

Effect of exotic interacting massive particles including long-lived colored particles on big bang nucleosynthesis

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Kawasaki¹, MK, PRD83, 055011, arXiv:1012.0435

MK, Kajino², Yoshida¹, Mathews³, PRD80, 103501, arXiv:0906.3516

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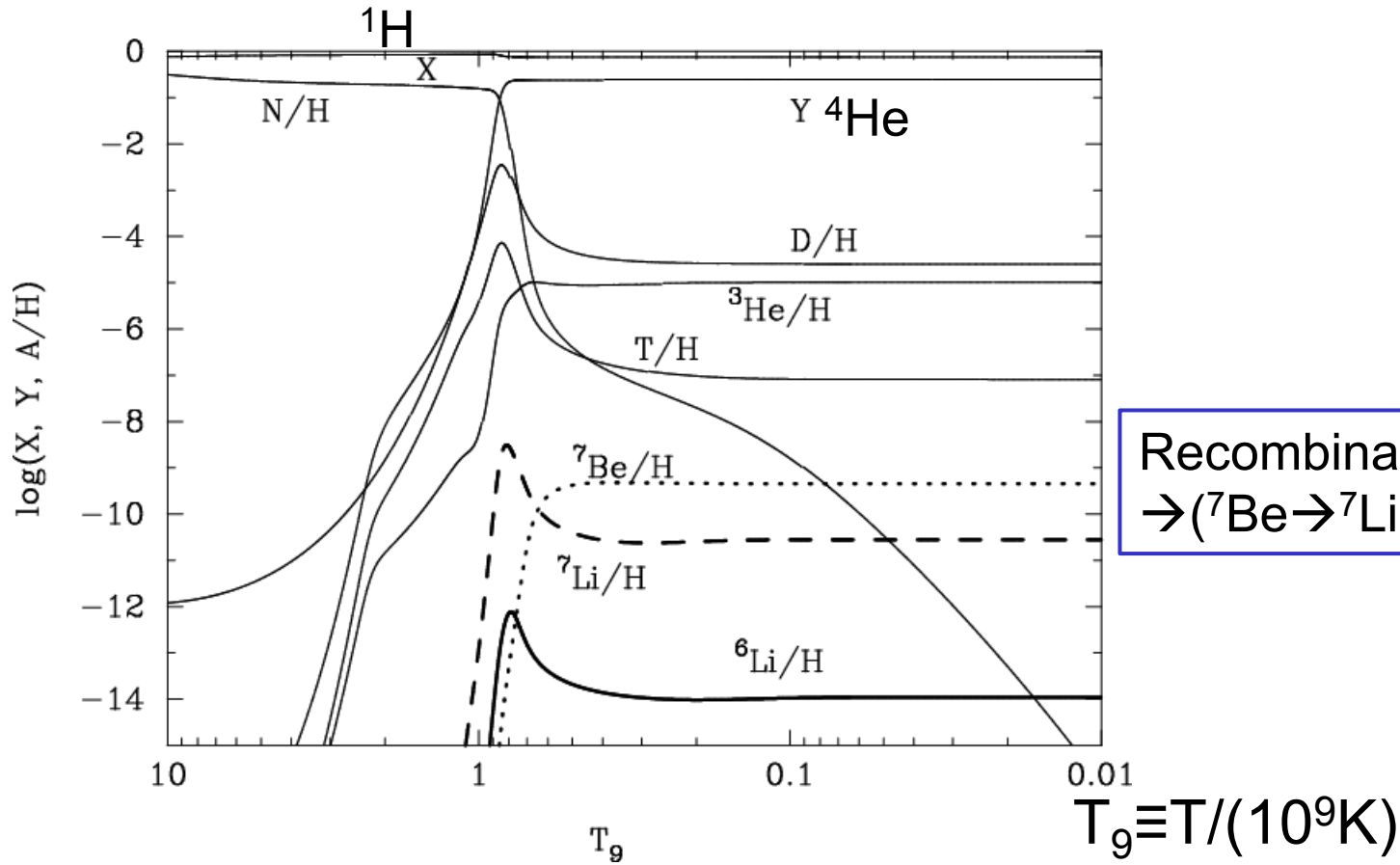
5) RIKEN Nishina Center

†) JSPS research fellow

2010/5/9

Introduction

Standard Big Bang Nucleosynthesis (BBN)



Observed abundances of light elements

WMAP7

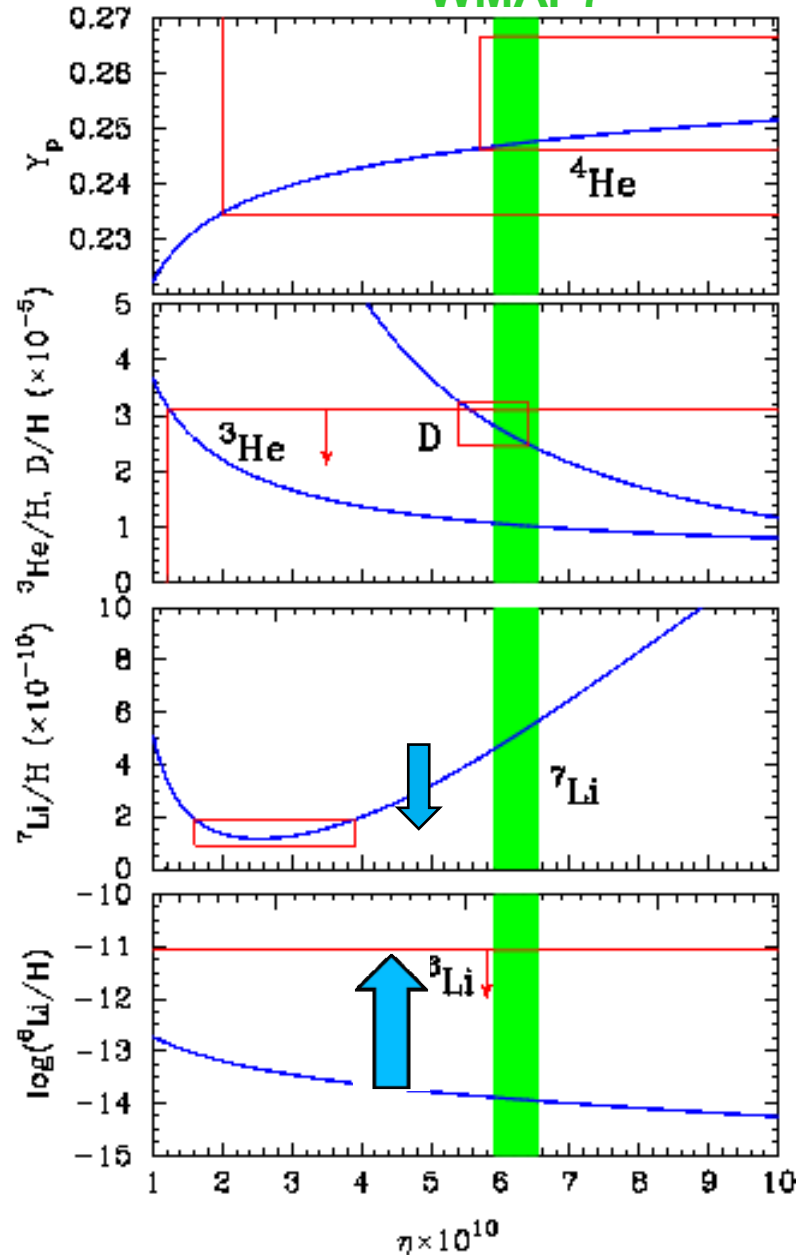
➤ SBBN: one parameter

baryon-to-photon ratio η

$$\leftarrow \eta = (6.225^{+0.157}_{-0.154}) \times 10^{-10}$$

(WMAP: Larson et al. 2010)

Li problems



Observed abundances of light elements

$$\log(^{6,7}\text{Li}/\text{H})+12$$

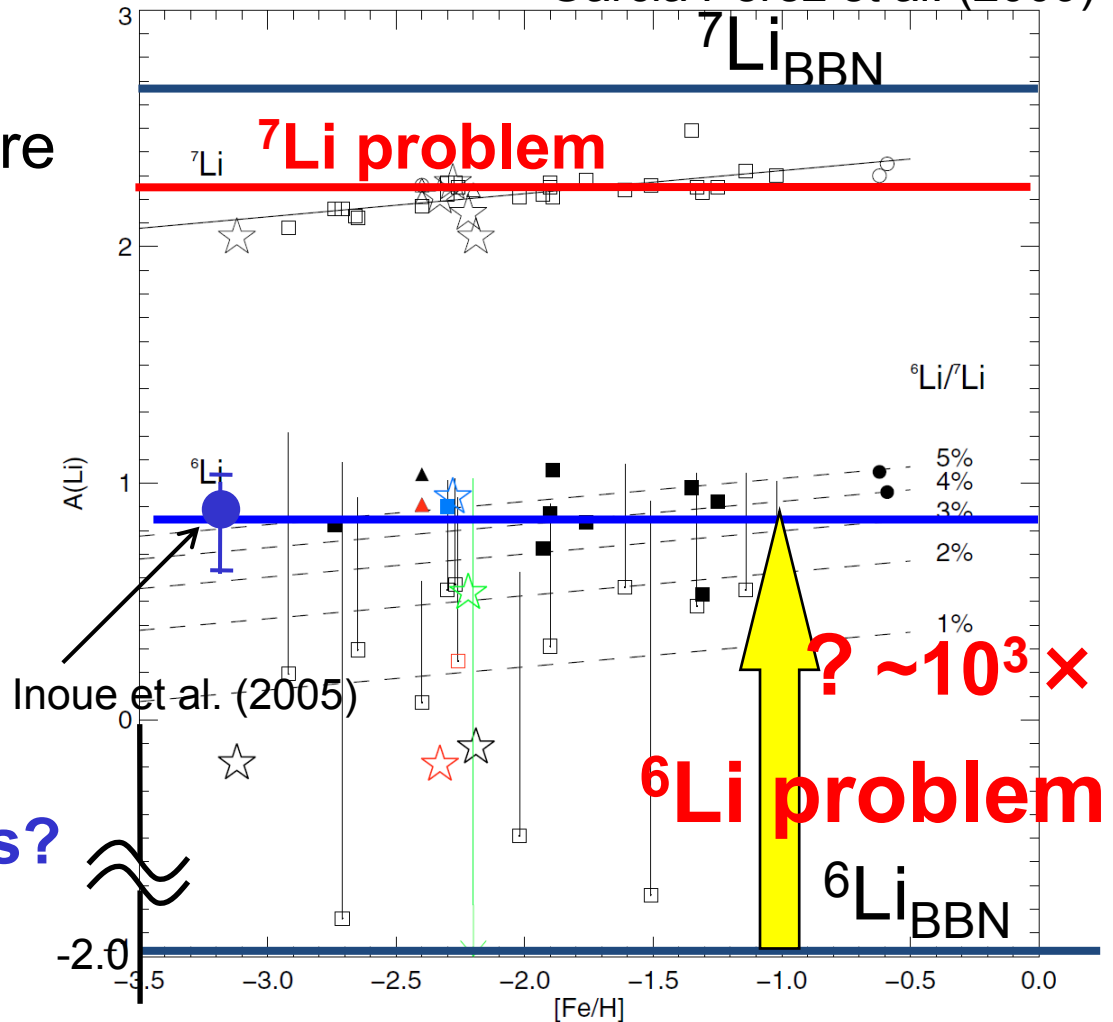
Garcia Perez et al. (2009)

- ^7Li in metal-poor stars are $\sim 1/3$ times (CMB+standard BBN prediction)

- ^6Li is abundant in some of the stars

Signature of new physics?

- ^9Be , B, C: plateau not seen

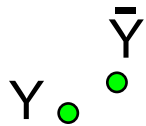


$$[\text{Fe}/\text{H}] = \log\left[\frac{(\text{Fe}/\text{H})}{(\text{Fe}/\text{H})_{\odot}}\right]$$

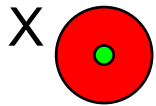
Long-lived Heavy Colored Particles

- long-lived heavy ($m \gg \text{GeV}$) colored particles appear in particle models beyond the standard model
e.g., Split SUSY (Arkani-Hamed & Dimopoulos 2005), ...

- History in the early universe (J. Kang et al. 2008)



- ✓ Long-lived exotic colored particles (Y)



- ✓ $T < T_c \sim 180 \text{ MeV}$

→ **Y particles get confined in hadrons X**

[strongly interacting massive particle (SIMP)]

(Walfran 1979; Dover et al. 1979)

- ✓ final abundance

$$\frac{n_X}{n_b} \approx 10^{-8} \left(\frac{R}{\text{GeV}^{-1}} \right)^{-2} \left(\frac{T_B}{180 \text{ MeV}} \right)^{-3/2} \left(\frac{m}{\text{TeV}} \right)^{1/2}$$

Studies on long-lived SIMP (X) in BBN

➤ NX force is the same as $N\Lambda$ force

(Dicus & Teplitz 1980, Plaga 1995, Mohapatra & Teplitz 1998)

✓ Mass of X is implicitly assumed to be $m_X \sim m_\Lambda = 1 \text{ GeV}$

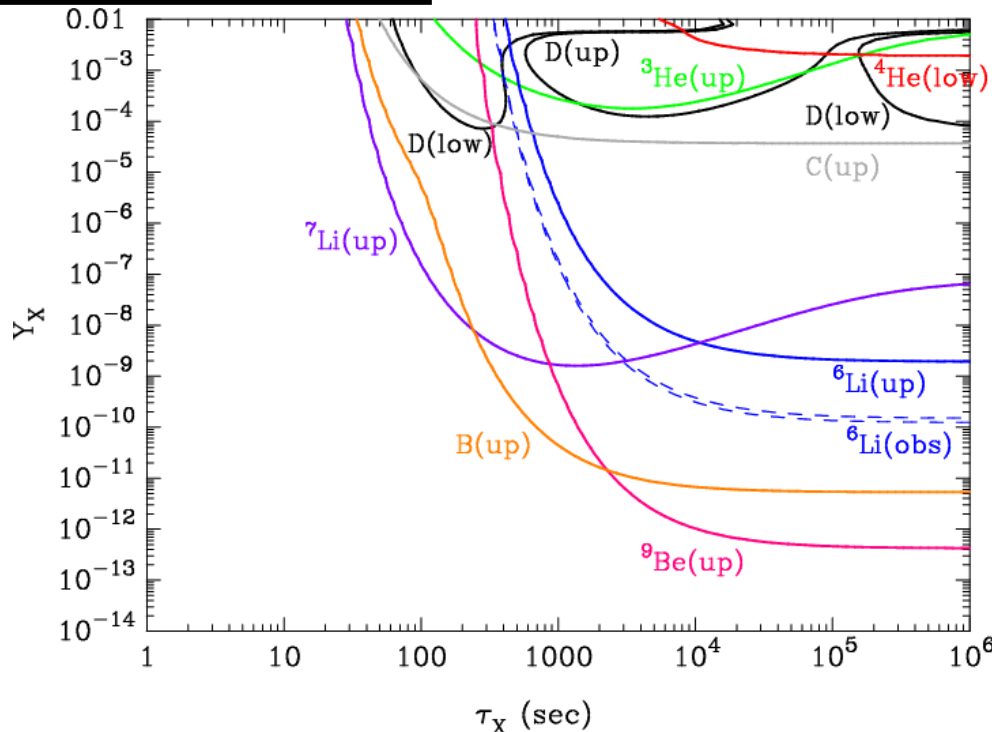
➤ NX force is same as the NN force, and $m_X \gg 1 \text{ GeV}$

(MK et al. 2009)

✓ Estimation of binding energies between A and X, and reaction rates

✓ Nonequilibrium calculation of abundances

Abundance n_X/n_b



✓ ${}^9\text{Be}$ and B can be produced more than in SBBN

✓ ${}^{10}\text{B}/{}^{11}\text{B} \sim 10^5$ **high ratio**
c.f. Galactic CR (${}^{10}\text{B}/{}^{11}\text{B} \sim 0.4$)
SN ν -process (${}^{10}\text{B}/{}^{11}\text{B} \ll 1$)

lifetime τ_X

Goal

- To investigate signatures of X particles on elemental abundances
- Interaction between X and N is unknown
→ studying multiple cases of interaction strength (δ)
- process of ${}^7\text{Be}$ destruction was found
→ solution to the ${}^7\text{Li}$ problem

Model

1. Binding energies of nuclides to an X

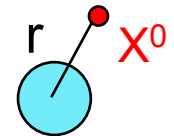
[Assumption]

➤ X (spin 0, **charge 0**, mass m_X)

➤ XN potential is Gaussian

parameter 1

Yahiro et al. (1986)



nucleon N

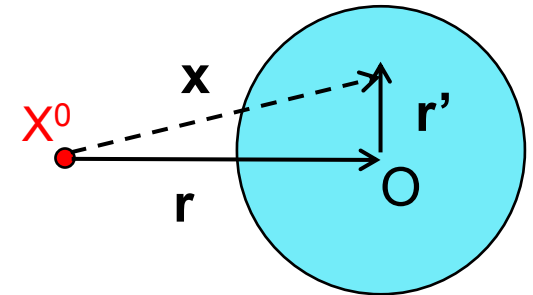
$$v(r) = v_0 \delta \exp \left[- (r/r_0)^2 \right]$$

$$v_0 = -72.15 \text{ MeV and } r_0 = 1.484 \text{ fm}$$

$\delta=1$ reproduces the binding energy of n+p

➤ XA potential: integration of (XN potential multiplied by nucleon densities)

$$V(r) = \int v(\mathbf{x}) \rho(\mathbf{r}') d\mathbf{r}'$$



nucleus A

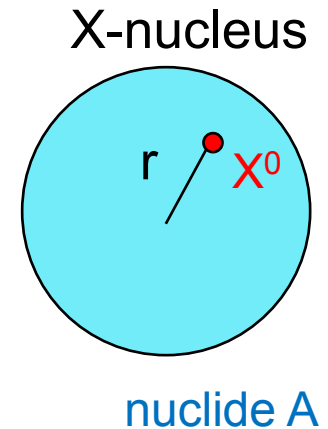
➤ $A \geq 2$ nuclei have Gaussian nucleon densities

Model

1. Binding energies of nuclides to an X

Schrödinger equation → binding energies and wave functions

$$\left[-\frac{\hbar^2}{2\mu} \nabla^2 + V(r) - E \right] \psi_{lm}(\mathbf{r}) = 0$$



2. Reaction rates

➤ Calculated binding energies → reaction Q-values

[notation] (1) $1+2 \rightarrow 3+4$ reaction: $1(2,3)4$

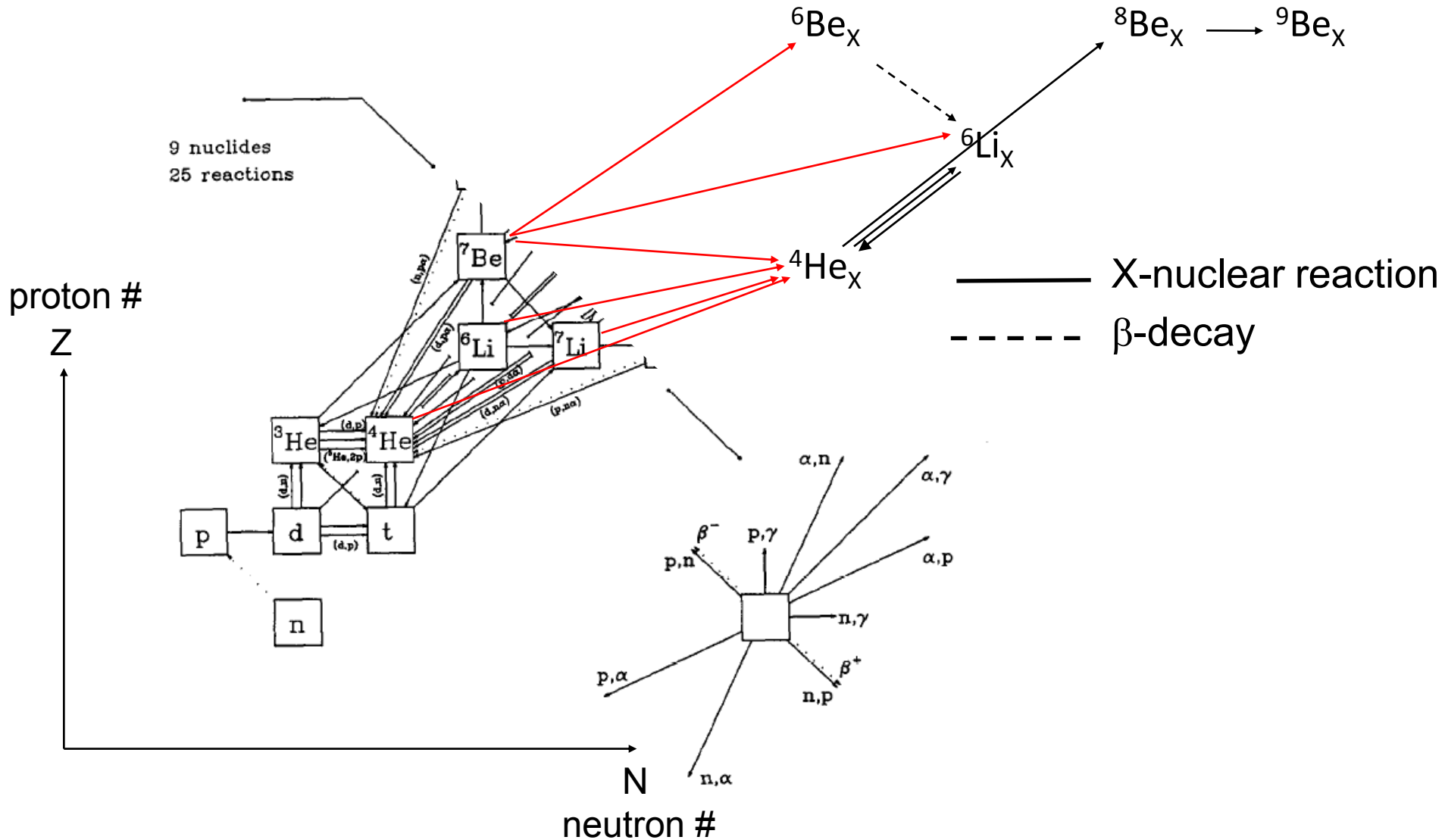
(2) bound state of A and X: A_X

See Kawasaki & MK (2011) for details

✓ **Uncertainties** in estimated binding energies and reaction rates

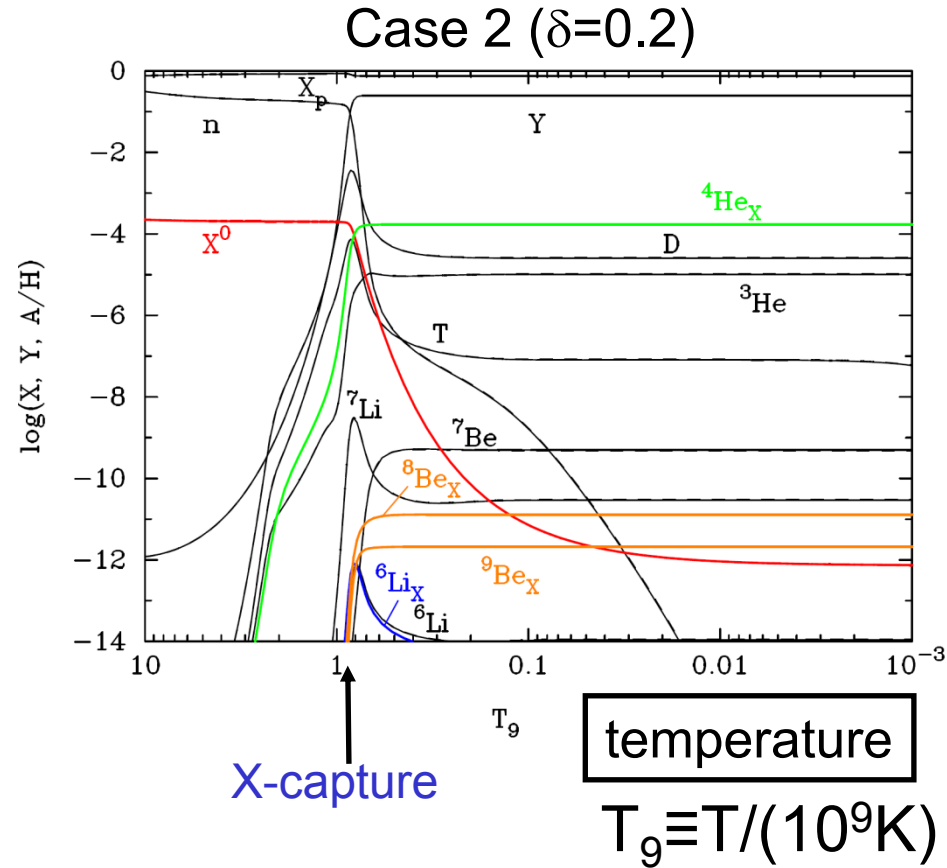
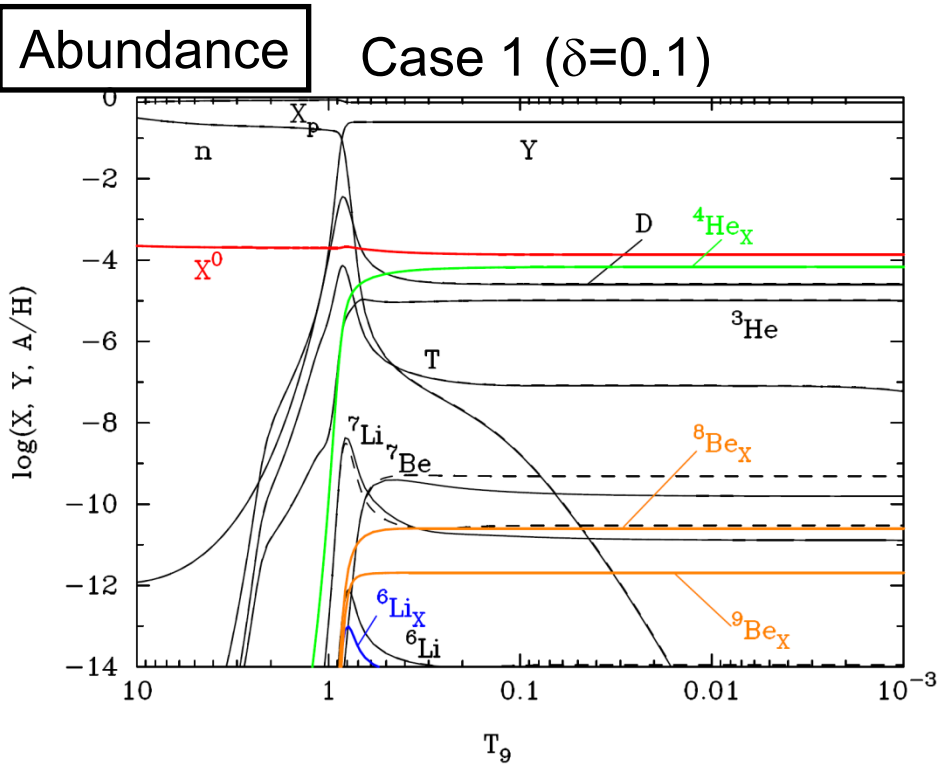
3. Reaction network

➤ Up to ${}^9\text{Be}_x$



Result 1: Nuclear flow

◆ $n_x = 1.7 \times 10^{-4} n_b$, $m_x = 100 \text{ GeV}$, $\tau_x \gg 200 \text{ s}$



✓ $T_9 \gtrsim 1$: Xs are in the free state.

✓ $T_9 \sim 1$: ^4He is produced

→ X is captured by ^4He [1/3 (Case1), ~ 1 (Case2)]

✓ ^7Be & ^7Li react with free X → destroyed



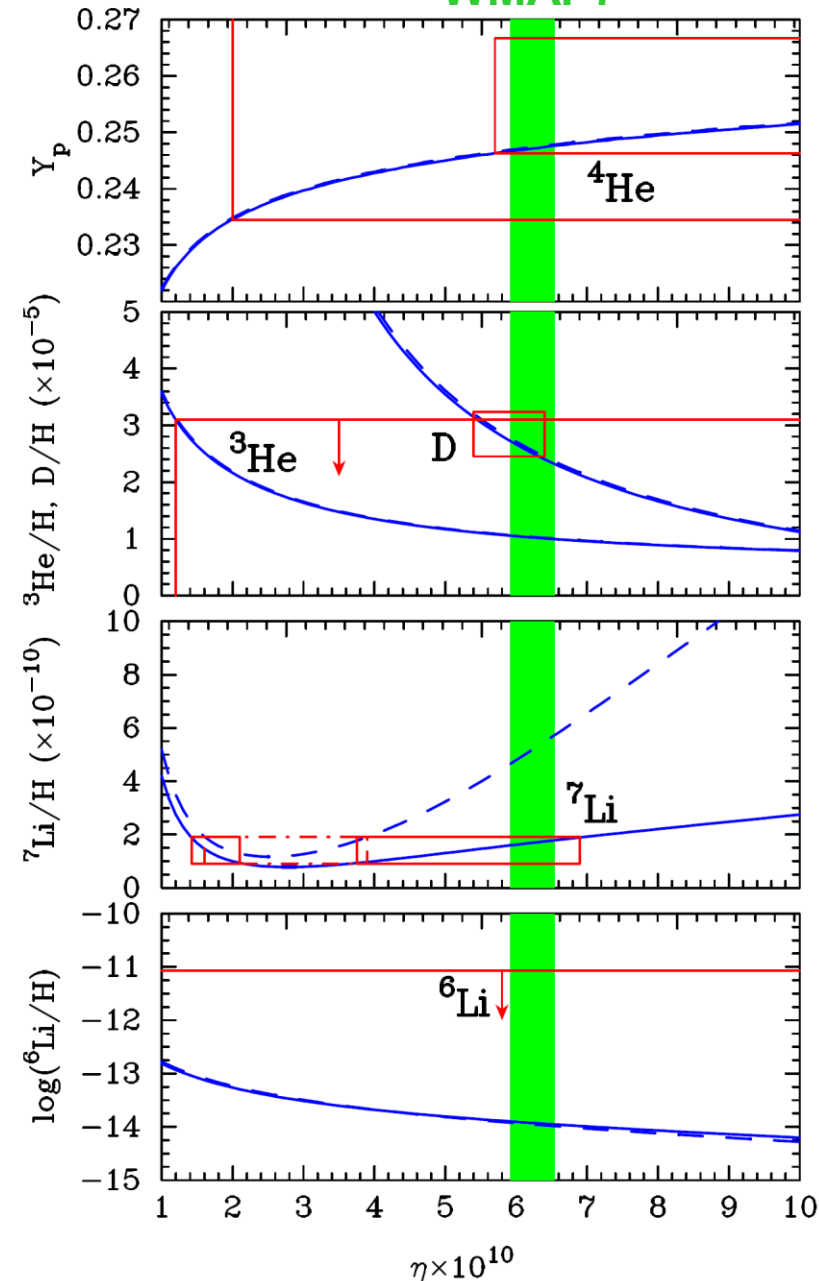
Result 2: reduction of Li abundance

WMAP7

Solid line : BBN including X

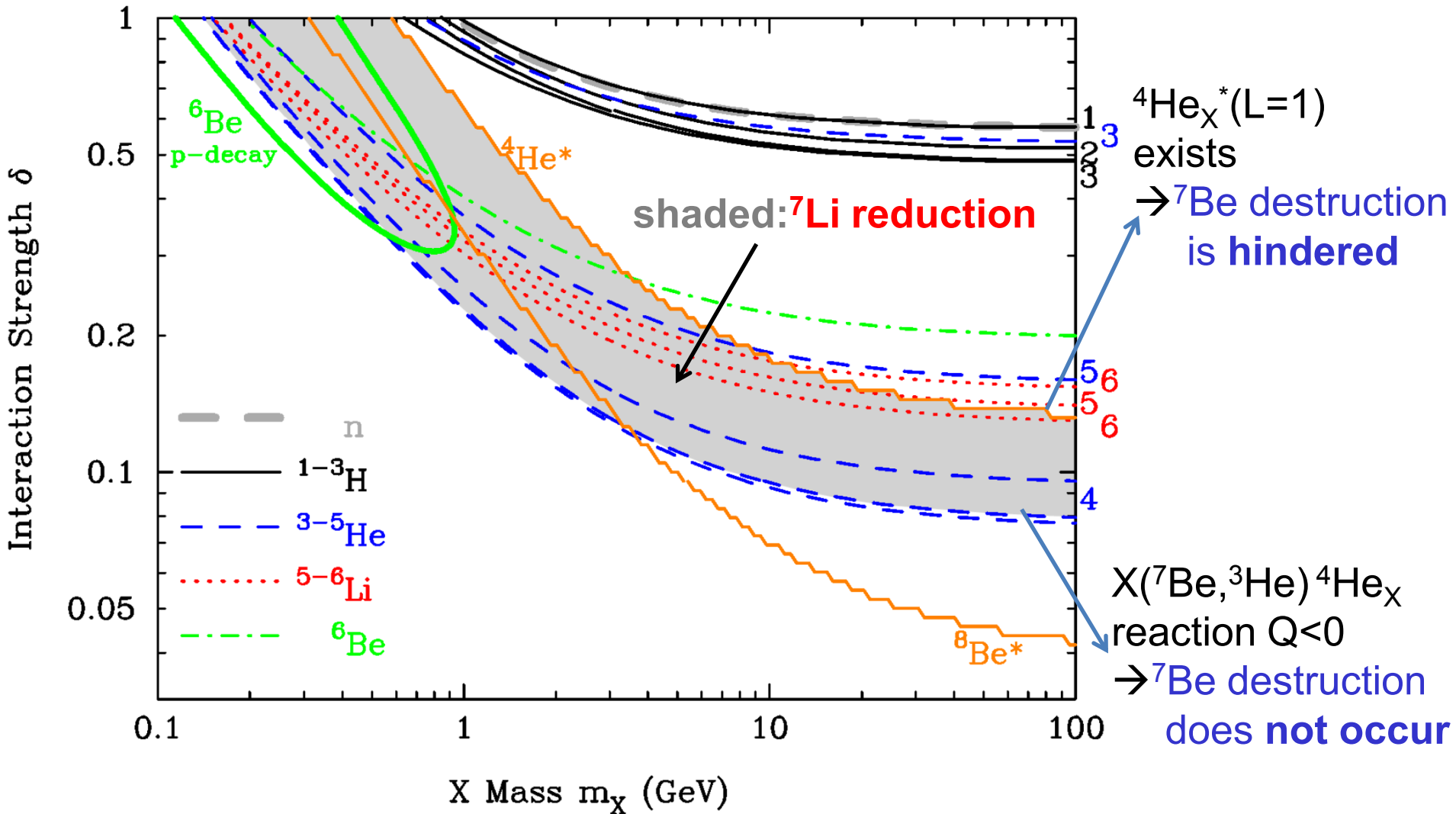
- ◆ $m_x = 100 \text{ GeV}$, $n_x = 1.7 \times 10^{-4} n_b$,
- ◆ $\tau_x \gg 200 \text{ s}$
- ◆ Case 1 ($\delta = 0.1$)

Dashed line : standard BBN



➤ ^7Li reduces
→ new solution to the ^7Li problem

Result 3: Parameter region for Li reduction



Summary

- We study **effects of long-lived strongly interacting massive particle X^0 on BBN**
- Evolutions of elemental abundances are calculated in cases of **XN force weaker than the NN force [$\delta \sim 0.1$]**
→ **sub-SIMP**
- ✓ **^7Be and ^7Li are destroyed via X^0 capture**
 $X(^7\text{Be}, ^3\text{He})^4\text{He}_x$, $X(^7\text{Li}, t)^4\text{He}_x$
- ✓ If there is no excited state of $^4\text{He}_x$ ($L=1$)
→ significant fraction of the X^0 s escape capture by ^4He
- ✓ We show a possibility of **resolving the ^7Li problem**
→ Constraint on parameter region is derived

Processes affecting elemental abundances

	Model	${}^6\text{Li}$ problem solved ?	${}^7\text{Li}$ problem solved ?	Signatures on other nuclides ?
Density fluctuation [$z \sim 10^9$]	Inhomogeneous BBN	no	no	${}^9\text{Be}$ [1]
Existence of particle [$z \sim 10^9$]	sub-SIMP X^0 [2]	?	✓	${}^9\text{Be}$?
	SIMP X^0 [3]	no	no	${}^9\text{Be}$ and/or ${}^{10}\text{B}$
	CHAMP X^- *	✓ [4]	✓ [5,6]	no *
Decay of particle [$z \lesssim 10^9$]	Hadronic decay	✓ [7]	✓ [7]	${}^9\text{Be}$? [8]
	Radiative decay	✓ [9]	no	no
Early stars [$z \sim O(10)$]	Early cosmic ray	✓ [10]	no	${}^9\text{Be}$ and ${}^{10,11}\text{B}$ [11]

* Latest calculation: MK et al. PRD 81, 083521 (2010)

[1] Boyd, Kajino (1989)

[2] Kawasaki, MK (2011)

[3] MK, Kajino, Yoshida, Mathews (2009)

[4] Pospelov (2007)

[5] Bird, Koopmans, Pospelov (2008) [stronger]

[6] MK, Kajino, Boyd, Yoshida, Mathews (2007)
[weaker]

[7] Dimopoulos, Esmailzadeh, Hall, Starkman (1988)

[8] Pospelov, Pradler (2010)

[9] Jedamzik (2000);

Kawasaki, Kohri, Moroi (2001)

[10] Rollinde, Vangioni, Olive (2006)

[11] Rollinde, Maurin, Vangioni, Olive, Inoue (2008);
MK (2008)

Model

1. Binding energies of nuclides to an X

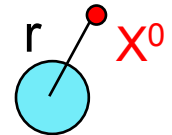
[Assumption]

➤ X (spin 0, **charge 0**, mass m_X)

➤ XN potential is Gaussian

reproducing the binding energy and scattering phase shift of n+p system

Yahiro et al. (1986)



nucleon N

parameter 1

$$v(r) = v_0 \delta \exp \left[- (r/r_0)^2 \right]$$

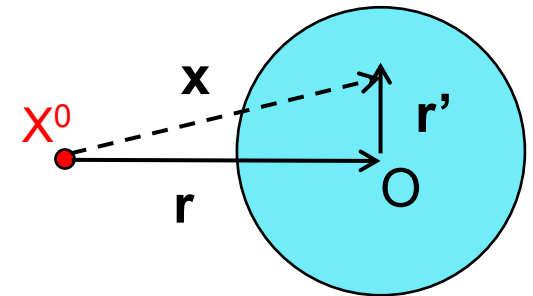
$$v_0 = -72.15 \text{ MeV and } r_0 = 1.484 \text{ fm}$$

parameter 2

➤ XA potential: integration of XN potential multiplied by nucleon densities

$$V(r) = \int v(\mathbf{x}) \rho(\mathbf{r}') d\mathbf{r}'$$

➔ nuclear deformation by interaction of X
is not considered



nucleus A

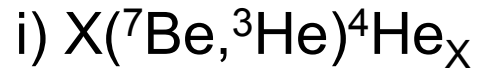
➤ $A \geq 2$ nuclei have Gaussian nucleon densities

2. Reaction rates

- Nonresonant components only
- $m_X = 100 \text{ GeV}$
- Calculated binding energies \rightarrow reaction Q-values
- In this scenario, **${}^7\text{Be}$ is destroyed via a reaction with X**
 - \rightarrow Destruction efficiency depends on
escape fraction of X from capture by ${}^4\text{He}$
 - \rightarrow Two cases for different fractions are considered
(interaction strength: $\delta = 0.1$ & 0.2)

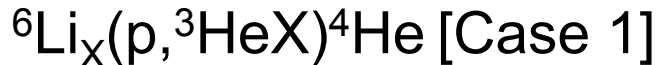
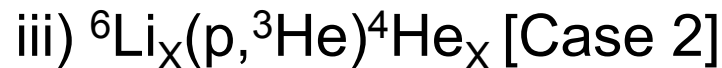
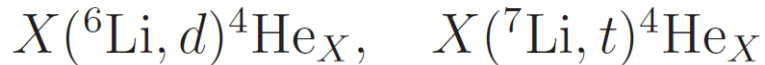
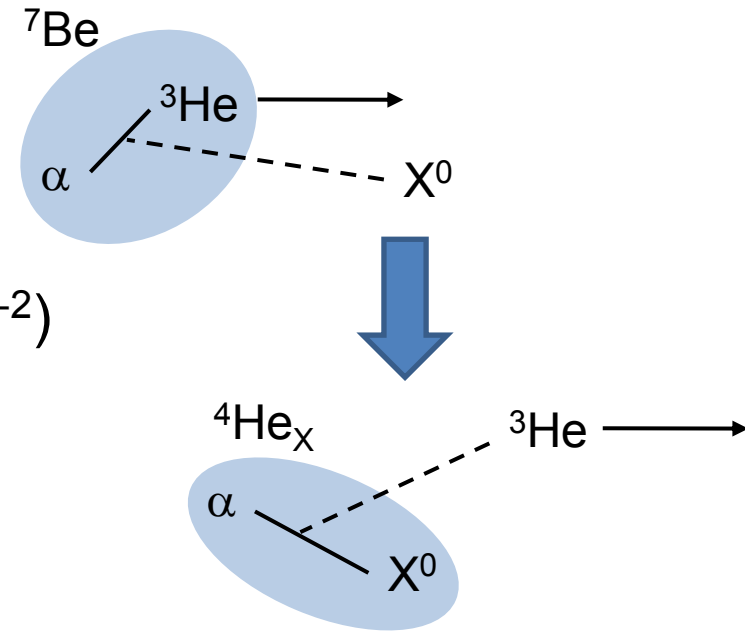
[notation] (1) $1+2 \rightarrow 3+4$ reaction: $1(2,3)4$
(2) bound state of A and X: A_X

2-1. Non-radiative reactions

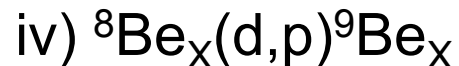


✓ $^6\text{Li}(n, \alpha)^3\text{H}$ rate is adopted
(nonresonant component)

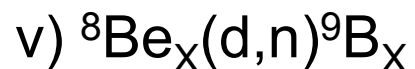
✓ Reduced mass dependence
of cross section is corrected ($\sigma \propto \mu^{-2}$)



✓ $^6\text{Li}(p, ^3\text{He})^4\text{He}$ rate is adopted



✓ $^7\text{Be}(d, p\alpha)^4\text{He}_X$ rate is adopted



✓ $Q < 0 \rightarrow$ neglected

2. Reaction rates

➤ Calculated binding energies → reaction Q-values

[notation] (1) $1+2 \rightarrow 3+4$ reaction: $1(2,3)4$

(2) bound state of A and X: A_X

2-1. Non-radiative reactions

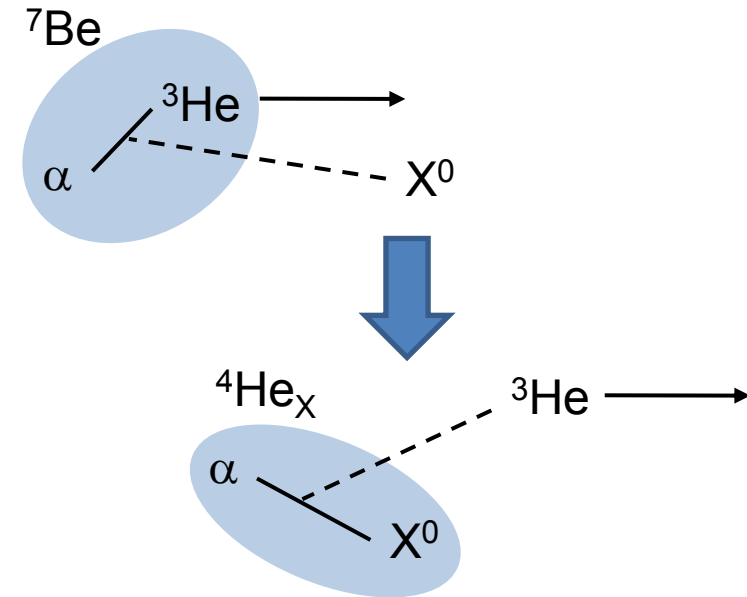
i) $X(^7\text{Be}, ^3\text{He})^4\text{He}_X$

✓ $^6\text{Li}(n, \alpha)^3\text{H}$ rate is adopted

✓ Reduced mass factor
is corrected ($\sigma \propto \mu^{-2}$)

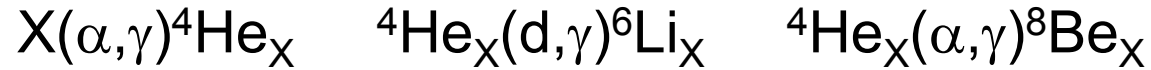
...

See Kawasaki & MK (2011) for details



✓ **Uncertainties** in estimated binding energies and reaction rates

2-2. Radiative reactions



➤ Wave functions of bound & scattering states are calculated with the code RADCAP(Bertulani 2003)

→ cross sections

→ reaction rates are fitted as a function of temperature

2-3. β -decay ($^6\text{Be}_X(,\beta^+\nu_e)^6\text{Li}_X$)

✓ $^6\text{He}(,\beta^-\bar{\nu}_e)^6\text{Li}$ rate is used after correction for the Q-value

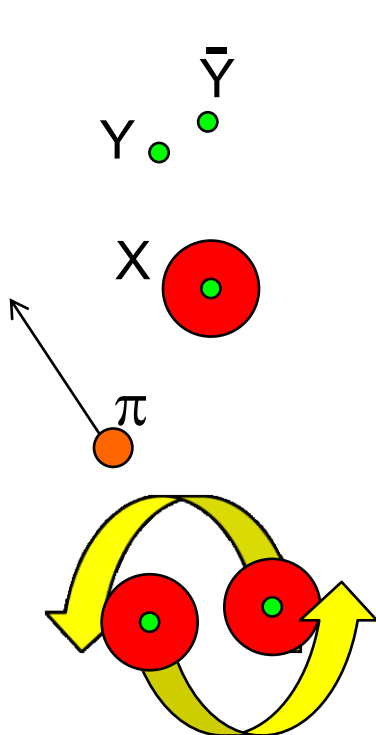
✓ **Large uncertainties in estimation of binding energies and reaction rates**

→ **realistic calculation with a quantum many-body models are needed**

Long-lived Heavy Colored Particles

➤ Existence of long-lived heavy ($m \gg \text{GeV}$) colored particles are suggested in particle models beyond the standard model e.g., Split SUSY (Arkani-Hamed & Dimopoulos 2005), ...

➤ History in the early universe: (J. Kang et al. JHEP 9, 86, 2008)



✓ In the early universe, hypothetical colored particles (Y) annihilate \rightarrow relic abundance $n_Y/n_b \sim 10^{-4}$

✓ $T < T_c \sim 180 \text{ MeV} \rightarrow$ heavy partons get confined in hadrons (X)

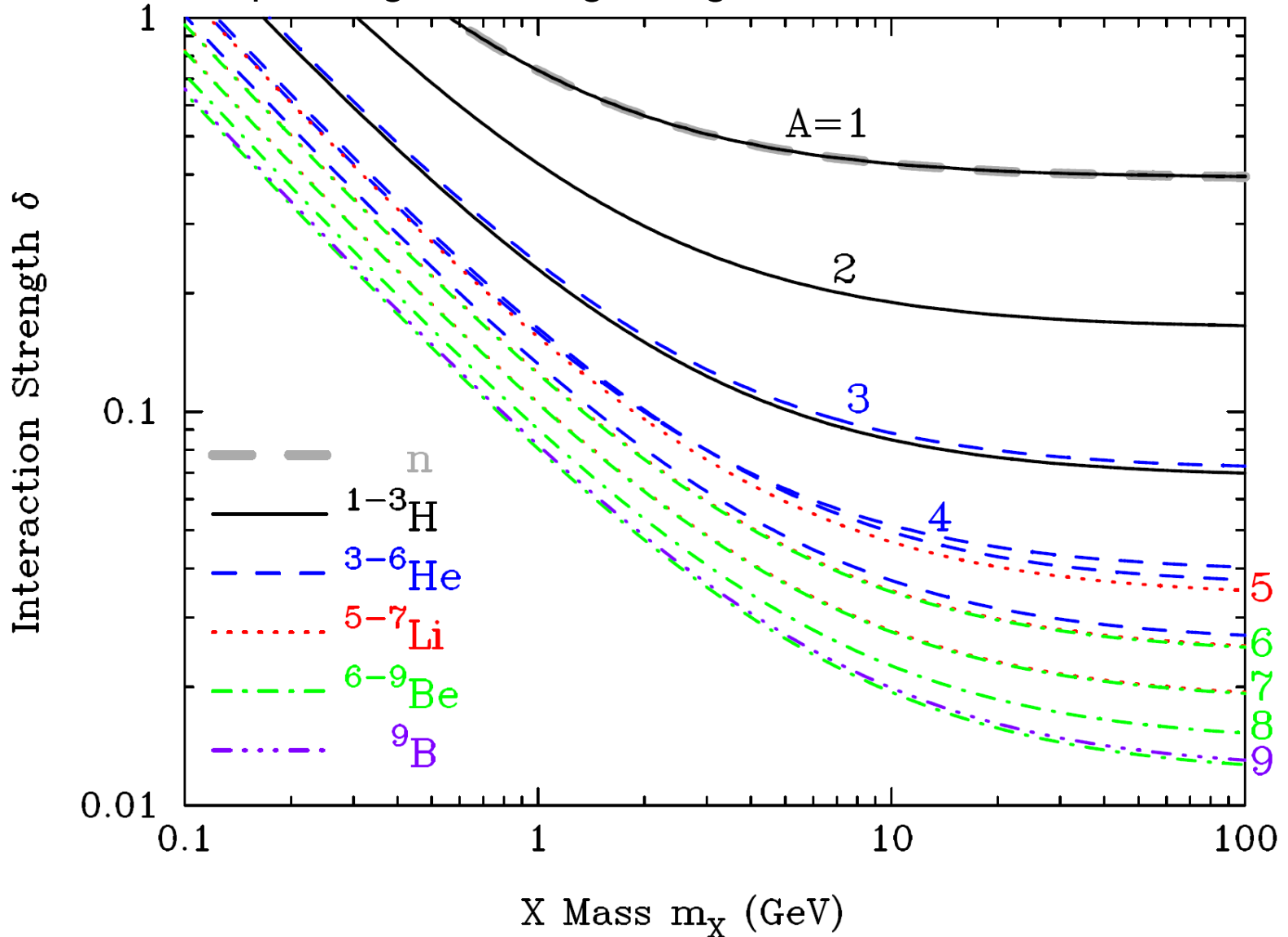
✓ $X+X$ form the bound state \rightarrow decay into lower energy states \rightarrow annihilate \rightarrow final abundance

$$\frac{n_X}{n_b} \approx 10^{-8} \left(\frac{R}{\text{GeV}^{-1}} \right)^{-2} \left(\frac{T_B}{180 \text{ MeV}} \right)^{-3/2} \left(\frac{m}{\text{TeV}} \right)^{1/2}$$

Model

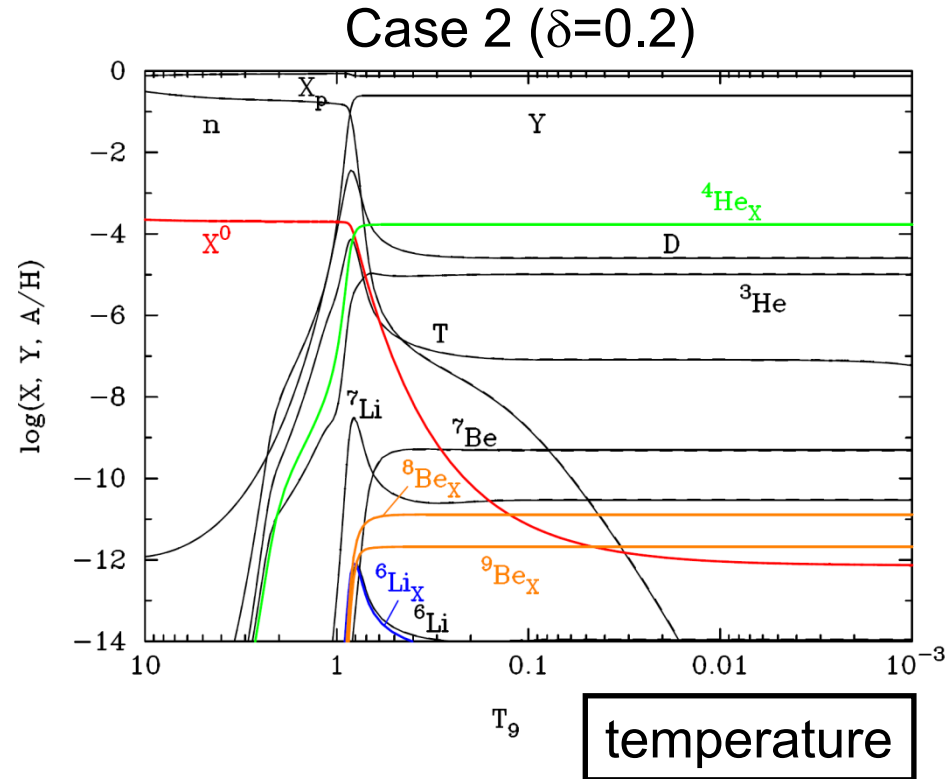
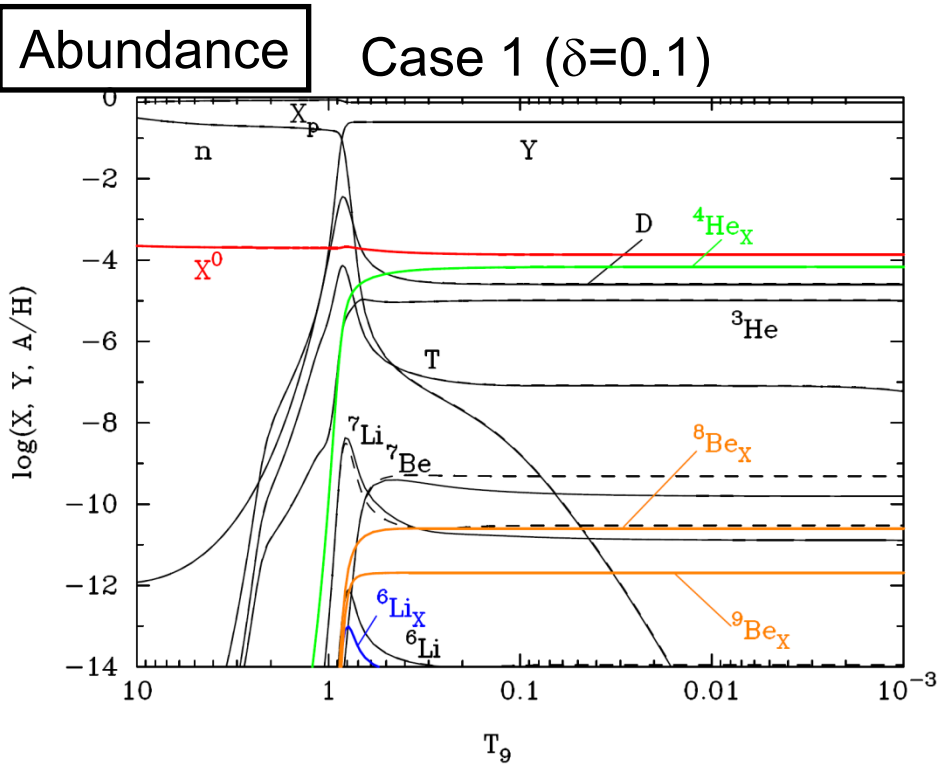
1. Binding energies of nuclides to an X

Contours corresponding to binding energies=0.1MeV



Result 1: Nuclear flow

◆ $n_x = 1.7 \times 10^{-4} n_b$, $m_x = 100 \text{ GeV}$, $\tau_x \gg 200 \text{ s}$



temperature

$T_9 \equiv T / (10^9 \text{ K})$

✓ **no excited ${}^4\text{He}_x(L=1)^*$**
 → E1 transition
 (${}^4\text{He}+X$) p-wave → ${}^4\text{He}_x$ ground state is dominant

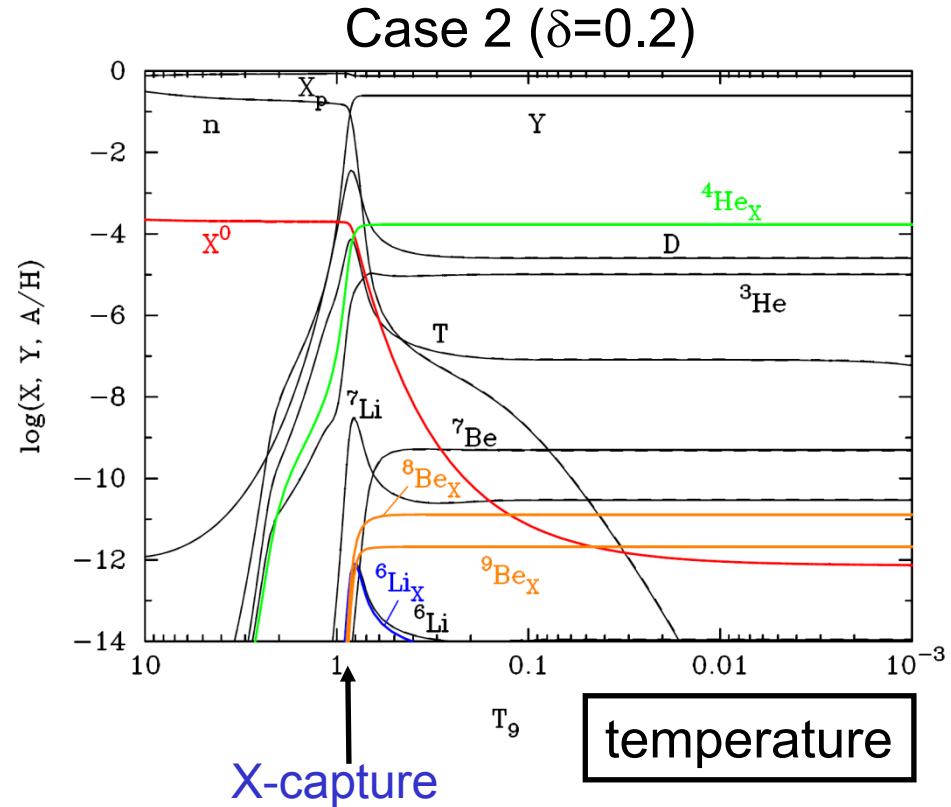
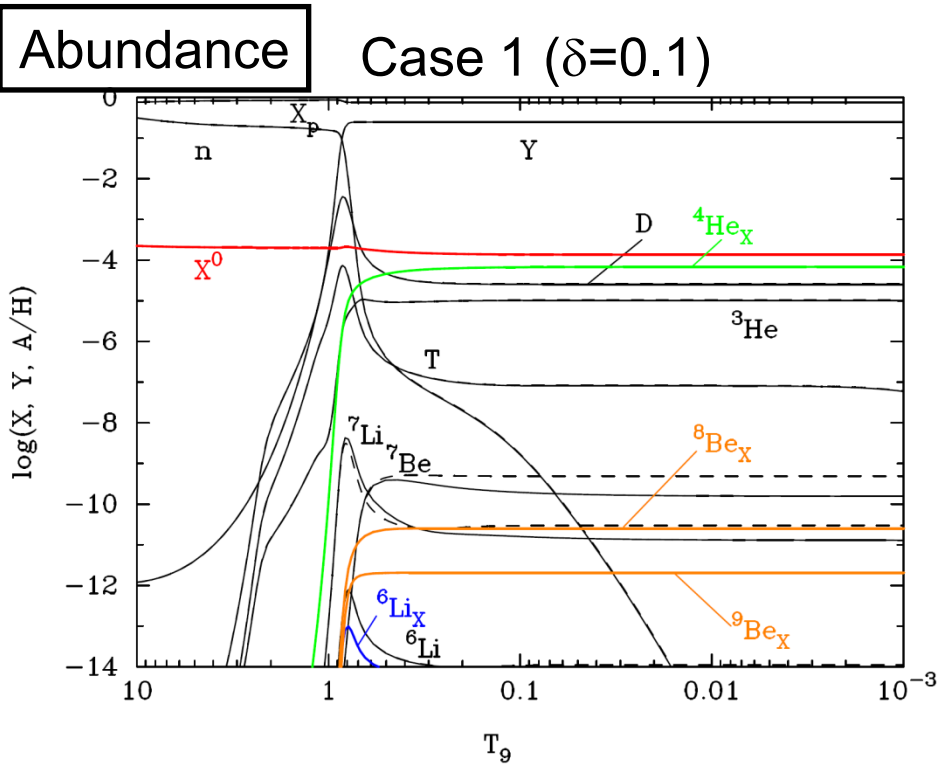
✓ **Excited ${}^4\text{He}_x(L=1)^*$**
 → E1 transition
 (${}^4\text{He}+X$) s-wave → ${}^4\text{He}_x$ excited state is dominant

→ **X-capture by ${}^4\text{He}$ is weak**

→ **X capture by ${}^4\text{He}$ is strong**

Result 1: Nuclear flow-1

◆ $n_x = 1.7 \times 10^{-4} n_b$, $\tau_x \gg 200s$



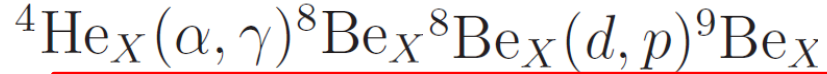
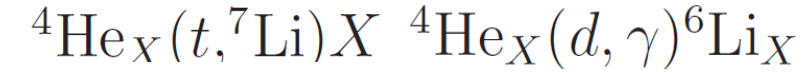
✓ $T_9 \geq 1$: Xs are in the free state.

✓ $T_9 \sim 1$: ^4He is produced

→ X is captured by ^4He [1/3 (Case1), large portion (Case2)]

✓ nuclear reaction of $^4\text{He}_x$

→ heavy X-nuclei are produced



✓ ^7Be & ^7Li react with free X → destroyed

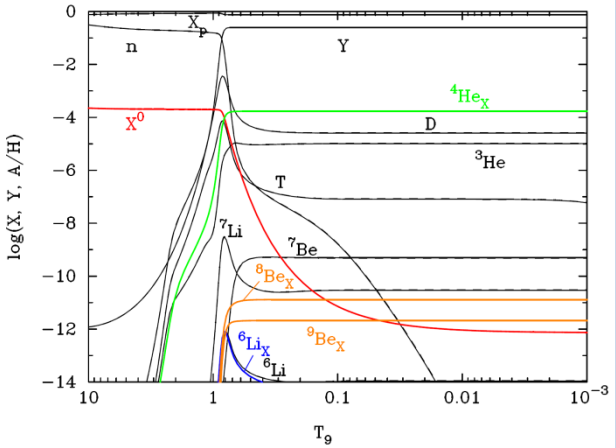
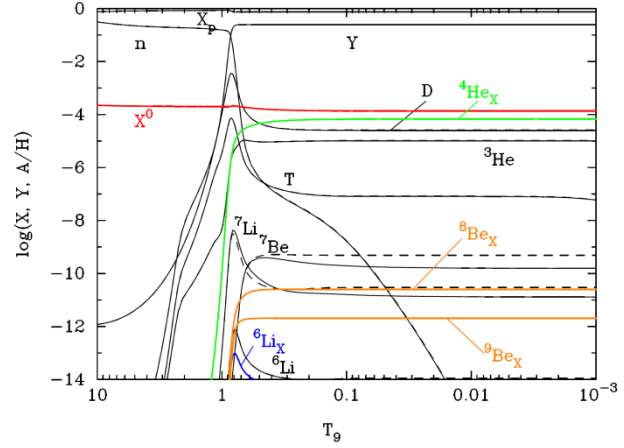


Result 1: Nuclear flow-2

Abundance

Case 1 ($\delta=0.1$)

Case 2 ($\delta=0.2$)



$T_9 = T / (10^9 \text{ K})$
temperature

✓ no excited ${}^4\text{He}_x(L=1)^*$

→ E1 transition from $({}^4\text{He}+X)$ p-wave to ${}^4\text{He}_x$ ground state is dominant

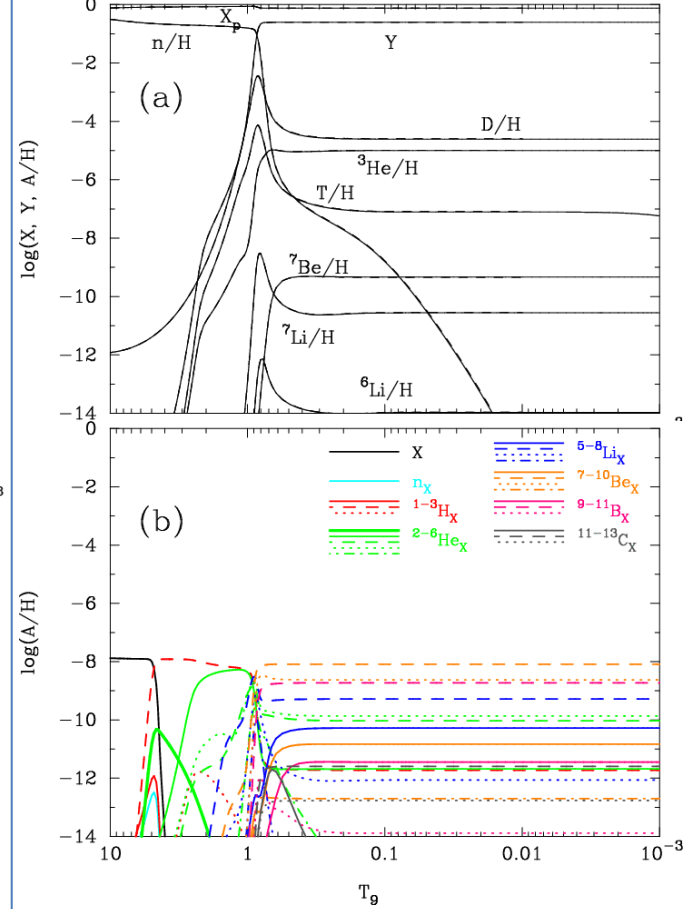
→ X-capture by ${}^4\text{He}$ is weak

✓ Excited ${}^4\text{He}_x(L=1)^*$

→ E1 transition from $({}^4\text{He}+X)$ s-wave to ${}^4\text{He}_x$ excited state is dominant

→ X capture by ${}^4\text{He}$ is strong

MK et al. 2009 ($\delta \sim 1$)



Bound states of X & ${}^5\text{A}$ exist

→ heavy X-nuclei efficiently form

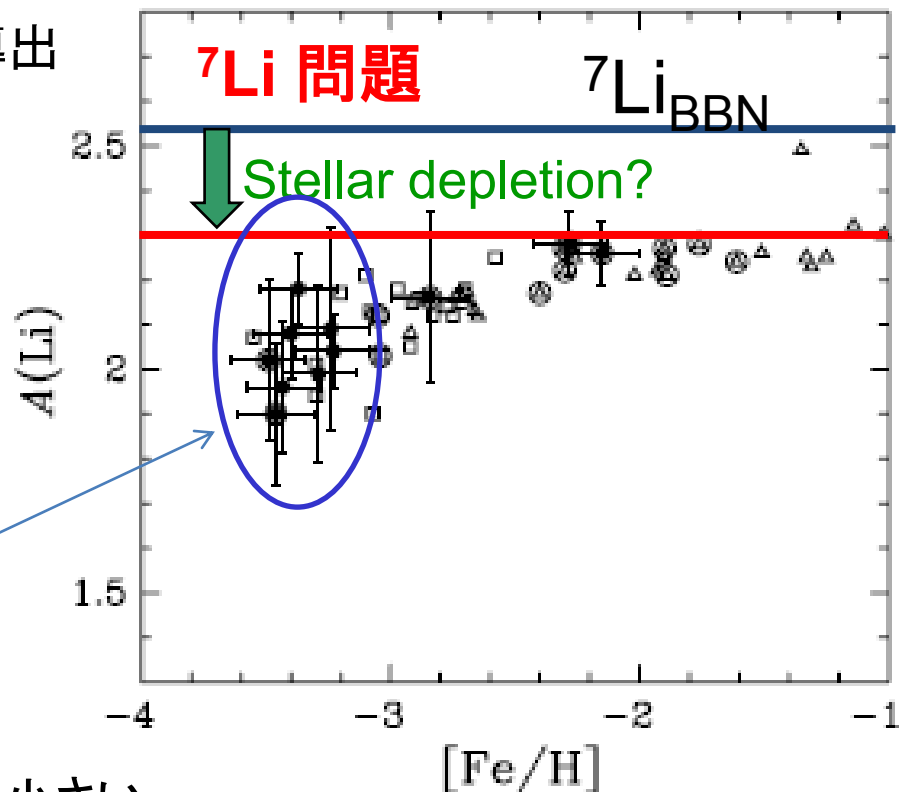
${}^7\text{Li}$ 問題

[Li組成分析]

- Li I 6708 Å lineのスペクトルfitでLi組成導出
- ドップラー補正はFe I線の位置から決定
- ${}^6\text{Li}$ の寄与は無視
- fitパラメーターはLi組成(と波長シフト)

$\log({}^7\text{Li}/\text{H})+12$

Aoki et al. (2009)



$[\text{Fe}/\text{H}] < -3$ で平均的にLi組成が低い

$$A(\text{Li}) = 2.03 \pm 0.09$$

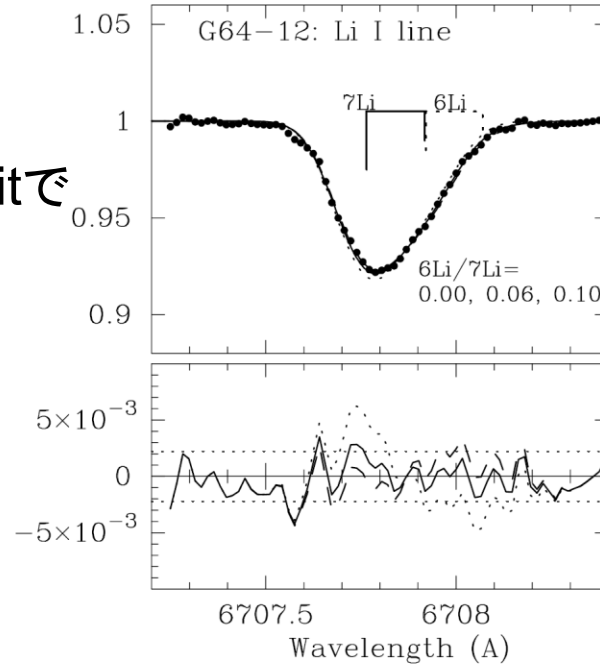
- 標準偏差が測定誤差(0.07-0.23 dex)より小さい
→ Li組成の分散を未検出
- high metallicity star $A(\text{Li}) = 2.27$ より0.24dex 小さい
(Asplund et al. (2006)の $[\text{Fe}/\text{H}] = -2.5$ -- -2.0 の星
 $A(\text{Li}) = 2.23$ より0.27dex 小さい)
→ この差は大きい

$$[\text{Fe}/\text{H}] = \log\left\{\frac{(\text{Fe}/\text{H})}{(\text{Fe}/\text{H})_{\odot}}\right\}$$

^6Li 問題

[Li組成分析]

- Li I 6708 Å lineのスペクトルfitでLi同位体比導出



Inoue et al. (2005)

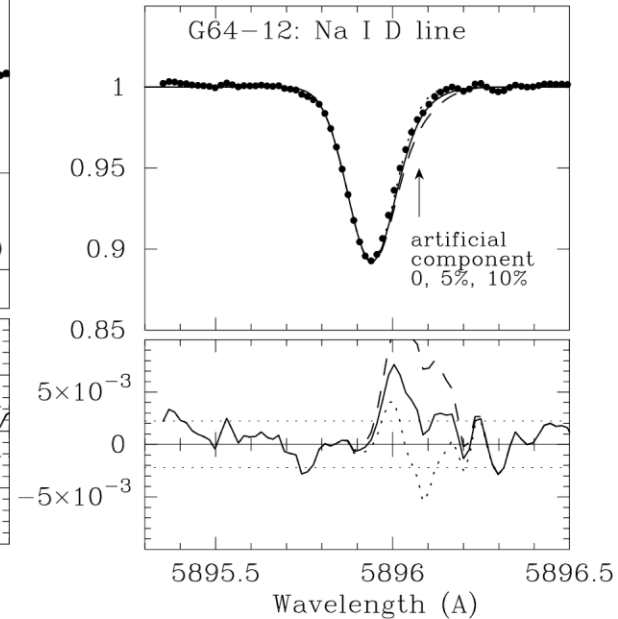


Figure 1. Left: Observed spectra (filled circles) of the Li I line in G 64-12 compared with synthetic spectra (curves) assuming $^6\text{Li}/^7\text{Li}=0$ (dotted), 0.06 (solid) and 0.1 (dashed). The bottom panel shows the difference between synthetic and observed spectra, with the horizontal dotted lines marking $\pm 1\sigma$ in the observed flux. Right: Same as left, except for the Na I D line. The synthetic spectra include artificial line asymmetries that mimic ^6Li in the Li I line at levels of 0 (dotted), 5 (solid) and 10 (dashed) %.

- 9つのMPHSsで ^6Li の検出 (Asplund et al. 2006)

- しかし、大気の3D効果を取り入れる必要がある (Cayrel et al. 2007)

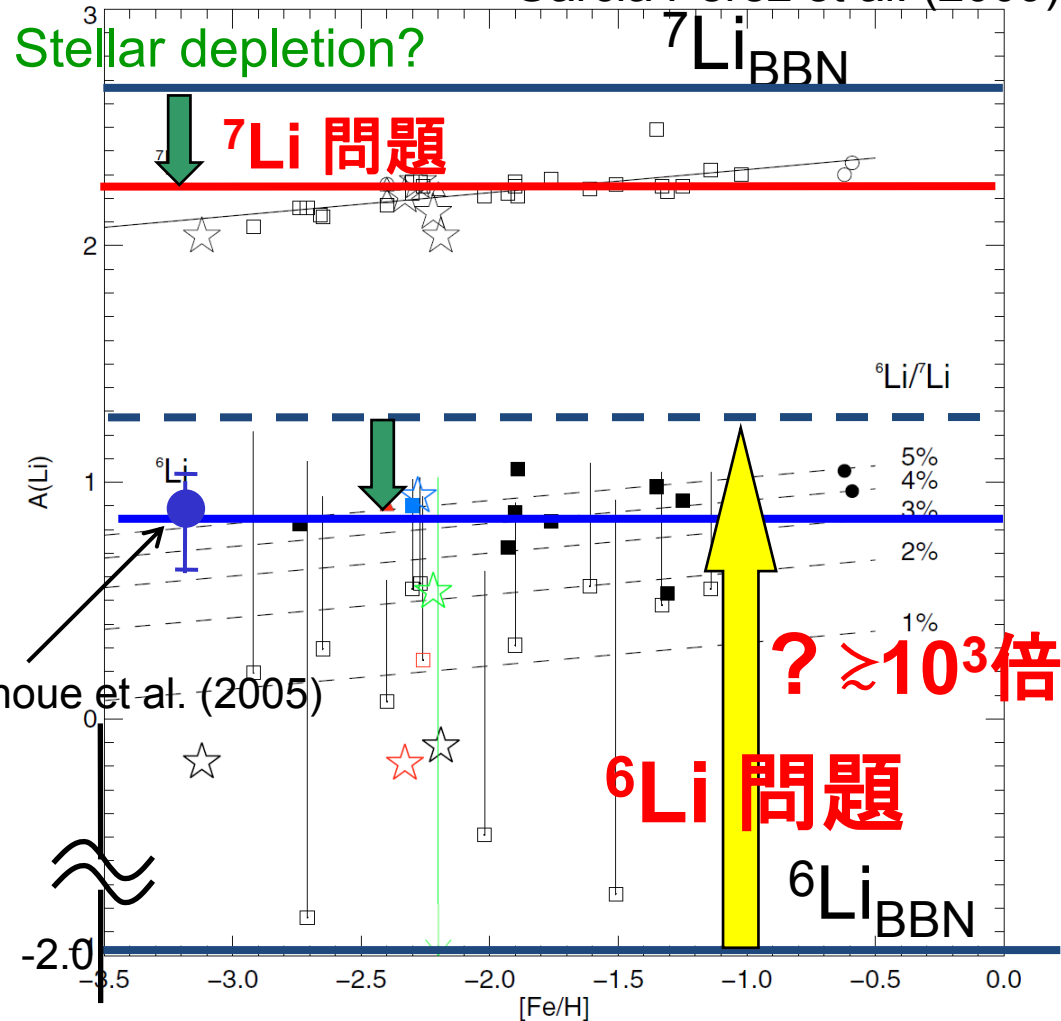
- 3D効果を考慮すると、検出とされた星の数が9→4に減る

- しかし、確かに ^6Li がある星も存在 (Garcia Perez et al. 2009)

天体物理的過程

$$\log(^{6,7}\text{Li}/\text{H})+12$$

Garcia Perez et al. (2009)



[⁶Li合成]

➤ 初期の宇宙線 α と背景 α の $\alpha+\alpha$
(Montmerle 1977)

✓ 銀河形成前の
超新星宇宙線

(Rollinde et al. 2005)

✓ 銀河形成中の構造形成
ショック宇宙線
(Suzuki & Inoue 2002)

➤ 超新星ejecta α +星周 α の $\alpha+\alpha$
(Nakamura et al. 2006)

➤ MPHSフレア³Heの⁴He(³He,p)
(Tatischeff & Thibaud 2007)

[^{6,7}Li減少]

➤ 銀河初期の超高効率化学進化(星内燃焼→ejection)(Piau et al. 2006)

➤ MPHS表面での原子拡散・乱流混合 (Richard et al. 2005)

始原組成の観測的制限1

- D : QSOの方向にある吸収系 (Pettini et al. 2008)

$$\log D/H = -4.55 \pm 0.03 (2\sigma)$$

- ^3He : 銀河系のHII region (Bania et al. 2002)

$$^3\text{He}/H = (1.9 \pm 0.6) \times 10^{-5} (2\sigma, \text{上限})$$

- ^4He : metal-poor outer galaxiesのHII region

(Izotov & Thuan, 2010) $Y = 0.2565 \pm 0.0051 (2\sigma)$

(Aver et al. 2010) $Y = 0.2561 \pm 0.0108 (2\sigma)$

- ^6Li : Metal-Poor Halo Star (Asplund et al. 2006)

[隕石の ^6Li 組成(Lodders 2003)を超えない]

$$^6\text{Li}/H = (7.1 \pm 0.7) \times 10^{-12} (2\sigma, \text{上限})$$

- ^7Li : Metal-Poor Halo Star (Ryan et al. 2000)

$$^7\text{Li}/H = (1.23^{+0.68}_{-0.32}) \times 10^{-10} (95\% \text{ CL})$$

始原組成の観測的制限2

● ${}^9\text{Be}$: Metal-Poor Halo Star (Ito et al. 2009)

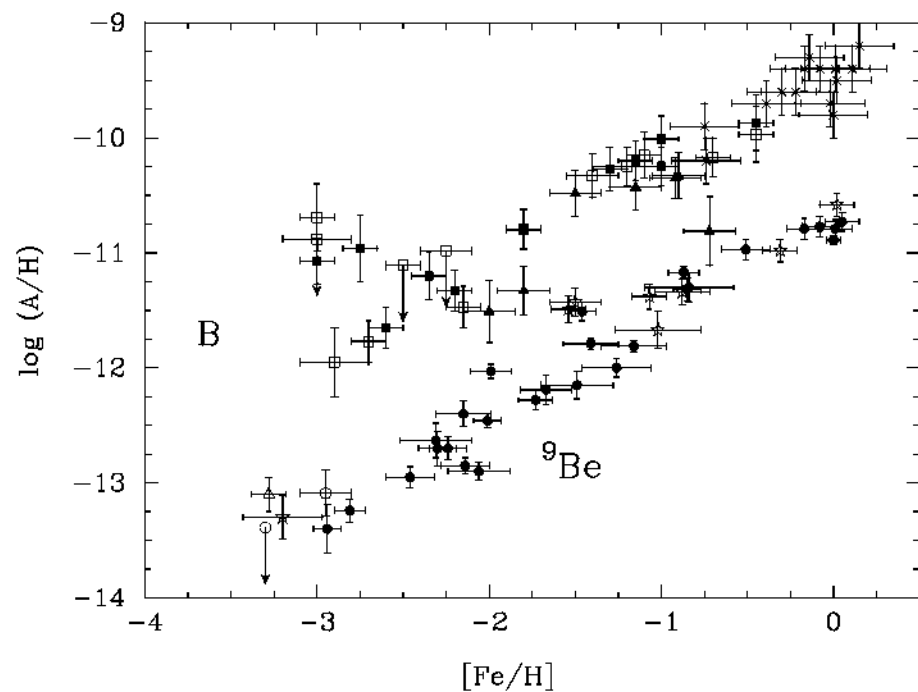
$${}^9\text{Be}/\text{H} < 10^{-14}$$

● B : Metal-Poor Halo Star (Duncan et al. 1997, Garcia Lopez et al. 1998)

$$\text{B}/\text{H} < 10^{-12}$$

● C : Metal-Poor Halo Star (Suda et al. compilation arXiv:0806.3697)

$$\text{C}/\text{H} < 10^{-8}$$



Suda et al. (2008)

