Top Quark Phenomenology:

ADD model and the Minimal Length Scenario

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NExT PhD school, Cosener's House 18/07/2011

Overview

- ADD Model
 - Background
 - Particles and coupling to SM
 - Virtual gravitons
- Minimal Length Scenario: natural regulator
- t-tbar channel at collider experiments
 - precision LHC discovery channel
- ADD+MLS in the t-tbar channel: LHC reach to set bounds
- Outlook & Summary

The ADD Model

- Alternative viewpoint on the gauge hierarchy problem proposed by Arkani-Hamed, Dimopoulos and Dvali (arXiv:hep-ph/9803315)
 - δ compact extra dimensions
 - Fundamental Planck scale is very low ~ O(TeV)
 - SM degrees of freedom localised on a 3-brane embedded in the bulk
- 'Large' extra dimensional volume allows the fundamental Planck scale to be as low as the electroweak scale

$$M_4^2 = M_D^{2+\delta} V_\delta$$

- Simplest realisation of a model with transverse extra dimensions
- Toroidal compactification: $\mathcal{M}_4 \times T_\delta$
- Common radius R

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The ADD Model Particles and Coupling

Linearised gravity in D dimensions – degrees of freedom described by fluctuations of the metric about flatness:

$$g_{MN} = \eta_{MN} + \kappa h_{MN}; \quad \kappa^2 = \frac{1}{M_D^{2+\delta}} \qquad h_{MN} = \begin{pmatrix} \frac{G_{\mu\nu}}{A_{i\nu}} & \frac{A_{\mu j}}{\phi_{ij}} \end{pmatrix}$$

Kaluza Klein reduction leads to a 4D effective action for the Lorentz group reps into which the D dimensional tensor decomposes

$$h_{MN}(x^{\mu}, y^{i}) = \sum_{\vec{n}} h_{MN}(x^{\mu}) \exp\left(i\frac{2\pi n \cdot y}{R}\right)$$

- Massless zero modes
- Each (non zero) KK level consists of:
 - 1 tensor field (graviton)
 - δ -1 vector fields (graviphotons)
 - $(\delta^2 \delta 2)/2 + 1$ scalars (graviscalars + radion)

•
$$m_n = \frac{2\pi |\vec{n}|}{R}$$

•
$$m_n = \frac{2\pi |\vec{n}|}{R}$$

The ADD Model KK Gravitons

Each KK level has universal coupling to matter via the energy-momentum tensor

$$S_{int} = \int d^4x \sqrt{\bar{g}} \mathcal{L}_{SM} \Rightarrow \sum_{\vec{n}} \int d^4x \, G^{\vec{n}}_{\mu\nu} T^{\mu\nu} + \phi^{\vec{n}} T^{\mu}_{\mu}$$

- Scalars and vectors decouple apart from the radion
- Although each mode couples with strength ~1/M₄, the sum of contributions from all modes increase the effective coupling up to 1/M_D
- Collider signals could arise from
 - real production (X + missing E_T)
 - virtual production (excesses in σ or differential distributions)
- Large extra dimensions = many closely spaced KK levels approximate the sum of propagators as an integral (Han et al. arXiv:hep-ph/0508097)

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$$D(p^{2}) = \sum_{\vec{n}} \frac{1}{p^{2} - m_{n}^{2} + i\epsilon} \Rightarrow \int^{\Lambda^{2}} \frac{\rho(m_{n})dm_{n}^{2}}{p^{2} - m_{n}^{2} + i\epsilon}$$

➡ Ad hoc cutoff...

Minimal Length Scenario

Natural Regulator

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- Perturbative string theory brings about the possibility of having a minimal length scale strings cannot probe distances smaller than the string scale: I_S~1/M_{pl}
- Generalised uncertainty principle, modified relation between p and k

$$[\hat{x}, \hat{p}] = i \frac{\partial p}{\partial k}$$

- Interesting from a phenomenological point of view in the framework of large extra dimensions because of the lowered fundamental Planck scale, M_D
- Modification of the momentum measure to account for the squeezing of high momentum modes – tool for regularising UV divergences

$$d^4p \to d^4p \prod_{\mu} \frac{\partial k_{\mu}}{\partial p_{\mu}}$$

Use MLS to incorporate a smooth UV cutoff for the sum over KK mode amplitudes

Minimal Length Scenario

Phenomenological Approach

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- Many issues which aren't addressed our approach is to parametrise the effects of MLS at I_S ~ 1/M_{D/S} and apply them to the ADD
- How this extension affects bounds on the parameters of the model in the light of early LHC data in the t-tbar channel
 - Regularisation of the KK graviton sum
 - Modification of the phase space integration (Hossenfelder et al. arXiv:hep-th/0305262)

$$\tilde{\rho}(m_n) = \rho(m_n) \frac{\partial \omega}{\partial E}; \qquad d\tilde{\sigma} = d\sigma \prod_i \frac{E_i}{\omega_i} \prod_{\mu} \frac{\partial k_{\mu}}{\partial p_{\mu}}$$
Unruh relations: $k_M = \frac{1}{l_S} \tanh\left(\frac{p_M}{M_S}\right)$

$$D(s) = \frac{R^d s^{d/2 - 1}}{(4\pi)^{d/2} \Gamma(d/2)} \left[\pi \operatorname{sech}^2\left(\frac{\sqrt{s}}{M_S}\right) + 2iI' \right]; \qquad I' = P\left[\int_0^\infty \frac{y^{d-1}}{1 - y^2} \operatorname{sech}^2\left(\frac{\sqrt{s}}{M_S}y\right) dy \right]$$

Finite! Numerically integrable

Why we love it...

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- One of the main LHC discovery channels NLO $\sigma \sim 830$ pb at 14 TeV
 - Low statistical uncertainty
- Dominated (~90%) by gluon fusion
 - PDF's are important always improving
- ➡ NLO fully known, NNLO on the way (NLL,NNLL)
 - Theoretical errors under control
 - Normalisation of SM background known
- Experimental uncertainty predominantly due to b-tagging and jet energy scale
- Possible to define asymmetries (Γ_{top} > Λ_{QCD}) spin information transferred to decay products A_{FB}, spin asymmetries
- 'Precision' channel, $m_t = 173.3 \pm 1.1 \text{ GeV/c}^2 (\text{CDF})$

NLO SM Corrections to t-tbar cross-section

100 QCD LO: solid/red EW LO: dashed/green 10^{-1} QCD/EW NLO: dotted/black $d\sigma/dM_{tt}~(pb/GeV)$ QCD NLO: dot-dashed/blue 10^{-2} LHC at 14 TeV 10^{-3} 10^{-4} 10^{-5} 100 80 Correction (%)60 40 20 0 -20 1000 2000 3000 4000 () M_{tt} (GeV)

 $p p \rightarrow t t$

- In general, KK graviton interactions are suppressed by powers of M_S
- Interference terms with the SM (QCD) will dominate if they exist

$$\sigma_T \sim \sigma_{SM} + (M_S)^{-4} \sigma_{int} + (M_S)^{-8} \sigma_{grav}$$

- For a hadron collider, the di-quark initial state is mediated by a gluon colour structure prevents an interference term with ADD
- gg initial state in the t and u channels are colour singlets and will therefore contribute an interference term

$$|\mathcal{M}_{int}|^{2} = \left(\sum_{\vec{n}} \frac{3}{3} - \frac{G_{\mu\nu}^{\vec{n}}}{2} - \frac{$$

 Promising way to search for evidence of large extra dimensions – excess in the ttbar cross-section, invariant mass distribution

- Look at the effect of the ADD model and MLS on the t-tbar cross section as a function of the fundamental scale, M_S, determine the reach of the LHC to set bounds on it
- Consider other observables such as differential distributions that could yield deviations from the SM
- Based on work done on the Drell-Yan channel (Bhattacharya et al. arXiv:hep-ph/0408295)
 - MLS: $M_S > 1.5 \text{ TeV}$ @ Tevatron, $M_S > 6 \text{ TeV}$ @ LHC (14 TeV)
 - Improve on LHC bounds, look at 7 TeV and look to compare with early data
- For consistency, compute tree-level ADD/MLS contribution and compare it to treelevel SM predictions
- Included q-qbar initial state as well as pure graviton contribution from gg diagram (Mathews et al. arXiv:hep-ph/9811501v4)

- Calculate the contribution to the t-tbar cross section of ADD with and without MLS over a region of parameter space
 - Vary fundamental Planck scale, M_S
 - Vary number of extra dimensions, δ =3-6
 - Compare ADD+MLS to pure ADD in terms of reach on M_S bounds, divergent behaviour
- Experimental setup: LHC at 14TeV (100 fb⁻¹) and Tevatron (4 fb⁻¹), adding LHC at 7 TeV (5 fb⁻¹) for possible data comparison.
 - CTEQ6 PDF sets
 - Madgraph for LO SM matrix elements: QCD + EW (v. small)
 - Make kinematical cuts (invariant mass) to minimise SM background and focus on regions where NLO corrections are not as important compared to ADD
- When determining reach to set bounds on M_S, fold in the t-tbar reconstruction efficiencies for various decay channels: ~4%(3%) at the LHC(Tevatron)

Differential Distributions

- Use differential cross-section wrt invariant mass to guide kinematic cuts and look at behaviour as one approaches the cutoff $M_{S}(M_{D})$
- Divergence of the ADD cross-section as one approaches the cutoff
- MLS tames this behaviour prevents overestimation of cross-section





Differential Distributions

- MLS competes better with ADD for larger δ
- Invariant mass cut of 1 TeV
- Integrate invariant mass up to 0.8Ms
- Within this range can expect some excesses of order 100% depending on parameters



10¹

10⁰

10⁻¹

10⁻²

10⁻³

10⁻⁴

400

600

800

1000

1200

1400

1600

1800

2000

dσ/dM_{tt} [pb/GeV]

SM

ADD MLS

 M_{tt} Distribution for M_S =2TeV, LHC at 14 TeV, δ =3

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- Determined appropriate invariant mass ranges for 3 experimental setups
- Plot total production cross-section as a function of M_S and compare against 95% CL upper and lower bounds for LO SM cross-section
- LHC at 14 TeV assumed 100 fb⁻¹
- Reconstruction efficiencies had a significant effect on attainable bounds need more integrated luminosity



Total Cross-section: LHC at 14 TeV

- Show that the LHC at 14 TeV has the potential to reach fundamental scales of up to ~4.2 TeV – lowered to 4 TeV in the MLS
- Relies on a large amount of data
- Drell-Yan study did not include acceptance or reconstruction efficiencies
- Tevatron does not improve bounds because of dominance of q-qbar initial state
- LHC at 7 TeV shows some promise of determining bounds assuming 5 fb⁻¹



Total Cross-sections: LHC at 7 TeV





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M _S Bounds	t-tbar (incl. ε _{reco.})		Drell-Yan		Lint
	ADD	MLS	ADD	MLS	fb⁻¹
LHC @ 14 TeV	4.2 TeV	4 TeV	7 TeV	6 TeV	100
LHC @ 7 TeV	1.5 TeV	1.3 TeV			5

- Overall, in taming the divergent nature of the cross-sections, MLS lowers the reach of experiments to set bounds on M_S
- More focused cuts could improve signal to background ratio, help make up for the losses due to poor reconstruction efficiencies
- In principle t-tbar channel at the LHC could improve on the reach of Drell-Yan because of the gg initial state and ADD interfering with QCD rather than EW at the Tevatron

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- As an example: significance of ADD/MLS signal in deviation from SM t-tbar cross section as a function of luminosity for M_S = 4 TeV
- LHC at design luminosity certainly has potential to probe the existence of large extra dimensions with fundamental scales of several TeV
- More focused event selection will certainly improve ability to set bounds



Results Significance plots: LHC @ 14 TeV



- These are rough numbers: what we really want is to compare with early data and get some real bounds perhaps combining the Drell-Yan and t-tbar channels
- Started by looking at simple quantities like total cross-section and invariant mass distributions but there are many more variables that one could look at
- Could exploit the potential of asymmetries at the LHC via interference with SM
 - EW or NLO QCD for t-tbar
 - Drell-Yan
- Recent paper has looked at centrality ratio $\chi = \exp(y_1 y_2)$ as a discerning variable: measure of angular distribution (Franceschini et al. arXiv:1101.4919v2 [hep-ph])
- A kinematical variable being measured at ATLAS and CMS to quantify angular distribution of dijet final state
- Expected to be flatter for QCD than for the exchange of a spin 2 particle

Outlook Centrality Ratio

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Outlook Centrality Ratio

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- Considered a simple extension of the ADD model whereby a minimal length is assumed applied to the t-tbar channel
- This scenario, as expected, regulates some of the divergent behaviour of the ADD model by suppressing high mass virtual KK modes
- Looking at the total cross-section and invariant mass distributions, investigated the reach of collider experiments to set bounds on the fundamental Planck scale of the model, M_D
- Considered LHC at design energy, Tevatron and also 'early LHC data' scenario with the aim of using data and obtaining bounds
- Reconstruction efficiency of the t-tbar final state limits the ability to set bounds on M_s without large amounts of data
- Potential to look at other quantities such as χ and asymmetries