

RHIC Non-linear Chromaticity and IR Correction

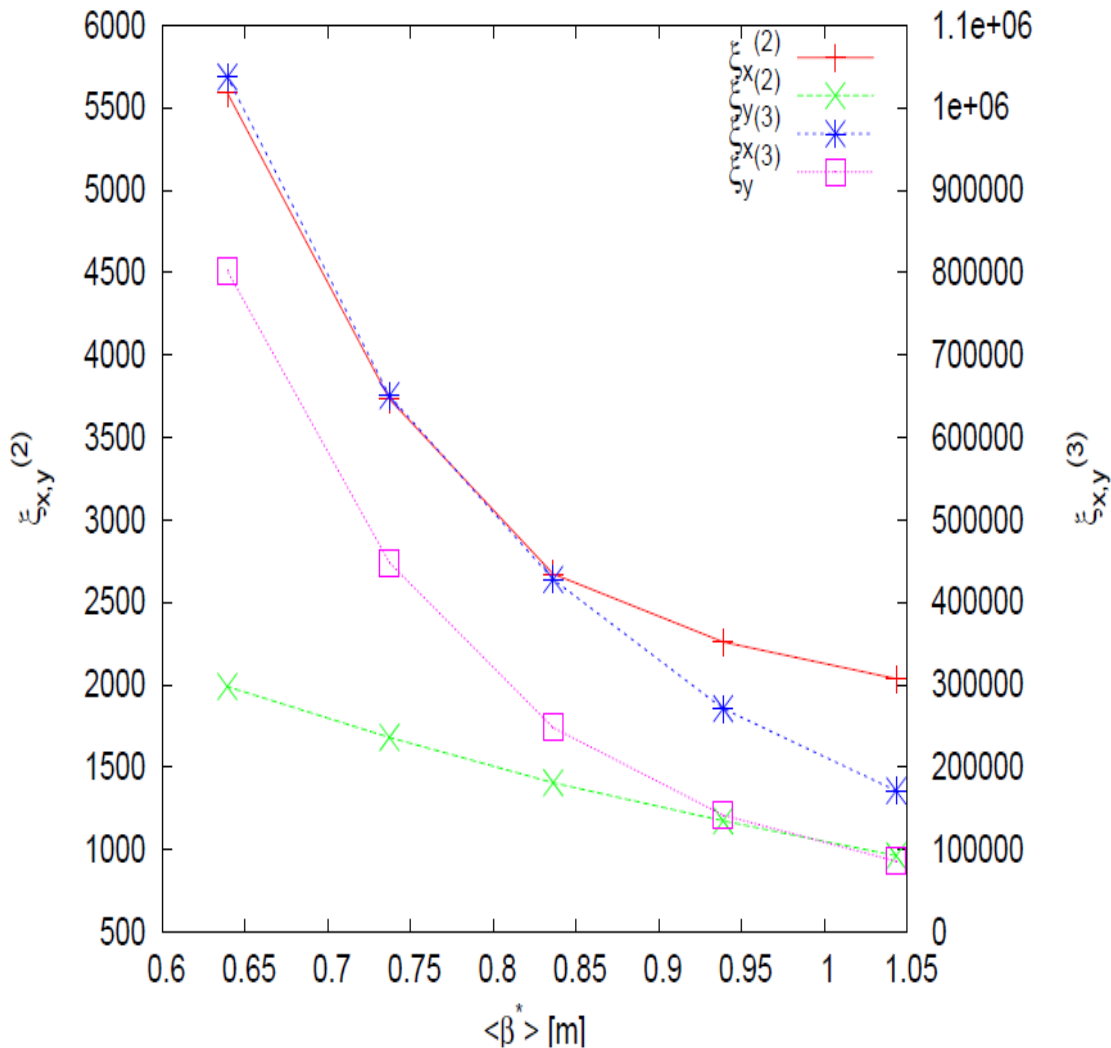
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and Modeling for High-Performance Storage Rings,
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Motivation for non-linear chromaticity correction

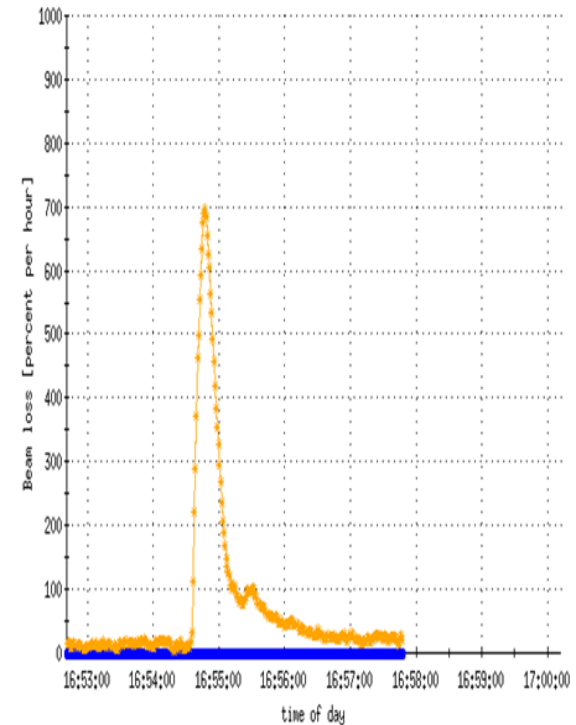
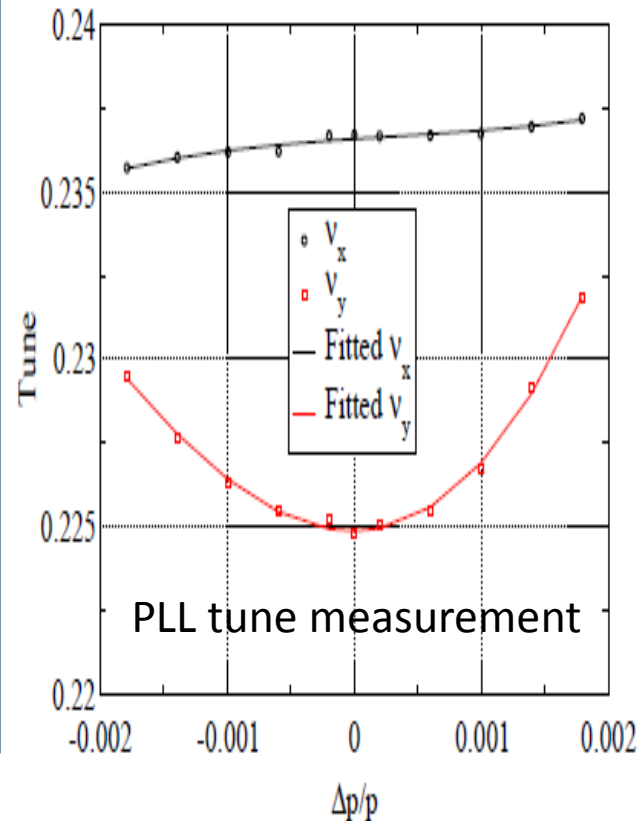
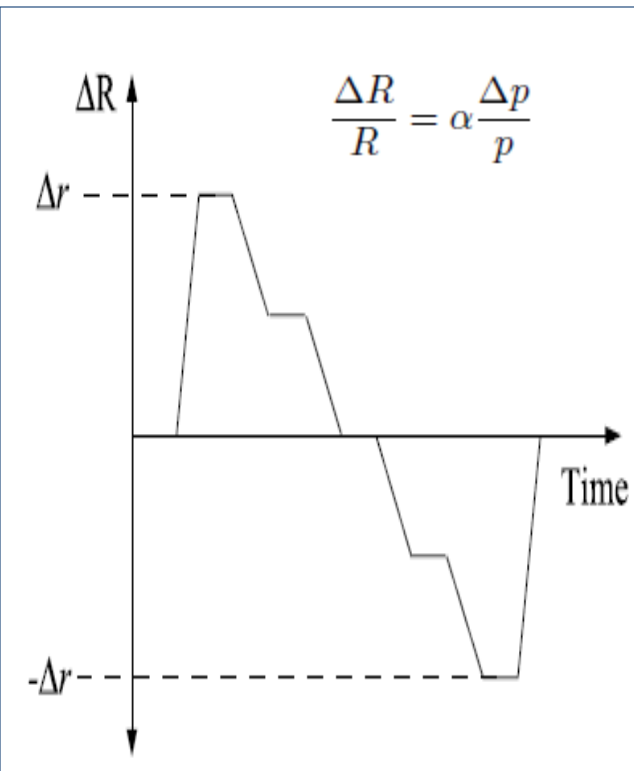


➤ Low β^* lattice will give large off-momentum tune shifts and off-momentum β -beat.

➤ Currently $\beta^* = 0.7m$ for RHIC proton and heavy ion runs. We are marching to or even below $\beta^*=0.5m$. Dynamic β^* squeezing to 0.58m demonstrated with Stochastic cooling.

➤ In the heavy ion runs, RF re-bucketing requires a enough momentum aperture. Before re-bucketing, $(dp/p_0)_{\max} = 0.0009$. After re-bucketing, $(dp/p_0)_{\max} = 0.0017$.

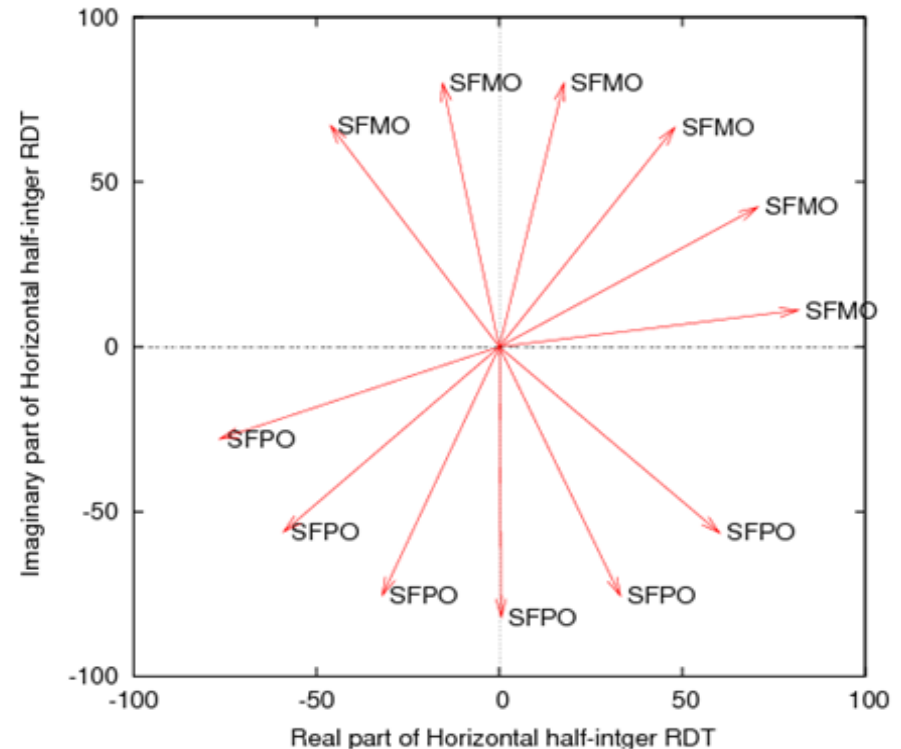
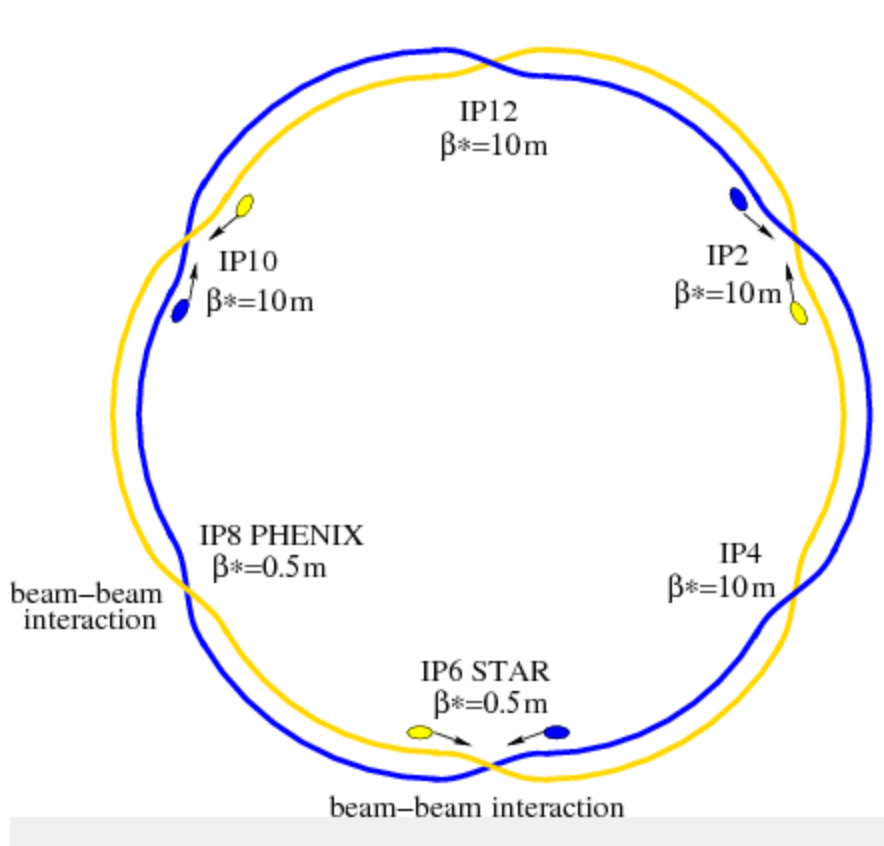
Measurement of linear & non-linear chromaticity



- Off-momentum tunes are measured with PLL and fitted up to 3rd order.
- Reliable measurement is very important for correction and modeling.
- Beam decay in this measurement is a good indicator of momentum aperture.

(S. Tepikian, et al, PAC'05)

RHIC arc chromatic sextupoles



➤ There are 144 arc sextupoles in 6 arcs. Before 2007, 2-family scheme was used for chromaticity correction. After 2007, the number of arc sextupole power supplies was doubled.

➤ The betatron phase advance per FODO cell is about 83degree. We split the SFs and SDs into two sub-families in each arc. There are totally 24 sub-families.

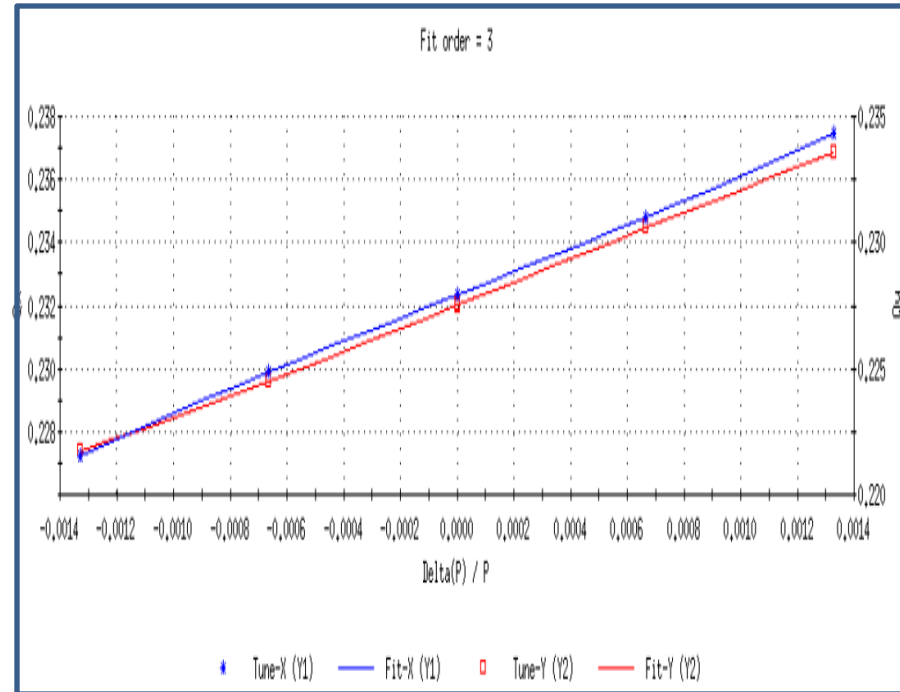
(Y. Luo, et al., PAC'07)

Lattice model based correction

Table 4: Blue store: linear chromaticity measurement and modeling.

	Tunes (Q_x, Q_y)	Chromaticities (Q'_x, Q'_y)	Chromatic sextupole strengths (K_2)
Case 5:			SF: 0.3310m^{-3} SD: -0.6575m^{-3}
Measurements	(28.228, 29.215)	(4.0, 4.3)	-
Modeling with $(K_2L)_{dipole,mean} = 0\text{m}^{-2}$	-	(0.8, 8.2)	-
Modeling with $(K_2L)_{dipole,mean} = 0.0095\text{m}^{-2}$	-	(4.0, 5.1)	-
Case 6:			SF: 0.3202m^{-3} SD: -0.6348m^{-3}
Measurements	(28.232, 29.227)	(1.5, 1.6)	-
Modeling with $(K_2L)_{dipole,mean} = 0\text{m}^{-2}$	-	(1.2, 5.1)	-
Modeling with $(K_2L)_{dipole,mean} = 0.0095\text{m}^{-2}$	-	(1.9, 2.0)	-

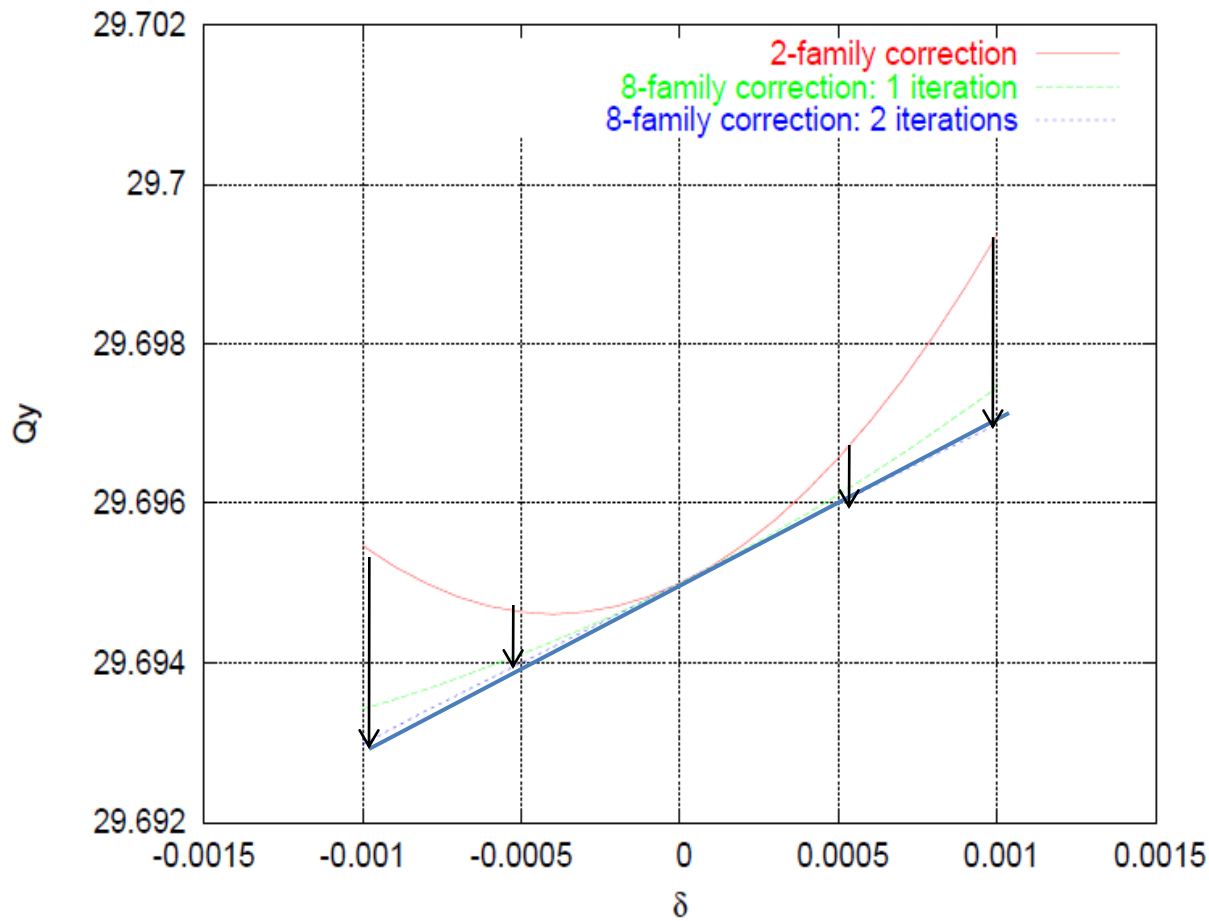
$\Delta Q'_{x,y} \sim 4$ at store.



- Non-linear chromaticity can be easily corrected if the optics model reflects real machine.
- One difficulty is discrepancy in Q' from machine and prediction. We model $\langle b_2 \rangle$ in arc dipoles from Q' split. On top of it, we applied model based correction.
- This method worked very well just by typing the correction strengths for the Blue ring in 2007 Au-Au run when $\beta^* = 0.8\text{m}$.

(Y. Luo, et al., BNL C-AD/AP/Note-276, 2007)

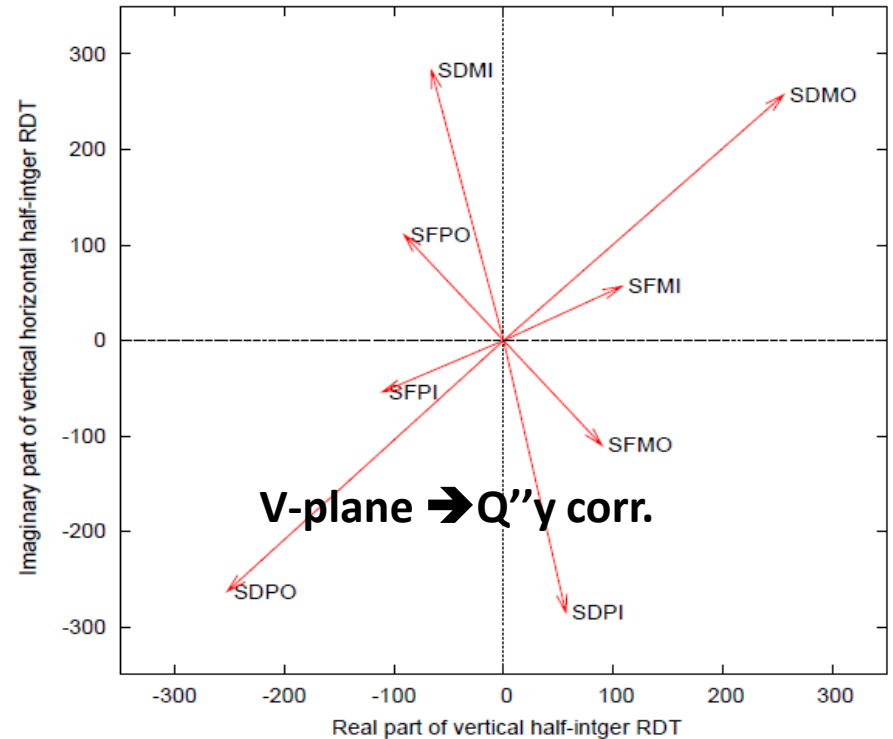
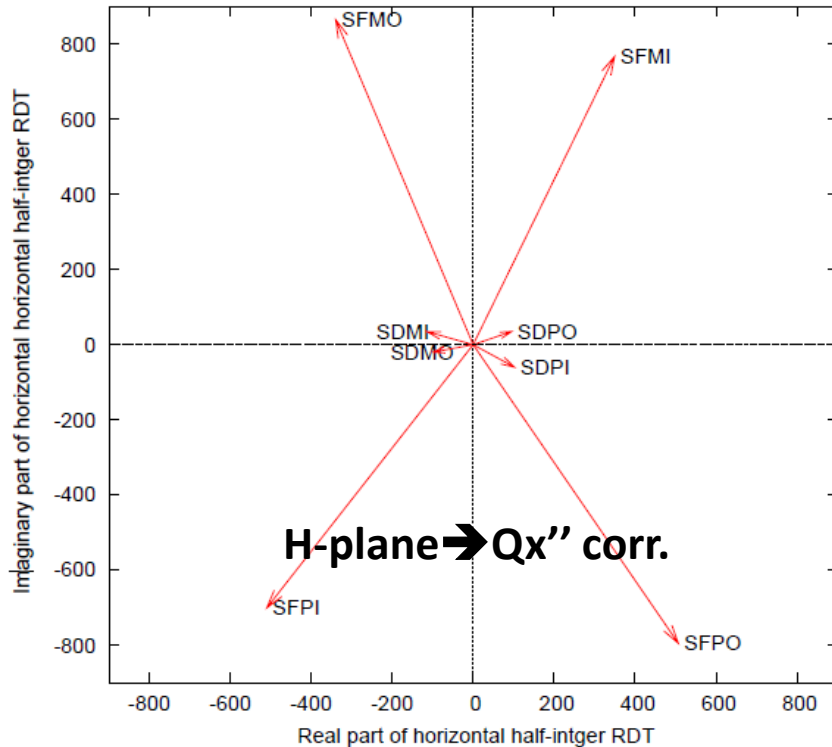
Response Matrix based on-line correction



(Y. Luo, et al. PAC'07.)

- Using the response matrix of off-momentum tunes w.r.t. the strengths of sextupole families to push the off-momentum tunes onto the projected off-momentum tunes by the linear chromaticity.
- This method corrects all order of non-linear chromaticities.
- This method was tested and didn't work very well due to SVD over-shooting and change of linear chromaticity with correction (c.o.d. in sexts).

4-Knobs for on-line correction



- According to their contribution to half-integer resonance driving terms, 144 arc chromatic sextupoles are sorted into four families (4-knobs) to minimize the second order chromaticities.
- This method avoids sextupole polarity reversal and unbalanced sextupole correction strengths. The correction will not change the first order chromaticity at least in principle.
- It was proved to be very effective to correct second order chromaticities and was implemented in RHIC control system.

(Y. Luo, et al., IPAC 2010)

Measurement vs. model calculation

Table 1: Second order chromaticity comparison

Run#	Ring	Δr [mm]	X [meas]	X [model]	Y [meas]	Y [model]
4207	Yellow	2.0	12 ± 26	137	1332 ± 28	1141
			-13 ± 44		1661 ± 119	
			94 ± 27		1563 ± 32	
			448 ± 33		1452 ± 25	
4454	Blue	2.0	410 ± 19	844	855 ± 24	940
			346 ± 27		838 ± 62	
			365 ± 17		924 ± 48	
			292 ± 31		871 ± 21	
Averaging measured values for a comparison						
4207	Yellow		134 ± 179	137	1460 ± 105	1141
4454	Blue		361 ± 40	844	870 ± 26	940

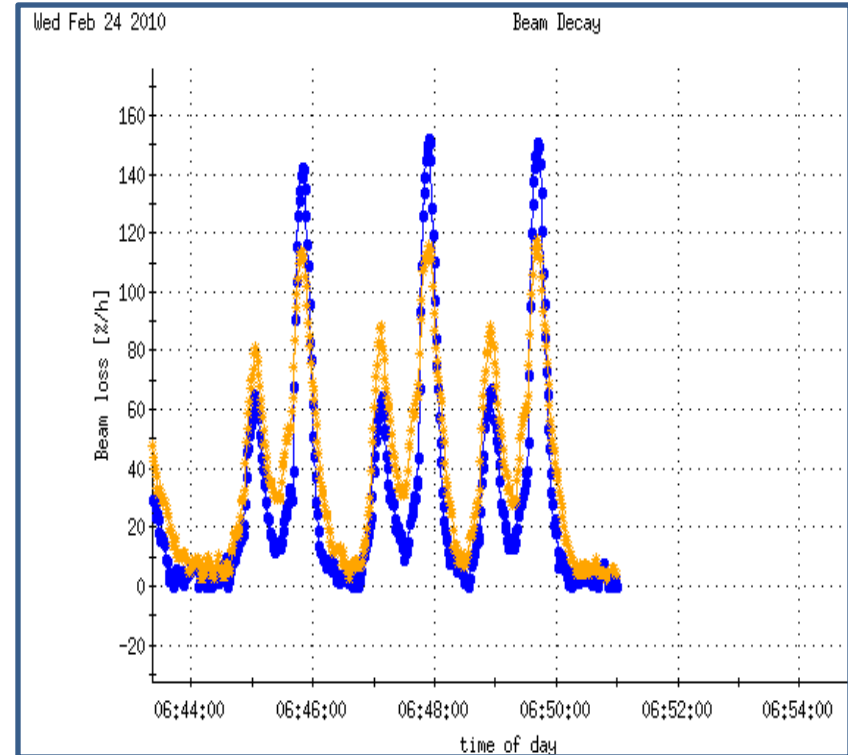
$\beta^*=1.0$ m
lattice

(Ref: S. Tepikian, et al, PAC'05)

- Measurement of second order chromaticity agreed with the model prediction very well before 2007 with $\beta^* > 0.8$ m lattices.
- Large discrepancies were observed with $\beta^* < 0.7$ m lattices.

Errors in the chromaticity measurement

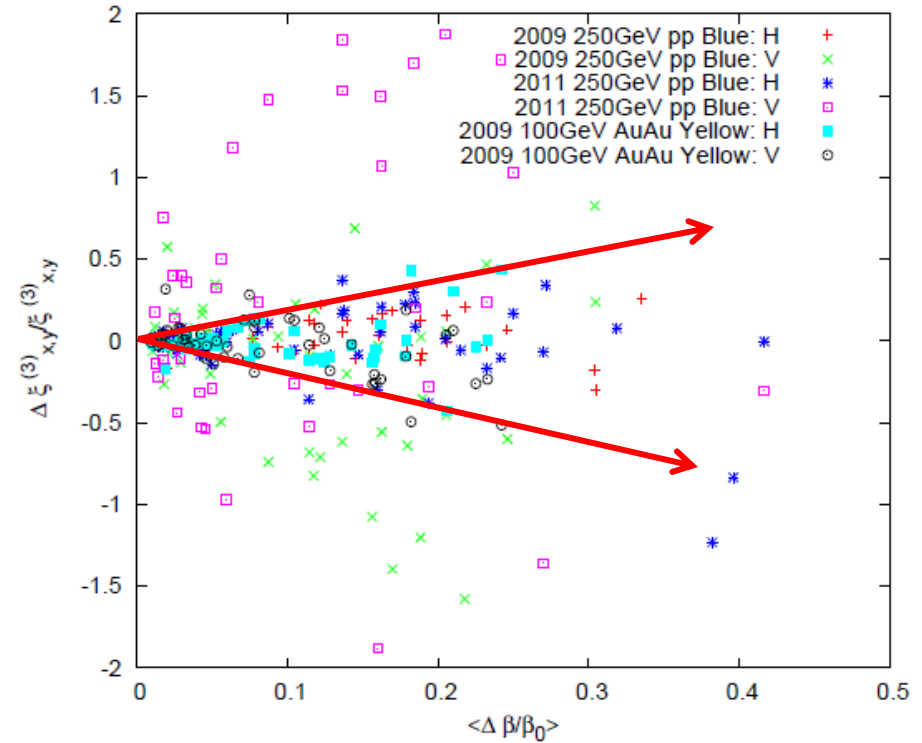
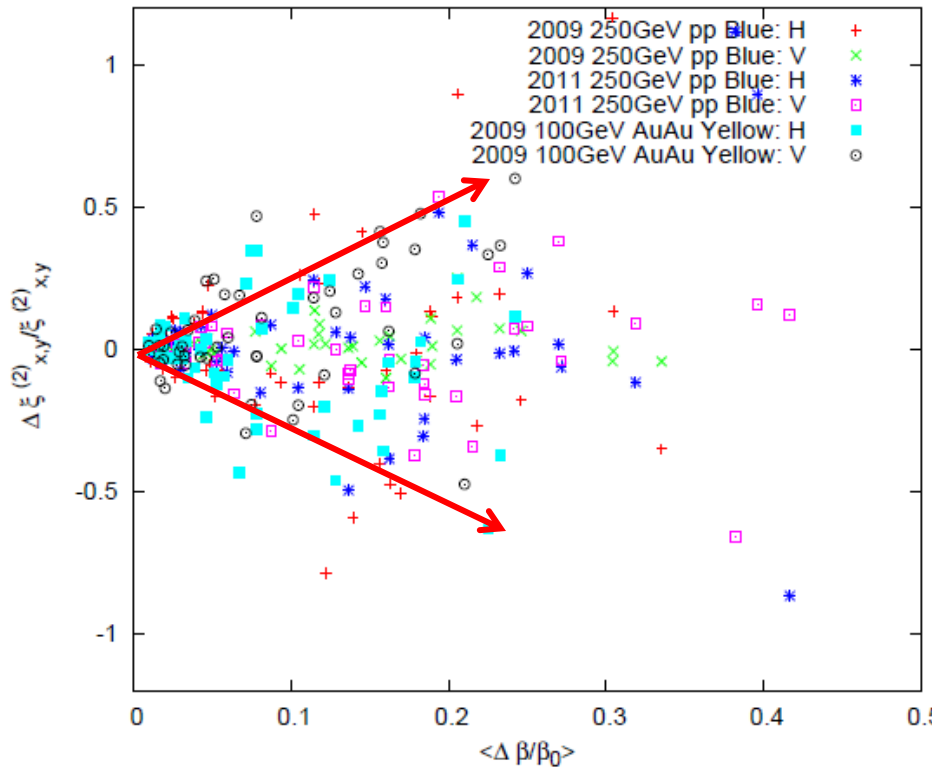
Radial (mm)	$Q''_x/2$	$Q''_y/2$
0.5	1089	-1645
	1112	-942
	1053	-1156
	870	-1201
	852	-1039
	779	-1377
Model	3077	4728



(Yellow 2010 Au-Au run, $\beta^*=0.7$ m)

- Measurement tunes and fitted chromaticities didn't repeat very well.
- Small momentum aperture prevented large (dp/p_0) scan.
- Un-balanced beam decay with positive and negative (dp/p_0) in this example.

Non-linear chromaticity and on-momentum β -beat



➤ Simulation shows that more than 10% on-momentum beta-beat will introduce 30% second order chromaticity uncertainty (1200 units here).

(Y. Luo, et al., BNL C-AD/AP/418, 2011)

Source of non-linear chromaticities

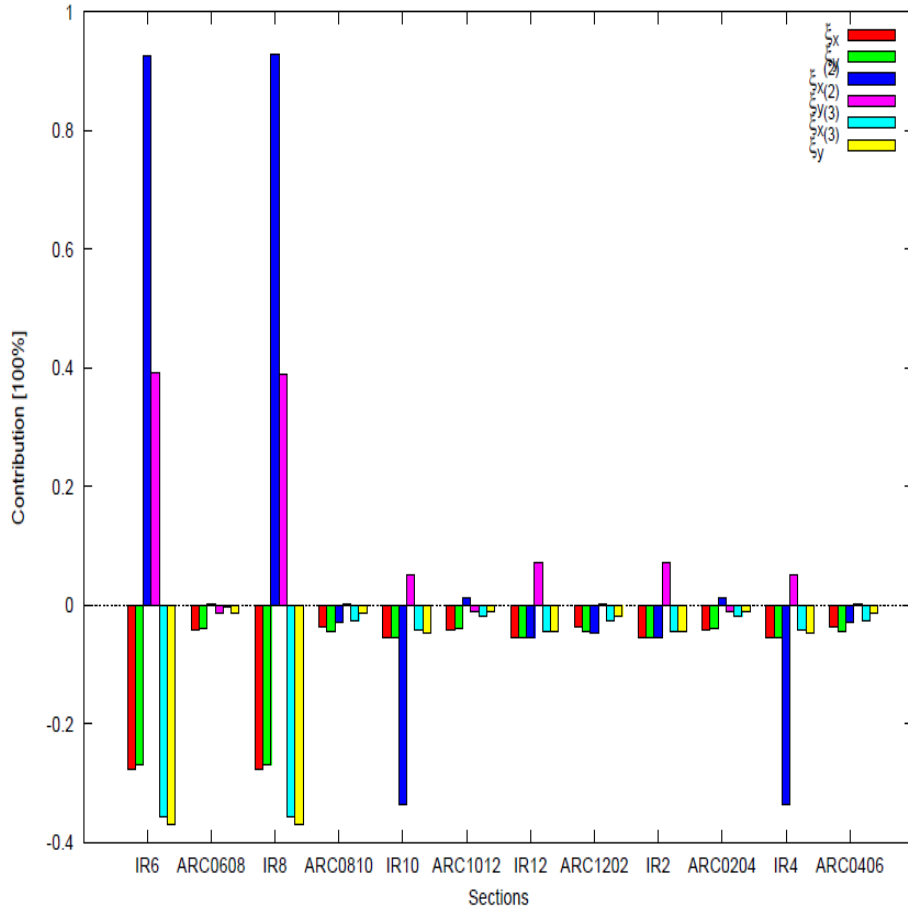


Table 4: Contributions to the linear and nonlinear chromaticities

Sections	$\xi_x^{(1)}$	$\xi_y^{(1)}$	$\xi_x^{(2)}$	$\xi_y^{(2)}$	$\xi_x^{(3)}$	$\xi_y^{(3)}$
2009-pp-Blue:						
IR6 and IR8	-0.55	-0.53	1.85	0.78	-0.71	-0.73
Other IRs	-0.21	-0.21	-0.77	0.24	-0.17	-0.18
Arcs	-0.22	-0.24	-0.07	-0.02	-0.11	-0.081
2011-pp-Blue:						
IR6 and IR8	-0.58	-0.57	1.20	0.84	-0.74	-0.78
Other IRs	-0.20	-0.20	-0.20	0.17	-0.15	-0.16
Arcs	-0.21	-0.22	0.00	-0.02	-0.10	-0.05
2010-AuAu-Yellow:						
IR6 and IR8	-0.47	-0.50	0.34	1.15	-0.74	-0.86
Other IRs	-0.26	-0.24	0.71	-0.20	-0.17	-0.03
Arcs	-0.25	-0.25	-0.06	0.04	-0.08	-0.10

IR6 and IR8 contributes most of non-linear chromaticities than the other sections.

(Y. Luo, et al., BNL C-AD/AP/Note-418)

Phase advances between IP6 and IP8

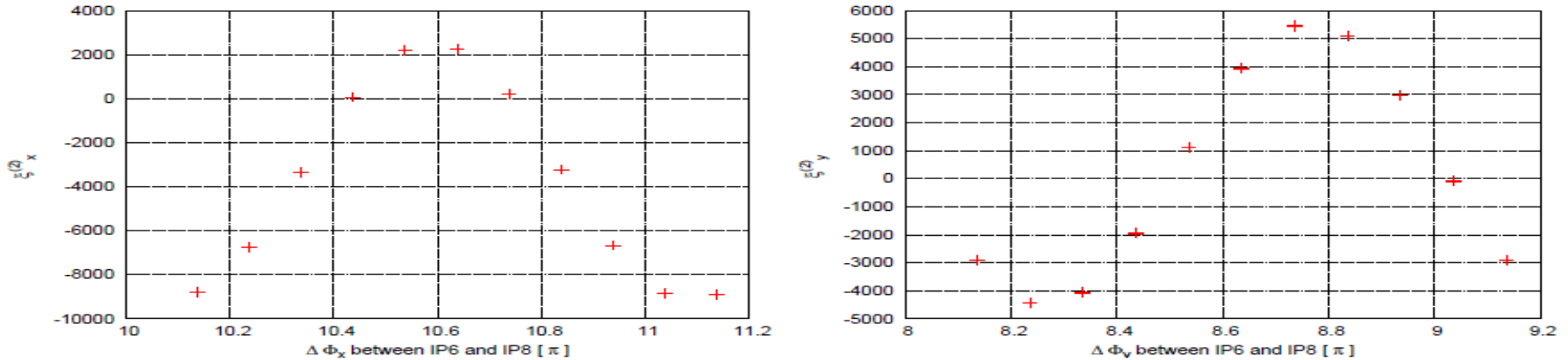
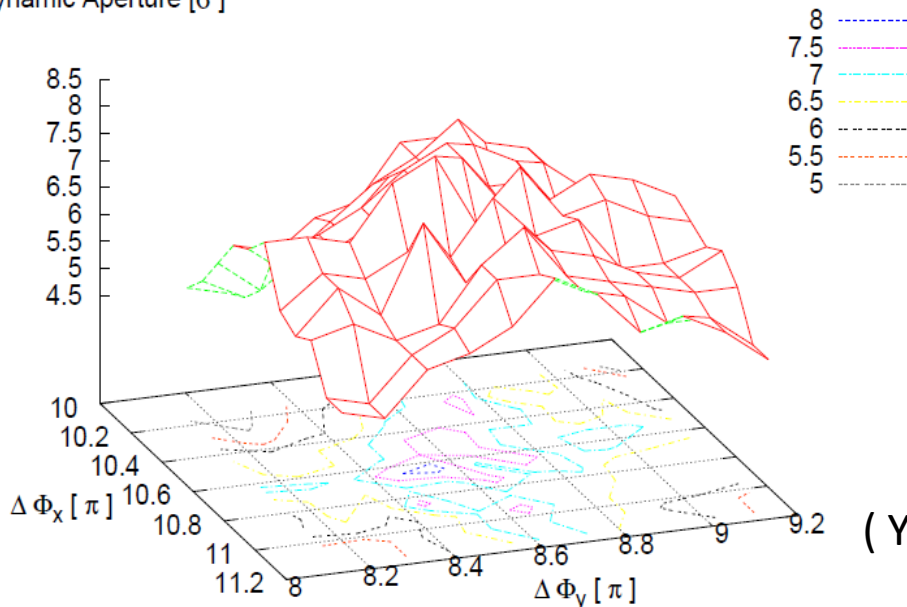


Figure 2: 2011pp-Blue: second order chromaticities versus phase advances between IP6 and IP8.

Dynamic Aperture [σ]

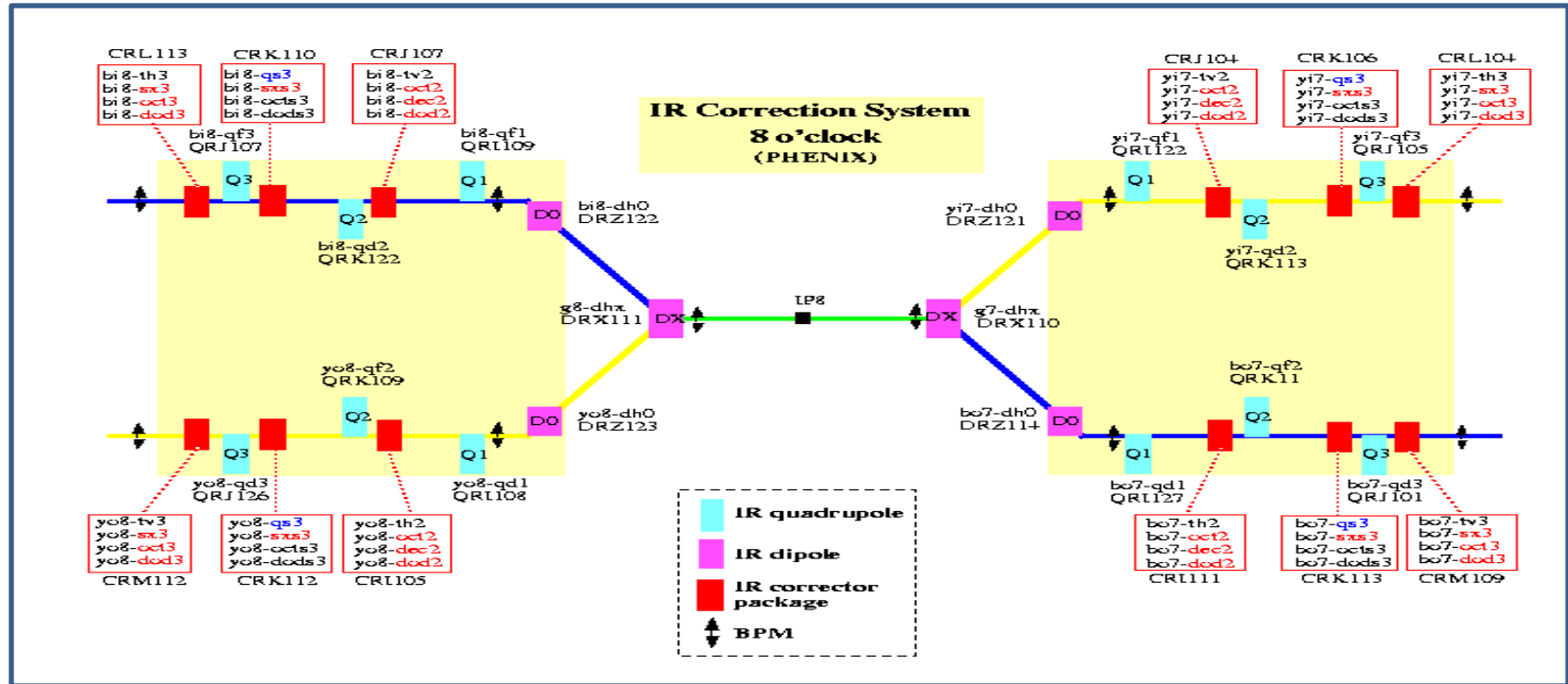


- It is possible to adjust the phase advances between IP6 and IP8 to cancel their the contribution to the global chromatic effects from IR6 and IR8.
- Unfortunately so far we have no enough arc quadrupole power supplies to do this job.

(Y. Luo, et al., BNL C-AD/AP/Note-426, 2011;
Y. Luo, et al., PAC 2011.)

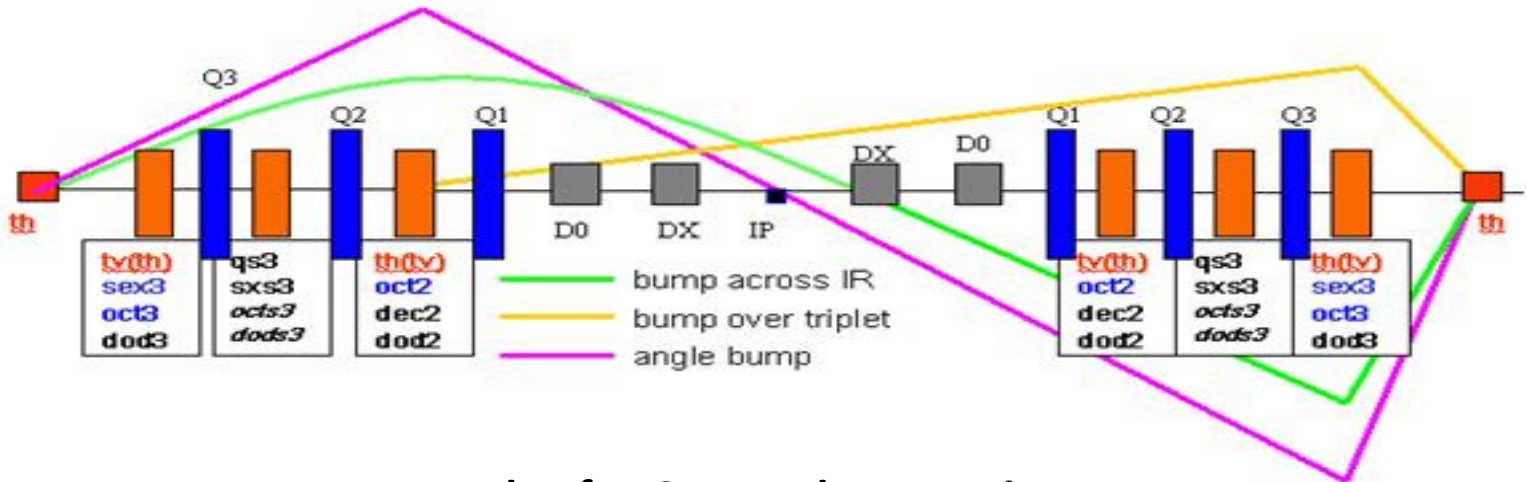
2011pp run:default phase advances between IPs: (10.65 π , 8.64 π)

IR Multipole Correction

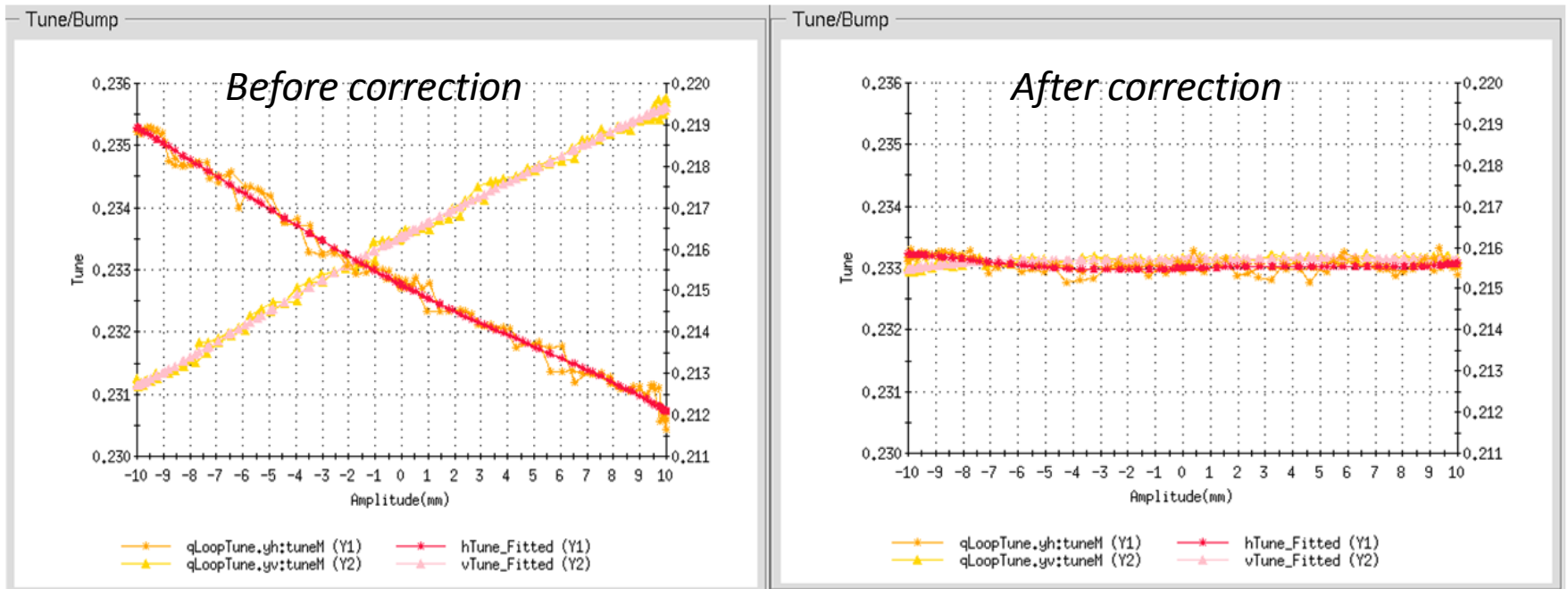


- 1) J. Wei, Particle Accelerators, Vol. 55, 1996 .-→Action-angle minimization
- 2) J-P. Koutchouk, F. Pilat, V. Ptitsyn, et al. In PAC 2001. →IR bump method
- 3) Y. Luo, et al., PAC 2005. →Compare different methods
- 4) F. Pilat, Y. Luo, et al., PAC 2005. →RHIC experiences with IR bump
- 5) W. Fischer, Y. Luo, et al., IPAC 2010. → 5th and 6th IR non-linear correction

On-line IR bump correction



Example of IR Sextupole correction

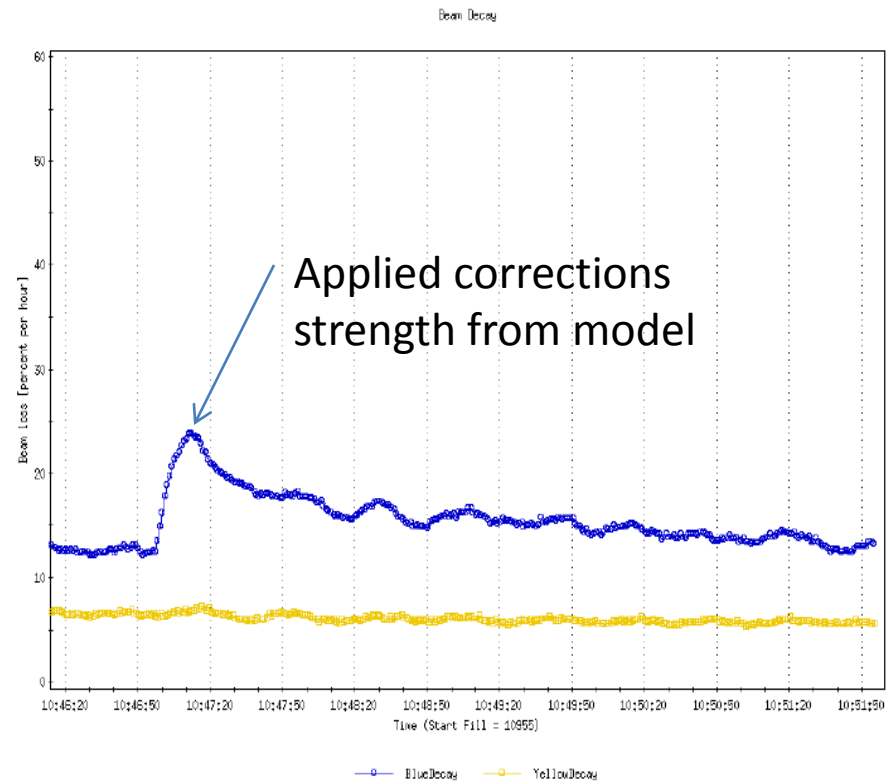


IR bump versus model correction

Blue ring:		
bi5-sx3	0.001439	-0.00040
bo6-sx3	-0.003861	-0.00050
bo7-sx3	-0.002804	-0.00400
bi8-sx3	0.001890	0.00180
Yellow ring:		
yo5-sx3	-0.000922	-0.00100
yi6-sx3	0.000288	0.00070
yi7-sx3	0.001679	0.00200
yo8-sx3	-0.003085	-0.00200

Model

IR-bump



- In June 2009, we compared the correction strengths from IR non-linear model and that from IR-bump.
- Experiment showed that the corrections from IR-bump method gave a better beam lifetime.

3Qx Resonance Correction

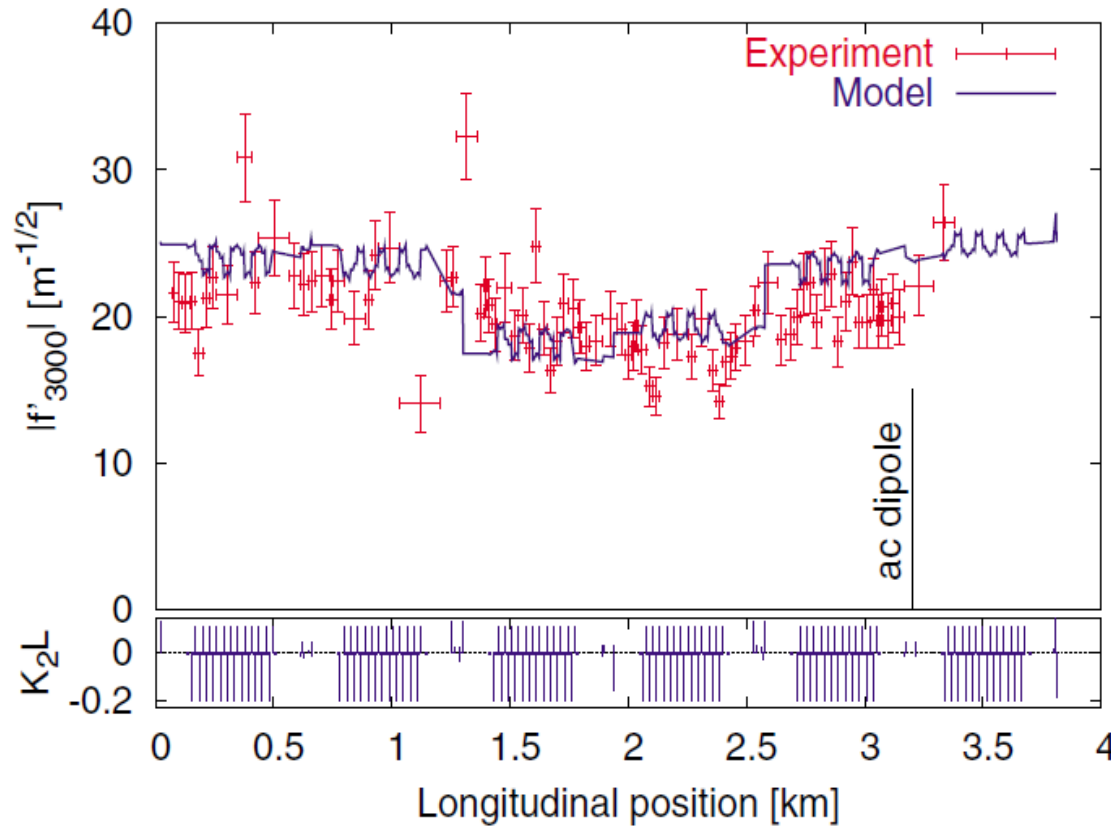
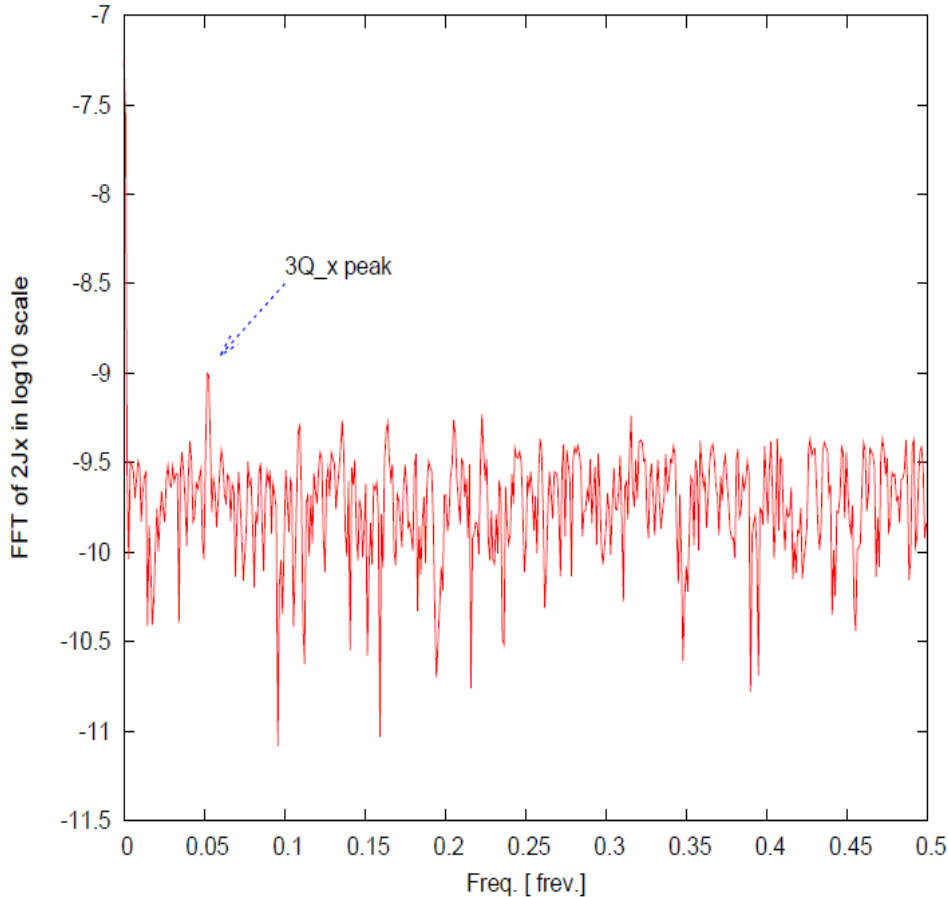


FIG. 4. (Color) Measurement of $|f'_{3000}|$ with an ac dipole. The bottom plot shows the sextupolar components of the ring.

- 3Qx resonance driving terms were measured with the TBT BPM data generated by AC dipole.
- Two methods to obtain the driving terms were tried in RHIC.
- This year the global correction of 3Qx was proposed but was not implemented due to tight machine time.

- 1) R. Tomas, et al., Phys. Rev. ST Accel. Beams 8, 024001 (2005).
- 2) R. Bartolini and F. Schmidt, Part. Accel. 59, 93(1998).

Extract h3000 from FFT of Jx



$$\begin{aligned}
 J_x(N) &= J_x + \frac{A_{21000}(2J_x)^{3/2}}{\sin(\pi\nu_x)} \cos(\hat{\phi}_{21000} + \phi_x + N2\pi\nu_x) \\
 &\quad + \frac{A_{10110}\sqrt{2J_x}2J_y}{\sin(\pi\nu_x)} \cos(\hat{\phi}_{10110} + \phi_x + N2\pi\nu_x) \\
 &\quad + \frac{3A_{30000}(2J_x)^{3/2}}{\sin(3\pi\nu_x)} \cos[\hat{\phi}_{30000} + 3(\phi_x + N2\pi\nu_x)] \\
 &\quad + \frac{A_{10020}\sqrt{2J_x}2J_y}{\sin[\pi(\nu_x - 2\nu_y)]} \cos[\hat{\phi}_{10020} + \phi_x - 2\phi_y + N2\pi(\nu_x - 2\nu_y)] \\
 &\quad + \frac{A_{10200}\sqrt{2J_x}2J_y}{\sin[\pi(\nu_x + 2\nu_y)]} \cos[\hat{\phi}_{10200} + \phi_x + 2\phi_y + N2\pi(\nu_x + 2\nu_y)] \\
 &\quad + O(b_3^2), \\
 J_y(N) &= J_y - \frac{2A_{10020}\sqrt{2J_x}2J_y}{\sin[\pi(\nu_x - 2\nu_y)]} \cos[\hat{\phi}_{10020} + \phi_x - 2\phi_y + N2\pi(\nu_x - 2\nu_y)] \\
 &\quad + \frac{2A_{10200}\sqrt{2J_x}2J_y}{\sin[\pi(\nu_x + 2\nu_y)]} \cos[\hat{\phi}_{10200} + \phi_x + 2\phi_y + N2\pi(\nu_x + 2\nu_y)] \\
 &\quad + O(b_3^2)
 \end{aligned} \tag{156}$$

- 1) Y. Luo, et al., in Proceedings of the 2007 PAC.
- 2) J. Bengtsson, SLS Note 9/97.

Backup slides: Perturbation Theory

Definitions:

$$\xi_{x,y}^{(1)} = \frac{\partial Q_{x,y}}{\partial \delta} \quad (4)$$

$$\xi_{x,y}^{(2)} = \frac{1}{2} \frac{\partial Q_{x,y}^2}{\partial \delta^2} \quad (5)$$

$$\xi_{x,y}^{(3)} = \frac{1}{6} \frac{\partial Q_{x,y}^3}{\partial \delta^3}. \quad (6)$$

The first order chromaticities:

$$\xi_{x,y}^{(1)} = \frac{1}{4\pi} \oint \beta_{x,y}(s) [\mp K_1(s) \pm K_2(s) D_x(s)] ds. \quad (7)$$

The second order chromaticities:

$$\xi_{x,y}^{(2)} = -\frac{1}{2} \xi_{x,y}^{(1)} + \frac{1}{8\pi} \oint [\mp K_1 \pm K_2 D_x] \frac{\partial \beta_{x,y}}{\partial \delta} ds + \frac{1}{8\pi} \oint \pm K_2 \beta_{x,y} D_x^{(2)} ds \quad (8)$$

The third order chromaticities:

$$\begin{aligned} \xi_{x,y}^{(3)} = & -\frac{1}{3} \xi_{x,y}^{(2)} \\ & + \frac{1}{24\pi} \oint [\mp K_1 \pm K_2 D_x] \frac{\partial^2 \beta_{x,y}}{\partial \delta^2} ds + \frac{1}{24\pi} \oint [\pm K_1 \mp K_2 D_x] \frac{\partial \beta_{x,y}}{\partial \delta} ds \\ & + \frac{1}{24\pi} \oint [\pm K_2 D_x^{(2)}] \frac{\partial \beta_{x,y}}{\partial \delta} ds + \frac{1}{24\pi} \oint [\pm K_2 D_x^{(3)}] \beta_{x,y} ds \\ & + \frac{1}{24\pi} \oint [\mp K_2 D_x^{(2)}] \beta_{x,y} ds + \frac{1}{24\pi} \oint \pm K_2 D_x^{(2)} \frac{\partial \beta_{x,y}}{\partial \delta} ds. \end{aligned} \quad (9)$$

where $D_x = \frac{\partial x_{\text{geo}}}{\partial \delta}$, $D_x^{(2)} = \frac{\partial x_{\text{geo}}^2}{\partial \delta^2}$, and $D_x^{(3)} = \frac{\partial x_{\text{geo}}^3}{\partial \delta^3}$.

Second order chromaticity, off-momentum β -beat and half-integer resonance driving terms

According to Ref. [2], the second order chromaticities are given by

$$\xi_{x,y}^{(2)} = -\frac{1}{2}\xi_{x,y}^{(1)} + \frac{1}{8\pi} \oint [\mp K_1 \pm K_2 D_x] \frac{\partial \beta_{x,y}}{\partial \delta} ds + \frac{1}{8\pi} \oint \pm K_2 \beta_{x,y} D_x^{(2)} ds, \quad (1)$$

where K_1 and K_2 are the strengths of quadrupoles and sextupoles. $\xi_{x,y}^{(1)}$ are the first order chromaticities. $\beta_{x,y}$ and $D_x^{(2)}$ are the betatron amplitude functions and horizontal dispersion. δ is the relative momentum deviation.

The second term in Eq. (1) is the dominant one which is determined by the off-momentum β -beat $\frac{\partial \beta_{x,y}}{\partial \delta}$. The relative off-momentum β -beat is

$$\frac{1}{\beta_{x,y}(s)} \frac{\partial \beta_{x,y}(s)}{\partial \delta} = \pm \frac{1}{2 \sin(2\pi Q_{x,y})} \oint [\beta_{x,y}(s') [\mp K_1(s') \pm K_2(s') D_x(s')] \cos(2|\phi_{x,y}(s) - \phi_{x,y}(s')| - 2\pi Q_{x,y}) ds'. \quad (2)$$

We define half integer resonance driving terms as

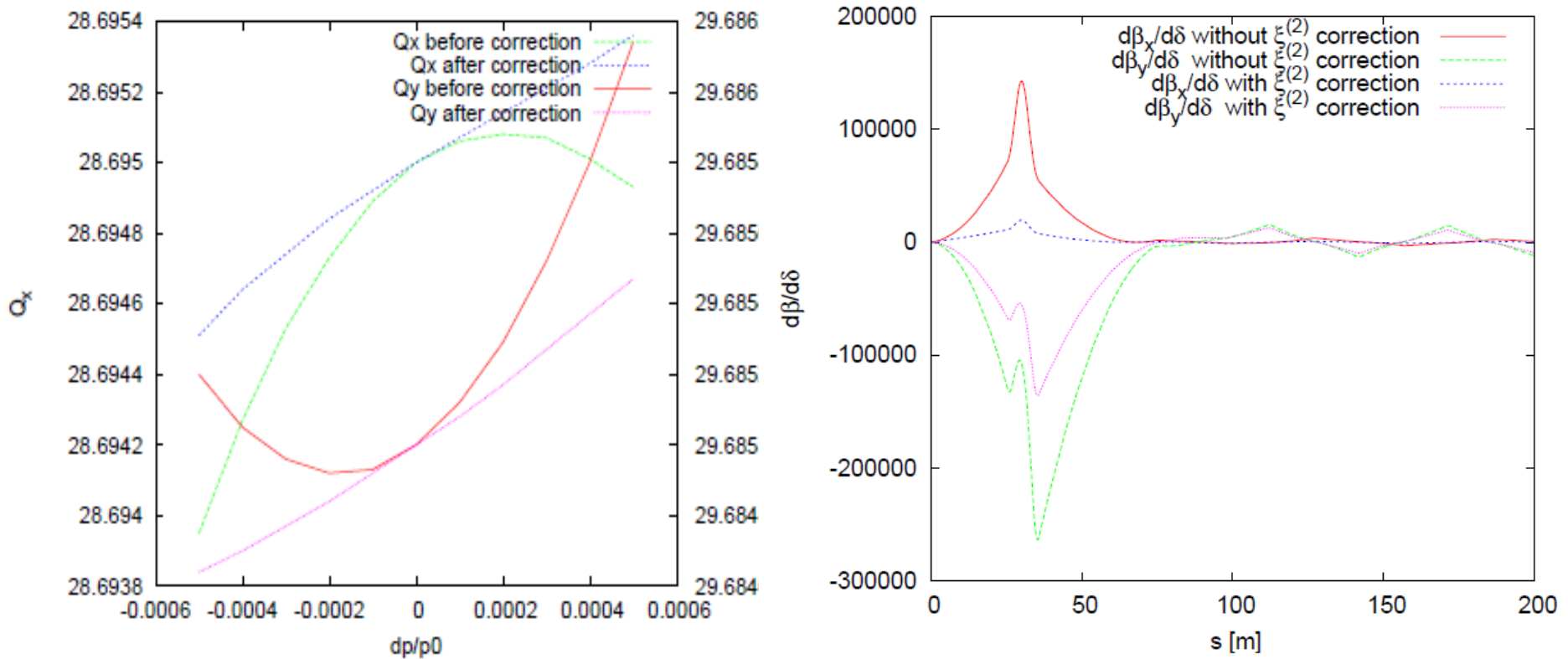
$$h_{20001} = \sum_i^N [-(K_1 L)_i + (K_2 D_x L)_i] \beta_{x,i} e^{-i2\phi_{x,i}}, \quad (3)$$

$$h_{00201} = \sum_i^N [(K_1 L)_i - (K_2 D_x L)_i] \beta_{y,i} e^{-i2\phi_{y,i}}. \quad (4)$$

Here h_{20001} and h_{00201} are the horizontal and vertical half integer resonance driving terms. Comparing Eq. (2) to Eqs. (3) and (4), the relative off-momentum β -beat is governed by the half integer resonance driving terms.

(Y. Luo, et al., BNL C-AD/AP/Note-426, 2011; Y. Luo, et al., PAC 2011.)

An example: Off-momentum β -beat reduced after second order chromaticity correction



(Y. Luo, et al., BNL C-AD/AP/Note-348, 2007)