

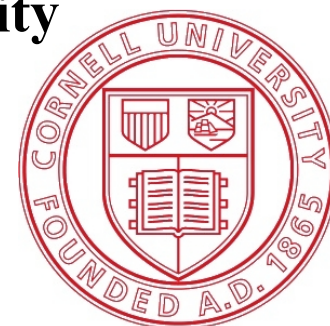


# Charged Particles from Dark Matter: Galactic Propagation Effects

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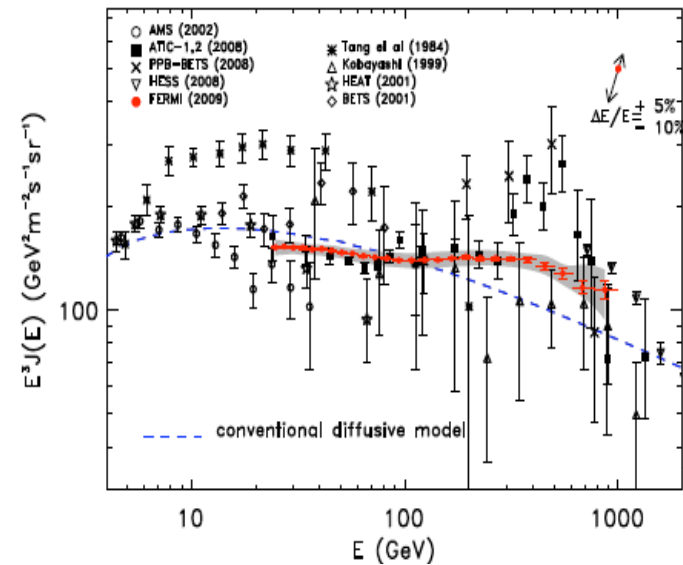
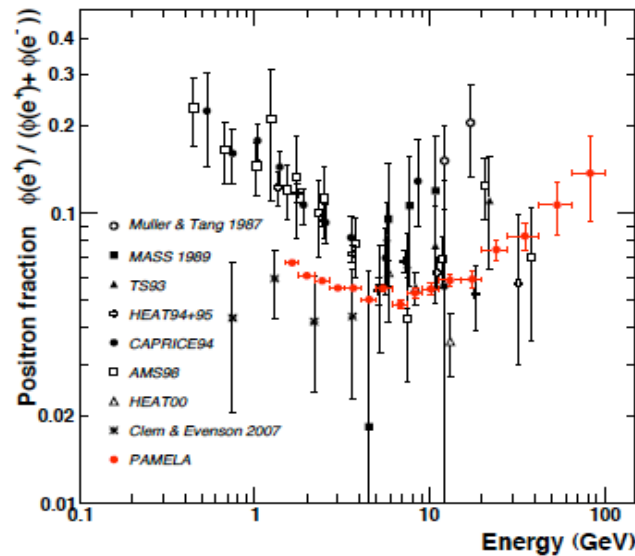
**Bibhushan Shakya**  
**Cornell University**

PHENO 2011  
May 9, 2011



Based on arXiv: 1012.3772, 1002.4588 with Maxim Perelstein

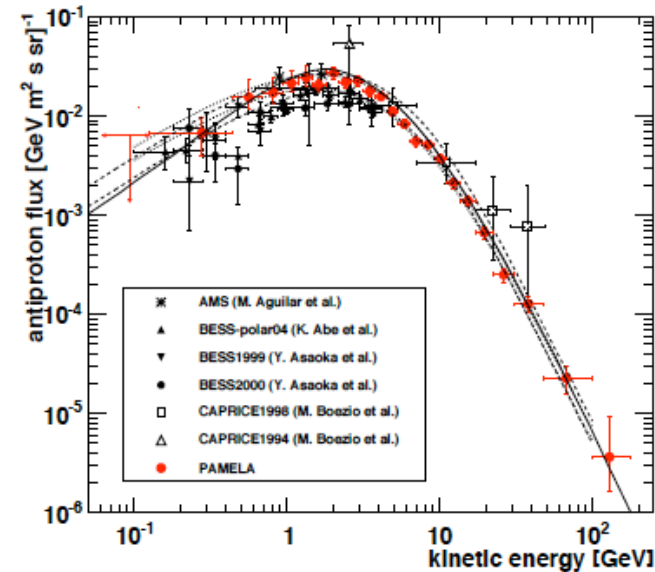
# Indirect Evidence of Dark Matter?



Data:

excess in  $e^+$  fraction and  $e^+e^-$  flux, no excess in antiprotons

Fit well to leptophilic dark matter annihilation with boosted cross sections in the galaxy



# Connecting Theory to Observation

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Spectrum  
observed  
experimentally

Annihilation cross  
section predicted by  
theory

$$\frac{d\phi}{dE} = \int_{E_{min}}^{m_\chi} dE' (\text{astrophysical correction}) \times \frac{d\sigma}{dE'}$$

# Connecting Theory to Observation

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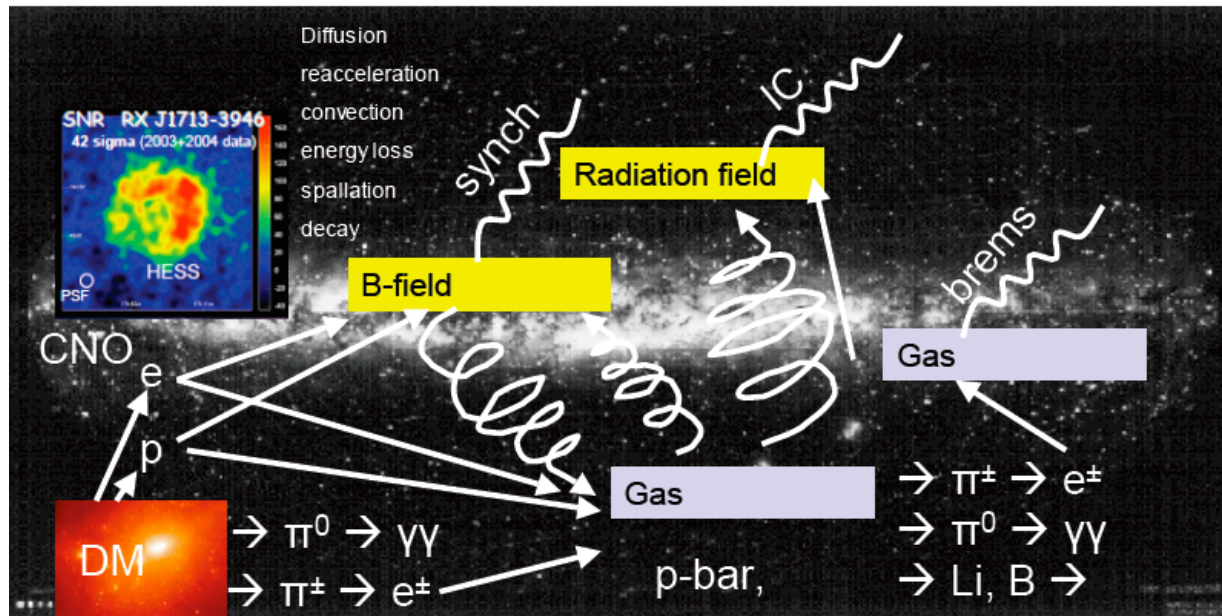
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Spatial distribution of dark matter sources,  
propagation effects, interaction with  
interstellar matter....

## (Nontrivial) Cosmic ray propagation in the galaxy

**A mess!** Electrons, positrons from dark matter annihilation interact with the galactic interstellar medium, losing energy and directional information.



From Jan Conrad's 08/07/09 talk at the SLAC Summer Institute

Galactic propagation significantly alters the spectrum

# The diffusion equation (positrons)

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$$-\nabla [K(\mathbf{x}, E) \nabla \psi] - \frac{\partial}{\partial E} [b(E)\psi] = q(\mathbf{x}, E) \quad \text{Source term}$$

position independent  
Diffusion coefficient  
 $K(E) = K_0 \epsilon^\delta$

$$\epsilon = E/E_0, E_0 = 1 \text{ GeV}$$

Energy loss rate

$$b(E) = \frac{E_0}{\tau_E} \epsilon^2$$

$$\tau_E = 10^{16} \text{ sec}$$

Positron density per unit volume per unit energy

Parameters defining a propagation model:

$K$ ,  $\delta$ ,  $L$  (size of region in which transport equation solved)

Model	$\delta$	$K_0$ [kpc <sup>2</sup> /Myr]	$L$ [kpc]
MED	0.70	0.0112	4
M1	0.46	0.0765	15
M2	0.55	0.00595	1

# The diffusion equation (antiprotons)

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Convective wind term

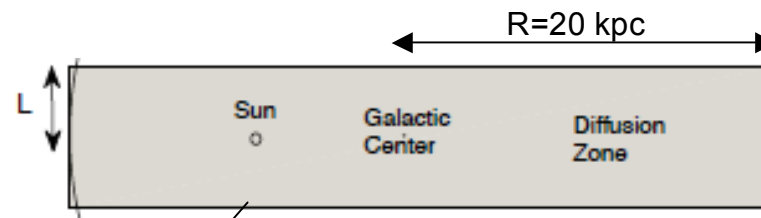
$$-\nabla [K(\mathbf{x}, E) \nabla n_{\bar{p}}] + \frac{\partial}{\partial z}(V_C(z)n_{\bar{p}}(E, \mathbf{x})) + 2h\delta(z)\Gamma_{ann}n_{\bar{p}}(E, \mathbf{x}) = q_{\bar{p}}(\mathbf{x}, E)$$

Antiproton interaction with interstellar medium, confined to galactic plane

Omitted: energy loss term  
(negligible for the more massive antiprotons)

Model	$\delta$	$K_0$ (kpc <sup>2</sup> /Myr)	$L$ (kpc)	$V_C$ (km/s)
MIN	0.85	0.0016	1	13.5
MED	0.70	0.0112	4	12
MAX	0.46	0.0765	15	5

Conventional approach: solve transport equation in a thin cylindrical disk (half-thickness  $L$ , radius  $R = 20\text{kpc}$ ), the *diffusion zone*, where galactic magnetic fields confine positrons. Outside this region, positrons are assumed to propagate freely and escape.



$\psi=0$  at this boundary



# Conventional Solution (positrons)

[T. Delahaye, R. Lineros, F. Donato, N. Fornengo and P. Salati, Phys. Rev. D 77, 063527 (2008)]

1. Write positron density as a Bessel-Fourier series

$$\psi(z, r, \epsilon) = \sum_i \sum_n P_{i,n}(\epsilon) J_0 \left( \frac{\alpha_i r}{R} \right) \sin \left( \frac{n\pi(z + L)}{2L} \right)$$

2. Change variables as  $t = \frac{\tau_E \epsilon^{\delta-1}}{1-\delta}$ ,  $\tilde{P}_{i,n} = \epsilon^2 P_{i,n}$

and take Bessel and Fourier transforms of the transport equation:

$$\frac{d\tilde{P}_{i,n}}{dt} + K_0 \left( \left( \frac{\alpha_i}{R} \right)^2 + \left( \frac{n\pi}{2L} \right)^2 \right) \tilde{P}_{i,n} = \epsilon^{2-\delta} Q_{i,n} \quad (Q_{i,n}: \text{Bessel-Fourier transform of source term } q)$$

3. Solve to get:  $\tilde{P}_{i,n}(t) = \int_0^t \tilde{Q}_{i,n}(t_S) \exp[-\omega_{i,n}(t - t_S)] dt_S$

$$\text{where } \tilde{Q}_{i,n} = \epsilon^{2-\delta} Q_{i,n} \text{ and } \omega_{i,n} = K_0 \left[ \left( \frac{\alpha_i}{R} \right)^2 + \left( \frac{n\pi}{2L} \right)^2 \right]$$

## Conventional Solution (antiprotons)

Expand antiproton density as a Bessel series

$$n_{\bar{p}}(\rho, z, E) = \sum_i N_i(z, E) J_0 \left( \frac{\zeta_i \rho}{R} \right)$$

Assuming position independent K and V and solving in the cylindrical disk, the solution is

$$N_i(z) = e^{a(|z|-L)} \frac{y_i(L)}{B_i \sinh(S_i L/2)} [\cosh(S_i z/2) + A_i \sinh(S_i z/2)] - \frac{y_i(z)}{K S_i}$$

where

$$S_i = 2 \left( a^2 + \frac{\zeta_i^2}{R^2} \right)^{1/2}, \quad A_i = \frac{V_C + 2h\Gamma_{ann}}{K S_i}; \quad B_i = K S_i [A_i + \coth(S_i L/2)]$$
$$y_i(z) = 2 \int_0^z e^{a(z-z')} \sinh [S_i(z - z')/2] q_i(z') dz'. \quad a = V_C/(2K)$$

BUT these are extremely simplified propagation models, with many clearly unphysical approximations:

- The diffusion coefficient is position independent, but the diffusion behavior should be weaker away from the galactic center
- There is a sharp cutoff at the diffusion zone boundary, where the diffusion coefficient essentially gets an infinite jump
- The dark matter halo extends significantly beyond this disk, but these models neglect everything outside the diffusion zone boundary. (eg. for  $L=1$  kpc and dark matter with an isothermal profile, the diffusion zone contains only  $\sim 10\%$  of the dark matter mass.)

Justification: Simple, semi-analytic, easy to implement solutions.

# Can do better: A more realistic setup

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For antiprotons, the absence of an energy loss term allows significant analytic progress

Diffusion: charged particles getting confined by galactic magnetic fields  $\longrightarrow$  diffusion coefficient should follow spatial variations of the galactic magnetic field strength

$$B(\rho, z) \approx (11\mu\text{G}) \times \exp\left(-\frac{\rho}{10 \text{ kpc}} - \frac{|z|}{2 \text{ kpc}}\right) \implies K(E, z) = K_e(E) \exp(|z|/z_t)$$

Has been studied numerically, have best fit parameters

(Evoli, Gaggero, Grasso, Maccione, JCAP 0810, 018 (2008) )

For antiprotons, if the convective wind term is assumed to have a similar exponential profile (or can be neglected), **CAN solve the diffusion equation analytically!**

# New Solution

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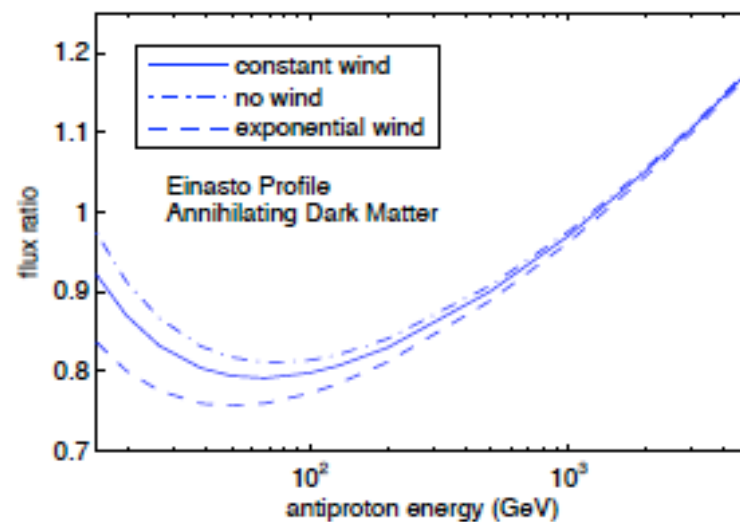
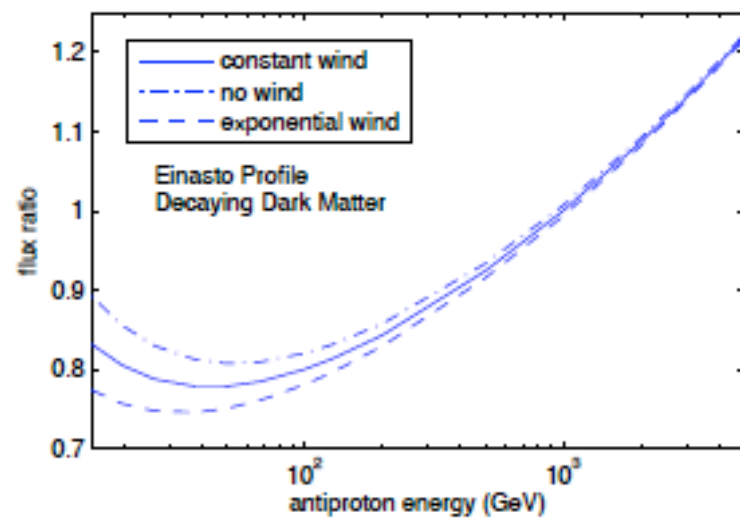
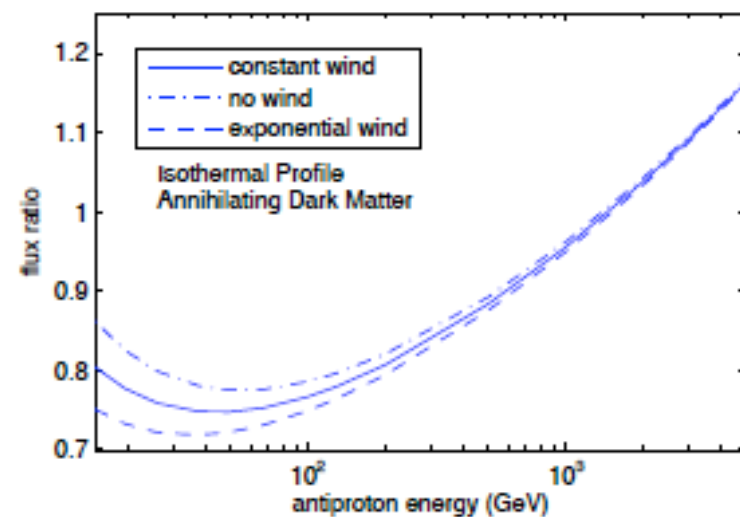
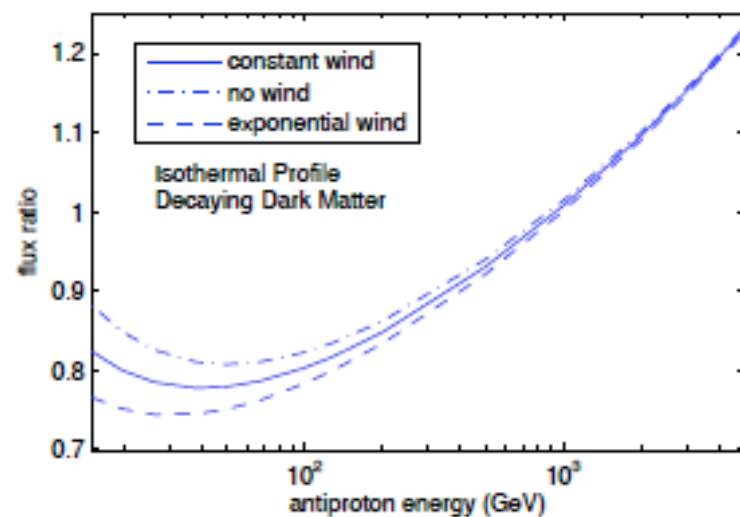
$$N_i(z) = e^{a(|z|-L)} \frac{y_i(L)}{B_i \sinh(S_i L/2)} [\cosh(S_i z/2) + A_i \sinh(S_i z/2)] - \frac{y_i(z)}{K_e S_i}$$

same form as the conventional solution, with slightly different definitions

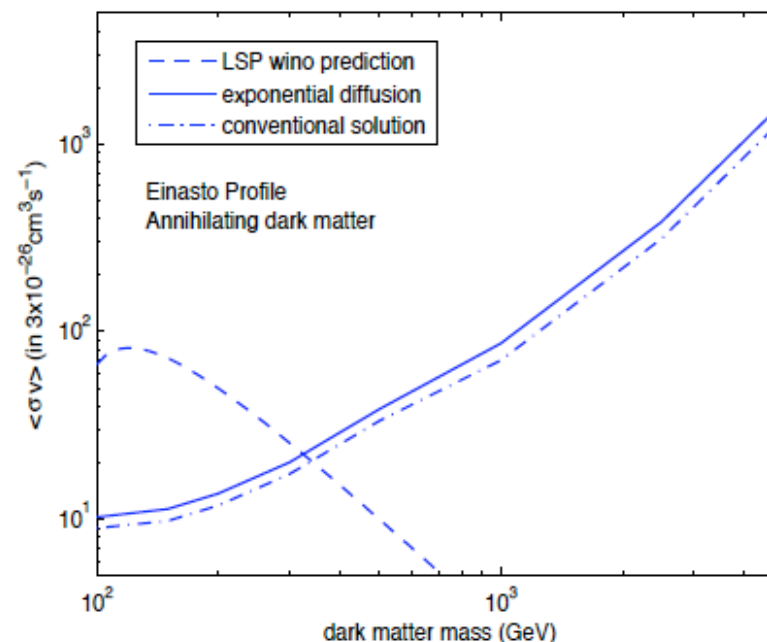
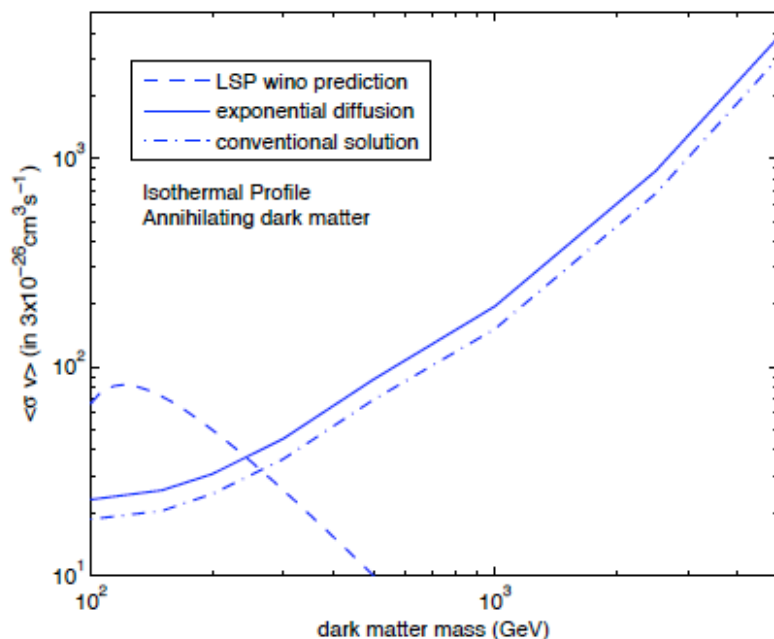
$$\begin{aligned} a &= \frac{V_e}{2K_e} - \frac{1}{2z_t}, \quad S_i = 2 \left( a^2 + \frac{\zeta_i^2}{R^2} + \frac{V_e}{z_t K_e} \right)^{1/2}, \\ A_i &= \frac{V_e + 2h\Gamma_{ann}}{K_e S_i} + \frac{1}{z_t S_i}; \quad B_i = K_e S_i [A_i + \coth(S_i L/2)] \\ y_i(z) &= 2 \int_0^z e^{a(z-z')} \sinh[S_i(z-z')/2] q_i(z') e^{-z'/z_t} dz'. \\ V_C(z) &= V_e \exp(|z|/z_t) \end{aligned}$$

As simple to evaluate as the conventional solution!

# Comparison with conventional solution



# Antiproton bounds for WIMP dark matter



Assume stable dark matter pair annihilating into  $W^+W^-$   
Bounds from conventional and new solutions agree to within  $\sim 20\%$

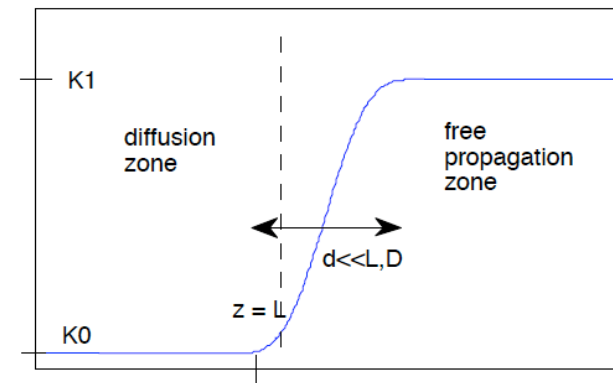
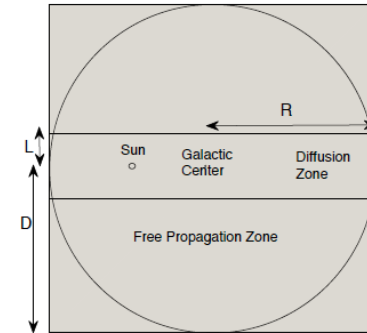
# Correction for positrons

1. Set boundary condition at  $|z|=D$ , not  $|z|=L$ .

2. Make diffusion coefficient position dependent to incorporate different behavior in diffusion and free propagation zones

$$K(z, \epsilon) = \left( K_0 + \tilde{K}(z) \right) \epsilon^\delta$$

- Different modes mix, equations no longer decoupled.
- Diagonalize matrices numerically, compute solution.
- 10-15% enhancement in for MIN model, percent level correction for MED model.





# Summary

- New, analytic, easy-to-use solution to antiproton flux from dark matter (valid at energies higher than several hundred GeV) in a more realistic propagation model, includes contributions from the full dark matter halo
- deviates from conventional solution by  $\sim 25\%$  for realistic parameters
- Smaller correction for positrons, negligible for MED propagation



“When is science going to explain the dark matter  
I find in my belly button every morning?”

# Questions