### **Modelling experience at Diamond**

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### Outline

- Introduction to the diamond storage ring
- Modelling and correcting the <u>linear optics</u>
- Modelling and correcting the <u>nonlinear optics</u>

Frequency Maps

Spectral lines analysis and resonance driving terms

Limits of the methods

Conclusions and ongoing work





### **Diamond aerial view**



Diamond is a third generation light source open for users since January 2007 100 MeV LINAC; 3 GeV Booster; 3 GeV storage ring 2.7 nm emittance – 300 mA – 18 beamlines in operation (10 in-vacuum small gap IDs)

### Diamond storage ring main parameters non-zero dispersion lattice



48 Dipoles; 240 Quadrupoles; 168 Sextupoles
(+ H and V orbit correctors + 96 Skew Quadrupoles)
3 SC RF cavities; 168 BPMs

**Quads + Sexts have independent power supplies** 

Energy	3 GeV		
Circumference	561.6 m		
lo. cells	24		
Symmetry	6		
Straight sections	6 x 8m, 18 x 5m		
nsertion devices	4 x 8m, 18 x 5m		
Beam current	300 mA (500 mA)		
Emittance (h, v)	2.7, 0.03 nm rad		
.ifetime	> 10 h		
/lin. ID gap	7 mm (5 mm)		

Beam size (h, v)123, 6.4  $\mu$ mBeam divergence (h, v)24, 4.2  $\mu$ rad(at centre of 5 m ID)178, 12.6  $\mu$ mBeam divergence (h, v)16, 2.2  $\mu$ rad(at centre of 8 m ID)

#### **Linear optics modelling with LOCO** Linear Optics from Closed Orbit response matrix – J. Safranek et al.



#### LOCO has solved the problem of the correct implementation of the linear optics

### Linear coupling correction with LOCO

Skew quadrupoles can be simultaneously zero the <u>off diagonal blocks</u> of the measured response matrix and the <u>vertical disperison</u>

$$\chi^{2}(\overline{Q},\overline{G}_{BPMs},\overline{S}_{q},\overline{K}_{BPMs},...) = \sum_{i,j} \left( R_{ij}^{measured} - R_{ij}^{model}(\overline{Q},\overline{S}_{q},\overline{G}_{BPMs},\overline{K}_{BPMs},...) \right)^{2}$$

$$Measured Response Matrix$$

$$Measu$$

### **Residual vertical dispersion after correction**

Without skew quadrupoles off

After LOCO correction

r.m.s. Dy = 14 mm r.m.s. Dy = 700 μm







### **Measured emittances**

Coupling without skew quadrupoles off K = 0.9%

(at the pinhole location; numerical simulation gave an average emittance coupling  $1.5\% \pm 1.0\%$ )

Emittance [2.78 - 2.74] (2.75) nm Energy spread [1.1e-3 - 1.0-e3] (1.0e-3)

After coupling correction with LOCO (2\*3 iterations)

1<sup>st</sup> correction K = 0.15% 2<sup>nd</sup> correction K = 0.08%

V beam size at source point 6 µm

Emittance coupling  $0.08\% \rightarrow V$  emittance 2.2 pm

Variation of less than 20% over different measurements





### Comparison machine/model and Lowest vertical emittance

	Model emittance	Measured emittance	$\beta$ -beating (rms)	Coupling* (ε <sub>y</sub> / ε <sub>x</sub> )	Vertical emittance
ALS	6.7 nm	6.7 nm	0.5 %	0.1%	4-7 pm
APS	2.5 nm	2.5 nm	1 %	0.8%	20 pm
ASP	10 nm	10 nm	1 %	0.01%	1-2 pm
CLS	18 nm	17-19 nm	4.2%	0.2%	36 pm
Diamond	2.74 nm	2.7-2.8 nm	0.4 %	0.08%	2.2 pm
ESRF	4 nm	4 nm	1%	0.1%	4.7 pm
SLS	5.6 nm	5.4-7 nm	4.5% H; 1.3% V	0.05%	2.0 pm
SOLEIL	3.73 nm	3.70-3.75 nm	0.3 %	0.1%	4 pm
SPEAR3	9.8 nm	9.8 nm	< 1%	0.05%	5 pm
SPring8	3.4 nm	3.2-3.6 nm	1.9% H; 1.5% V	0.2%	6.4 pm
SSRF	3.9 nm	3.8-4.0 nm	<1%	0.13%	5 pm

\* best achieved

## Comparison real lattice to model linear and nonlinear optics



The calibrated nonlinear model is meant to <u>reproduce all the measured</u> <u>dynamical quantities</u>, giving us insight in which resonances affect the beam dynamics and possibility to <u>correct</u> them





### Comparison real lattice to model linear and nonlinear optics

Frequency Maps and amplitudes and phases of the spectral line of the betatron motion can be used to compare and correct the real accelerator with the model



Combining the complementary information from FM and spectral line should allow the calibration of the nonlinear model and a full control of the nonlinear resonances

## Frequency map and detuning with momentum comparison machine vs model (I)

Using the measured **Frequency Map** and the measured **detuning with momentum** we can build a fit procedure to calibrate the nonlinear model of the ring



- tracking data
- build FM and detuning with momentum



- BPMs data with kicked beams
- measure FM and detuning with momentum

$$A_{target} = (Q_{x}[(x, y)_{1}], ..., Q_{x}[(x, y)_{n}], Q_{y}[(x, y)_{1}], ..., Q_{y}[(x, y)_{n}], ....$$
$$..., Q_{x}(\delta_{1}), ..., Q_{x}(\delta_{m}), Q_{y}(\delta_{1}), ..., Q_{y}(\delta_{m}))$$

The distance between the two vectors

$$\chi^{2} = \sum_{k} \left( A_{Model} (j) - A_{Measured} (j) \right)^{2}$$

can be used for a Least Square Fit of the sextupole gradients to minimise the distance  $\chi 2$  of the two vectors

## Frequency map and detuning with momentum comparison machine vs model (I)



Sextupole strengths variation less than 3%

#### The most complete description of the nonlinear model is mandatory !

Measured multipolar errors to dipoles, quadrupoles and sextupoles (up to b10/a9)

- Correct magnetic lengths of magnetic elements
- Fringe fields to dipoles and quadrupoles
- Substantial progress after correcting the frequency response of the Libera BPMs

# Frequency map and detuning with momentum comparison machine vs model (II)



The fit procedure based on the reconstruction of the measured FM and detunng with momentum describes well the **dynamic aperture**, the **resonances excited** and the dependence of the **synchrotron tune vs RF frequency** 

R. Bartolini et al. Phys. Rev. ST Accel. Beams 14, 054003





### **Frequency Analysis of betatron motion**

**Example: Spectral Lines for tracking data for the Diamond lattice** 



Each spectral line can be associated to a resonance driving term

J. Bengtsson (1988): CERN 88–04, (1988). R. Bartolini, F. Schmidt (1998), Part. Acc., **59**, 93, (1998). R. Tomas, PhD Thesis (2003)

### All diamond BPMs have turn-by-turn capabilities







## Spectral line (-1, 1) in V associated with the sextupole resonance (-1,2)







### Frequency Analysis of Betatron Motion and Lattice Model Reconstruction

Using the measured amplitudes and phases of the spectral lines of the betatron motion we can build a fit procedure to calibrate the nonlinear model of the ring



Least Square Fit of the sextupole gradients to minimise the distance  $\chi^2$  of the two Fourier coefficients vectors

### Simultaneous fit of (-2,0) in H and (1,-1) in V



Both resonance driving terms are decreasing

### **Sextupole variation**



Now the sextupole variation is limited to < 5%

#### **Both resonances are controlled**

We measured a slight improvement in the lifetime (10%)





### **Limits of the Frequency Analysis techniques**

**<u>BPMs precision</u>** in turn by turn mode (+ gain, coupling and non-linearities)

10  $\mu$ m with ~10 mA

very high precision required on turn-by-turn data (not clear yet is few tens of  $\mu$ m is sufficient); Algorithm for the precise determination of the betatron tune lose effectiveness quickly with noisy data. R. Bartolini et al. Part. Acc. 55, 247, (1995)

**Decoherence** of excited betatron oscillation reduce the number of turns available Studies on oscillations of beam distribution shows that lines excited by resonance of order m+1 decohere m times faster than the tune lines. This decoherence factor m has to be applied to the data R. Tomas, PhD Thesis, (2003)

The machine <u>tunes are not stable</u>! Variations of few 10<sup>-4</sup> are detected and can spoil the measurements

**BPM gain and coupling** can be corrected by LOCO,

**BPM nonlinearities** corrected as per R. Helms and G. Hofstaetter PRSTAB 2005 **BPM frequency response** can be corrected with a proper deconvolution of the time filter used to built t-b-t data form the ADC samples R. Bartolini subm. to PRSTAB

### Conclusions



J.A.I.

