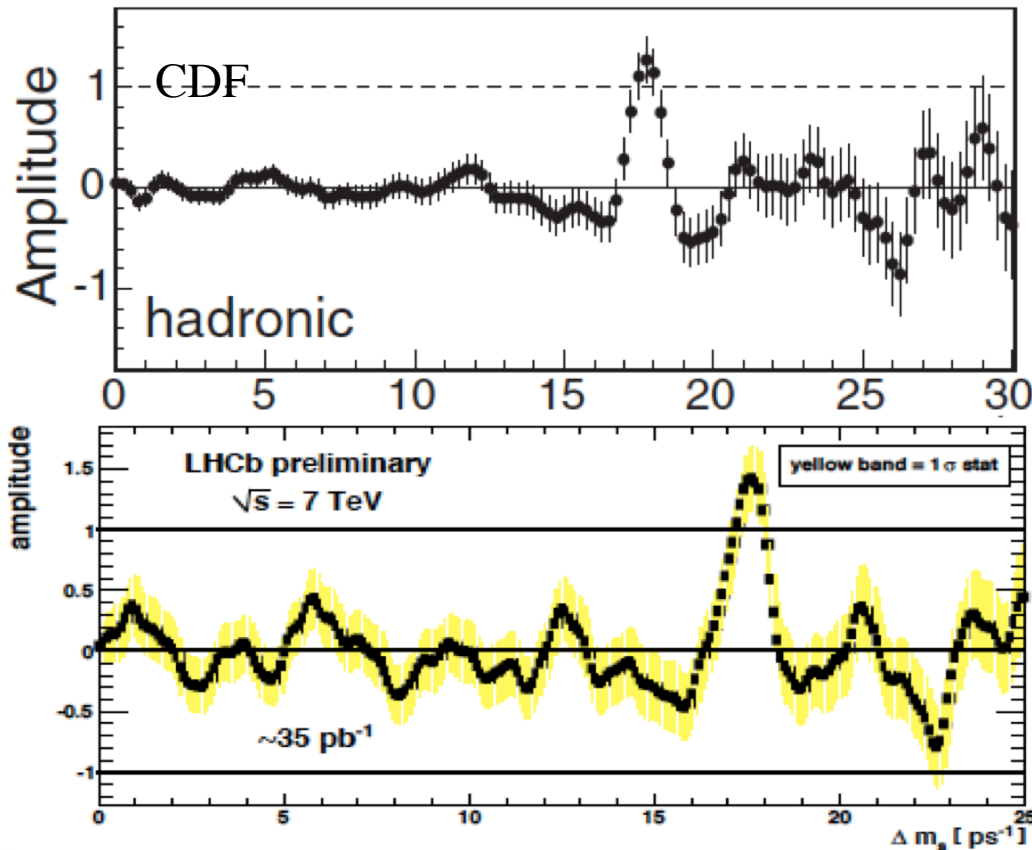
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$\beta\gamma c\tau_B \approx 1 \text{ cm}$

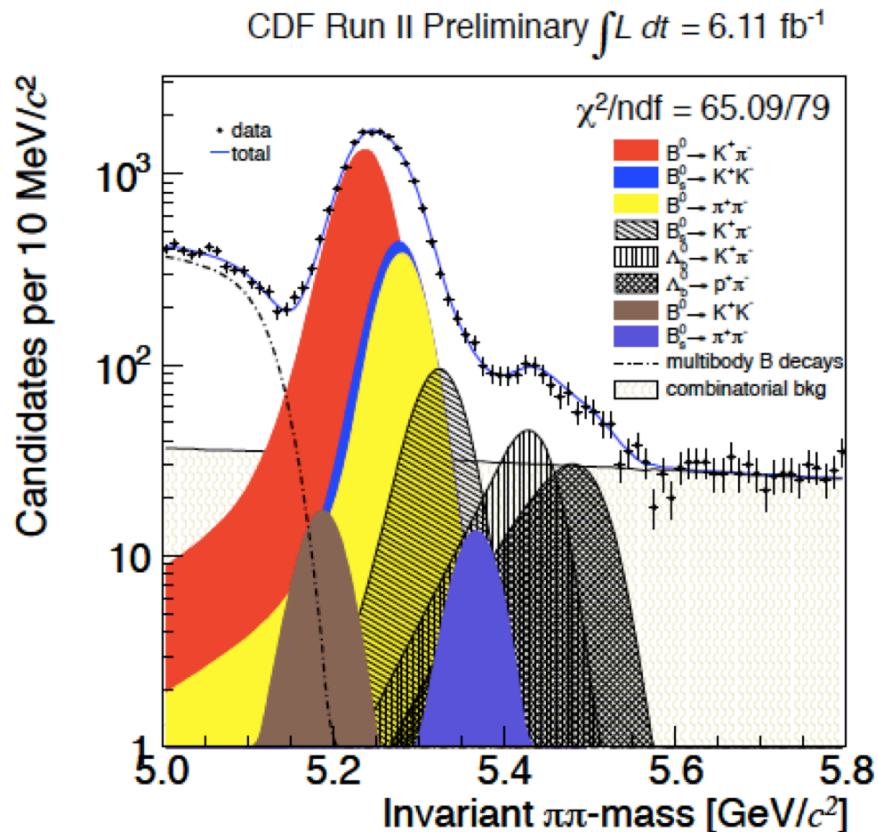
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Advantage of LHCb over CDF and D0 are:

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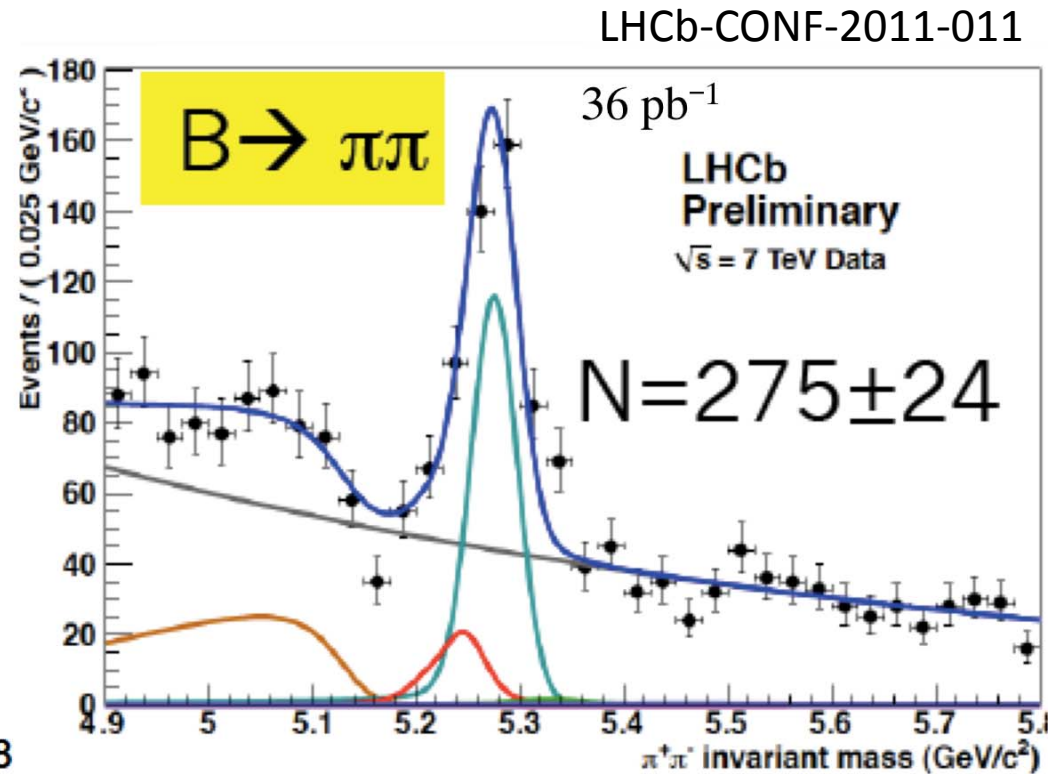
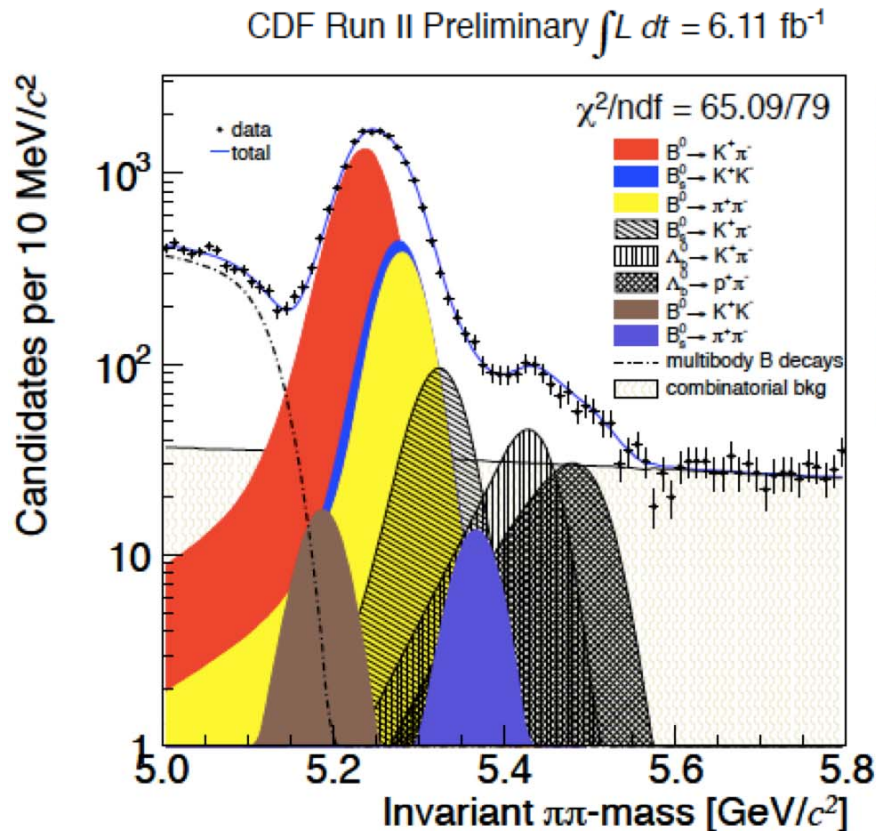


Public Note 10498

Flavour Physics at Hadron Machines

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Flavour Physics at Hadron Machines

Advantage of LHCb over CDF and D0 are:

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Reflects on the systematic errors

$A_{CP}(B_d \rightarrow K^+\pi^-)$

CDF $-0.086 \pm 0.023 \pm 0.009$ 1 fb^{-1}

LHCb $-0.074 \pm 0.033 \pm 0.008$ 36 pb^{-1}

σ_{stat} for BABAR: ± 0.016 BELLE: ± 0.018

Flavour Physics at Hadron Machines

Advantage of LHCb over CDF and D0 are:

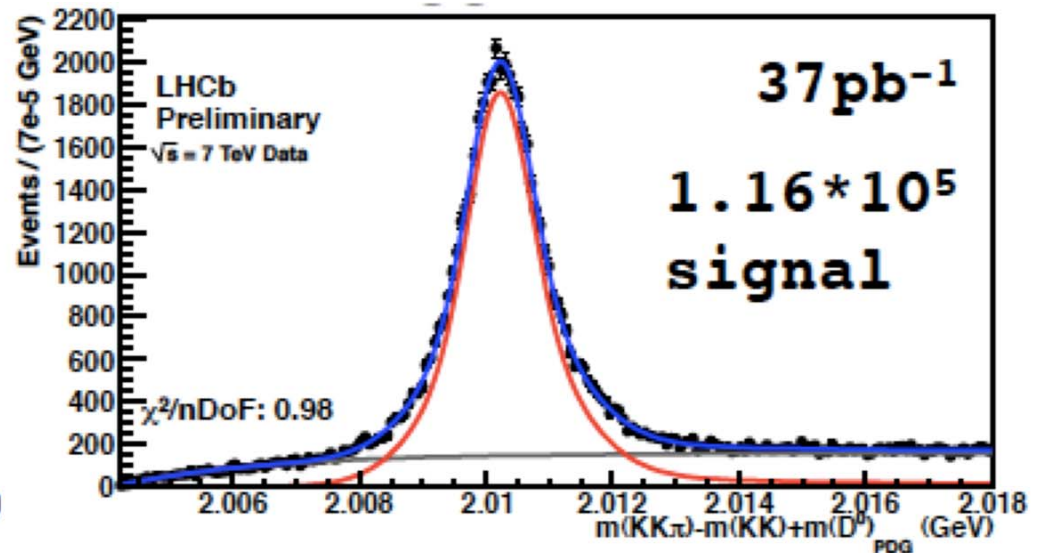
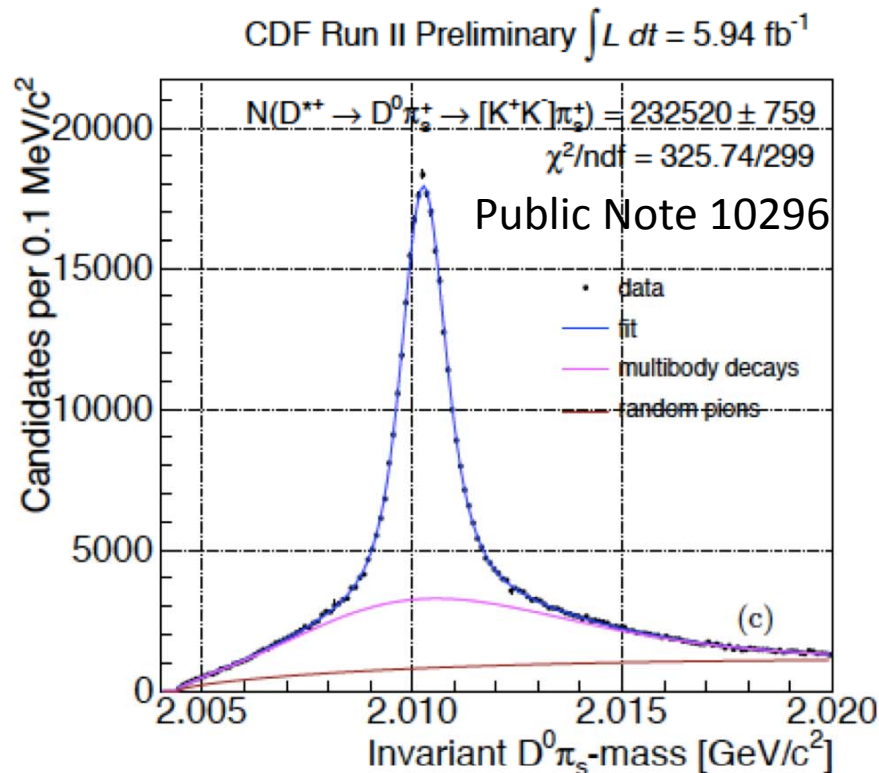
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CP violation in D decays is getting attention and trigger efficiency is an issue!

Flavour Physics at Hadron Machines

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CP violation in D decays is getting attention and trigger efficiency is an issue!

Flavour tagged D^0 and $\bar{D}^0 \rightarrow K^+K^-$

CDF	0.48×10^6	6 fb^{-1}
LHCb	0.12×10^6	36 pb^{-1}

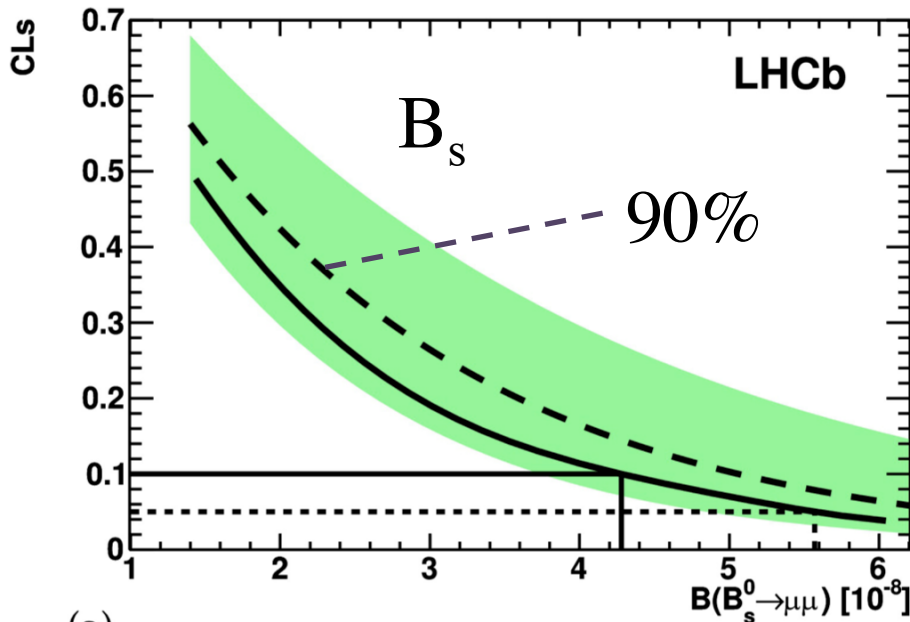
LHCb get 40 times more events for the same $\int L dt$

Flavour Physics at Hadron Machines

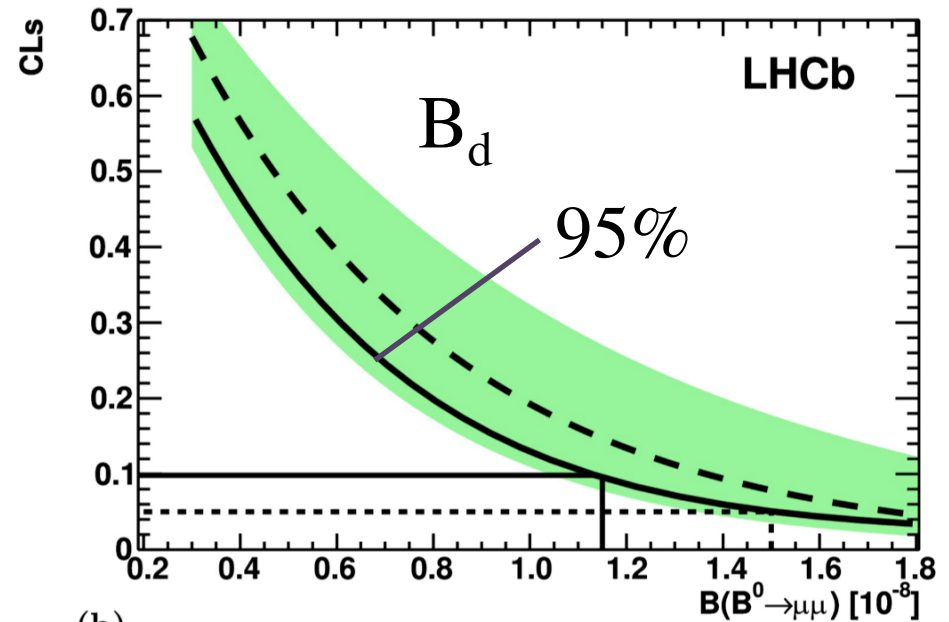
- By the end of 2012, LHCb will correct $\gtrsim 2 \text{ fb}^{-1}$ of data
 - Exclude NP physics contribution to $B_s \rightarrow \mu^+ \mu^-$ up to a level of SM (currently up to ~ 10 times SM) or find one (current best limit: 4.3×10^{-8} by CDF 3.7 fb^{-1})

SM prediction: $(0.32 \pm 0.02) \times 10^{-8}$
 $(0.010 \pm 0.001) \times 10^{-8}$

36 pb^{-1}



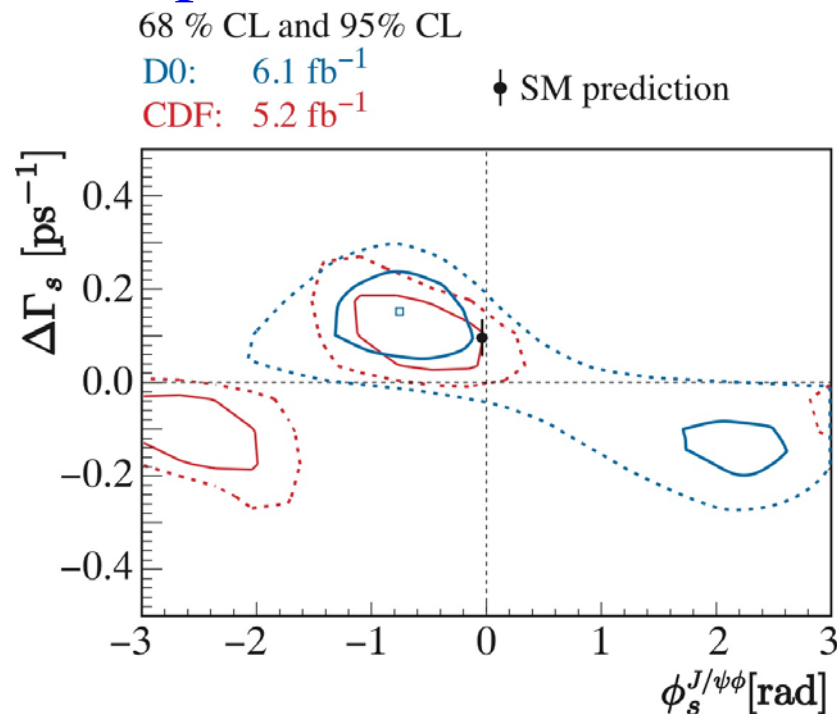
$< 5.6 \times 10^{-8}$ (95% CL)



$< 1.5 \times 10^{-8}$ (95% CL)

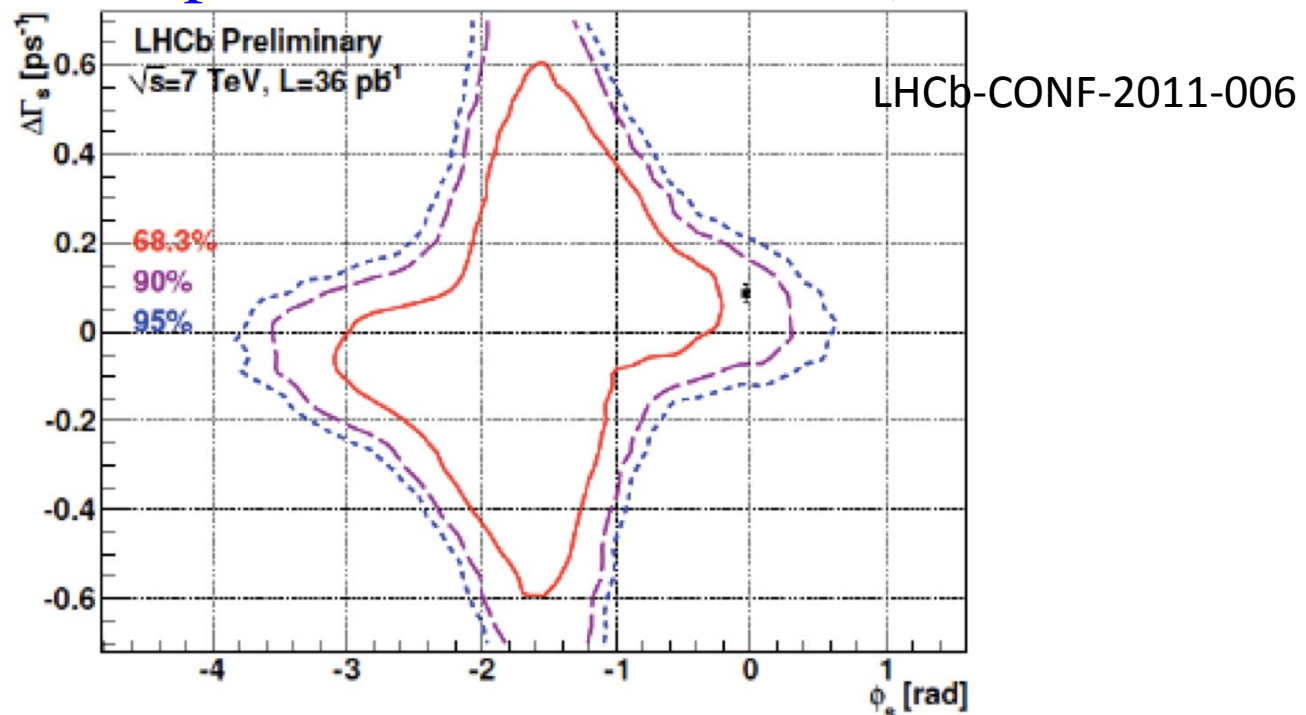
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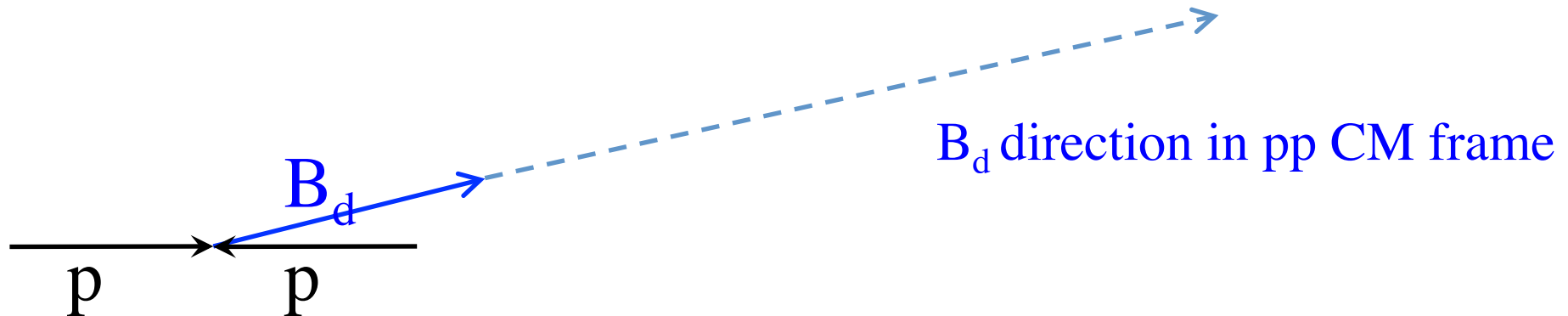


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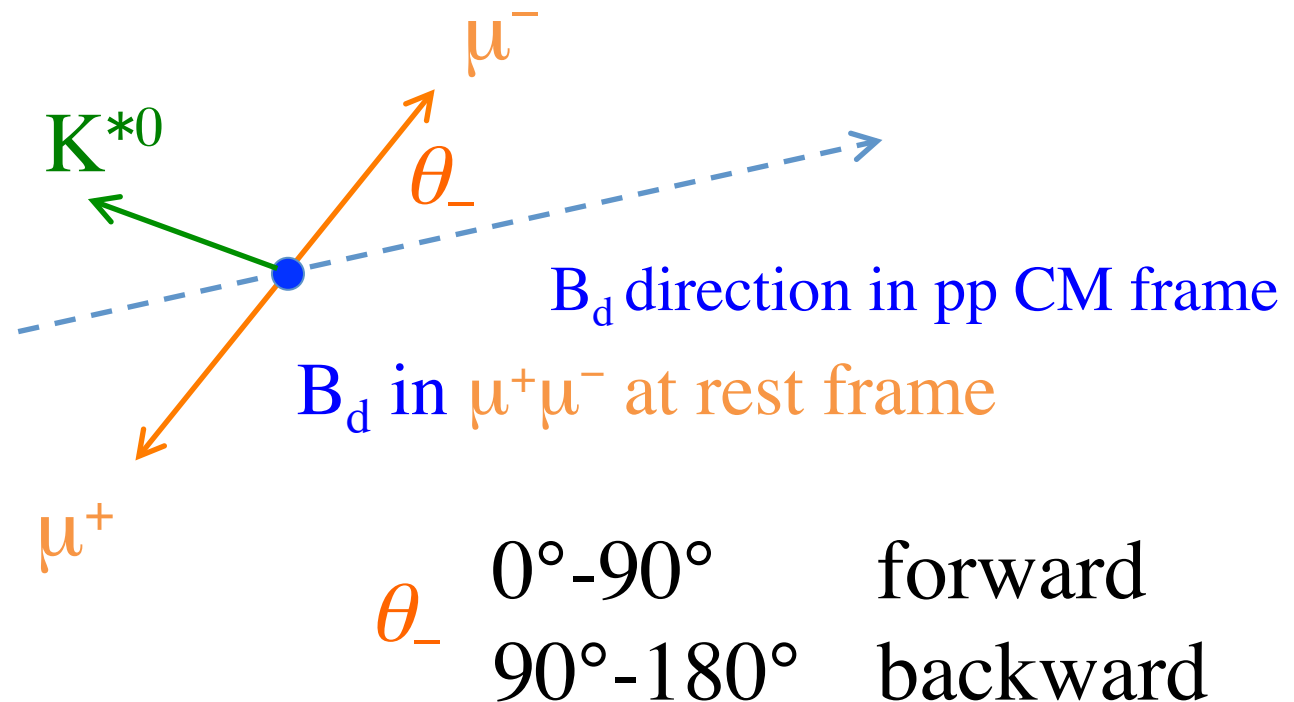
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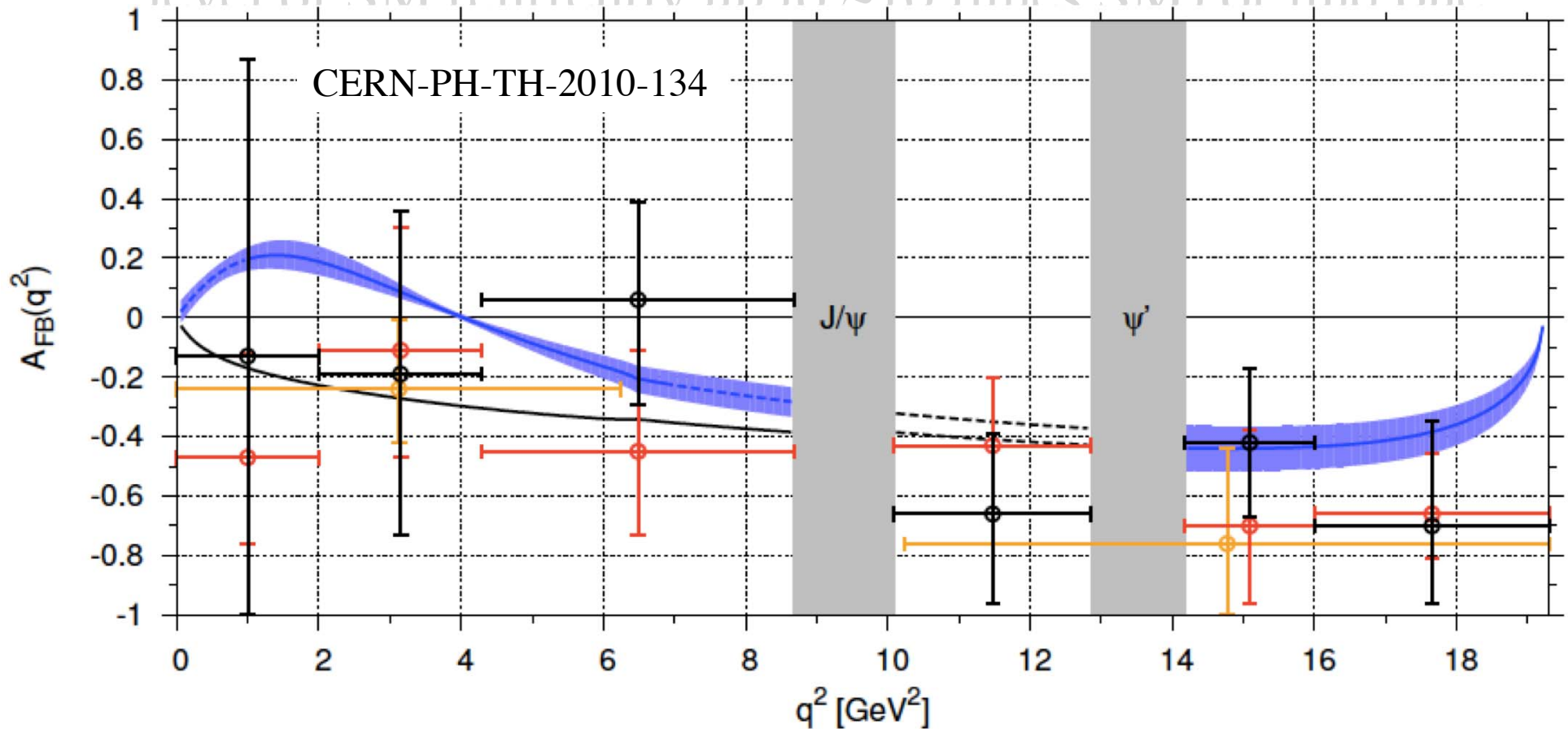
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BABAR(75%) BELLE(80%) CDF (4.4 fb^{-1}) SM prediction

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And much more

Flavour Physics at Hadron Machines

- And more than 5 fb^{-1} by the end of 2017!

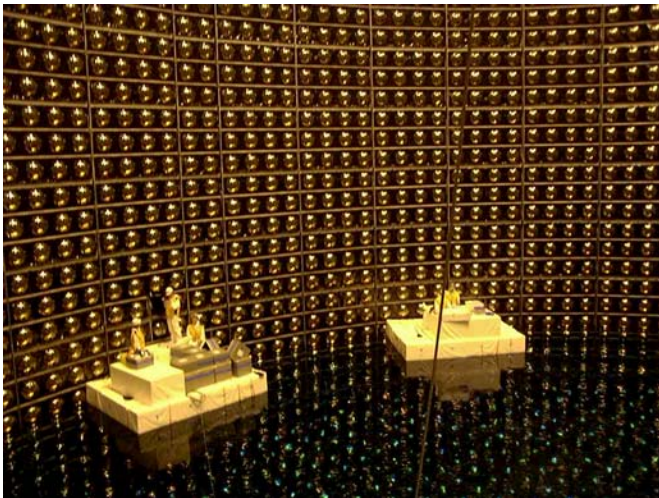
Epilogue

- There exists solid observations for physics beyond the Standard Model

Epilogue

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Neutrino oscillations

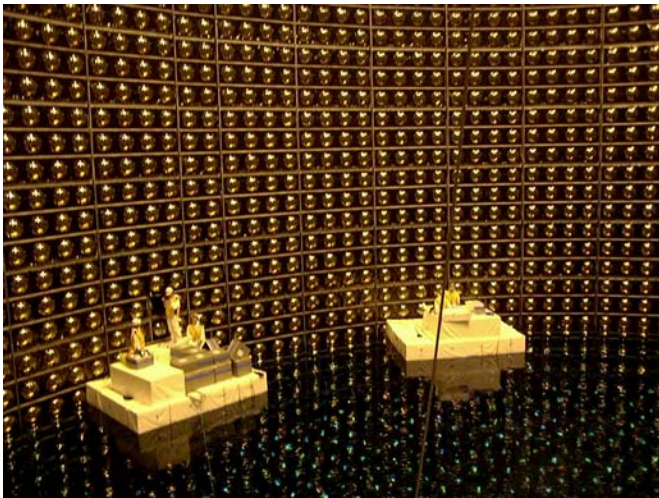
S-KAMIOKANDE



Epilogue

- There exists solid observations for physics beyond the Standard Model
 - Neutrino oscillations
 - Dark matter

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Bullet Galaxy Clusters



Epilogue

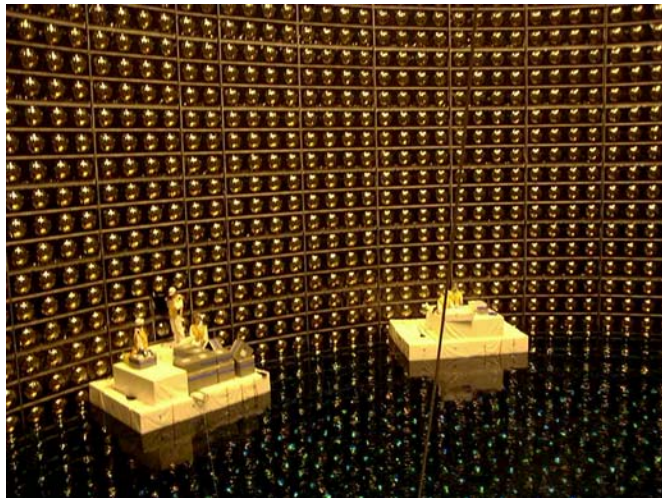
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Neutrino oscillations

Dark matter

$$N_B / N_\gamma = 10^{-10}$$

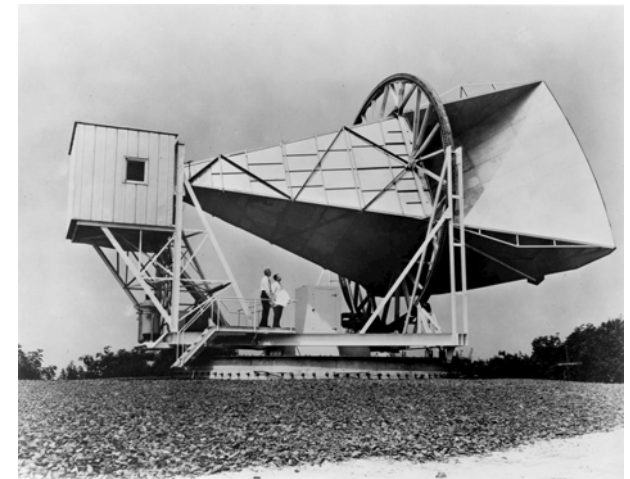
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Bullet Galaxy Clusters



The Horn Antenna
Bell Telephone Laboratory



Epilogue

- There exists a **strong anticipation** that new physics is just around the corner...


Epilogue

By the middle of 2012...

ATLAS CMS high p_T physics
LHCb flavour physics
Particle Physics



Epilogue

By the middle of 2012...

ATLAS CMS high p_T physics	BSM
LHCb flavour physics	Only SM
Particle Physics	




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Particle Physics			

Epilogue

By the middle of 2012...

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LHCb flavour physics	Only SM	BSM	BSM	
Particle Physics				

Oh, no more space left...

Epilogue

In any case,



Exciting time is ahead of us all.

Epilogue

T2K preprint,

(Dated: June 13, 2011)

the expected number of such events is $1.5 \pm 0.3(\text{syst.})$. Under this hypothesis, the probability to observe six or more candidate events is 7×10^{-3} , equivalent to 2.5σ significance. At 90% C.L., the data are consistent with $0.03(0.04) < \sin^2 2\theta_{13} < 0.28(0.34)$ for $\delta_{\text{CP}} = 0$ and normal (inverted) hierarchy.

Exciting time is ahead of us all.