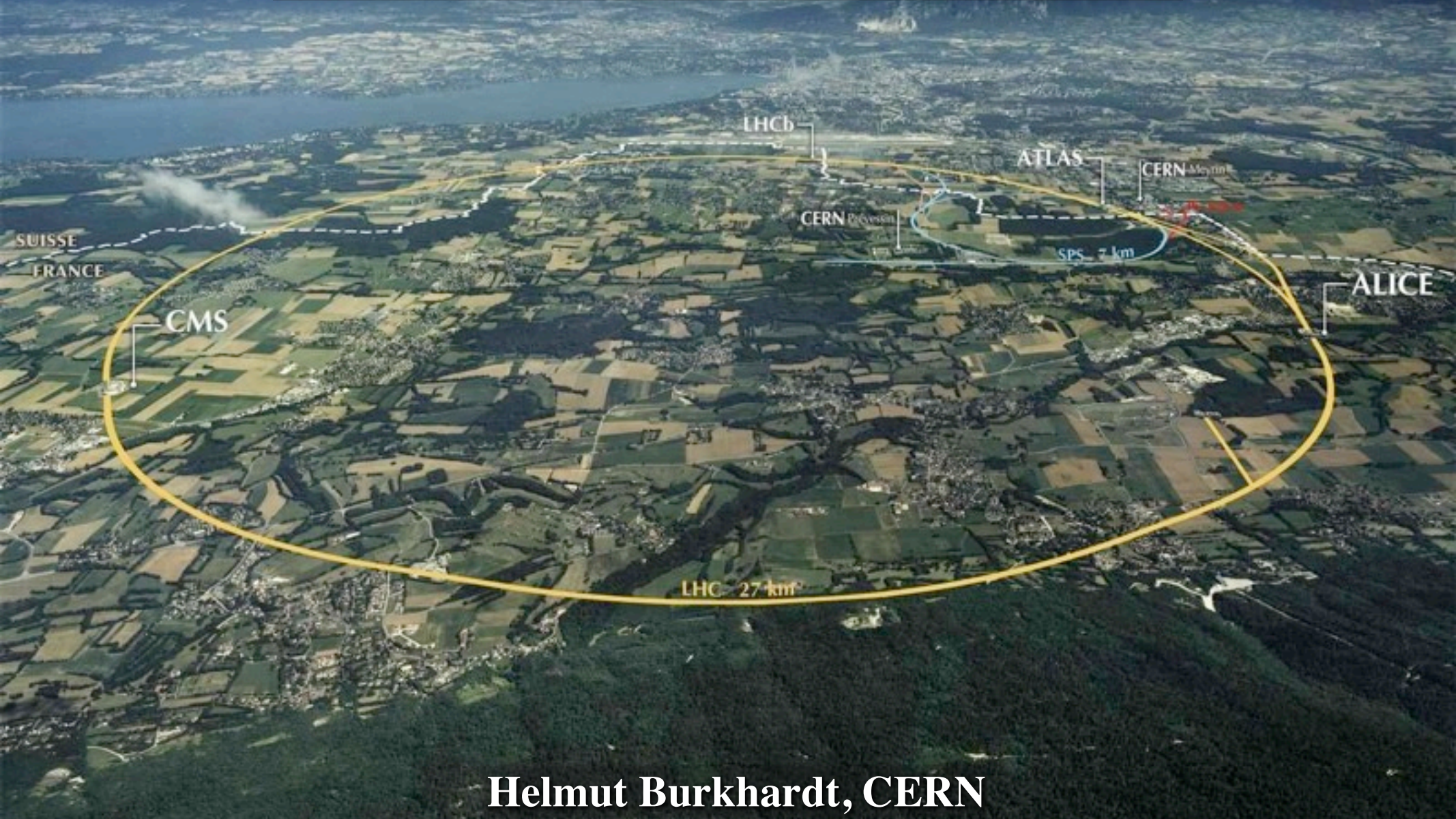


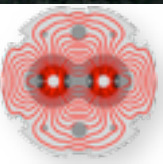
Introduction to Accelerators



Helmut Burkhardt, CERN



[ISEF 2013](#) 24 June 2012



Contents

- **Concepts: Energy Gain, E / B field. Units**
- **Types of accelerators : Ring, Collider, Linac, e+e-, pp ; Cosmic**
- **Components: Source, Magnets, resonant Cavities**
- **Basic machine optics**
- **Energy and Luminosity**
- **Synchrotron Radiation**
- **Limitations, current and future challenges**

- **Mixed with examples - mostly from CERN machines and in particular the LHC**

General, introductory refs. and books on Accelerators :

E. D. Courant and H. S. Snyder, *Theory of the Alternating-Gradient Synchrotron*, [pdf](#)

M. Sands, *Physics of Electron Storage Rings*, [SLAC Report No. 121](#); Wiedemann, *Particle Accelerator Physics* Bd. I,II

S.Y. Lee, *Accelerator Physics*, [World Scientific](#); M. Conte, W. MacKay, *Physics of Particle Accelerators*, [World Scientific](#)

CERN CAS yellow reports ; K. Wille, *The physics of particle accelerators*, Oxford University Press, 1996

Accelerators for Particle Physics, H. Burkhardt, in Handbook of Particle Detection and Imaging, [Ed. C. Grupen](#), Oct. 2011

The Large Hadron Collider : O. Brüning, H. Burkhardt, S. Myers, [10.1016/j.pnnp.2012.03.001](#), [CERN-ATS-2012-064](#)

Accelerators and Colliders, Landolt-Börnstein New Series I/21C, [Springer 2013](#)

Accelerators at the Energy Frontier

Livingston plot

Exponential growth
of E_{cm} in **time**

Starting in 60's
with e^+e^- at about 1 GeV

Factor 4 every 10 y

$pp, p\bar{p}$: $E_{cm} / 6$
still **5 ×** above e^+e^- at
same time

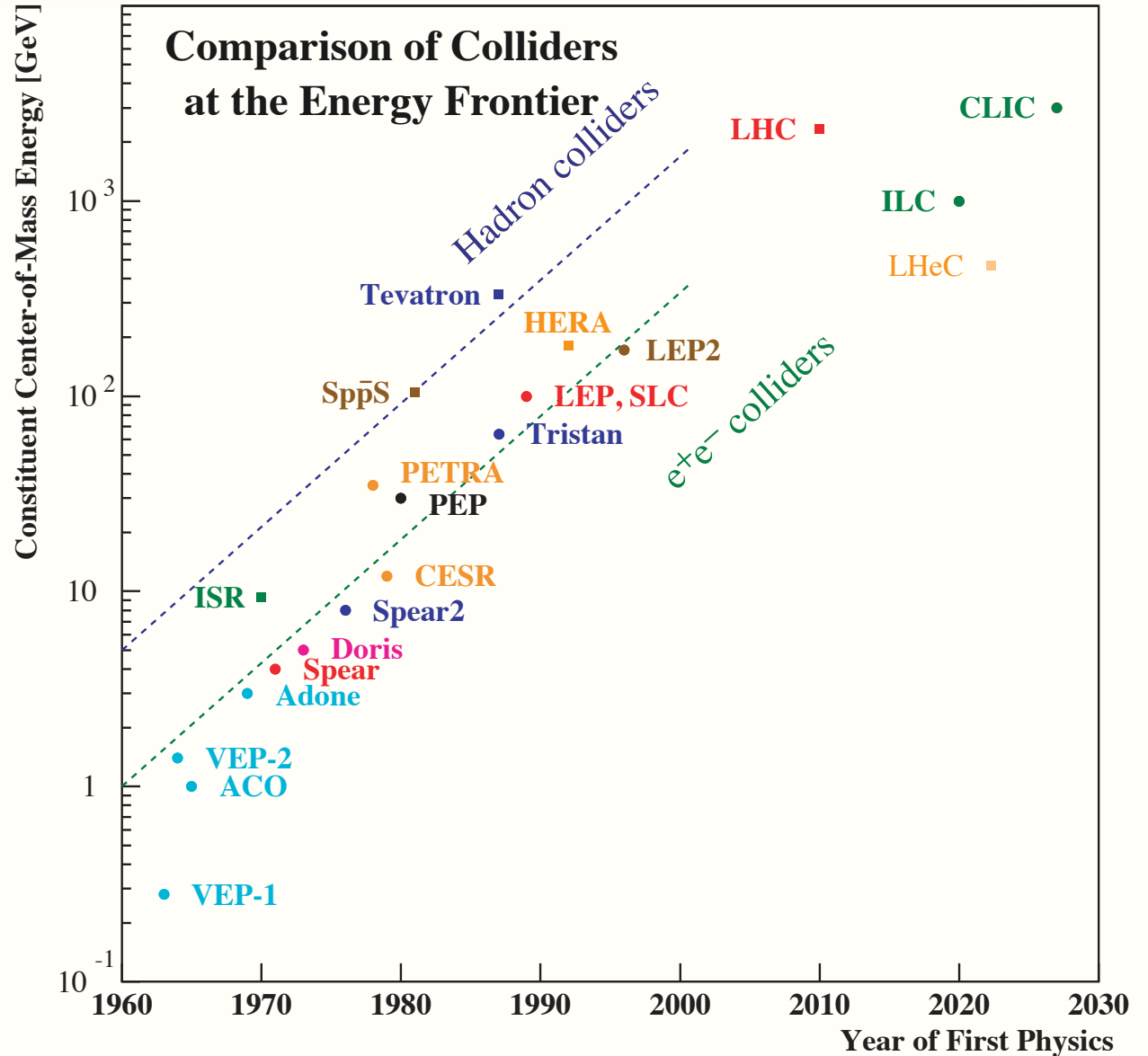
$pp, p\bar{p}$: **discovery**

e^+e^- : **precision**

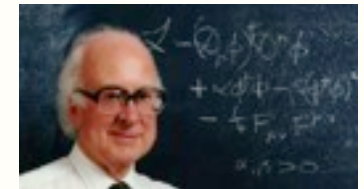
both required machines

+ ep : hadron structure, QCD

HERA, LHeC



The LHC is a major step forward
Discovery machine : Higgs ...



Basic concepts and units

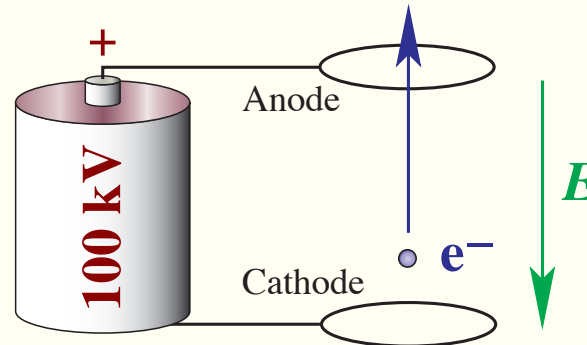
Electric field :

Acceleration

or rather

Energy gain

100 keV



Electric charge **e**
and electric field **E**

Special relativity, Lorentz transformation

$$E = \gamma m c^2 \quad p = \beta \gamma m c \quad \beta = \frac{v}{c} \quad \gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

$$m_e \approx 0.511 \text{ MeV}/c^2 \quad m_p \approx 938 \text{ MeV}/c^2 \quad e \approx 1.602 \times 10^{-19} \text{ C}$$

For $E = 10 \text{ GeV}$:

Electron $\beta = 0.999\,999\,9987 \quad \gamma = 19569.5$

Proton $\beta = 0.995\,588\,4973 \quad \gamma = 10.6579$

Unit conversion

$$\frac{e^2}{4\pi\epsilon_0} = \alpha \hbar c = r_{\text{part}} m_{\text{part}} c^2 = 1.43996 \times 10^{-18} \text{ GeV m}$$

$$\hbar c = 197.327 \times 10^{-18} \text{ GeV m}$$

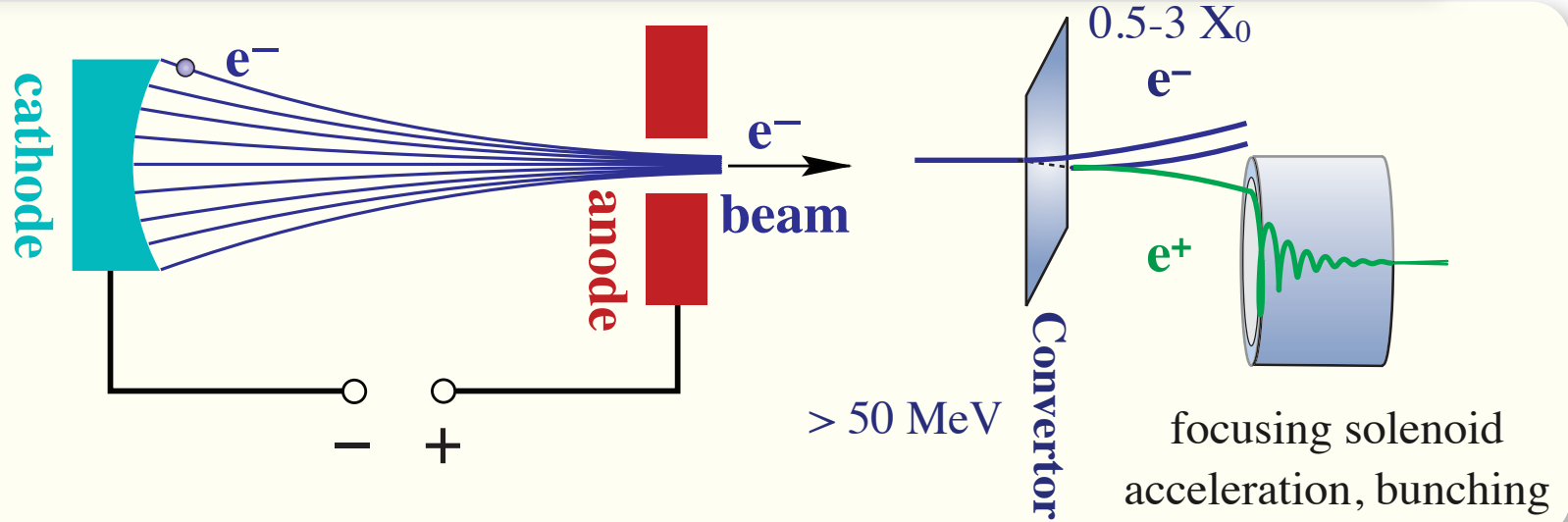
$$(\hbar c)^2 = 3.8938 \times 10^{-32} \text{ GeV}^2 \text{ m}^2 = 3.8938 \times 10^5 \text{ GeV}^2 \text{ nb}$$

for precise numbers see [PDG](#)

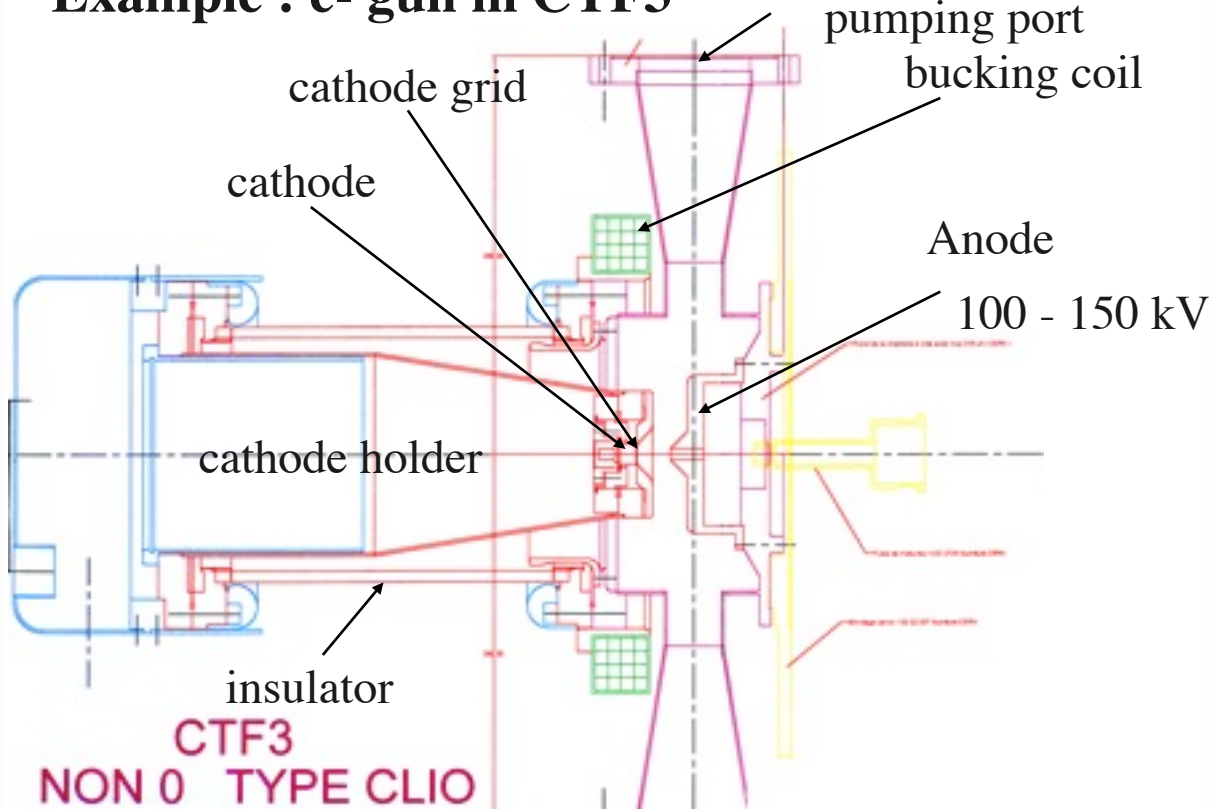
giga G = 10^9 tera T = 10^{12} peta P = 10^{15} exa E = 10^{18} zetta Z = 10^{21} yotta Y = 10^{24}

Particle sources

Thermionic electron source
principle same as cathode ray tube



Example : e- gun in CTF3



challenges :

high intensity

polarized e^- sources

damping rings for minimum emittance

undulator polarized e^+ sources

Proton and ion sources

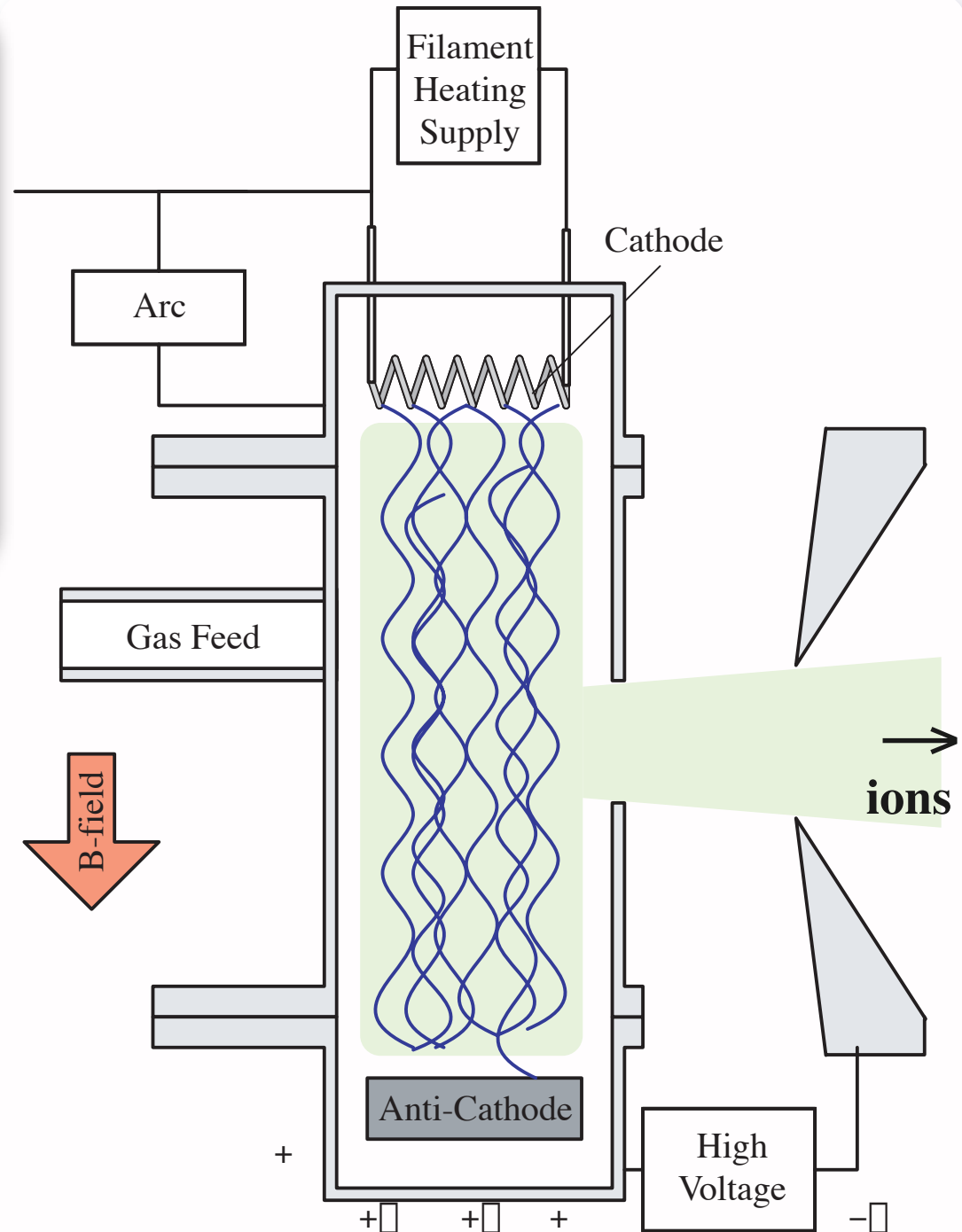
Various methods exist to produce p (H^+), H^- (p with 2 e^-) and heavy ions - heavier atoms, most electrons removed

Typically involves : **low pressure heated gas ionized gas / plasma**, inject H_2 to get protons, **or surface sputtering** and **electric and magnetic fields** to keep the electrons

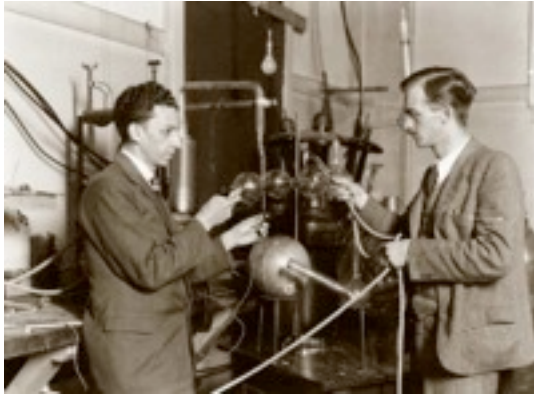
CERN p-source and 50 MeV Linac



special techniques
H⁻ injection
RadioFrequency
Quadrupole

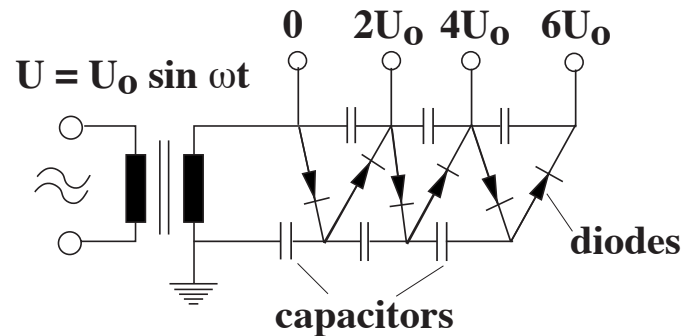


Linear Acceleration with Electrostatic Field



Cockcroft Walton
voltage multiplier

allows for DC, 100 % duty factor
limited by HV-breakdown $\sim 1 \text{ MV / m}$

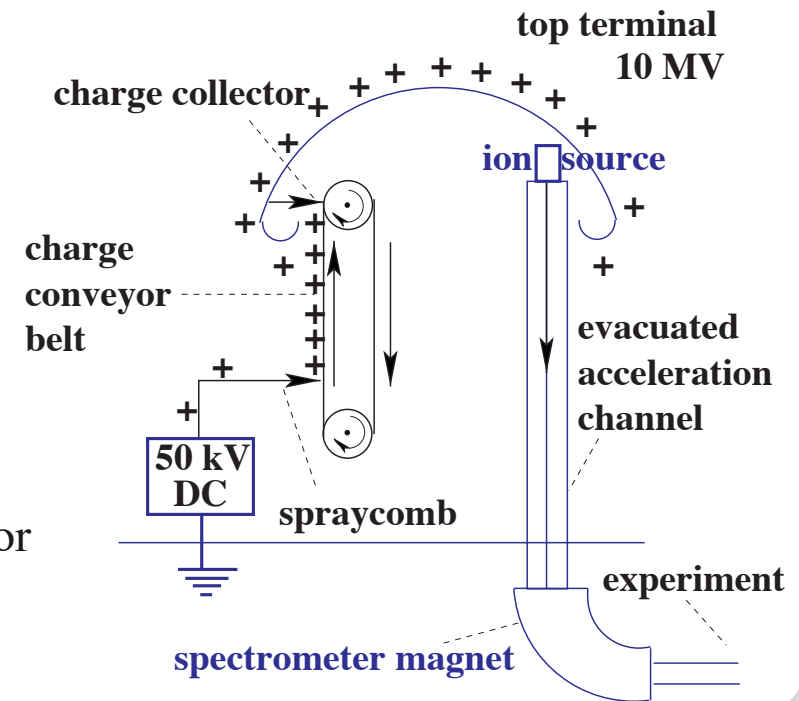


800 kV
proton pre-
injector
used at
CERN
until 1993



Van de Graaff generator
static electricity from belts

Oak Ridge Tandem Van de Graaff generator
reached 25.5 MV using pressurised SF₆



Time Varying Fields

Radio-frequency or short RF acceleration

- allows for multiple passages
- bunched beams, reduced duty cycle
- higher RF frequencies allow for higher acceleration gradients

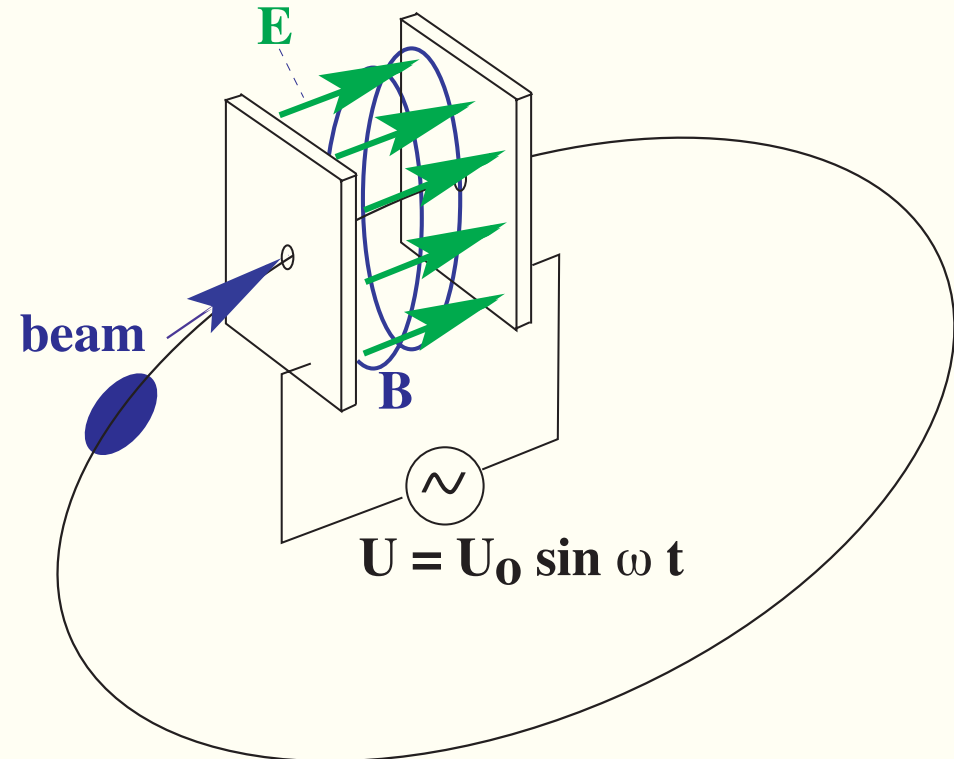
no time for breakdown / flashover

LEP , SC	8 MV / m at 352 MHz
Tesla / ILC, SC	31.5 MV / m at 1.3 GHz
CLIC	100 MV / m at 12 GHz

little gain above 12 GHz

SC limit ~ 50 MV/m, reached for single cell surface gradients higher than acceleration gradients, smooth structures

high f : shorter bunches - collective effects (peak current) and alignment more difficult
less energy stored in structure



Basic parameters, Lorentz Force

$$\mathbf{F} = q (\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

charge q , normally $q = e$; $q = Z e$ for ions

- Electric field \mathbf{E} provides the acceleration or rather energy gain
- The magnetic field \mathbf{B} keeps the particles on their path

ρ is the radius of curvature for motion perpendicular to the static magnetic field. Often called

- gyromagnetic or Larmor radius in astroparticle physics
 - bending radius for accelerators
- $B\rho$ known as magnetic rigidity, units Tm

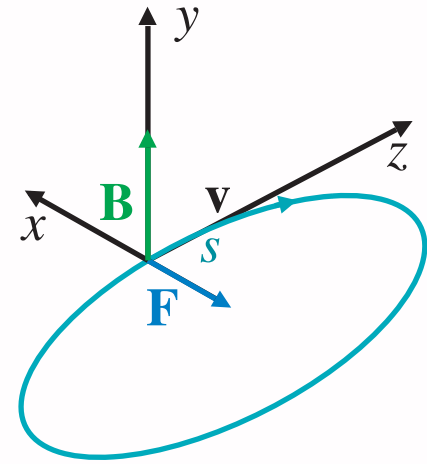
LHC

- Momentum $\mathbf{p} = 7 \text{ TeV}/c$
- LHC bending radius $\rho = 2804 \text{ m}$
- Bending field $\mathbf{B} = 8.33 \text{ Tesla}$
- magnets at 1.9 K , super-fluid He

Circular motion for

$$\mathbf{E} = 0$$

$$\mathbf{v} \perp \mathbf{B}$$



$$\mathbf{B} = \frac{\mathbf{p}}{q \rho}$$

for $q = e$ numerically
 $B \text{ [T]} = p \text{ [GeV}/c] \cdot 3.336 \text{ m} / \rho$
high energy, $v = c$ “ $p = E$ ”
 $E < E_H = q B \rho$ Hillas criterion

Astroparticle

units $10^{-4}\text{T} = 1\text{Gauss}$; a.u. = $1.5 \times 10^{11}\text{m}$

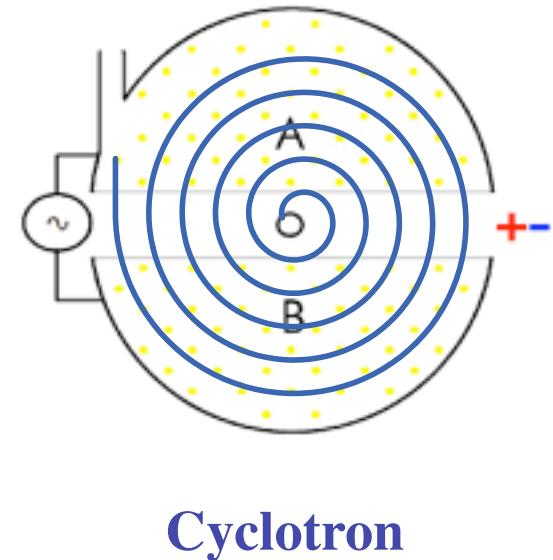
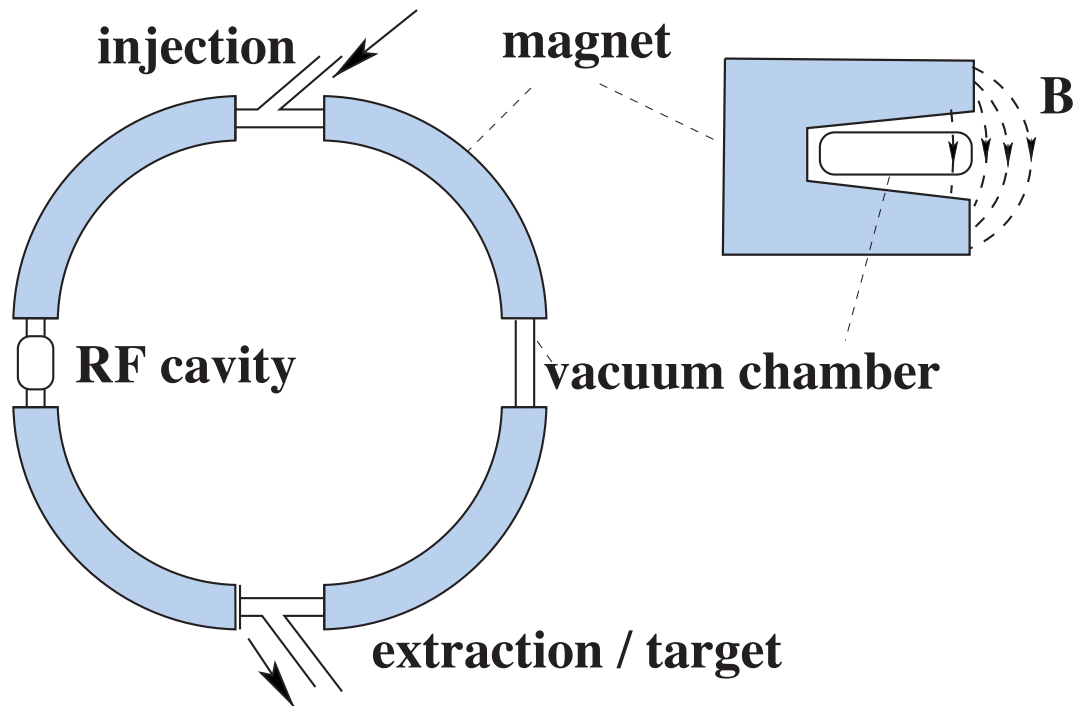
Solar system $\mathbf{B} = 10 \mu\text{G}$ $E = 5 \text{ TeV}$ $\rho = 11 \text{ a.u.}$

Intergalactic $\mathbf{B} = 1 \text{ nG}$ $E = 5 \text{ PeV}$ (knee)

$\rho = 1.7 \times 10^{19}\text{m}$ (4 % of galaxy-radius)

Circular Accelerator

- **Cyclotron** : constant rf-frequency. Magnetic field radius ρ increases with energy. Used for smaller machines



- **Synchrotron** : $\rho = \text{const.}$ **B increased with energy.** RF-frequency adjusted slightly ($\beta = 0.999 \dots 1.0$). Most HEP and all CERN ring accelerators PS, SPS, LEP, LHC of this type. Principle same for e, p, heavy-ion – PS, SPS – accelerate(d) all of these, in some cases switching within seconds

Phase stability I

acceleration,
ramping up in energy :

- allow for enough RF-voltage
- ramp up magnets
- particle adjust themselves in radius and phase to gain on average the right amount of energy

LHC nominal RF parameters

Voltage at injection 8 MV

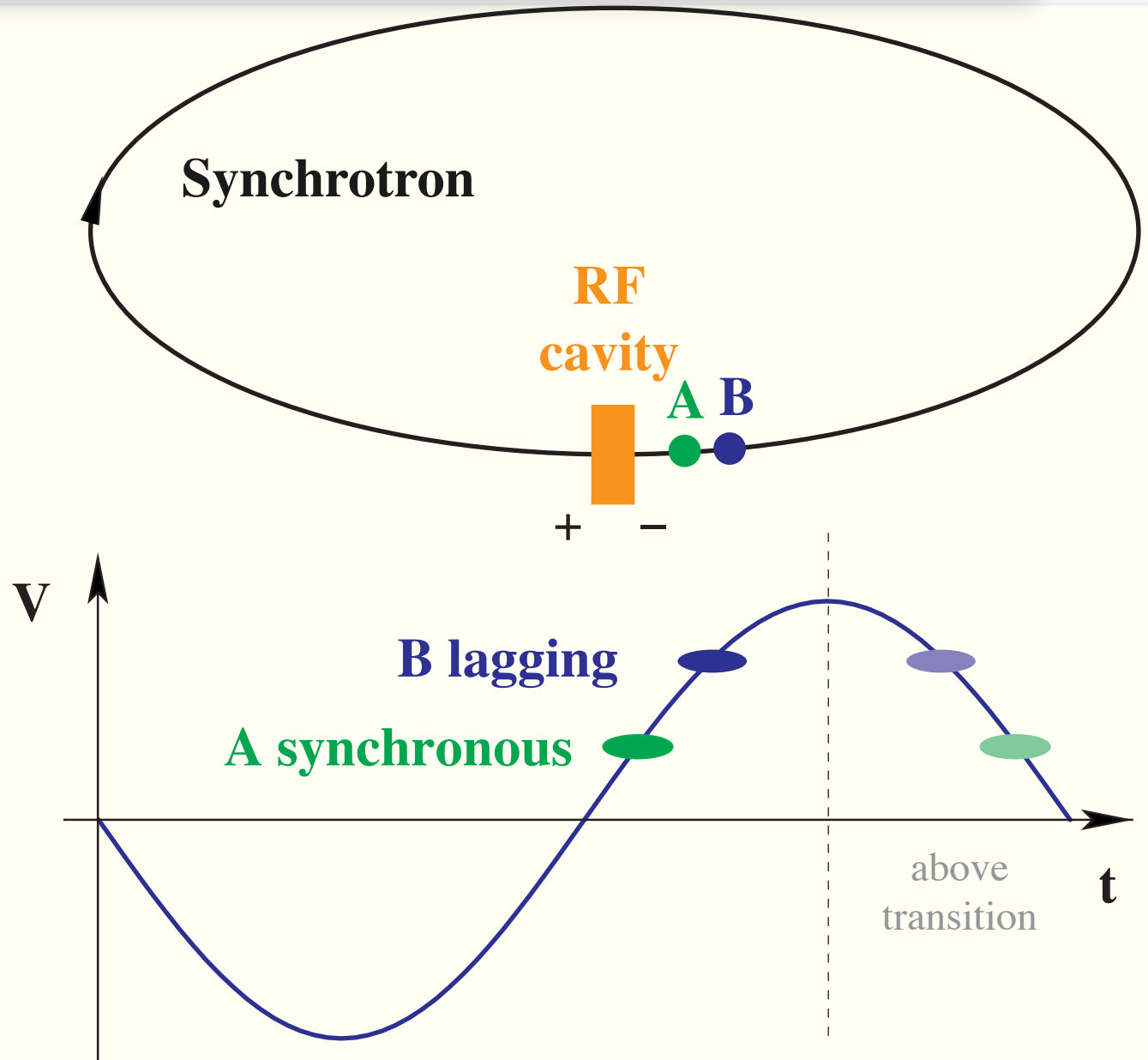
top energy 16 MV

Revolution frequency $f_{rf} = h f_{rev}$

Circumference $L = v / f_{rev} = \beta c / f_{rev}$

$h = 35\,640$ $f_{rf} = 400.7896$ MHz $L = 26658.864$ m

$f_{rev} = 11.2455$ kHz 1 turn in 88.92446 μ s



Magnets and Power Consumption

Why super conducting magnets ?

$$P = R I^2$$

LEP

B = 0.1 T LEP2 ~ 100 GeV

(half) cells with each three 11.55 m long dipole magnets

I = 4500 A together **R = 1 mΩ** **P = 20 kW / cell**

488 cells

P = 10 MW

if we would have kept the same magnets for the LHC

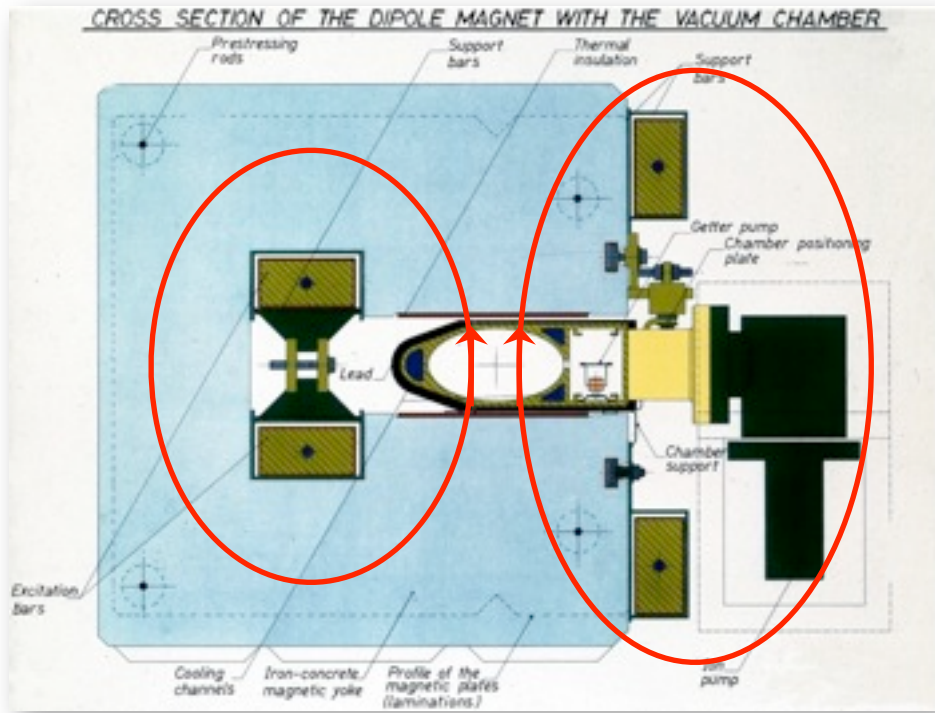
LHC **B ∝ I** **B = 8.38 T**

would need now **I = 280 kA** with LEP magnets **R = 1 mΩ**

P = 78 MW / cell × **488 cells** **total power P = 38 GW**

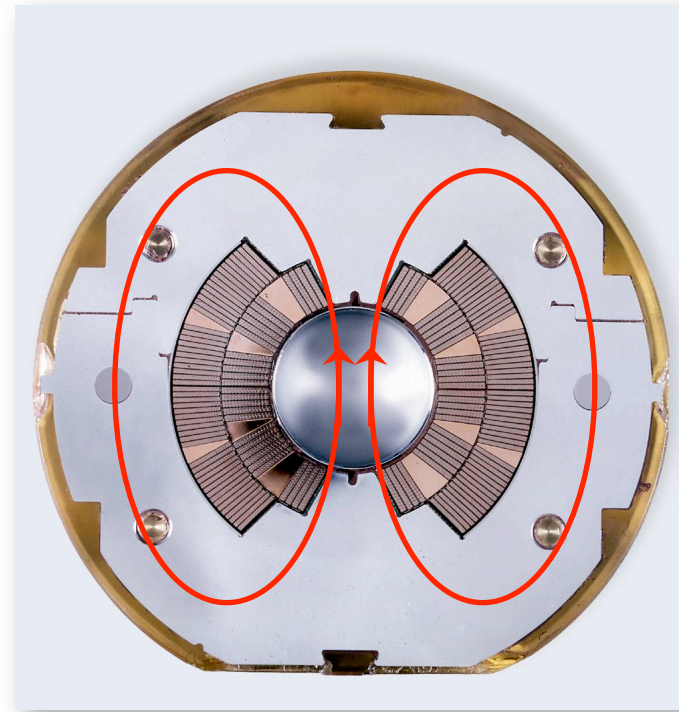
Magnet technology

warm



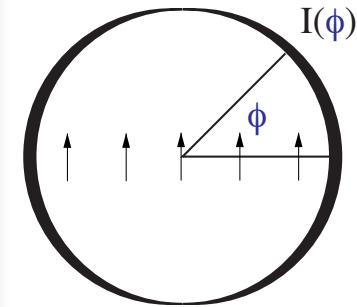
- field quality given by pole face geometry
- field amplified by Ferromagnetic material
- hysteresis and saturation ~ 2 T
- Ohmic losses for high magnet currents

cold



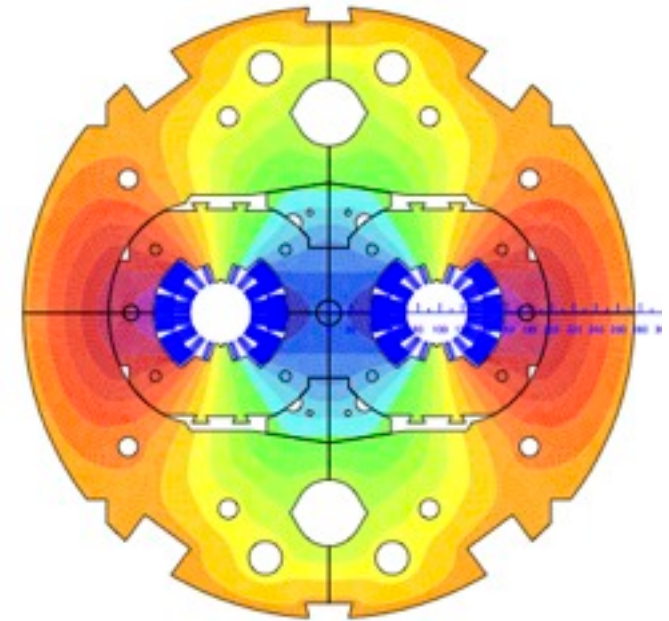
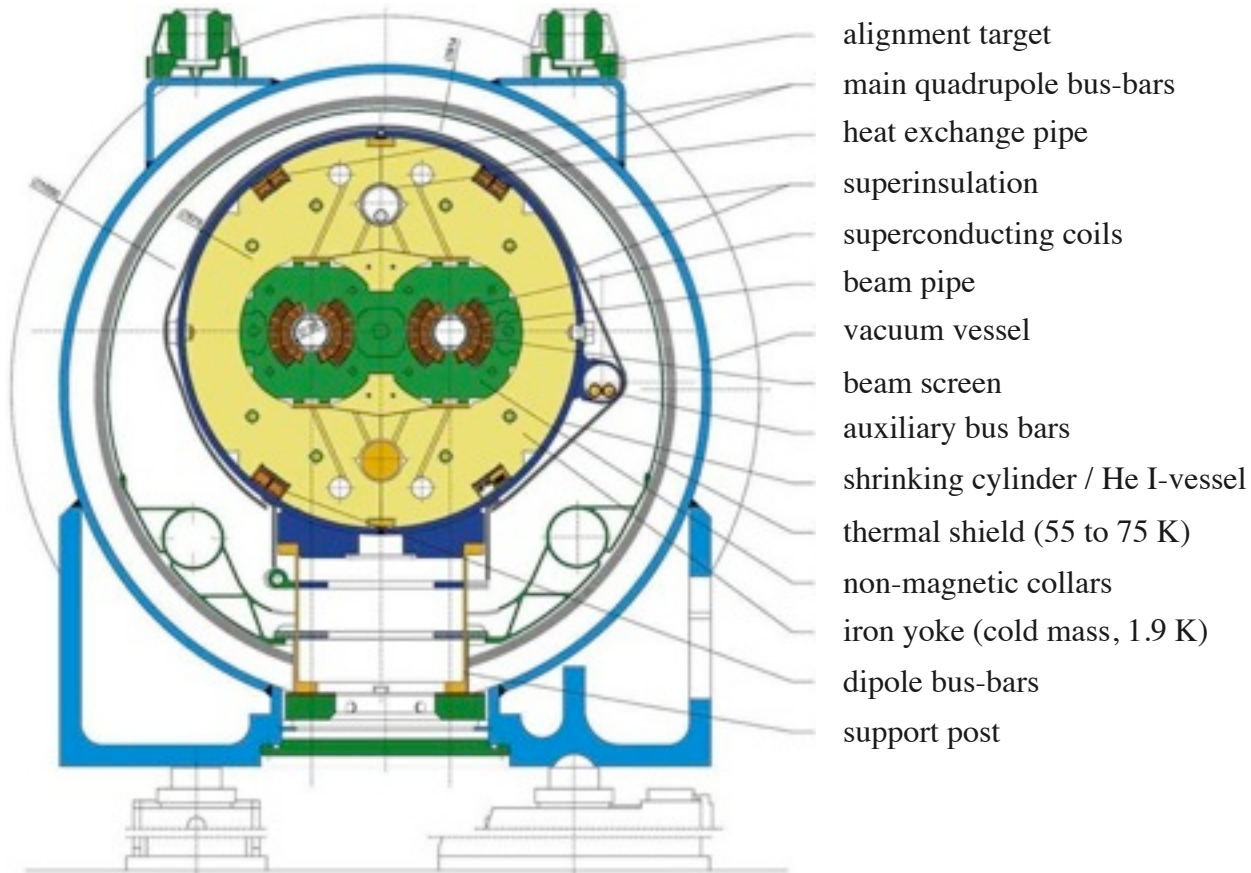
- field quality given by coil geometry
- requires cooling to cryogenic temperatures
- persistent currents and snap back
- risk of magnet quenches

Dipole current distribution
 $I(\Phi) = I_0 \cos(\Phi)$



LHC dipole magnet

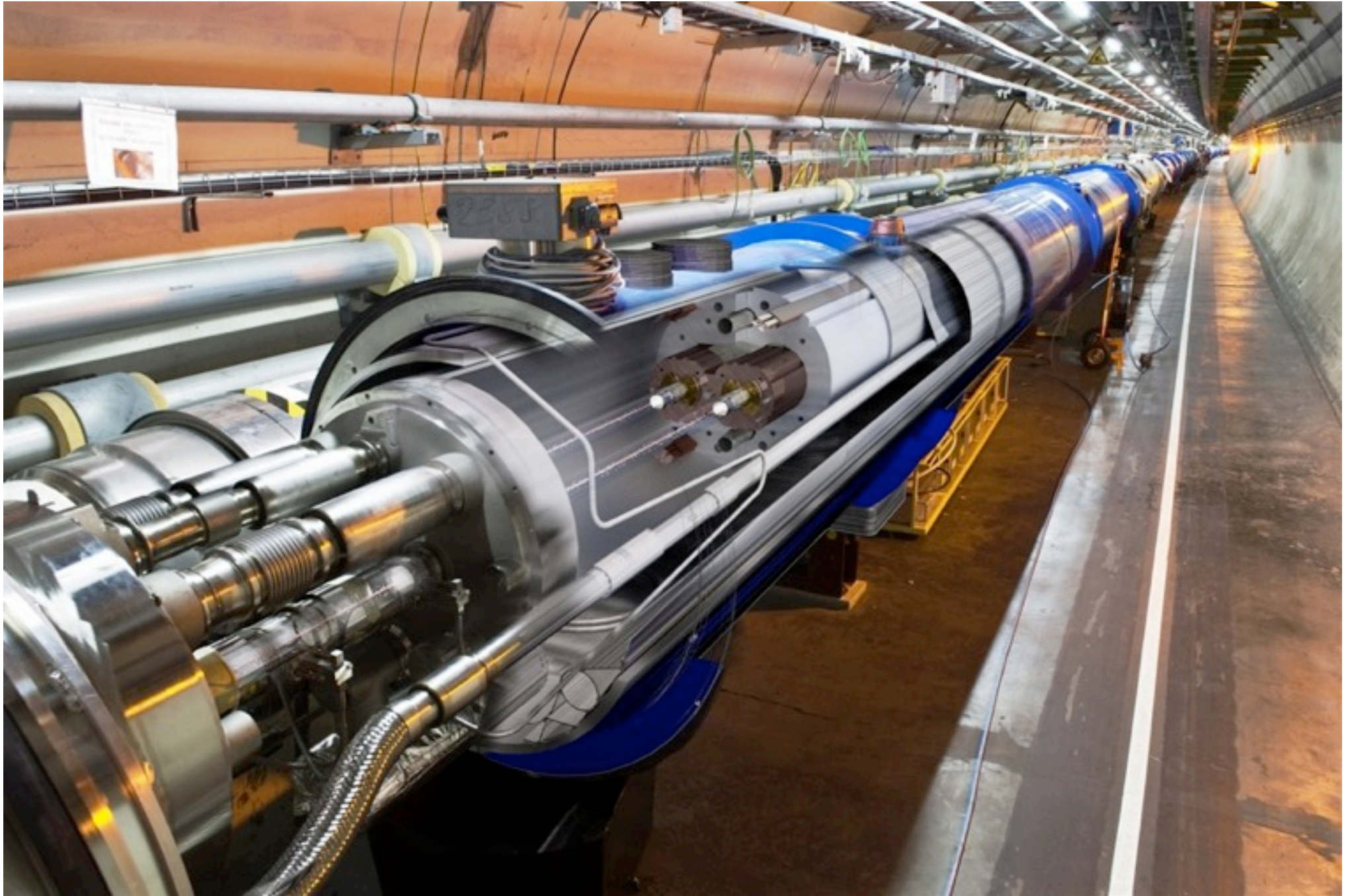
2-in-1 dipole magnet, 8.33 T field, 15 m long, mass 30 ton



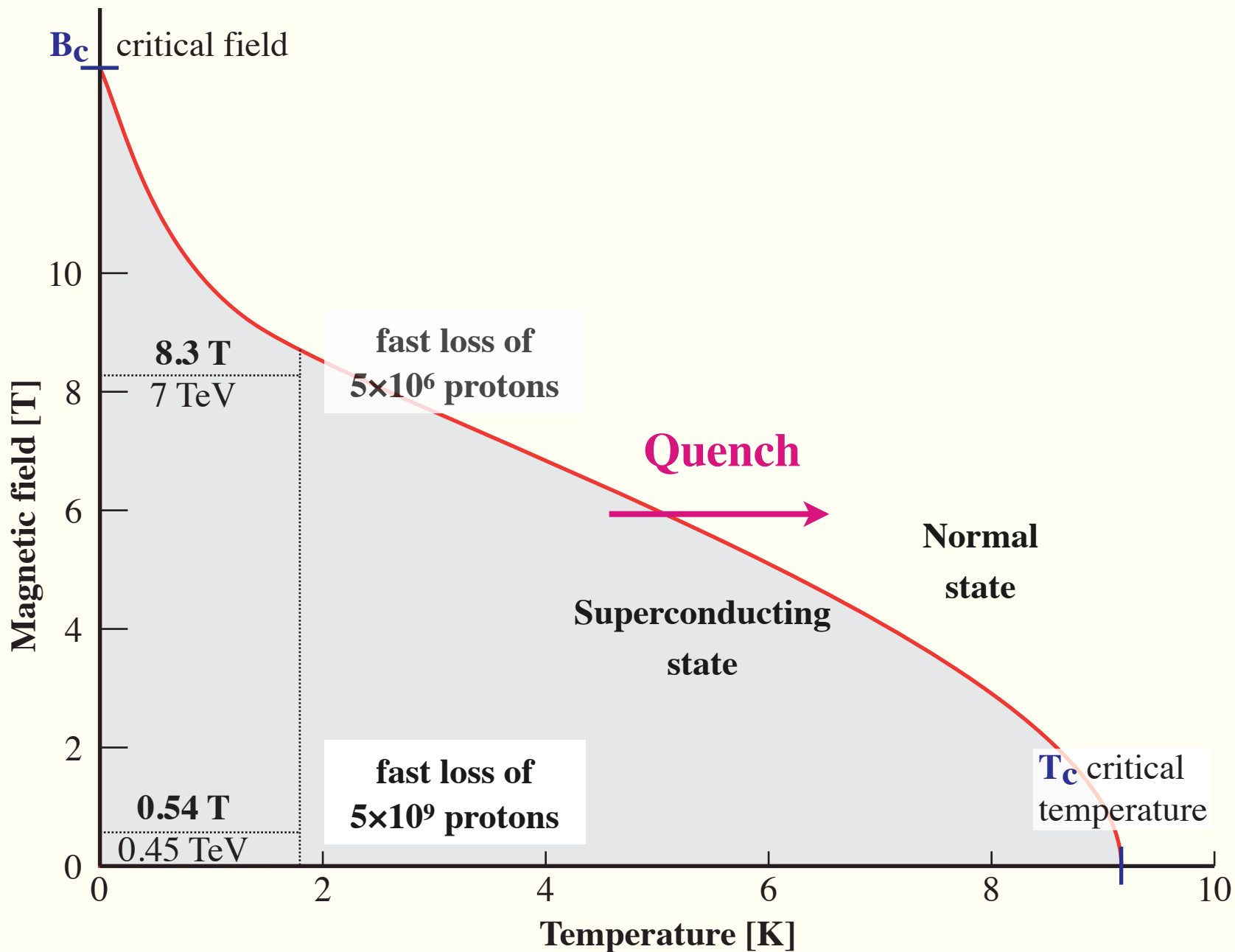
current distribution

LHC dipole magnet cross-section

LHC magnets installed in the tunnel



Operational margin of a superconducting LHC dipole



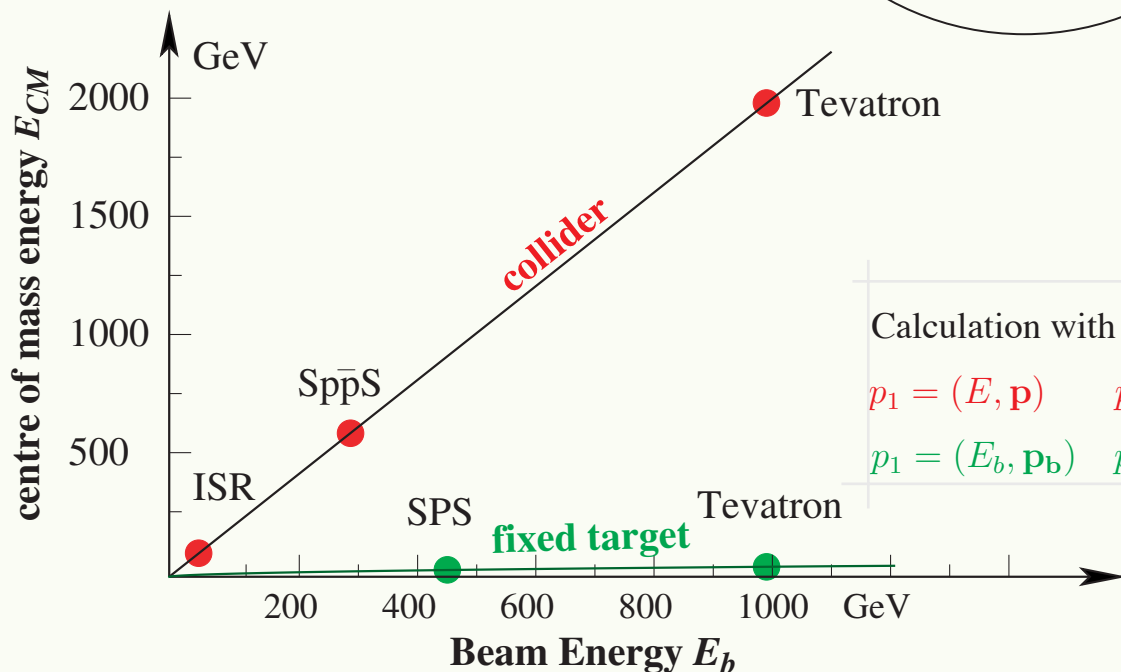
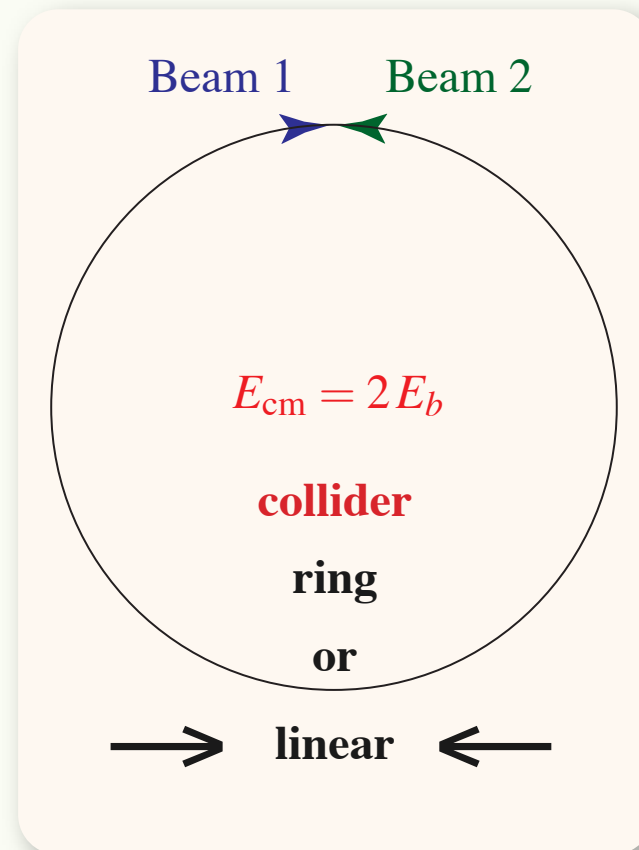
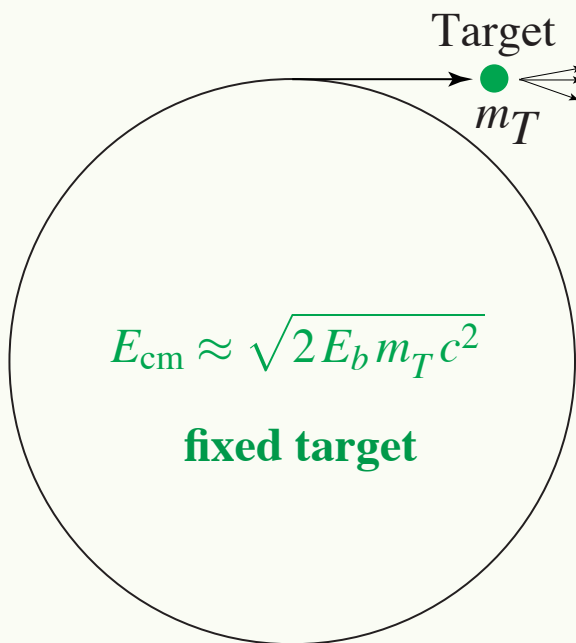
Fixed Target vs Collider

Fixed target, high energy collisions :

Energy “lost” as kinetic energy

High Energy $e+e^-$ and
very high energy pp gain
a lot from **colliders**

Gain for LHC is by $\times 122$
(14 TeV / 114.6 GeV)



Calculation with four vectors for $c = 1$ $E_{CM} = \sqrt{s}$ $s = (p_1 + p_2)^2$

$p_1 = (E, \mathbf{p})$ $p_2 = (E, -\mathbf{p})$ $s = 2m^2 + 2E^2 + 2p^2 = 4E^2$ **collider**

$p_1 = (E_b, \mathbf{p}_b)$ $p_2 = (m_T, \mathbf{0})$ $s = m_b^2 + m_T^2 + 2m_T E_b$ **fixed target**

Primary cosmic ray spectrum

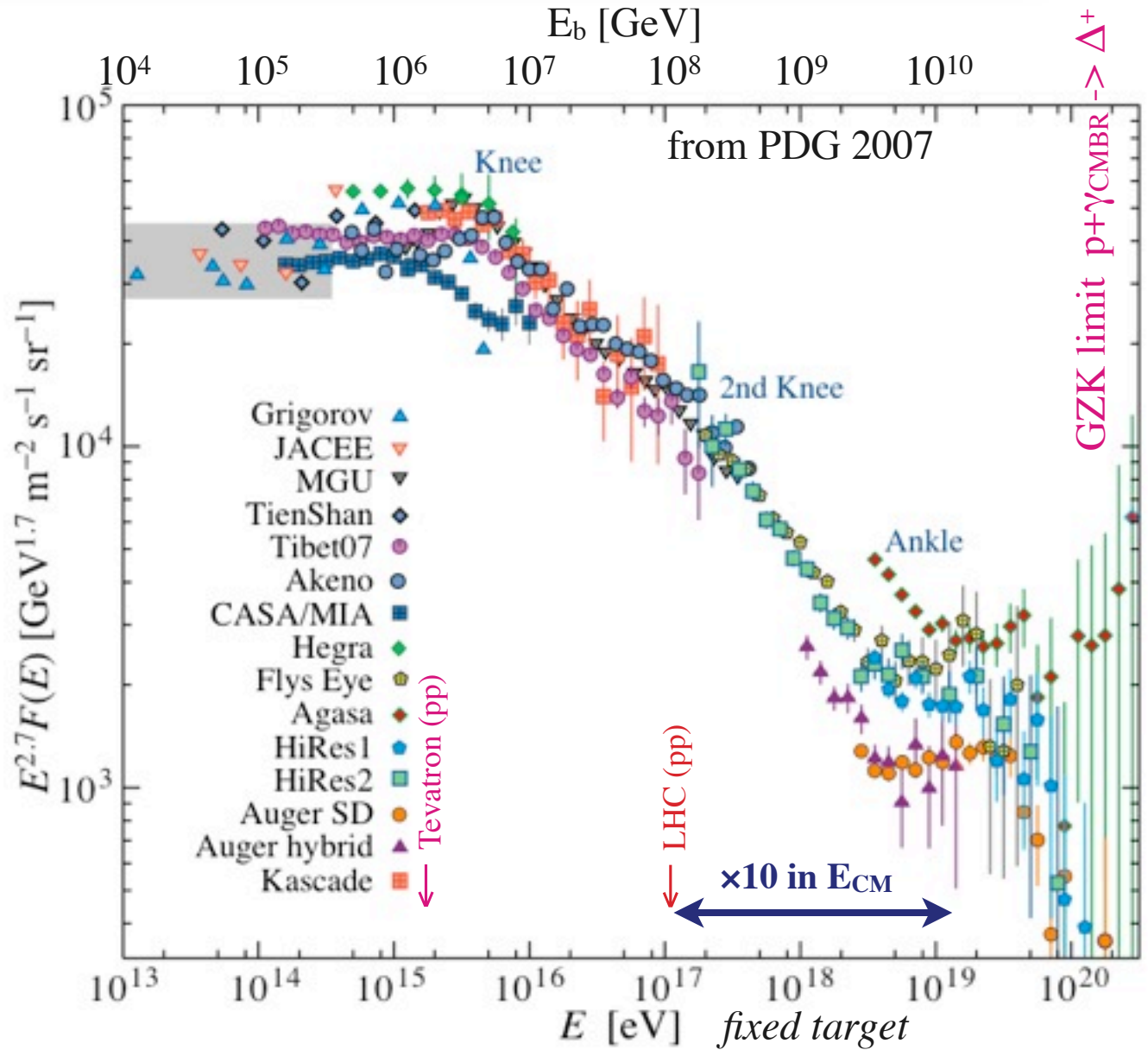
E spectrum falls as $E^{-2.7}$
 to knee at $E \approx 5e15 \text{ eV}$
 $= 5 \times 10^6 \text{ GeV}$
 $\sim 1 \text{ particle/m}^2 \text{ and year}$
 origin galactic

above $\sim E^{-3}$

back to $E^{-2.7}$ at very
 highest energies

conversion to E_{cm}

E_b [eV]	E_{cm} [TeV]
10^{13}	0.137
10^{15}	1.370
10^{17}	13.70 \approx LHC pp
10^{19}	137.0 \leftarrow LHC ions
10^{21}	1370.



Nature has much larger and more powerful **cosmic accelerators** than we can ever built.

With colliders we can get to these collision energies in clean laboratory conditions.

The LHC already gets us to within 1-2 orders of magnitude of the very highest cosmic rays.

Luminosity and collision rates

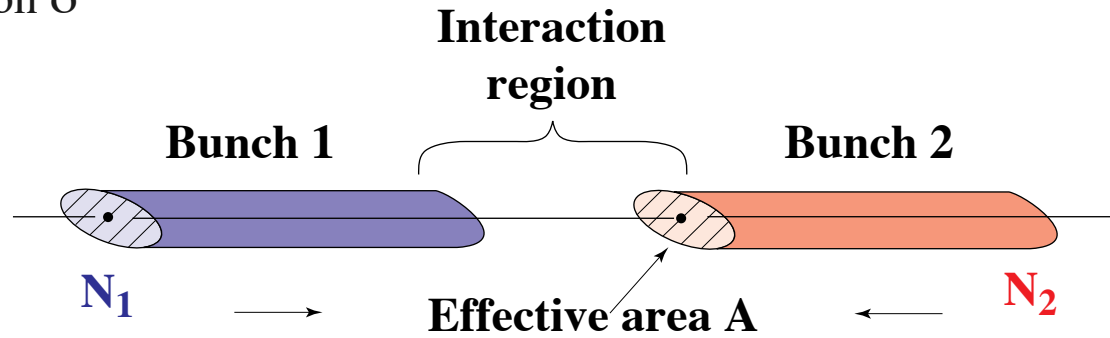
Event rate for process with cross section σ

$$\dot{n} = \mathcal{L} \sigma$$

Luminosity from bunch

crossings at frequency $f = f_{\text{rev}} n_b$

$$\mathcal{L} = \frac{N_1 N_2 f}{A}$$



for Gaussian bunches with rms sizes $\sigma_x \sigma_y$ $A = 4 \pi \sigma_x \sigma_y$

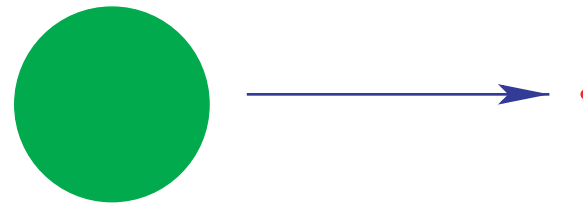
High **Luminosity** : $N \uparrow$ collide many particles, $A \downarrow$ squeezed in small bunches

LHC 1.15×10^{11} protons, $n_b = 2808$ ($f \uparrow$ crossings at 25 ns intervals)

Beams squeezed using strong large aperture quadrupoles around the interaction points

from ~ 0.2 mm to

$$\sigma_x = \sigma_y = 17 \mu\text{m}$$



$$\langle \beta \rangle_{\text{arc}} = 80 \text{ m}$$

$$\beta_{\text{IP}} = 0.5 \text{ m}$$

Rare new processes, like Higgs production can have very small cross section, like $1 \text{ fb} = 10^{-39} \text{ cm}^2$. LHC designed for very high Luminosity $\mathbf{L} = 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

Event rate for such rare processes : ~ 1 new particle every 28h.

Instead pp $\sigma_{\text{tot}} \approx 0.1$ barn 30 / crossing

Alternate gradient focusing

**Quadrupole lens
focusing in x,
defocusing in y
or vice versa**

$$\mathbf{F} = e (\mathbf{v} \times \mathbf{B})$$

here

$$\mathbf{F} = e (0, 0, v) \times (B_x, B_y, 0)$$

$$= e (-v B_y, +v B_x, 0)$$

Combine F D

Defocusing when at

small amplitude

Overall focusing

Normal (light) optics :

Focal length of two lenses

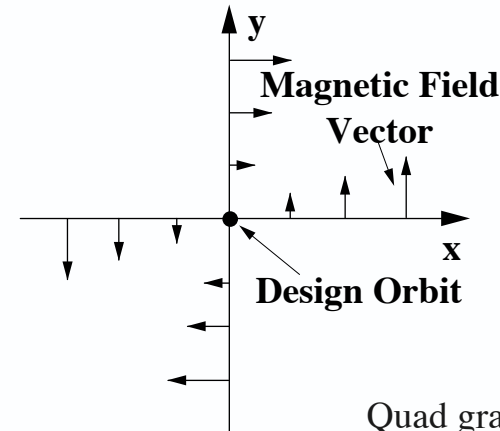
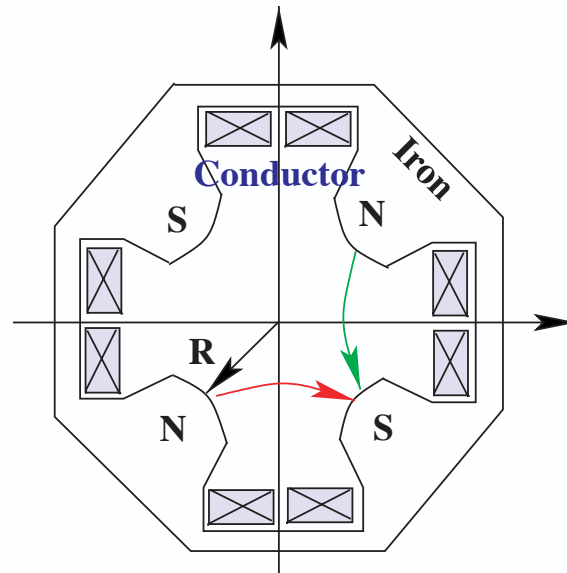
at distance D

$$1/f = 1/f_1 + 1/f_2 - D/f_1 f_2$$

is overall focusing

with $1/f = D/f^2$

for $f = f_1 = -f_2$



$$B_x = k y$$

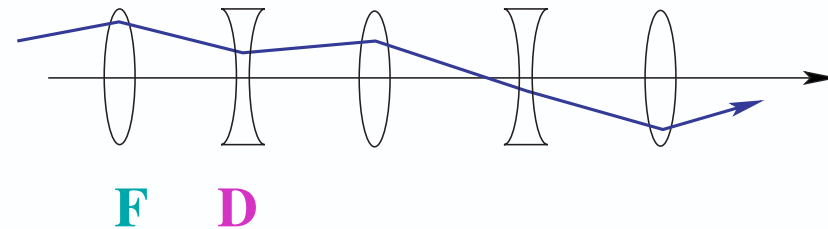
$$B_y = k x$$

$$B_z = 0$$

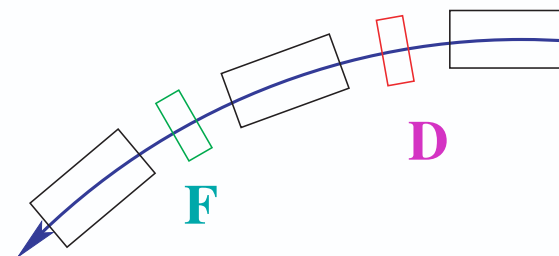
$$\nabla \times \mathbf{B} = \mathbf{0}$$

Quad gradients in the LHC
 $K = 1/B_0 \partial B_y / \partial x \approx 200 \text{ T/m}$

**alternate gradient
focusing**



**together with
bending magnets
FODO lattice**



N. C. Christofilos, unpublished manuscript in 1950 and patent

Courant, Snyder in 1952, Phys. Rev. 88, pp 1190 - 1196 + longer review in [Annals of Physics 3 \(1958\)](#)

Betatron motion

Equation of motion of particles in a ring (with bending fields) **and quadrupoles** (field gradients $\propto \partial B / \partial r$)

In both transverse planes, here written with x for x, y : known as Mathieu-Hill equation

$$x''(s) + k(s) x(s) = 0, \quad \text{derived in 1801 to describe planetary motion}$$

Generalised oscillator equation with position dependent, periodic restoring force $k(L+s) = k(s)$ given by the quadrupole gradients (+ the small weakly focusing bending term in the ring plane)

Solution : $x(s) = \sqrt{\epsilon \beta(s)} \cos(\mu(s) + \phi)$

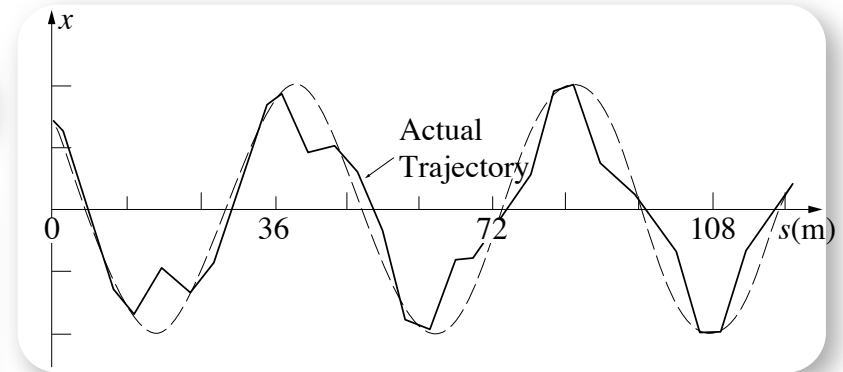
Phase advance

Lyapunov-Floquet Transformation

Tune # of betatron oscillations

$$\mu(s) = \int_0^s \frac{ds}{\beta(s)}$$

$$Q = \mu / 2\pi$$



*motion $x/\sqrt{\beta}$ plotted with phase advance
normalised coordinates - becomes simple cos*

$\beta(s)$ **beta function**, describes the focusing properties of the magnetic lattice

\mathcal{E} invariant, together with $\beta(s)$ amplitude. “single particle emittance”

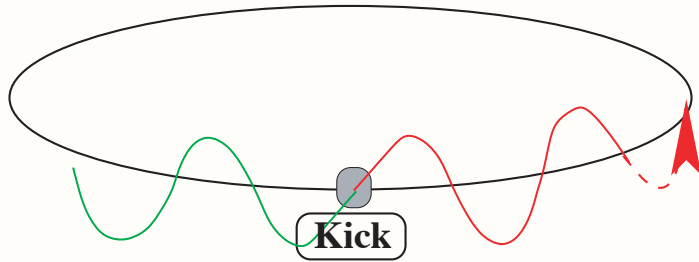
Motion conveniently described in phase space (x, x') where $x' = p_x / p$

and linear optics elements as matrices ; with simple case for M, applies for IP to IP

$$\begin{pmatrix} x(s) \\ x'(s) \end{pmatrix} = \mathbf{M} \begin{pmatrix} x(s_0) \\ x'(s_0) \end{pmatrix} \quad \mathbf{M} = \begin{pmatrix} \cos 2\pi Q & \beta \sin 2\pi Q \\ -\frac{1}{\beta} \sin 2\pi Q & \cos 2\pi Q \end{pmatrix}$$

Accelerator design : starts with magnet lattice based on linear beam optics ; MAD program

Orbit stability and tune



Misalignments and dipole field errors

→ **orbit perturbations**

would add up on successive turns

for integer tune $Q = N$

Higher order field errors,

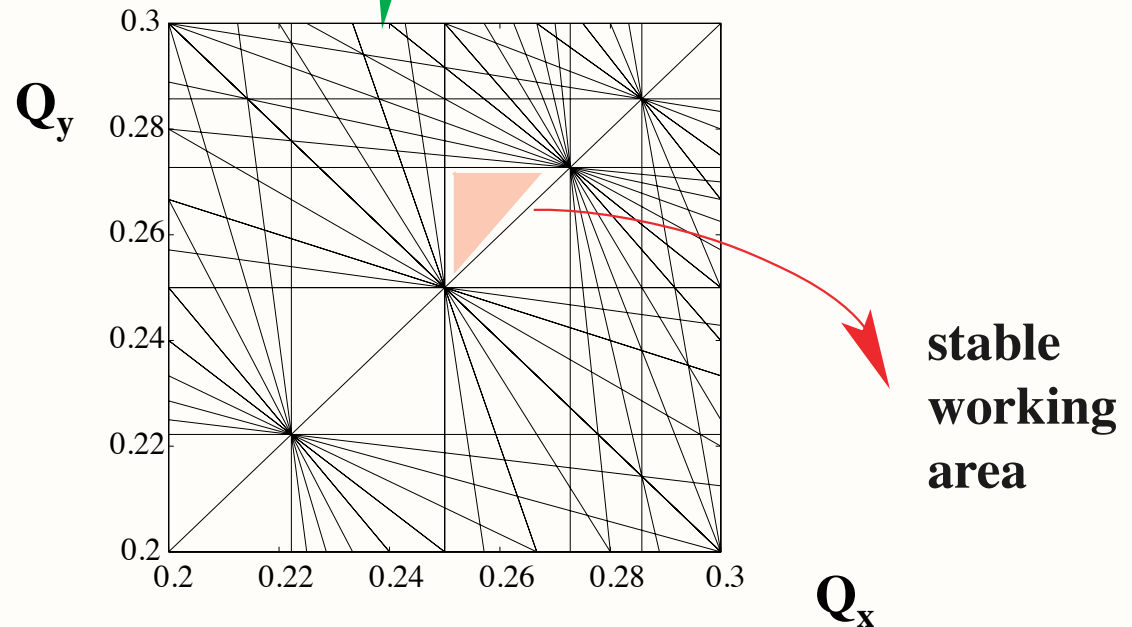
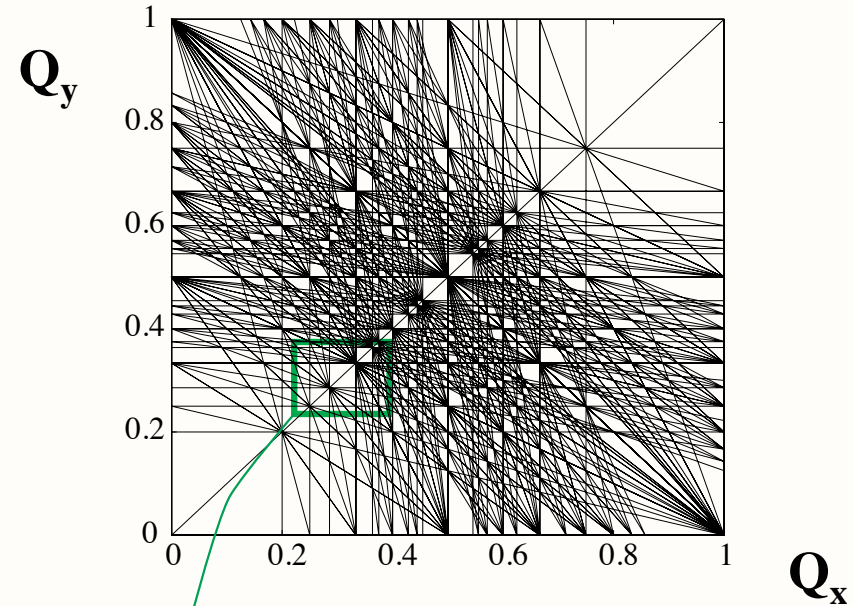
Quad., Sext. perturbations.

Avoid simple fractional tunes

$nQ_x + mQ_y + mQ_s = \text{int.}$

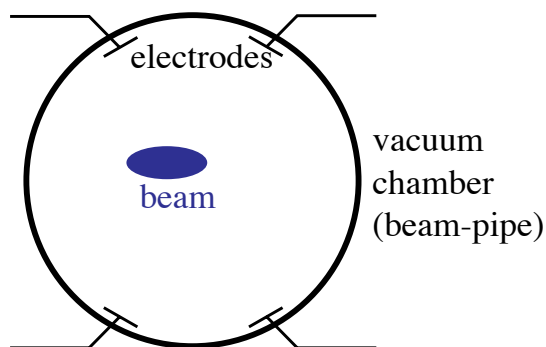
Minimise field and alignment

errors

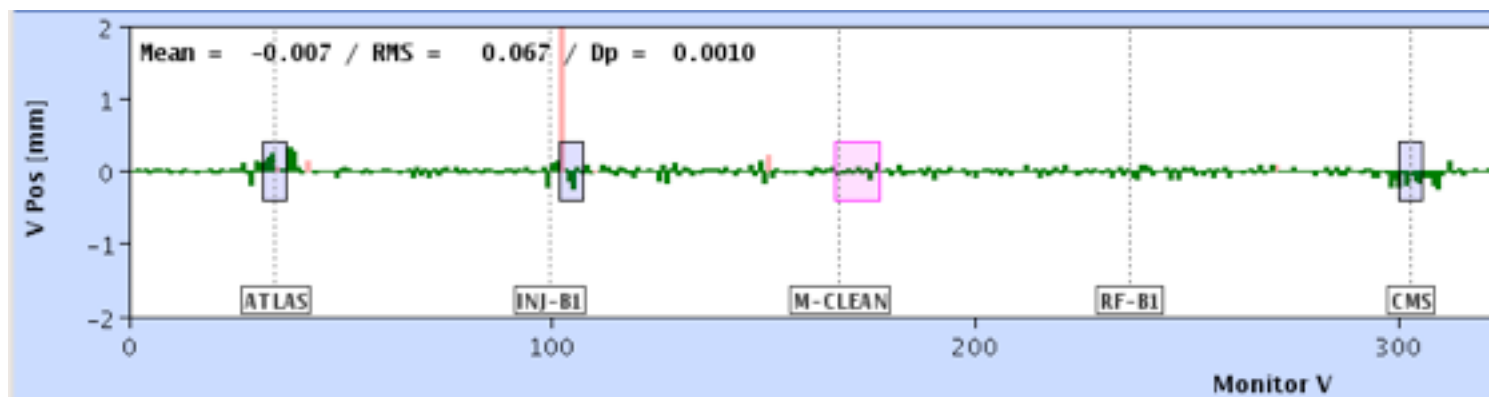


Orbit, tune measurement and peak beam current

vertical orbit, June 2011, 1st half of LHC shown



Beam Pickup Monitor



$\langle I_b \rangle$ average ring
and
 \hat{I} local peak
current

$$\hat{I} = \frac{\langle I_b \rangle L}{\sqrt{2\pi} \sigma_z}$$

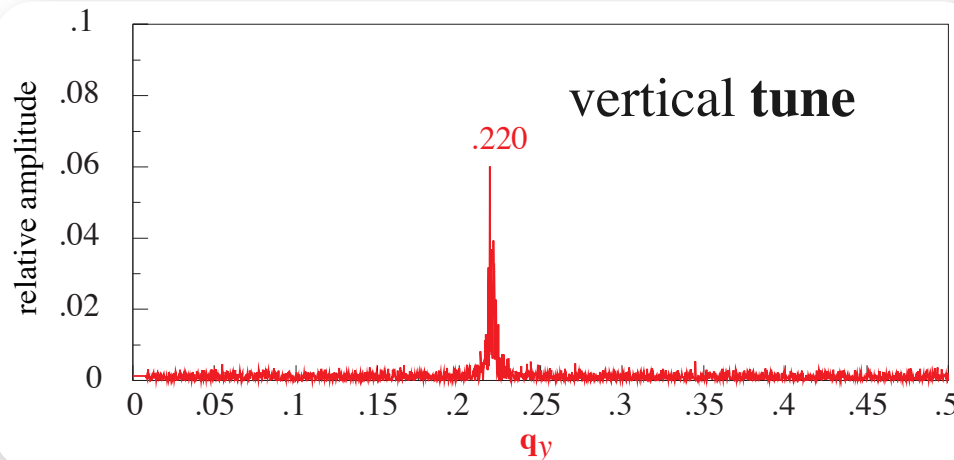
Typical numbers, for a single bunch $\langle I_b \rangle = n e f_{\text{rev}}$

LEP	$n = 4 \times 10^{11}$	$\langle I_b \rangle = 0.72 \text{ mA}$	$\sigma_z = 2 \text{ cm}$	$\hat{I} = 960 \text{ A}$
LHC	$n = 1.15 \times 10^{11}$	$\langle I_b \rangle = 0.21 \text{ mA}$	$\sigma_z = 7.55 \text{ cm}$	$\hat{I} = 73.2 \text{ A}$

$f_{\text{rev}} = 11245 \text{ kHz}, \quad L = 26658.9 \text{ m}$

Bunch peak currents are many Amperes !
Strong signals, used to monitor beam position and oscillations

Also source of undesirable effects :
wake fields, heating, instabilities



Transverse beam size and emittance

consider : beam of many particles on stable orbit and

simple case : dispersion and slope $\beta' = 0$ by default at IP - relevant for experiments

beam size, r.m.s. $\sigma(s) = \sqrt{\varepsilon \beta(s)}$

beam divergence, r.m.s. $\theta(s) = \sqrt{\varepsilon / \beta(s)}$

product $\varepsilon = \sigma(s) \theta(s)$

β - function : local machine quantity - focusing of lattice

Emittance ε : beam quantity - the average action

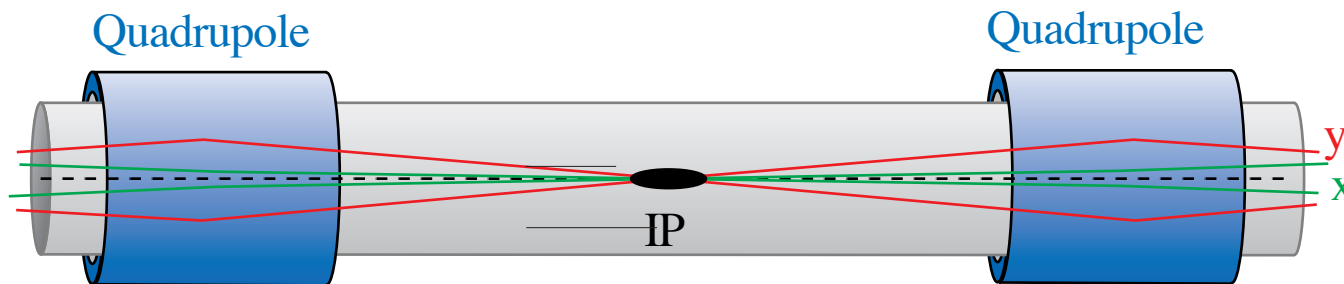
related to phase space density or kind of beam temperature

given by initial conditions (injected beam)

or equilibrium of quantum excitation and damping - 2nd lecture

in ideal machine : x, y, z motion uncoupled, 3 emittances $\varepsilon_x, \varepsilon_y, \varepsilon_z$

IP: squeeze β to a minimum, called β^* \Rightarrow maximum of divergence, needs aperture



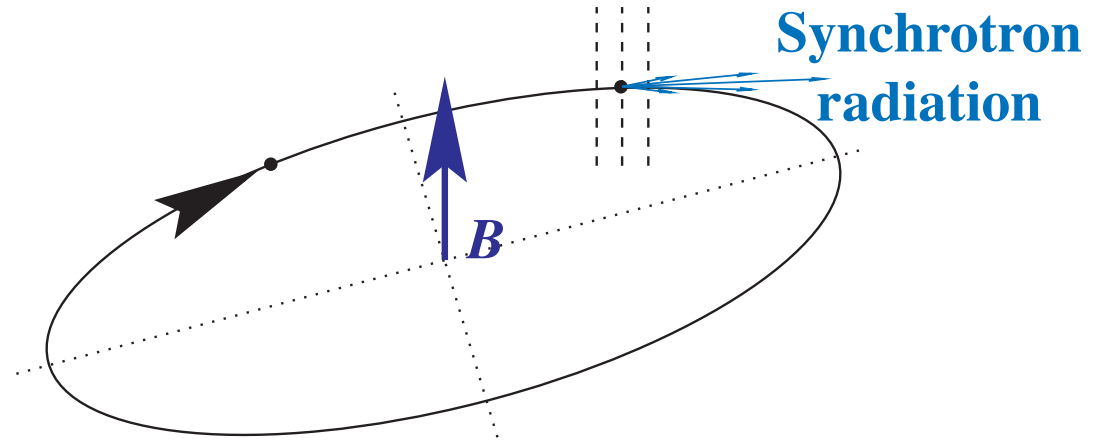
LHC $\varepsilon_N = \varepsilon \beta \gamma = 3.75 \mu\text{m}$, at top $E_b = 7 \text{ TeV}$: $\varepsilon = 0.503 \text{ nm}$, $\beta^* = 0.55 \text{ m}$, $\sigma^* = 16.63 \mu\text{m}$, $\theta^* = 30 \mu\text{rad}$

Standard Synchrotron Radiation

$$E_c = \frac{3}{2} \frac{\hbar c \gamma^3}{\rho} = 2.96 \times 10^{-7} \text{ eV m} \frac{\gamma^3}{\rho}$$

$$U_0 = \frac{e^2}{3\epsilon_0} \frac{\gamma^4}{\rho} \approx 6.0317 \cdot 10^{-9} \text{ eV m} \frac{\gamma^4}{\rho}$$

$$P_b = \frac{U_0 I_b}{e}$$



		E GeV	γ	ϱ m	U_0 MeV	E_c keV	τ_d s	N 10^{12}	I mA	P_b MW	B T
RHIC	Au	A×100	107.4	242.8	21×10^{-6}	1.5×10^{-6}	4.9×10^6	0.06	60	1.3×10^{-12}	3.42
LHC	p	7000	7460.5	2804	0.0067	0.044	61729	646	1163	0.0072	8.33
LEP1	e	45.6	89237	3026	126	69.5	23×10^{-3}	2.22	4	0.5	0.05
LEP2	e	104.5	204501	3026	3490	836	1.9×10^{-3}	2.8	5	18	0.115

Same beam energy E and radius ϱ : electron instead of proton $U_0 \sim \gamma^4 : (m_p/m_e)^4 = 1.13 \times 10^{13}$

Electrons, $E \gg 100$ GeV needs linear collider (ILC / CLIC)

Damping time E / U_0 turns or $\tau_d = t_{\text{rev}} E / U_0$ revolution time LEP/LHC $t_{\text{rev}} = 88.9 \mu\text{s}$

Gold ions Au⁷⁹⁺ A=197 $\langle E_\gamma \rangle = 8/(15\sqrt{3}) E_c \quad 8/(15\sqrt{3}) \approx 0.308$

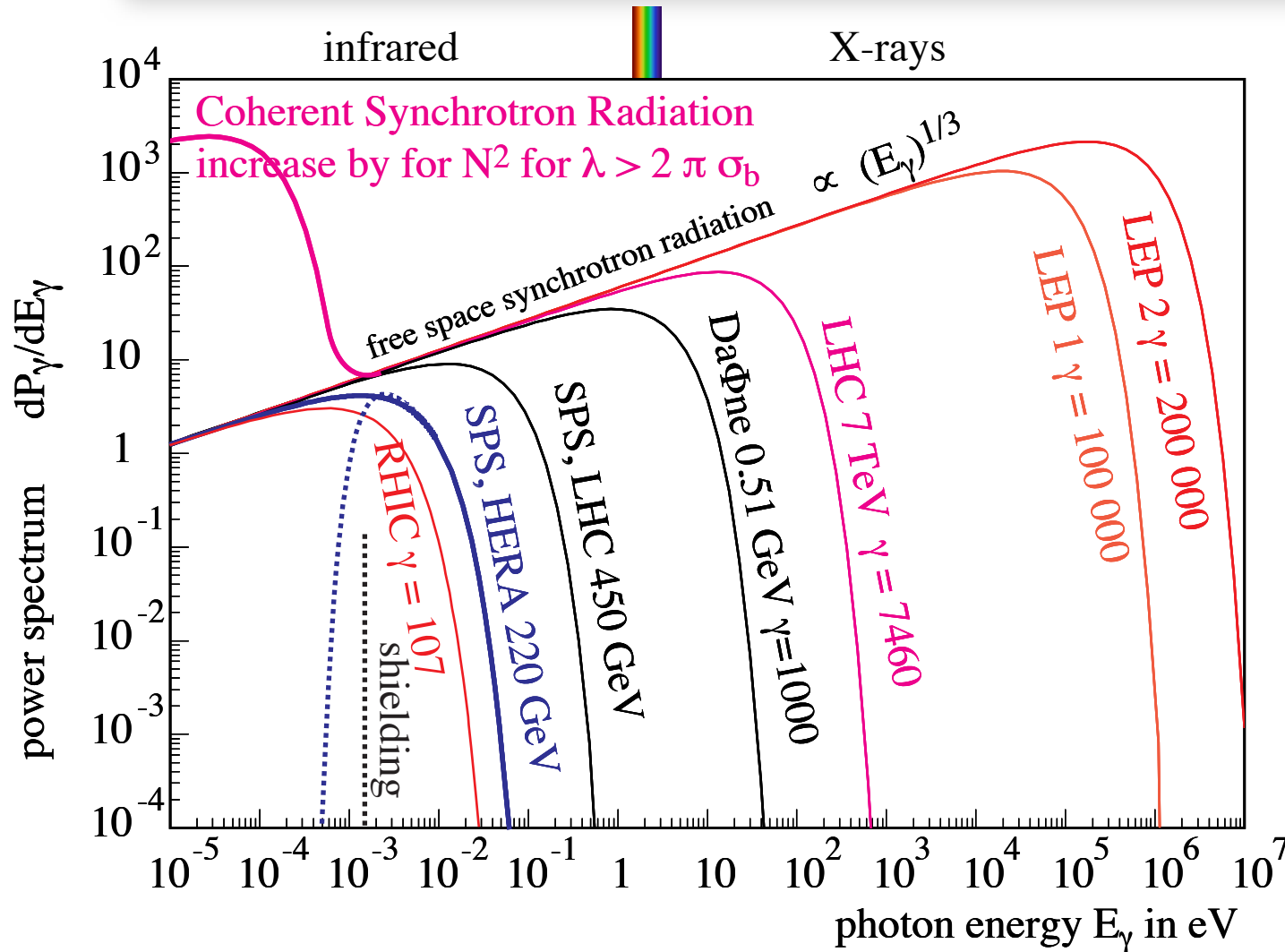
Synchrotron light monitor

Picture from
LEP. Typical
transverse
rms beam size
0.15 mm vertical
1.5 mm horiz.



Mirror, small slit, telescope and camera : beams continuously visible.
Now also used for protons in the LHC.

Power Spectrum, Free space, Cutoff and CSR



$$\frac{f_{\text{cutoff}}}{f_{\text{rev}}} = \sqrt{\frac{2}{3}} \left(\frac{\pi\rho}{h} \right)^{3/2}$$

ρ bending radius
 h chamber height
 cutoff relevant
 for $\gamma \approx 100$

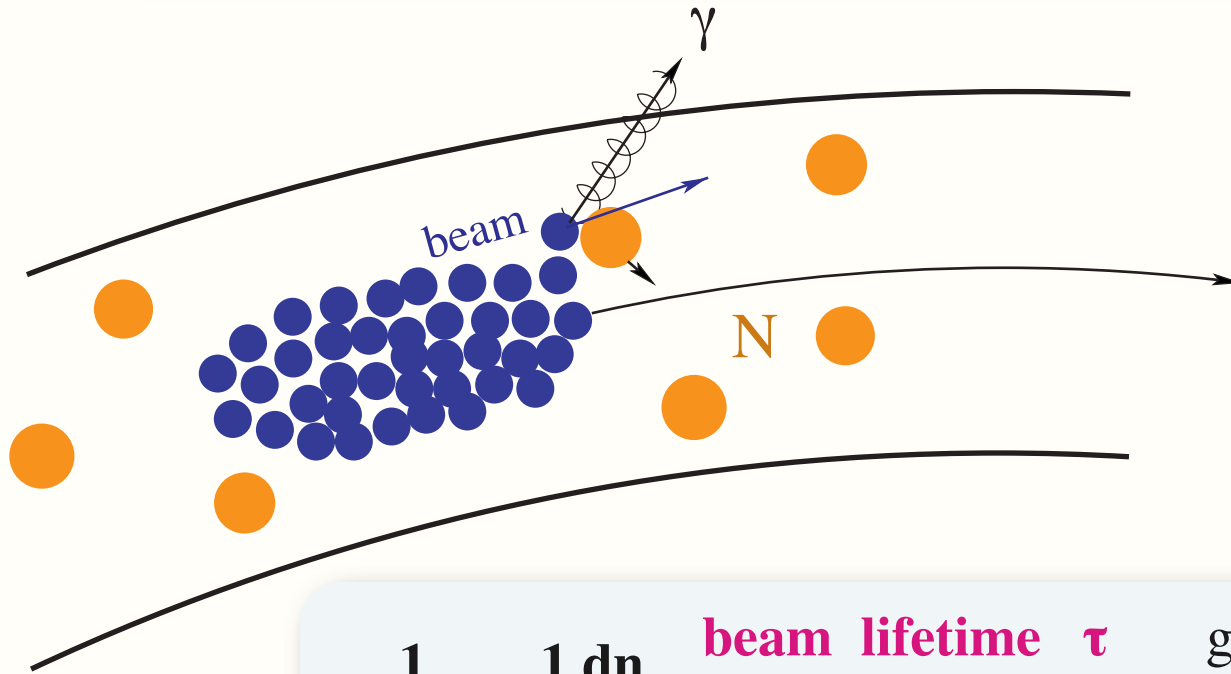
12 orders of magnitude
 in E_γ and λ
 10^{-5} eV $\lambda = 0.124$ m
 10^{+7} eV $\lambda = 124$ fm

Effects which can modify the low energy, long wavelength spectrum :

- i) **Coherent Synchrotron Radiation CSR** increases radiation and loss
- ii) **Boundary conditions - cutoff by conducting chamber** decreases radiation and loss

Energy Loss of Gold Ions in RHIC, [EPAC 2008](#)

Vacuum, beam Gas - lifetime



Beam blow up, core + halo
Background to experiments
loss, radiation, beam and
Luminosity lifetime

Minimize effect :
Good vacuum
O(nTorr or 10^{-9} mb)
Collimation

$$\frac{1}{\tau} = - \frac{1}{n} \frac{dn}{dt}$$

beam lifetime τ general expression
 average time between collisions leading to beam loss
 inverse normalised loss rate

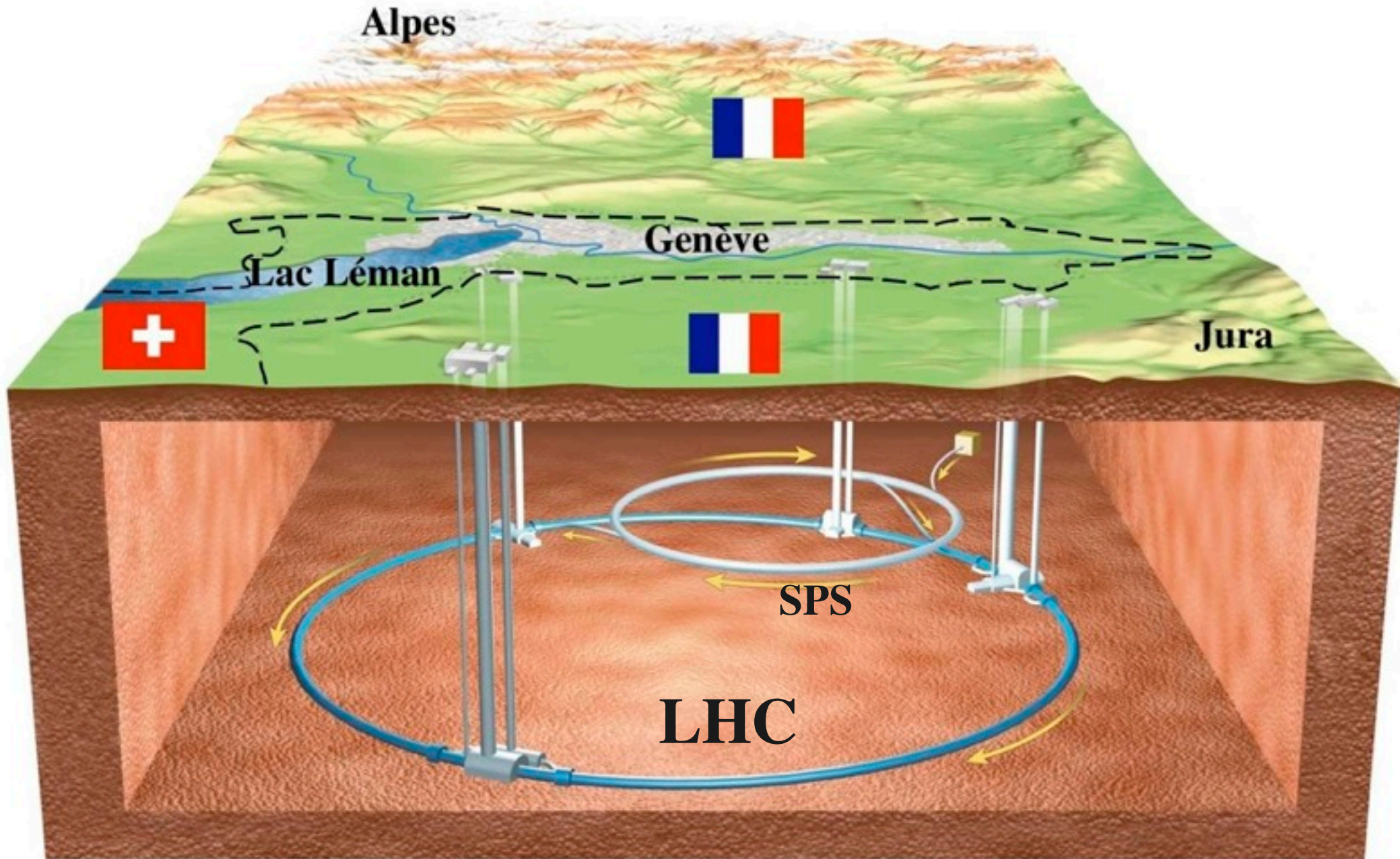
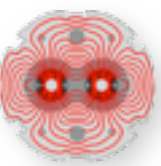
$$p = 1 \text{ ntorr} = 1.33 \times 10^{-7} \text{ Pa}$$

$$\rho_m = \frac{p}{kT} = 3.26 \times 10^{13} \text{ molecules / m}^3$$

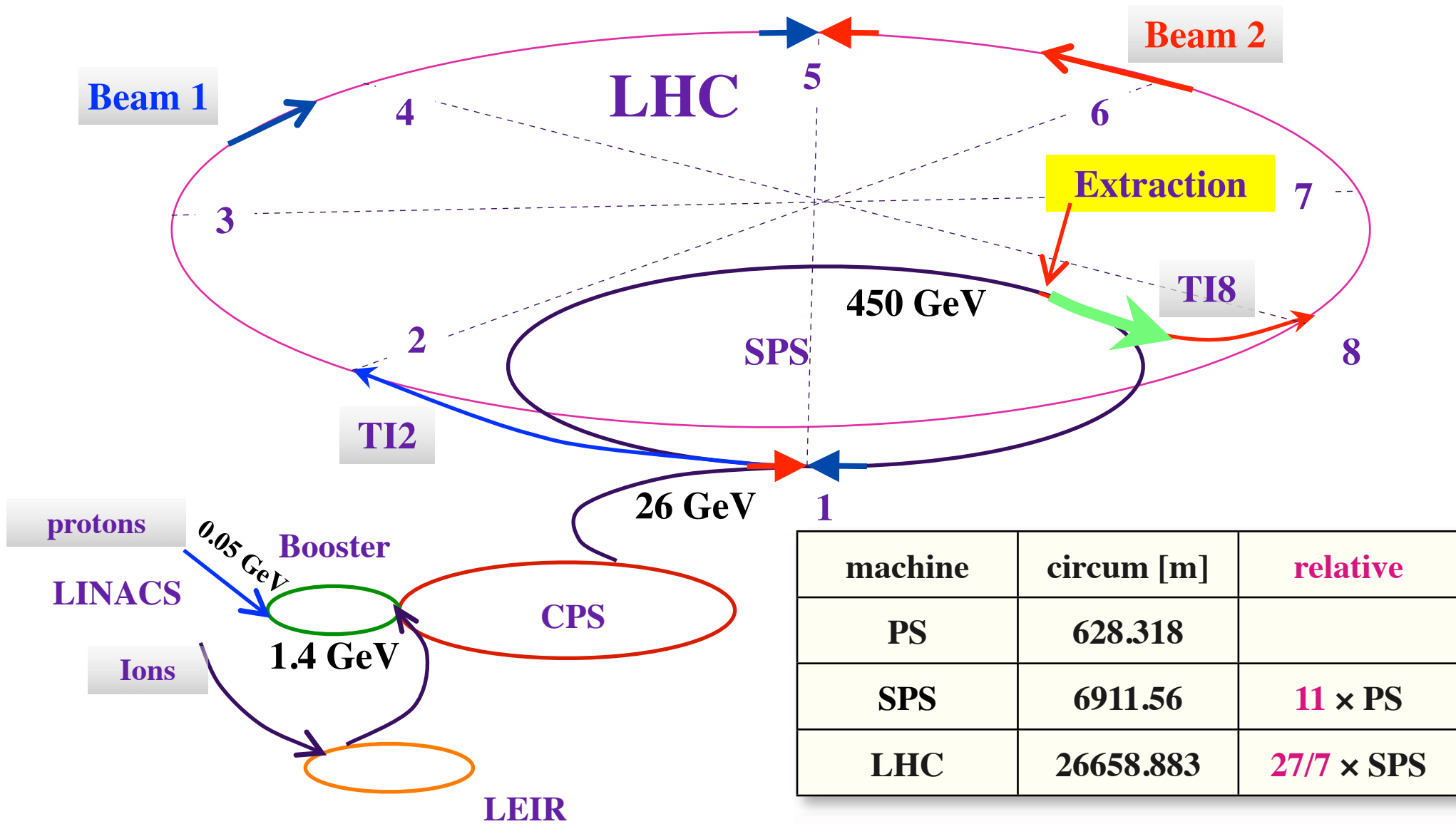
$$\text{typical cross section } \sigma = 6 \text{ barn} = 6 \times 10^{-28} \text{ m}^2$$

$$\text{collision probability } P_{\text{coll}} = \sigma \rho_m = 1.96 \times 10^{-14} / \text{m}$$

$$\tau = \frac{1}{P_{\text{coll}} c} = 1.7 \times 10^5 \text{ s} = 47 \text{ hours} \quad \text{for } v \approx c$$



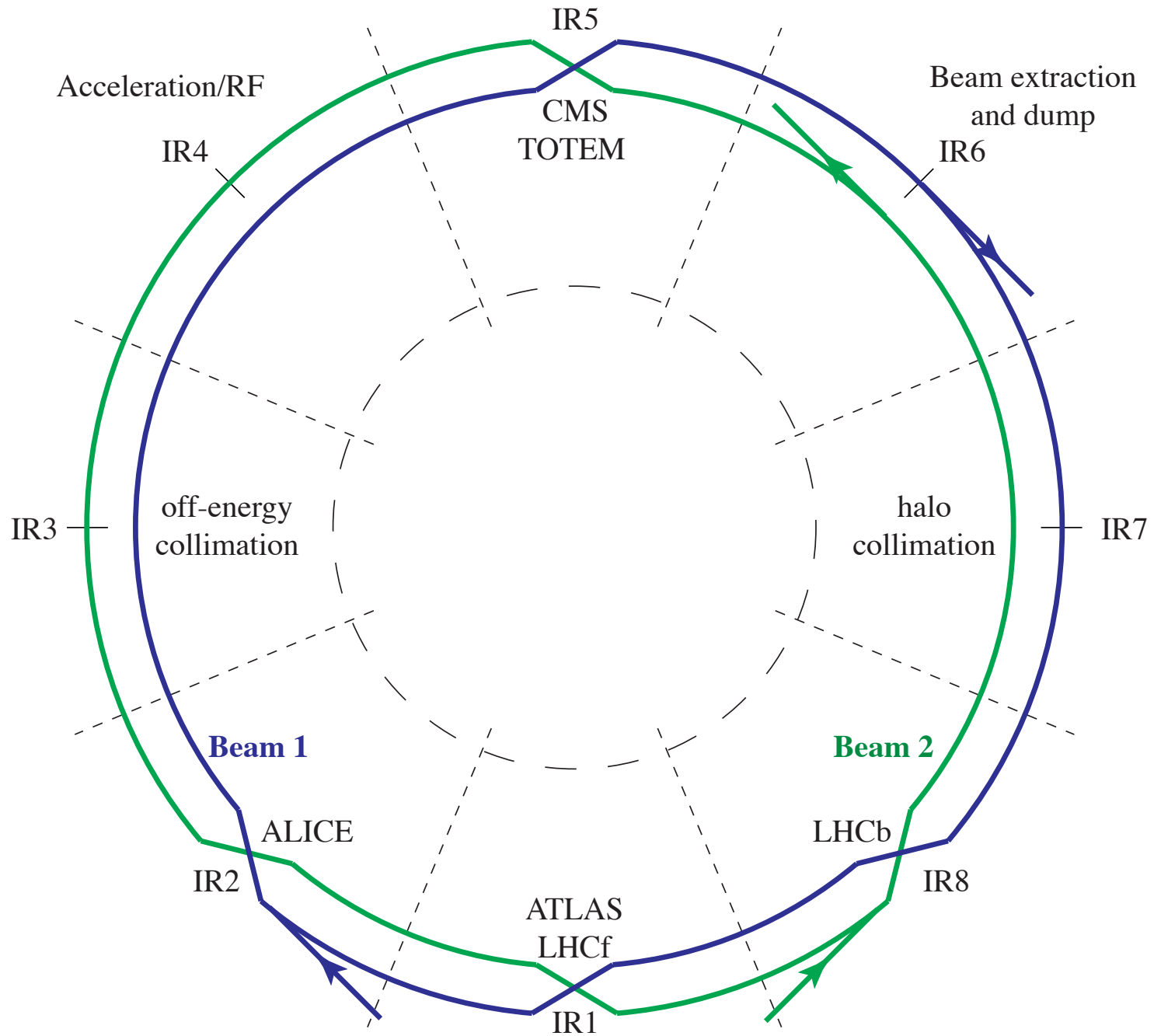
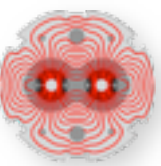
The CERN accelerator complex : injectors and transfer



Beam size of protons decreases with energy : area $\sigma^2 \propto 1 / E$
 Beam size largest at injection, using the full aperture

simple rational fractions for **synchronization**
 based on a single frequency
 generator at injection

Layout of the LHC



10 September 2008

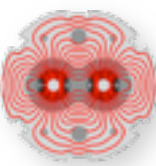


10:30 beam 1 3 turns
15:00 beam 2 3 turns
22:00 beam 2 several 100 turns





LHC status



- main LHC challenge : damage potential --- **increase safely (slowly) the intensity**
- **enormous stored energy** : nominal is 10 GJ in magnets, 362 MJ in beam; 0.7 MJ melts 1kg Cu
- currently 3.3 GJ in magnets, 130 MJ in beam

LHC :

2009 first collisions, mostly at injection energy 2x450 GeV

2010 2x3.5 TeV, $\beta^* = 3.5$ m, $L_{\text{peak}} = 0.2 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$ $\int L dt = 0.044 \text{ fb}^{-1}$ 368 bunches

2011 2x3.5 TeV, $\beta^* = 1.0$ m, $L_{\text{peak}} = 3.5 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$ $\int L dt = 6.1 \text{ fb}^{-1}$ 1380 bunches

2012 2x4.0 TeV, $\beta^* = 0.6$ m, $L_{\text{peak}} = 7.7 \times 10^{33} \text{ cm}^{-2}\text{s}^{-1}$ $\int L dt = 23.3 \text{ fb}^{-1}$ 1380 bunches

2013 Pb-p run, **shutdown**, magnet interconnects, restart in 2015 at 2x6.5 TeV, increase #bunches

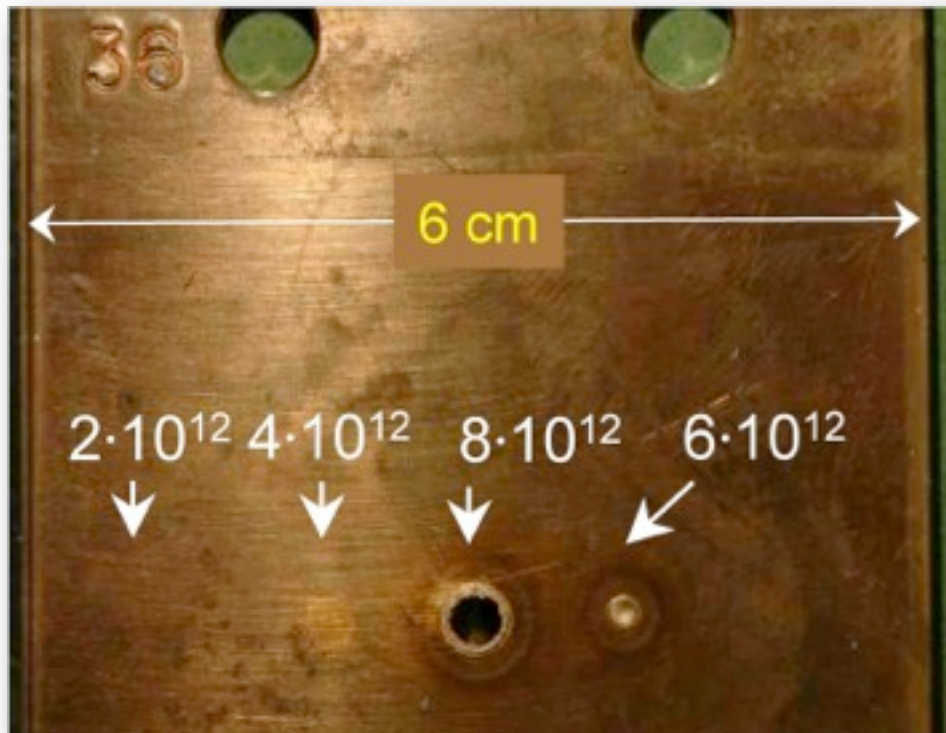
	LHC design	achieved
Momentum at collision, TeV/c	7	4
Luminosity, $\text{cm}^{-2}\text{s}^{-1}$	1.0E+34	7.7E+33
Dipole field at top energy, T	8.33	4.8
Number of bunches, each beam	2808	1380
Particles / bunch	1.15E+11	1.70E+11
Typical beam size in ring, μm	200 – 300	~300
Beam size at IP, μm	17	20

Damage potential : confirmed in controlled SPS experiment

controlled experiment with beam extracted from SPS at 450 GeV in a single turn, with perpendicular impact on Cu + stainless steel target

450 GeV protons →

r.m.s. beam sizes $\sigma_{x/y} \approx 1$ mm



SPS results confirmed :

8×10^{12} clear damage

2×10^{12} below damage limit

for details see V. Kain et al., PAC 2005 [RPPE018](#)

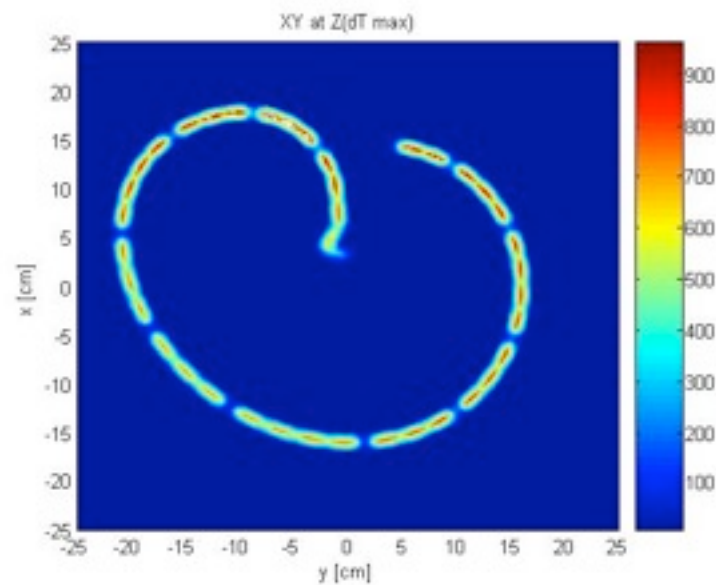
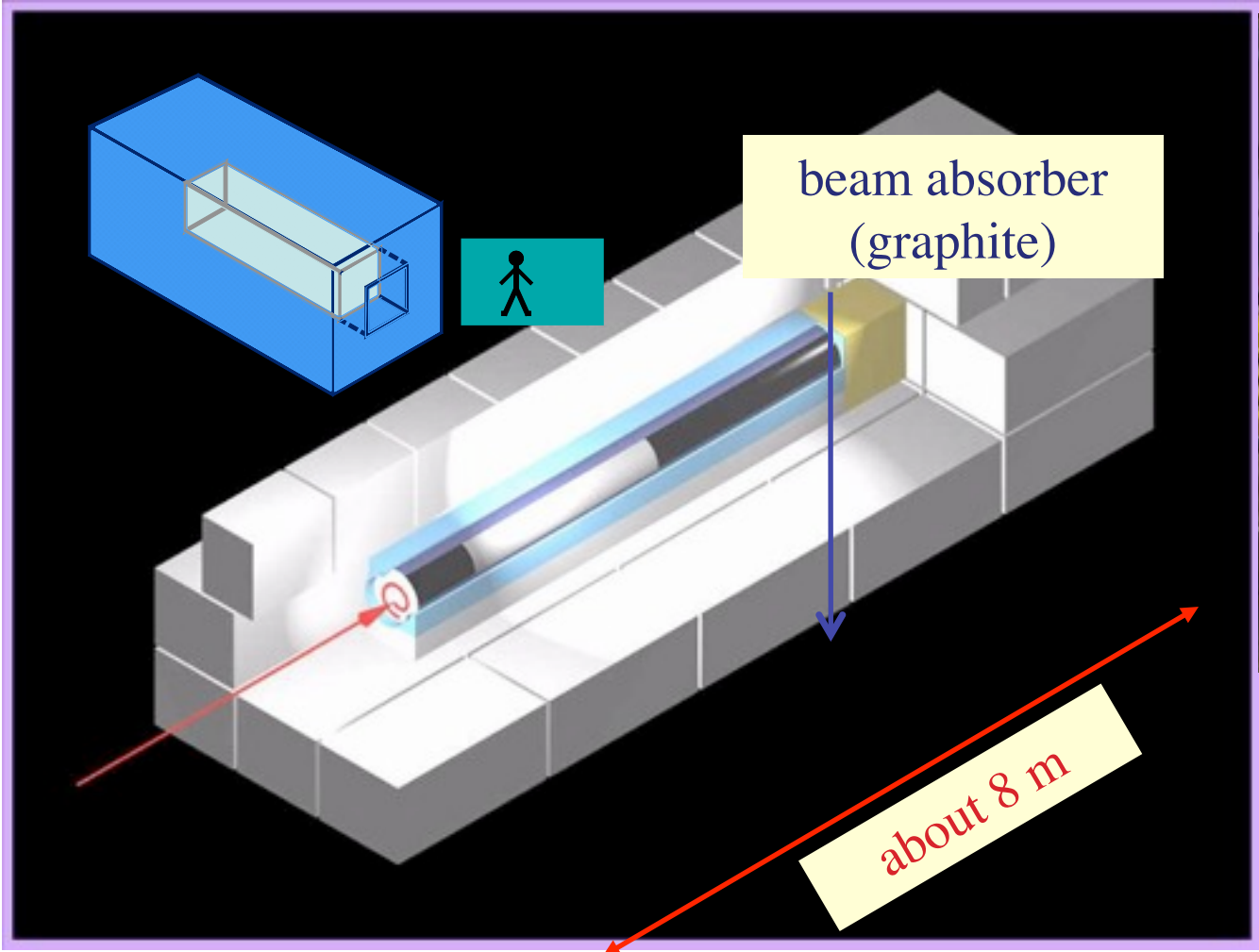
For comparison, the LHC nominal at 7 TeV :

$2808 \times 1.15 \times 10^{11} = 3.2 \times 10^{14}$ p/beam

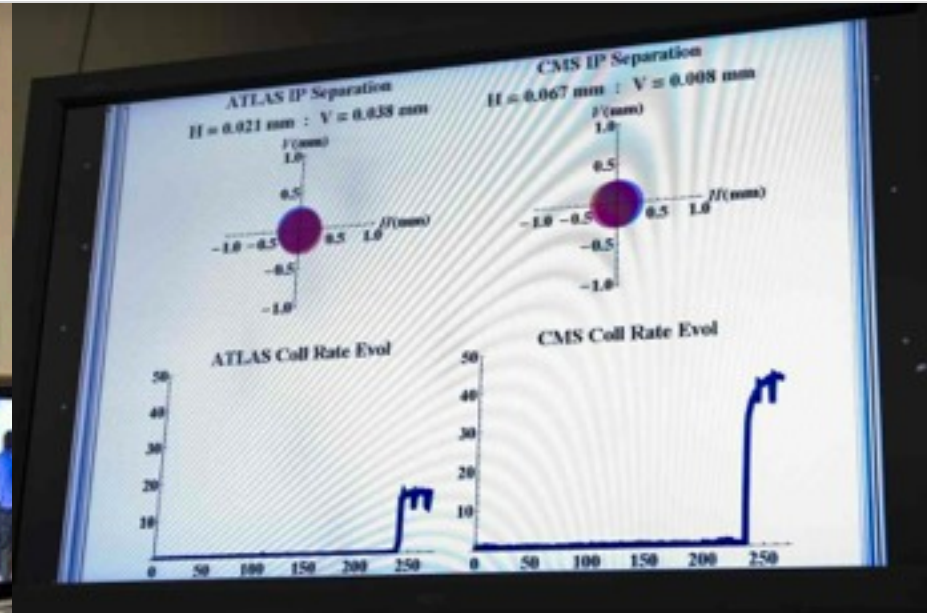
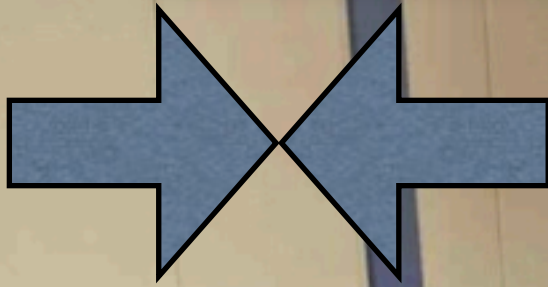
at $\langle \sigma_{x/y} \rangle \approx 0.2$ mm

over 3 orders of magnitude above damage level for perpendicular impact

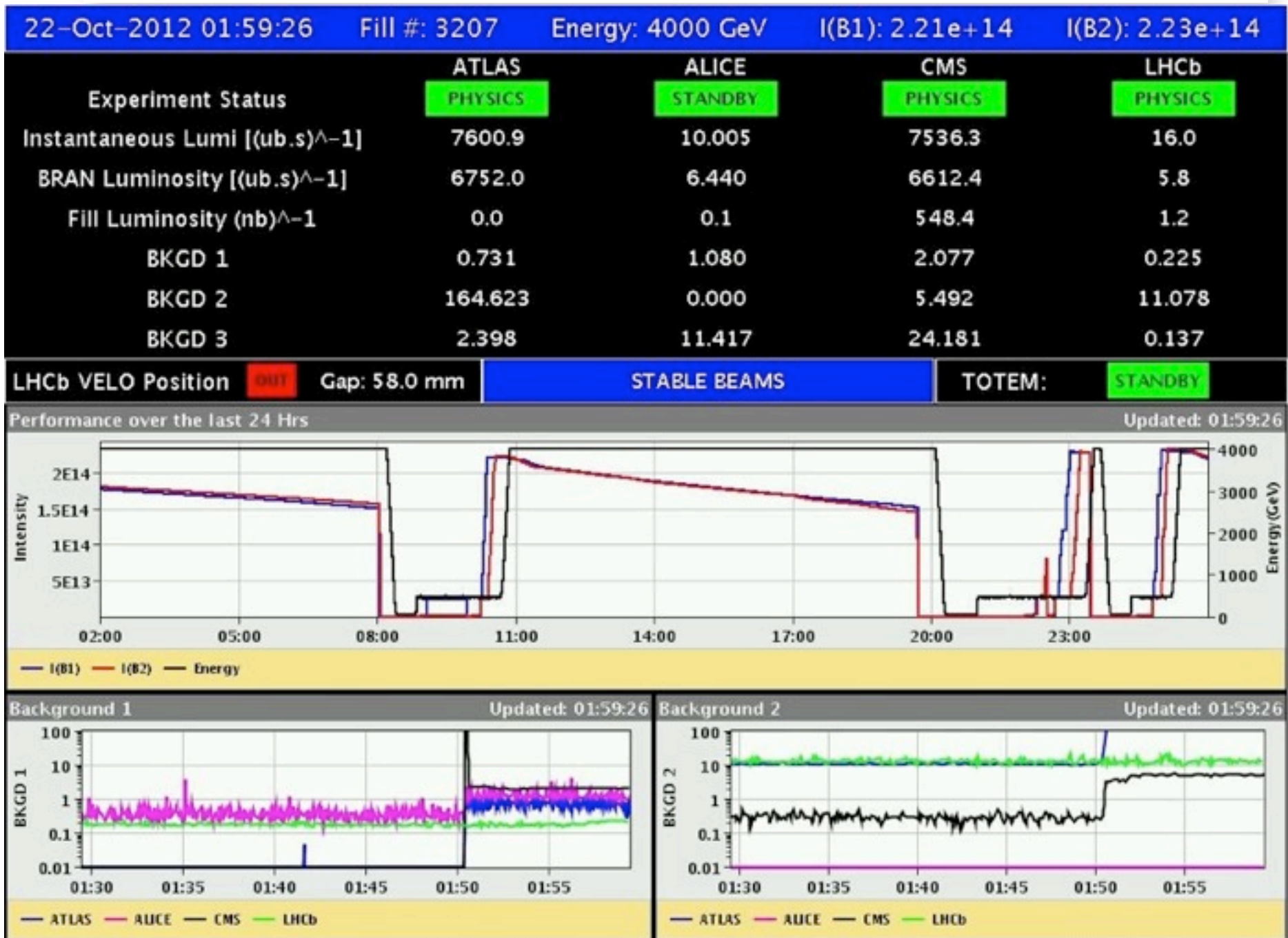
Dumping the LHC beam



First high energy 3.5TeV+3.5TeV collisions, 30 March 2010

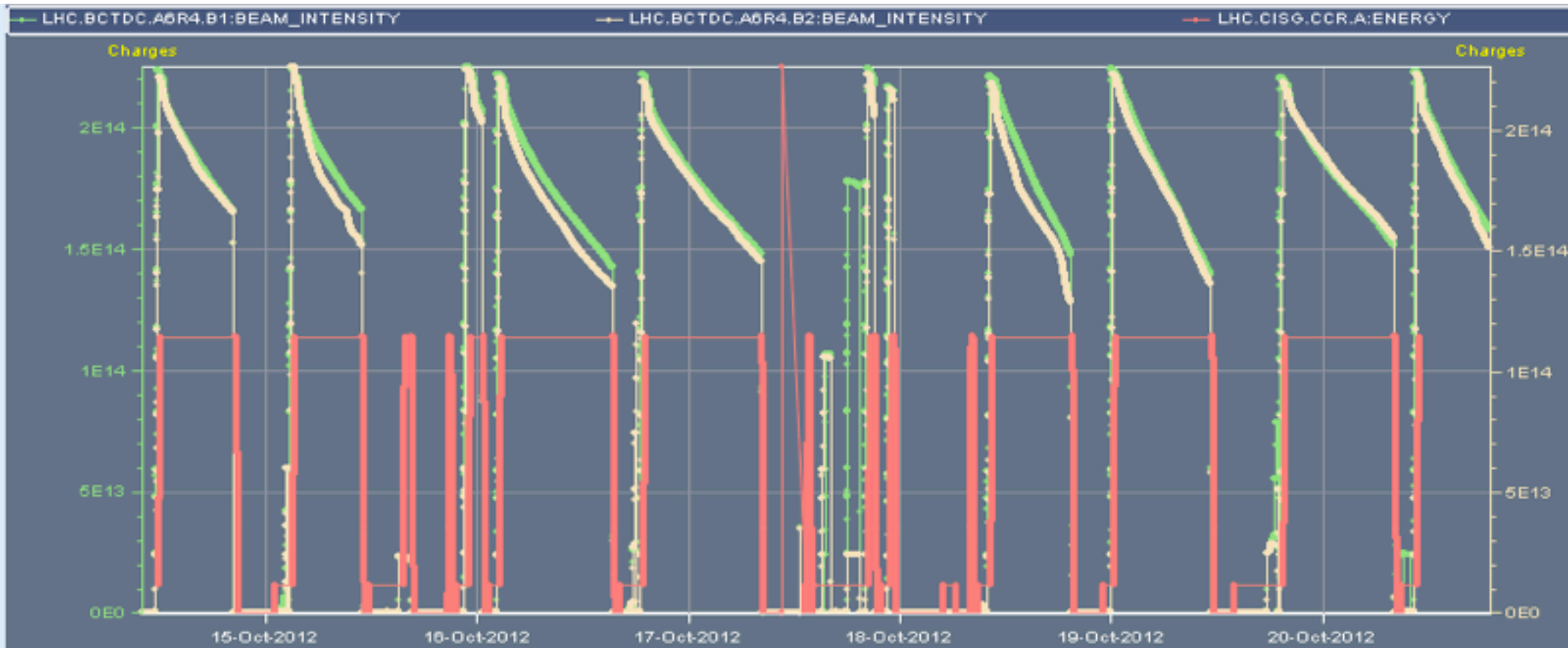


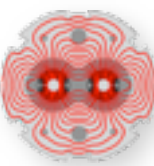
LHC running very well



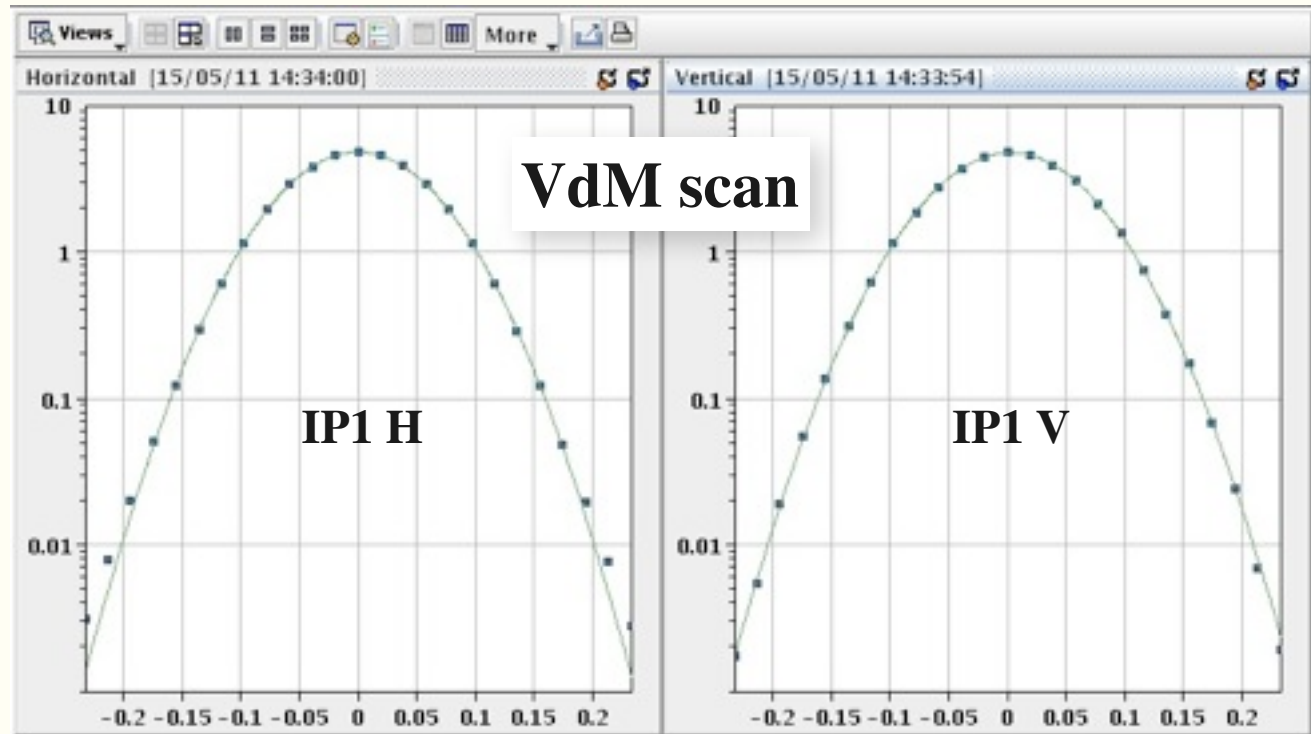
peak Luminosity $7.8 \times 10^{33} \text{cm}^{-2}\text{s}^{-1}$

LHC typical week, Oct. '12, 1.2 pb-1



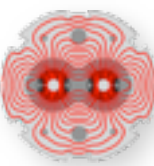


- absolute luminosity normalization
 - low, well understood backgrounds
 - precision optics for ATLAS-ALFA and TOTEM
- $\beta^* = 1000$ m, Oct.'12



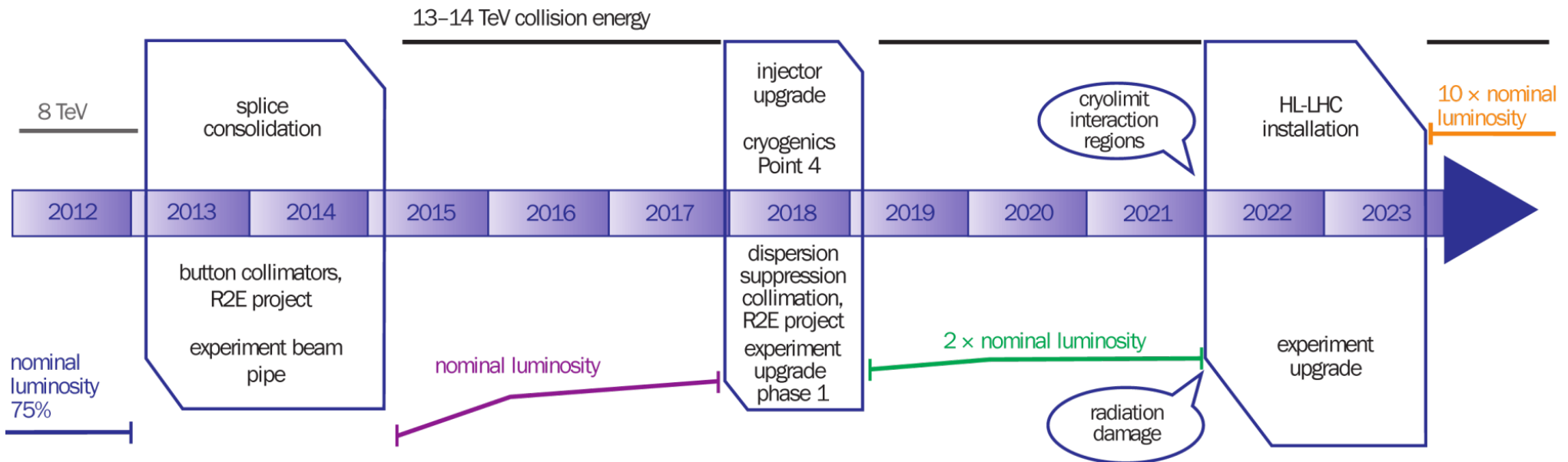
precise measurement of the luminous region + beam intensity --> absolute luminosity and cross section calibration

currently ~ 3 % level (Tevatron had ~ 15 %)



The LHC is still a rather young machine

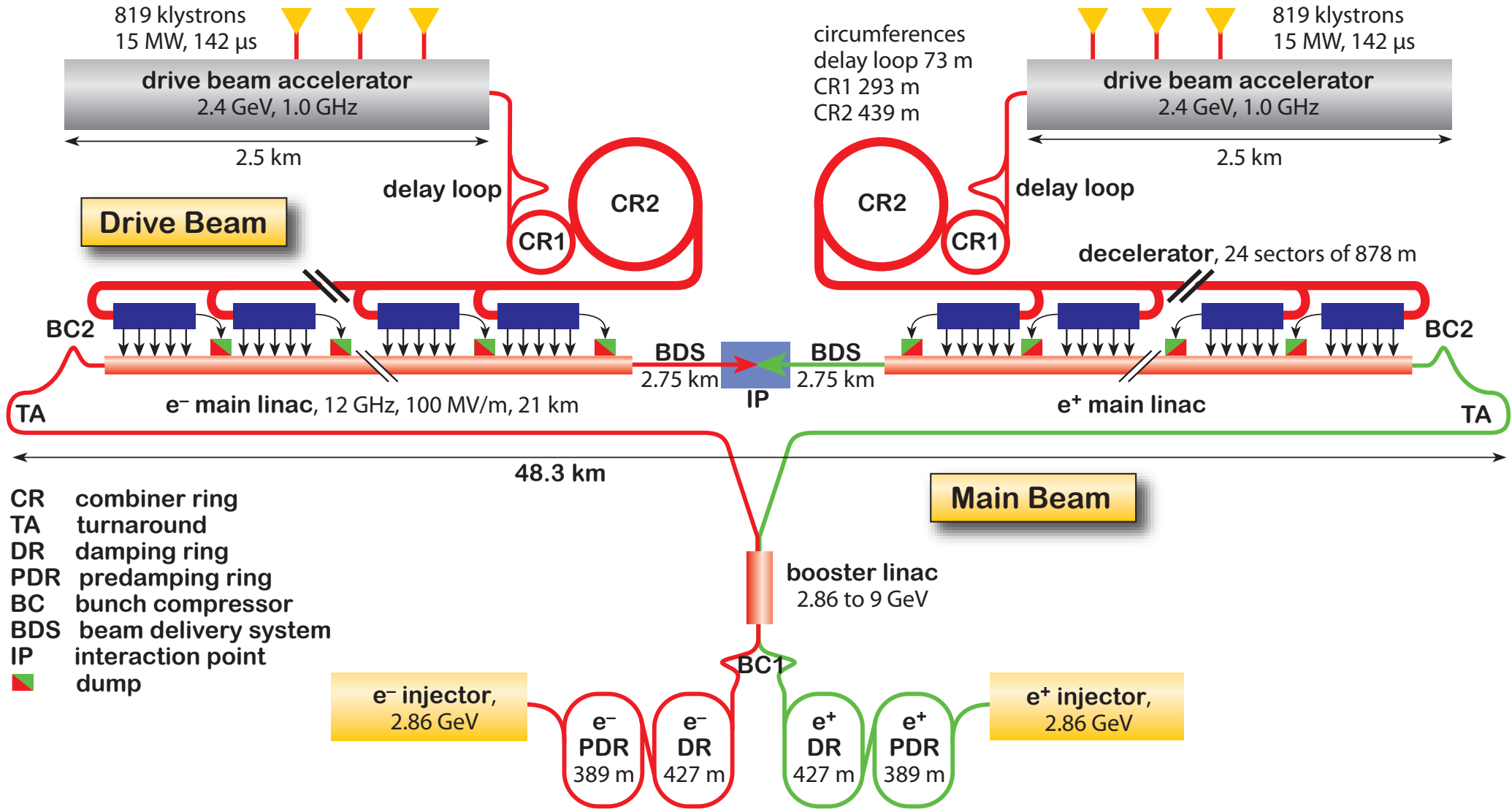
Operation planning + upgrade studies (HL-LHC) extend to ~ 2030



Further ideas already exist (HE-LHC, LHeC, TLEP)

We also study other machines, and in particular CLIC →

CLIC



Overview of the CLIC layout at $\sqrt{s} = 3$ TeV

The machine requires only one drive beam complex for stages 1 and 2.

- **The largest flag-ship accelerator is the LHC here at CERN**
- **By now many more accelerators outside particle physics**
#Accelerators in the world : O (30 000) mostly smaller for medical and industrial applications
- **Broad range of particle accelerator types and applications**

Large research facilities for :

Synchrotron light, UV, X-Ray (electron accelerators)

High intensity proton accelerators + neutron spallation sources

condensed matter, material science and biology research,

accelerator driven subcritical fission (energy production & radioactive waste incineration)

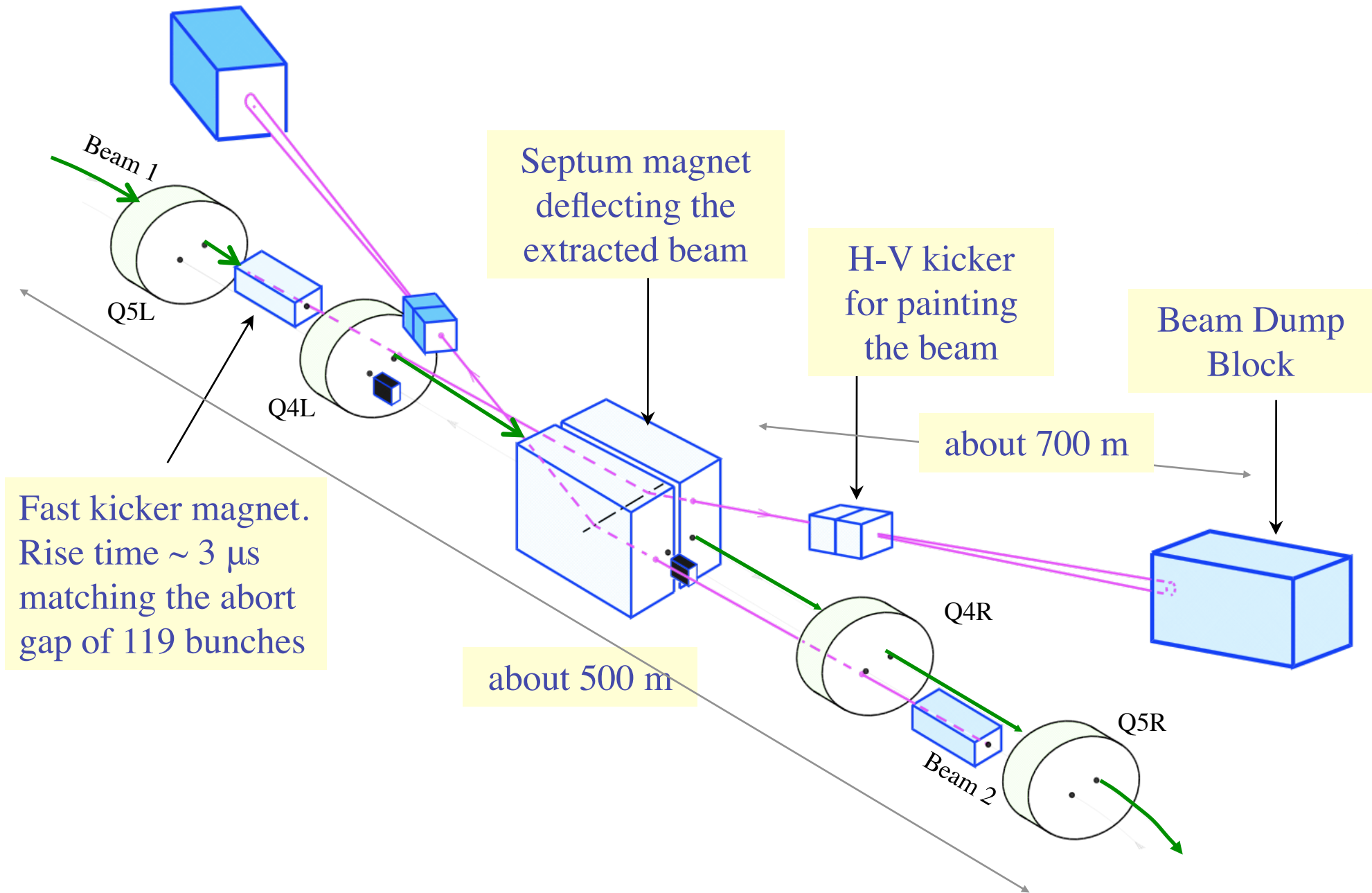
Yearly international accelerator conferences IPAC, last one in May'13 in [Shanghai](#)

Some of the hot-subjects and keywords :

- **Free electrons lasers FEL, X-FEL, Laser induced coherent SR**
- **Advanced LINACS -- including recirculation and energy recovery ERL**
- **New acceleration techniques :**
 - **Dielectric, LASER, Plasma driven**

Reserve

Schematic layout of beam dump system in IR6



Radiation of an accelerated Charge

General concept - power radiated by an accelerated charge. Relativistic version of Lamor's formula, derived by Lienard in 1898, before relativity was known.

Photon spectrum : J. Schwinger Phys. Rev. 75 (1949) pp. 1912-1925

Here written with formulas in SI units. More info + references in my paper on MC generation of [SynRad](#) CERN-OPEN-2007-018

power radiated by an accelerated charge

$$P = \frac{e^2 \gamma^2}{6\pi\epsilon_0 m^2 c^3} \left[\left(\frac{d\mathbf{p}}{dt} \right)^2 - \beta^2 \left(\frac{dp}{dt} \right)^2 \right]$$

relativistic
Lamor formula

results in a major energy loss for a ring at high γ

$\mathbf{v} \perp \dot{\mathbf{v}}$

$$\left(\frac{d\mathbf{p}}{dt} \right)^2 - \underbrace{\beta^2 \left(\frac{dp}{dt} \right)^2}_0 = \dot{\mathbf{p}}^2 \quad P = \frac{e^2}{6\pi\epsilon_0 m^2 c^3} \gamma^2 \dot{\mathbf{p}}^2$$

Perpendicular acceleration, B-field (or E_{\perp} field). Motion in circular machine.

$\mathbf{v} \parallel \dot{\mathbf{v}}$

$$\left(\frac{d\mathbf{p}}{dt} \right)^2 = \left(\frac{dp}{dt} \right)^2 \quad \left(\frac{d\mathbf{p}}{dt} \right)^2 - \beta^2 \left(\frac{dp}{dt} \right)^2 = \dot{p}^2 (1 - \beta^2) = \frac{\dot{p}^2}{\gamma^2}$$

$$P = \frac{e^2}{6\pi\epsilon_0 m^2 c^3} \dot{p}^2$$

Parallel acceleration, E-field, Linac case
cancellation, $1/\gamma^2$

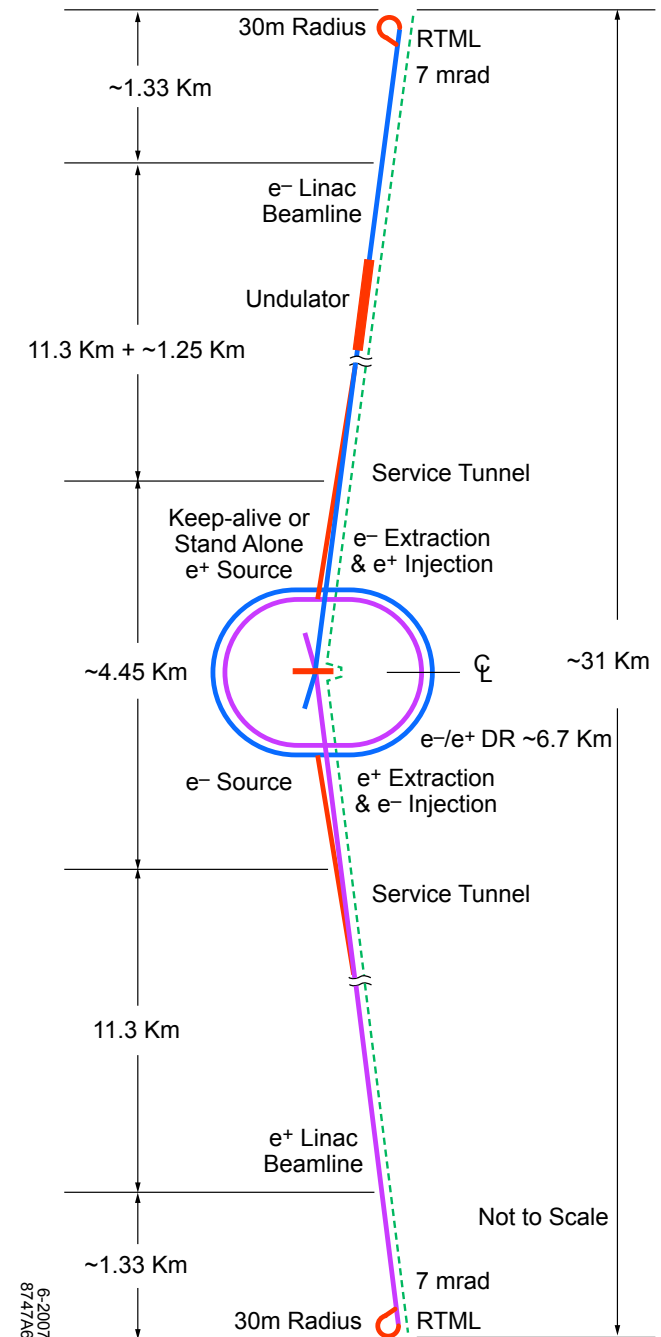
The energy loss for linear acceleration is very small.

Example: CLIC gradient 100 MV/m. Loss is 11 keV/s or only 0.4 eV for a 1 TeV 10 km Linac

ILC

ILC TDR Handover, 12 June 2013

- 200-500 GeV centre-of-mass, 31 km long
- Luminosity: $2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
- Based on accelerating gradient of 31.5 MV/m
1.3 GHz superconducting RF



Two Beam Scheme

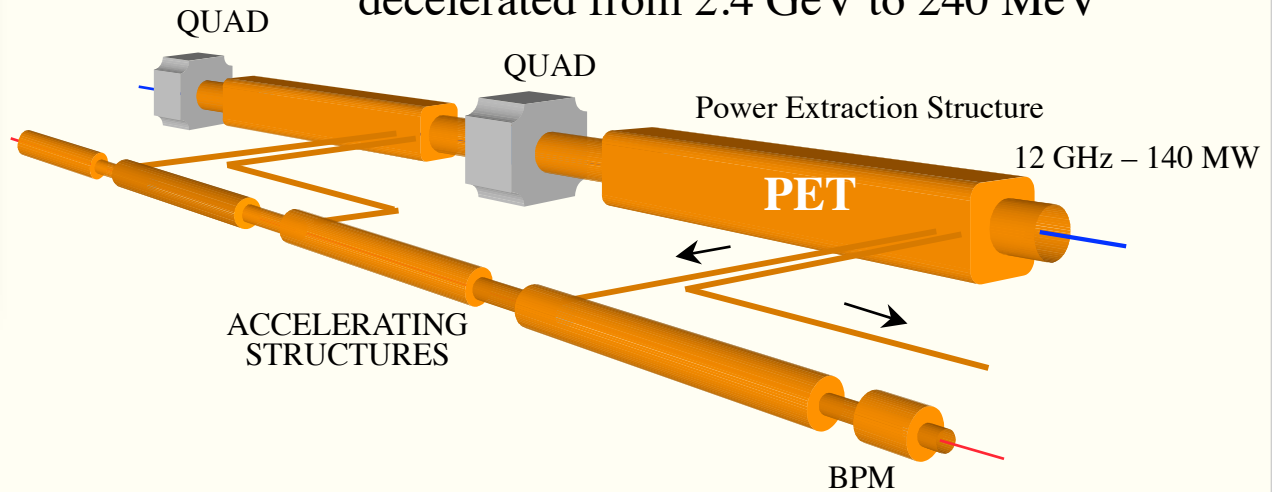
Drive Beam supplies RF power

- 12 GHz bunch structure
- low energy (2.4 GeV - 240 MeV)
- high current (100A)

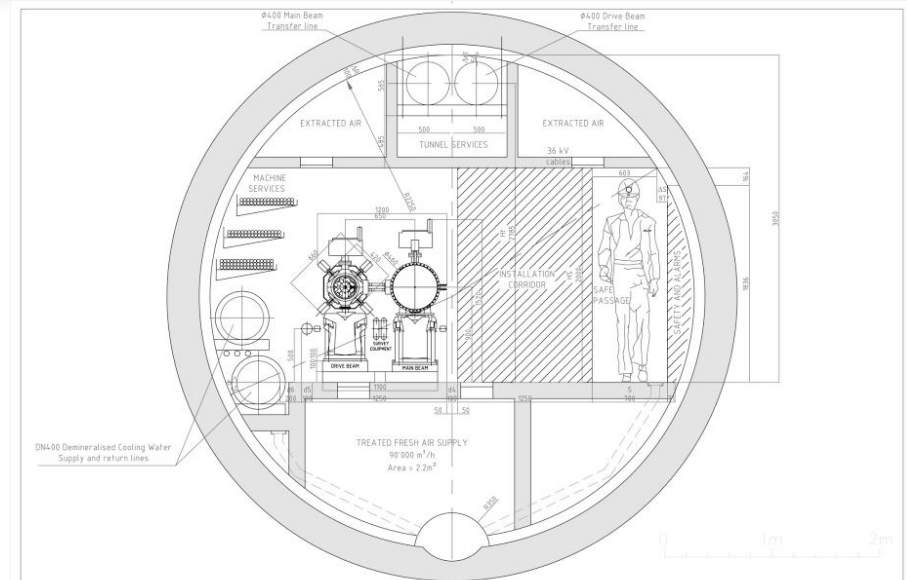
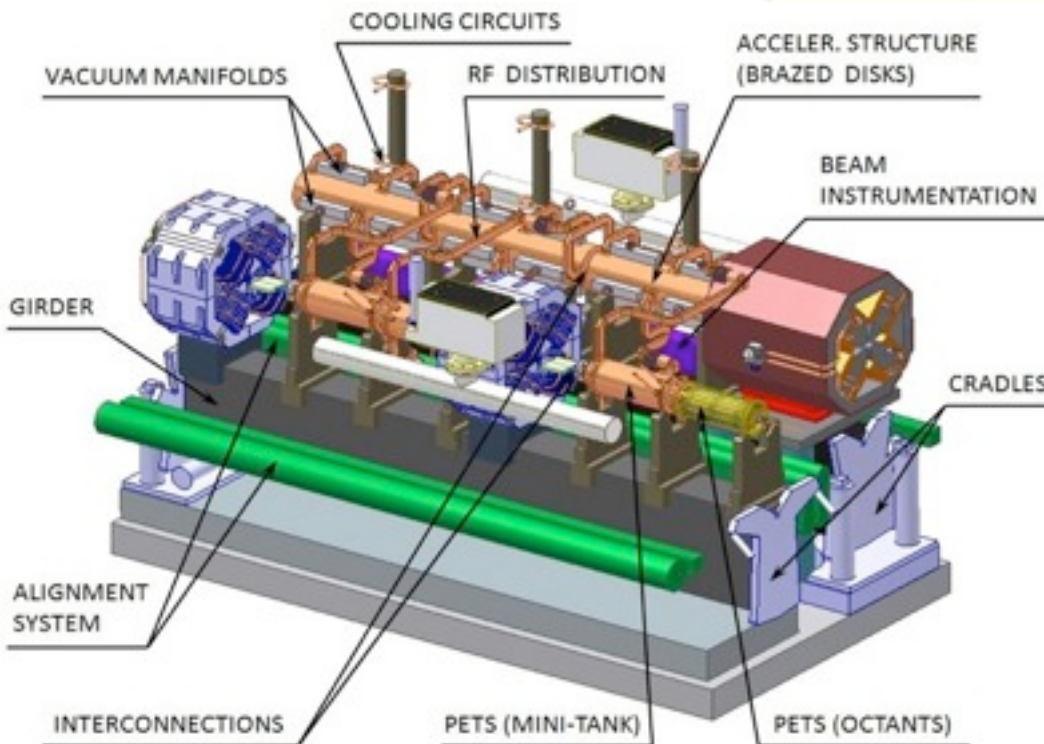
warm (not superconducting) RF

Drive beam - 100 A, 240 ns

decelerated from 2.4 GeV to 240 MeV



Main beam - 1.2 A, 156 ns bunch trains
accelerated from 9 GeV to 1.5 TeV



ILC and CLIC parameters

ILC: Superconducting RF

CLIC: normal conducting copper RF

500 GeV

3 TeV

accelerating gradient:

31.5 MV/m

100 MV/m

35 MV/m target

RF Peak power:

0.37 MW/m , 1.6 ms, 5 Hz

275 MW/m, 240 ns, 50 Hz

RF average power:

2.9 kW/m

3.7 kW/m

total length:

31 km

48.4 km

site power :

230 MW

392 MW

Beam structure:

particles per bunch:

20×10^9

3.7×10^9

2625 bunches / pulse of 0.96 ms

312 bunches / pulse of 156 ns

bunch spacing

369 ns

0.5 ns



The main 2013-14 LHC consolidations

1695 Openings and final reclosures of the interconnections

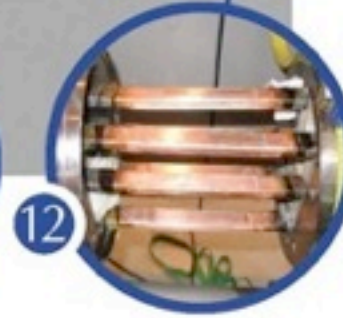
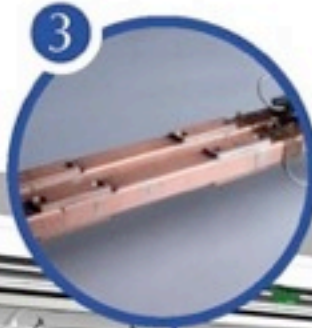
Complete reconstruction of 1500 of these splices

Consolidation of the 10170 13kA splices, installing 27 000 shunts

Installation of 5000 consolidated electrical insulation systems

300 000 electrical resistance measurements

10170 orbital welding of stainless steel lines



18 000 electrical Quality Assurance tests

10170 leak tightness tests

4 quadrupole magnets to be replaced

15 dipole magnets to be replaced

Installation of 612 pressure relief devices to bring the total to 1344

Consolidation of the 13 kA circuits in the 16 main electrical feed-boxes