

Testing Bulk Matter RS Model
through
Rare Decays of Smuon and Scharm

Toshifumi Yamada

KEK

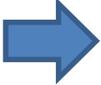
arXiv:1105.XXXX [hep-ph]

Introduction

- Origin of Yukawa coupling hierarchy in SM ?
- “5D RS spacetime with matter fields in the bulk”

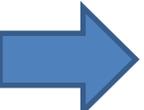
Weaknesses :

i. Severe constraints on KK scale :

e.g. $K^0-\bar{K}^0$ mixing  1st KK gluon mass > 21 TeV

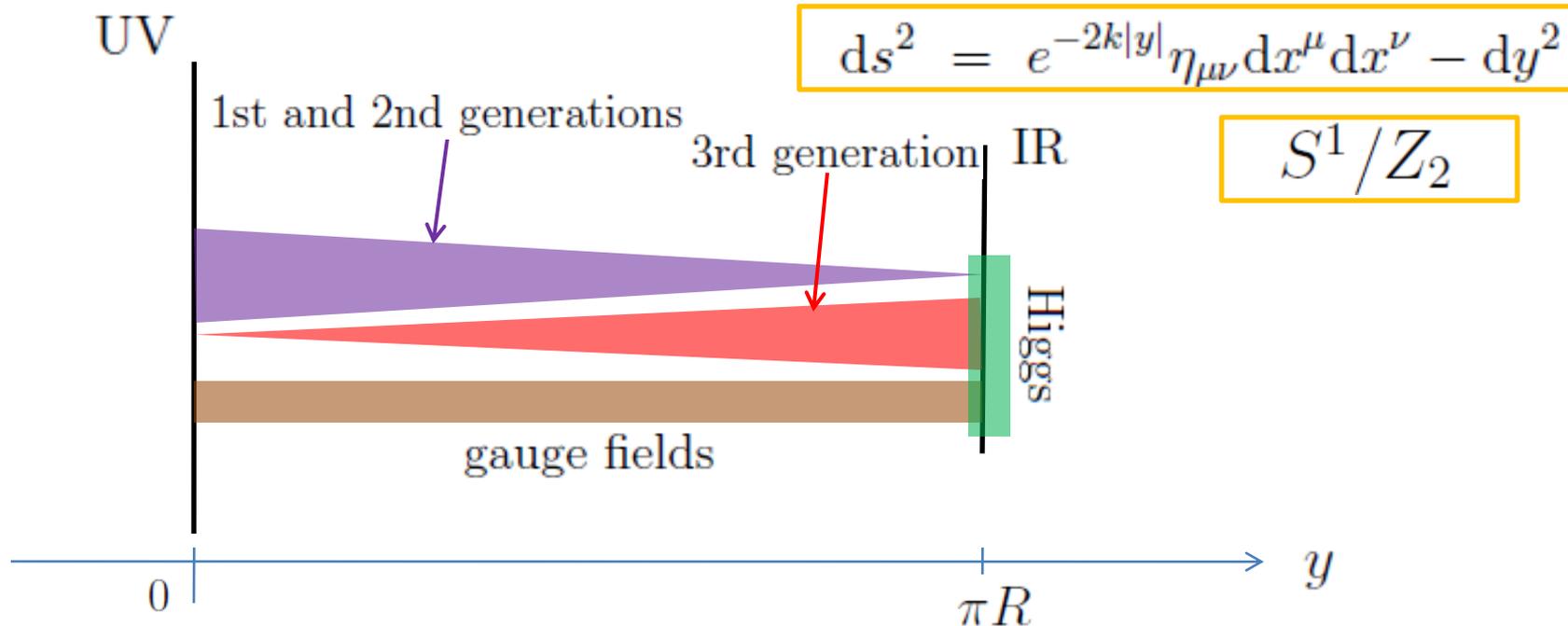
ii. KK scale need not be \sim TeV to explain the Yukawa hierarchy.

- How can we confirm bulk matter RS model if **KK is far above TeV** scale?

 If **SUSY particles** are accessible at colliders, we can do that by studying their **rare decays**.

Setup

Bulk Matter RS Model



- **Geometrical overlap** between bulk matters and IR-localized Higgs is given by

$$\sqrt{\frac{1 - 2c_i}{2\{1 - e^{-(1-2c_i)kR\pi}\}}} \sqrt{\frac{1 - 2c_j}{2\{1 - e^{-(1-2c_j)kR\pi}\}}} \quad (c_i : 5D \text{ Dirac mass, } i, j : \text{ flavor indices})$$

- Can be exponentially suppressed with $O(1)$ c_i 's .

Geometrical Overlap

- Determine the geometrical overlap factors so that they explain the Yukawa hierarchy.

Define $\alpha_i \equiv \sqrt{\frac{1 - 2c_i}{2\{1 - e^{-(1-2c_i)kR\pi}\}}}$ for SU(2) doublet quark Q_i ,
 $\beta_i, \gamma_i, \delta_i, \epsilon_i$ resp. for U_i, D_i, L_i, E_i .

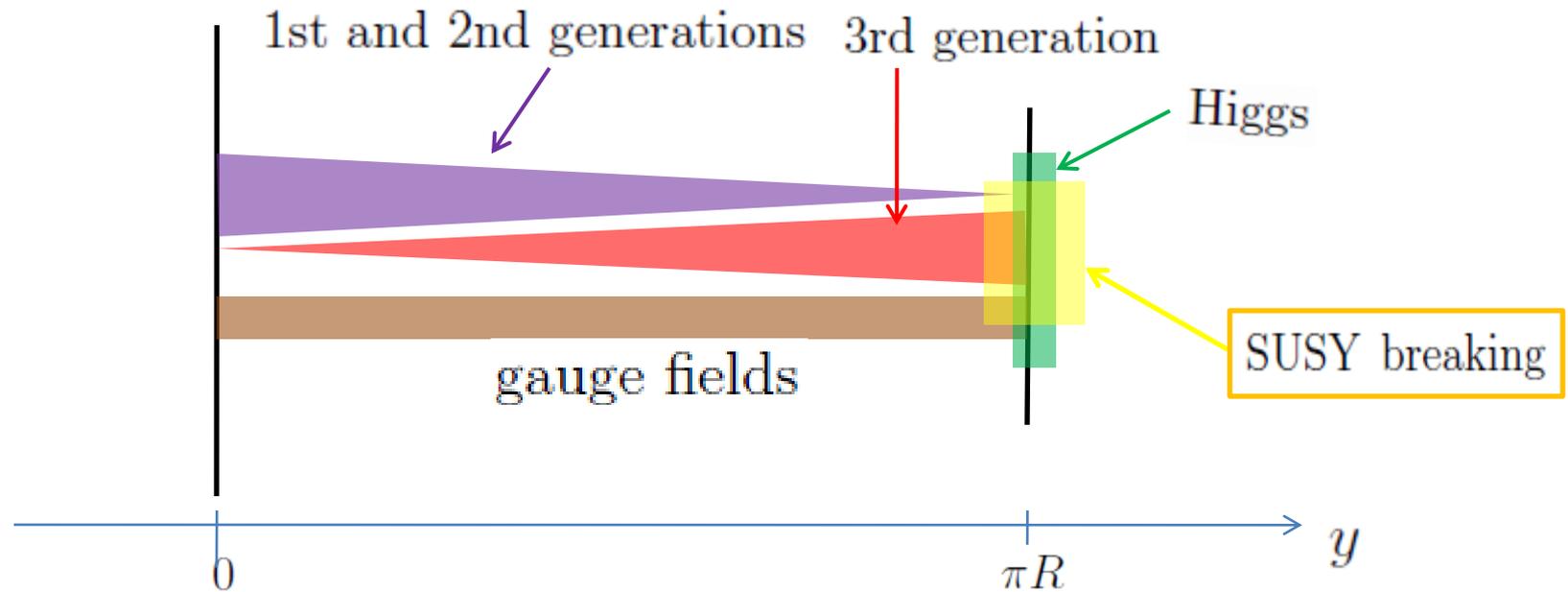
- From the data on SM fermion mass, CKM matrix and neutrino oscillations, the **orders** of $\alpha_i, \beta_i, \gamma_i, \delta_i, \epsilon_i$ are fixed as :

$$\alpha_1 \sim \lambda^3, \quad \alpha_2 \sim \lambda^2, \quad \alpha_3 \sim 1; \quad \beta_1 \sim \lambda^{-3} \frac{m_u}{v}, \quad \beta_2 \sim \lambda^{-2} \frac{m_c}{v}, \quad \beta_3 \sim 1; \quad \dots$$

$$3\delta_1 \sim \delta_2 \sim \delta_3; \quad \epsilon_1 \sim \delta_1^{-1} \frac{m_e}{v} \tan \beta, \quad \epsilon_2 \sim \delta_2^{-1} \frac{m_\mu}{v} \tan \beta, \quad \epsilon_3 \sim \delta_3^{-1} \frac{m_\tau}{v} \tan \beta;$$

Note : Absolute scale of δ_3 's is a free parameter because we do not know the seesaw scale.

SUSY Bulk Matter RS Model



- Gravity mediation for 1st and 2nd generation matter fields is geometrically suppressed. ➡ 5D sequestering
- Introduce 0 to several “messenger” fields.

➡ Hybrid of gravity mediation (on IR brane), gaugino mediation and (optional) gauge mediation.

Two Scales of Soft Terms

- Gravity mediation contributions on IR brane
- Gaugino mediation contributions

(Gaugino masses arise from contact terms, then their RG effects give rise to matter soft masses)

➔ The mass scale of soft terms :

$$\sim \frac{|\langle F \rangle|}{M_5 e^{-kR\pi}} \equiv M_X$$

F-term SUSY breaking

- Gauge mediation contributions

➔ The mass scale of soft terms :

$$\sim \frac{1}{16\pi^2} \frac{|\langle F \rangle|}{M_{mess}} \equiv M_G$$

typical messenger mass

Sparticle Mass Spectrum

Flavor-Violating Soft Terms

- Gravity mediation violates flavor.

Geometrical overlap between matters and IR brane controls the **flavor structure** of its contributions.

➡ Matter soft mass and A-terms from grav. med. are expressed in terms of $\alpha_i, \beta_i, \gamma_i, \delta_i, \epsilon_i$:

$$(m_Q^2) \sim \alpha_i \alpha_j M_X^2, \quad (m_U^2) \sim \beta_i \beta_j M_X^2, \quad (m_D^2) \sim \gamma_i \gamma_j M_X^2,$$

$$(m_L^2) \sim \delta_i \delta_j M_X^2, \quad (m_E^2) \sim \epsilon_i \epsilon_j M_X^2;$$

$$(A_u)_{ij} \sim \beta_i \alpha_j M_X, \quad (A_d)_{ij} \sim \gamma_i \alpha_j M_X, \quad (A_e)_{ij} \sim \epsilon_i \delta_j M_X$$

- Gaugino med. and gauge med. only generate flavor-universal terms.

➡ **Grav. med. contributions are detectable thru flavor-violating soft terms.**

Gravity Mediation vs. MFV

- Unfortunately, RG of Yukawa couplings also generate flavor-violating soft terms.
(Minimal Flavor Violation)
- Compare the predictions of “grav.med. in bulk matter RS model” with those of MFV.

Orders of MFV effects

up-type quark Yukawa



- In Y_u -diagonal flavor basis, MFV effects on $(m_U^2)_{ij}$, $(A_u)_{ij}$ are given by :

$$\begin{aligned} (m_U^2)_{ij} &\sim \beta_i(\alpha_i)^2(\alpha_j)^2\beta_j \max\{M_X, M_G\}^2, \\ (A_u)_{ij} &\sim \beta_i(\beta_i)^2\alpha_j M_X \end{aligned}$$

- $(m_D^2)_{ij}$, $(A_d)_{ij}$ in Y_d -diagonal basis:

$$\begin{aligned} (m_D^2)_{ij} &\sim \gamma_i(\alpha_i)^2(\alpha_j)^2\gamma_j \max\{M_X, M_G\}^2 \\ (A_d)_{ij} &\sim \gamma_i(\gamma_i)^2\alpha_j M_X \end{aligned}$$

- $(m_E^2)_{ij}$, $(A_e)_{ij}$ in Y_e -diagonal basis, if singlet neutrinos exist:

$$(m_E^2)_{ij} \sim \epsilon_i(\delta_i)^2(\delta_j)^2\epsilon_j \max\{M_X, M_G\}^2, \quad (A_e)_{ij} \sim \epsilon_i(\epsilon_i)^2\delta_j M_X$$



If $M_X \gtrsim M_G$, grav. med. induces

larger amount of flavor violation than MFV for

$$(m_U^2)_{ij}, (m_D^2)_{ij}, (m_E^2)_{ij}, (A_u)_{ij}, (A_d)_{ij}, (A_e)_{ij} \quad (i = 1, 2)$$

in their natural bases.

Experimental Test

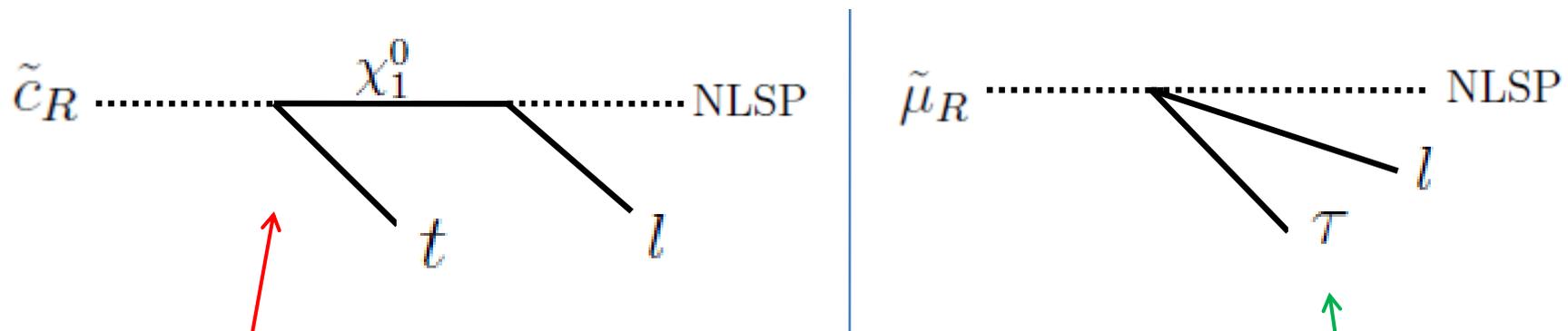
- The model predicts that flavor-violating soft terms :
 $(m_U^2)_{ij}, (m_D^2)_{ij}, (m_E^2)_{ij}, (A_u)_{ij}, (A_d)_{ij}, (A_e)_{ij} \quad (i = 1, 2)$
 are much larger than MFV effects if $M_X \gtrsim M_G$.
- **Measure the orders of these terms** to test bulk matter RS model.
- In our model, LSP is gravitino, and NLSP is charged slepton and long-lived, so that it is **possible to reconstruct the masses** of all SUSY matter particles.

SUSY Particles of Interest

- **2nd generation, SU(2) singlet** sparticles give the best channels to test the model;
 - i. Large mixings between 2nd and 3rd generations.
 - ii. Much larger mixings among SU(2) singlet sparticles than MFV.
 - iii. Doublet-singlet mixings for 2nd gen. sparticles are small, hence doublet and singlet sparticles separately form mass eigenstates.

Rare Decay Events

- \tilde{c}_R and $\tilde{\mu}_R$ mainly decay into $c \chi_1^0$ and $\mu + \text{lepton}(s) + \text{NLSP}$ resp. They can be identified by their reconstructed masses.
- Focus on the following rare decay events :



- Their rates are predicted to be :
 - $Br(\tilde{c}_R \rightarrow t \chi_1^0) \sim (\beta_2 \beta_3)^2 \sim 1/300$
 - $Br(\tilde{\mu}_R \rightarrow \tau + \dots) \sim (\epsilon_2 \epsilon_3)^2 \sim 3 \cdot 10^{-11} (\delta_3)^{-4} \tan^4 \beta$
 when $M_X \sim M_G$.

- The scale of δ_3 can be obtained by measuring the rate of MFV-dominant event: $Br(\tilde{\mu}_L \rightarrow \tau + \dots) \sim (\delta_2 \delta_3)^2 \sim (\delta_3)^4$ ¹⁵

In Other Scenarios

- If $M_X \ll M_G$, impossible to test the model because MFV effects surpass grav. med. effects.
- If $M_X < M_G$, $Br(\tilde{c}_R \rightarrow t \chi_1^0)$ and $Br(\tilde{\mu}_R \rightarrow \tau + \dots)$ are smaller, but the model can be tested through their ratio.
- If there is no 5D structure but grav. med. effects are sizeable, we will observe

$$Br(\tilde{\mu}_R \rightarrow \tau + \dots) \sim Br(\tilde{\mu}_R \rightarrow e + \dots) .$$

Conclusions

- Discussed the test of bulk matter RS model in case KK scale is far above TeV.
- Flavor-violating soft terms reflect the 5D structure of bulk matter RS model.
- Distinguishable from MFV effects.
- To test bulk matter RS model at colliders, measure the rates of rare decays in which \tilde{C}_R and $\tilde{\mu}_R$ resp. decay into SM t and τ .