

Lectures on Extra Dimensions

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- ① Motivations and mass hierarchy
- ② Strings, branes and extra dimensions (flat and warped)
- ③ Main experimental signatures
- ④ Extra U(1)'s
- ⑤ Gravity scale and number of species
- ⑥ Low string coupling and linear dilaton background

BSM physics: driven by mass hierarchy problem

Higgs mass: very sensitive to high energy physics $m_H \sim \text{UV cutoff } \Lambda$

why gravity is so weak compared to the other interactions? $\Lambda = M_P$ [5]

Possible answer (alternative to supersymmetry): Low UV cutoff $\Lambda \sim \text{TeV}$

- low scale gravity \Rightarrow large extra dimensions, warped dimensions
- low string scale \Rightarrow low scale gravity, ultra weak string coupling [6]

Experimentally testable framework:

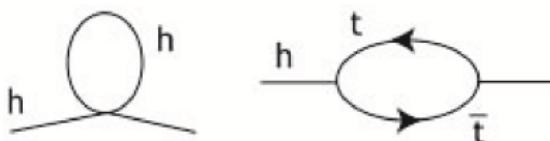
- spectacular model independent predictions
 - radical change of high energy physics at the TeV scale
- explicit model building is not necessary at this moment

Mass hierarchy problem

Higgs mass: very sensitive to high energy physics

1-loop radiative corrections:

dominant contributions:



$$\mu_{\text{eff}}^2 = \mu_{\text{bare}}^2 + \left(\frac{\lambda}{8\pi^2} - \frac{3\lambda_t^2}{8\pi^2} \right) \Lambda^2 + \dots$$

UV cutoff: $\int^\Lambda \frac{d^4 k}{k^2}$ scale of new physics

High-energy validity of the Standard Model : $\Lambda \gg \mathcal{O}(100) \text{ GeV} \Rightarrow$

“unatural” fine-tuning between μ_{bare}^2 and radiative corrections

order by order

Mass hierarchy problem

example: $\Lambda \sim \mathcal{O}(M_{\text{Planck}}) \sim 10^{19} \text{ GeV}$, loop factor $\sim 10^{-2}$

$$\Rightarrow \mu_{1-\text{loop}}^2 \sim 10^{-2} \times 10^{38} = \pm 10^{36} \text{ (GeV)}^2$$

$$\text{need } \mu_{\text{bare}}^2 \sim \mp 10^{36} \text{ (GeV)}^2 - 10^4 \text{ (GeV)}^2$$

- adjustment at the level of 1 part per 10^{32} $\mu_{\text{bare}}^2 / \mu_{1-\text{loop}}^2 = -1 \mp 10^{-32}$
- new adjustment at the next order, etc

$$\text{highest order } N: (10^{-2})^N \times 10^{38} \lesssim 10^4 \Rightarrow N \gtrsim 18 \text{ loops !}$$

- no fine tuning : $10^{-2}\Lambda^2 \lesssim 10^4 \text{ (GeV)}^2 \Rightarrow \Lambda \lesssim 1 \text{ TeV}$

→ new physics within LHC range ! [2]

Newton's law

$$m \bullet \xleftarrow{r} \bullet m \quad F_{\text{grav}} = G_N \frac{m^2}{r^2} \quad G_N^{-1/2} = M_{\text{Planck}} = 10^{19} \text{ GeV}$$

Compare with electric force: $F_{\text{el}} = \frac{e^2}{r^2} \Rightarrow$

effective dimensionless coupling $G_N m^2$ or in general $G_N E^2$ at energies E

$$E = m_{\text{proton}} \Rightarrow \frac{F_{\text{grav}}}{F_{\text{el}}} = \frac{G_N m_{\text{proton}}^2}{e^2} \simeq 10^{-40} \Rightarrow \text{Gravity is very weak !}$$

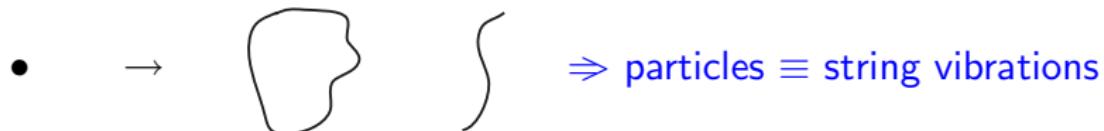
At what energy gravitation becomes comparable to the other interactions?

$$M_{\text{Planck}} \simeq 10^{19} \text{ GeV} \rightarrow \text{Planck length: } 10^{-33} \text{ cm}$$

$10^{15} \times$ the LHC energy! [2] [12]

String theory: Quantum Mechanics + General Relativity

point particle → extended objects



Framework for unification of all interactions

Mass scale: string tension $M_s \leftrightarrow$ string size: l_s

Consistent theory : 9 spatial dimensions !

six new dimensions of space

matter and gauge interactions may be localized in less than 9 dims ⇒

our universe on a p -brane ? [10]

extended in p spatial dimensions $p = 0$: particle, $p = 1$: string, . . .

Extra Dimensions

how they escape observation?

finite size R

Kaluza and Klein 1920

energy cost to send a signal:

$$E > R^{-1} \leftarrow \text{compactification scale}$$

experimental limits on their size

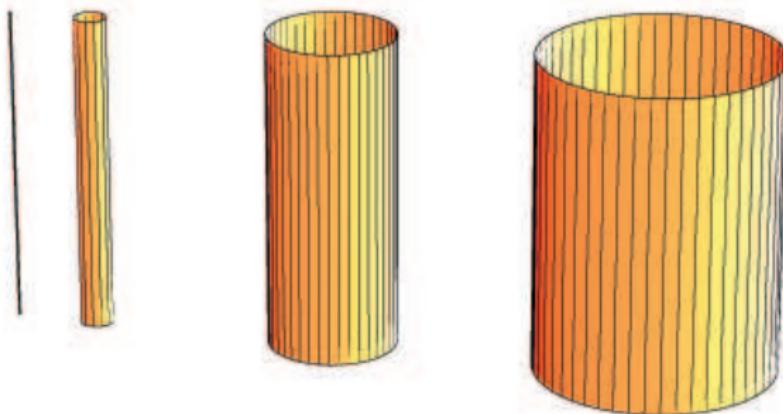
light signal : $E \gtrsim 1 \text{ TeV}$

$$R \lesssim 10^{-16} \text{ cm}$$

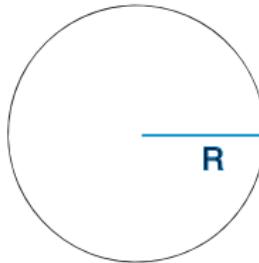
how to detect their existence?

motion in the internal space \Rightarrow mass spectrum in 3d

Dimensions D=??



- example:
- one internal circular dimension
 - light signal



plane waves e^{ipy} periodic under $y \rightarrow y + 2\pi R$

\Rightarrow quantization of internal momenta: $p = \frac{n}{R}; n = 0, 1, 2, \dots$

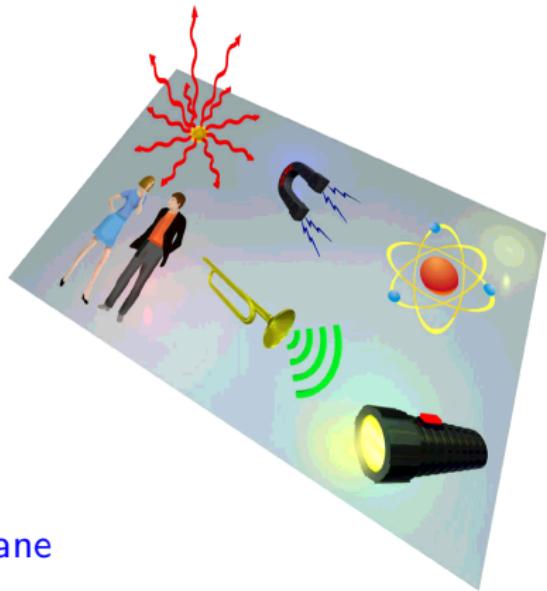
\Rightarrow 3d: tower of Kaluza Klein particles with masses $M_n = n/R$

$$p_0^2 - \vec{p}^2 - p_5^2 = 0 \Rightarrow p^2 = p_5^2 = \frac{n^2}{R^2}$$

$E \gg R^{-1}$: emission of many massive photons

\Leftrightarrow propagation in the internal space [6]

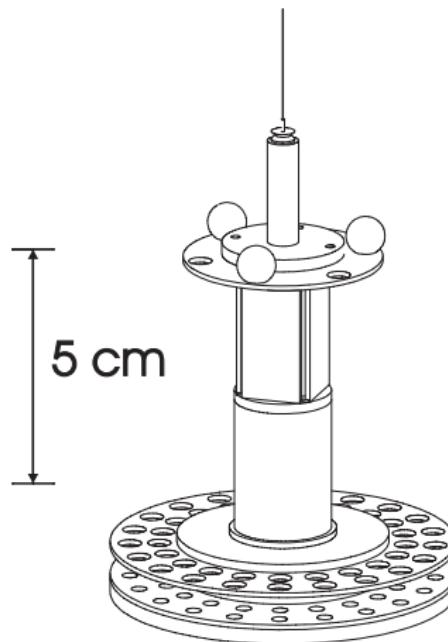
Our universe on a membrane



Two types of new dimensions:

- longitudinal: along the membrane
- transverse: “hidden” dimensions

only gravitational signal $\Rightarrow R_{\perp} \lesssim 1 \text{ mm} !$



$R_{\perp} \lesssim 45 \mu\text{m}$ at 95% CL

- dark-energy length scale $\approx 85 \mu\text{m}$ [21]

Low scale gravity

Extra large \perp dimensions can explain the apparent weakness of gravity

total force = observed force \times volume \perp [5]

total force $\simeq \mathcal{O}(1)$ at 1 TeV n dimensions of size R_\perp

$n = 1 : R_\perp \simeq 10^8$ km excluded [53]

$n = 2 : R_\perp \simeq 0.1$ mm $(10^{-12}$ GeV) possible

$n = 6 : R_\perp \simeq 10^{-13}$ mm $(10^{-2}$ GeV)

- distances $> R_\perp$: gravity 3d

however for $< R_\perp$: gravity $(3+n)d$ [14]

- strong gravity at 10^{-16} cm $\leftrightarrow 10^3$ GeV

10³⁰ times stronger than thought previously! [15]

Low scale gravity

Extra large \perp dimensions can explain the apparent weakness of gravity

total force = observed force \times volume \perp

$$G_N^* E^{2+n} = G_N E^2 \times V_\perp E^n$$

$$G_N^* = M_*^{-(2+n)} : (4+n)\text{-dim gravitational constant}$$

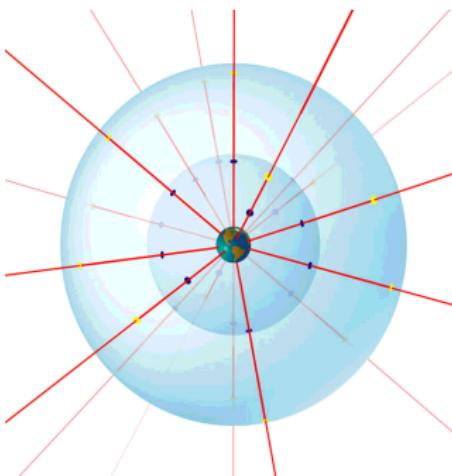
total force $\simeq \mathcal{O}(1)$ at 1 TeV n dimensions of size R_\perp

$$\Rightarrow V_\perp = R_\perp^n$$

$$\Rightarrow 1 = E^2/M_P^2 \times (R_\perp E)^n \text{ for } E = M_* \simeq 1 \text{ TeV}$$

Gravity modification at submillimeter distances

Newton's law: force decreases with area



$$3d: \text{force} \sim 1/r^2$$

$$(3+n)d: \text{force} \sim 1/r^{2+n}$$

observable for $n = 2$: $1/r^4$ with $r \ll .1 \text{ mm}$ [12]

Connect string theory to the real world

- Are there low energy string predictions testable at LHC ?
- What can we hope to learn from LHC on string phenomenology ?



Very different answers depending mainly on the value of the string scale M_s

- arbitrary parameter : Planck mass $M_P \longrightarrow \text{TeV}$

- physical motivations \Rightarrow favored energy regions:

- High :
$$\begin{cases} M_P^* \simeq 10^{18} \text{ GeV} & \text{Heterotic scale} \\ M_{\text{GUT}} \simeq 10^{16} \text{ GeV} & \text{Unification scale} \end{cases}$$
- Intermediate : around 10^{11} GeV ($M_s^2/M_P \sim \text{TeV}$)
SUSY breaking, strong CP axion, see-saw scale
- Low : TeV (hierarchy problem)

High string scale

perturbative heterotic string : the most natural for SUSY and unification

gravity and gauge interactions have same origin

massless excitations of the closed string

But mismatch between string and GUT scales:

$$M_s = g_H M_P \simeq 50 M_{\text{GUT}} \quad g_H^2 \simeq \alpha_{\text{GUT}} \simeq 1/25 \quad [21]$$

in GUTs only one prediction from 3 gauge couplings unification: $\sin^2 \theta_W$

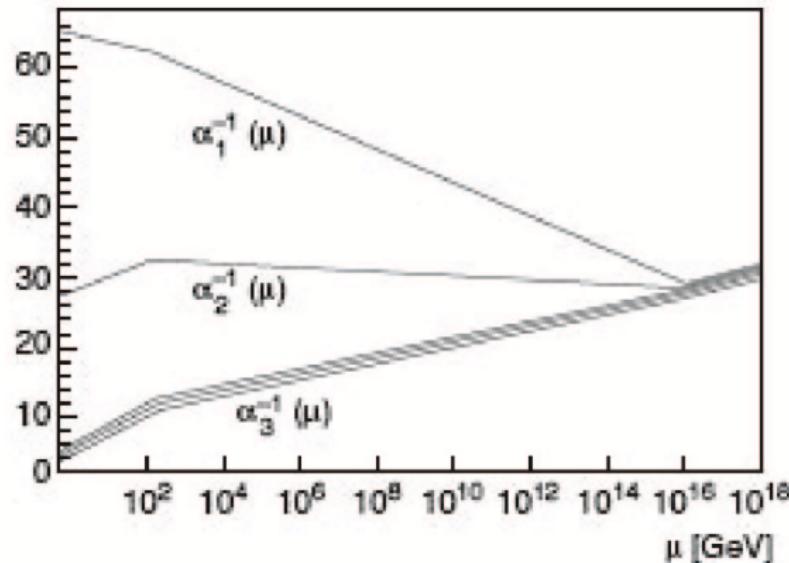
introduce large threshold corrections or strong coupling $\rightarrow M_s \simeq M_{\text{GUT}}$

but loose predictivity

Gauge coupling unification

Energy evolution of gauge couplings $\alpha_i = g_i^2 / 4\pi \Rightarrow$

low energy data → extrapolation at high energies:



unification at $M_{GUT} \simeq 10^{15} - 10^{16}$ GeV

Intersecting branes: ‘perfect’ for SM embedding

product of unitary gauge groups (**brane stacks**) and bi-fundamental reps

but no unification: no prediction for M_s , independent gauge couplings

however GUTs: problematic:

- no perturbative $SO(10)$ spinors
- no top-quark Yukawa coupling in $SU(5)$: $10 \ 10 \ 5_H$

$SU(5)$ is part of $U(5) \Rightarrow U(1)$ charges : 10 charge 2 ; 5_H charge ± 1

\Rightarrow cannot balance charges with $SU(5)$ singlets

can be generated by D-brane instantons but ...

→ Non-perturbative M/F-theory models:

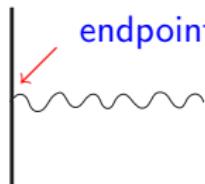
combine good properties of heterotic and intersecting branes

but lack exact description for systematic studies

D-brane embedding of the Standard Model

Generic spectrum: N coincident branes $\Rightarrow U(N)$

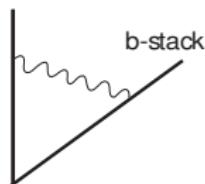
a-stack



endpoint transformation: N_a or \bar{N}_a $U(1)_a$ charge: +1 or -1
 \Rightarrow "baryon" number

- open strings from the same stack \Rightarrow adjoint gauge multiplets of $U(N_a)$
- stretched between two stacks \Rightarrow bifundamentals of $U(N_a) \times U(N_b)$

a-stack



non-oriented strings \Rightarrow also:

- orthogonal and symplectic groups $SO(N), Sp(N)$
- matter in antisymmetric + symmetric reps

Type I string theory \Rightarrow D-brane world

I.A.-Arkani-Hamed-Dimopoulos-Dvali '98

- gravity: closed strings propagating in 10 dims
- gauge interactions: open strings with their ends attached on D-branes

Dimensions of finite size: n transverse $6 - n$ parallel

calculability $\Rightarrow R_{\parallel} \simeq l_{\text{string}}$; R_{\perp} arbitrary

$$M_p^2 \simeq \frac{1}{g_s^2} M_s^{2+n} R_{\perp}^n \quad g_s = \alpha : \text{weak string coupling [17]}$$

Planck mass in $4 + n$ dims: M_*^{2+n}

$$M_s \sim 1 \text{ TeV} \Rightarrow R_{\perp}^n = 10^{32} l_s^n \quad \text{small } M_s/M_P : \text{extra-large } R_{\perp}$$

[49] [50]

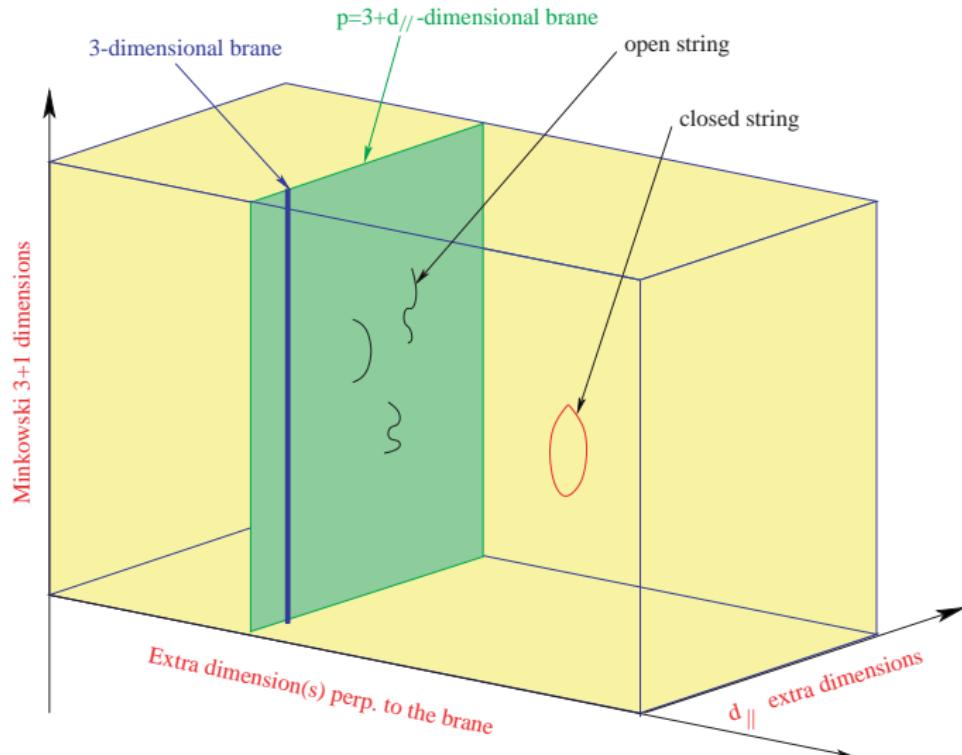
$$R_{\perp} \sim .1 - 10^{-13} \text{ mm for } n = 2 - 6 \quad \text{[11]}$$

distances $< R_{\perp}$: gravity $(4+n)$ -dim \rightarrow strong at 10^{-16} cm [23]

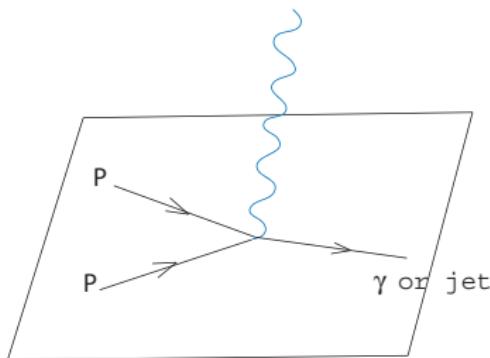
Braneworld

2 types of compact extra dimensions:

- parallel (d_{\parallel}): $\lesssim 10^{-16}$ cm (TeV) [26]
- transverse (\perp): $\lesssim 0.1$ mm (meV) [40]



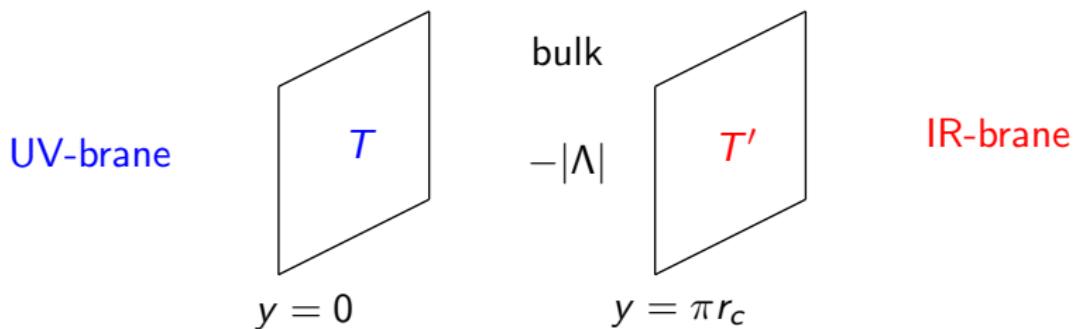
Gravitational radiation in the bulk \Rightarrow missing energy



Collider bounds on R_{\perp} in mm			
	$n = 2$	$n = 4$	$n = 6$
LEP 2	4.8×10^{-1}	1.9×10^{-8}	6.8×10^{-11}
Tevatron	5.5×10^{-1}	1.4×10^{-8}	4.1×10^{-11}
LHC	4.5×10^{-3}	5.6×10^{-10}	2.7×10^{-12}
NLC	1.2×10^{-2}	1.2×10^{-9}	6.5×10^{-12}

Randal Sundrum models

spacetime = slice of AdS₅ : $ds^2 = e^{-2k|y|} \eta_{\mu\nu} dx^\mu dx^\nu + dy^2$ $k^2 \sim \Lambda/M_5^3$



- exponential hierarchy: $M_W = M_P e^{-2kr_c}$ $M_P^2 \sim M_5^3/k$ $M_5 \sim M_{GUT}$
- 4d gravity localized on the UV-brane, but KK gravitons on the IR

$$m_n = c_n k e^{-2kr_c} \sim \text{TeV} \quad c_n \simeq (n + 1/4) \text{ for large } n$$

\Rightarrow spin-2 TeV resonances in di-lepton or di-jet channels [51] [52]

- weakly coupled for $m_n < M_5 e^{-2kr_c}$ $\Rightarrow k < M_5$
- viable models: SM gauge bosons in the bulk, Higgs on the IR-brane
- AdS/CFT duals to strongly coupled 4d field theories
composite Higgs models, technicolor-type $g_{YM} = M_5/k > 1$

Other accelerator signatures

- Large TeV dimensions seen by SM gauge interactions
⇒ KK resonances of SM gauge bosons [22] I.A. '90

$$M_n^2 = M_0^2 + \frac{n^2}{R^2} ; \quad n = \pm 1, \pm 2, \dots$$

- string physics and possible strong gravity effects

Massive string vibrations ⇒ e.g. resonances in dijet distribution [34] [36]

$$M_j^2 = M_0^2 + M_s^2 j ; \quad \text{maximal spin : } j+1$$

higher spin excitations of quarks and gluons with strong interactions

Anchordoqui-Goldberg-Lüst-Nawata-Taylor-Stieberger '08

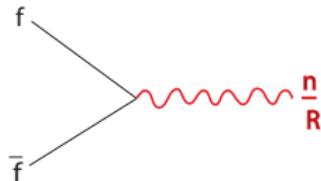
production of micro-black holes? [38]

Giddings-Thomas, Dimopoulos-Landsberg '01

Localized fermions (on 3-brane intersections) [40]

⇒ single production of KK modes

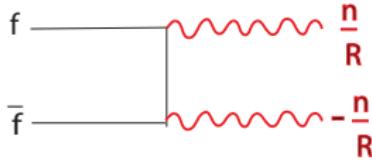
I.A.-Benakli '94



- strong bounds indirect effects: $R^{-1} \gtrsim 3 \text{ TeV}$
- new resonances but at most $n = 1$

Otherwise KK momentum conservation [30]

⇒ pair production of KK modes (universal dims)

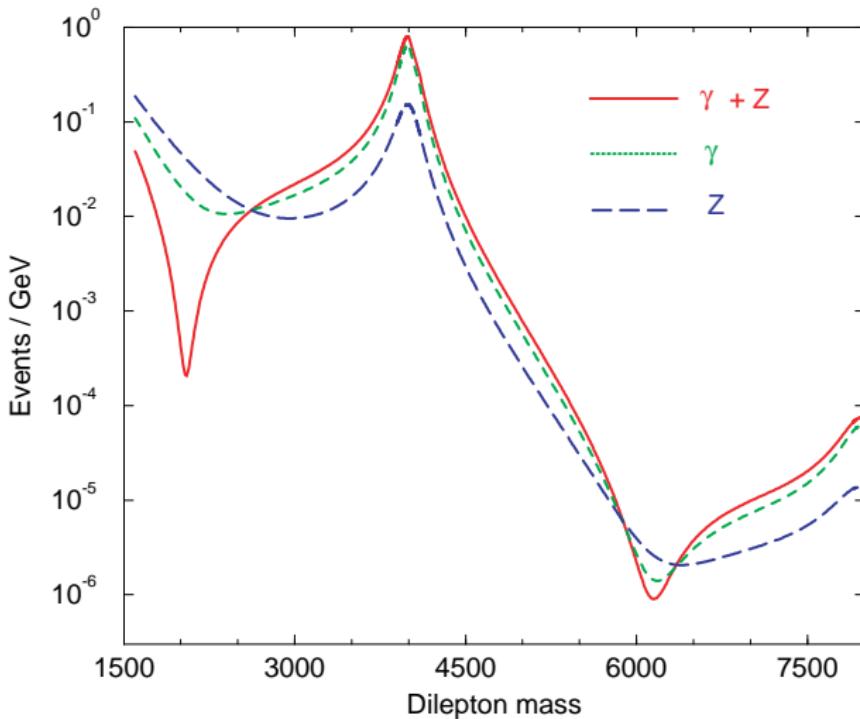


- weak bounds $R^{-1} \gtrsim 300\text{-}500 \text{ GeV}$
- no resonances
- lightest KK stable : dark matter candidate

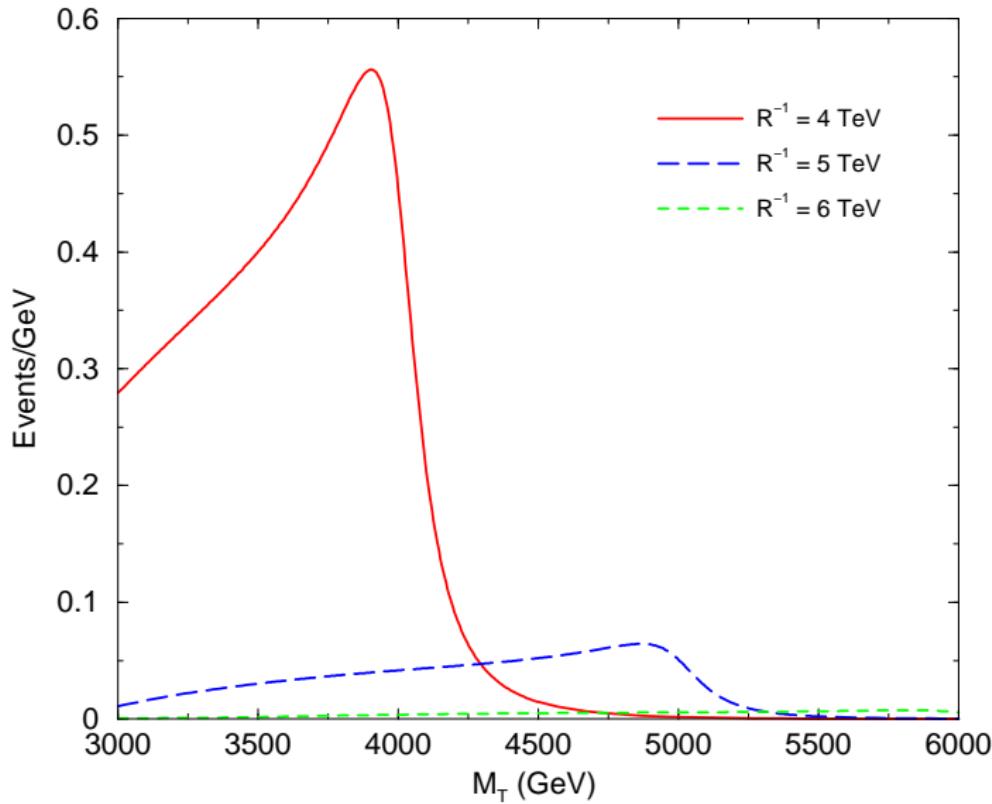
Servant-Tait '02

$R^{-1} = 4 \text{ TeV}$

I.A.-Benakli-Quiros '94, '99



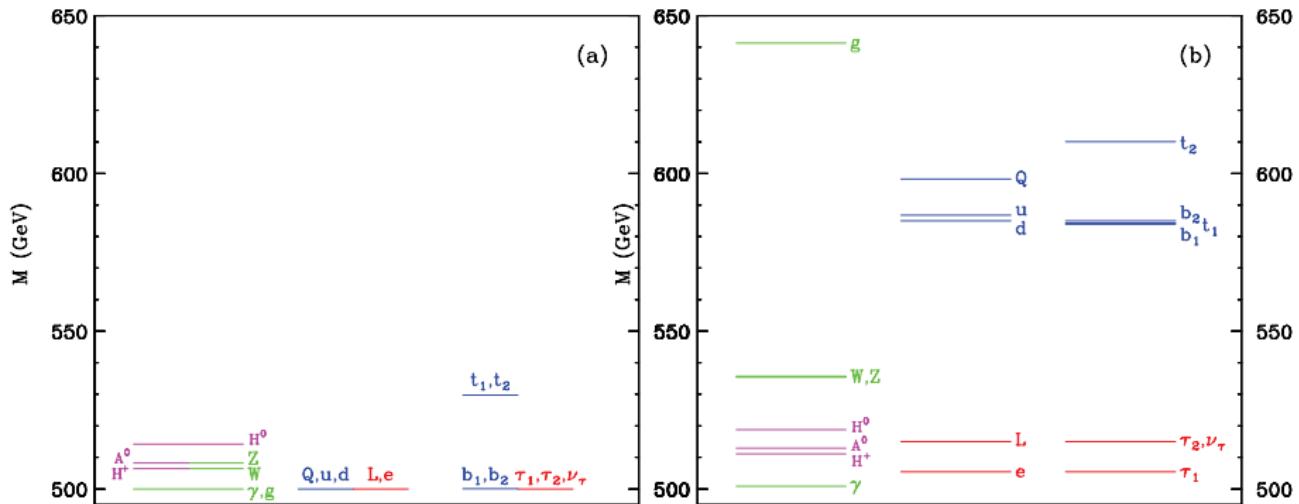
KK W-production at LHC in the $l\nu$ channel [27]



Mass spectrum

Radiative corrections \Rightarrow mass shifts that lift degeneracy at lowest KK level
divergent sum over KK modes in the loop \Rightarrow cutoff scale $\Lambda \simeq 10/R$

Cheng-Matchev-Schmaltz '02

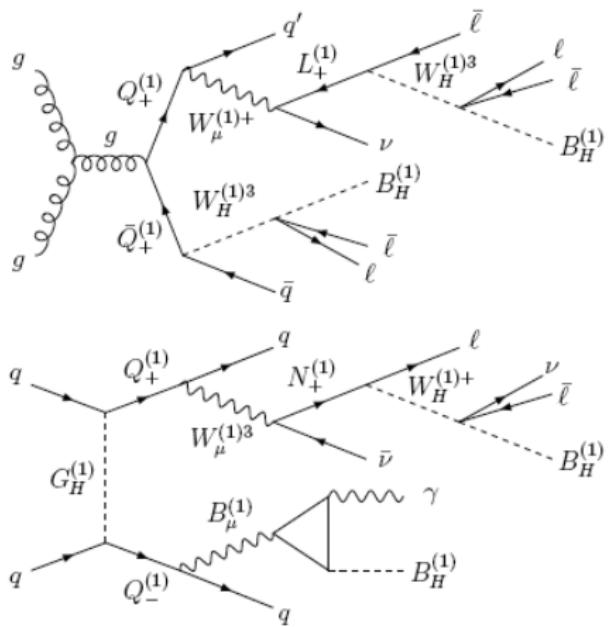
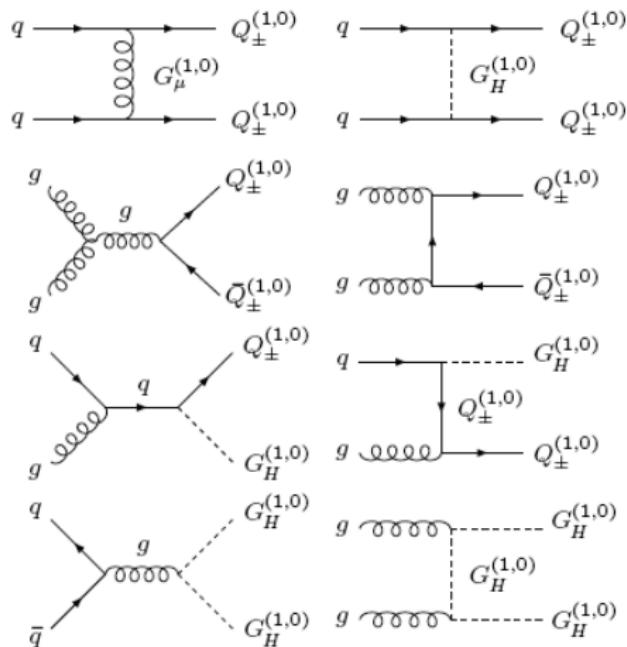


UED hadron collider phenomenology

- large rates for KK-quark and KK-gluon production
LHC: 1-100 pb for $R^{-1} \lesssim 800$ GeV
- cascade decays via KK- W bosons and KK-leptons
determine particle properties from different distributions
- missing energy from LKP: weakly interacting escaping detection
- phenomenology similar to supersymmetry [33]
spin determination important for distinguishing SUSY and UED [26]

gluino	1/2	KK-gluon	1
squark	0	KK-quark	1/2
chargino	1/2	KK- W boson	1
slepton	0	KK-lepton	1/2
neutralino	1/2	KK- Z boson	1

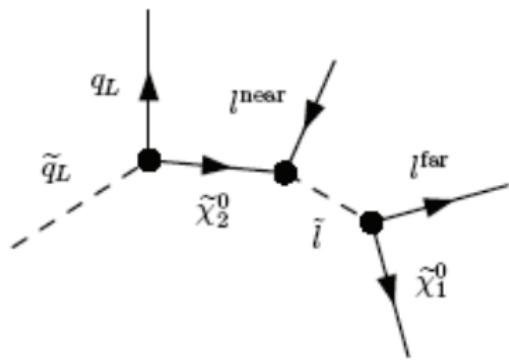
Production at LHC



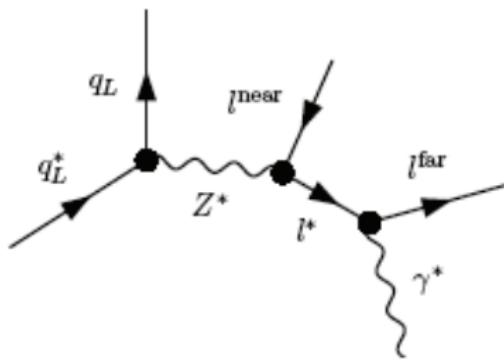
SUSY vs UED signals at LHC

Example: jet dilepton final state [31]

SUSY



UED



Massive string vibrations

indirect effects: virtual exchanges \Rightarrow effective interactions

e.g. four-fermion operators

Actual limits: Matter fermions on

- same set of branes $\Rightarrow M_s \gtrsim 500$ GeV dim-8: $\frac{g^2}{M_s^4}(\bar{\psi}\partial\psi)^2$
- brane intersections : $M_s \gtrsim 2 - 3$ TeV dim-6: $\frac{g^2}{M_s^2}(\bar{\psi}\psi)^2$

Cullen-Perelstein-Peskin, I.A.-Benakli-Laugier '00

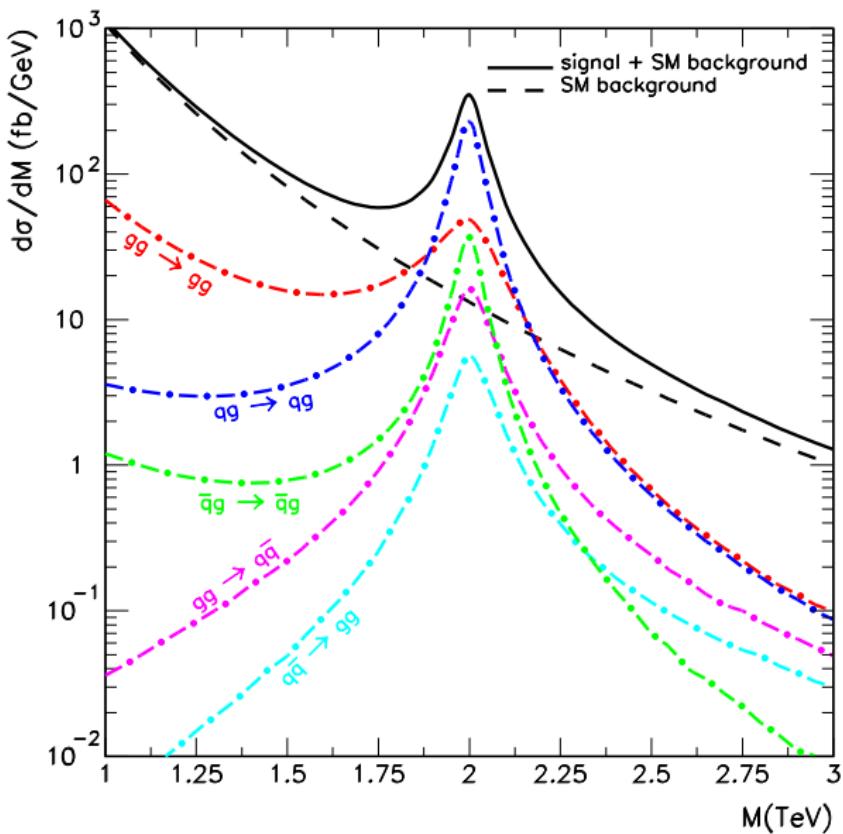
High energies \Rightarrow direct production: string physics

Universal deviation from Standard Model in dijet distribution

$M_s = 2 \text{ TeV}$

Width = 15-150 GeV

Anchordoqui-Goldberg-Lüst-Nawata-Taylor-Stieberger '08 [26]



present LHC limits (2010 data): $M_s \gtrsim 2.5 \text{ TeV}$

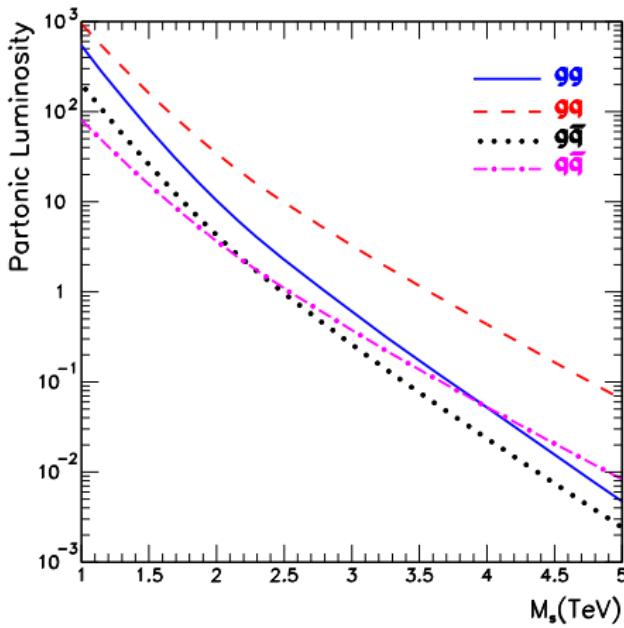
Tree level superstring amplitudes involving at most 2 fermions and gluons:
model independent for any compactification, # of susy's, even none
no intermediate exchange of KK, windings or graviton emmission
Universal sum over infinite exchange of string (Regge) excitations [26]

Parton luminosities in pp above TeV

are dominated by gq, gg

⇒ model independent

$gq \rightarrow gq, gg \rightarrow gg, gg \rightarrow q\bar{q}$



Cross sections

$$\left. \begin{array}{l} |\mathcal{M}(gg \rightarrow gg)|^2 , \quad |\mathcal{M}(gg \rightarrow q\bar{q})|^2 \\ |\mathcal{M}(q\bar{q} \rightarrow gg)|^2 , \quad |\mathcal{M}(qg \rightarrow qg)|^2 \end{array} \right\}$$

model independent
for any compactification

Lüst-Stieberger-Taylor '08

$$|\mathcal{M}(gg \rightarrow gg)|^2 = g_{YM}^4 \left(\frac{1}{s^2} + \frac{1}{t^2} + \frac{1}{u^2} \right)$$

$$\times \left[\frac{9}{4} (s^2 V_s^2 + t^2 V_t^2 + u^2 V_u^2) - \frac{1}{3} (s V_s + t V_t + u V_u)^2 \right]$$

$$|\mathcal{M}(gg \rightarrow q\bar{q})|^2 = g_{YM}^4 \frac{t^2 + u^2}{s^2} \left[\frac{1}{6} \frac{1}{tu} (t V_t + u V_u)^2 - \frac{3}{8} V_t V_u \right] \quad M_s = 1$$

$$V_s = -\frac{tu}{s} \quad B(t, u) = 1 - \frac{2}{3}\pi^2 tu + \dots \quad V_t : s \leftrightarrow t \quad V_u : s \leftrightarrow u$$

YM limits agree with e.g. book "Collider Physics" by Barger, Phillips

Black hole production

String-size black hole energy threshold : $M_{\text{BH}} \simeq M_s/g_s^2$

Horowitz-Polchinski '96, Meade-Randall '07

- string size black hole: $r_H \sim l_s = M_s^{-1}$
- black hole mass: $M_{\text{BH}} \sim r_H^{d-3}/G_N \quad G_N \sim l_s^{d-2} g_s^2$

weakly coupled theory \Rightarrow strong gravity effects occur much above M_s, M_*

$$g_s \sim 0.1 \text{ (gauge coupling)} \quad \Rightarrow \quad M_{\text{BH}} \sim 100M_s$$

Comparison with Regge excitations : $M_n = M_s \sqrt{n} \Rightarrow$

production of $n \sim 1/g_s^4 \sim 10^4$ string states before reach M_{BH}

Other accelerator signatures

- extra $U(1)$'s and anomaly induced terms
 - masses suppressed by a loop factor
 - new Chern-Simons type interactions
 - usually associated to known global symmetries of the SM
 - (anomalous or not) such as (combinations of)
 - Baryon and Lepton number, or PQ symmetry

Two kinds of massive $U(1)$'s:

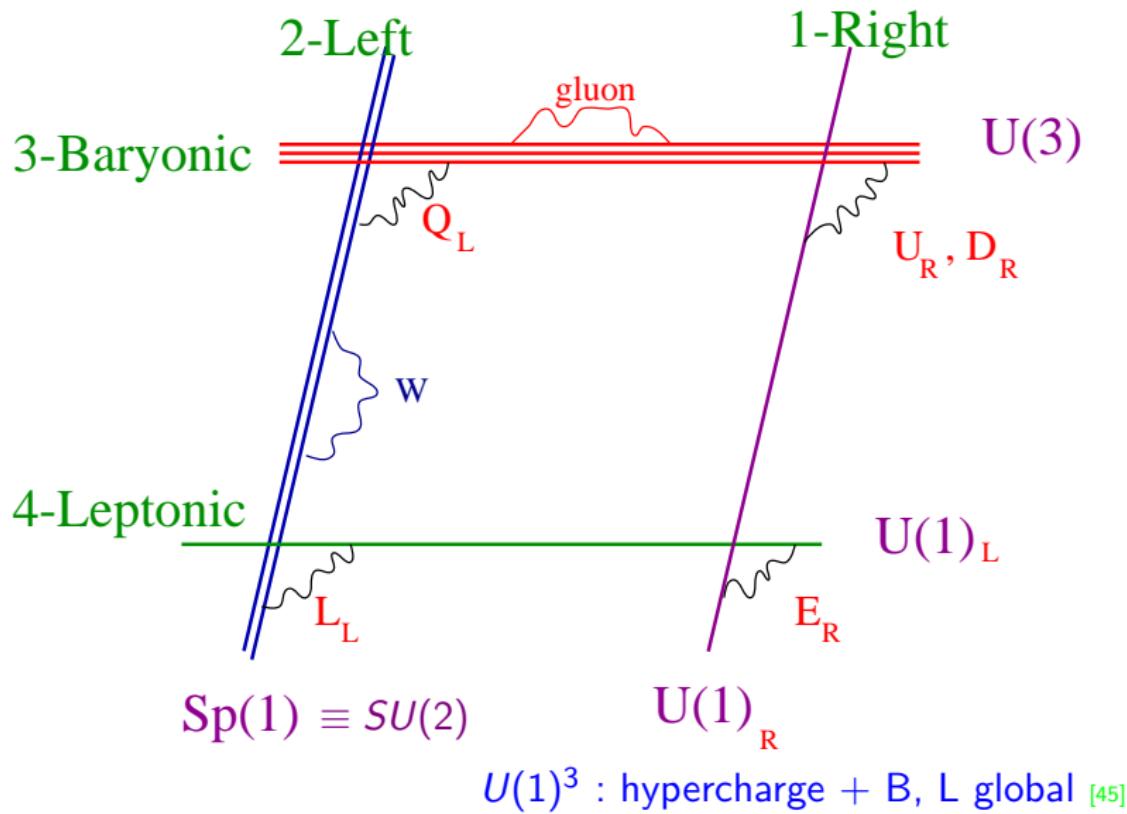
I.A.-Kiritsis-Rizos '02

- 4d anomalous $U(1)$'s: $M_A \simeq g_A M_s$

- 4d non-anomalous $U(1)$'s: (but masses related to 6d anomalies)

$$M_{NA} \simeq g_A M_s V_2 \leftarrow (6d \rightarrow 4d) \text{ internal space} \Rightarrow M_{NA} \geq M_A$$

Standard Model on D-branes



global symmetries

- B and L become massive due to anomalies
Green-Schwarz terms
- the global symmetries remain in perturbation
 - Baryon number \Rightarrow proton stability
 - Lepton number \Rightarrow protect small neutrino masses

no Lepton number $\Rightarrow \frac{1}{M_s} LLHH \rightarrow$ Majorana mass: $\frac{\langle H \rangle^2}{M_s} LL$

\sim GeV

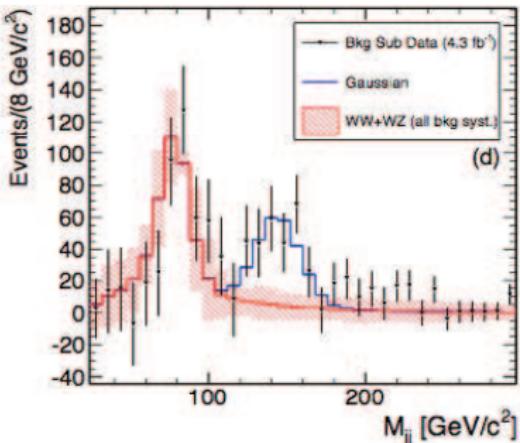
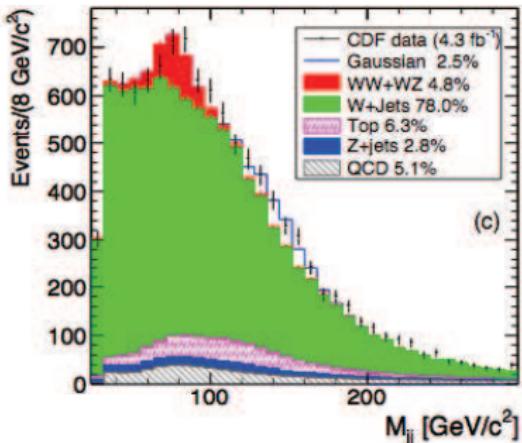
- $B, L \Rightarrow$ extra Z' s
with for instance leptophobic couplings leading to CDF-type events [21]

Anchordoqui-I.A.-Goldberg-Huang-Lüst-Taylor to appear



"Bump" in W+2jet events - CDF

- Routine measurement of WW to $\ell\nu jj$ cross section led to deviations between predicted by ALPGEN background and data around 145 GeV
- "Excess" could be described by signal with Gaussian shape with cross section of $\sim 4 \text{ pb}$

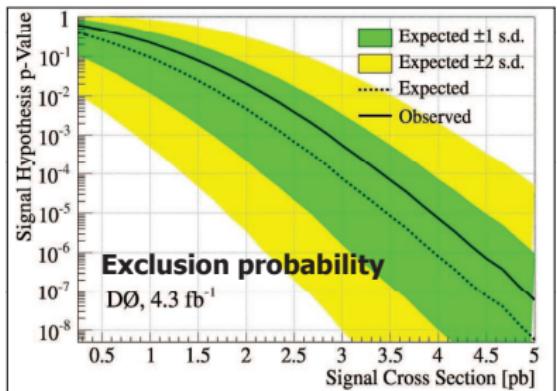
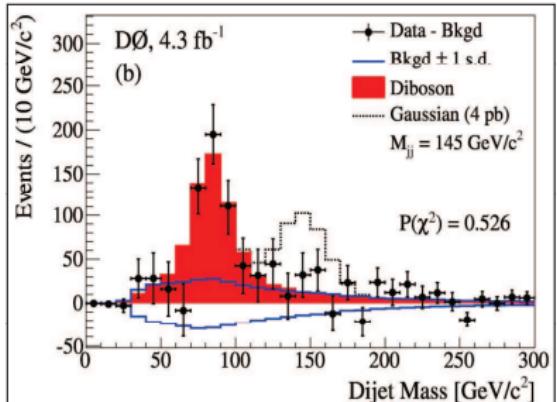
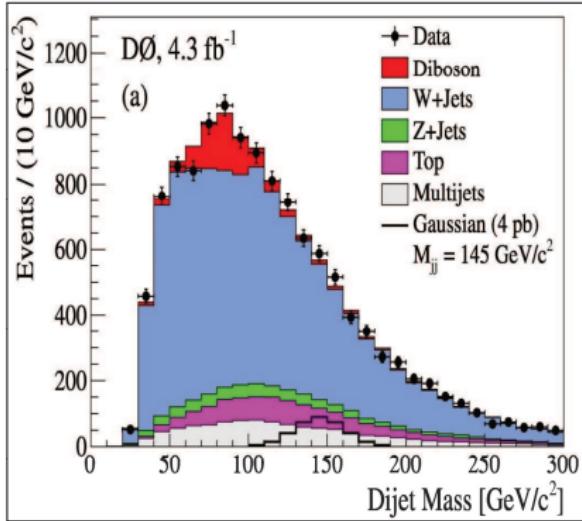


- What is this
 - Mis-modelling of Standard Model background?
 - Hints of new particle with mass of $\sim 145 \text{ GeV}$ decaying to two jets?



"Bump" in W+2jet events - DØ

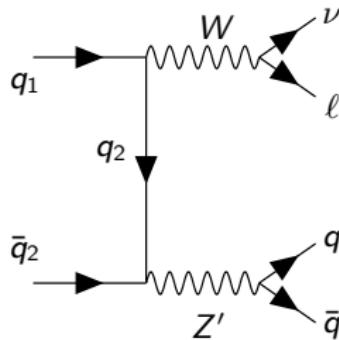
- DØ repeated analysis with exactly the same selections as CDF
- No evidence of resonance near to 145 GeV observed
 - Limits set
- Importance of the two experiments



Dmitri Denisov, Erice, June 2011

possible explanation: Z' in D-brane models

Anchordoqui-Goldberg-Huang-Lüst-Taylor '11



Required properties: leptophobic, no mixing with Z after EW symmetry breaking, much lighter than M_s

Good candidate: anomalous Baryon number $U(1)_B$

$U(3) \times SU(2) \times U(1)_L \times U(1)_R$ D-brane model [40]

3 natural $U(1)$'s: $Y, B, B - L$

R-neutrino $\Rightarrow B - L$ anomaly free \Rightarrow natural heavier than B

In general: 3×3 rotation matrix \Rightarrow 6 parameters: 3 couplings + 3 angles

4 parameters are fixed by the SM: 2 couplings and 2 angles

from hypercharge combination $Y = \frac{1}{2}(B - L + Q_{1R})$

remaining 2 fixed by phenomenological constraints:

leptophobic $\lesssim 1\%$, no mixing with Z ,

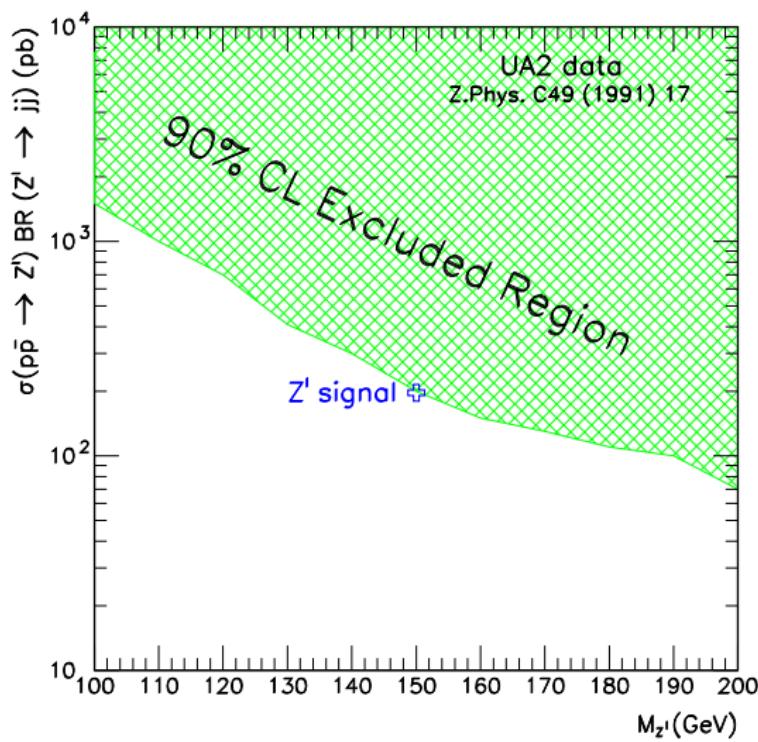
$\sigma(p\bar{p} \rightarrow WZ') \times \text{BR}(Z' \rightarrow jj) \lesssim 1.9 \text{ pb}$ at $\sqrt{s} = 1.96 \text{ TeV}$,

UA2 upper limit on $\sigma(p\bar{p} \rightarrow Z') \times \text{BR}(Z' \rightarrow jj)$ at $\sqrt{s} = 630 \text{ GeV}$

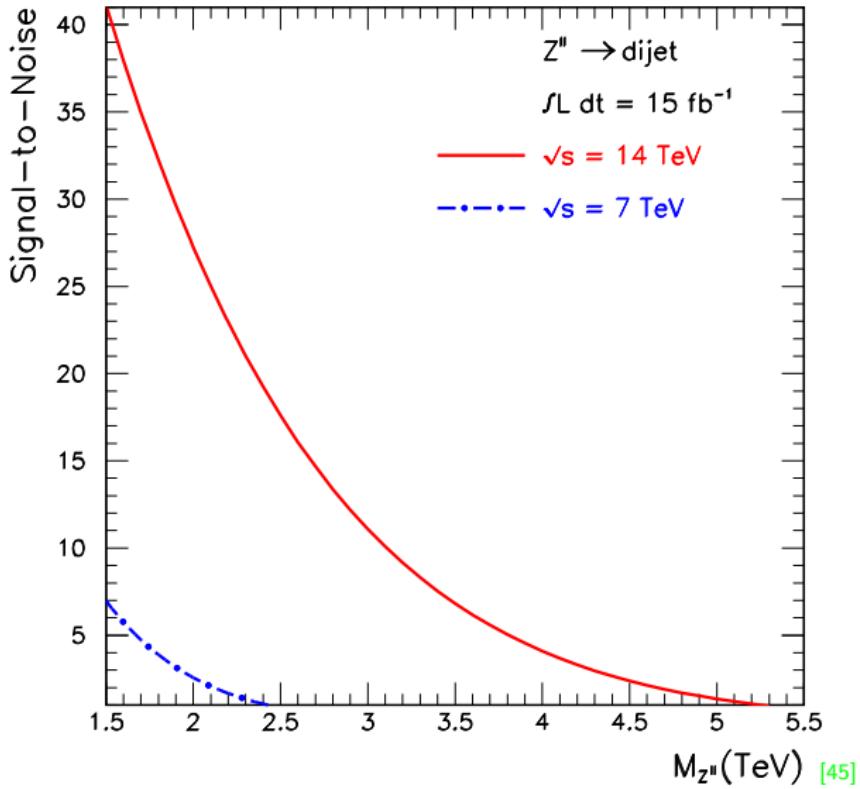
further predictions: Z'' [47] $\gamma + \text{jet}$ ($gg \rightarrow \gamma g$ from Y component on color stack)

Bounds on Z' 's from direct production at UA2

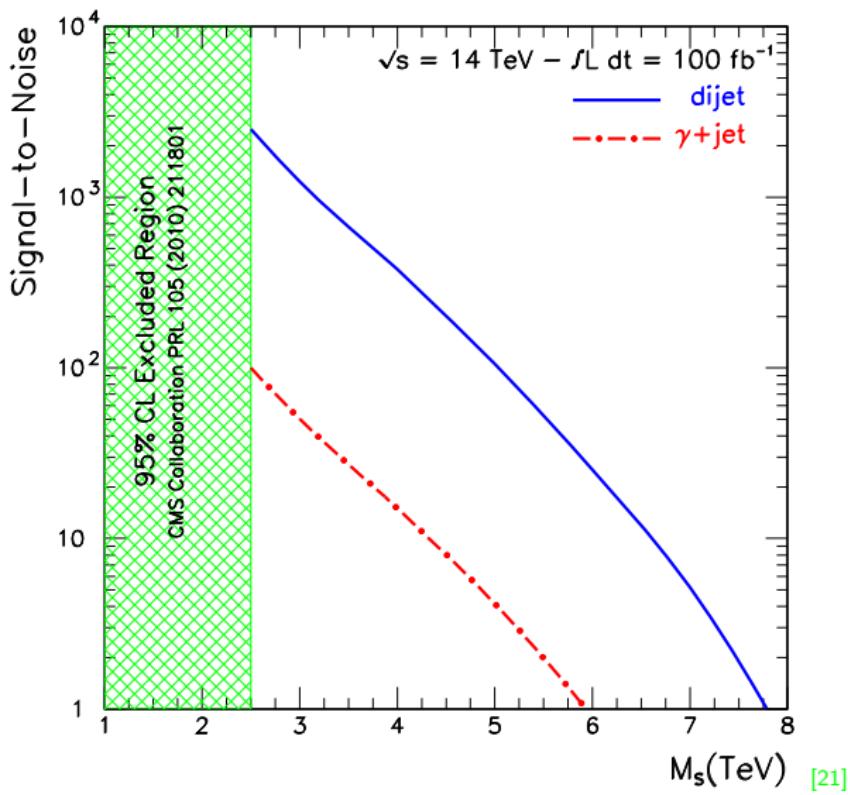
$p\bar{p} \rightarrow Z' \rightarrow jj$ at $\sqrt{s} = 640$ GeV



Dijet signals of Z'' at the LHC



$pp \rightarrow$ dijet vs $pp \rightarrow \gamma$ jet at LHC



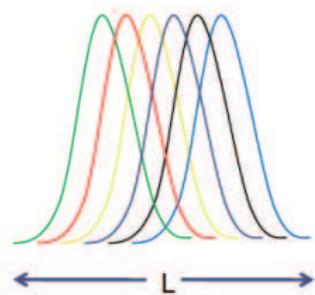
More general framework: large number of species

N particle species \Rightarrow lower quantum gravity scale : $M_*^2 = M_p^2/N$

Dvali '07, Dvali, Redi, Brustein, Veneziano, Gomez, Lüst '07-'10

derivation from: black hole evaporation or quantum information storage

Pixel of size L containing N species storing information:



localization energy $E \gtrsim N/L \rightarrow$

Schwarzschild radius $R_s = N/(LM_p^2)$

no collapse to a black hole : $L \gtrsim R_s \Rightarrow L \gtrsim \sqrt{N}/M_p = 1/M_*$

$M_* \simeq 1 \text{ TeV} \Rightarrow N \sim 10^{32} \text{ particle species !}$

2 ways to realize $N = 10^{32}$ lowering the string scale

- ① Large volume compactifications

SM on D-branes [21]

$$N = R_{\perp}^n I_s^n : \text{number of KK modes up to energies of order } M_* \simeq M_s$$

- ② $N \sim$ effective number of string modes contributing to the BH bound

Dvali-Lüst '09, Dvali-Gomez '10

$$N_s = \frac{1}{g_s^2} \text{ with } g_s \simeq 10^{-16} \quad \text{SM on NS5-branes}$$

I.A.-Pioline '99, I.A.-Dimopoulos-Giveon '01

in this case gravity does NOT become strong at M_s

Both ways are compatible with the general string relation:

$$M_p^2 = \frac{1}{g_s^2} V_6 M_s^8 \quad V_6 : \text{internal 6d compactification volume}$$

Gauge/Gravity duality \Rightarrow toy 5d bulk model

Gravity background : near horizon geometry (holography) Maldacena '98

Analogy from D3-branes : AdS_5

NS-5 branes : $(\mathcal{M}_6 \otimes \mathbb{R}_+)$



linear dilaton background in 5d flat string-frame metric $\Phi = -\alpha|y|$

Aharony-Berkooz-Kutasov-Seiberg '98

"cut" the space of the extra dimension \Rightarrow gravity on the brane

$$S_{bulk} = \int d^4x \int_0^{r_c} dy \sqrt{-g} e^{-\Phi} (M_5^3 R + M_5^3 (\nabla \Phi)^2 - \Lambda)$$

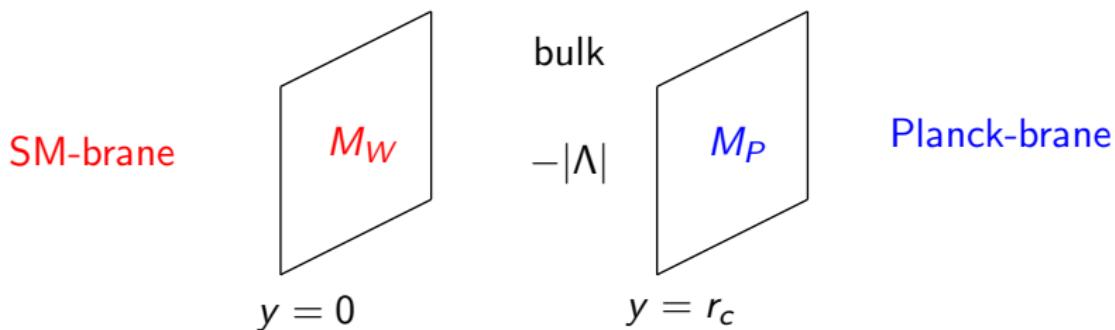
$$S_{vis(hid)} = \int d^4x \sqrt{-g} e^{-\Phi} (L_{SM(hid)} - T_{vis(hid)})$$

Tuning conditions: $T_{vis} = -T_{hid} \leftrightarrow \Lambda < 0$ [24]

Linear dilaton background IA-Arvanitaki-Dimopoulos-Giveon '11

$$g_s^2 = e^{-\alpha|y|} ; \quad ds^2 = e^{\frac{2}{3}\alpha|y|} (\eta_{\mu\nu} dx^\mu dx^\nu + dy^2) \leftarrow \text{Einstein frame [24]}$$

$z \sim e^{\alpha y/3} \Rightarrow$ polynomial warp factor + log varying dilaton

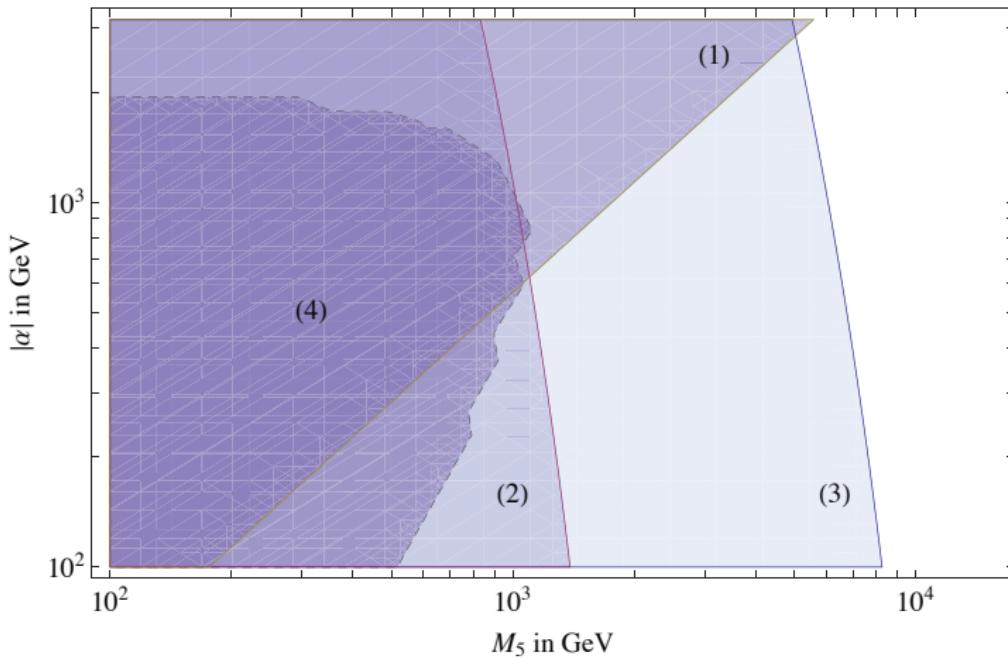


- exponential hierarchy: $g_s^2 = e^{-\alpha|y|} \quad M_P^2 \sim \frac{M_5^3}{\alpha} e^{\alpha r_c} \quad \alpha \equiv k_{RS}$
- 4d graviton flat, KK gravitons localized near SM
- SM particles cannot be in the bulk
 - bulk gauge bosons: exp suppressed couplings

LST KK graviton phenomenology

- KK spectrum : $m_n^2 = \left(\frac{n\pi}{r_c}\right)^2 + \frac{\alpha^2}{4}$; $n = 1, 2, \dots$
⇒ mass gap + dense KK modes $\alpha \sim 1 \text{ TeV}$ $r_c^{-1} \sim 30 \text{ GeV}$
- couplings : $\frac{1}{\Lambda_n} \sim \frac{1}{(\alpha r_c) M_5}$
⇒ extra suppression by a factor $(\alpha r_c) \simeq 30$
- width : $1/(\alpha r_c)^2$ suppression $\sim 1 \text{ GeV}$
⇒ narrow resonant peaks in di-lepton or di-jet channels
- extrapolates between RS and flat extra dims ($n = 1$) [12]
 $\alpha \gtrsim (0.1 \text{ mm})^{-1} \sim 10^{-2} \text{ eV}$ from microgravity experiments
⇒ distinct experimental signals

Bounds on the LST parameter space



exclusion by (1) perturbativity (2) Tevatron with 5.4 fb^{-1} data

(3) LHC 14 TeV with 10 fb^{-1} (4) diphoton at Tevatron 5.4 fb^{-1}

Conclusions

TeV strings and large extra dimensions: Physical reality or imagination?

- Well motivated theoretical framework
 - with many testable experimental predictions
 - new resonances, missing energy
- Stimulus for micro-gravity experiments
 - look for new forces at short distances
 - higher dim graviton, scalars, gauge fields

But: - unification has to be dropped

- physics is radically changed above string scale

LHC: will explore the physics beyond the Standard Model