

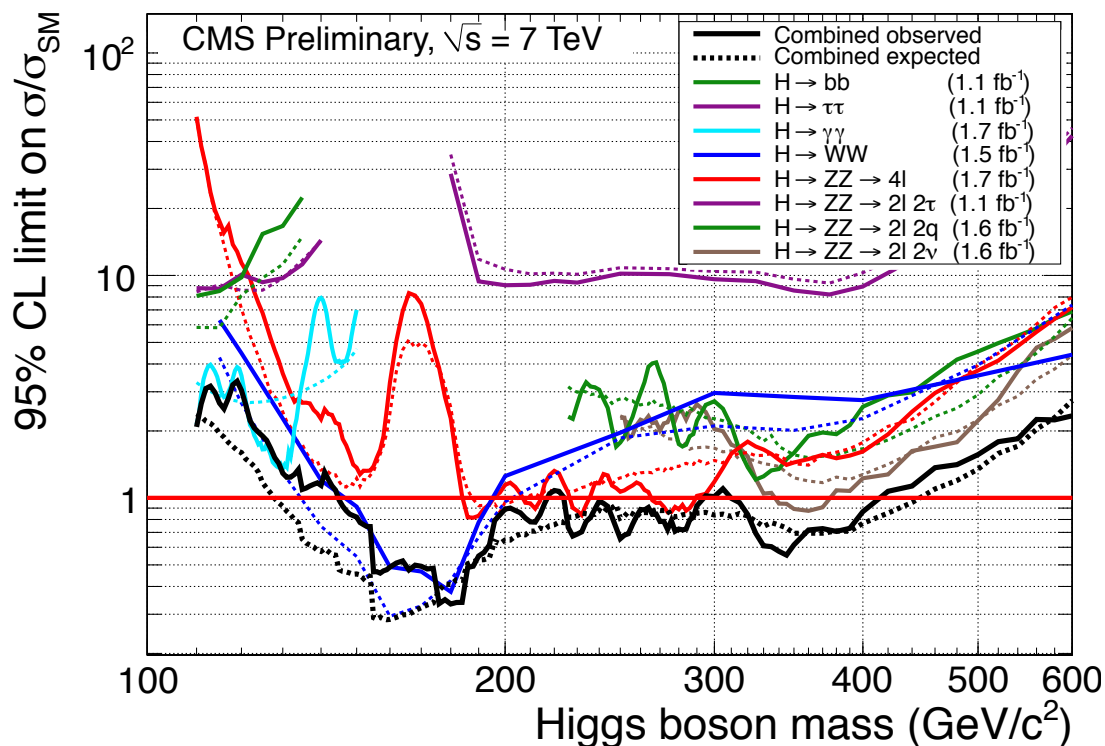
BSM: Implications from/for Higgs searches

G. Bélanger
LAPTH

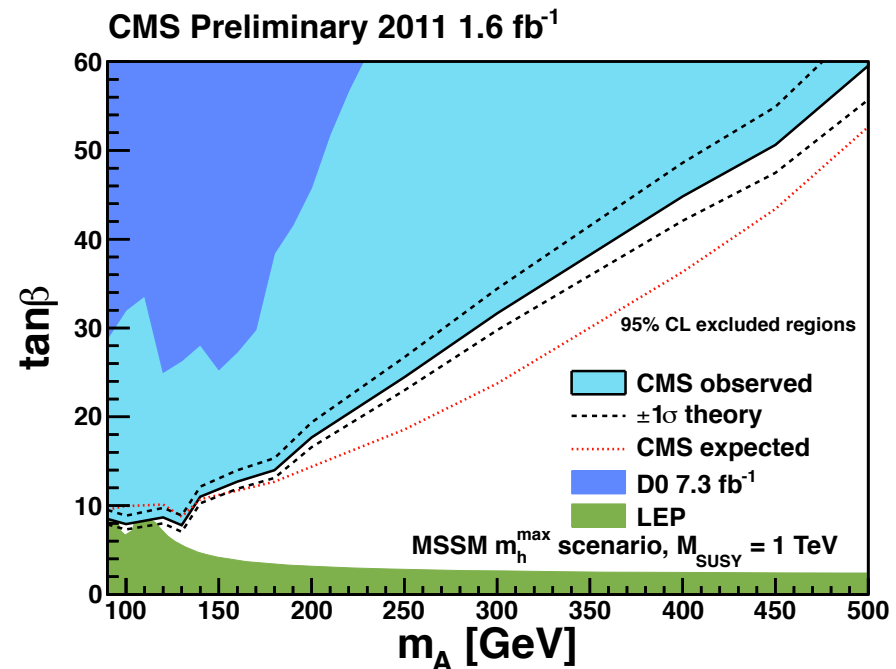
Contribution to Les Houches -
A .Albornoz-Vasquez, GB, A. Belyaev, M. Brown, F.
Boudjema, G. DaSilva, G. Drieu La Rochelle, R. Godbole, F.
Mahmoudi, F. Moortgat, A. Pukhov ...

Introduction

- ATLAS and CMS have excluded a SM - Higgs in the range $144\text{GeV} < m_h < 440\text{GeV}$ with 1.7fb^{-1}
- With $\sim 10\text{fb}^{-1}$ expect to close the range $114 < m_h < 144\text{GeV}$



- In extensions of the SM, these limits cannot be applied directly, new physics can affect the Higgs couplings to $WW, bb, tt, \gamma\gamma, gg$
- Higgs can decay invisibly thus reducing all signatures in SM particles
- Higgs can have new decay channels
- In some extensions of SM, Higgs searches in different channels are important, e.g. search for heavy Higgs doublet leads to constrain in the m_A - $\tan\beta$ plane in SUSY



Higgs searches and BSM

- The SUSY case
 - Comparison with SM
 - $gg-h-\gamma\gamma$
 - Pseudoscalar and light DM
- Invisible Higgs (MSSM and extensions)
- BSM Higgs
- Conclusion

The SUSY case

- Modification of couplings
 - low M_A
 - stop sector
 - chargino sector
- Pseudoscalar Higgs search and the link to dark matter

Higgs couplings in MSSM

- Non decoupling case - low m_A
- Ignoring radiative corrections

$$\cos^2(\beta - \alpha) = \frac{M_h^2(M_Z^2 - M_h^2)}{M_A^2(M_H^2 - M_h^2)}$$

- Higgs coupling to SM particles`

$$\frac{g_{W^+W^-h}}{g_{SM}} = \sin(\beta - \alpha), \quad \frac{g_{\bar{t}th}}{g_{SM}} = \frac{\cos \alpha}{\sin \beta}, \quad \frac{g_{\bar{b}bh}}{g_{SM}} = -\frac{\sin \alpha}{\cos \beta}.$$

- Modification of couplings

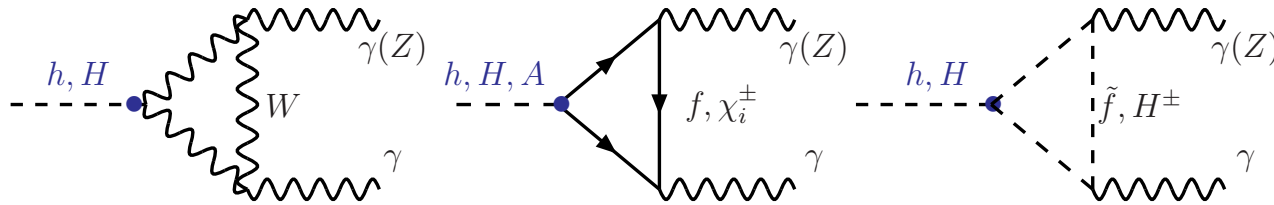
$$\tan \alpha \tan \beta = -(1 + r)$$

$$r \simeq \frac{2M_Z^2}{M_A^2} \frac{\tan^2 \beta - 1}{\tan^2 \beta + 1} \geq 0$$

$$R_{VVh} \simeq 1 - \frac{r^2}{2} \frac{\tan^2 \beta}{(1 + \tan^2 \beta)^2} \quad R \simeq 1 - \frac{r}{1 + \tan^2 \beta} \quad R_{bbh} \simeq 1 + r \frac{\tan^2 \beta}{1 + \tan^2 \beta}$$

- Increase in bbh , more modest corrections (decrease) in tth, WWh

Loop-induced couplings



- $h\gamma\gamma$: dominant contribution: W loop , top loop opposite sign
- If hWW coupling not modified, $h\gamma\gamma$ not much affected
- Total width dominated by bb , larger than in SM

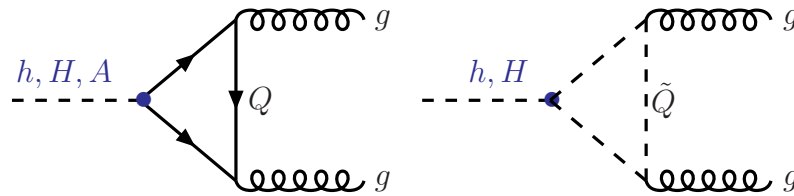
$$b - \frac{m_b/v}{1 + \Delta m_b} \bar{b} \left[g_b^h \left(1 - \frac{\Delta m_b}{\text{tg}\alpha \text{tg}\beta} \right) h \right]$$

- branching $h\gamma\gamma$ can drop

$$R_{\gamma\gamma} = \frac{BR^{SUSY}(h \rightarrow \gamma\gamma)}{BR^{SM}(h \rightarrow \gamma\gamma)}$$

- inNMSSM light Higgs can have large singlet component and much reduced coupling to b-quarks **Br($\gamma\gamma$) increases** up to factor 6
 - Ellwanger 1012.1201

ggH



- hgg dominant contribution top loop
- Expect small decrease at low m_A
- Stop contribution can be large

– Djouadi, PLB345(98) 101

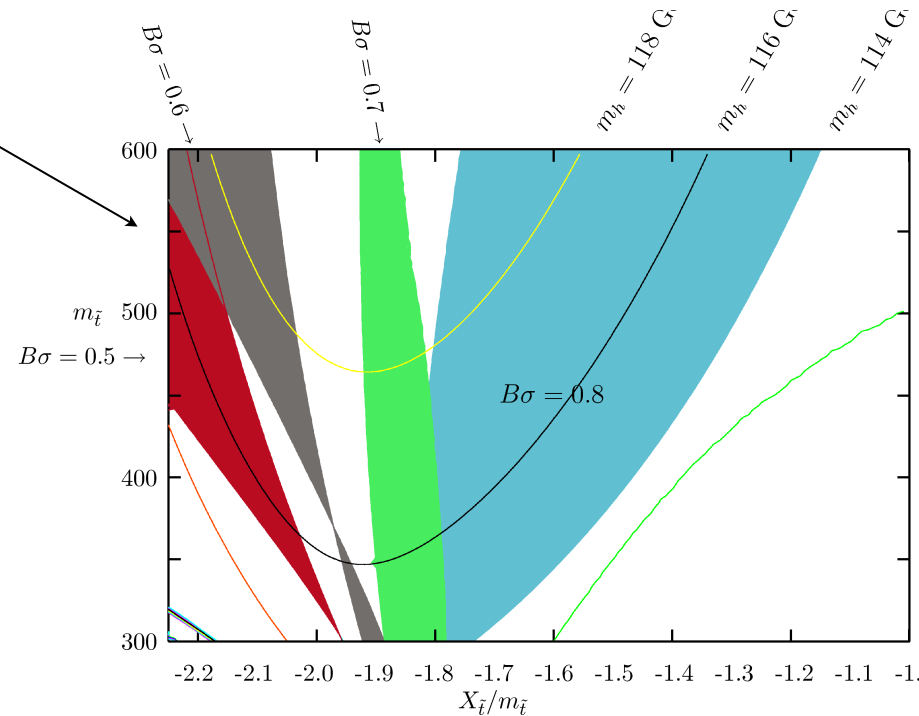
$$V_{\tilde{t}_1 \tilde{t}_1 h} \simeq \frac{g}{M_W} \left(\sin^2(2\theta_{\tilde{t}}) \frac{(m_{\tilde{t}_1}^2 - m_{\tilde{t}_2}^2)}{4} + m_t^2 \right) + M_Z^2 \cos(2\beta) \left(\left(\frac{1}{2} - \frac{2}{3} \sin^2 \theta_W \right) \cos^2 \theta_{\tilde{t}} + \frac{2}{3} \sin^2 \theta_W \sin^2 \theta_{\tilde{t}} \right)$$

- no-mixing case: stop interferes **constructively** with top
--> increase in hgg, decrease(modest) in $h\gamma\gamma$
- large mixing : stop interferes **destructively** with top
--> decrease in hgg

gg-H- $\gamma\gamma$ and SUSY

- Large corrections for light stop
 - GB, Boudjema, Sridhar, hep-ph/9904348
- Effect of top/stop on gg $\gamma\gamma$
- Effect of chargino on h $\gamma\gamma$ less important
 - GB, Boudjema, Donato, Godbole, Rosier, hep-ph/0002039
- Decoupling limit (large m_A) : hWW, htt, hbb SM-like (ignoring radiative correction effects which can lead to large modifications in hbb)

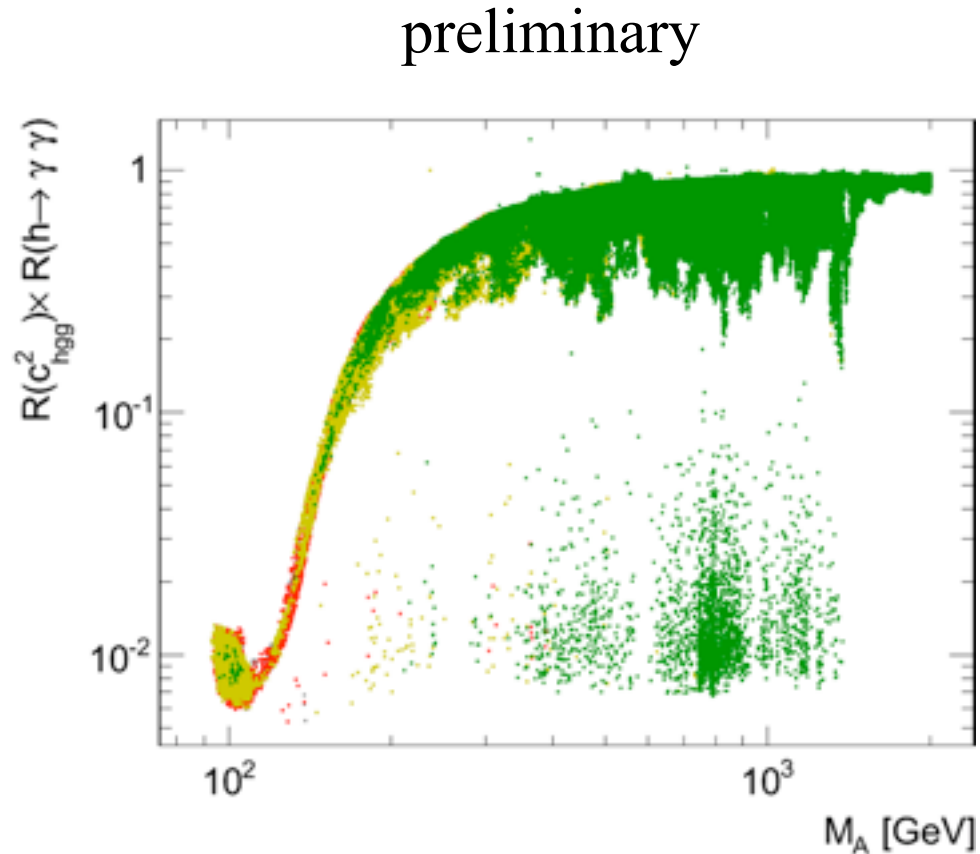
- $m_A=400\text{GeV}$, $\tan\beta=30$
 $M_{\text{susy}}=1\text{TeV}$



Low, Shalgar, 0901.0266

Higgs in MSSM

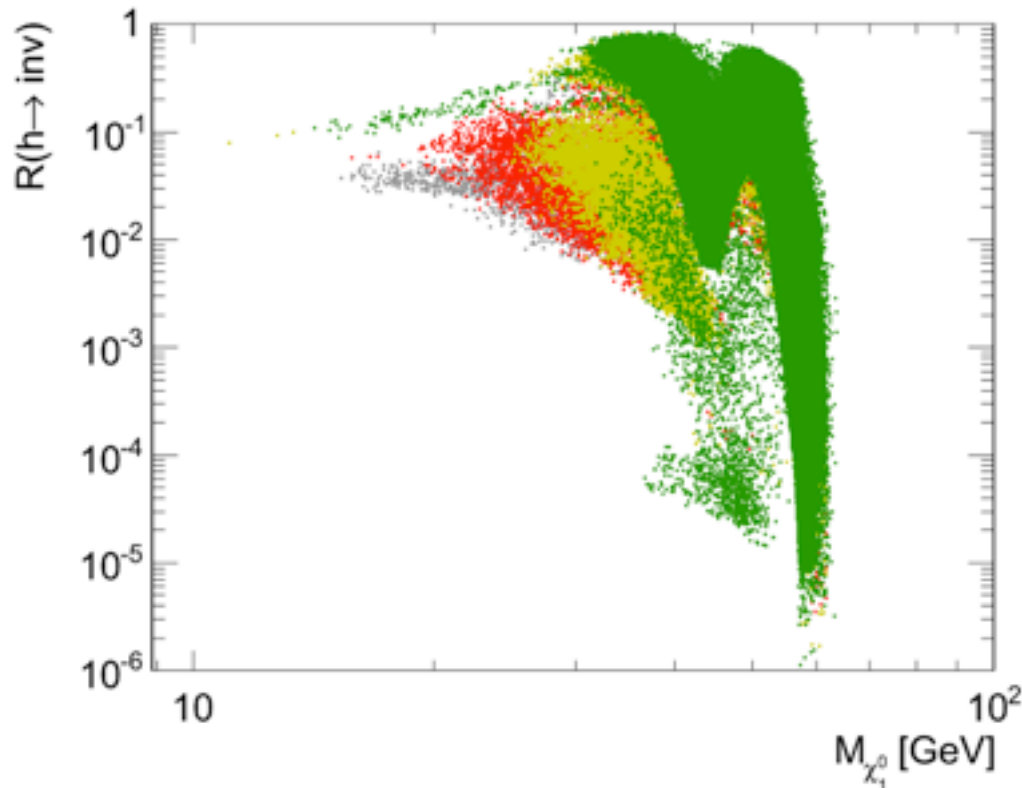
- Scan of MSSM with 11 free parameters
 - $M_1, M_2, M_3, \mu, \tan\beta, M_A, A_t, M_{IL}, M_{IR}, M_{q1}, M_{q3}$
 - MCMC approach
 - Higgs searches (HiggsBounds)
 - Flavour physics
 - DM observables
- LHC exclusion squarks not included
- Drop in $gg-h-\gamma\gamma$ signal at low M_A
 - partially excluded by DM constraints



D. Albornoz Vasquez

Higgs in MSSM

- Can have large invisible width \rightarrow reduction in all channels
- Possible modest enhancement of $WW \rightarrow h \rightarrow \tau\tau$

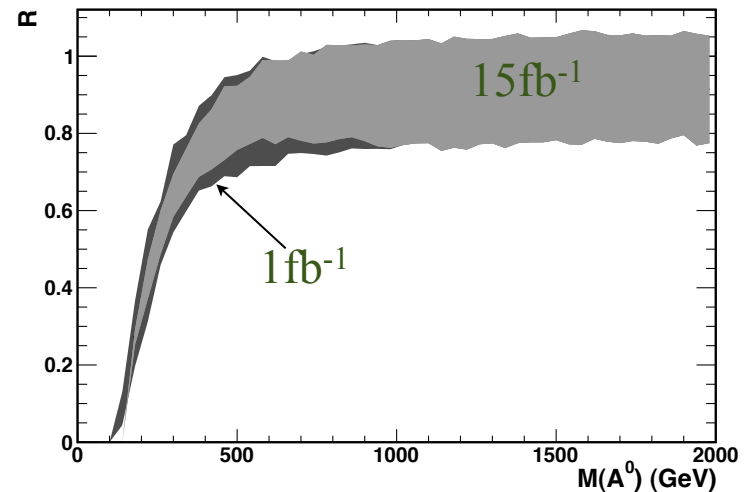


Yellow/red
Direct detection+
gamma-ray Fermi

Light Higgs in pMSSM

- Global scan in pMSSM
 - Mass limits (includes LHC 1fb^{-1})
 - Higgs searches (HiggsBounds)
 - Flavour physics
 - DM relic density
- If $m_A > 500\text{ GeV}$, $R > 0.7$
 - to rule out/explore large fraction of MSSM parameter space need slightly more luminosity than SM

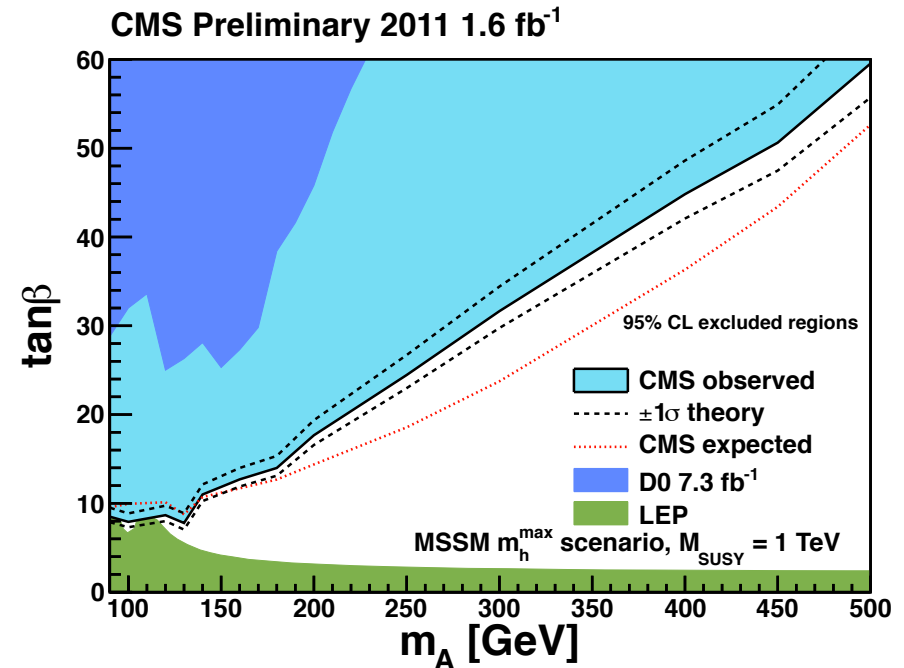
$$R = \frac{\sigma_{gg\gamma\gamma}^{\text{MSSM}}}{\sigma_{gg\gamma\gamma}^{\text{SM}}}$$



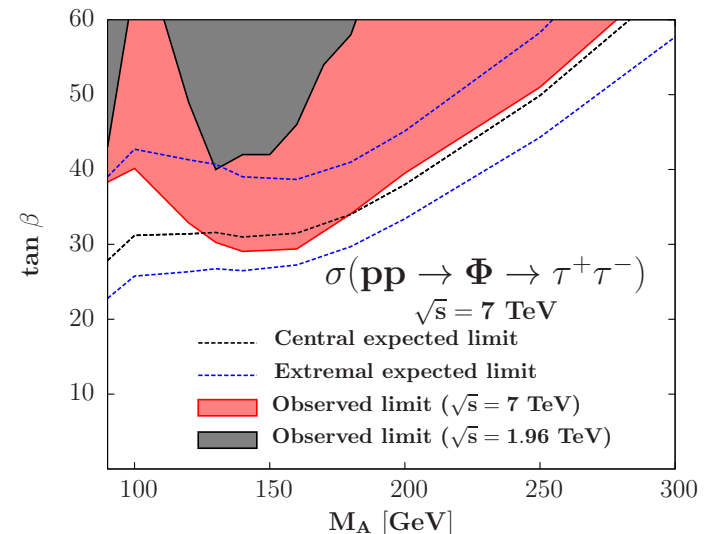
Arbey, Battaglia, Mahmoudi, arXiv:1110.3726

Impact of heavy Higgs search

- Search for $pp-H \rightarrow \tau\tau$ in CMS and ATLAS
- Production $gg \rightarrow H$ (through b loops) and bbH
- Strong enhancement for MSSM heavy doublet at large values of $\tan\beta$



- Theoretical uncertainties
 - Baglio, Djouadi, arXiv:1103.6247
 - At a given luminosity weakens the limit on $\tan\beta$

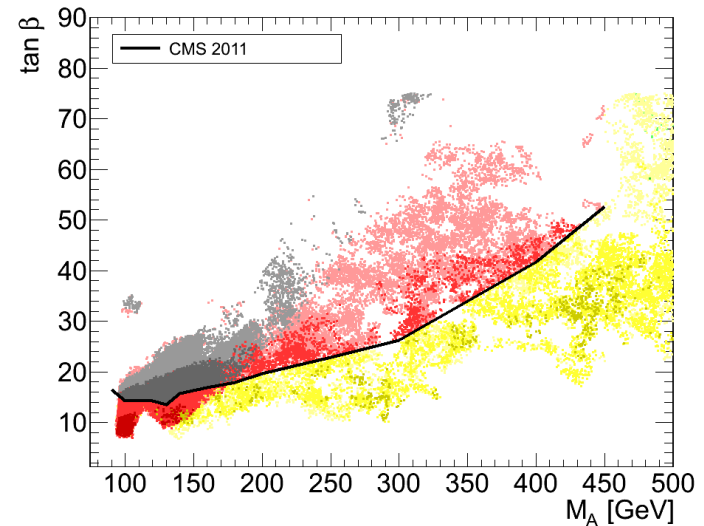


Impact of heavy Higgs search

- One example, MSSM with light neutralino DM ($<30\text{GeV}$)
- Possible only for non-universal gaugino mass
- Annihilation of LSP pairs:
 - through $\chi\chi \rightarrow A \rightarrow bb$
 - into fermion pairs through sfermion exchange (stau)
 - Both require particular conditions on spectrum.
- Consider Higgs case, for annihilation in early universe to be efficient m_A should be light and coupling large - $\tan\beta$ enhanced (not possible to have resonance effect)
- These scenarios constrained by heavy Higgs searches

Higgs in MSSM(2)

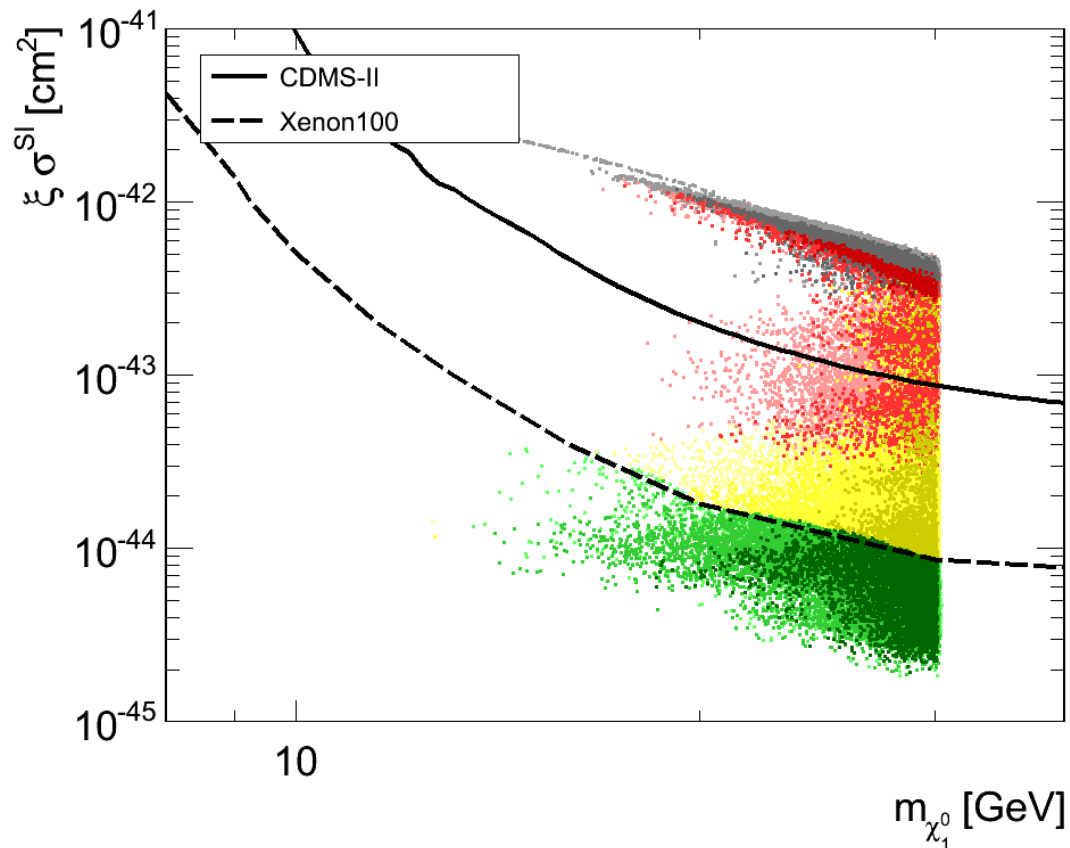
- Scan of MSSM focus on light neutralino ($<30\text{GeV}$)
 - 11 free parameters
 - $M_1, M_2, M_3, \mu, \tan\beta, M_A, A_t, M_{IL}, M_{IR}, M_{q1}, M_{q3}$
 - MCMC approach
 - Higgs searches (36pb^{-1}) (HiggsBounds)
 - Flavour physics
 - DM observables
- Pseudoscalar search major impact on allowed points



Albornoz Vasquez, GB, Boehm, arXiv:1108.1338

Light neutralino and A/H

- All light neutralino scenarios with large direct detection cross section constrained from negative Higgs searches + other astrophysical constraints



Invisible Higgs in BSM

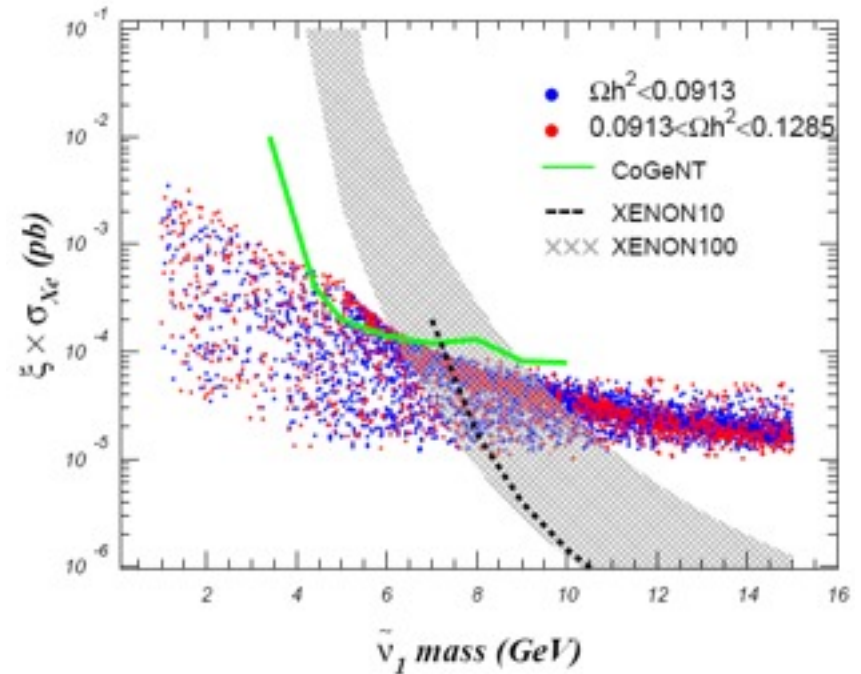
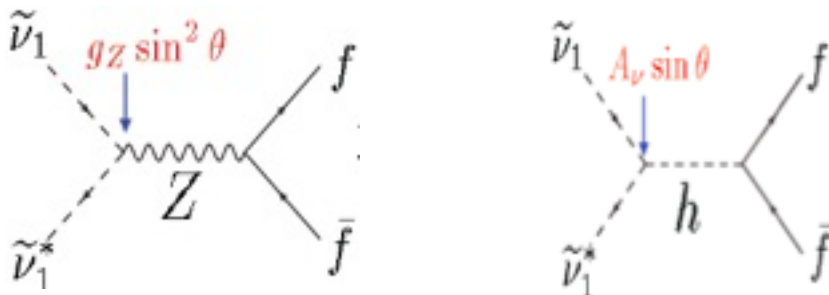
- Many BSM models have significant Br_{inv} , in particular models with light DM
 - MSSM with non-universal gaugino masses (as shown previously)
 - NMSSM
 - MSSM+sneutrino with large L/R mixing
 - UMSSM
 - Extended scalar sector
 - Little Higgs model
- In general invisible mode significant only in corner of parameter space due to kinematics
- DM 10GeV motivated by hints in direct detection experiments - whether or not these are confirmed, important to consider invisible mode
- Some models also have new decay modes for the Higgs, best example NMSSM with new singlet higgses, higgs to higgs decays

$$h \rightarrow A_1 A_1 \rightarrow 4 \tau$$

- Ellwanger et al, Phys. Rep. 486(2010)1

Example: sneutrino DM

- Extension of MSSM with RH neutrino+ sneutrino
 - Arkani-Hamed et al hep-ph/0006312
- Large L/R mixing makes sneutrino good DM candidate
- Annihilation

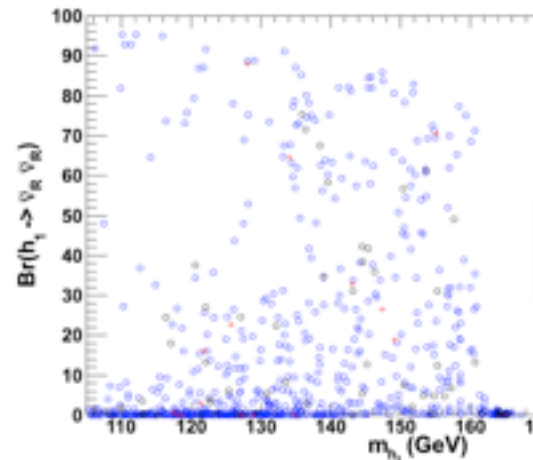
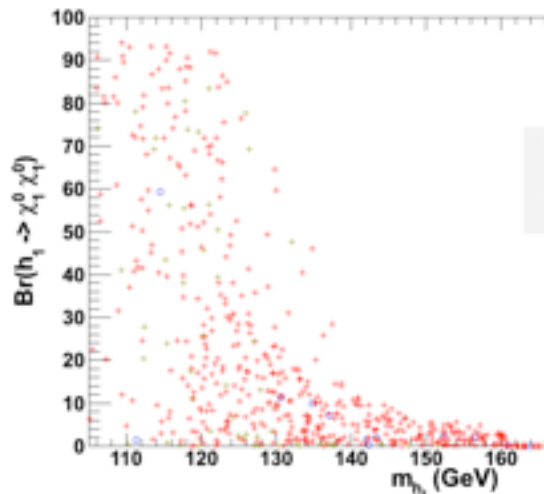


- Sneutrino couples to Higgs, if light nearly 100% Br_{inv}

GB, Kakizaki, Kraml, Park, Pukhov
JCAP 1011(2010)017

Example 2: UMSSM

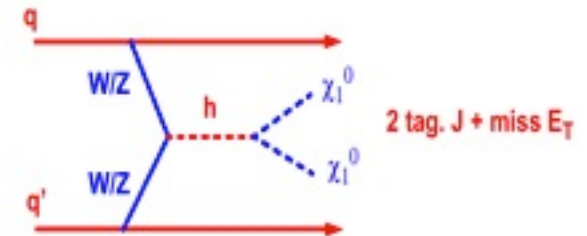
- MSSM with extended gauge symmetry U(1)
- Higgs sector h_1, h_2, h_3, A, H
- Scale of new gauge boson $Z' > 1.3-1.7\text{TeV}$ (depend on couplings)
- Light Higgs can be heavier than MSSM, SM-like $WW h_1$ coupling
- When LSP (neutralino or sneutrinoR) light, $h \rightarrow \text{inv}$



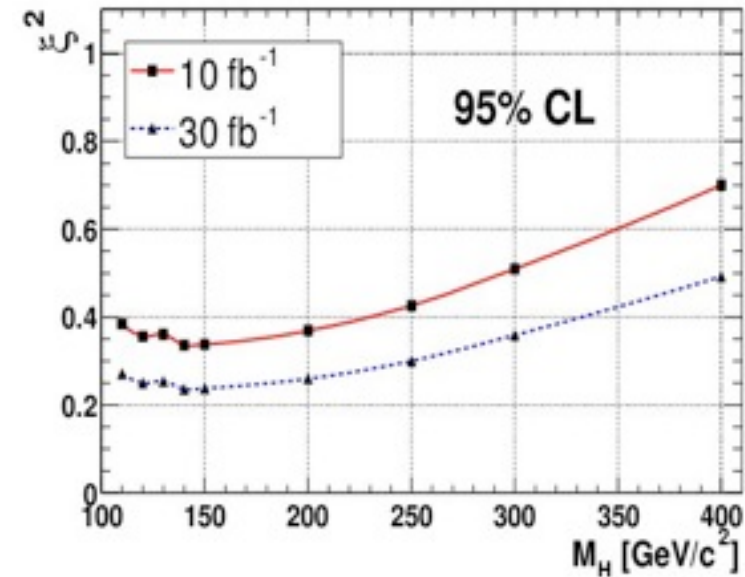
J. DaSilva, GB, in prep

Invisible Higgs at LHC

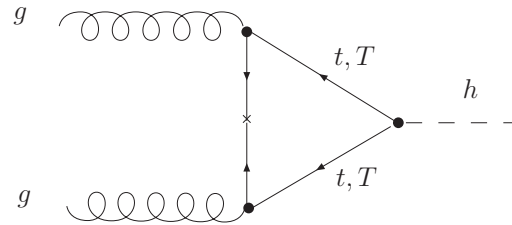
- Invisible decays --> reduction in all visible channels
- Search for invisible Higgs at LHC, important to complement visible channels and probe models with large invisible width
 - Vector boson fusion
 - Cavalli et al, Les Houches
 - hep-ph/ 0203056
 - Associated production with t quarks
 - Gunion, PRL72(94)199
 - Assoc. production with gauge bosons
 - ZH-> dilepton + missing E_t
 - sensitive to $Br_{inv} > .42$ for $m_h = 120$ at LHC14 ($100fb^{-1}$)
 - Godbole et al, hep-ph/030413



$$\xi^2 = \frac{\sigma(H_{jj})}{\sigma(H_{jj})_{SM}} \times Br(H \rightarrow inv)$$



BSM Higgs



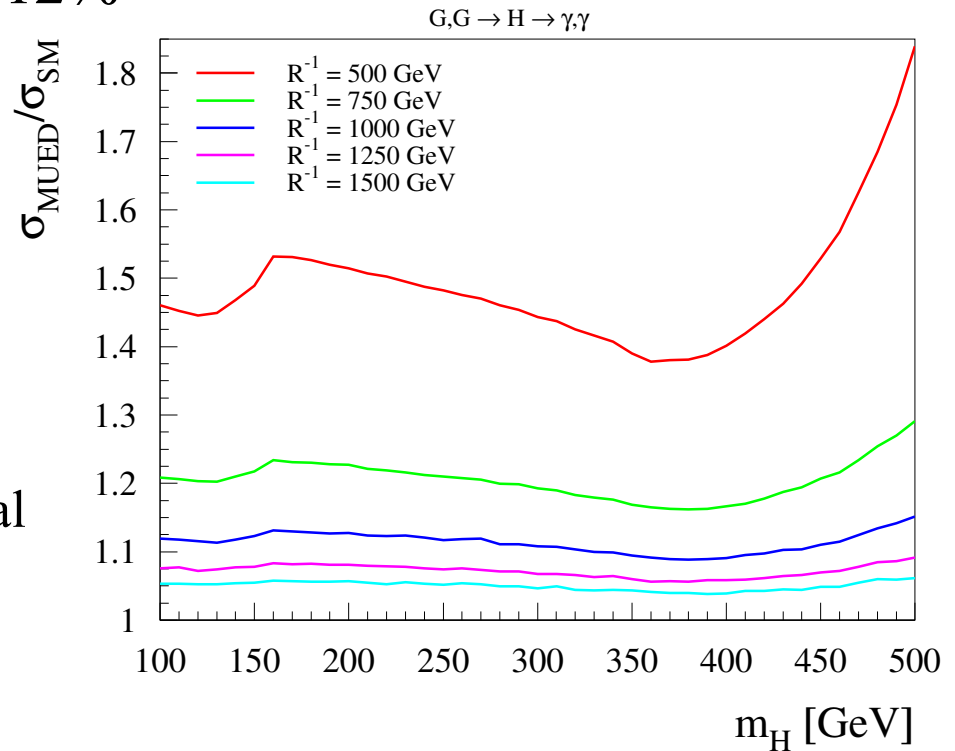
- New quark in the loop, contribution same sign as top, leads to increase in hgg
- Contribution also to $h\gamma\gamma$, lead to small decrease (W contribution dominant)
- Found in many BSM
 - Warped extra dimension model
 - Djouadi, Moreau, hep-ph 0707.3800
 - UED models (5D and 6D)
 - see G. Cacciapaglia et al 0901.0927
 - UED 6D - Nishiwaki et al, 1108.1764

Higgs in UED

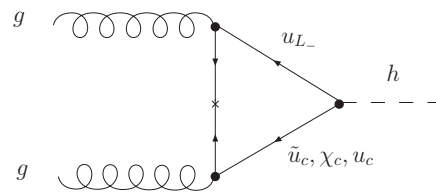
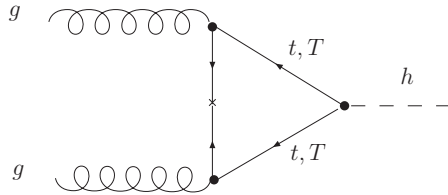
- KK top quarks contribution to ggh same sign as top \rightarrow increase in hgg width
- $m_{KK} < 500 \text{ GeV}$ \rightarrow constrained by Higgs search with 1.7 fb^{-1} if $m_h \sim 125 \text{ GeV}$ (similar to indirect limits)
- $m_{KK} \sim 1 \text{ TeV}$ \rightarrow $\gamma\gamma$ increase $\sim 12\%$

mUED

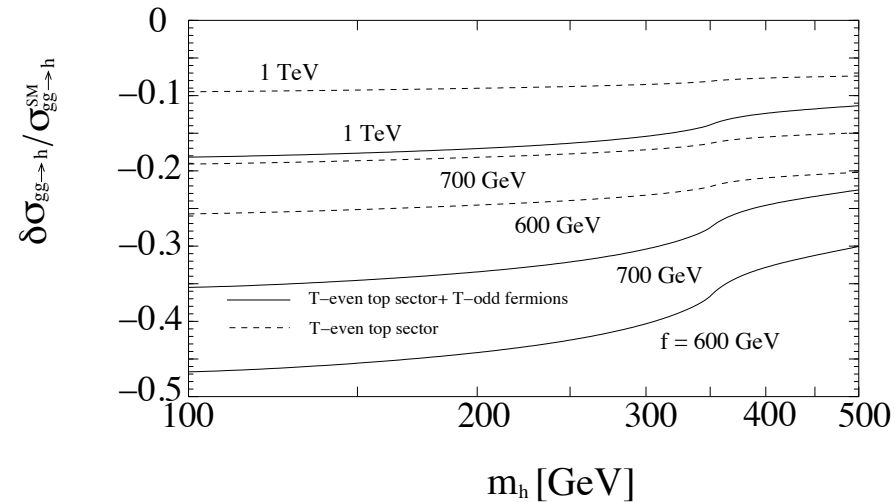
M. Brown, A. Belyaev et al



Little Higgs



- T-parity
- New T-even and T-odd fermions
- T-odd fermions: contribution opposite sign as top quark
- Significant decrease in ggh (20% for $f=1\text{TeV}$)
- Littlest Higgs model also allows for invisible width, $H \rightarrow A_H A_H$ significant only $f \sim 500\text{GeV}$
 - Hundi et al hep-ph/0611116



Chen, Tobe, Yuan, hep-ph/0602211

Conclusion

- In BSM, Higgs signal can be enhanced or reduced relative to SM - channel dependent
- Need exclusion in each channel independently
- SUSY Higgs not probed with 1.7fb^{-1} except in the $m_A/\tan\beta$ plane
- Search for Invisible Higgs important to cover all scenarios