

# What is “Discovering SUSY” ?

- E.g. – what makes Supersymmetry different to Universal Extra Dimensional models with Kaluza-Klein particles.
- One part of the answer:

**SPIN**



QUACK !



QUACK !

Not all things that quack are ducks!

# We will see two important themes:

- Mass measurements will precede<sup>(\*)</sup> spin determinations
- “Spin measurement”<sup>(\*\*)</sup> should not be confused with “sensitivity to spin”

(\*) or will at best be simultaneous with

(\*\*) Here “spin measurement” means “determining unambiguously the correct nature (scalar, fermion, vector) of one or more particles in a decay chain or model”

(more info at)

# A REVIEW OF SPIN DETERMINATION AT THE LHC

Lian-Tao Wang and Itay Yavin

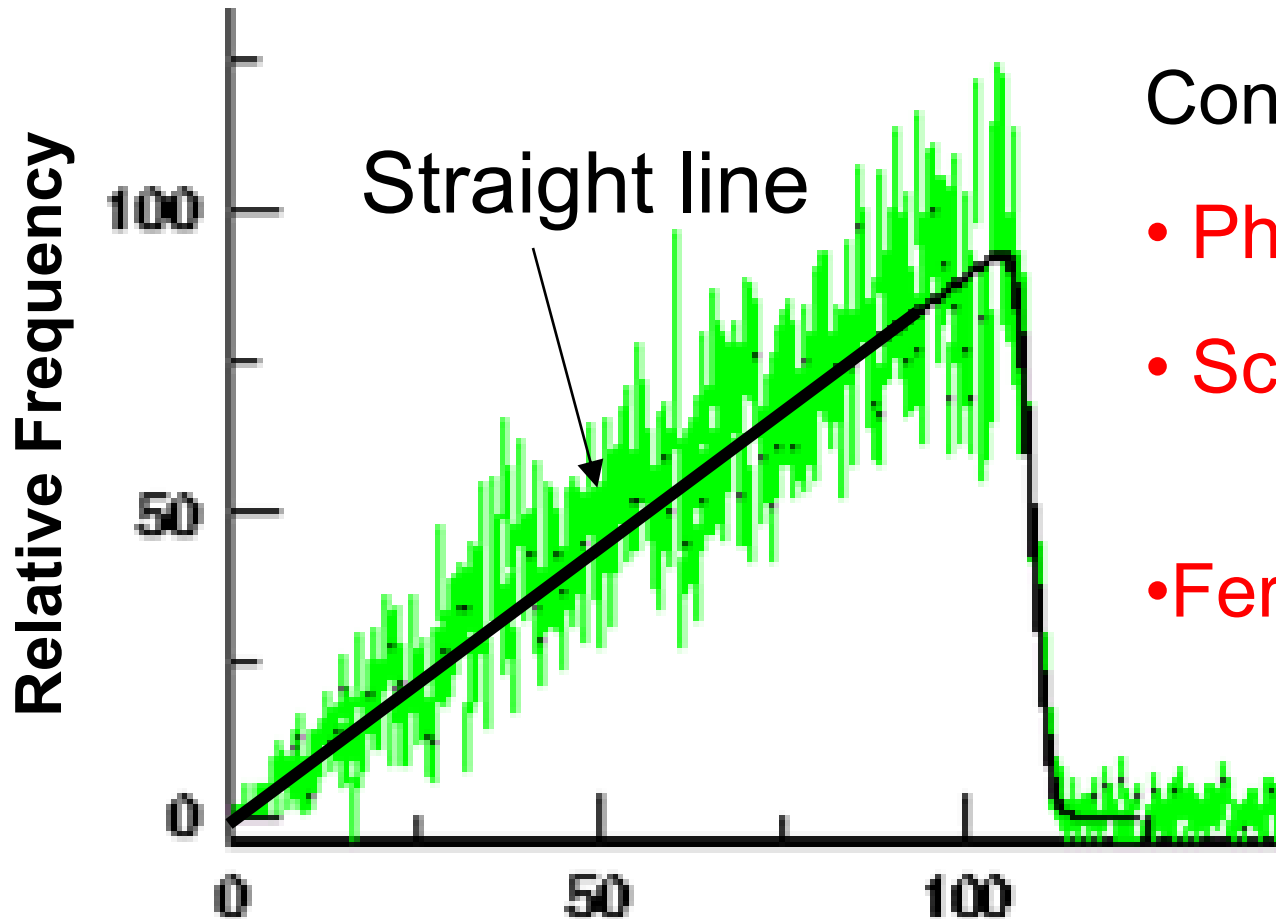
[arXiv:0802:2726](https://arxiv.org/abs/0802.2726)

# Spin determination topics

- Consistency checks
- Spins in “QLL chain”
  - A.Barr hep-ph/0405052
  - Smillie et al hep-ph/0605286
  - Florida etc [arXiv:0808.2472](https://arxiv.org/abs/0808.2472)
  - Biglietti et al ATL-PHYS-PUB-2007-004
- Slepton Spin (production)
  - A.Barr hep-ph/0511115
- MAOS method
  - Cho, Kong, Kim, Park arXiv:0810.4853
- Gluino chain spin
  - Alvez, Eboli, Plehn hep-ph/0605067
- Spins in chains with charginos
  - Wang and Yavin hep-ph/0605296
  - Smillie hep-ph/0609296
- Spins in chains radiating photons
  - Ehrenfeld et al arXiv:0904.1293



# Spin Consistency Check

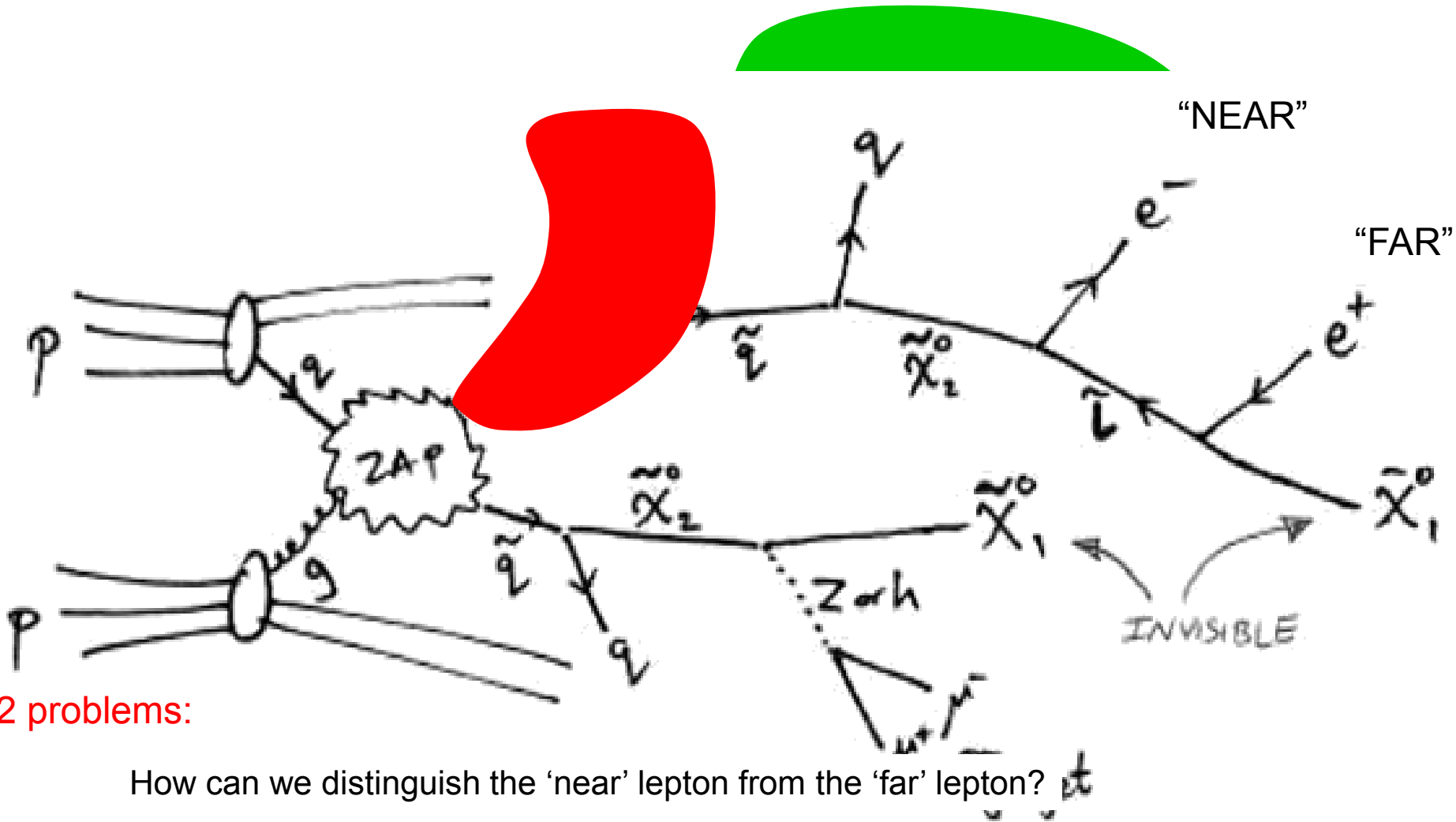


Consistent with:

- Phase-space
- Scalar slepton (SFSF)
- Fermion KK lepton (FVFFV)

**Di-Lepton Invariant Mass (GeV)**

# QL Spin Determination (A.Barr)



2 problems:

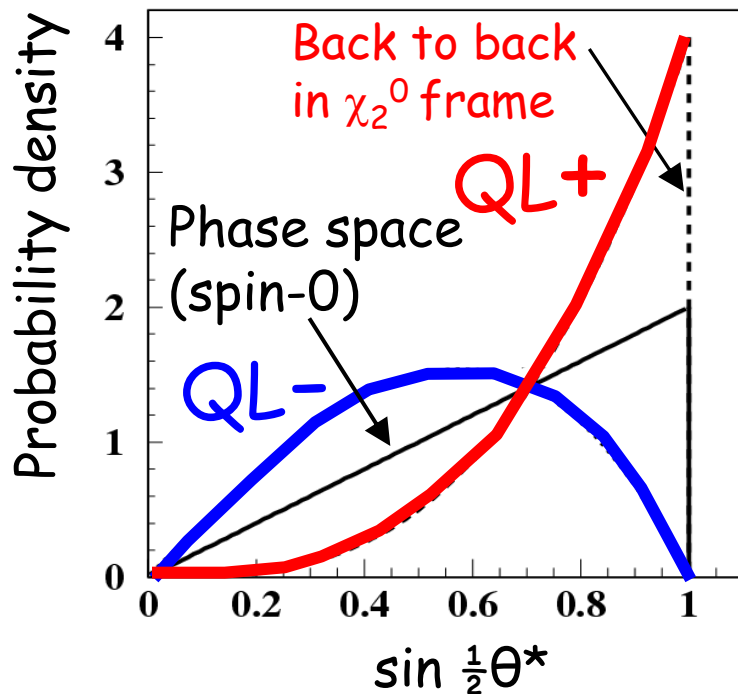
How can we distinguish the ‘near’ lepton from the ‘far’ lepton?

How can we tell  $l^+ q$  from  $l^+ \bar{q}$ ?

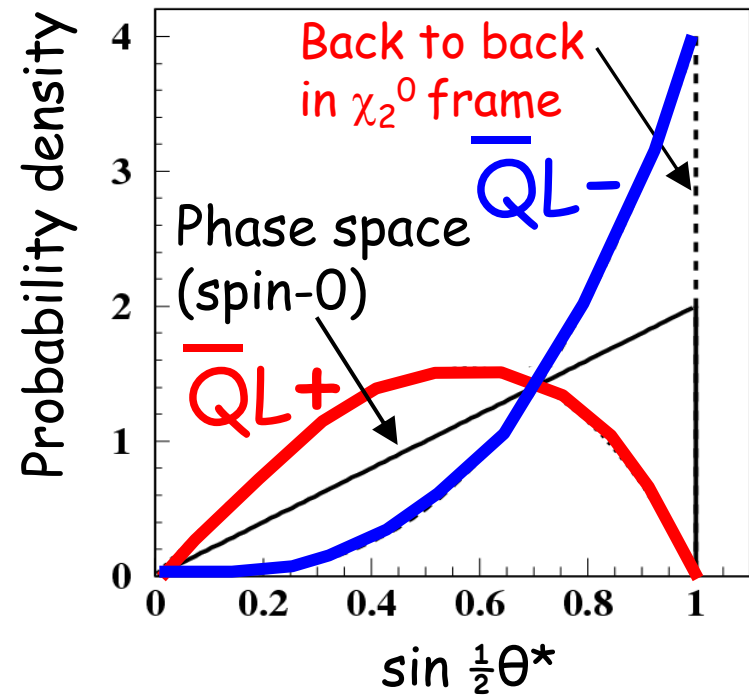


# Quark+NearLepton invariant mass distributions for:

**L+** **L-** and  
**QUARKS**



**L+** **L-** and  
**ANTI-QUARKS**

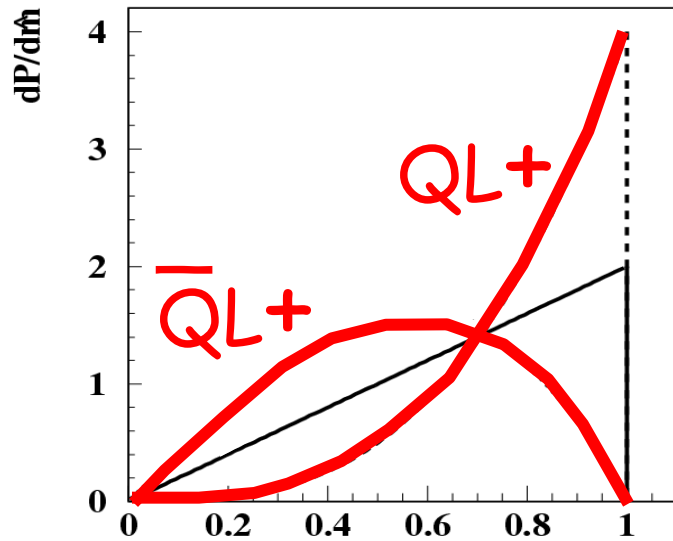


# Experimental problem

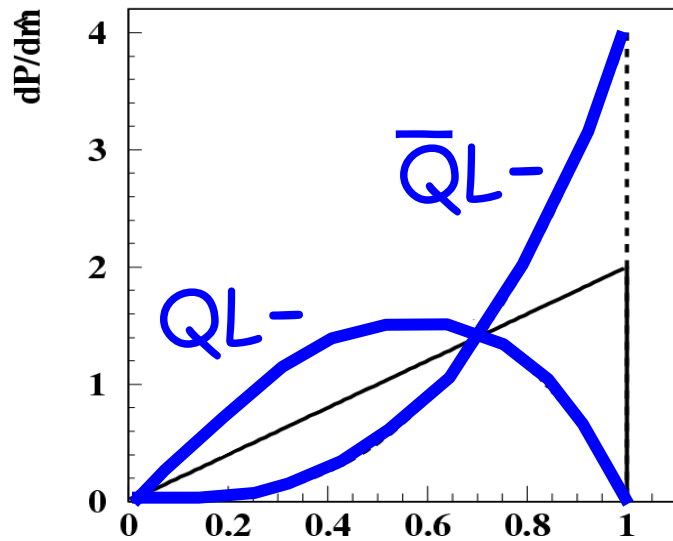
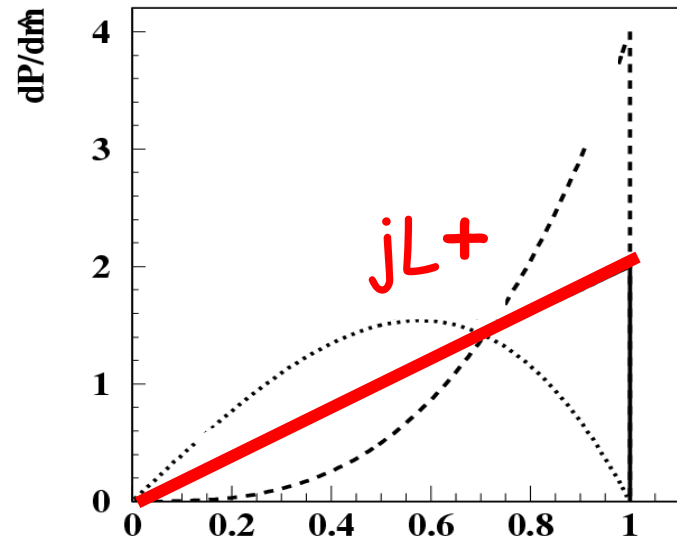
- Cannot reliably distinguish **QUARKs** from **ANTI-QUARKs**

Can only distinguish lepton charge  
**RED**( $Q_L^+$ ,  $\bar{Q}_L^+$ ) from **BLUE**( $Q_L^-$ ,  $\bar{Q}_L^-$ )

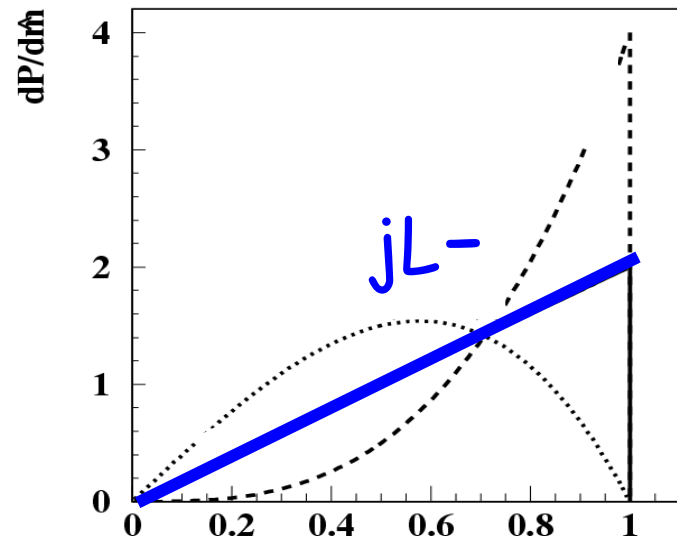
# Expect QUARK and ANTI-QUARK contributions to cancel:



SUM  $\rightarrow$

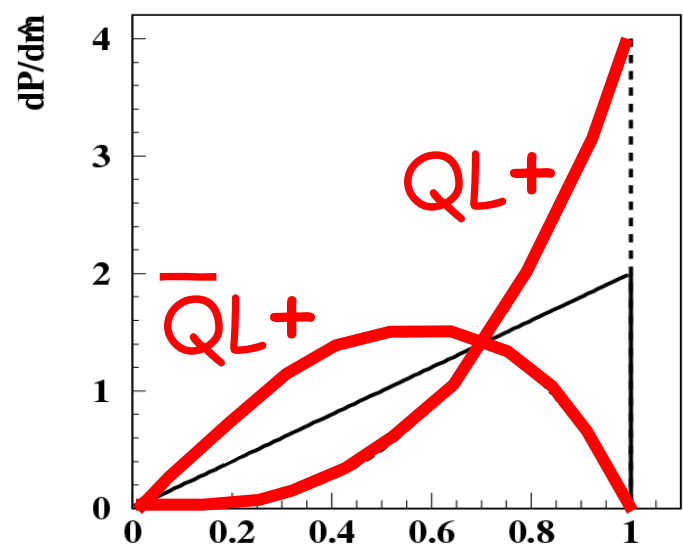


SUM  $\rightarrow$

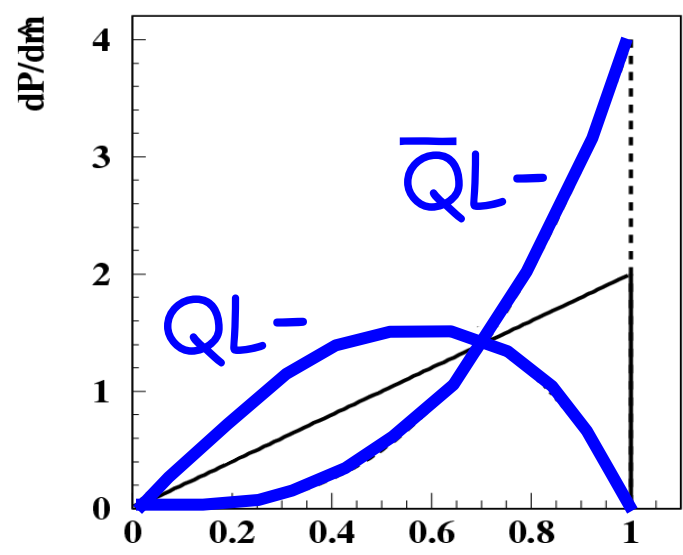
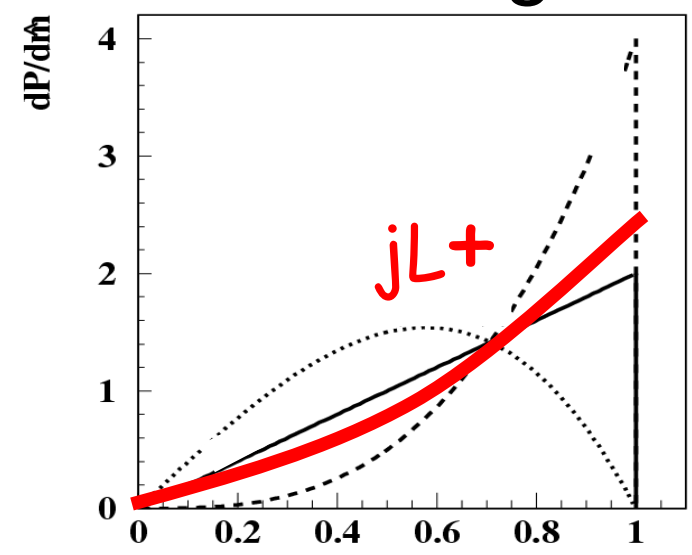


# But LHC is Proton-Proton machine

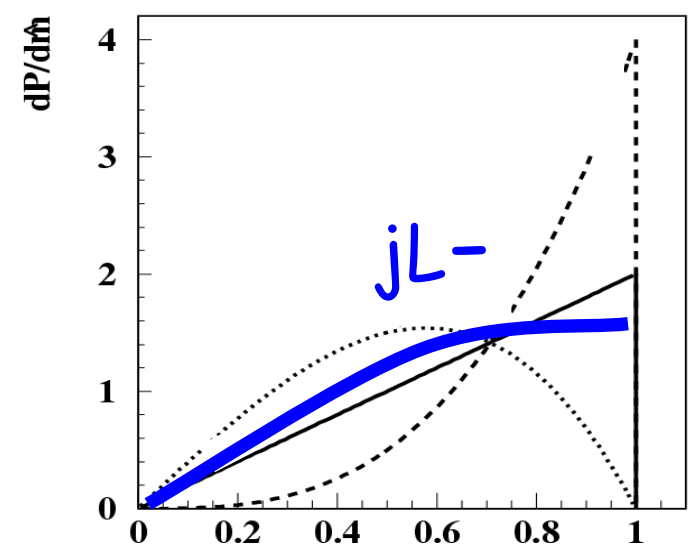
- More Quarks than Anti-Quarks! So get:



SUM →

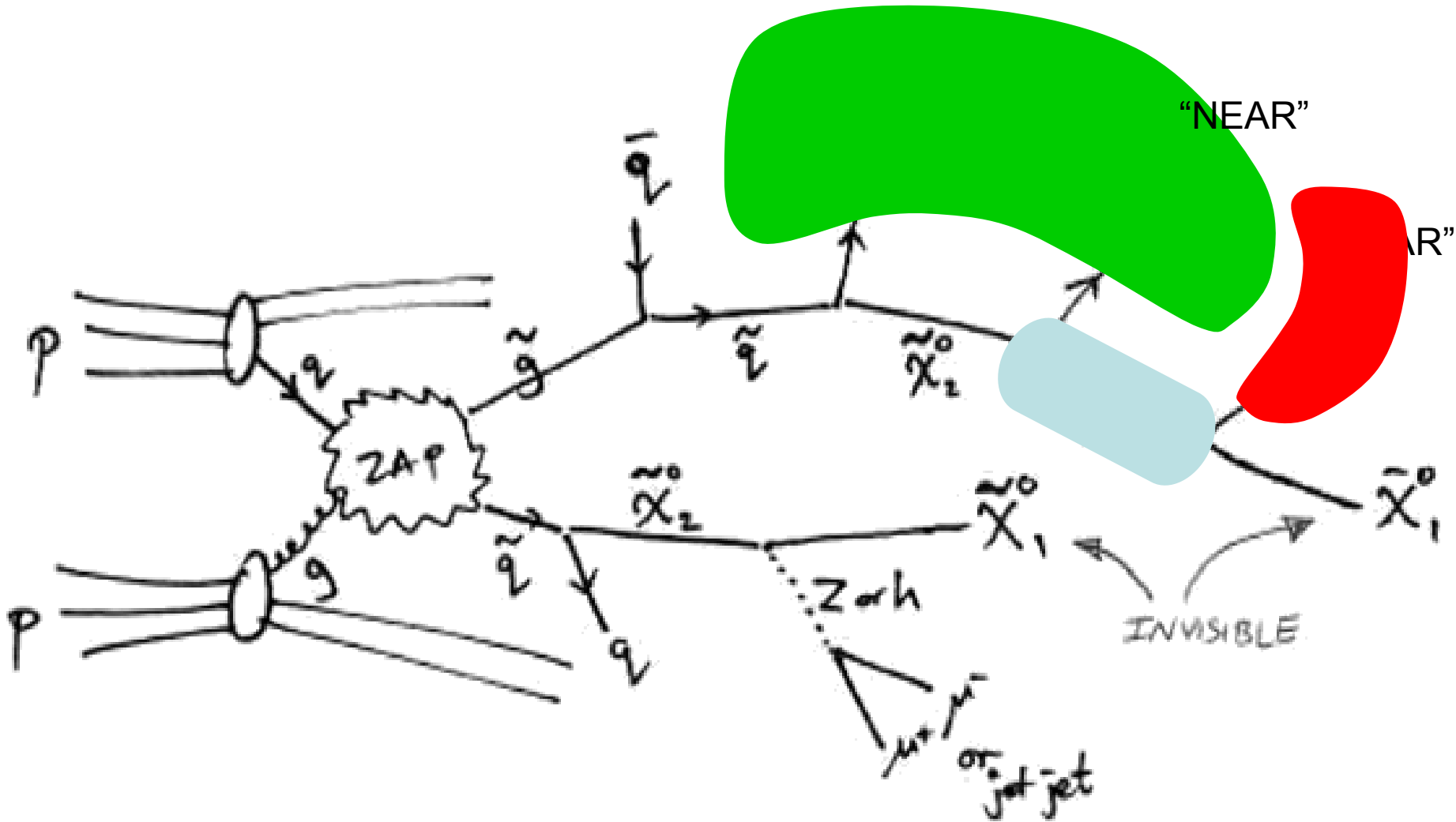


SUM →

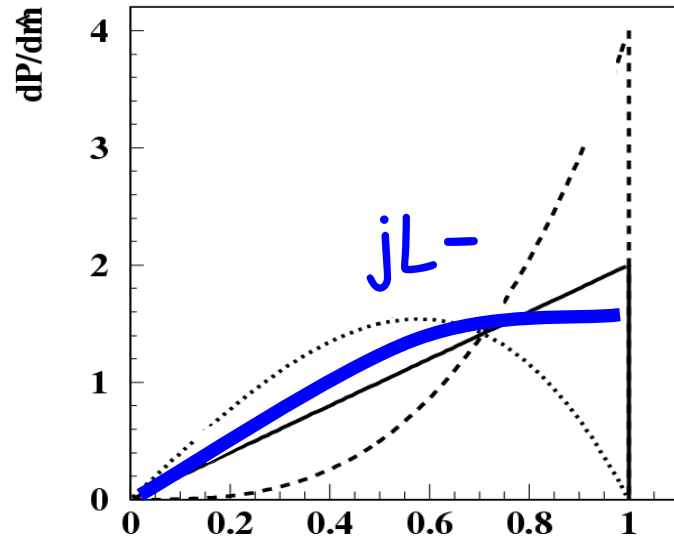
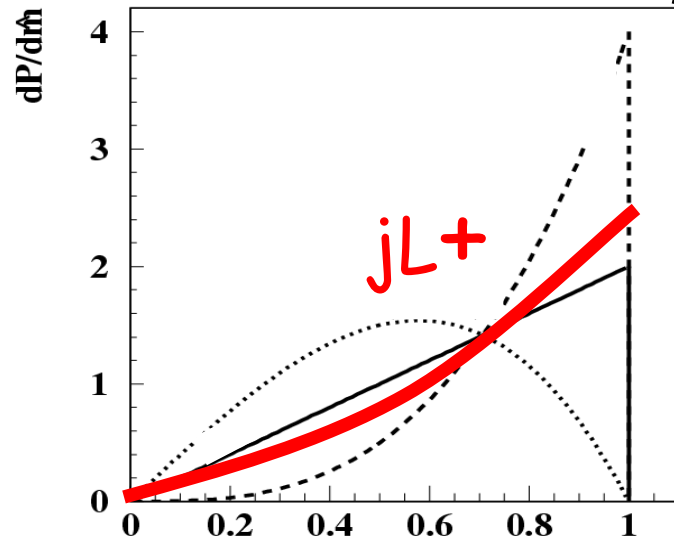


Asymmetry!

# “Far” Lepton washout?

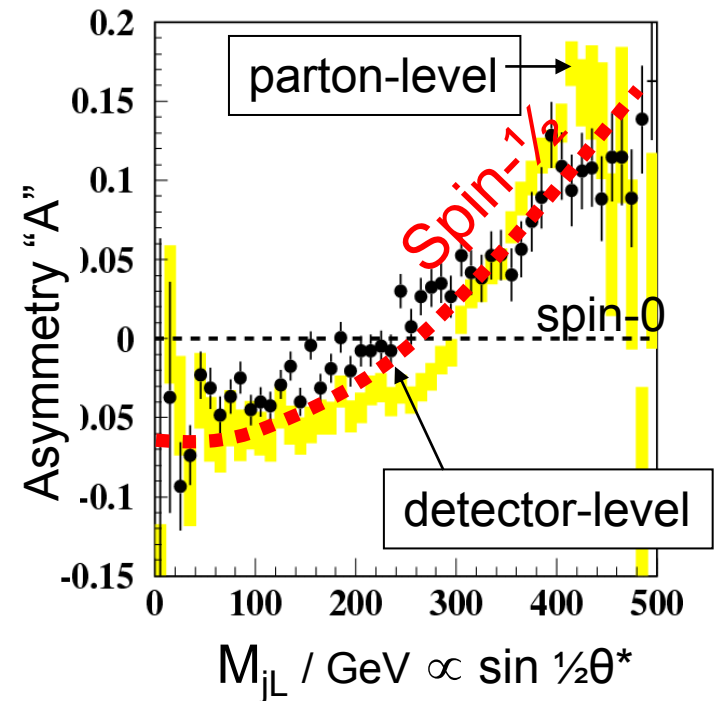


# So define $m_{jL^+}$ , $m_{jL^-}$ asymmetry



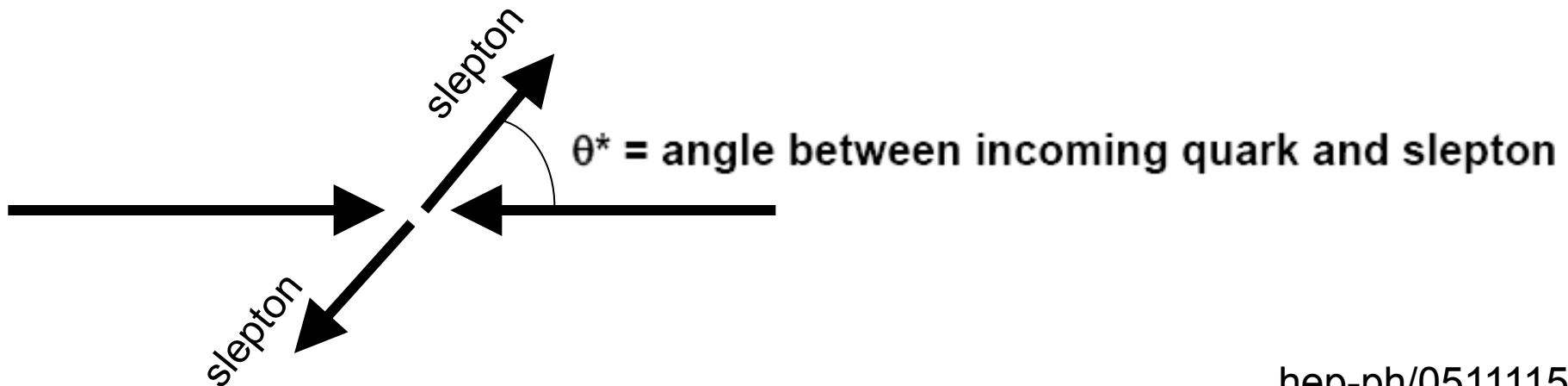
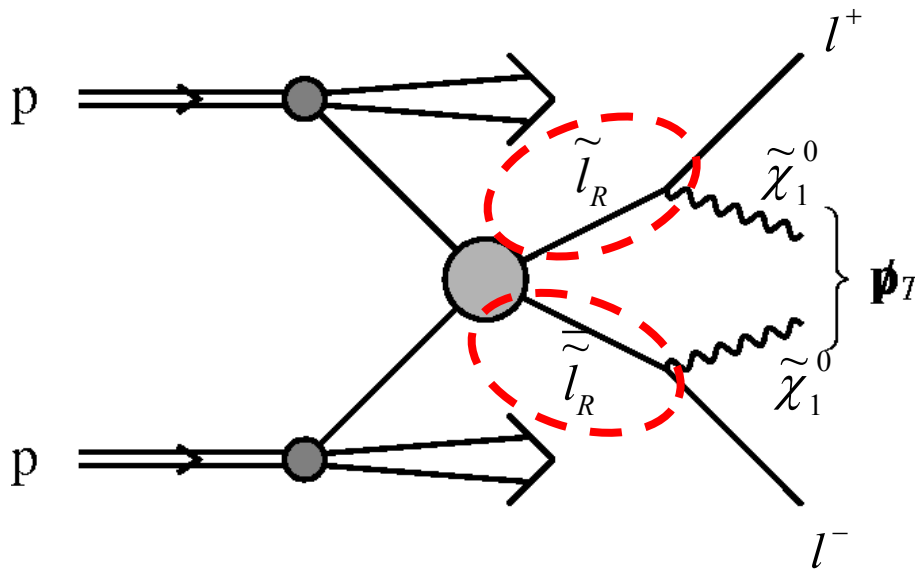
$$A = \frac{s^+ - s^-}{s^+ + s^-}$$

where  $s^\pm = \frac{d\sigma}{dm_{jL^\pm}}$



Different method altogether

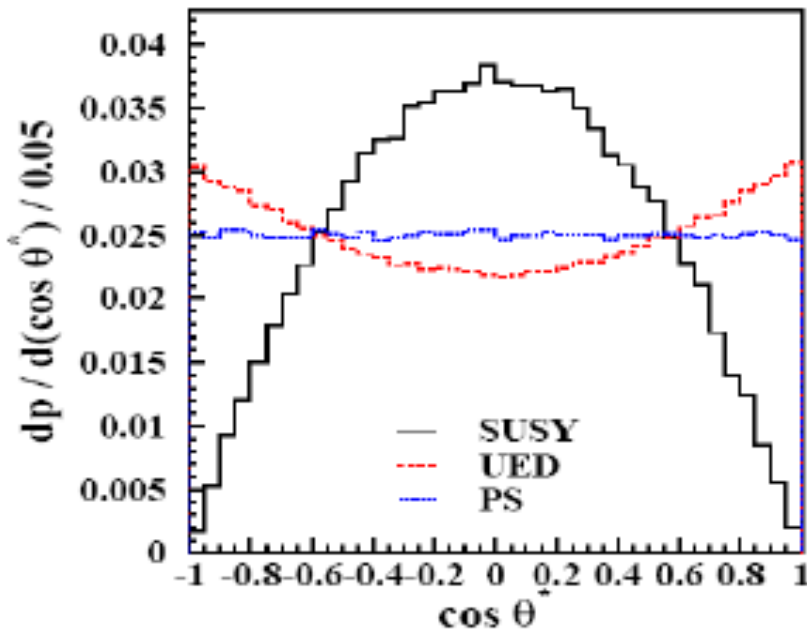
# Direct slepton spin detection: $qq \rightarrow Z\gamma^* \rightarrow \text{slepton slepton}$





# Look at slepton production angle in c.o.m.

- $\theta^*$  = angle between incoming quark and slepton



$$\left(\frac{d\sigma}{d\cos\theta^*}\right)_{\text{SUSY}} \propto 1 - \cos^2\theta^* \quad \text{SUSY = LHC point 5}$$

$$\left(\frac{d\sigma}{d\cos\theta^*}\right)_{\text{PS}} \propto \text{constant}$$

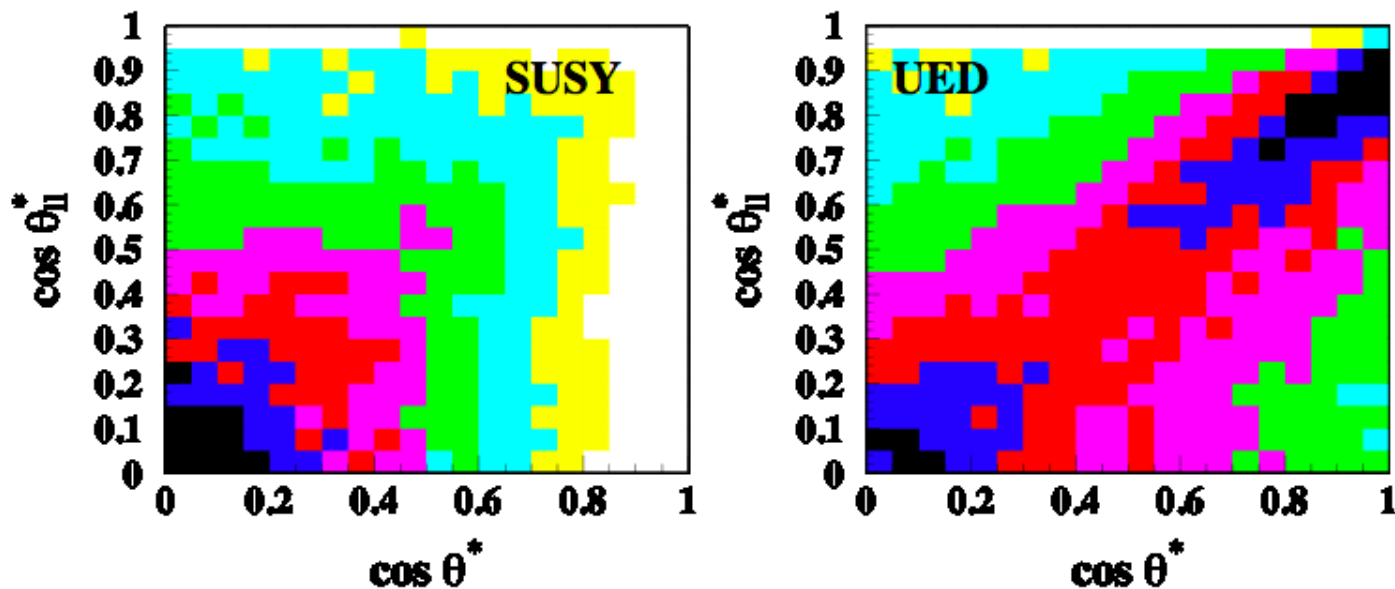
$$\left(\frac{d\sigma}{d\cos\theta^*}\right)_{\text{UED}} \propto 1 + \left(\frac{E_{\ell_1}^2 - M_{\ell_1}^2}{E_{\ell_1}^2 + M_{\ell_1}^2}\right) \cos^2\theta^*$$

Sensitive to spin, but can we measure  $\theta^*$ ?

**Define:**  $\cos\theta_u^* \equiv \cos\left(2 \tan^{-1} \exp(\Delta\eta_{e^+e^-}/2)\right) = \tanh(\Delta\eta_{e^+e^-}/2)$

# Have some access to desired angle

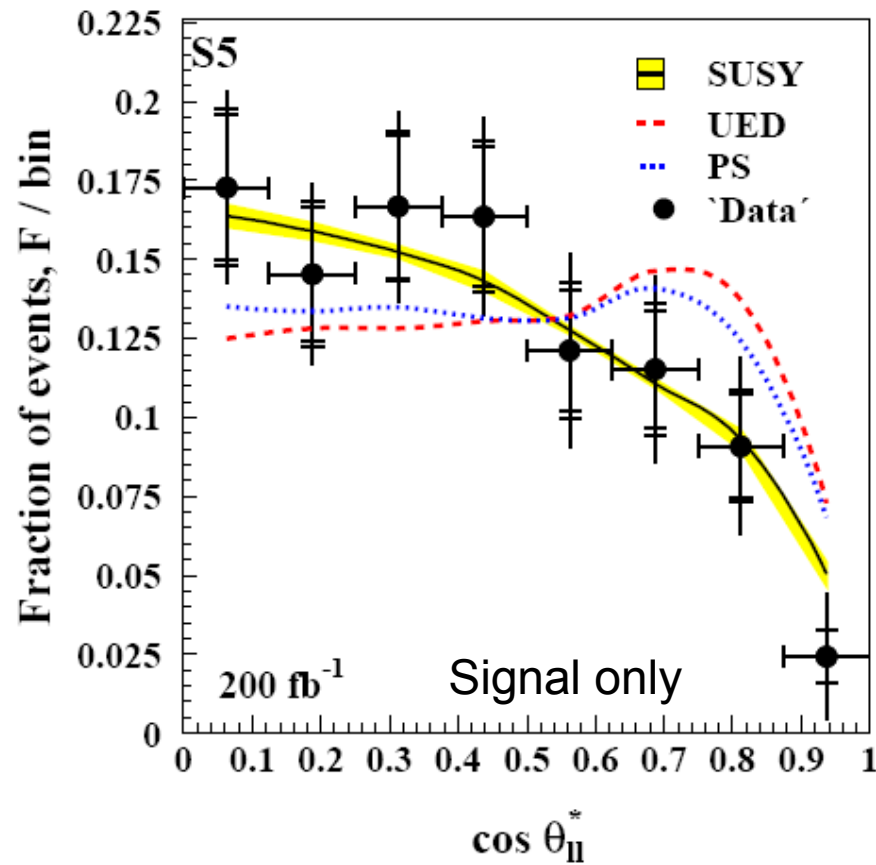
Distribution of  $\cos \theta_{ll}^* \equiv \tanh(\Delta\eta_{\ell^+\ell^-}/2)$   
is correlated with  $Z^0/\gamma$  decay angle  $\theta^*$



# Direct slepton spin (A.Barr)

hep-ph/0511115

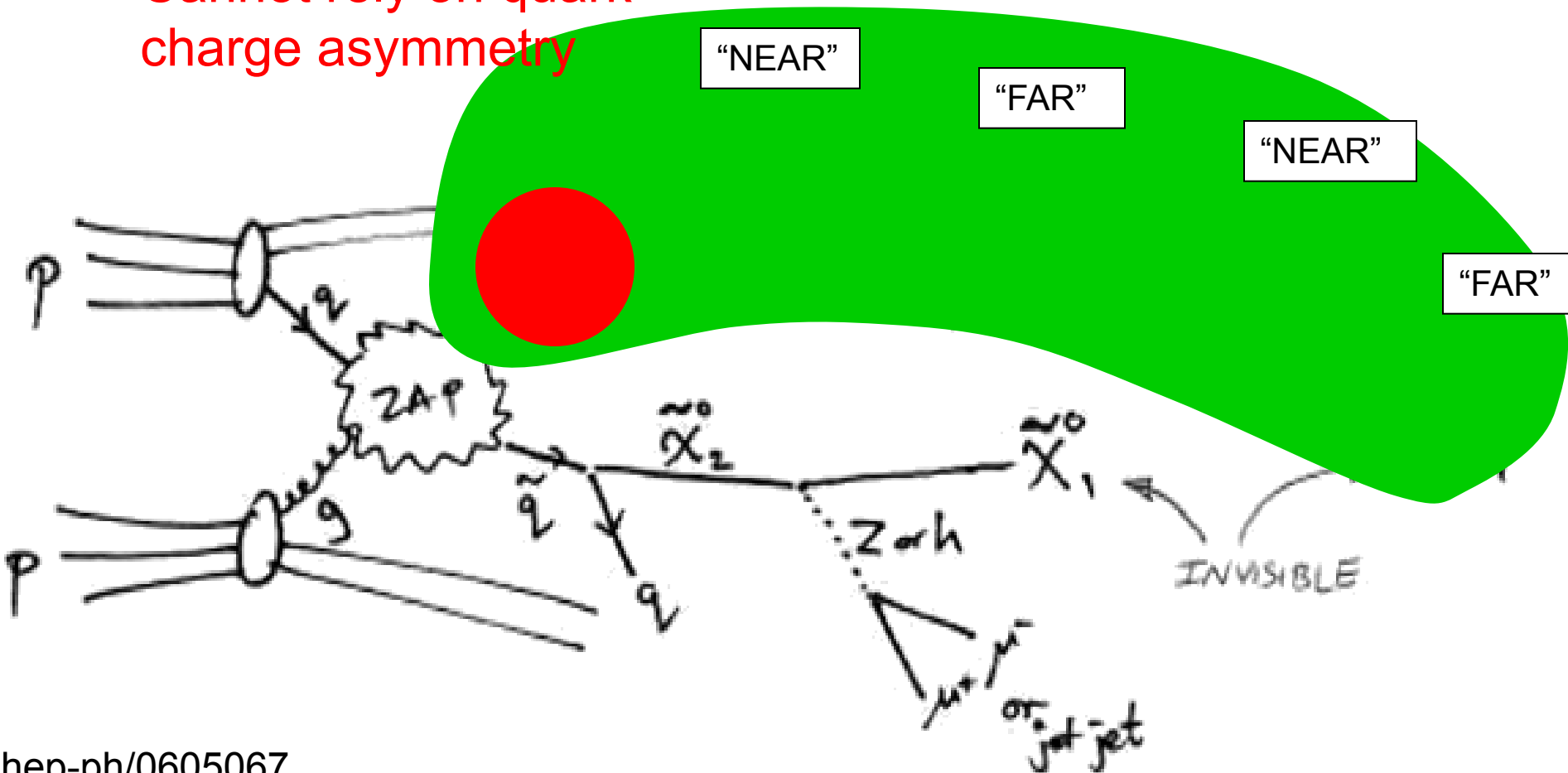
2 years high luminosity?



Different again

# Spin Determination (T.Plehn et.al.)

- What if we **want to investigate chain from gluino?**
- Crucial to test gluino nature
- **Cannot rely on quark charge asymmetry**

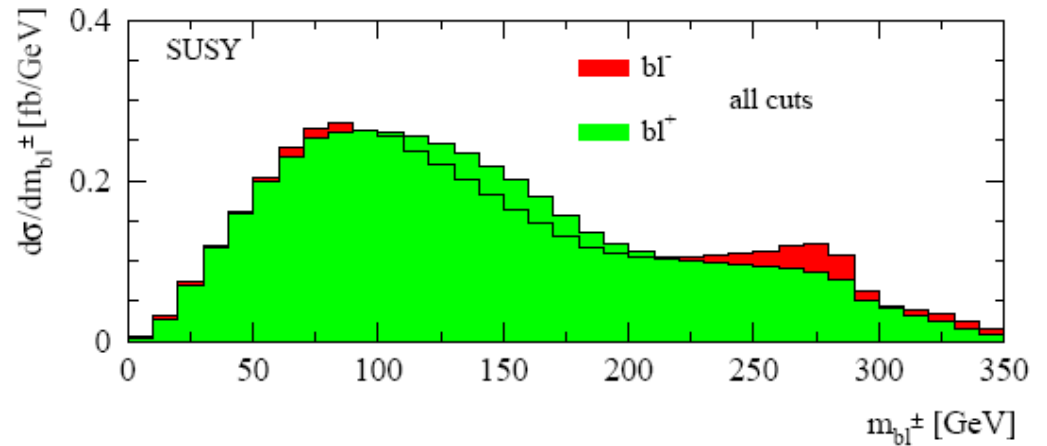




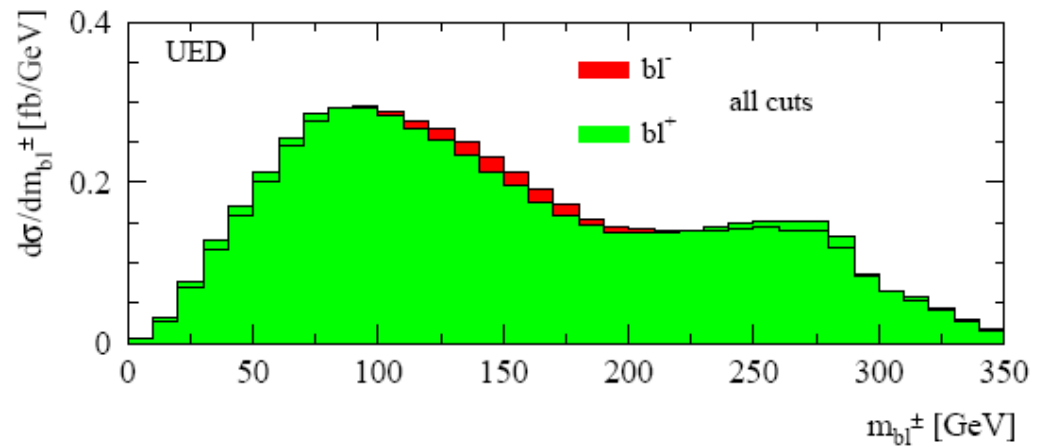


# $M_{BL}^+$ and $M_{BL}^-$ distributions

SUSY



UED



Room for an asymmetry!



# So define asymmetry

$$A^\pm(m_{bl}) = \frac{d\sigma/dm_{bl+} - d\sigma/dm_{bl-}}{d\sigma/dm_{bl+} + d\sigma/dm_{bl-}}$$

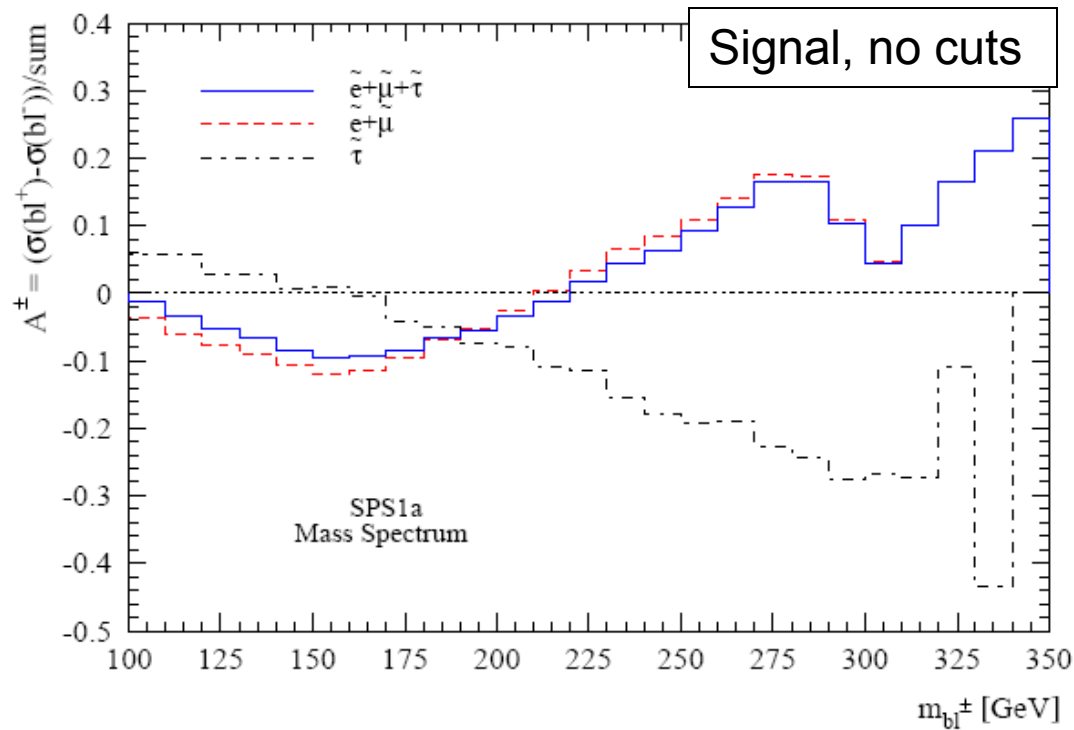
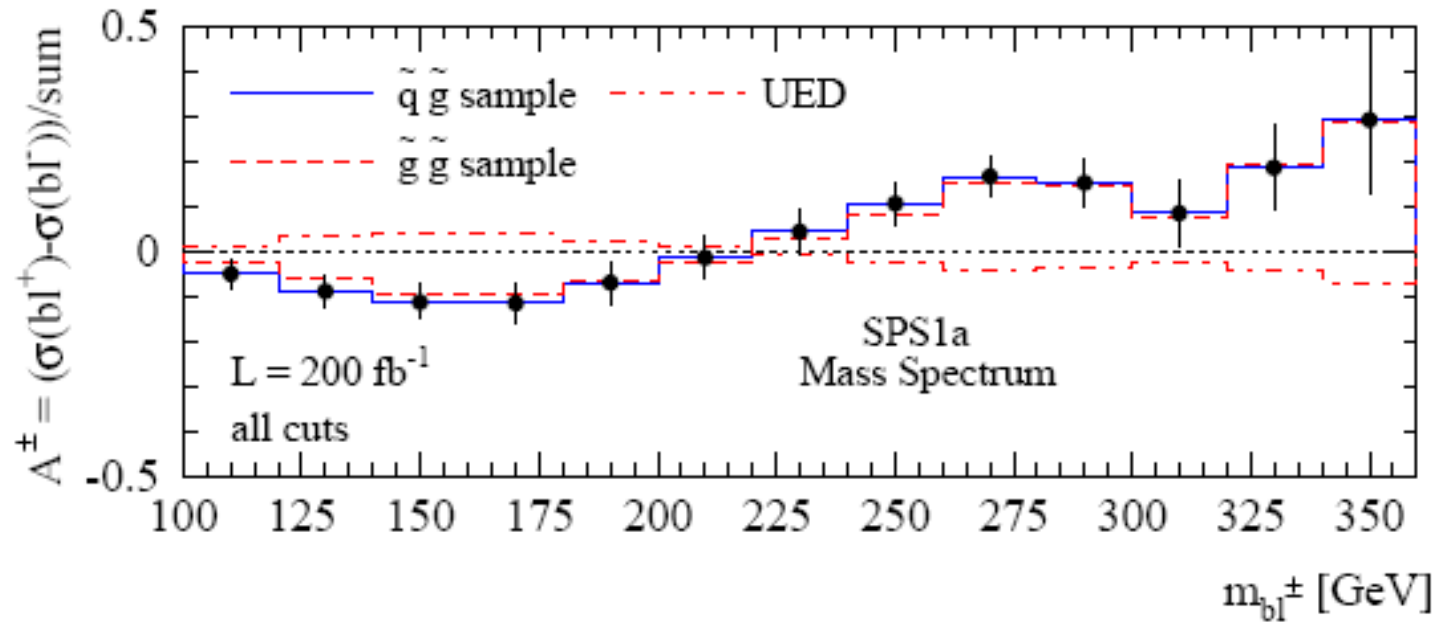


Figure 3: Bottom-lepton asymmetry for the SUSY signal only. The curves shown are for the first and second generation sleptons and for leptons coming from an intermediate  $\tilde{\tau}$ .

# After realistic cuts, SPS1A, 200 fb<sup>-1</sup>

Asymmetry  
still  
observable



Acceptance cuts:

$$\begin{aligned}
 p_{T,b} &> 50 \text{ GeV} & p_{T,\ell} &> 10 \text{ GeV} \\
 p_{T,j}^{\min} &> 40 \text{ GeV} & p_{T,j}^{\max} &> 150 \text{ GeV} \\
 |\eta_i| &< 2.4 & \Delta R_{ik} &> 0.4 \quad (i, k = b, j, \ell)
 \end{aligned}$$

Cuts to reject Standard Model

$$m_{\ell\ell} < 80 \text{ GeV} \quad M_{\text{eff}} > 450 \text{ GeV} \quad m_{jj} < 300 \text{ GeV}$$

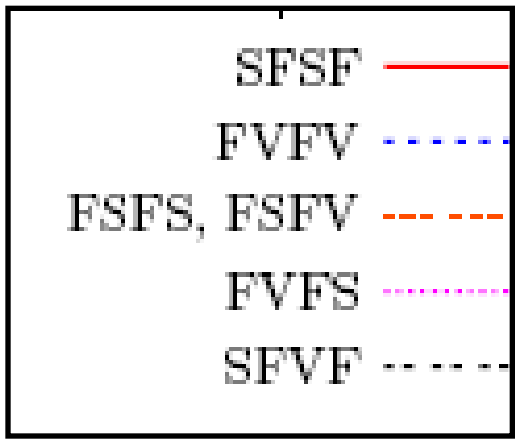
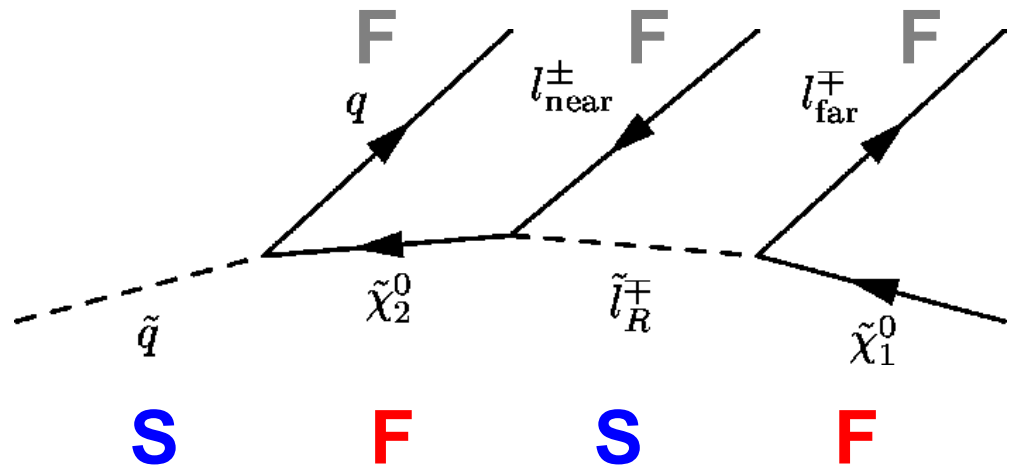
For a quantitative study we choose the (collider friendly) parameter point SPS1a. The masses in the gluino decay cascade are  $m_{\tilde{g}} = 608 \text{ GeV}$ ,  $m_{\tilde{b}_1} = 517 \text{ GeV}$ ,  $m_{\tilde{b}_2} = 547 \text{ GeV}$ ,  $m_{\tilde{\chi}_2^0} = 181 \text{ GeV}$ ,  $m_{\tilde{\ell}_1} = 145 \text{ GeV}$ ,  $m_{\tilde{\ell}_2} = 202 \text{ GeV}$ ,  $m_{\tilde{\tau}_1} = 136 \text{ GeV}$ ,  $m_{\tilde{\tau}_2} = 208 \text{ GeV}$ , and  $m_{\tilde{\chi}_1^0} = 97 \text{ GeV}$ . The NLO production cross sections are 7.96 pb for  $\tilde{g}\tilde{g}$ , 8.02 pb for  $\tilde{q}\tilde{q}^*$ , 26.6 pb for  $\tilde{q}\tilde{g}$ , and 7.51 pb for  $\tilde{q}\tilde{q}$ . For the SPS1a parameter choice the lighter of the

Back to long chains

# Spin sensitivity elsewhere in the llq chain (Smillie et.al.)

## Later more general follow-up (Matchev, Kong, et al)

$$\bar{\Psi}_F (g_L P_L + g_R P_R) \Psi_f \Phi + h.c.$$



Scalar	Fermion	Scalar	Fermion
Fermion	Vector	Fermion	Vector
Fermion	Scalar	Fermion	Scalar
Fermion	Vector	Fermion	Scalar
Fermion	Scalar	Fermion	Vector
Scalar	Fermion	Vector	Fermion

Cannot distinguish:  
 {FSFS, FSFV} and {FVFS, FVFV}

hep-ph/0605286 arXiv:0808.2472

# But masses matter

**SPS1a** mass spectrum: (GeV)

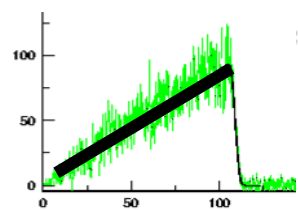
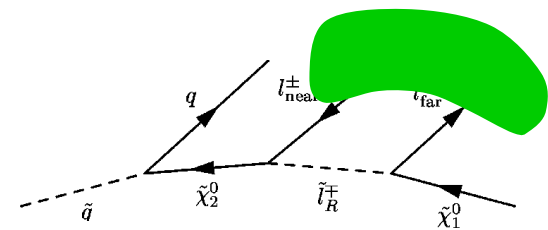
$A$	$B$	$C$	$D$
$\tilde{\chi}_1^0$	$\tilde{e}_R$	$\tilde{\chi}_2^0$	$\tilde{u}_L$
96	143	177	537

**UED-type** mass spectrum: (GeV)

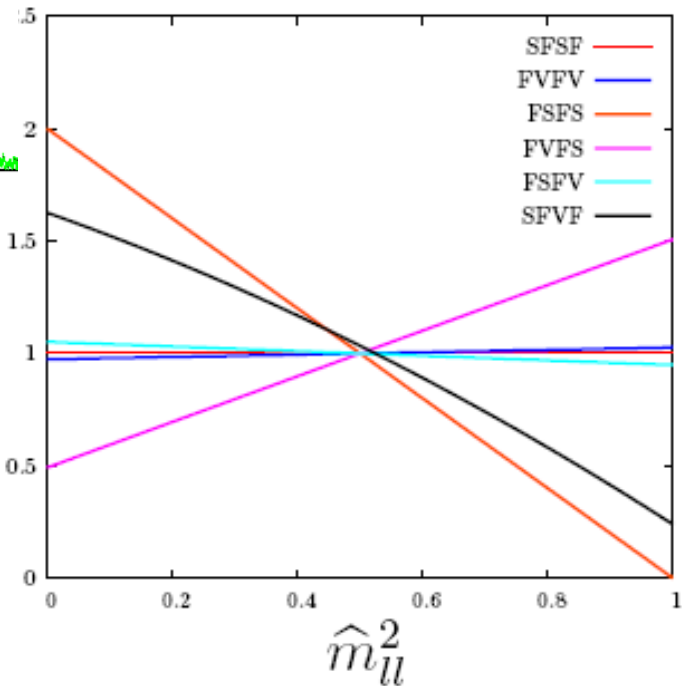
( $R^{-1} \sim 800$  GeV)

$A$	$B$	$C$	$D$
$\gamma^*$	$l_L^*$	$Z^*$	$q_L^*$
800	824	851	956

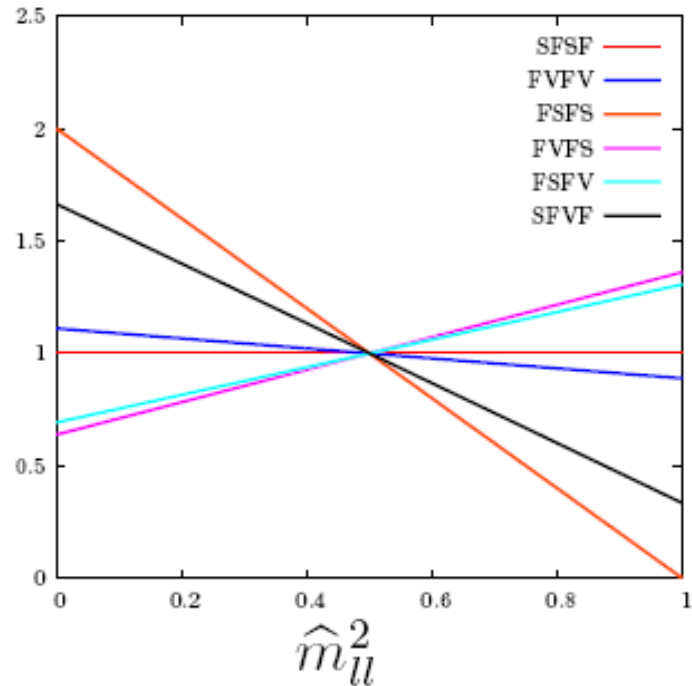
# Maybe masses are not too important for $m_{ll}$ distribution



$$\frac{dP}{d\hat{m}_{ll}^2}$$



SPS1a masses



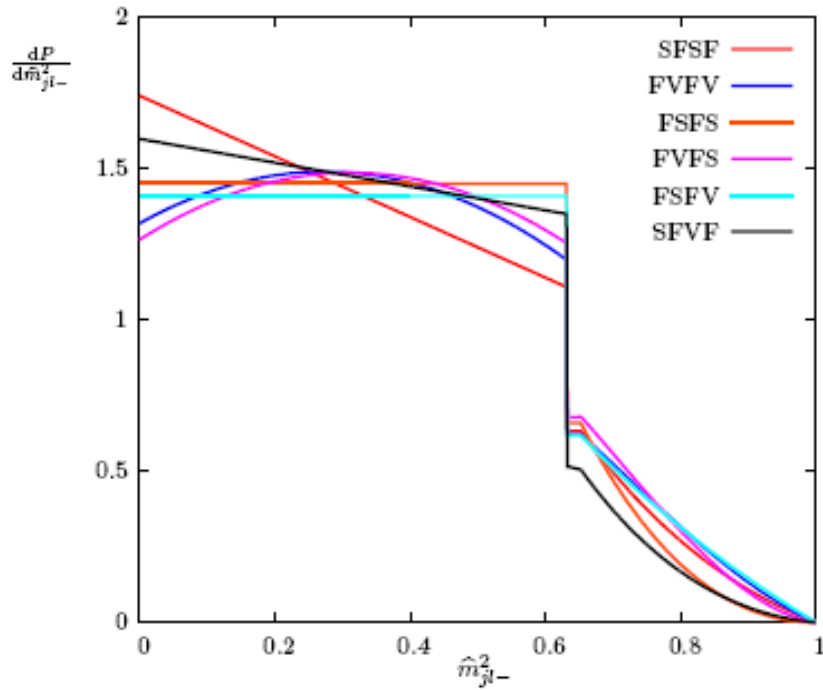
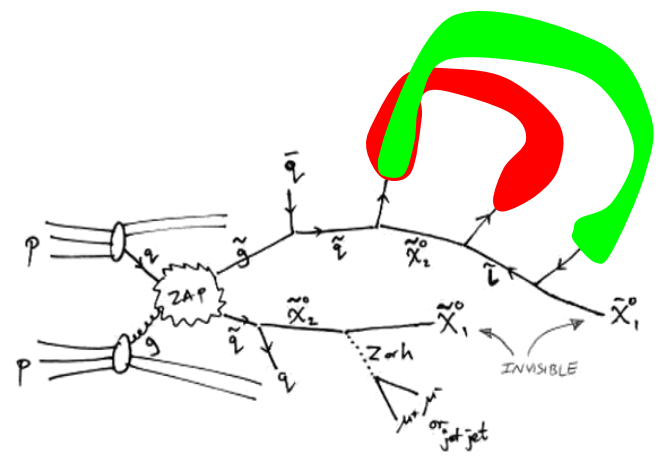
UED type masses

hep-ph/0605286

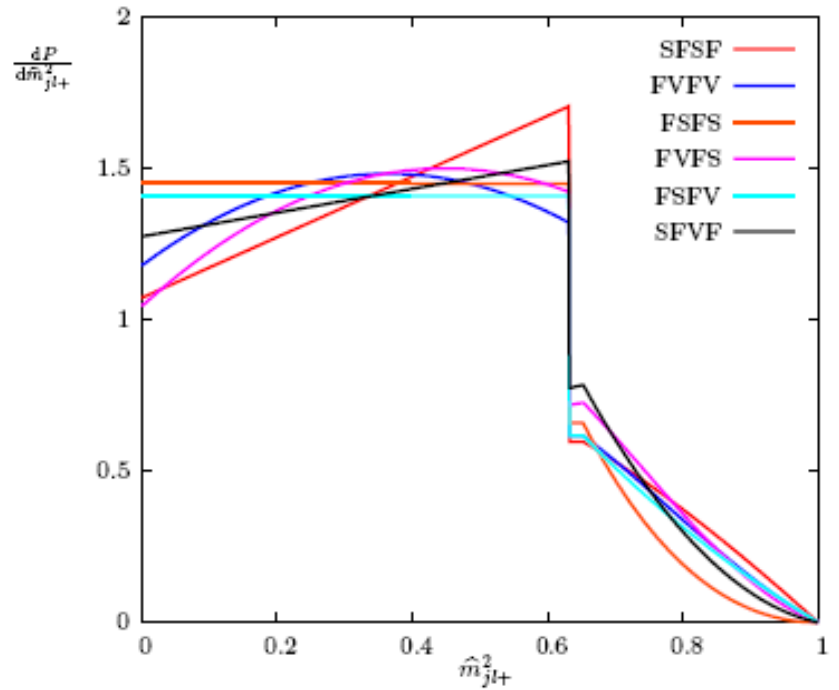
... but this fun ....

jet +  $l^\pm$

At SPS 1a:

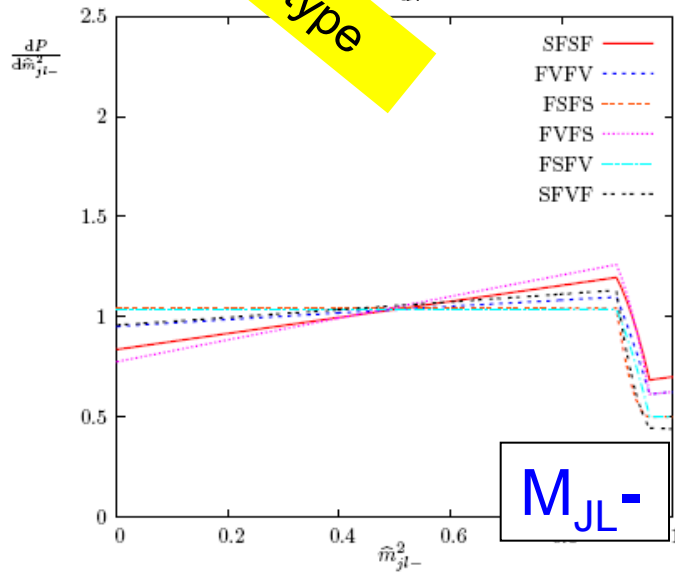
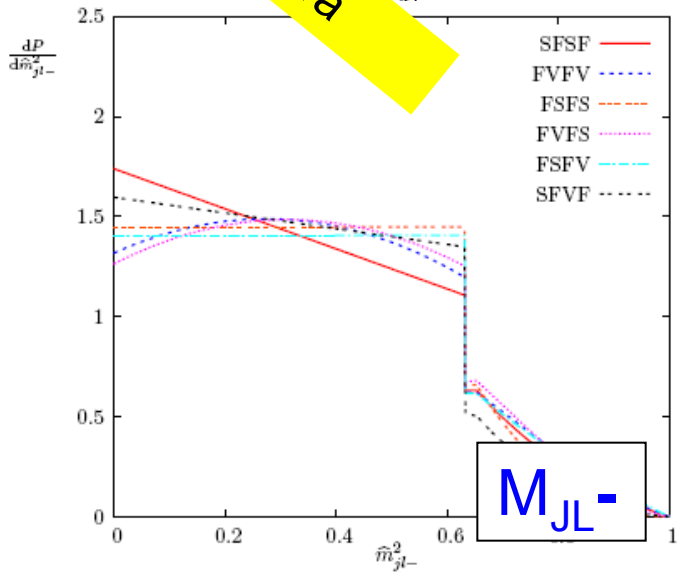
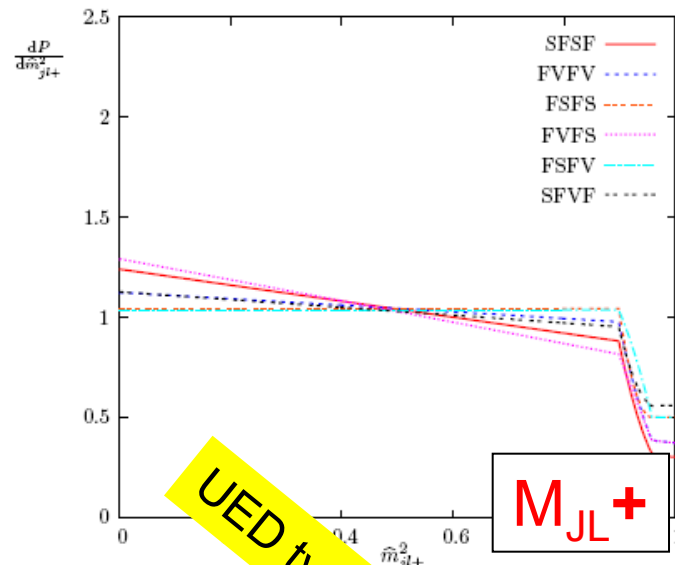
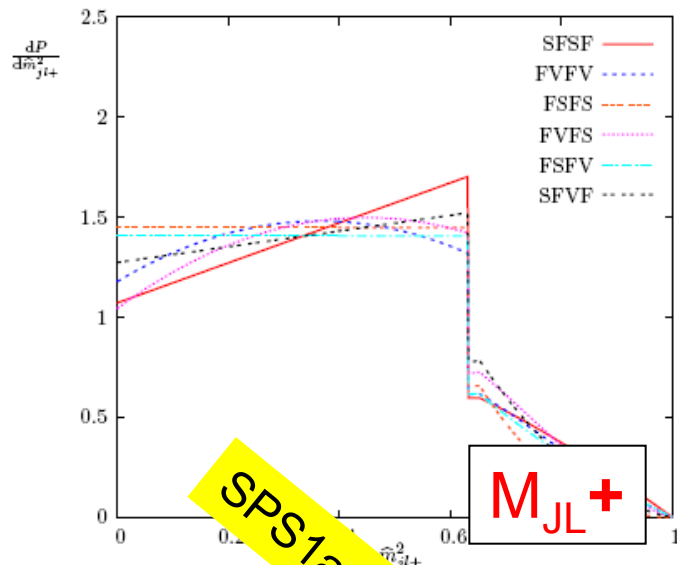


jet +  $l^-$



jet +  $l^+$

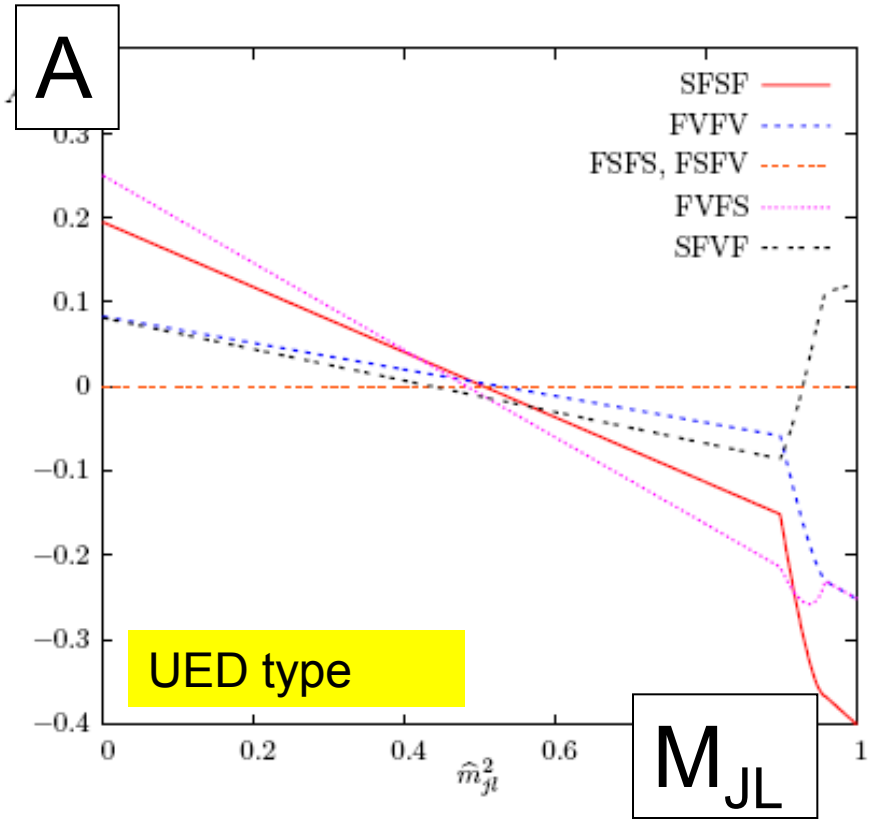
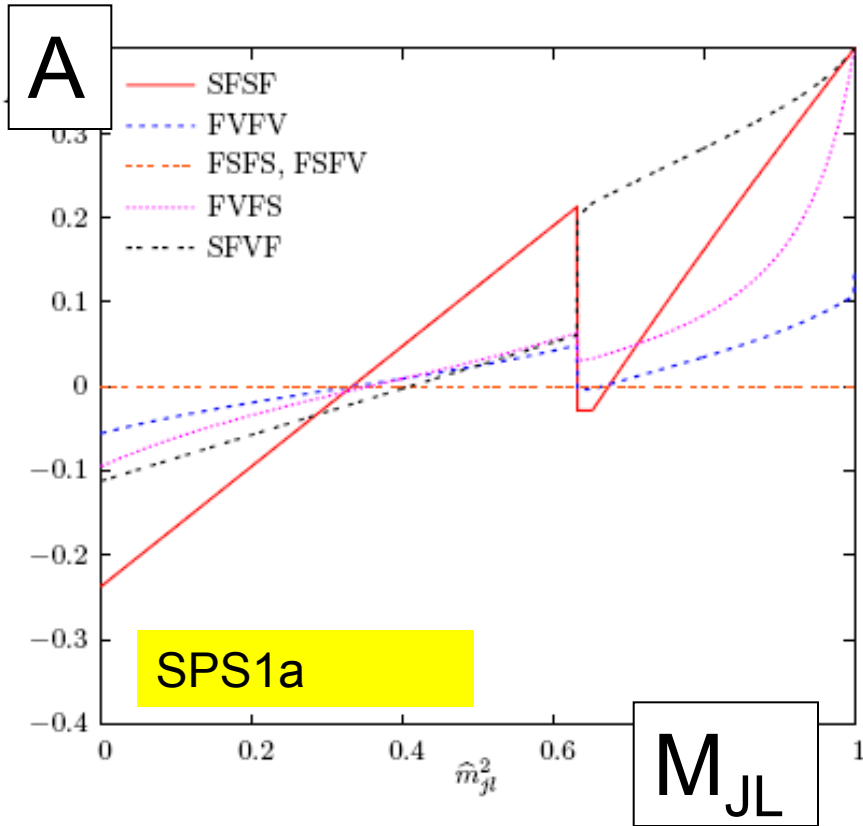
.... is spoiled. ☹️





# Example asymmetries:

(a big mix of spin and mass spectrum) ☹️



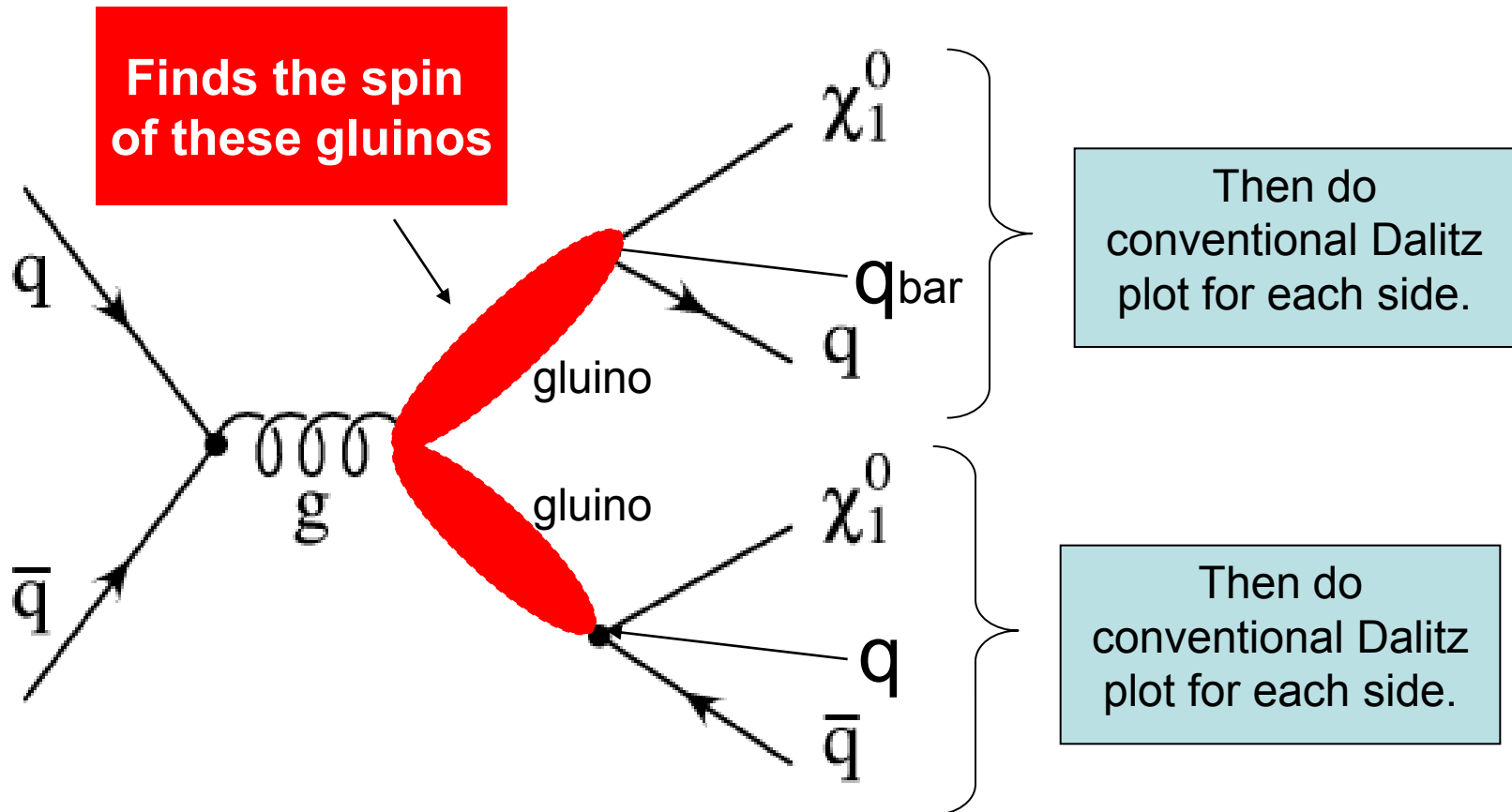
$$A = \frac{\text{Red Box} - \text{Blue Box}}{\text{Red Box} + \text{Blue Box}}$$

Yet another game one can play

# $M_{T2}$ -assisted (MAOS) spin determination

$$pp \rightarrow Y(1) + \bar{Y}(2) \rightarrow V(p_1)\chi(k_1) + V(p_2)\chi(k_2), \quad Y \rightarrow q(p_q)\bar{q}(p_{\bar{q}})\chi(k):$$

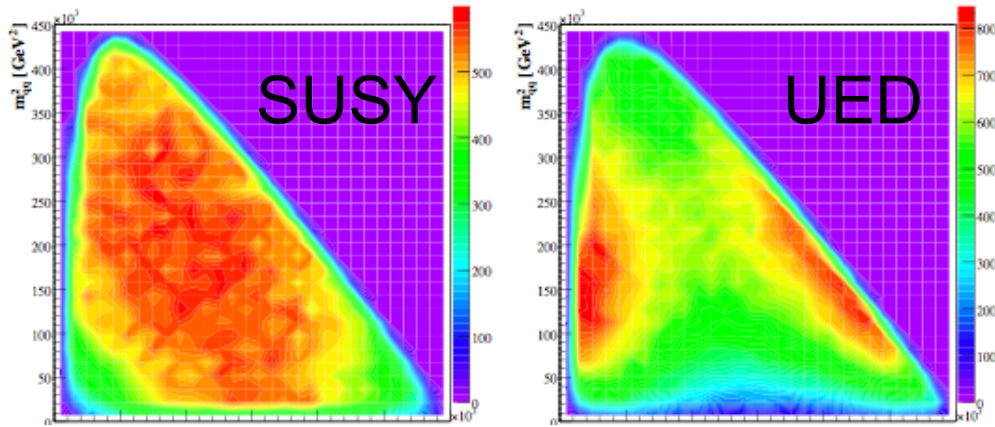
Use splitting for which leads to  $M_{T2}$  solution to assign 4-momenta to invisible particles:



# $M_{T2}$ -assisted (MAOS) spin determination

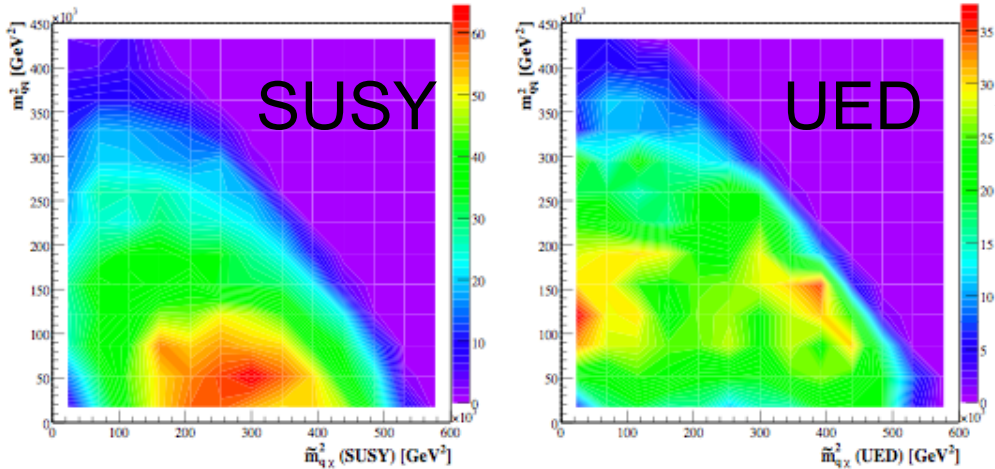
$$pp \rightarrow Y(1) + \bar{Y}(2) \rightarrow V(p_1)\chi(k_1) + V(p_2)\chi(k_2), \quad Y \rightarrow q(p_q)\bar{q}(p_{\bar{q}})\chi(k).$$

$$M_{T2}(p_i, m_\chi) \equiv \min_{\mathbf{k}_{1T} + \mathbf{k}_{2T} = \mathbf{p}_T^{\text{miss}}} [\max\{M_T^{(1)}, M_T^{(2)}\}] \rightarrow \text{assign 4-momenta}$$



$$m_{\chi, Y} = m_{\chi, Y}^{\text{true}}$$

$$\mathcal{L} = \infty$$

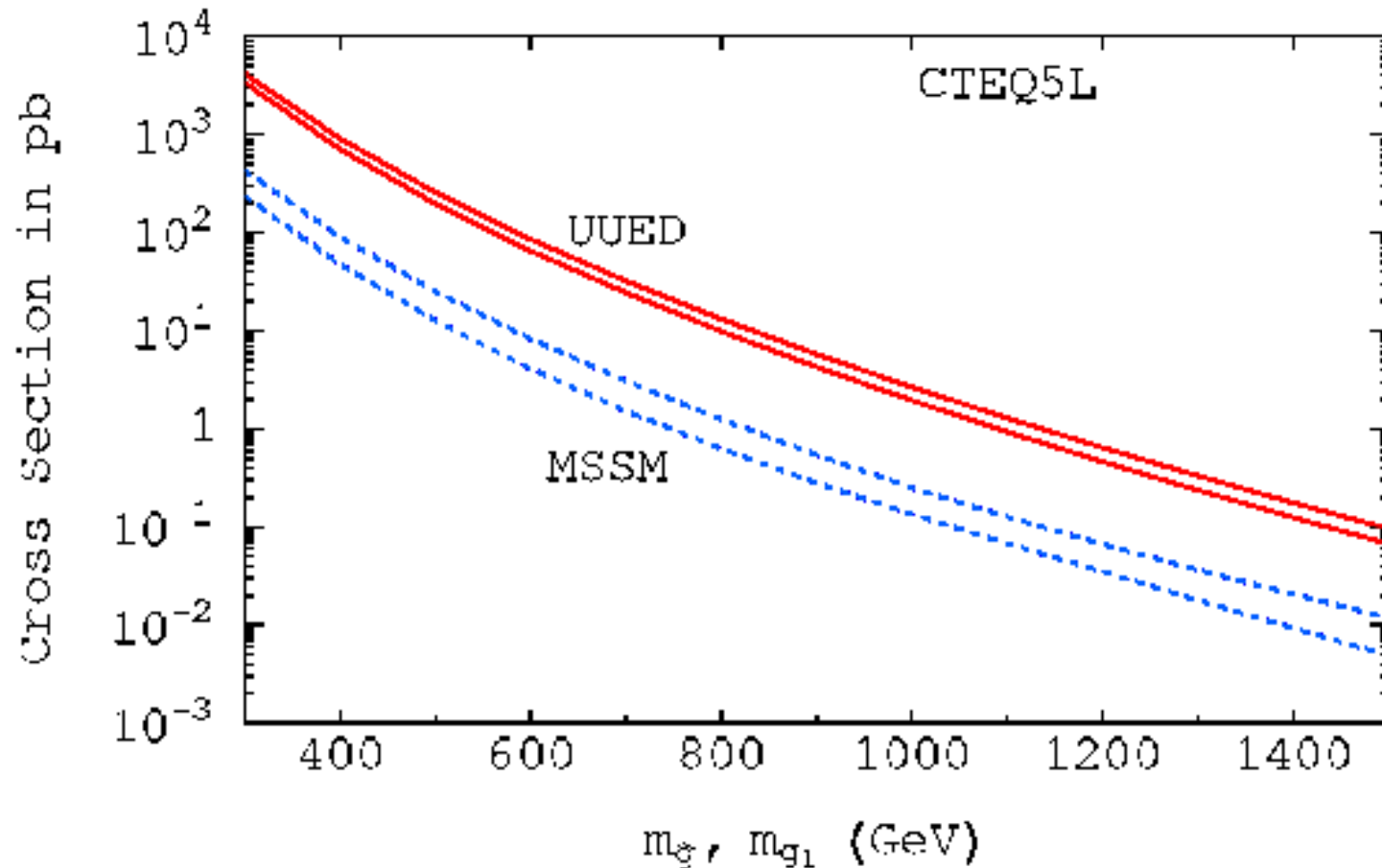


$$m_\chi = 0, m_Y = M_{T2}^{\text{max}}(m_\chi = 0)$$

$$\mathcal{L} = 300 \text{ fb}^{-1}$$

Cho, Choi, Kim, Park, 0810.4853

# Reminder: cross sections reveal spins



➔ Higher spins mean higher cross sections  
(for given masses)

# End Notes

- QLL chain
  - Some spin “sensitivity” – but no strong UED/SUSY separation
  - Reduced discriminatory power when considering general couplings (Matchev/Kong).
- Di-slepton production
  - Better chance of separating UED/SUSY
  - Still model dependent
- Both require large cross sections
- Masses inextricably intertwined.

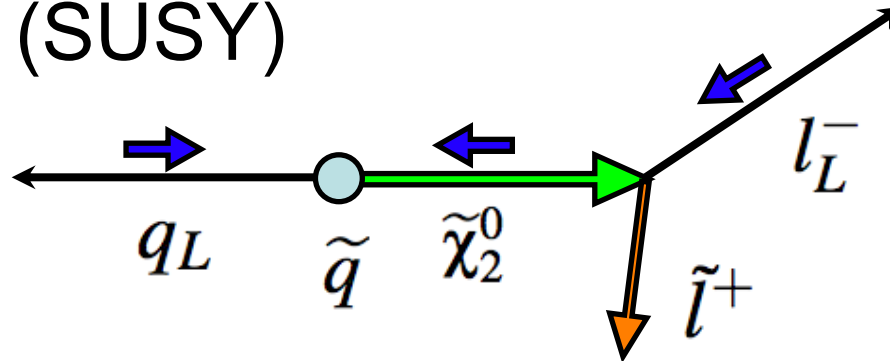


**Backup slides**

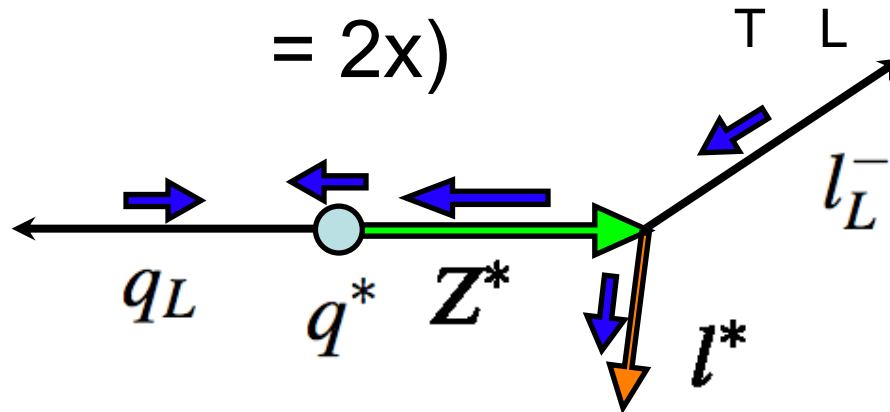


# Helicity dependence

○ Process 1 (SUSY)



○ Process 1 (UED, transverse  $Z^*$ :  $P/P = 2x$ )



⇒ Both prefer high  $(ql^-)^{near}$  invariant mass