... a personal review on

Leptogenesis

Joshelle Wesins (CERN)

TH-Institute on Flavor, CERN, 08/05/08

Based on some works with:

V. Cirigliano, A. De Simone, G. Isidori, F.R. Joaquim, A. Riotto, C.A. Savoy

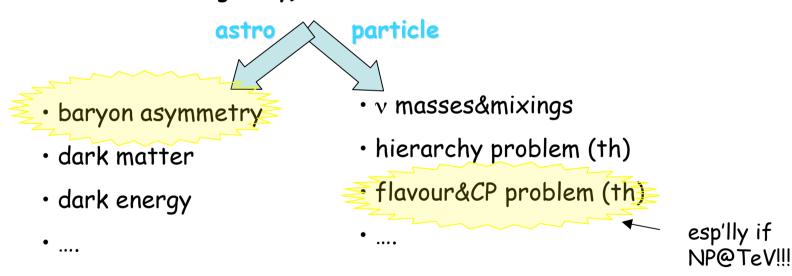
FOR A RECENT GENERAL REVIEW: Davidson Nardi Nir, arXiv:0802.2962

...leptogenesis is maybe the most promising mechanism to explain B-asymmetry

BUT WHY A TALK ABOUT IT IN A "FLAVOR AS A WINDOW TO NEW PHYSICS AT LHC" WORKSHOP?

New Physics

necessary because (even postponing the problem of unification with gravity) it must account for:



→ leptogenesis and flavor are both related to New Physics bSM, so that flavor might play an important role in the generation of B-asymmetry

PLAN

- 1- What is and how to measure the B asymmetry
- 2- Dynamical mechanisms for Baryogenesis (beyond SM) > leptogenesis
- 3- Various leptogenesis → seesaw type I, II, III, ...
- 4- Leptogenesis via type I seesaw
 - a- effects: flavor, resonant, quantum
 - b- related phenomenology
 - i- connection with CPV in v masses
 - ii- range of RH ν masses
 - iii- embedding in susy: gravitino problem
 - iv- embedding in susy: LFV, EDM,
 - v- embedding in GUT/flavor models: 2 examples

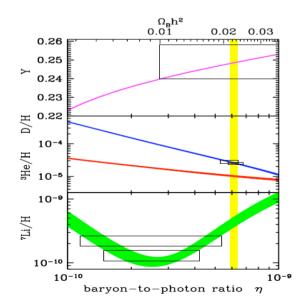
5- Conclusions

Baryon asymmetry of the universe

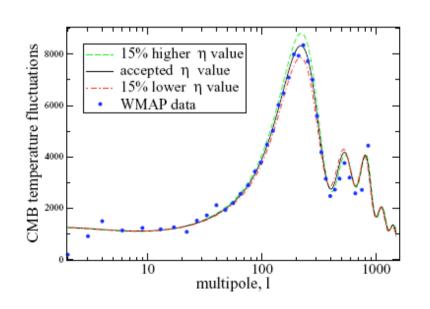
 n_X = # density of X

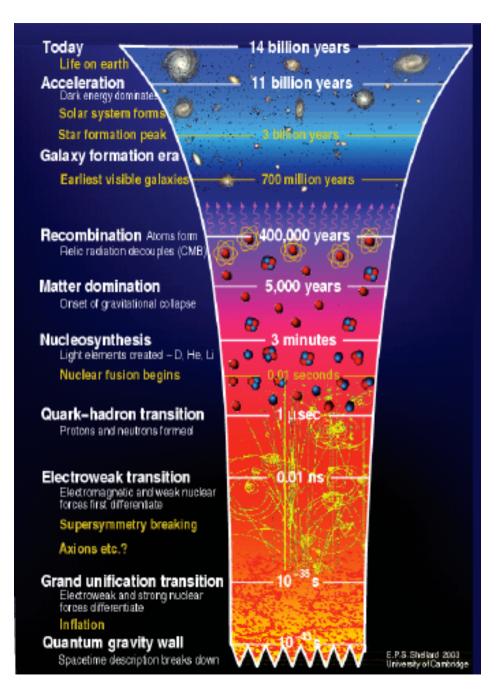
$$\eta_B = \frac{n_B - n_{\bar{B}}}{n_{\gamma}} = 274 \times 10^{-10} \Omega_B h^2 = (6.21 \pm .16) \times 10^{-10}$$
 @68% How to measure?

BBN - Historically



CMB - At present more precise





←Today

←CMB

←BBN

A NAIVE estimation of the B asymmetry of the universe TODAY

1.5 x
$$10^{79}$$
 Hatoms in observed universe = 1.25 10^{11} galaxies x 10^{11} stars/galaxy x 1.2 10^{57} Hatoms/star

Hubble Space Telescope as in our galaxy as in the Sun $\frac{2\times10^{30} \text{ kg/star}}{1.67\times10^{-27} \text{ kg/Hatom}}$

3.6 x 10^{80} m³ = volume of observed universe

Hatom density =
$$\frac{4.2 \times 10^{-2} / \text{m}^3}{10^{-10}}$$
 = $\frac{10^{-10}}{10^{-10}}$ = $\frac{4.1 \times 10^8}{10^{-10}}$ = $\frac{10^{-10}}{10^{-10}}$ eq at T=2.73K

Baryon asymmetry of the universe

A <u>non-trivial</u> value! If the universe were B-antiB symmetric

$$\frac{n_B}{n_\gamma} = \frac{n_{ar{B}}}{n_\gamma} \sim 10^{-20}$$
annihilation catastrophe!

From fine-tuned (1/109) initial conditions?

No: with inflation, any preexisting B-asymm is diluted to a negligible value, due to entropy production during reheating





Need a <u>dynamical mechanism</u> to generate B-asymmetry after inflation!

Baryogenesis

Sakharov's 3 conditions ['67] to dynamically create the asymmetry

- 1 B Violation2 C & CP Violations
- 3- out of thermal equil'm

<u>Baryogenesis</u>

Sakharov's 3 conditions ['67] to dynamically create the asymmetry

1- B Violation → at quantum-level (triangle anomaly)
 2- C & CP Violations → maximal & TOO tiny

3- out of thermal equil'm > NOT so STRONG 1° EWPT

<u>SM</u>

possess all ingredients but does not work...

[Gavela Hernandez Orloff Pene]

Topological trans'n: sphalerons
$$\begin{array}{c} \text{Q}_1\\ \text{B-L conserved:}\\ \Delta_B = \Delta_L = \pm 3 \begin{array}{c} \text{L}_{\text{p}}\\ \text{L}_{\text{t}} \end{array}$$
 [Kuzmin Rubakov Shaposhnikov '85]

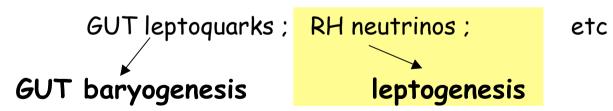
Mechanisms for Baryogenesis

Sakharov's 3 conditions ['67] to dynamically create the asymmetry

```
1- B Violation
2- C & CP Violations
3- out of thermal equil'm
```

A) out of eq decay (3-) of heavy particles whose int's violate C, CP (2-)

and B-L (1-) so that SM sphalerons do not erase the B asymmetry:



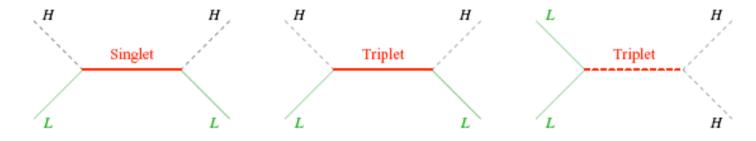
B) OTHERS: EW baryogenesis (modification of EWPT): 2HDM, MSSM; spontaneous baryogenesis; Affleck-Dine; gravitational leptogenesis; etc

Various Leptogensis

...according to type of seesaw inducing (

$$\mathscr{L} = \mathscr{L}_{SM} + \frac{(LH)^2}{2\Lambda_I}$$

The neutrino Majorana mass operator (LH)² can be mediated by tree level exchange of:



- a fermion singlet ('see-saw');
- II) a fermion triplet;
- III) a scalar triplet.



The seesaw (I)

[P.Minkowski '77]

Dirac-Yukawa Majorana-mass if complex
$$\rightarrow$$
 CPV Δ L=2
$$\mathcal{L}_{ss} = \bar{\nu}_R Y_{\nu} \nu_L H + \bar{\nu}_R^c M_{\nu} \nu_R + h.c.$$

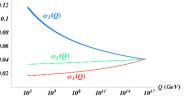
Mass eigenstates:

Heavy:
$$Npprox
u_R$$
 with $M_
u=(M_1,M_2,M_3)$ (where M1< M2< M3)

Light:
$$\nu \approx \nu_L \quad \text{with} \quad m_\nu = U^* m_\nu^d U^\dagger = Y_\nu^T M_R^{-1} Y_\nu v^2$$

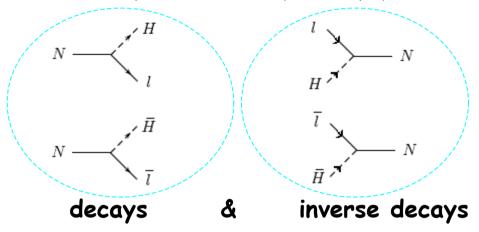
MNS mixing matrix

$$m_v = O(eV) \rightarrow M_v = O(10^{15} GeV)$$
, near SUSY g.c.u.! $\frac{au}{au}$

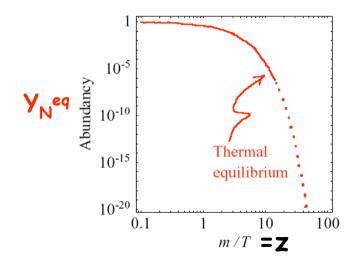


0. In the thermal bath with $\Gamma(N o \ell H, ar{\ell} H) > H(T)$

 \leftarrow tipically T>M_R

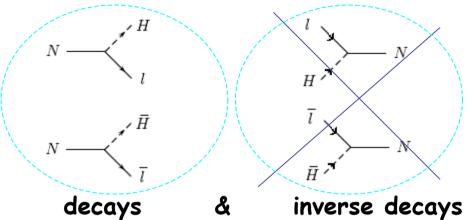


are in thermal equilibrium



1. In the thermal bath when $\Gamma(N o \ell H, ar{\ell} ar{H}) extbf{<} H(T)$

← tipically T<M_R

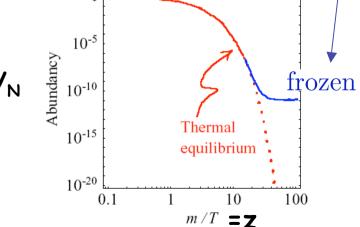


are no more effective in diluting Y_N

go out of equilibrium!

sooner (i.e. N's more abundant) the more

$$K = \Gamma/H(T = M_R) < 1$$

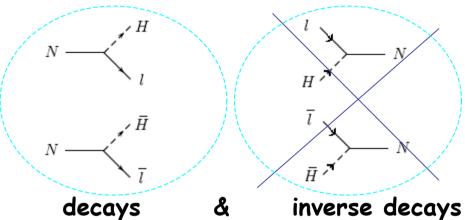


<u>Leptogenesis</u>

[Fukugita Yanagida '86]

1. In the thermal bath when $\Gamma(N o \ell H, \bar{\ell} \bar{H}) \blacktriangleleft H(T)$

← tipically T<M_R



are no more effective in diluting Y_N

go out of equilibrium!

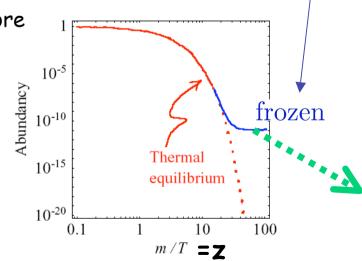
sooner (i.e. N's more abundant) the more

$$K = \Gamma/H(T = M_R) < 1$$

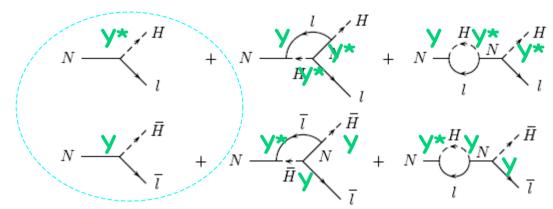
HOWEVER: later N's decay

$$\frac{dY_{N_i}}{dz} = -D_i \left(Y_{N_i} - Y_{N_i}^{eq} \right)$$

Boltzmann eq'n



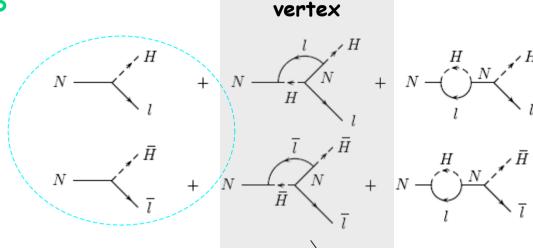
2. Violating C&CP



lepton asymm

$$\epsilon_{i\alpha} \equiv \frac{\Gamma(N_i \to \ell_\alpha \bar{H}) - \Gamma(N_i \to \bar{\ell}_\alpha H)}{\Gamma(N_i \to \ell_\alpha \bar{H}) + \Gamma(N_i \to \bar{\ell}_\alpha H)} = \sum_{j \neq i} \frac{1}{8\pi} \frac{\mathrm{Im} \left[(\lambda_\nu)_{i\alpha} (\lambda_\nu)^\dagger_{\alpha j} (\lambda_\nu \lambda_\nu^\dagger)_{ij} \right]}{\left(\lambda_\nu \lambda_\nu^\dagger \right)_{ii}} \left(g_{\mathbf{S}}^{(j,i)} + g_{\mathbf{S}}^{(j,i)} \right)$$

2. Violating C&CP

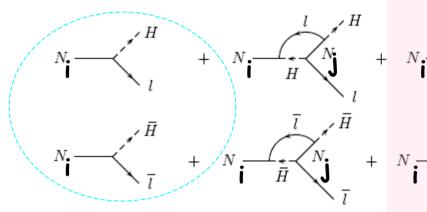


lepton asymm

$$\epsilon_{i\alpha} \equiv \frac{\Gamma(N_i \to \ell_\alpha \bar{H}) - \Gamma(N_i \to \bar{\ell}_\alpha H)}{\Gamma(N_i \to \ell_\alpha \bar{H}) + \Gamma(N_i \to \bar{\ell}_\alpha H)} = \sum_{j \neq i} \frac{1}{8\pi} \frac{\mathrm{Im} \left[(\lambda_\nu)_{i\alpha} (\lambda_\nu)_{\alpha j}^\dagger (\lambda_\nu \lambda_\nu^\dagger)_{ij} \right]}{\left(\lambda_\nu \lambda_\nu^\dagger \right)_{ii}} \stackrel{(g^{(j,i)}_{\mathbf{S}})}{(g^{(j,i)}_{\mathbf{S}})} + g^{(j,i)}_{\mathbf{S}})$$

most important for hierarchical RH ν 's \rightarrow lower limit on M_1

2. Violating C&CP



lepton asymm

$$\epsilon_{i\alpha} \equiv \frac{\Gamma(N_i \to \ell_\alpha \bar{H}) - \Gamma(N_i \to \bar{\ell}_\alpha H)}{\Gamma(N_i \to \ell_\alpha \bar{H}) + \Gamma(N_i \to \bar{\ell}_\alpha H)} = \sum_{j \neq i} \frac{1}{8\pi} \frac{\mathrm{Im} \left[(\lambda_\nu)_{i\alpha} (\lambda_\nu)^\dagger_{\alpha j} (\lambda_\nu \lambda_\nu^\dagger)_{ij} \right]}{\left(\lambda_\nu \lambda_\nu^\dagger \right)_{ii}} \left(g_{\mathbf{J},i}^{(j,i)} + g_{\mathbf{S},i}^{(j,i)} \right)$$

Resonant effects

[Pilaftsis, Covi Roulet, Flanz... '97]

relevant for nearly deg RH v's

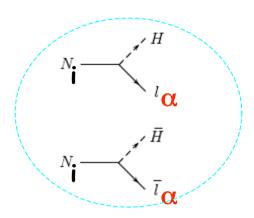
resonantly enhanced for

self-energy

$$\Delta M_{ij} = M_i - M_j \lesssim \Gamma_j$$
 (i-decays, j-internal)

 \rightarrow M₁ as low as TeV

<u>Leptogenesis</u>

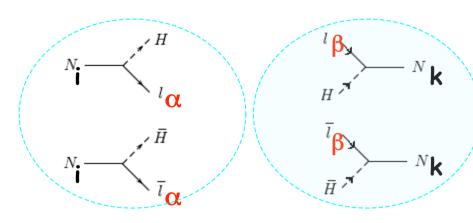


lepton asymm

$$\mathbf{E}_{\mathbf{i}\alpha} \equiv \frac{\Gamma(N_i \to \ell_\alpha \bar{H}) - \Gamma(N_i \to \bar{\ell}_\alpha H)}{\Gamma(N_i \to \ell_\alpha \bar{H}) + \Gamma(N_i \to \bar{\ell}_\alpha H)} = \sum_{j \neq i} \frac{1}{8\pi} \frac{\mathrm{Im} \left[(\lambda_\nu)_{i\alpha} (\lambda_\nu)_{\alpha j}^\dagger (\lambda_\nu \lambda_\nu^\dagger)_{ij} \right]}{\left(\lambda_\nu \lambda_\nu^\dagger \right)_{ii}} \; (g_{\mathbf{i}}^{(j,i)} + g_{\mathbf{S}}^{(j,i)})$$

3. BV- Sphaleron conserve $Y_{\Delta\alpha} = B/3 - L_{\alpha}$

$$\eta_B \frac{n_\gamma}{s} = Y_B = \frac{12}{37} \sum_{\alpha} Y_{\Delta_{\mathbf{Q}}}(z \to \infty)$$



inverse decays

wash out $\varepsilon_{i\alpha}$ only if $\beta = \alpha$

lepton asymm

$$\mathbf{E}_{\mathbf{i}\alpha} \equiv \frac{\Gamma(N_i \to \ell_\alpha \bar{H}) - \Gamma(N_i \to \bar{\ell}_\alpha H)}{\Gamma(N_i \to \ell_\alpha \bar{H}) + \Gamma(N_i \to \bar{\ell}_\alpha H)} = \sum_{j \neq i} \frac{1}{8\pi} \frac{\mathrm{Im} \left[(\lambda_\nu)_{i\alpha} (\lambda_\nu)_{\alpha j}^\dagger (\lambda_\nu \lambda_\nu^\dagger)_{ij} \right]}{\left(\lambda_\nu \lambda_\nu^\dagger \right)_{ii}} \; \left(g_{\mathbf{i}}^{(j,i)} + g_{\mathbf{S}}^{(j,i)} \right)$$

3. BV- Sphaleron conserve $Y_{\Delta\alpha} = B/3 - L_{\alpha}$

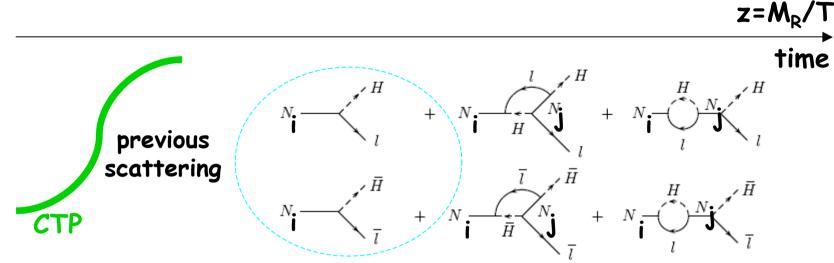
Boltzmann eq'n:
$$\frac{dY_{\Delta \mathbf{C}}}{dz} = -\sum_{i} \epsilon \mathbf{E}_{i} D_{i} \left(Y_{N_{i}} - Y_{N_{i}}^{eq} \right) - W_{\mathbf{C}} A_{\alpha\alpha} | Y_{\Delta_{\mathbf{C}}} | Y_{\Delta_{\mathbf{C}$$

$$\eta_B \frac{n_\gamma}{s} = Y_B = \frac{12}{37} \sum_{\alpha} Y_{\Delta_{\mathbf{Q}}}(z \to \infty)$$

Flavor effects

[Barbieri et al; Abada et al;...]

 $A_{\alpha\alpha}$ depend on which $Y_{ch,lept}$ interins are in equilim



lepton asymm

$$\mathbf{E}_{\mathbf{i}\alpha} \left(\mathbf{Z} \right) = \frac{\Gamma(N_i \to \ell_\alpha \bar{H}) - \Gamma(N_i \to \bar{\ell}_\alpha H)}{\Gamma(N_i \to \ell_\alpha \bar{H}) + \Gamma(N_i \to \bar{\ell}_\alpha H)} = \sum_{j \neq i} \frac{1}{8\pi} \frac{\mathrm{Im} \left[(\lambda_\nu)_{i\alpha} (\lambda_\nu)^\dagger_{\alpha j} (\lambda_\nu \lambda_\nu^\dagger)_{ij} \right]}{\left(\lambda_\nu \lambda_\nu^\dagger \right)_{ii}} \left(g_{\mathbf{S}}^{(j,i)} + g_{\mathbf{S}}^{(j,i)} \right) m^{(i,j)} \left(\mathbf{Z} \right)$$

Quantum effects

[De Simone & Riotto '07]

relevant if t'scale-
$$\epsilon_{\rm i\alpha}$$
> t'scale-Y $_{\rm Ni}$ namely $\Delta M_{ij}=M_i-M_j\lesssim \Gamma_i$

$$2 \sin^2 \left(\frac{1}{2} \frac{M_j - M_i}{2H(M_1)} z^2\right) - \frac{\Gamma_j}{M_j - M_i} \sin \left(\frac{M_j - M_i}{2H(M_1)} z^2\right)$$

typical timescale $1/\Delta M_{ij}$

from solving quantum

Leptogenesis-related phenomenology

I-Connection with CPV in v masses

From seesaw: given M_R , reconstruction of Y_v ambiguous up to R (\sim =diag)

$$m_
u=U^*\hat{m}_
u U^\dagger=Y_
u^T\hat{M}_R^{-1}Y_
u v^2$$
 Casas-Ibarra parameterization $Y_
u=\sqrt{\hat{M}_R}R\sqrt{\hat{m}_
u}U^\dagger/v$

I-Connection with CPV in v masses

From seesaw: given M_R , reconstruction of Y_v ambiguous up to R (^=diag)

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u v^2$$
 Casas-Ibarra parameterization $Y_
u=\sqrt{\hat{M}_R}R\sqrt{\hat{m}_
u}U^\dagger/v$

(hierarchical case)

Neglecting flavor:
$$\epsilon_1 = -\frac{3M_1}{16\pi v^2} \frac{\operatorname{Im}\left(\sum_{\rho} m_{\rho}^2 R_{1\rho}^2\right)}{\sum_{\beta} m_{\beta} \left|R_{1\beta}\right|^2}$$

function of R and m, only!

[Davidson Ibarra, Branco IM et al, etc]

Including flavor:
$$\epsilon_{\alpha} = -\frac{3M_{1}}{16\pi v^{2}} \frac{\operatorname{Im}\left(\sum_{\beta\rho}m_{\beta}^{1/2}m_{\rho}^{3/2}U_{\alpha\beta}^{*}U_{\alpha\rho}R_{1\beta}R_{1\rho}\right)}{\sum_{\beta}m_{\beta}\left|R_{1\beta}\right|^{2}}$$

introduce dep on U

[Davidson, Ibarra, Petcov et al, etc]

II-Range for RH v masses

Neglecting flavor:

efficiency factor $\kappa < 0.1$

$$6 \times 10^{-10} \approx \eta_B \approx 10^{-2} \epsilon \, \kappa \lesssim 10^{-3} \epsilon$$
 \rightarrow $\epsilon \gtrsim 10^{-6}$

$$\rightarrow$$
 $\epsilon \gtrsim 10^{-6}$

II-Range for RH v masses

Neglecting flavor:

efficiency factor K<0.1

$$6 \times 10^{-10} \approx \eta_B \approx 10^{-2} \epsilon \, \kappa \lesssim 10^{-3} \epsilon$$
 \rightarrow $\epsilon \gtrsim 10^{-6}$

$$\rightarrow$$
 $\epsilon \gtrsim 10^{-6}$

Hierarchical (enough)

[Davidson Ibarra '02]

$$|\epsilon| < \frac{3}{16\pi} \frac{(m_{\rm max} - m_{\rm min}) M_1}{v_u^2}$$



$$10^9 {\rm GeV} \lesssim M_1 \lesssim T_{RH}$$

Nearly degenerate

$$|\epsilon_{N_1}(\text{resonance})| \simeq \frac{1}{2} \frac{|\mathcal{I}m[(\lambda^{\dagger}\lambda)_{12}^2]|}{(\lambda^{\dagger}\lambda)_{11}(\lambda^{\dagger}\lambda)_{22}} < O(1)$$



$$\text{TeV} \lesssim M_1 \lesssim T_{RH}$$

Including flavor: lower bound on M_1 relaxed in general by O(2)

III-Embedding in SUSY: gravitinos

TRH has to be small enough to avoid overproduction of gravitinos during reheating. Being only gravitationally coupled to MSSM particles, they decay late destroying successful BBN. [Moroi, etc]

For
$$m_{3/2} pprox 10^2 - 10^3 {
m GeV}$$
 gravitino bound

$$T_{RH} \lesssim 10^5 - 10^7 \text{GeV}$$

Tension with lower bound

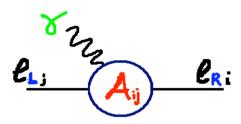
on M
$$_{\scriptscriptstyle 1}$$
 for hier RH v's $10^9 {
m GeV} \lesssim M_1$

→thermal production of RH v's is inefficient

WAYS OUT:

non-thermal production, e.g. via flat directions in scalar potential [Giudice etc]; stable gravitinos (but then model dependence); resonant leptogenesis; etc etc

IV-Embedding in SUSY: LFV & EDM



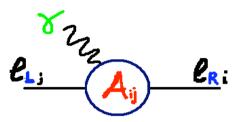
Loops w/ Sleptons & Gauginos

LFVdecays

$$BR(\ell_i \to \ell_j \gamma) \propto |A_{ij}|^2 = f_{LL} |\delta_{ji}^{LL}|^2 + f_{RR} |\delta_{ji}^{RR}|^2 + f_{LR} |\delta_{ji}^{LR}|^2 + f_{RL} |\delta_{ji}^{RL}|^2 + ...$$

$$d_{\ell_i} = ImA_{ii} = f_a m_{\ell_i} Ima_i + f_{LLRR} Im(\delta^{LL} m_\ell \delta^{RR})_{ii} + ...$$

IV-Embedding in SUSY: LFV & EDM



Loops w/ Sleptons & Gauginos

From **mSUGRA** at M_{Pl,} susy seesaw induce FV and CPV at I.e. via **RGE**

$$C_{ij}^k = Y_{\nu ki}^* Y_{\nu kj} \ln \frac{M_{Pl}}{M_k}$$

[BorzumatiMasiero '86]

$$-\frac{1}{(4\pi)^2} \frac{6m_0^2 + 2a_0^2}{\bar{m}_L^2} \sum_k C_{ij}^k$$

$$BR(\ell_i \to \ell_j \gamma) \propto |A_{ij}|^2 = f_{LL} |\delta_{ji}^{LL}|^2 + f_{RR} |\delta_{ji}^{RR}|^2 + f_{LR} |\delta_{ji}^{LR}|^2 + f_{RL} |\delta_{ji}^{RL}|^2 + ...$$

EDMs

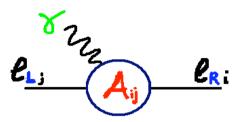
$$d_{\ell_i} = Im A_{ii} = f_a m_{\ell_i} Im a_i + f_{LLRR} Im (\delta^{LL} m_\ell \delta^{RR})_{ii} + ...$$
 FC
 $\delta^{RR} = \int_{\{1, 1\}^k} \frac{\ln_{k'}^k}{\ln^{2l}} \operatorname{Im}(C^k C^{k'})_{ii}$
 $\delta^{RR} = \int_{\{1, 1\}^k} \frac{\ln_{k'}^k}{\ln^{2l}} \operatorname{Im}(C^k C^{k'})_{ii}$
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[EllisHisanoLolaRaidalShimizu '01]

$$\frac{8m_{\ell_i}}{(4\pi)^6} \frac{(6m_0^2 + 2a_0^2)(6m_0^2 + 3a_0^2)}{\bar{m}_L^2 \bar{m}_R^2} \frac{m_\tau^2 \tan^2 \beta}{v^2} \sum_{k>k'} \tilde{\ln}_{k'}^k \operatorname{Im} \left(C^k \frac{m_\ell^2}{m_\tau^2} \, C^{k'} \right)_{ii}$$

[IM '03]

IV-Embedding in SUSY: LFV & EDM



Loops w/ Sleptons & Gauginos

From **mSUGRA** at M_{Pl,} susy seesaw induce FV and CPV at I.e. via **RGE**

$$C_{ij}^k = Y_{\nu ki}^* Y_{\nu kj} \ln \frac{M_{Pl}}{M_k}$$

at present (future)

LFV decays BR(
$$\mu \rightarrow e \gamma$$
) < 10⁻¹¹⁽⁻¹³⁾ $\rightarrow C_{21}$ < 10⁻¹⁽⁻²⁾

strong impact on see-saw models

[Buchmuller et al; Sato, Tobe, Yanaqida; Casas Ibarra; King et al; Lavignac I.M. Savoy; Masiero et al;]

EDMs

 $d_e < 10^{-27(-30)} e cm \rightarrow$

seesaw contribution to de observable in future only if $tg\beta$ large, RH v's are Hi and various yukawas are O(1) [IM'04]

$$egin{aligned} \mathbf{\mathcal{C}}.\mathbf{\mathcal{G}}.\ Y_{
u} &= \left(egin{array}{cccc} \lesssim \mathbf{i}ar{\mathbf{o}^2} &\lesssim \mathbf{i}ar{\mathbf{o}^2} \ \mathcal{O}(1) & oldsymbol{lpha} oldsymbol{\mathcal{O}} & \mathcal{O}(1) \ \mathcal{O}(1) & oldsymbol{lpha} oldsymbol{\mathcal{O}} & \mathcal{O}(1) \end{array}
ight) \end{aligned}$$

e.g. $Y_{\nu} = \begin{pmatrix} \lesssim_{1}\bar{o}^{2} & \lesssim_{1}\bar{o}^{2} & \lesssim_{1}\bar{o}^{2} \\ \mathcal{O}(1) & \text{\approx0} & \mathcal{O}(1) \\ \mathcal{O}(1) & \text{\approx0} & \mathcal{O}(1) \end{pmatrix}$ Th leptogenesis & observable seesaw-induced $d_{e} \rightarrow M_{1} > 10^{11} \text{ GeV} \rightarrow \text{can't explain both with seesaw}$

[IM Riotto Joaquim, ph/0701270]

V-Embedding in flavor models/GUTs

To look for correlations among observables:

2 examples with A) hierarchical B) degenerate RH v's

A SU(5)XU(1)_F with q>0 [Froggatt Nielsen]

$$\begin{cases} \eta_B: & M_1 = O(10^{11}) \text{ GeV} \\ \mu \to e \gamma: & M_3 \le 5 \times 10^{12-13} \text{ GeV} \times 1/30 \text{ in future} \end{cases}$$

→ WHOLE CLASS could be TESTED!! [IM Savoy '05]

B Minimal Lepton Flavor Violation

Quantum effects [DeSimone Riotto Blanchet DiBari Raffelt] might be as large as $O(10^3)$ [Cirigliano De Simone Isidori IM Riotto '07] ...despite "Minimal", not a very predictive model for η_B

Conclusions

New TeV-scale physics should not Violate too much F&CP

New physics must explain B-asymmetry \rightarrow e.g. via leptogenesis

New Phys = seesaw

As happens for leptogenesis

(and depending on the details of the model),

the flavor structure of New Physics

can play a significant role in the generation of the B-asymmetry