

Non-SUSY scenarios with MET signatures

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CERN TH

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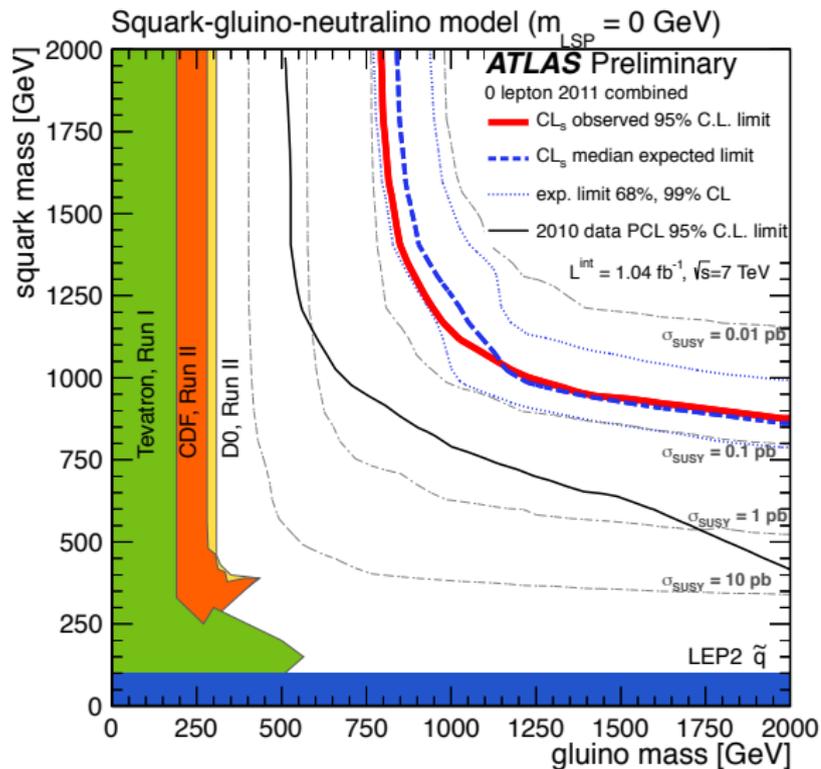
Outline

- ▶ SUSY, but not as we know it
- ▶ everything else
- ▶ a proposal: counting DM particles

SUSY, but not as we know it

SUSY, but not as we know it

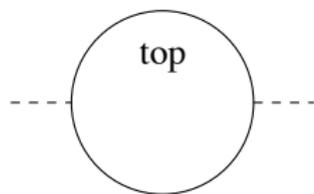
Isn't SUSY **dead** already?



SUSY, but not as we know it

There are now **strong bounds** on

- ▶ **gluino** mass
- ▶ **common** squark mass



SUSY, but not as we know it

Naturalness at the LHC \implies

- ▶ stop quarks
- ▶ a higgsino

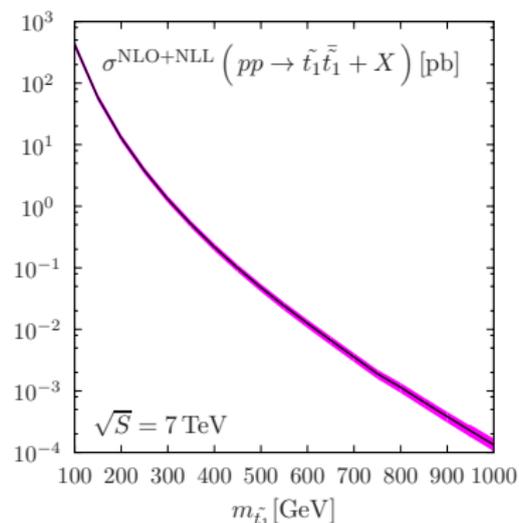
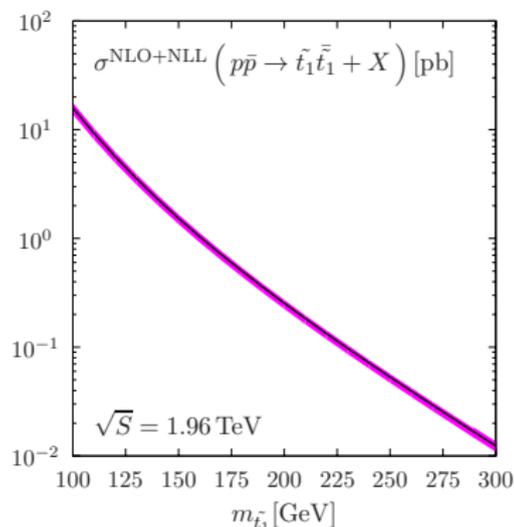
$$m_Z^2 \sim -m_{H_u}^2 - |\mu|^2$$

So could see only

- ▶ $2(\tilde{t} \rightarrow t\tilde{\chi}^0) \implies t\bar{t} + \cancel{E}_T$
- ▶ $2(\tilde{t} \rightarrow b\tilde{\chi}^\pm) \implies b\bar{b} + \cancel{E}_T$

SUSY, but not as we know it

Stop bounds are (and always will be) **weak**



Beenakker et al., 1006.4771

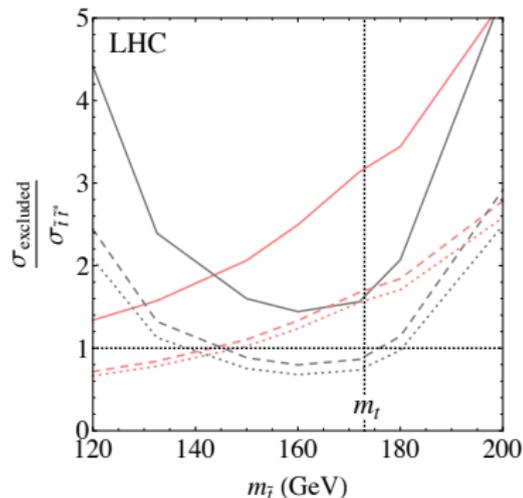
Small x-section at $m_{\tilde{t}} - m_{\tilde{\chi}} \gg m_t$

SUSY, but not as we know it

Stop bounds are (and always will be) weak
CDF, 2.7/fb: $m_{\tilde{t}} > 150\text{GeV}$

CDF, 0912.1308

Kats & Shih, 1106.0030



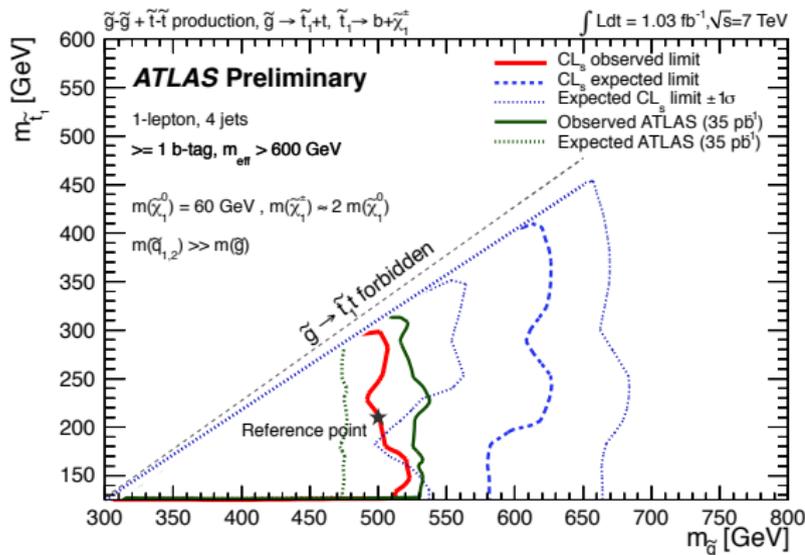
ATLAS, 3/fb: $m_{\tilde{t}} > 180\text{GeV}$

Kats & Shih, 1106.0030

SUSY, but not as we know it

Glauino: $\delta m_{\tilde{t}}^2 \sim \frac{8\alpha_s}{3\pi} M_3^2 \log \frac{\Lambda}{m_{\tilde{t}}}$

- ▶ Either M_3 or Λ not large
- ▶ $2(\tilde{g} \rightarrow \tilde{t}t)$



SUSY, but not as we know it

Are there **models** for this?

- ▶ More minimal SUSY
- ▶ Flavourful SUSY
- ▶ Partial SUSY

Cohen, Kaplan, & Nelson, 9607394

Dimopoulos & Giudice, 1995

Barbieri & al., 1004.2256, 1105.2296

Craig, Green, & Katz, 1103.3708

Gherghetta and Pomarol, 0302001

Sundrum, 0909.5430

BMG & Redi, 1004.5114

Gherghetta & al., 1104.3171

SUSY, but not as we know it

Is *R-parity* sufficient to prevent proton decay?

▶ $W \supset \frac{QQQL}{\Lambda}$

Yanagida & Sakai, 1982

Weinberg, 1982

▶ Consequences for *colliders* and *DM*

The LHC will not kill SUSY ...

... but it is unlikely to be the SUSY we* know and love.

* Well, some of us.

Everything else

Everything else

Focus on the **DM** motivation

Not so hard to build a **model** with a DM candidate

- ▶ Need a **neutral**, colourless particle
- ▶ Need it to be **long-lived**
- ▶ Need the right **relic density**

Everything else

Not so hard to build a model with a DM candidate

- ▶ Need a neutral, colourless particle. [Fiat](#)
- ▶ Need it to be long-lived. [Fiat](#)
- ▶ Need the right relic density. [Fiat](#)

Everything else

A **symmetry**, exact or accidental, makes the LSP long-lived or stable

- ▶ DM/SM charged/uncharged
- ▶ Exact \implies stable
- ▶ Accidental \implies long-lived

Everything else

Two particles are long-lived, but shouldn't be

- ▶ DM
- ▶ The **proton**

There may be an interesting interplay.

Everything else, e.g. I

In SUSY

- ▶ *R*-parity stabilizes DM
- ▶ *R*-parity makes the proton long-lived

Everything else, e.g. II

In the SM plus a **singlet**

- ▶ a Z_2 stabilizes DM
- ▶ unnaturalness ($\implies B$) makes the proton long-lived

McDonald, 1994

Everything else, e.g. III

In a **composite Higgs** model

- ▶ approximate B makes the proton long-lived
- ▶ $3B - n_c + \bar{n}_c \bmod 3$ (a $Z_3 \subset U(1)_B \times SU(3)_c$) stabilizes DM
- ▶ Z_2 vs. Z_3

Agashe & Servant, 0403143

Walker, 0907.3142

Agashe & al., 1003.0899

Agashe & al., 1012.4460

Everything else, e.g. IV

Non-Abelian symmetry, S_3

- ▶ Fields $\in \mathbf{1}', \mathbf{2}$
- ▶ Multiple stable components

Adulpravitchai, Batell & Pradler, 1103.3053

Conjecture: most models could be extended to have a DM candidate ...

... so what should we look for at the LHC & beyond?

Guess a TeV-scale Lagrangian

- ▶ Post-dict relic density
- ▶ Tricky at LHC: Strong vs. weak interactions.

Baltz & al., 0602187

Guess a TeV-scale Simplified Model

Arkani-Hamed & al., 0703088

Alves & al., 1105.2838

Guess a sub-TeV, effective Lagrangian

- ▶ Relate direct/indirect detection to collider searches
- ▶ There are a lot of operators

Goodman & al., 1005.1286

Bai, Fox & Harnik, 1005.3797

Goodman & al., 1008.1783

Davoudiasl & al., 1106.4320

How about **counting** invisible particles in collider **events**?

Giudice, BMG, & Mahbubani, 1108.1800

Why count DM?

- ▶ Multiple production \implies evidence for a symmetry
- ▶ Counting \implies nature of the symmetry
- ▶ e.g. Count mod 2
- ▶ Odd $\implies \mathbb{Z}_2$
- ▶ Odd \implies non-Abelian or $DM \neq \overline{DM}$

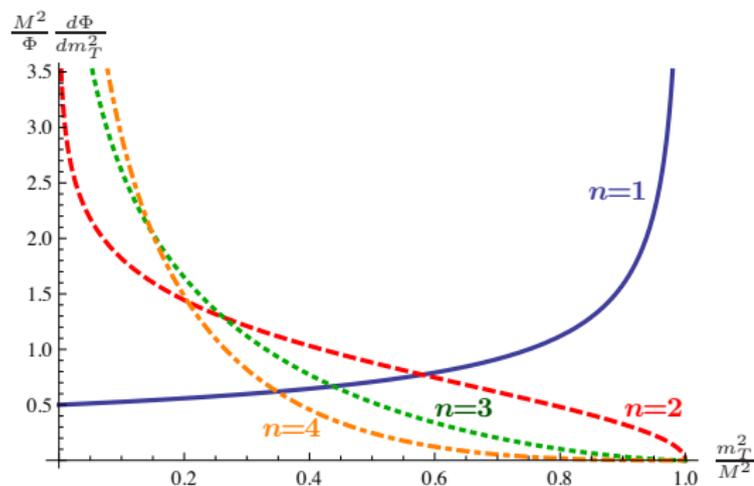
How to count DM?

- ▶ # invisibles, n , sets dimension of phase space
- ▶ observables are projections thereof
- ▶ strong dependence on n

Toy example

$M \rightarrow P + nX$, massless X

- ▶ m_T is maximised \implies invisibles **parallel** and **transverse**
- ▶ $\frac{d\Phi}{dm_T^2} \propto \left(1 - \frac{m_T^2}{M^2}\right) n^{-\frac{3}{2}}$



How to count DM?

- ▶ Strong dependence on n
- ▶ Other dependencies should be **small** or **known**
- ▶ Can tolerate **errors** up to $O(0.5)$

Other dependencies

- ▶ pdfs
- ▶ decay widths
- ▶ detector effects
- ▶ topology
- ▶ masses
- ▶ spins

e. g. Mass dependence

$M \rightarrow P + nX$, massive X

- ▶ m_T is maximised \implies invisibles transverse and relatively at rest
- ▶ $\frac{d\Phi}{dm_T^2} \propto \left(1 - \frac{m_T^2}{M^2}\right)^{\frac{3n}{2}-2}$
- ▶ Allow mass to **float** as a nuisance parameter

e. g. Spin dependence

e.g. $n=2$

▶ e.g. $h \rightarrow WW \rightarrow 2l2\nu, m_h = 2m_W$

▶ \implies measure $n < 2$

▶ e.g. $\mathcal{L} \supset A\Psi P_L \psi + B\Psi P_R \psi$

▶ $A \rightarrow \psi\Psi \rightarrow \psi\bar{\psi}B$

▶ \implies measure $n > 2$

or other spin effects

Wang & Yavin, 0802.2726

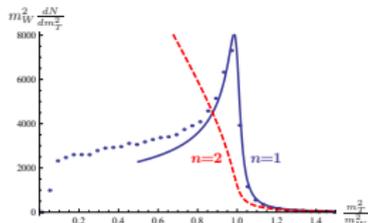
Rare **pathologies** more a blessing than a curse.

Strategy

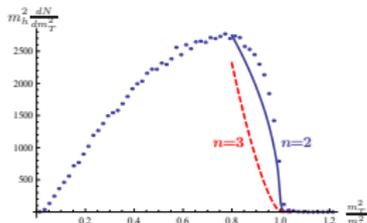
- ▶ Generate **phase space** with n invisibles
- ▶ Float invisible **mass** and **topology**
- ▶ Convolve **width/detector** effects
- ▶ **Fit endpoint** behaviour of appropriate observables (invariant masses, m_T , m_{T2})

SM examples

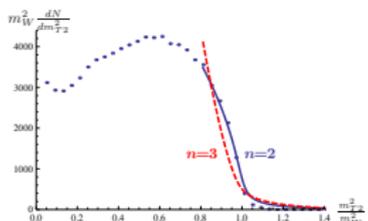
$$W \rightarrow l\nu, h \rightarrow 2l2\nu, 2t \rightarrow 2b2l2\nu$$



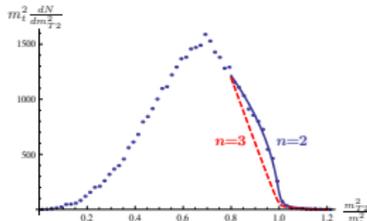
(a) $m_{l\nu}^2$ distribution for W production and leptonic decay at the Tevatron.



(b) $m_{h\nu}^2$ distribution for $h \rightarrow W^+W^-$ and leptonic W decay at the LHC, with $m_h = 180$ GeV.



(c) $m_{h\nu}^2$ distribution for $h \rightarrow W^+W^-$ and leptonic W decay at the LHC, with $m_h = 180$ GeV.



(d) $m_{t\nu}^2$ distribution for $t\bar{t}$ production and leptonic decay at the LHC.

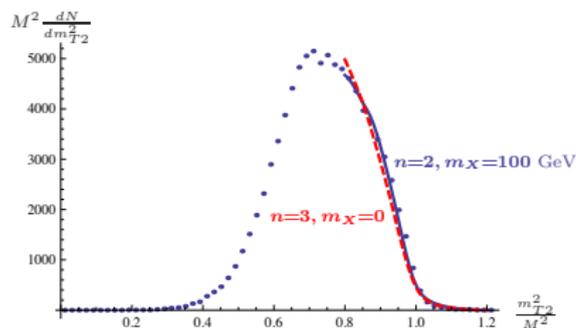
- ▶ No showering or detector effects
- ▶ Best fit correct in all cases

BSM example

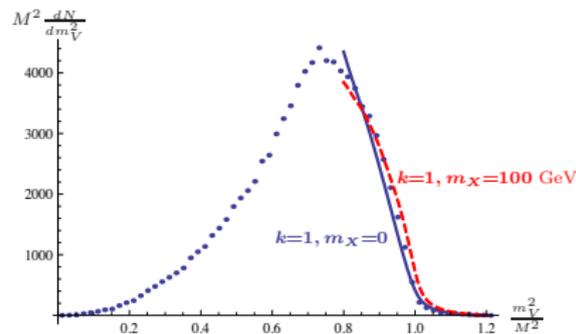
- ▶ $2(\tilde{u}_L \rightarrow u_L \tilde{\chi}_2^0 \rightarrow u_L \tilde{\ell}_R^+ \tilde{\ell}_R^- \rightarrow u_L \ell_R^+ \ell_R^- \tilde{\chi}_1^0)$
- ▶ Spin effects present

Barr, 0405052

- ▶ Float mass, not topology



(a) $m_{T_2}^2$ distribution for SUSY cascade process at the LHC.



(b) $m_{\tilde{\nu}}^2$ distribution for one leg of same SUSY cascade process at the LHC. Similar results are obtained for the other leg.

Summary

- ▶ Reports of the death of SUSY are greatly exaggerated
- ▶ The alternatives are **countless**
- ▶ **Count** invisible particles instead?