

Theory perspective on precision Higgs measurements

*EUCARD LEP3
CERN, 18 June 2012*

Christophe Grojean
CERN-TH

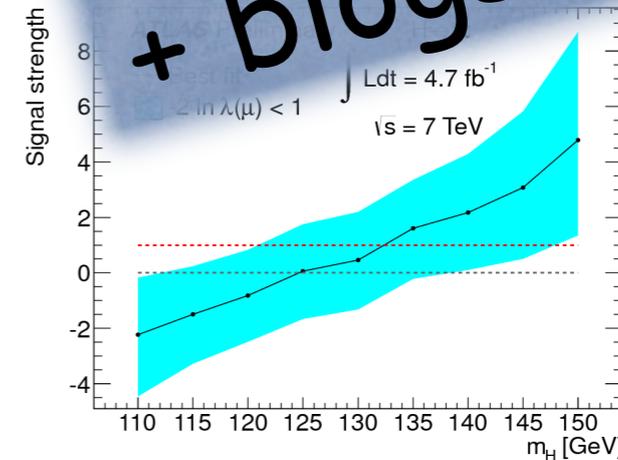
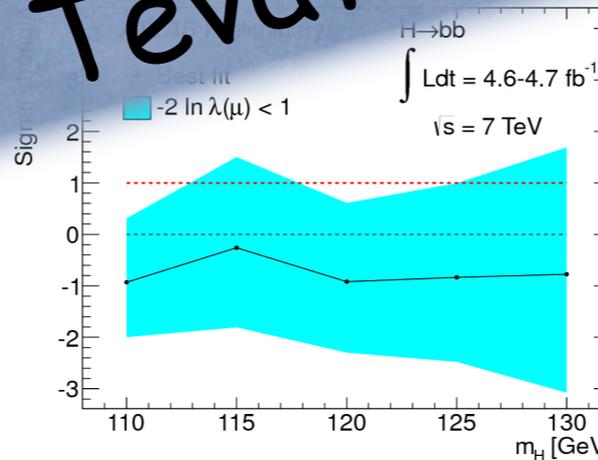
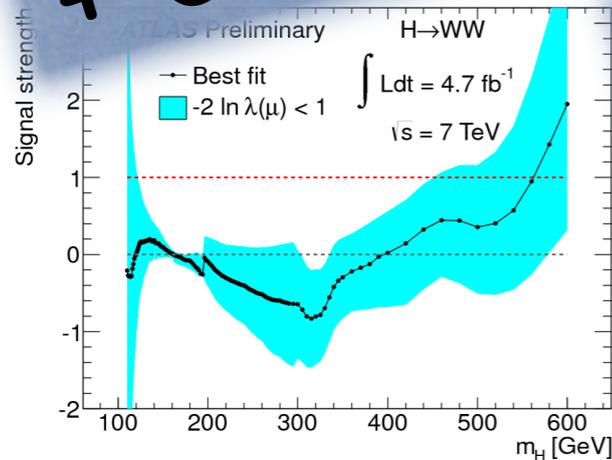
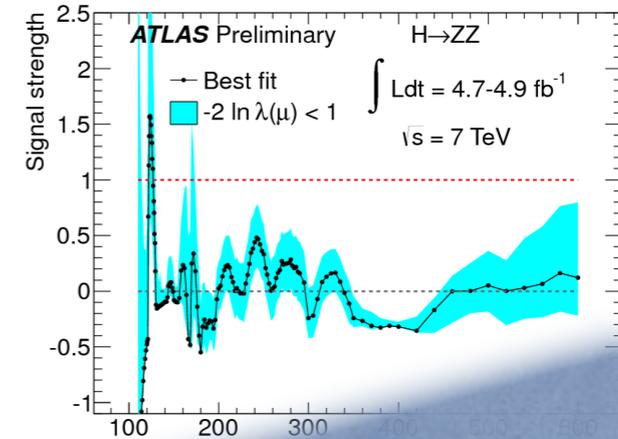
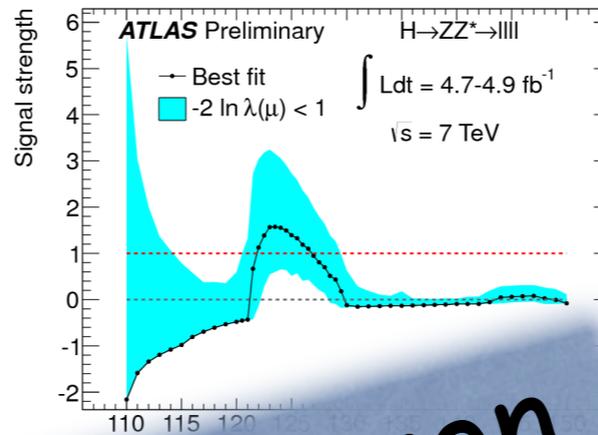
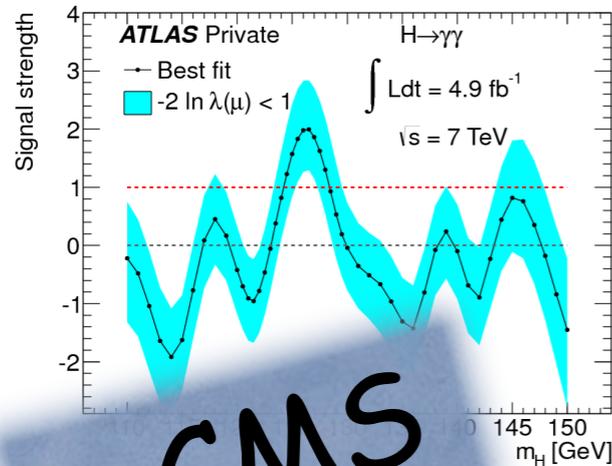
(christophe.grojean@cern.ch)



Facts that we have to live with

signal strength

$$\mu_i = \frac{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i)}{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i) |_{\text{SM}}}$$



+ CMS

+ Tevatron

+ blogs...

Still not enough information:

Correlations?
 Exact likelihoods?

Chiral Lagrangian for a light Higgs

$$\begin{aligned}
 \mathcal{L} = & \frac{1}{2}(\partial_\mu h)^2 - \frac{1}{2}m_h^2 h^2 - \frac{d_3}{6} \left(\frac{3m_h^2}{v} \right) h^3 - \frac{d_4}{24} \left(\frac{3m_h^2}{v^2} \right) h^4 \dots \\
 & - \left(m_W^2 W_\mu W_\mu + \frac{1}{2}m_Z^2 Z_\mu Z_\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right) \\
 & - \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left(1 + c_\psi \frac{h}{v} + c_{2\psi} \frac{h^2}{v^2} + \dots \right) \\
 & + \frac{g^2}{16\pi^2} \left(c_{WW} W_{\mu\nu}^+ W_{\mu\nu}^- + c_{ZZ} Z_{\mu\nu}^2 + c_{Z\gamma} Z_{\mu\nu} \gamma_{\mu\nu} \right) \frac{h}{v} + \dots \\
 & + \frac{g^2}{16\pi^2} \left[\gamma_{\mu\nu}^2 \left(c_{\gamma\gamma} \frac{h}{v} + \dots \right) + G_{\mu\nu}^2 \left(c_{gg} \frac{h}{v} + c_{2gg} \frac{h^2}{v^2} \dots \right) \right] \\
 & + \frac{g^2}{16\pi^2} \left[\frac{c_{hhgg}}{\Lambda^2} G_{\mu\nu}^2 \frac{(\partial_\rho h)^2}{v^2} + \frac{c'_{hhgg}}{\Lambda^2} G_{\mu\rho} G_{\rho\nu} \frac{\partial_\mu h \partial_\nu h}{v^2} + \dots \right] \\
 & + \dots
 \end{aligned}$$

A few (reasonable) assumptions:

spin-0 & CP-even



 $\gamma\gamma$ WW & ZZ

custodial symmetry



 EWPD

no Higgs FCNC



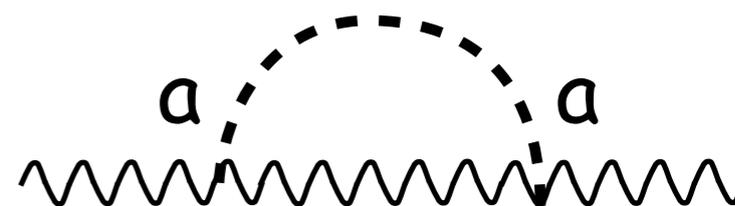
 Flavor

Contino, Grojean, Moretti, Piccinini, Rattazzi '10

Azatov, Contino, Galloway '12

EW constraints

The parameter 'a' controls the size of the one-loop IR contribution to the LEP precision observables



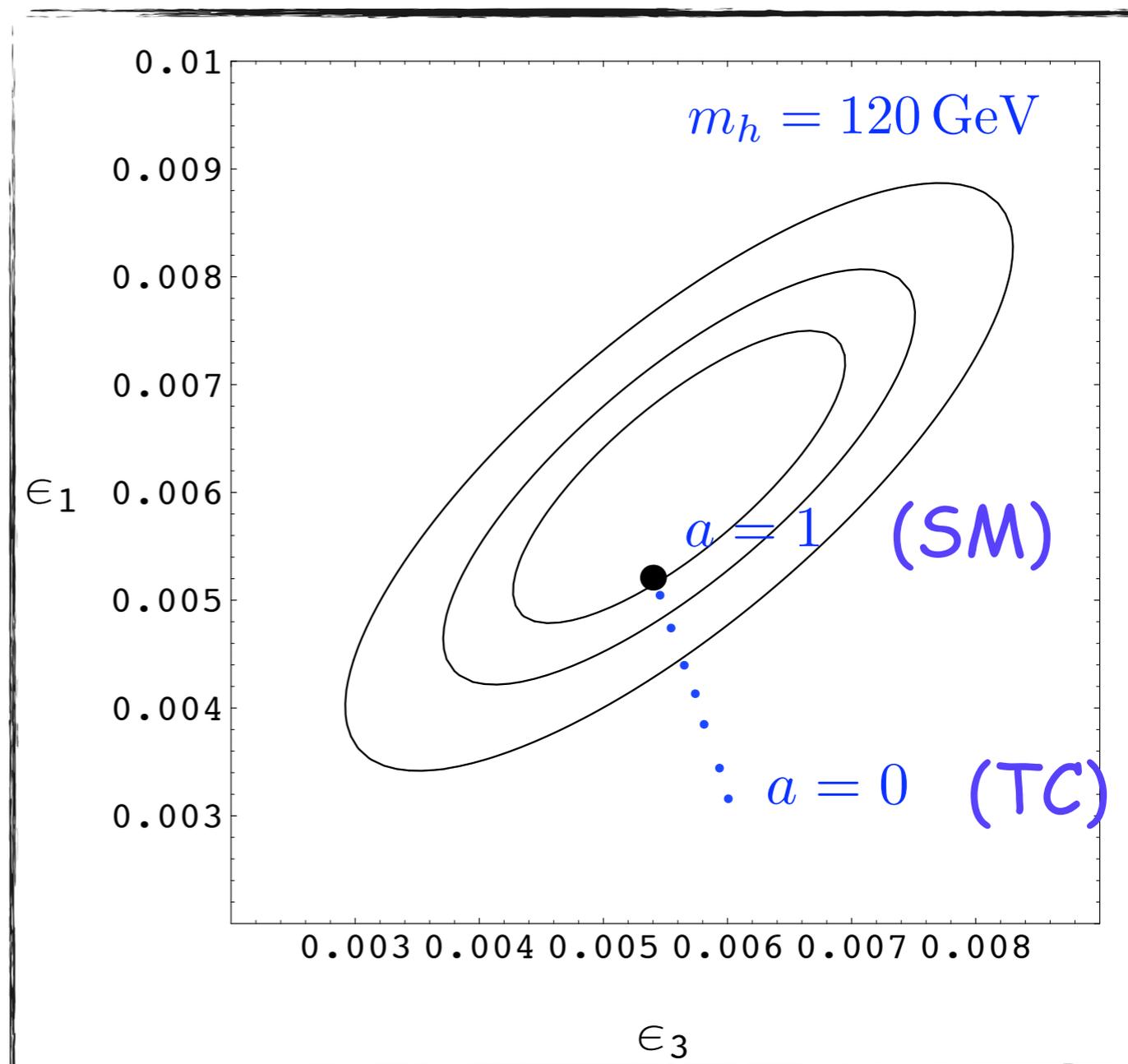
$$\epsilon_{1,3} = c_{1,3} \log(m_Z^2/\mu^2) - c_{1,3} a^2 \log(m_h^2/\mu^2) - c_{1,3} (1 - a^2) \log(m_\rho^2/\mu^2) + \text{finite terms}$$

$$c_1 = + \frac{3}{16\pi^2} \frac{\alpha(m_Z)}{\cos^2 \theta_W}$$

$$c_3 = - \frac{1}{12\pi} \frac{\alpha(m_Z)}{4 \sin^2 \theta_W}$$

$$\Delta\epsilon_{1,3} = -c_{1,3} (1 - a^2) \log(m_\rho^2/m_h^2)$$

Barbieri, Bellazzini, Rychkov, Varagnolo '07



EW data constraints on 'a'

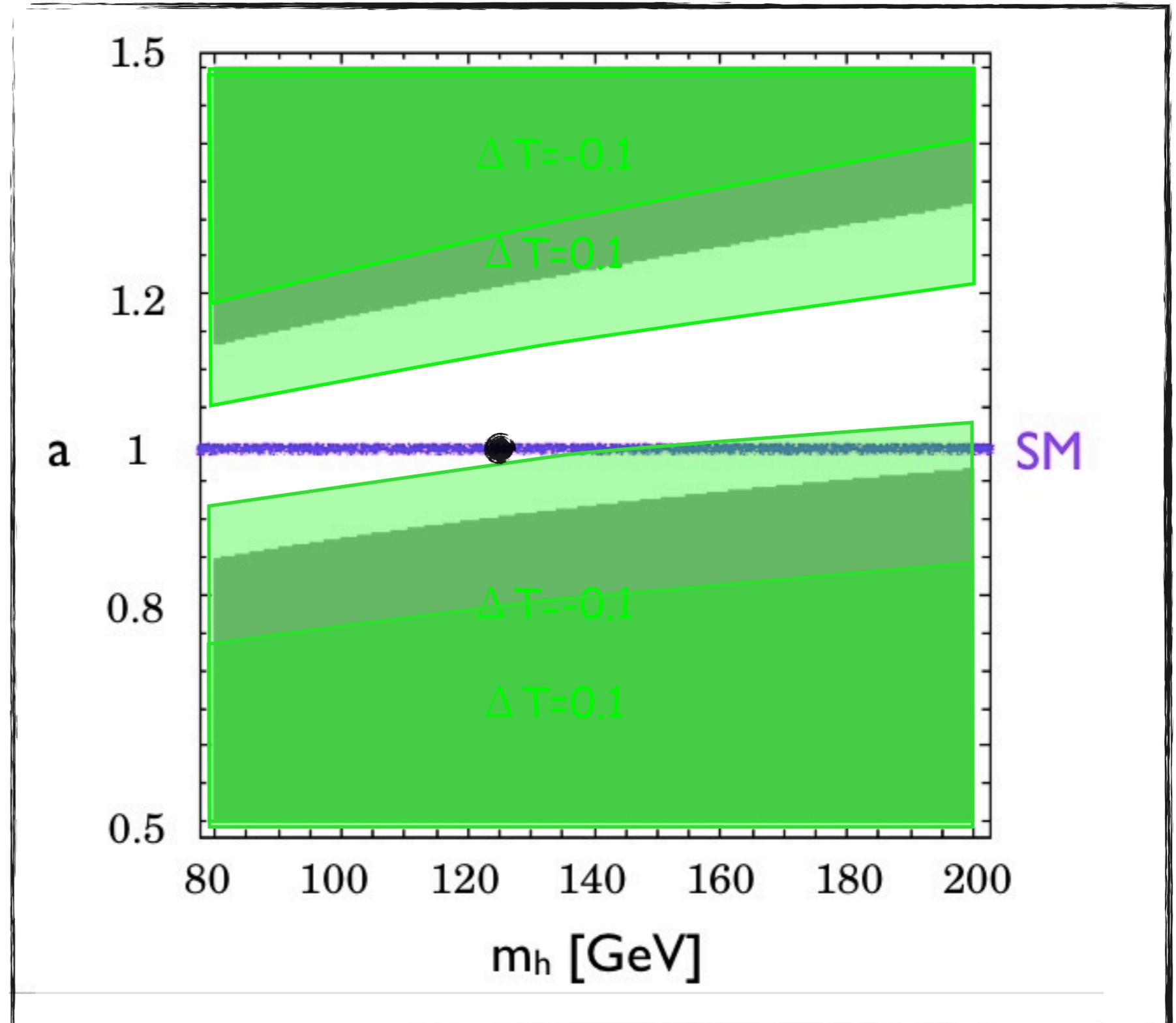
EW fit with SM degrees of freedom + (composite) Higgs

- EW data require less than 15-20% deviations in the couplings of the Higgs to gauge bosons

- EW data don't constraint the other Higgs couplings

note:

additional UV contributions to S and T can modify the preferred values of couplings



Chiral Lagrangian for a light Higgs @ LHC

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \frac{1}{2}m_h^2 h^2 - \frac{d_3}{6} \left(\frac{3m_h^2}{v} \right) h^3 - \frac{d_4}{24} \left(\frac{3m_h^2}{v^2} \right) h^4 \dots$$

$$- \left(m_W^2 W_\mu W_\mu + \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right)$$

$$- \sum_{\psi=u,d,l} m_{\psi^{(i)}} \bar{\psi}^{(i)} \psi^{(i)} \left(1 + c_\psi \frac{h}{v} + \dots \right)$$

still too much freedom
 \Downarrow
 dynamical assumptions needed
 \Downarrow
 to explore deformations of the SM

$$+ G_{\mu\nu}^2 \left(c_{gg} \frac{h}{v} + c_{2gg} \frac{h^2}{v^2} \dots \right)$$

$$+ \frac{g^2}{16\pi^2} \left[\frac{c_{hhgg}}{\Lambda^2} G_{\mu\nu}^2 \frac{(\partial_\rho h)^2}{v^2} + \frac{c'_{hhgg}}{\Lambda^2} G_{\mu\rho} G_{\rho\nu} \frac{\partial_\mu h \partial_\nu h}{v^2} + \dots \right]$$

+ ...

A few (reasonable) assumptions:

spin-0 & CP-even



$\gamma\gamma$



WW & ZZ

custodial symmetry



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no Higgs FCNC



Flavor

Contino, Grojean, Moretti, Piccinini, Rattazzi '10

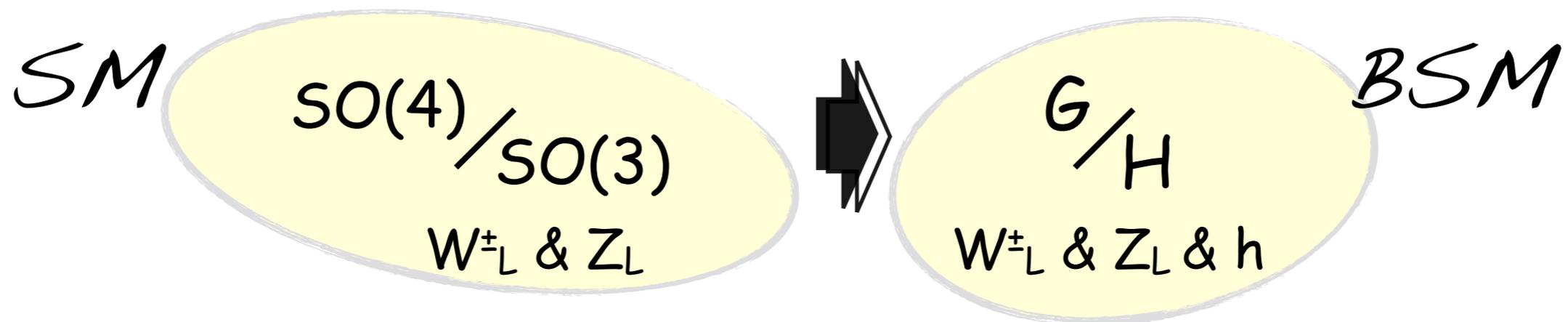
Azatov, Contino, Galloway '12

The New Physics Mass Gap

One solution to the hierarchy pb:

Higgs transforms non-linearly under some global symmetry

Higgs=Pseudo-Goldstone boson (PGB)



Examples: $SO(5)/SO(4)$: 4 PGBs = W^\pm_L, Z_L, h

Minimal Composite Higgs Model

Agashe, Contino, Pomarol '04

$SO(6)/SO(5)$: 5 PGBs = H, a

Next MCHM

Gripaios, Pomarol, Riva, Serra '09

$SU(4)/Sp(4, \mathbb{C})$: 5 PGBs = H, s

$SO(6)/SO(4) \times SO(2)$: 8 PGBs = $H_1 + H_2$

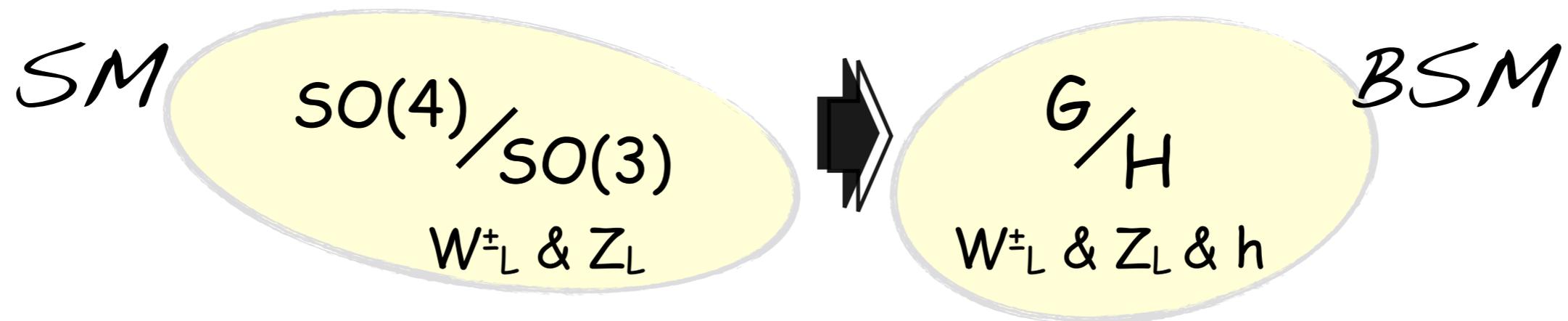
Minimal Composite Two Higgs Doublets

Mrazek, Pomarol, Rattazzi, Serra, Wulzer '11

The New Physics Mass Gap

One solution to the hierarchy pb:
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How can we tell the difference with the SM Higgs?

What are the experimental constraints?

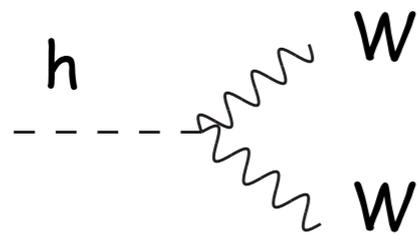
2 parameter Higgs physics @ LHC 2011-2012

$$\frac{c_H}{2f^2} (\partial^\mu |H|^2)^2$$

$$\xi = v^2 / f^2$$

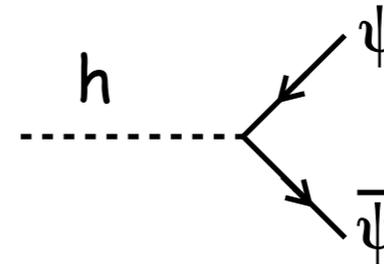
$$\frac{c_\psi y_\psi}{f^2} |H|^2 \bar{\psi}_L H \psi_R$$

Controls the hWW , hZZ couplings



$$a g_{hVV}^{SM}$$

Controls the $h\psi\psi$ couplings



$$c g_{h\psi\psi}^{SM}$$

Explicit (and calculable) models built in AdS_5 spacetimes

Agashe, Contino, Pomarol '04

Contino, Da Rold, Pomarol '06

MCHM5

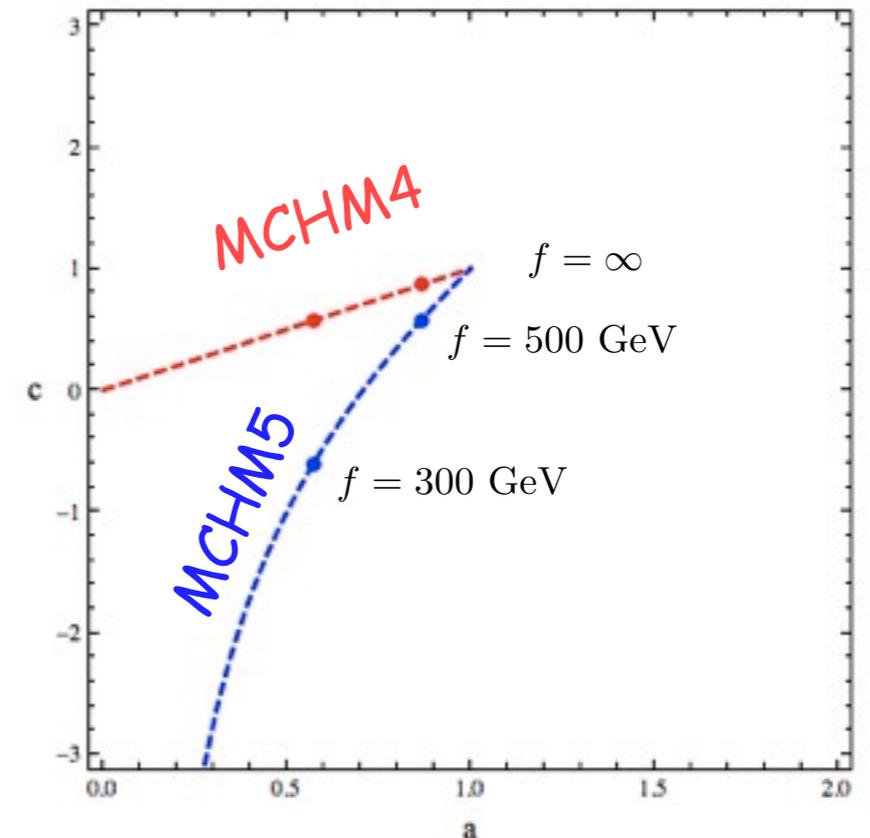
$$a = \sqrt{1 - \xi} \quad c = \frac{1 - 2\xi}{\sqrt{1 - \xi}}$$

MCHM4

$$a = \sqrt{1 - \xi} \quad c = \sqrt{1 - \xi}$$

disfavored by EW data (Zbb)

SM is recovered as a limit when the compositeness scale is well above weak scale



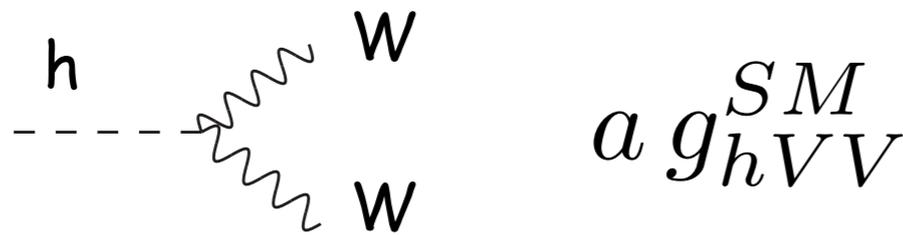
2 parameter Higgs physics @ LHC 2011-2012

$$\frac{c_H}{2f^2} \left(\partial^\mu |H|^2 \right)^2$$

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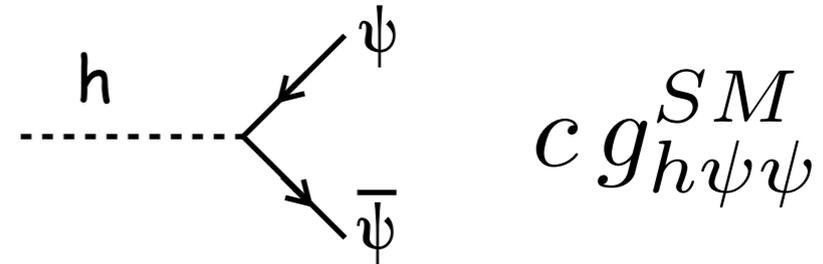
$$\frac{c_y y_\psi}{f^2} |H|^2 \bar{\psi}_L H \psi_R$$

Controls the hWW, hZZ couplings



$$a = 1 - c_H \xi / 2$$

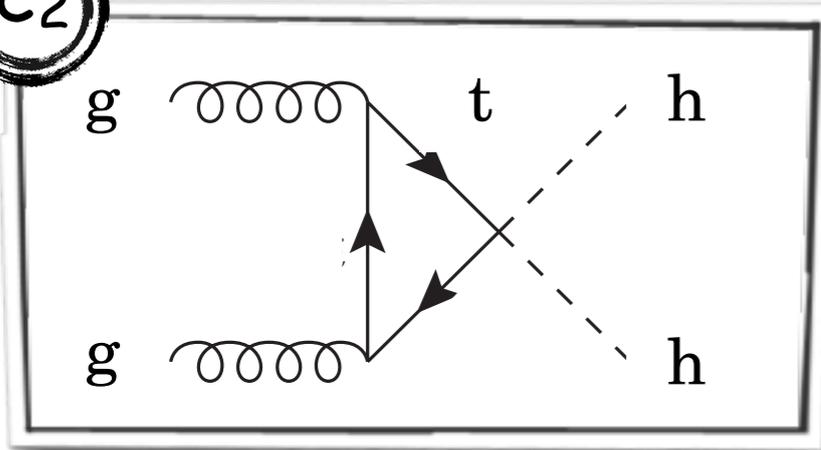
Controls the $h\psi\psi$ couplings



$$c = 1 - (c_H + 2c_y) \xi / 2$$

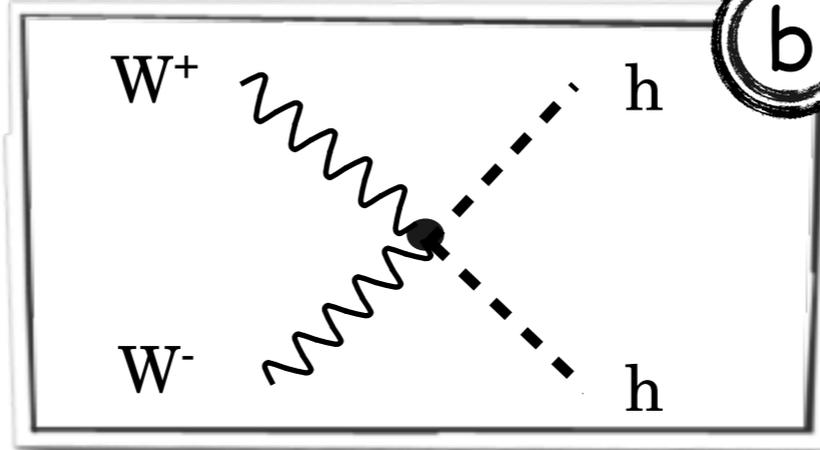
other couplings are very interesting as they are directly testing non-linearities/strong interactions of the Higgs

C2



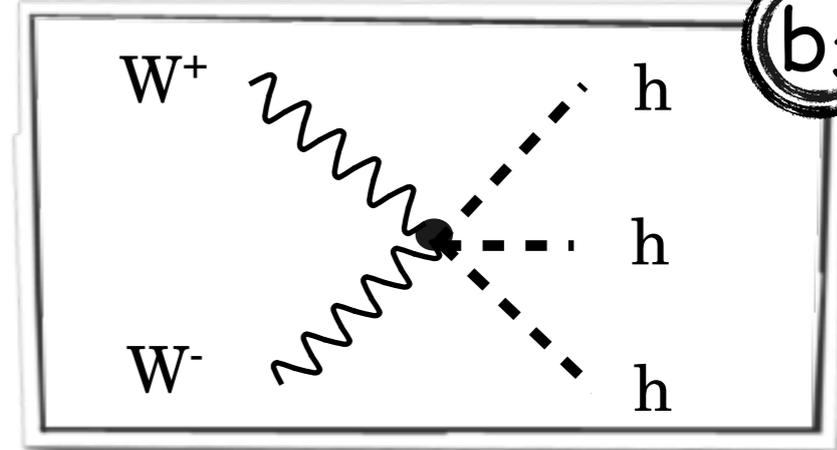
Gröber, Mühlleitner '10
Contino et al '12
Gillioz et al 'to appear

b



Contino, Grojean,
Moretti, Piccinini, Rattazzi '10

b3



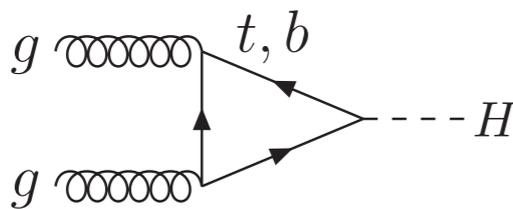
Contino, Grojean, Pappadopulo,
Rattazzi, Thamm 'to appear

but they are not on agenda of the current LHC run

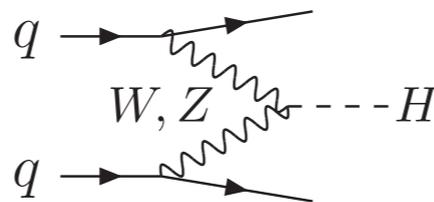
2 parameter Higgs physics @ LHC₂₀₁₁₋₂₀₁₂

- Higgs couplings modified w.r.t. SM but same kinematics
(particular to single Higgs process - with more than one Higgs, sensitive to derivative couplings)
- Background processes unaffected

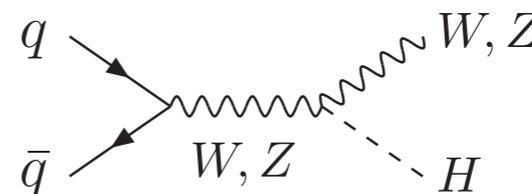
\Downarrow \Downarrow \Downarrow
 simple rescaling of SM searches



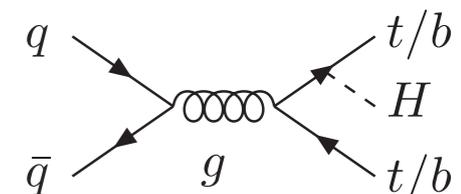
c^2



a^2



a^2



c^2

$$\frac{\sigma_{NLO}^{SM}}{\sigma_{NLO}}$$



The QCD NLO rescale trivially in the flavor universal limit.
Not the EW NLO



$$\Gamma(H \rightarrow f\bar{f}) = c^2 \Gamma^{SM}(H \rightarrow f\bar{f}),$$

$$\Gamma(H \rightarrow VV) = a^2 \Gamma^{SM}(H \rightarrow VV),$$

$$\Gamma(H \rightarrow gg) = c^2 \Gamma^{SM}(H \rightarrow gg),$$

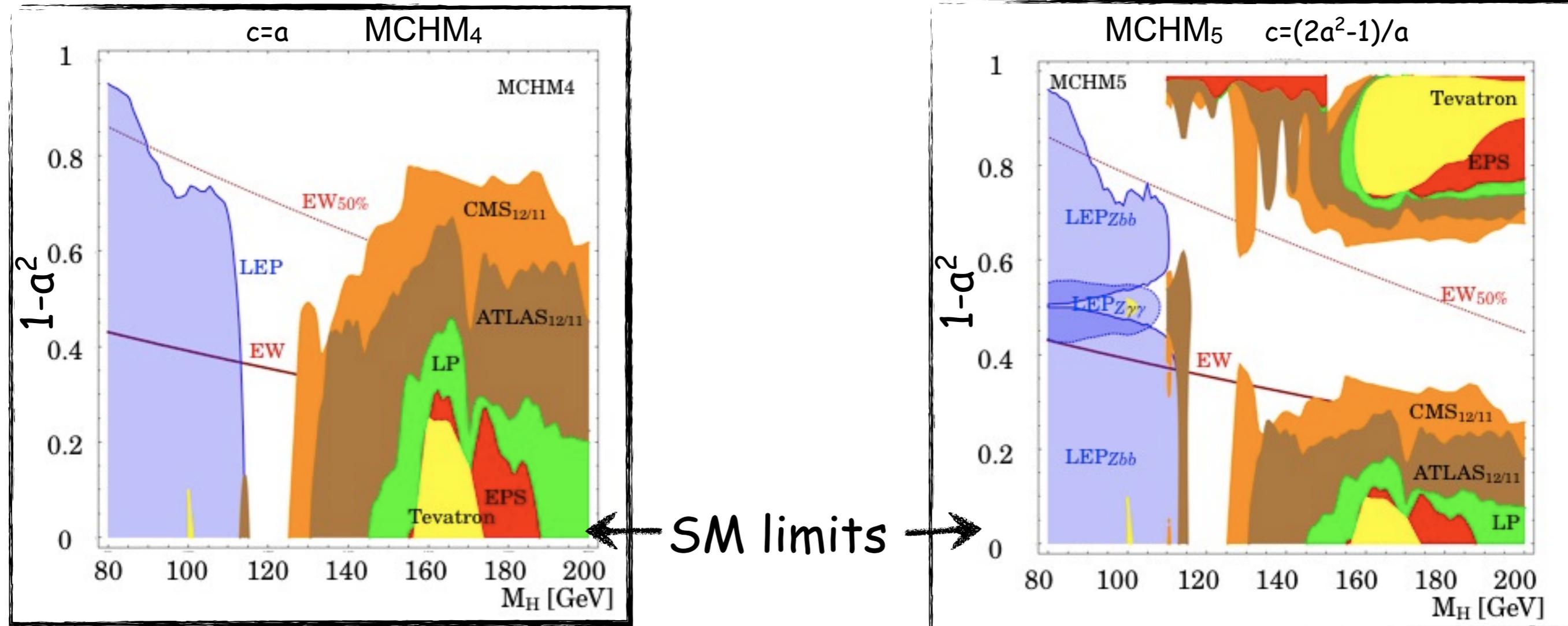
$$\Gamma(H \rightarrow \gamma\gamma) = \frac{(cI_\gamma + aJ_\gamma)^2}{(I_\gamma + J_\gamma)^2} \Gamma^{SM}(H \rightarrow \gamma\gamma),$$

$$\simeq (1.26a - 0.26c)^2 \text{ for } m_h = 125 \text{ GeV}$$

Deformation of the SM Higgs: current constraints

the SM exclusion bounds are easily rescaled in the (m_H, a) plane

Espinosa, Grojean, Muehlleitner '11

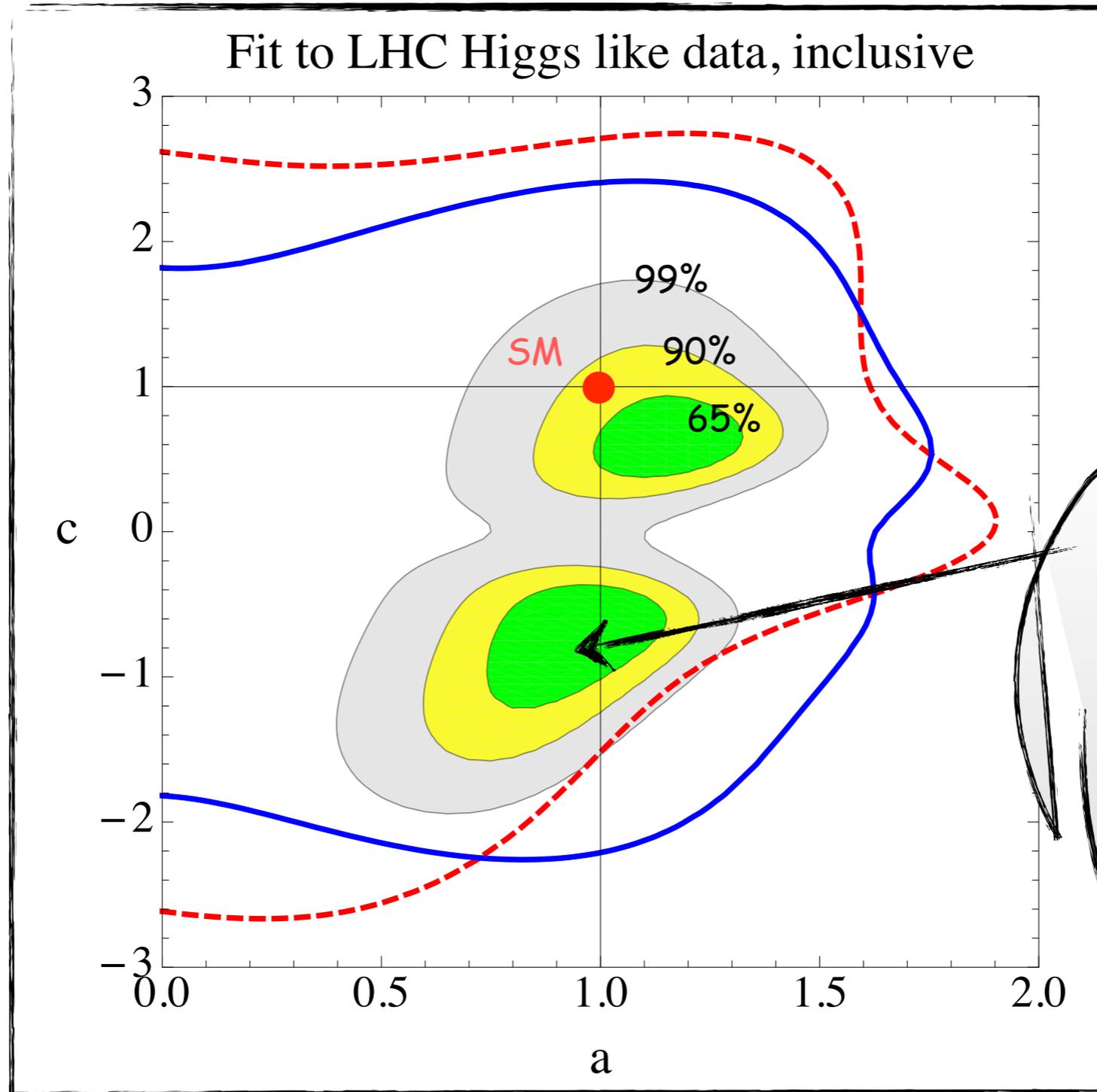


the LHC can do much more than simply excluding the SM Higgs

for similar analysis, see also [Azatov, Contino, Galloway '12](#)

Model independent χ^2 fit to LHC excess @ 125

Espinosa, Grojean, Muhlleitner, Trott '12



note: a fermiophobic Higgs is disfavored by data (mostly VBF channels)

"disfermiophilia"

the current data prefers "negative" coupling to fermions
 \approx
 positive interference between top and W in $\gamma\gamma$ channel

 Atlas 95%CL exclusion

—
 CMS 95%CL exclusion

SM 82%CL
 away from
 best fit point

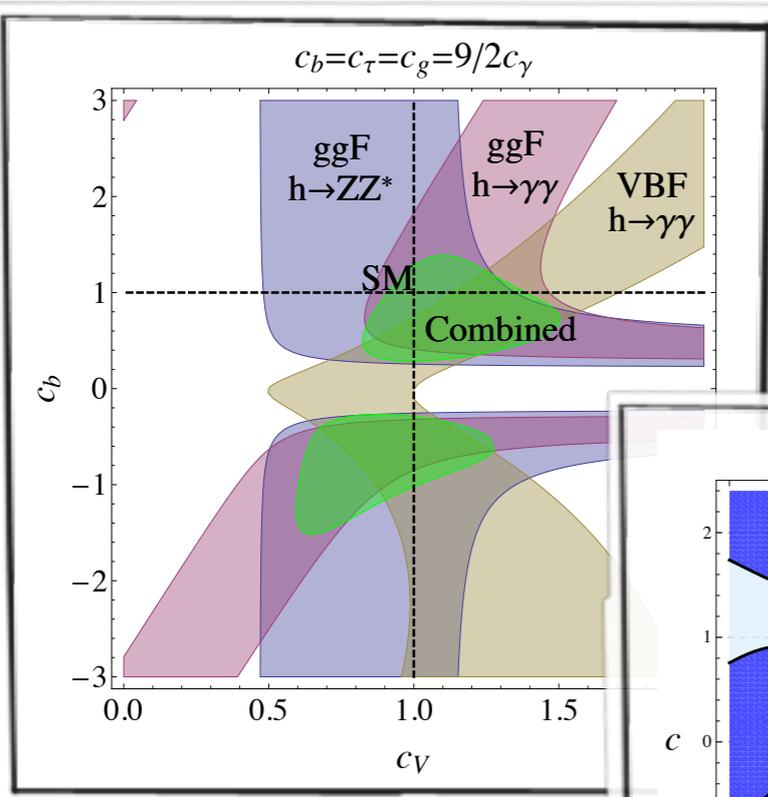
Two minima:

$(a,c)=(1.13,0.58)$
 $\chi^2=2.86$

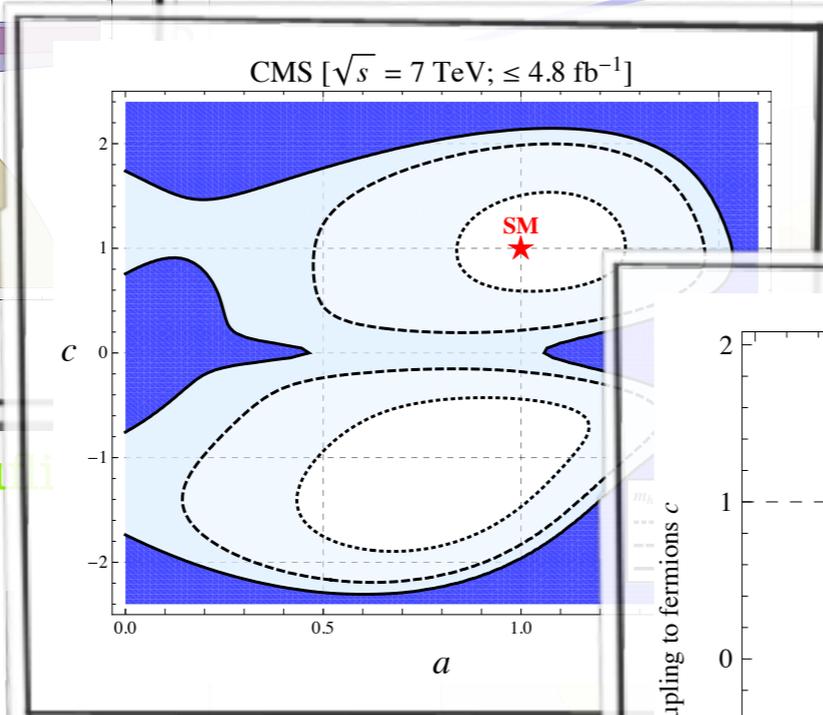
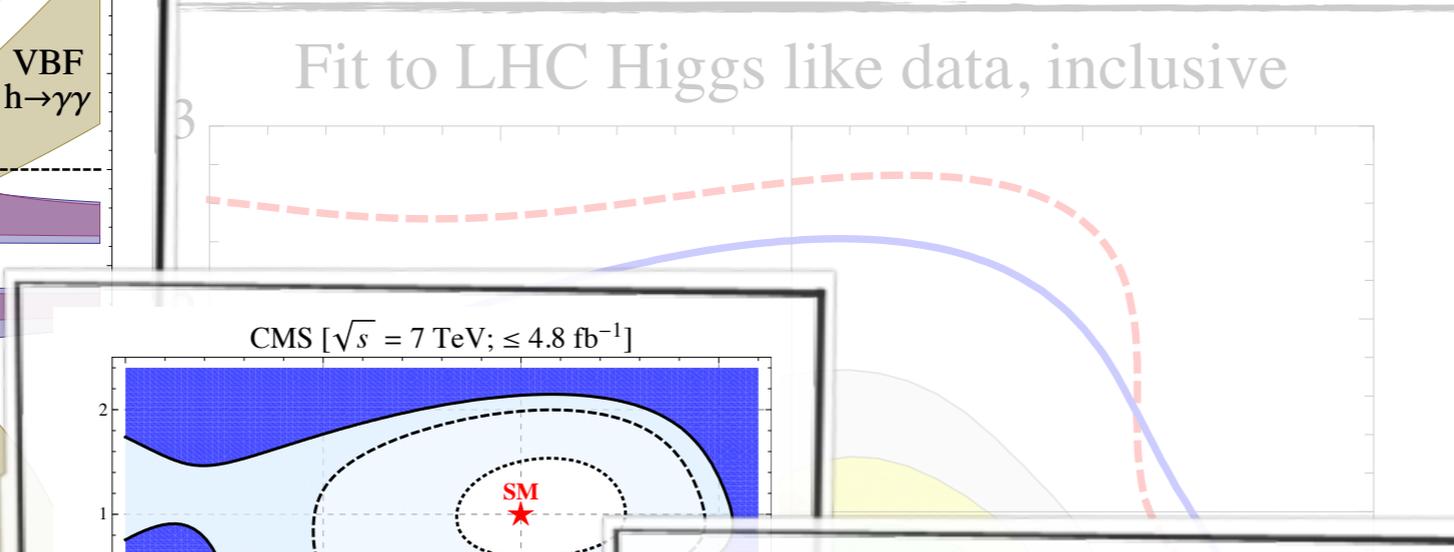
$(a,c)=(0.96,-0.64)$
 $\chi^2=1.96$

Model independent χ^2 fit to LHC excess @ 125

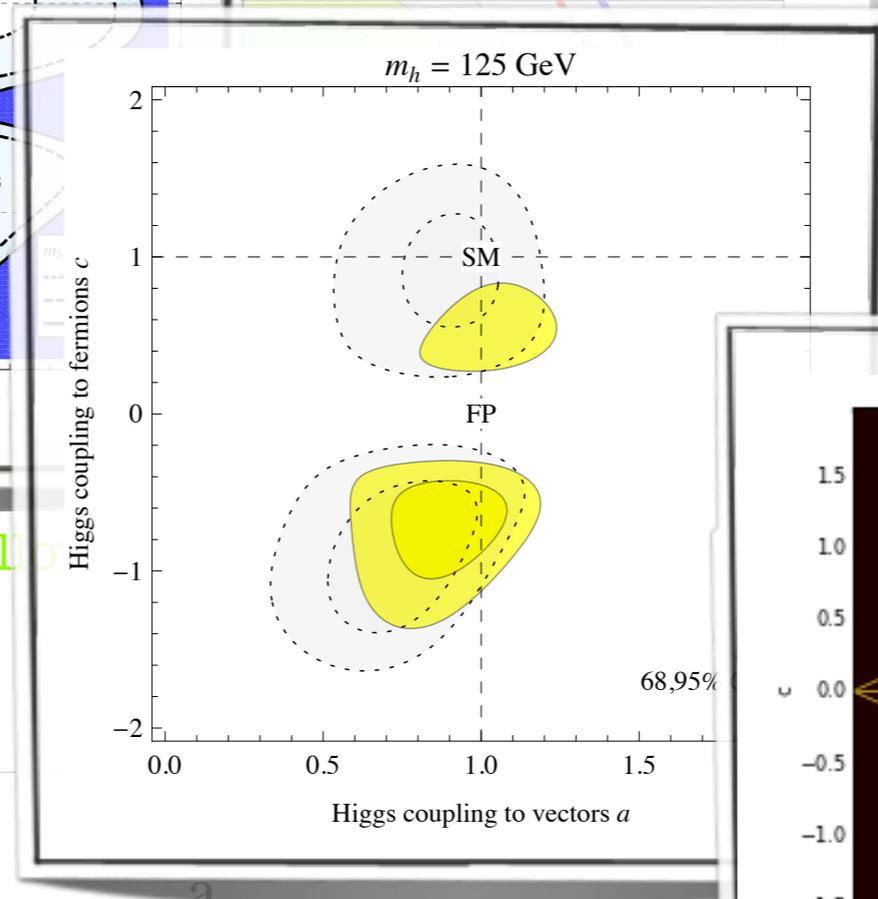
Espinosa, Grojean, Muhlleitner, Trott '12



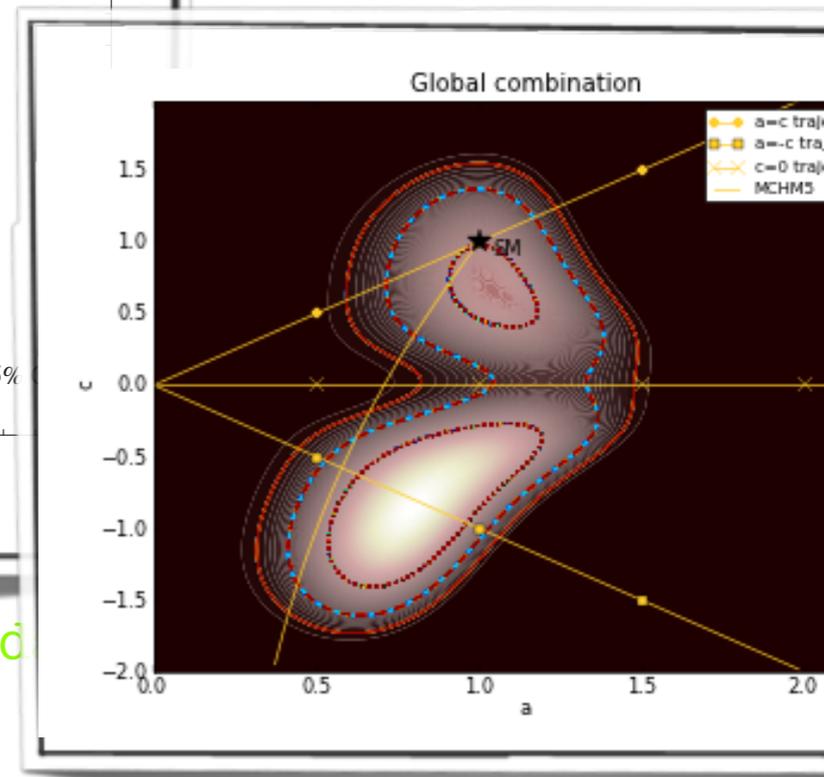
Carni, Falkowski, Kulesh, Volansky '12



Azatov, Contino, Gallo '12



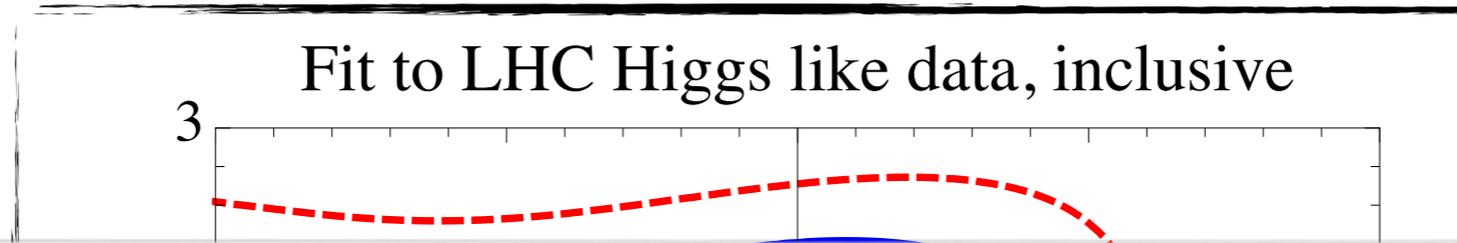
Giardino, Kannike, Raidal, Strumia '12



Ellis, You '12

Model independent χ^2 fit to LHC excess @ 125

Espinosa, Grojean, Muhlleitner, Trott '12



note: a fermiophobic Higgs is disfavored by

many issues to "validate" this kind of fits:

1. exact likelihood (departure from Gaussians...)
2. correlations among channels
3. combination of rescaled channels
4. cut efficiencies
5. ...
6. missing parameters: $c_t \neq c_b$, $a_W \neq a_Z$...

(a,c)=(0.96,-0.64)

$\chi^2=1.96$

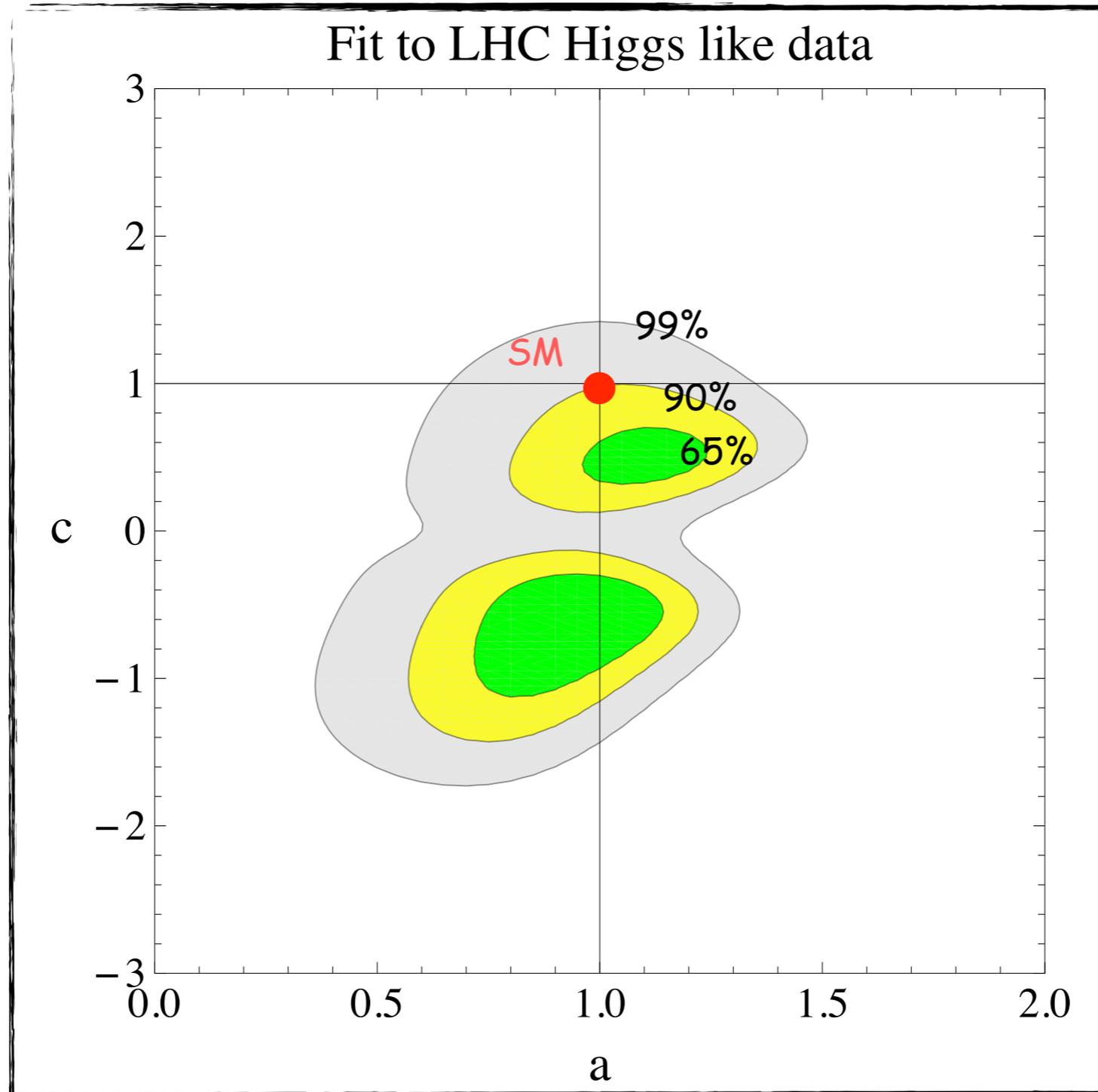
for similar analyses, see also

Azatov, Contino, Galloway '12

Carni, Falkowski, Kuflik, Volansky '12

Model independent χ^2 fit to (Moriond) LHC data

Espinosa, Grojean, Muhlleitner, Trott '12



note: a fermiophobic Higgs is disfavored by data (mostly VBF channels) at 97%CL

SM 88%CL
away from
best fit point
($\sim 2\sigma$)

Two minima:

$(a,c)=(1.18,0.55)$
 $\chi^2=7.5$

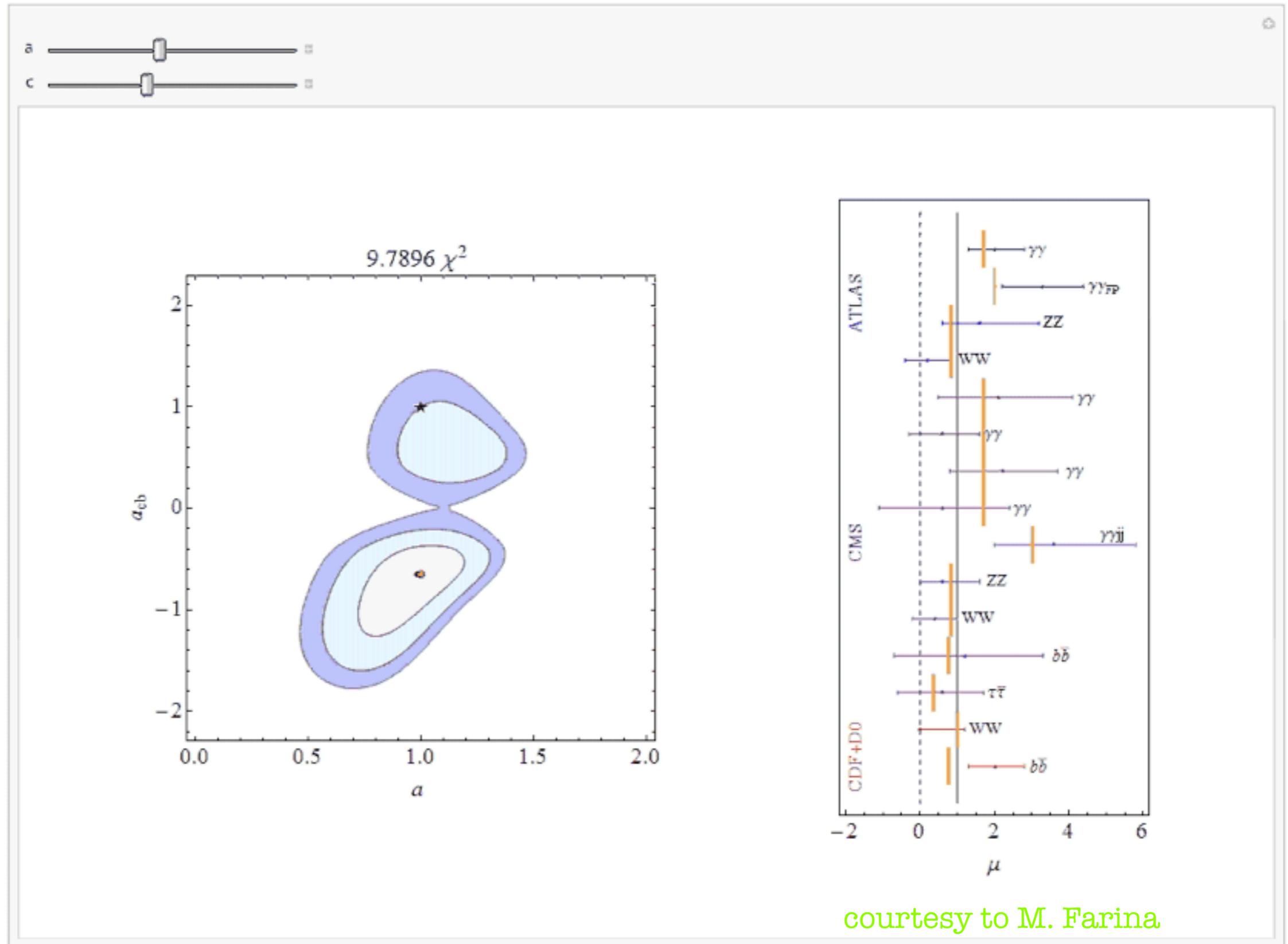
$(a,c)=(0.99,-0.64)$
 $\chi^2=6.3$

Azatov, Contino, Galloway '12

Carni, Falkowski, Kuflik, Volansky '12

for similar analyses, see also

Model independent χ^2 fit to LHC data



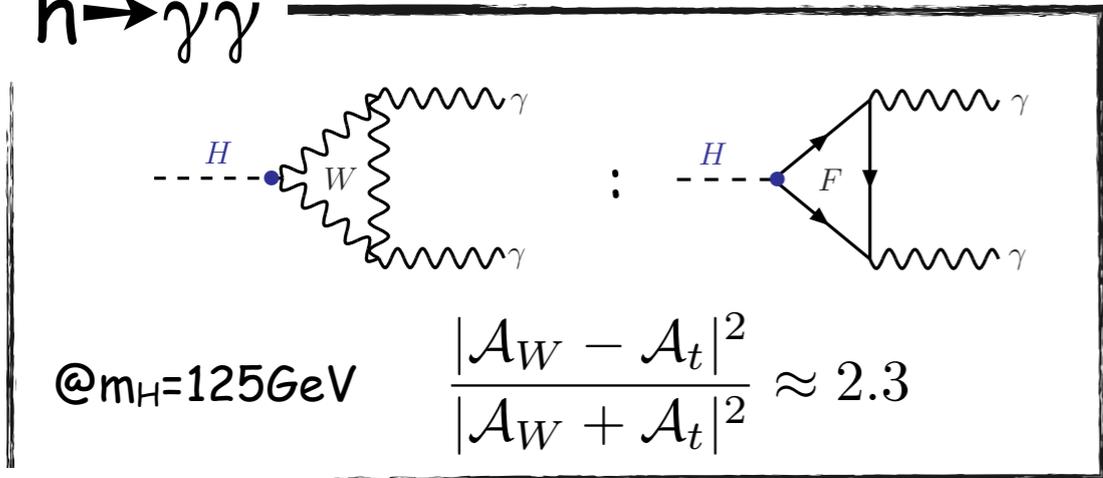
Fermiophilia or Disfermiophilia?

Farina, Grojean, Maltoni,
Salvioni, Thamm 'in progress

difficult!

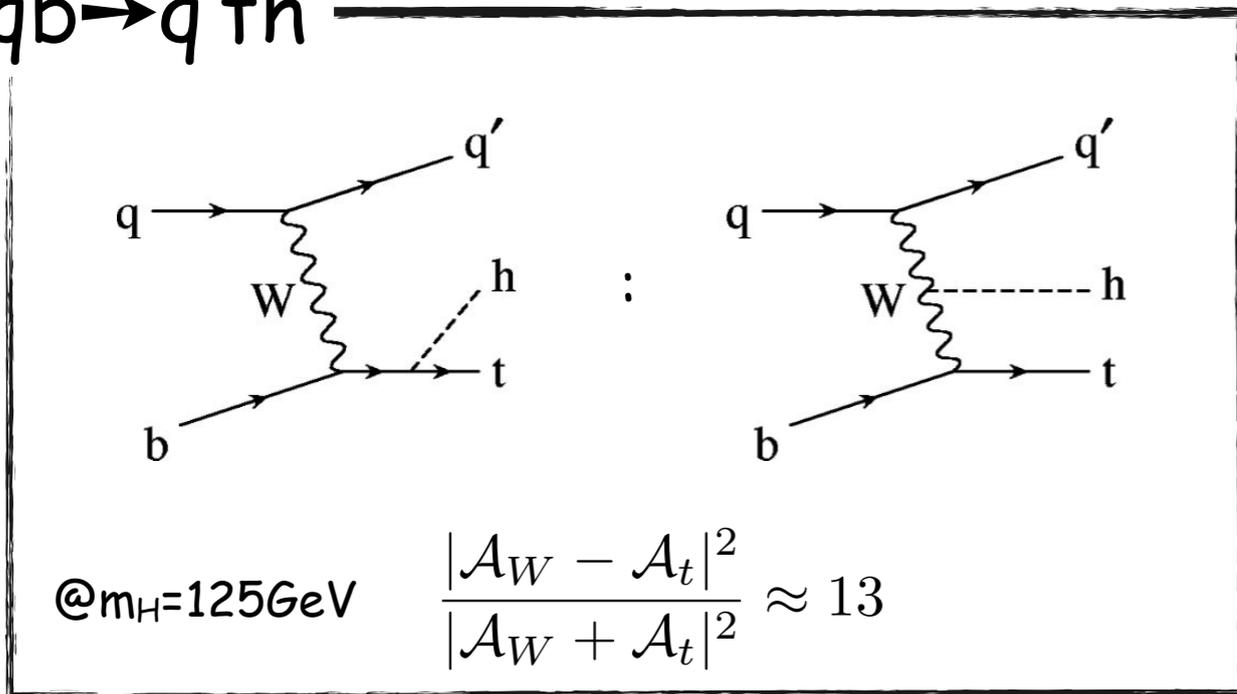
difference is physically relevant only in the presence of
strong interference with single $h\psi\psi$ coupling

$h \rightarrow \gamma\gamma$

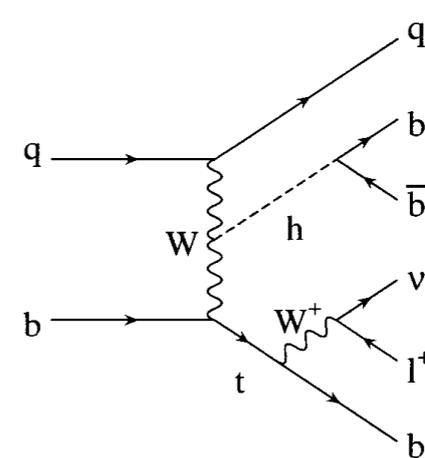


rare decay and one needs
some largish luminosity to
be sensitive to the sign of c

$qb \rightarrow q'th$



look at final state: $3b + 1 \text{ fwd jet} + l^\pm + p^T$.



Maltoni, Stelzer, Willenbrock '01

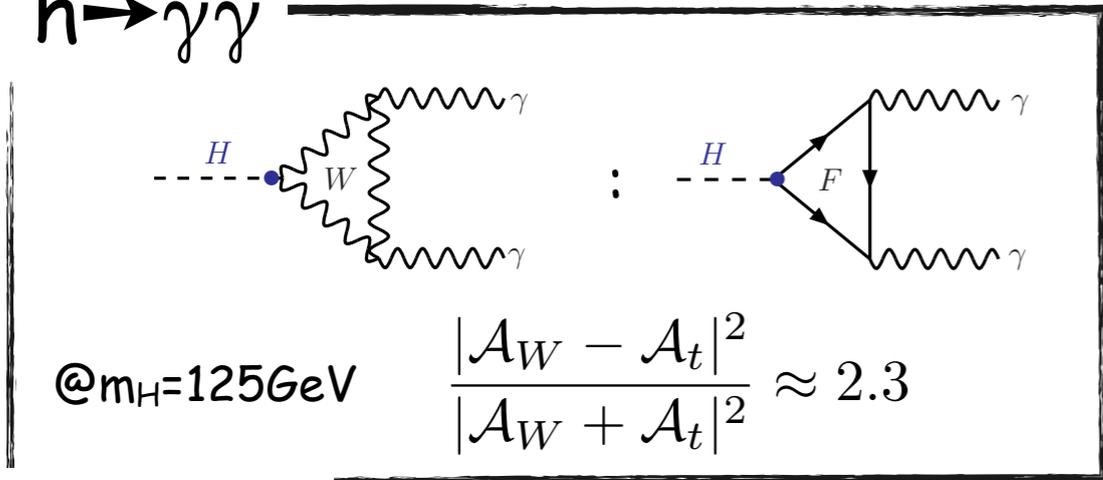
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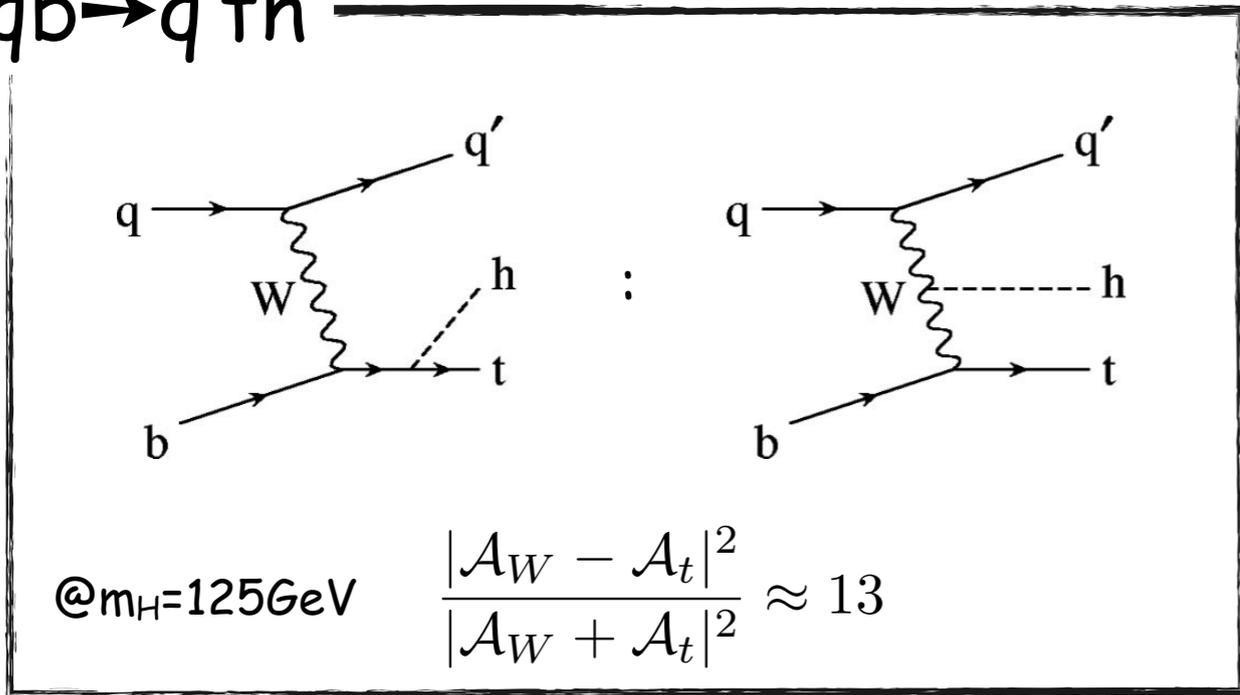
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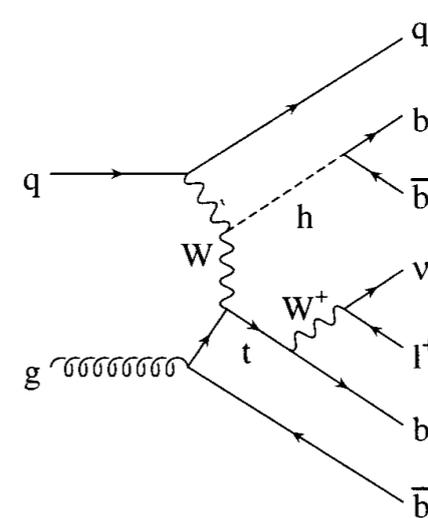


rare decay and one needs
some largish luminosity to
be sensitive to the sign of c

$qb \rightarrow q'th$



look at final state: $4b + 1 \text{ fwd jet} + l^\pm + p^T$.

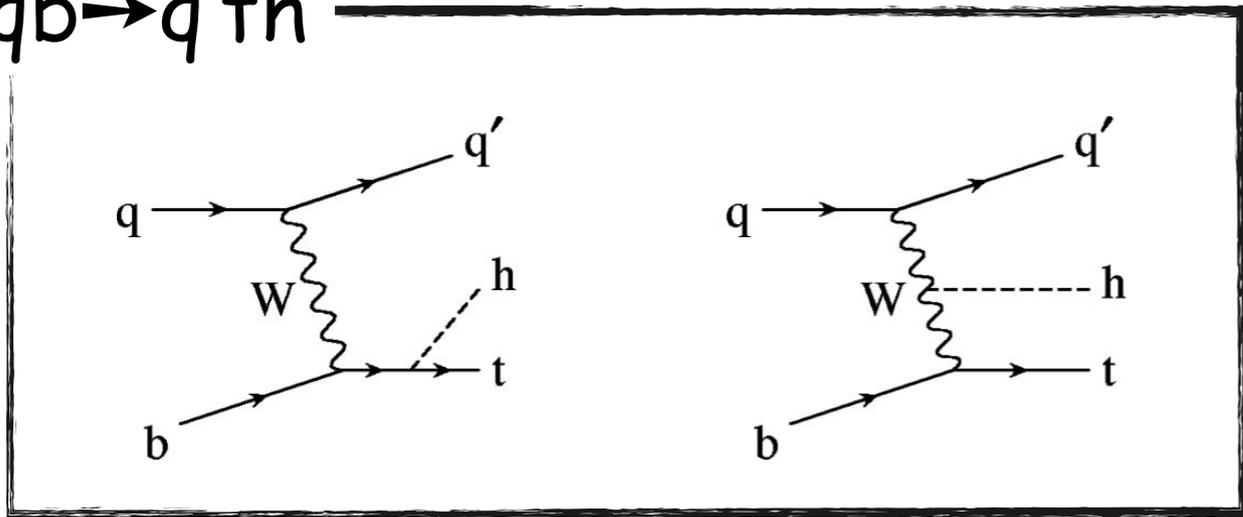


Maltoni, Stelzer, Willenbrock '01

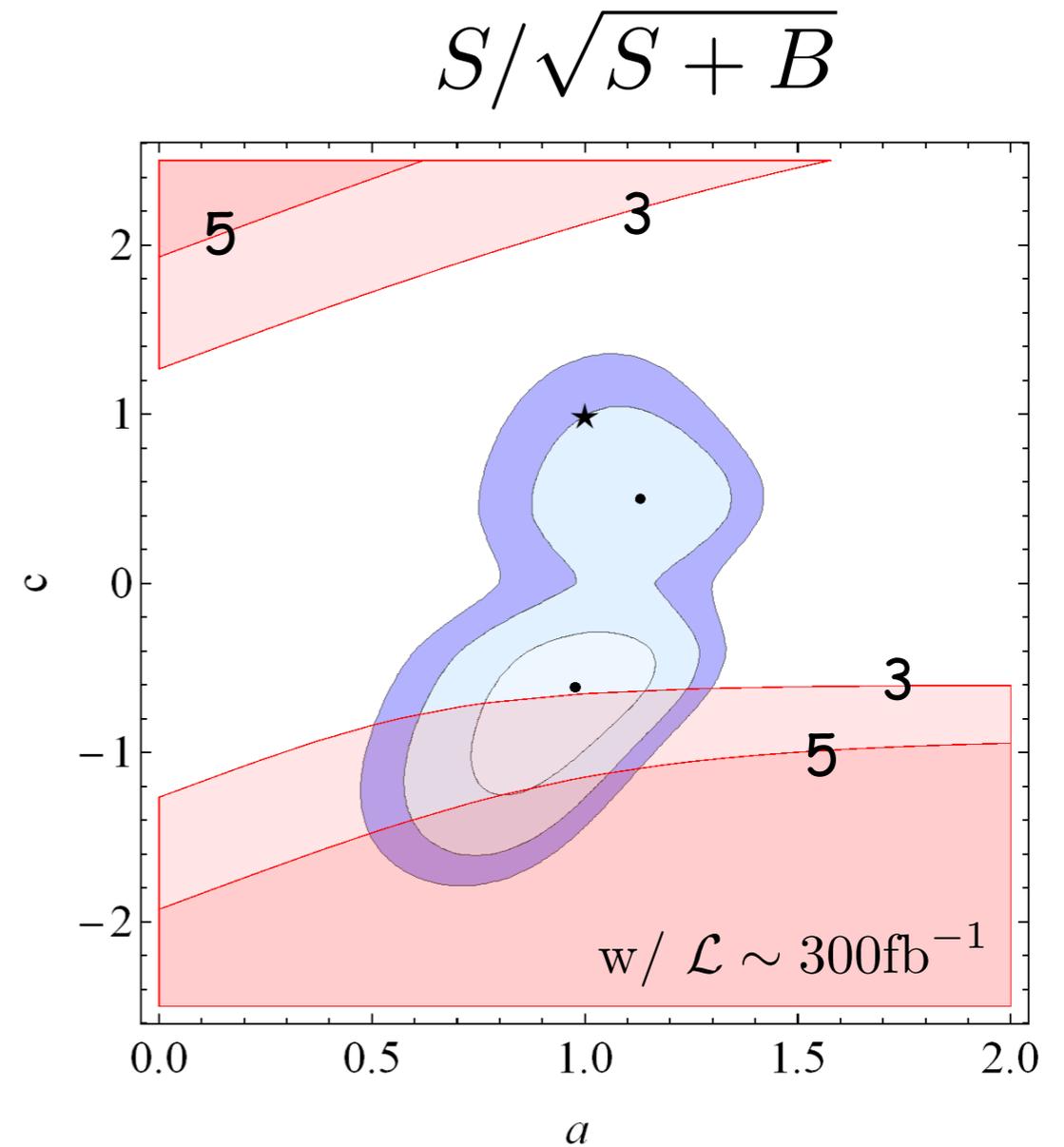
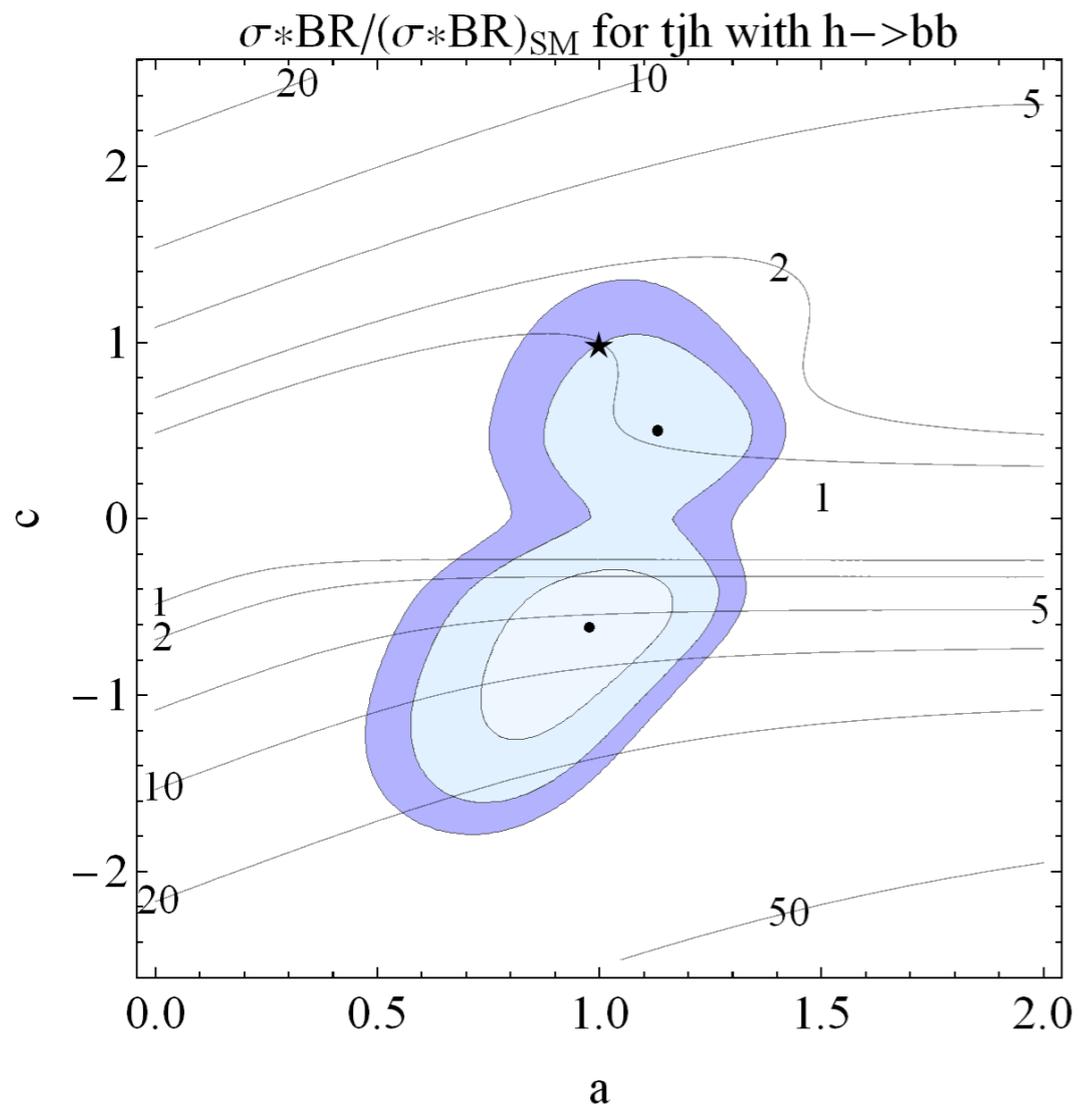
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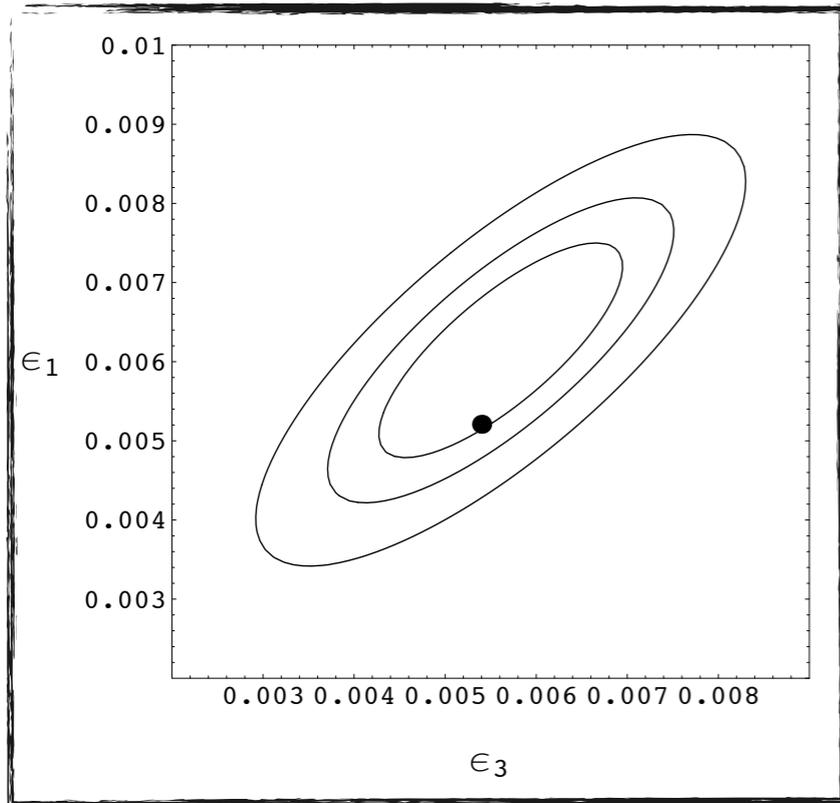
$qb \rightarrow q'th$



*the sign ambiguity will remain
with us for a long time!*



A tension between LHC and EW data?



EW fit strongly suggests custodial symmetry

$$\Sigma = e^{i\sigma^a \pi^a / v}$$

Goldstone of
 $SU(2)_L \times SU(2)_R / SU(2)_V$

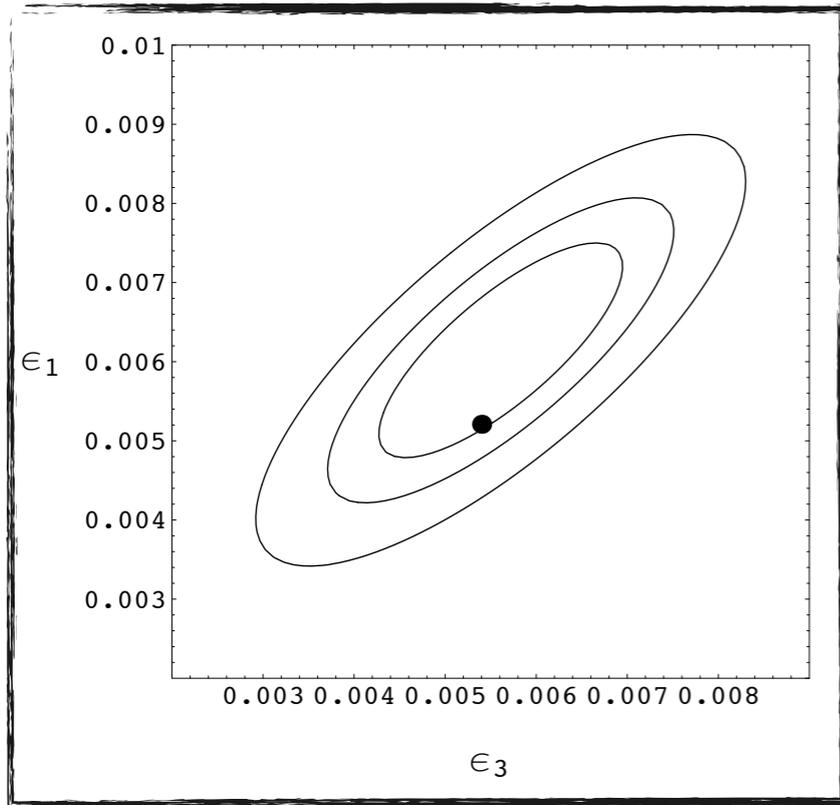
$$\frac{v^2}{4} \text{Tr} (D_\mu \Sigma^\dagger D^\mu \Sigma) \Rightarrow \rho = 1 \text{ ie } \epsilon_1 = \hat{T} = 0 \quad \checkmark$$

$$\text{also } \Rightarrow \mu_{ZZ} = \mu_{WW}$$

$$\left(\mu_i = \frac{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i)}{\sum_j \mathcal{A}_{ji} \sigma(j \rightarrow h) \times \text{Br}(h \rightarrow i) |_{\text{SM}}} \right)$$

$$\frac{v^2}{8} \text{Tr}^2 (\Sigma^\dagger D_\mu \Sigma \sigma^3) \Rightarrow \rho = 2 \text{ ie } \epsilon_1 = \hat{T} = 1 \quad \text{strongly disfavored}$$

A tension between LHC and EW data?



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$$\Sigma = e^{i\sigma^a \pi^a / v}$$

Goldstone of
 $SU(2)_L \times SU(2)_R / SU(2)_V$

$$\frac{v^2}{4} \text{Tr} (D_\mu \Sigma^\dagger D^\mu \Sigma) \Rightarrow \rho = 1 \text{ ie } \epsilon_1 = \hat{T} = 0 \quad \checkmark$$

$$\text{also } \Rightarrow \mu_{ZZ} = \mu_{WW} \quad \times$$

but

| Channel [Exp] | $\mu_{119.5} (\mu_{119.5}^L)$ | $\mu_{124} (\mu_{124}^L)$ | $\mu_{125} (\mu_{125}^L)$ |
|--|-------------------------------|---------------------------|---------------------------|
| $pp \rightarrow Z Z^* \rightarrow \ell^+ \ell^- \ell^+ \ell^-$ [ATLAS] | $-0.5^{+0.5??} (5.1)$ | $1.6^{+1.4}_{-0.8} (4.7)$ | $1.4^{+1.3}_{-0.8} (4.1)$ |
| $pp \rightarrow W W^* \rightarrow \ell^+ \nu \ell^- \bar{\nu}$ [ATLAS] | $0.0^{+1.2}_{-1.3} (2.4)$ | $0.1^{+0.7}_{-0.7} (1.6)$ | $0.1^{+0.7}_{-0.6} (1.4)$ |

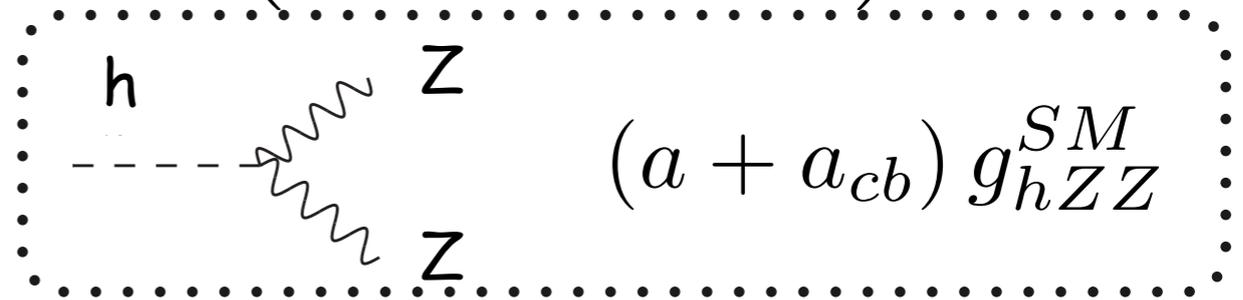
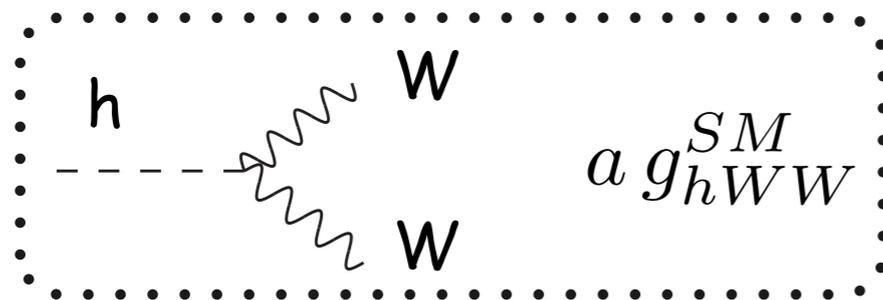
1. has LHC identified a violation of the custodial symmetry?
2. if yes, how to reconcile LHC data with EW data?

$$\frac{v^2}{8} \text{Tr}^2 (\Sigma^\dagger D_\mu \Sigma \sigma^3) \Rightarrow \rho = 2 \text{ ie } \epsilon_1 = \hat{T} = 1 \quad \text{strongly disfavored}$$

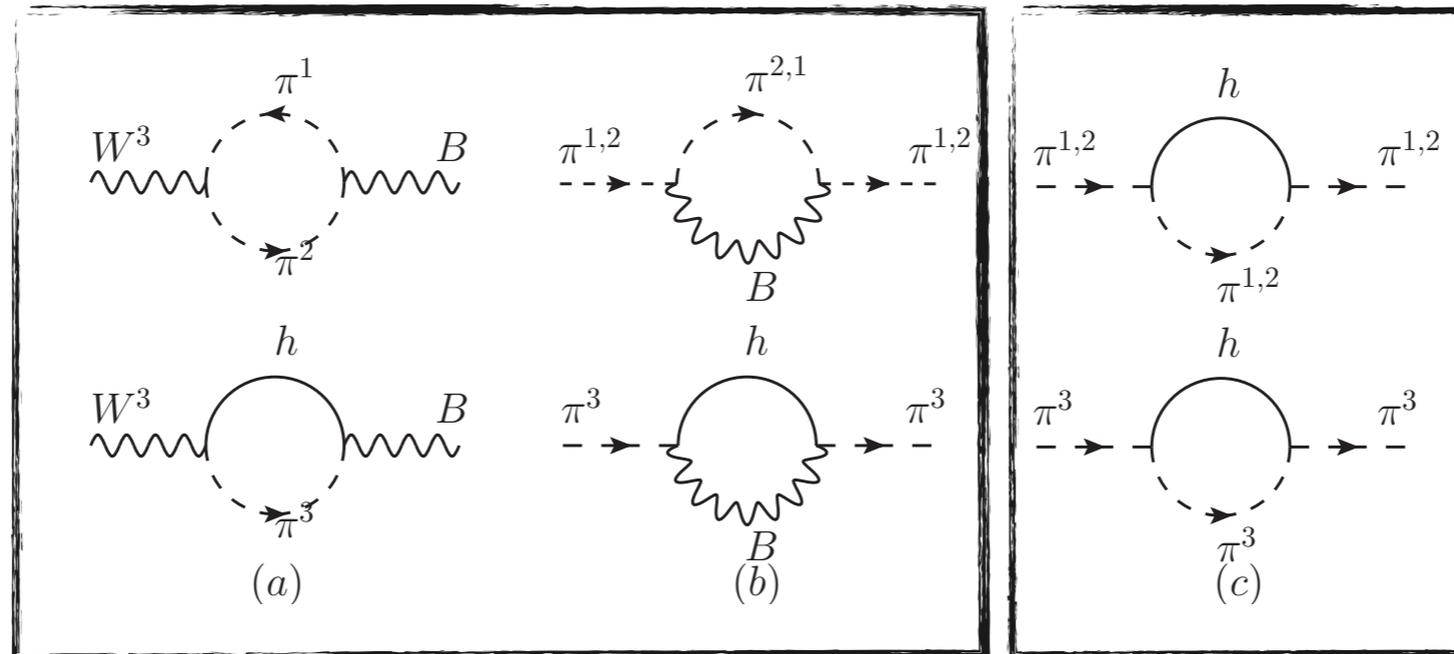
DisZphilia or how to live with custodial breaking

Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left(\text{Tr} [\Sigma^\dagger D_\mu \Sigma \sigma^3] \right)^2 \left(t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$



EWPT



Log(Λ) UV sensitivity

Λ^2 UV sensitivity

$$\Delta\epsilon_1 = -\frac{3}{16\pi} \frac{\alpha(m_Z)}{\cos^2 \theta_W} \left[1 - (a + a_{cb})^2 + \left(\frac{g}{g'} \right)^2 (a^2 - (a + a_{cb})^2) \right] \log \left(\frac{\Lambda^2}{m_h^2} \right)$$

$$\Delta\epsilon_3 = +\frac{1}{48\pi} \frac{\alpha(m_Z)}{\sin^2 \theta_W} \left[1 - (a + a_{cb})^2 \right] \log \left(\frac{\Lambda^2}{m_h^2} \right)$$

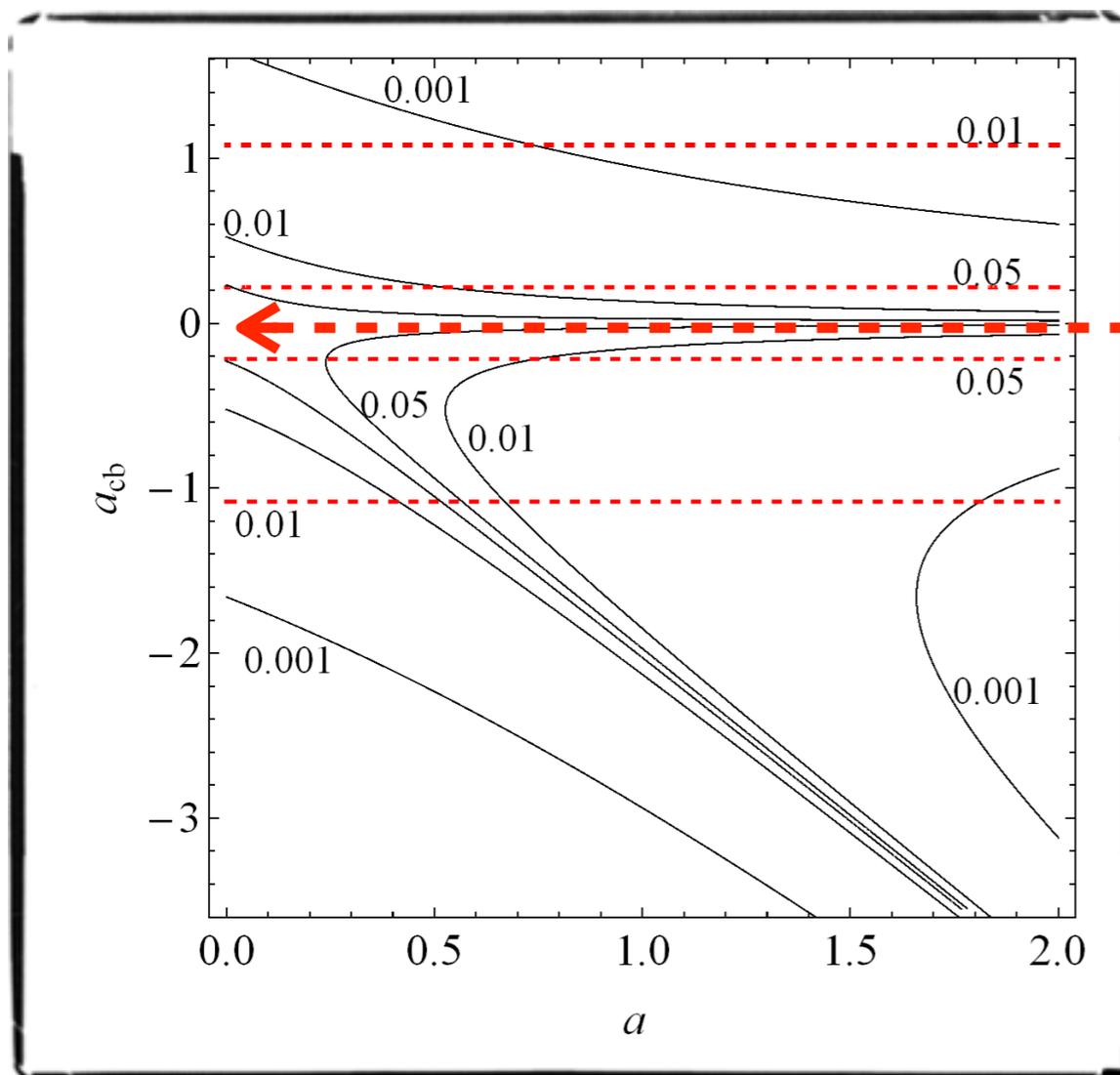
$$\Delta\epsilon_1 = \left((a + a_{cb})^2 - a^2 \right) \frac{\Lambda^2}{16\pi^2 v^2}$$

DisZphilia or how to live with custodial breaking

Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left(\text{Tr} \left[\Sigma^\dagger D_\mu \Sigma \sigma^3 \right] \right)^2 \left(t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$

EWPT highly model-dependent
tuning between tree-level and loop contributions?
new light states?



Amount of fine-tuning in EWPT

← ——— → custodial invariance

Λ^2 UV sensitivity

ie

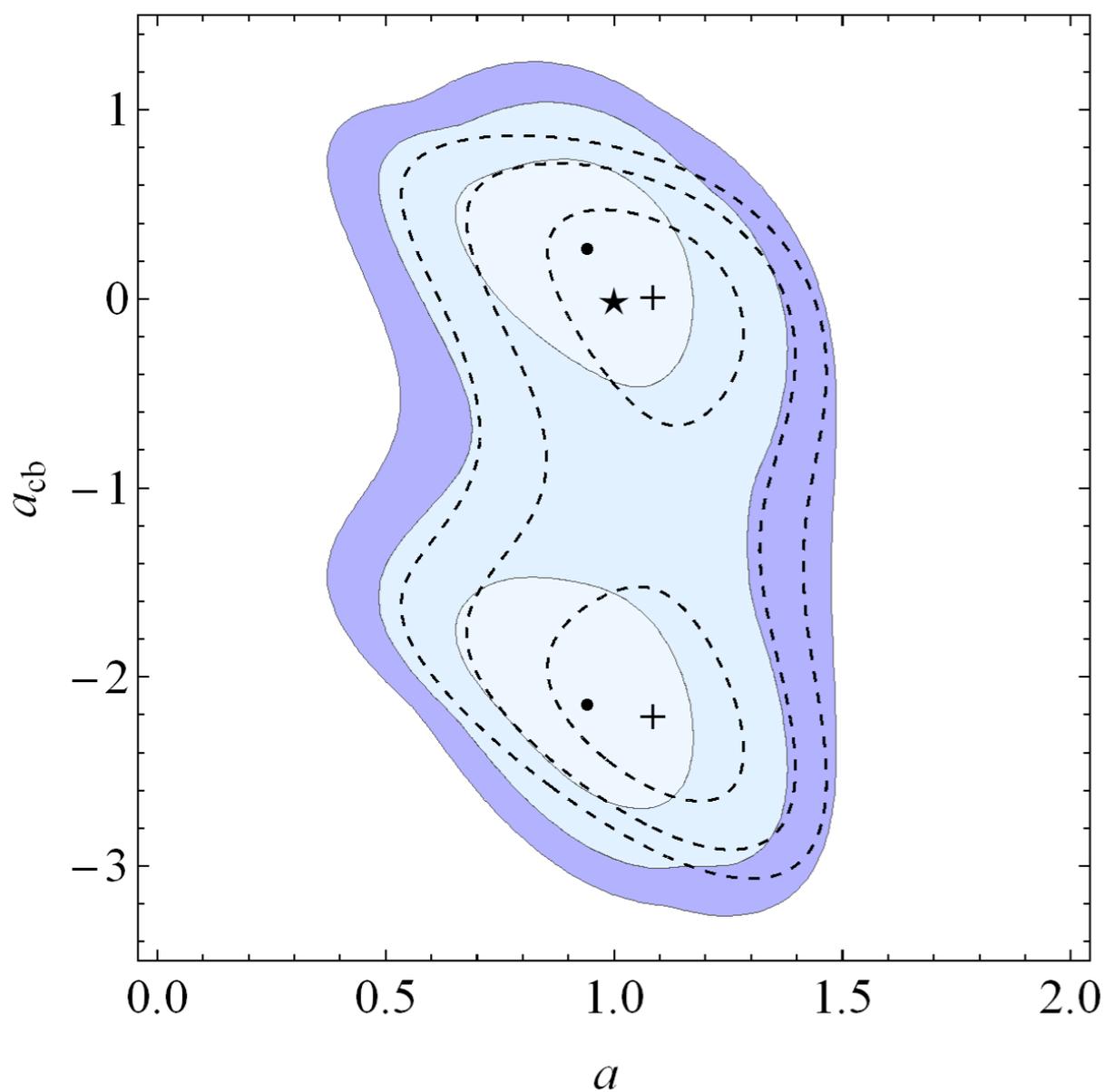
could be as bad as the hierarchy problem

DisZphilia or how to live with custodial breaking

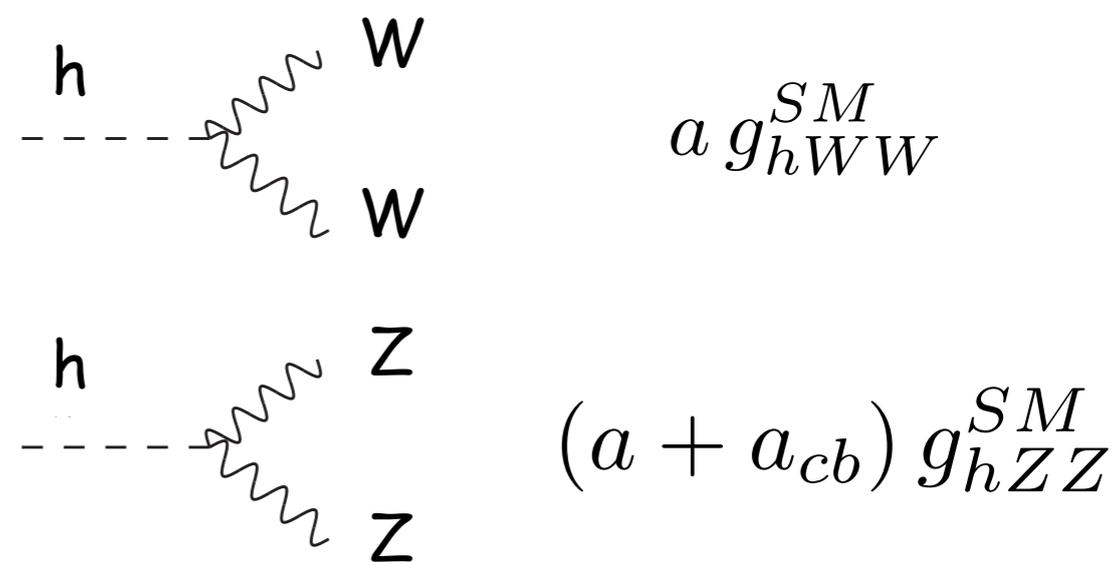
Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left(\text{Tr} [\Sigma^\dagger D_\mu \Sigma \sigma^3] \right)^2 \left(t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$

Fit to LHC data



----- $c = 1$
 ————— marginalization over c



LHC data are symmetric under

$$(a + a_{cb}) \leftrightarrow -(a + a_{cb})$$

i.e.

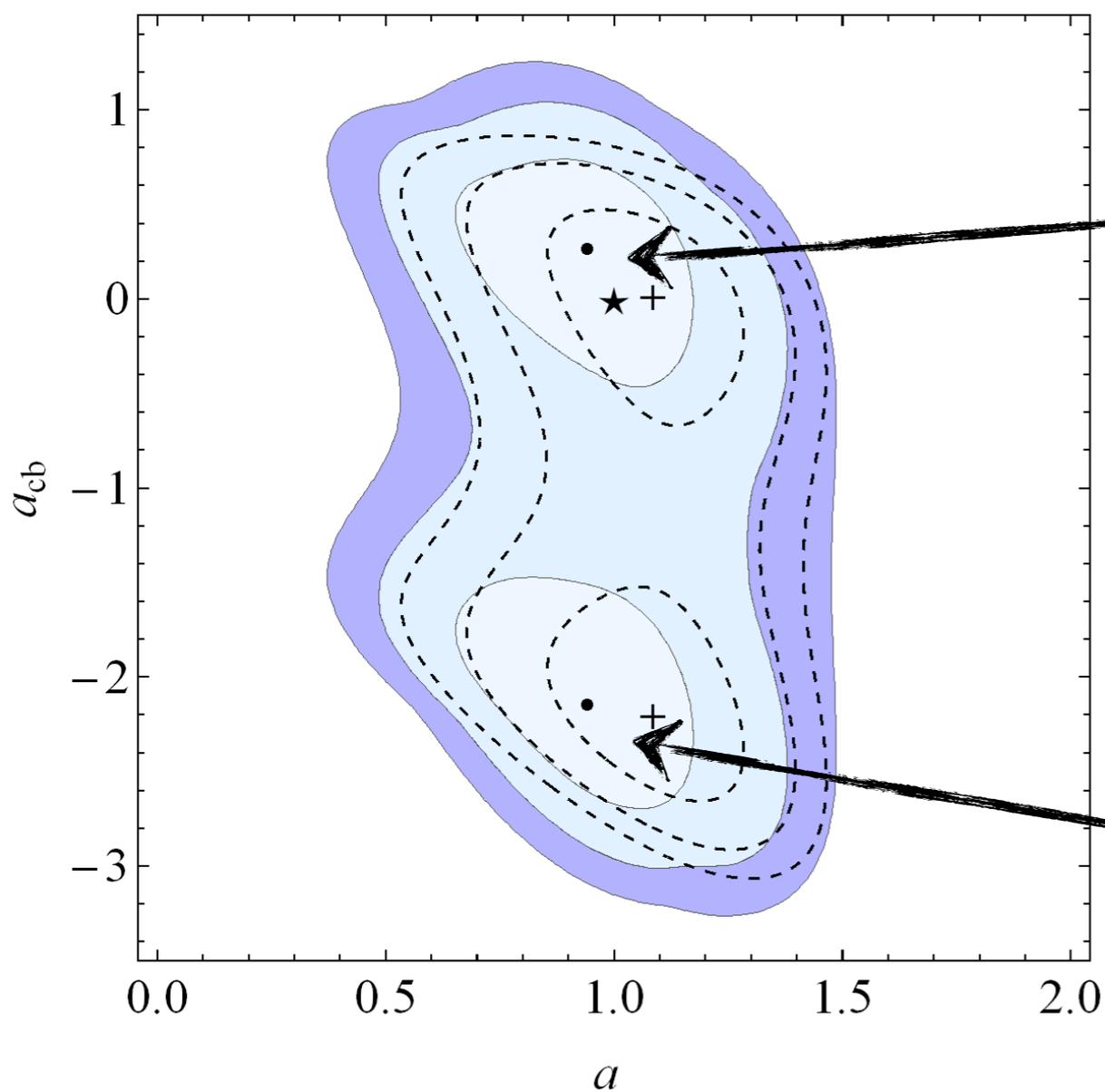
$$(a, a_{cb}) \leftrightarrow (a, -(2a + a_{cb}))$$

DisZphilia or how to live with custodial breaking

Farina, Grojean, Salvioni '12

$$\mathcal{L}_{cb} = -\frac{v^2}{8} \left(\text{Tr} [\Sigma^\dagger D_\mu \Sigma \sigma^3] \right)^2 \left(t_{cb} + 2a_{cb} \frac{h}{v} + \dots \right)$$

Fit to LHC data



----- c = 1
 ————— marginalization over c

$$2 \frac{h}{v} \left(m_W^2 W_\mu^+ W_\mu^- + \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right)$$

the two solutions can only be distinguished in the presence of interference with a single hZZ vertex

$$2 \frac{h}{v} \left(m_W^2 W_\mu^+ W_\mu^- - \frac{1}{2} m_Z^2 Z_\mu Z_\mu \right)$$

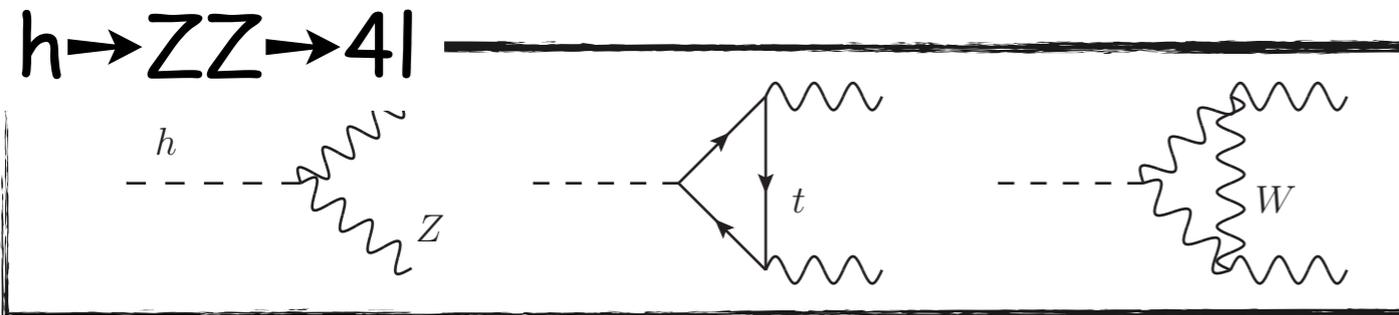
↑
 "disZphilia"

Zphilia or DisZphilia?

difficult!

Farina, Grojean, Salvioni '12

difference is physically relevant only in the presence of interference with single hZZ coupling

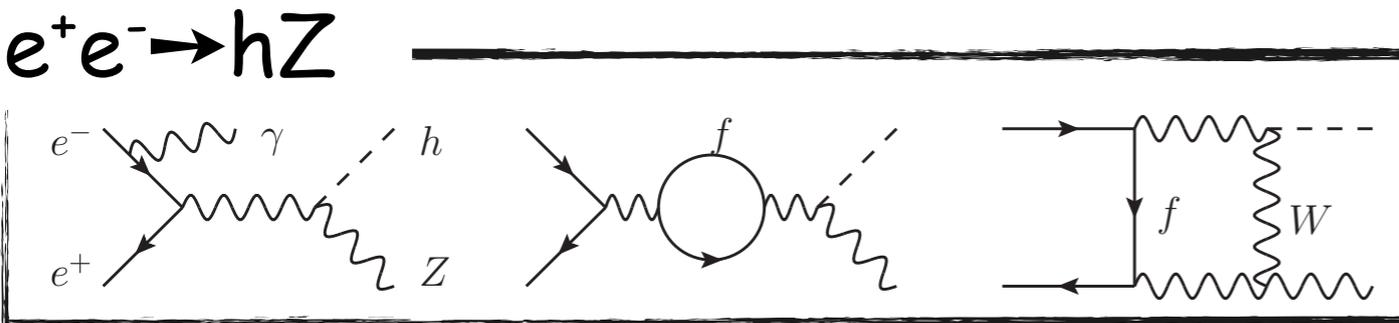


TH prediction

$$\Delta = \left| \frac{\Gamma_Z^+ - \Gamma_Z^-}{\Gamma_Z^+ + \Gamma_Z^-} \right| = \delta \approx 1\%$$

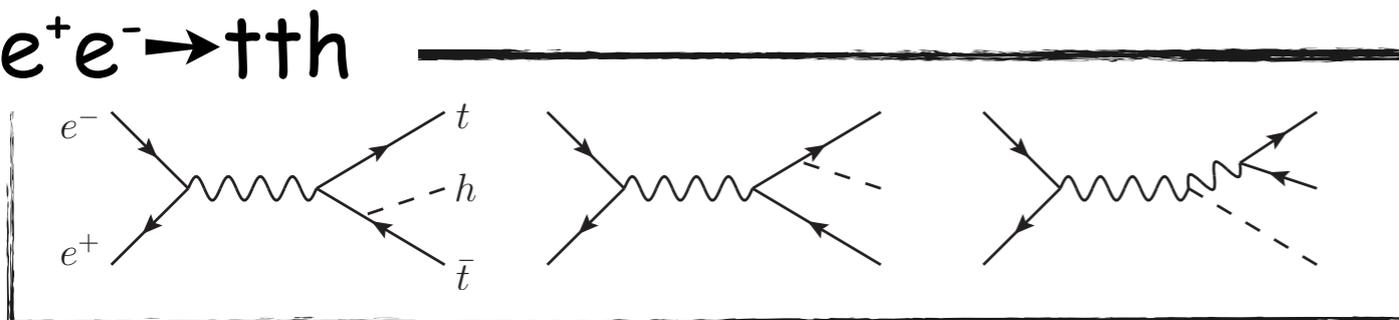
ILC ($\sqrt{s}=800\text{GeV}$ and 1ab^{-1})

$\approx 1\%$



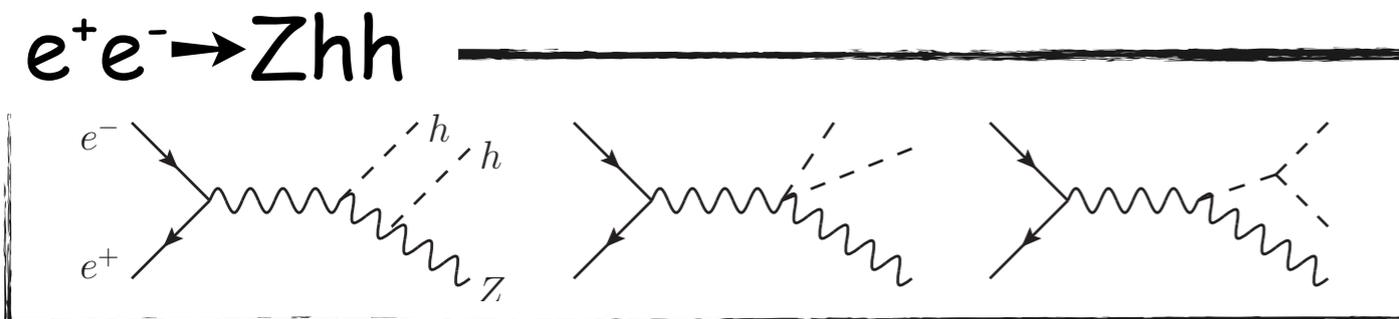
$$\Delta = \left| \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \right| \approx 15\%$$

$\approx 5\%$



$$\Delta = \left| \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \right| \lesssim 4\%$$

$\approx 10\%$



$$\Delta = \left| \frac{\sigma_+ - \sigma_-}{\sigma_+ + \sigma_-} \right| \approx 50\%$$

$\approx 10\%$



Signs of New Particles?

Espinosa, Grojean, Muhlleitner, Trott '12

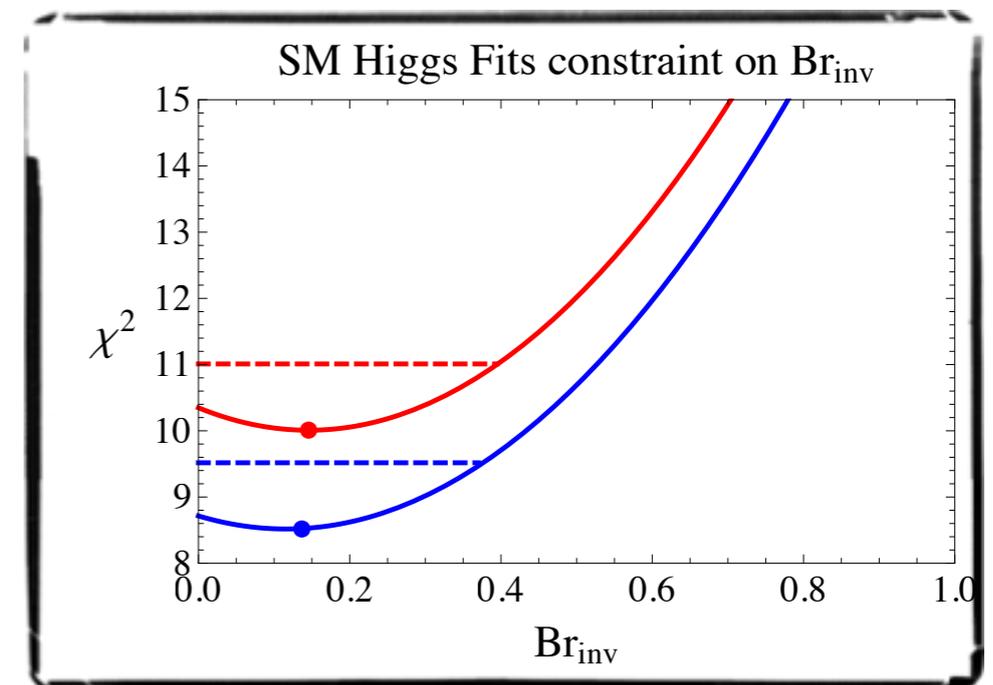
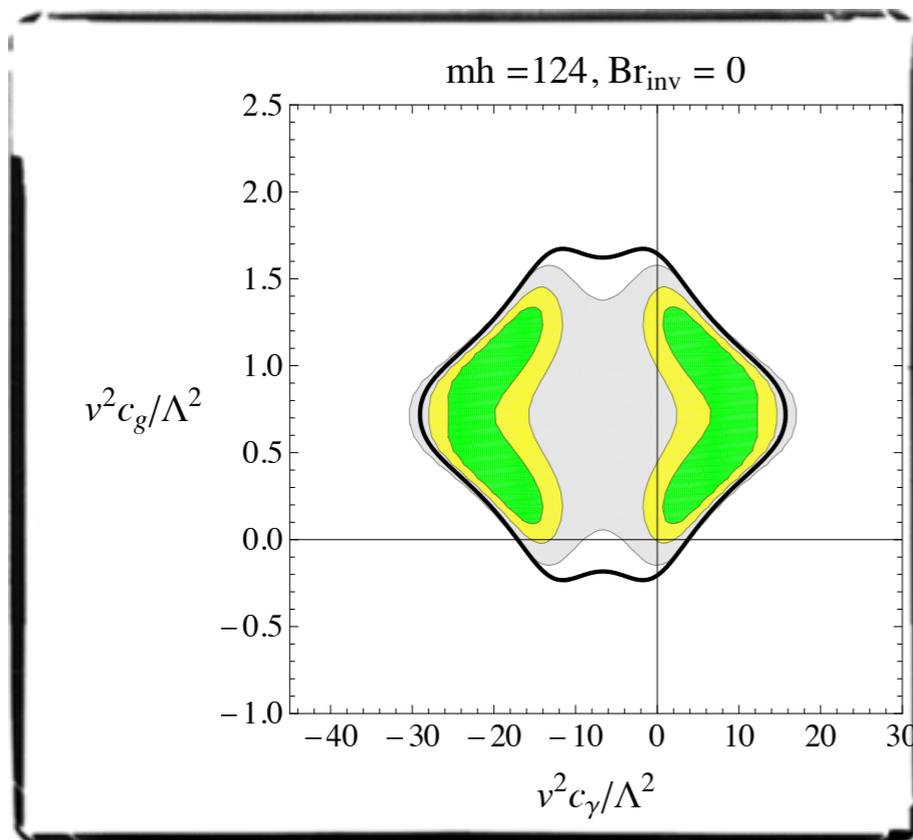
The Higgs can couple to new particles

charged under
SM gauge group

neutral under
SM gauge group

$$\mathcal{L} = -\frac{\tilde{c}_\gamma e^2}{32\pi^2 \Lambda^2} H^\dagger H F_{\mu\nu} F^{\mu\nu} - \frac{\tilde{c}_g g_s^2}{32\pi^2 \Lambda^2} H^\dagger H G_{\mu\nu}^a G^{a\mu\nu}$$

$$\text{Br}(h \rightarrow f) \equiv \frac{\Gamma(h \rightarrow f)}{\Gamma_{\text{SM}} + \Gamma_{\text{inv}}} = (1 - \text{Br}_{\text{inv}}) \times \text{Br}_{\text{SM}}(h \rightarrow f)$$

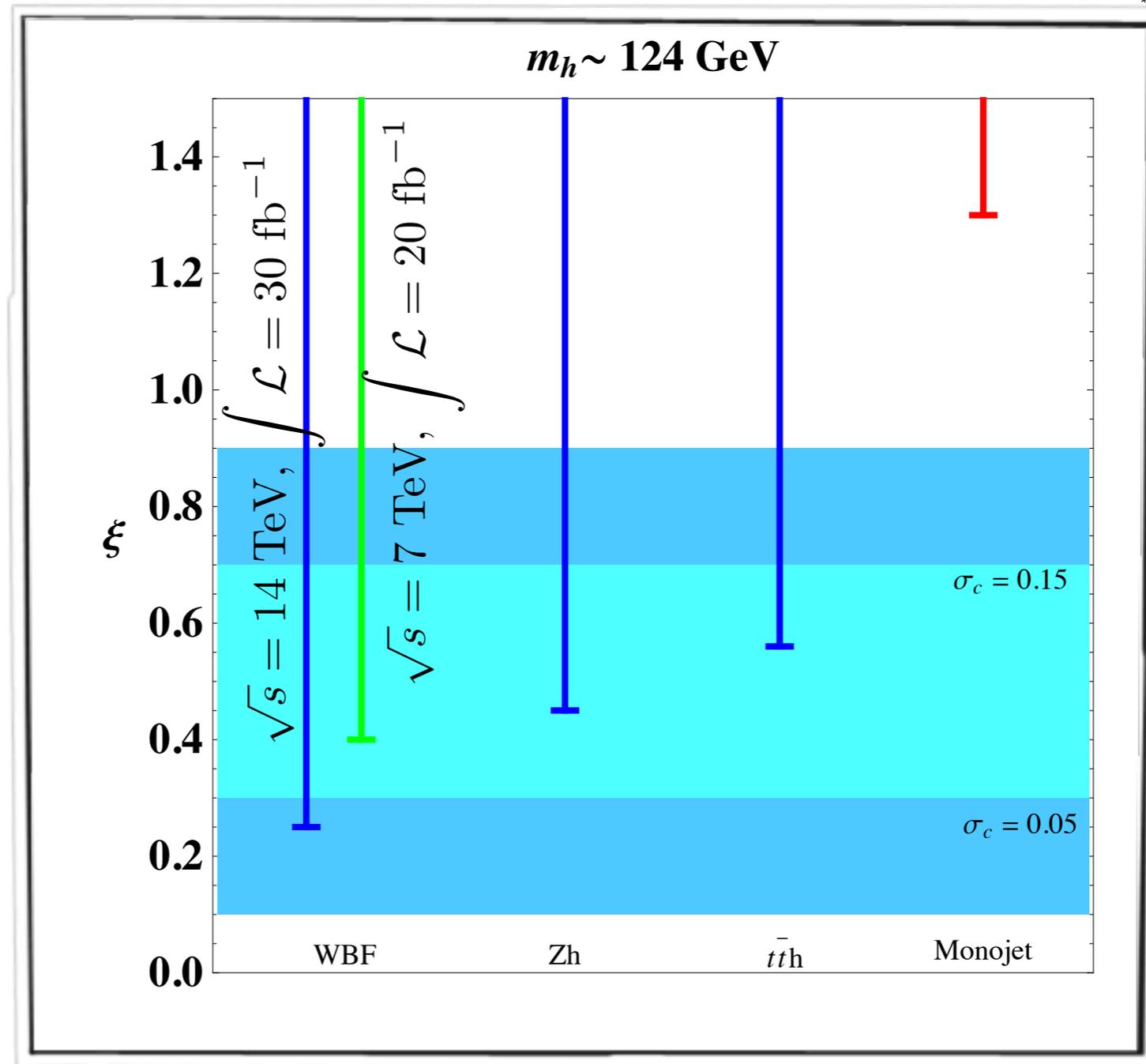


Search for Invisible Decays with Visible Channels

Espinosa, Grojean, Muhlleitner, Trott '12

direct (vertical) vs indirect (horizontal) searches

Values of
 $\xi = \sigma / \sigma_{SM} BR(h \rightarrow inv)$
 for which a 95%CL limit can be imposed



The Question of the next Decade(s)

What is really this Higgs boson that might have been discovered at $\sim 125\text{GeV}$?

"Higgs = emergency tire of the SM"

Altarelli @ Blois'10



[picture courtesy to Andreas Weiler]