

# Lower limits on gamma lines from WIMP annihilations

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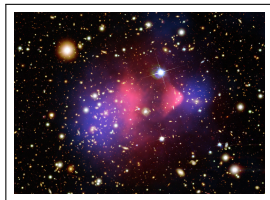
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May 9, 2011

K. Abazajian, Z. Chacko, C. Kilic  
to appear

## Weakly Interacting Massive Particles

There are many gravitational observations which point towards existence of dark matter.

- ▶ Galactic rotation curves
- ▶ Cosmic Microwave Background
- ▶ Gravitational lensing observations



The WIMP miracle

- ▶ The relic abundance of a particle with a weak-scale mass and weak-scale interactions is consistent with observations.

$$\langle \sigma_{Av} \rangle \sim \frac{g^4}{4\pi(1 \text{ TeV})^2} \sim 10^{-26} \text{ cm}^3/\text{s}$$

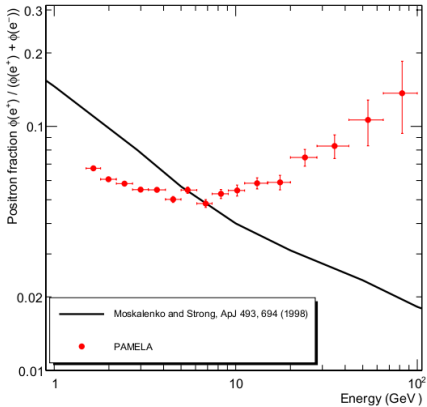
## Indirect detection of WIMPs

Dark matter particles in galactic haloes can find each other and annihilate into SM states.

- ▶ Tests robust prediction of the WIMP hypothesis,  
 $\langle \sigma_{Av} \rangle \sim 3 \times 10^{-26} \text{ cm}^3/\text{s}$
- ▶ Complementary to direct detection
  - ▶ coupling to other standard model states
  - ▶ different set of astrophysical parameters
- ▶ Background estimation is hard.

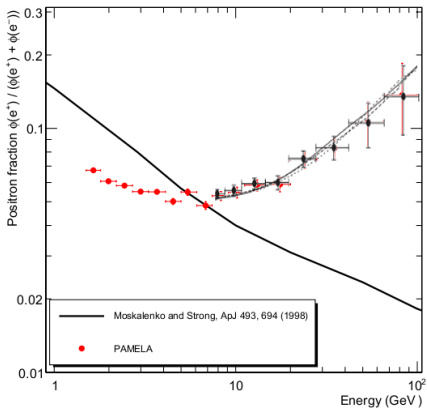
# Example: Positron excess in PAMELA

Detection of dark matter?



## Example: Positron excess in PAMELA

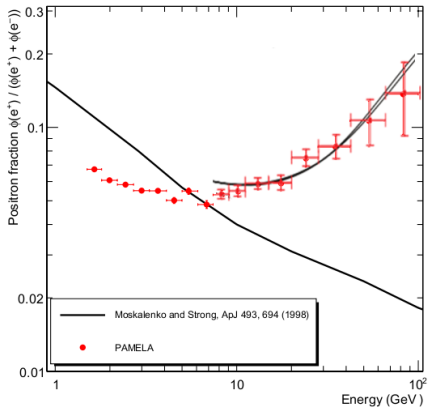
Detection of dark matter?



► Annihilation through a light boson

# Example: Positron excess in PAMELA

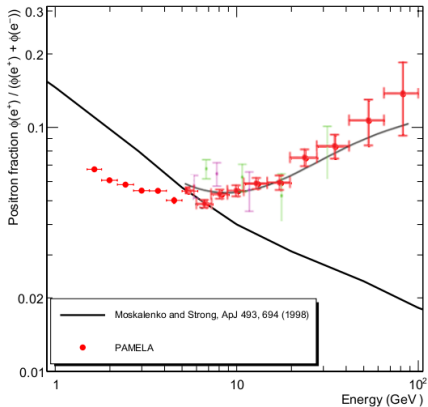
Detection of dark matter?



► Kaluza-Klein dark matter

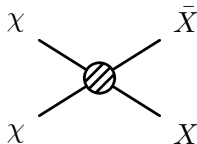
# Example: Positron excess in PAMELA

Detection of dark matter?



► Geminga pulsar

## Continuum photon spectrum

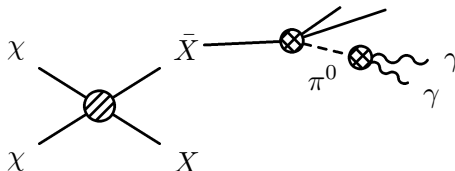
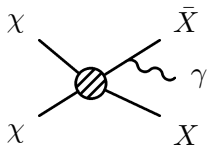


- ▶ Tree level annihilation to two-body SM final states arising from a renormalizable theory
- ▶ Cannot annihilate to photons directly



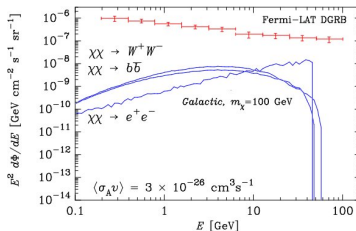
# Continuum photon spectrum

Bremsstrahlung/Hadronization

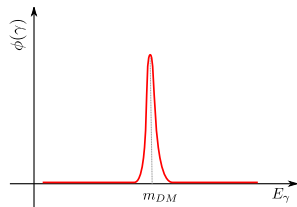
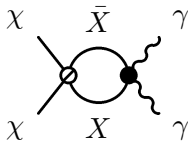


- ▶ Independent of the form of the coupling
- ▶ Depends on
  - ▶ dark matter mass
  - ▶ SM final state

Model independent!



## Photon lines

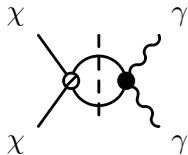


- ▶ No known astrophysical background
- ▶ Model dependent

$$|\mathcal{M}|^2 \sim \left| \begin{array}{c} \text{Diagram 1} \\ + \\ \text{Diagram 2} \\ + \\ \text{Diagram 3} \\ + \\ \dots \end{array} \right|^2$$

- ▶ Depends on detailed form of the couplings
- ▶ New physics often dominates

## Photon lines



Use unitarity relations to put lower bounds on the line strength relative to continuum strength

- ▶ Continuum  $\gamma$ -rays produced at tree-level from annihilation to SM states
- ▶ Photon lines arise at one loop
- ▶ Limit to  $\gamma\gamma$  mode
- ▶ Assume single SM channel dominates the imaginary part

# Unitarity

- ▶ We use the  $|J, M; L, S\rangle$  basis

$J$ : total angular momentum	} conserved
$M$ : $z$ -component of $J$	
$L$ : orbital angular momentum	} not conserved
$S$ : total spin	

- ▶ From unitarity, assuming time-reversal (CP) invariance,

$$2 \operatorname{Im}\langle f|T|i\rangle = \sum_X \langle f|T^\dagger|X\rangle \langle X|T|i\rangle$$

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 contribution to  
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$$\sigma_{IM}(\chi\chi \rightarrow \gamma\gamma) = \rho \sigma(\chi\chi \rightarrow X\bar{X}) \sigma(X\bar{X} \rightarrow \gamma\gamma)$$

↑	↑	↑	↑
Imaginary part contribution to $\sigma(\chi\chi \rightarrow \gamma\gamma)$	phase space factor	DM cross section to SM	SM cross section

We have assumed a unique state  $X$  with unique  $L, S$  values.

## Model-independent limits

$$\frac{\phi_{\text{line}}}{\phi_{\text{continuum}}} \propto \frac{\sigma(\chi\chi \rightarrow \gamma\gamma)}{\sigma(\chi\chi \rightarrow X\bar{X})} \geq \frac{\sigma\left(\begin{array}{c} \chi \\ \diagdown \\ \text{---} \\ \diagup \\ \chi \end{array} \begin{array}{c} \gamma \\ \diagup \\ \text{---} \\ \diagdown \\ \gamma \end{array}\right)}{\sigma\left(\begin{array}{c} \chi \\ \diagdown \\ \text{---} \\ \diagup \\ \chi \end{array} \begin{array}{c} \bar{X} \\ \diagup \\ \text{---} \\ \diagdown \\ X \end{array}\right)} = \rho \sigma\left(\begin{array}{c} \bar{X} \\ \diagdown \\ \text{---} \\ \diagup \\ X \end{array} \begin{array}{c} \gamma \\ \diagup \\ \text{---} \\ \diagdown \\ \gamma \end{array}\right)$$

When can we find unique intermediate states?

## Example 1: Majorana Fermion

- ▶ Dark matter particles are non-relativistic  $\Rightarrow$  Initial  $L = 0$
- ▶ Wavefunction is anti-symmetric  $\Rightarrow$  Initial  $S = 0$

Dark Matter	Initial spin	Annihilation		Bound
		channel	mode	
Majorana Fermion	$J = 0 \quad -$	$WW$	$L = 0, S = 0 \quad +$	
			$L = 1, S = 1 \quad -$	
$L = 2, S = 2 \quad +$				
		$f\bar{f}$	$L = 0, S = 0 \quad -$	
			$L = 1, S = 1 \quad +$	

Majorana fermion dark matter will always give a photon line with at least a certain strength.

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Majorana fermion dark matter will always give a photon line with at least a certain strength.

## Example 2: Dirac Fermion

Dark Matter	Initial spin	Annihilation		Bound
		channel	mode	
Dirac Fermion	$J = 0 \quad -$	$WW$	$L = 0, S = 0 \quad +$	
			$L = 1, S = 1 \quad -$	
	$L = 2, S = 2 \quad +$			
	$ff$	$L = 0, S = 0 \quad -$		
$L = 1, S = 1 \quad +$				
	$J = 1 \quad +$			

## Example 2: Dirac Fermion

Dark Matter	Initial spin	Annihilation		Bound
		channel	mode	
Dirac Fermion	$J = 0 \quad -$	$WW$	$L = 0, S = 0 \quad +$	✓
			$L = 1, S = 1 \quad -$	
	$L = 2, S = 2 \quad +$			
Dirac Fermion	$J = 0 \quad -$	$f\bar{f}$	$L = 0, S = 0 \quad -$	✓
			$L = 1, S = 1 \quad +$	
	$J = 1 \quad +$	Forbidden by Landau-Yang theorem		

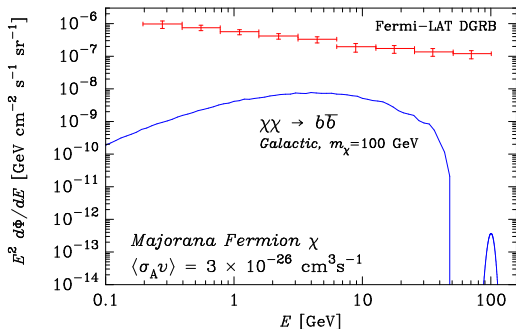


## Other cases

Dark Matter	Initial spin	Annihilation		Bound
		channel	mode	
Scalar	$J = 0$	$WW$	$L = 0, S = 0$ $L = 2, S = 2$	In non-relativistic, ultra-relativistic limits.
		$f\bar{f}$	$L = 1, S = 1$	✓
Majorana Fermion	$J = 0$	$WW$	$L = 1, S = 1$	✓
		$f\bar{f}$	$L = 0, S = 0$	✓
Dirac Fermion	$J = 0$	$WW$	$L = 1, S = 1$	✓
		$f\bar{f}$	$L = 0, S = 0$	✓
	$J = 1$	Forbidden		
Real Vector Boson	$J = 0$	$WW$	$L = 0, S = 0$ $L = 2, S = 2$	In non-relativistic, ultra-relativistic limits.
		$f\bar{f}$	$L = 0, S = 0$	✓
	$J = 2$	$WW$	$L = 2, S = 0$ $L = \{0, 1, 2, 3, 4\}, S = 2$	In non-relativistic limit.
		$f\bar{f}$	$L = \{1, 2, 3\}, S = 1$	In non-relativistic, ultra-relativistic limits.

# Phenomenology of Majorana Fermion

$b\bar{b}$  spectrum

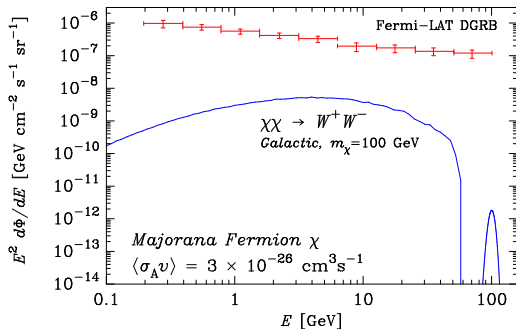


$$\frac{\sigma(\chi\chi \rightarrow \gamma\gamma)}{\sigma(\chi\chi \rightarrow b\bar{b})} \geq \frac{3e^4 m_b^2}{32\pi^2 m_\chi^2} \frac{1}{\beta} [\tanh^{-1} \beta]^2$$

$$\beta = \sqrt{1 - \frac{m_b^2}{m_\chi^2}}$$

# Phenomenology of Majorana Fermion

## $WW$ spectrum

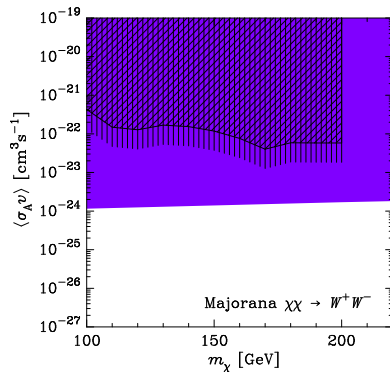
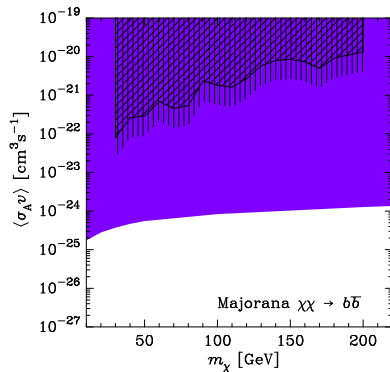


$$\frac{\sigma(\chi\chi \rightarrow \gamma\gamma)}{\sigma(\chi\chi \rightarrow WW)} \geq \frac{e^4}{8\pi^2} \beta^{3/2} [\tanh^{-1} \beta]^2$$

$$\beta = \sqrt{1 - \frac{m_W^2}{m_\chi^2}}$$

# Phenomenology of Majorana Fermion

## Limits



## Comments

Bounds obtained from lower limit of lines are currently weaker than the diffuse isotropic spectrum

- ▶ Bounds are largely statistics limited and will improve with more data.
- ▶ Future experiments with better resolution will be more sensitive to lines.
- ▶ If these bounds are comparable, the overall bound from lines in a given model are expected to dominate.