

# Tuning of Generator LUARLW for R Scan at BESIII

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### **Outline**

- **Motivation**
- Some related problems need to solve
- Main LUARLW parameters need to tune
- Preliminary tuning results
- Work next to do

## Motivation

#### R scan at BESIII

#### BESIII finished two phases of R scan data taking:

First phase: (2012.05.28 – 06.08)

 $J/\Psi$  scan: 3.050 - 3.12 GeV

R scan: 2.23 - 3.40 GeV

Second phase: (2013.12.09 – 2014.01.26)

R scan: 3.85 - 4.60 GeV

Two scans cover the whole BEPCII energy region

#### Main goals:

- R value
- Resonant line shape and resonance parameters
- .....

### R measurement in experiment

Theoretical definition:

$$R = \frac{\sigma_{had}^{0}(e^{+}e^{-} \rightarrow \gamma^{*} \rightarrow \text{hadrons})}{\sigma_{\mu\mu}^{0}(e^{+}e^{-} \rightarrow \gamma^{*} \rightarrow \mu^{+}\mu^{-})}$$

Experimental expression:

$$R = \frac{N_{had} - N_{bg}}{\sigma_{\mu\mu}^{0} L \epsilon_{trg} \epsilon_{had} (1 + \delta)}$$

#### The basic tasks of R measurement:

- ➤ Data taking (raw data)
- > Data analysis (BG removing, hadronic selection, luminosity measurement)
- Theoretical studies (ISR calculation, generator tuning, efficiency)
- > Error analysis

#### MC and ISR in R measurement

#### R value experiment is to measure the hadronic cross section.

Observed total cross section:  $\sigma_{obs}^{tot} = \frac{N_{evt}}{r}$ 

$$\sigma_{obs}^{tot} = \frac{N_{evt}}{L}$$

Not physical quantity, experiment dependent.

 $\text{Total cross section in physics:} \quad \sigma_{phys}^{tot} = \frac{N_{evt}}{L\bar{\epsilon}} \quad \begin{array}{c} \text{Physical quantity,} \\ \text{experiment independent.} \end{array}$ 

$$\sigma_{phys}^{tot} = \frac{N_{evt}}{L\bar{\epsilon}}$$

Hadronic detection efficiency: 
$$\bar{\epsilon} = \frac{N_{exp}^{obs}}{N_{exp}^{gen}} \approx \frac{N_{MC}^{obs}}{N_{MC}^{gen}}$$

Born cross section in physics:

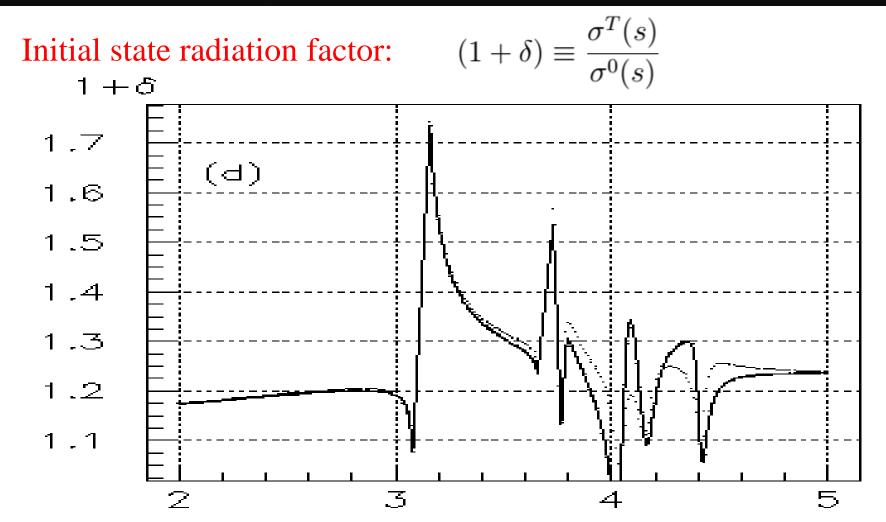
$$\sigma_{phys}^0 = \frac{N_{evt}}{L\bar{\epsilon}(1+\delta)}$$

ISR factor: 
$$\sigma(s)_{the}^{tot} = \int dx \sigma_{the}^{0}(s') F(x;s) \equiv \sigma_{the}^{0}(s) (1+\delta)$$

Main error:

$$\frac{\Delta\sigma_{phys}^{0}}{\sigma_{phys}^{0}} = \sqrt{(\frac{\Delta N_{evt}}{N_{evt}})^{2} + (\frac{\Delta L}{L})^{2} + (\frac{\Delta \bar{\epsilon}}{\bar{\epsilon}})^{2} + (\frac{\Delta (1+\delta)}{(1+\delta)})^{2}}$$

### Energy dependence of $(1+\delta)$



The ISR factor  $(1+\delta)$  is energy dependent. Energy dependence of  $(1+\delta)$  reflects the energy dependence of the Born cross section  $\sigma_7^0(s)$ .

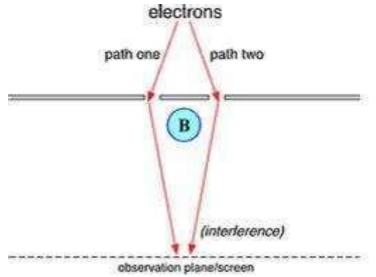
### Problems need to solve

#### **Interference:**

According to quantum mechanics principle, if an exclusive hadronic state can produce via *N* channels, it has *N* Feynman diagrams with amplitude *A*i, the differential cross section:

$$\frac{d\sigma}{d\Omega} \propto |\sum A_i|^2 \neq \sum |A_i|^2$$

The measured distributions certainly contain interference effects.



#### **Interference:**

In R scan above 3.85 GeV, DD states can produce via both continuous and resonant modes, the differences between them are the propagators

$$e^{+}e^{-}\Rightarrow \underline{\gamma^{*}+\psi}\Rightarrow \begin{array}{c} D\bar{D};\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s};\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s},D_{s}\bar{D}_{s}^{*};\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s},D_{s}\bar{D}_{s}^{*},D_{s}^{*}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D^{*}\bar{D}^{*},D\bar{D}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D_{s}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D^{*}\bar{D}^{*},D\bar{D}^{*},D^{*}\bar{D}^{*},D^{*}\bar{D}^{*},D^{*}\bar{D}_{s}^{*},D^{*}\bar{D}_{s}^{*},D\bar{D}_{1}^{*},D\bar{D}_{2}^{*},\\ D\bar{D},D^{*}\bar{D}^{*},D^{$$

In QCD, the expression of amplitude *Ac* of continuous channel are function of form factors. But in most case, form factors are unknown.

Amplitude of resonance *Ar* is described by Breit-Wigner form:

$$Ar \equiv \mathcal{T} = \frac{M\sqrt{\Gamma^e \Gamma^h}}{W^2 - M^2 + iM\Gamma_{tot}} e^{i\delta}$$

For wide resonance, like  $\Psi(4040)$ ,  $\Psi(4160)$ ,  $\Psi(4415)$ , the hadronic widths are energy dependent.

#### **Energy dependence of hadronic width:**

for unstable particle, the width  $\Gamma = 1$ / lifetime  $\tau$ . the calculation of  $\Gamma$  relay on nonperturbative QCD, one has to employ phenomenological models. For example:

• Simple squared-potential well model:

final state decay momentum

$$\Gamma_{\rm R}^{
m f_c}({
m s}) = {
m G_R} \cdot \sum_{
m L} rac{Z_{
m f_c}^{
m 2L+1}}{B_{
m L}^{
m f_c}} \qquad Z_{
m f_c} = {
m r} \cdot {
m P_{
m f_c}} = rac{1}{2\sqrt{{
m s}}} \, \lambda({
m s}, {
m M_D^2}, {
m M_D^2})$$
 $B_0 = 1, \qquad \qquad {
m strong interaction scale} \sim 1\text{-}3 {
m fm}$ 
 $B_1 = 1 + Z^2, \qquad \qquad B_2 = 9 + 3Z^2 + Z^4, \qquad \qquad B_3 = 225 + 45Z^2 + 6Z^4 + Z^6$ 

 $\Gamma$  increase with the energy  $s = E \text{cm}^2 = W^2$  smoothly. This model is too simple to describe the details of the hadronic decay dynamics!

• Hadronic width in equivalent interaction theory 
$$1^{--} \rightarrow 0^{-C} + 0^{-C} \quad (VPP)$$
Heavy  $\Psi$  family has three decay modes: 
$$1^{--} \rightarrow 1^{-C} + 0^{-C} \quad (VVP)$$
Effective Hamiltonians: 
$$1^{--} \rightarrow 1^{-C} + 1^{-C} \quad (VVV)$$

$$\mathcal{H}_{eff} = g_1 \psi_{\mu} [(\partial_{\mu} D) \bar{D}' - D(\partial_{\mu} \bar{D}')] \quad (VPP)$$

$$\mathcal{H}_{eff} = g_2 \epsilon_{\mu\nu\lambda\sigma} (\partial_{\mu} \psi_{\nu}) (\partial_{\lambda} D_{\sigma}^*) \cdot \bar{D} \quad (VVP)$$

$$\mathcal{H}_{eff} = g_4 \psi_{\mu} [(\partial_{\mu} D_{\nu}^*) \bar{D}_{\nu}^* - D_{\nu}^* (\partial_{\mu} \bar{D}_{\nu}^*)] + g_5 (\partial_{\lambda} \psi_{\mu}) [(\partial_{\mu} D_{\nu}^*) (\partial_{\nu} \bar{D}_{\lambda}^*) - (\partial_{\nu} D_{\lambda}^*) (\partial_{\mu} \bar{D}_{\nu}^*)] \quad (VVV)$$
Energy dependent width: 
$$\Gamma_R^{fc}(s) = G_R^{fc} \cdot \frac{\lambda^3 (s, M_{D^*}^2, M_{D^*}^2)}{s^{5/2}}, \quad (V \rightarrow PP)$$

$$\Gamma_R^{fc}(s) = G_R^{fc} \cdot \frac{\lambda^3 (s, M_{D^*}^2, M_{D^*}^2)}{s^{3/2}}, \quad (V \rightarrow PV)$$

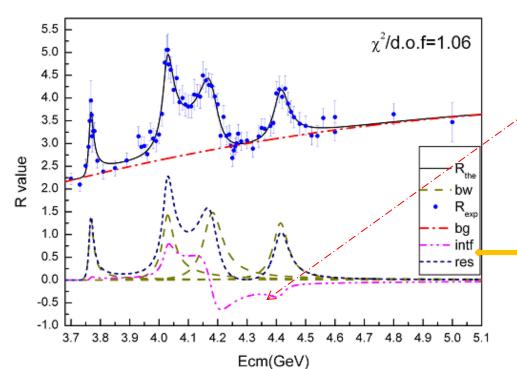
$$\Gamma_R^{fc}(s) = G_R^{fc} \cdot \frac{\lambda^3 (s, M_{D^*}^2, M_{D^*}^2)}{s^{5/2}} \cdot [3 + \frac{\lambda^2 (s, M_{D^*}^2, M_{D^*}^2)}{4M_{D^*}^2, M_{D^*}^2}] \quad (V \rightarrow VV)$$

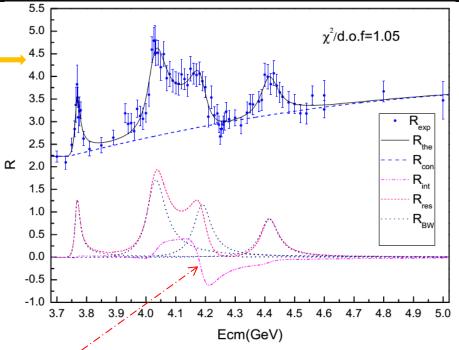
Where,  $G_R$  are the generalized form-factors, and are the scalar functions of the 4-momentum of the initial and final particles. But no enough 12 knowledge about them, they have to be treated as the free parameters.

### Model dependence of line shape



Physics Letters B 660 (2008) 315–319 PDG 2008 – 2014





Different model cause different resonant line shape.

Hadronic width: equivalent interaction theory

#### PDG cites BESII measurements

 $\psi$ (3770)

$$I^{G}(J^{PC}) = 0^{-}(1^{-})$$

 $\psi$ (4160)

$$I^{G}(J^{PC}) = 0^{-}(1^{-})$$

#### $\psi(3770)$ MASS

OUR FIT includes measurements of  $m_{\psi(2S)}$ ,  $m_{\psi(3770)}$ , and  $m_{\psi(3770)} - m_{\psi(2S)}$ . VALUE (MeV) EVTS DOCUMENT ID TECN COMMENT

3772.92  $\pm$  0.35 OUR FIT Error includes scale factor of 1.1.

3775.2  $\pm$  1.7 OUR AVERAGE Error includes scale factor of 1.4. See the ideogram below.

3772.0  $\pm$  1.9 1 ABLIKIM 08D BES2  $e^+e^- \rightarrow \text{hadrons}$ 

 $\psi$ (4040)

$$I^{G}(J^{PC}) = 0^{-}(1^{-})$$

#### $\psi(4040)$ MASS

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT
4039 ± 1 OUR ESTIMATE			
4039.6± 4.3	ABLIKIM 08D	BES2	$e^+e^- \rightarrow \text{hadrons}$

 $\psi$ (4415)

$$I^G(J^{PC}) = 0^-(1^{--})$$

#### $\psi$ (4415) MASS

VALUE (MeV)	DOCUMENT ID		TECN	COMMENT
4421 ± 4 OUR ESTI	MATE			
4415.1± 7.9	<sup>1</sup> ABLIKIM	08D	BES2	$e^+e^- \rightarrow hadrons$
• • • We do not use the	following data for av	verages	, fits, li	mits, etc. • • •
4411 ± 7	<sup>2</sup> PAKHLOVA	08A	BELL	10.6 $e^+e^- \rightarrow D^0D^-\pi^+\gamma$
4425 ± 6	3 SETH	05A	RVUE	$e^+e^- \rightarrow hadrons$
4429 ± 9	<sup>4</sup> SETH	05A	RVUE	$e^+e^- \rightarrow hadrons$
4417 ±10	BRANDELIK	78c	DASP	e+ e-
4414 ± 7	SIEGRIST	76	MRK1	e+ e-
1 Reanalysis of data pre	cented in RAL02c Fr	om a m	lobal fit	over the center-of-mass energy

<sup>&</sup>lt;sup>1</sup> Reanalysis of data presented in BAI 02c. From a global fit over the center-of-mass energy region 3.7–5.0 GeV covering the  $\psi(3770)$ ,  $\psi(4040)$ ,  $\psi(4160)$ , and  $\psi(4415)$  resonances. Phase angle fixed in the fit to  $\delta = (234 \pm 88)^{\circ}$ .

#### $\psi$ (4160) MASS

VALUE			DOCUMENT II				
4153	± 3	OUR ESTIMATE					
4191.	7± 6.5	1	ABLIKIM	08D	BES2	$e^+e^- \rightarrow$	hadrons
• • •	We do	not use the following	data fer averag	es, fits,	iimits,	et <del>c.</del> • • •	
4151	± 4	2	SETH	05A	RVUE	$e^+e^- \rightarrow$	hadrons
4155	± 5	3	SETH	05A	RVUE	$e^+e^- \rightarrow$	hadrons
4159	±20		BRANDELIK	78c	DASP	e+ e-	

- <sup>1</sup> Reanalysis of data presented in BAI 02c. From a global fit over the center-of-mass energy region 3.7–5.0 GeV covering the  $\psi(9770)$ ,  $\psi(4040)$ ,  $\psi(4160)$ , and  $\psi(4415)$  resonances. Phase angle fixed in the fit  $\delta \delta = (293 \pm 57)^{\circ}$ .
- <sup>2</sup> From a fit to Crystal Ball (OSTERHELD 86) data.
- 3 From a fit to BES (BAI 02c) data.

#### $\psi$ (4160) WIDTH

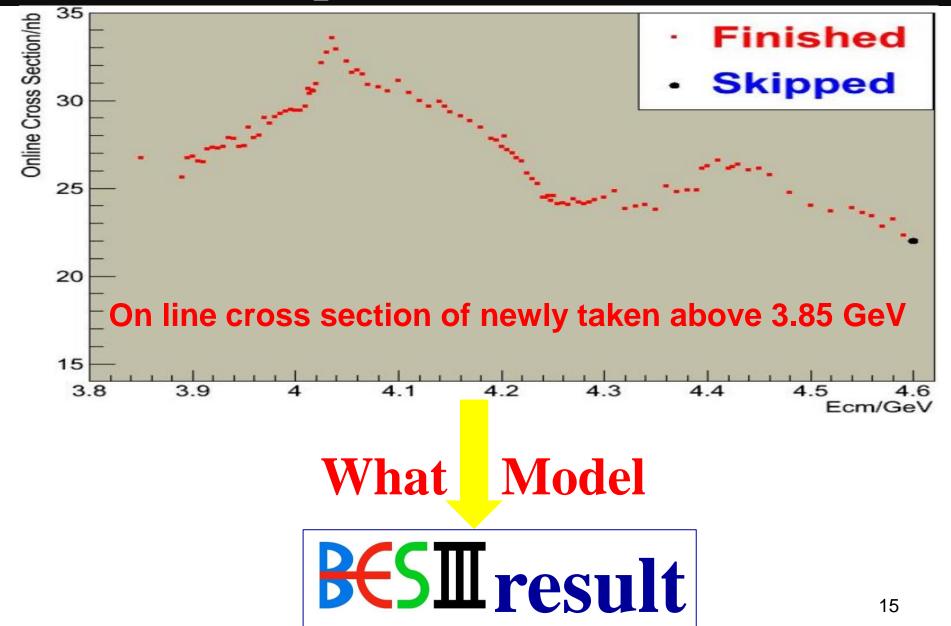
	E (MeV)		DOCUMENT ID		TECN	COMMENT	
103	± 8. Q	IR ESTIMATE -	. — . —				<u> </u>
	8±12.3		<sup>4</sup> ABLIKIM	08D	BES2	$e^+e^- \rightarrow$	hadrons
• •	<ul> <li>We do n</li> </ul>	ot use the following	data for averag	es, rits,	limits,	etc. • • •	
107	±10		<sup>5</sup> SETH	05A	RVUE	$e^+e^- \rightarrow$	hadrons
107	±16		<sup>6</sup> SETH	05A	RVUE	$e^+e^- \rightarrow$	hadrons
78	±20		BRANDELIK	78C	DASP	e+e-	

- <sup>4</sup> Reanalysis of data presented in BAI 02C. From a global fit over the center-of-mass energy region 3.7–5.0 GeV covering the  $\psi(3770)$ ,  $\psi(4040)$ ,  $\psi(4160)$ , and  $\psi(4415)$  resonances. Phase angle fixed in the fit  $(6.6 \pm 0.23) \pm 0.00$
- <sup>5</sup> From a fit to Crystal Ball (OSTERHELD 86) data.
- 6 From a fit to BES (BAI 02c) data.

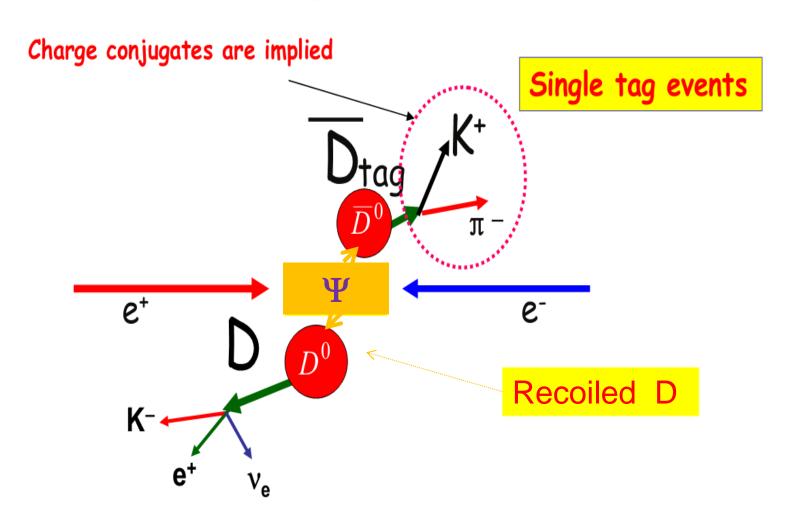
#### $\psi$ (4160) PARTIAL WIDTHS

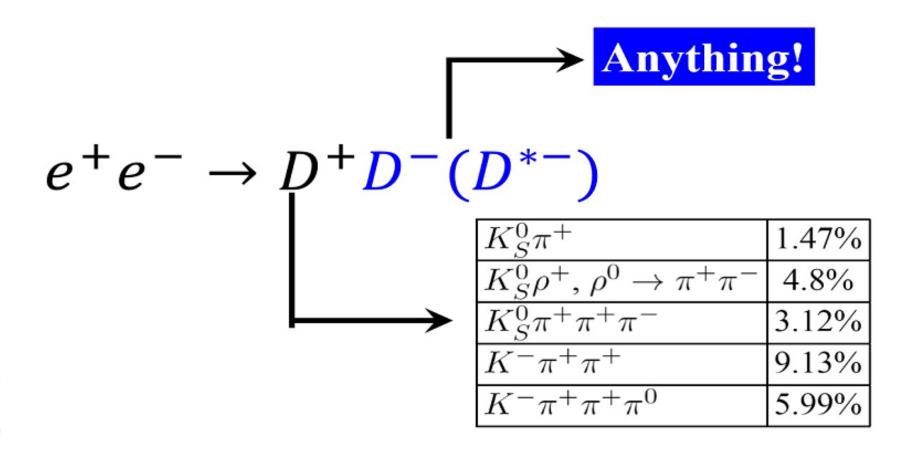
Γ(e <sup>+</sup> e <sup>-</sup> )				Γ
VALUE (keV)	DOCUMENT ID	TECN	COMMENT	
0.83±0.07 OUR ESTIMATE 0.48±0.22	7	BES2	$e^+e^- \rightarrow hadrons$	-

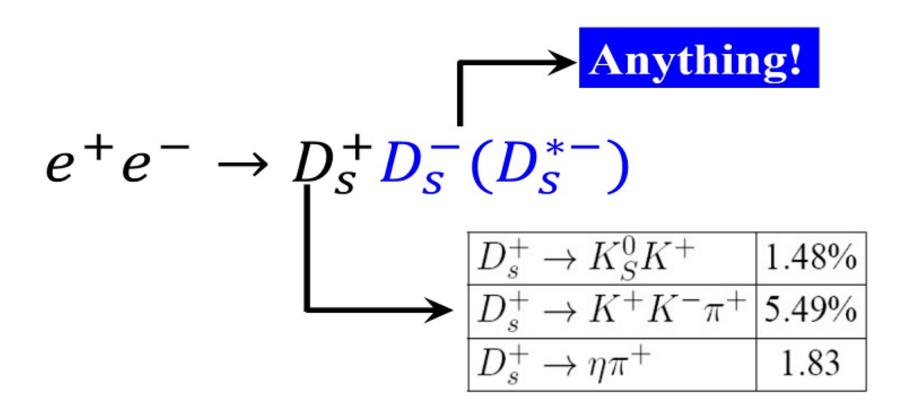
### Line shape measured at BESIII



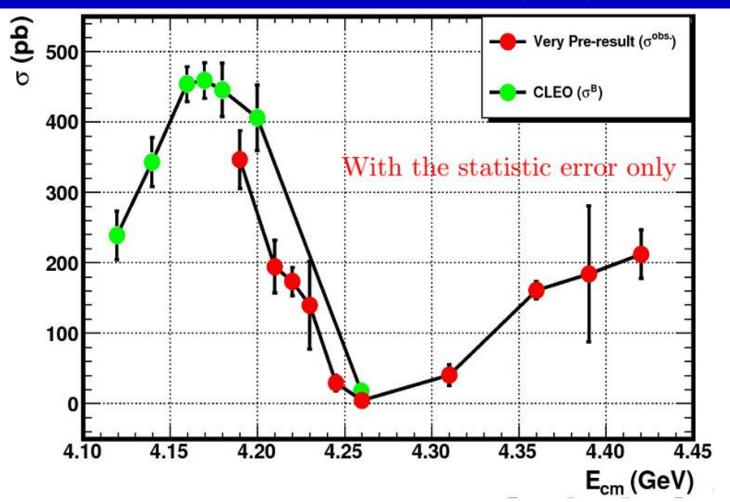
Determination of  $\Gamma$  need measure the branch ratio of DD states. We are going to use D tag package selecting DD events.







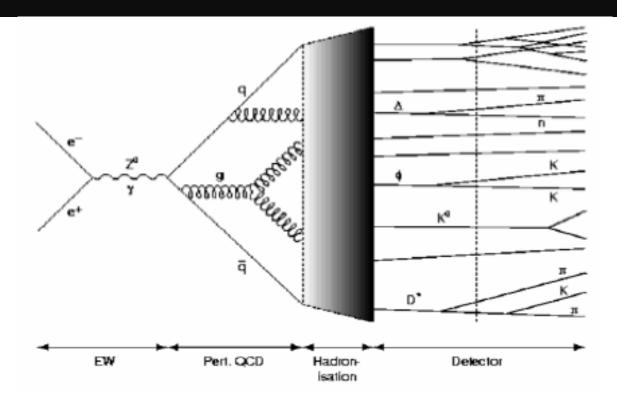
#### Very preliminary results for $D_s^+D_s^{*-}$ mode



At present, our analysis is different from the measurements<sub>19</sub> by CLEO, we have to check our method carefully.

# Main parameters to tune

#### **Pictures of Hadronic Production**



Hadronization belongs to QCD nonperturbative problem, and one has to employ phenomenological models.

#### Two typical models:

- Cluster decay model (HERWIG)
- ➤ String fragmentation model (JETSET/PYTHIA) √

### **Motivation for Developing LUARLW**

#### hep-ph/9910285

#### Few-Body States in Lund String Fragmentation Model

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#### Abstract

The well-known Monte Carlo simulation packet JETSET is not built in order to describe few-body states (in particular at the few GeV level in  $e^+e^-$  annihilation as in BEPC). In this note we will develop the formalism to use the basic Lund Model area law directly for Monte Carlo simulations.

#### **LUARLW** simulation

LUARLW can simulate the productions and decays of inclusive continuous channels and  $J^{PC} = 1^{--}$  resonances with ISR return from 2 - 5 GeV to meet the needs of R scan.

$$e^+e^-\Rightarrow\gamma^*\Rightarrow\begin{cases} V\left(\rho,\omega,\phi\right)\\ q\bar{q}\Rightarrow \text{ string}\Rightarrow \text{ hadrons}\\ gq\bar{q}\Rightarrow \text{ string}\Rightarrow \text{ hadrons} \end{cases}$$
 
$$e^+e^-\Rightarrow\gamma^*\Rightarrow J/\psi\Rightarrow\begin{cases} \gamma^*\Rightarrow e^+e^-,\ \mu^+\mu^-\\ \gamma^*\Rightarrow \bar{q}\Rightarrow \text{ string}\Rightarrow \text{ hadrons}\\ \gamma qg\Rightarrow \text{ hadrons}\\ \gamma qg\Rightarrow \text{ hadrons}\\ \gamma qg\Rightarrow \text{$$

$$e^{+}e^{-} \Rightarrow \gamma^{*} \Rightarrow \begin{cases} \psi(4040) \Rightarrow DD, D^{*}D^{*}, DD^{*}, D^{*}D, D_{s}D_{s}, \text{ other decay modes} \\ \psi(4160) \Rightarrow D\bar{D}, D^{*}\bar{D}^{*}, D\bar{D}^{*}, D^{*}\bar{D}, D_{s}\bar{D}_{s}, D_{s}\bar{D}_{s}^{*}, \text{ other decay modes} \\ \psi(4415) \Rightarrow D\bar{D}, D^{*}\bar{D}^{*}, D\bar{D}^{*}, D^{*}\bar{D}, D_{s}\bar{D}_{s}, D_{s}\bar{D}_{s}^{*}, D_{s}^{*}\bar{D}_{s}^{*}, \text{ other decay modes} \\ X(4260) \Rightarrow \text{ possible decay modes} \end{cases}$$

#### Simulation of radiative events

In MC, hadronic events are classed into two types:

1 Nonradiative events

Born, soft (k < ko) and virtual radiation. The effective c.m.s. energy of the hadronic system equals to the initial  $e^+e^-$  energy

Weight:  $\sigma^{VSB} = \sigma^0(s)[1 + \beta \ln k_0 + \delta_{AR}]$ 

**2** Radiative events

Hard photon with  $k>k_0$ , the effective c.m.s. energy of the hadronic system smaller than the initial  $e^+e^-$  energy, s'=s(1-k)

Weight:  $\sigma^{HB} = \int_{k_0}^{k_m} dk \frac{\partial \sigma^{HB}}{\partial k}$ 

Momentum and polar angle distribution of radiative photon:

$$d\sigma^{HB}(s) = \frac{\alpha}{\pi^2} \frac{\sin^2 \theta}{(1 - a^2 \cos^2 \theta)} \frac{dk d\Omega_{\gamma}}{k} (1 - k + \frac{k^2}{2}) d\sigma^0(s')$$

Note: for the narrow resonance, the effect of energy spread must be considered 24

### Multiplicity distribution

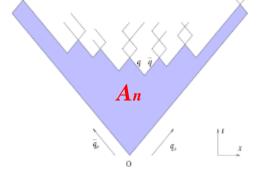
In MC simulation, various events are sampled according to multiplicity distribution of preliminary hadrons  $P_n$ , which can be derived from the Lund area law.

Define partition function for string  $\rightarrow n$  hadrons:

$$Z_n = s \int d\Phi_n \exp(-b\mathcal{A}_n)$$

Multiplicity distribution for string fragmentations:

$$P_n = Z_n / \sum Z_n$$



At some approximations:

$$P_n(s) = \frac{\mu^n}{n!} \exp[c_0 + c_1(n-\mu) + c_2(n-\mu)^2].$$

μ may use QCD exponential form

$$\mu = a + b \exp\{c \cdot [\ln(s/Q_0^2)]^{1/2}\}\$$

where  $c_0$ ,  $c_1$ ,  $c_2$ , a, b, c are parameters to be tuned using data.

### Main parameters to be tuned

• Parameters  $\sigma$  and  $\xi$  in Lund area law

$$\mathcal{M}_{\perp} = \exp(-\sum_{j=1}^{n} \vec{k}_{j}^{2})$$
,  $\vec{k}_{j} \equiv \frac{\vec{p}_{\perp j}}{2\sigma}$   
 $\mathcal{M}_{//} = \exp(i\xi \mathcal{A}_{n})$ ,  $\xi = \frac{1}{2\kappa} + i\frac{b}{2}$ 

Parameters PARJ(\*) in JETSET

```
PARJ(01)=0.15
                !p(qq)/p(q) suppression of diquark-antidiquark production
                !p(ss)/p(dd,uu)
PARJ(02)=0.30
                                            suppression of s quark production
PARJ(03)=0.75
                !P(us)/P(ud))/(P(s)/P(d)) extra suppression of strange diquark
                !P(ud1)/P(ud0) suppression of spin 1 diquark to spin 0 ones
PARJ(04)=0.10
PARJ(05)=0.50
                !P(BMB)/[P(BB)+P(BMB)] relative probability for BMB to BB
PARJ(06)=0.50
                !extra suppression for having a ssbar
PARJ(07)=0.50
                 !extra suppression for having a strange meson
PARJ(11)=0.55
                !ratio of light meson with spin 1/0
                 !ratio of strange meson with spin 1/0
PARJ(12)=0.55
PARJ(13)=0.75
                 !ratio of charm meson with spin 1/0
PARJ(14)=0.05
                 !ratio of S=0 meson with L=1, J=1
PARJ(15)=0.05
                 !ratio of S=1 meson with L=1, J=0
PARJ(16)=0.05
                 !ratio of S=1 meson with L=1, J=1
                                                                       26
PARJ(17)=0.10
                 !ratio of S=1 meson with L=1, J=2
```

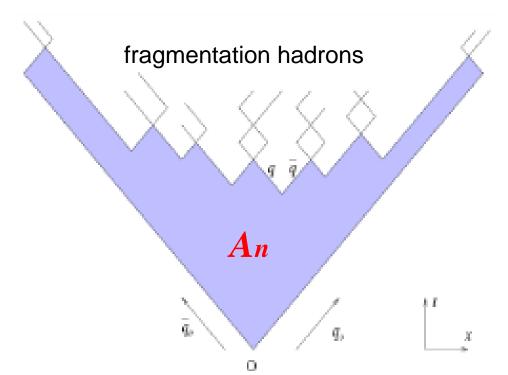
### Functions of the main parameters

$$P_n(s) = \frac{\mu^n}{n!} \exp[c_0 + c_1(n-\mu) + c_2(n-\mu)^2]. \longrightarrow \text{Multiplicity distribution}$$

Parameters PARJ(\*\*)

$$\mathcal{M}_{\perp} = \exp(-\sum_{j=1}^{n} \vec{k}_{j}^{2})$$
,  $\vec{k}_{j} \equiv \frac{\vec{p}_{\perp j}}{2\sigma}$   
 $\mathcal{M}_{//} = \exp(i\xi \mathcal{A}_{n})$ ,  $\xi = \frac{1}{2\kappa} + i\frac{b}{2}$ 

Momentum distribution

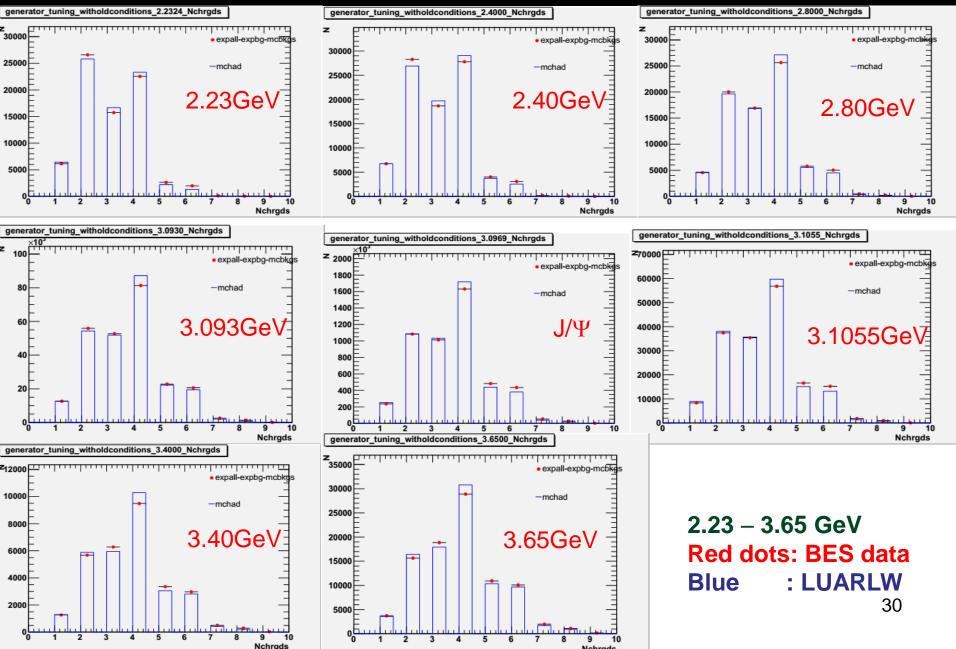


### Goal of tuning

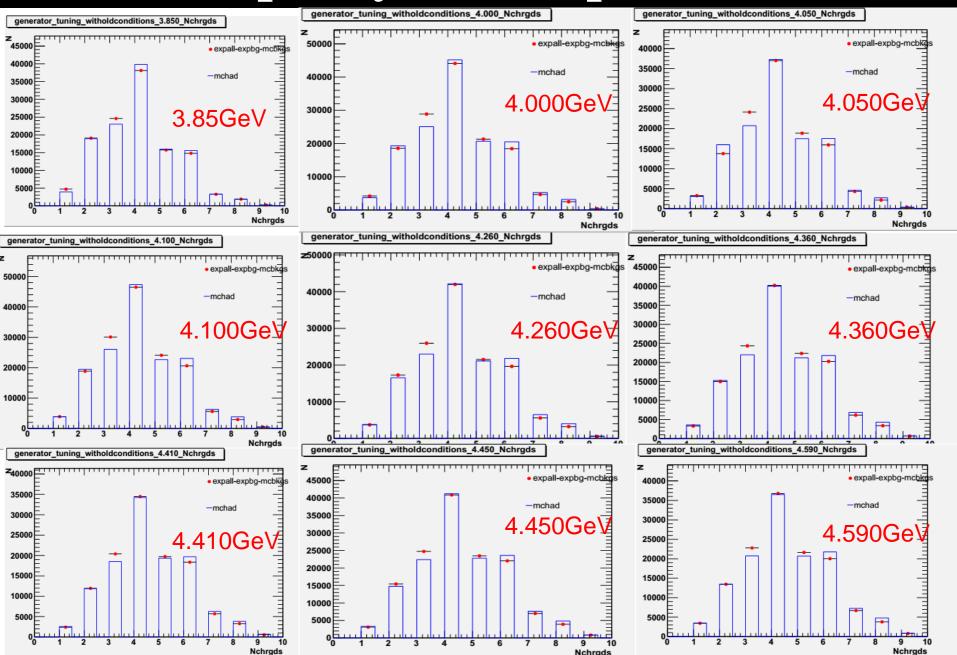
- Use **two sets** of parameters to control the simulations of continuous and  $J^{PC} = 1^{--}$  resonant states.
- LUARLW could agree with data well for most of the inclusive distributions and DD final states at all energies.

# Preliminary tuning

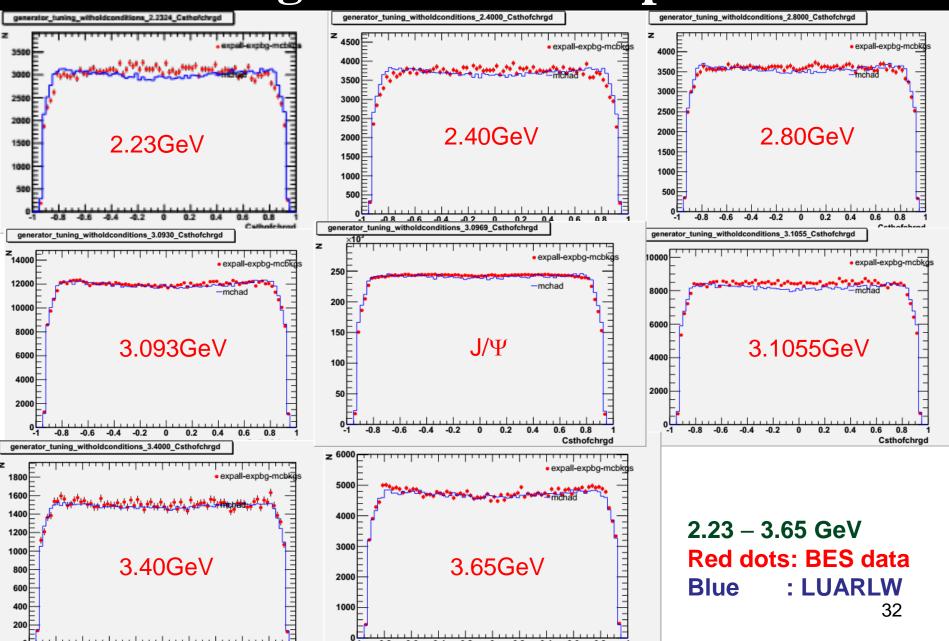
### Multiplicity below open charm



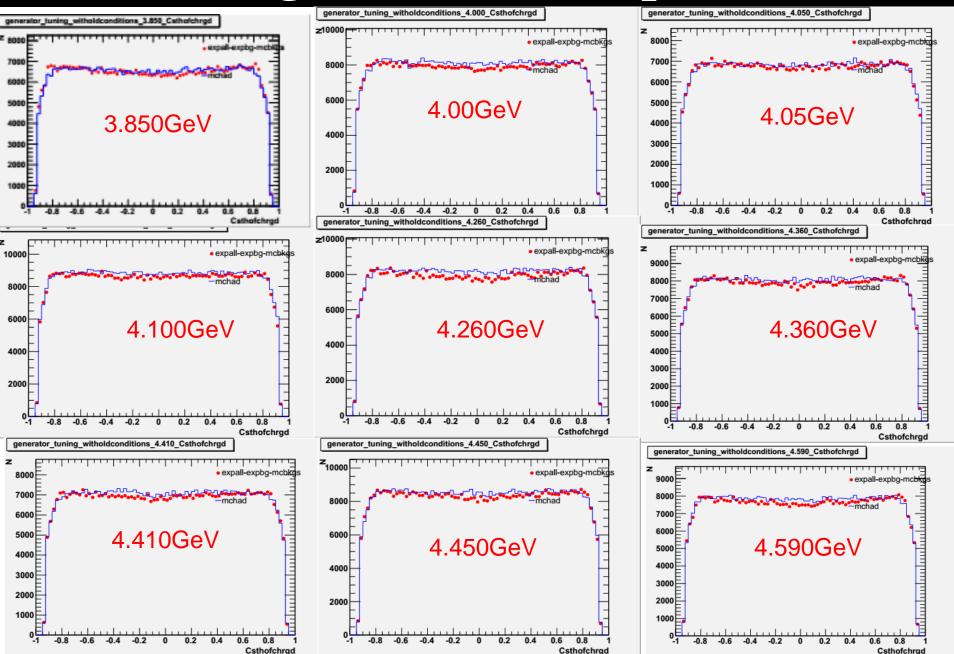
Multiplicity above open charm



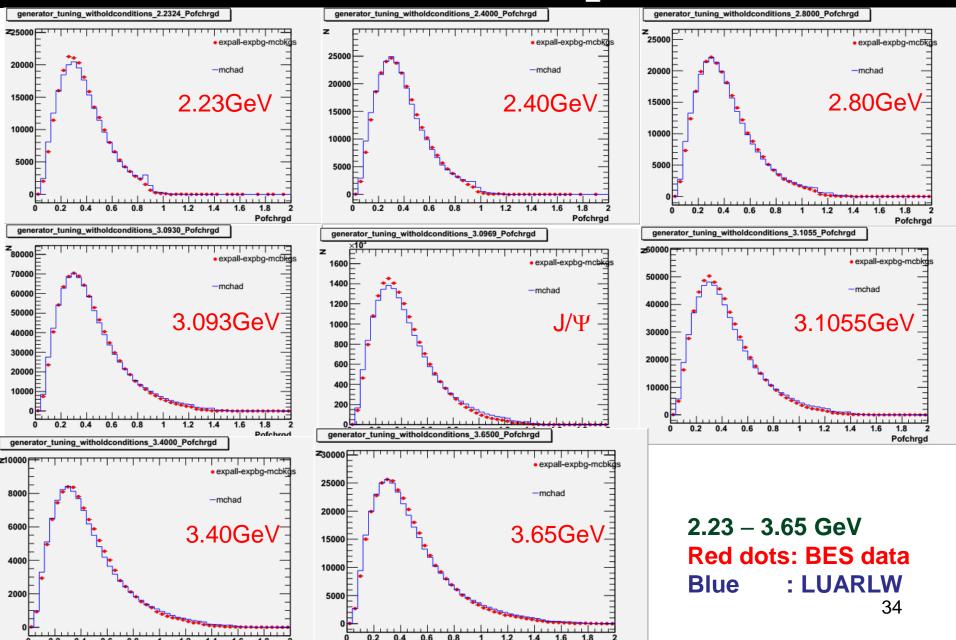
### Polar Angle cosθ below open charm



### Polar Angle cosθ above open charm

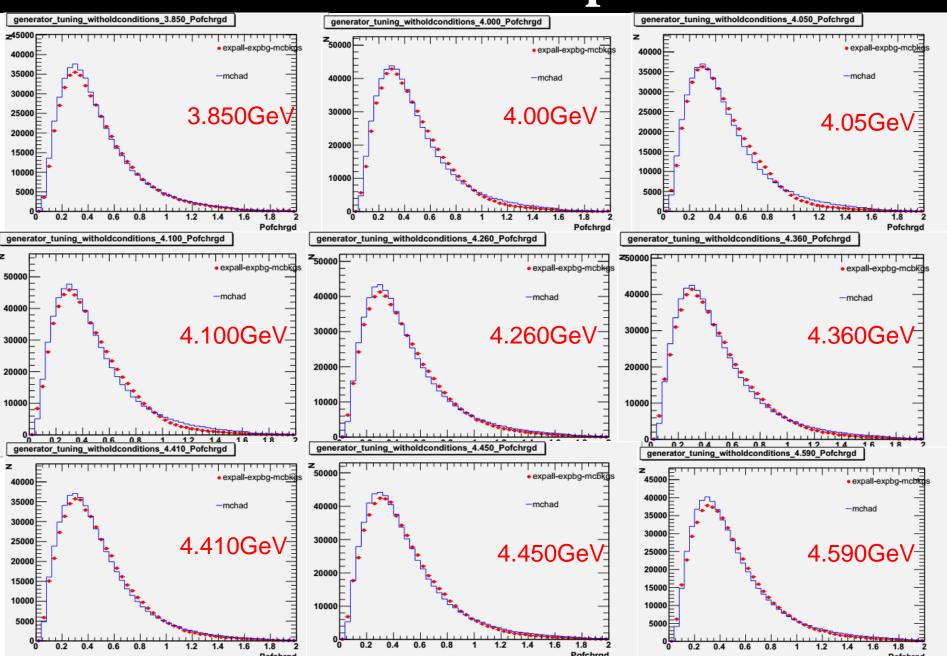


### Momentum below open charm



Pofchrgd

Momentum above open charm



#### Summary

The parameters in LUARLW are tuning based on the data samples taken at BESIII from 2.23 - 4.59 GeV.

#### Present status:

- \* LUARLW is tuned using two sets of parameters corresponding to the below and above open charm thresholds data separately, but now the parameters of LUARLW are far from the well tuned.
- \* Criteria for inclusive hadronic events need to fix.
- \* D tag package using for selecting DD final states has to study further.

#### Next to do:

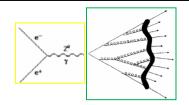
- \* Continue tuning the LUARLW parameters.
- Study the theoretical calculations for interference effects.
- \* Study the width energy dependence of  $\Psi(4040), \Psi(4160)$  and  $\Psi(4415)$ .
- Calculate ISR factor employing above improvements.
- \* Check if the related QED generator work well.

# Back Up

# **Key Points of Lund Area Law**

#### Hadron production process:

$$e^+e^- \Rightarrow q\bar{q} \Rightarrow \text{string} \Rightarrow m_1 + m_2 + \dots + m_n$$



#### "Matrix element":

$$\mathcal{M} \equiv \mathcal{M}_{\text{QED}}(e^+e^- \to q\bar{q})\mathcal{M}_{\text{LUND}}(q\bar{q} \to m_1, m_2, \cdots m_n)$$

Lund model: 
$$\mathcal{M}_{\text{LUND}}(q\bar{q} \to m_1, m_2, \cdots m_n) = C_n \mathcal{M}_{\perp} \mathcal{M}_{//}$$

Transverse momentum (Gaussian): 
$$\mathcal{M}_{\perp} = \exp(-\sum_{j=1}^{n} \vec{k}_{j}^{2})$$
,  $\vec{k}_{j} \equiv \frac{\vec{p}_{\perp j}}{2\sigma}$ 

Longitudinal momentum (area law): 
$$\mathcal{M}_{//} = \exp(i\xi \mathcal{A}_n)$$
,  $\xi = \frac{1}{2\kappa} + i\frac{b}{2}$ 

#### Probability for string fragmenting into *n* hadrons :

$$d\wp_n(q\bar{q}\to m_1,m_2,\cdots m_n) = (2\pi)^4\delta(1-\sum_{j=1}^n\frac{m_{\perp j}^2}{sz_j})\delta(1-\sum_{j=1}^nz_j)\delta(\sum_{j=1}^n\vec{k}_j)\sum |\mathcal{M}_{LUND}|^2d\Phi_n$$

Fraction of light-cone momentum: 
$$z_j \equiv (E_j \pm p_{zj})/(E_0 \pm P_{z0})$$

### **Solutions of Lund Area Law**

#### String → 2 hadrons

$$\wp_{2} = \frac{C_{2}}{\sqrt{\lambda(s, m_{\perp 1}^{2}, m_{\perp 2}^{2})}} [exp(-b\mathcal{A}_{2}^{(1)}) + exp(-b\mathcal{A}_{2}^{(2)})]$$

$$\mathcal{A}_{2}^{(1,2)} = \frac{1}{2}(s + m_{\perp 1}^{2} + m_{\perp 2}^{2} \mp \sqrt{\lambda}) \qquad \lambda(a, b, c) \equiv a^{2} + b^{2} + c^{2} - 2ab - 2bc - 2ca$$

#### String → 3 hadrons

$$d\wp_3 = \frac{C_3}{\sqrt{\Lambda}} exp(-b\mathcal{A}_3) d\mathcal{A}_3$$

$$\Lambda = [(s - \mathcal{A}_3)(\mathcal{A}_3 - m_{\perp 1}^2 - m_{\perp 2}^2 - m_{\perp 3}^2) - m_{\perp 1}^2 m_{\perp 2}^2 - m_{\perp 2}^2 m_{\perp 3}^2 - m_{\perp 3}^2 m_{\perp 1}^2]^2 - 4s m_{\perp 1}^2 m_{\perp 2}^2 m_{\perp 3}^2$$

#### • String $\rightarrow n$ hadrons $(n \ge 4)$

$$d\wp_{n}(s;s_{1},s_{2}) = \frac{ds_{1}ds_{2}}{\sqrt{\lambda(s,s_{1},s_{2})}} [exp(-b\Gamma^{(1)}) + exp(-b\Gamma^{(2)})]\wp_{n_{1}}(s_{1})\wp_{n_{2}}(s_{2})$$

$$s_{1} = (\sum_{j=1}^{n_{1}} p_{\circ j})^{2} = Z_{1}(1 - Z_{2})s, \quad s_{2} = (\sum_{j=n_{1}+1}^{n} p_{\circ j})^{2} = Z_{2}(1 - Z_{1})s$$

$$A_{n} = A_{n_{1}} + A_{n_{2}} + \Gamma, \quad \Gamma = (1 - Z_{1})(1 - Z_{2})s$$

$$\Gamma^{(1,2)} = \frac{1}{4s} \left[ s - s_{1} + s_{2} \pm \sqrt{\lambda(s,s_{1},s_{2})} \right] \left[ s + s_{1} - s_{2} \mp \sqrt{\lambda(s,s_{1},s_{2})} \right]$$

$$s_{1} = (\sum_{j=1}^{n} p_{\circ j})^{2} = Z_{1}(1 - Z_{2})s$$

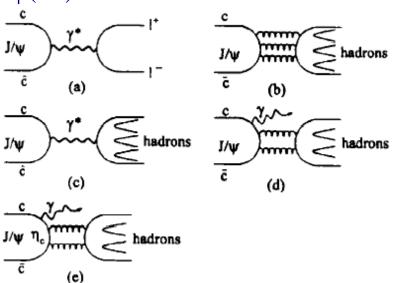
$$s_{2} = (\sum_{j=1}^{n} p_{\circ j})^{2} = Z_{2}(1 - Z_{1})s$$

$$s_{3} = (\sum_{j=1}^{n} p_{\circ j})^{2} = Z_{1}(1 - Z_{2})s$$

$$s_{4} = \sum_{j=1}^{n} p_{\circ j} + \sum_{j$$

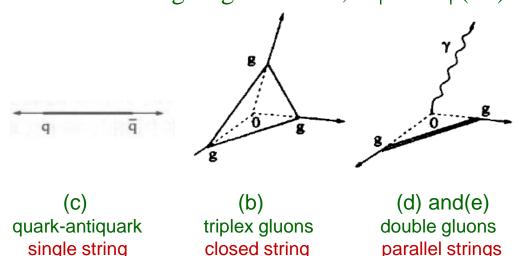
### **Charmonium Decay Modes**

 $J/\psi$  and  $\psi(2S)$  are charmonia with  $J^{PC}=1^{--}$ . Decay modes predicts by QCD



- (a) Electric-magnetic lepton decay
- (b) Strong interaction decay
- (c) Electric-magnetic hadronic decay
- (d) Radiative hadronic decay
- (e) Radiative M1 transision

In the Lund string fragmentation,  $J/\psi$  and  $\psi(2S)$  contain three types of string structures



Due to Lund area law has the invariant relativistic property, it can be used to treat substring fragmentation independently.

#### **Continum states**

#### Continuum state at $E_{cm} = 1 \text{ GeV}$

I particle		KS	KF	Of	^i9	P_X	P_9	p_z	Е	M
1 !e-!		21	11	1	0	0,000	0,000	0,500	0,500	0,001
2 !e+!		21	-11	1	0	0.000	0.000	-0.500	0.500	0.001
3 gamma		1	22	2	2	0.000	0.000	0.174	0.174	0.000
4 [gamma!		21	22	2	1	0.000	0.000	-0.174	0.826	0.808
5 (ū)	Ĥ	12	2	2	4	0.000	0.000	-0.087	0.413	0.006
6 (u~)	٧	11	-2	2	4	0.000	0.000	-0.087	0.413	0.006
7 (string)		11	92	2	5	0.000	0.000	-0.174	0.826	0.808
8 (rho0) ¯		11	113	3	- 7	0.000	0.000	-0.174	0.826	0.808
9 pi+		1	211	1	8	-0.355	0.080	0.023	0.390	0.140
10 pi-		1	-211	1	8	0.355	-0.080	-0,197	0.436	0.140

#### Continuum state at $E_{cm} = 2 \text{ GeV}$

I particle		KS	KF	ori	ig	P_X	P_9	P_Z	Ε	M
1 (e-)		12	11		0	0.000	0.000	1,000	1,000	0.001
2 (e+)		11	-11		0	0.000	0.000	-1.000	1.000	0.001
3 !gamma!		21	22		1	0.000	0.000	0.000	2,000	2,000
4 (š)	Ĥ	12	3		3	0.071	-0,006	-0.977	1,000	0.199
5 (s~)	٧	11	-3		3	-0.071	0.006	0.977	1.000	0.199
6 (string)		11	92		4	0.000	0.000	0.000	2,000	2,000
7 (K0)		11	311		6	-0,267	0.022	-0,650	0.861	0.498
8 (K*~0)		11	-313		6	0.267	-0.022	0.650	1,139	0.897
9 (K_SO)		11	310		7	-0.267	0.022	-0.650	0.861	0.498
10 K-		1	-321		8	0.237	-0.235	0.213	0.633	0.494
11 pi+		1	211		8	0.030	0.213	0.436	0.506	0.140
12 (pi0)		11	111		9	-0,207	-0,186	-0.387	0.495	0.135
13 (pi0)		11	111		9	-0.060	0.207	-0,263	0.366	0.135
14 gamma		1	22	1	12	-0,190	-0.181	-0,268	0.375	0.000
15 gamma		1	22	1	12	-0.017	-0.004	-0,118	0.120	0.000
16 gamma		1	22	1	L3	-0.042	0.023	-0.129	0.138	0.000
17 gamma		1	22	1	L3	-0.018	0.184	-0.134	0,228	0.000

#### Continuum state at $E_{cm} = 3 \text{ GeV}$

					- 0116			
I particle	KS	KF	orig	P_X	P_9	p_z	Ε	M
1 (e-)	12	11	0	0.000	0.000	1,500	1,500	0.001
2 (e+)	11	-11	0	0.000	0.000	-1,500	1,500	0.001
3 gamma	1	22	1	0.000	0.000	0.034	0.034	0.000
4 !gamma!	21	22	1	0.000	0.000	-0.034	2,966	2,966
5 (ŭ)	A 12	2	4	-1,088	0.832	0.553	1.477	0.006
6 (u~)	V 11	-2	4	1.088	-0.832	-0.587	1,490	0.006
7 (string)	11	92	5	0.000	0.000	-0.034	2,966	2,966
8 (eta)	11	221	7	-0.229	0.066	0.084	0.603	0.548
9 (K*-)	11	-323	7	0.148	-0.031	0.507	0.935	0.771
10 (KO)	11	311	7	-0.185	0.182	-0.832	1,004	0.498
11 pi+	1	211	7	0.266	-0.216	0,208	0.424	0,140
12 (K~0)	11	-311	9	0.129	-0.082	0.570	0.772	0.498
13 pi-	1	-211	9	0.019	0.051	-0.063	0.163	0.140
14 K_L0	1	130	10	-0.185	0.182	-0.832	1,004	0.498
15 K_L0	1	130	12	0,129	-0.082	0.570	0.772	0.498
16 (pi0)	11	111	8	0.045	0.090	0.053	0.177	0.135
17 (pi0)	11	111	8	-0.064	-0.049	0.071	0.172	0.135
18 (pi0)	11	111	8	-0,209	0.024	-0.040	0.253	0.135
19 gamma	1	22	16	0.091	0.077	0.026	0,122	0.000
20 gamma	1	22	16	-0.046	0.013	0.027	0.055	0.000
21 gamma	1	22	17	-0,002	0.034	0.060	0.069	0.000
22 gamma	1	22	17	-0.062	-0.082	0.011	0.104	0.000
23 gamma	1	22	18	-0.020	-0.020	0.035	0.045	0.000
24 gamma	1	22	18	-0,189	0.044	-0, <u>07</u> 5	0.208	0.000
								— ———.

#### J/Ψ production and decay

$\mathrm{J}/\psi[$		J/ $\psi$
I particle KS	KF orig p_x p_y p_z E m	I particle KS KF orig p_x p_y p_z E m
1 !e-! 21 2 !e+! 21 3 gamma 1 4 !gamma! 21 5 (c) A 12 6 (c") V 11 7 !J/psi! 21 8 !gamma! 21 9 mu+ 1 10 mu- 1	11 0 0.000 0.000 1.550 1.550 0.00 -11 0 0.000 0.000 -1.550 1.550 0.00 22 1 0.000 0.000 -0.003 0.003 0.00 22 1 0.000 0.000 0.003 3.097 3.09 4 4 0.000 0.000 0.001 1.549 1.39 -4 4 0.000 0.000 0.001 1.549 1.39 443 6 0.000 0.000 0.001 1.549 1.39 22 7 0.000 0.000 0.003 3.097 3.09 22 7 0.000 0.000 0.003 3.097 3.09 -13 8 -0.231 -1.489 -0.339 1.548 0.10 13 8 0.231 1.489 0.342 1.549 0.10	1 2 !e+! 21 -11 0 0.000 0.000 -1.550 1.550 0.001 0 3 !gamma! 21 22 1 0.000 0.000 0.000 3.100 3.100 7 4 (c) A 12 4 3 0.000 0.000 0.000 1.550 1.350 0 5 (c") V 11 -4 3 0.000 0.000 0.000 1.550 1.350 0 6 !J/psi! 21 443 5 0.000 0.000 0.000 3.100 3.097 7 7 (g) A 11 21 6 -0.539 -0.223 0.264 0.641 0.000 7 8 (g) V 11 21 6 1.332 0.113 -0.737 1.528 0.000 6 9 (g) A 11 21 6 -0.793 0.110 0.473 0.931 0.000 6 10 (d) I 12 1 7 -0.082 -0.024 0.030 0.091 0.010 11 (d") I 12 -1 7 -0.457 -0.199 0.234 0.550 0.010
$\mathrm{J}/\psi$	c ys Eladrons	12 (u) I 12 2 8 0.044 0.101 -0.076 0.134 0.006 13 (u") I 12 -2 8 1.289 0.012 -0.661 1.394 0.006 14 (d) I 12 1 9 -0.235 0.213 0.189 0.369 0.010 15 (d") V 11 -1 9 -0.558 -0.103 0.284 0.562 0.010 16 (string) 12 92 10 1.206 -0.012 -0.632 1.485 0.593 17 (string) 12 92 12 -0.514 -0.002 0.208 0.695 0.419
I particle KS  1 !e-! 21 2 !e+! 21 3 gamma 1 4 !gamma! 21 5 (c) A 12 6 (c") V 11 7 !J/psi! 21 8 !gamma! 21 9 (u) A 12 10 (u") V 11 11 (string) 11 12 pi- 1 13 (eta') 11	KF orig p_x p_y p_z E m  11 0 0.000 0.000 1.550 1.550 0.0  -11 0 0.000 0.000 -1.550 1.550 0.0  22 1 0.000 0.000 -0.003 0.003 0.0  22 1 0.000 0.000 0.003 3.097 3.0  4 4 0.000 0.000 0.001 1.549 1.3  -4 4 0.000 0.000 0.001 1.549 1.3  443 6 0.000 0.000 0.001 1.549 1.3  443 6 0.000 0.000 0.001 3.097 3.0  22 7 0.000 0.000 0.003 3.097 3.0  22 8 1.523 -0.039 -0.276 1.548 0.0  -2 8 -1.523 0.039 0.278 1.549 0.0  92 9 0.000 0.000 0.003 3.097 3.0  -211 11 -0.160 0.078 -0.035 0.229 0.1  331 11 1.327 0.134 -0.182 1.652 0.5	18 (string) 11 92 14 -0.692 0.014 0.423 0.919 0.432 01 19 pi- 1 -211 16 0.455 0.033 -0.054 0.480 0.140 01 20 pi+ 1 211 16 0.172 0.010 -0.188 0.291 0.140 00 21 pi- 1 -211 16 0.579 -0.055 -0.390 0.714 0.140 07 22 pi+ 1 211 17 -0.332 -0.131 0.204 0.434 0.140 050 23 (pi0) 11 111 17 -0.182 0.130 0.005 0.261 0.135 0.24 pi- 1 -211 18 -0.342 -0.147 0.142 0.422 0.140 07 25 pi+ 1 211 18 -0.350 0.161 0.202 0.497 0.140 0.66 06 06 06 07 40
14 pi+ 1 15 pi+ 1 16 pi- 1 17 (eta) 11 18 gamma 1	211 11 -1.168 -0.212 0.224 1.216 0.1 211 13 0.127 0.006 -0.178 0.260 0.1 -211 13 0.200 -0.002 -0.023 0.245 0.1 221 13 1.000 0.129 0.019 1.148 0.5 22 17 0.033 -0.067 0.135 0.154 0.0	40 40 40 48 42

0.000

19 gamma

0.967

		A SECONDARY		<b>J/</b> Ψ	proc	lucti	on and d	eca	y Yast	<b>M</b> 1	radiativ	/e trans	ition	
	$J/\psi_{\overline{c}}^{c}$	,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,		drons				$\mathrm{J}/\psi$		*******	$\equiv$	adrons		
I particle	KS	KF orig	P_X P.	y p_z	Ε	M	I particle	KS	KF ori	9 P_X	P_9	P_Z	Ε	M
1 !e-! 2 !e+! 3 gamma 4 !gamma! 5 (c) 6 (c") 7 !J/psi! 8 gamma 9 (g) 10 (g) 11 (d) 12 (d") 13 (s) 14 (s") 15 (string) 16 (string) 17 pi- 18 K+ 19 gamma 20 (K"0)	21 21 1 21 A 12 V 11 21 A 11 V 11 A 12 I 12 I 12 V 11 11 11	11 0 -11 0 22 1 22 1 4 4 -4 4 443 6 22 7 21 19 -1 9 -1 9 -1 19 -2 11 -2 11 -2 11 -2 15 -3 10 -3	0.000 0.000	00 -1.550 00 -0.003 00 0.002 00 0.002 00 0.003 78 -0.518 17 -0.570 61 1.091 79 -0.346 62 -0.224 10 0.349 71 0.743 50 0.397 72 0.102 37 0.500 71 -0.004	0.355 0.481 0.890 1.542 0.836 0.747 0.796 0.171	0.001 0.001 0.000 3.097 1.350 1.350 3.097 0.000 0.000 0.010 0.100 0.199 0.199 1.139 0.614 0.140 0.494 0.000 0.498,	1 !e-! 2 !e+! 3 gamma 4 !gamma! 5 (c) 6 (c") 7 !J/psi! 8 gamma 9 (eta_c) 10 (g) 11 (g) 12 (u) 13 (u") 14 (d) 15 (d") 16 (string) 17 (string) 18 pi+ 19 (K0) 20 (K"0) 21 (pi0) 22 pi- 23 (pi0)	21 21 21 21 A 12 V 11 21 11 A 12 V 11 A 12 I 12 I 12 V 11 11 11 11 11	-2 1 1 1 -1 1 92 1 92 1 211 1	1 -0.169 2 -0.023 3 0.001 6 0.129 6 0.083 6 -0.235 7 0.029 7 -0.158	0.000 0.000 0.000 0.000 0.000 0.000 -0.105 0.105 0.490 -0.385 -0.094 0.584 -0.066 -0.319 -0.413 0.518 -0.147 -0.170 -0.096 0.317	1,549 -1,551 0,003 -0,006 -0,003 -0,006 -0,004 0,036 1,435 -1,399 0,478 0,957 -0,235 -1,164 -0,686 0,722 -0,565 0,021 -0,142 0,476 0,352 -0,105	1.549 1.552 0.003 3.097 1.549 1.549 3.097 0.115 2.982 1.523 1.459 0.509 1.014 0.245 1.214 1.723 1.258 0.614 0.533 0.576 0.588 0.415 0.255	0.001 0.001 0.000 3.097 1.350 1.350 3.097 0.000 0.000 0.000 0.006 0.006 0.010 0.010 1.526 0.891 0.140 0.498 0.498 0.498 0.135 0.140 0.135

# $\Psi \mbox{(2S)}$ production and decay $\psi(2S) \to \pi^+\pi^- J/\psi$

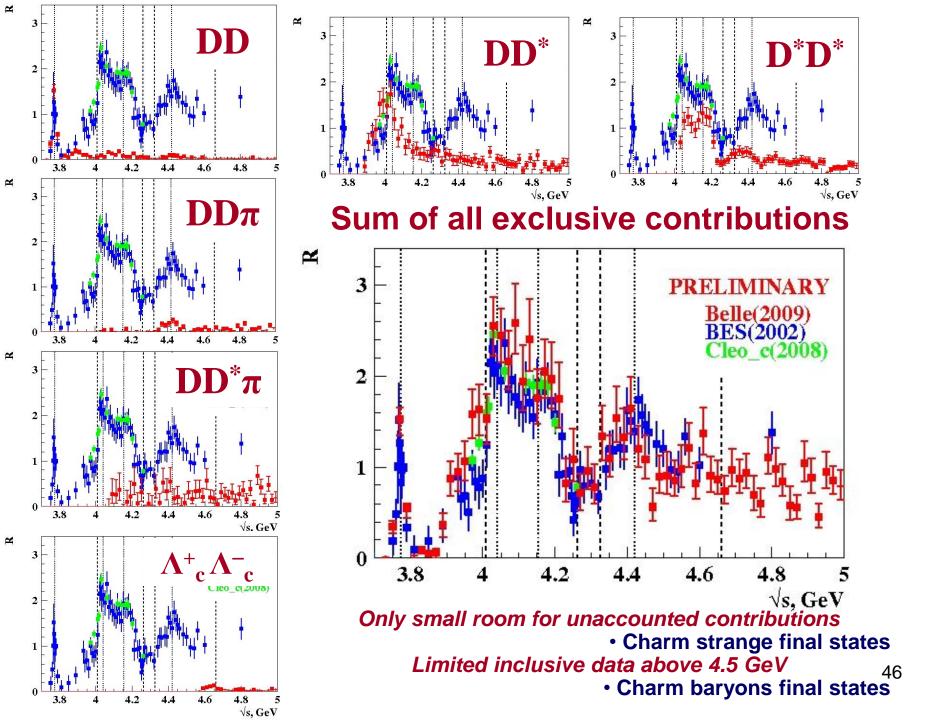
$$\psi(2S) \to \pi^+\pi^- J/\psi$$

$$\psi(2S) \to \eta J/\psi$$

	, ,	/	, ,		$\psi(2S)  o \eta J/\psi$
I particle KS	KF ori	9 P_X P_9	p_z E	M	I particle KS KF orig p_x p_y p_z E m
1 !e-! 21 2 !e+! 21 3 gamma 1 4 !gamma! 21 5 (c) A 12 6 (c") V 11 7 !psi3686! 21 8 pi+ 1 9 pi- 1 10 (J/psi) 11 11 (g) A 11 12 (g) V 11 13 (g) A 11 14 (s) I 12 15 (s") I 12 15 (s") I 12 16 (u) I 12 17 (u") I 12 18 (u) I 12 17 (u") I 12 18 (u) I 12 17 (u") V 11 20 (string) I 12 21 (string) I 12 22 (string) I 12 23 (K"0) I 1 24 pi- 1 25 (pi0) I 1 26 (pi0) I 1 27 gamma 1 28 K+ 1	11	0 1.400 0.164 0 -0.593 0.691 1 -0.230 -0.436 1 -0.185 -0.172 2 0.718 -0.424 2 0.682 0.588 3 -0.102 0.795 3 -0.491 -0.104 4 0.452 0.152 6 0.226 -0.528 8 -0.287 0.623 0 -0.072 0.172 0 0.525 -0.020 1 0.424 -0.401 1 -0.020 0.069 2 -0.279 0.185	1.844 1.844 -1.844 1.844 -0.003 0.003 0.003 3.686 0.002 1.843 0.002 1.843 0.003 3.686 -0.074 0.279 -0.010 0.275 0.087 3.133 -0.243 0.775 0.074 1.412 0.256 0.946 -0.136 0.548 -0.108 0.226 0.044 0.835 0.031 0.577 0.225 0.832 0.031 0.114 -0.105 1.125 0.075 0.948 0.118 1.059 -0.211 0.572 0.106 0.554 -0.152 0.618 0.283 0.321 0.206 0.393 -0.088 0.666	0.001 0.001 0.000 3.686 1.350 3.686 0.140 0.140 0.140 0.000 0.000 0.000 0.000 0.199 0.199 0.006 0.006 0.006 0.006 1.014 0.751 0.798 0.498 0.140 0.135 0.135	1  e-  21

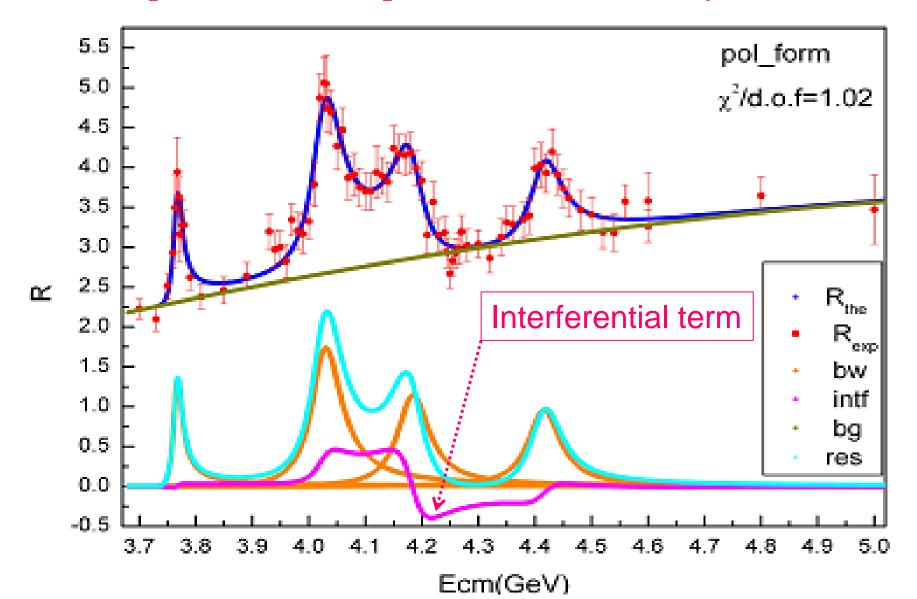
## Main problems in generator tuning

- Some of free parameters in LUARLW need to tune via comparing the related experimental data.
- Before parameter tuning, one has to have a set of event selection criteria and obtain a group of hadronic spectra.
- For data below open charm threshold, hadronic events selection conditions are almost fixed. But for the data newly taken above 3.85 GeV, we are studying the criteria for DD states, there are many problems to discuss.



## Interferential in resonant line shape

For example, the line shape of heavier Ψ family



### D tag

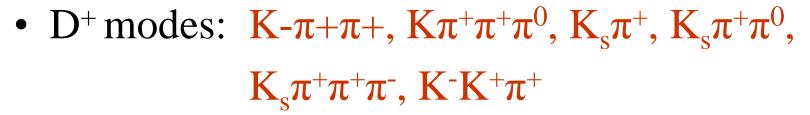
- Hadron D decay reconstruction to find all possible D candidates (tags)
- Check the D tag information if no tag is found, the event is discarded
- Only the events with tags are kept, and compare data with MC
- Tag methods: single D tag and Double D tag

### D tag

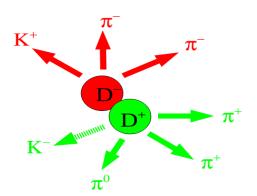
### D-Tag Candidates (1)

- $\Delta E = E_D E_{beam}$
- $m_{BC} = \sqrt{E_{beam}^2 P_D^2}$





- $D_s^+$  modes:  $K_sK^+, K^-K^+\pi^+, K_sK^+\pi^0$
- Charge-conjugate modes implied



### D tag

### D-Tag Candidates (2)

- Ingredients:
  - charged  $K/\pi$  pass vertex and PID cuts
  - Ks from  $\pi^+\pi^+$ ,  $\pi^0$  from  $\gamma\gamma$ , invariant mass cuts
- Cut on  $\Delta E$  and fit  $m_{BC}$  distributions to extract yields
- Analyze MC samples with same procedure
- Compare Data's distribution with MC's

### Summary for D tag prestudy

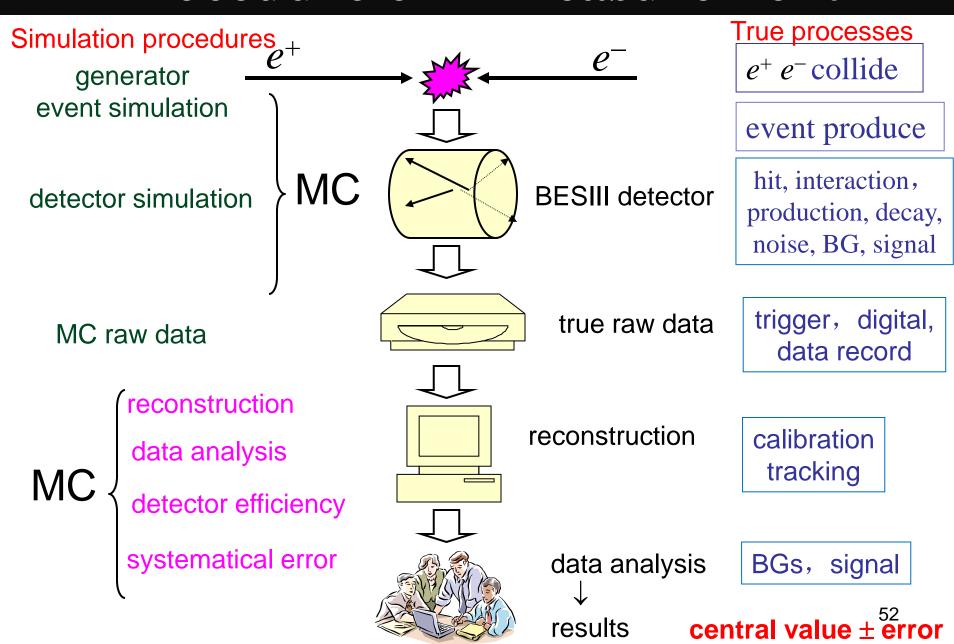
#### Summary

- Observe obviously signals for  $e^+e^- \to D\bar{D}, D\bar{D}^*$  with D tag at energy Y(4260), Y(4230), Y(4360) and scan data at Y(4260), Y(4360).
- A very preliminary observed cross section for  $D_s^+D_s^{*-}$  mode is presented with the statistic error only (without implying the charged-conjugate mode).

#### 2 Next to do

- Carefully check for every tag channel
- Detailed background study
- Improve fitting method (the current fit is very preliminary).
- Understand ISR effect
- Systematic error study

# Procedure of R measurement



### **Parameters Tuning Flow Chart**

Compare true data with MC simulated distributions

