



# Confinement and 4-manifolds

based on:

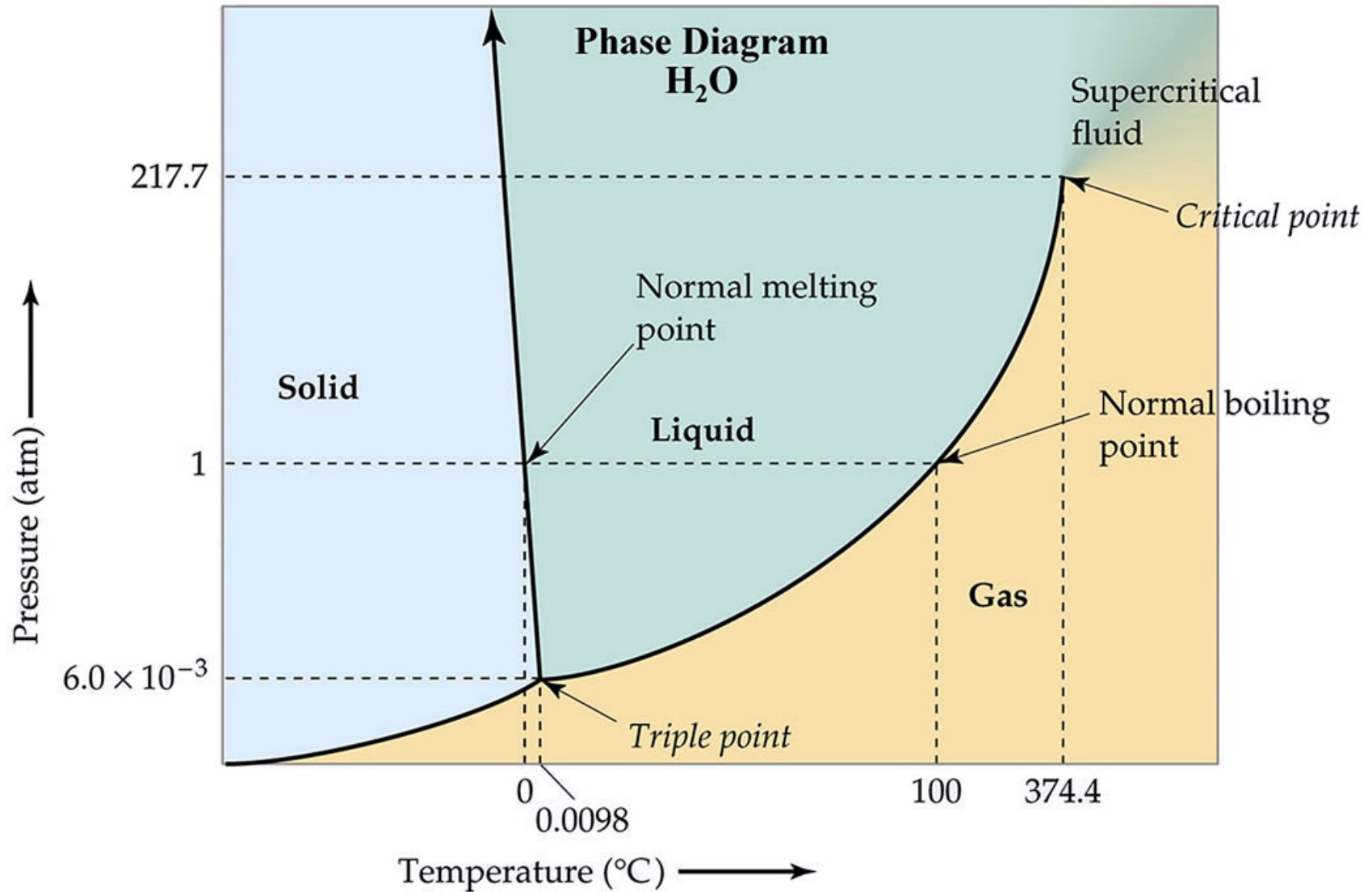
arXiv:1306.4320 (2d  $\mathcal{N} = (0,2)$  theories labeled by 4-manifolds)

arXiv:1404.2929 (duality defects and Lefschetz fibrations)

arXiv:1404.5314 (exact solutions of  $\mathcal{N} = (0,2)$  gauge theories)

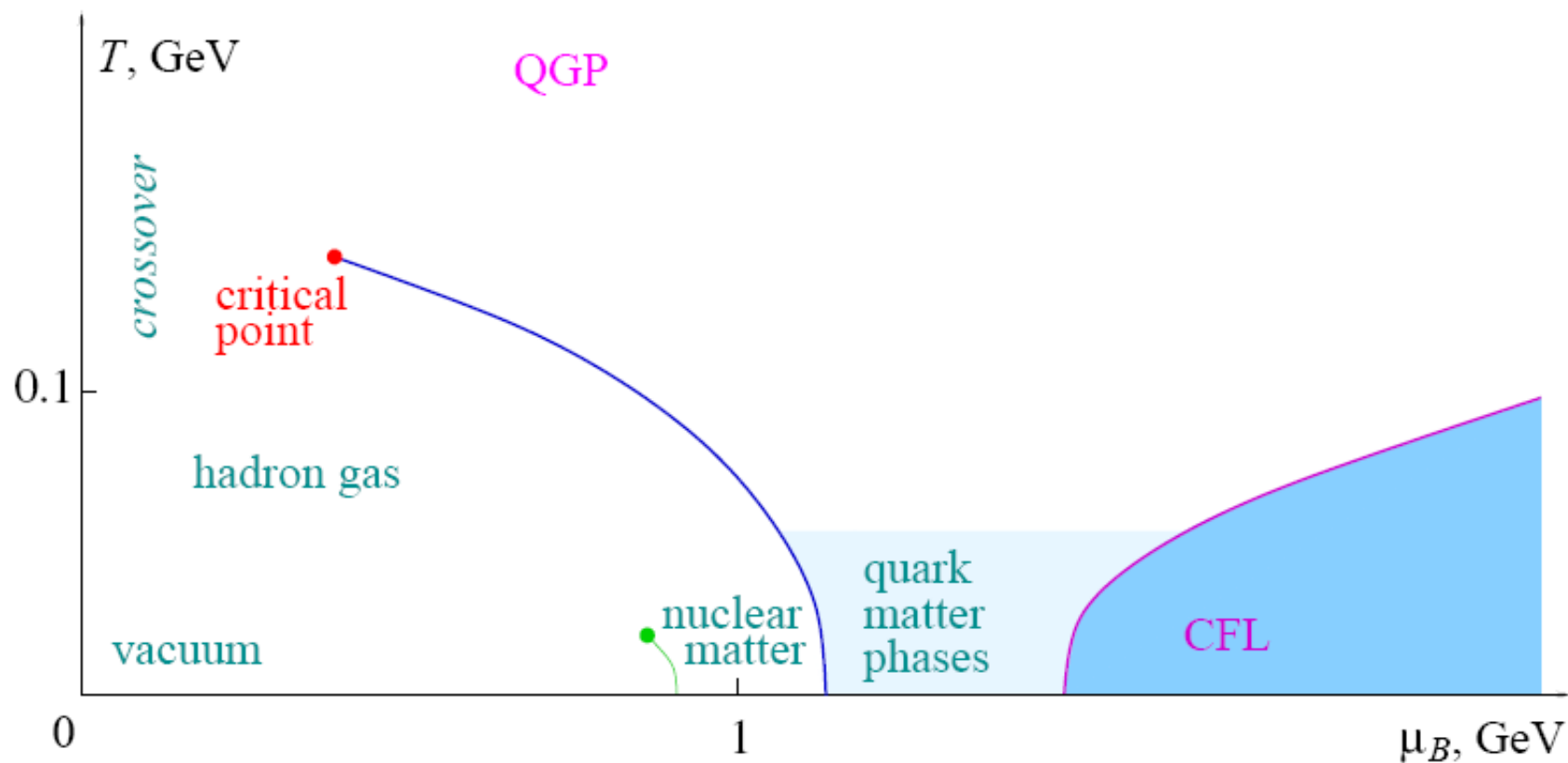
with **A.Gadde** and **P.Putrov**

# Phases

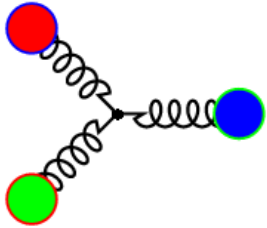


# Phases

[M.Stephanov]



# Can we quantitatively understand confinement and the mass gap?



- Extensively tested in computer simulations
- Paper-and-pencil computation?

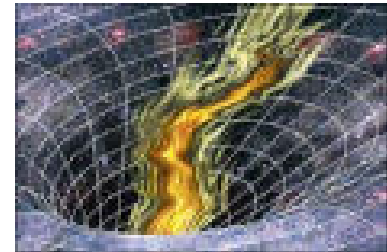


CLAY  
MATHEMATICS  
INSTITUTE

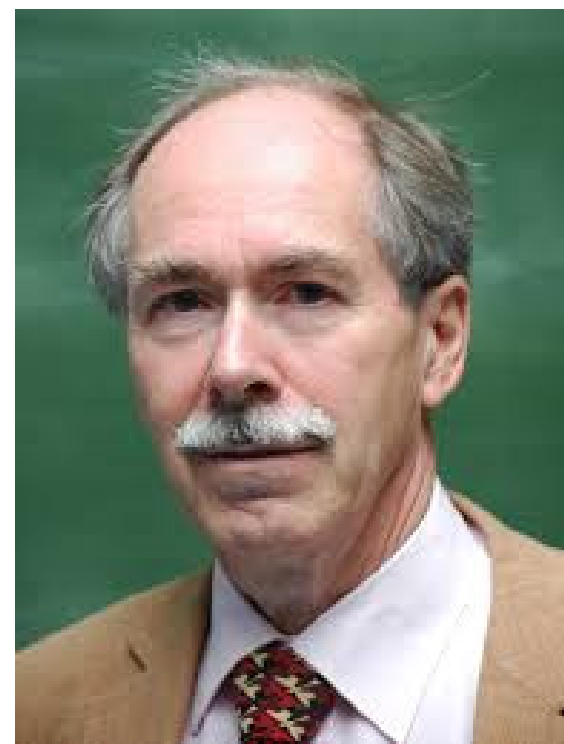
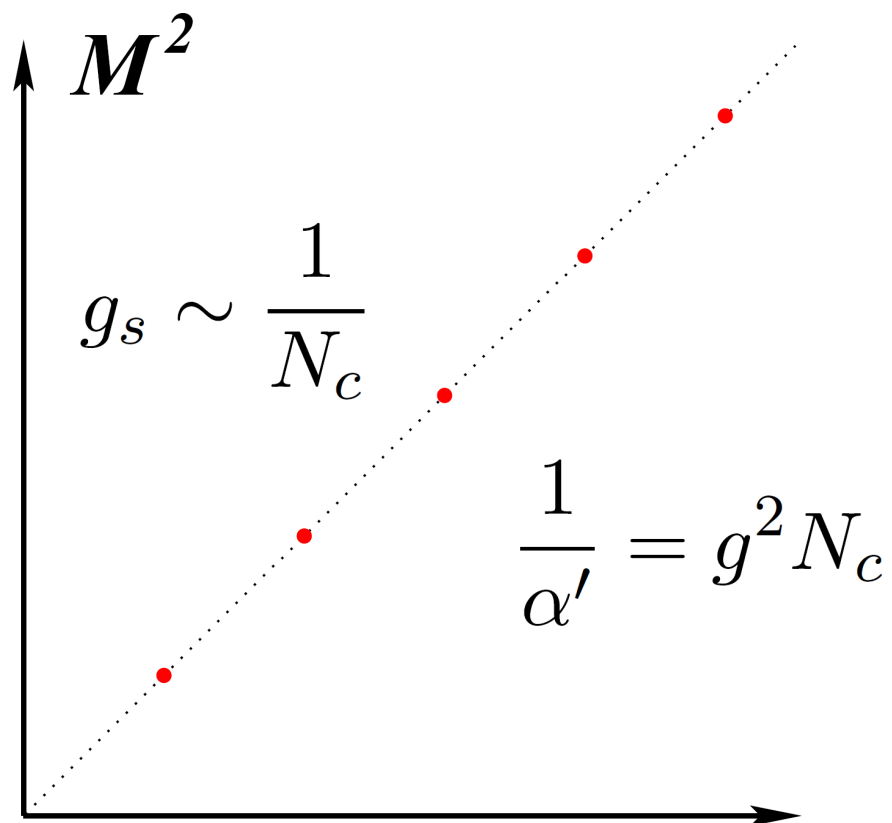
\$1,000,000 Prize



The answer may involve gravity!

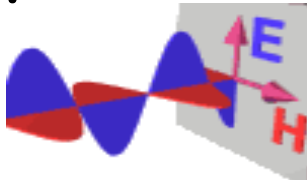


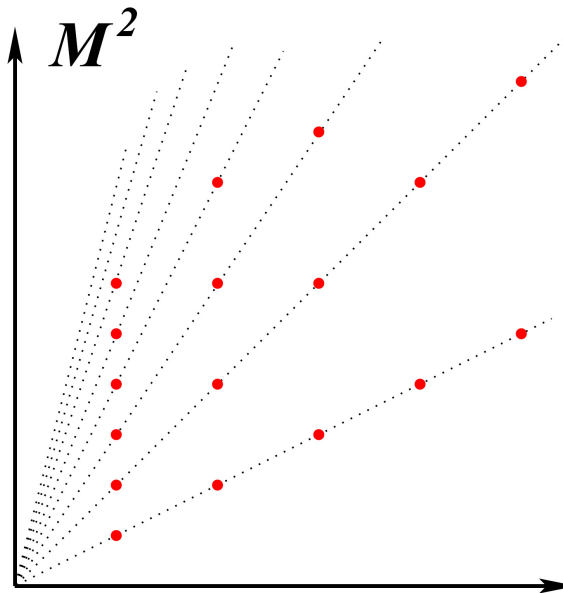
# Solvable Gauge Theories



't Hooft

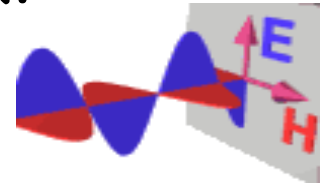
# Solvable Gauge Theories

- 2d Yang-Mills: almost “topological”
- QED: confinement / screening 
- 2d  $\mathcal{N} = 0$  QCD: one Regge trajectory
- QCD with a massive adjoint: higher Regge

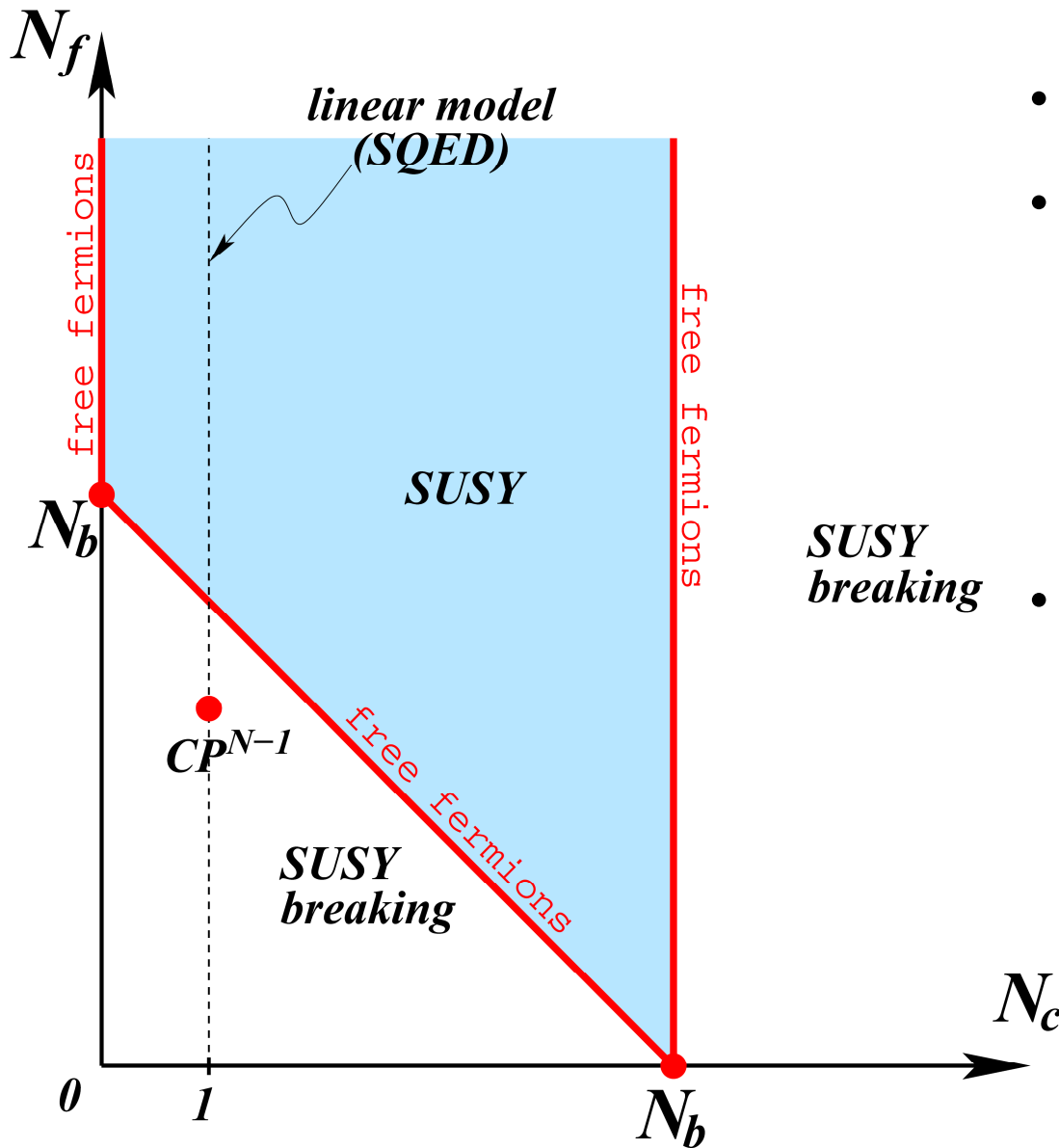


# Solvable Gauge Theories

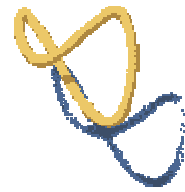
- 2d Yang-Mills: almost “topological”
- QED: confinement / screening
- 2d  $\mathcal{N} = 0$  QCD: one Regge trajectory
- QCD with a massive adjoint: higher Regge
- 2d  $\mathcal{N} = (1,1)$  SQCD: very similar
- 2d  $\mathcal{N} = (0,2)$  SQCD: ?
- 2d  $\mathcal{N} = (2,2)$  SQCD: ?
- :



# 2d $\mathcal{N} = (0,2)$ SQCD



- New  $(0,2)$  SCFTs
- Exactly solvable string phenomenology



*SUSY breaking*

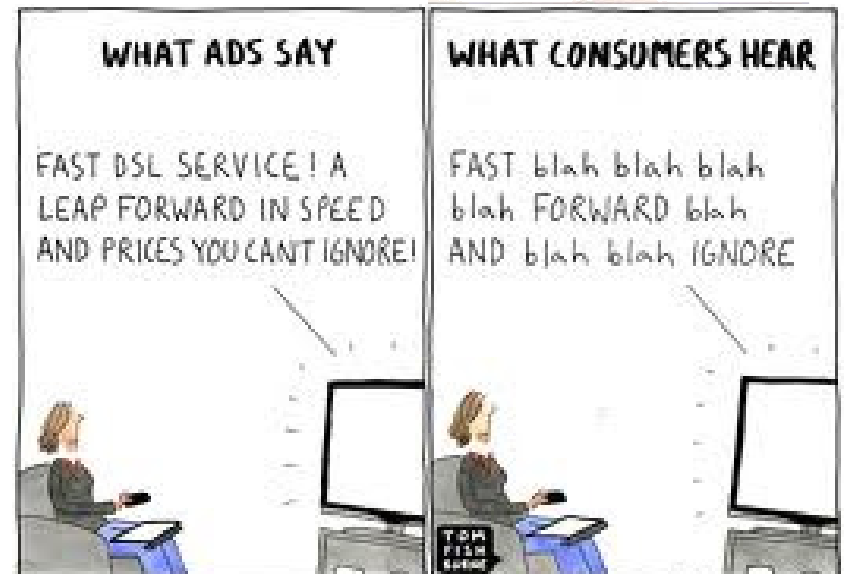
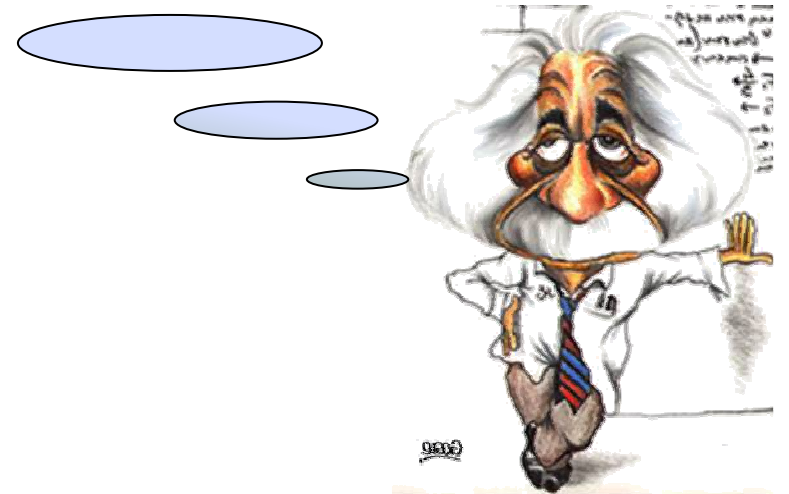
- Holographic dual of the Veneziano limit



higher spin theory

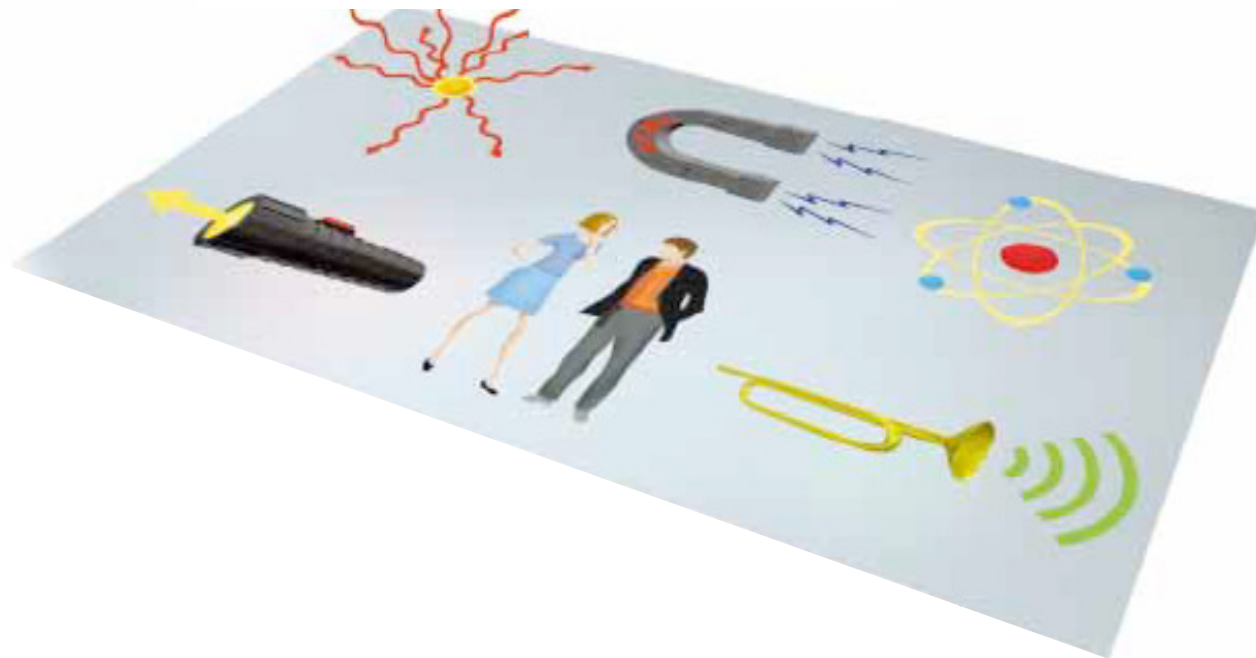


Okay, let us consider fivebranes on a  
coassociative 4-manifold in a  $G_2$   
holonomy space ...



# What are Fivebranes?

- 6-dimensional submanifolds in 11-dimensional space-time of M-theory



# Kaluza-Klein compactification

6d fivebrane theory

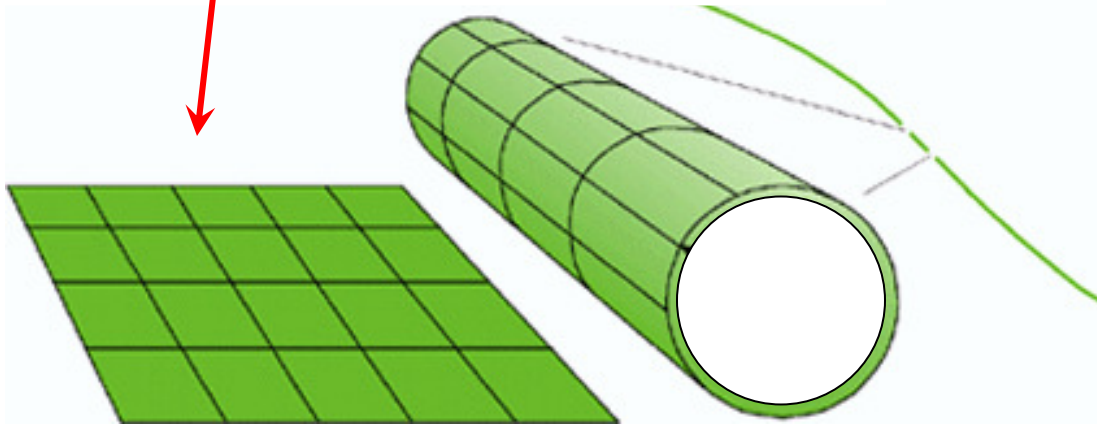
on  $\mathbb{R}^{6-n} \times M_n$



"effective" theory

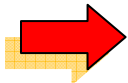
$T[M_n]$

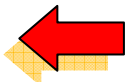
in  $6-n$  dimensions



depends on  
topology and  
geometry of  $M_n$

$$6 = 3 + 3$$

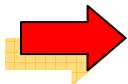
3-manifold  $M_3$   3d  $\mathcal{N} = 2$  theory  
 $T[M_3]$

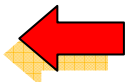
complex flat connections  supersymmetric vacua

$$dA + A \wedge A = 0 \qquad \frac{\partial \mathcal{W}}{\partial \sigma_i} = 0$$

$$\mathcal{M}_{\text{flat}}(M_3; G_{\mathbb{C}}) = \mathcal{M}_{\text{SUSY}}(T[M_3])$$

$$6 = 3 + 3$$

3-manifold  $\mathcal{M}_3$   3d  $\mathcal{N} = 2$  theory  
 $\mathcal{T}[\mathcal{M}_3]$

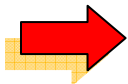
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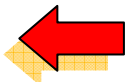
Example:  $\mathcal{M}_3 = L(k, 1)$

$$\frac{\partial \mathcal{W}}{\partial \sigma_i} = 0$$

$\rho$  = highest weight integrable representation of  
the loop group  $\mathcal{L}G$  at level  $k$

$$6 = 3 + 3$$

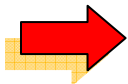
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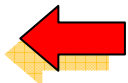
complex flat connections  supersymmetric vacua

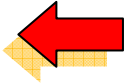
- Knot complements: A-polynomial
- generalized / quantum / homological Volume Conjecture ...

$$\hat{A}(\hat{x}, \hat{y}) Z(\mathcal{M}) = 0$$

$$6 = 3 + 3$$

3-manifold  $M_3$   3d  $\mathcal{N} = 2$  theory  
 $T[M_3]$

complex flat connections  supersymmetric vacua


complex Chern-Simons partition function  vortex partition function

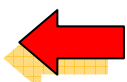
$$Z_{CS}(M_3; G_{\mathbb{C}}) = Z_{\text{vortex}}(T[M_3])$$

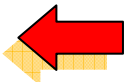
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3-manifold  $M_3$   3d  $\mathcal{N} = 2$  theory  
 $\mathcal{T}[M_3]$

complex flat connections  supersymmetric vacua

complex Chern-Simons  
partition function  vortex partition  
function

Knot homology  Q-cohomology

$$\mathcal{H}(\mathcal{O}) = \mathcal{H}_{\text{BPS}}$$



# Embedding in 11 dimensions

$q$ -grading

homological grading

$$U(1)_P \times U(1)_F$$



space-time:

$$\mathbb{R} \times \mathbb{R}_{q,t}^4 \times T^*M_3$$

$$\parallel \quad \cup \quad \cup$$

$N$  M5-branes:

$$\mathbb{R} \times \mathbb{R}_q^2 \times M_3$$

$$M_4 = \mathbb{R} \times M_3$$

$$\mathcal{H}(\mathcal{O}) = \mathcal{H}_{\text{refined BPS}}^{(\text{open})}$$

# Embedding in 11 dimensions

$q$ -grading

homological grading

$$U(1)_P \times U(1)_F$$



space-time:

$$\mathbb{R} \times \mathbb{R}_{q,t}^4 \times T^*M_3$$

$\parallel$

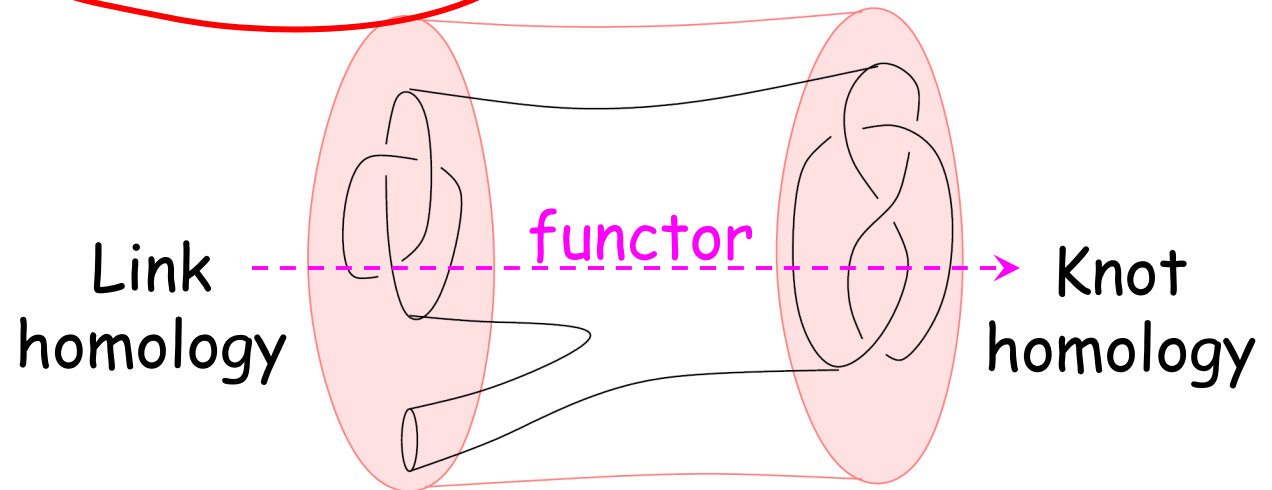
$\cup$

$\cup$

$$M_4 = \mathbb{R} \times M_3$$

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# Embedding in 11 dimensions

space-time:

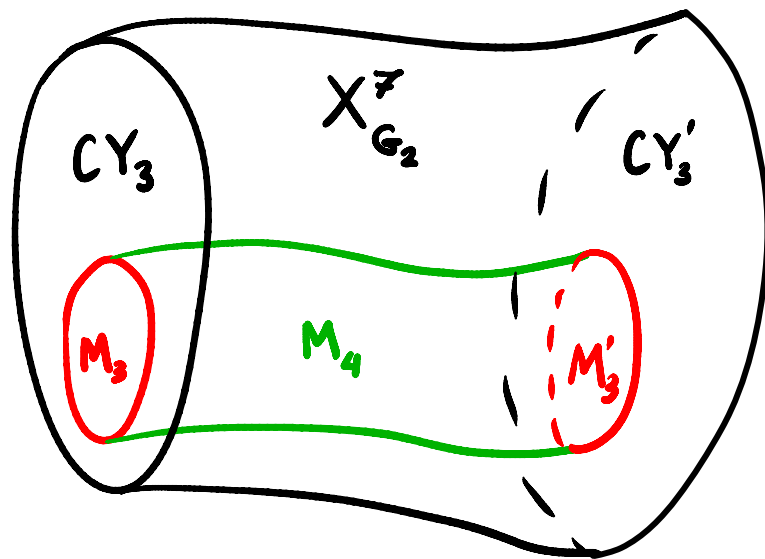
$$\mathbb{R} \times \mathbb{R}_{q,t}^4 \times T^*M_3$$

$$X^7 = \mathbb{R} \times T^*M_3$$

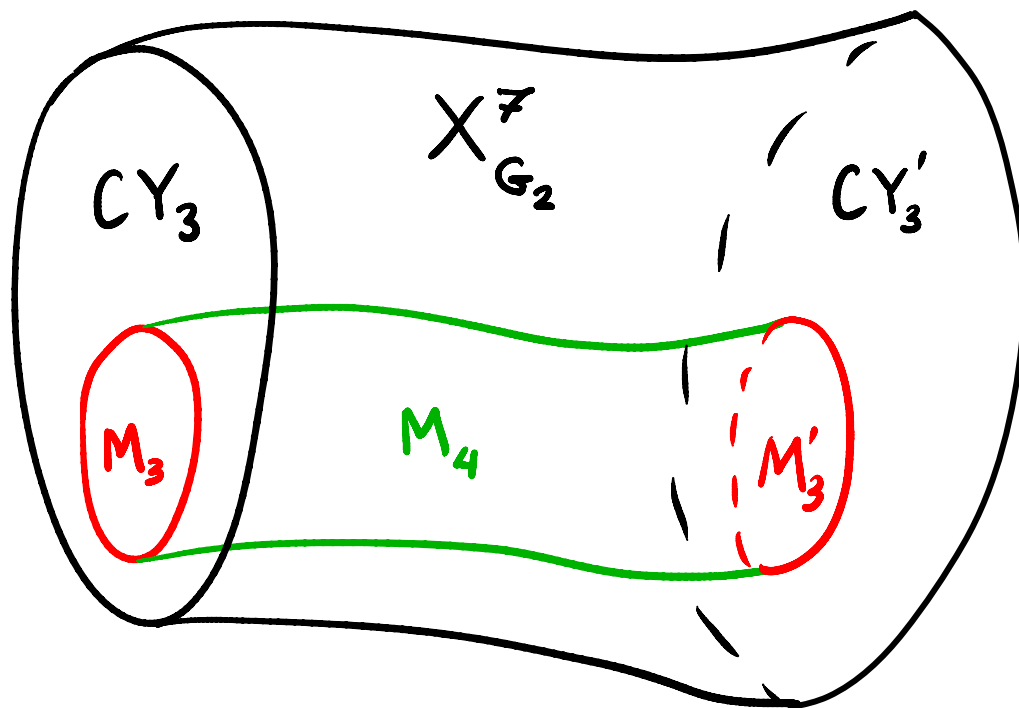
$N$  M5-branes:

$$\mathbb{R} \times \mathbb{R}_q^2 \times M_3$$

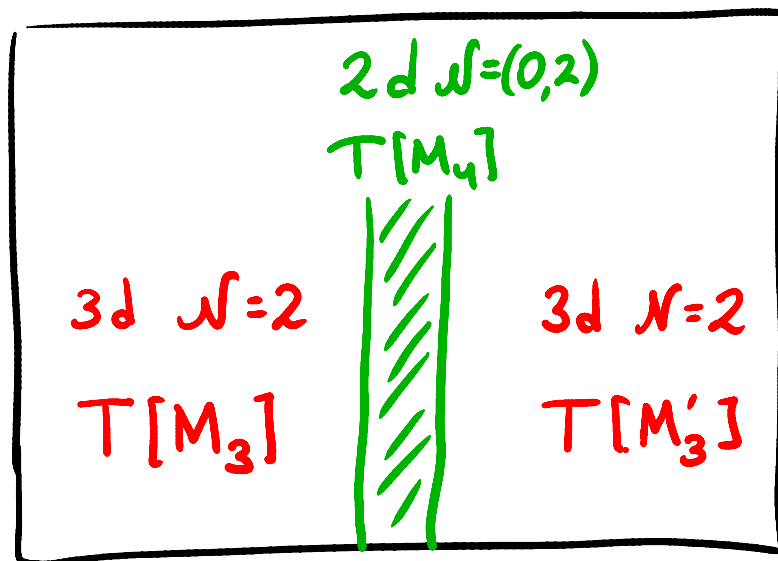
$$M_4 = \mathbb{R} \times M_3$$



cobordism



half-BPS  
domain wall

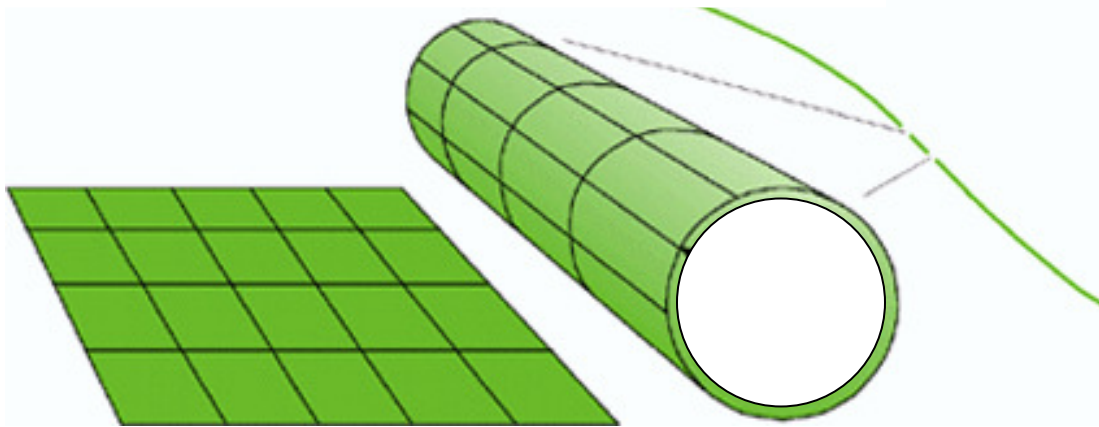


$$6 = 2 + 4$$

4-manifold  $\mathcal{M}_4$  

2d  $\mathcal{N} = (0,2)$  theory  
 $\mathcal{T}[\mathcal{M}_4]$

6d fivebrane theory  
on  $\mathbb{R}^2 \times \mathcal{M}_4$



depends on  
topology and  
geometry of  $\mathcal{M}_4$

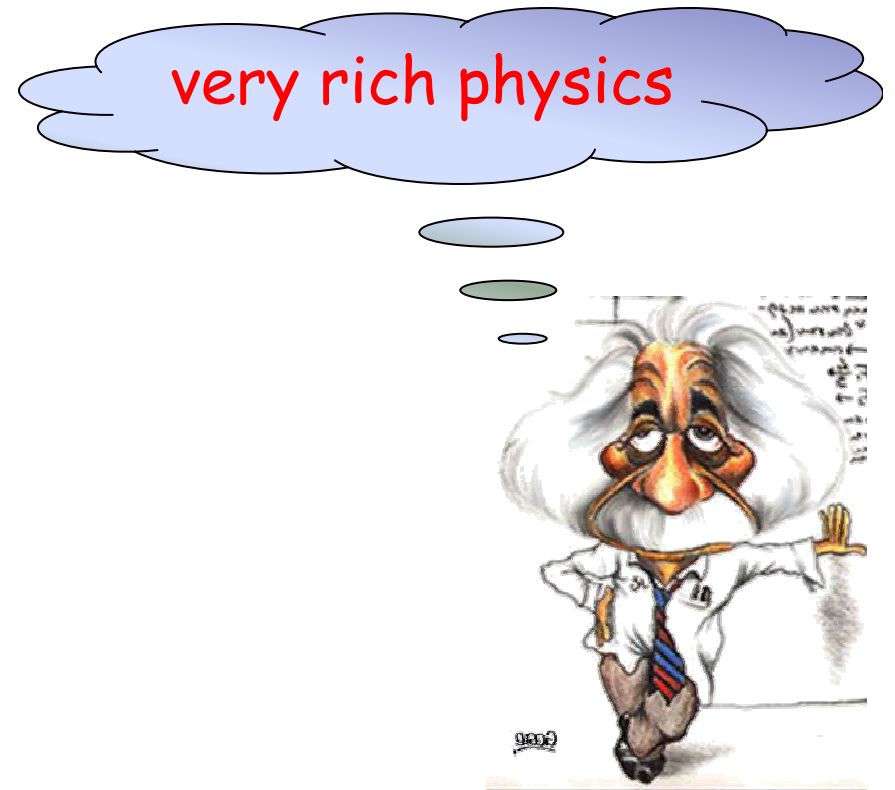
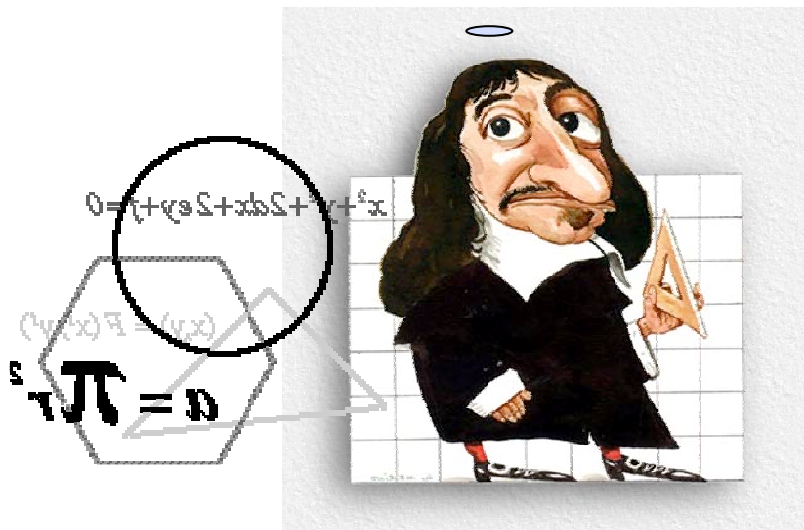
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4-manifold  $M_4$  


2d  $\mathcal{N} = (0,2)$  theory  
 $T[M_4]$

very rich math

very rich physics



$$6 = 2 + 4$$

4-manifold  $M_4$   2d  $\mathcal{N} = (0,2)$  theory  
 $T[M_4]$

$Z[M_4] = \sum_n (\mathbf{x}^c) q^n \chi(\mathcal{M}_{n,c}) =$  equivariant elliptic genus  
of the 2d  $\mathcal{N} = (0,2)$  theory

moduli space  
of instantons 

$$G_{\text{flavor}}(T[M_4]) = U(1)^{b_2}$$

$$F_A^+ = 0$$

modular !

$$6 = 2 + 4$$

6d fivebrane

on  $T^2 \times M_4$

[C.Vafa, E.Witten]

$\mathcal{N} = 4$  super-Yang-Mills  
on  $M_4$

2d (0, 2) theory  $T[M_4]$   
on  $T^2$

$Z[M_4] = \sum_n (x^q) q^n \chi(\mathcal{M}_{n,c}) =$  equivariant elliptic genus  
of the 2d  $\mathcal{N} = (0, 2)$  theory

moduli space  
of instantons

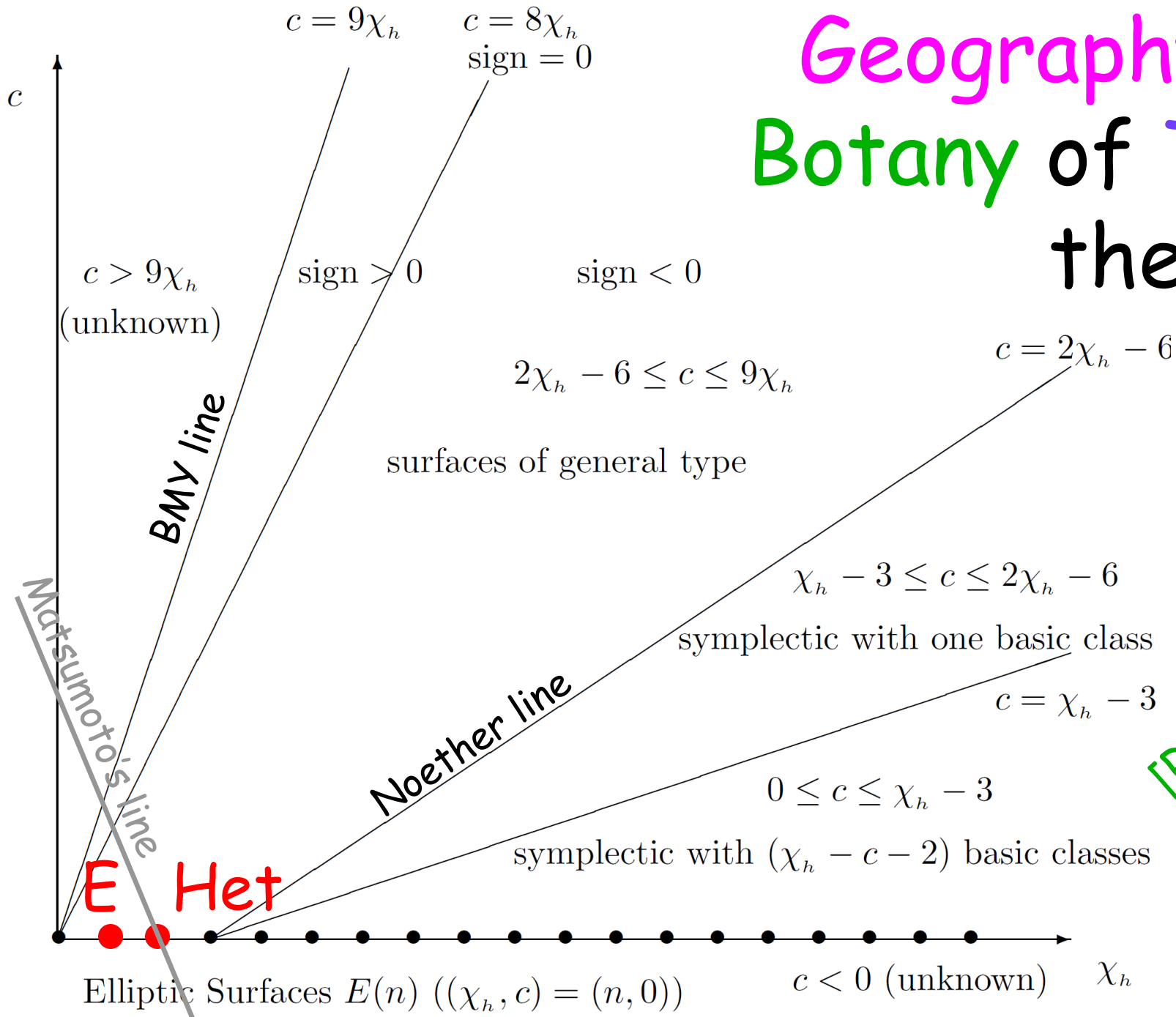
$$G_{\text{flavor}}(T[M_4]) = U(1)^{b_2}$$

$$F_A^+ = 0$$

modular !



# Geography and Botany of $T[M_4]$ theories



[R. Fintushel]

# Instantons on ALE spaces

$\mathcal{M}_4 = A_k$  ALE space, bounded by  $\mathcal{M}_3 = L(k+1, 1)$



elliptic genus of  $T[A_k, U(N)]$

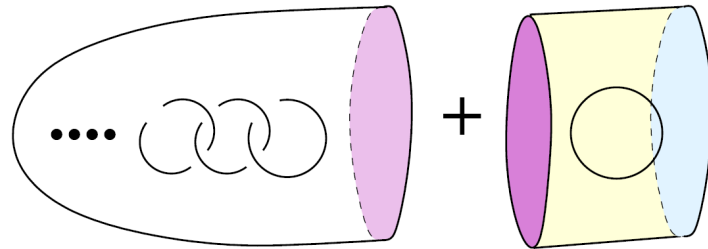
$\sum_n x^c q^n \chi(\mathcal{M}_{n,c})$  = character of a level  $N$  representation of affine  $SU(k+1)$  labeled by flat connection  $\rho$  on  $L(k+1, 1)$



[H.Nakajima]

# First New Results

$\mathcal{M}_4$  = cobordism between  $L(k,1)$  and  $L(k+1,1)$



$\sum_n x^c q^n \chi(\mathcal{M}_{n,c}) =$  branching function of a  $G/H$  coset

$$\chi_\rho^G = \sum_{\rho'} C_\rho^{\rho'} \chi_{\rho'}^H$$

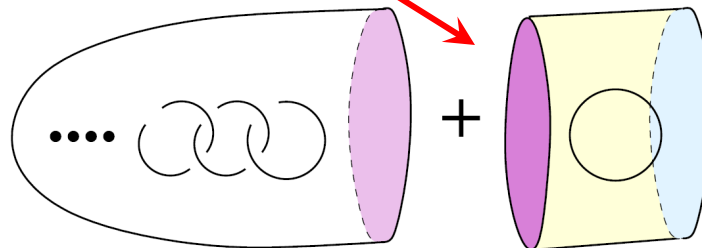
# First New Results



$$\chi_{\rho}^G = \sum_{\rho'} C_{\rho}^{\rho'} \chi_{\rho'}^H$$

flat connection on  $L(k,1)$

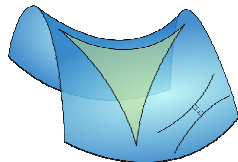
flat connection on  $L(k+1,1)$



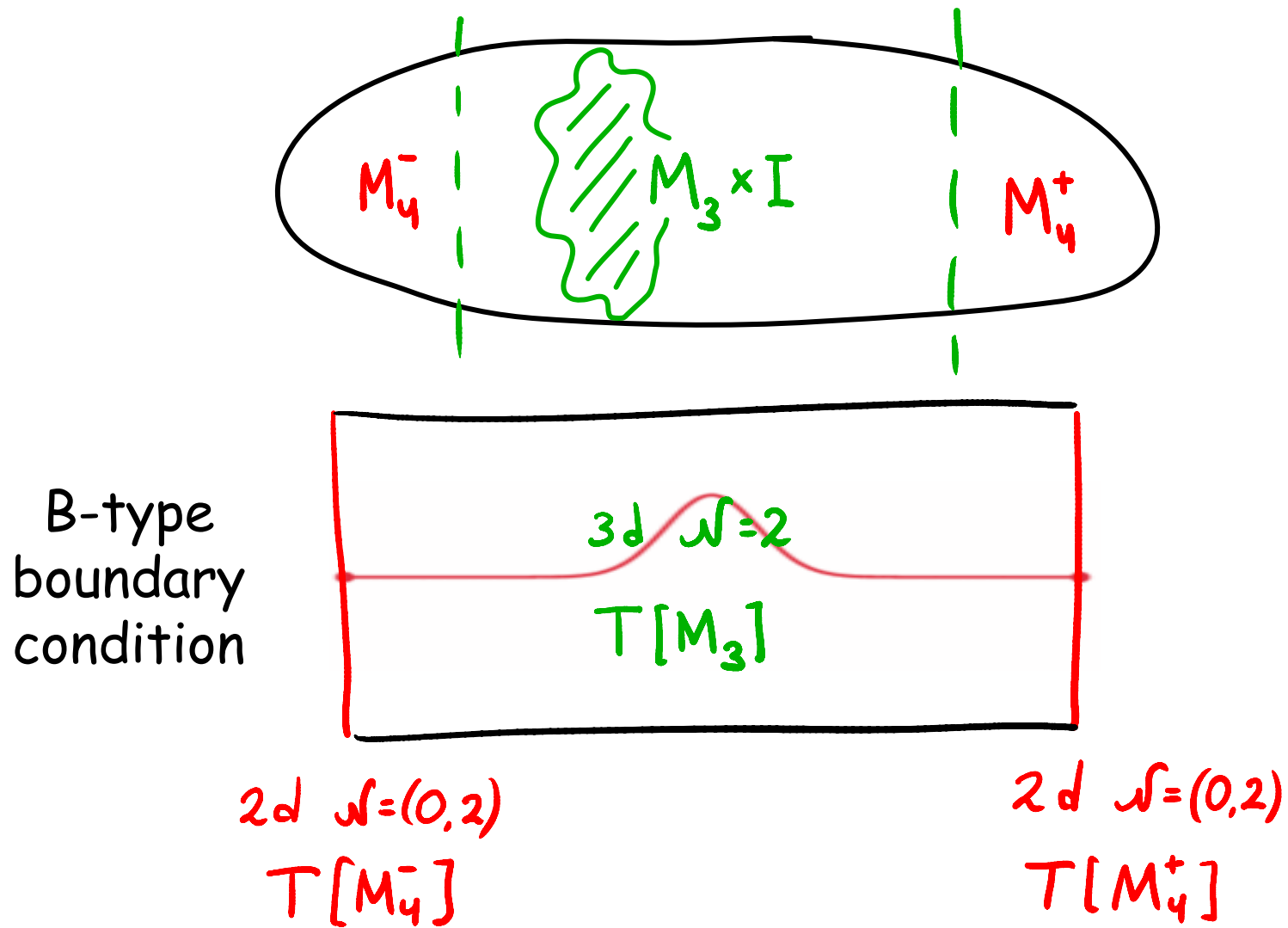
# New 4-manifold invariants from 2d $\mathcal{N} = (0,2)$ theories

$$M_4 \xrightarrow{\text{red arrow}} T[M_4; G] \xrightarrow{\text{red arrow}} Z_{T[M_4; G]} = \text{4-manifold invariant}$$

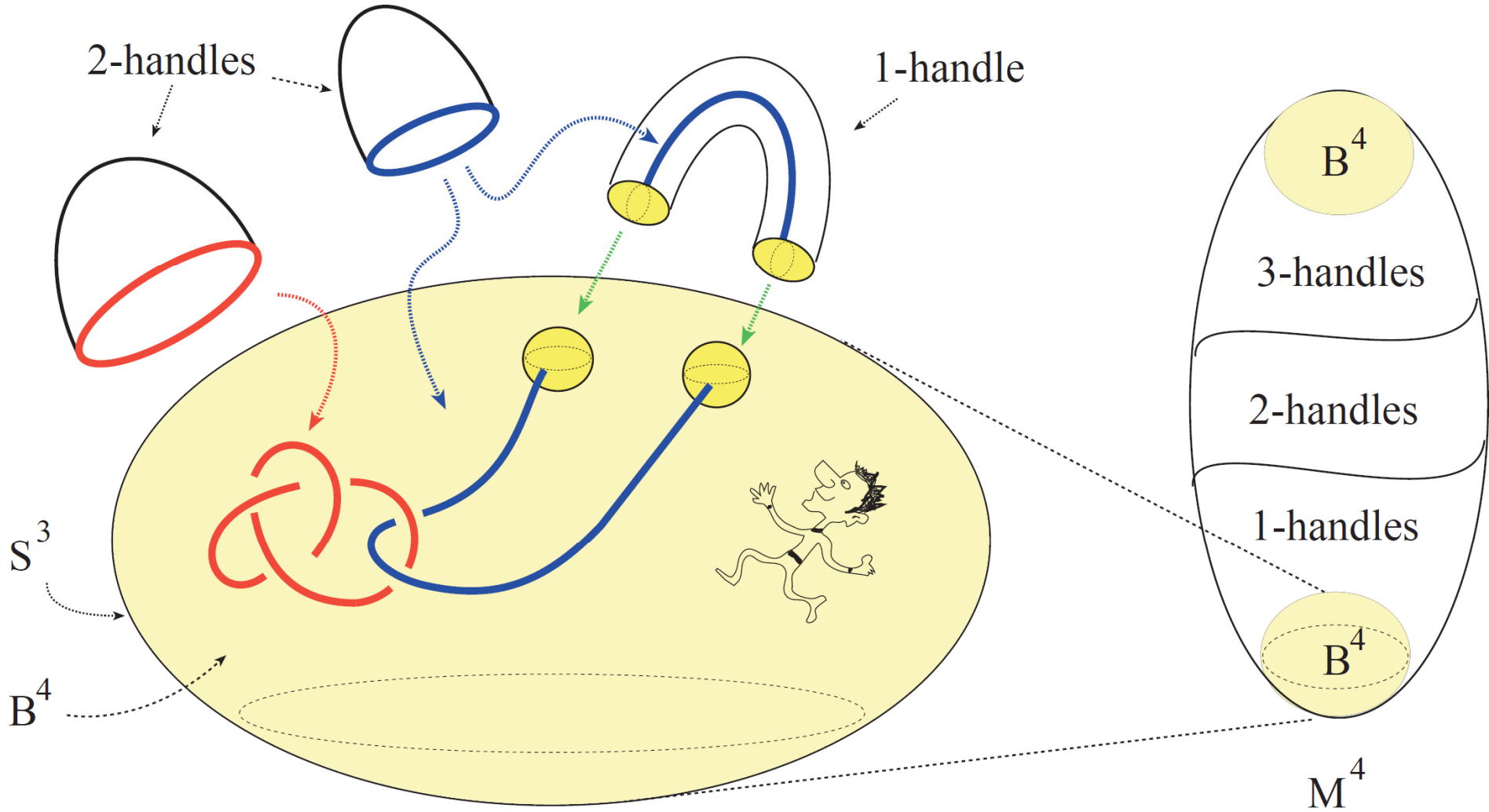
- equivariant elliptic genus =  $\chi(\mathcal{M}_{\text{inst}})$
- Q-cohomology = Donaldson invariants
- moduli space of marginal couplings, etc.



# Gluing

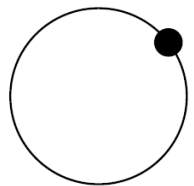


# Building blocks

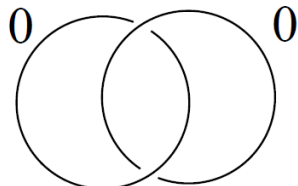


S. Akbulut, 2012

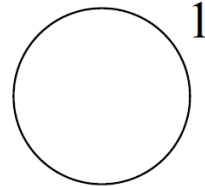
# Kirby diagrams



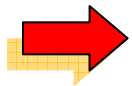
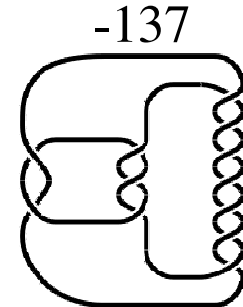
$S^1 \times S^3$



$S^2 \times S^2$



$\mathbb{C}P^2$



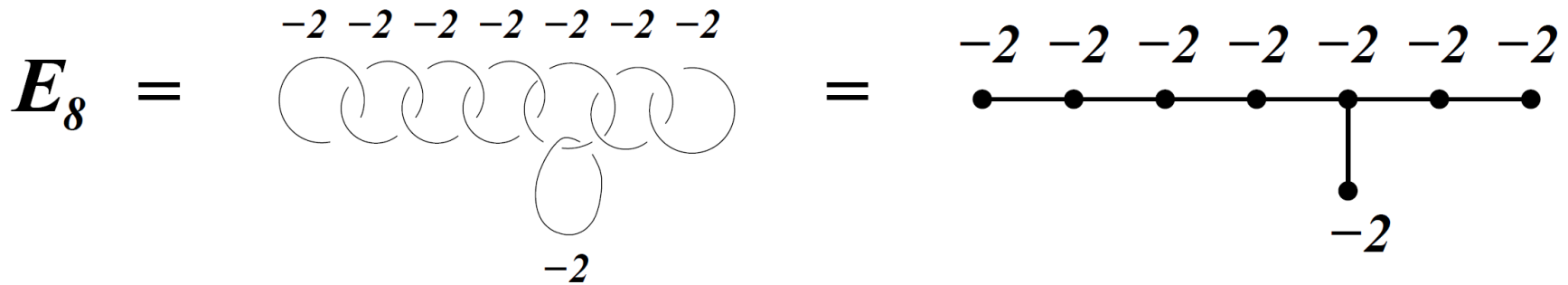
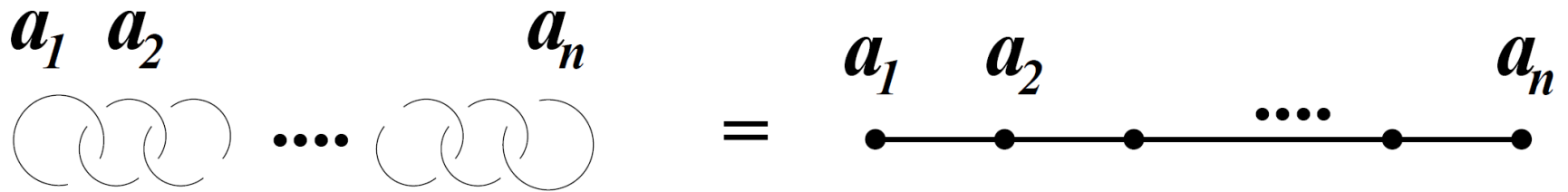
$$M_4 : K_1^{a_1} K_2^{a_2} \dots K_n^{a_n}$$

Intersection form on  $H_2(M_4; \mathbb{Z})$ :

$$Q_{ij} = \begin{cases} \text{lk}(K_i, K_j), & \text{if } i \neq j \\ a_i, & \text{if } i = j \end{cases}$$

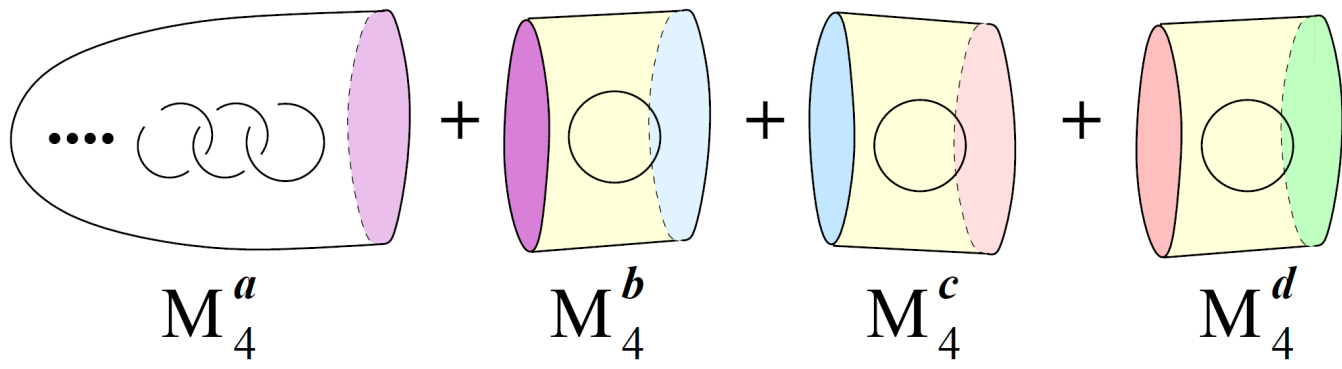


# Plumbing graphs

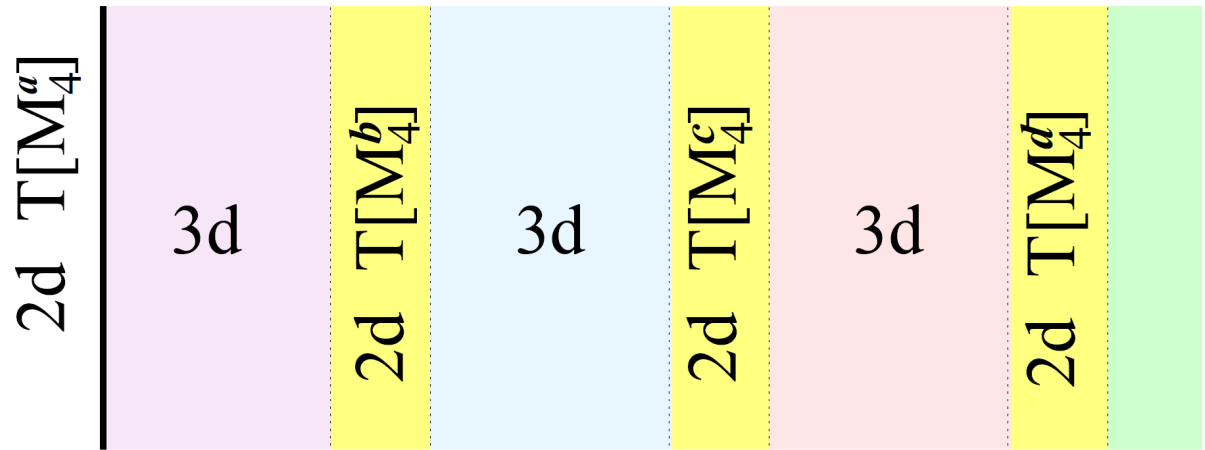
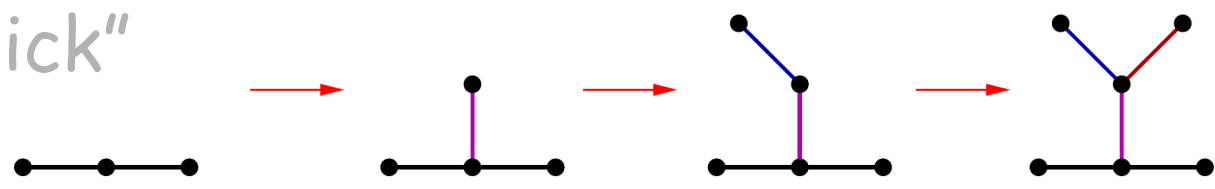


does not always work:

4-manifold bounded by a 3-torus



"Norman trick"



$$\chi_\rho^G = \sum_{\rho'} C_\rho^{\rho'} \chi_{\rho'}^H$$

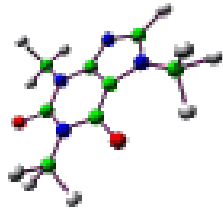


coset branching function

<b>4-manifold <math>M_4</math></b>	<b>2d <math>(0, 2)</math> theory <math>T[M_4]</math></b>
handle slides	dualities of $T[M_4]$
boundary conditions	vacua of $T[M_3]$
3d Kirby calculus	dualities of $T[M_3]$
cobordism from $M_3^-$ to $M_3^+$	domain wall (interface) between $T[M_3^-]$ and $T[M_3^+]$
gluing	fusion
Vafa-Witten partition function	flavored (equivariant) elliptic genus
$Z_{VW}$ (cobordism)	branching function
instanton number	$L_0$
embedded surfaces	chiral operators
Donaldson polynomials	chiral ring relations

# Quiver Chern-Simons theory

vertex  $\bullet$   $\longleftrightarrow$  **U(1)** Chern-Simons at level  $a$



$$S = \frac{a}{4\pi} \int (A \wedge dA + \dots)$$



edge



$$S = \frac{1}{2\pi} \int (A_i \wedge dA_j + \dots)$$

cf. [D.Belov, G.Moore]  
 [A.Kapustin, N.Saulina]  
 [J.Fuchs, C.Schweigert, A.Valentino]

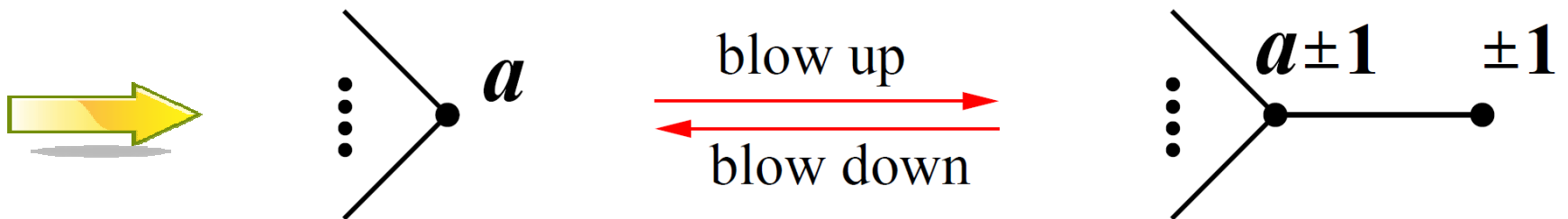
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# Quiver Chern-Simons theory

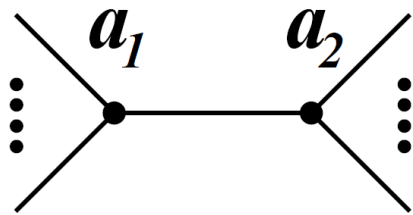
$$\begin{array}{c} \vdots \\ \diagup \\ \bullet \\ \diagdown \end{array} \begin{array}{c} \mathbf{a \pm 1} \\ \text{---} \\ \bullet \\ \text{---} \\ \mathbf{\pm 1} \end{array} = \frac{1}{4\pi} \int \left( \pm A \wedge dA + 2B \wedge dA + (a \pm 1)B \wedge dB + \dots \right)$$

integrate out  $A$   $= \frac{1}{4\pi} \int \left( \pm B \wedge dB \mp 2B \wedge dB + (a \pm 1)B \wedge dB + \dots \right)$

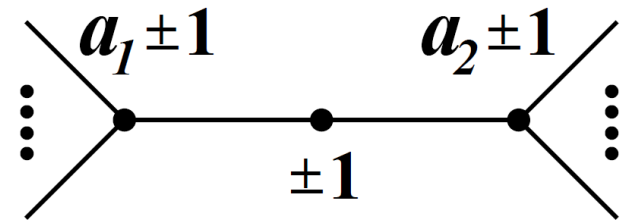
$$\begin{array}{c} \vdots \\ \diagup \\ \bullet \\ \diagdown \end{array} \mathbf{a} = \frac{1}{4\pi} \int (aB \wedge dB + \dots)$$



# 3d Kirby moves

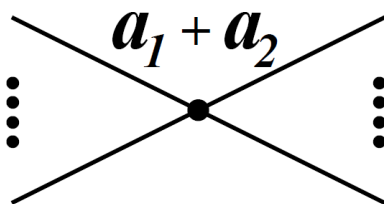
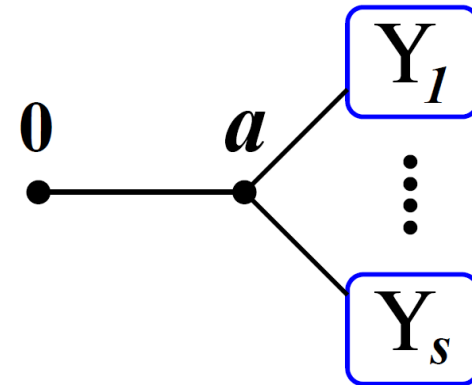


blow up  
 $\rightleftarrows$   
 blow down

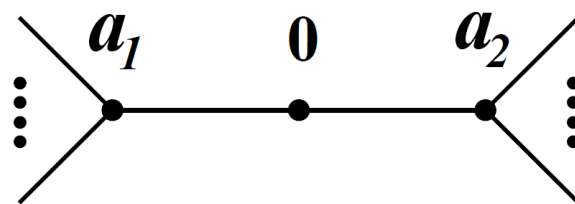


(disjoint union)

$\rightleftarrows$



$\rightleftarrows$



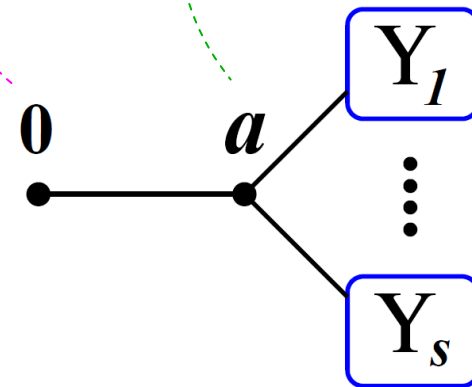
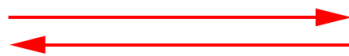
# 3d Kirby moves

$$\mathcal{L} = \frac{1}{4\pi} \int (2A \wedge dB + aB \wedge dB + \dots)$$

$A$  is Lagrange multiplier

$$\boxed{Y_1} + \dots + \boxed{Y_s}$$

(disjoint union)

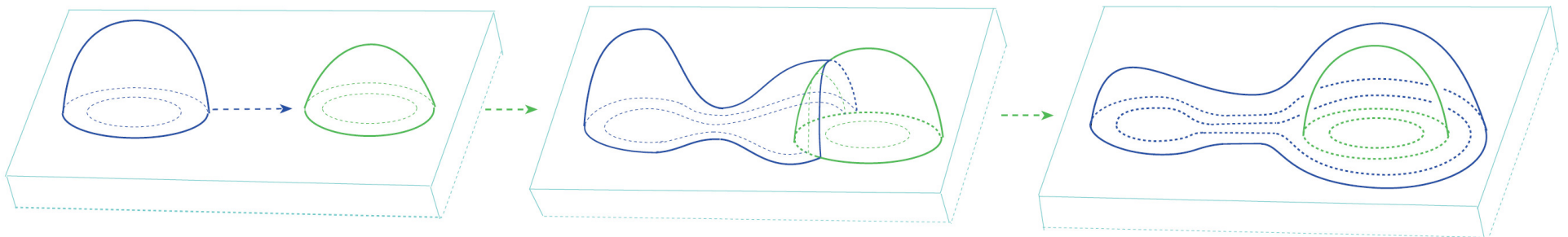


Integrating out  $A$  makes  $B$  pure gauge  
and removes all its Chern-Simons couplings

# 4d Kirby moves

➔ identity for instanton partition function (= equivariant elliptic genus):

$$(q; q)_\infty \int dz \frac{1}{\theta(z^{-1}w)} \prod_{i=1}^{N_f} \theta(q^{\frac{1}{2}} z x_i) = \prod_{i=1}^{N_f} \theta(q^{\frac{1}{2}} x_i w)$$





# 4d Kirby moves

➔ identity for instanton partition function (= equivariant elliptic genus):

$$(q; q)_\infty \int dz \frac{1}{\theta(z^{-1}w)} \prod_{i=1}^{N_f} \theta(q^{\frac{1}{2}} z x_i) = \prod_{i=1}^{N_f} \theta(q^{\frac{1}{2}} x_i w)$$

2d  $\mathcal{N} = (0, 2)$  SQED

	$\Phi$	$\Psi_{i=1, \dots, N_f}$	$\cong$
$U(1)_{\text{gauge}}$	-1	+1	
$U(1)_{\text{flavor}}$	+1	0	

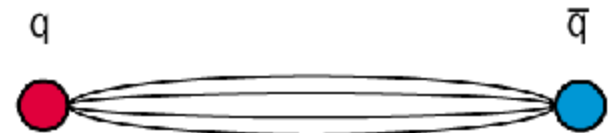
$\left\{ \begin{array}{l} N_f \text{ Fermi multiplets } \Psi'_{i=1, \dots, N_f} \\ \text{with charge } +1 \text{ under } U(1)_{\text{flavor}} \end{array} \right\}$

2d (0,2) "twisted superpotential"

$$\tilde{\mathcal{J}} = -\frac{i}{8\pi} (N_f - 1) \log(\Phi)$$

"mesons"

$$\Psi'_i = \Phi \Psi_i$$



# New 2d $\mathcal{N} = (0,2)$ dualities from 4-manifolds

2d  $\mathcal{N} = (0, 2)$  SQCD

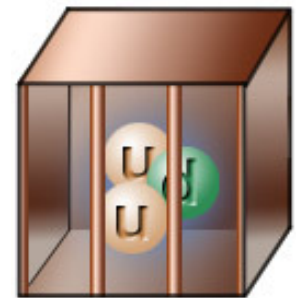
	$\Phi$	$\Psi$	$P$	$\Gamma$
$U(N_c)_{\text{gauge}}$	$\square$	$\overline{\square}$	$\overline{\square}$	$\mathbf{1}$
$U(N_b)_{\text{flavor}}$	$\overline{\square}$	$\mathbf{1}$	$\mathbf{1}$	$\square$
$U(N_f)_{\text{flavor}}$	$\mathbf{1}$	$\square$	$\mathbf{1}$	$\mathbf{1}$
$U(N_f - N_b + 2N_c)$	$\mathbf{1}$	$\mathbf{1}$	$\square$	$\overline{\square}$

$\cong$

dual 2d  $\mathcal{N} = (0, 2)$  SQCD

	$\Phi'$	$\Psi'$	$P'$	$\Gamma'$
$U(N_b - N_c)_{\text{gauge}}$	$\square$	$\overline{\square}$	$\overline{\square}$	$\mathbf{1}$
$U(N_b)$	$\overline{\square}$	$\mathbf{1}$	$\mathbf{1}$	$\square$
$U(N_f - N_b + 2N_c)$	$\mathbf{1}$	$\square$	$\mathbf{1}$	$\mathbf{1}$
$U(N_f)$	$\mathbf{1}$	$\mathbf{1}$	$\square$	$\overline{\square}$

- 2d (0,2) analogue of Seiberg duality
- the **very first** non-abelian 2d (0,2) duality



**MATH**



**The End**

**PHYSICS**

