

NLO QCD corrections to Vector Boson Pair Production via Vector Boson Fusion

Giuseppe Bozzi

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Universität Karlsruhe

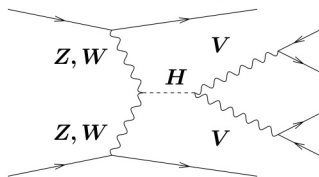
IFAE 2007
Napoli, 12.4.2007

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - NLO: real contributions
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

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Higgs production via VBF



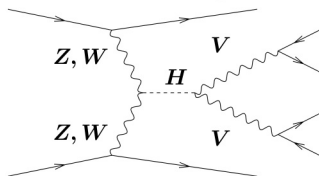
- $\sigma(qq \rightarrow qqH) \sim 0.2 \cdot \sigma(gg \rightarrow H)$ at the LHC
- Clean experimental signature
 - *two highly energetic outgoing jets*
 - *large rapidity interval between jets*
 - *no hadronic activity in the rapidity interval between jets*
- NLO QCD corrections moderate (5-10%)

[Han, Valencia, Willenbrock (1991)]

[Figy, Oleari, Zeppenfeld (2003)]

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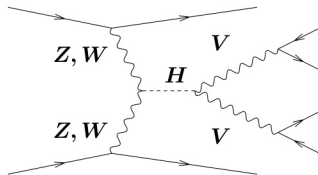
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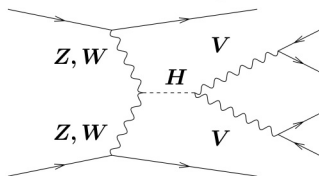
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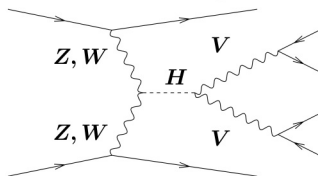
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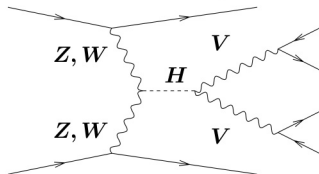
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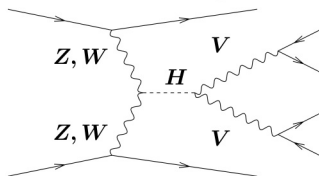
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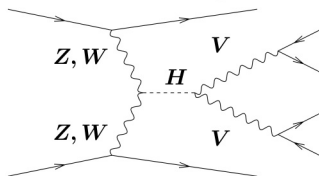
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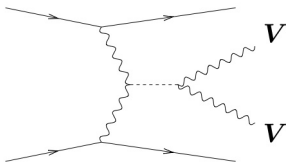
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VV production via VBF ($V = W^\pm, Z$)



- Background to Higgs production via VBF

- $\sigma(qq \rightarrow qqW^+W^-)$ between 3.5% and 15% of the Higgs signal for $115 \text{ GeV} \leq M_H \leq 160 \text{ GeV}$

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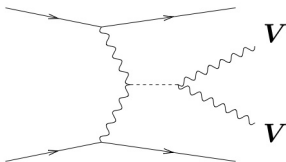
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- New Physics

- possible signal: enhancement of $qq \rightarrow qqVV$ over SM predictions at high \sqrt{s}
 - subprocess $V_L V_L \rightarrow V_L V_L$ intimately related to EWSB

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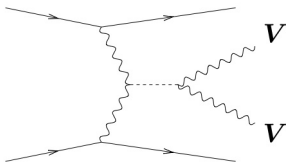
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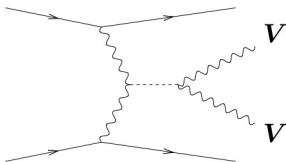
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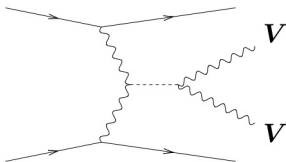
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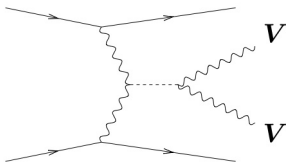
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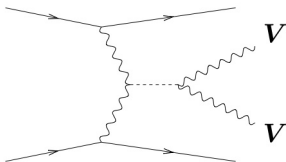
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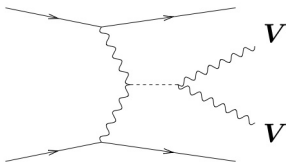
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- Multi-parton process: huge number of Feynman diagrams

- $2 \rightarrow 4$ for $qq \rightarrow qqVV$

- $2 \rightarrow 6$ for $qq \rightarrow qq l^+ l^- \nu_l \bar{\nu}_l, qq l^+ l^- l^+ l^-, qq l^+ l^- l^+ \nu_l$

→ how to speed up the evaluation?

- Suitable treatment of pentagon contributions

→ how to solve numerical instabilities?

→ Build a fully-flexible partonic Monte Carlo program allowing for

- computation of jet observables at NLO-QCD accuracy
- straightforward implementation of cuts

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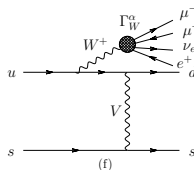
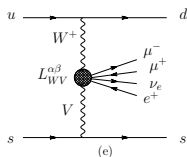
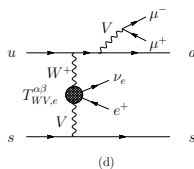
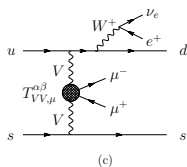
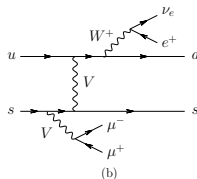
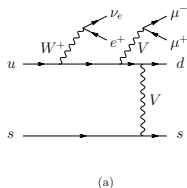
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- Up to 200 diagrams at LO!

- only $us \rightarrow ds$ shown

→ add $u\bar{s} \rightarrow d\bar{s}, \dots$

(a) same quark line

(b) different quark lines

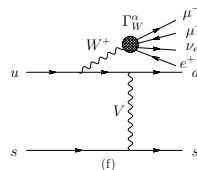
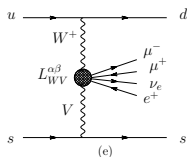
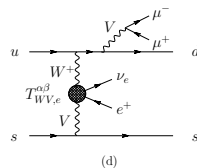
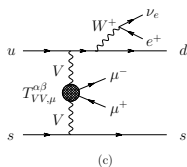
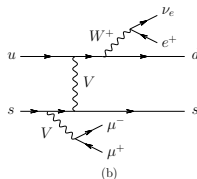
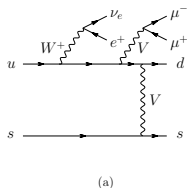
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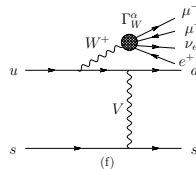
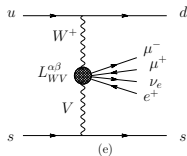
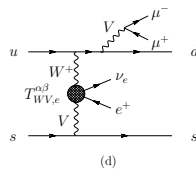
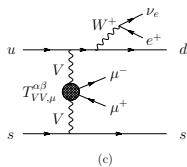
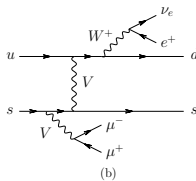
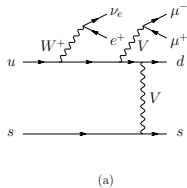
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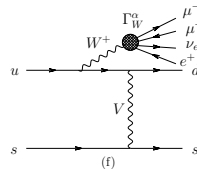
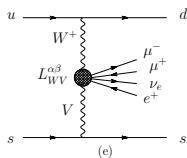
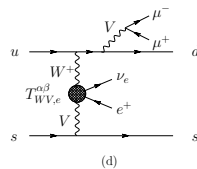
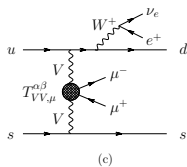
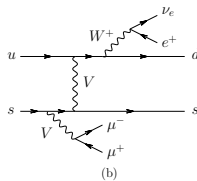
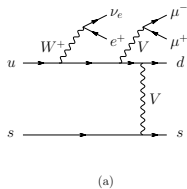
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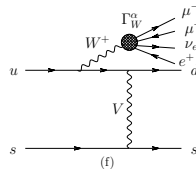
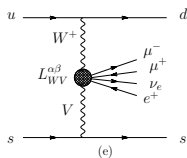
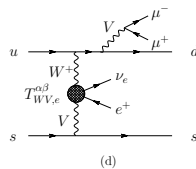
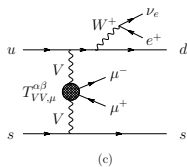
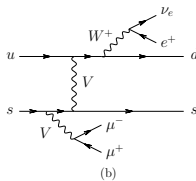
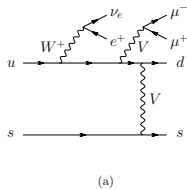
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Topologies @ LO (WZ case)



- Up to 200 diagrams at LO!

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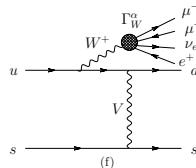
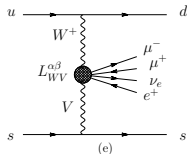
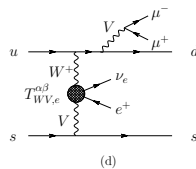
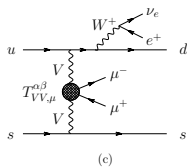
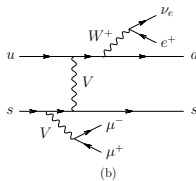
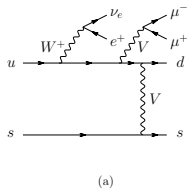
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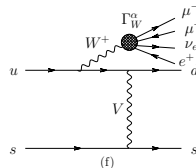
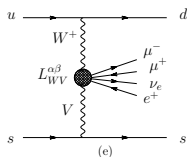
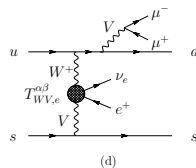
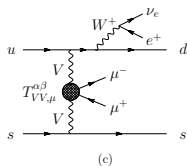
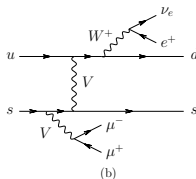
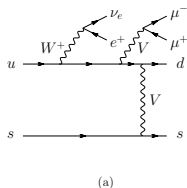
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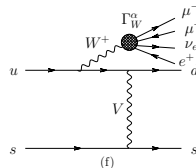
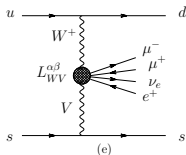
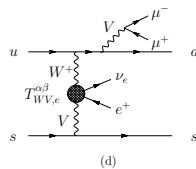
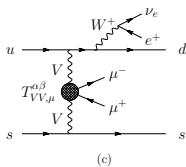
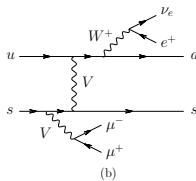
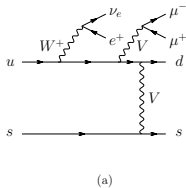
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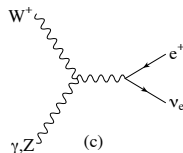
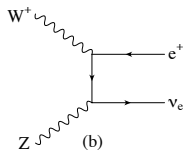
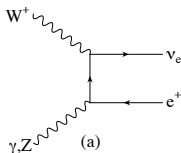
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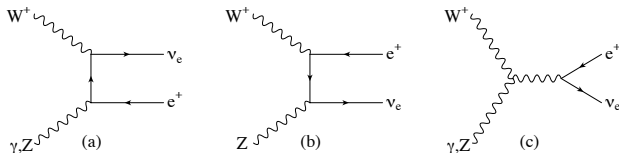
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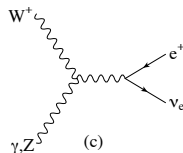
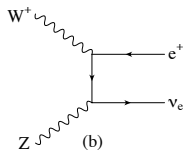
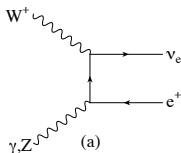
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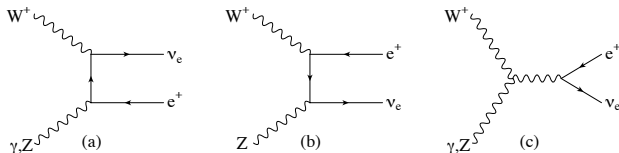
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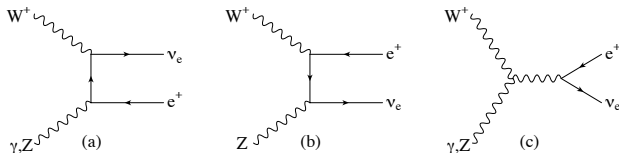
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Real corrections

- Attach a gluon to the quark lines in all possible ways
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 - **Soft** and **collinear** singularities
 - standard *Catani-Seymour* dipole subtraction
- [Catani, Seymour (1997)]
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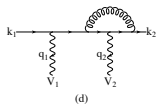
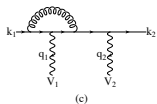
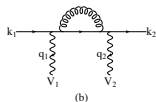
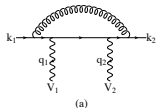
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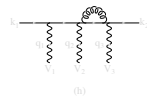
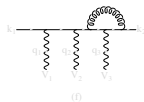
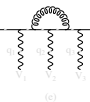
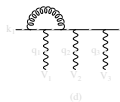
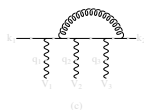
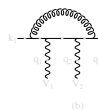
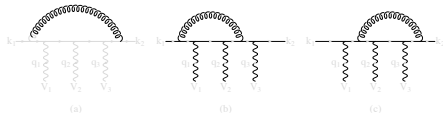
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

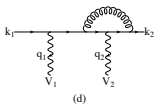
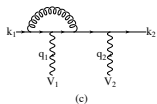
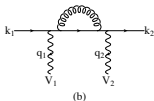
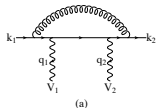
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- pentagons

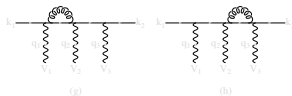
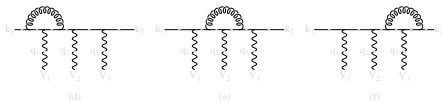
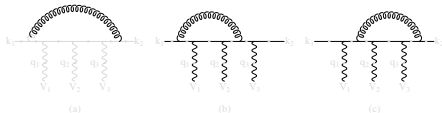
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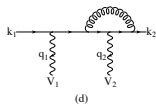
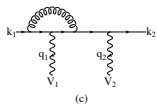
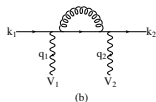
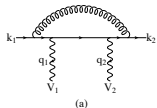
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

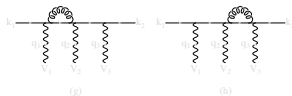
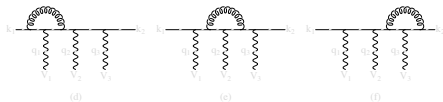
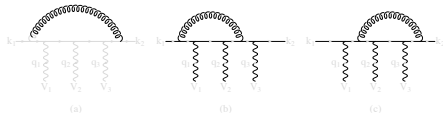
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

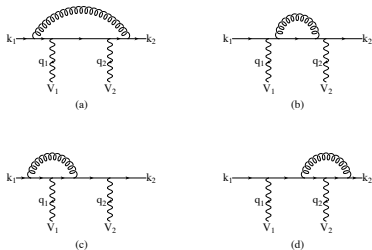
quark line with 3 bosons attached



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- triangles
- boxes
- pentagons

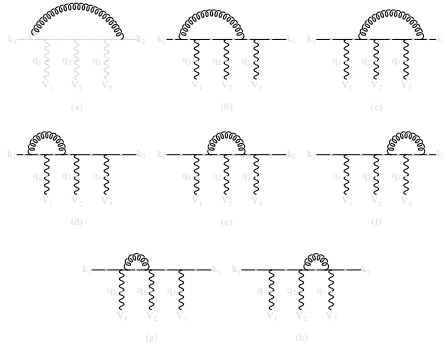
Topologies

quark line with 2 bosons attached



- self-energies
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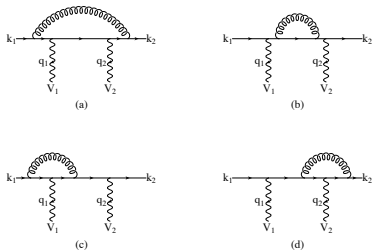
quark line with 3 bosons attached



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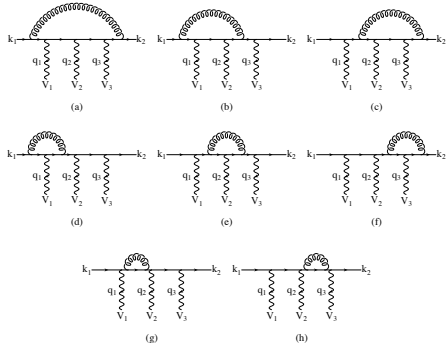
Topologies

quark line with 2 bosons attached



- self-energies
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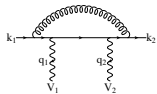
quark line with 3 bosons attached



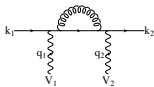
- self-energies
- triangles
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Topologies

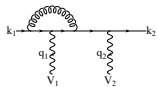
quark line with 2 bosons attached



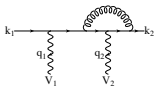
(a)



(b)



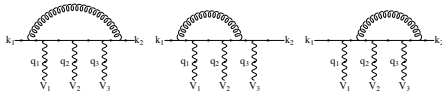
(c)



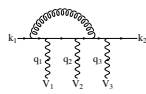
(d)

- self-energies
- triangles
- boxes

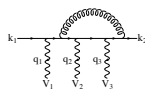
quark line with 3 bosons attached



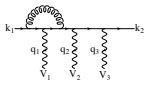
(a)



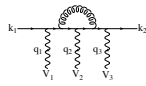
(b)



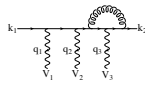
(c)



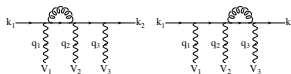
(d)



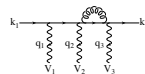
(e)



(f)



(g)

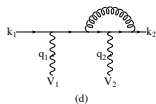
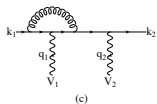
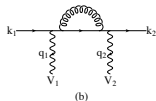
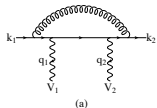


(h)

- self-energies
- triangles
- boxes
- pentagons

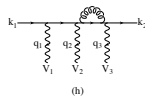
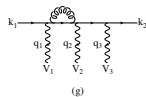
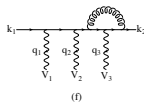
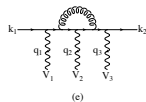
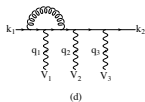
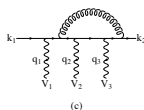
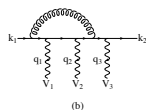
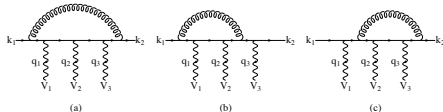
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

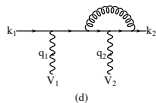
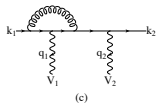
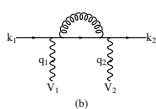
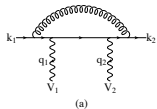
quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- pentagons

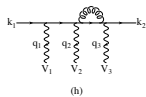
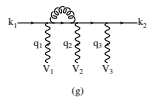
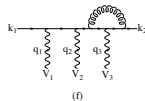
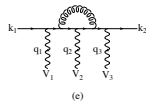
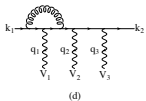
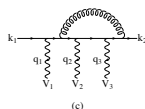
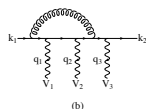
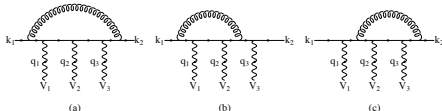
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

quark line with 3 bosons attached

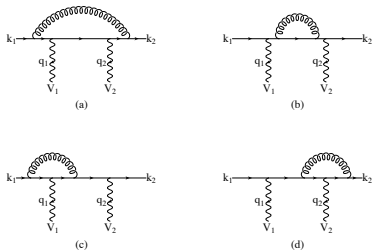


- self-energies
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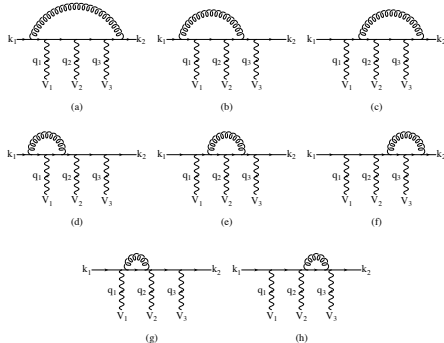
Topologies

quark line with 2 bosons attached



- self-energies
- triangles
- boxes

quark line with 3 bosons attached



- self-energies
- triangles
- boxes
- **pentagons**

Finite contributions

Summing up:

$$\mathcal{M}_V = \mathcal{M}_B \frac{\alpha_s(\mu_R)}{4\pi} C_F \left(\frac{4\pi\mu_R^2}{Q^2} \right)^\epsilon \Gamma(1 + \epsilon) \left[-\frac{2}{\epsilon^2} - \frac{3}{\epsilon} + c_{\text{virt}} \right] + \widetilde{\mathcal{M}}_V$$

- divergent part proportional to Born amplitude
 - exactly cancels the phase-space integral of the dipole terms
- finite term proportional to Born amplitude
- finite *non-universal* term $\widetilde{\mathcal{M}}_V$
 - can be computed in $d = 4$ dimensions
 - given in terms of the **finite parts** of the Passarino-Veltman $B_{ij}, C_{ij}, D_{ij}, E_{ij}$ coefficient functions

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Evaluation of $\widetilde{\mathcal{M}}_V$

- Divergent contributions in the expression of scalar integrals can generate finite terms by multiplying tensor coefficients with a $(d - 4)$ in the numerator

→ keep track of how the divergent contributions feed into the expression of tensor coefficients

[Oleari, Zeppenfeld (2003)]

- Two-, Three-, Four-point tensor integrals

→ computed through Passarino-Veltman reduction procedure

→ numerically stable in phase-space regions relevant for VBF

[Passarino, Veltman (1979)]

- Five-point tensor integrals

→ numerical instabilities if kinematical invariants (Gram determinant) become small

→ use Denny-Dittmaier reduction formalism

[Denny, Dittmaier (2005)]

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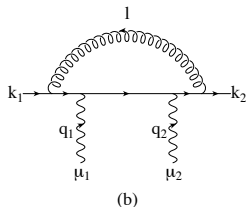
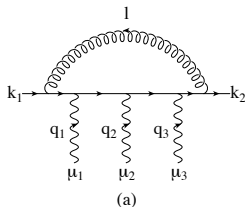
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Electromagnetic Ward Identities

Gauge invariance: a pentagon can be reduced to box integrals



$$\mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) \equiv \int \frac{d^d l}{(2\pi)^d} \gamma^\alpha \frac{1}{l + k_1 + \not{q}_{123}} \gamma_{\mu_3} \frac{1}{l + k_1 + \not{q}_{12}} \gamma^{\mu_2} \frac{1}{l + k_1 + \not{q}_1} \gamma_{\mu_1} \frac{1}{l + k_1} \gamma_\alpha \frac{1}{l}$$

$$q_1^{\mu_1} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_2 \mu_3}(k_1, q_1 + q_2, q_3) - \mathcal{D}_{\mu_2 \mu_3}(k_1 + q_1, q_2, q_3)$$

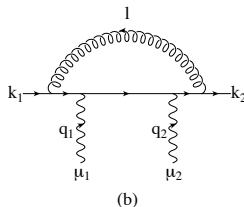
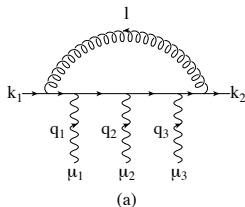
$$q_2^{\mu_2} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1 \mu_3}(k_1, q_1, q_2 + q_3) - \mathcal{D}_{\mu_1 \mu_3}(k_1, q_1 + q_2, q_3)$$

$$q_3^{\mu_3} \mathcal{E}_{\mu_1 \mu_2 \mu_3}(k_1, q_1, q_2, q_3) = \mathcal{D}_{\mu_1 \mu_2}(k_1, q_1, q_2) - \mathcal{D}_{\mu_1 \mu_2}(k_1, q_1, q_2 + q_3)$$

Express $\mathcal{E}_{\mu_1 \mu_2 \mu_3}$ ($\mathcal{D}_{\mu_1 \mu_2}$) as a sum of coefficients up to E_{ij} (D_{ij}) and verify the Ward identities \rightarrow **strong check** of the code!

Electromagnetic Ward Identities

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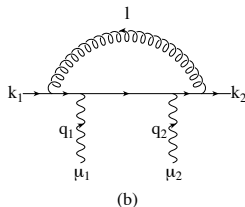
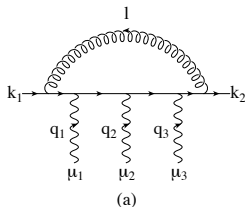
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Electromagnetic Ward Identities

Gauge invariance: a pentagon can be reduced to box integrals



$$\mathcal{E}_{\mu_1\mu_2\mu_3}(k_1, q_1, q_2, q_3) \equiv \int \frac{d^d l}{(2\pi)^d} \gamma^\alpha \frac{1}{l + k_1 + \not{q}_{123}} \gamma_{\mu_3} \frac{1}{l + k_1 + \not{q}_{12}} \gamma^{\mu_2} \frac{1}{l + k_1 + \not{q}_1} \gamma_{\mu_1} \frac{1}{l + k_1} \gamma_\alpha \frac{1}{l^2}$$

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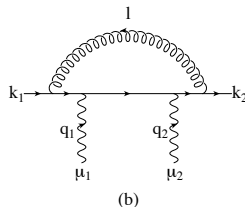
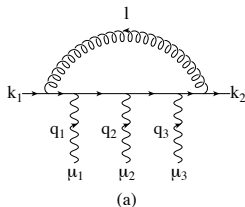
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“True” pentagons

- Loop amplitudes eventually contracted with leptonic currents

- Example: $W^+(q_+)$, $W^-(q_-)$, $\gamma/Z(q_0)$ with leptonic decays J_+ , J_-

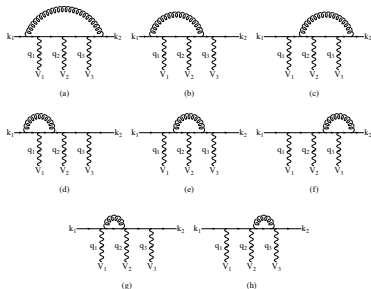
- $M_5 = J_+^{\mu_1} J_-^{\mu_2} \tilde{\mathcal{P}}_{\mu_1\mu_2\mu_3}(k_1, q_+, q_-, q_0)$

- Project J_{\pm} on the respective momenta ($J_{\pm}^{\mu} = x_{\pm} q_{\pm}^{\mu} + r_{\pm}^{\mu}$), so that the vectors r_{\pm} , in the center-of-mass system of the W pair, have zero time component ($r_{\pm} \cdot (q_+ + q_-) = 0$)

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- Ward Identities reduce magnitude of coefficients multiplying pentagon loops and thus, the overall pentagon

contribution to cross section



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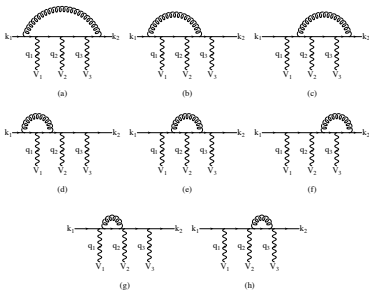
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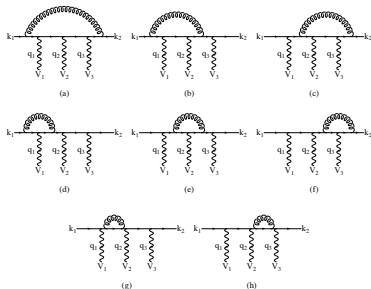
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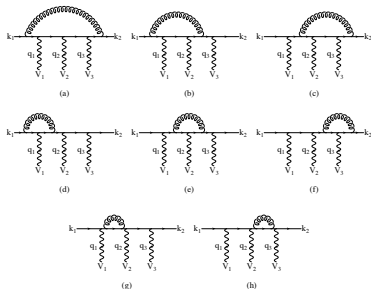
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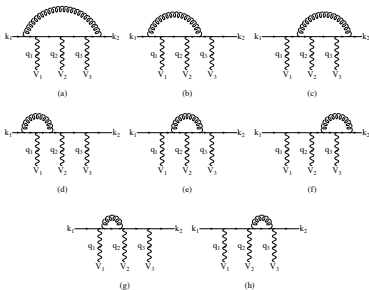
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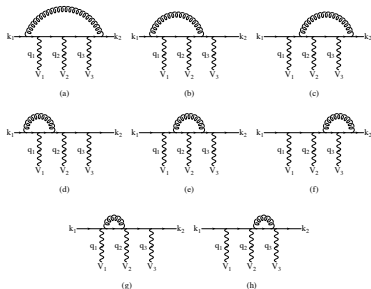
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Numerical stability of pentagon contributions

- Gauge-check procedure

- Identify the fraction f of events for which numerical pentagon reduction violates Ward identity by more than 10%
- Discard these points for the calculation of the finite parts
- Correct the remaining pentagon contributions by a factor $1/(1 - f)$
- Error induced by this approximation: far below numerical accuracy of Monte Carlo program

- PV formalism: $f \sim 15\%$

- DD formalism: $f \sim 0.1\%$

→ Pentagons under control using DD formalism!

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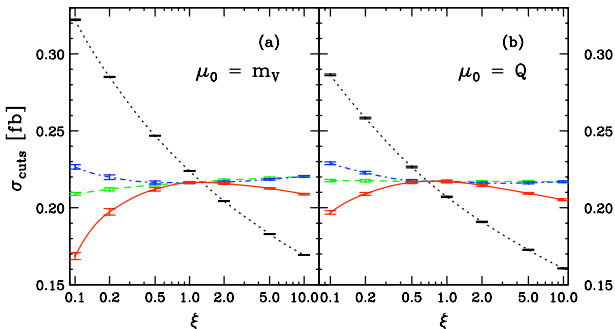
Outline

- 1 Motivation
 - Why Vector Boson Fusion?
- 2 Elements of the calculation
 - Tree-level features
 - NLO: real contributions
 - NLO: virtual contributions
- 3 Selected results
 - Differential distributions at the LHC

VBF cuts

Tagging Jets	$p_{Tj} \geq 20 \text{ GeV}, \quad y_j \leq 4.5$ $\Delta y_{jj} = y_{j_1} - y_{j_2} > 4,$ $y_{j_1} \cdot y_{j_2} < 0$
Charged Leptons	$p_{Tl} > 20 \text{ GeV}, \quad \eta_l \leq 2.5$ $y_{j,min} < \eta_l < y_{j,max}$ $\Delta R_{jl} \geq 0.4$
Higgs on/off	$M_{VV} > M_H + 10 \text{ GeV}$ (WW,ZZ continuum only)

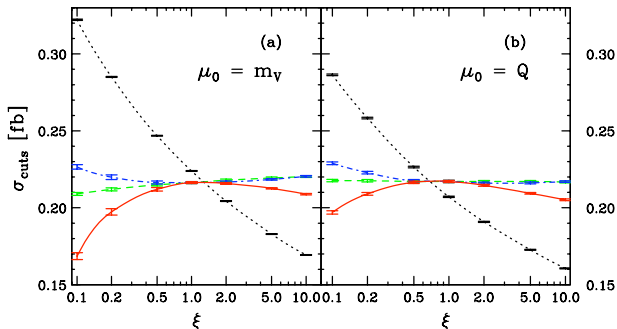
Scale dependence - total σ (WZ case)



• Two possible scales

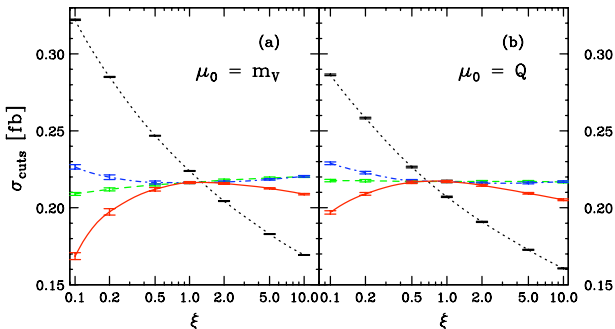
- $m_V = (m_Z + m_W)/2$
- $Q = \text{momentum-transfer of exchanged vector boson in VBF graphs}$
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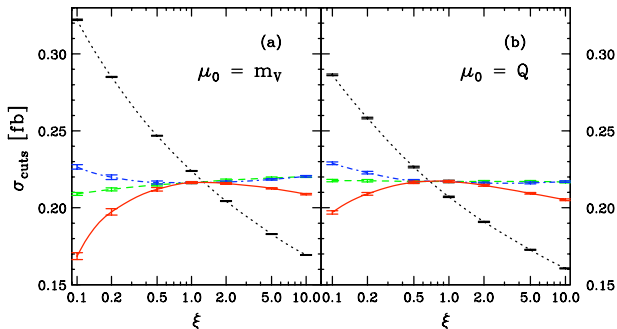
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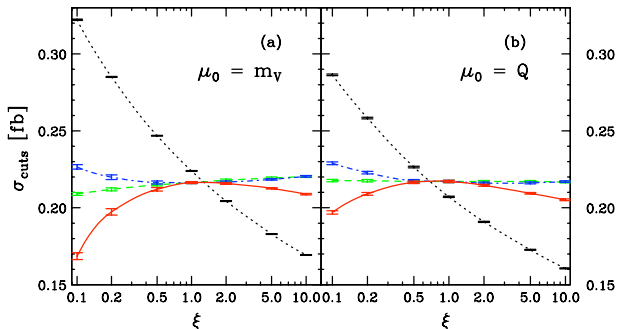
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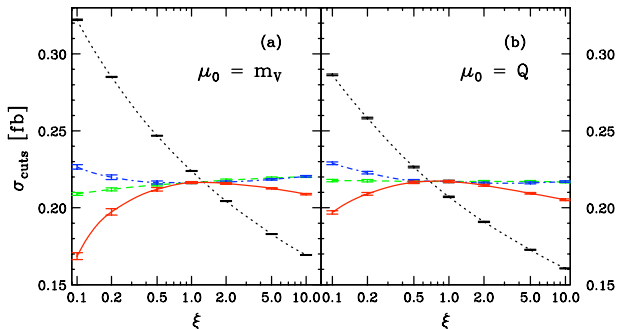
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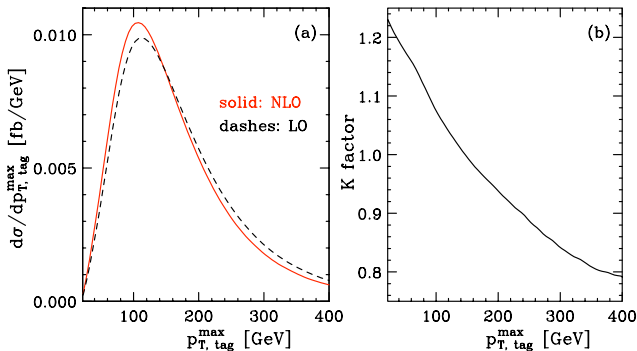
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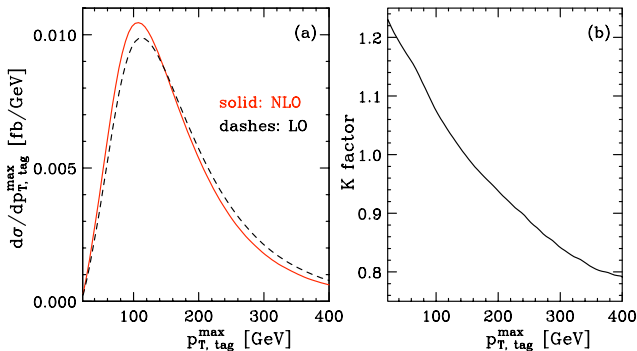
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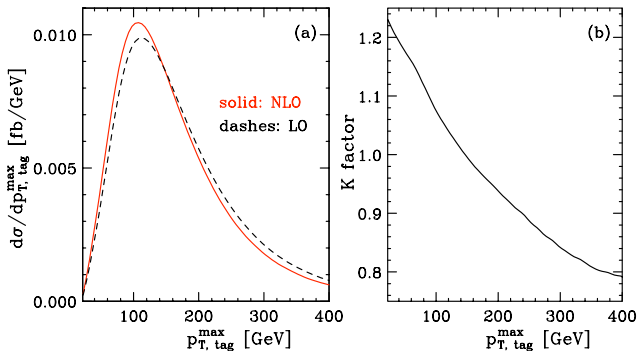
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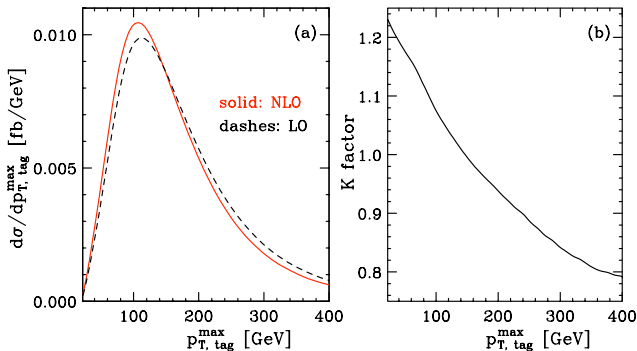
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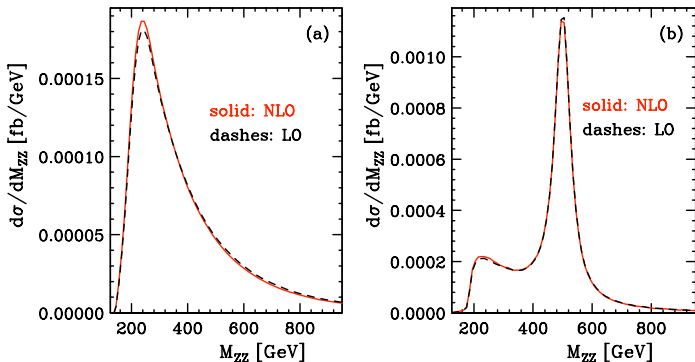
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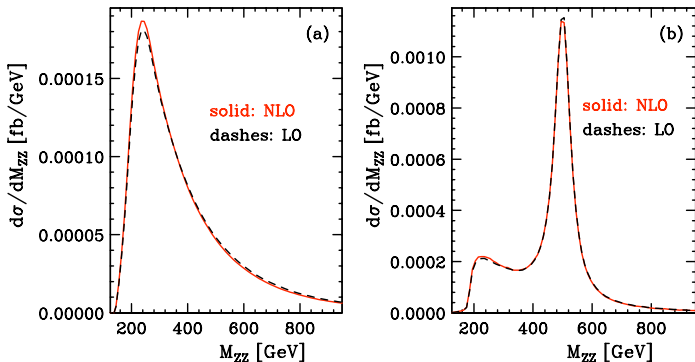
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Invariant mass - lepton pairs (ZZ case)



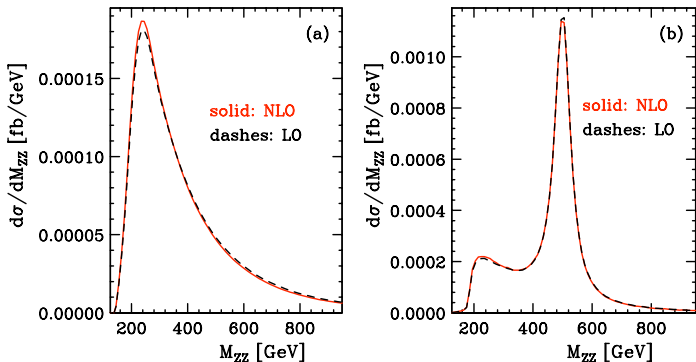
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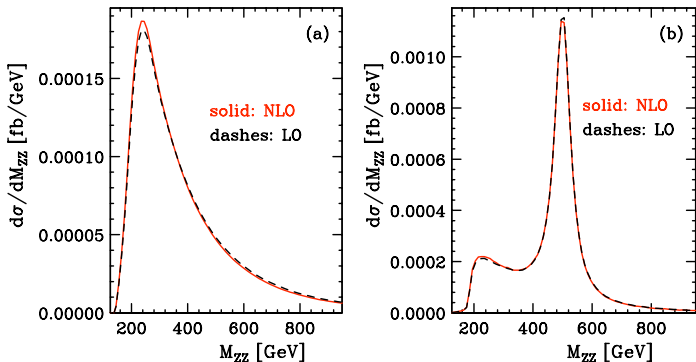
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- NLO corrections under excellent control (modest K-factors and scale dependence)

- Outlook: VBFNLO*

Monte Carlo program of relevant VBF processes at NLO QCD

Summary

- Fully-flexible parton-level Monte Carlo program with NLO QCD cross sections and distributions for

$$pp \rightarrow W^+ W^- jj \quad pp \rightarrow ZZjj \quad pp \rightarrow W^\pm Zjj$$

including leptonic decays

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- Outlook: **VBFNLO**
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