

The Local Dark Matter Density

New constraints on the Milky Way's dark disc and the shape of the halo

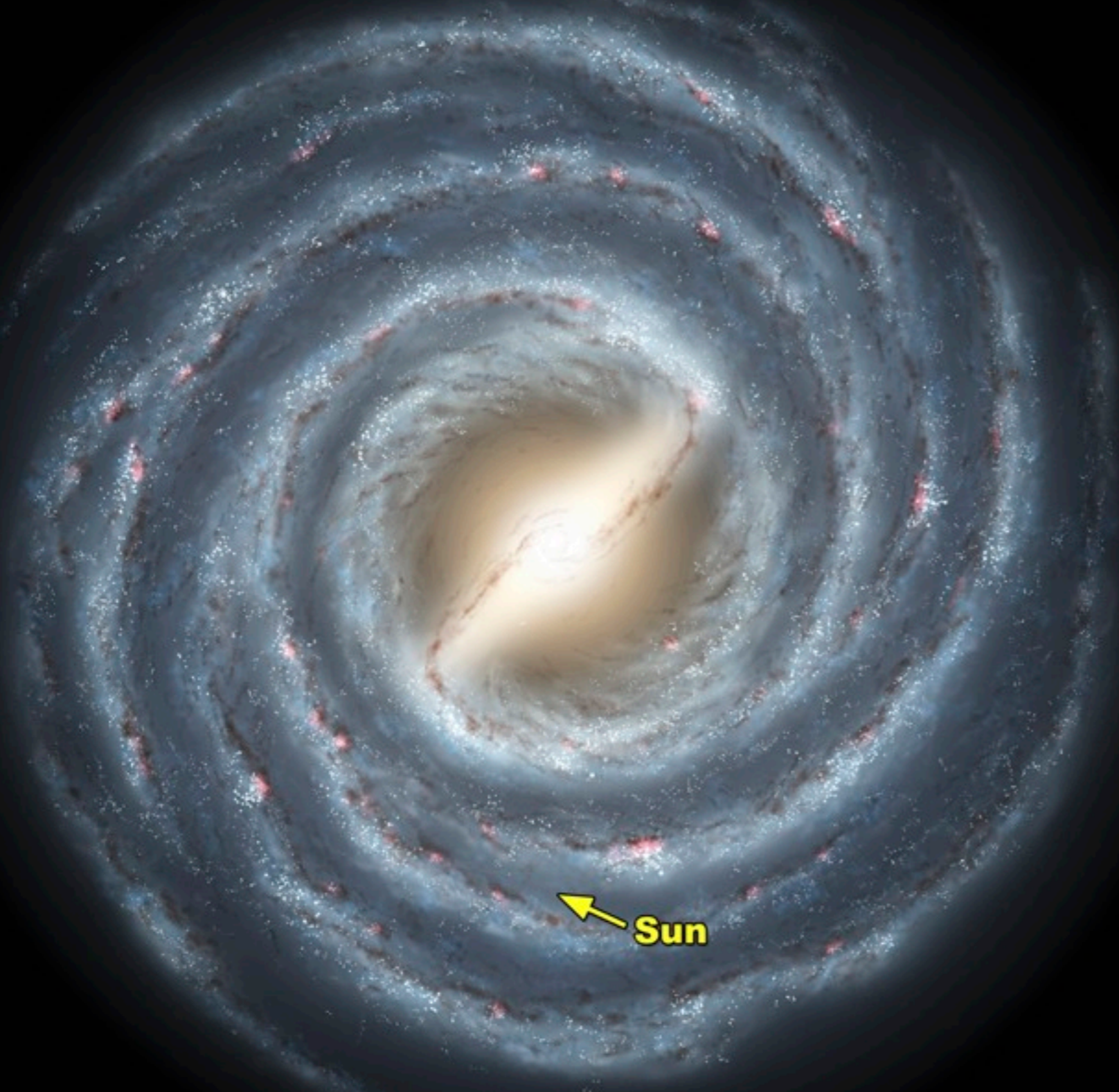
Prof. Justin Read | University of Surrey

Silvia Garbari; George Lake; Greg Ruchti; Oscar Agertz

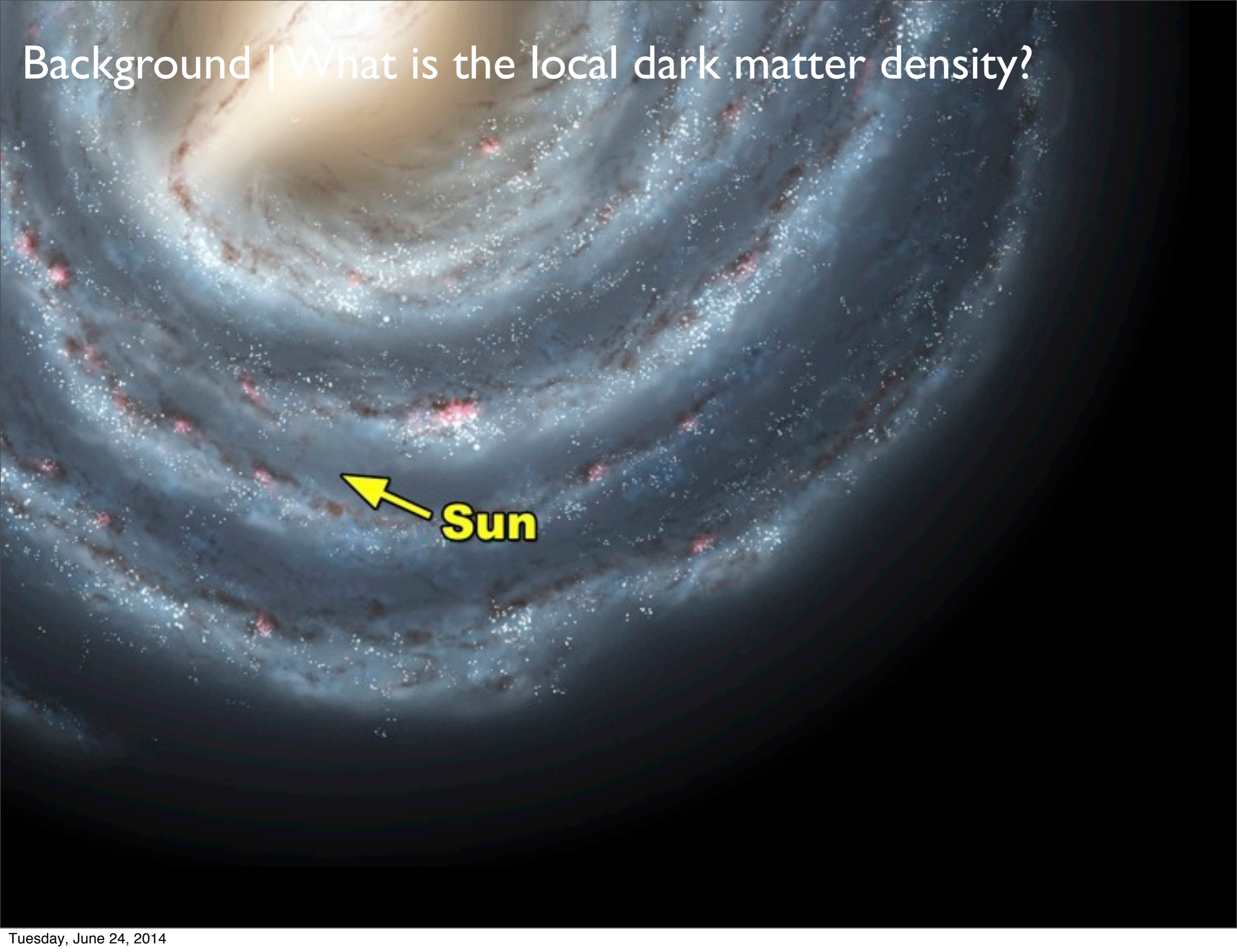
Background

[What is ρ_{dm} ? How do we measure it? Why is it interesting?]

Background | What is the local dark matter density?



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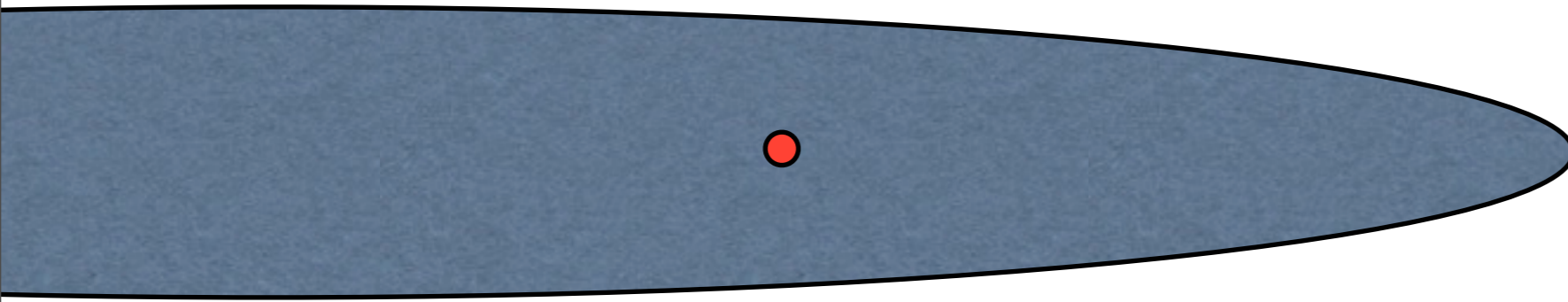
Sun

Background | What is the local dark matter density?



Background | How can we measure the local DM density?

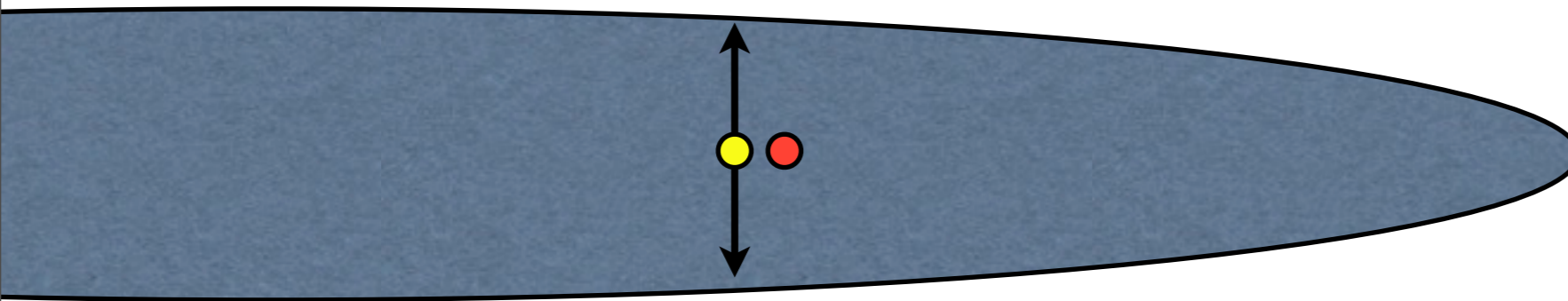
I. Local measure:



ρ_{dm}

Background | How can we measure the local DM density?

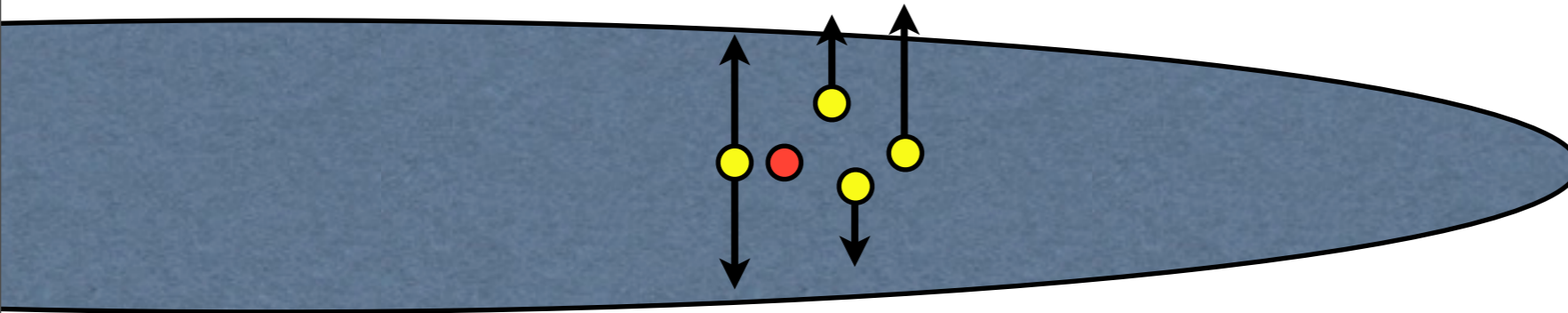
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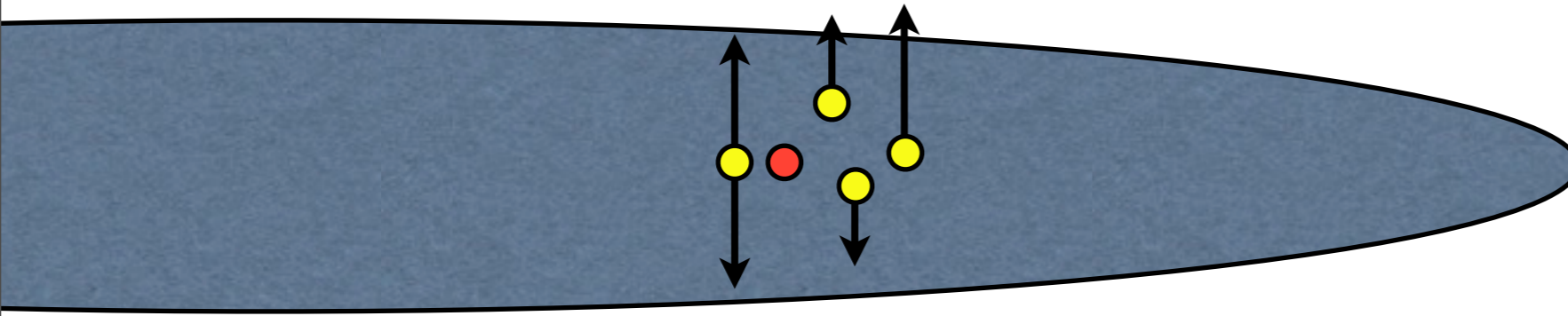
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$$\rho_{\text{dm}}$$

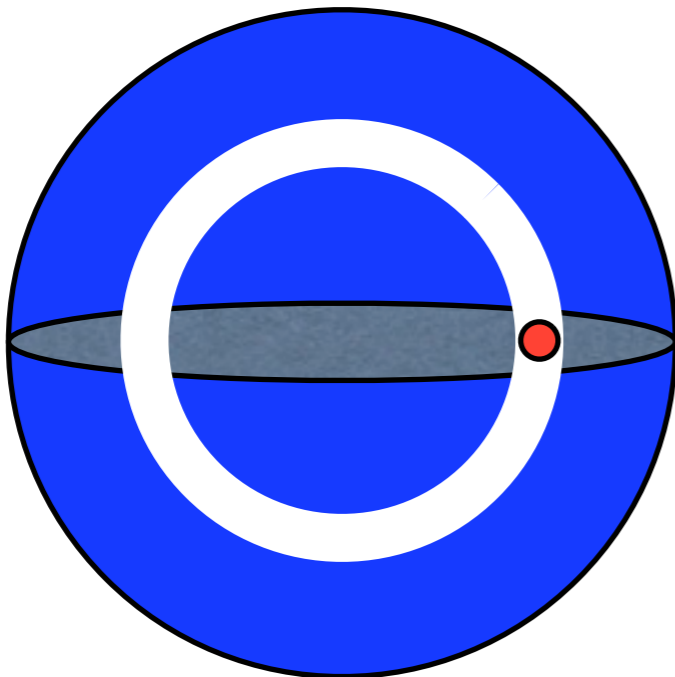
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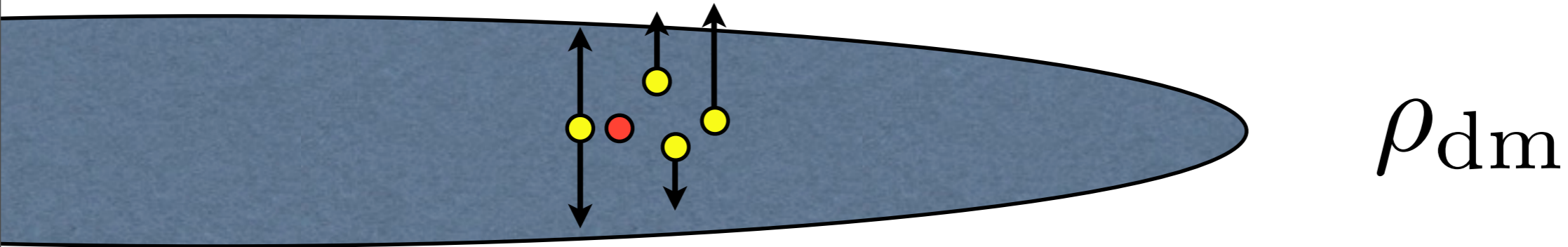
2. Global measure:



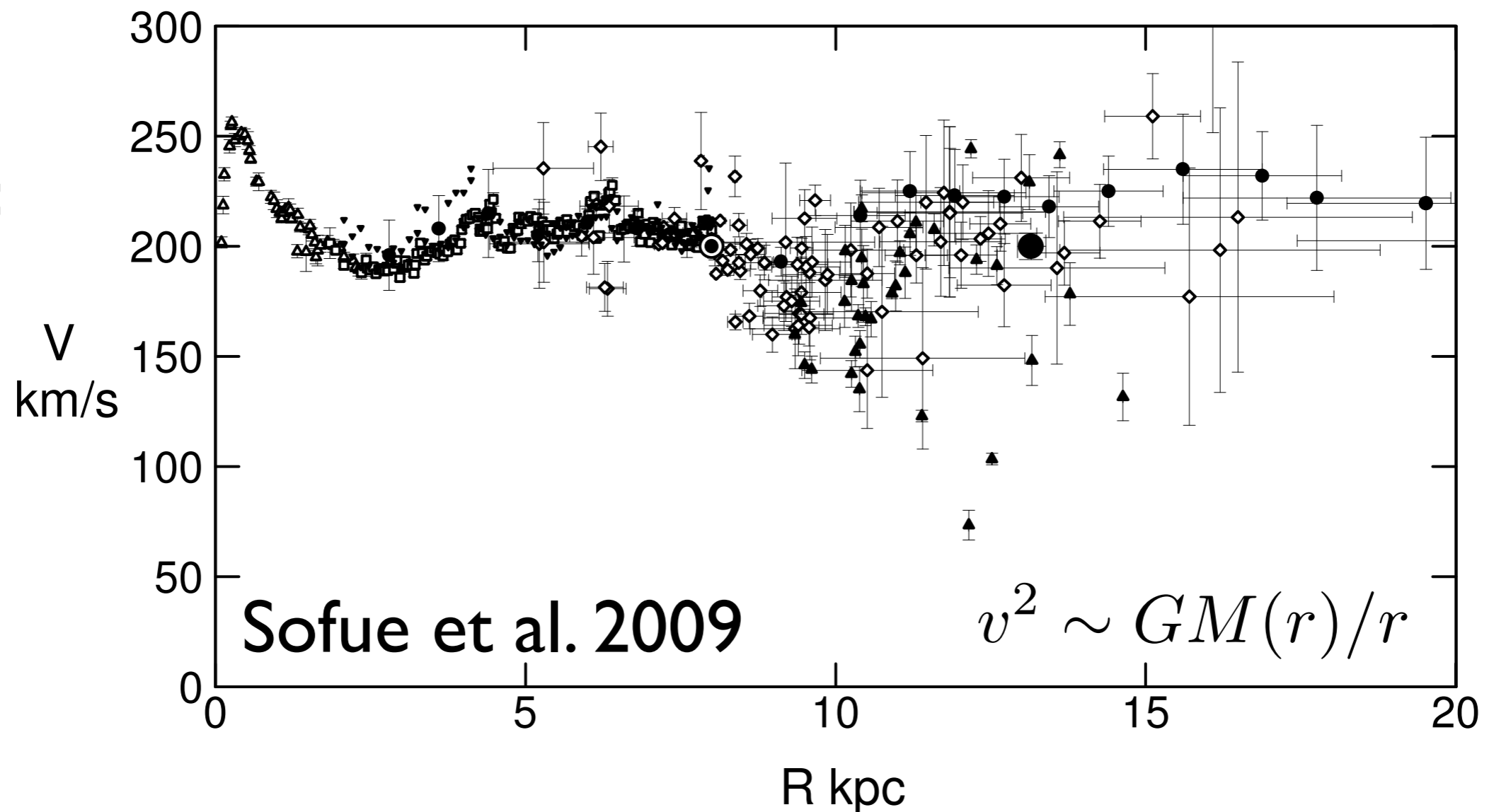
$$\rho_{\text{dm,ext}}$$

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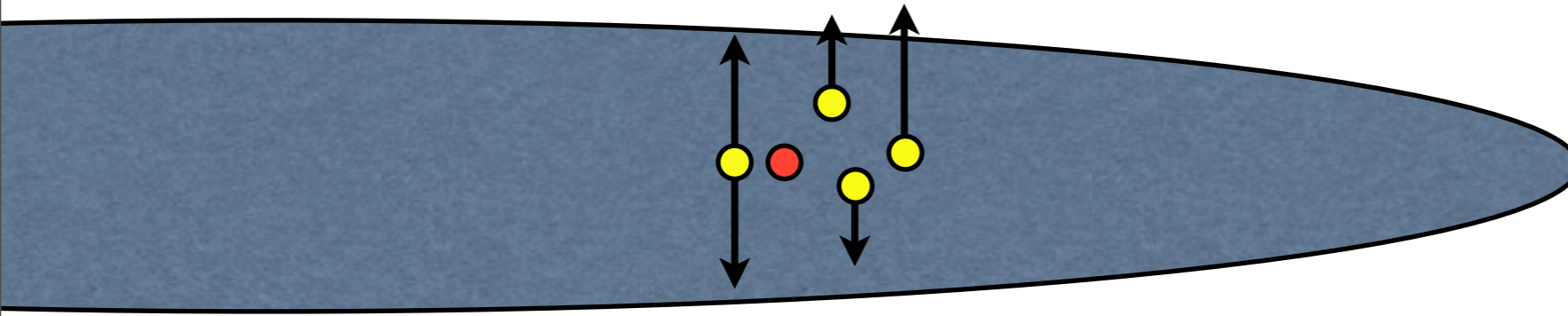


2. Glc



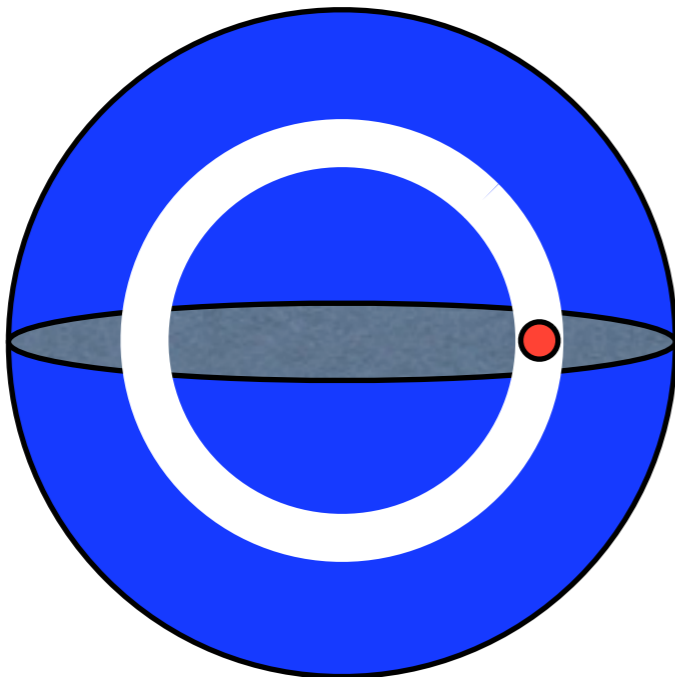
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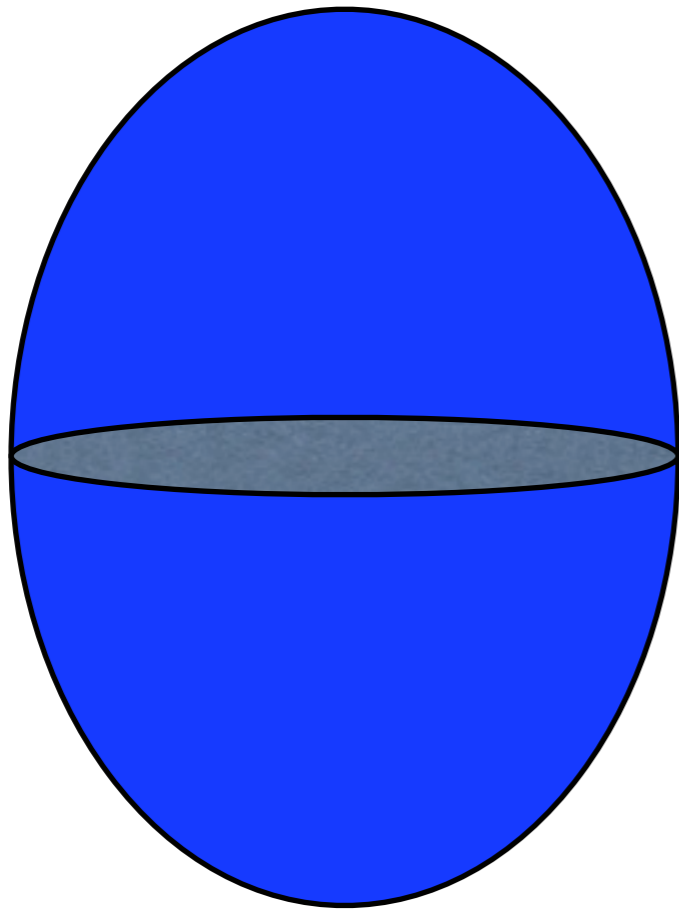


$$\rho_{\text{dm,ext}}$$

Background | Why measure the local dark matter density?

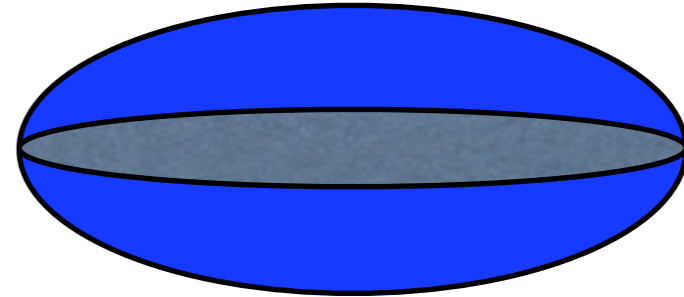
I. Halo shape ...

$$\rho_{\text{dm}} < \rho_{\text{dm,ext}}$$



Prolate

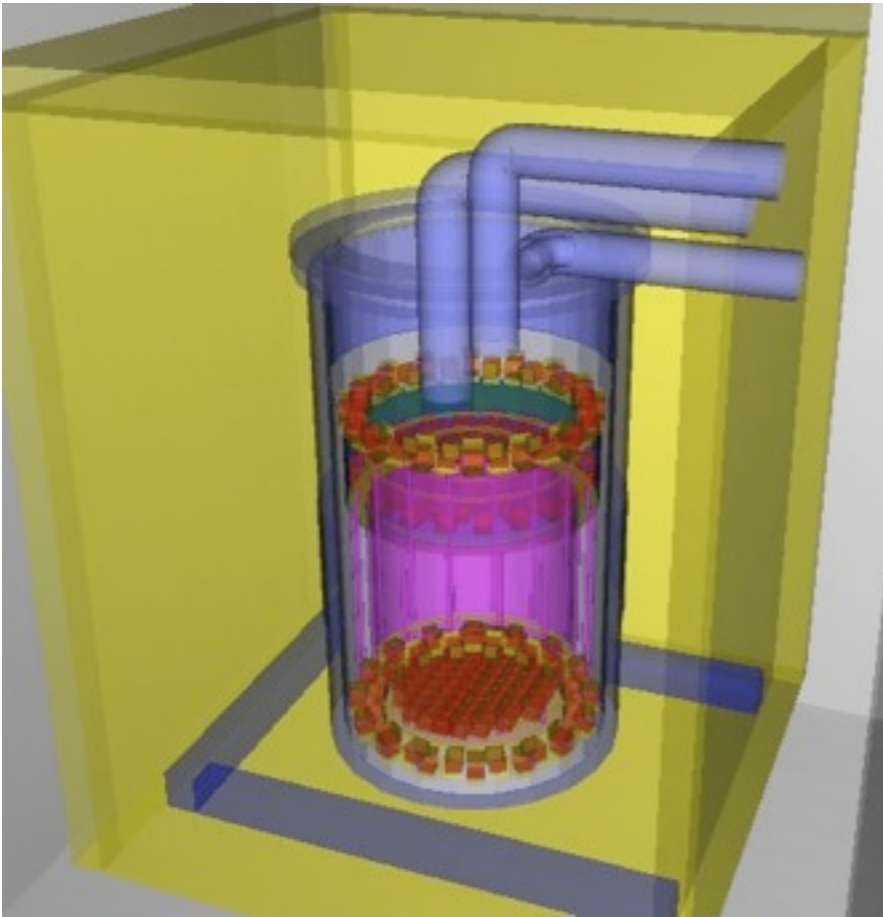
$$\rho_{\text{dm}} > \rho_{\text{dm,ext}}$$



Oblate/dark disc

Background | Why measure the local dark matter density?

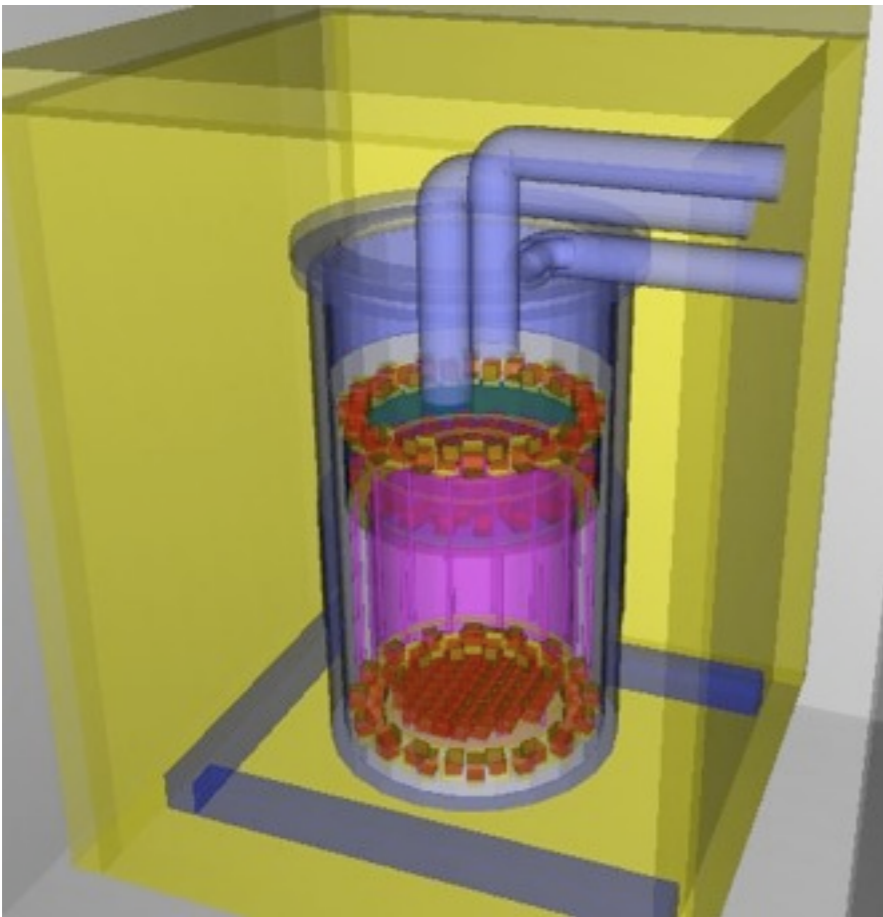
2. Detecting dark matter



- Big tub of **inert** material
- Deep underground
- Wait for rare event
- Need to know very local phase space distribution

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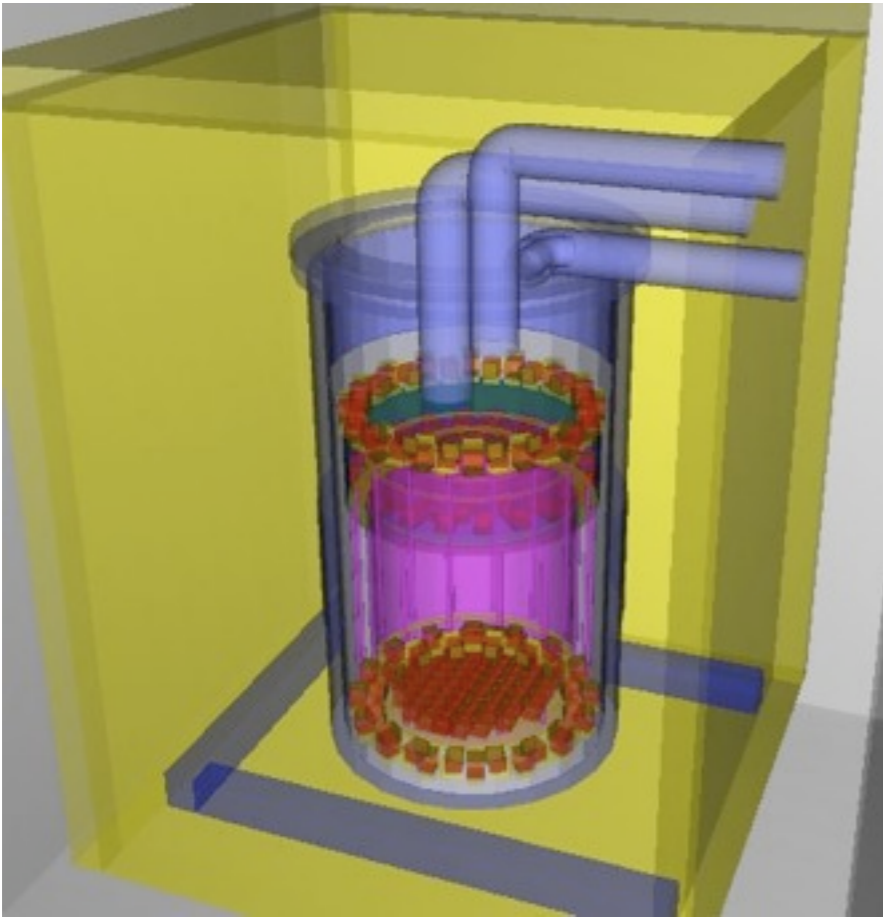


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$$\frac{dR}{dE} = \frac{\rho \sigma_{\text{wn}} |F(E)|^2}{2m\mu^2} \int_{v > \sqrt{ME/2\mu^2}}^{v_{\text{max}}} \frac{f(\mathbf{v}, t)}{v} d^3v$$

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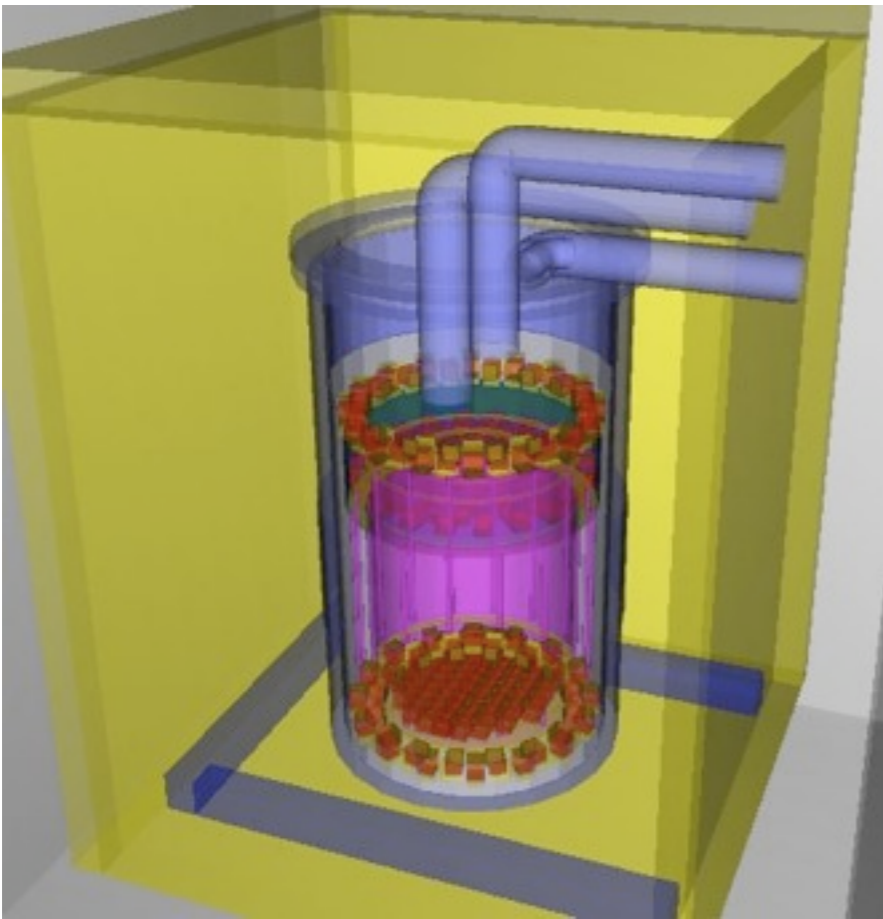


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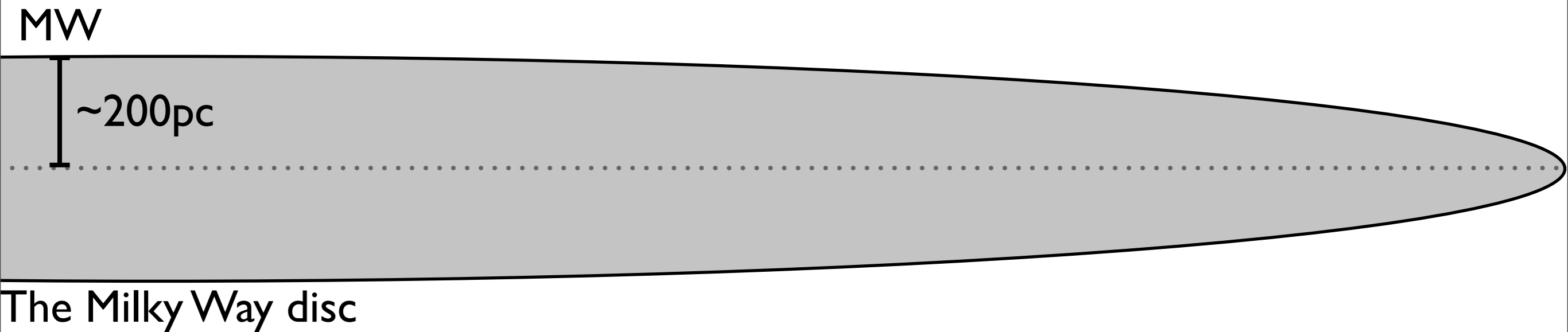
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Background | The need for simulations

1. $\rho_{\text{lab}} \neq \rho_{\text{dm}} (< 1 \text{ kpc})$

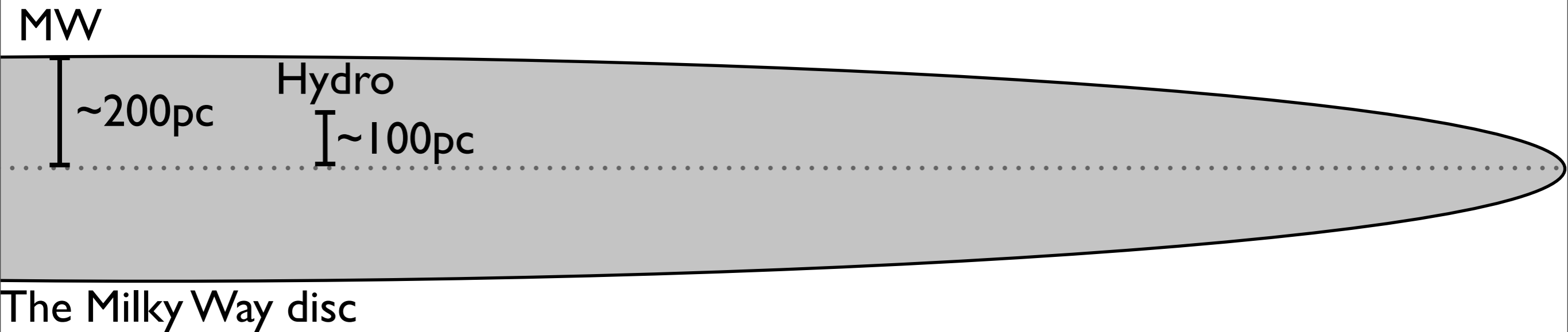
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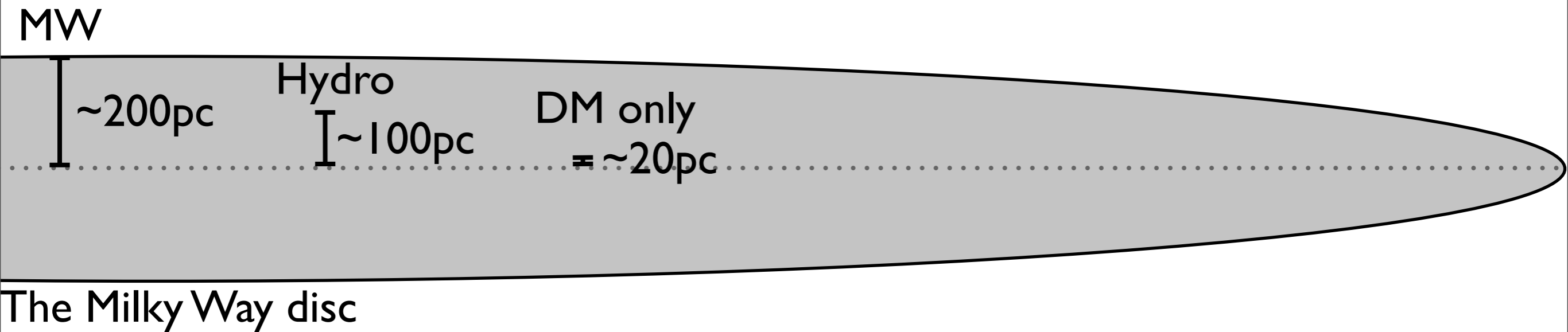
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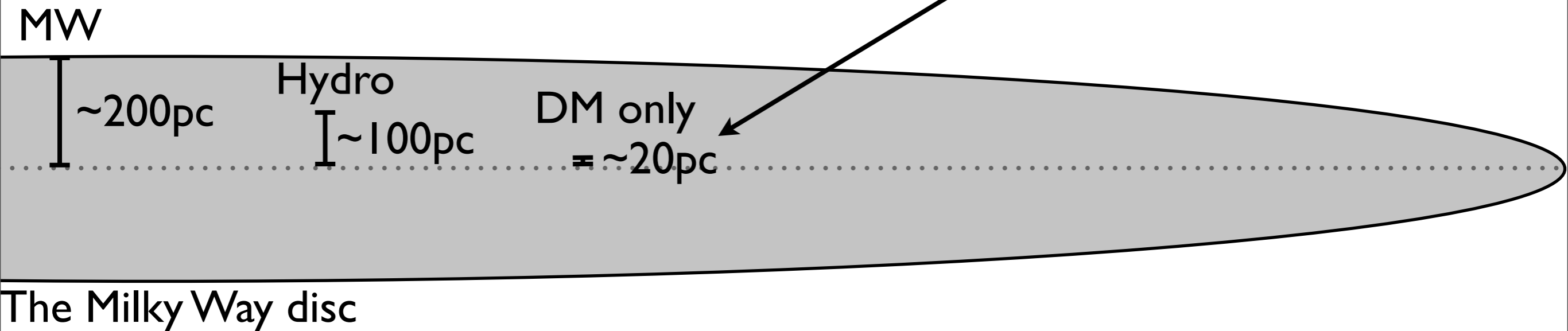
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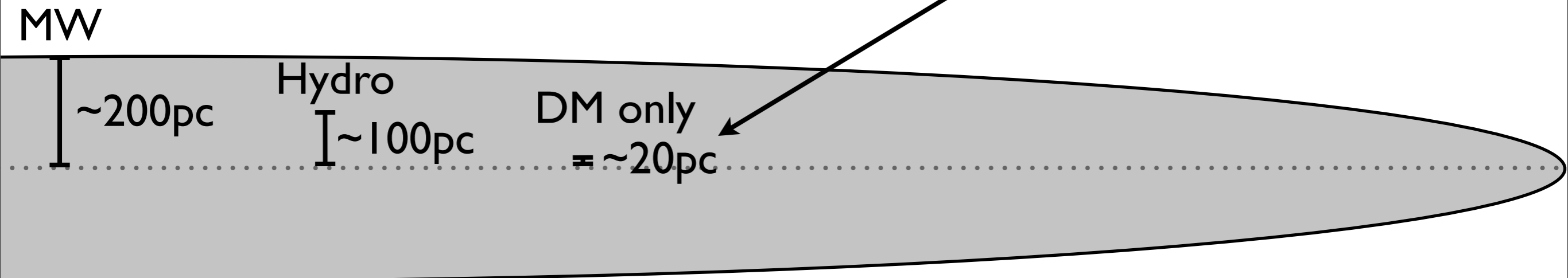
Solar system is a million times smaller than this!



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The Milky Way disc

2. Need $f(\mathbf{v}, t)$

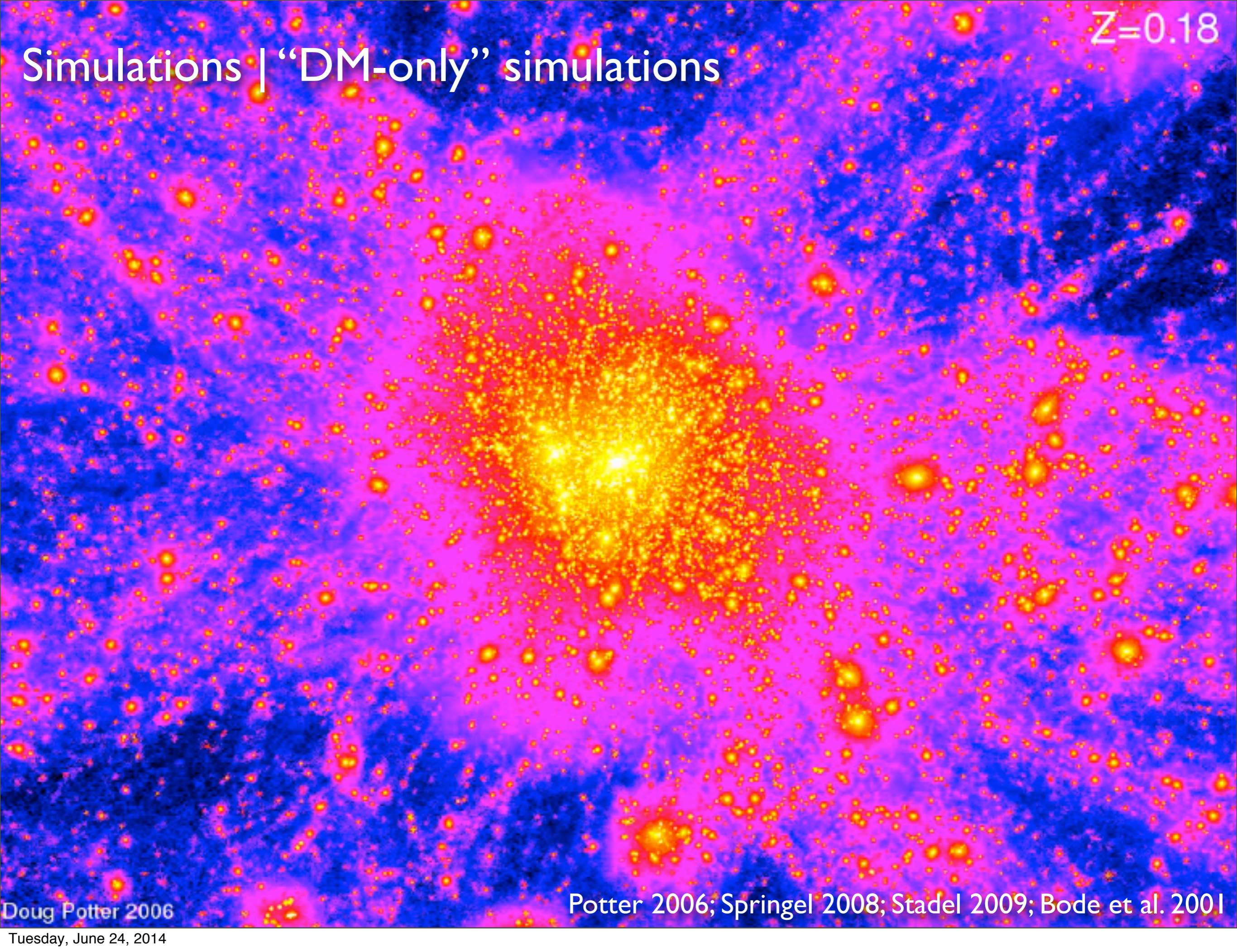
Simulations | “DM-only” simulations

Potter 2006; Springel 2008; Stadel 2009; Vogelsberger et al. 2009

Simulations | “DM-only” simulations

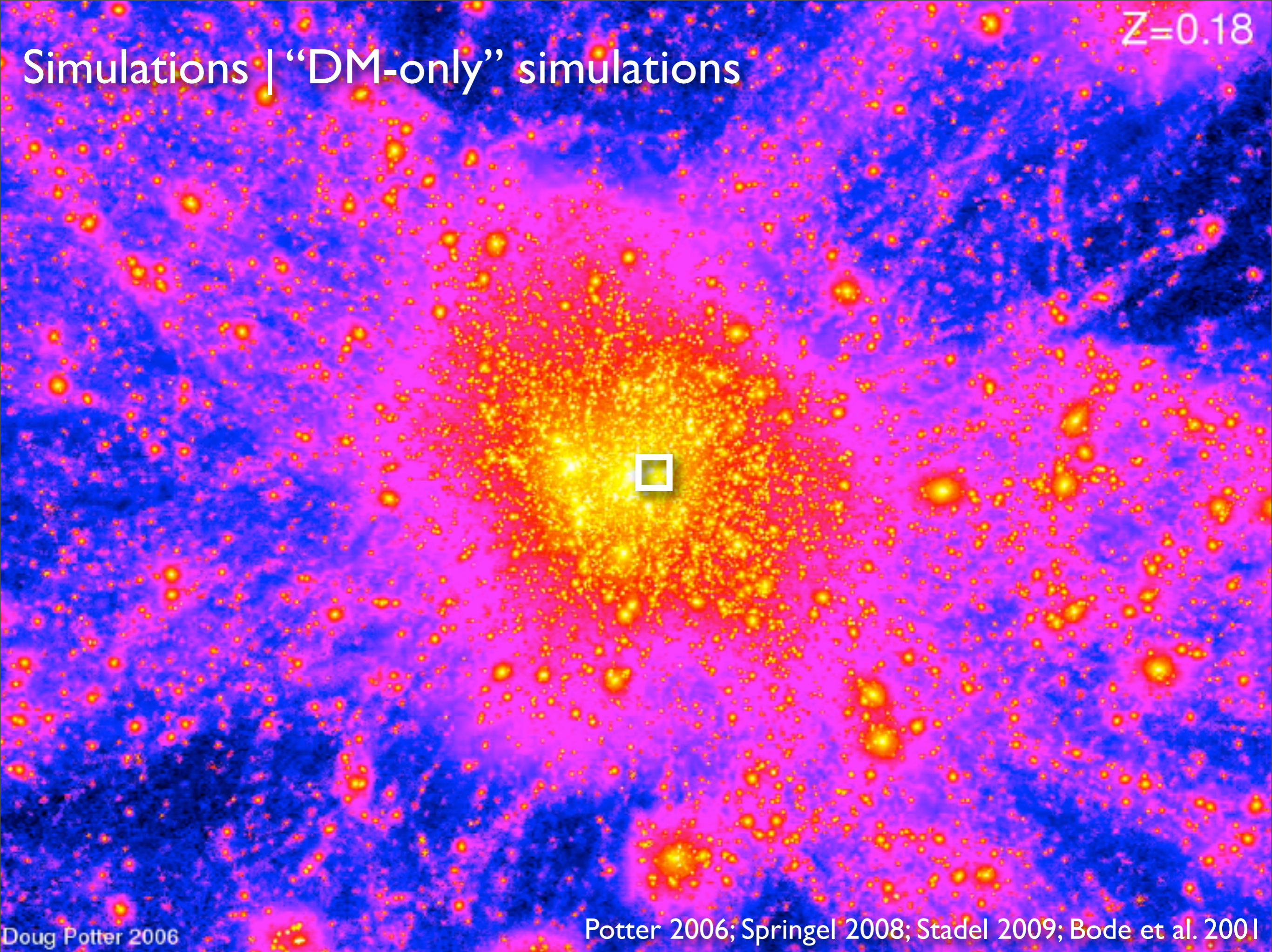
$Z=0.18$

Simulations | “DM-only” simulations



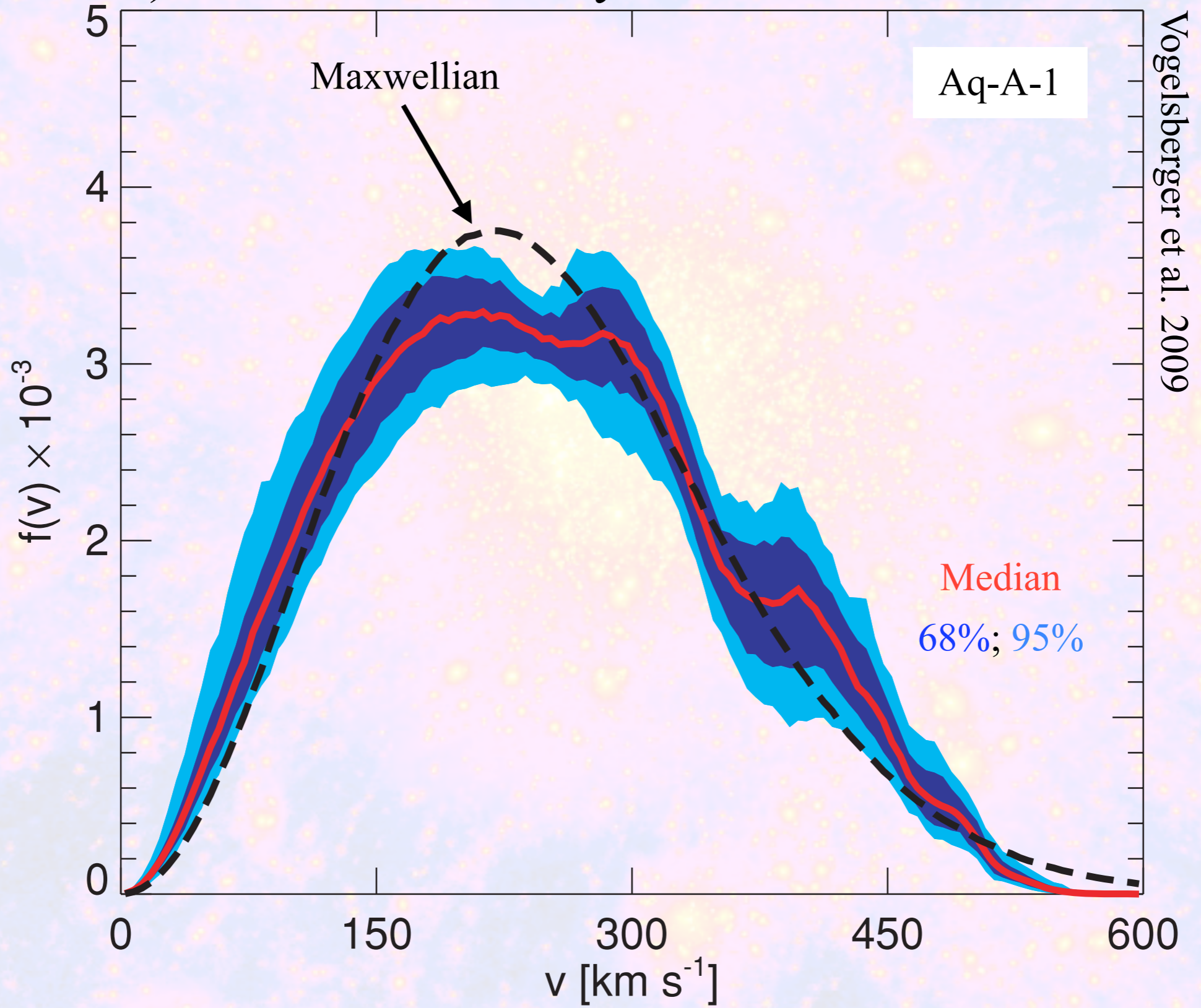
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b) Local DM velocity PDF



Vogelsberger et al. 2009

Simulations | “DM-only” simulations | Fine structure

- **Unresolved substructure** | **not likely important**
[Vogelsberger et al. 2009; Zemp et al. 2009; Kamionkowski et al. 2008]
- **Unresolved streams** | **not likely important**
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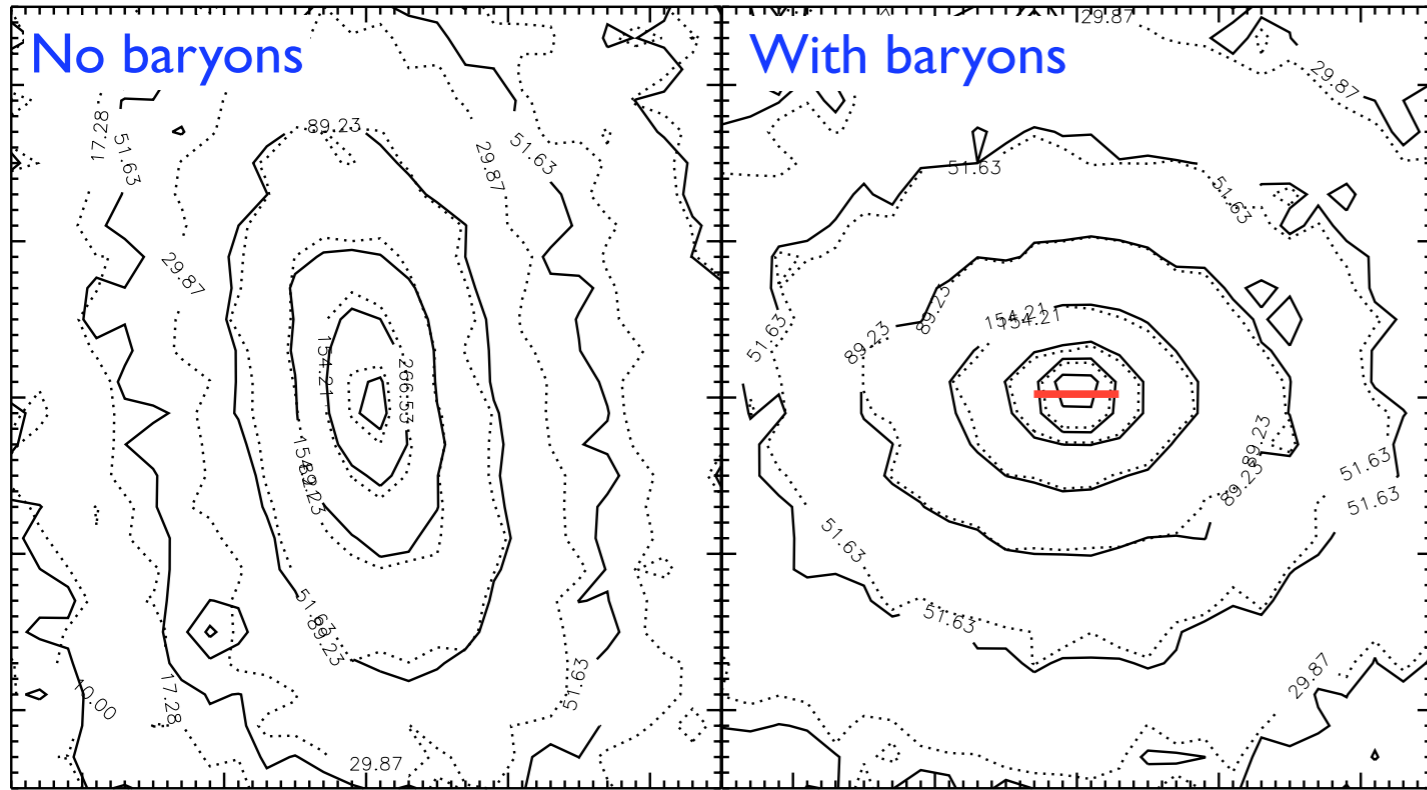
$$\rho_{\text{dm}} \Rightarrow \rho_{\text{lab}}$$

~600 light years ~metres

Simulations | The importance of baryons

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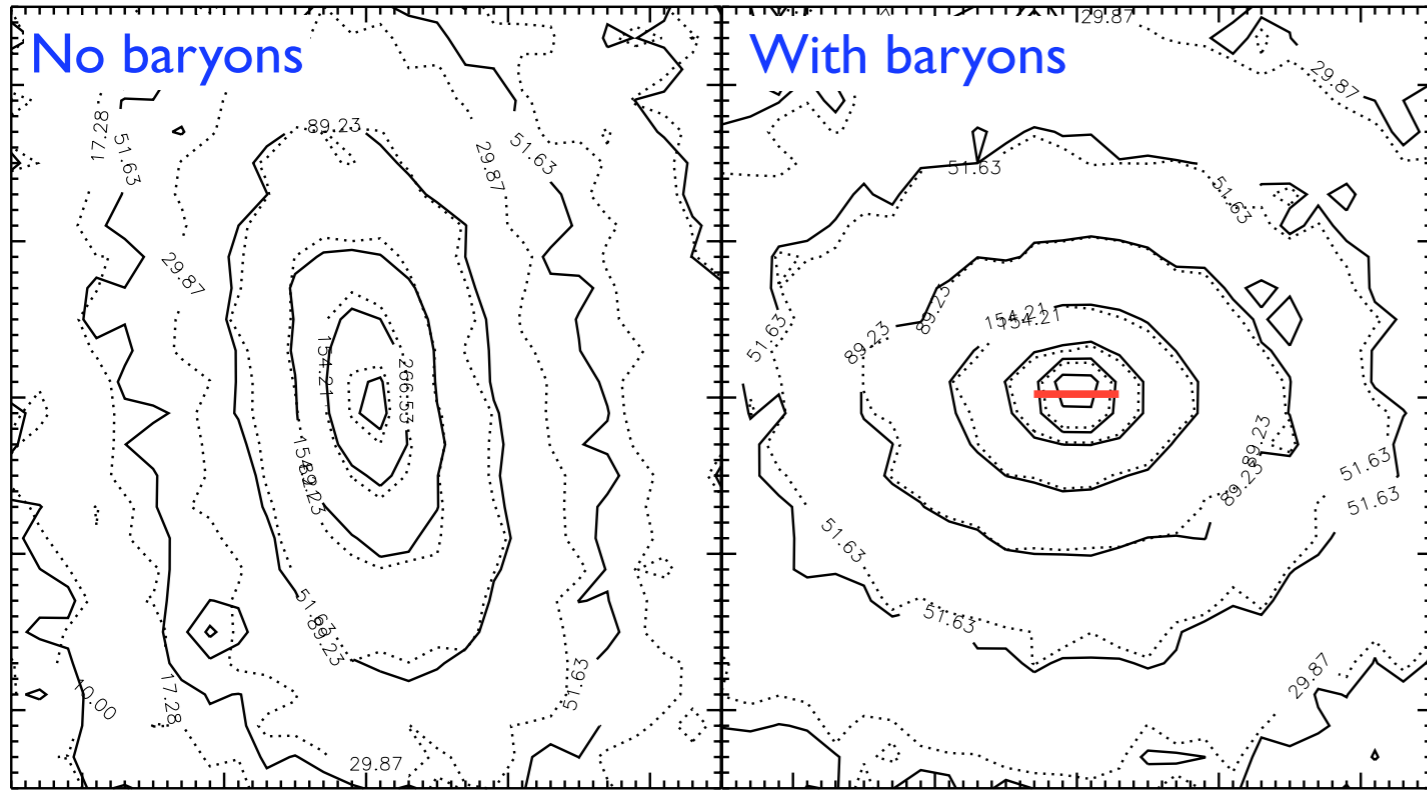
Shape change



Katz & Gunn 1991; Dubinski 1994; Debattista et al. 2008; Read et al. 2009

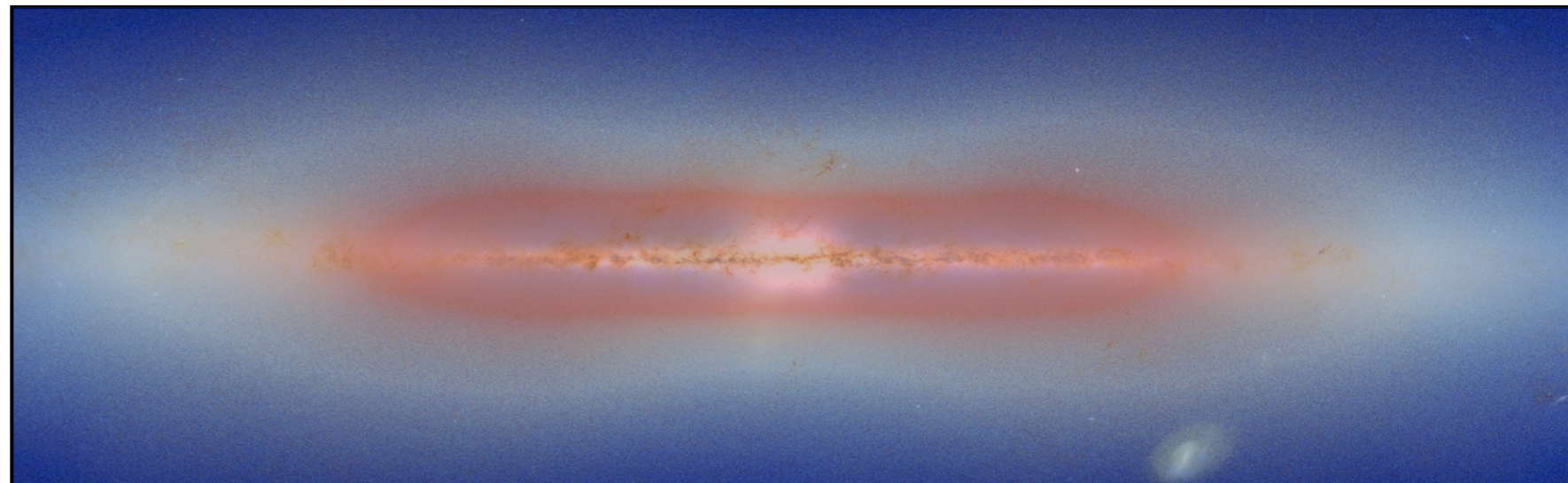
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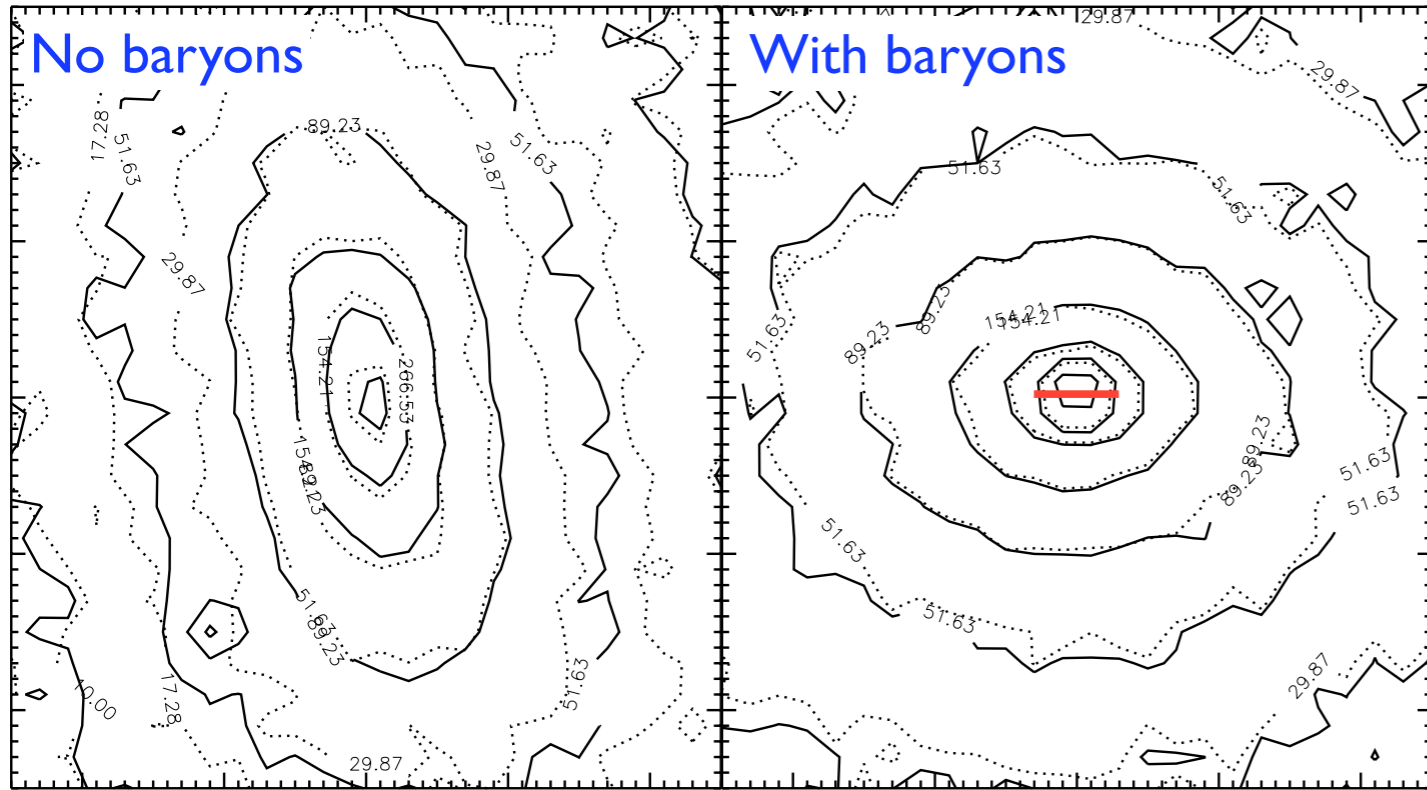
Dark discs



Lake 1989; Read et al. 2008/9

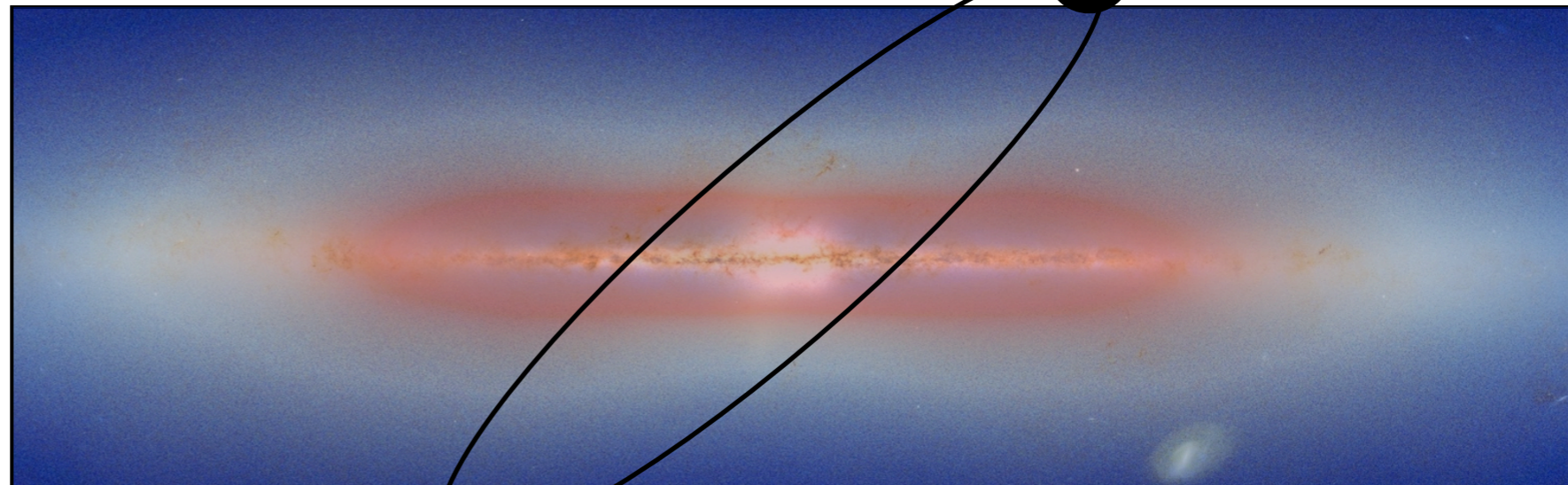
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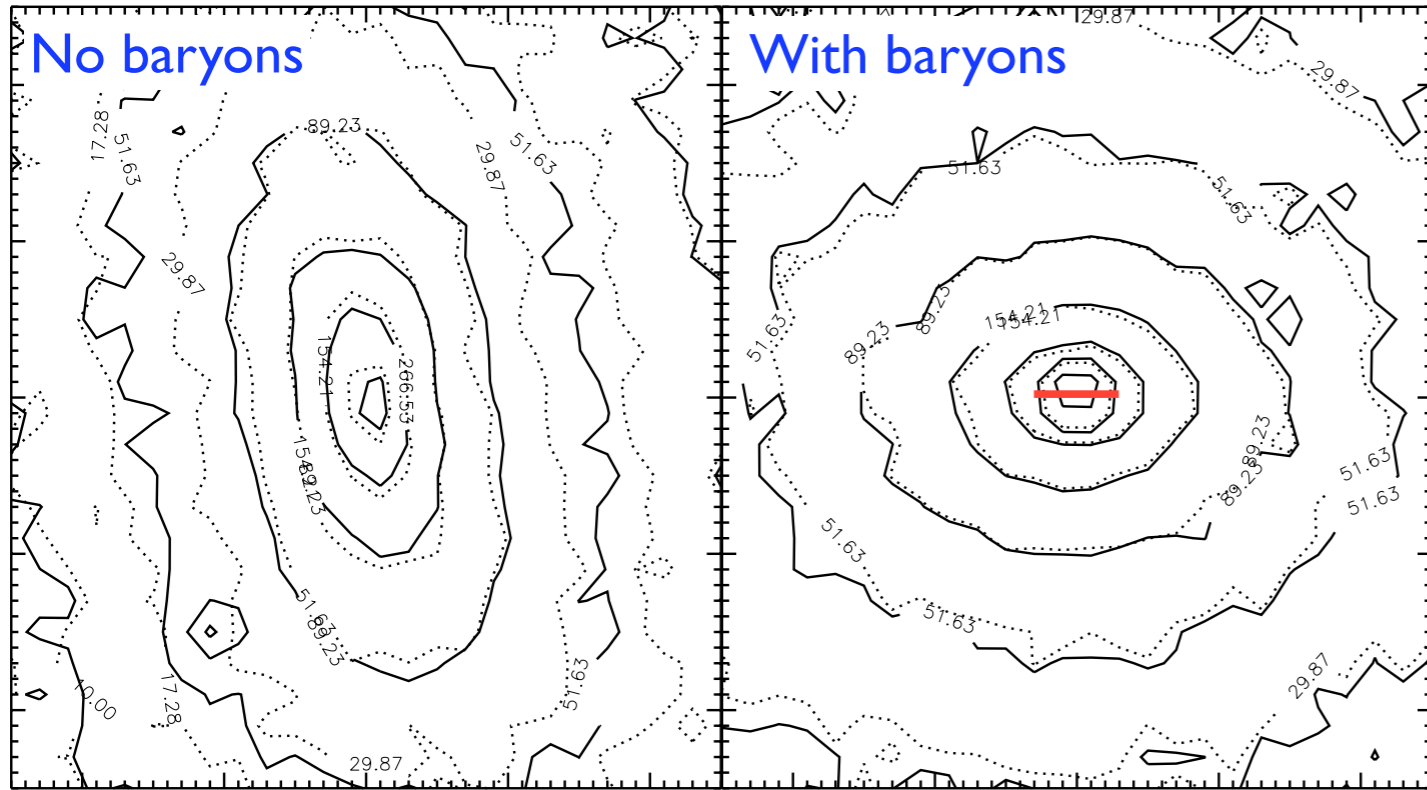
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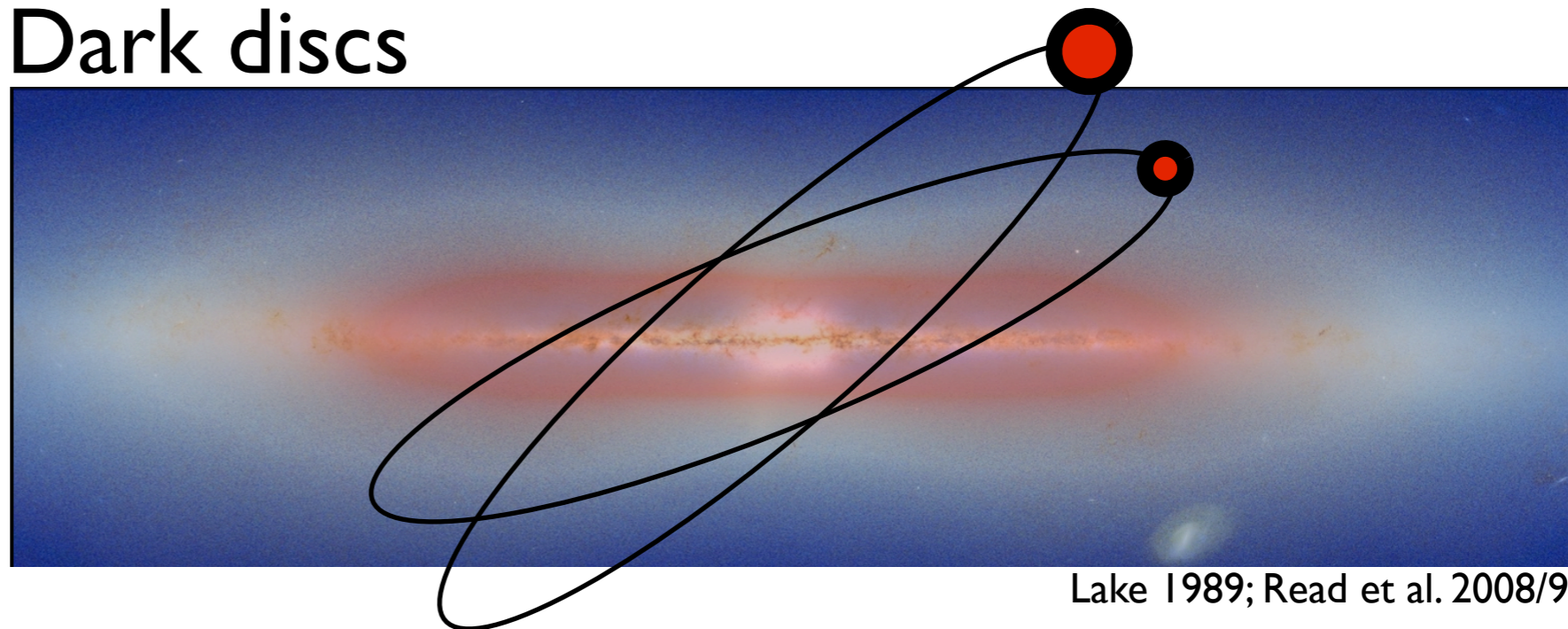
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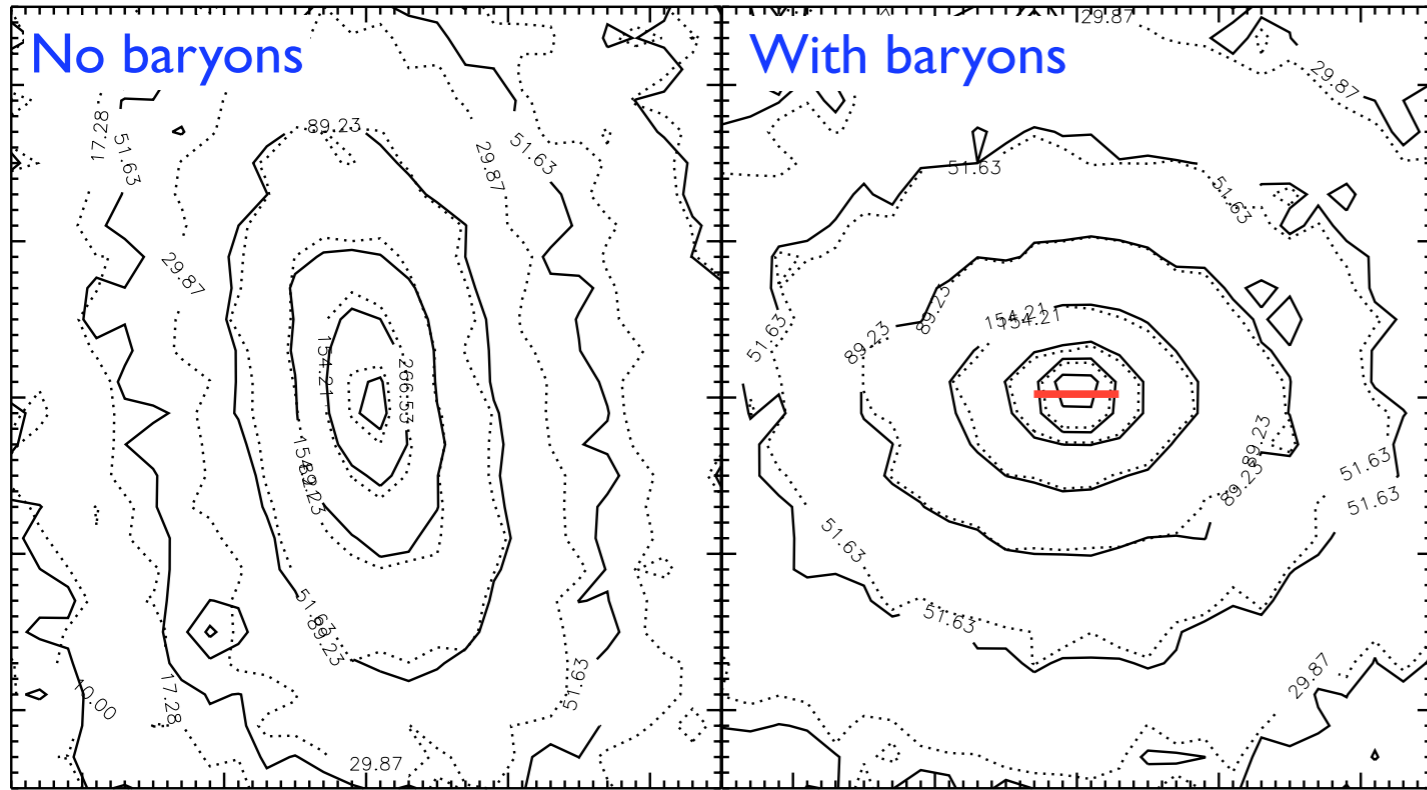
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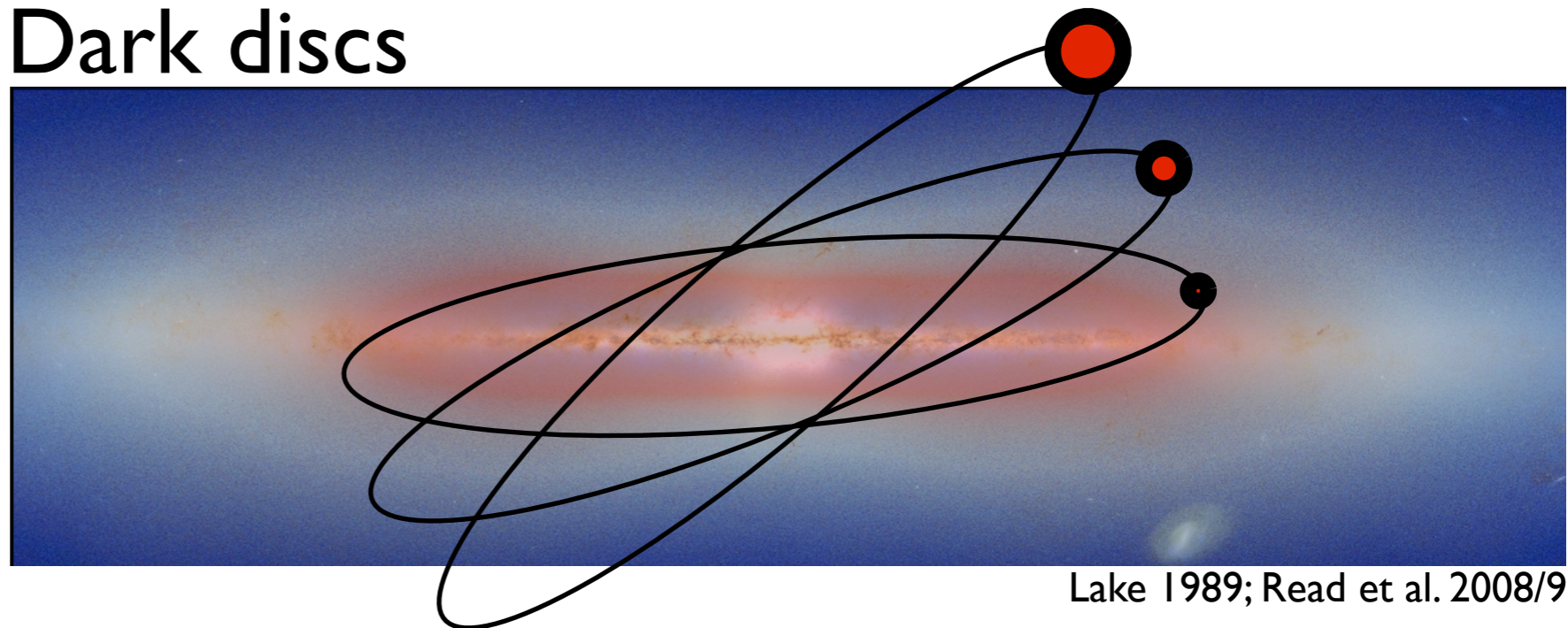
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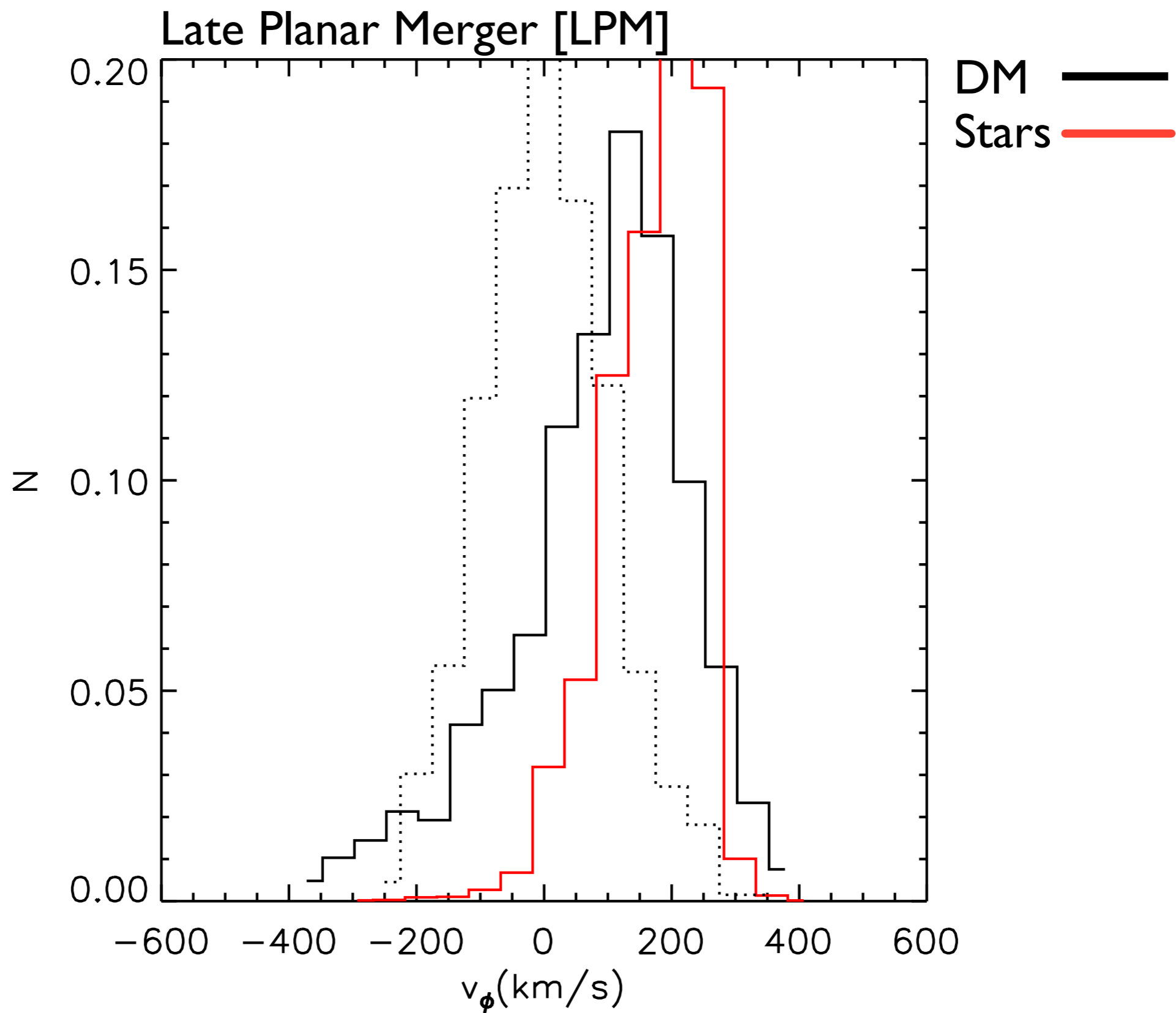


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Simulations | Dark discs

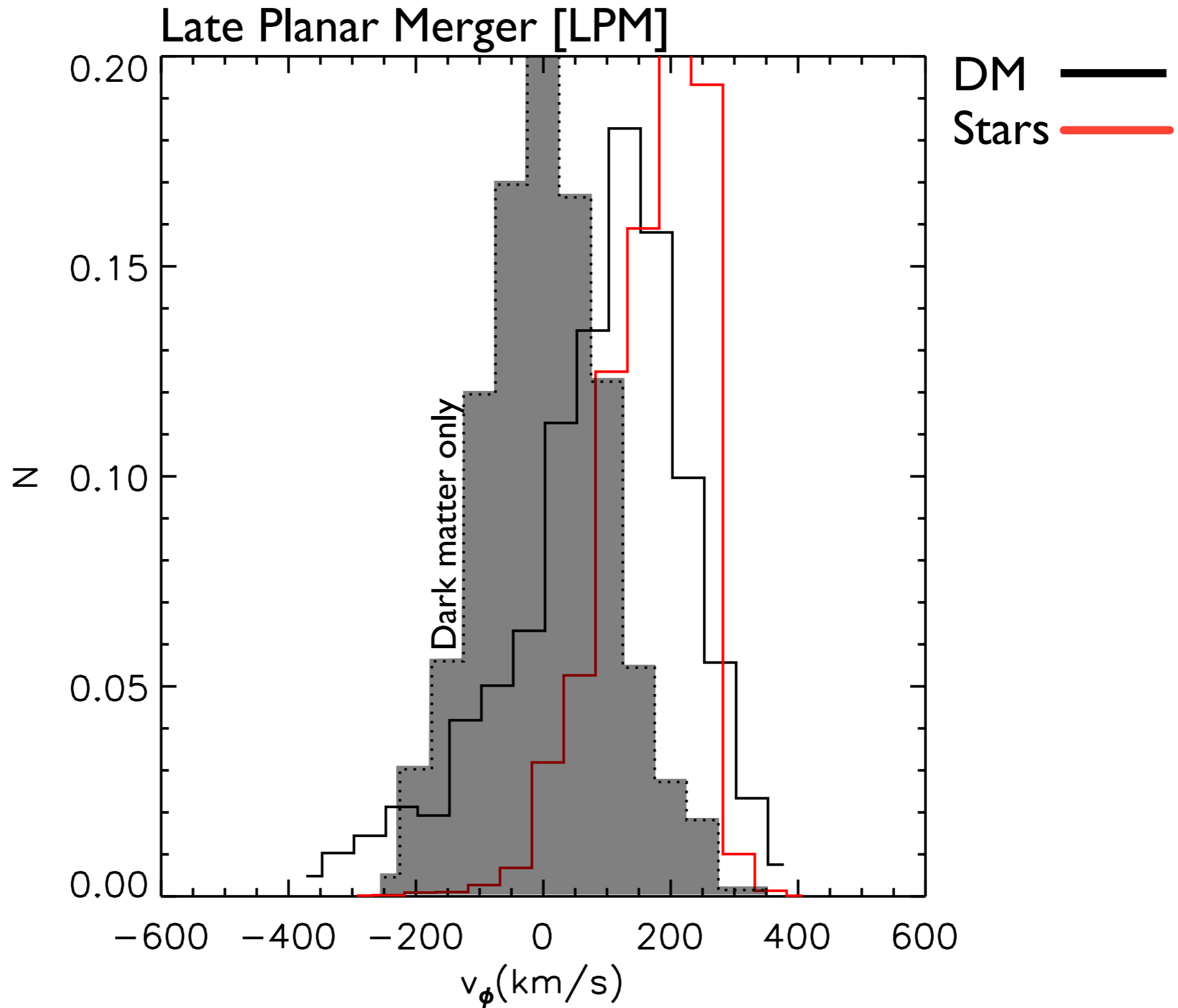
Stellar disc

Simulations | Dark discs



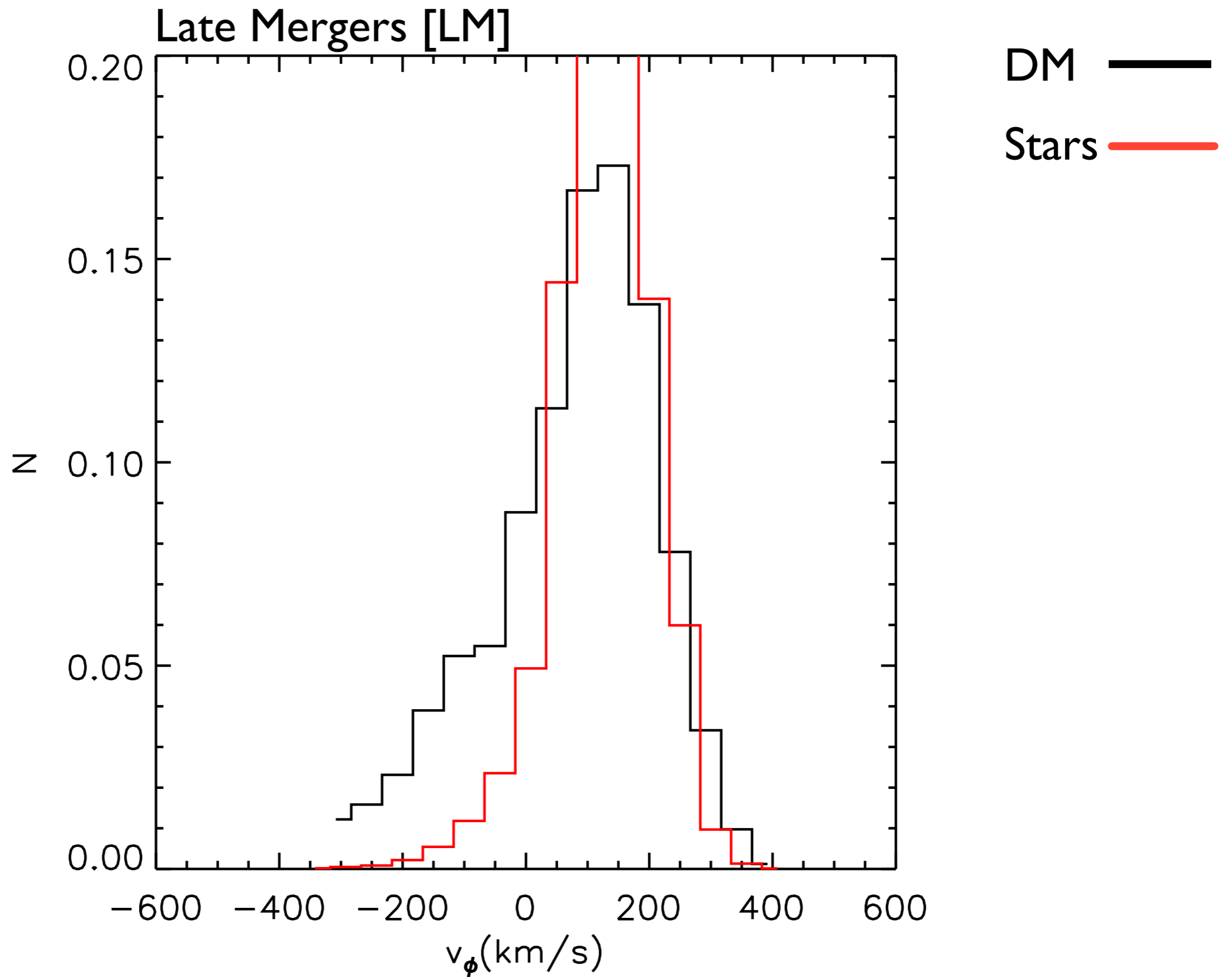
Read et al., 2008/9; Bruch et al. 2009a/b.

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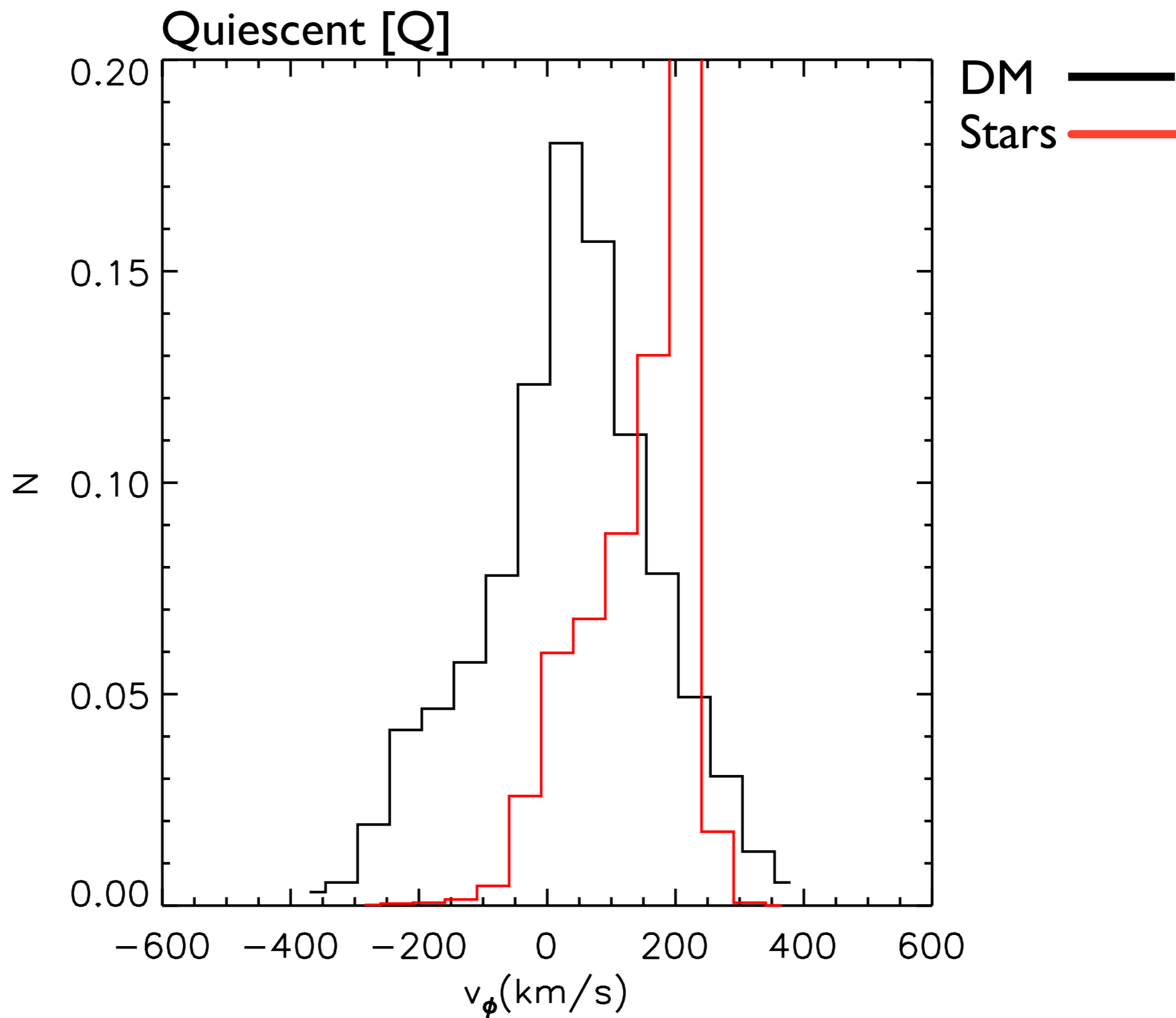
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Measurement

[ρ_{dm} ; the local halo shape; and the MW's dark disc]

Measurement | Theory

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$$\frac{df}{dt} = 0 = \frac{\partial f}{\partial t} + \mathbf{v} \cdot \frac{\partial f}{\partial \mathbf{r}} - \nabla \Phi \cdot \frac{\partial f}{\partial \mathbf{v}}$$

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But! hard to measure $f(\mathbf{r}, \mathbf{v})$

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$$\frac{1}{R} \frac{\partial}{\partial R} (R \nu_i \overline{v_R v_z}) + \frac{\partial}{\partial z} \left(\nu_i \overline{v_z^2} \right) + \nu_i \frac{\partial \Phi}{\partial z} = 0$$

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“tilt term”

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$h \ll R_d$

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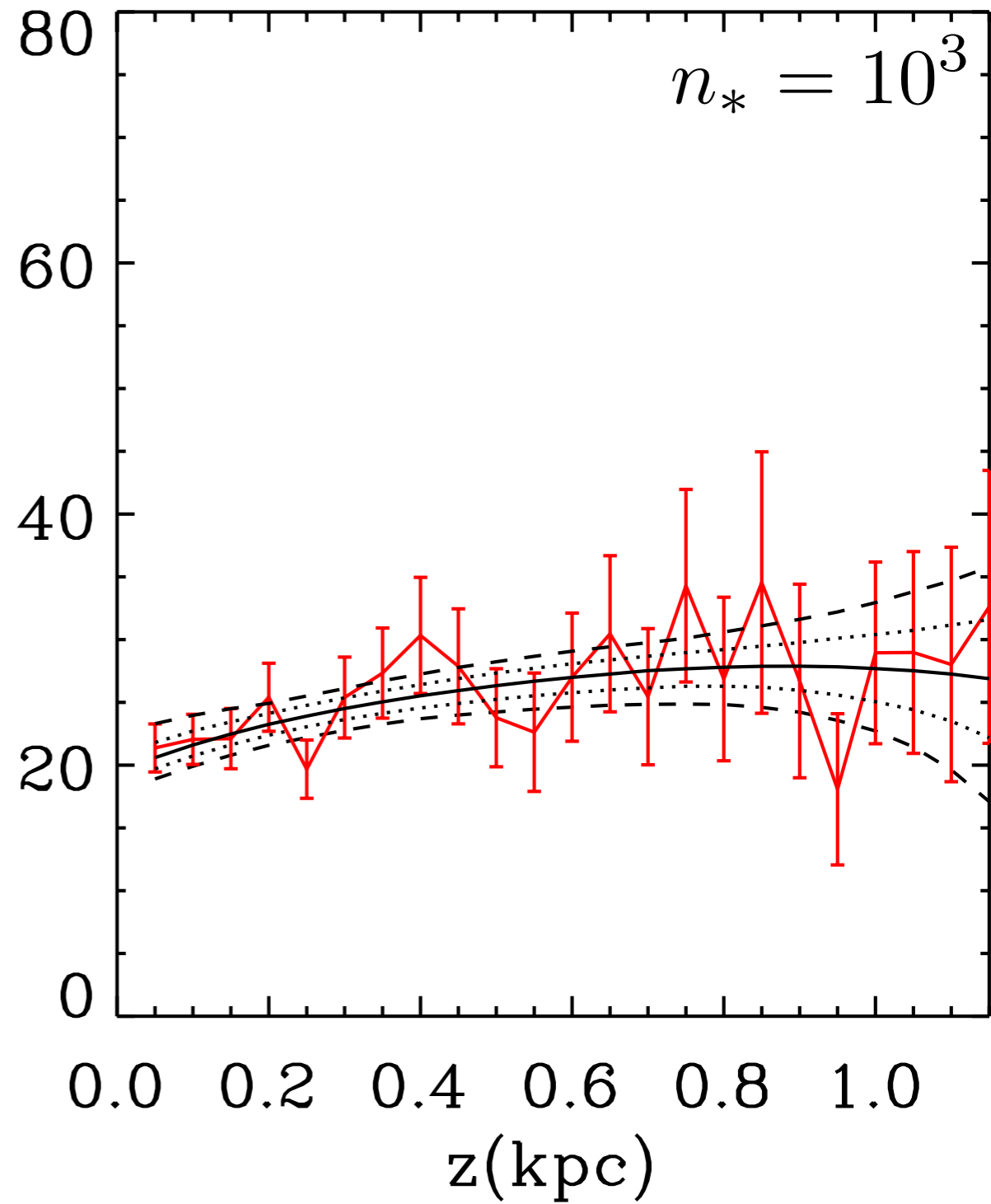
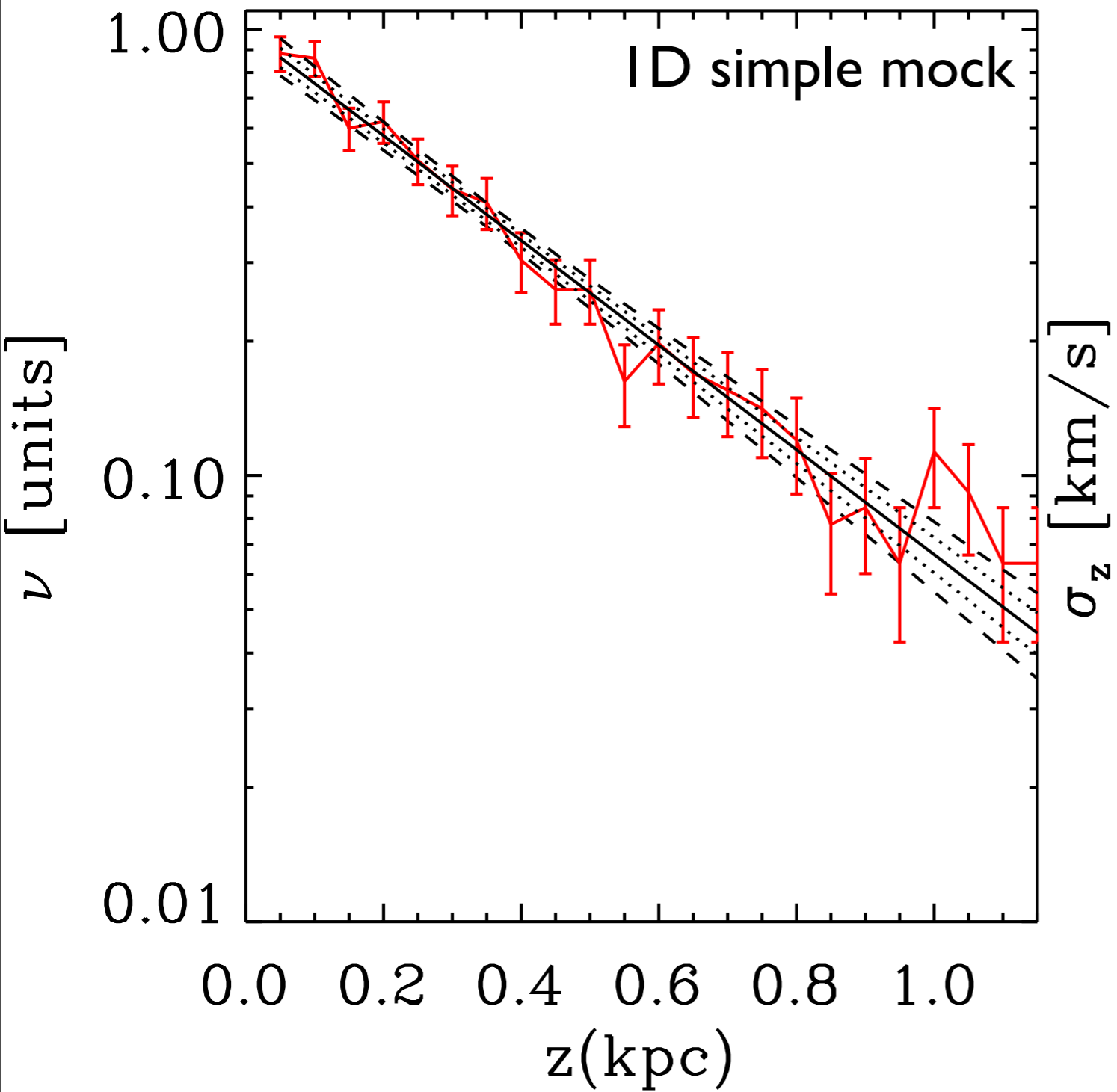
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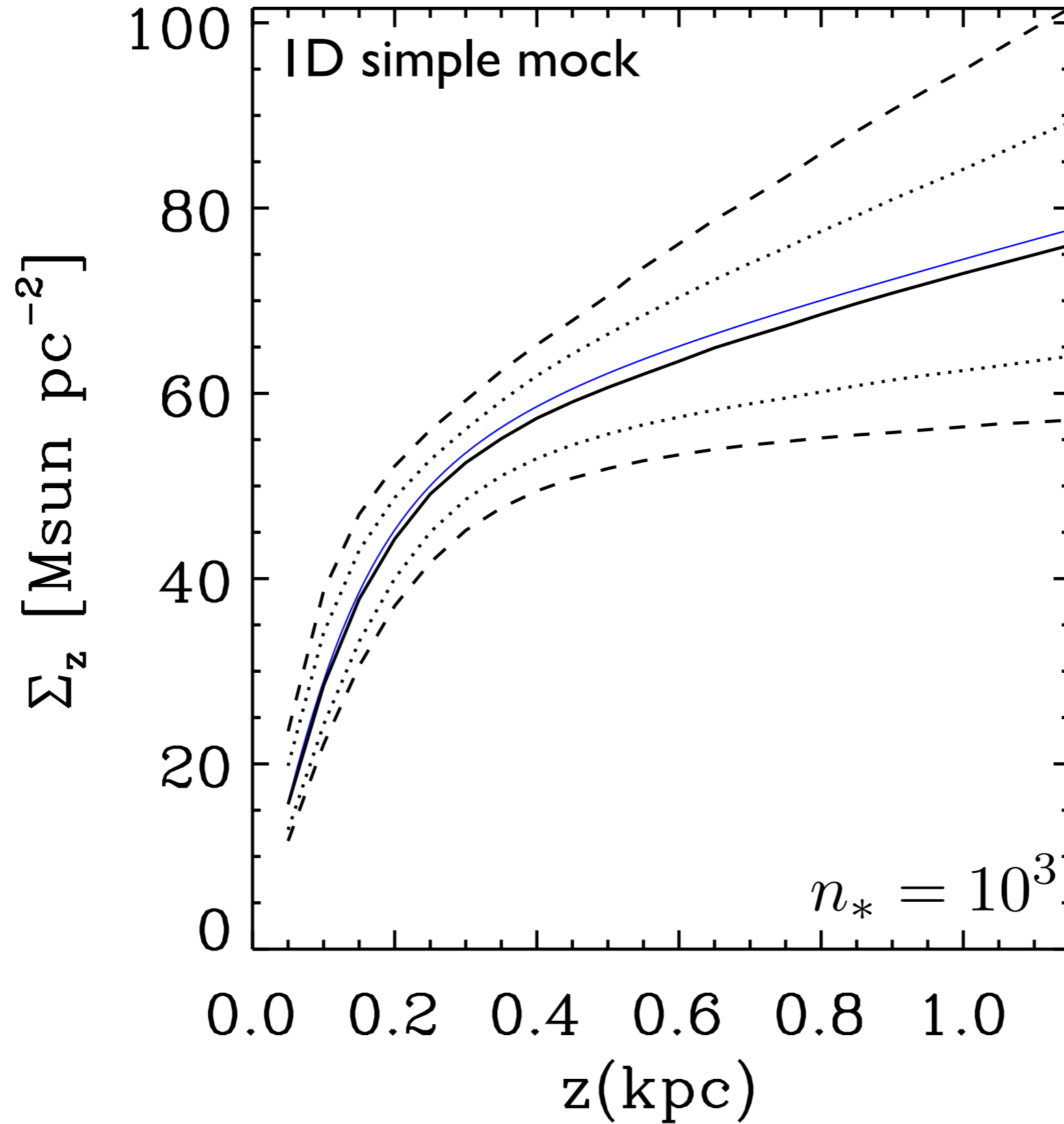
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$\propto \Sigma_z(z)$

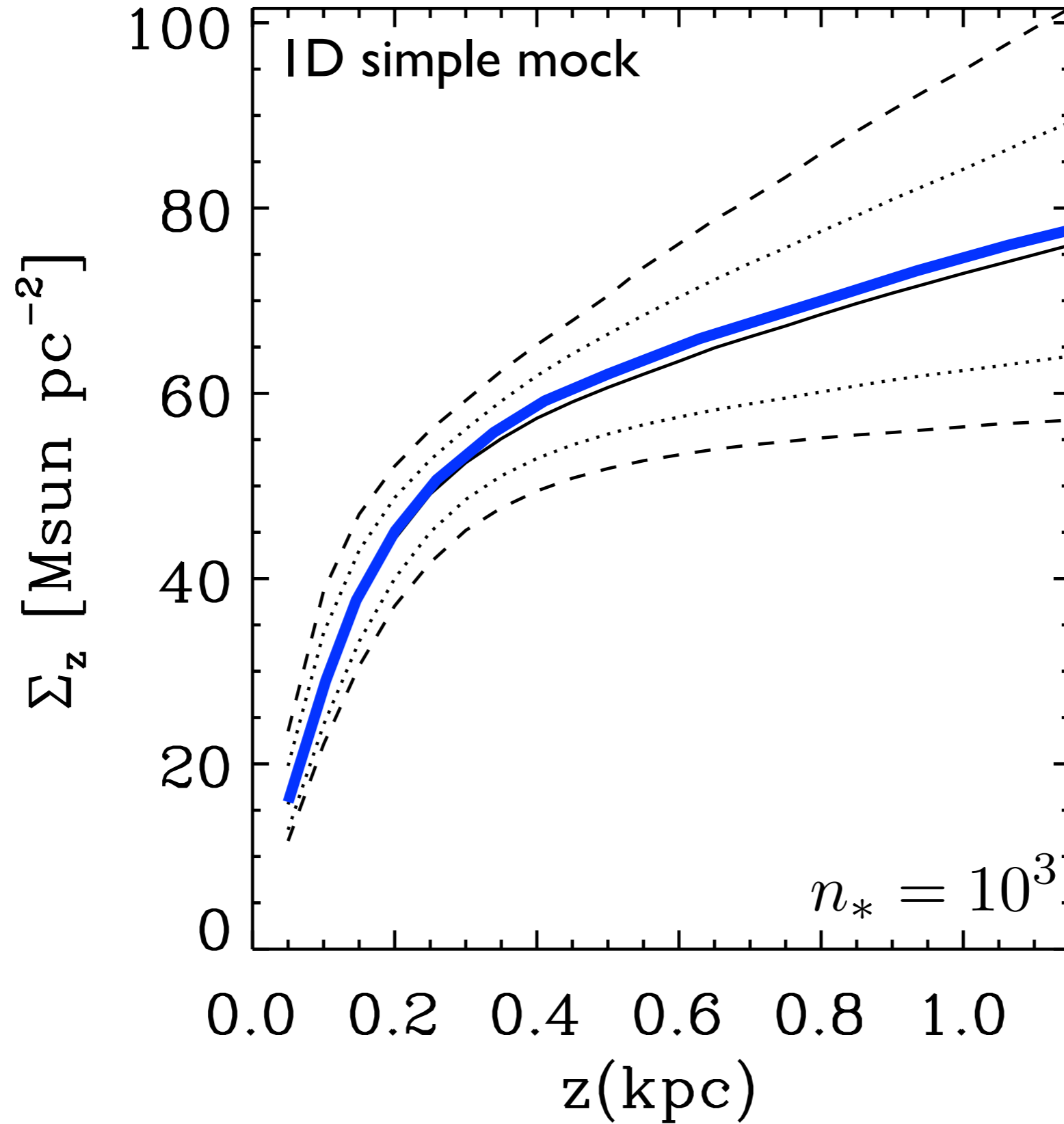
Measurement | Mock data



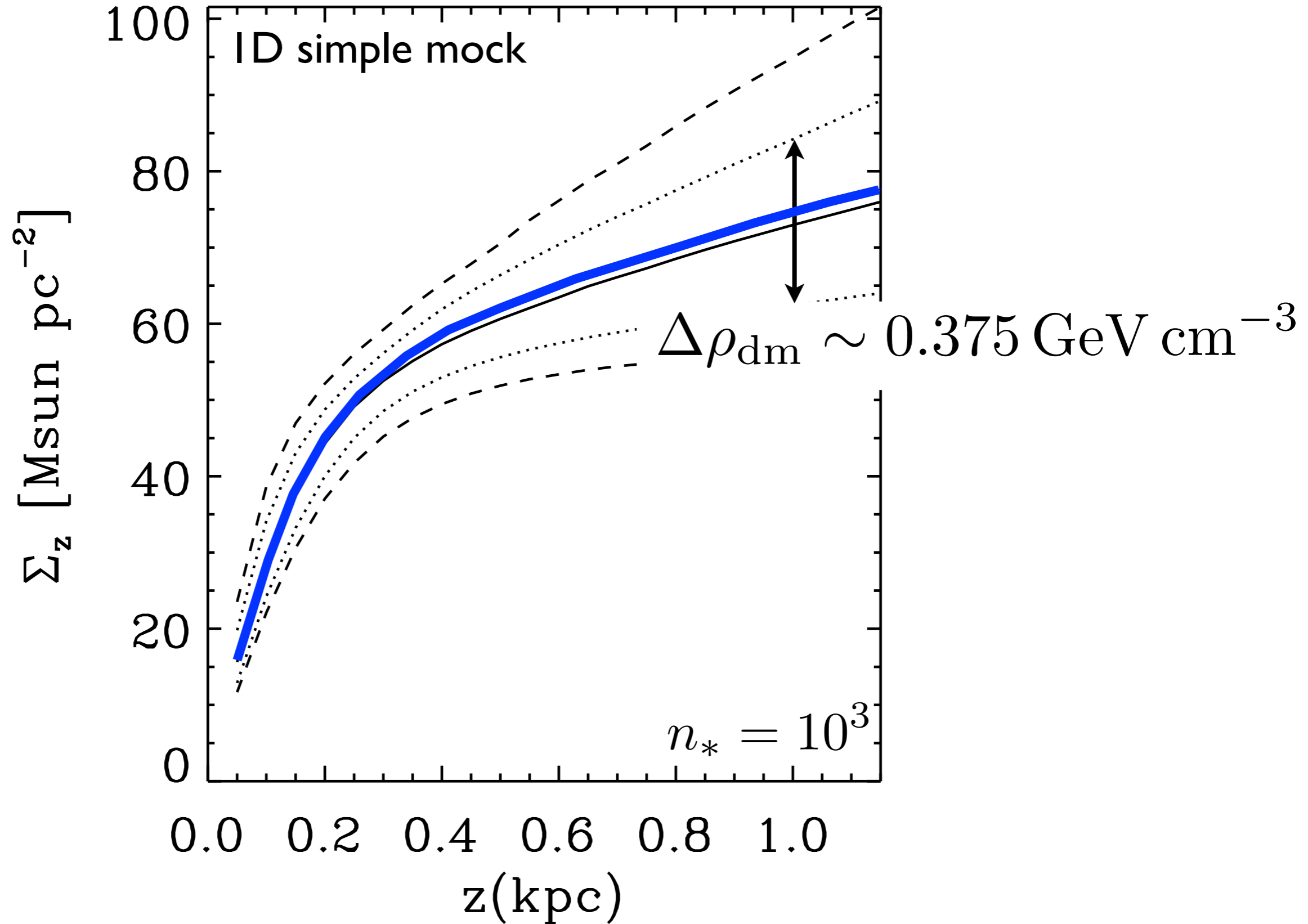
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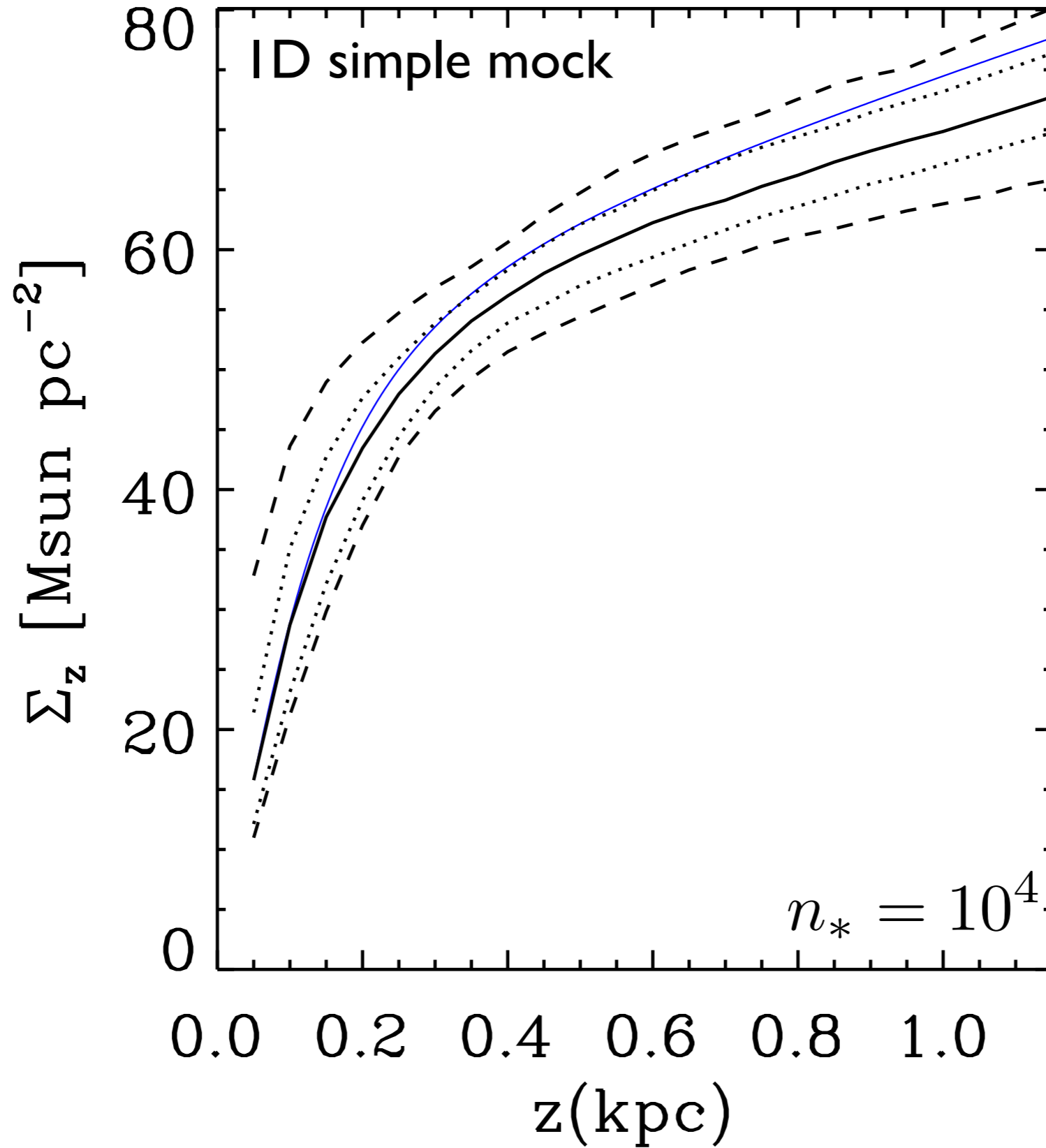
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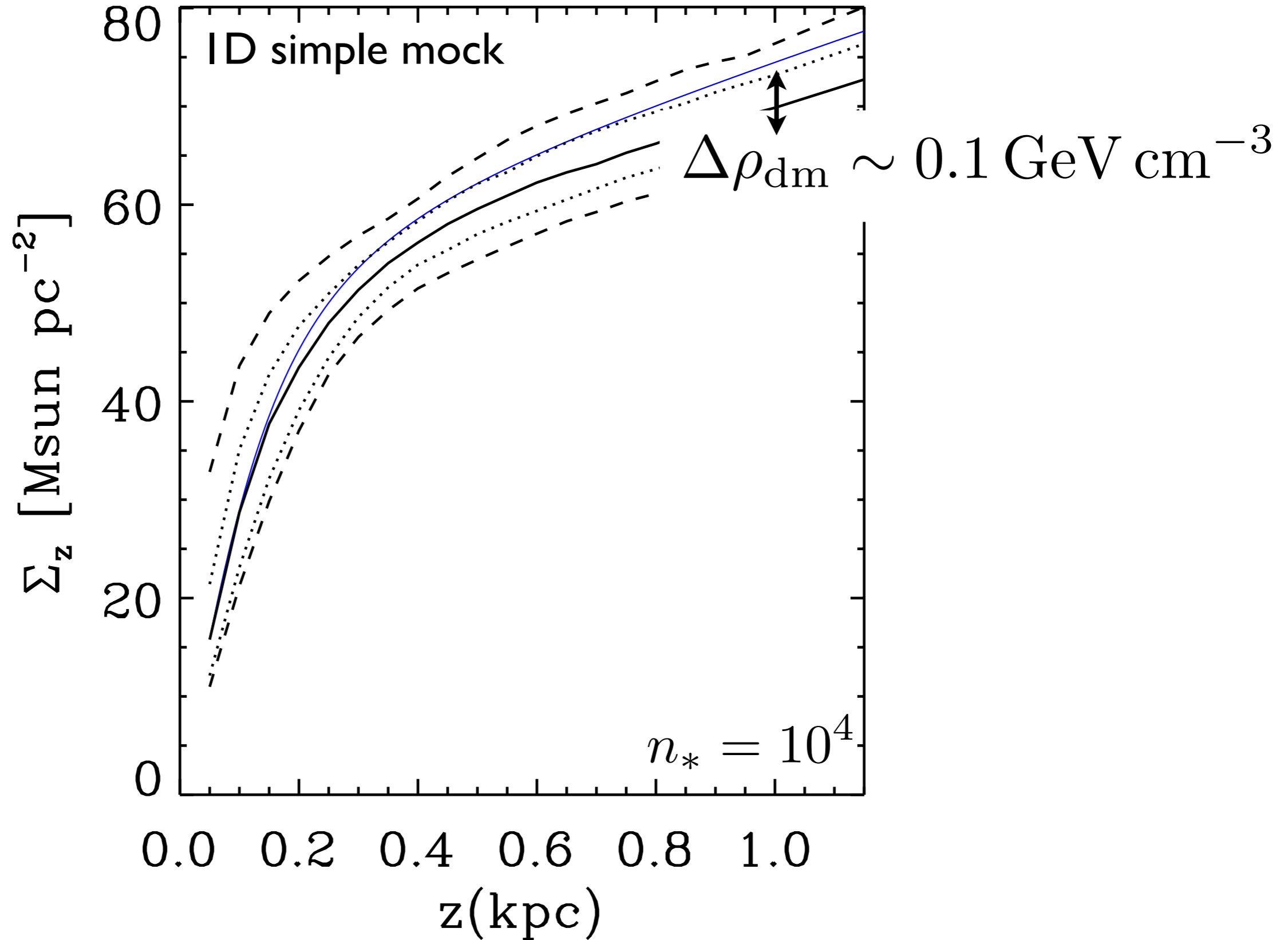
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Measurement | Mock data



Measurement | Real data

Need a good tracer:

Measurement | Real data

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- Well mixed \Rightarrow equilibrium
- Well populated \Rightarrow good statistics (at high z !)
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- Velocity data (v_z)
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a) 2000 stars

Volume Complete

[Kuijken & Gilmore 1989]

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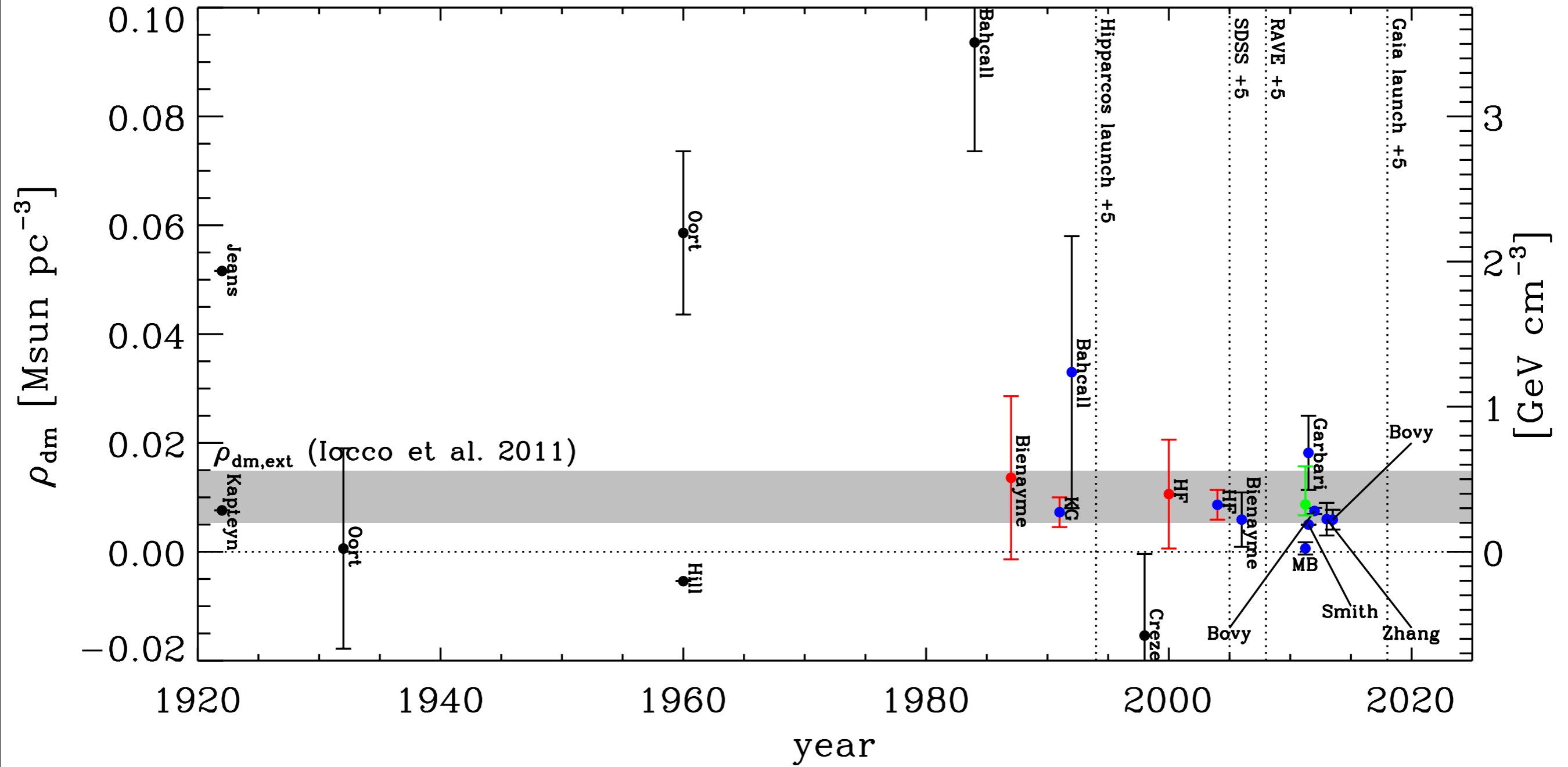
[Kuijken & Gilmore 1989]

b) 10,000 stars

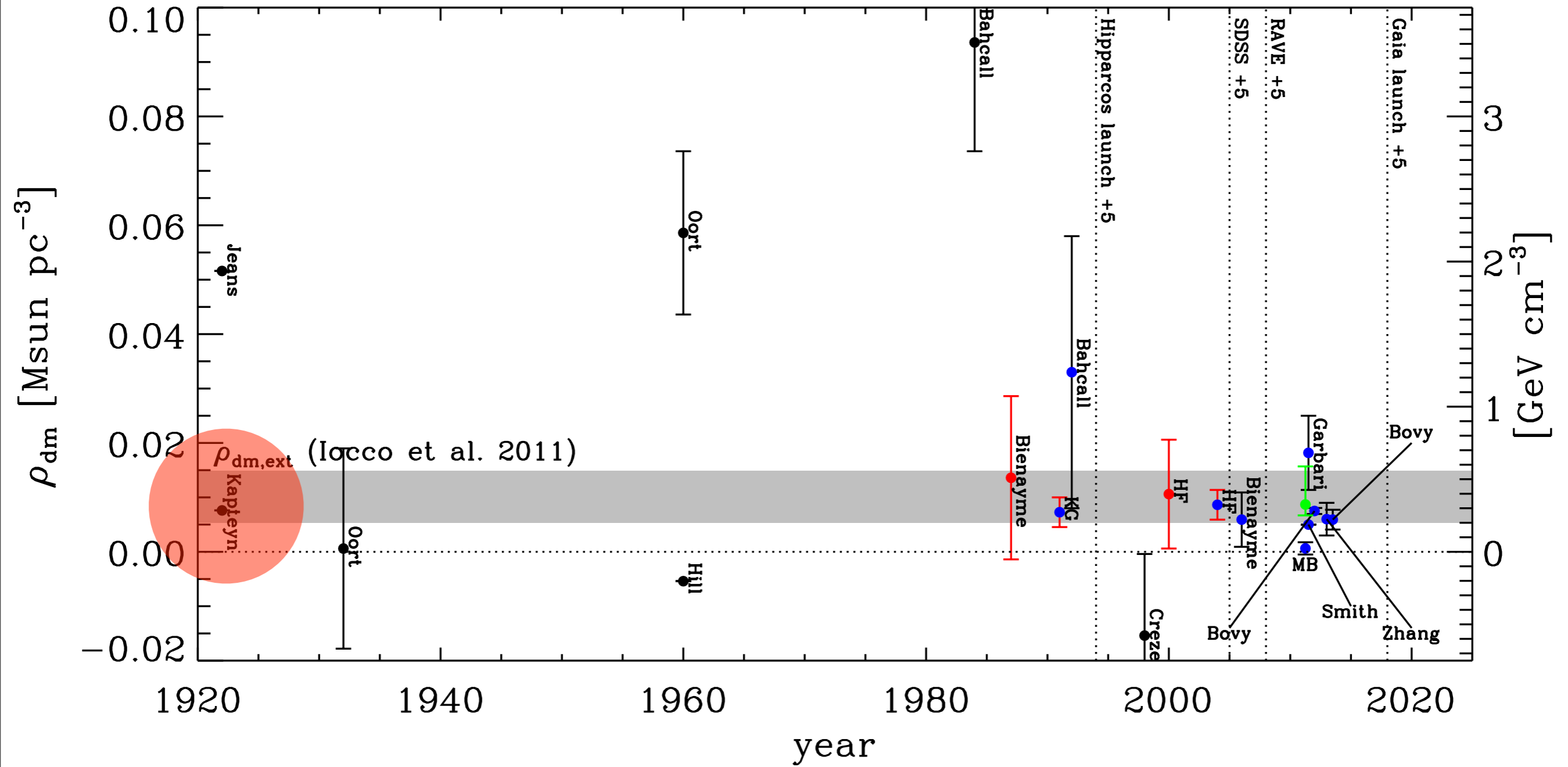
Complex SF

[Zhang et al. 2013]

Measurement | Historic measures



Measurement | Historic measures



Measurement | Historic measures

FIRST ATTEMPT AT A THEORY OF THE ARRANGEMENT AND MOTION OF THE SIDEREAL SYSTEM¹

BY J. C. KAPTEYN²

ABSTRACT

First attempt at a general theory of the distribution of masses, forces, and velocities in the stellar system.—(1) *Distribution of stars.* Observations are fairly well represented, at least up to galactic lat. 70° , if we assume that the equidensity surfaces are similar ellipsoids of revolution, with axial ratio 5.1, and this enables us to compute quite readily (2) *the gravitational acceleration at various points due to such a system*, by summing up the effects of each of ten ellipsoidal shells, in terms of the acceleration due to the average star at a distance of a parsec. The total number of stars is taken as 47.4×10^9 . (3) *Random and rotational velocities.* The nature of the equidensity surfaces is such that the stellar system cannot be in a steady state unless there is a general rotational motion around the galactic polar axis, in addition to a random motion analogous to the thermal agitation of a gas. In the neighborhood of the axis, however, there is no rotation, and the behavior is assumed to be like that of a gas at uniform temperature, but with a gravitational acceleration ($G\eta$) decreasing with the distance ρ . Therefore the density Δ is assumed to obey the barometric law: $G\eta = -\bar{u}^2(\delta\Delta/\delta\rho)/\Delta$; and taking the mean random velocity \bar{u} as 10.3 km/sec., the author finds that (4) *the mean mass of the stars* decreases from 2.2 (sun = 1) for shell II to 1.4 for shell X (the outer shell), the average being close to 1.6, which is the value independently found for the average mass of both components of visual binaries. In the galactic plane the resultant acceleration—gravitational minus centrifugal—is again put equal to $-\bar{u}^2(\delta\Delta/\delta\rho)/\Delta$, \bar{u} is taken to be constant and the average mass is assumed to decrease from shell to shell as in the direction of the pole. The angular velocities then come out such as to make the linear rotational velocities about constant and equal to 19.5 km/sec. beyond the third shell. If now we suppose that part of the stars are rotating one way and part the other, the relative velocity being 39 km/sec., we have a quantitative explanation of the phenomenon of star streaming, where

Measurement | Historic measures

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First attempt at a general theory of the distribution of the stars in the stellar system.—(1) Distribution of stars. Observed up to least up to galactic lat. 70° , if we assume the stars to be distributed in ellipsoids of revolution, with axial ratio 5.1, and readily (2) the gravitational acceleration at various distances, taking into account the effects of each of ten ellipsoidal shells, and adding to the average star at a distance of a parsec. 47.4×10^9 . (3) Random and rotational velocities. The general motion of the surfaces is such that the stellar system cannot be in a state of general rotational motion around the galactic axis, however, there is no rotation, and the behavior is like a gas at uniform temperature, but with a gravitation that varies with the distance ρ . Therefore the density Δ is determined by $G\eta = -\bar{u}^2(\delta\Delta/\delta\rho)/\Delta$; and taking the mean radial velocity \bar{u} as constant, the author finds that (4) the mean mass of the stars is 1.4 for shell X (the outer shell), the average being close to 1.6, which is the value independently found for the average mass of the components of the spiral arms. In the galactic plane the resultant acceleration—gravitational minus centrifugal—is again put equal to $-\bar{u}^2(\delta\Delta/\delta\rho)/\Delta$, \bar{u} is taken to be constant and the average mass is assumed to decrease from shell to shell as in the direction of the pole. The angular velocities then come out such as to make the linear rotational velocities about constant and equal to 19.5 km/sec. beyond the third shell. If now we suppose that part of the stars are rotating one way and part the other, the relative velocity being 39 km/sec., we have a quantitative explanation of the phenomenon of star streaming, where



Jacobus Cornelius Kapteyn
[1851 - 1922]

with the distance ρ . Therefore the density Δ is assumed to obey the barometric law. $G_{\eta} = -\bar{u}^2(\delta\Delta/\delta\rho)/\Delta$; and taking the mean random velocity \bar{u} as 10.3 km/sec., the **Measurement of the mean mass of the stars** decreases from 2.2 (sun = 1) for shell II to 1.4 for shell X (the outer shell), the average being close to 1.6, which is the value independently found for the average mass of both components of visual binaries. In the galactic plane the resultant acceleration—gravitational minus centrifugal—is again put equal to $-\bar{u}^2(\delta\Delta/\delta\rho)/\Delta$, \bar{u} is taken to be constant and the average mass is assumed to decrease from shell to shell as in the direction of the pole. The angular velocities then come out such as to make the linear rotational velocities about constant and equal to 19.5 km/sec. beyond the third shell. If now we suppose that part of the stars are rotating one way and part the other, the relative velocity being 39 km/sec., we have a quantitative explanation of the phenomenon of star-streaming, where the relative velocity is also in the plane of the Milky Way and about 40 km/sec. It is incidentally suggested that when the theory is perfected it may be possible to determine *the amount of dark matter* from its gravitational effect. (5) The *chief defects of the theory* are: That the equidensity surfaces assumed do not agree with the actual surfaces, which tend to become spherical for the shorter distances; that the *position of the center of the system* is not the sun, as assumed, but is probably some 650 parsecs away in the direction galactic long. 77° , lat. — mass of the stars was assumed to be the same in all shells in derivation for the variation of G_{η} with ρ on the basis of which the variation from shell to shell and the constancy of the rotational velocity were either the assumption or the conclusions are wrong; and that no comparison has been made between stars of different types.



1. *Equidensity surfaces supposed to be similar*

Mount Wilson Contribution No. 188³ a provisional derivation was given of the star-density in the stellar system. The question was

Measurement | Historic measures

It is incidentally suggested that when the theory is perfected it may be possible to determine *the amount of dark matter* from its gravitational effect.

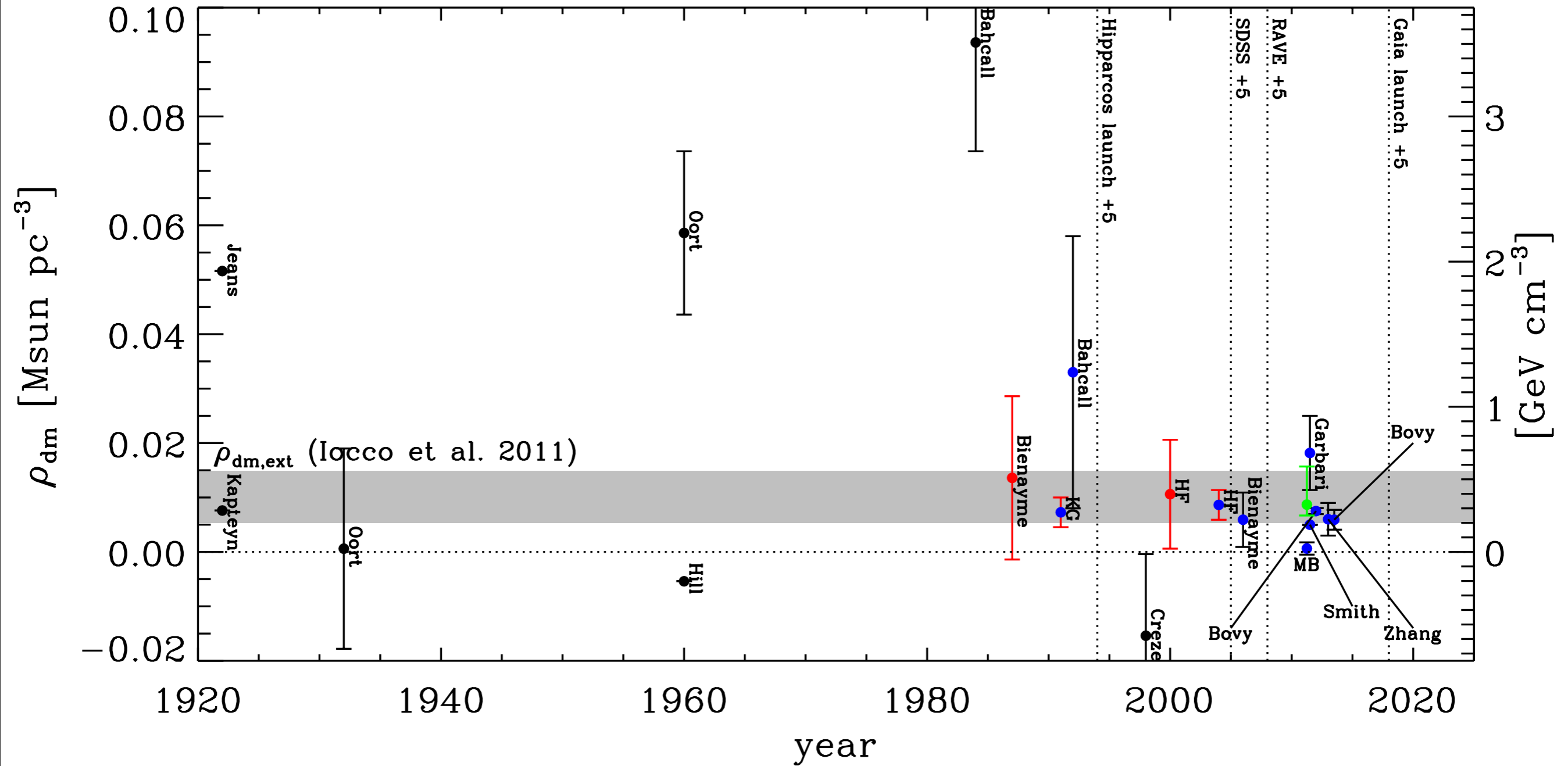


Measurement | Historic measures

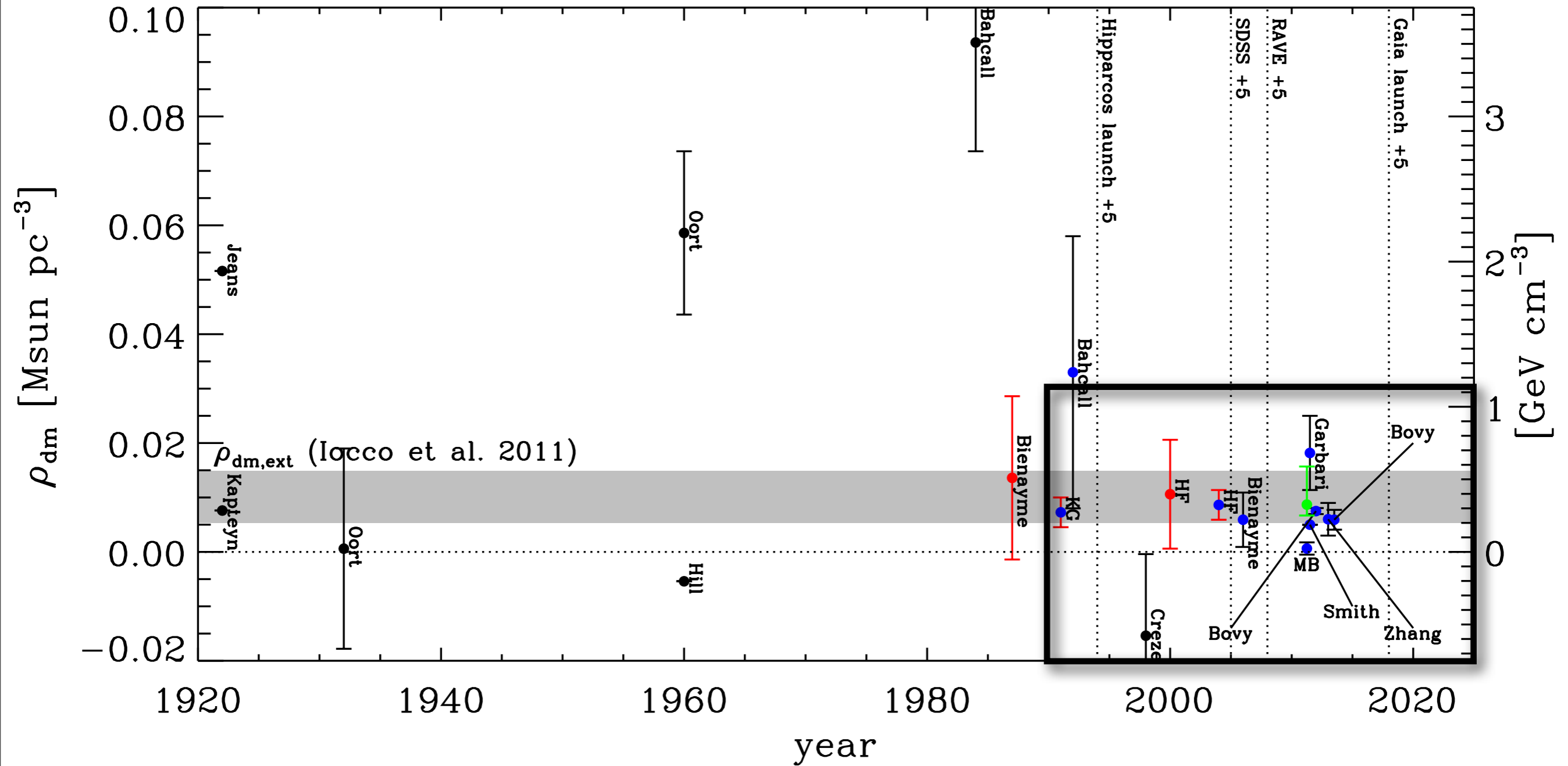
It is incidentally suggested that when the theory is perfected it may be possible to determine *the amount of dark matter* from its gravitational effect.



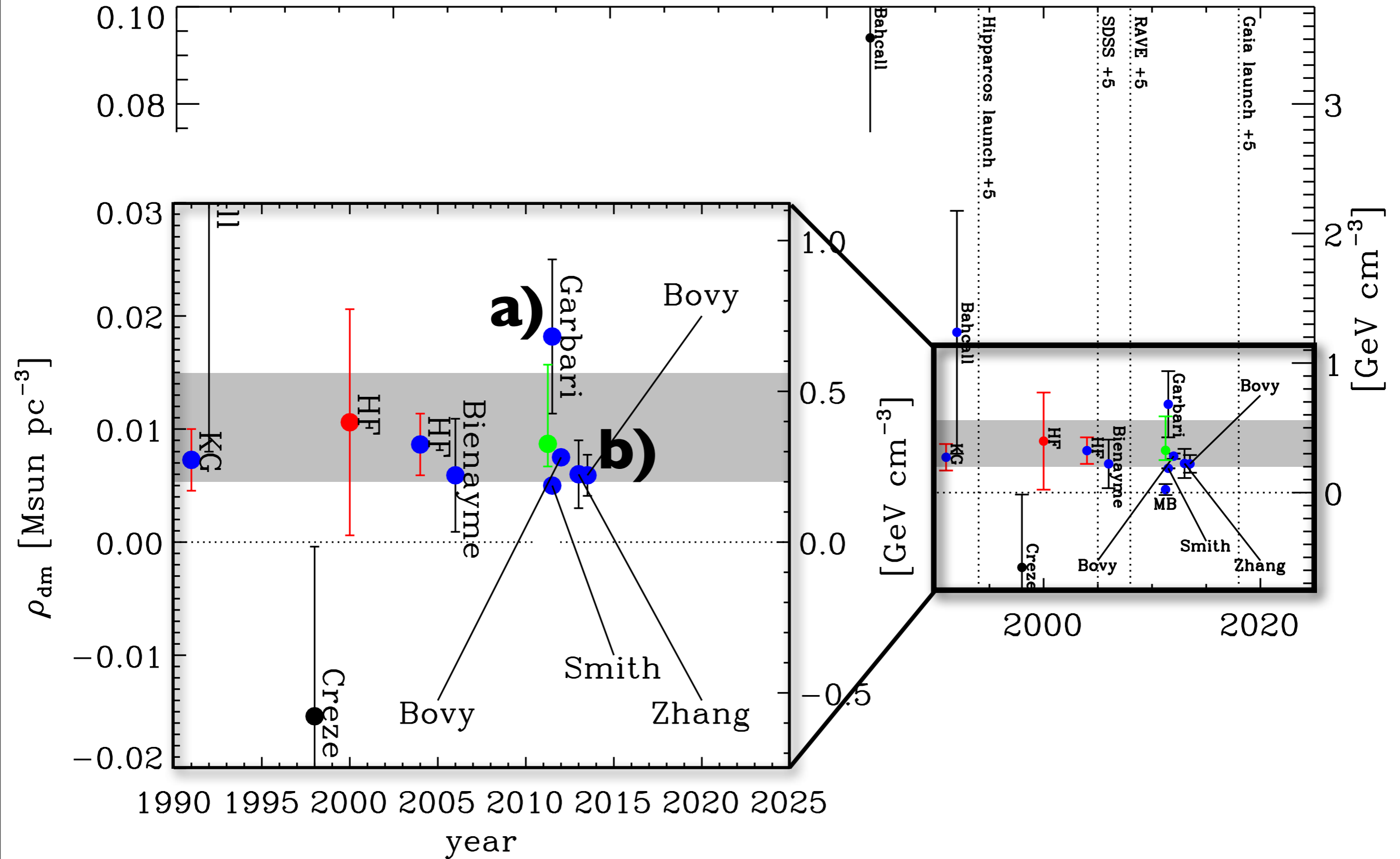
Measurement | Comparison of recent measures



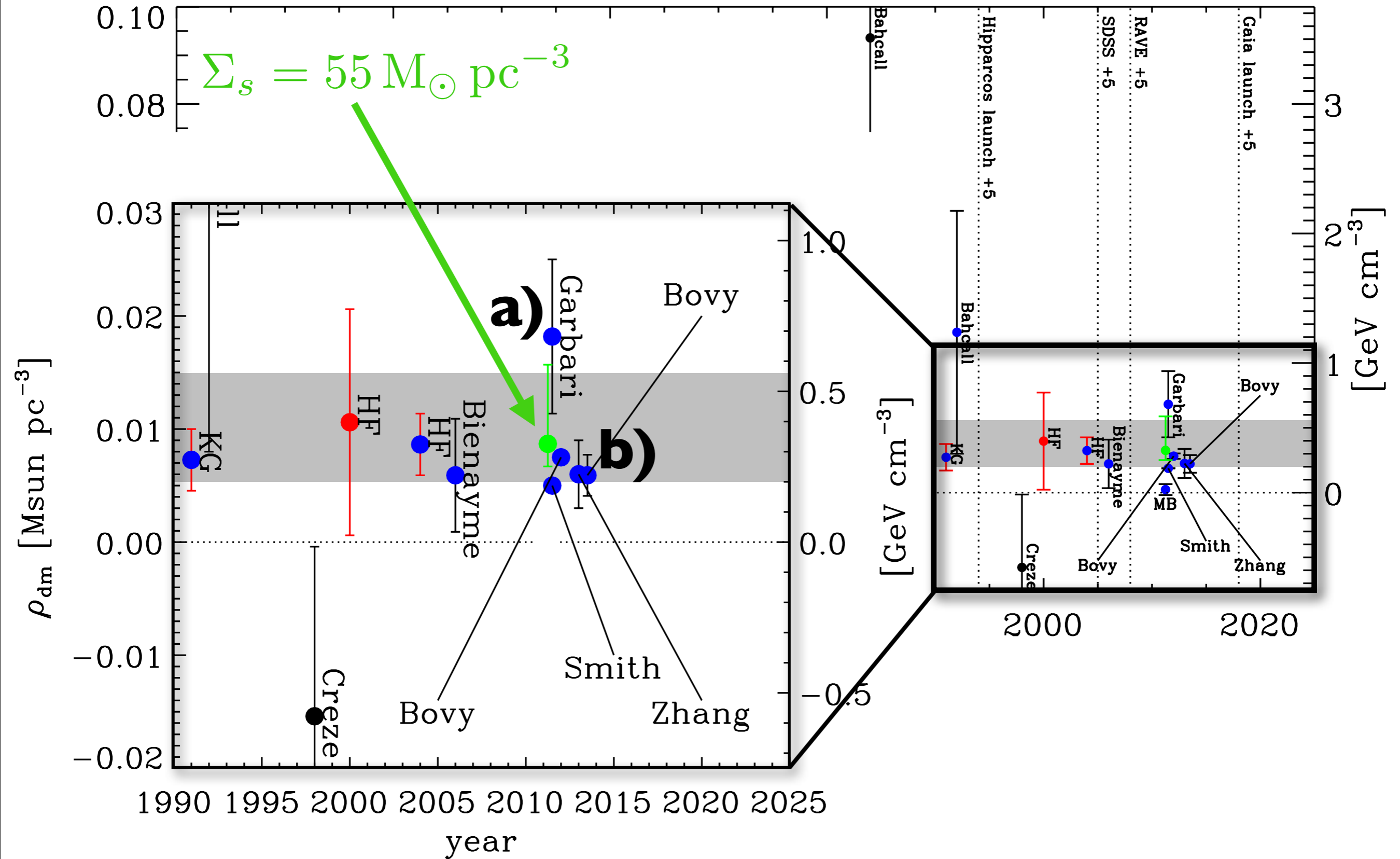
Measurement | Comparison of recent measures



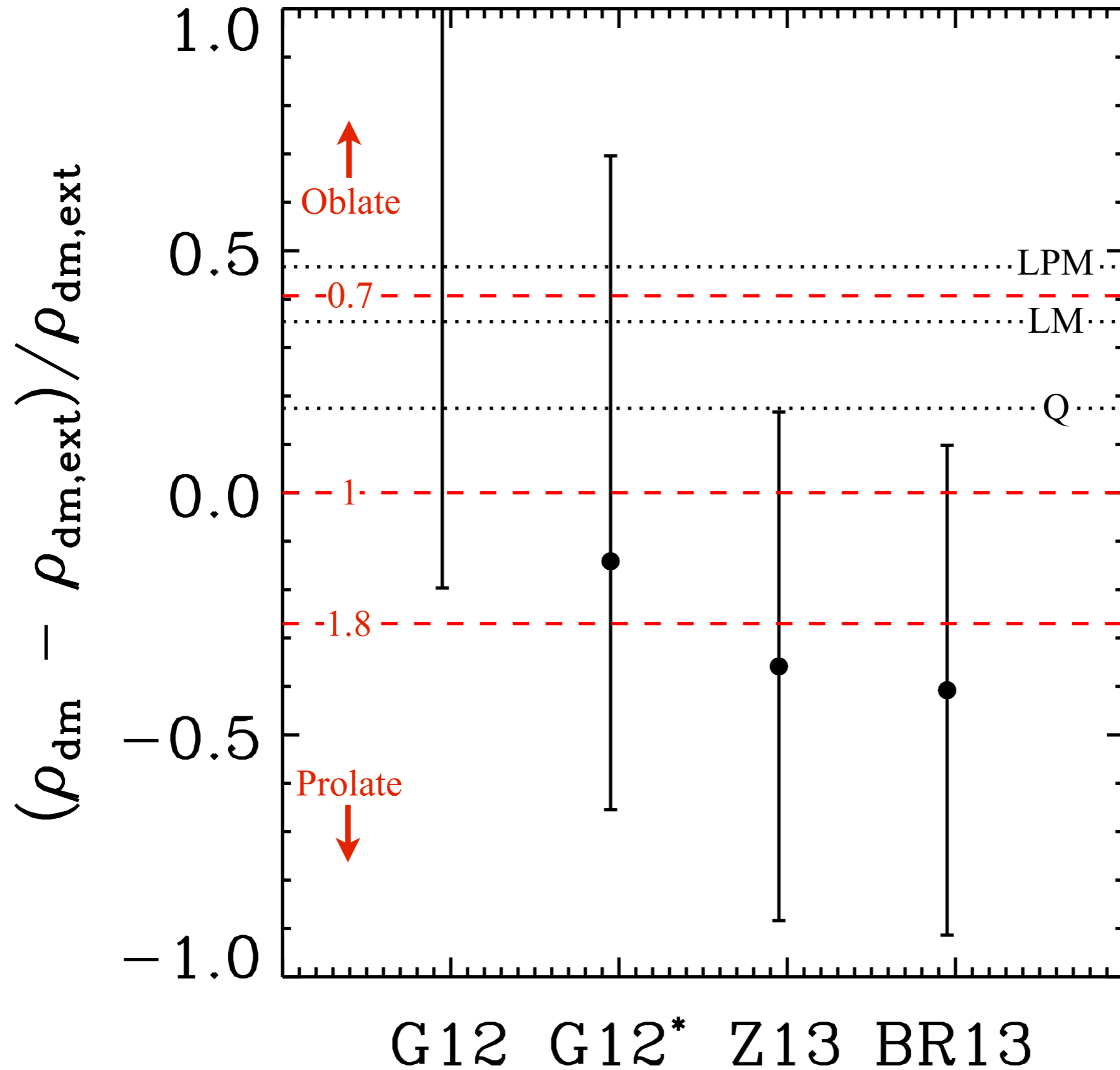
Measurement | Comparison of recent measures



Measurement | Comparison of recent measures

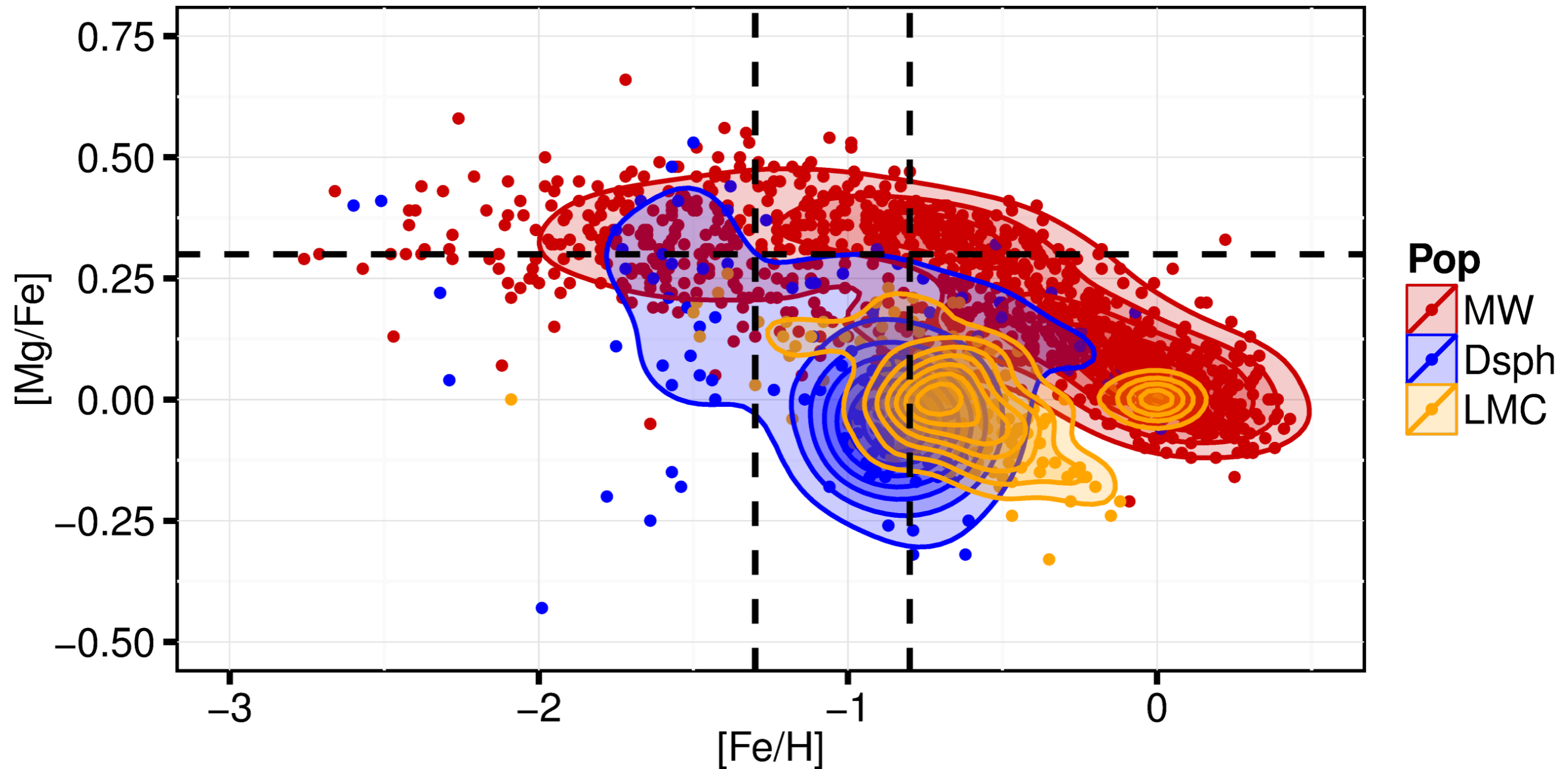


Measurement | The dark disc



Measurement | The dark disc | The hunt for accreted stars

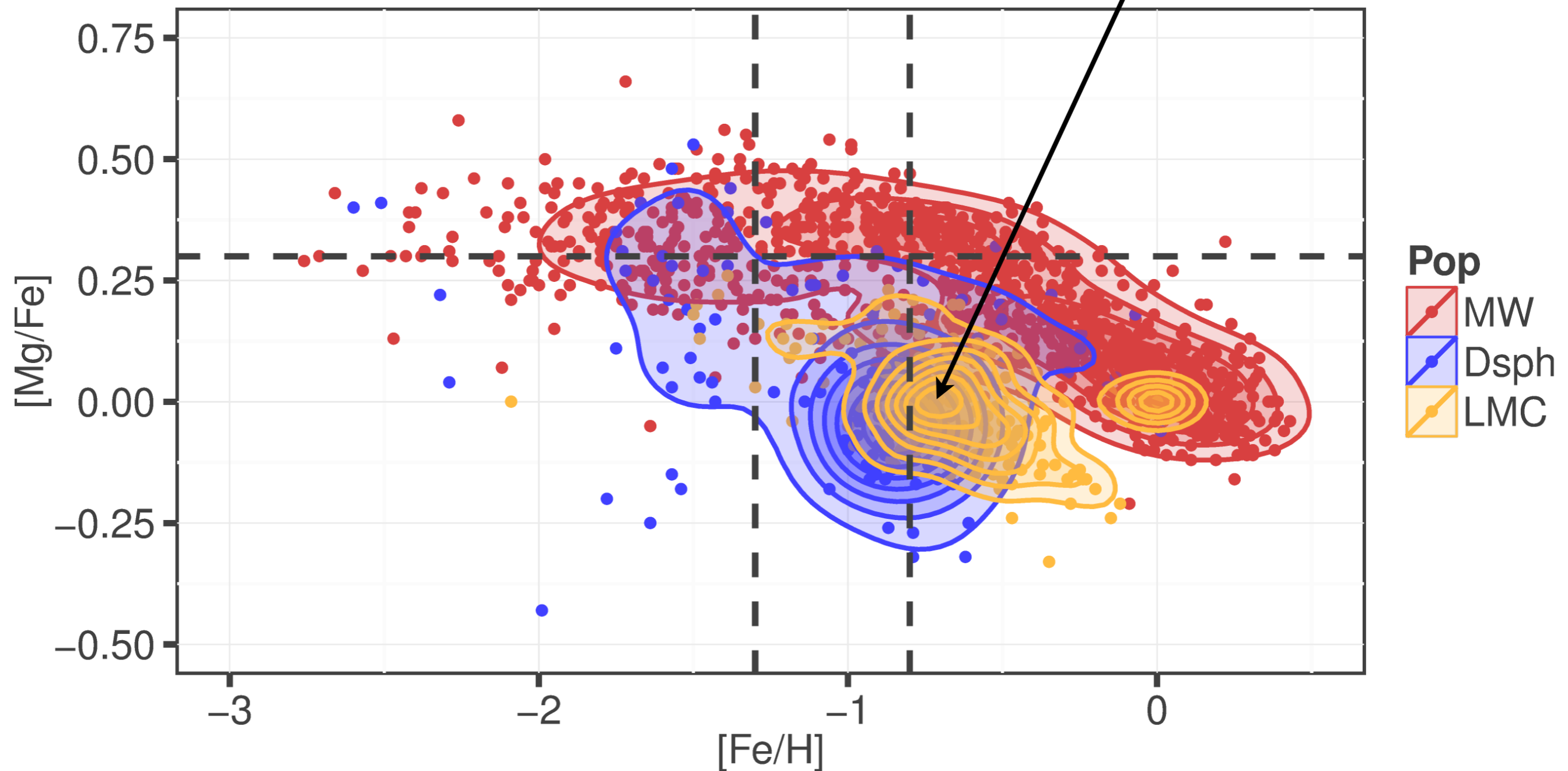
I. A chemical template ...



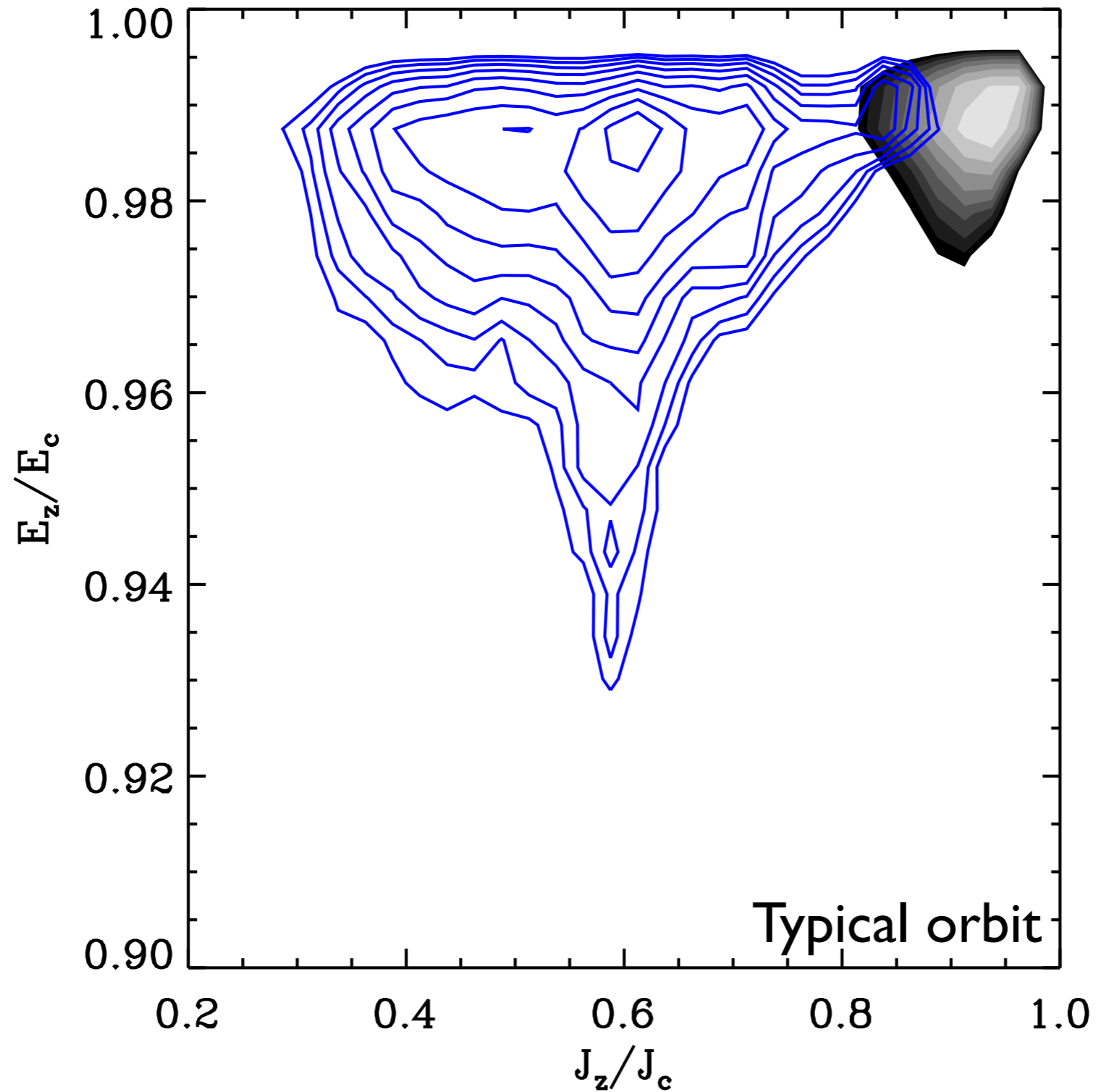
Measurement | The dark disc | The hunt for accreted stars

I. A chemical template ...

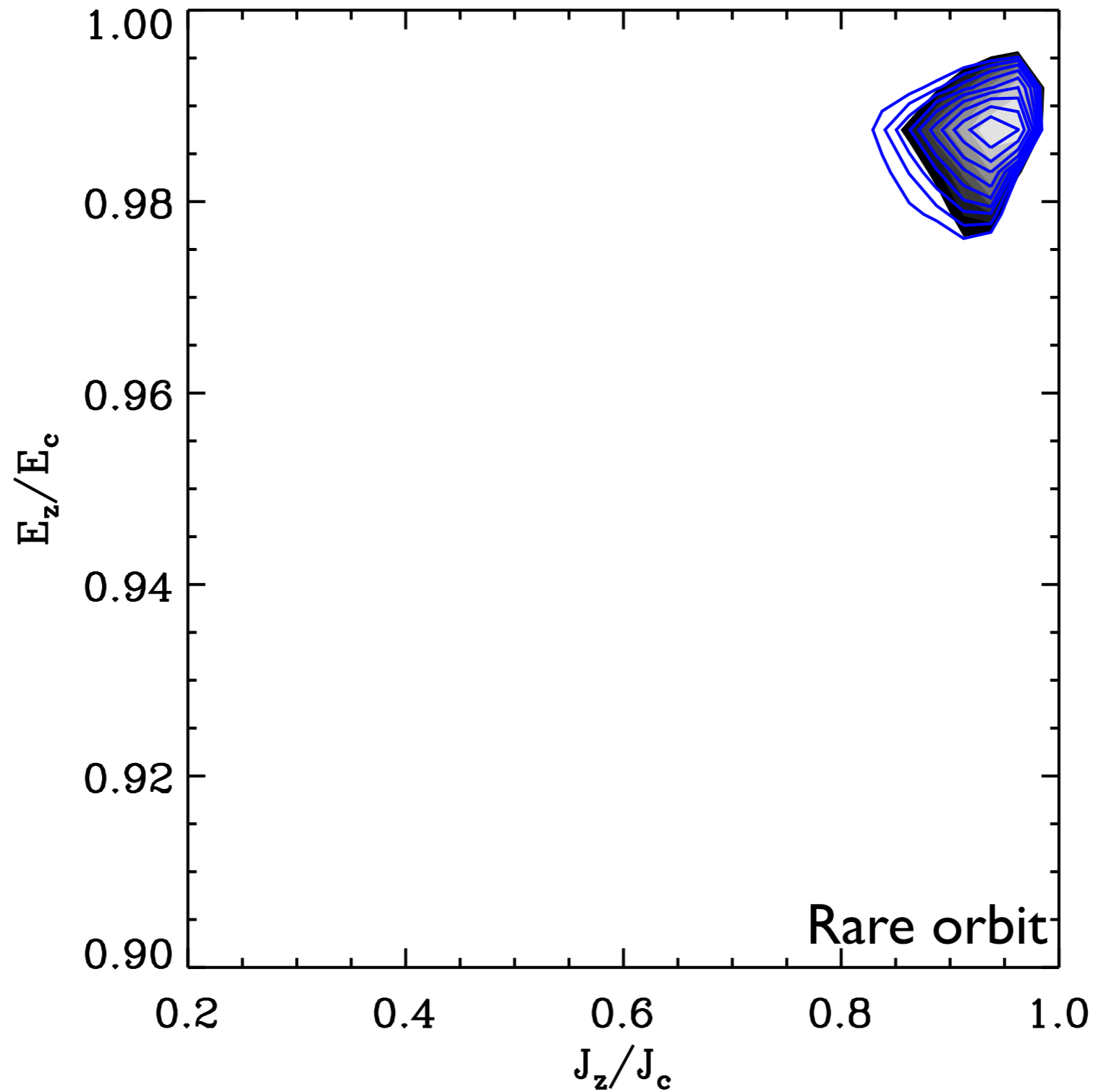
Born in dwarf or outer disc



2.A kinematic template ...

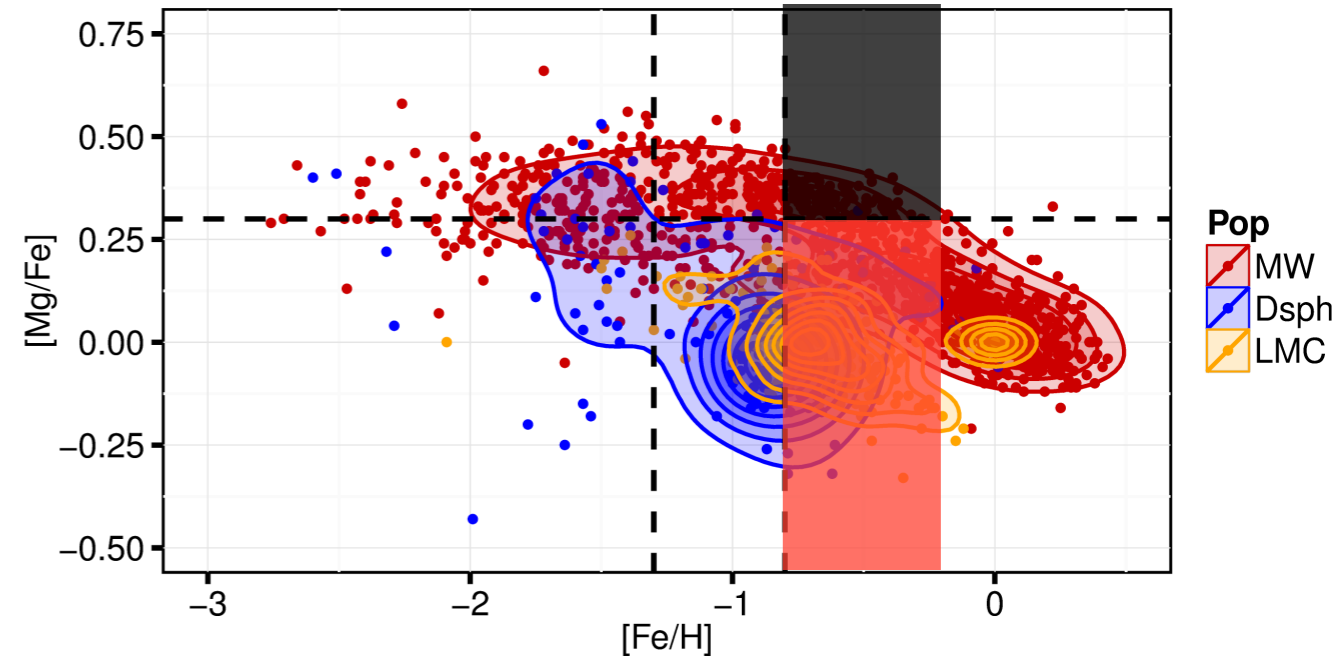
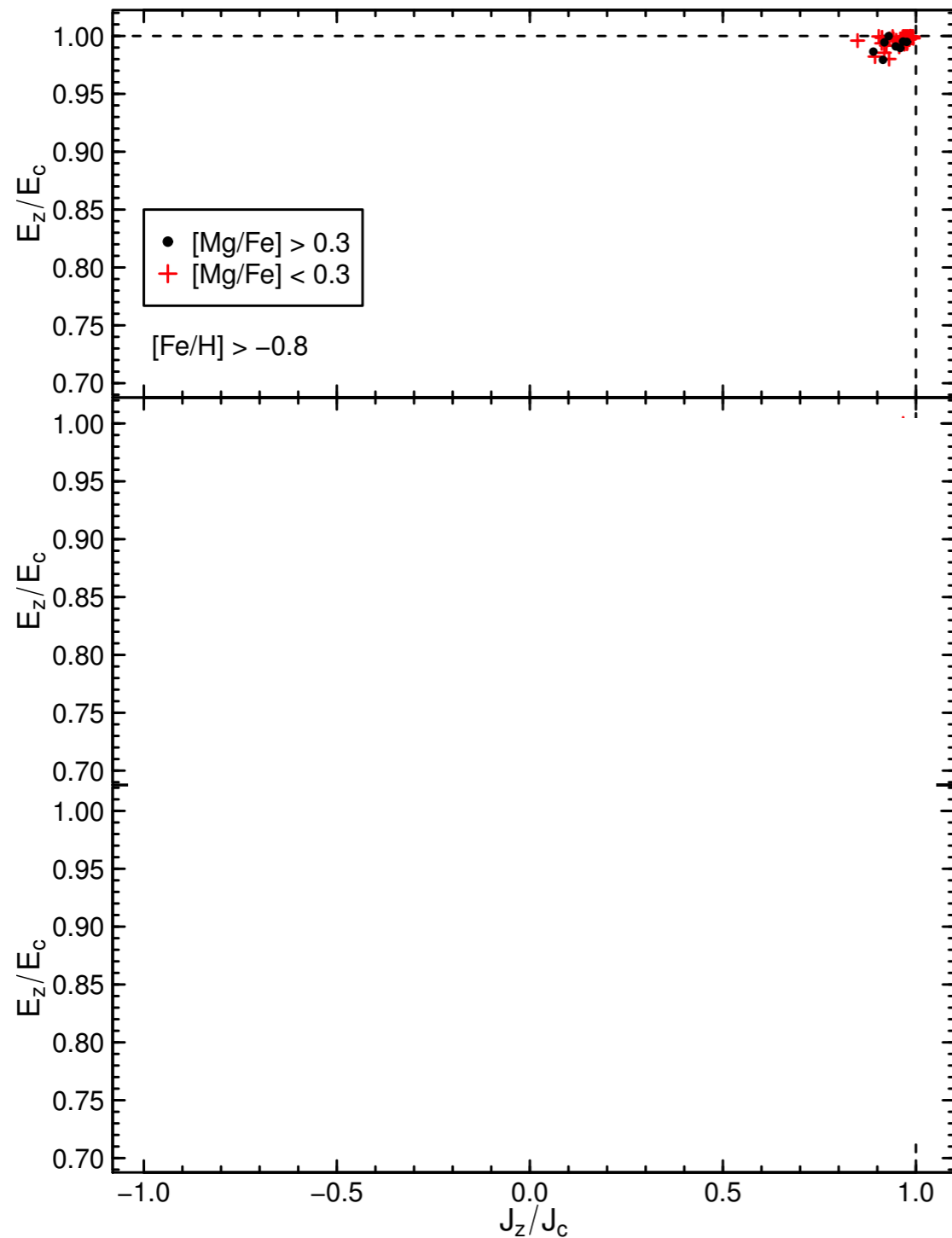


2.A kinematic template ...



Measurement | The dark disc | The hunt for accreted stars

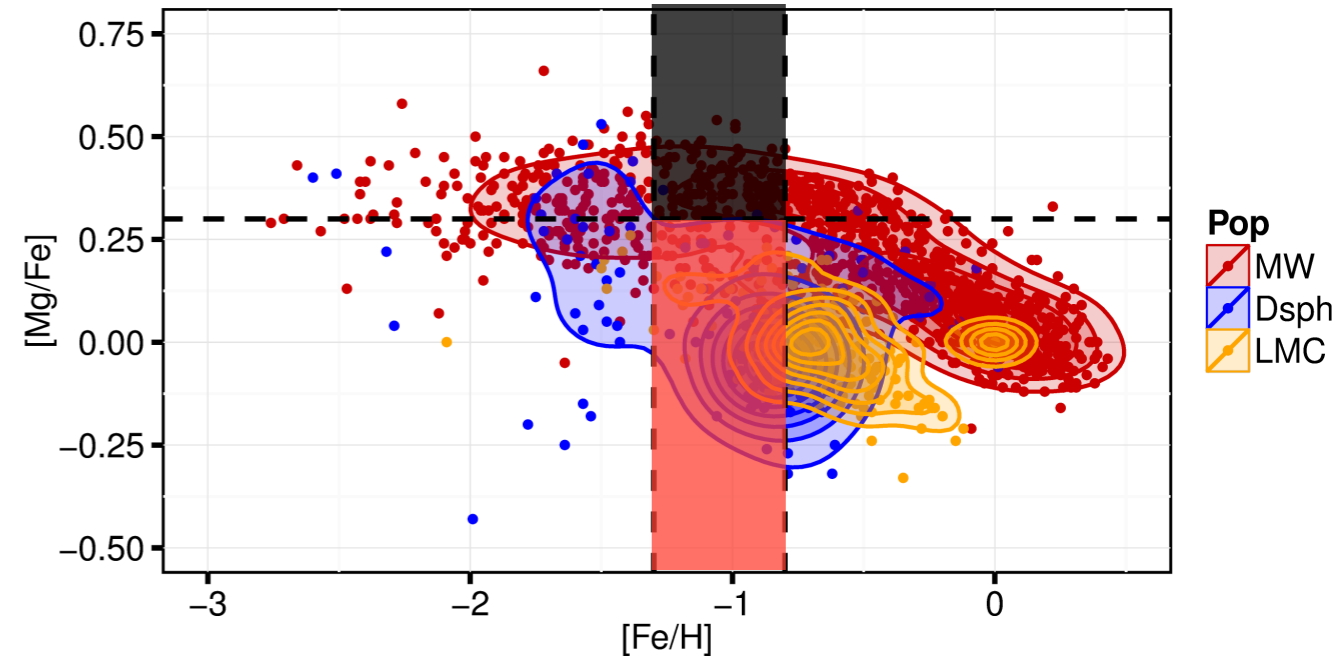
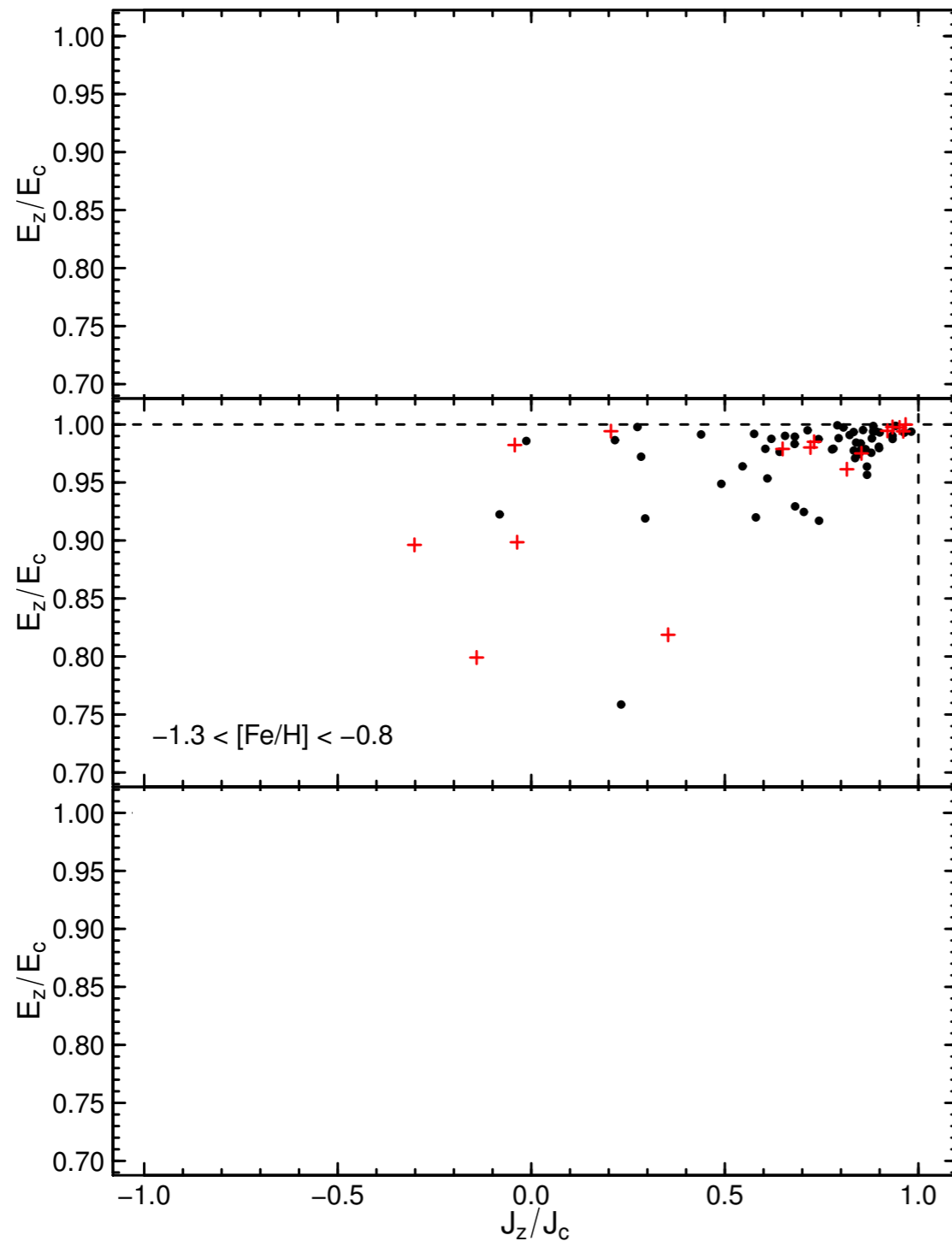
Real data [Ruchti et al. 2014 submitted]



Ruchti, Read et al. 2014, submitted

Measurement | The dark disc | The hunt for accreted stars

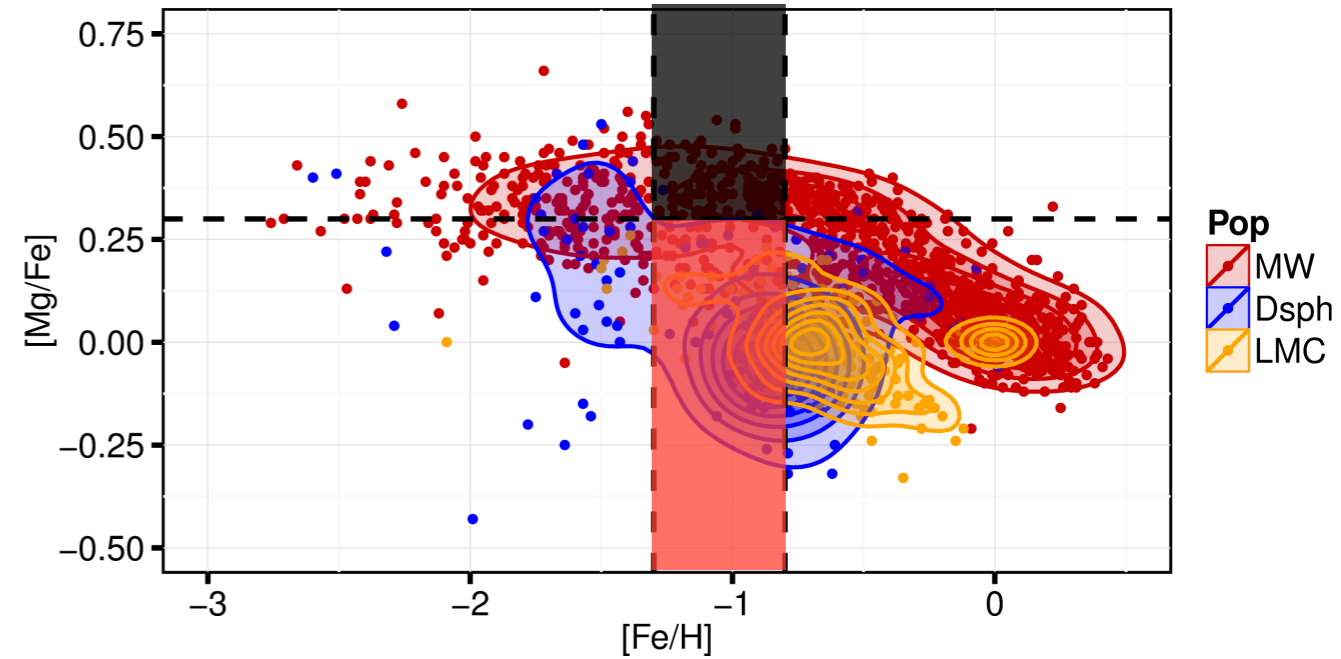
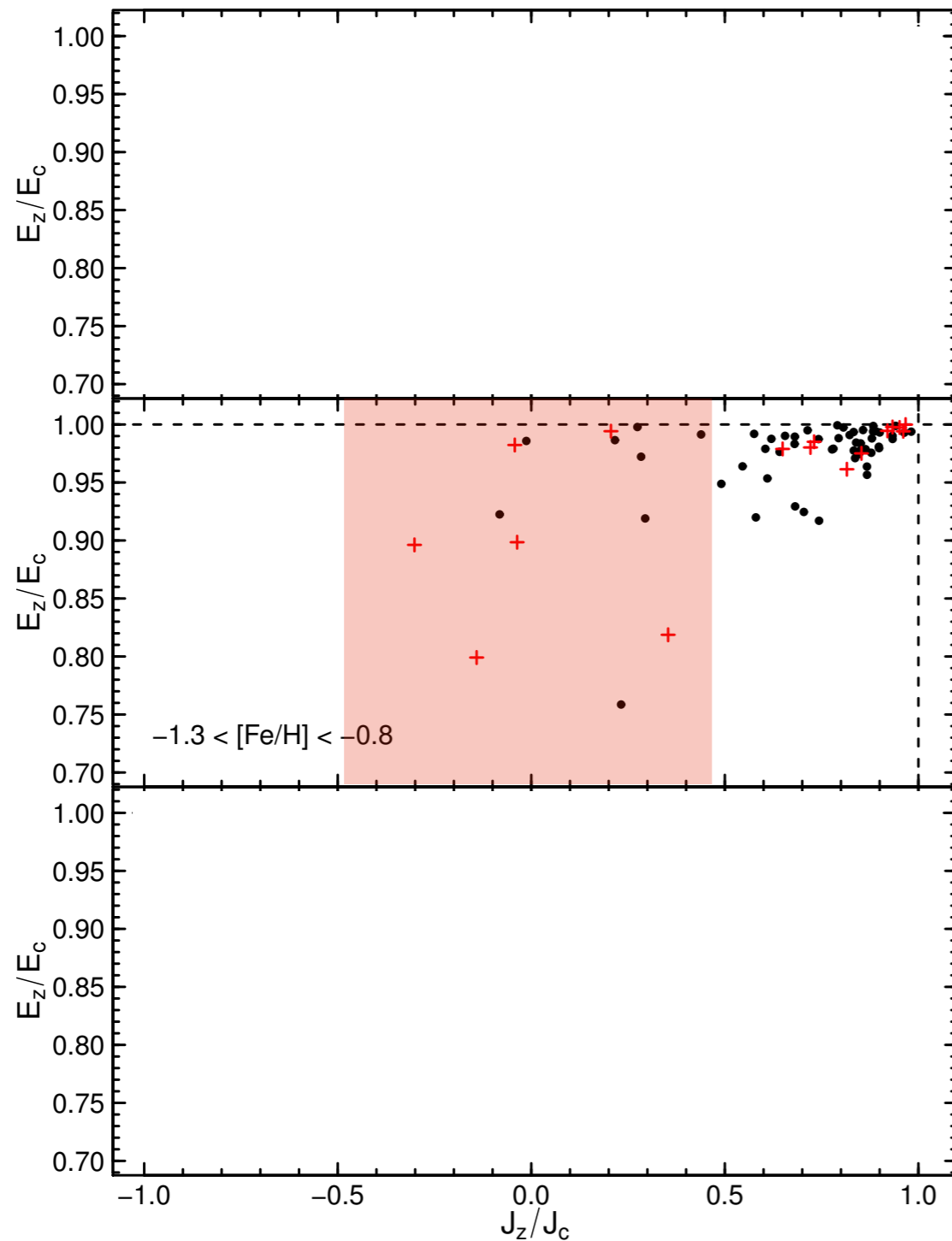
Real data [Ruchti et al. 2014 submitted]



Ruchti, Read et al. 2014, submitted

Measurement | The dark disc | The hunt for accreted stars

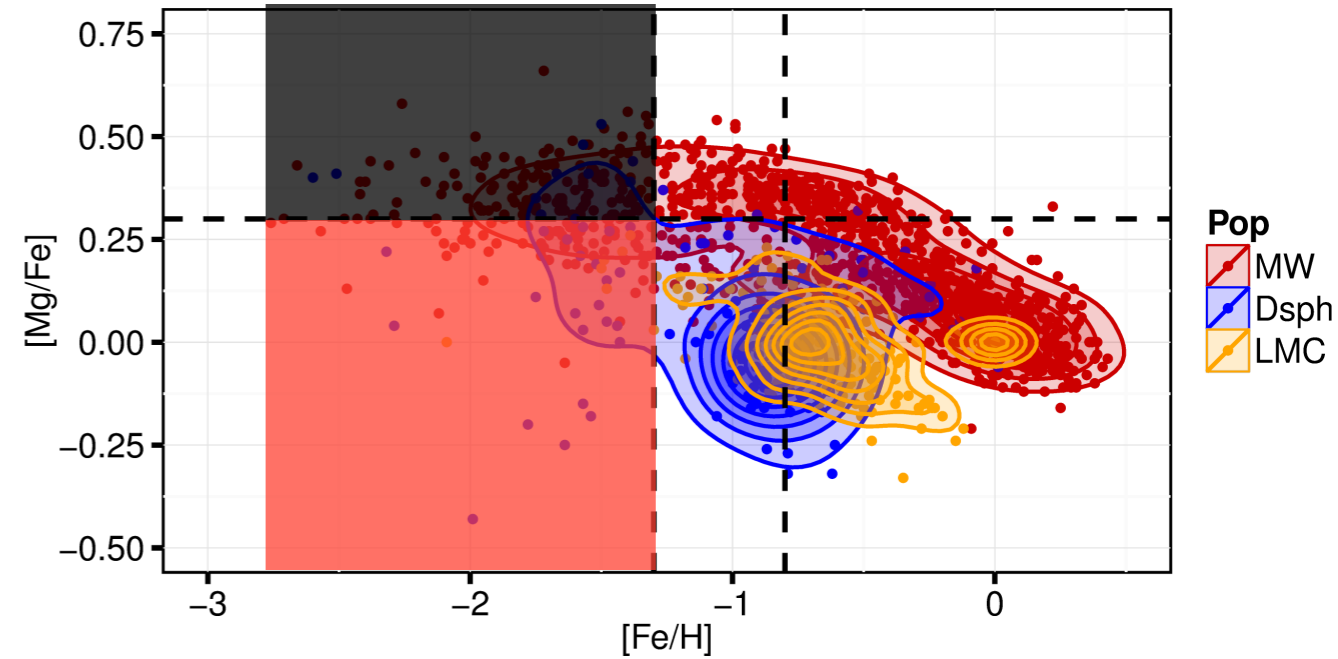
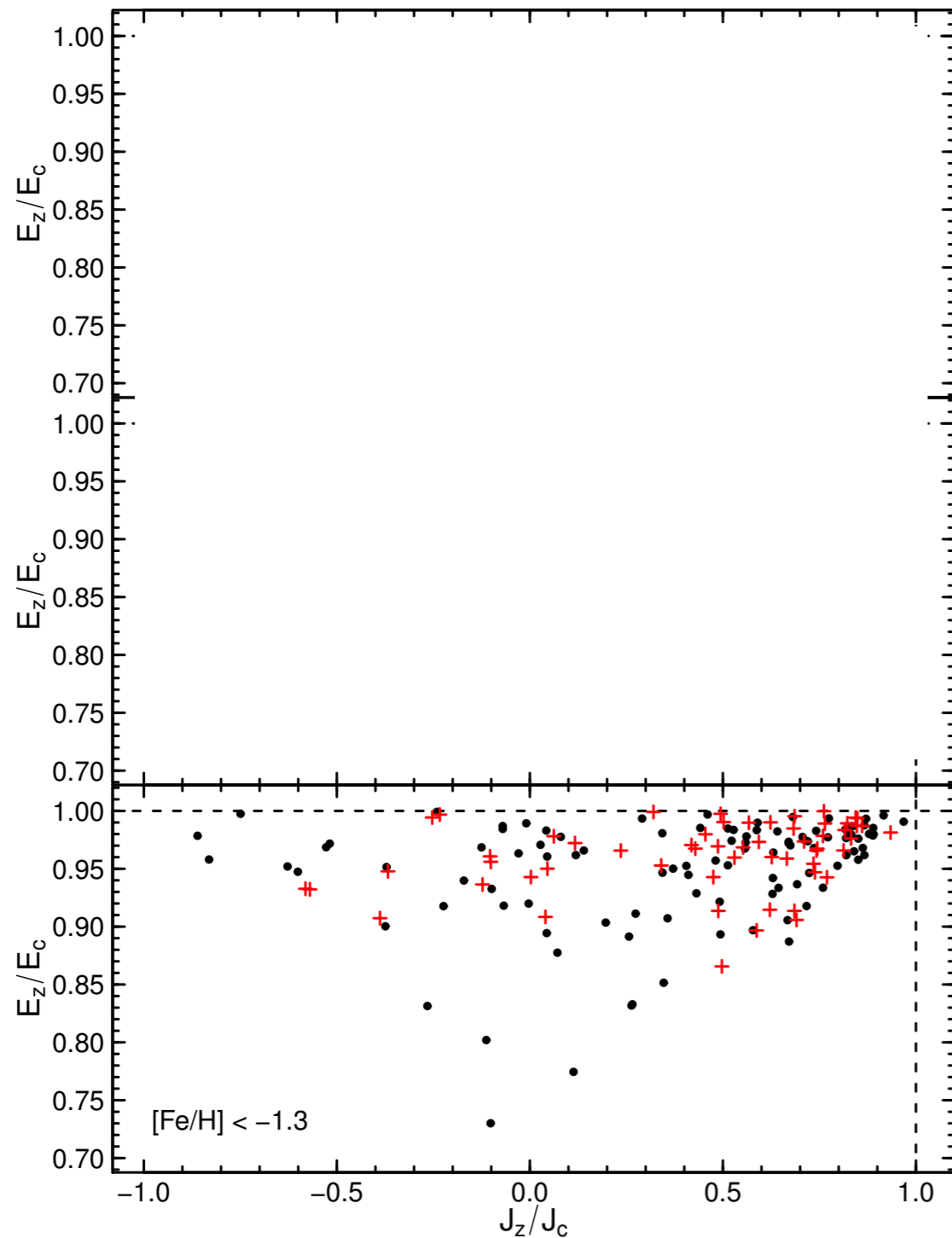
Real data [Ruchti et al. 2014 submitted]



Ruchti, Read et al. 2014, submitted

Measurement | The dark disc | The hunt for accreted stars

Real data [Ruchti et al. 2014 submitted]



Ruchti, Read et al. 2014, submitted

Conclusions

- The latest constraints on the local dark matter density give:

$$\rho_{\text{dm}} = 0.33^{+0.26}_{-0.075} \text{ GeV cm}^{-3} \quad \rho_{\text{dm}} = 0.25 \pm 0.09 \text{ GeV cm}^{-3}$$

[volume complete; G12*;R14] [SDSS; Z13]

- Comparing these with the rotation curve implies a near-spherical MW halo at $\sim 8\text{kpc}$, little dark disc, and a quiescent merger history.
- We have searched for stars accreted along with the dark disc, finding none so far; this supports the “quiescent MW” scenario.
- Gaia will move us into the realm of truly precise measurements of the Local Dark Matter Density.