

Implications of Higgs Searches and Discovery

Michele Papucci
LBNL

SLAC Summer Institute 2012

Aug 1st, 2012

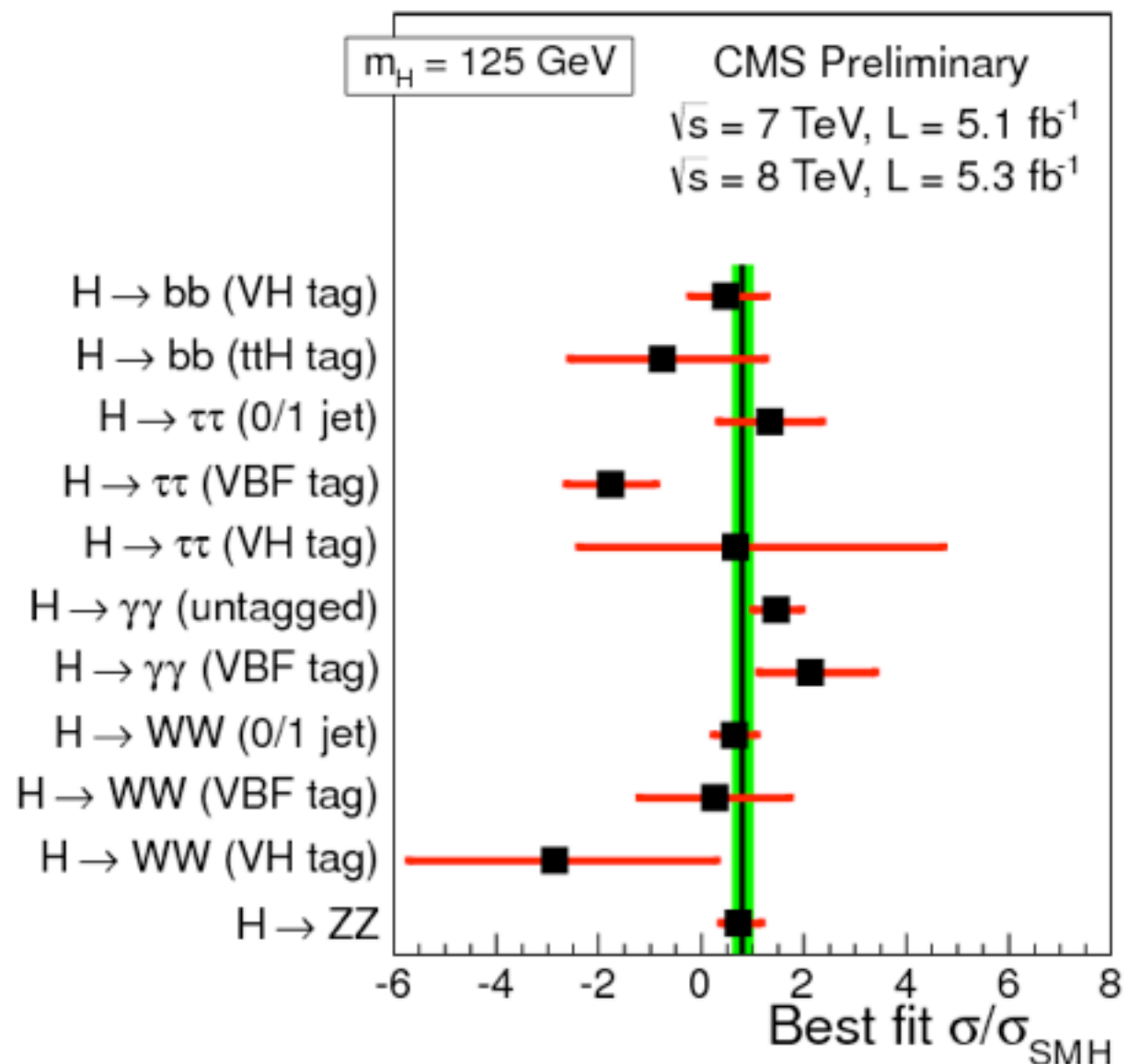
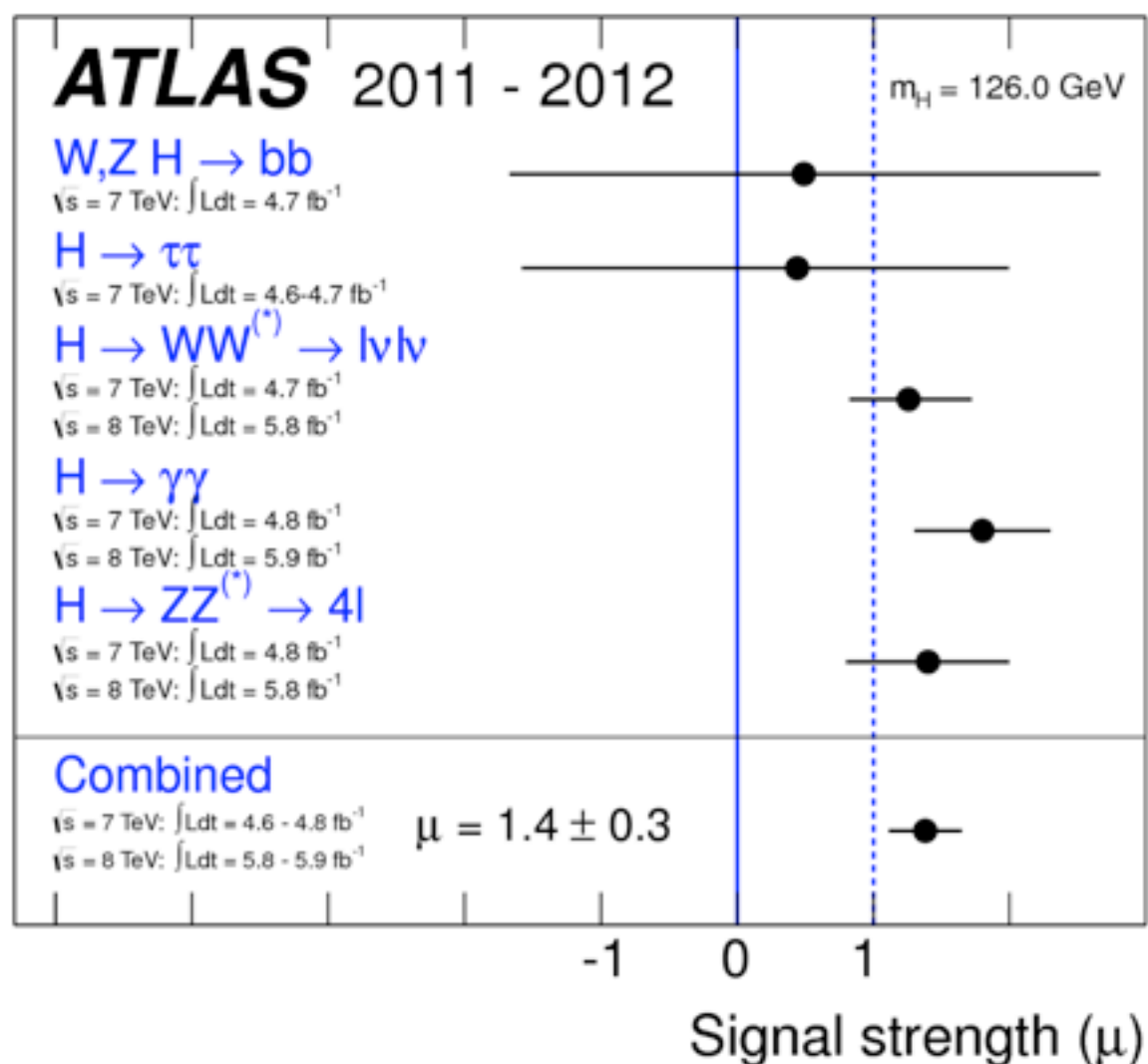
Implications of ~~Higgs~~ *a new particle?* Searches and Discovery

Michele Papucci
LBNL

SLAC Summer Institute 2012

Aug 1st, 2012

- On July 4th, a new particle has been discovered by ATLAS and CMS around 125-126 GeV
- It's a boson, and can decay to $\gamma\gamma$, ZZ^* , WW^*
- Couplings to fermions ($bb, \tau\tau$) are still not very constrained



Outline

- Is it a Higgs or an impostor? Does it participate in EWSB?
- If it is an Higgs, is it fundamental or composite?
- Implications for specific models?
- ...

Outline

- Is it a Higgs or an impostor? Does it participate in EWSB?
- If it is an Higgs, is it fundamental or composite?
- Implications for specific models?
- ...

This kind of questions will occupy us for a while...

Outline

- Is it a Higgs or an impostor? Does it participate in EWSB?
- If it is an Higgs, is it fundamental or composite?
- Implications for specific models?
- ...

Today

This kind of questions will occupy us for a while...

Two ways to proceed

Two ways to proceed

- Top - down:

Two ways to proceed

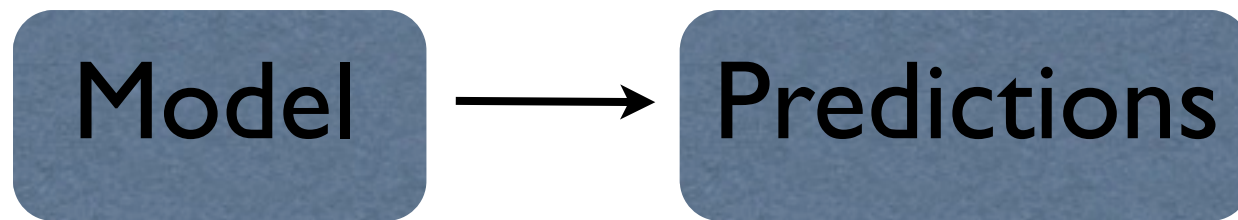
- Top - down:



Model

Two ways to proceed

- Top - down:



Two ways to proceed

- Top - down:



Two ways to proceed

- Top - down:



ex: SUSY (MSSM in particular) → M. Carena's lecture
(+ other examples tomorrow)

Two ways to proceed

- Top - down:



ex: SUSY (MSSM in particular) \rightarrow M. Carena's lecture
(+ other examples tomorrow)

- Bottom - up: start from effective Lagrangian for EWSB
(see H. Haber's lecture)

Two ways to proceed

- Top - down:



ex: SUSY (MSSM in particular) → M. Carena's lecture
(+ other examples tomorrow)

- Bottom - up: start from effective Lagrangian for EWSB (see H. Haber's lecture)

Look at what data says about the 125GeV particle being an impostor or being composite, or ...

Two ways to proceed

- Top - down:



ex: SUSY (MSSM in particular) → M. Carena's lecture
(+ other examples tomorrow)

- Bottom - up: start from effective Lagrangian for EWSB (see H. Haber's lecture)

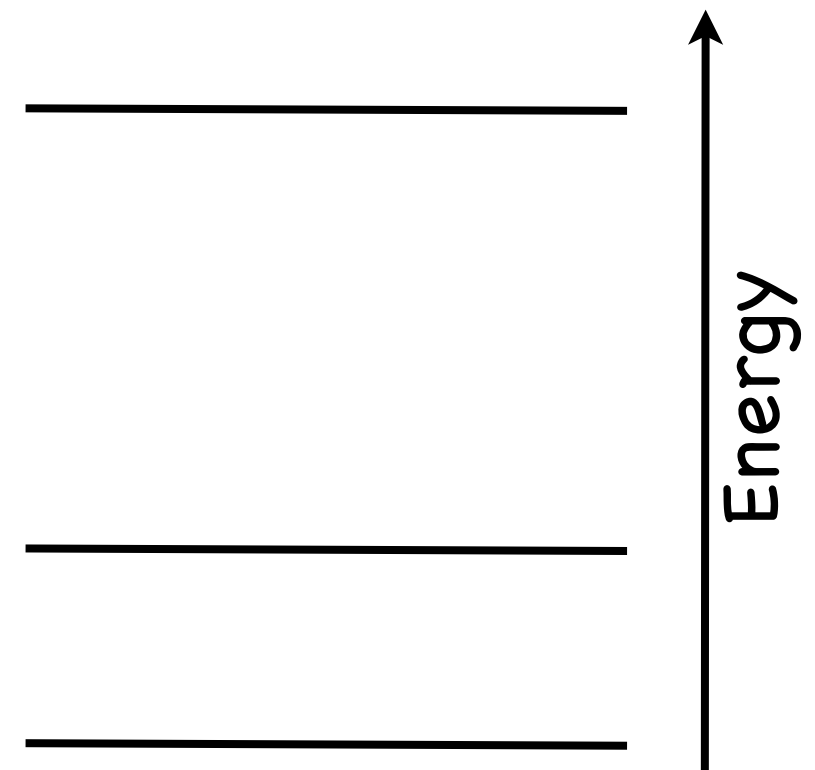
Look at what data says about the 125GeV particle being an impostor or being composite, or ...

maybe too early to play these games, but we've been waiting for so long...

Choices for EWSB

Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

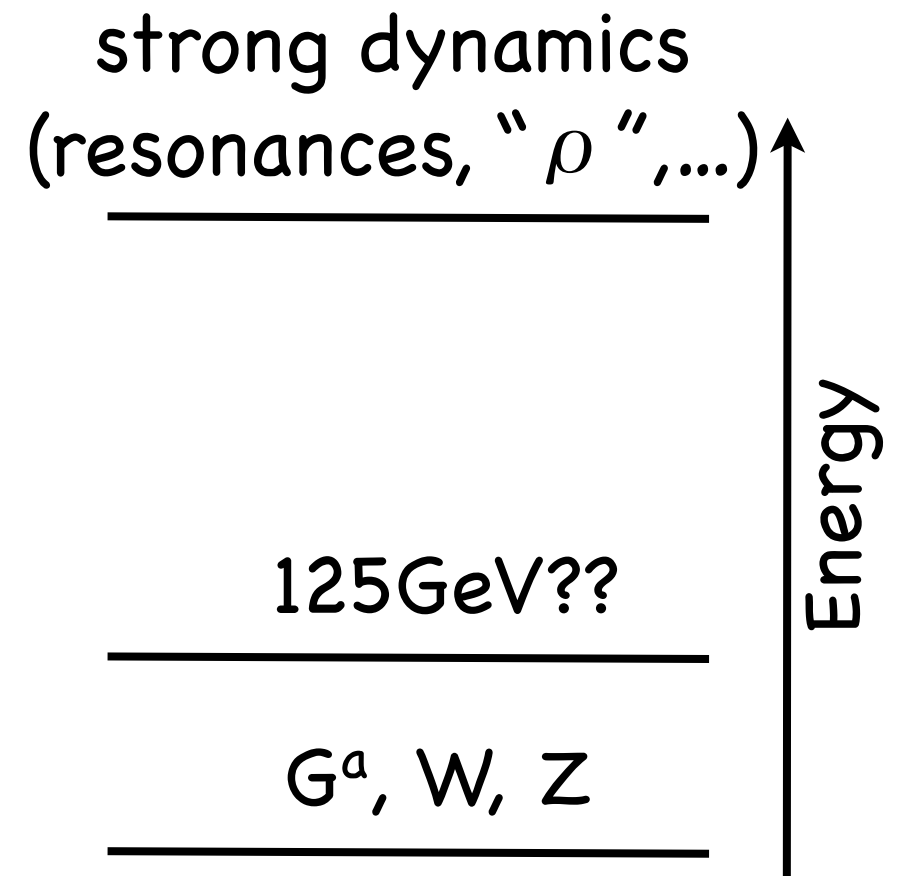


Choices for EWSB

Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

I. Nothing (“Higgsless”)



Choices for EWSB

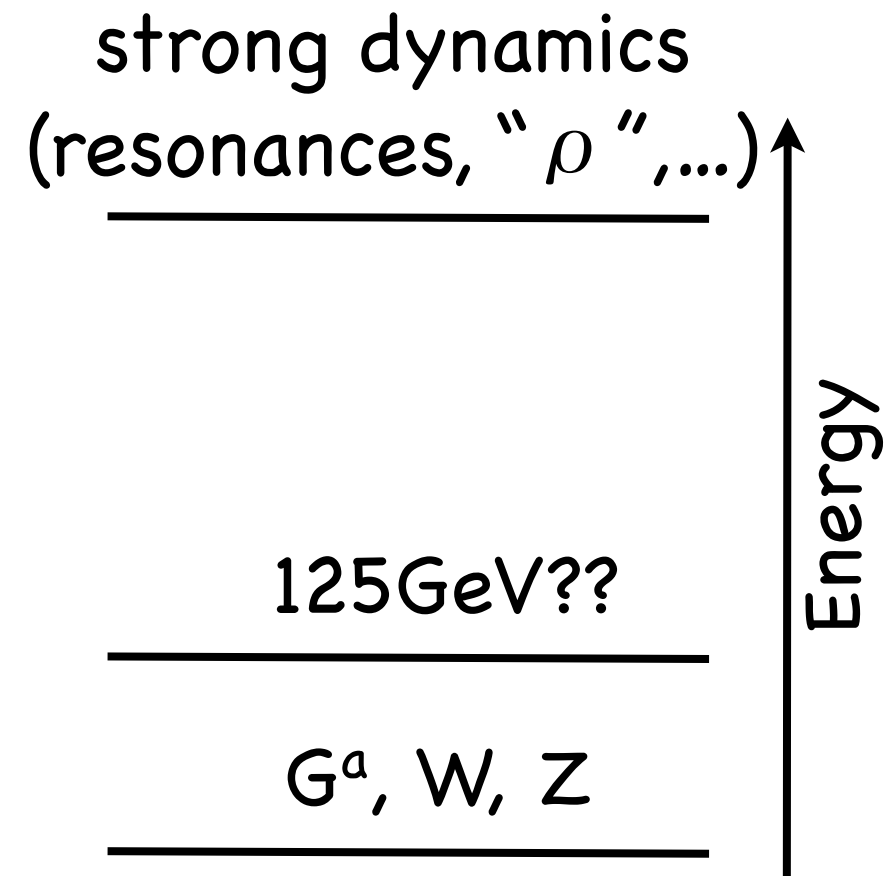
Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

I. Nothing (“Higgsless”)

125 GeV scalar is an “impostor”
(more on it later)

Strong dynamics takes care of UV problems of the Goldstone scattering amplitudes

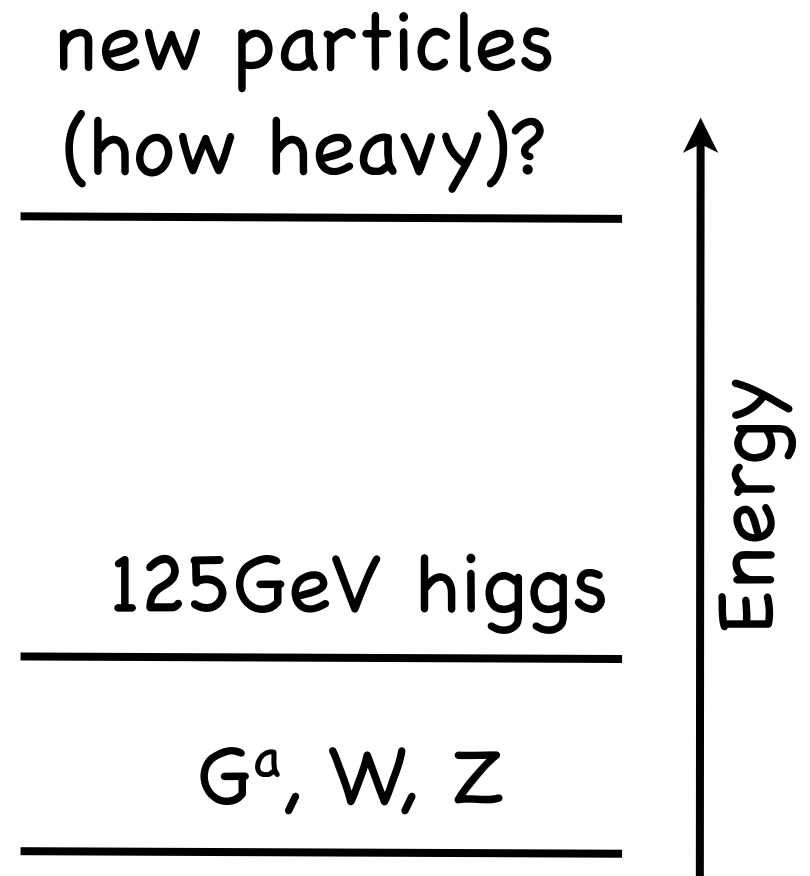


Choices for EWSB

Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

1. Nothing (“Higgsless”)
2. $SU(2)_L \times U(1)_Y$ linearly realized and weakly coupled: (a bunch of) fundamental Higgs(es)



Choices for EWSB

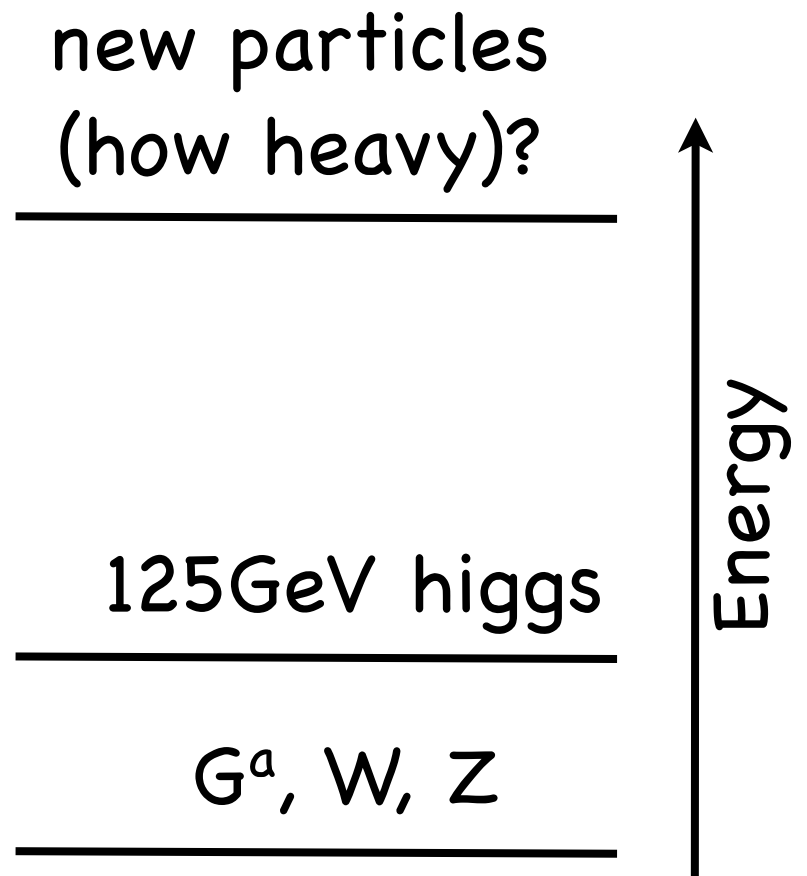
Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

1. Nothing (“Higgsless”)
2. $SU(2)_L \times U(1)_Y$ linearly realized and weakly coupled: (a bunch of) fundamental Higgs(es)

Higgs takes care of UV problems of the Goldstone scattering amplitudes

New particles may be present to fix the hierarchy problem (naturalness)



Choices for EWSB

Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

-
- The diagram illustrates the energy spectrum for three different scenarios of electroweak symmetry breaking (EWSB). A vertical arrow on the right is labeled "Energy". Three horizontal lines represent different energy levels:
- 1. Nothing ("Higgsless")**: This scenario corresponds to the highest energy level, labeled "strong dynamics (resonances, ' ρ ', ...)".
 - 2. $SU(2)_L \times U(1)_Y$ linearly realized and weakly coupled: (a bunch of) fundamental Higgs(es)**: This scenario corresponds to the middle energy level, labeled "125 GeV composite Higgs (like a pion)".
 - 3. $SU(2)_L \times U(1)_Y$ linearly realized + strong dynamics ("composite Higgs", etc...)**: This scenario corresponds to the lowest energy level, labeled " G^a, W, Z ".

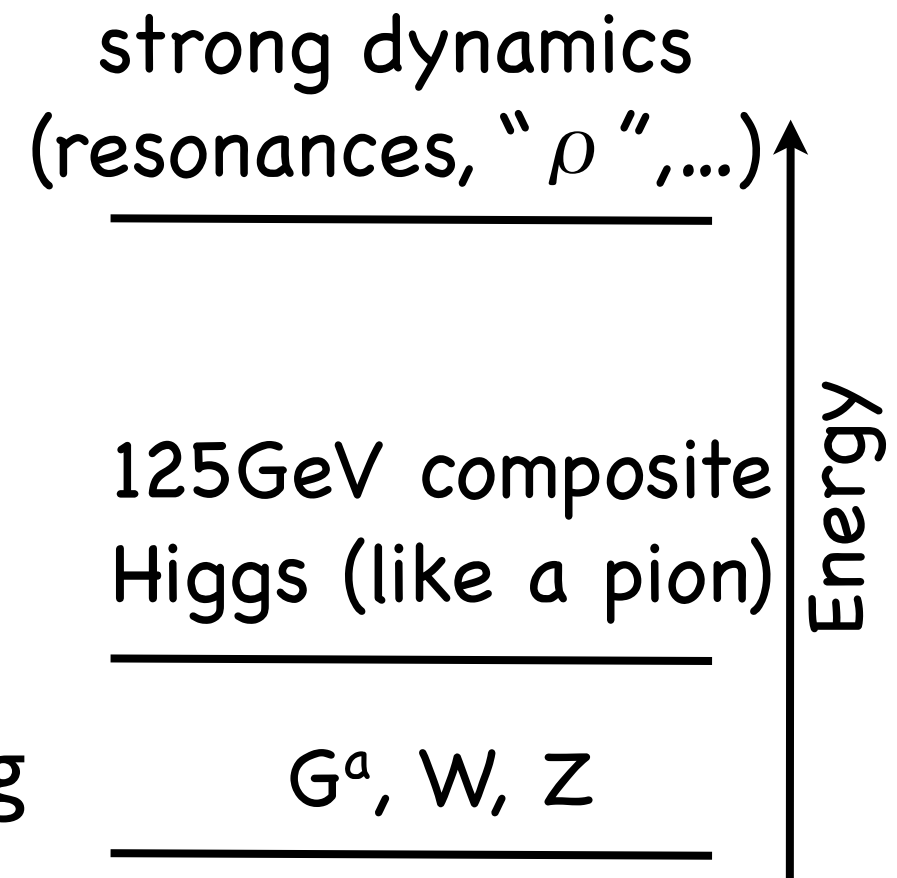
Choices for EWSB

Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

1. Nothing (“Higgsless”)
2. $SU(2)_L \times U(1)_Y$ linearly realized and weakly coupled: (a bunch of) fundamental Higgs(es)
3. $SU(2)_L \times U(1)_Y$ linearly realized + strong dynamics (“composite Higgs”, etc...)

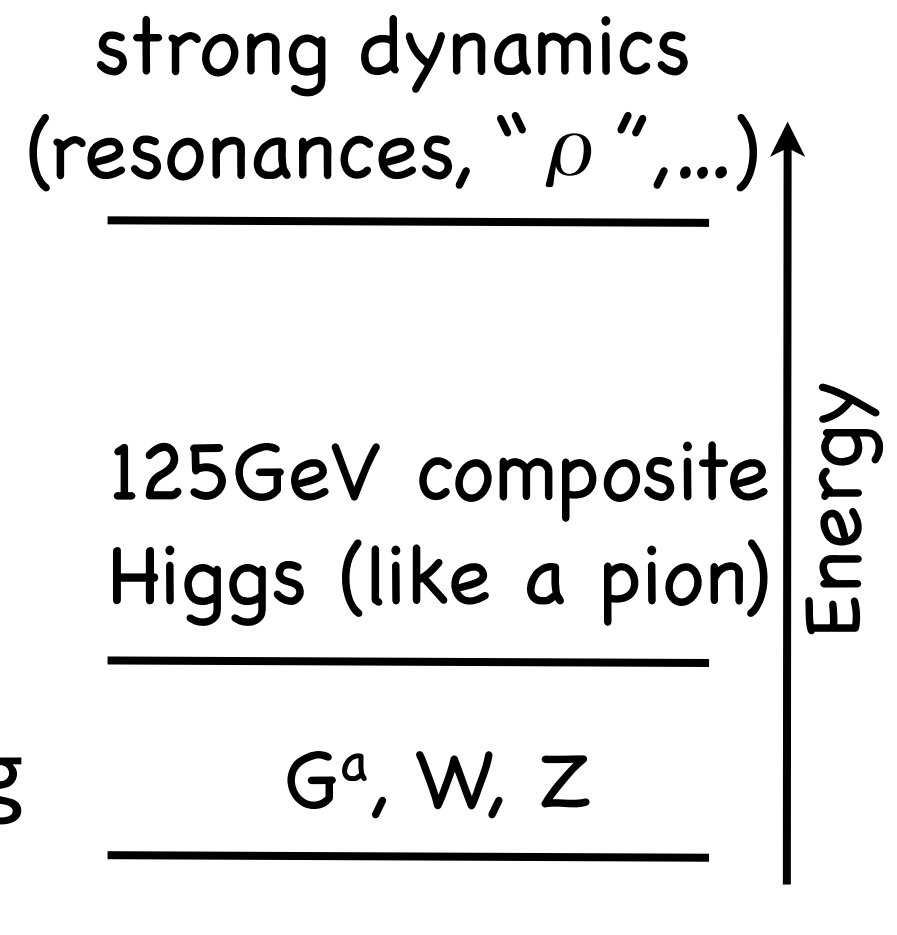
Higgs only partially takes care of UV problems of the Goldstone scattering amplitudes, rest is done by strong dynamics



Choices for EWSB

Non-linear realization of $SU(2)_L \times U(1)_Y$ (describes Goldstone bosons G^a , breaks down around 1 TeV)

+ what??

- 
- The diagram illustrates the energy spectrum for three different EWSB scenarios. A vertical arrow on the right is labeled 'Energy'. Three horizontal lines represent different energy levels:
- 1. Nothing ("Higgsless")**: This scenario corresponds to the highest energy level, labeled "strong dynamics (resonances, " ρ ", ...)".
 - 2. $SU(2)_L \times U(1)_Y$ linearly realized and weakly coupled: (a bunch of) fundamental Higgs(es)**: This scenario corresponds to the middle energy level, labeled "125 GeV composite Higgs (like a pion)".
 - 3. $SU(2)_L \times U(1)_Y$ linearly realized + strong dynamics ("composite Higgs", etc...)**: This scenario corresponds to the lowest energy level, labeled " G^a, W, Z ".

Implication of current data on 1 vs. 2 vs. 3?

Bottom-up approach

Let's count parameters to fit to Higgs data...

Effective (chiral) lagrangian for the Higgs

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light
- keep custodial symmetry (LEP data)

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light
- keep custodial symmetry (LEP data)
- no FCNC's due from Higgs interactions (Minimal Flavor Viol')

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light
- keep custodial symmetry (LEP data)
- no FCNC's due from Higgs interactions (Minimal Flavor Viol')
- $U(1)_{\text{e.m.}} \times SU(3)_C$ invariance

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light
- keep custodial symmetry (LEP data)
- no FCNC's due from Higgs interactions (Minimal Flavor Viol')
- $U(1)_{\text{e.m.}} \times SU(3)_C$ invariance

Motivated by
experimental results

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light
- keep custodial symmetry (LEP data)
- no FCNC's due from Higgs interactions (Minimal Flavor Viol')
- $U(1)_{\text{e.m.}} \times SU(3)_C$ invariance
- Cutoff at $\Lambda \rightarrow$ expansion in $(p_\mu/\Lambda)^n$

Motivated by
experimental results

Effective (chiral) lagrangian for the Higgs

- One higgs: spin-0, neutral
- no other particle is light
- keep custodial symmetry (LEP data)
- no FCNC's due from Higgs interactions (Minimal Flavor Viol')
- $U(1)_{\text{e.m.}} \times SU(3)_C$ invariance
- Cutoff at $\Lambda \rightarrow$ expansion in $(p_\mu/\Lambda)^n$

Motivated by
experimental results

(spin-2 unlikely) (story can be adapted to multiple higgses if necessary)

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

see H.Haber lecture

Take the $U(1)_{\text{e.m.}} \times SU(3)_C$ lagrangian with massive W, Z and restore $SU(2)_L \times U(1)_Y$ invariance by using:

$$\Sigma(x) = e^{i T^a G^a(x)/v}$$

example:

$$\begin{aligned} W_\mu^a W^{\mu a} &\rightarrow \text{Tr} [(D_\mu \Sigma)^\dagger (D^\mu \Sigma)] \\ &= (W_\mu^a - \partial_\mu G^a)(W^{a\mu} - \partial^\mu G^a) \end{aligned}$$

+ add the most generic function for the “radial” mode
(the Higgs boson)

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$\mathcal{O}(p^2)$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$+ \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger D^\mu \Sigma \right] \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$+ \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger D^\mu \Sigma \right] \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$+ \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger D^\mu \Sigma \right] \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

$\mathcal{O}(p^4)$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$+ \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger D^\mu \Sigma \right] \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$\mathcal{O}(p^4)$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$+ \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger D^\mu \Sigma \right] \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$\mathcal{O}(p^4)$

$$+ \frac{g^2}{32\pi^2} \left(k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu} \right) \frac{h}{v} + \dots$$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$+ \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger D^\mu \Sigma \right] \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) \Sigma \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$\mathcal{O}(p^4)$

$$+ \frac{g^2}{32\pi^2} \left(k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu} \right) \frac{h}{v} + \dots$$

$+ \mathcal{O}(p^6) + \dots$

e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

Coleman, Wess, Zumino PRD 117 (1969) 2239

Callan, Coleman, Wess, Zumino PRD 117 (1969) 2247

$$\Sigma(x) = e^{i\sigma^a G^a(x)/v}$$

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

$\mathcal{O}(p^2)$

$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$\mathcal{O}(p^4)$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

$+ \mathcal{O}(p^6) + \dots$

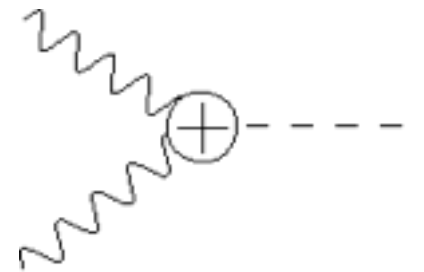
e.g. Contino, Grojean, Moretti, Piccinini, Rattazzi, JHEP 1005 (2010) 089

Effective (chiral) lagrangian for the Higgs

$$\begin{aligned}
 \mathcal{L} = & \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right) \\
 & - \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right) \\
 & - \frac{v}{\sqrt{2}} (\bar{u}_L^i d_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.} \\
 & \frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right) \\
 & + \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots \\
 & + \mathcal{O}(p^6) + \dots
 \end{aligned}$$

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$



WWh, ZZh

$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i d_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

$$+ \mathcal{O}(p^6) + \dots$$

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

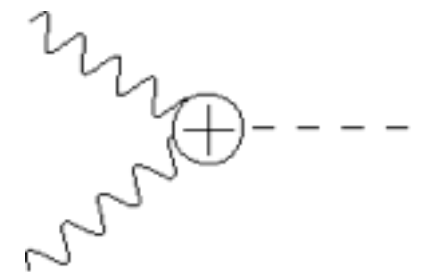
$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i d_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

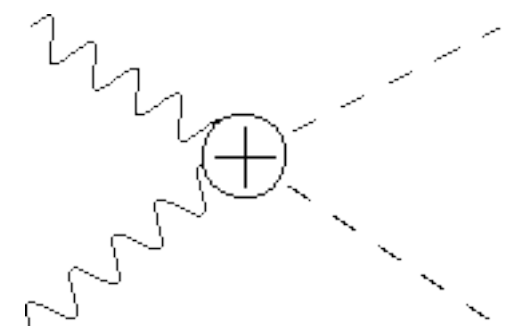
$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

$$+ \mathcal{O}(p^6) + \dots$$



WWh, ZZh



WWhh, ZZhh

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

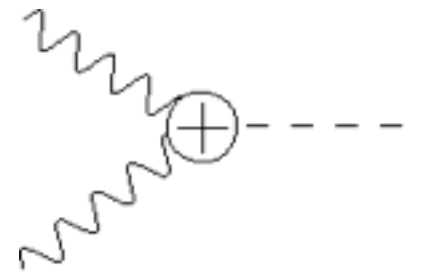
$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i d_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

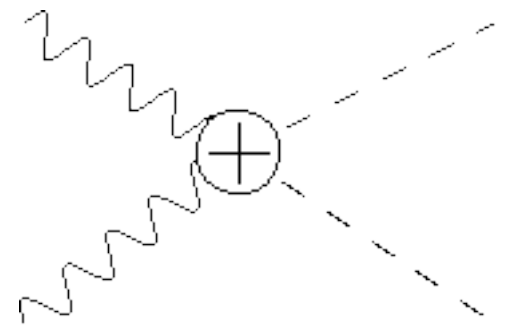
$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

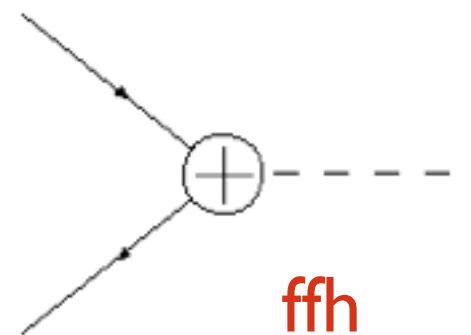
$$+ \mathcal{O}(p^6) + \dots$$



WWh, ZZh



WWhh, ZZhh



ffh

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

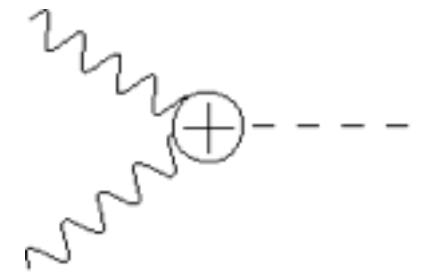
$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i d_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

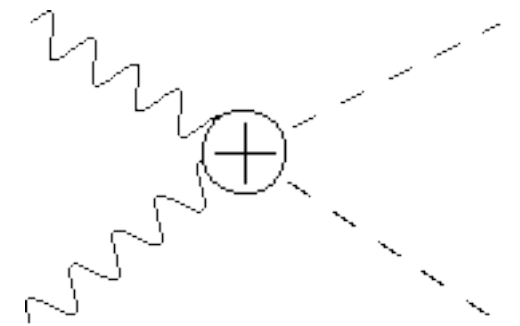
$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

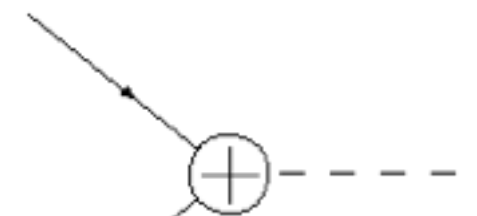
$$+ \mathcal{O}(p^6) + \dots$$



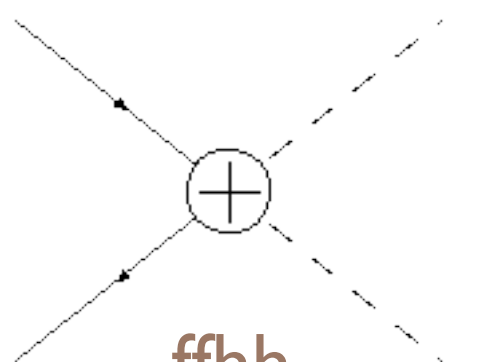
WWWh, ZZh



WWhh, ZZhh



ffh



ffhh

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

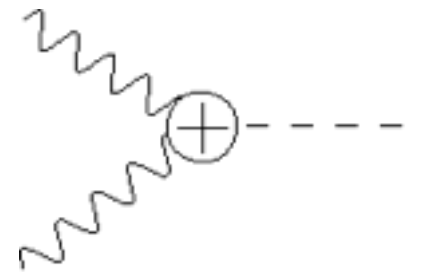
$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

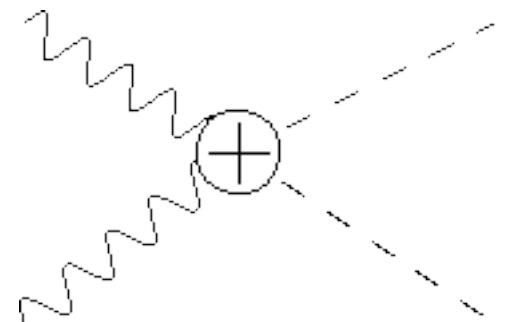
$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

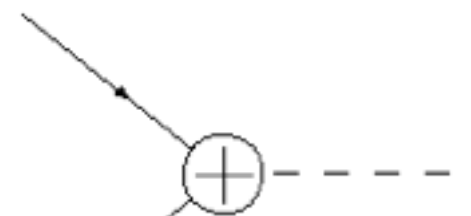
$$+ \mathcal{O}(p^6) + \dots$$



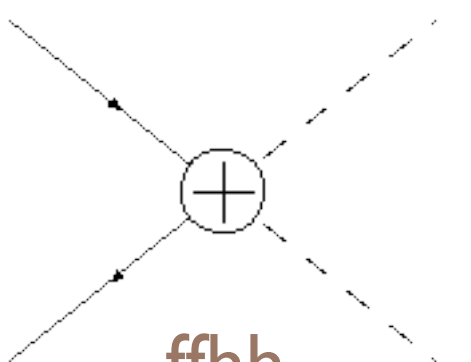
WWh, ZZh



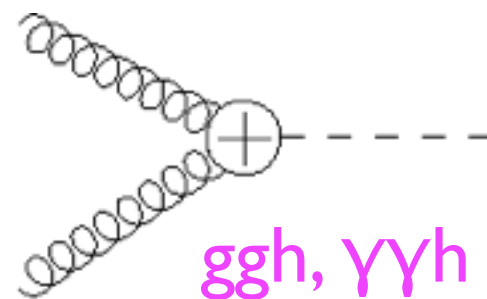
WWhh, ZZhh



ffh



ffhh



ggh, $\gamma\gamma h$

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

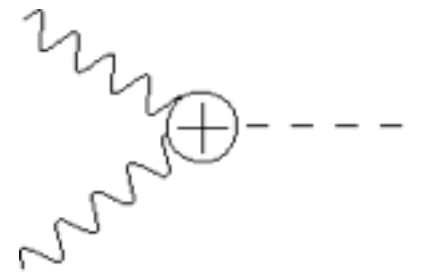
$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i d_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

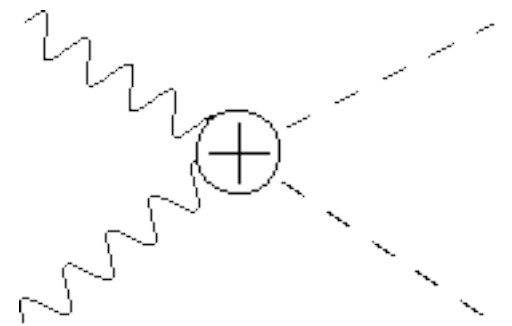
$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} (k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu}) \frac{h}{v} + \dots$$

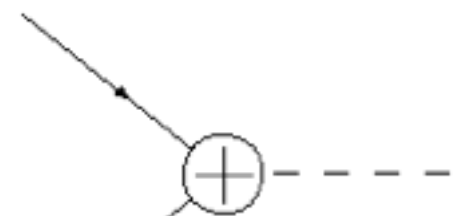
$$+ \mathcal{O}(p^6) + \dots$$



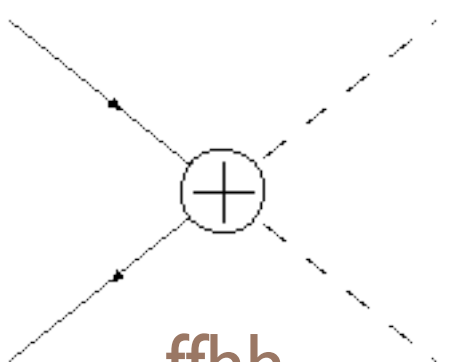
WWZh, ZZh



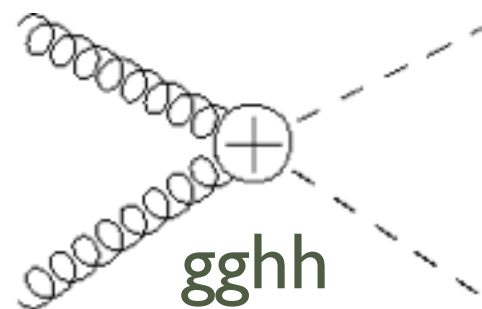
WWhh, ZZhh



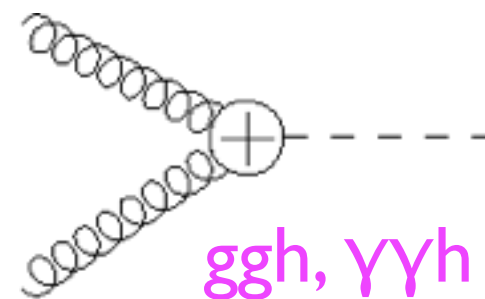
ffh



ffhh



gghh



ggh, $\gamma\gamma h$

Effective (chiral) lagrangian for the Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \left(\frac{1}{2} m_h^2 h^2 + d_3 \left(\frac{m_h^2}{2v} \right) h^3 + d_4 \left(\frac{m_h^2}{8v^2} \right) h^4 + \dots \right)$$

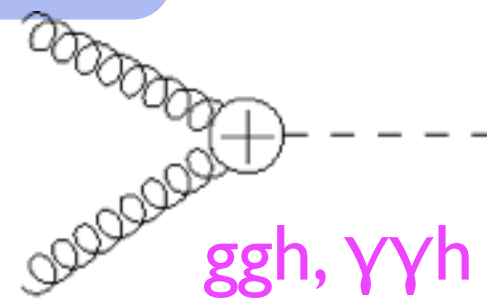
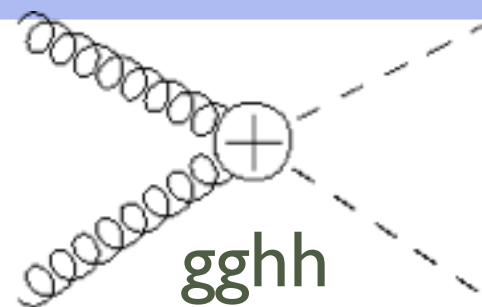
$$- \left(m_W^2 W_\mu W^\mu + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + b_3 \frac{h^3}{v^3} + \dots \right)$$

$$- \frac{v}{\sqrt{2}} (\bar{u}_L^i \bar{d}_L^i) - \left[1 + c \frac{h}{v} + c_2 \frac{h^2}{v^2} + \dots \right] \begin{pmatrix} y_{ij}^u u_R^j \\ y_{ij}^d d_R^j \end{pmatrix} + \text{h.c.}$$

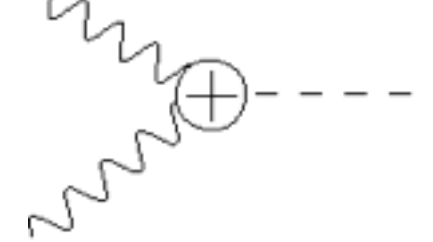
$$\frac{g_s^2}{48\pi^2} G^{\mu\nu a} G_{\mu\nu}^a \left(k_g \frac{h}{v} + \frac{1}{2} k_{2g} \frac{h^2}{v^2} + \dots \right) + \frac{e^2}{32\pi^2} F_{\mu\nu} F^{\mu\nu} \left(k_\gamma \frac{h}{v} + \dots \right)$$

$$+ \frac{g^2}{32\pi^2} \left(k_{WW} W_{\mu\nu}^+ W^{-\mu\nu} + k_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + k_{Z\gamma} Z_{\mu\nu} F^{\mu\nu} \right) \frac{h}{v} + \dots$$

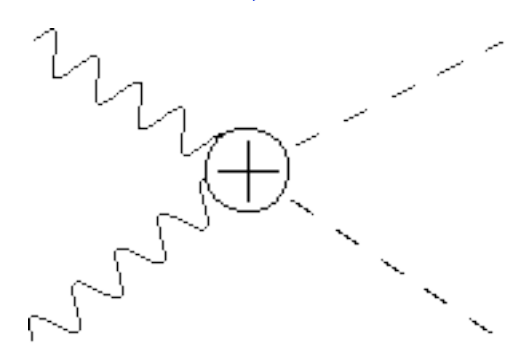
$$+ \mathcal{O}(p^6) + \dots$$



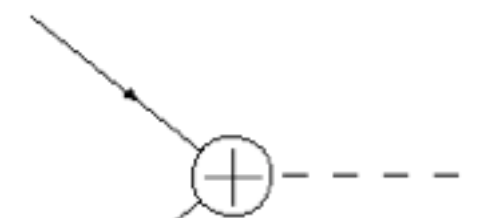
(higher order)



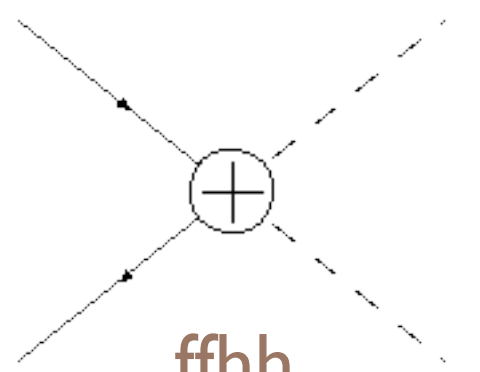
WW h, ZZ h



WW hh, ZZ hh



ff h



ff hh

Effective (chiral) lagrangian for the Higgs

- Most general Lagrangian:

Effective (chiral) lagrangian for the Higgs

- Most general Lagrangian:
- sufficient parameterization for all the EWSB models with 1 light Higgs

Effective (chiral) lagrangian for the Higgs

- Most general Lagrangian:
 - sufficient parameterization for all the EWSB models with 1 light Higgs
 - also sufficient for the “impostors” models with one light impostor and a heavy Higgs (or Higgsless)

Effective (chiral) lagrangian for the Higgs

- Most general Lagrangian:
 - sufficient parameterization for all the EWSB models with 1 light Higgs
 - also sufficient for the “impostors” models with one light impostor and a heavy Higgs (or Higgsless)
 - SM is $a=1$, $b=1$, $c=1$, other c 's and b 's $=0$, etc. [in this case $H(x)/v=\Sigma(x)(1+h(x)/v)$]

Effective (chiral) lagrangian for the Higgs

- Most general Lagrangian:
 - sufficient parameterization for all the EWSB models with 1 light Higgs
 - also sufficient for the “impostors” models with one light impostor and a heavy Higgs (or Higgsless)
 - SM is $a=1$, $b=1$, $c=1$, other c 's and b 's $=0$, etc. [in this case $H(x)/v = \Sigma(x)(1+h(x)/v)$]

but too many parameters!!

for early data need more constraints

→ need to add more assumptions

Composite Higgs

Recall:

Composite Higgs

Recall:

Technicolor:

QCD:

Composite Higgs

Recall:

Technicolor:

QCD:

Goldstone bosons \longleftrightarrow “pions”

Composite Higgs

Recall:

Technicolor:

QCD:

Goldstone bosons



“pions”

v



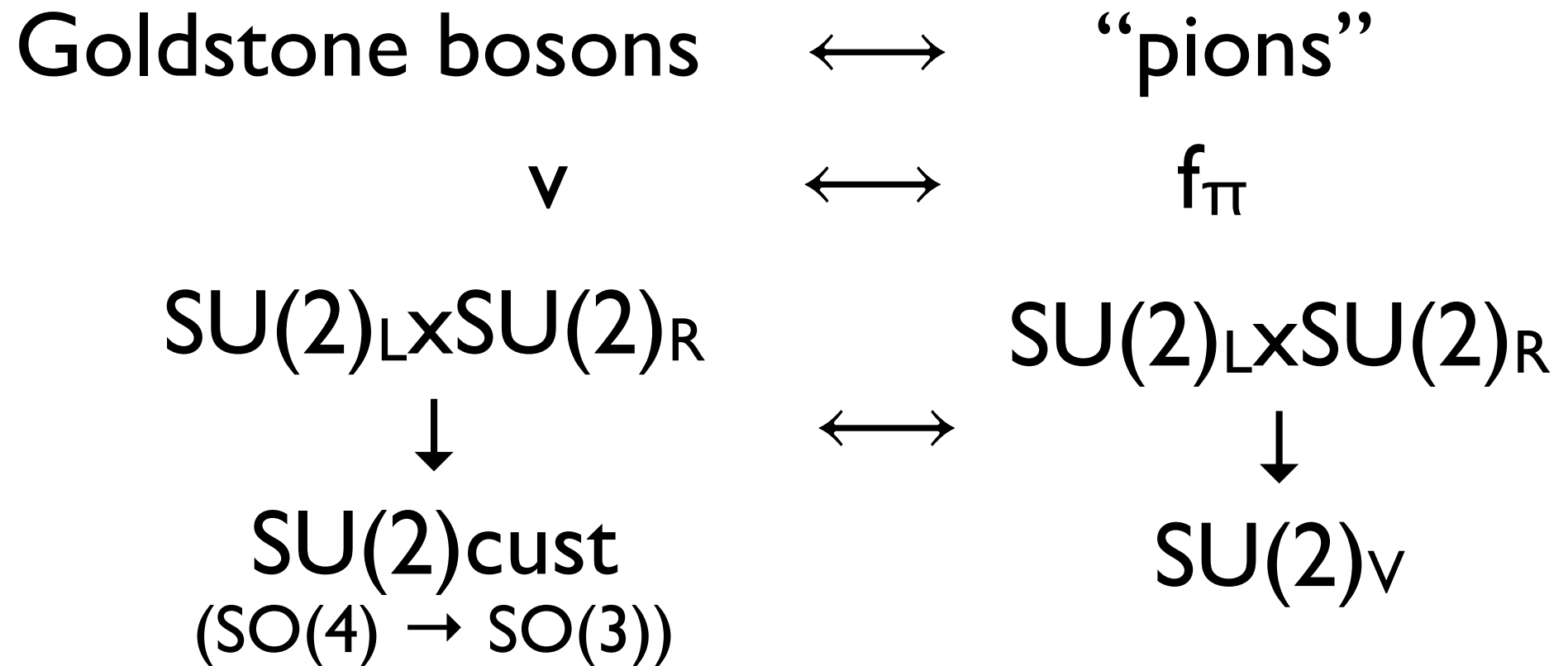
f_π

Composite Higgs

Recall:

Technicolor:

QCD:



Composite Higgs

Recall:

Technicolor:

QCD:

Goldstone bosons

\longleftrightarrow

“pions”

v

\longleftrightarrow

f_π

$SU(2)_L \times SU(2)_R$

\downarrow

$SU(2)_{\text{cust}}$
($SO(4) \rightarrow SO(3)$)

\longleftrightarrow

$SU(2)_L \times SU(2)_R$

\downarrow

$SU(2)_V$

$\Lambda = 4\pi v$

\longleftrightarrow

$\Lambda = 4\pi f_\pi$

Composite Higgs

Recall:

Technicolor:

QCD:

Goldstone bosons

\longleftrightarrow

“pions”

v

\longleftrightarrow

f_π

$SU(2)_L \times SU(2)_R$

\downarrow

\longleftrightarrow

$SU(2)_L \times SU(2)_R$

\downarrow

$SU(2)_{\text{cust}}$
($SO(4) \rightarrow SO(3)$)

$SU(2)_V$

$\Lambda = 4\pi v$

\longleftrightarrow

$\Lambda = 4\pi f_\pi$

$m_{\rho_{TC}} \sim \Lambda$

$m_\rho \sim \Lambda$

Composite Higgs

To reduce params \rightarrow have the Higgs boson as another “pion” of some strongly coupled theory
howto: just enlarge the symmetry group to have more Goldstone bosons (and symmetry breaking at above the EW scale)

Composite Higgs

To reduce params \rightarrow have the Higgs boson as another “pion” of some strongly coupled theory

howto: just enlarge the symmetry group to have more Goldstone bosons (and symmetry breaking at above the EW scale)

ex: $SO(5) \rightarrow SO(4)$ broken at scale $f (>v)$:

4 massless d.o.f.: $G^a + h$

$$\Lambda_c = 4\pi f > \Lambda = 4\pi v$$

$$m_\rho \sim \Lambda_c$$

Composite Higgs

To reduce params \rightarrow have the Higgs boson as another “pion” of some strongly coupled theory

howto: just enlarge the symmetry group to have more Goldstone bosons (and symmetry breaking at above the EW scale)

ex: $SO(5) \rightarrow SO(4)$ broken at scale $f (>v)$:

4 massless d.o.f.: $G^a + h$

$$\Lambda_c = 4\pi f > \Lambda = 4\pi v$$

$$m_\rho \sim \Lambda_c$$

later $SO(4)$ (and $SU(2)_L \times U(1)_Y$) broken by v :
 G^a “eaten” by W, Z, h massive with $m_h \sim O(v)$

Composite Higgs

To reduce params \rightarrow have the Higgs boson as another “pion” of some strongly coupled theory

howto: just enlarge the symmetry group to have more Goldstone bosons (and symmetry breaking at above the EW scale)

ex: $SO(5) \rightarrow SO(4)$ broken at scale $f (>v)$:

4 massless d.o.f.: $G^a + h$

$$\Lambda_c = 4\pi f > \Lambda = 4\pi v$$

$$m_\rho \sim \Lambda_c$$

later $SO(4)$ (and $SU(2)_L \times U(1)_Y$) broken by v :
 G^a “eaten” by W, Z , h massive with $m_h \sim O(v)$

($\Lambda_c > \Lambda$ b/c h partially unitarize $G^a G^a$ scattering)

Composite Higgs

To reduce params \rightarrow have the Higgs boson as another “pion” of some strongly coupled theory

howto: just enlarge the symmetry group to have more Goldstone bosons (and symmetry breaking at above the EW scale)

ex: $SO(5) \rightarrow SO(4)$ broken at scale $f (>v)$:

4 massless d.o.f.: $G^a + h$

$$\Lambda_c = 4\pi f > \Lambda = 4\pi v$$

$$m_\rho \sim \Lambda_c$$

later $SO(4)$ (and $SU(2)_L \times U(1)_Y$) broken by v :
 G^a “eaten” by W, Z , h massive with $m_h \sim O(v)$

($\Lambda_c > \Lambda$ b/c h partially unitarize $G^a G^a$ scattering)

Improvement of
 $O(\xi = v^2/f^2)$

Composite Higgs

“Strongly Interacting Light Higgs” (SILH) Giudice, Grojean, Pomarol,
Rattazzi arXiv:hep-ph/0703164

most general lagrangian for Higgs boson as a pGB,
expanded at $O(\xi)$

Composite Higgs

“Strongly Interacting Light Higgs” (SILH) Giudice, Grojean, Pomarol,
Rattazzi arXiv:hep-ph/0703164

most general lagrangian for Higgs boson as a pGB,
expanded at $O(\xi)$

Parameters	SILH
a	$1 - (c_H - c_r/2) \xi/2$
b	$1 + (c_r - 2c_H) \xi$
b_3	$(c_r - 2c_H)2\xi/3$
c	$1 - (c_H/2 + c_y) \xi$
c_2	$-(c_H + 3c_y + c_r/4) \xi/2$
d_3	$1 + (c_6 - c_r/4 - 3c_H/2) \xi$
d_4	$1 + (6c_6 - 25c_H/3 - 11c_r/6) \xi$
$k_g = k_{2g}$	$3c_g(y_t^2/g_\rho^2)\xi$
k_γ	$2c_\gamma(g^2/g_\rho^2)\xi$

Gillioz, Grober, Grojean, Muhlleitner,
Salvioni arXiv:1206.7120

Composite Higgs

“Strongly Interacting Light Higgs” (SILH) Giudice, Grojean, Pomarol,
Rattazzi arXiv:hep-ph/0703164

most general lagrangian for Higgs boson as a pGB,
expanded at $O(\xi)$

- 2 parameters control
a, b, b₃, c, c₂

Parameters	SILH
a	$1 - (c_H - c_r/2) \xi/2$
b	$1 + (c_r - 2c_H) \xi$
b_3	$(c_r - 2c_H) 2\xi/3$
c	$1 - (c_H/2 + c_y) \xi$
c_2	$-(c_H + 3c_y + c_r/4) \xi/2$
d_3	$1 + (c_6 - c_r/4 - 3c_H/2) \xi$
d_4	$1 + (6c_6 - 25c_H/3 - 11c_r/6) \xi$
$k_g = k_{2g}$	$3c_g(y_t^2/g_\rho^2)\xi$
k_γ	$2c_\gamma(g^2/g_\rho^2)\xi$

Gillioz, Grober, Grojean, Muhlleitner,
Salvioni arXiv:1206.7120

Composite Higgs

“Strongly Interacting Light Higgs” (SILH)

Giudice, Grojean, Pomarol,
Rattazzi arXiv:hep-ph/0703164

most general lagrangian for Higgs boson as a pGB,
expanded at $O(\xi)$

- 2 parameters control
a, b, b_3 , c, c_2
- + 1 parameter for higgs
self interactions

Parameters	SILH
a	$1 - (c_H - c_r/2) \xi/2$
b	$1 + (c_r - 2c_H) \xi$
b_3	$(c_r - 2c_H) 2\xi/3$
c	$1 - (c_H/2 + c_y) \xi$
c_2	$-(c_H + 3c_y + c_r/4) \xi/2$
d_3	$1 + (c_6 - c_r/4 - 3c_H/2) \xi$
d_4	$1 + (6c_6 - 25c_H/3 - 11c_r/6) \xi$
$k_g = k_{2g}$	$3c_g(y_t^2/g_\rho^2)\xi$
k_γ	$2c_\gamma(g^2/g_\rho^2)\xi$

Gillioz, Grober, Grojean, Muhlleitner,
Salvioni arXiv:1206.7120

Composite Higgs

“Strongly Interacting Light Higgs” (SILH)

Giudice, Grojean, Pomarol,
Rattazzi arXiv:hep-ph/0703164

most general lagrangian for Higgs boson as a pGB,
expanded at $O(\xi)$

- 2 parameters control a, b, b_3, c, c_2
- + 1 parameter for higgs self interactions
- + 2 params for k_g, k_{g^2} and k_γ (parametrically suppressed by $O(g^2/g_\rho^2)$)

Parameters	SILH
a	$1 - (c_H - c_r/2) \xi/2$
b	$1 + (c_r - 2c_H) \xi$
b_3	$(c_r - 2c_H) 2\xi/3$
c	$1 - (c_H/2 + c_y) \xi$
c_2	$-(c_H + 3c_y + c_r/4) \xi/2$
d_3	$1 + (c_6 - c_r/4 - 3c_H/2) \xi$
d_4	$1 + (6c_6 - 25c_H/3 - 11c_r/6) \xi$
$k_g = k_{2g}$	$3c_g(y_t^2/g_\rho^2)\xi$
k_γ	$2c_\gamma(g^2/g_\rho^2)\xi$

Gillioz, Grober, Grojean, Muhlleitner,
Salvioni arXiv:1206.7120

Composite Higgs

“Strongly Interacting Light Higgs” (SILH)

Giudice, Grojean, Pomarol,
Rattazzi arXiv:hep-ph/0703164

most general lagrangian for Higgs boson as a pGB,
expanded at $O(\xi)$

- 2 parameters control a, b, b_3, c, c_2
- + 1 parameter for higgs self interactions
- + 2 params for k_g, k_{g2} and k_γ (parametrically suppressed by $O(g^2/g_\rho^2)$)

Parameters	SILH
a	$1 - (c_H - c_r/2) \xi/2$
b	$1 + (c_r - 2c_H) \xi$
b_3	$(c_r - 2c_H) 2\xi/3$
c	$1 - (c_H/2 + c_y) \xi$
c_2	$-(c_H + 3c_y + c_r/4) \xi/2$
d_3	$1 + (c_6 - c_r/4 - 3c_H/2) \xi$
d_4	$1 + (6c_6 - 25c_H/3 - 11c_r/6) \xi$
$k_g = k_{2g}$	$3c_g(y_t^2/g_\rho^2)\xi$
k_γ	$2c_\gamma(g^2/g_\rho^2)\xi$

Gillioz, Grober, Grojean, Muhlleitner,
Salvioni arXiv:1206.7120

SILH: for Higgs decays still 4 free parameters: $ggh, \gamma\gamma h, VVh, f\bar{f}h$

Composite Higgs

Composite Higgs

- Current Higgs phenomenology is encoded in g_{hVV} , g_{ffh} (for each fermion), g_{hgg} , $g_{h\gamma\gamma}$

Composite Higgs

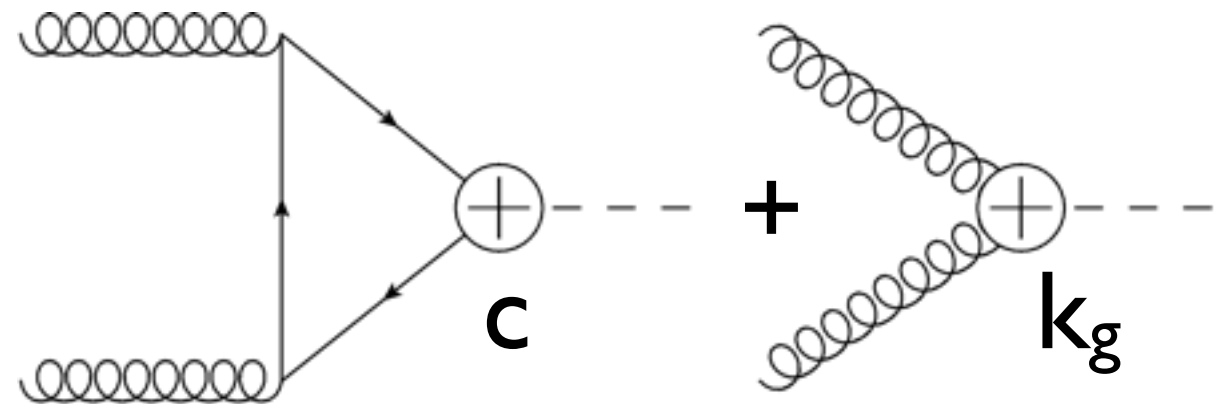
- Current Higgs phenomenology is encoded in g_{hVV} , g_{ffh} (for each fermion), g_{hgg} , $g_{h\gamma\gamma}$
- SILH+Flavor Universality: $a, c_{\text{univ}}, k_g, k_\gamma$

Composite Higgs

- Current Higgs phenomenology is encoded in g_{hVV} , g_{ffh} (for each fermion), g_{hgg} , $g_{h\gamma\gamma}$
- SILH+Flavor Universality: $a, c_{\text{univ}}, k_g, k_\gamma$
- Forced to 4-dim fit to check for compositeness from the very beginning? :(

Composite Higgs

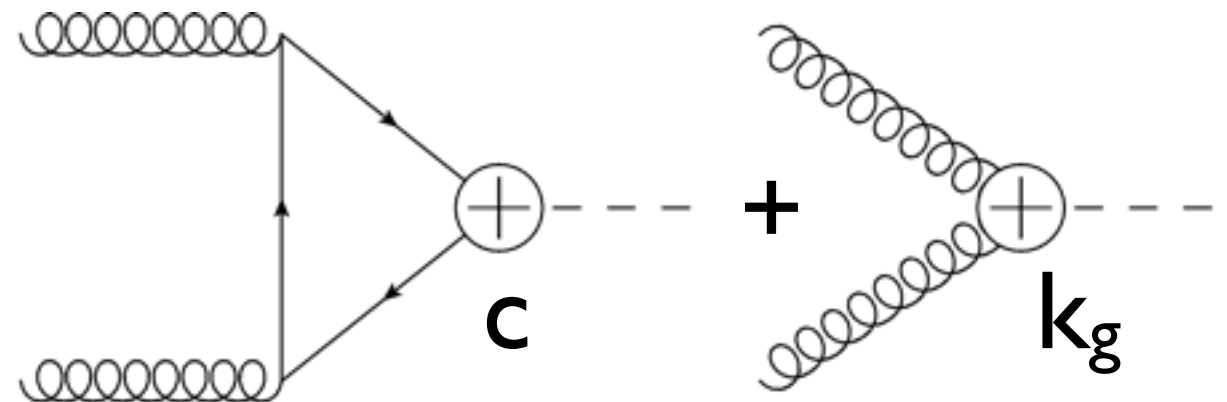
- Current Higgs phenomenology is encoded in g_{hVV} , g_{ffh} (for each fermion), g_{hgg} , $g_{h\gamma\gamma}$
- SILH+Flavor Universality: $a, c_{\text{univ}}, k_g, k_\gamma$
- Forced to 4-dim fit to check for compositeness from the very beginning? :(



Composite Higgs

- Current Higgs phenomenology is encoded in g_{hVV} , g_{ffh} (for each fermion), g_{hgg} , $g_{h\gamma\gamma}$
- SILH+Flavor Universality: $a, c_{\text{univ}}, k_g, k_\gamma$
- Forced to 4-dim fit to check for compositeness from the very beginning? :(

In “minimal” pGB models
the contribution from heavy
fermions (like top partners)
cancel out at zero
momentum \rightarrow drop k_g, k_γ



Falkowski, PRD 77 (2008) 055018

Low, Rattazzi, Vichi, JHEP 1004 (2010) 126

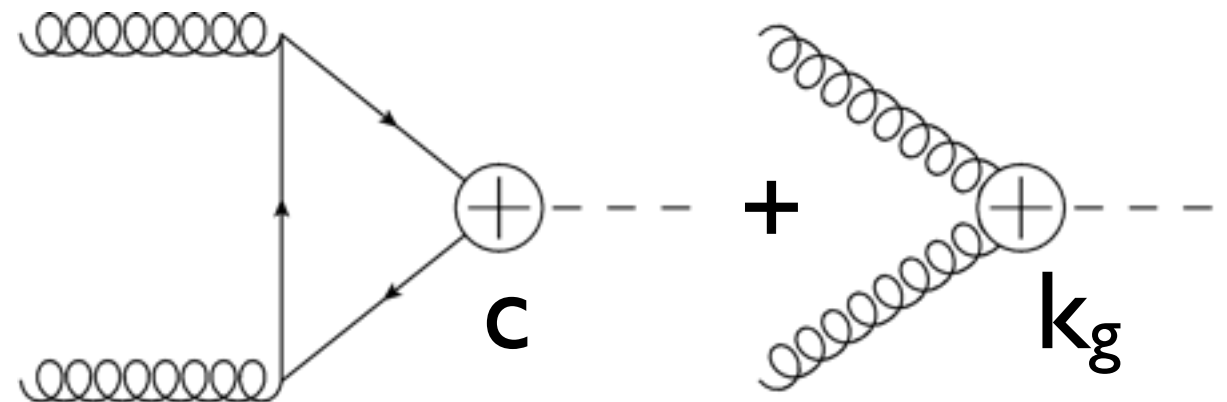
Azatov, Galloway, PRD 85 (2012) 055013

Composite Higgs

- Current Higgs phenomenology is encoded in g_{hVV} , g_{ffh} (for each fermion), g_{hgg} , $g_{h\gamma\gamma}$
- SILH+Flavor Universality: $a, c_{\text{univ}}, k_g, k_\gamma$
- Forced to 4-dim fit to check for compositeness from the very beginning? :(

In “minimal” pGB models
the contribution from heavy
fermions (like top partners)
cancel out at zero

momentum \rightarrow drop k_g, k_γ



Minimal Composite Higgs Models:
2-paramters Higgs pheno: a & c

Falkowski, PRD 77 (2008) 055018

Low, Rattazzi, Vichi, JHEP 1004 (2010) 126

Azatov, Galloway, PRD 85 (2012) 055013

Composite Higgs

Composite Higgs

- Many composite Higgs models can be parameterized in terms of $a=g_{VVh}$ and $c=g_{ffh}$ as long as there is one Higgs boson and a mass gap with the partners/resonances

Composite Higgs

- Many composite Higgs models can be parameterized in terms of $a=g_{VVh}$ and $c=g_{ffh}$ as long as there is one Higgs boson and a mass gap with the partners/resonances
- deviation from unity is function of $\xi=v^2/f^2$

Composite Higgs

- Many composite Higgs models can be parameterized in terms of $a=g_{VVh}$ and $c=g_{ffh}$ as long as there is one Higgs boson and a mass gap with the partners/resonances
- deviation from unity is function of $\xi=v^2/f^2$
- In specific models a and c tend to be related

Composite Higgs

- Many composite Higgs models can be parameterized in terms of $a=g_{VVh}$ and $c=g_{ffh}$ as long as there is one Higgs boson and a mass gap with the partners/resonances
- deviation from unity is function of $\xi=v^2/f^2$
- In specific models a and c tend to be related

MCHM = Minimal Composite Higgs Model (w/ $SO(5)/SO(4)$)

Composite Higgs

- Many composite Higgs models can be parameterized in terms of $a=g_{VVh}$ and $c=g_{ffh}$ as long as there is one Higgs boson and a mass gap with the partners/resonances
- deviation from unity is function of $\xi=v^2/f^2$
- In specific models a and c tend to be related

“MCHM5”: $a = \sqrt{1 - \xi}$, $c = \frac{1 - 2\xi}{\sqrt{1 - \xi}}$ (also LR Twin Higgs)

MCHM = Minimal Composite Higgs Model (w/ $SO(5)/SO(4)$)

Composite Higgs

- Many composite Higgs models can be parameterized in terms of $a=g_{VVh}$ and $c=g_{ffh}$ as long as there is one Higgs boson and a mass gap with the partners/resonances
- deviation from unity is function of $\xi=v^2/f^2$
- In specific models a and c tend to be related

“MCHM5”: $a = \sqrt{1 - \xi}$, $c = \frac{1 - 2\xi}{\sqrt{1 - \xi}}$ (also LR Twin Higgs)

“MCHM4”: $a = \sqrt{1 - \xi}$, $c = \sqrt{1 - \xi}$ (also simplest little Higgs with $(SU(3)/SU(2))^2$)

MCHM = Minimal Composite Higgs Model (w/ $SO(5)/SO(4)$)

Looking at the data

DISCLAIMER

- In the following I will use fits to Higgs parameters done by theorists, to illustrate the main points

DISCLAIMER

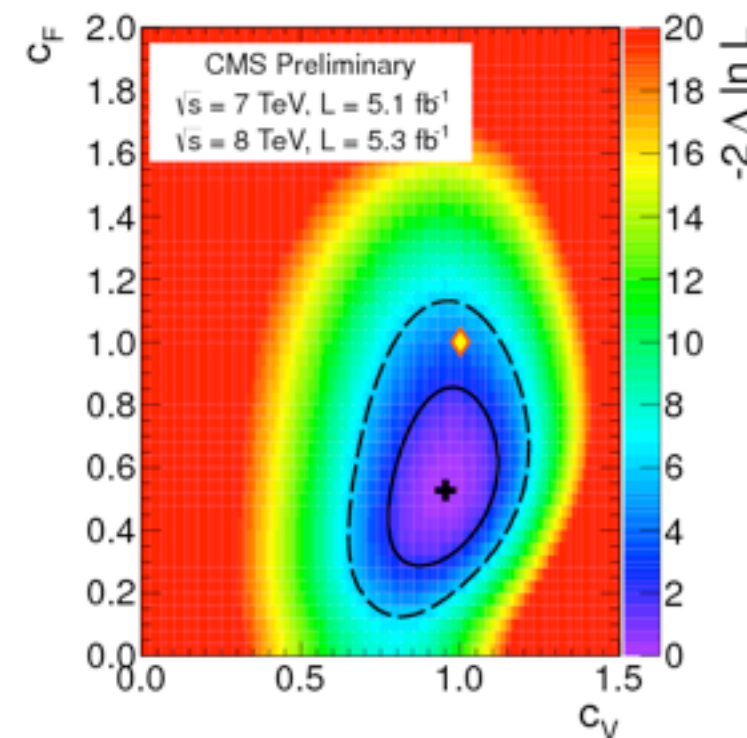
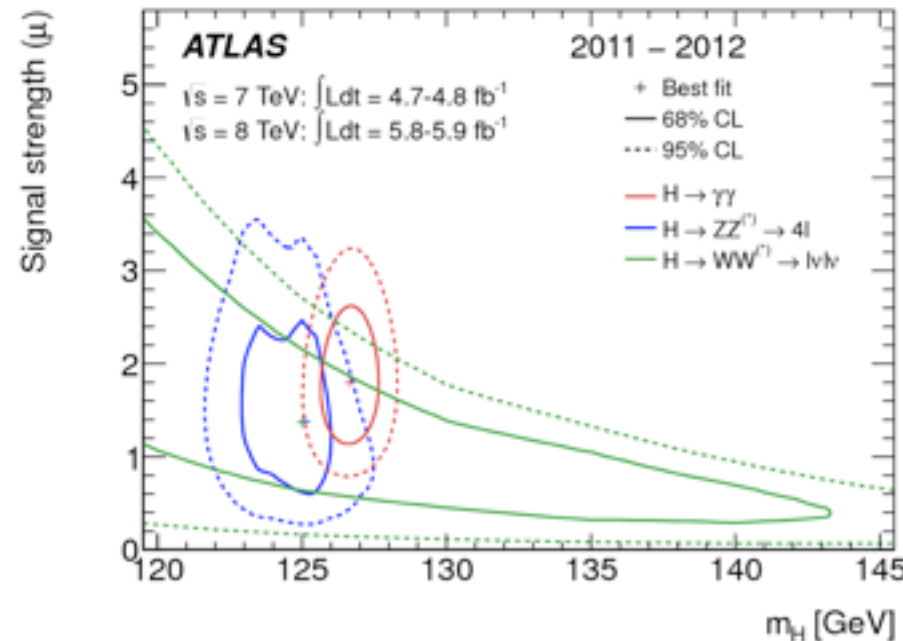
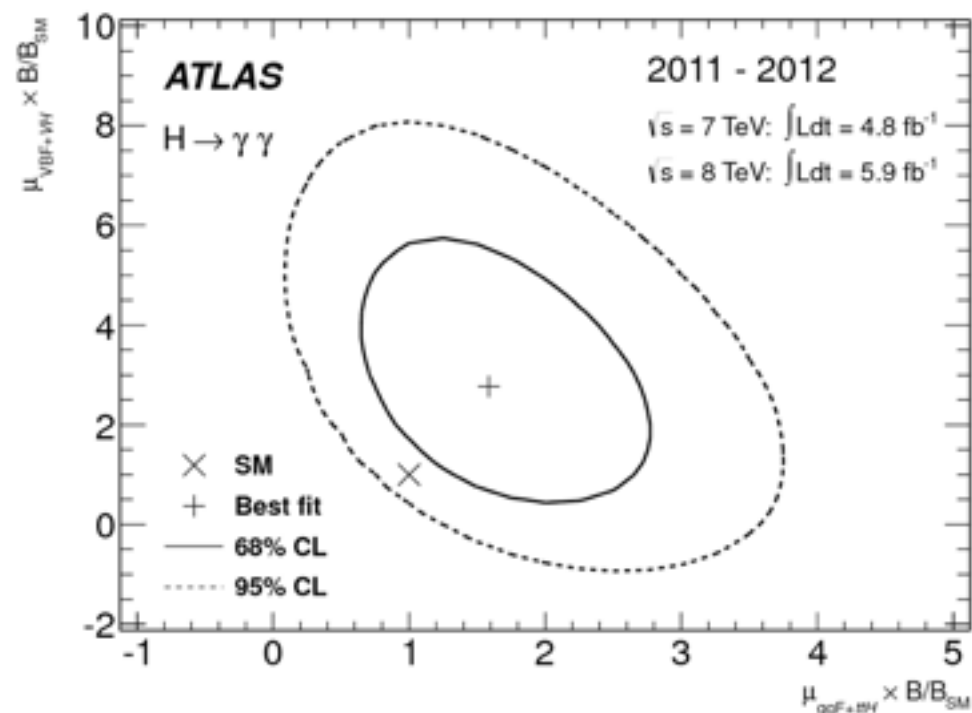
- In the following I will use fits to Higgs parameters done by theorists, to illustrate the main points
- Higgs data has still large uncertainties and the likelihoods are not public → results are approximate

DISCLAIMER

- In the following I will use fits to Higgs parameters done by theorists, to illustrate the main points
- Higgs data has still large uncertainties and the likelihoods are not public → results are approximate
- Ultimately these fits will/should be performed/checked by the experiments (already started)

DISCLAIMER

- In the following I will use fits to Higgs parameters done by theorists, to illustrate the main points
- Higgs data has still large uncertainties and the likelihoods are not public → results are approximate
- Ultimately these fits will/should be performed/checked by the experiments (already started)

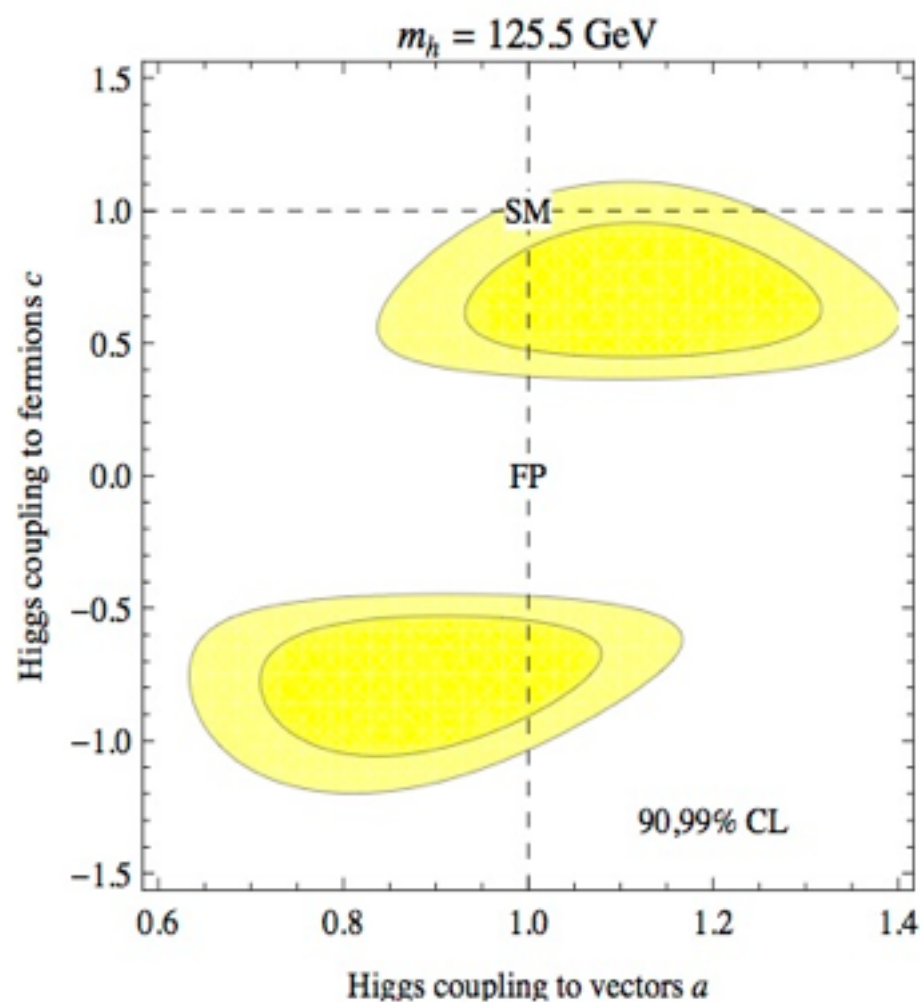


DISCLAIMER

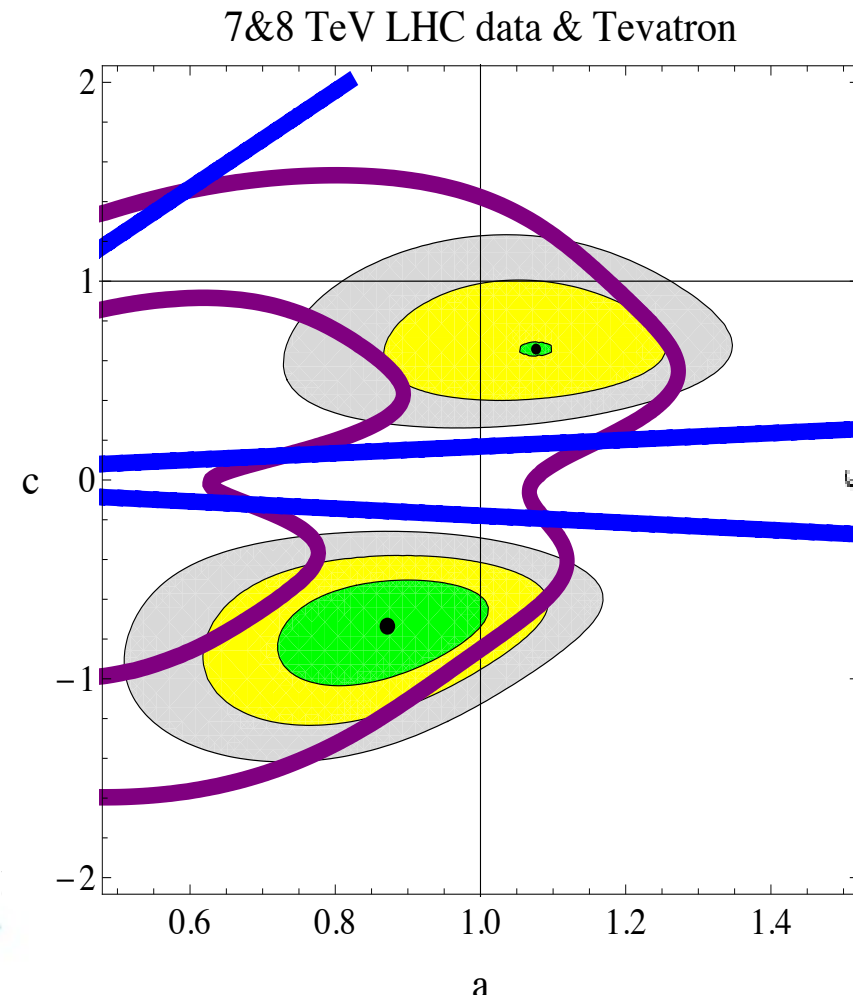
- In the following I will use fits to Higgs parameters done by theorists, to illustrate the main points
- Higgs data has still large uncertainties and the likelihoods are not public → results are approximate
- Ultimately these fits will/should be performed/checked by the experiments (already started)
- Caveat: “sigmas” are “theorist sigmas”, but qualitative conclusions should hold

Looking at the data

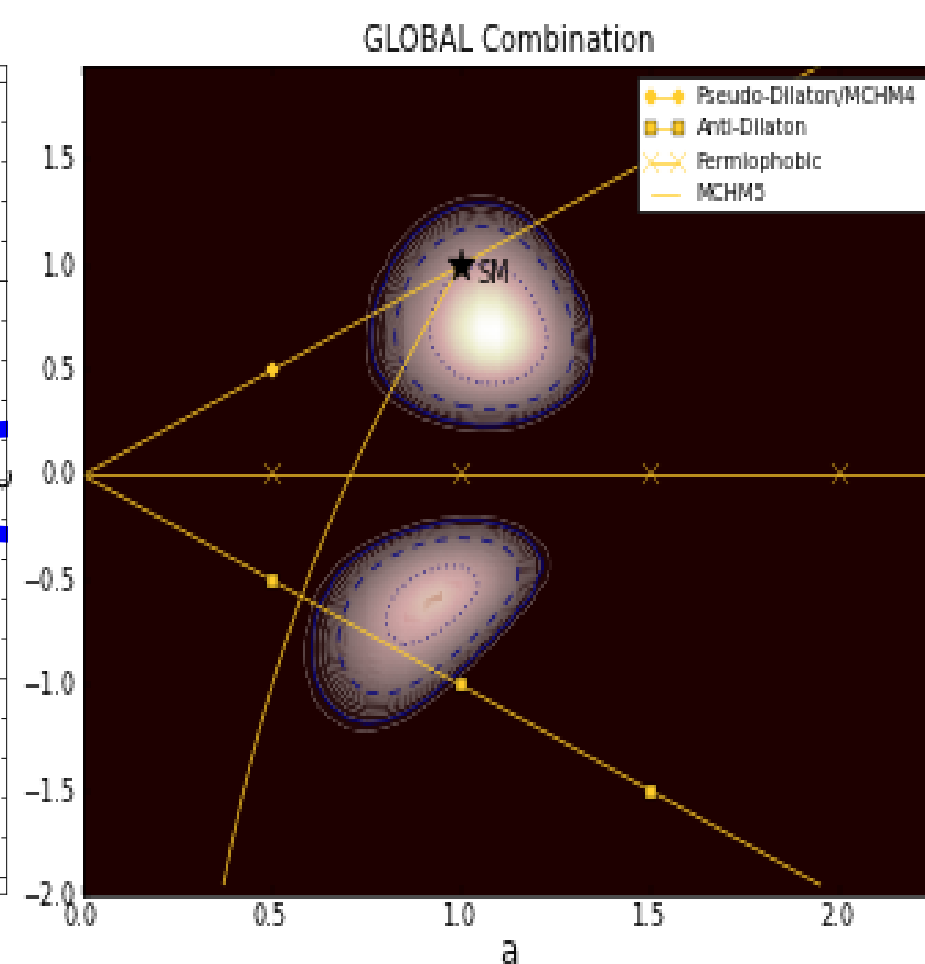
- Theorists views at the (a,c) plane after July 4th (ATLAS+CMS+TeVatron)



Giardino, Kannike, Raidal,
Strumia arXiv:1207.1347



Espinosa, Grojean, Muhlleitner,
Trott arXiv:1207.1717

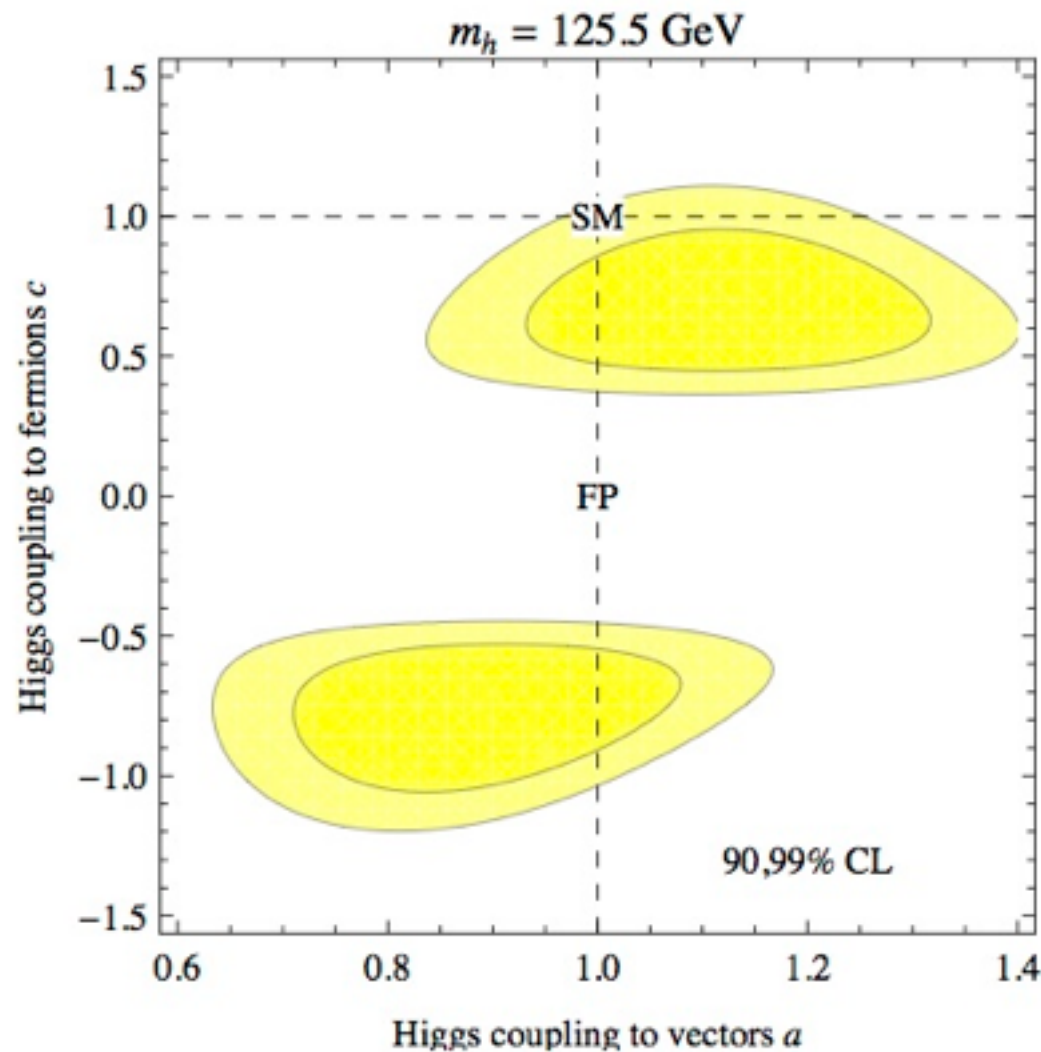


Ellis, You arXiv:1207.1693

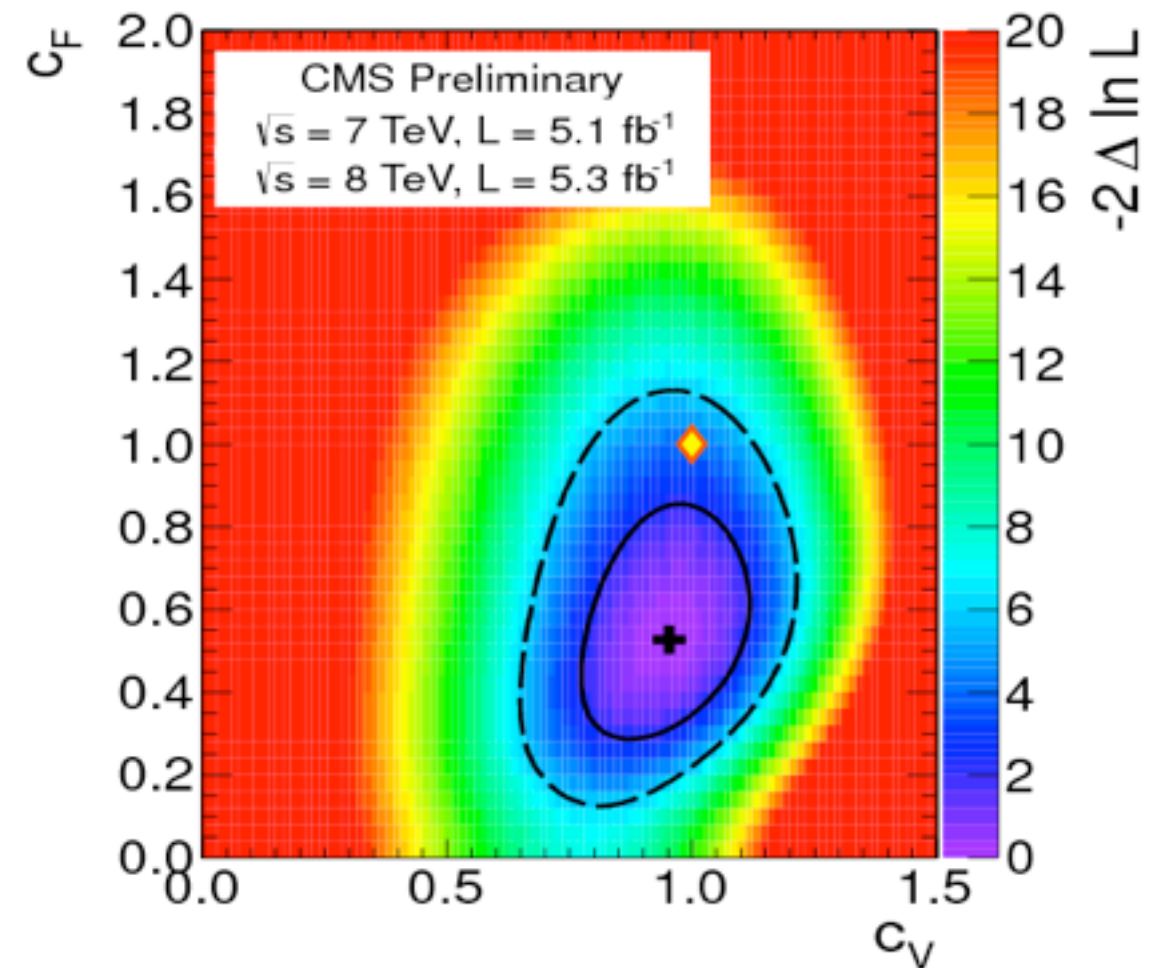
The second solution is allowed by a degeneracy in $h \rightarrow \gamma\gamma$: $\sim |8.3a - 1.8c|^2$

Looking at the data

- Sanity check with experimental fit (CMS):

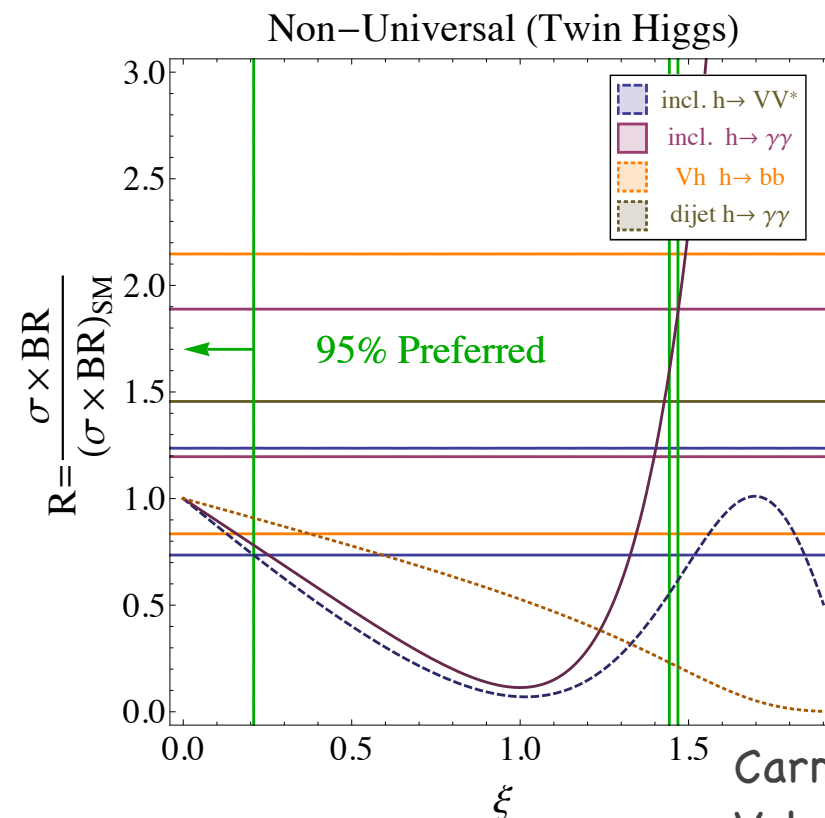
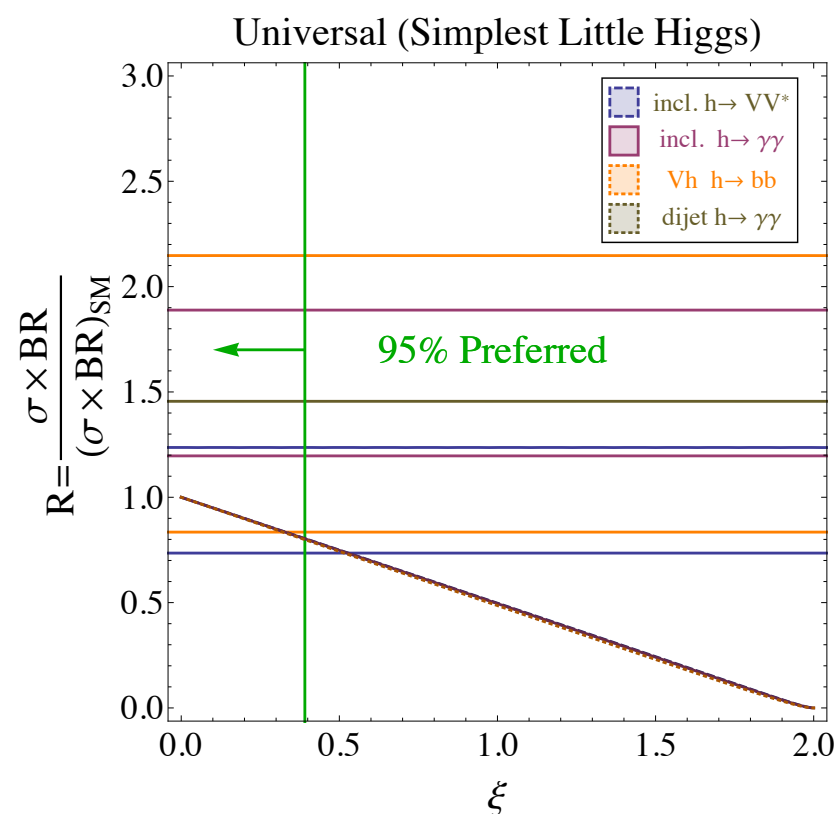


Giardino, Kannike, Raidal,
Strumia arXiv:1207.1347



Looking at the data

- The fact that the new particle walks and talks as the SM Higgs boson constrains the amount of compositeness of the Higgs boson
- $\xi < 0.3-0.4$ is still allowed (and already expected by EW precision data) $\rightarrow f > 1.6-1.8 v$



Carmi, Falkowski, Kuflik,
Volansky, Zupan arXiv:1207.1718

wait a sec...

Testing compositeness?

wait a sec...

Testing compositeness?

- Clearly many models can produce deviations in c & a (or even in c, a, k_Y, k_g)

wait a sec...

Testing compositeness?

- Clearly many models can produce deviations in c & a (or even in c, a, k_γ, k_g)
- Seeing a deviation in c & a doesn't point to Higgs being “composite”

wait a sec...

Testing compositeness?

- Clearly many models can produce deviations in c & a (or even in c, a, k_Y, k_g)
- Seeing a deviation in c & a doesn't point to Higgs being “composite”
- To test compositeness one need to check predictions from pGB nature of the Higgs

wait a sec...

Testing compositeness?

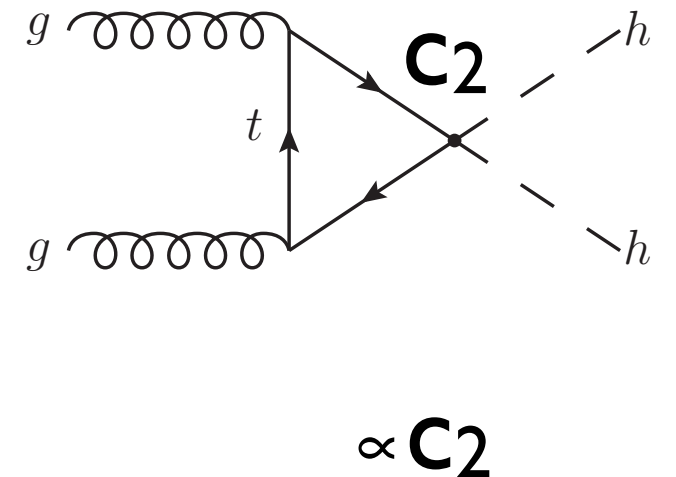
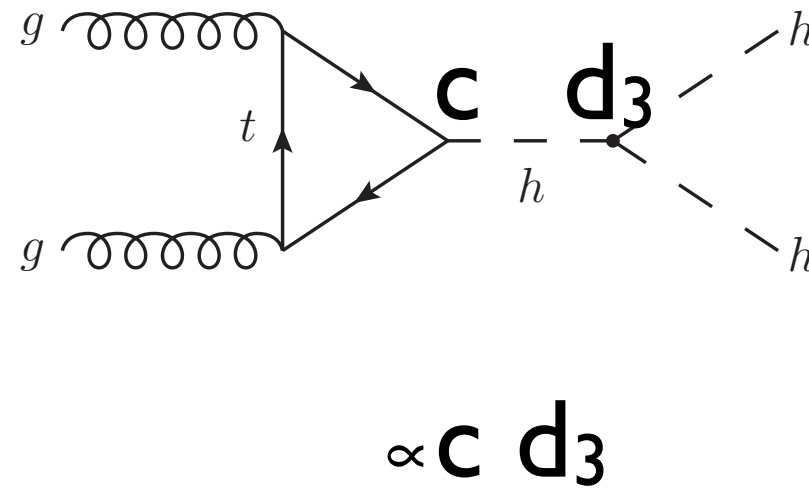
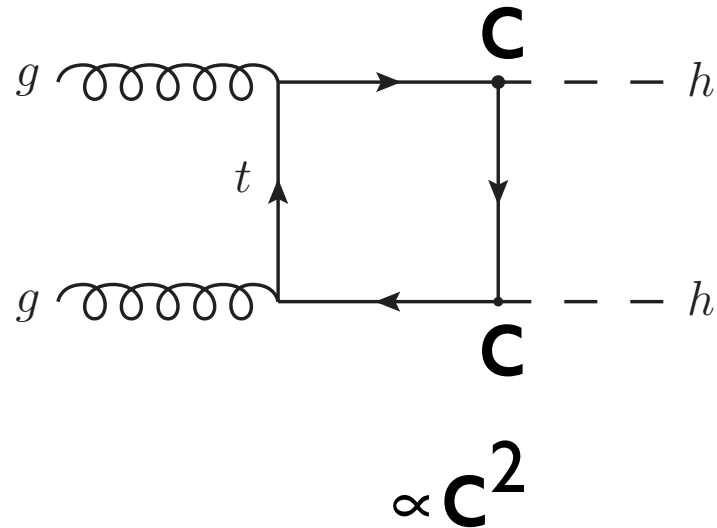
- Clearly many models can produce deviations in c & a (or even in c, a, k_Y, k_g)
- Seeing a deviation in c & a doesn't point to Higgs being “composite”
- To test compositeness one need to check predictions from pGB nature of the Higgs
- $pGB \rightarrow “\exp(i h/v)” \rightarrow$ relations between lowest order in h and higher orders

wait a sec...

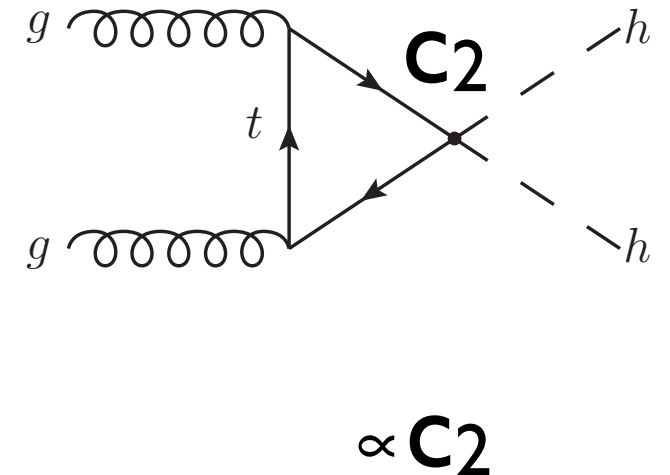
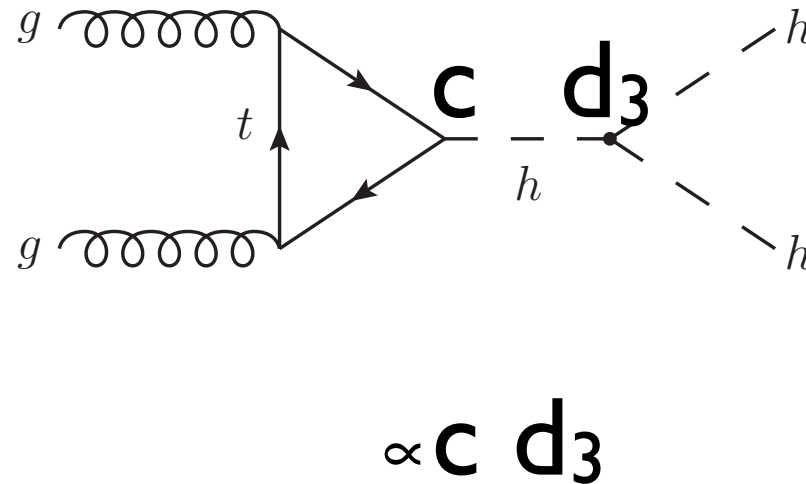
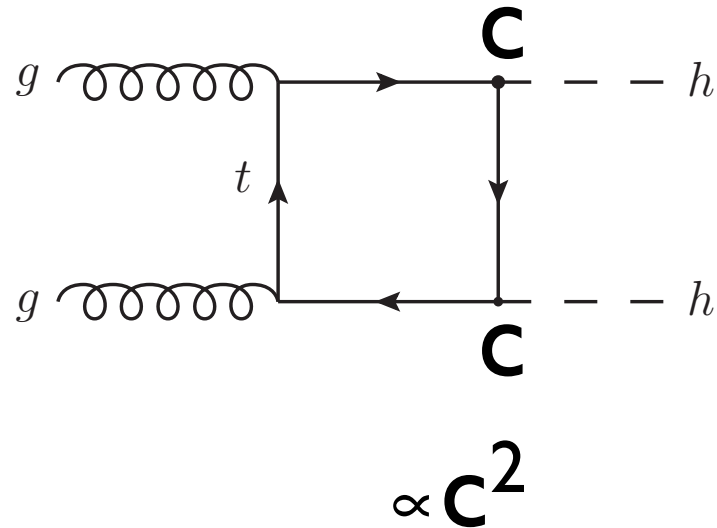
Testing compositeness?

- Clearly many models can produce deviations in c & a (or even in c, a, k_Y, k_g)
- Seeing a deviation in c & a doesn't point to Higgs being “composite”
- To test compositeness one need to check predictions from pGB nature of the Higgs
- $pGB \rightarrow “\exp(i h/v)” \rightarrow$ relations between lowest order in h and higher orders
- need to look (at least) at higgs pair production or WW scattering (and s dependence of σ)...

Testing compositeness?

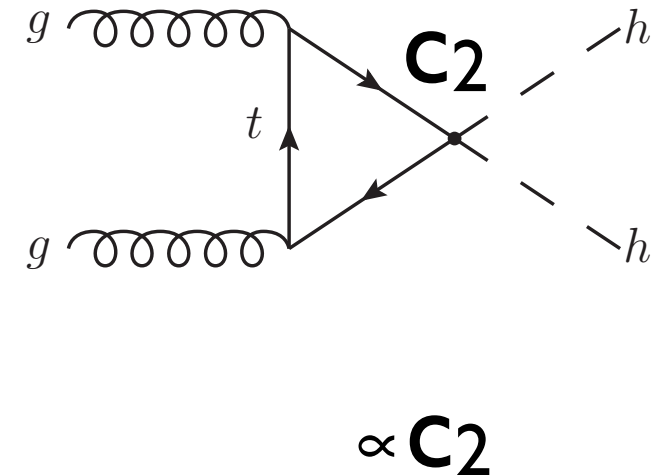
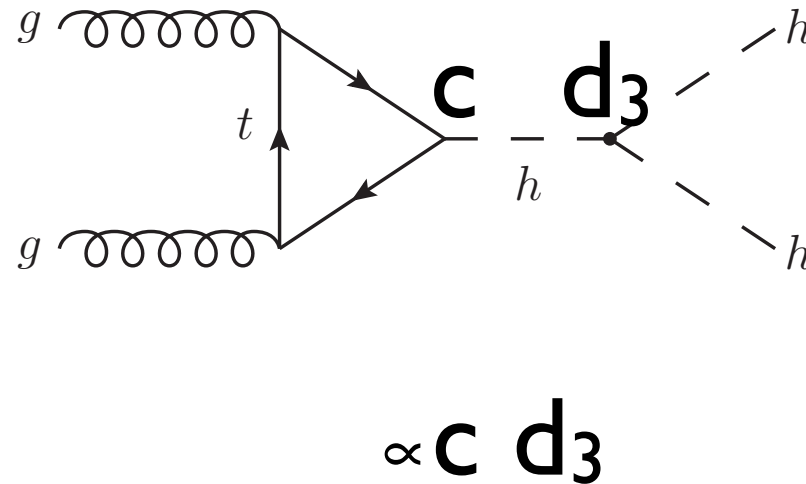
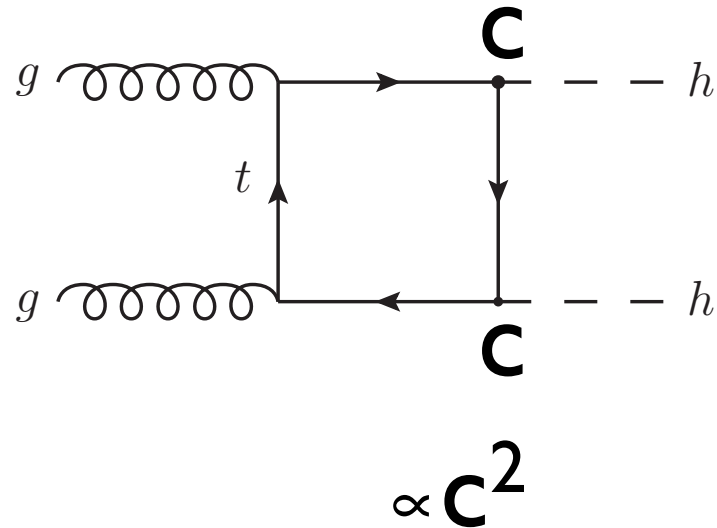


Testing compositeness?



- SILH: relation between c and c_2

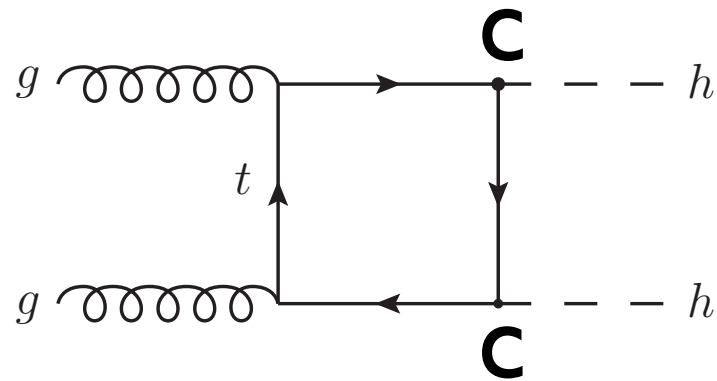
Testing compositeness?



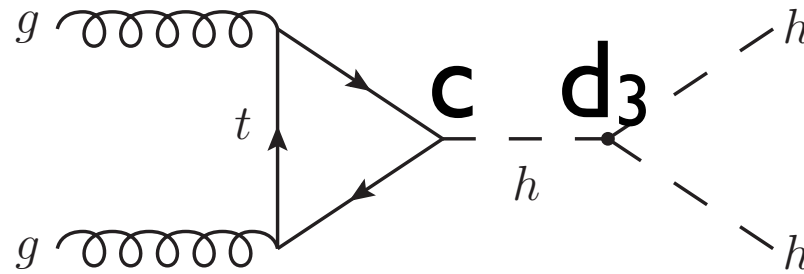
- SILH: relation between c and c_2
- process is very little sensitive to d_3 (good news)

Dib, Rosenfeld, Zerwekh, JHEP 0605 (2006) 074, Grober and Muhlleitner, JHEP 1106 (2011) 020

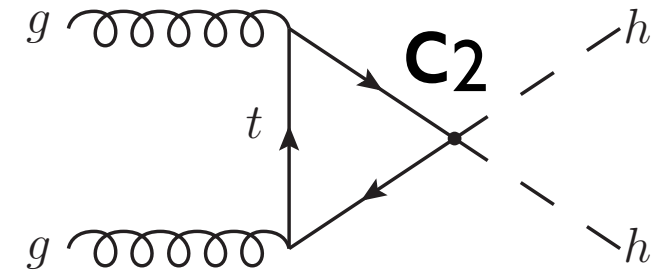
Testing compositeness?



$$\propto c^2$$



$$\propto c d_3$$



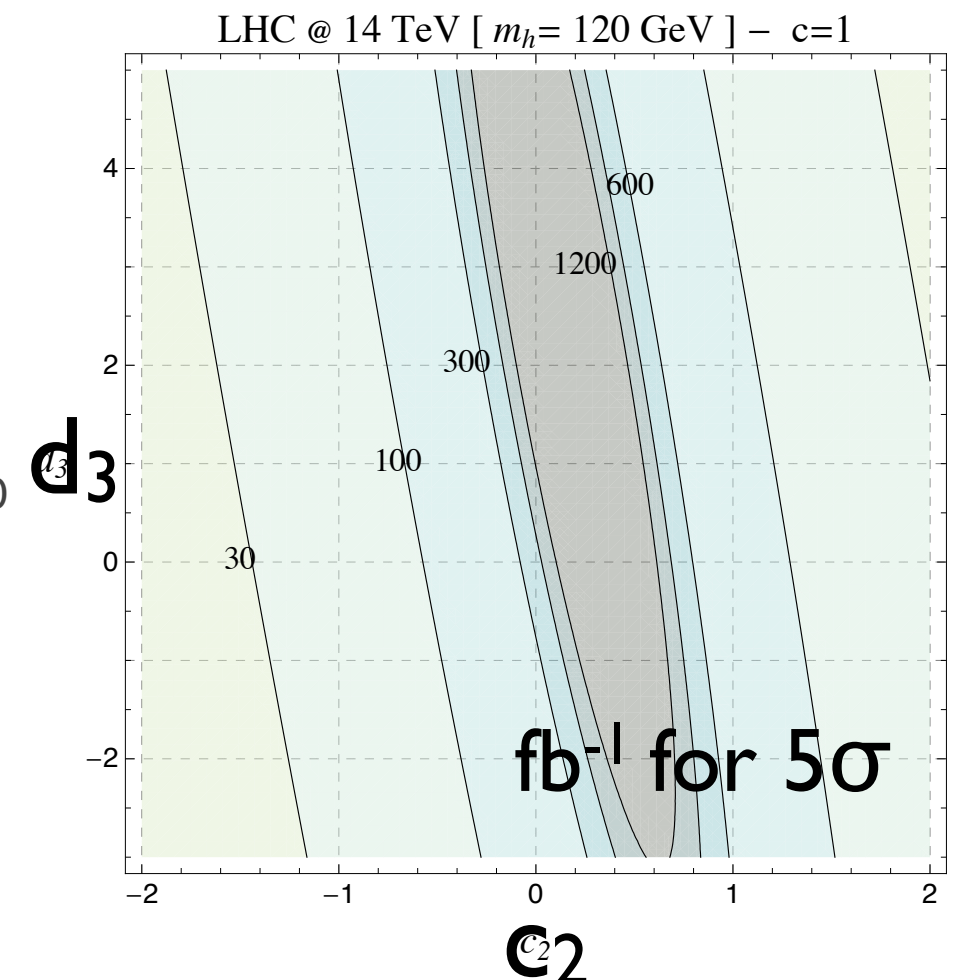
$$\propto c_2$$

- SILH: relation between c and c_2
- process is very little sensitive to d_3 (good news)

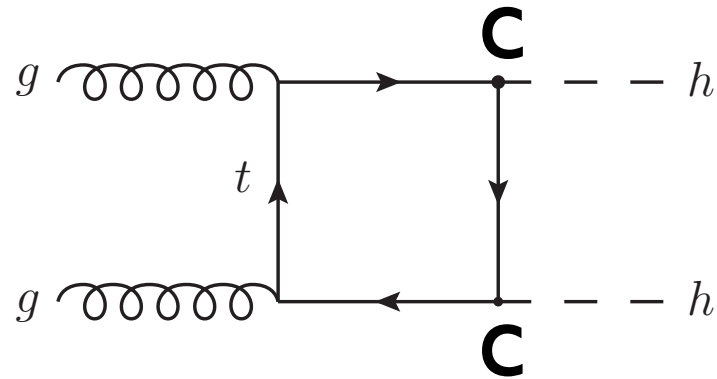
Dib, Rosenfeld, Zerwekh, JHEP 0605 (2006) 074, Grober and Muhlleitner, JHEP 1106 (2011) 020

- very high luminosity required

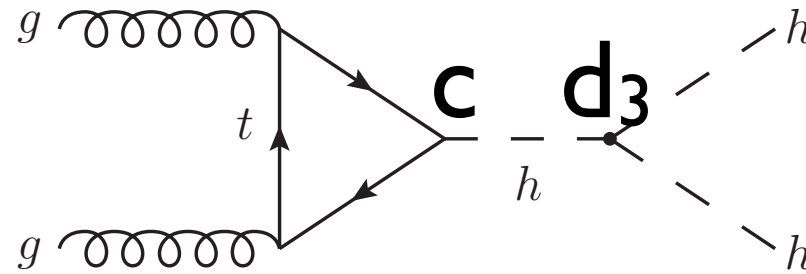
Contino, Ghezzi, Moretti, Panico, Piccinini, Wulzer, arXiv:1205.5444



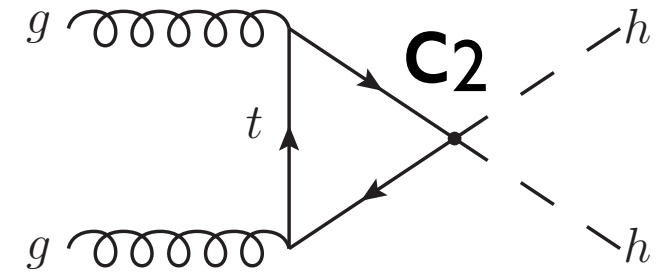
Testing compositeness?



$$\propto c^2$$



$$\propto c d_3$$



$$\propto c_2$$

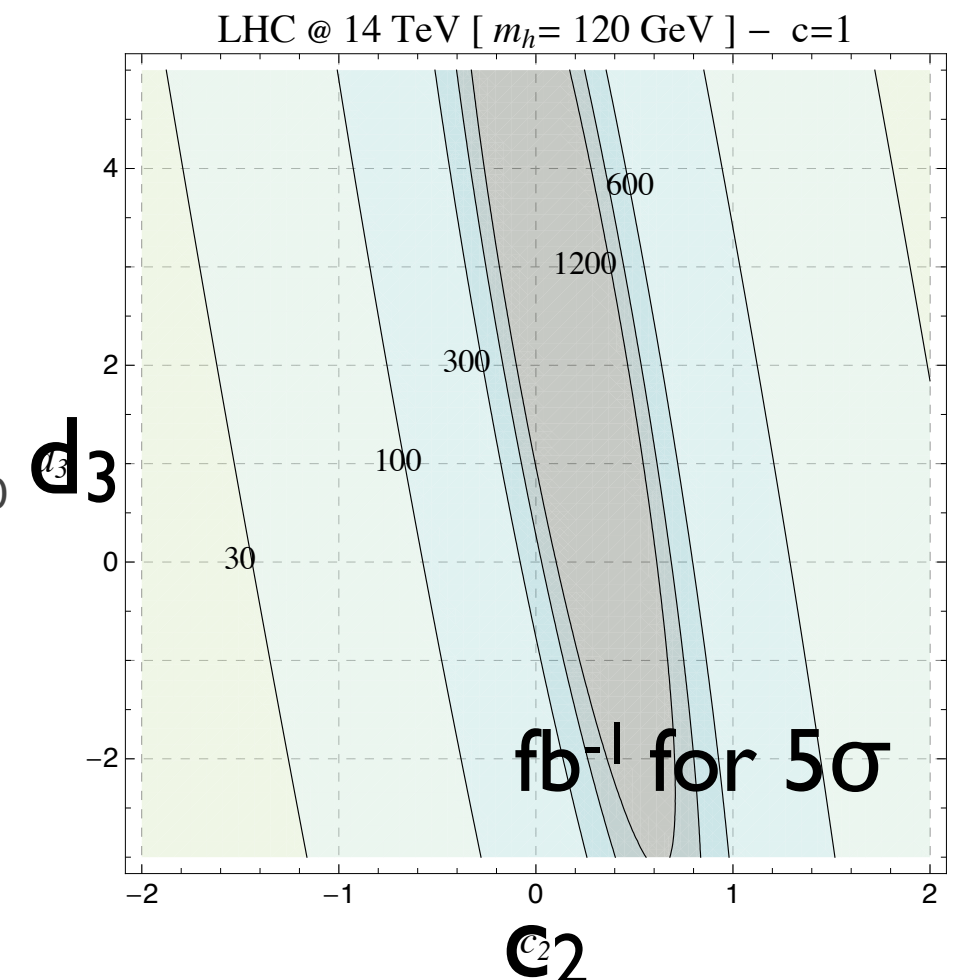
- SILH: relation between c and c_2
- process is very little sensitive to d_3 (good news)

Dib, Rosenfeld, Zerwekh, JHEP 0605 (2006) 074, Grober and Muhlleitner, JHEP 1106 (2011) 020

- very high luminosity required

Contino, Ghezzi, Moretti, Panico, Piccinini, Wulzer, arXiv:1205.5444

Hard to test compositeness
for a long time...



What about “impostors”?

here just meant to be a scalar particle
that couple to WW , ZZ ($\gamma\gamma$, gg) but
has nothing to do with EWSB

if it is an impostor, the Higgs is somewhere
else or it is an Higgsless scenario

How to couple a scalar to VV (not Higgs)

Low, Lykken, Shaughnessy, arXiv:1207.1093

How to couple a scalar to VV (not Higgs)

Low, Lykken, Shaughnessy, arXiv:1207.1093

- Look at $SU(2)_L \times U(1)_Y$ singlets (need to get couplings to both WW and ZZ)

How to couple a scalar to VV (not Higgs)

Low, Lykken, Shaughnessy, arXiv:1207.1093

- Look at $SU(2)_L \times U(1)_Y$ singlets (need to get couplings to both WW and ZZ)
- dim-5 couplings? (i.e. “ $\Phi W_{\mu\nu} W^{\mu\nu}$ ”) \rightarrow generic singlet scalar

$$\begin{aligned}\mathcal{L}_{sV_1V_2} = & \kappa_W \frac{\alpha}{8\pi m_s s_w^2} s W_{\mu\nu}^+ W^{-\mu\nu} + \left(\kappa_W \frac{c_w^2}{s_w^2} + \kappa_B \frac{s_w^2}{c_w^2} \right) \frac{\alpha}{16\pi m_s} s Z_{\mu\nu} Z^{\mu\nu} \\ & + \kappa_g \frac{\alpha_s}{16\pi m_s} s G_{\mu\nu}^a G^{a\mu\nu} + (\kappa_W + \kappa_B) \frac{\alpha}{16\pi m_s} s F_{\mu\nu} F^{\mu\nu} \\ & + \left(\kappa_W \frac{c_w}{s_w} - \kappa_B \frac{s_w}{c_w} \right) \frac{\alpha}{8\pi m_s} s F_{\mu\nu} Z^{\mu\nu},\end{aligned}$$

How to couple a scalar to VV (not Higgs)

Low, Lykken, Shaughnessy, arXiv:1207.1093

- Look at $SU(2)_L \times U(1)_Y$ singlets (need to get couplings to both WW and ZZ)
- dim-5 couplings? (i.e. “ $\Phi W_{\mu\nu} W^{\mu\nu}$ ”) \rightarrow generic singlet scalar

$$\begin{aligned} \mathcal{L}_{sV_1V_2} = & \kappa_W \frac{\alpha}{8\pi m_s s_w^2} s W_{\mu\nu}^+ W^{-\mu\nu} + \left(\kappa_W \frac{c_w^2}{s_w^2} + \kappa_B \frac{s_w^2}{c_w^2} \right) \frac{\alpha}{16\pi m_s} s Z_{\mu\nu} Z^{\mu\nu} \\ & + \kappa_g \frac{\alpha_s}{16\pi m_s} s G_{\mu\nu}^a G^{a\mu\nu} + (\kappa_W + \kappa_B) \frac{\alpha}{16\pi m_s} s F_{\mu\nu} F^{\mu\nu} \\ & + \left(\kappa_W \frac{c_w}{s_w} - \kappa_B \frac{s_w}{c_w} \right) \frac{\alpha}{8\pi m_s} s F_{\mu\nu} Z^{\mu\nu}, \end{aligned}$$

- dim-4 couplings? (i.e. “ $\Phi m_W W_\mu W^\mu$ ”) \rightarrow “dilaton” or “radion”

$$\begin{aligned} \mathcal{L} = c_\phi \frac{\phi}{v} & \left(2m_W^2 W_\mu^+ W_\mu^- + m_Z^2 Z_\mu Z_\mu - \sum_{f \in SM} m_f \bar{f} f \right) \sim \phi T_\mu^\mu \\ & c_g \frac{\alpha_s}{12\pi} \frac{\phi}{v} G_{\mu\nu}^a G^{a\mu\nu} + c_\gamma \frac{\alpha}{\pi} \frac{\phi}{v} F_{\mu\nu} F^{\mu\nu} \end{aligned}$$

c_g, c_γ different from SM due to conformal properties of UV

How to couple a scalar to VV (not Higgs)

How to couple a scalar to VV (not Higgs)

- Dilaton: simplifying assumption: UV is conformal invariant and SM is part of the CFT: $\beta_{UV} + \beta_{IR(SM)} = 0$
→ additional contribution proportional to $-\beta_{SM}$

$$\delta c_g = \frac{21}{2} c_\phi + (c_\phi - 1) A_f(\tau_t), \quad \delta c_\gamma = -\frac{11}{24} c_\phi + (c_\phi - 1) A_f(\tau_t).$$

How to couple a scalar to VV (not Higgs)

- Dilaton: simplifying assumption: UV is conformal invariant and SM is part of the CFT: $\beta_{UV} + \beta_{IR(SM)} = 0$
→ additional contribution proportional to $-\beta_{SM}$

$$\delta c_g = \frac{21}{2} c_\phi + (c_\phi - 1) A_f(\tau_t), \quad \delta c_\gamma = -\frac{11}{24} c_\phi + (c_\phi - 1) A_f(\tau_t).$$

- Radion: mode associated to the stabilization of the 5dim bulk size (distance between the two branes). In Randall Sundrum models the bulk is supposed to be dual to a conformal field theory in 4D, with the IR providing a dynamical breaking of conformal invariance: radion is the “dilaton” of the CFT

How to couple a scalar to VV (not Higgs)

How to couple a scalar to VV (not Higgs)

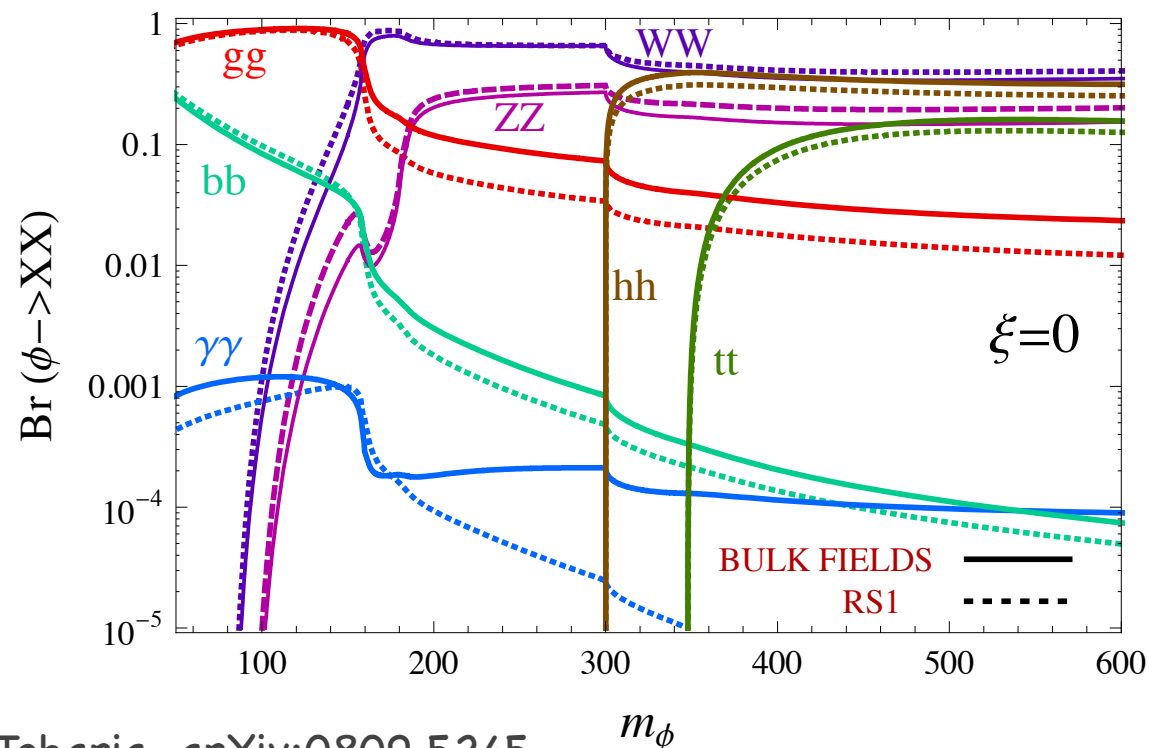
- Both dilaton and radion couplings depend on anomalous dimensions of fields above the weak scale (for radion bulk masses and profiles) \rightarrow some model dependence

How to couple a scalar to VV (not Higgs)

- Both dilaton and radion couplings depend on anomalous dimensions of fields above the weak scale (for radion bulk masses and profiles) \rightarrow some model dependence
- General features: $B(r \rightarrow gg)$ dominates. $B(r \rightarrow ZZ^*)/B(r \rightarrow \gamma\gamma)$ is model dependent

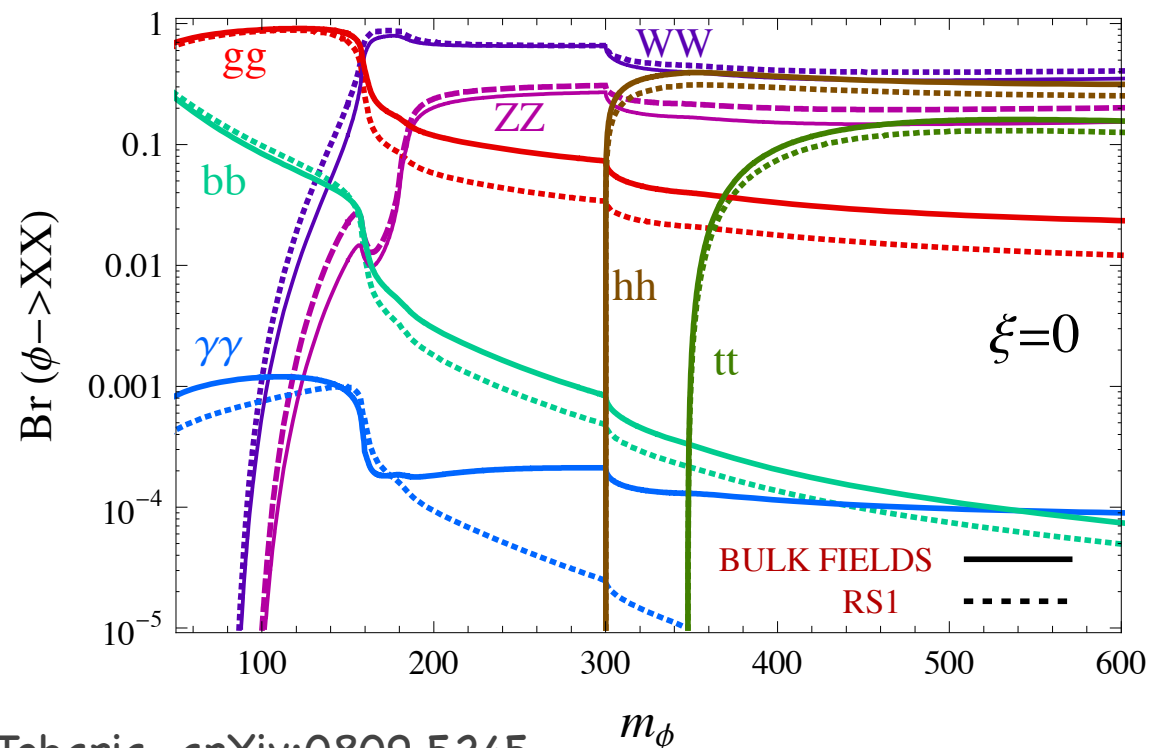
How to couple a scalar to VV (not Higgs)

- Both dilaton and radion couplings depend on anomalous dimensions of fields above the weak scale (for radion bulk masses and profiles) \rightarrow some model dependence
- General features: $B(r \rightarrow gg)$ dominates. $B(r \rightarrow ZZ^*)/B(r \rightarrow \gamma\gamma)$ is model dependent

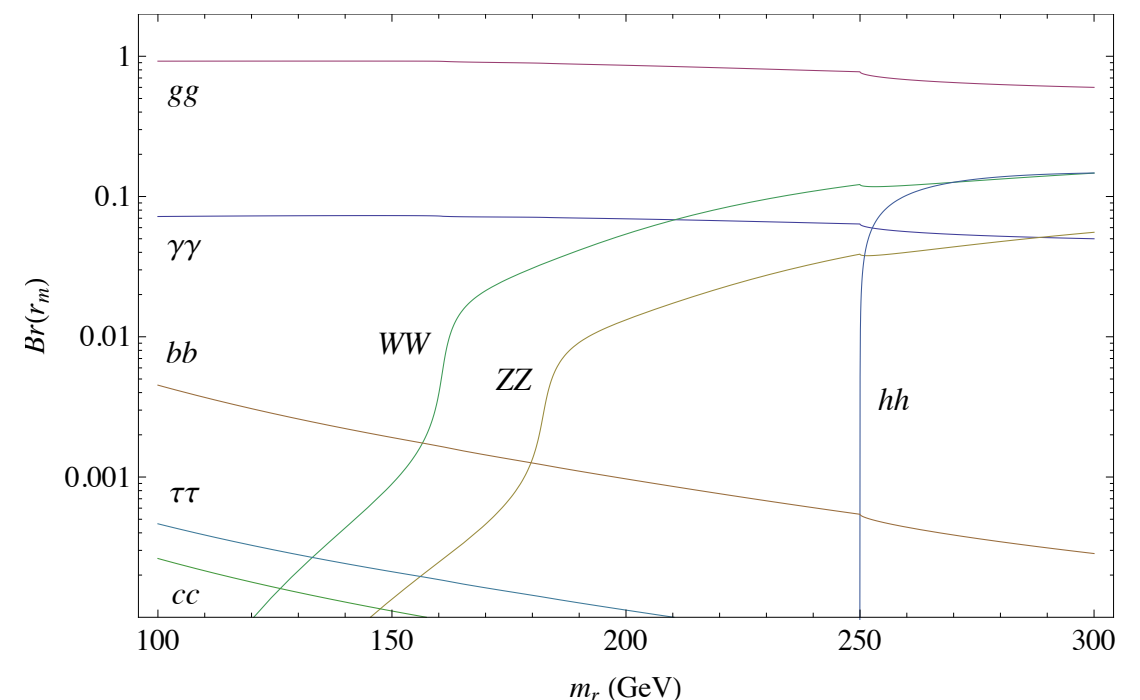


How to couple a scalar to VV (not Higgs)

- Both dilaton and radion couplings depend on anomalous dimensions of fields above the weak scale (for radion bulk masses and profiles) \rightarrow some model dependence
- General features: $B(r \rightarrow gg)$ dominates. $B(r \rightarrow ZZ^*)/B(r \rightarrow \gamma\gamma)$ is model dependent



Toharia, arXiv:0809.5245



Cox, Gherghetta, arXiv:1203.5870

Looking at the data

The “minimal” dilaton is still allowed, although not preferred

Looking at the data

The “minimal” dilaton is still allowed, although not preferred

also allowing higgs-dilaton mixing:

$$\begin{aligned}\delta c_g &= \frac{21}{2} c_\phi \sin \alpha + (c_\phi \sin \alpha + \cos \alpha - 1) A_f(\tau_t), \\ \delta c_\gamma &= -\frac{11}{24} c_\phi \sin \alpha + (c_\phi \sin \alpha + \cos \alpha - 1) A_f(\tau_t).\end{aligned}$$

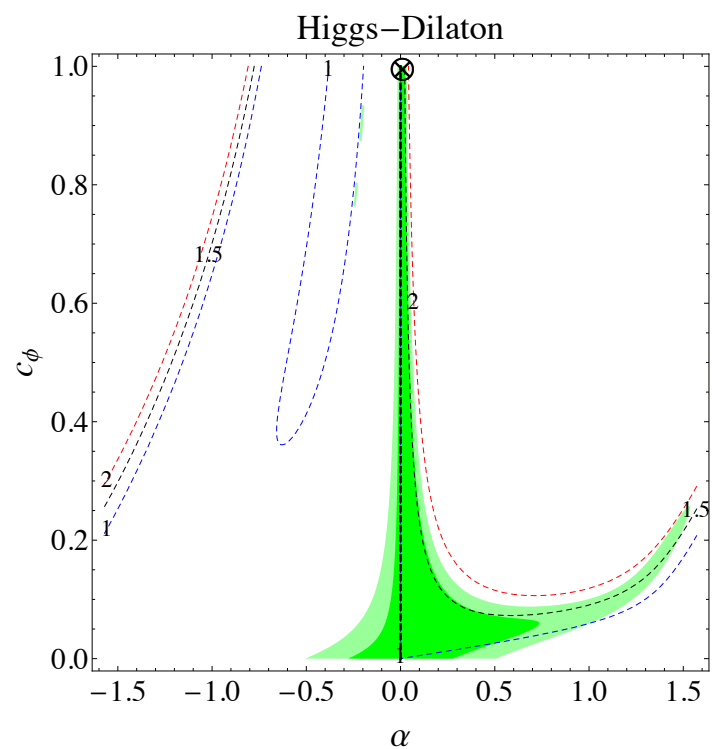
Looking at the data

The “minimal” dilaton is still allowed, although not preferred

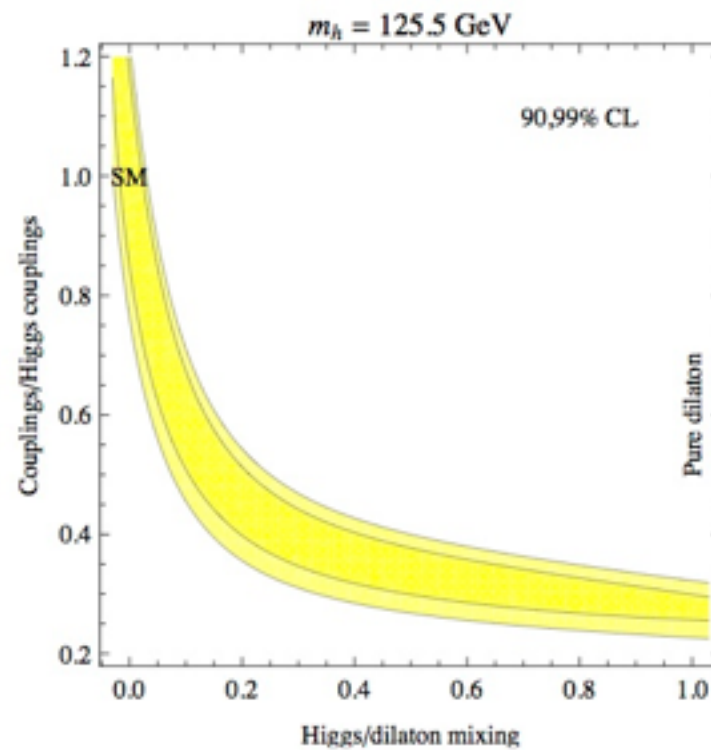
also allowing higgs-dilaton mixing:

$$\delta c_g = \frac{21}{2} c_\phi \sin \alpha + (c_\phi \sin \alpha + \cos \alpha - 1) A_f(\tau_t),$$

$$\delta c_\gamma = -\frac{11}{24} c_\phi \sin \alpha + (c_\phi \sin \alpha + \cos \alpha - 1) A_f(\tau_t).$$



Carmi, Falkowski, Kuflik,
Volansky, Zupan arXiv:1207.1718



Giardino, Kannike, Raidal,
Strumia arXiv:1207.1347

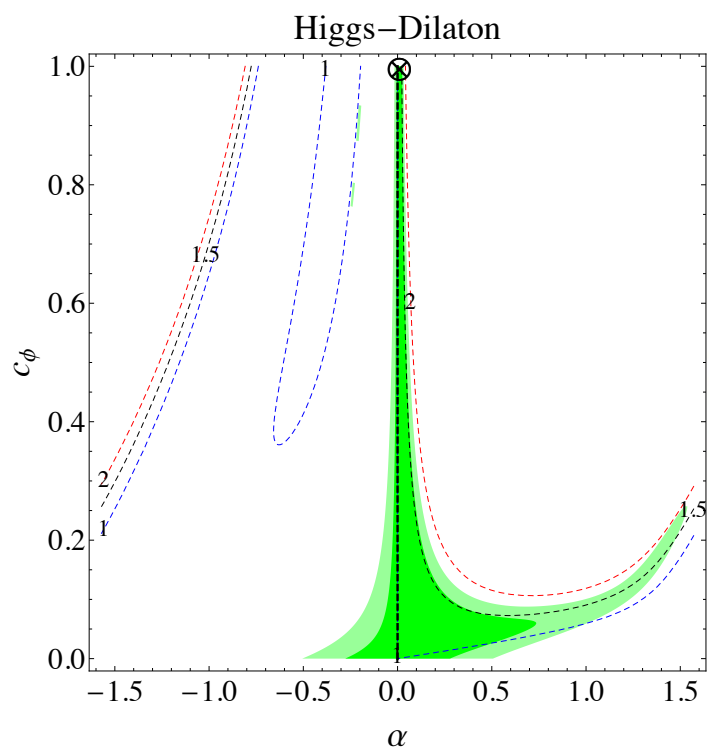
Looking at the data

The “minimal” dilaton is still allowed, although not preferred

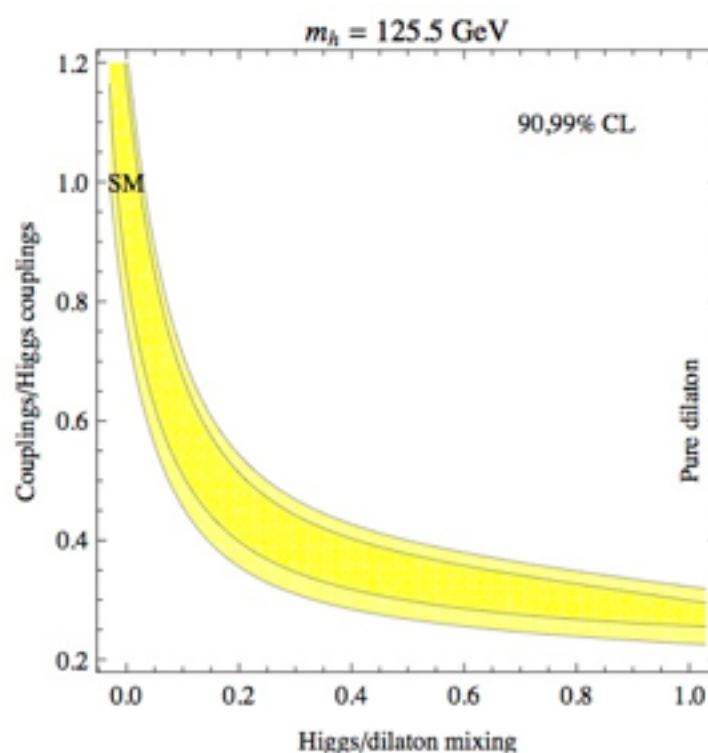
also allowing higgs-dilaton mixing:

$$\delta c_g = \frac{21}{2} c_\phi \sin \alpha + (c_\phi \sin \alpha + \cos \alpha - 1) A_f(\tau_t),$$

$$\delta c_\gamma = -\frac{11}{24} c_\phi \sin \alpha + (c_\phi \sin \alpha + \cos \alpha - 1) A_f(\tau_t).$$



Carmi, Falkowski, Kuflik,
Volansky, Zupan arXiv:1207.1718



Giardino, Kannike, Raidal,
Strumia arXiv:1207.1347

non-dilaton singlet “impostor”
is excluded by $\gamma Z/ZZ$

Low, Lykken, Shaughnessy, arXiv:1207.1093

