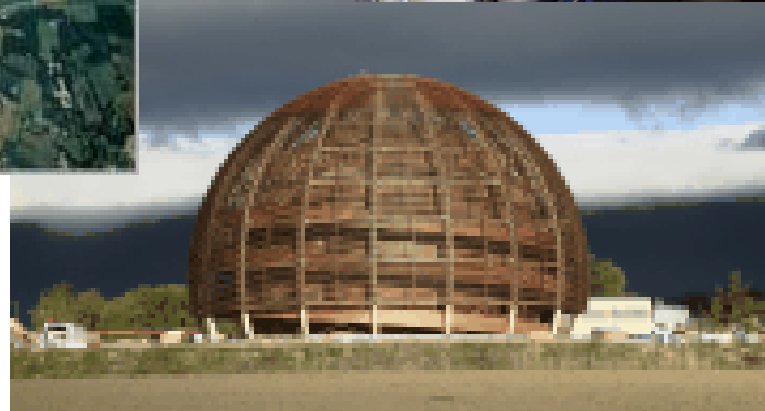
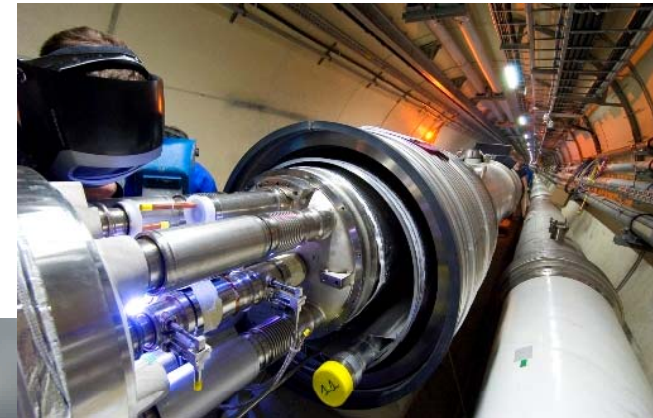


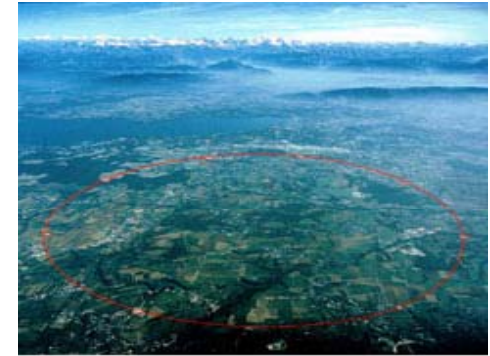
# Introduction to Accelerators

## Elena Wildner AT/MCS



# Contents

1. INTRODUCTION
2. THE ACCELERATOR CHAIN
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  1. Targets, Colliders
  2. Luminosity
5. ACCELERATOR TECHNOLOGI
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6. REFERENCES

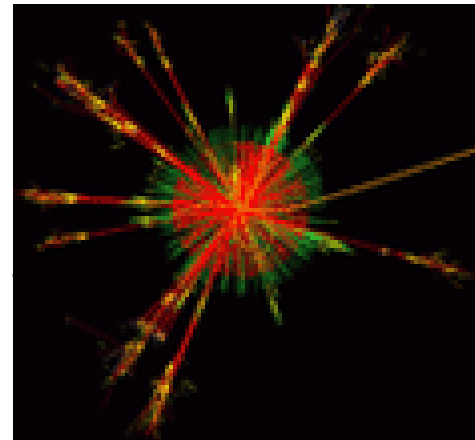


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# Application Areas

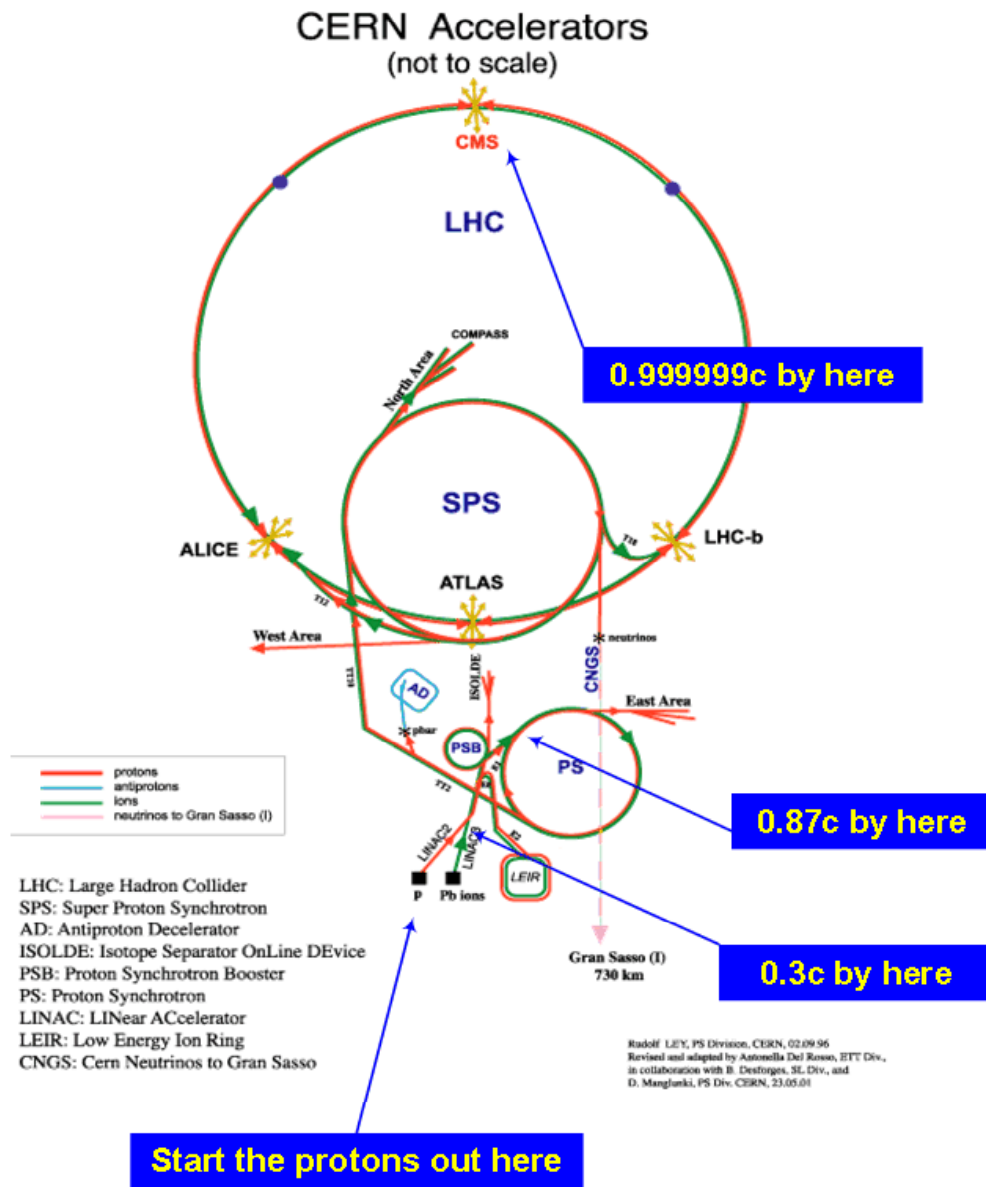
INTRODUCTION

- In your old TV set: Cathode Tube
- Material Physics
  - Photons from Electrons, Synchrotron Light
  - Material Surface
- Medicine
  - X-rays, synchrotron Radiation
  - Protons and Ions
- Food treatment
- Physics
  - Nuclear physics
  - Isotope production
  - High energy physics
- Etc.



# Accelerators and LHC experiments at CERN

INTRODUCTION



Energies:

Linac 50 MeV

PSB 1.4 GeV

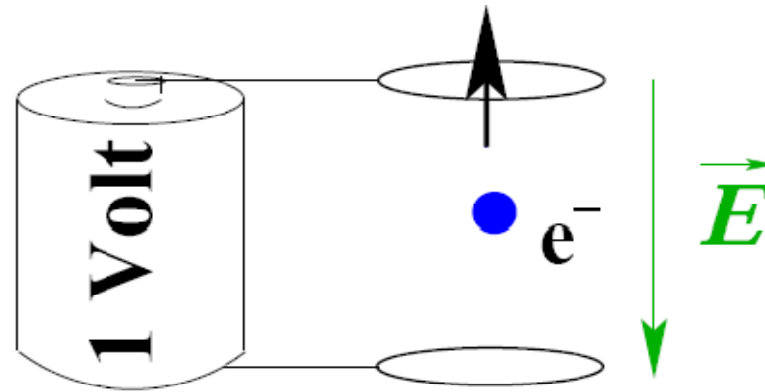
PS 28 GeV

SPS 450 GeV

LHC 7 TeV

Units?

# Units: Electronvolt



*Electronvolt, unit for energy denoted by  $eV$ , is used for small energies (joule)*

*1  $eV$  is defined as the energy needed to move one electron, with charge  $e$  (around  $1.602 \cdot 10^{-19}$  C) in an electric field with the strength 1 V/m a distance of 1 meter:*

$$1 \text{ eV} = 1.602 \cdot 10^{-19} \text{ joule.}$$

Acceleration

*In particle physics the unit  $eV$  is also used as a unit for mass since mass and energy are closely coupled through the relationship:*

$$E = mc^2, \quad m = \gamma \cdot m_0$$

*$m$  is the particle mass and  $c$  the speed of light in vacuum.*

*The mass of one electron, having a speed of  $v \ll c$  is around 0.5 MeV.*

Total energy

*From Wikipedia*

# Relativity

When particles are accelerated to velocities ( $v$ ) coming close to the velocity of light ( $c$ ):

then we must consider relativistic effects

$$\gamma = 1/\sqrt{1 - \beta^2}; \quad \beta = v/c$$

$$E = mc^2; \quad m = \gamma * m_0$$

Total Energy

Rest Mass

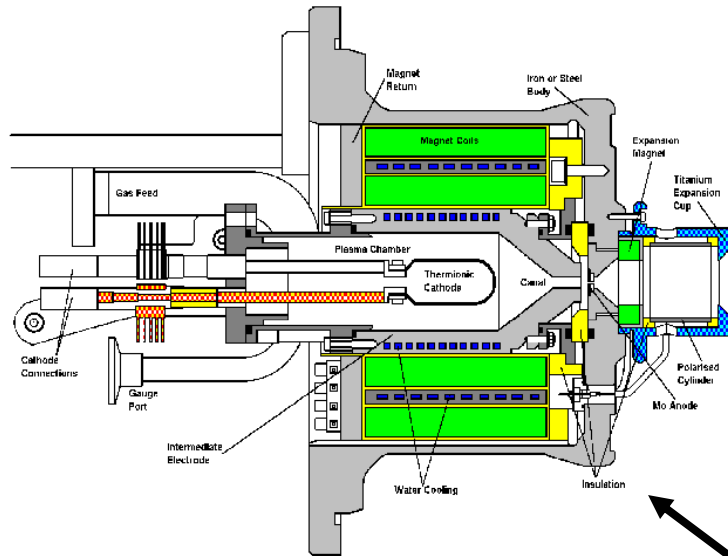
# Particle Sources and acceleration

## THE ACCELERATOR CHAIN

- Natural Radioactivity: alfa particles and electrons. Alfa particles have an energy of around 5 MeV (corresponds to a speed of  $\sim 15,000$  km/s).
- Production of particles: Particle sources
- Electrostatic fields are used for the first acceleration step after the source
- Linear accelerators accelerate the particles using Radio Frequency (RF) Fields
- Circular accelerators use RF and electromagnetic fields. Protons are today (2007+) accelerated to an energy of 7 TeV
- The particles need to circulate in vacuum (tubes or tanks) not to collide with other particles disturbing their trajectories.

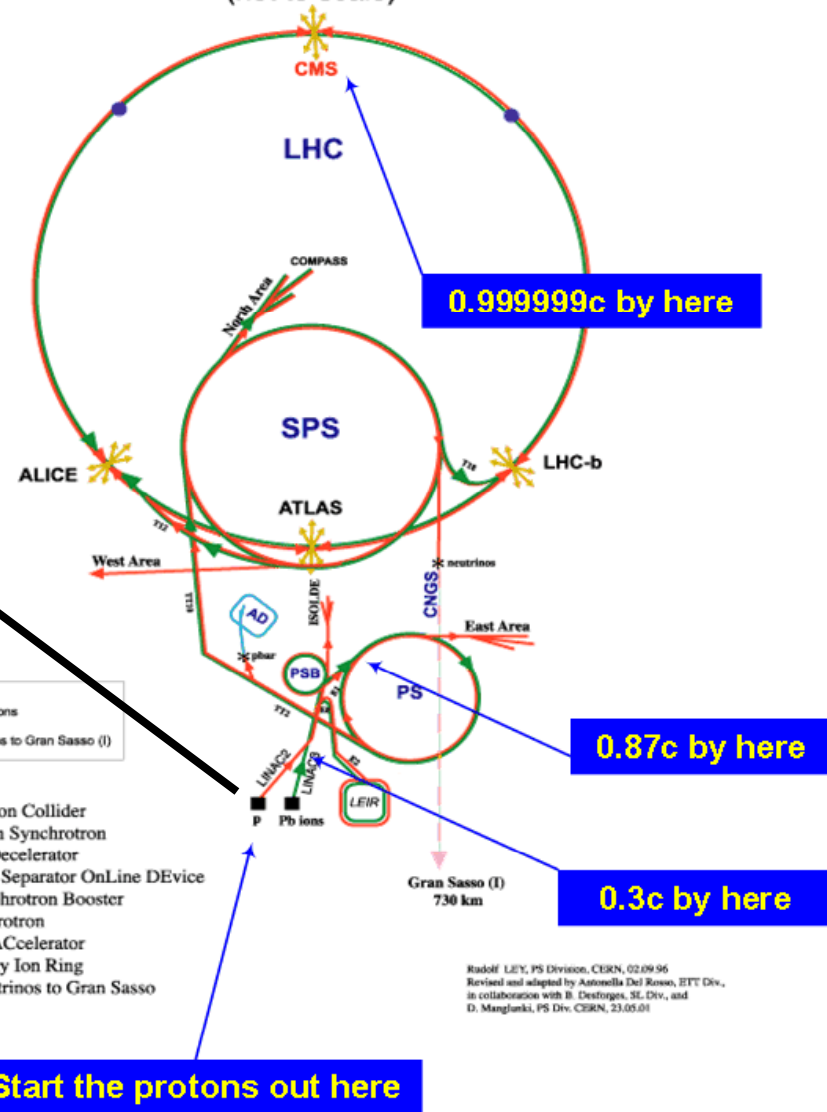
# Particle Sources 1

THE ACCELERATOR CHAIN



Duoplasmatron for proton production

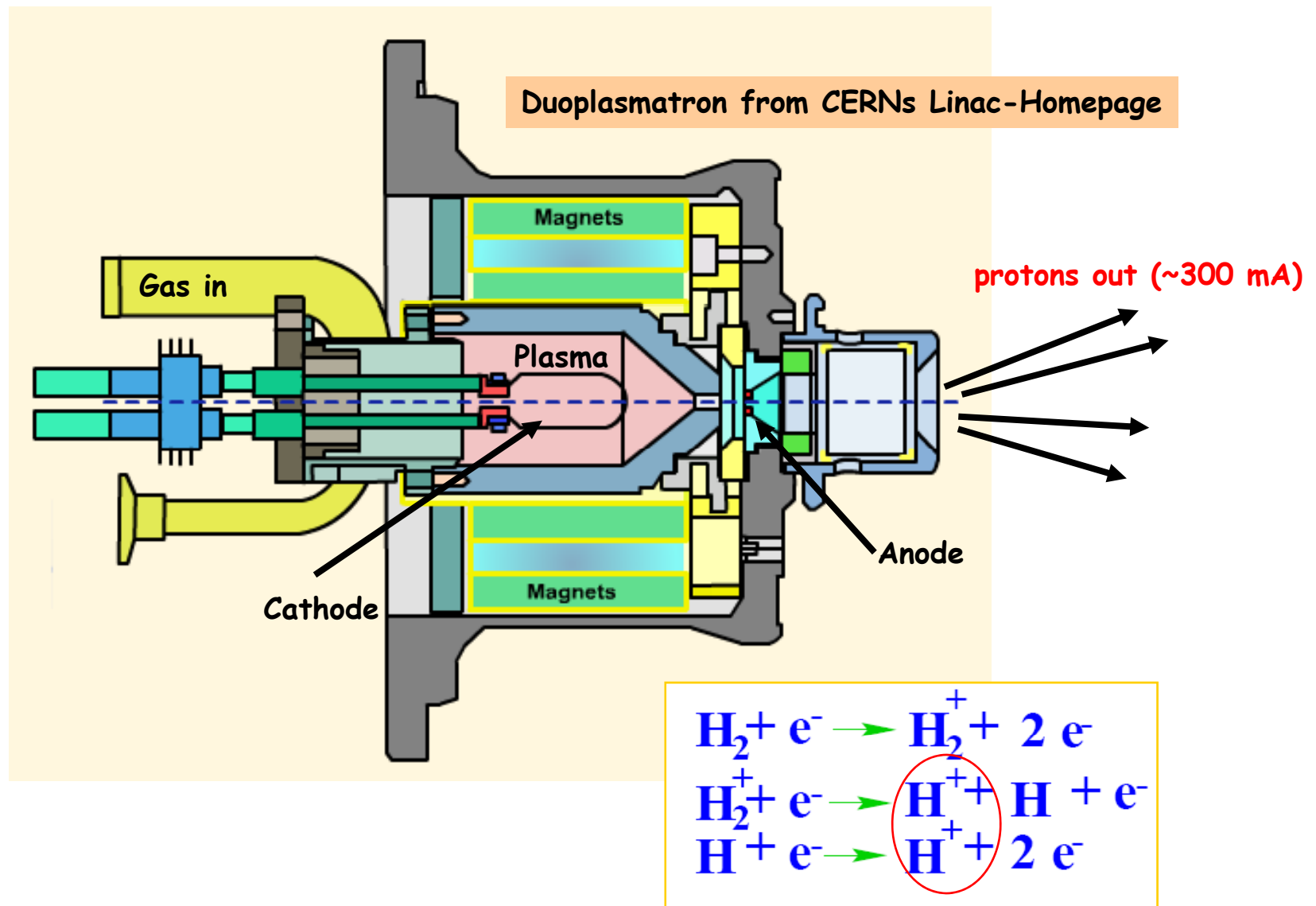
## CERN Accelerators (not to scale)





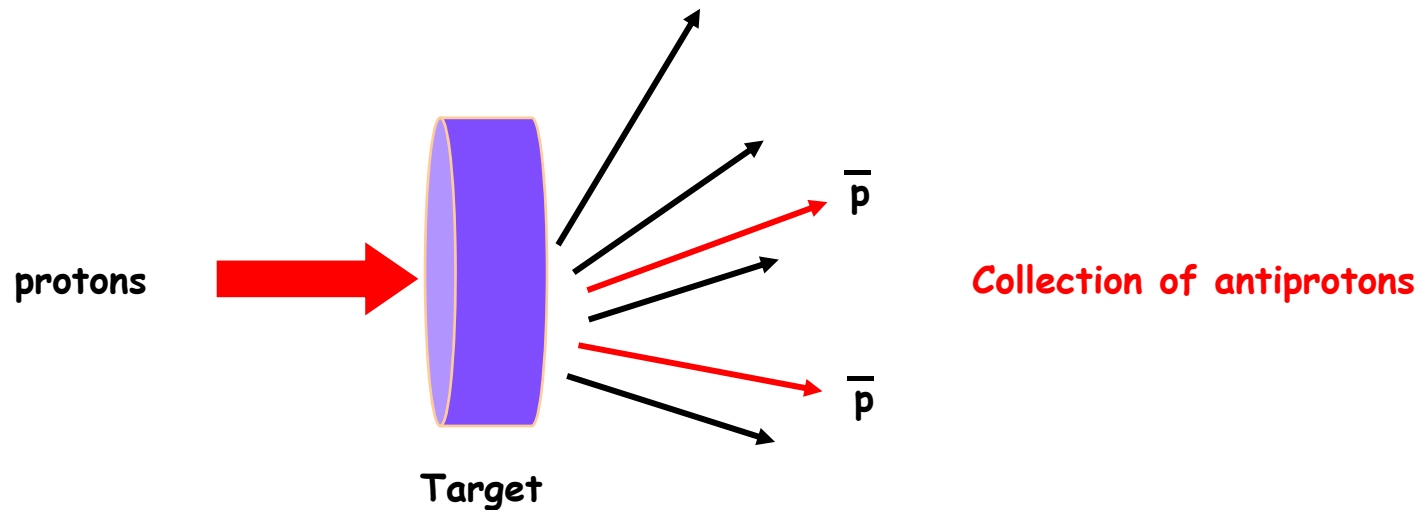
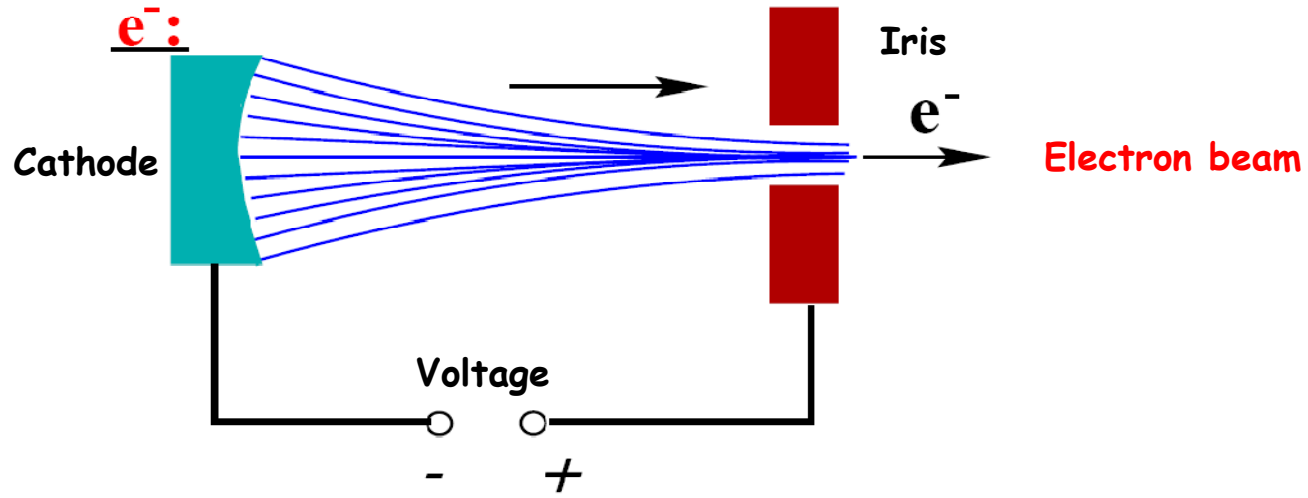
# Particle Sources 2

THE ACCELERATOR CHAIN



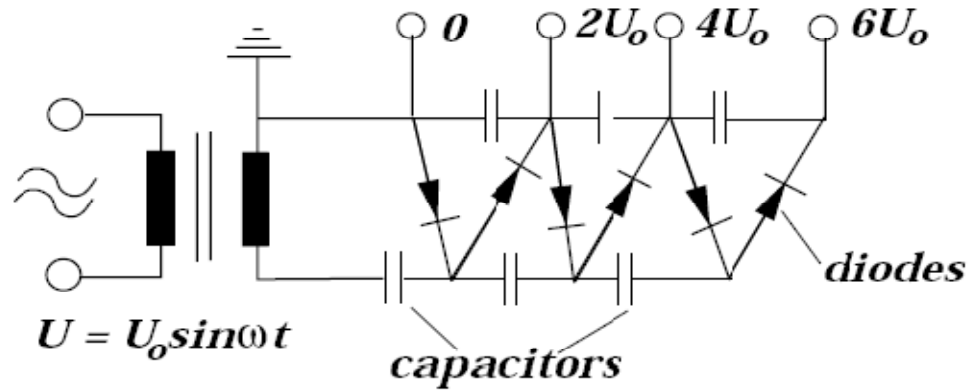
# Particle Sources 3

THE ACCELERATOR CHAIN

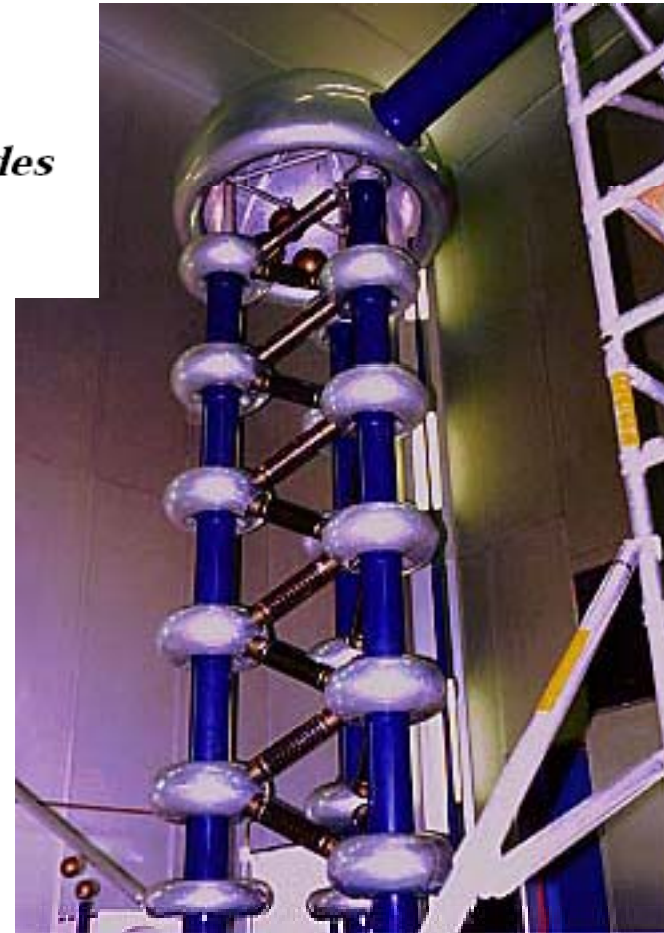


# The Cascade Generator

THE ACCELERATOR CHAIN



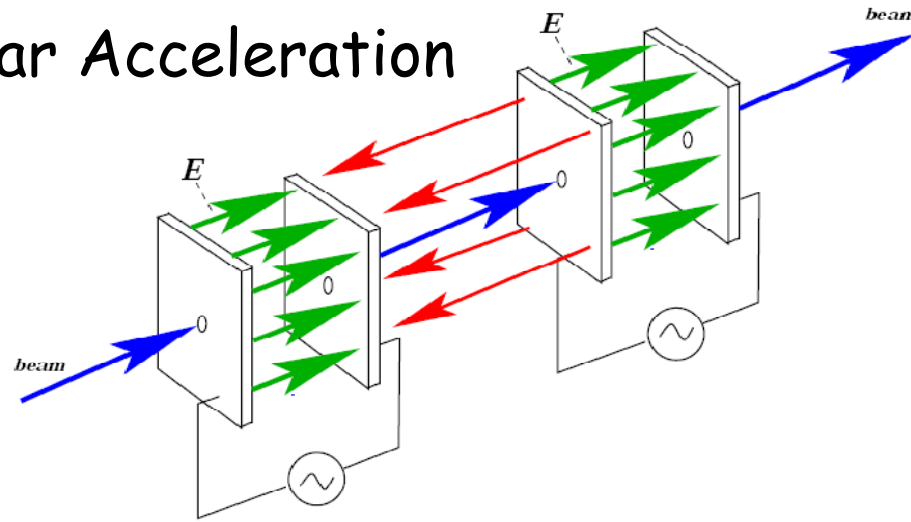
**Cockroft Walton: 4MV**



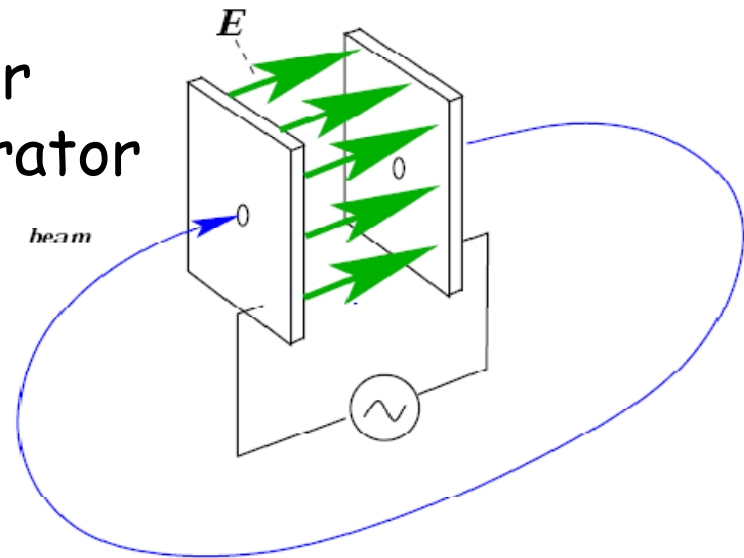
# Time Varying Electrical Fields

THE ACCELERATOR CHAIN

## Linear Acceleration



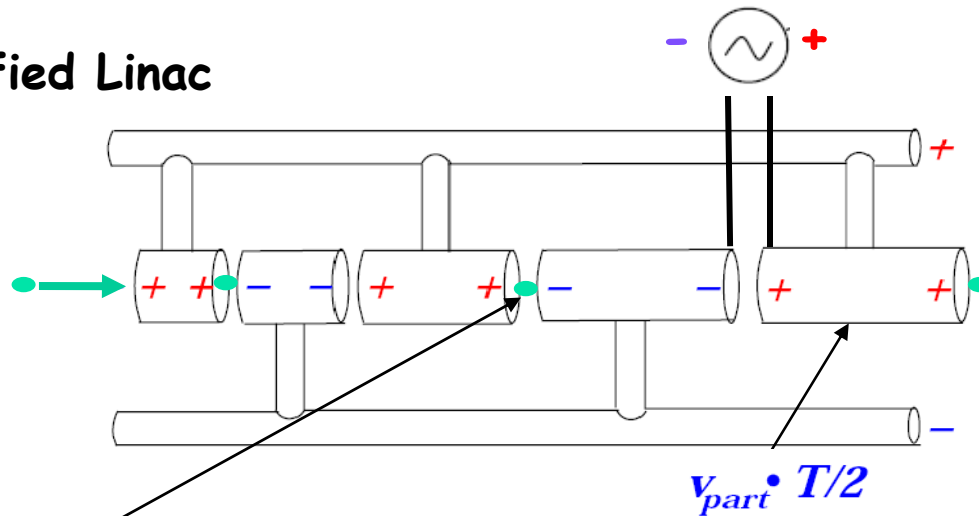
## Circular accelerator



# Linear accelerators

THE ACCELERATOR CHAIN

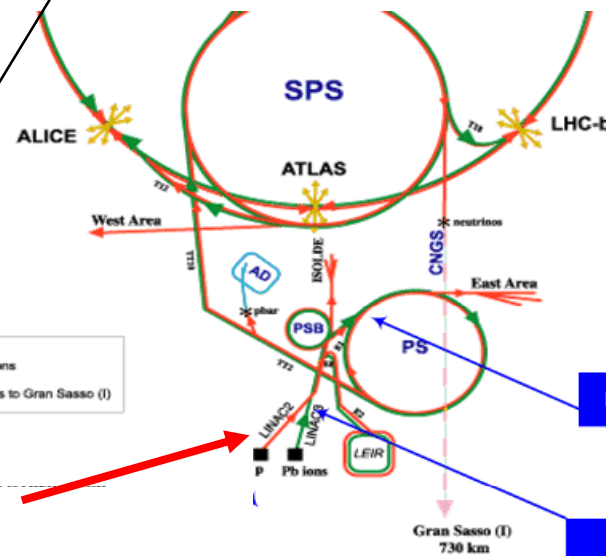
Simplified Linac



**Wideroe**  
**1928.**

The particles are grouped together to make sure that the field has the correct direction at the time the particle group passes the gap.

The speed of the particles increases and the length of the modules change so that the particle's arrival in the gap is synchronized with the field direction in the gap



Alvarez: Resonance tank

Linac

# The Cyclotron

THE ACCELERATOR CHAIN

Centripetal force = -Centrifugal force:

$$\frac{mv^2}{r} = Bqv$$

Reorganizing:

$$\frac{v}{r} = \frac{Bq}{m}$$

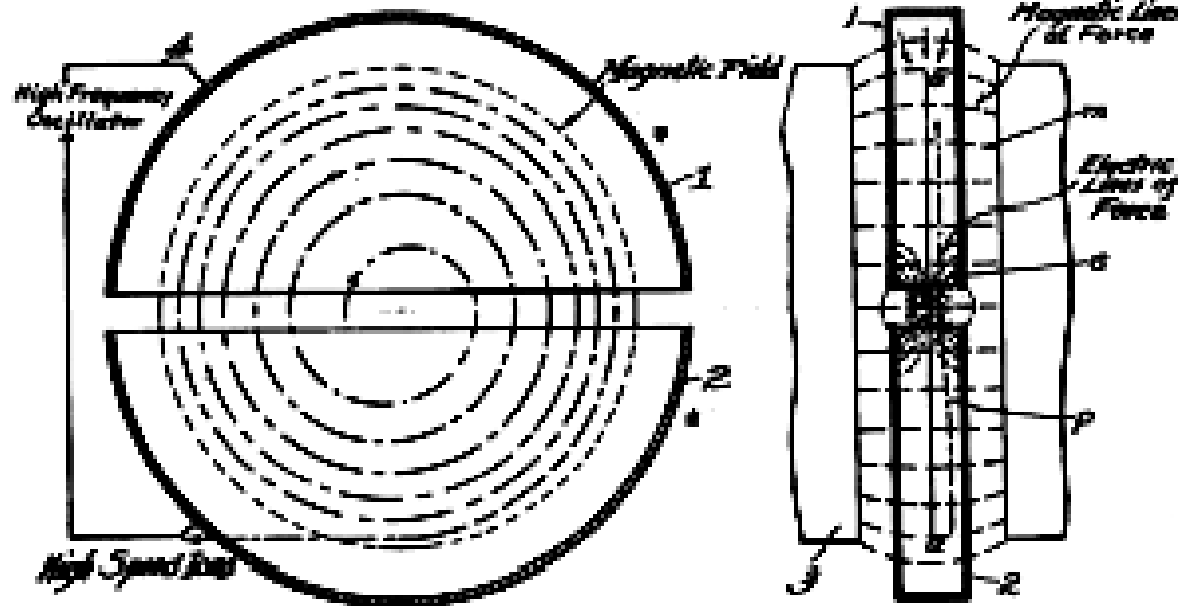
$$\downarrow$$

$$\omega = \frac{Bq}{m}$$

$$f = \frac{\omega}{2\pi}$$

$$f = \frac{Bq}{2m\pi}$$

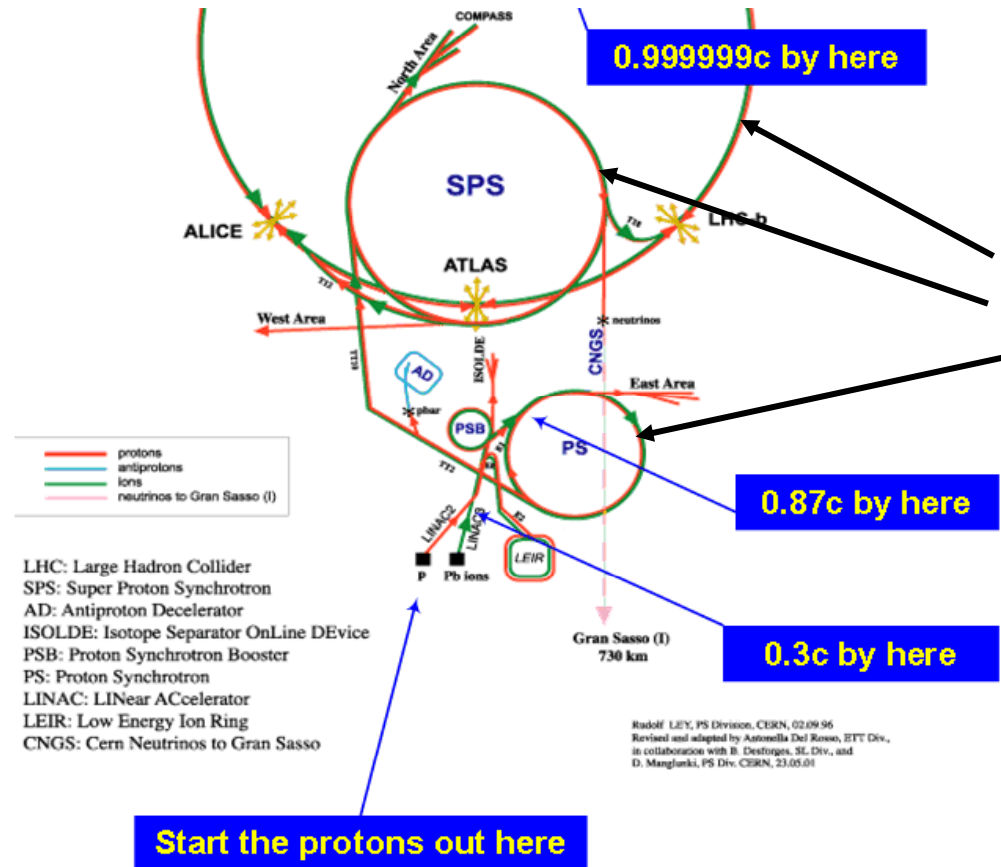
Continuous particle flux

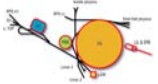


The frequency does not depend on the radius, if the mass is constant. When the particles become relativistic this is not valid any more. The frequency must change with the particle velocity: synchrocyclotron. The field can also change with the radius: isochronous cyclotron

# Synchrotrons at CERN

## THE ACCELERATOR CHAIN



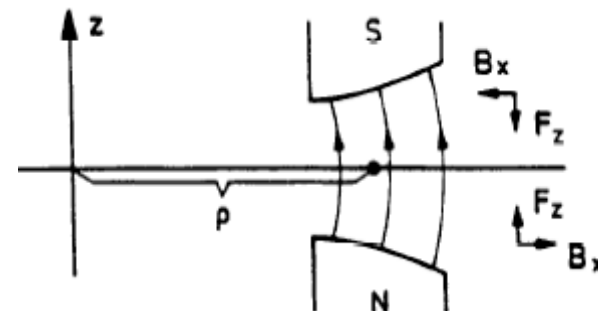
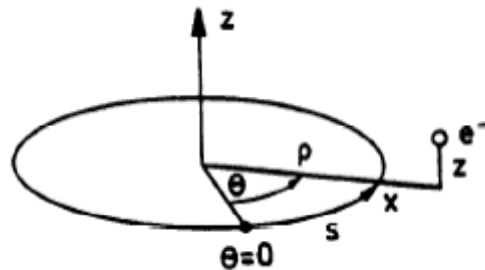
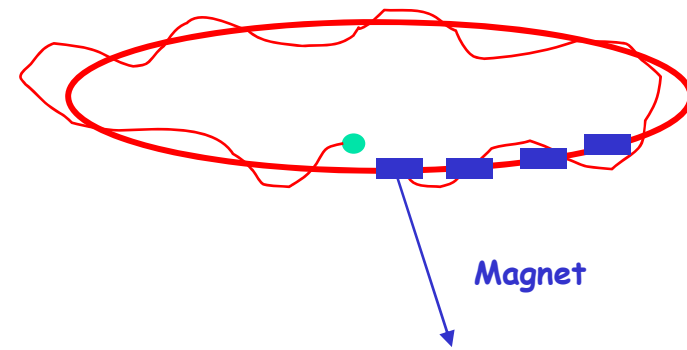
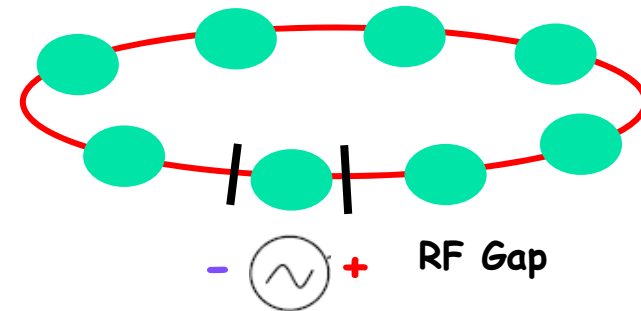


# The Synchrotron

HOW TO KEEP THE BEAM IN PLACE

Groups of particles are circulating synchronously with the RF field in the accelerating cavities

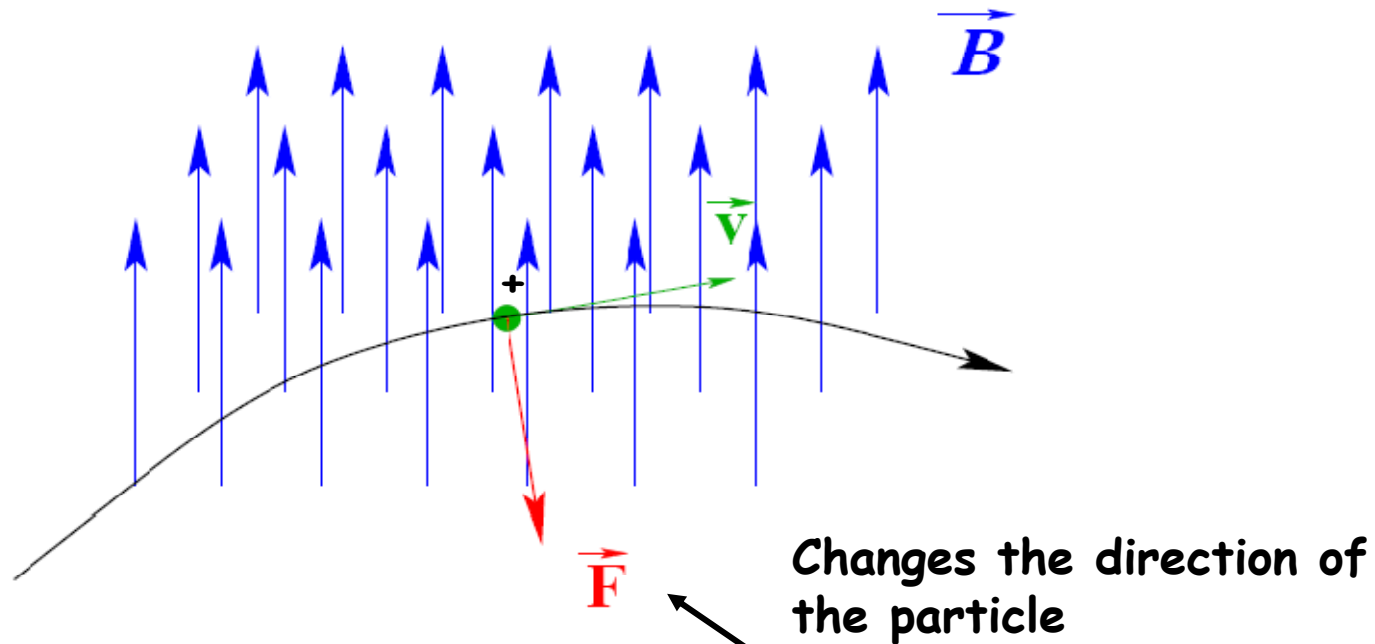
Each particle is circulating around an ideal (theoretical) orbit: for this to work out, acceleration and magnet fields must obey stability criteria!!





# Forces on the particles

STEERING



Lorentz:

$$\frac{d\vec{p}}{dt} = Q * \left( \vec{E} + \vec{v} \times \vec{B} \right)$$

Acceleration of the particles, field in the same direction as the velocity

# The Dipole

STEERING

Dipole Magnet, bends the particle trajectory in the horizontal plane (vertical field). Exception: correctors...

$$F_x = -ev_s B_y$$

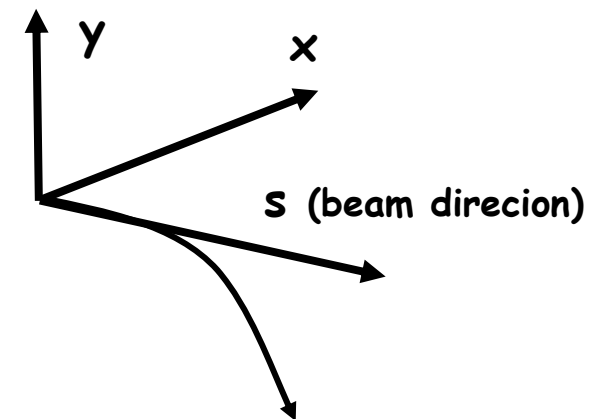
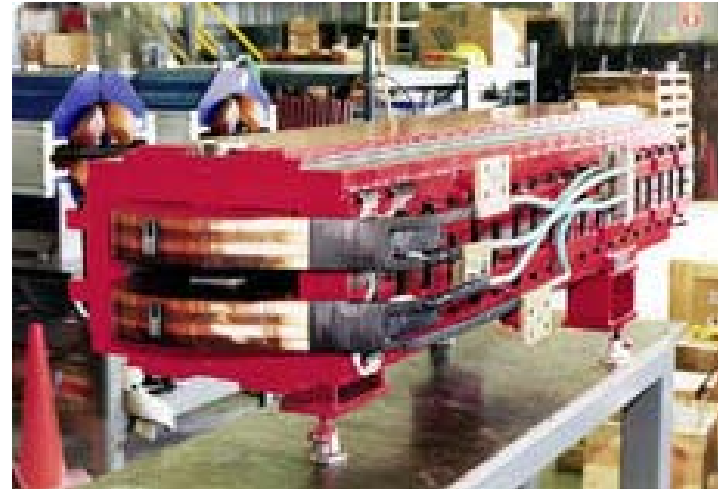
$$F_r = mv_s^2 / \rho$$

$$p = mv_s$$

$$\frac{1}{\rho(x, y, s)} = \frac{e}{p} B_y(x, y, s)$$

$$B\rho = \frac{p}{e}$$

**"Magnetic rigidity"**





## The Quadrupole 2

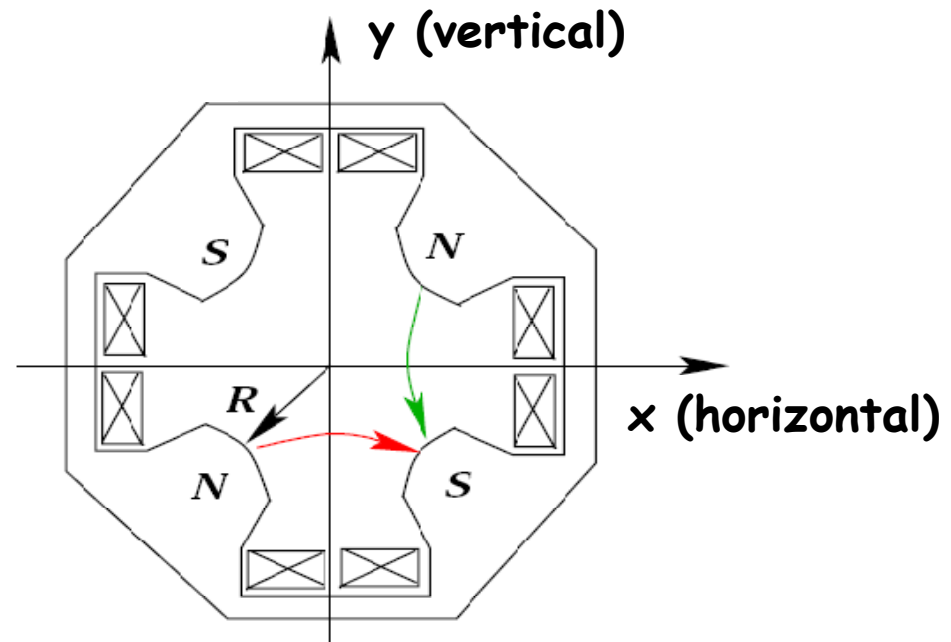
FOCUSING

$$B_x = -g \cdot y$$

$$B_y = -g \cdot x$$

$$F_x = g \cdot x$$

$$F_y = -g \cdot y$$

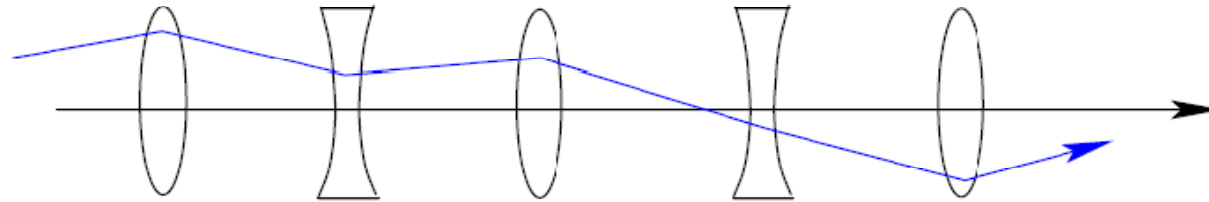


The force is proportional to  $x$  and to  $y$ :

Particles far from the center of the magnet are bent more, they get a more important correction.

# The Focusing System

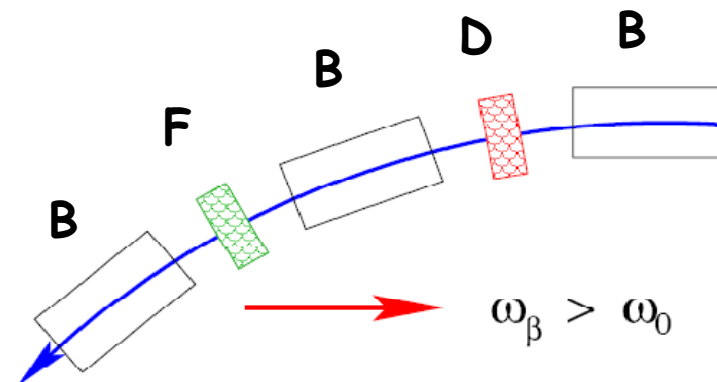
FOCUSING

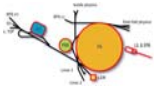


**"Alternate gradient focusing" gives an overall focusing effect (compare for example optical systems in cameras)**

**The beam takes up less space in the vacuum chamber, the amplitudes are smaller and for the same magnet aperture the field quality is better (cost optimization)**

**Synchrotron design: The magnets are of alternating field (focusing-defocusing)**





# The Oscillating Particles

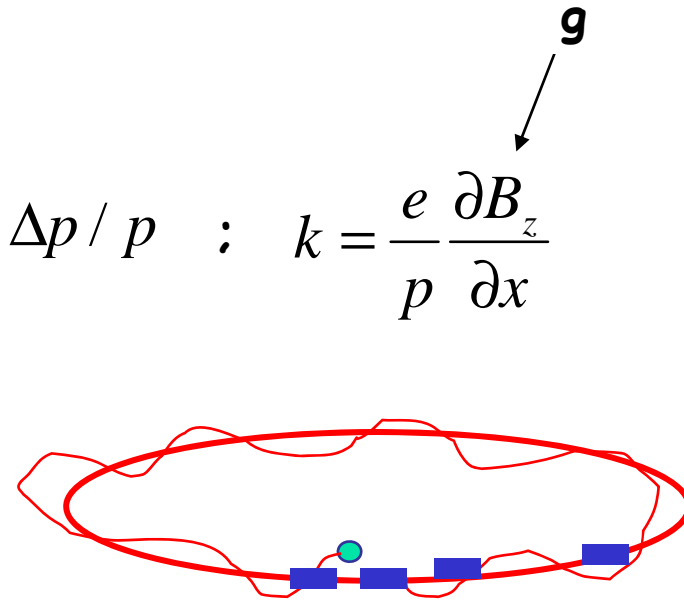
FOCUSING

The following kind of differential equations can be derived, compare the simple pendulum:

$$x''(s) + \left( \frac{1}{\rho^2(s)} - k(s) \right) \cdot x(s) = \frac{1}{\rho(s)} \Delta p / p \quad ; \quad k = \frac{e}{p} \frac{\partial B_z}{\partial x}$$

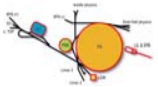
$$z''(s) + k(s) \cdot z(s) = 0$$

$$x(s) = \sqrt{\varepsilon \beta_x} (s) \cos\left(\frac{2\pi}{L} Q \cdot s + \delta\right)$$



**Oscillating movement with varying amplitude!**

The number of oscillations the particle makes in one turn is called the "tune" and is denoted Q. The Q-value is slightly different in two planes (the horizontal and the vertical planes). L is the circumference of the ring.

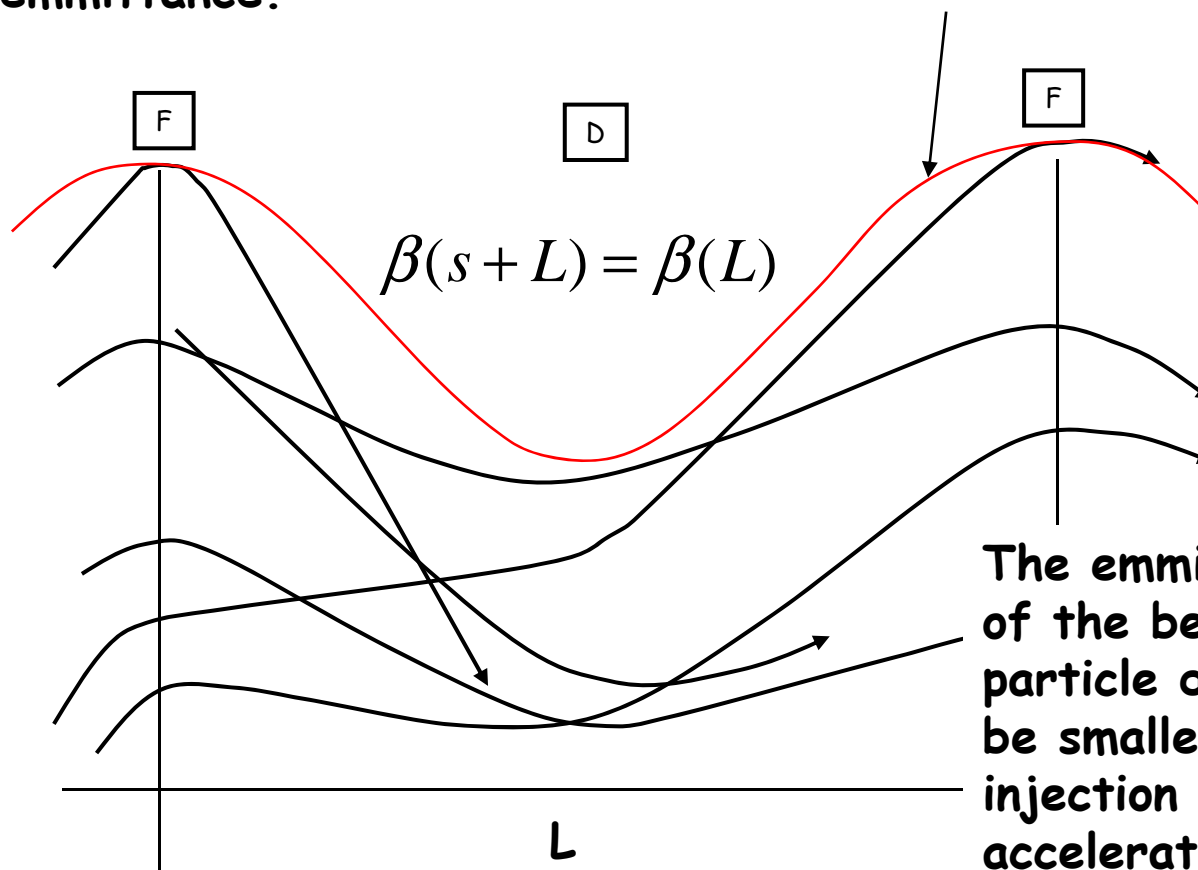


# The Beta Function

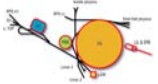
All particle excursions are confined by a function: the square root of the the beta function and the emittance.

$$x(s) = \sqrt{\varepsilon \beta_x(s)} \cos\left(\frac{2\pi}{L} Q \cdot s + \delta\right)$$

FOCUSING



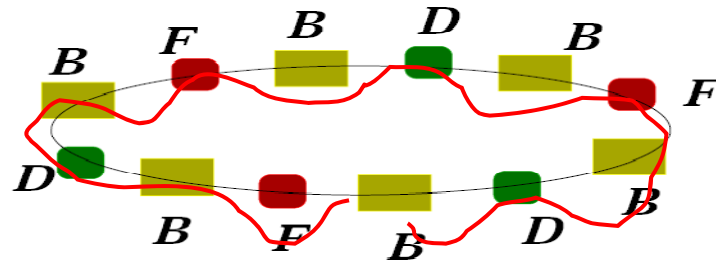
The emittance, a measure of the beam size and the particle divergences, cannot be smaller than after injection into the accelerator (normalized)



# Closed orbit, and field errors

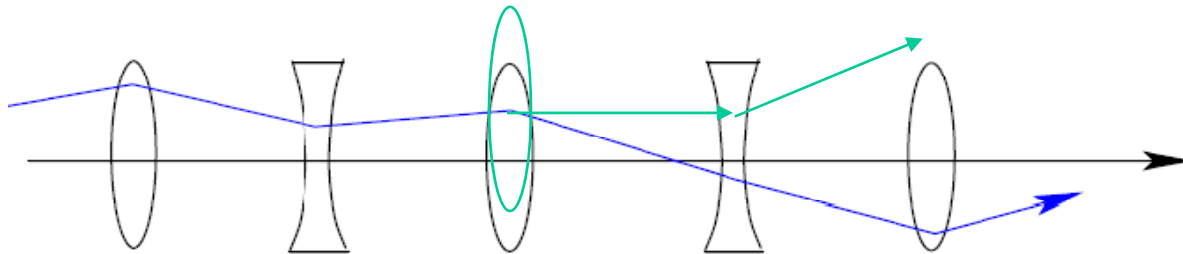
STEERING AND FOCUSING

Theoretically the particles oscillate around a nominal, calculated orbit



The magnets are not perfect, in addition they cannot be perfectly aligned.

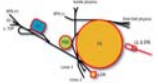
For the quadrupoles for example this means that the force that the particles feel is either too large or too small with respect to the theoretically calculated force. Effect: the whole beam is deviated.



$$F_x = g \cdot x$$

$$F_y = -g \cdot y$$

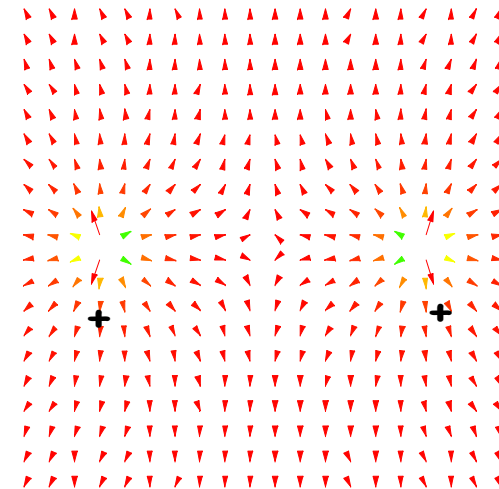
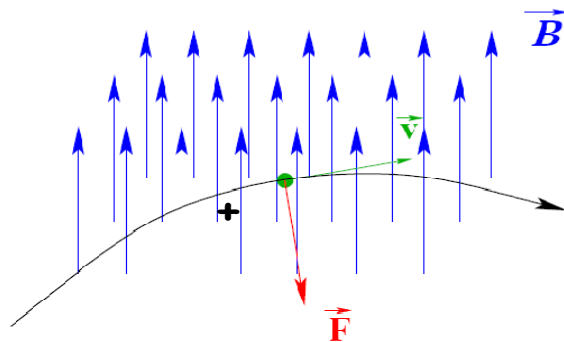




## Effects from other particles: example, space charge

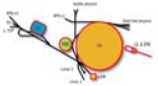
The field felt by a particle may come not only from the magnetic elements but also from the other particles.

If the beam is dense the influence of the coulomb field of other particles may be important.



$$\frac{d\vec{p}}{dt} = Q * \left( \vec{E} + \vec{v} \times \vec{B} \right)$$

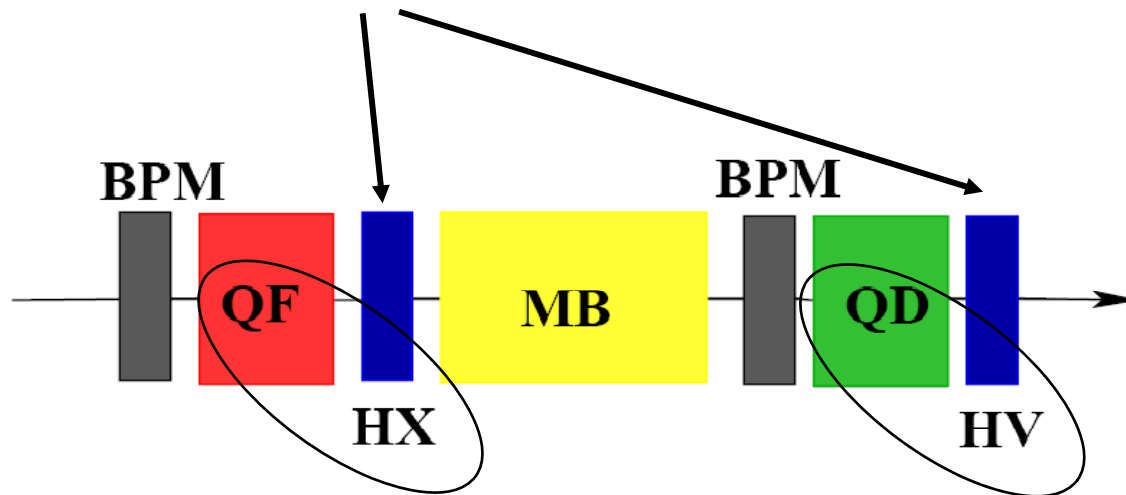
Correction of the field in the quadrupoles may be needed.



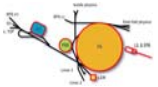
# Correctors

Beam Position Monitors are used to measure the center of the beam near a quadrupole, the beam should be in the center at this position.

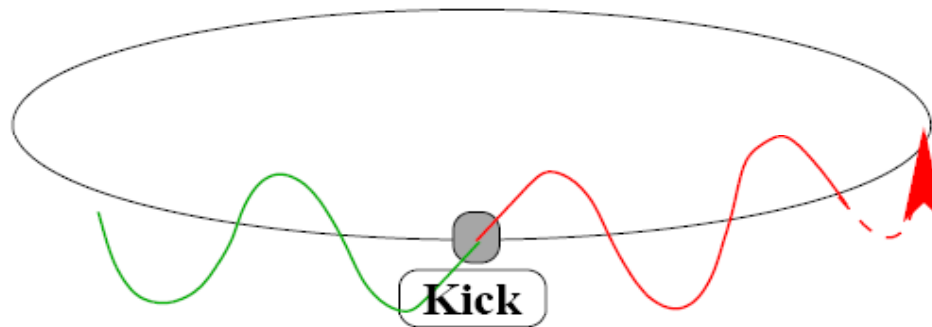
Small dipole magnets are used to correct possible beam position errors.



Other types of magnets are used to correct other types of errors for example non perfect magnetic fields.

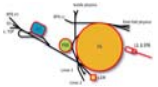


## Possible errors 1



The Q-value gives the number of oscillations the particles make in one turn. If this value is an integer, the beam "sees" the same magnet-error over and over again and we may have a resonance phenomenon. (Resonance) Therefore the Q-value is not an integer.

The magnets have to be good enough so that resonance phenomena do not occur. Non wanted magnetic field components (sextupolar, octupolar etc.) are comparable to  $10^{-4}$  relative to the main component of a magnet (dipole in a bending magnet, quadrupole in a focussing magnet etc.). This is valid for LHC



## Possible errors 2

Types of effects that may influence the accelerator performance and has to be taken into account:

Movement of the surface of the earth

Trains

The moon

The seasons

Construction work

...

Calibration of the magnets is important

Current regulation in the magnets

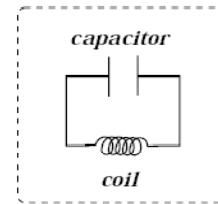
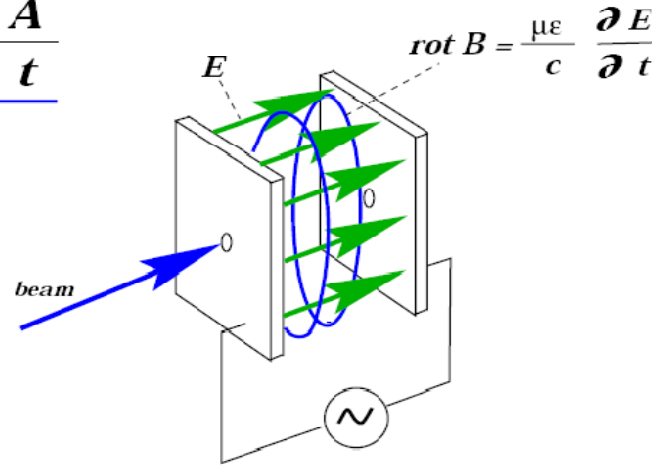
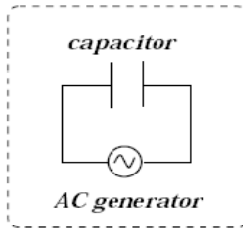
...

The energy of the particles must correspond to the field in the magnets, to permit the particle to stay in their orbits. Control of the acceleration!

# Electrical Fields for Acceleration

ACCELERATION

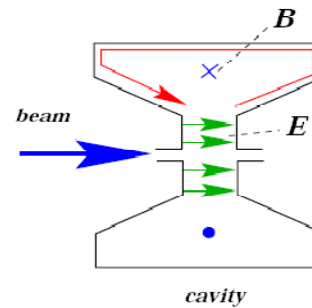
● 
$$\underline{E = - \frac{1}{c} \frac{\partial A}{\partial t}}$$



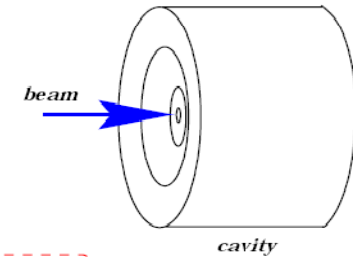
$$L = \frac{\mu_0 \cdot N^2 \cdot A}{l}$$

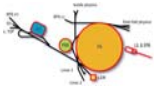
$$C = \frac{\epsilon_0 \cdot A}{d}$$

Resonance circuit  
Cavity for acceleration



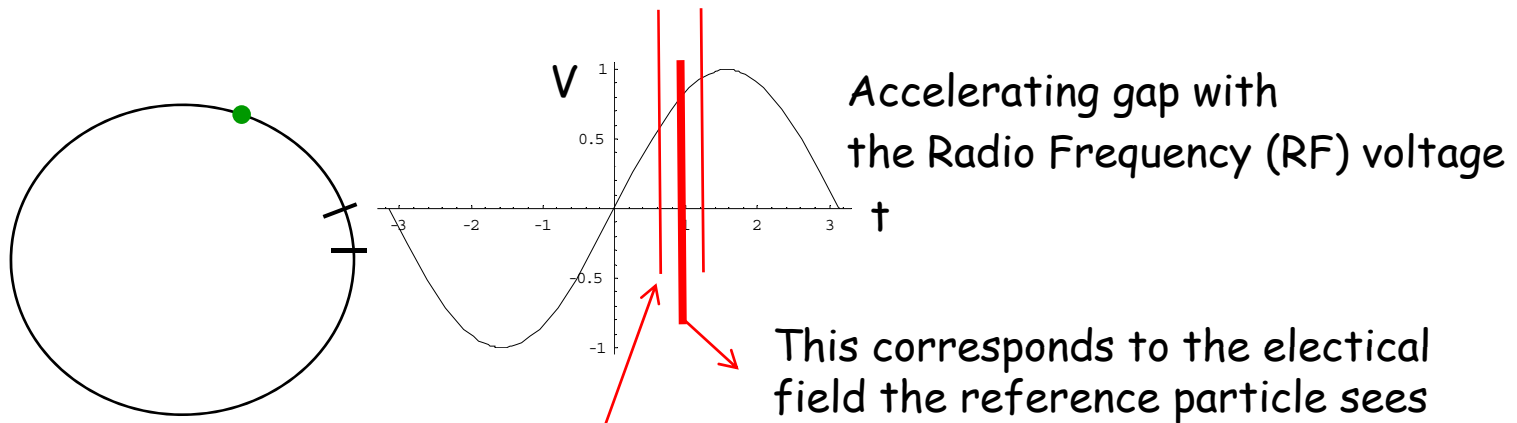
$f; Q; R$



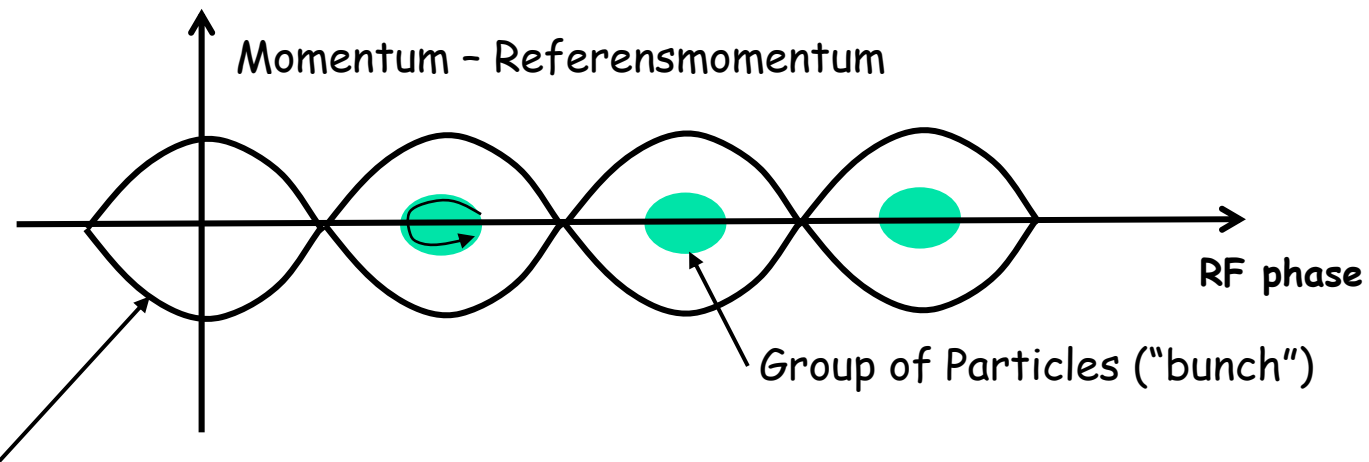


# The Synchrotron: grouping particles

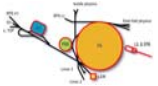
ACCELERATION



An early particle gets less energy increase



"Bucket": Energy/phase condition for stability

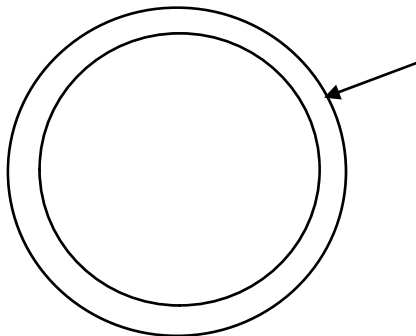


# The Synchrotron: Acceleration

ACCELERATION

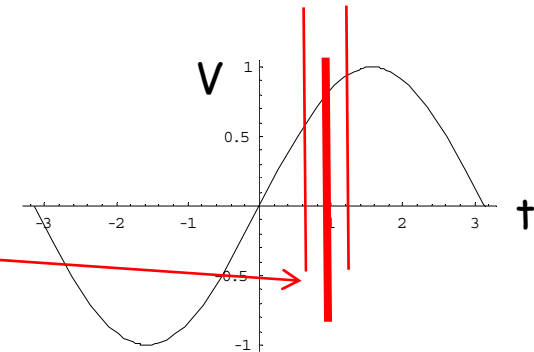
Magnetic field  $B$  and the velocity  $v$

- > revolution path, the reference trajectory
- > revolution time
- > the time at which the particle enters the RF gap



If  $v$  increases ->  
Radius larger than the ref

Revolution time larger  
Not adapted energy  
correction



To accelerate the particles, the magnetic field has to increase and the frequency has to be adjusted to keep the particles on the reference trajectory.

$$f = v / (2\pi R_0)$$



# Experiment

EXPERIMENT

## Targets:

Bombarding material with a beam directed out of the accelerator.

Bubbelchamber

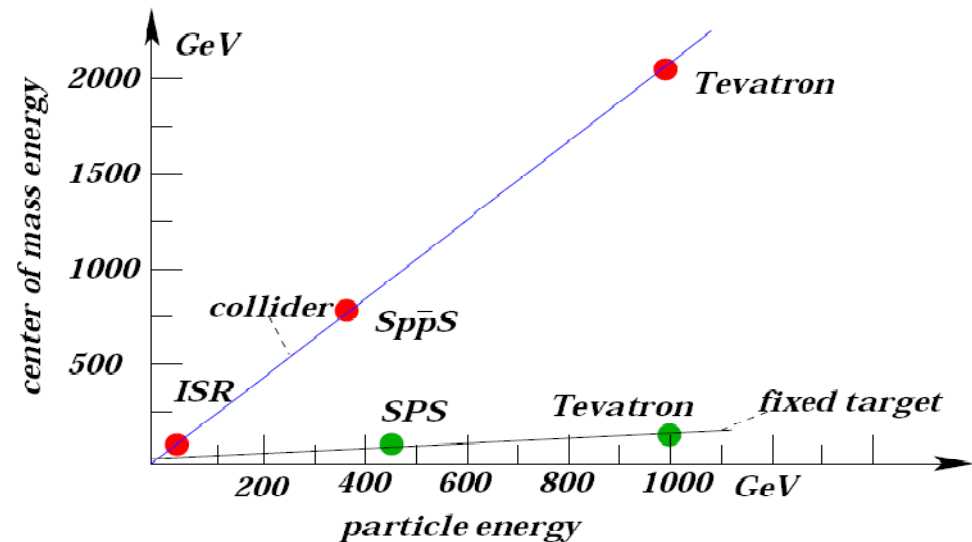
Available energy is calculated in the center of mass of the system (colliding objects)

To collide particle more interesting

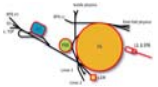
1960: electron/positron collider

1970: proton antiproton collider

2000: ions, gold

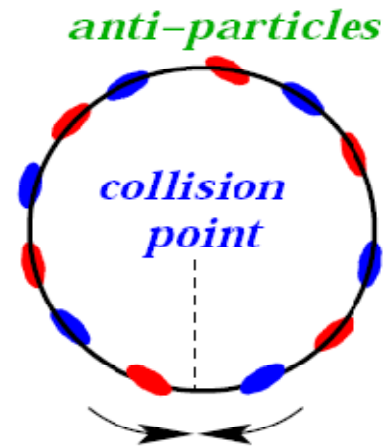
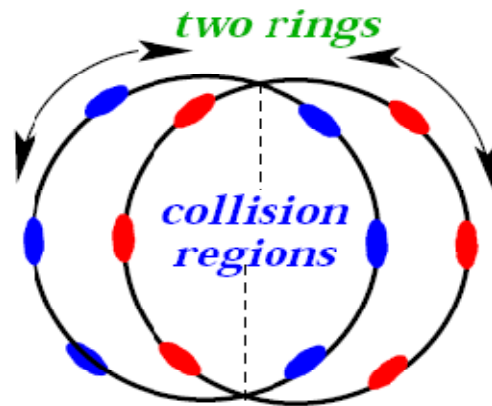






# Colliders

EXPERIMENT



- All particles do not collide at the same time -> long time is needed
- Two beams are needed
- Antiparticles are difficult (expensive) to produce (~1 antiproton/ $10^6$  protons)
- The beams affect each other: the beams have to be separated when not colliding



# Leptons/Hadrons

EXPERIMENT

## Lepton versus Hadron Collider

● Leptons: ( $e^+ / e^-$ )

■ elementary particles

→ well defined energy

→ precision experiments

● Hadrons: ( $p^+ / p^-$ )

■ multi particle collisions

→ energy spread

→ discovery potential

● Example:

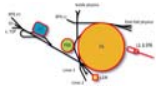
$Z_0$

1985 Sp $\bar{p}$ S

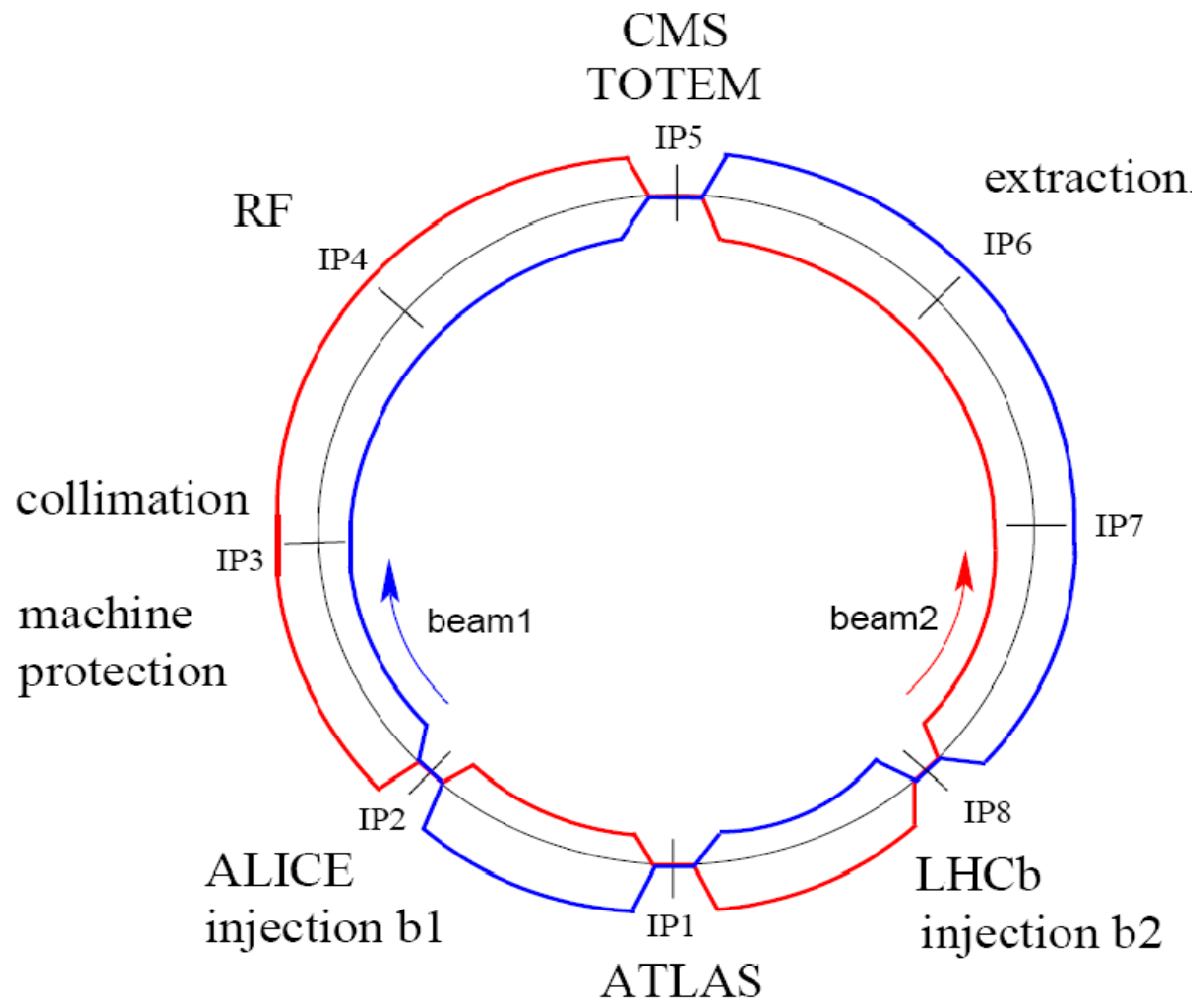
$p^+p^-$

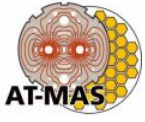
1990 LEP

$e^+e^-$



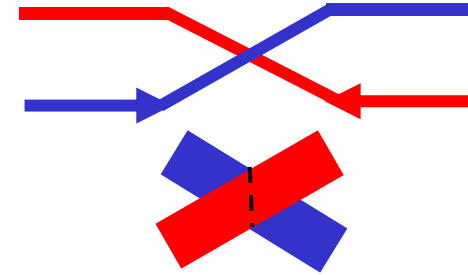
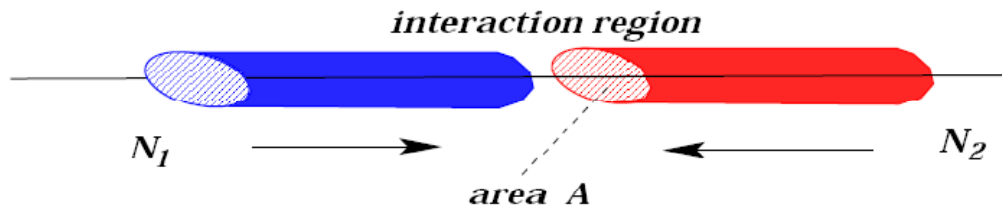
# The LHC





# Luminosity

EXPERIMENT



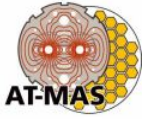
$$A = \pi \epsilon \beta^*$$

$$N_{ev}/sec = \sigma \cdot L$$

$$x(s) = \sqrt{\epsilon \beta_x(s)} \cos\left(\frac{2\pi}{L} Q \cdot s + \delta\right)$$

$$L = \frac{N_b^2 n_b f_{rev}}{4\pi \epsilon \beta^*} F$$

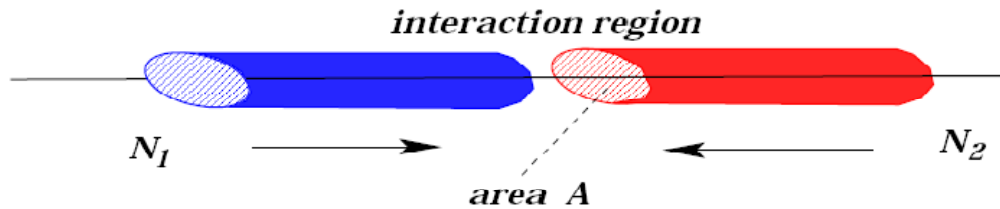
Number of particles per bunch (two beams) →  $N_b^2$   
 Number of bunches per beam →  $n_b$   
 Revolution frequency →  $f_{rev}$   
 Formfactor from the crossing angle →  $F$   
 Emittance →  $\epsilon$   
 Optical beta function →  $\beta^*$



# Luminosity: the beam size

EXPERIMENT

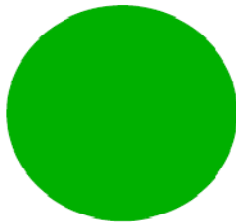
We need a small beam in the collision point



$$L = \frac{N_b^2 n_b f_{rev}}{4\pi\epsilon\beta^*} F$$

$$\sigma = \sqrt{\epsilon\beta^*}$$

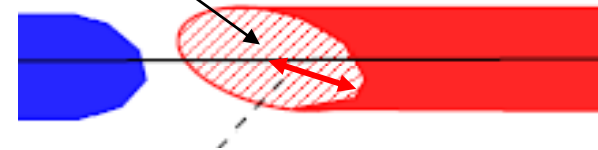
LHC:



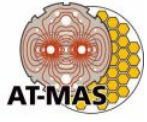
$\langle\beta\rangle_{arc} = 80 \text{ meter}$



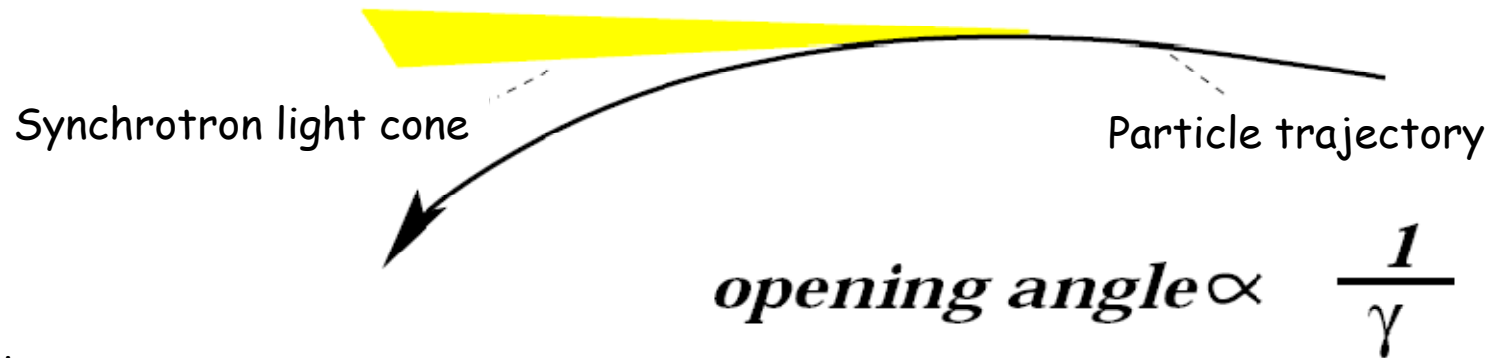
$\beta_{IP} = 0.5 \text{ meter}$



Limitation: Available magnetic field  
Magnet aperture



# Synchrotron light



Electromagnetic waves

Accelerated charged particles emit photons

Radio signals and x-ray

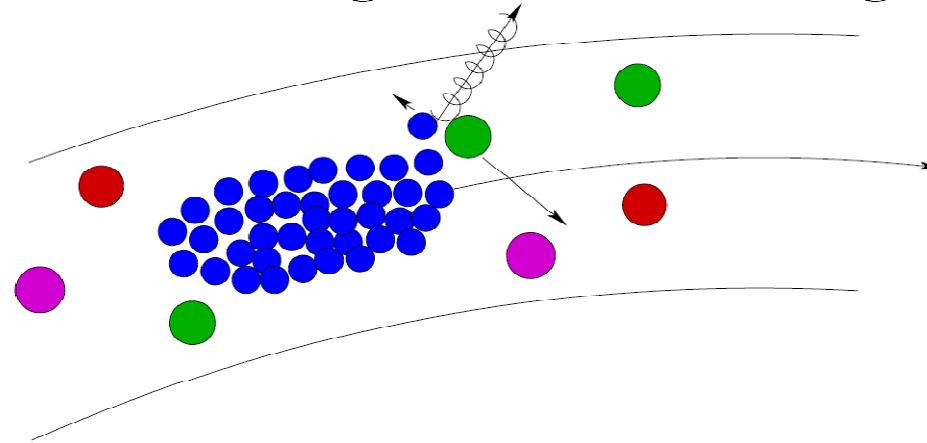
$$P \propto \frac{\gamma^4}{\rho^2}$$

$$E \propto \frac{\gamma^3}{\rho}$$

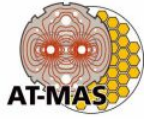
LEP:  $\gamma = 200000$

LHC:  $\gamma = 7000$

## Bremsstrahlung + Coulomb Scattering



- "Blow up" of the beam
- Particle losses
- Non wanted collisions in the experiments
- Limits the Luminosity

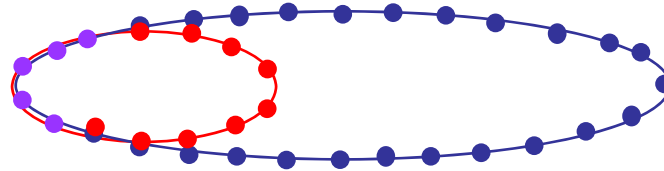


# Superconducting Technology 1

TECHNOLOGY

Why superconducting magnets?

Small radius, less number of particles in the machine, smaller machine



Energy saving, BUT infrastructure very complex

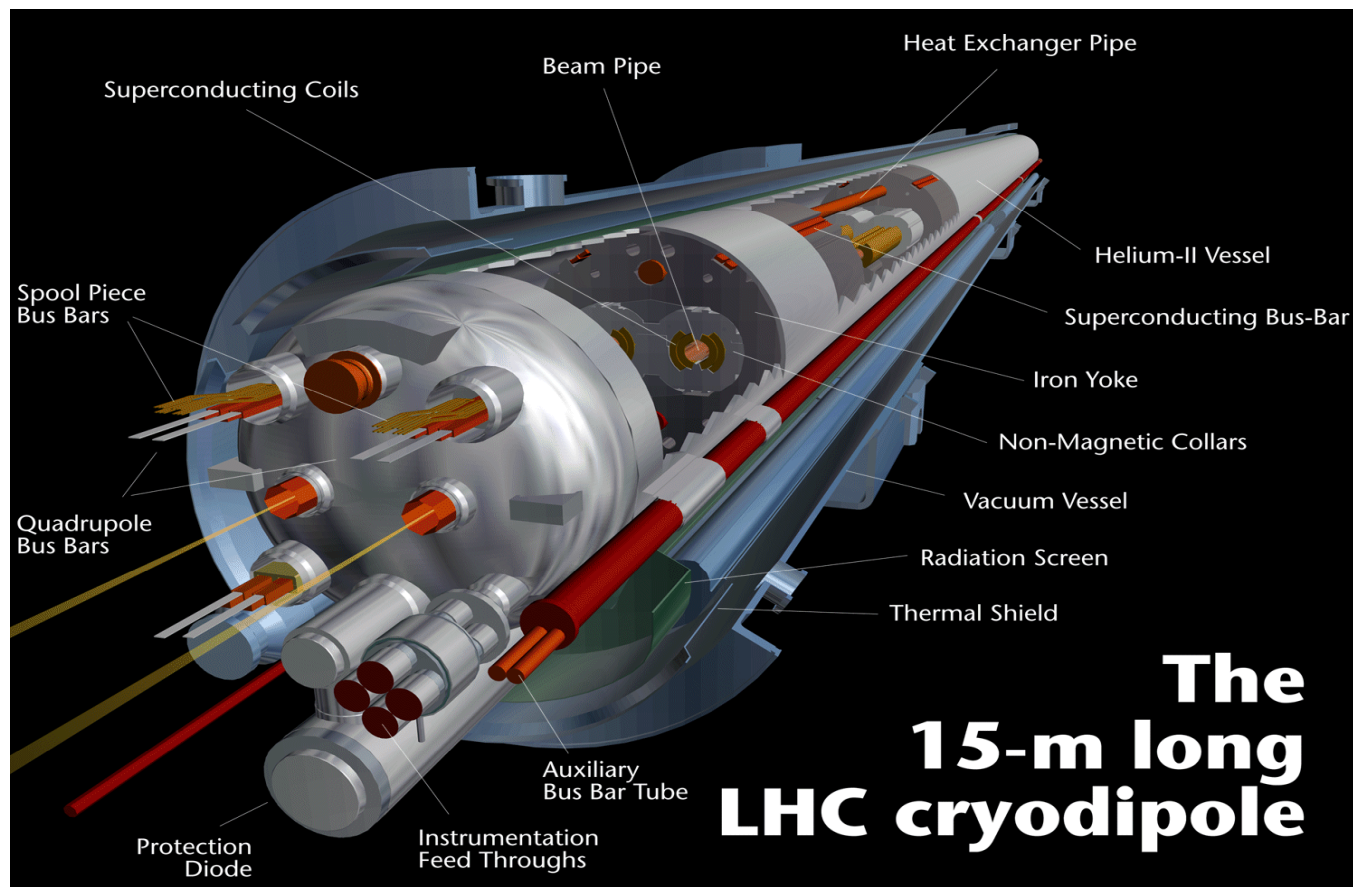


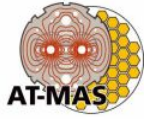


# The Superconducting Dipole for LHC

LHC dipole (1232 + reserves) built in 3 firms (Germany France and Italy, very large high tech project)

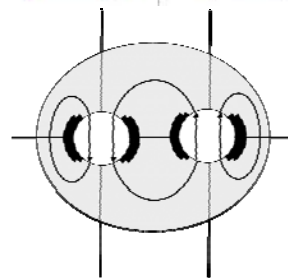
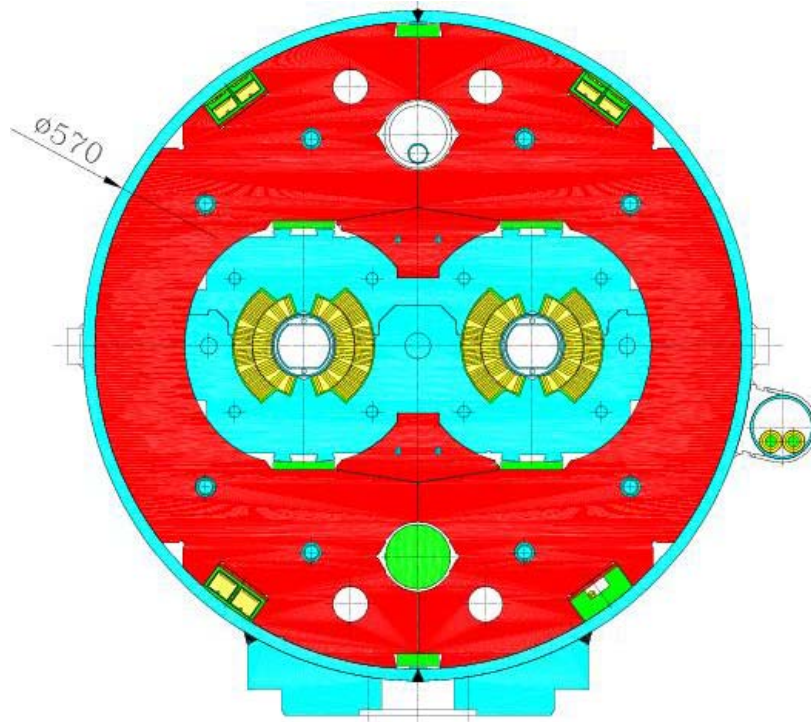
TECHNOLOGY





# The LHC Dipole

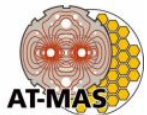
TECHNOLOGY



"Two in one"  
construction

Working  
temperature  
1.9 K!  
Coldest spot i the  
universe...





# Future accelerators

TECHNOLOGY

High Energy physics applications:

## LHC upgrade

Luminosity one order of magnitude higher

Insertion regions will change

Deposition in magnets of energy and debris

New superconductors

## Linear Colliders

Cost proportional to beam energy, for circular machines proportional to the square of the beam energy

International Linear Collider (ILC), 35 km, 500 GeV, electron-positron

Compact Linear Collider (CLIC), 38 km, 3 TeV, electron-positron

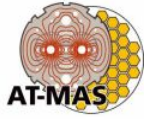


# Example: LHC 1

EXAMPLE: LHC

Table 2.1: LHC beam parameters relevant for the peak luminosity

		Injection	Collision
<b>Beam Data</b>			
Proton energy	[GeV]	450	7000
Relativistic gamma		479.6	7461
Number of particles per bunch		$1.15 \times 10^{11}$	
Number of bunches		2808	
Longitudinal emittance ( $4\sigma$ )	[eVs]	1.0	2.5 <sup>a</sup>
Transverse normalized emittance	[ $\mu\text{m rad}$ ]	3.5 <sup>b</sup>	3.75
Circulating beam current	[A]	0.582	
Stored energy per beam	[MJ]	23.3	362
<b>Peak Luminosity Related Data</b>			
RMS bunch length <sup>c</sup>	cm	11.24	7.55
RMS beam size at the IP1 and IP5 <sup>d</sup>	$\mu\text{m}$	375.2	16.7
RMS beam size at the IP2 and IP8 <sup>e</sup>	$\mu\text{m}$	279.6	70.9
Geometric luminosity reduction factor $F^f$		-	0.836
Peak luminosity in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$1.0 \times 10^{34}$
Peak luminosity per bunch crossing in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$3.56 \times 10^{30}$



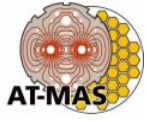
# Exempel: LHC 2

EXEMPEL: LHC

Table 2.1: LHC beam parameters relevant for the peak luminosity

		Injection	Collision
<b>Beam Data</b>			
Proton energy	[GeV]	450	7000
Relativistic gamma		479.6	7461
Number of particles per bunch		$1.15 \times 10^{11}$	
Number of bunches		2808	
Longitudinal emittance ( $4\sigma$ )	[eVs]	1.0	2.5 <sup>a</sup>
Transverse normalized emittance	[ $\mu\text{m rad}$ ]	3.5 <sup>b</sup>	3.75
Circulating beam current	[A]	0.582	
Stored energy per beam	[MJ]	23.3	362
<b>Peak Luminosity Related Data</b>			
RMS bunch length <sup>c</sup>	cm	11.24	7.55
RMS beam size at the IP1 and IP5 <sup>d</sup>	$\mu\text{m}$	375.2	16.7
RMS beam size at the IP2 and IP8 <sup>e</sup>	$\mu\text{m}$	279.6	70.9
Geometric luminosity reduction factor $F^f$		-	0.836
Peak luminosity in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$1.0 \times 10^{34}$
Peak luminosity per bunch crossing in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$3.56 \times 10^{30}$



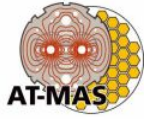


# Exempel: LHC 3

EXEMPEL: LHC

Table 2.1: LHC beam parameters relevant for the peak luminosity

		Injection	Collision
<b>Beam Data</b>			
Proton energy	[GeV]	450	7000
Relativistic gamma		479.6	7461
Number of particles per bunch		$1.15 \times 10^{11}$	
Number of bunches		2808	
Longitudinal emittance ( $4\sigma$ )	[eVs]	1.0	2.5 <sup>a</sup>
Transverse normalized emittance	[ $\mu\text{m rad}$ ]	3.5 <sup>b</sup>	3.75
Circulating beam current	[A]	0.582	
Stored energy per beam	[MJ]	23.3	362
<b>Peak Luminosity Related Data</b>			
RMS bunch length <sup>c</sup>	cm	11.24	7.55
RMS beam size at the IP1 and IP5 <sup>d</sup>	$\mu\text{m}$	375.2	16.7
RMS beam size at the IP2 and IP8 <sup>e</sup>	$\mu\text{m}$	279.6	70.9
Geometric luminosity reduction factor $F^f$		-	0.836
Peak luminosity in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$1.0 \times 10^{34}$
Peak luminosity per bunch crossing in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$3.56 \times 10^{30}$

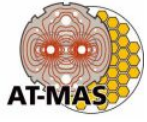


# Exempel: LHC 4

EXEMPEL: LHC

Table 2.1: LHC beam parameters relevant for the peak luminosity

		Injection	Collision
<b>Beam Data</b>			
Proton energy	[GeV]	450	7000
Relativistic gamma		479.6	7461
Number of particles per bunch		$1.15 \times 10^{11}$	
Number of bunches		2808	
Longitudinal emittance ( $4\sigma$ )	[eVs]	1.0	2.5 <sup>a</sup>
Transverse normalized emittance	[ $\mu\text{m rad}$ ]	3.5 <sup>b</sup>	3.75
Circulating beam current	[A]	0.582	
Stored energy per beam	[MJ]	23.3	362
<b>Peak Luminosity Related Data</b>			
RMS bunch length <sup>c</sup>	cm	11.24	7.55
RMS beam size at the IP1 and IP5 <sup>d</sup>	$\mu\text{m}$	375.2	16.7
RMS beam size at the IP2 and IP8 <sup>e</sup>	$\mu\text{m}$	279.6	70.9
Geometric luminosity reduction factor $F^f$		-	0.836
Peak luminosity in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$1.0 \times 10^{34}$
Peak luminosity per bunch crossing in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$3.56 \times 10^{30}$



# Exempel: LHC 5

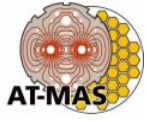
EXEMPEL: LHC

Table 2.1: LHC beam parameters relevant for the peak luminosity

		Injection	Collision
<b>Beam Data</b>			
Proton energy	[GeV]	450	7000
Relativistic gamma		479.6	7461
Number of particles per bunch		$1.15 \times 10^{11}$	
Number of bunches		2808	
Longitudinal emittance ( $4\sigma$ )	[eVs]	1.0	2.5 <sup>a</sup>
Transverse normalized emittance	[ $\mu\text{m rad}$ ]	3.5 <sup>b</sup>	3.75
Circulating beam current	[A]	0.582	
Stored energy per beam	[MJ]	23.3	362
<b>Peak Luminosity Related Data</b>			
RMS bunch length <sup>c</sup>	cm	11.24	7.55
RMS beam size at the IP1 and IP5 <sup>d</sup>	$\mu\text{m}$	375.2	16.7
RMS beam size at the IP2 and IP8 <sup>e</sup>	$\mu\text{m}$	279.6	70.9
Geometric luminosity reduction factor $F^f$		-	0.836
Peak luminosity in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$1.0 \times 10^{34}$
Peak luminosity per bunch crossing in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$3.56 \times 10^{30}$

Arcs:  $\sim 1$  mm transversellt



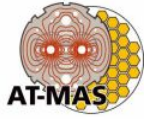


# Exempel: LHC 6

EXEMPEL: LHC

Table 2.1: LHC beam parameters relevant for the peak luminosity

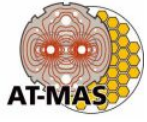
		Injection	Collision
<b>Beam Data</b>			
Proton energy	[GeV]	450	7000
Relativistic gamma		479.6	7461
Number of particles per bunch		$1.15 \times 10^{11}$	
Number of bunches		2808	
Longitudinal emittance ( $4\sigma$ )	[eVs]	1.0	2.5 <sup>a</sup>
Transverse normalized emittance	[ $\mu\text{m rad}$ ]	3.5 <sup>b</sup>	3.75
Circulating beam current	[A]	0.582	
Stored energy per beam	[MJ]	23.3	362
<b>Peak Luminosity Related Data</b>			
RMS bunch length <sup>c</sup>	cm	11.24	7.55
RMS beam size at the IP1 and IP5 <sup>d</sup>	$\mu\text{m}$	375.2	16.7
RMS beam size at the IP2 and IP8 <sup>e</sup>	$\mu\text{m}$	279.6	70.9
Geometric luminosity reduction factor $F^f$		-	0.836
Peak luminosity in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$1.0 \times 10^{34}$
Peak luminosity per bunch crossing in IP1 and IP5	[ $\text{cm}^{-2}\text{sec}^{-1}$ ]	-	$3.56 \times 10^{30}$



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Special thanks to Oliver Bruning for the reference list and for some material