

A complex visualization of particle detector data, likely from ATLAS, showing a dense network of tracks and energy deposits. The tracks are represented by thin, overlapping lines in shades of yellow and orange, radiating from a central point. The background is a deep blue, with numerous small white and red dots scattered throughout, representing individual particles or energy deposits. The overall appearance is that of a high-energy collision event.

Higgs Bosons and b Quarks

October, 2007

SLAC ATLAS Forum

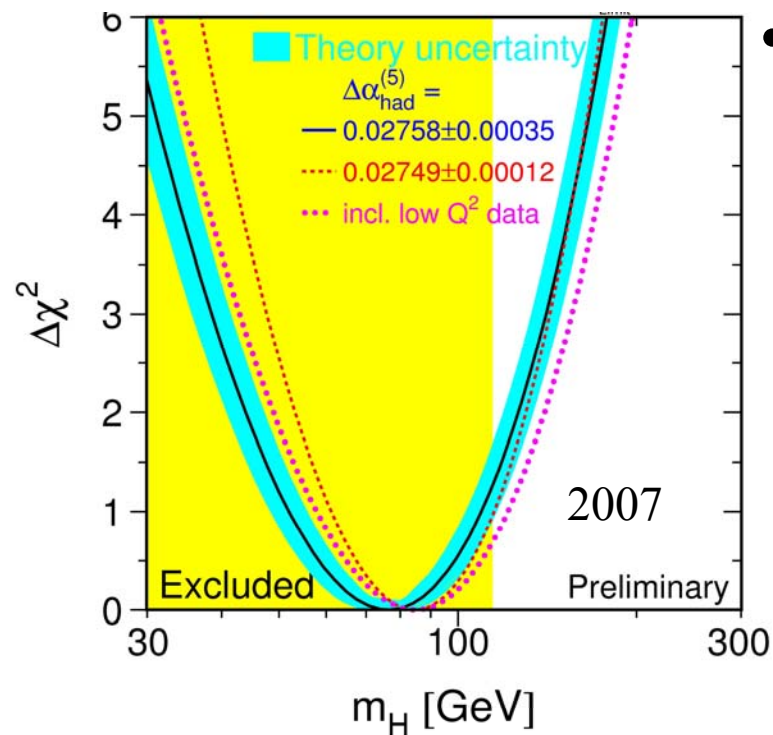
Sally Dawson

Laura Reina, Chris Jackson, Doreen Wackerath

Plan:

- Lightning review of SM Higgs physics
 - Discussion of $gg \rightarrow b\bar{b}h$ vs $bg \rightarrow bh$
 - Emphasize understanding of theoretical assumptions
- MSSM results for $bg \rightarrow bh$
 - Status of current (Summer, 2007) limits
- Effects of squark/gluino loops on $bg \rightarrow bh$
 - Why are these effects interesting?

Precision measurements limit Higgs Mass

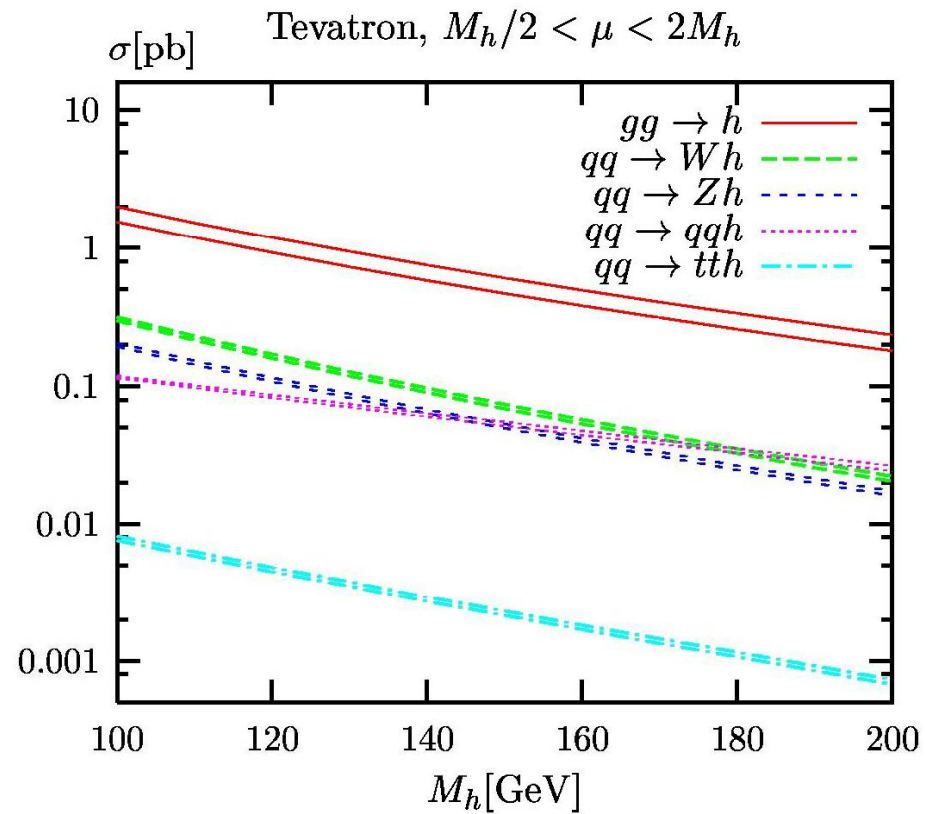


- LEP EWWG (July, 2007):

- $M_t = 170.9 \pm 1.8$ GeV
- $M_h = 76^{+36}_{-24}$ GeV
- $M_h < 144$ GeV (one-sided 95% cl)
- $M_h < 182$ GeV (Precision measurements plus direct search limit)

Best fit in region excluded from direct searches

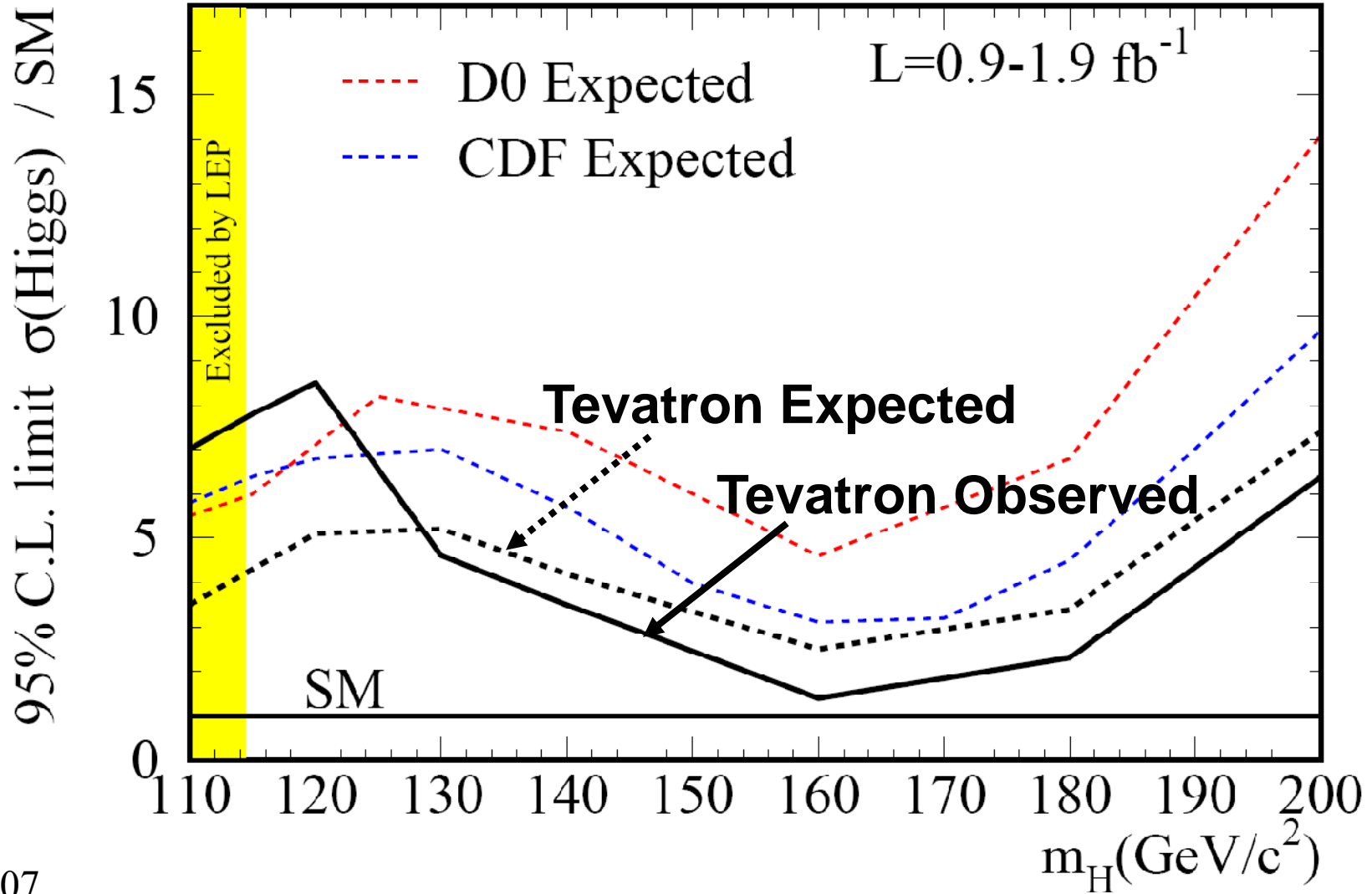
Producing the Higgs at the Tevatron



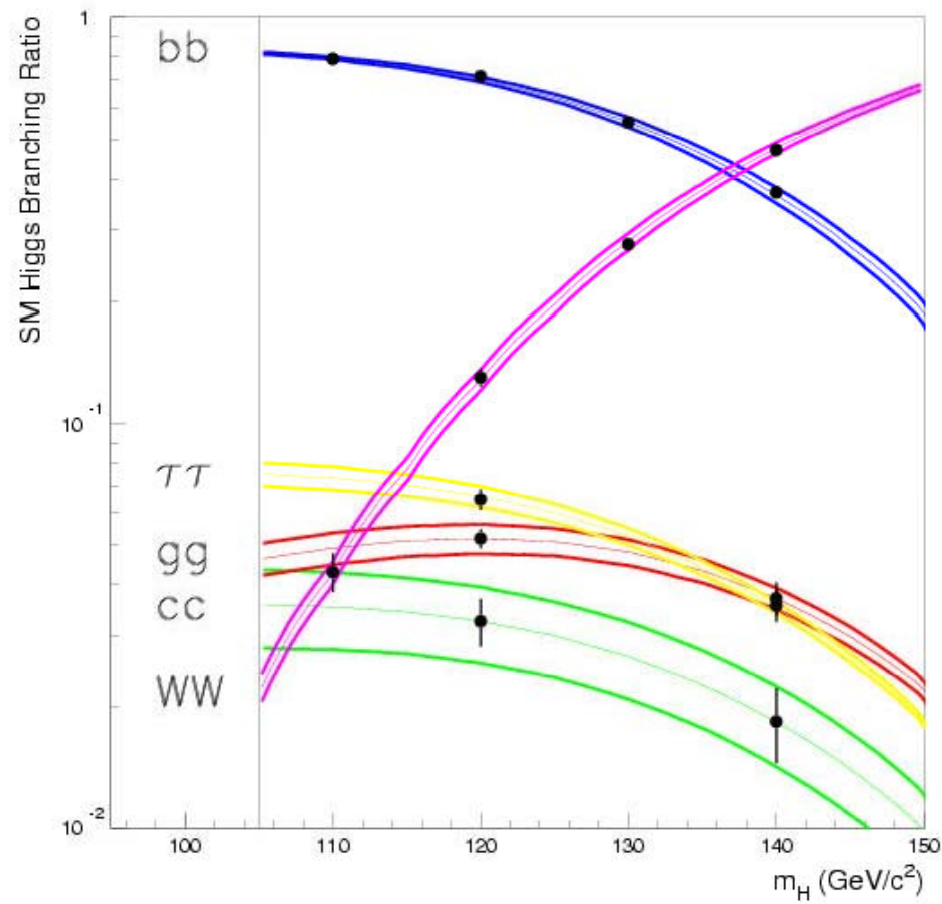
NNLO or NLO rates

$M_h/2 < \mu < M_h/4$

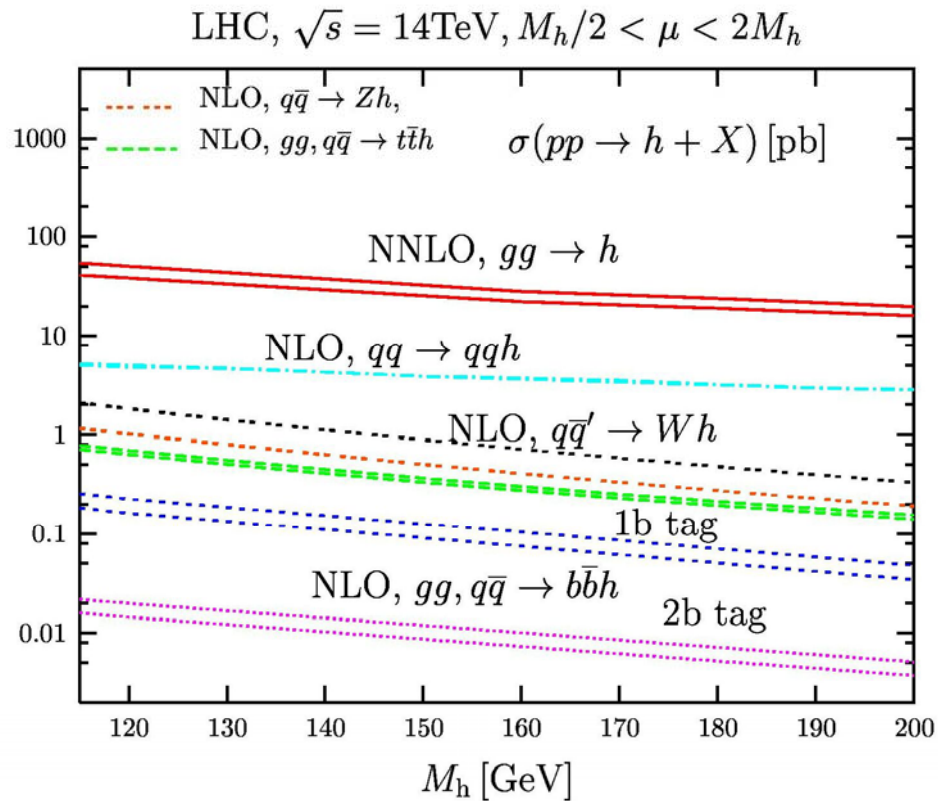
Tevatron Run II Preliminary



Limits understood from Branching Ratios



SM Production Mechanisms at LHC

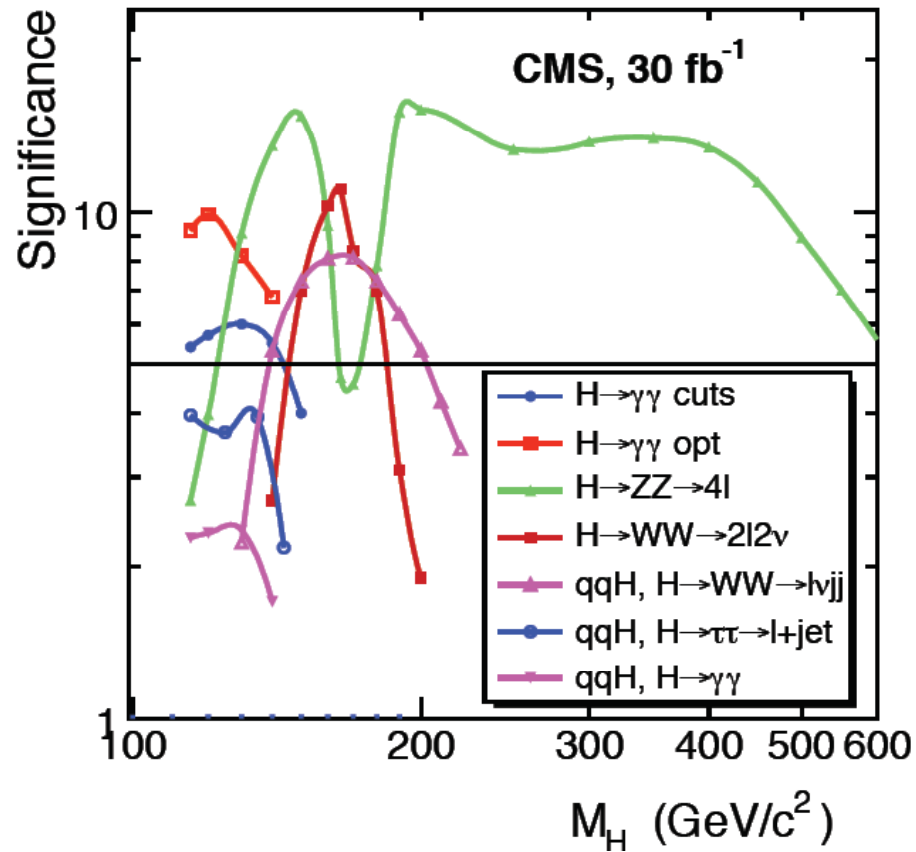


Bands show scale dependence

All important channels
calculated to NLO or NNLO

Production with b's
very small in SM

SM Higgs, CMS 2007

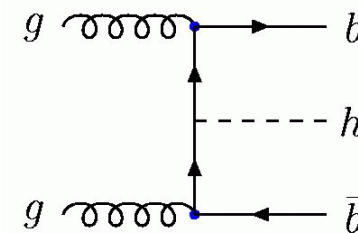


Includes radiative corrections

Higgs + b's aren't discovery mode for SM Higgs

$$pp \rightarrow b\bar{b}h$$

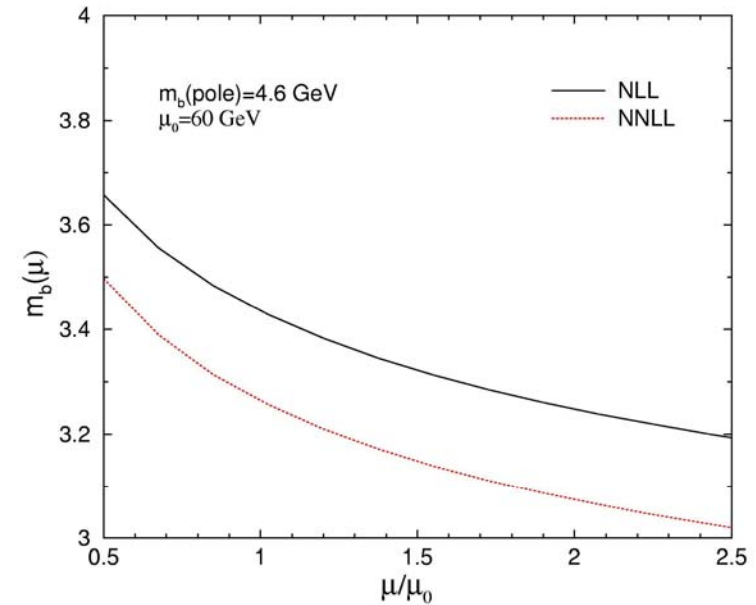
- Why is $b\bar{b}h$ interesting?
 - Direct measurement of b quark Yukawa coupling
(enhanced in MSSM at large $\tan \beta$)
 - Higgs discovery mode in SUSY models at large $\tan \beta$
 - Theoretical questions about b quark parton distribution functions (PDFs)
- Why do NLO corrections?
 - Improved theoretical reliability
 - Often find large numerical results



Which b mass?

- $\sigma(\text{bbh}) \approx (m_b/v)^2$
- Pole mass (from Υ decays):
 $m_b = 4.62 \text{ GeV}$
- $\overline{\text{MS}}$ bar mass:

$$\overline{m}_b(\mu) = m_b \left[1 - \frac{\alpha_s}{3\pi} \left(4 + 3 \ln \left\{ \frac{\mu^2}{m_b^2} \right\} \right) \right]$$



*Makes a big numerical difference
which b mass you use*

Use \overline{MS} Renormalization

- Compute the $O(\alpha_s)$ corrections:

$$\Gamma_1(h \rightarrow b\bar{b}) = \frac{3M_h}{8\pi} \left(\frac{m_b}{v}\right)^2 \left\{ 1 + \frac{2\alpha_s}{3\pi} \left[\frac{9}{2} - 3 \log\left(\frac{M_h^2}{m_b^2}\right) \right] \right\}$$

- Define the running b mass

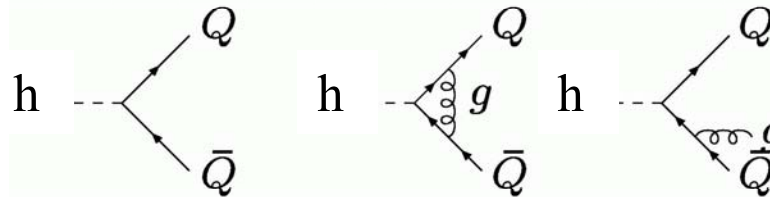
$$\bar{m}_b(\mu) = m_b \left[1 - \frac{\alpha_s}{3\pi} \left(4 + 3 \ln\left(\frac{\mu^2}{m_b^2}\right) \right) \right]$$

- Large logarithms absorbed to 2-loops

$$\Gamma(h \rightarrow b\bar{b}) = \Gamma_0 \left\{ 1 + 5.67 \frac{\alpha_s(M_h)}{\pi} + (36 - 1.4n_f) \frac{\alpha_s(M_h)^2}{\pi^2} \right\}$$

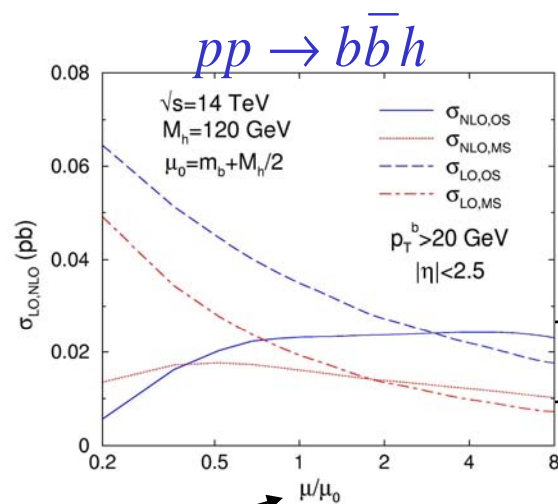
$$\Gamma_0(h \rightarrow b\bar{b}) = \frac{3M_h}{8\pi} \left(\frac{\bar{m}_b(M_h)}{v} \right)^2$$

Lore: \overline{MS} is best!



Scale and Scheme Dependence at NLO

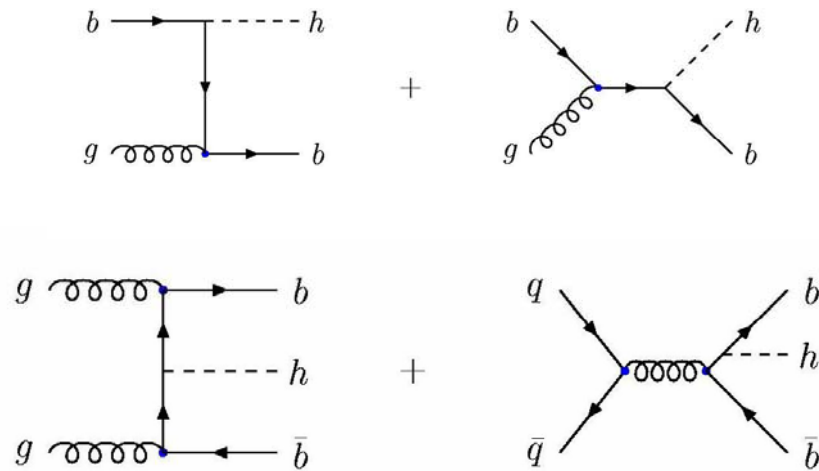
- NLO calculations improve scale dependence
 - Scale dependence enters in running of $\alpha_s(\mu)$ and PDFs, $g(\mu)$, as well as $\alpha_s^3 \log(\mu)$ contributions
 - Formally, scale dependence is $O(\alpha_s^4)$ but may be numerically large



- Large remaining scale/scheme dependence between OS and $\overline{\text{MS}}$ at NLO
- Effect $\approx 10\text{-}20\%$

Scale dependence

What is the dominant process for Higgs + b Production?



➤ Answer depends on whether you tag outgoing b's

➤ Is there double counting when including b initial state?

The b quark as a parton

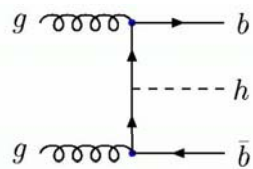
Absorb collinear logs in b quark distributions

$$b(x, \mu) = \frac{\alpha_s}{2\pi} \ln\left(\frac{\mu^2}{m_b^2}\right) \int_x^1 \frac{dz}{z} P_{bg}\left(\frac{x}{z}\right) g(z, \mu)$$

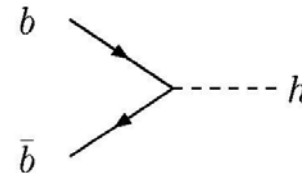
- Altarelli-Parisi evolution of PDFs sums $\alpha_s^n \ln^n(\mu^2/m_b^2)$
- b quark PDF $\approx \alpha_s \ln(\mu^2/m_b^2)$ relative to gluon PDF

Two Schemes for PDFs:

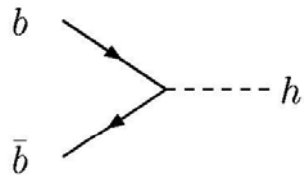
- 4 flavor number scheme (also called fixed flavor number scheme)
 - No b quarks in initial state
 - Lowest order process involving Higgs and b's is $gg \rightarrow b\bar{b}h$
- 5 flavor number scheme (also called variable flavor number scheme)
 - Define b quark PDFs (absorbs large logarithms)
 - Higgs produced with no p_T at lowest order ($b\bar{b} \rightarrow h$)
 - Higgs p_T generated at higher orders in expansion



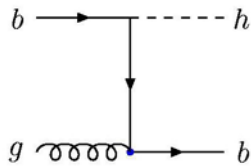
VS



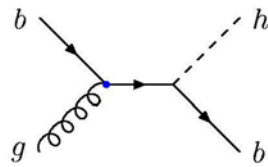
Counting Rules with b PDFs: Reordering of perturbation expansion



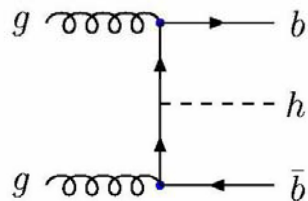
$$(\alpha_s \ln(M_h^2/m_b^2))^2 \approx .4$$



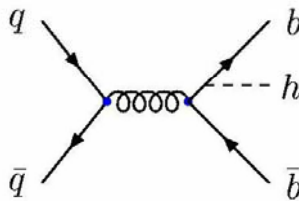
+



$$\alpha_s^2 \ln(M_h^2/m_b^2) \approx .06$$



+



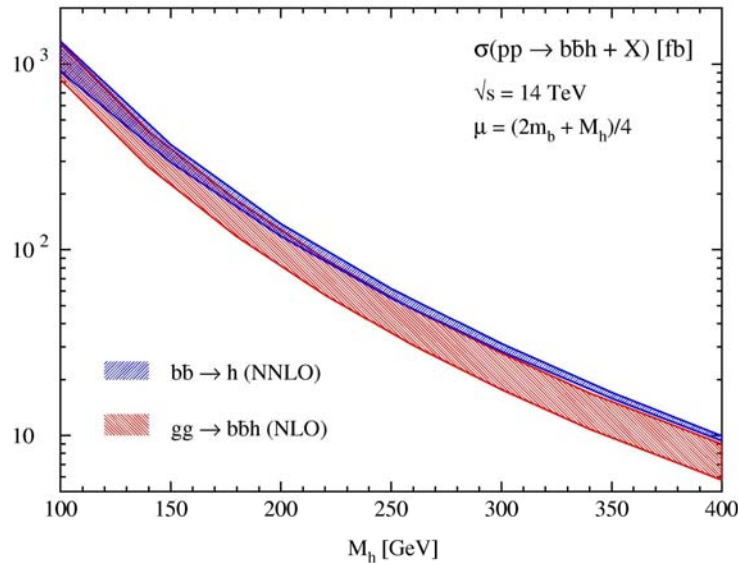
$$\alpha_s^2 \approx .01$$

Re-ordering of Perturbation Theory

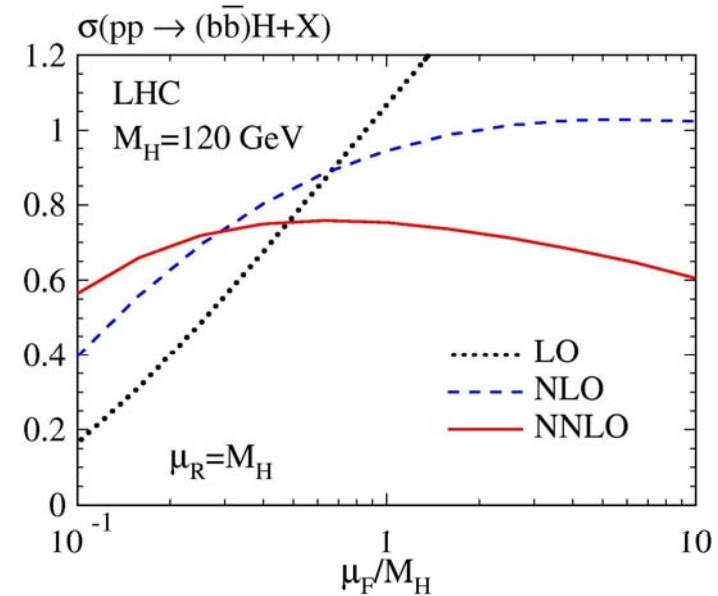
- 0 b tag process in 5FNS:
 - LO: $b\bar{b} \rightarrow h$ $O(\alpha_s^2 \Lambda_b^2)$
 - NLO: Virtual+real corrections $O(\alpha_s^3 \Lambda_b^2)$
 - NLO: $bg \rightarrow bh$ $O(\alpha_s^2 \Lambda_b)$, correction of $O(1/\Lambda_b)$ to tree level
 - NNLO: $gg \rightarrow b\bar{b}h$ $O(\alpha_s^2)$, correction of $O(1/\Lambda_b^2)$ to tree level
- 1 b tag process in 5FNS:
 - LO process is $bg \rightarrow bh$: Tree level, $O(\alpha_s^2 \Lambda_b)$
 - NLO includes new subprocess: $gg \rightarrow b\bar{b}h$, $O(1/\Lambda_b)$ correction to LO

$$\Lambda_b = \log(M_h^2/m_b^2)$$

Inclusive Cross Section for $b\bar{b} \rightarrow h$: 0 b tags



$b\bar{b} \rightarrow h$ vs $gg \rightarrow b\bar{b}h$

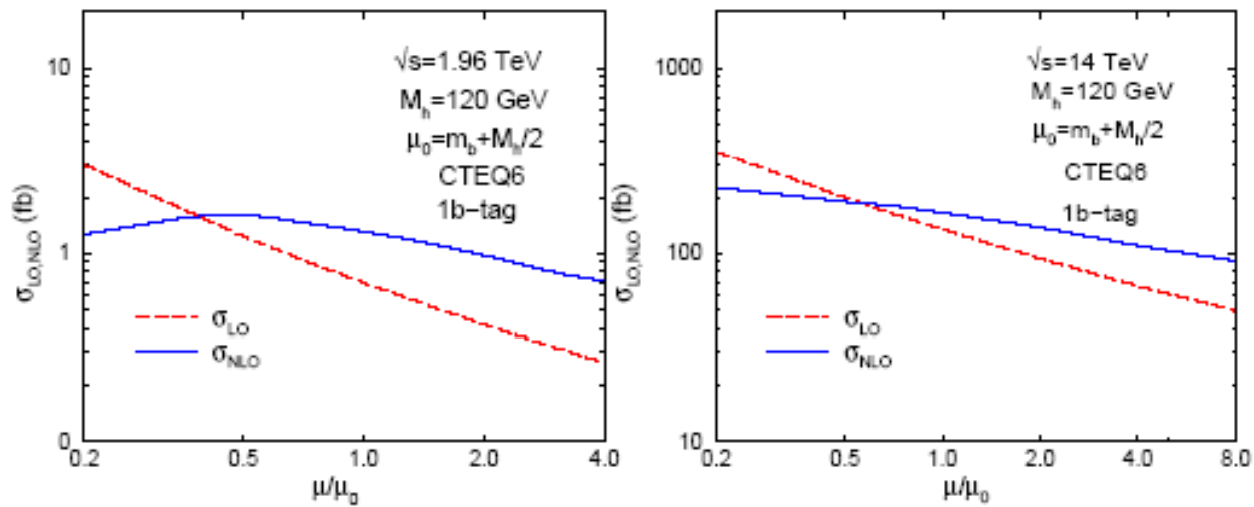


Almost no scale dependence at NNLO

It really doesn't matter which PDF scheme you use

What if only 1 b is tagged?

4FNS: $gg \rightarrow b\bar{b}h$



$$p_T^{b} > 20 \text{ GeV}$$

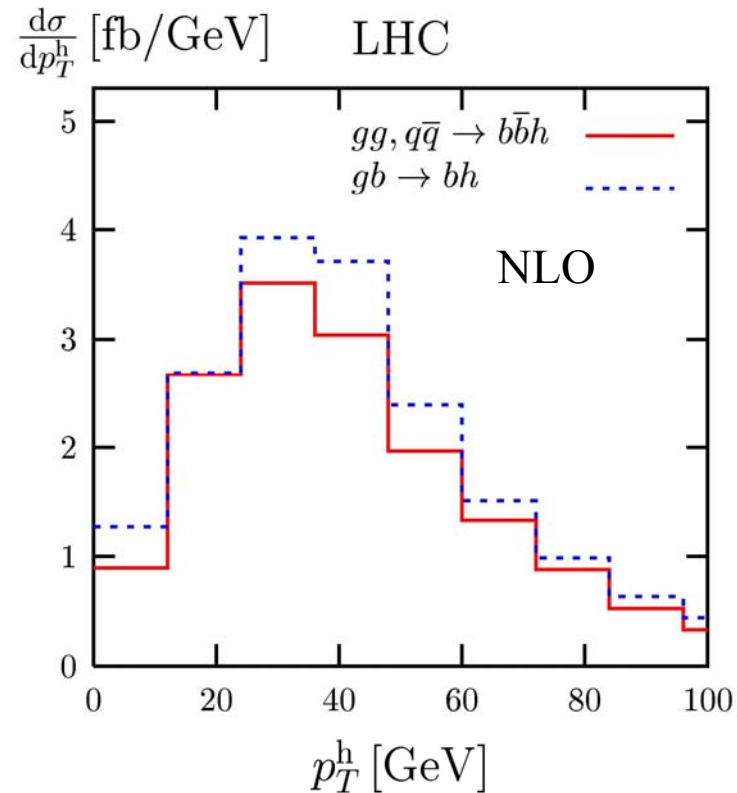
$$|\eta| < 2 \text{ (Tevatron), } 2.5 \text{ (LHC)}$$

$$\Delta R > 0.4$$

What about distributions?

Compare 4 and 5 Flavor Number PDF Schemes

Higgs plus single b at LHC:

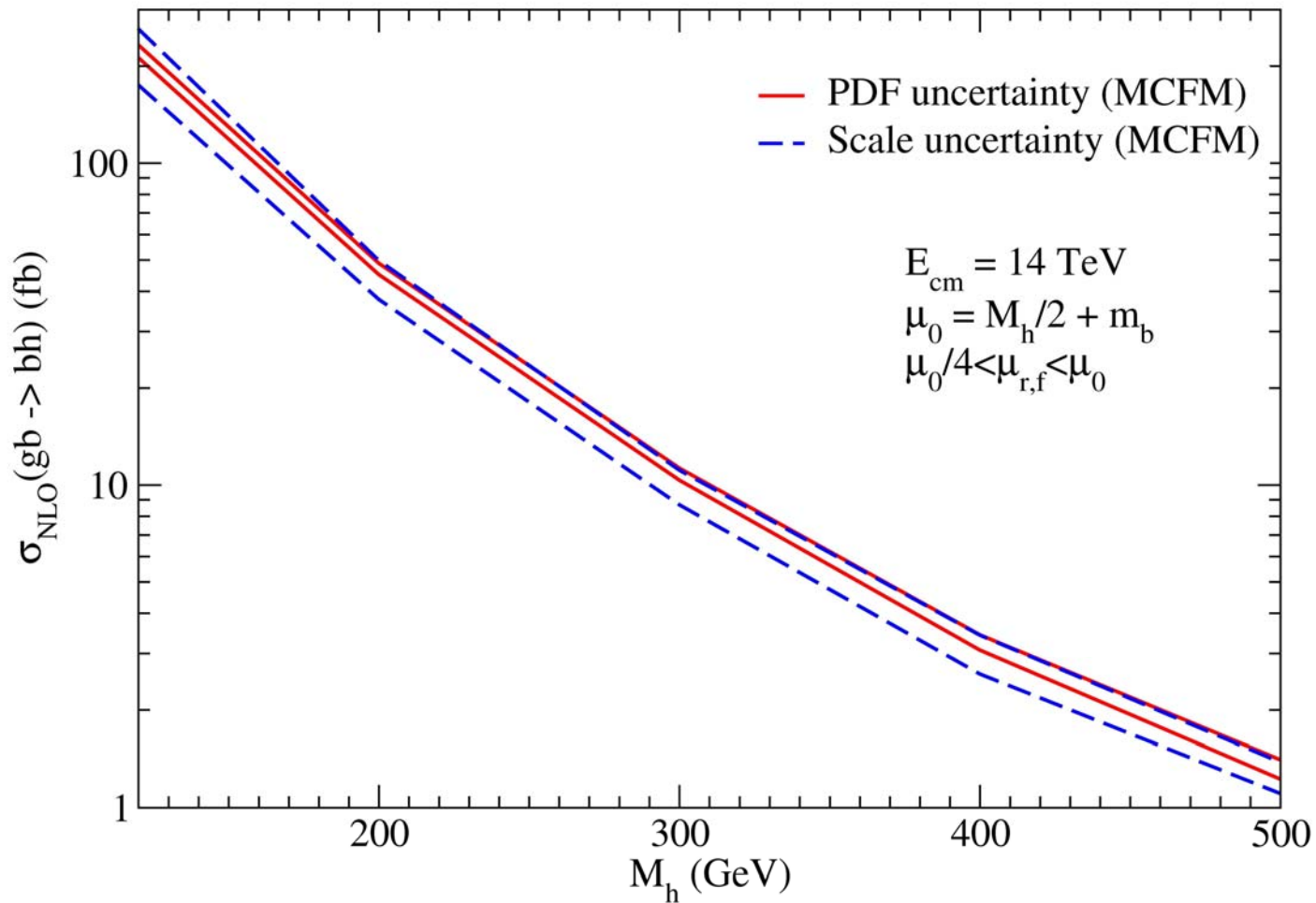


$p_T^{b} > 20$ GeV

$|\eta| < 2.5$

$\Delta R > 0.4$

Good Theoretical Understanding of Uncertainties



Higgs in the MSSM

- MSSM has 2 Higgs doublets: H_d and H_u

$$H_d = \begin{pmatrix} \phi_d^+ \\ \phi_d^0 \end{pmatrix}$$

$$H_u = \begin{pmatrix} \phi_u^0 \\ -\phi_u^- \end{pmatrix}$$

$$\phi_d^0 = \frac{1}{\sqrt{2}}(v_1 + h_d^0)$$

$$\phi_u^0 = \frac{1}{\sqrt{2}}(v_2 + h_u^0)$$

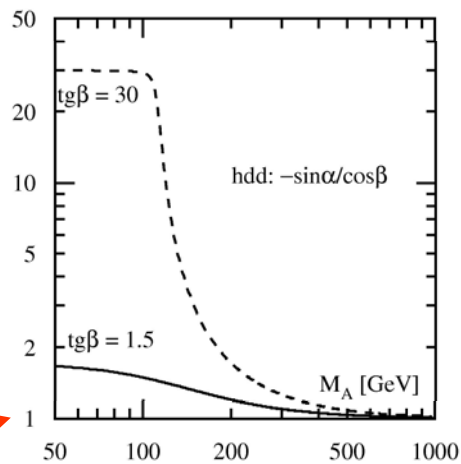
- Physical CP-Even Higgs bosons

$$\begin{pmatrix} h^0 \\ H^0 \end{pmatrix} = \begin{pmatrix} c_\alpha & -s_\alpha \\ s_\alpha & c_\alpha \end{pmatrix} \begin{pmatrix} h_u^0 \\ h_d^0 \end{pmatrix}$$

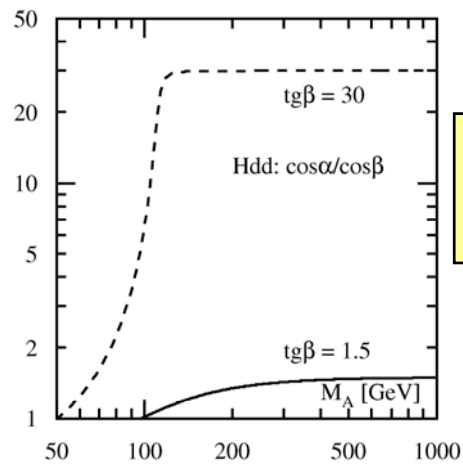
- Pseudoscalar, A^0 , and two charged Higgs, H^\pm

Higgs Couplings very different in MSSM

Light Higgs

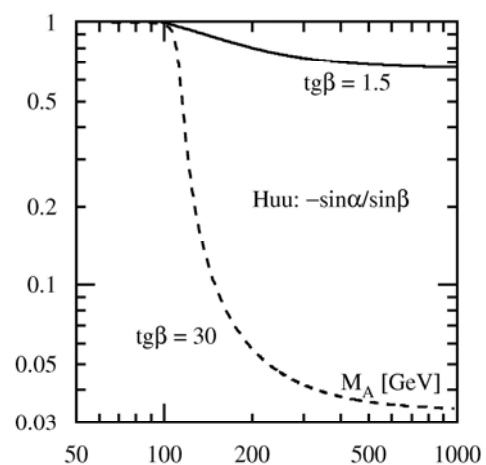
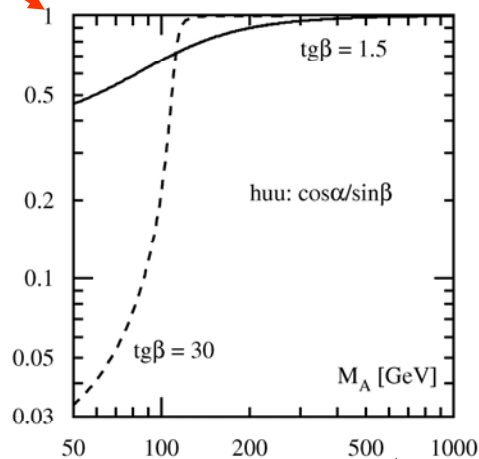
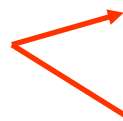


Heavy Higgs



Couplings to d, s, b enhanced at large $\tan\beta$

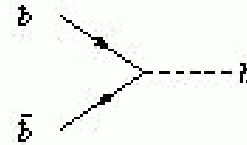
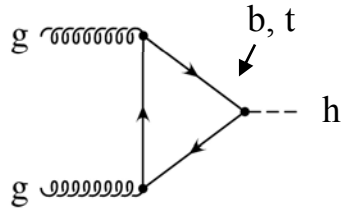
SM



Couplings to u, c, t suppressed at large $\tan\beta$

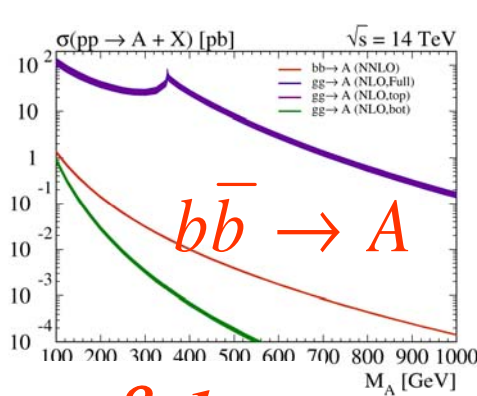
Decoupling limit

Large $\tan \beta$ Changes Relative Importance of Production Modes

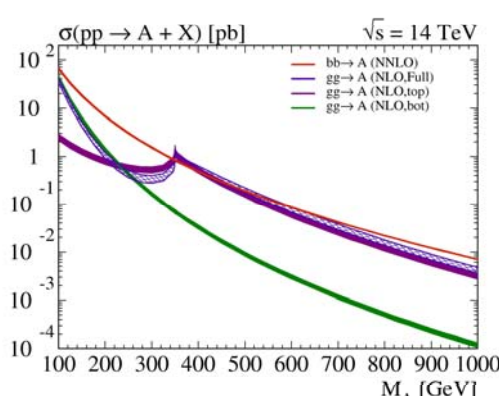


$$\sigma_{gg} = \frac{1}{M_h^2} \left(c_1 \cot^2 \beta + c_2 \frac{m_b^2}{M_h^2} + c_3 \frac{m_b^4}{M_h^4} \tan^2 \beta \right)$$

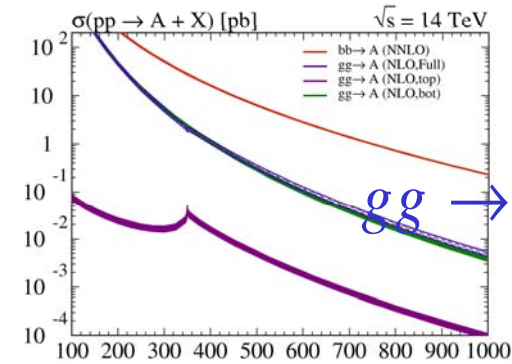
$$\sigma_{bb} = \frac{m_b^2}{M_h^4} c_4 \tan^2 \beta$$



$\tan \beta = 1$



$\tan \beta = 7$



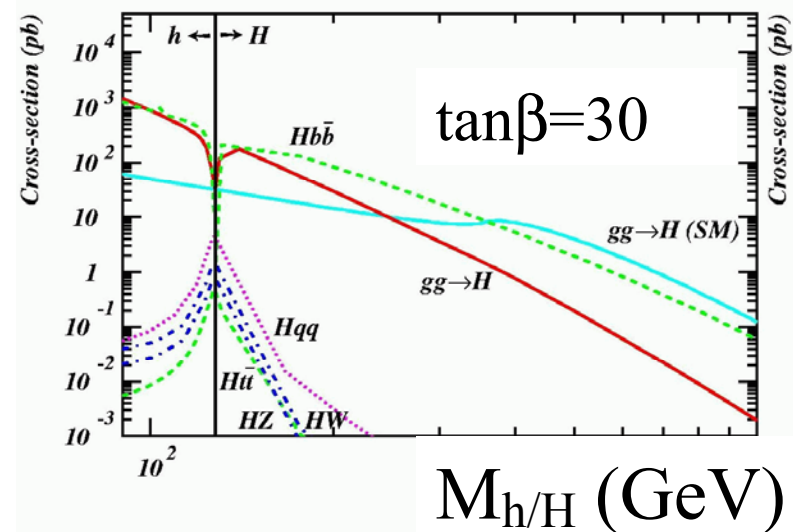
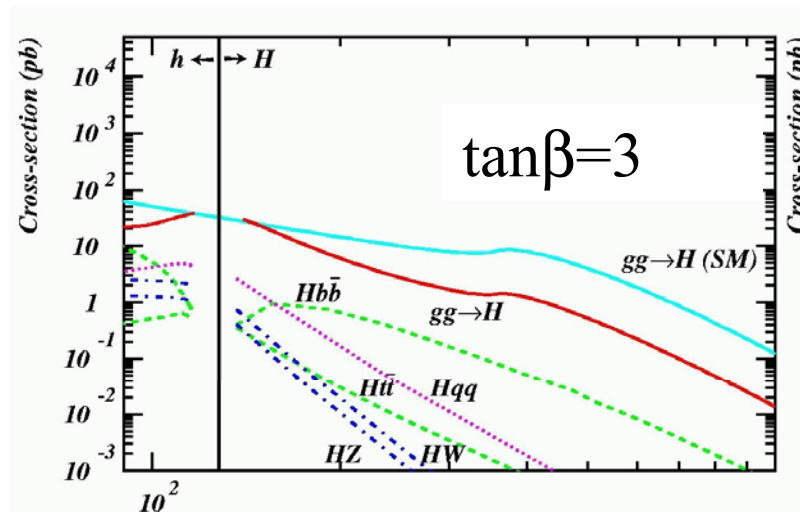
$\tan \beta = 40$

$\tan \beta \geq 7$, $b\bar{b}$ production mode dominates

Production of SUSY Higgs Bosons

- For large $\tan\beta$, dominant production mechanism is with b's
- bbh can be 10x's SM Higgs rate in SUSY for large $\tan\beta$

LHC



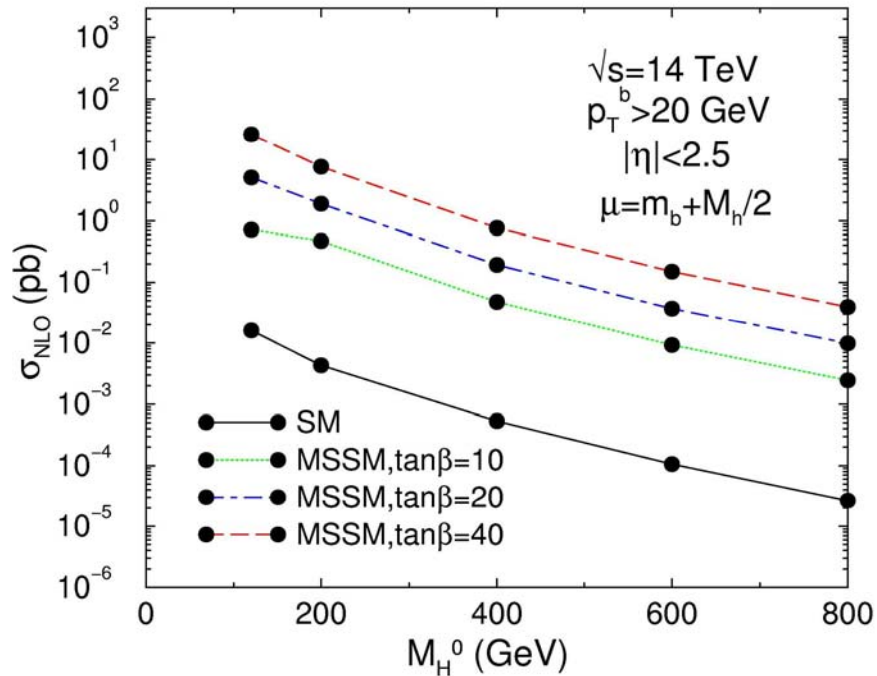
SUSY Higgs are produced with b's!

*Maximal Mixing

$pp, p\bar{p} \rightarrow b\bar{b}H$

Enhancement in MSSM

Note log scale!



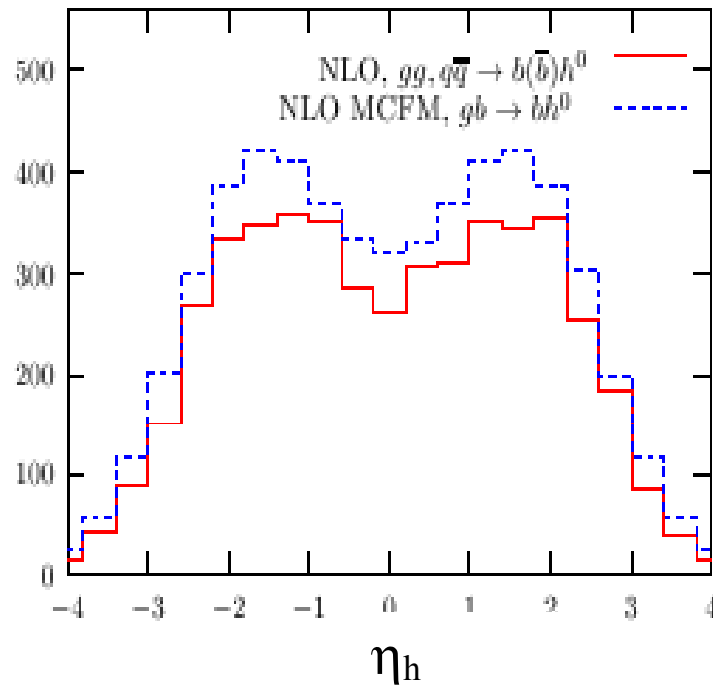
This is why the calculation is interesting!

α_{eff} from FeynHiggs with
 $M_{\text{SUSY}}=M_g=\mu=M_2=1 \text{ TeV},$
 $A_b=A_t=25 \text{ GeV}$

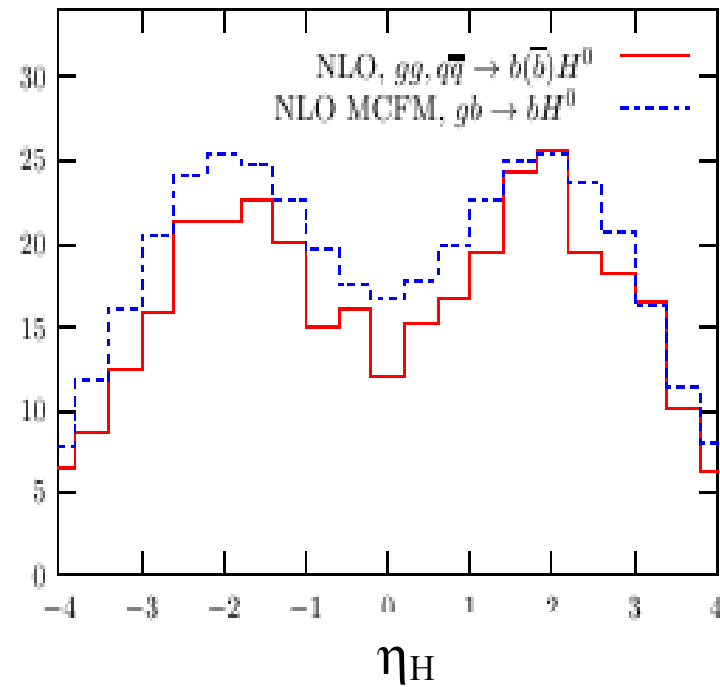
Can observe heavy MSSM scalar Higgs boson

Single b tag

$d\sigma/d\eta_h$ (fb/GeV) Tevatron



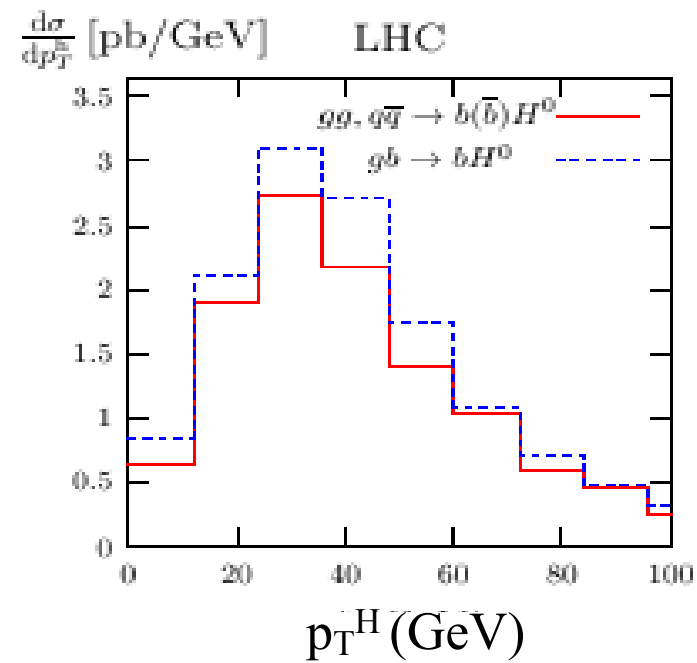
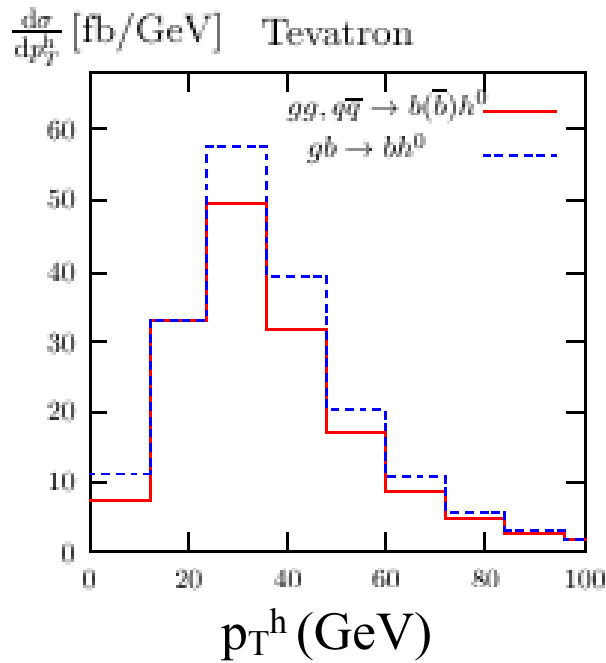
$d\sigma/d\eta_H$ (pb/GeV) LHC



MSSM with $M_h=M_H=120$ GeV, $\tan \beta=40$

Single b tag

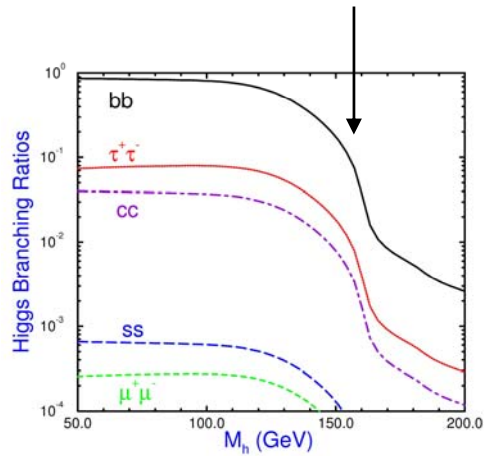
NLO



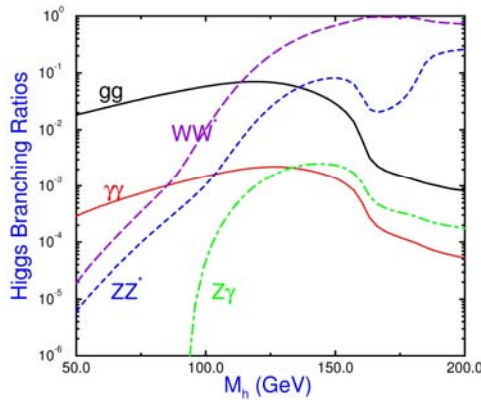
MSSM with $M_h=M_H=120$ GeV, $\tan \beta=40$

Higgs Decays also affected at large $\tan \beta$

- SM: Higgs branching rates to bb and $\tau^+\tau^-$ turn off as rate to W^+W^- turns on ($M_h > 160$ GeV)

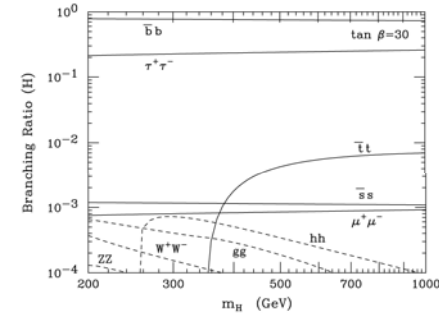


SM



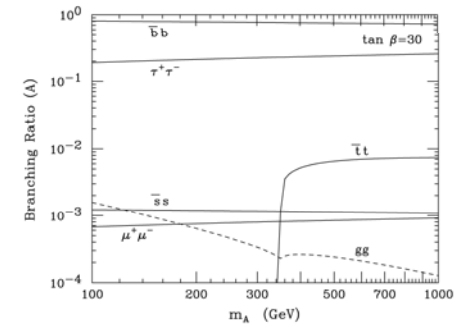
- MSSM: At large $\tan \beta$, rates to bb and $\tau^+\tau^-$ stay large

Heavy H^0 MSSM BRs



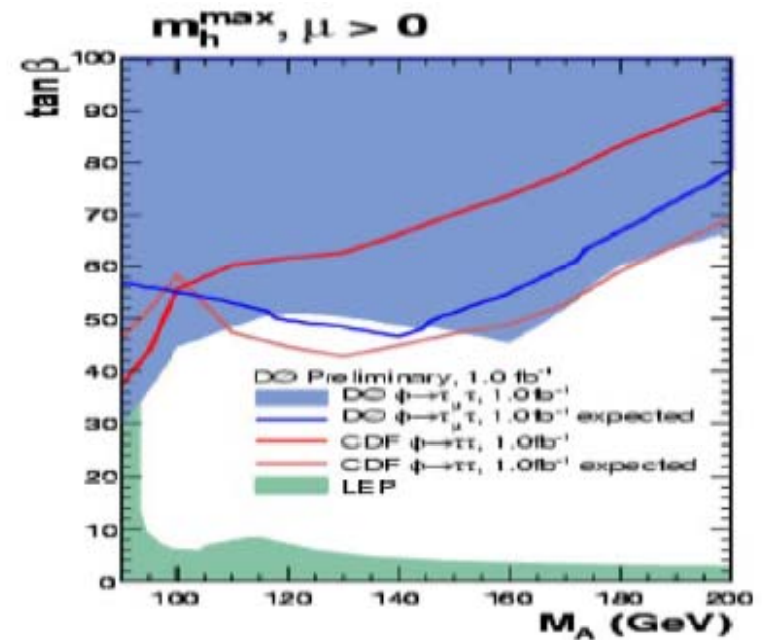
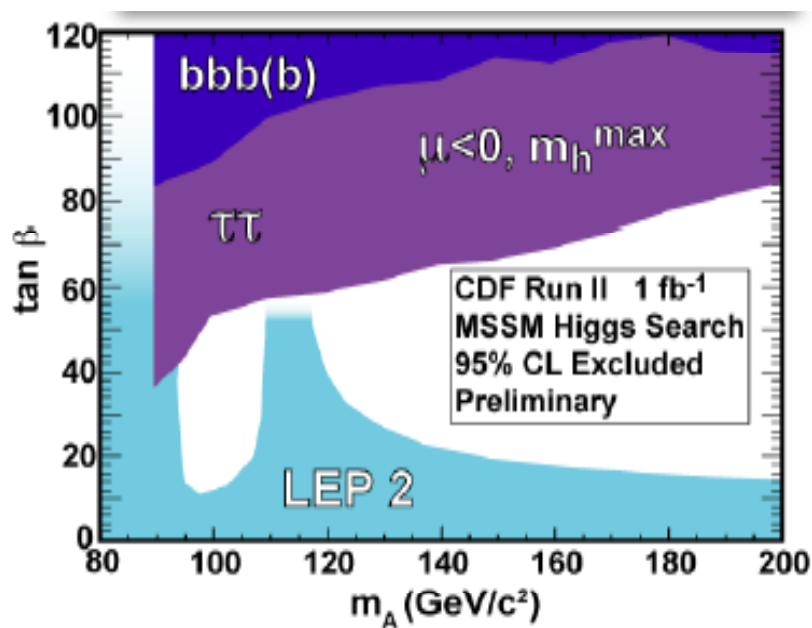
Rate to bb and $\tau^+\tau^-$ almost constant in MSSM

A^0 MSSM BRs



MSSM limits from $bg \rightarrow bh$ (1 fb^{-1})

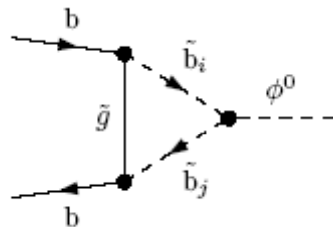
LP, 2007



30 fb^{-1} CMS expects to get to $\tan \beta \sim 15$ through bh ; $h \rightarrow \tau^+\tau$, $h \rightarrow bb$

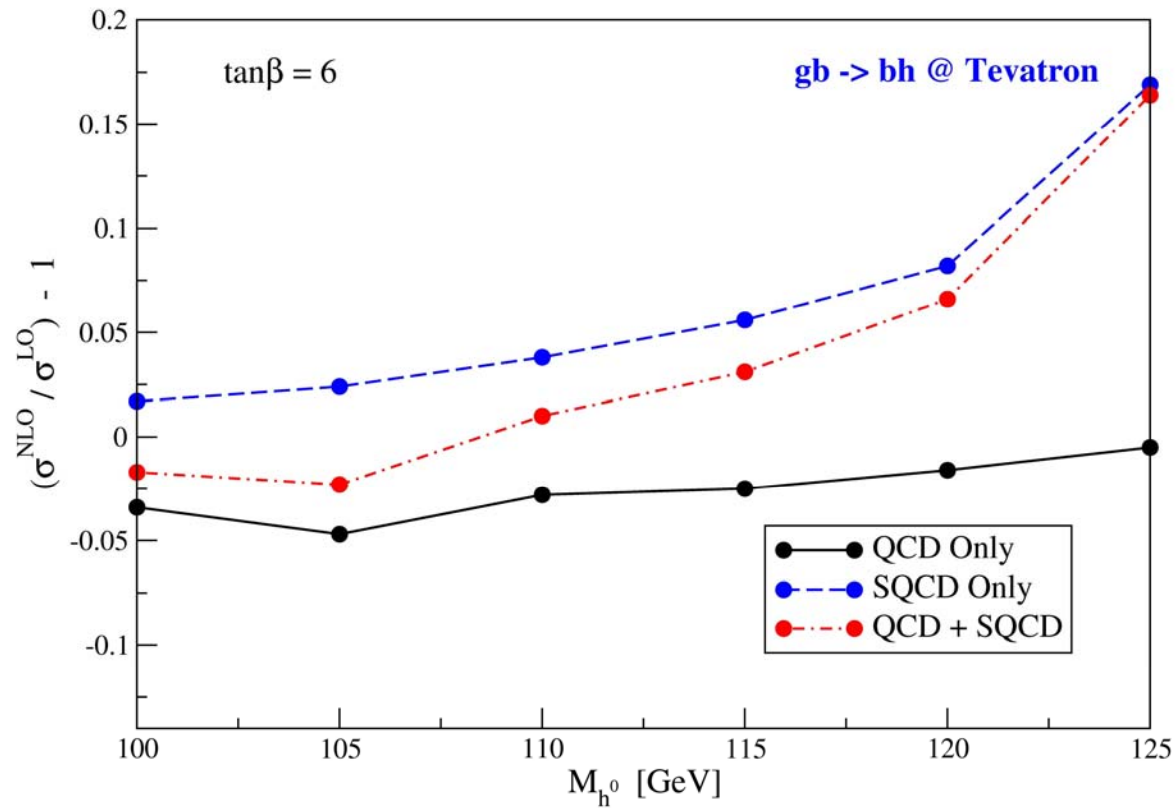
A Reliable Prediction

- We have $bg \rightarrow bh$ at QCD NLO
 - PDF/scale/scheme uncertainties $\sim 10\text{-}20\%$
- Are squark/gluino contributions relevant?
 - Important for $bb \rightarrow H, A$ at LHC
 - For some parameters as large as -50% effects
 - Squark/gluino effects almost completely described by *Improved Born Approximation*

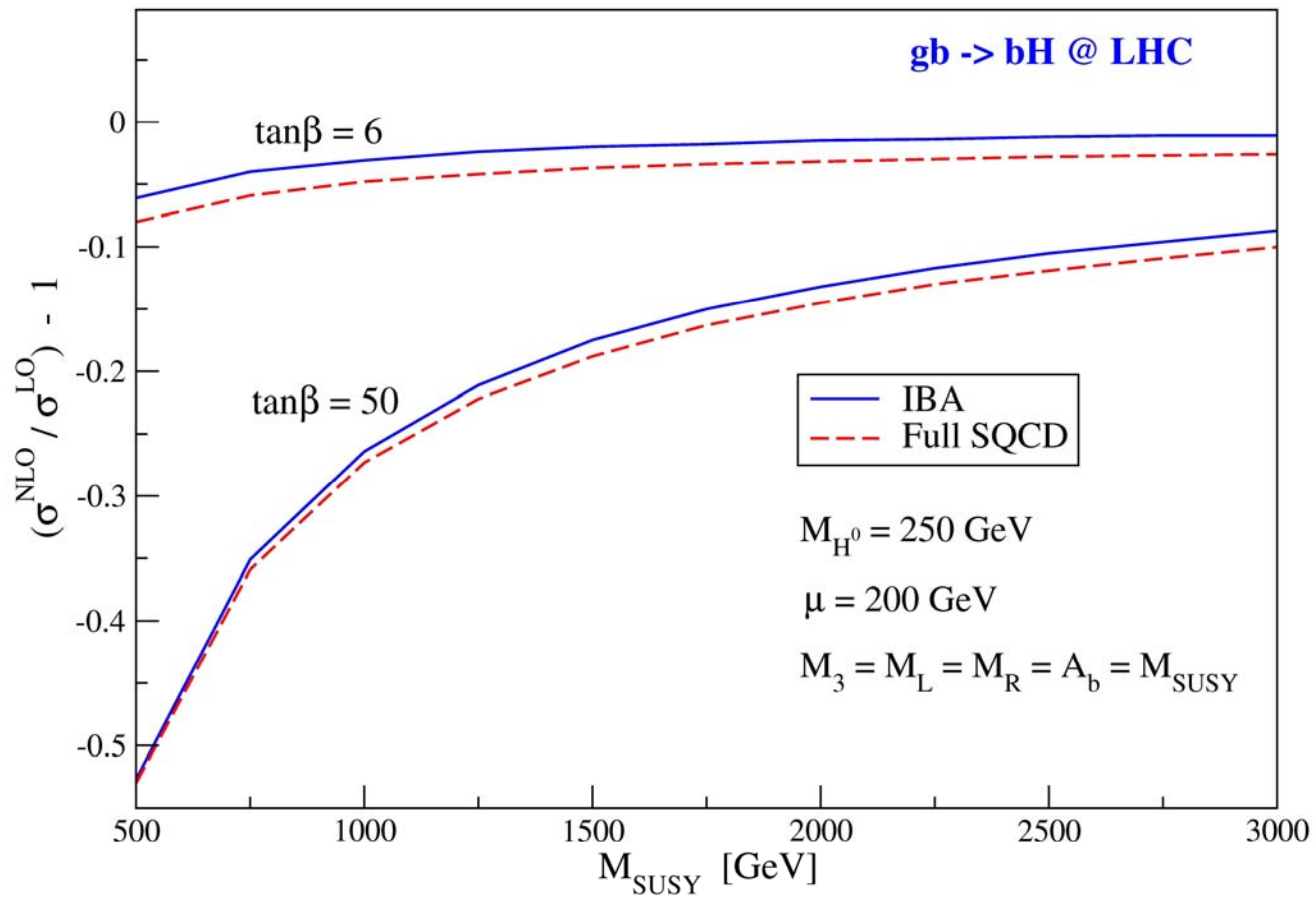


Dittmaier et al, hep-ph/0611353

Need SQCD for Reliable Predictions

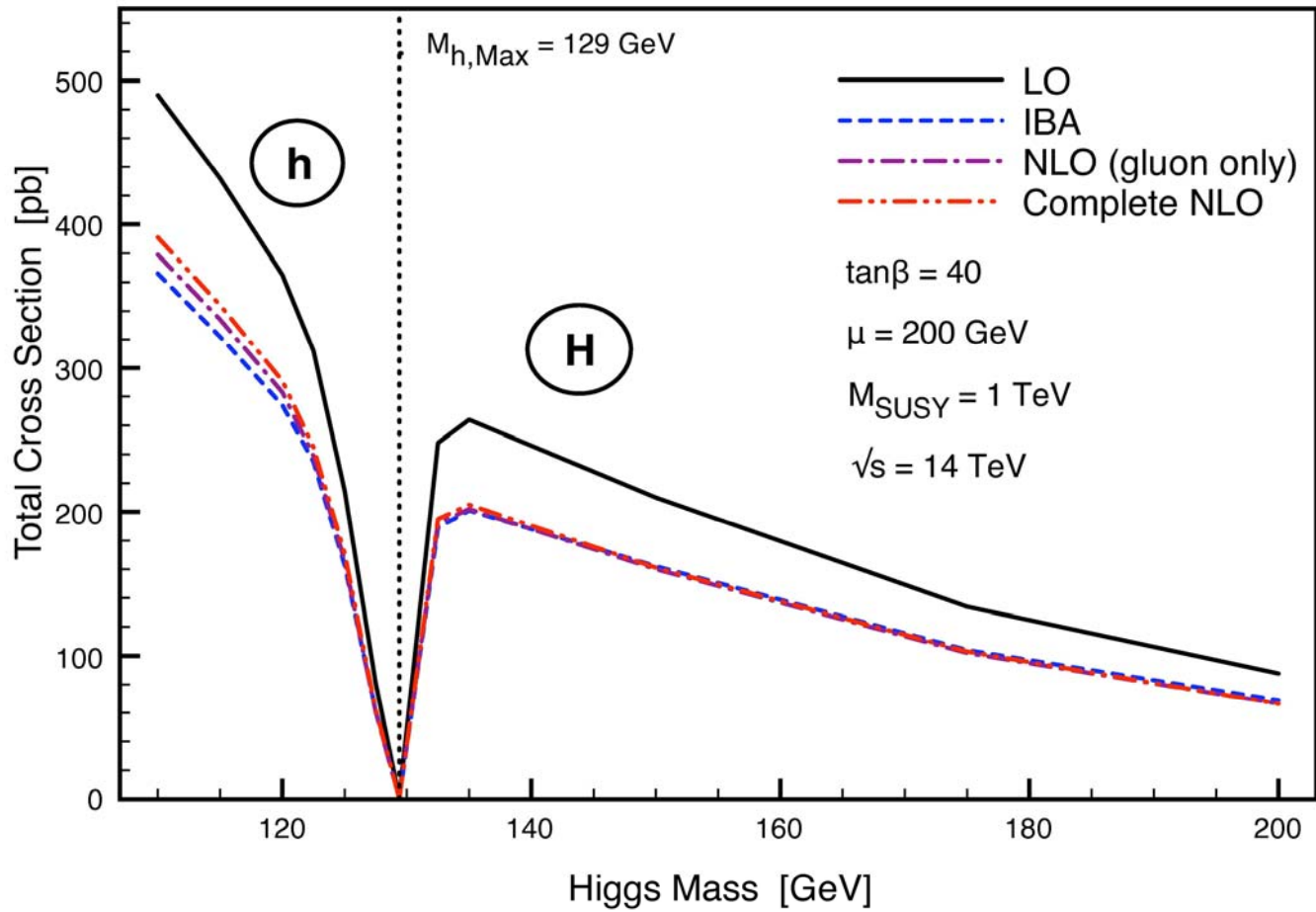


Squark/gluino loops important for large $\tan\beta$ and small M_{SUSY}



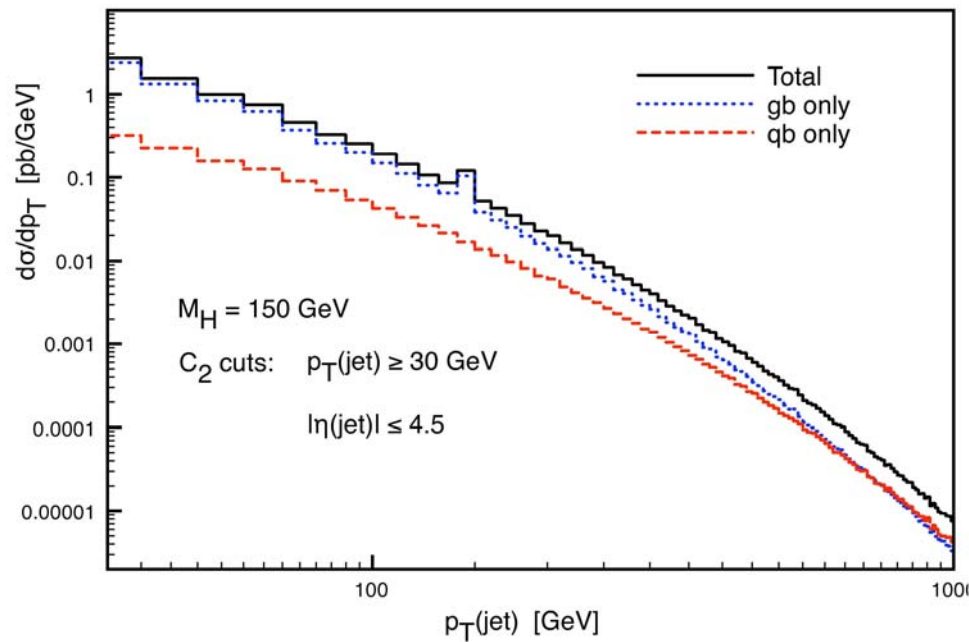
Note slow approach to decoupling limit for large $\tan\beta$

$gb \rightarrow b\phi$



Can $gb \rightarrow b\phi + \text{jet}$ be useful?

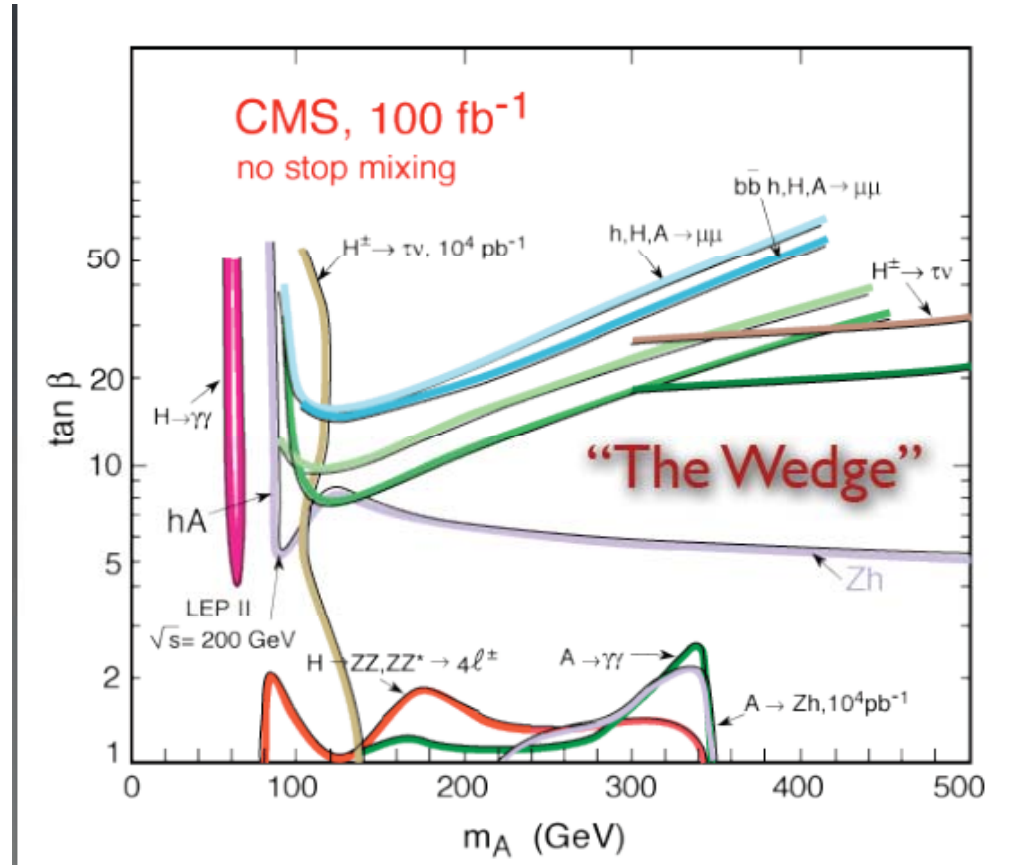
- LHC



More soon.....

Conclusions

- In the MSSM Higgs and b quarks go together at large $\tan \beta$
- Higgs production with b's is dominant mechanism for $\tan \beta > 7$
- Theoretical understanding of b PDFs: compatible answers in 4FNS and 5FNS for PDFs
- SUSY QCD corrections can be the same size as QCD corrections



Jan, 2007