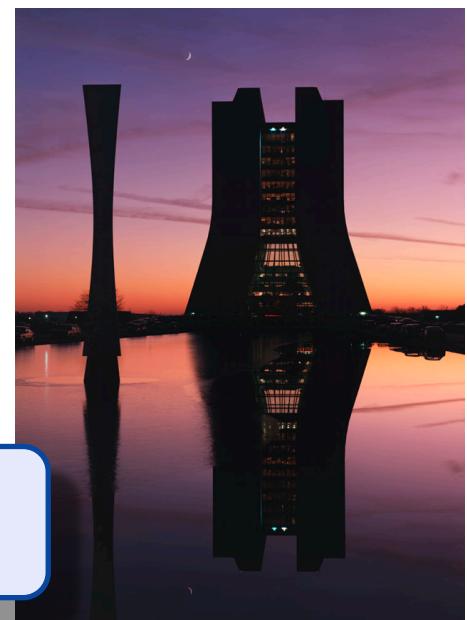
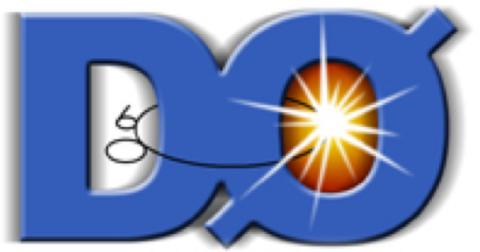




# Precision measurement of the mass of the top quark in l+jets final states using $9.7 \text{ fb}^{-1}$ of DØ data

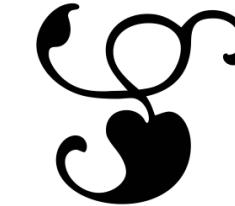


*Oleg Brandt (U Heidelberg)  
on behalf of the DØ collaboration*



- Motivation
- Selection
- The matrix element technique
- Calibration of method
- Result in data
- Data/MC comparisons
- Systematic uncertainties
  - Systematic uncertainties
    - Systematic uncertainties
    - Systematic uncertainties
      - Systematic uncertainties
- Conclusion
- 2014 Tevatron combination
- Anything to be learned?

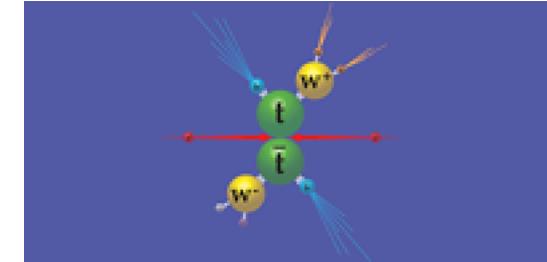
PRL 113, 032002 (2014)



Precision Measurement  
of the Top Quark Mass  
in *Lepton + Jets* Final States  
(D0 Collaboration)



Synopsis:  
Top Quark Mass Gets an Update

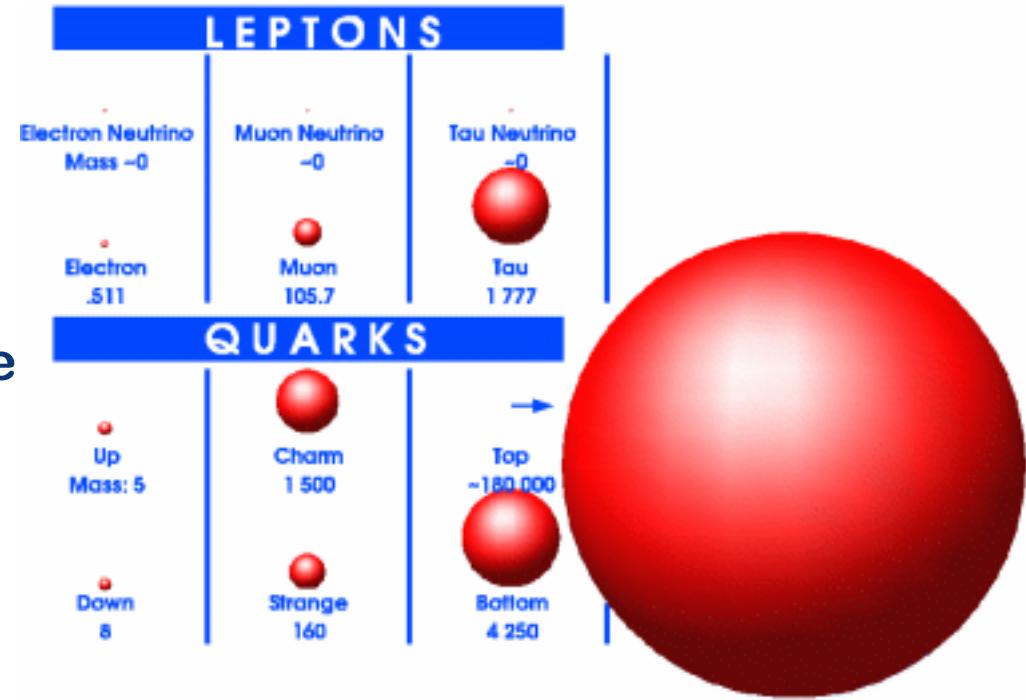




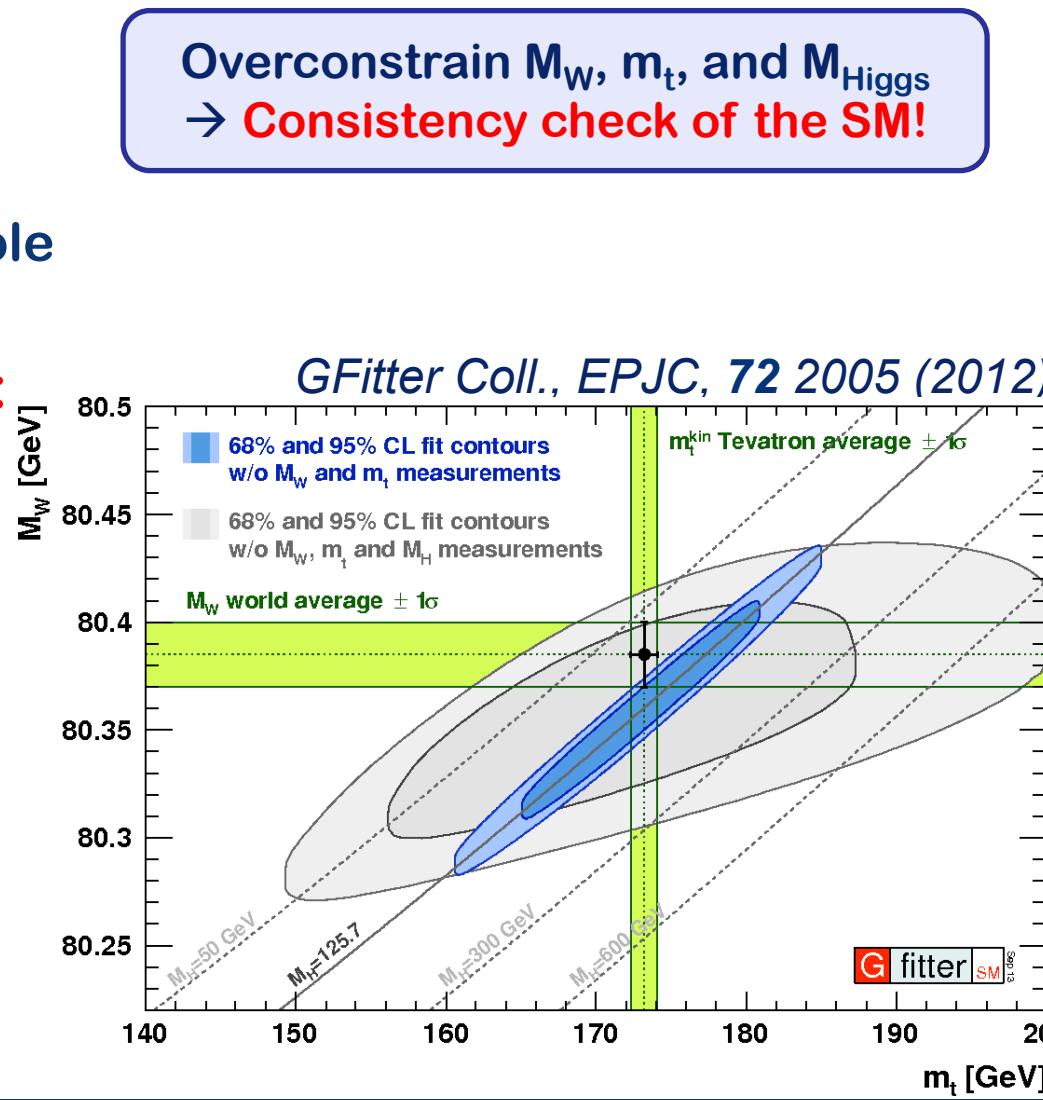
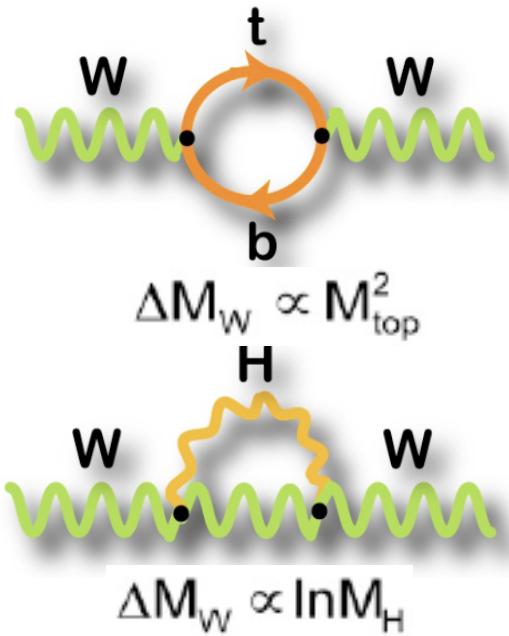
Events/(20 GeV/c<sup>2</sup>)



- The top quark is special:
  - It is the heaviest quark of the SM!
    - Why is it so heavy?
    - Does it play a special role in EWSB?



- The top quark is special:
  - It is the heaviest quark of the SM!
  - Why is it so heavy?
  - Does it play a special role in EWSB?
- $M_W$  related to  $m_t$  &  $M_{\text{Higgs}}$ :

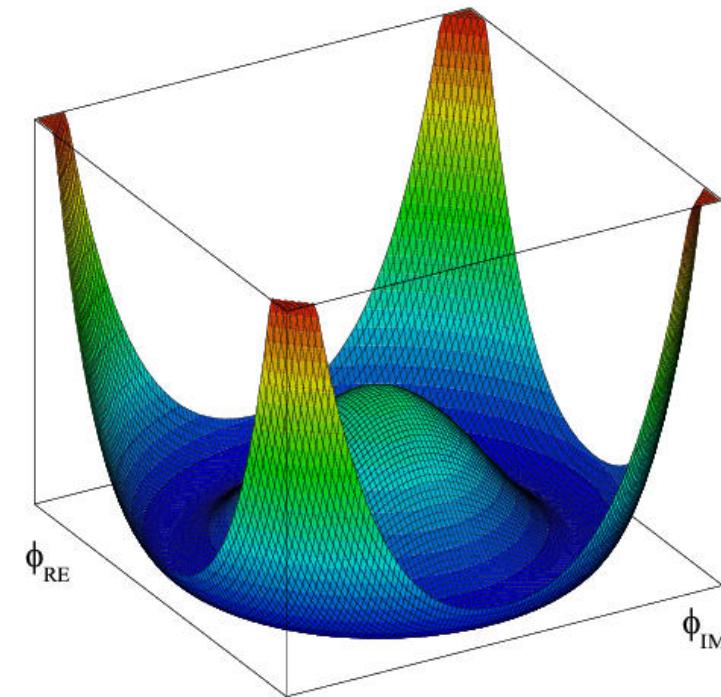


- If this is not enough, the top quark mass is a fundamental parameter of the SM
- The fate of our Universe depends on  $m_t$ !

- Consider the Higgs Lagrangian:

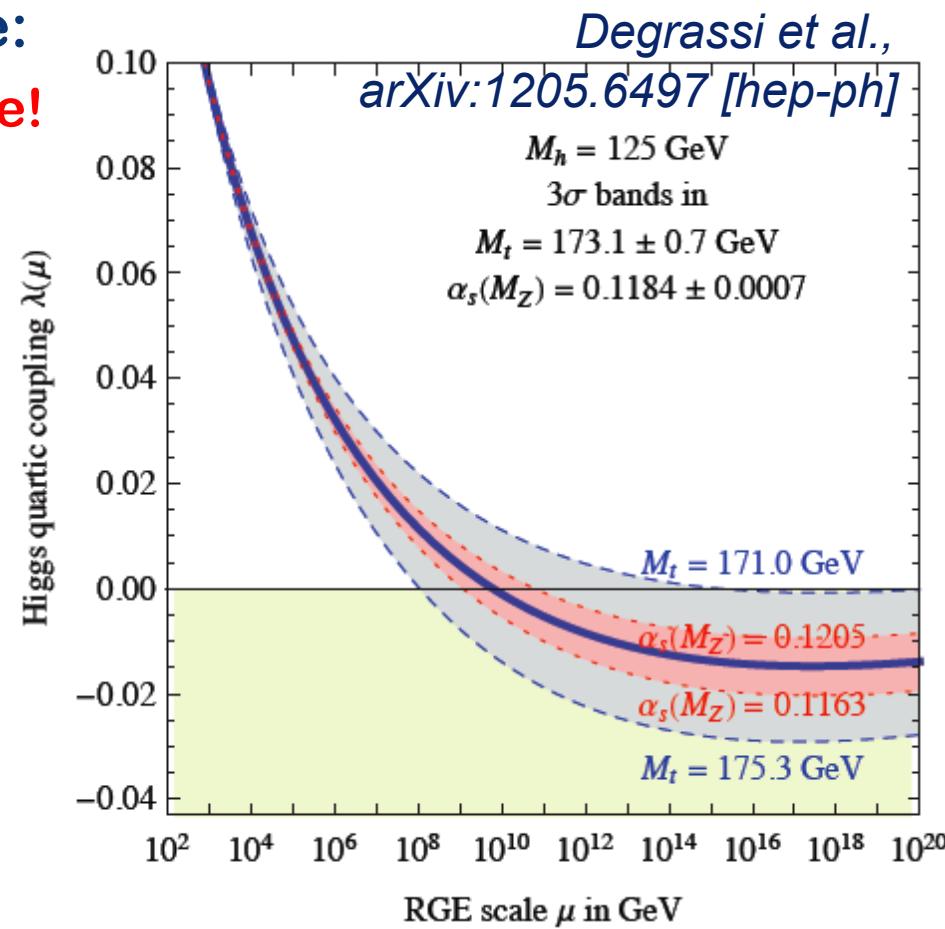
$$\mathcal{L}_H = \left| \left( \partial_\mu - igW_\mu^a \tau^a - i\frac{g'}{2} B_\mu \right) \phi \right|^2 + \mu^2 \phi^\dagger \phi - \lambda(\phi^\dagger \phi)^2,$$

- The quartic Higgs self-coupling term  $\lambda(\phi^\dagger \phi)^2$  is responsible for the mexican-hat shape of the potential
  - This works only if  $\lambda$  is positive...



- $\lambda$  receives **radiative corrections** from all particles of the SM, mostly from the top quark!
  - We can evolve these corrections using running group equation to Planck scale:
    - $\lambda$  should remain positive!

With the current world's best values for  $m_t$  and  $m_{\text{Higgs}}$ :  
 → Our Universe is only metastable!



The calculation includes NNLO effects,  
 RG equation at NNNLO



Will the top quark save our Universe?

Stay tuned!

We'll know more in  $10^{12}$  years...



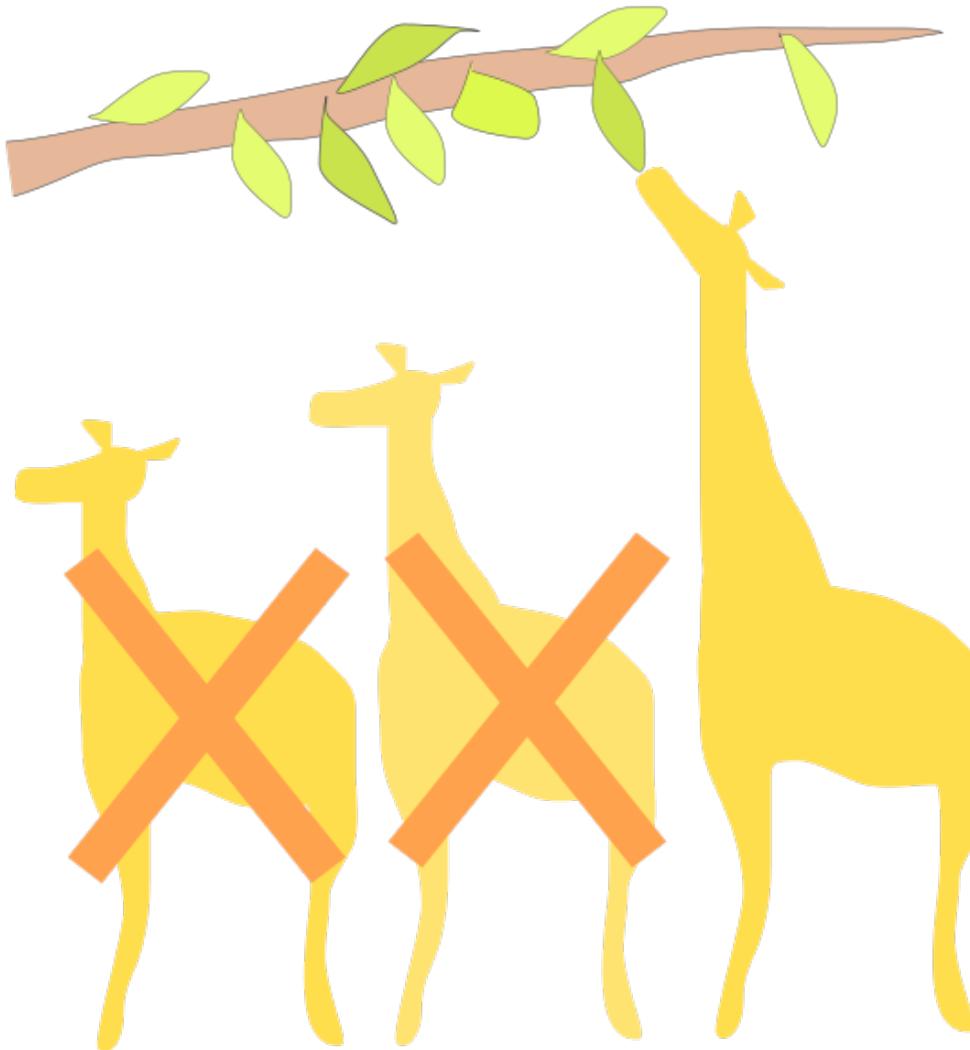
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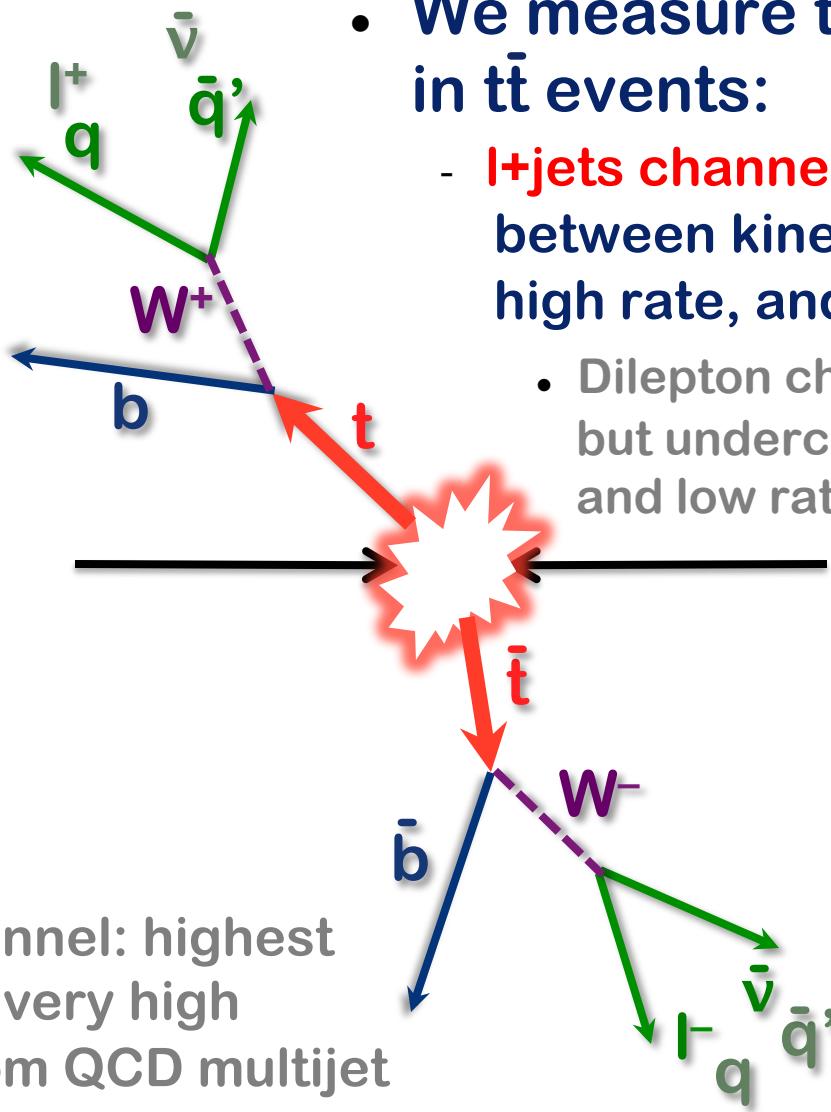
Stay tuned!

We will measure  $m_t$ ,  $M_{\text{Higgs}}$ , and  $\alpha_s$  with higher precision!



# Selection





- We measure the top mass in  $t\bar{t}$  events:
  - **I+jets channel:** good compromise between kinematic reconstruction, high rate, and backgrounds
    - Dilepton channel: low backgrounds, but underconstrained kinematics and low rate
- All-hadronic channel: highest branching ratio, very high backgrounds from QCD multijet production



- Use full  $9.7 \text{ fb}^{-1}$  of integrated luminosity & require:
  - Exactly one tight isolated electron or muon with:
    - $p_T > 20 \text{ GeV}$ ,  $|\eta_{\text{muon}}| < 2$ ,  $|\eta_{\text{electron}}| < 1.1$
  - Exactly four jets with cone parameter  $R=0.5$  and:
    - $p_T > 20 \text{ GeV}$ , leading jet  $p_T > 40 \text{ GeV}$
  - Reject multijet events with
    - $\cancel{p}_T > 20 \text{ GeV}$
    - Topological cuts
  - One or more b-tagged jet ( $\varepsilon_b \approx 65\%$ ,  $\varepsilon_{\text{light}} \approx 10\%$ ) to further reject backgrounds



- We split our data sample into **four epochs**:
  - Run IIa ( $1.1 \text{ fb}^{-1}$ )
  - Run IIb1 ( $1.2 \text{ fb}^{-1}$ )
  - Run IIb2 ( $3.0 \text{ fb}^{-1}$ )
  - Run IIb3 ( $4.4 \text{ fb}^{-1}$ )
- This is done to accurately take into account the **detector response variations** due to factors like:
  - Time dependence due to instantaneous luminosity
    - E.g. track reconstruction efficiency, b-tagging
  - Varying detector configuration
    - E.g. addition of Layer 0 to the silicon tracker after Run IIa
- We use **dedicated MC simulations** to model the detector response in **each of those epochs**



- **Simulations are done with:**
  - **t̄t signal:** alpgen+pythia (0,1,2 light partons)
  - **W+jets:** alpgen+pythia (0,1,...,5 light partons)
  - **W+cc, W+bb:** alpgen+pythia (0,1,2,3 light partons)
  - **Multijet events:** from data
    - Use “matrix method” to predict its yield [1]
    - Use events with inverted lepton isolation requirements to model the shape

[1] DØ Coll., submitted to PRD, arXiv:1401.5785 [hep-ex].



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  - Multijet events: from data
    - Use “matrix method” to predict its yield [1]
    - Use events with inverted lepton isolation requirements to model the shape
- To make sure we are on the same page:
  - We measure the so-called **MC mass:**
    - In our case  $m_t$  as implemented in alpgen
  - Close, but not necessarily identical to the pole mass

[1] DØ Coll., submitted to PRD, arXiv:1401.5785 [hep-ex].



- Sample composition assuming  $\sigma_{t\bar{t}} = 7.24 \text{ pb}$  [1]:

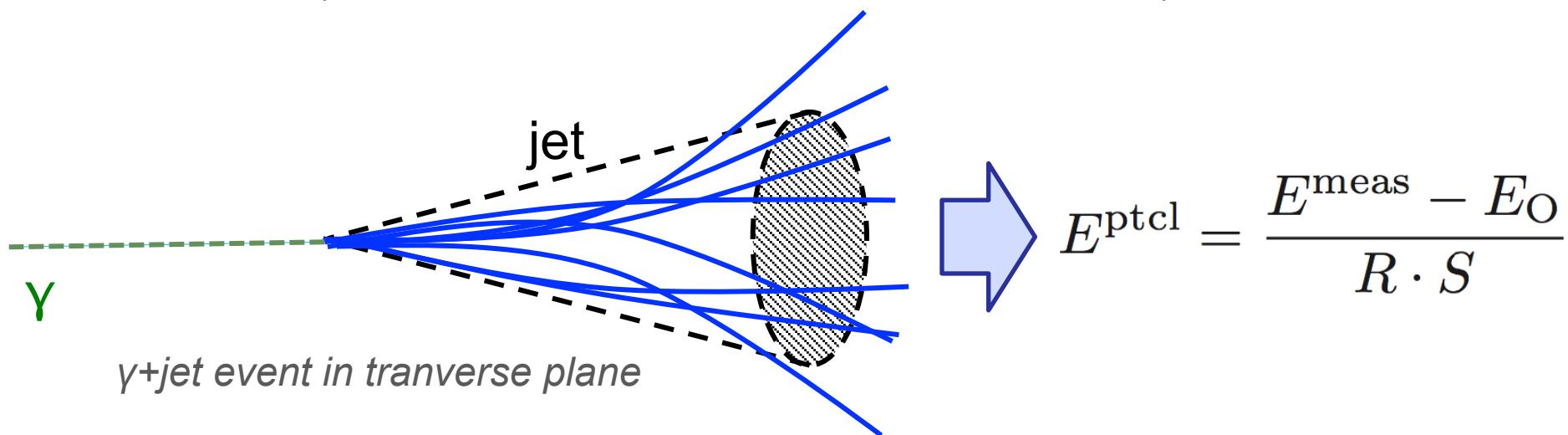
Contribution	$e + \text{jets}$			$\mu + \text{jets}$		
Data	1502.00	$\pm$	38.76	1286.00	$\pm$	35.86
$t\bar{t}$	918.11	$\pm$	3.63	824.88	$\pm$	3.48
$W + \text{jets}$	77.85	$\pm$	2.13	101.03	$\pm$	2.93
$W + \text{HF}$	125.98	$\pm$	2.12	162.21	$\pm$	2.81
Multijet	144.41	$\pm$	24.19	48.17	$\pm$	16.11
Other backgrounds	97.75	$\pm$	0.51	79.24	$\pm$	0.94
Expected	1364.10	$\pm$	24.65	1215.53	$\pm$	17.00

- We expect a signal fraction (f) of
  - 67% ( $e+\text{jets}$ )
  - 68% ( $\mu+\text{jets}$ )
- (will show data/simulations comparison plots later in the talk)

[1] Czakon et al, *PRL* **109**, 132001 (2012).



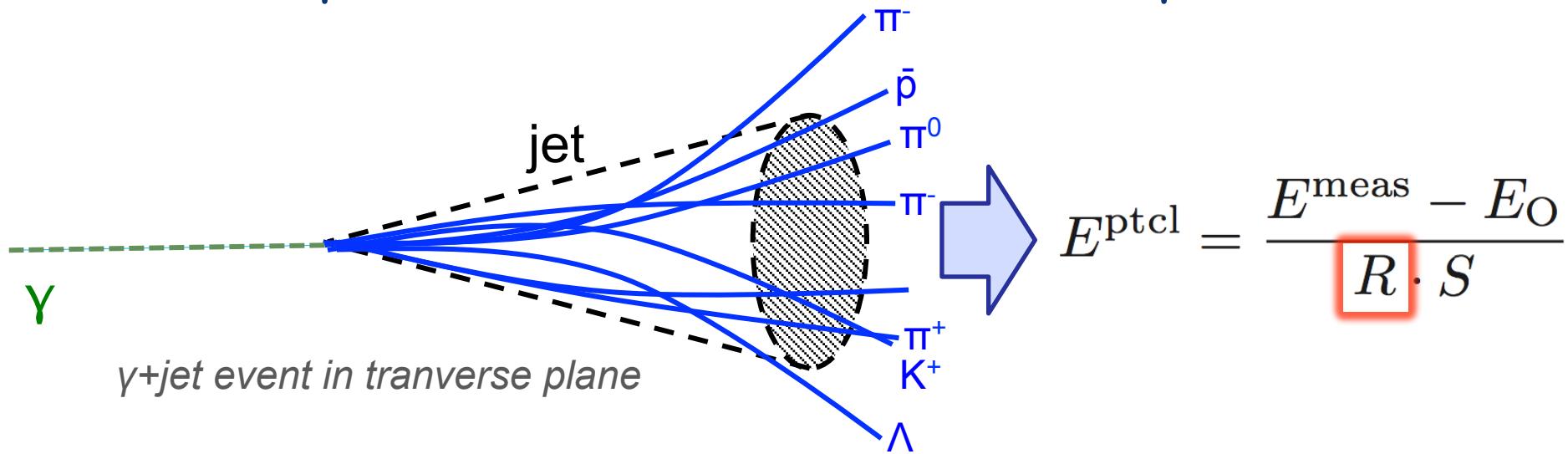
- We **calibrate jet energies at detector level to particle level** (in data and MC)
- Calibration procedure in a nutshell:
  - Calibrate EM energy scale with  $Z \rightarrow e^+e^-$
  - Correct energy scale for electrons to that of photons
  - Use  **$\gamma + \text{jet}$  events to calibrate major components of JES**
    - Expect momentum balance in transverse plane



- Use  $\gamma + \text{jet}$  and dijet events to extend calibration in  $pT, \eta$

# Jet energy scale (JES) calibration

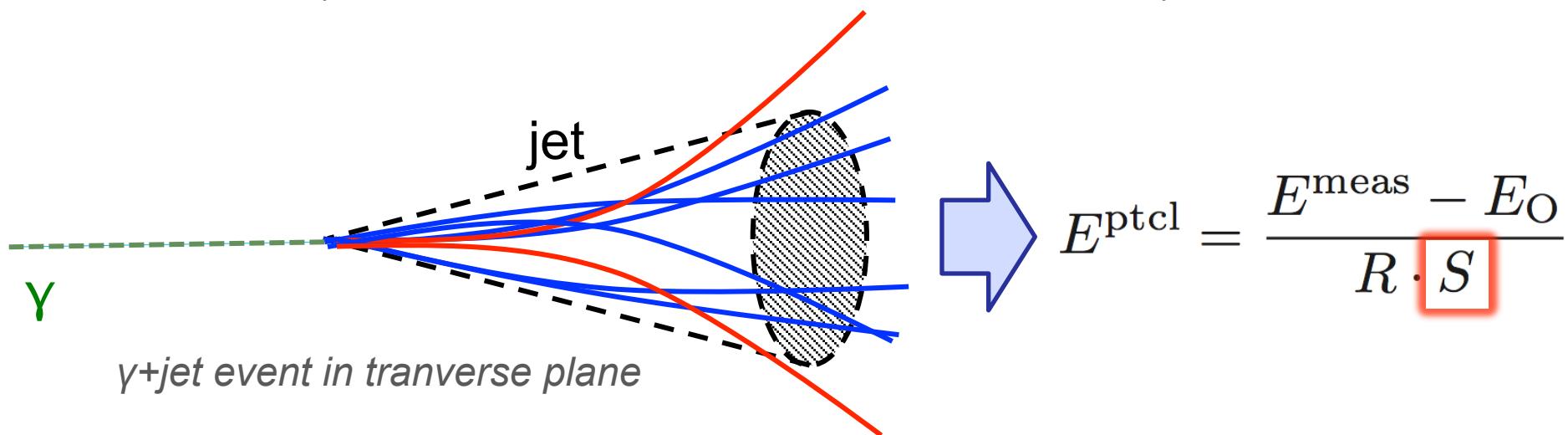
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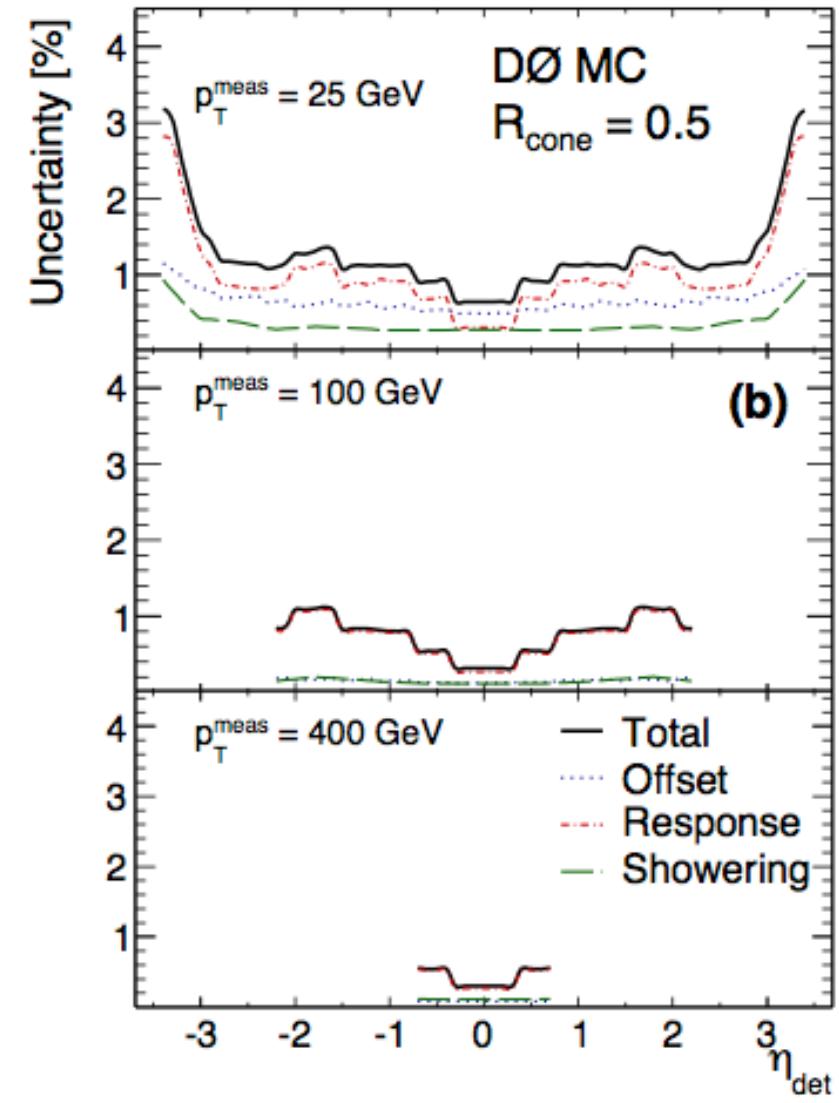
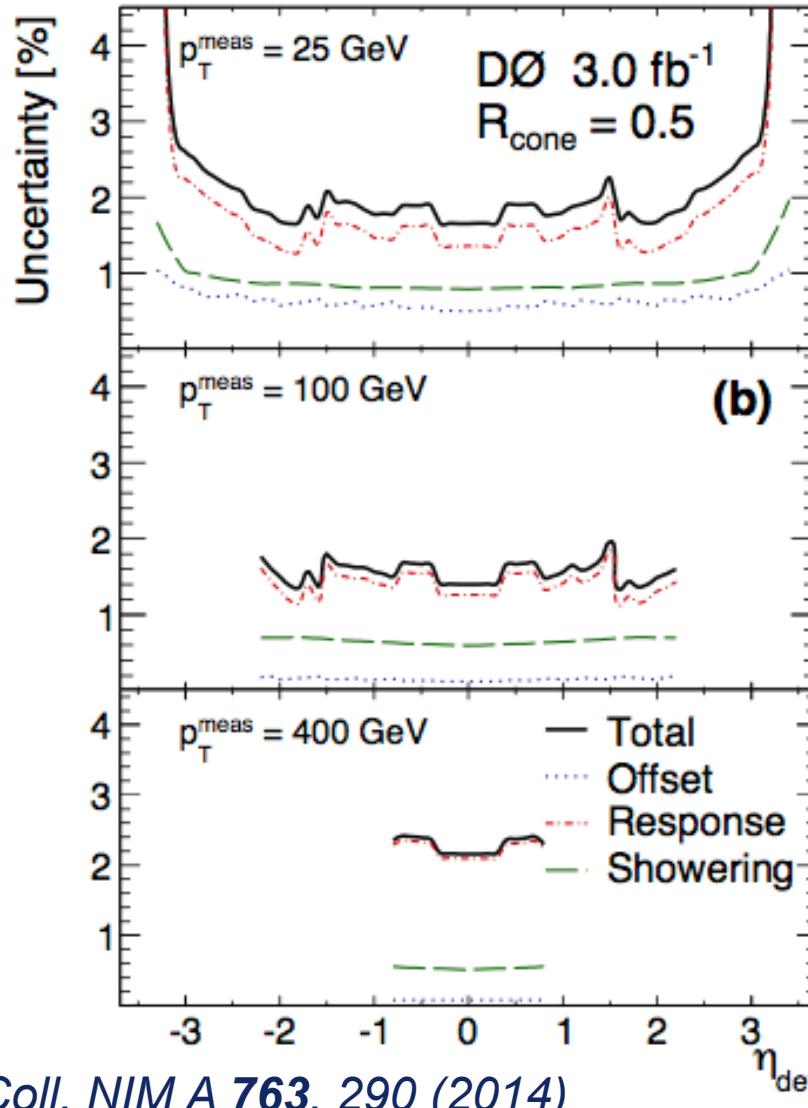
$$E^{\text{ptcl}} = \frac{E^{\text{meas}} - E_0}{R \cdot S}$$

*$\gamma + \text{jet}$  event in transverse plane*

- Use  $\gamma + \text{jet}$  and dijet events to extend calibration in  $pT, \eta$

# Refined jet energy scale calibration

- We use the new jet energy scale (JES) calibration:



$D\bar{\Omega}$  Coll, NIM A 763, 290 (2014)

Figures are representative of all Run II



- Apply dedicated corrections for:

- **u, d, c, s** quark jets
- **b** quark jets
- **gluon** jets

- The correction is given by:

$$F_{\text{corr}} = \frac{1}{\langle F \rangle_{\gamma+\text{jet}}} \cdot \frac{\sum_i E_i \cdot R_i^{\text{data}}}{\sum_i E_i \cdot R_i^{\text{MC}}}$$

*preserves default JES  
per constructionem*

- Derive single particle responses  $R_i$  in data/MC for:

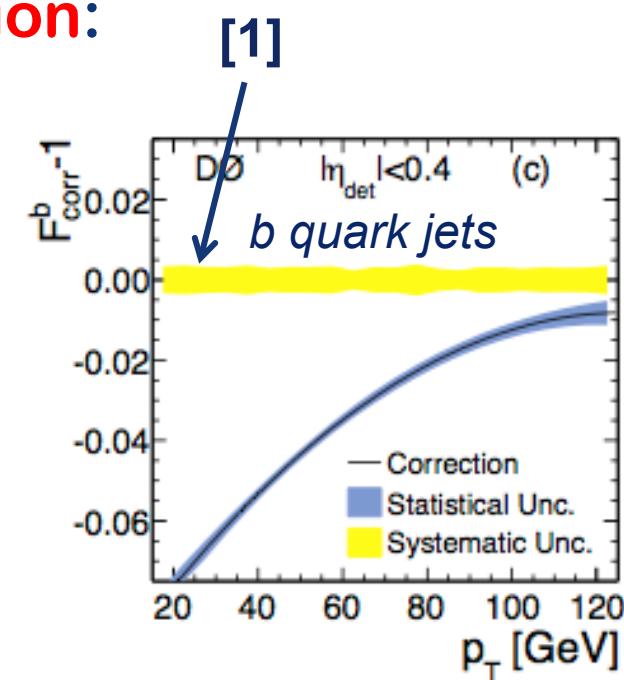
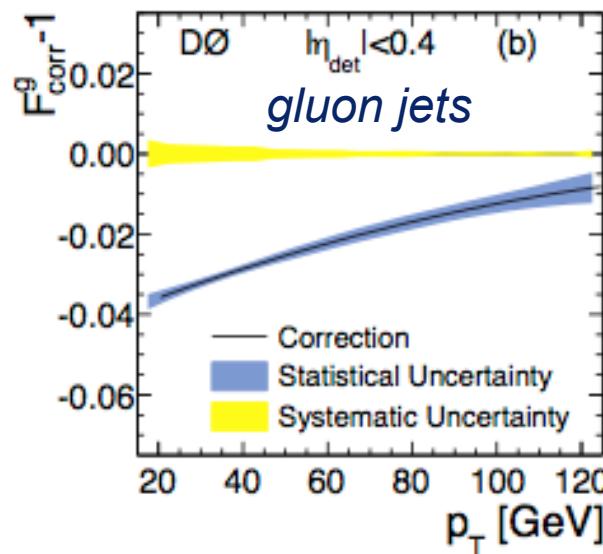
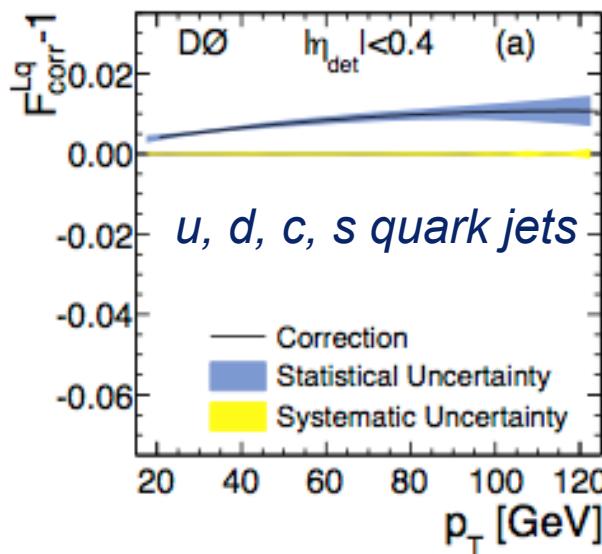
- $\gamma, e^\pm, \mu^\pm, \pi^\pm, K^\pm, K_0^S, K_0^L, p^\pm, n$  and  $\Lambda$ 
  - Use  $\gamma+\text{jet}$  sample  $\rightarrow$  **quark-dominated**
  - Use **dijet** sample  $\rightarrow$  **gluon-dominated**
- Take flavour composition of the samples from MC [1]

DØ Coll, NIM A 763, 290 (2014)

[1] X-check: uncertainty from this assumption covered by systematic uncertainty assigned

# Refined jet energy scale calibration

- The final flavour-dependent correction:

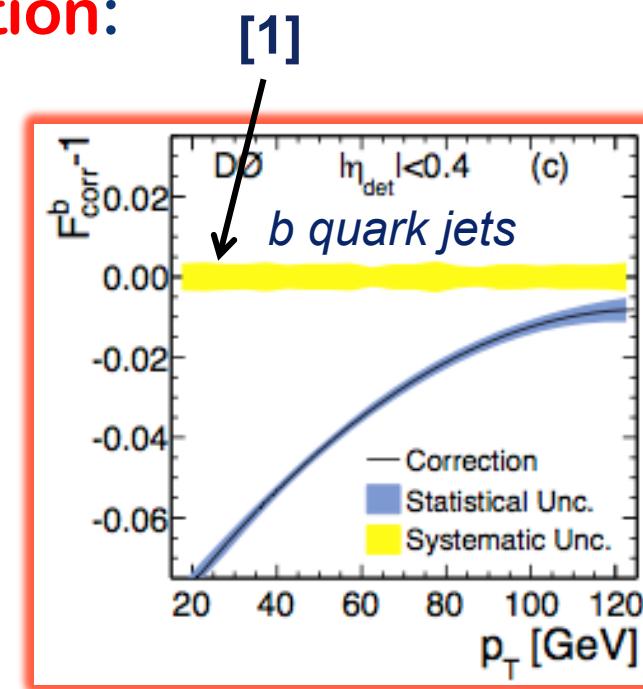
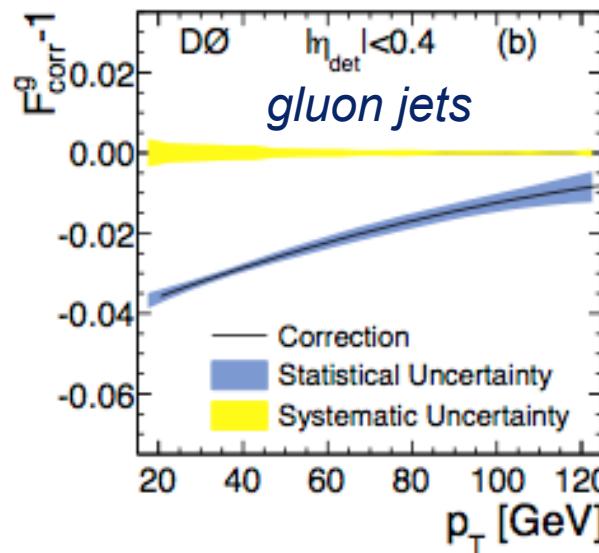
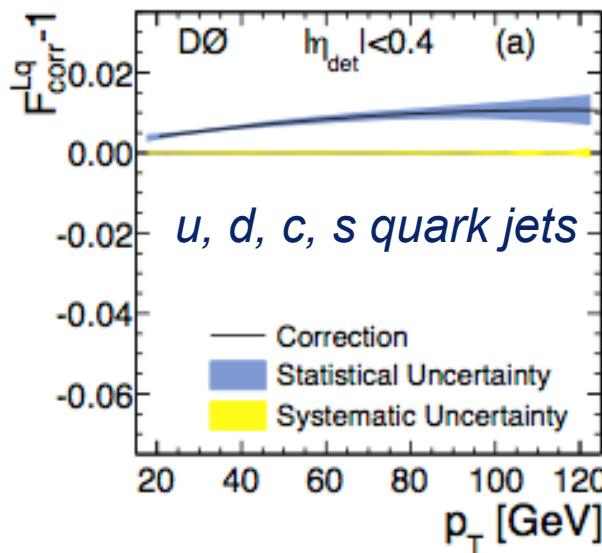


DØ Coll, NIM A 763, 290 (2014)

[1] Dominant source of systematic uncertainty: single particle response shapes in MC

# Refined jet energy scale calibration

- The final flavour-dependent correction:



- The correction accounts for the difference in JES for b quark jets and light quark jets:
  - Substantial reduction of one of the dominant systematic uncertainties!

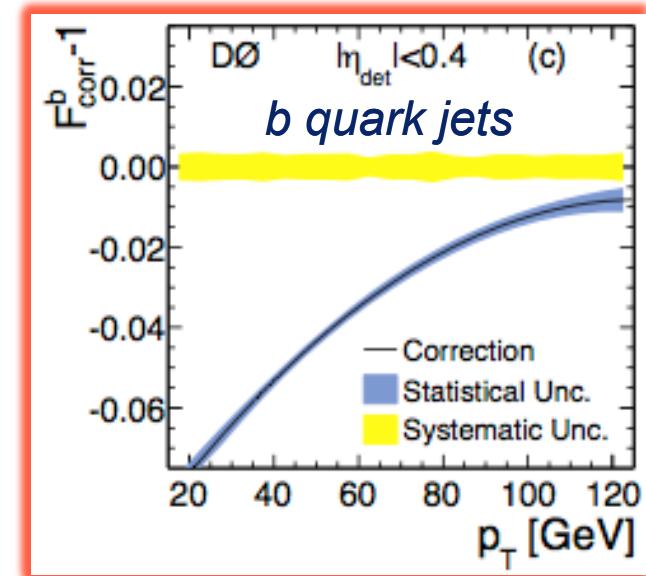
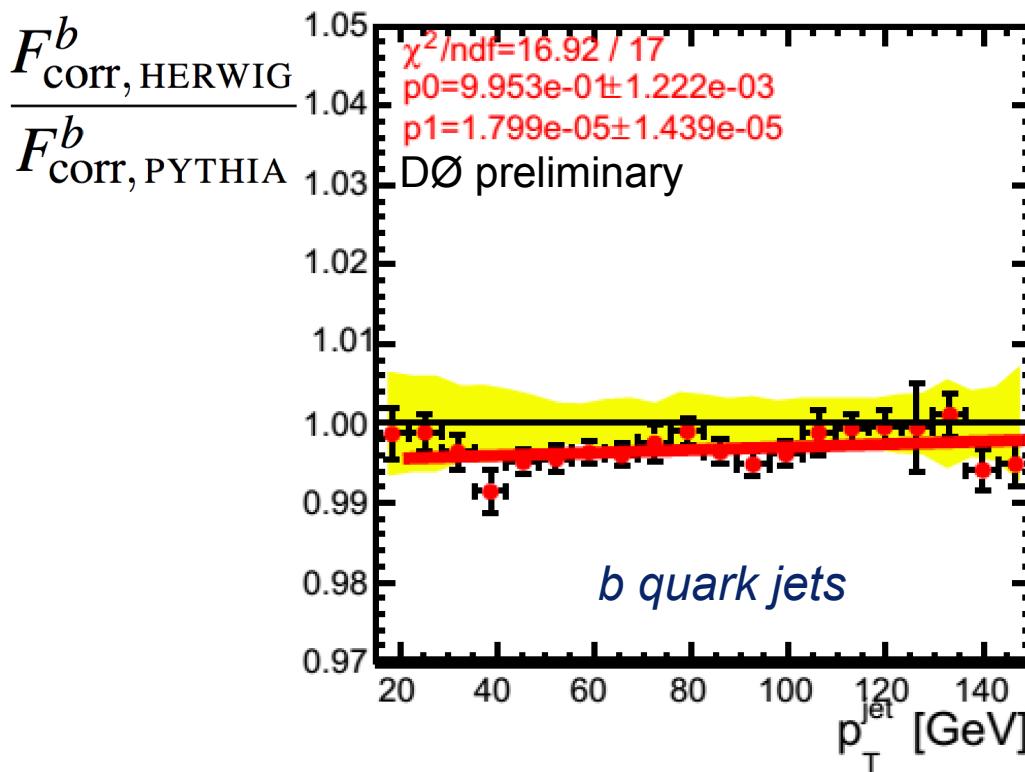
DØ Coll, NIM A 763, 290 (2014)

[1] Dominant source of systematic uncertainty: single particle response shapes in MC

- The final flavour-dependent correction:

## Cross-check #3:

Study impact of shower composition:



Excursion of double-ratio from unity covered by systematic uncertainties

Double-ratio plot preliminary (some weights for efficiencies from ID, trigger etc. missing)



# The matrix element technique



- **Matrix Element (ME) technique:**
  - Calculate the event probability on an event-by-event basis:

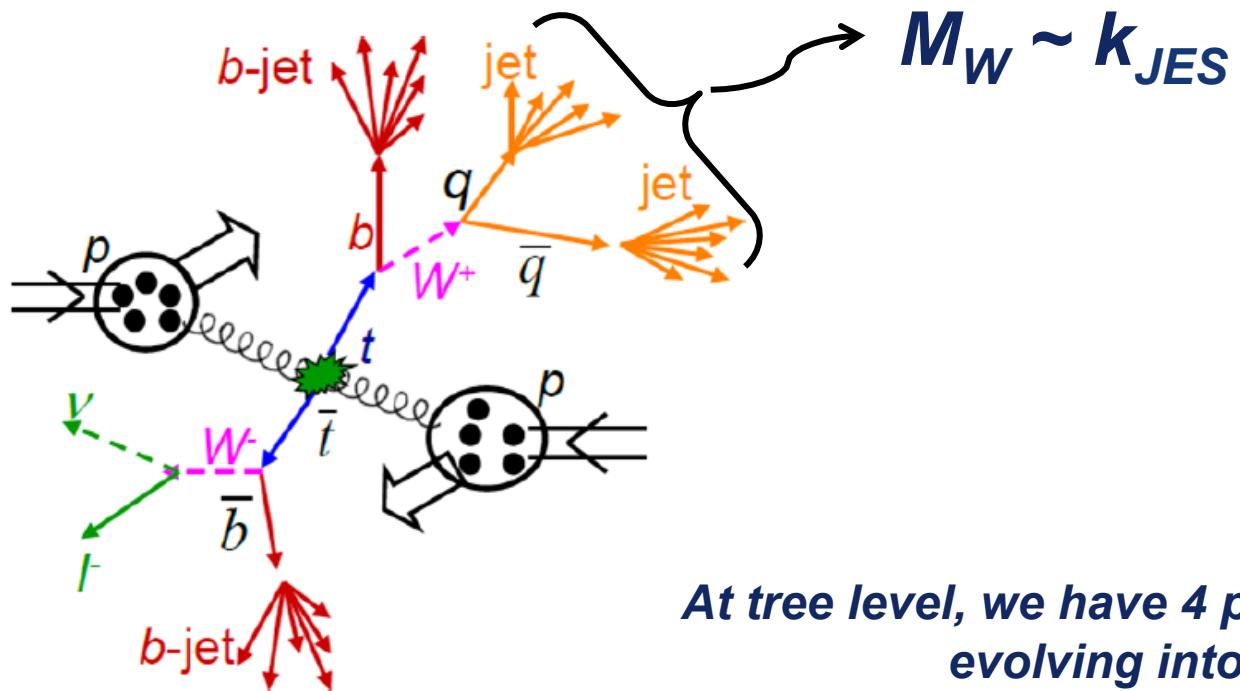
$$P_{\text{evt}}(m_{\text{top}}) \propto f P_{\text{sig}}(m_{\text{top}}) + (1 - f) P_{\text{bgr}}$$

$$P_{\text{sig}}(m_{\text{top}}) \propto \int \dots \underline{d\sigma_{t\bar{t}}(m_{\text{top}})}$$

$$d\sigma_{t\bar{t}} \propto |\mathcal{M}_{t\bar{t}}|^2(m_{\text{top}})$$

- **Advantages:**
  - Use 4-vectors with maximal kinematic and topological information → maximal statistical power
- **Disadvantages:**
  - High computational demand + theory assumptions

- perform an **in-situ calibration of the JES**:
  - Constrain energies of the two jets from W to be consistent with  $M_W$
  - This allows a **simultaneous extraction of  $m_t$  and the overall JES factor  $k_{JES}$ !**



*At tree level, we have 4 partons  
evolving into 4 jets:  
→ 24 jet-parton assignments*



## DØ matrix element technique in l+jets final states

b tagging-based weight to identify relevant jet-parton assignments

$$P_{\text{sig}} = \frac{1}{\sigma_{\text{obs}}^{t\bar{t}}} \sum_{i=1}^{24} w_i \int d\rho dm_1^2 dM_1^2 dm_2^2 dM_2^2 d\rho_\ell dq_1^x dq_1^y dq_2^x dq_2^y$$

**Integration over phase space (10 dim)**

$\sum_{\text{flavors}, \nu} |\mathcal{M}_{t\bar{t}}|^2 \frac{f'(q_1)f'(q_2)}{\sqrt{(\eta_{\alpha\beta} q_1^\alpha q_2^\beta)^2 - m_{q_1}^2 m_{q_2}^2}}$   $\Phi_6 W(x, y; k_{\text{JES}})$

**LO matrix element**  
PRD 53, 4886 (1996)  
PLB 411, 173 (1997)

**Phase space factor**

DØ Coll.,  
PRD 84, 032004 (2011)

Transfer functions (TFs) to map  
parton level quantities  $y$  to reco level quantities  $x$



## DØ matrix element technique in l+jets final states

Normalisation by observed cross section using the same LO ME

$$P_{\text{sig}} = \frac{1}{\sigma_{\text{obs}}^{t\bar{t}}} \sum_{i=1}^{24} w_i \int d\rho dm_1^2 dM_1^2 dm_2^2 dM_2^2 d\rho_\ell dq_1^x dq_1^y dq_2^x dq_2^y$$
$$\sum_{\text{flavors}, \nu} |\mathcal{M}_{t\bar{t}}|^2 \frac{f'(q_1)f'(q_2)}{\sqrt{(\eta_{\alpha\beta} q_1^\alpha q_2^\beta)^2 - m_{q_1}^2 m_{q_2}^2}} \Phi_6 W(x, y; k_{\text{JES}})$$

Sum over all 24 possible jet-parton assignments

Sum over incoming parton flavours and all neutrino  $p_z$  solutions

PDFs for Björken-x and PD for transverse momenta of incoming partons

DØ Coll.,  
PRD 84, 032004 (2011)



- We numerically calculate a 10 dimensional integral using MC integration techniques
  - Identical to the  $3.6 \text{ fb}^{-1}$  result [1] except:
    - Use low-discrepancy sequences for the MC integration
      - Deterministic sequence of points in our 10-dim parameter space providing optimal convergence
    - Factorise the JES factor  $k_{\text{JES}}$  from the ME calculation
      - Include it via the transfer function
    - Reduction of calculation time by  $\text{o}(100)!$

[1]  $D\bar{\nu}$  Coll., PRD 84, 032004 (2011)



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    - Factorise the JES factor  $k_{\text{JES}}$  from the ME calculation
      - Include it via the transfer function
    - Reduction of calculation time by  $\text{o}(100)$
- Increase the size of calibration samples!
  - Typical statistical uncertainty from size of MC samples:
    - $\approx 0.25 \text{ GeV} \rightarrow 0.01\text{--}0.05 \text{ GeV}$

Uncertainty (GeV)

$\pm 0.25$   
 $\pm 0.26$   
 $\pm 0.58$   
 $\pm 0.28$   
 $\pm 0.07$   
 $\pm 0.16$   
 $\pm 0.07$   
 $\pm 0.09$   
 $\pm 0.24$   
 $\pm 0.21$   
 $\pm 0.28$

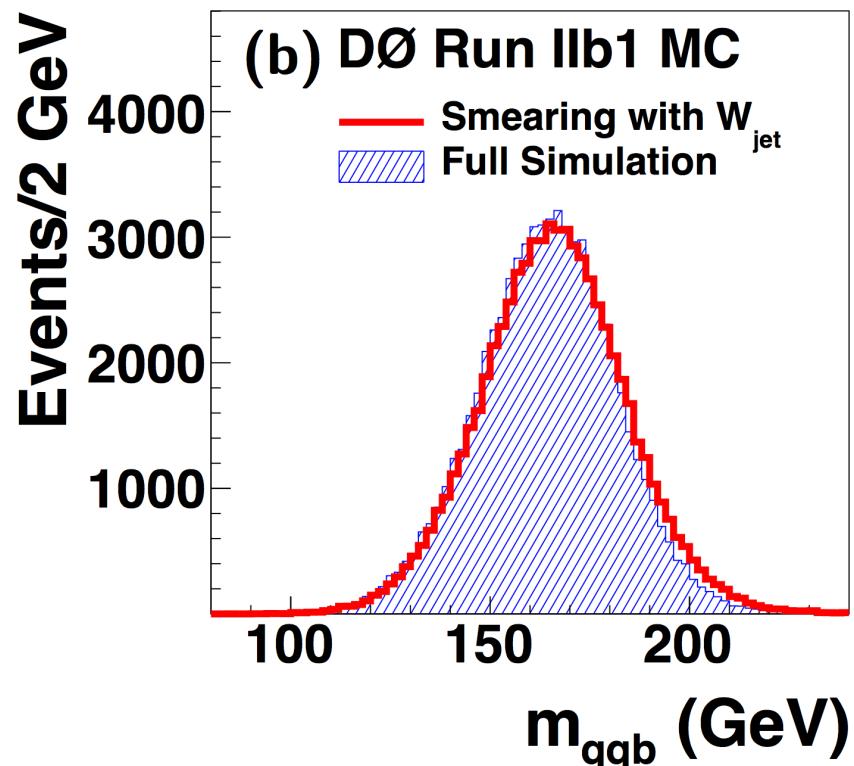
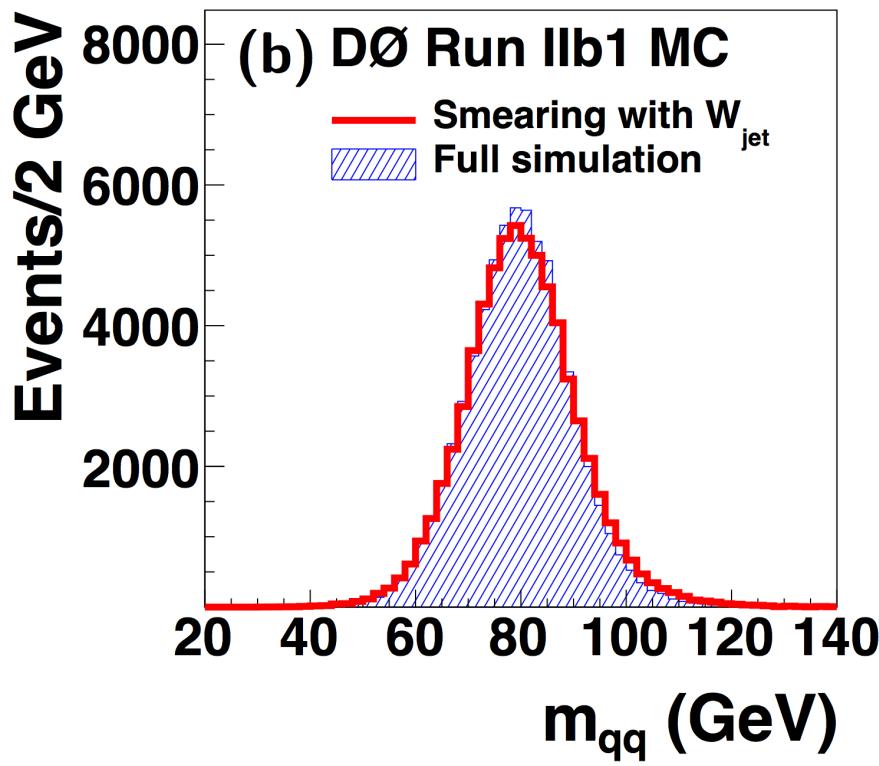
Excerpt from the table of systematic uncertainties of the  $3.6 \text{ fb}^{-1}$  analysis [1]

[1] DØ Coll., PRD 84, 032004 (2011)



- The Transfer Functions  $W(x, y; k_{\text{JES}})$  relate parton-level quantities to detector-level ones
- Parametrise the detector response:
  - For jets, we use the sum of two Gaussians:
- Parametrise jet energies:
  - treat separately: light quark jets, b-tagged jets with soft muon tag, all other b-jets
  - $\times 4 \mid \eta \mid$  regions for each
- Direction of jets and leptons in  $\eta \times \phi$  well-measured:
  - → use  $\delta$ -functions as transfer function!

- Compare parton momenta smeared with transfer functions to jet momenta in full simulation in:
  - Invariant mass of dijet system matched to W boson
  - Invariant mass of trijet system matched to top quark

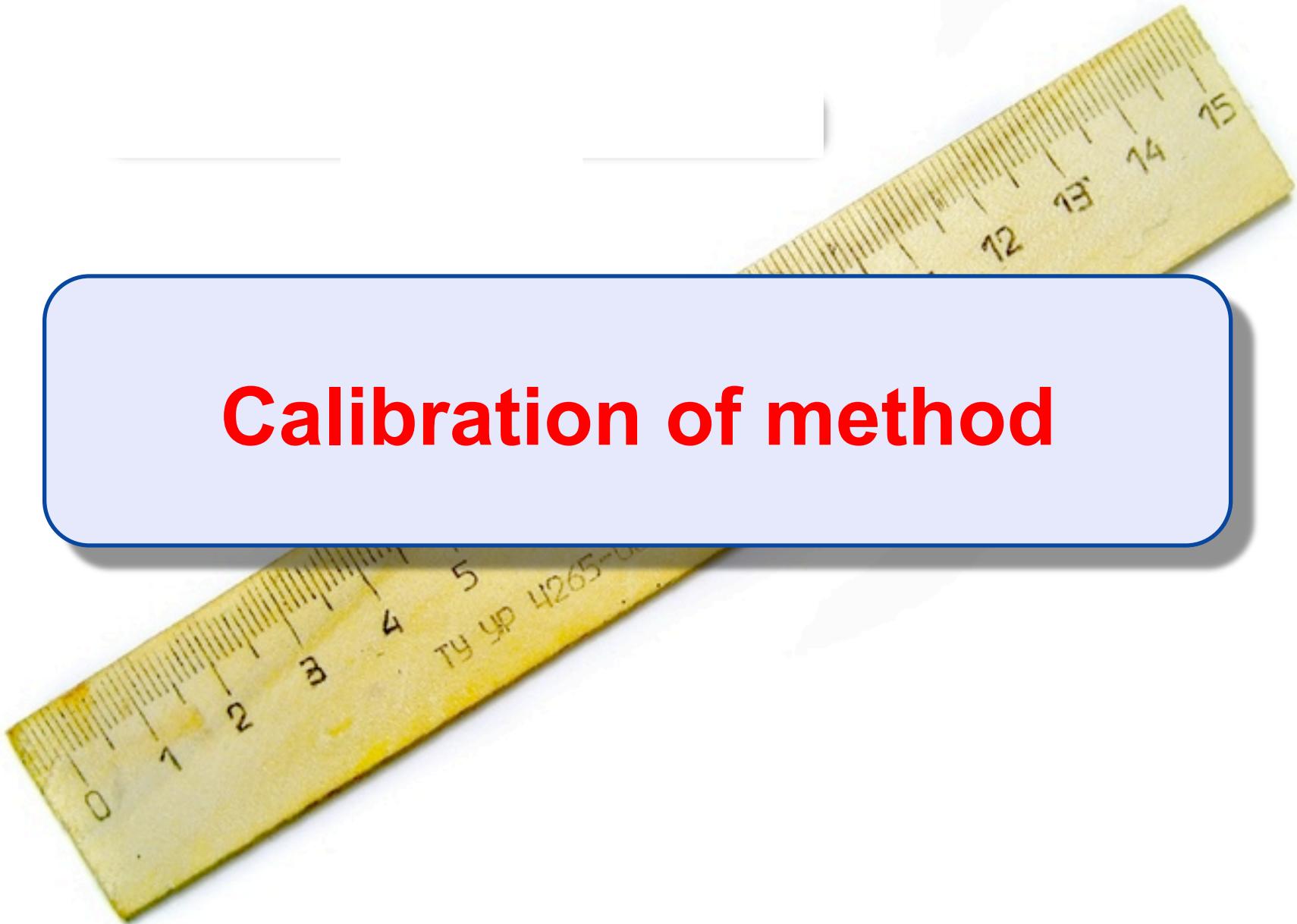


- Very good description of detector response!

Representative MC simulations → cf. backup for all



# Calibration of method

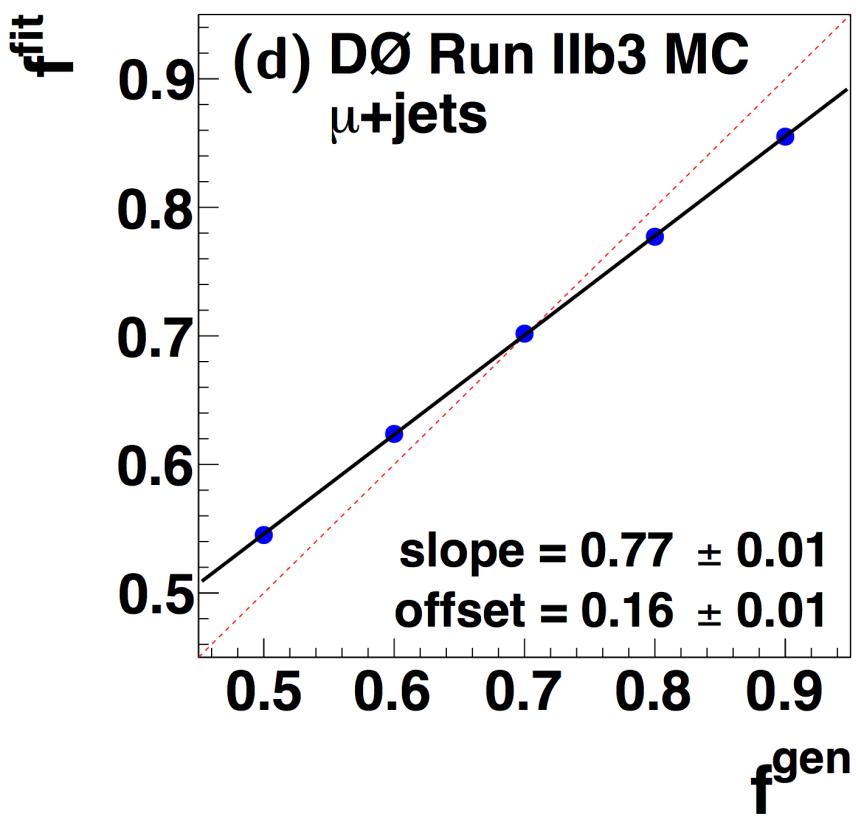
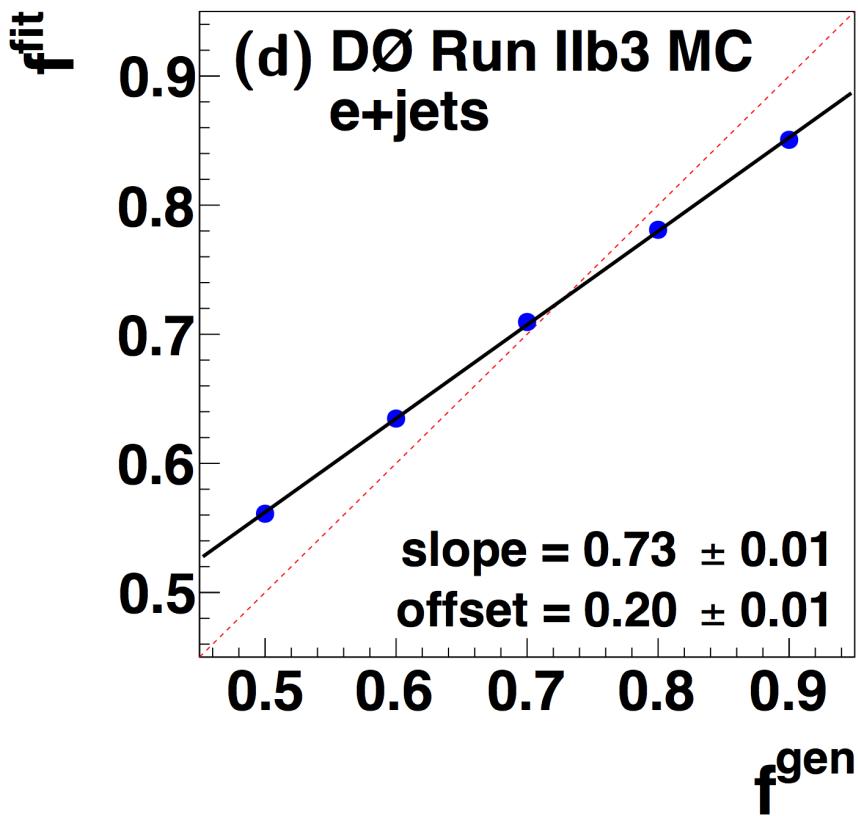




- Calibrate the method with **pseudo-experiments (PE)**
  - Keep in mind we use  $P_{\text{evt}}$  obtained from **first principles** with a LO ME and parametrised detector response
    - → calibration imperative
    - (in template methods this is merely a consistency check)
  - Dedicated calibration for **each epoch & channel**
    - → total of  $4 \times 2 = 8$  calibrations
  - Each PE consists of  $N_{\text{data}}$
  - PEs include:
    - **W+jets background** (dominant, adjusted according to  $f$ )
    - **MJ background** (11% and 4% for e and mu+jets)
    - (Other backgrounds contribute a total of ~ 5%)
  - Construct PEs according to signal fraction ( $f$ ) measured from data
  - 1000 PEs performed at each calibration point

# Extracting signal fraction from data

- We extract the signal fraction from data
  - Integrate  $L(f, m_t, k_{JES})$  over  $m_t$  and  $k_{JES} \rightarrow$  maximise in  $f$
  - Calibrate method response in  $f$ :



(Perfect method response to  $f$  is achieved at parton level with events generated according to the LO matrix element for the  $2 \rightarrow 2$  process)

Representative MC simulations, see backup for all



- We measure (after calibration):

Epoch	Final state	Signal fraction	$\sigma_{t\bar{t}}$ (pb)
Run IIa	$e + \text{jets}$	0.72	8.9
	$\mu + \text{jets}$	0.65	7.8
Run IIb1	$e + \text{jets}$	0.77	7.6
	$\mu + \text{jets}$	0.66	6.8
Run IIb2	$e + \text{jets}$	0.68	7.8
	$\mu + \text{jets}$	0.66	7.5
Run IIb3	$e + \text{jets}$	0.56	7.6
	$\mu + \text{jets}$	0.75	8.0
Run II	$e + \text{jets}$	0.63	7.8
	$\mu + \text{jets}$	0.70	7.6

- Values in good agreement with  $\sigma_{t\bar{t}} = 7.78^{+0.77}_{-0.64} \text{ pb}$  [1]

Typical statistical+calibration uncertainty on signal fraction: 1%, on  $\sigma_{t\bar{t}}$ : about 0.1 pb  
[1] DØ Collaboration, Phys. Rev. D 84, 012008 (2011).

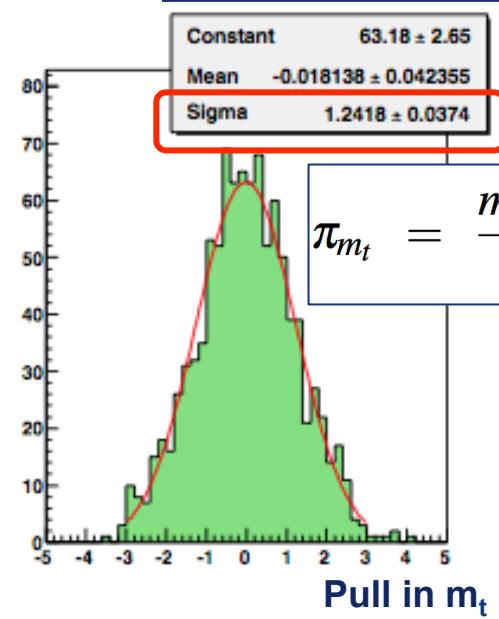
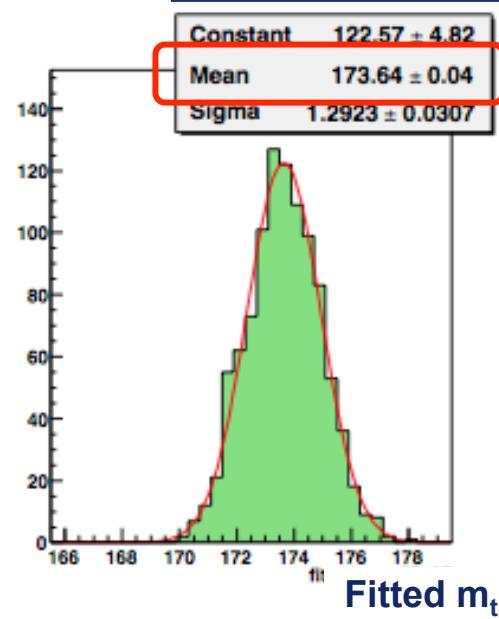
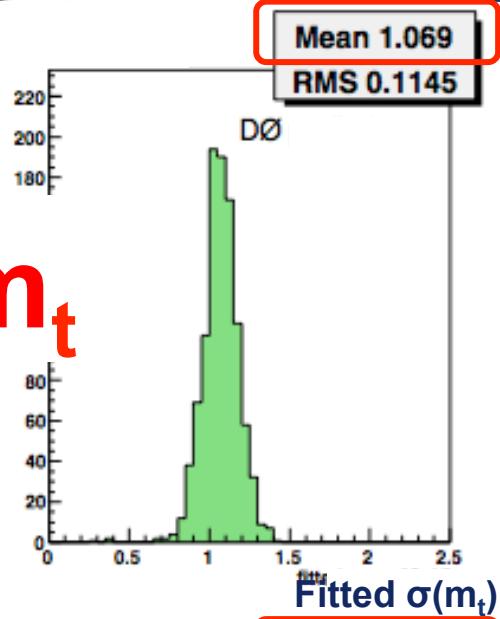


- For calibration of method response in  $m_t$  and  $k_{JES}$ :
  - Construct PEs according to the signal fractions measured for given epoch+channel
  - Vary  $N_{sig}$  in PEs according to binomial statistics
  - 1000 PEs performed at each calibration point
- Calibrate  $m_t$  with samples at
  - $m_t = 165, 170, 172.5, 175, 180 \text{ GeV}$
  - $k_{JES} = 1$
- Calibrate  $k_{JES}$  with samples at
  - $k_{JES} = 0.95, 1, 1.05$
  - $m_t = 172.5 \text{ GeV}$
  - $k_{JES}$  variation is done for background as well



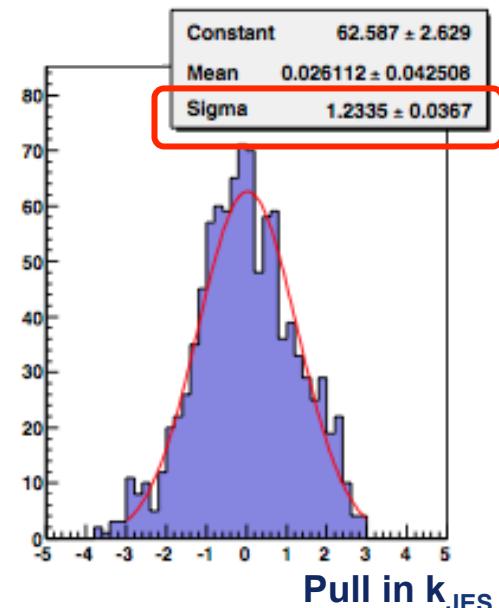
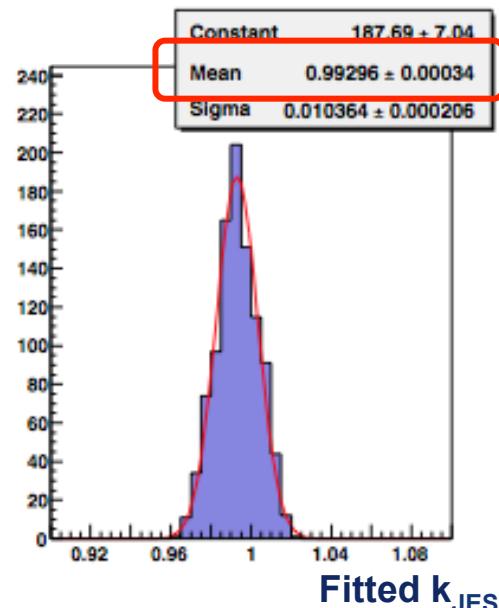
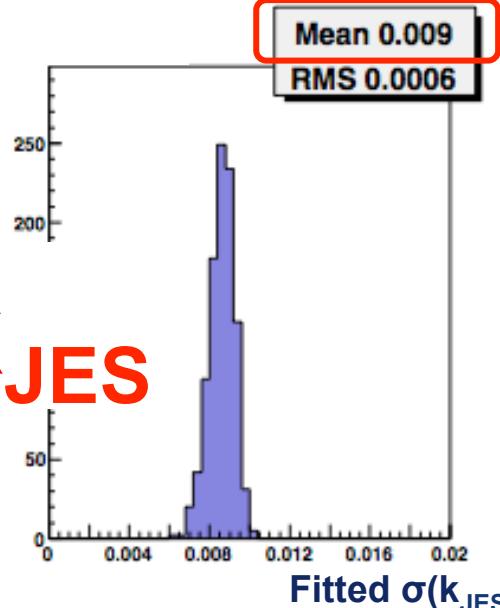
# Calibration of method response in $m_t$ , $k_{\text{JES}}$

$m_t$



$$\pi_{m_t} = \frac{m_t^{\text{fit}} - \langle m_t^{\text{fit}} \rangle}{\sigma_{m_t}^{\text{fit}}}$$

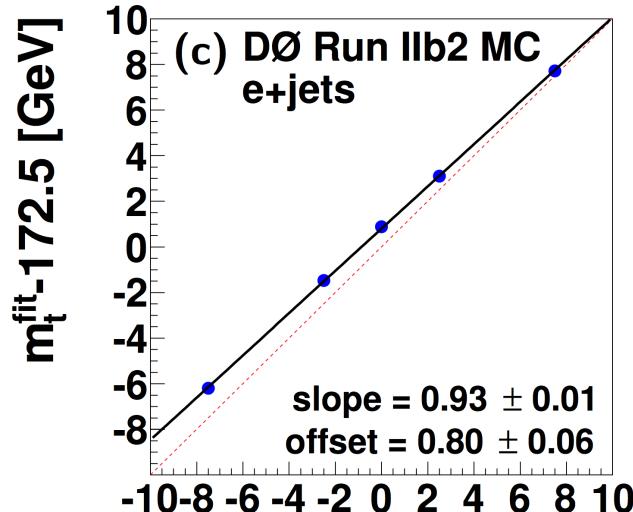
$k_{\text{JES}}$



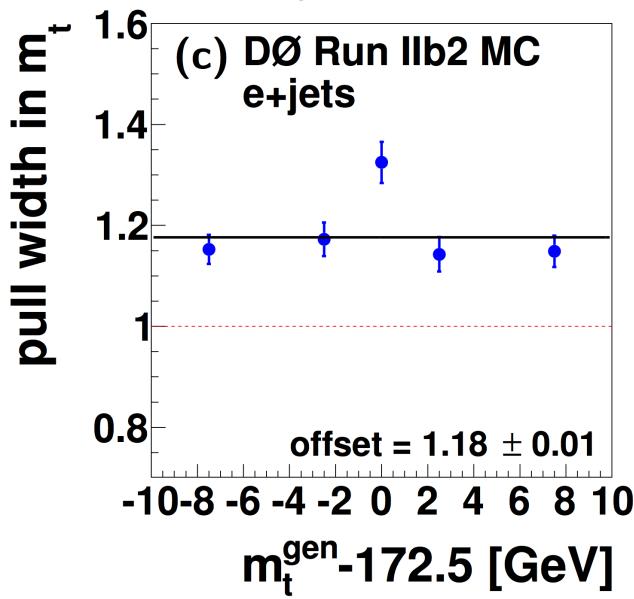
$e + \text{jets}$   
Run IIb2,  $m_t^{\text{gen}} = 172.5 \text{ GeV}$ ,  $k_{\text{JES}} = 1$



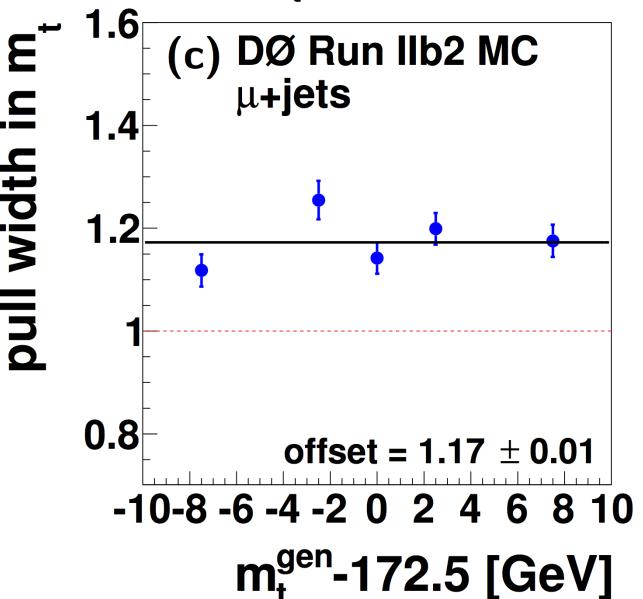
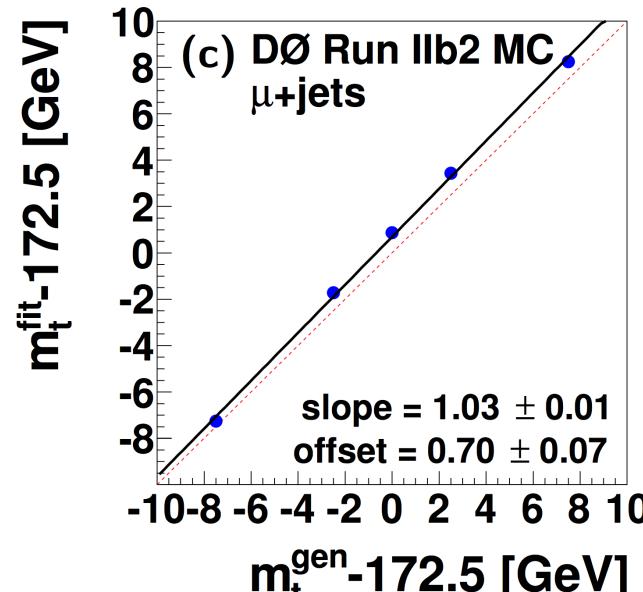
# Calibration of method response in $m_t$



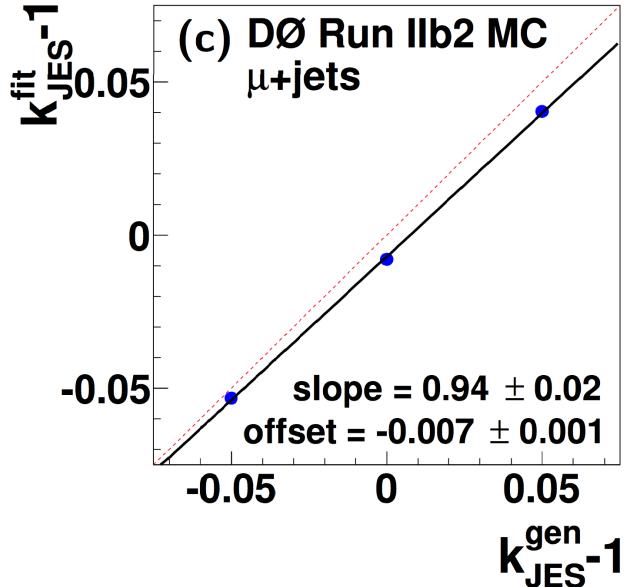
Calibrate  
 $m_t$  &  $\sigma(m_t)$



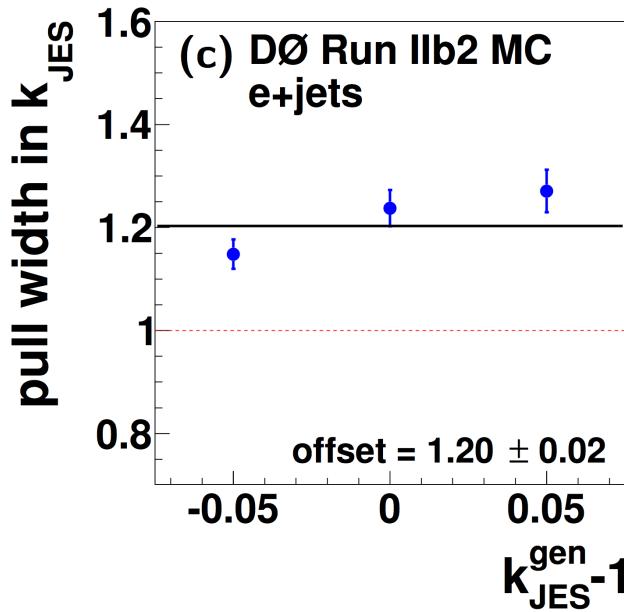
Calibrate  
 $\sigma(m_t)$



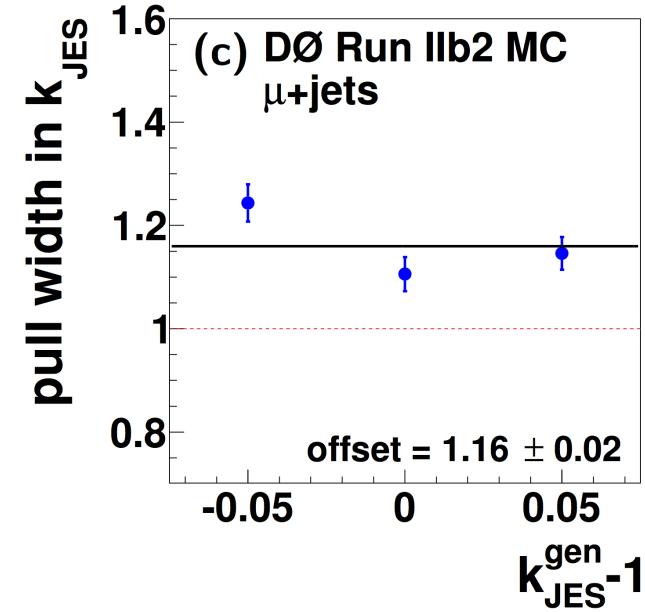
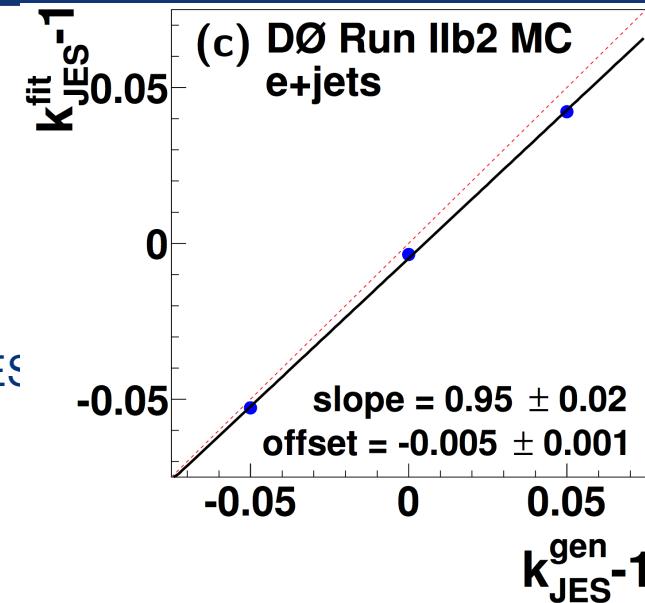
Representative MC simulations for Run IIb2 data  
→ cf. backup for others



Calibrate  
 $k_{\text{JES}}$  &  $\sigma(k_{\text{JES}})$



Calibrate  
 $\sigma(k_{\text{JES}})$

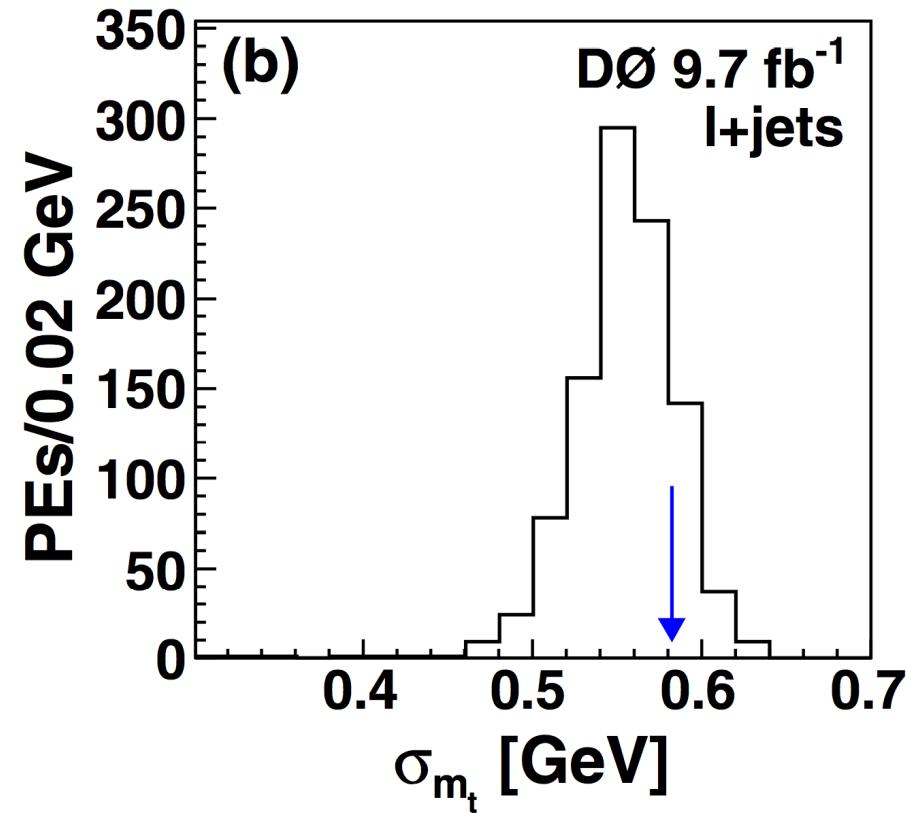
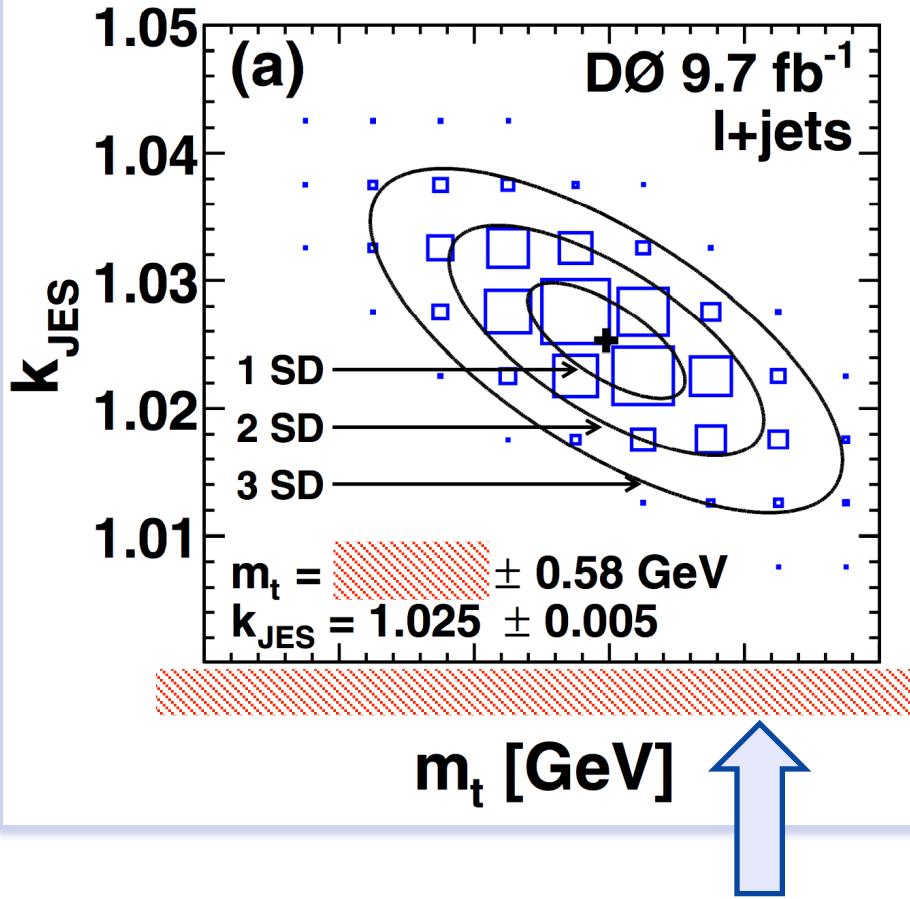


Representative MC simulations for Run IIb2 data  
→ cf. backup for others

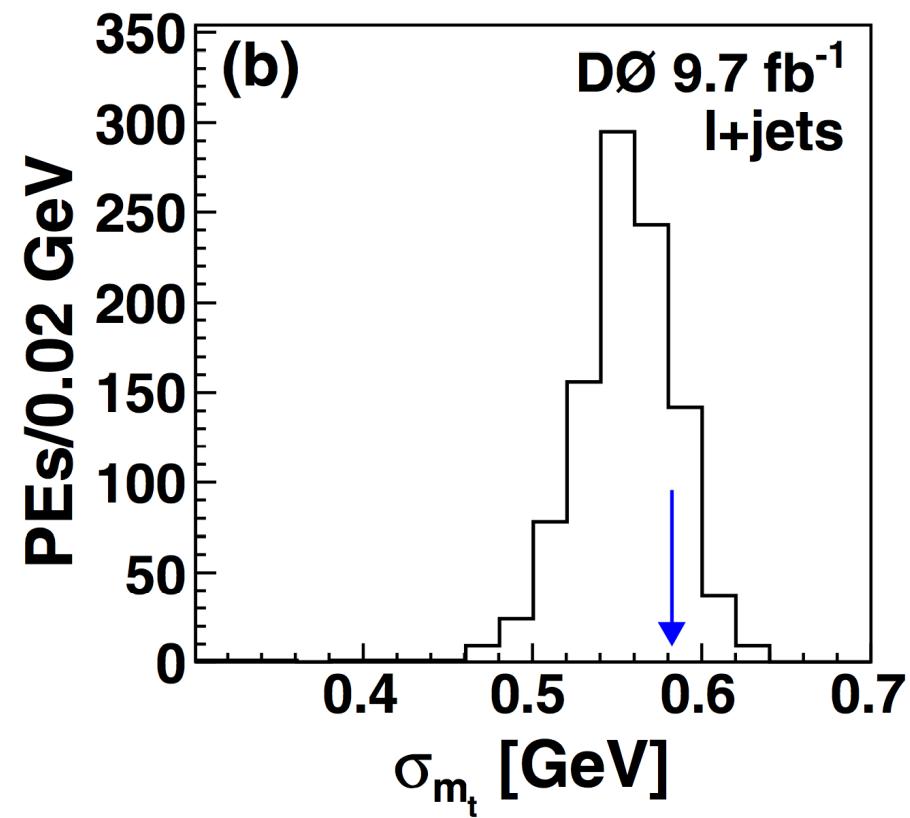
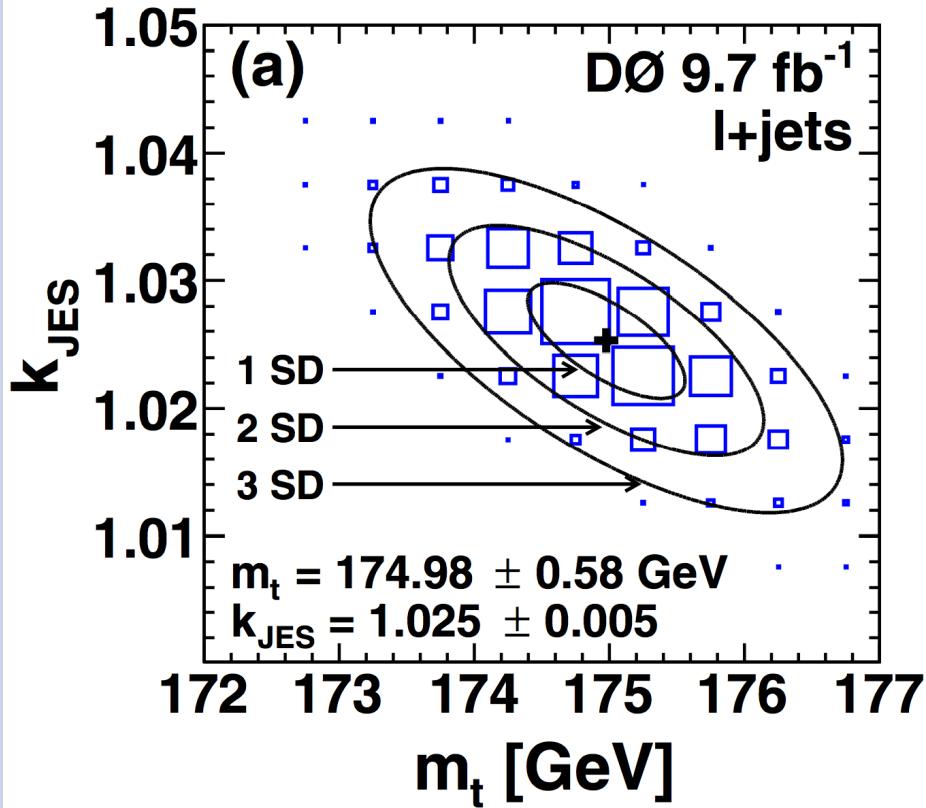


# Result





Data analysis was performed **blinded in  $m_t$** ,  
by adding an offset drawn from [-2 GeV, 2 GeV] around  
172.3 GeV (World average) according to a uniform prior



Data analysis was **unblinded** in last stages of review



# Final result in $\ell$ +jets final states using the full Run II dataset of $9.7 \text{ fb}^{-1}$ (statistical uncertainty only)

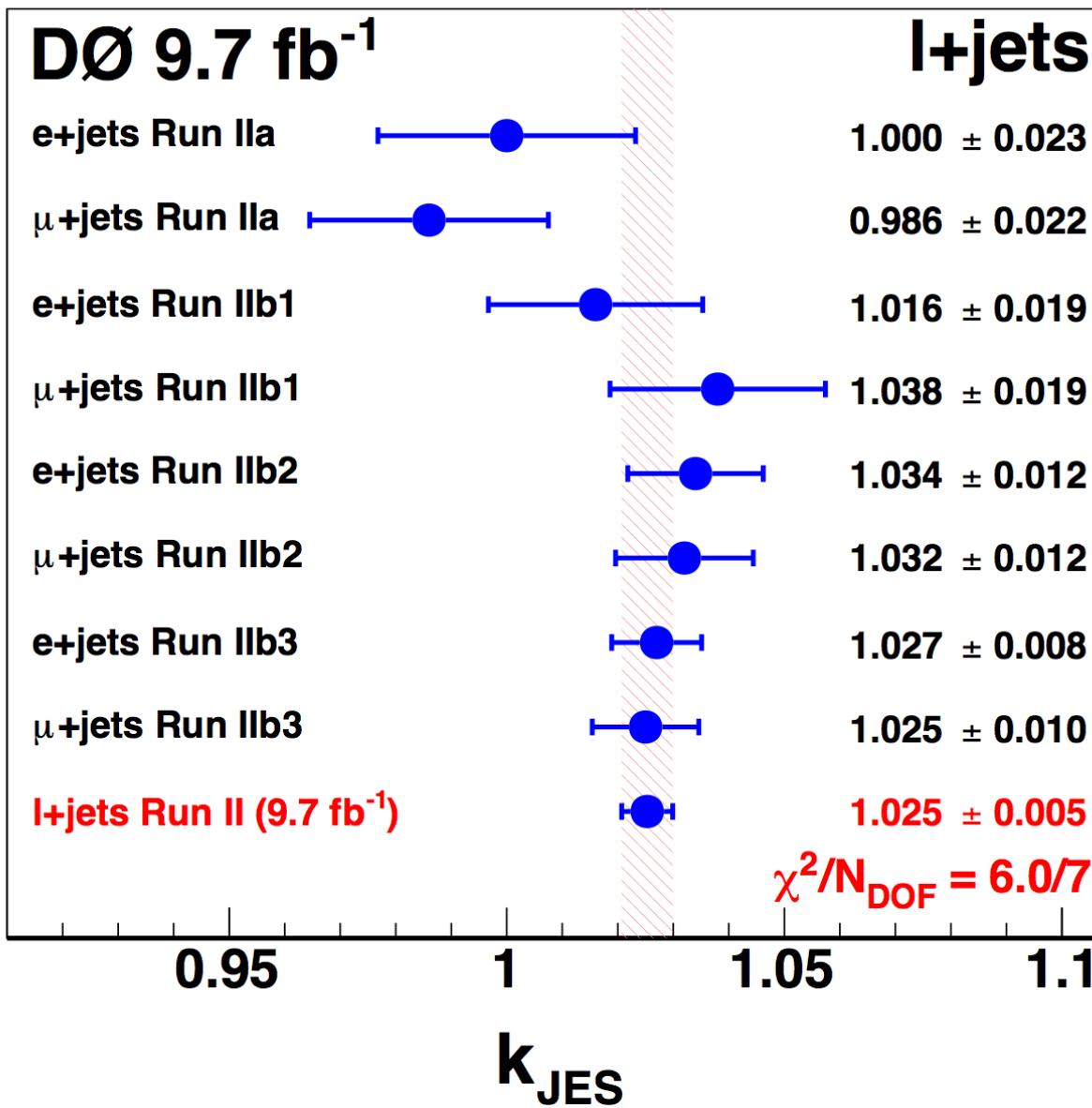
$$m_t = 174.98 \pm 0.41(\text{stat}) \pm 0.41(\text{JES}) \text{ GeV}.$$

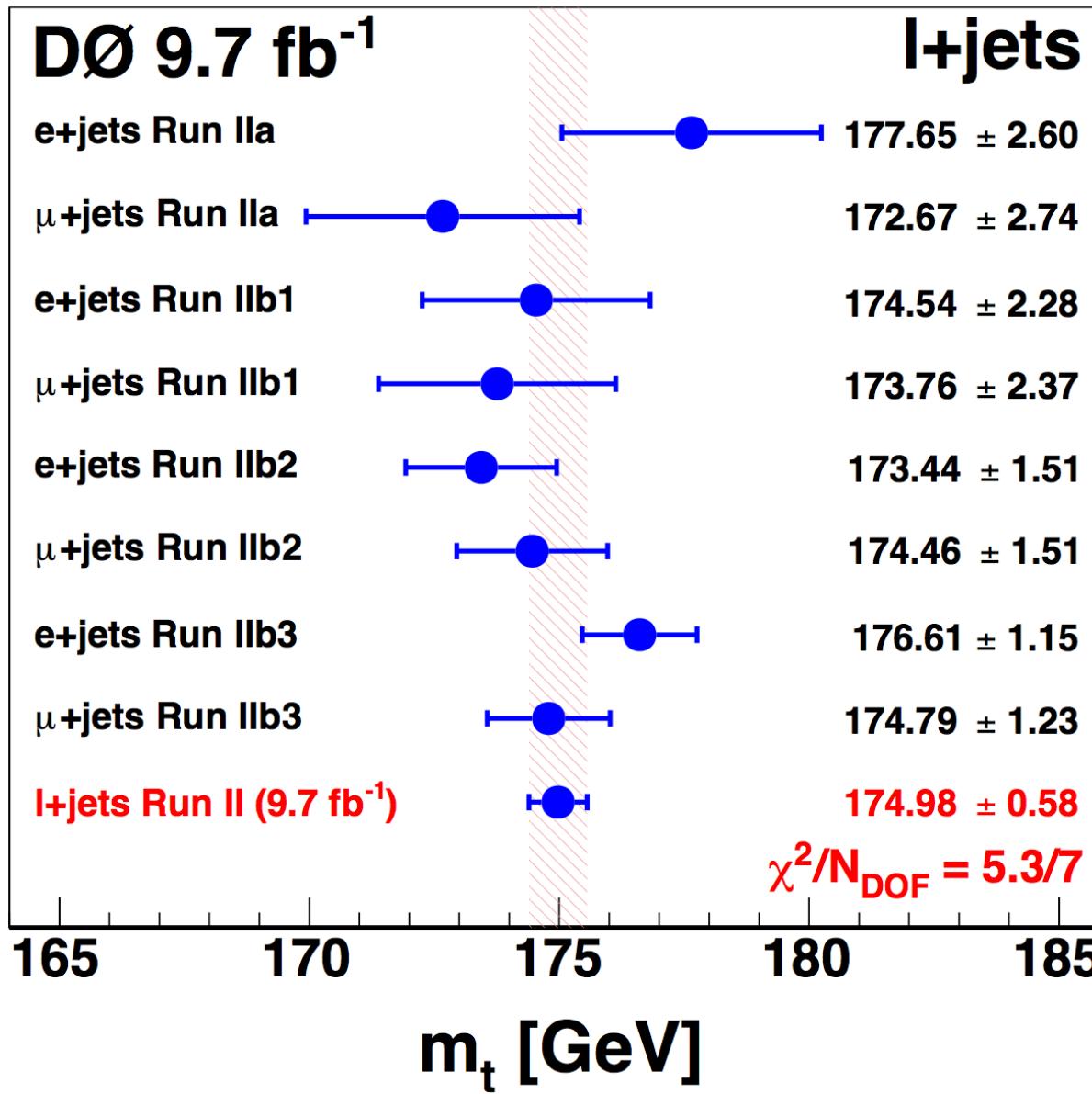
## Summary per final state

$\ell + \text{jets, final}$	:	$m_t$	=	$174.98 \pm 0.58 \text{ (stat+JES) GeV}$ ,
		$k_{\text{JES}}$	=	$1.025 \pm 0.005 \text{ (stat)}$ ;
$e + \text{jets, final}$	:	$m_t$	=	$175.55 \pm 0.81 \text{ (stat+JES) GeV}$ ;
		$k_{\text{JES}}$	=	$1.026 \pm 0.006 \text{ (stat)}$ ;
$\mu + \text{jets, final}$	:	$m_t$	=	$174.36 \pm 0.84 \text{ (stat+JES) GeV}$ .
		$k_{\text{JES}}$	=	$1.025 \pm 0.007 \text{ (stat)}$ ;



# Results per data taking epoch & channel



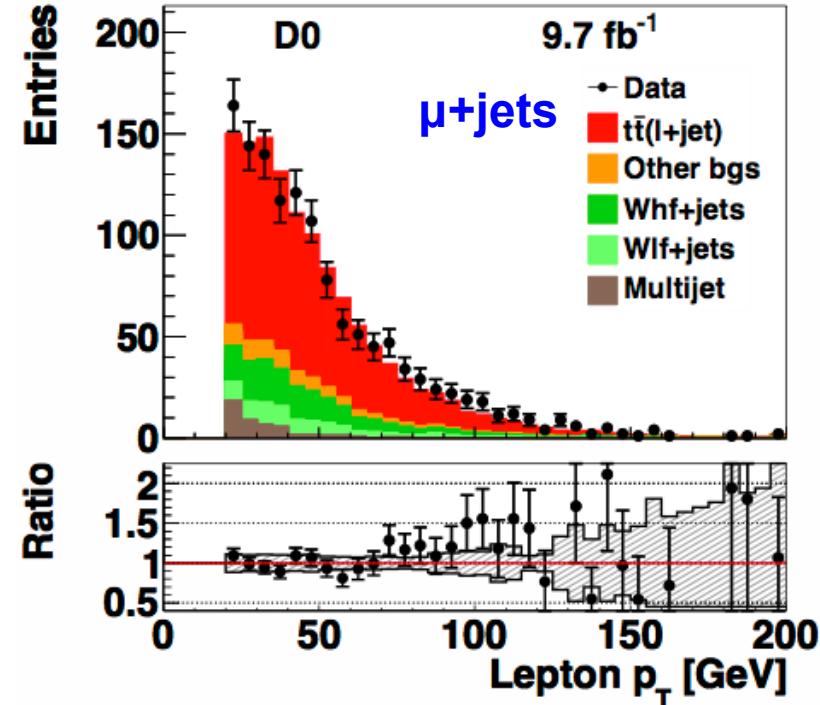
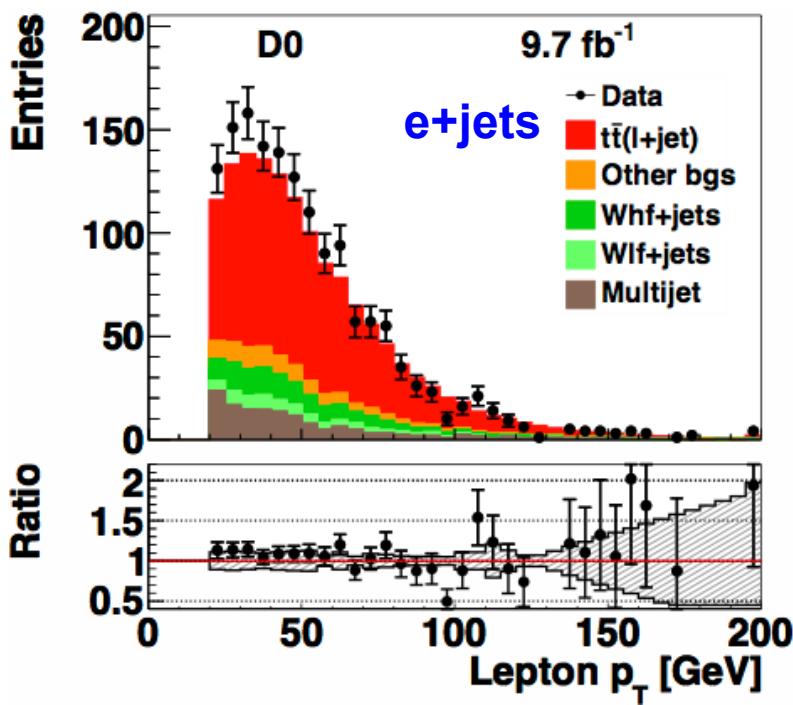




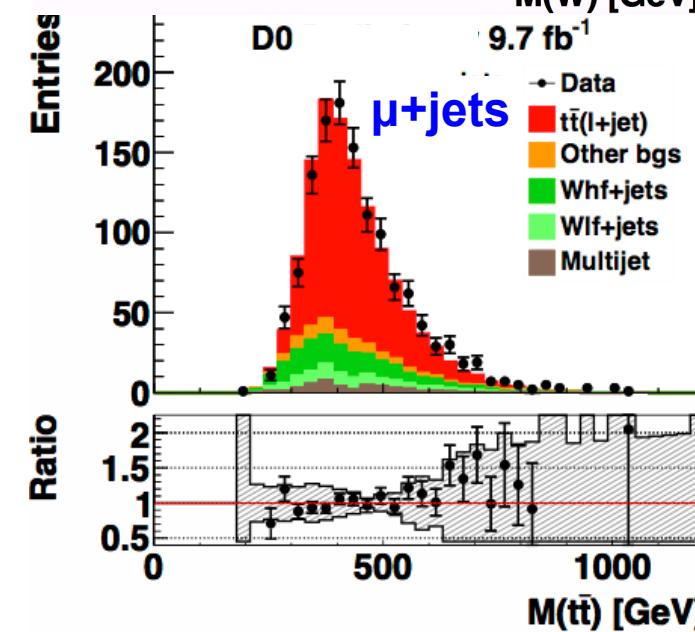
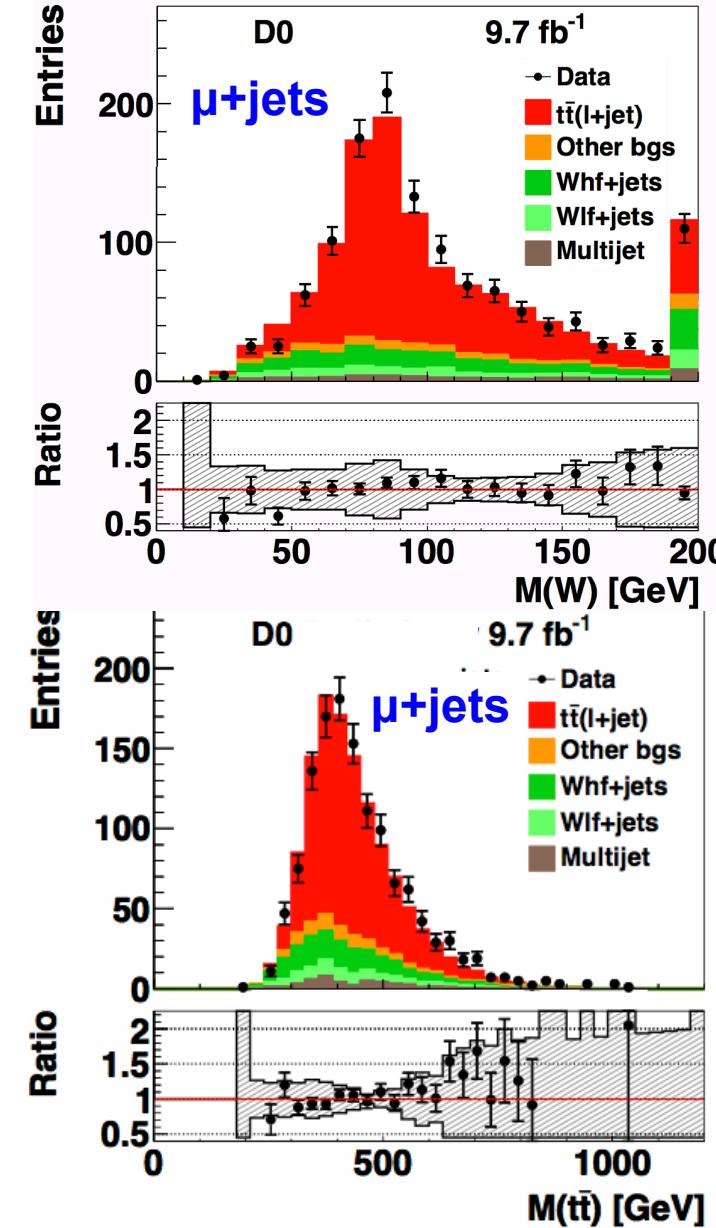
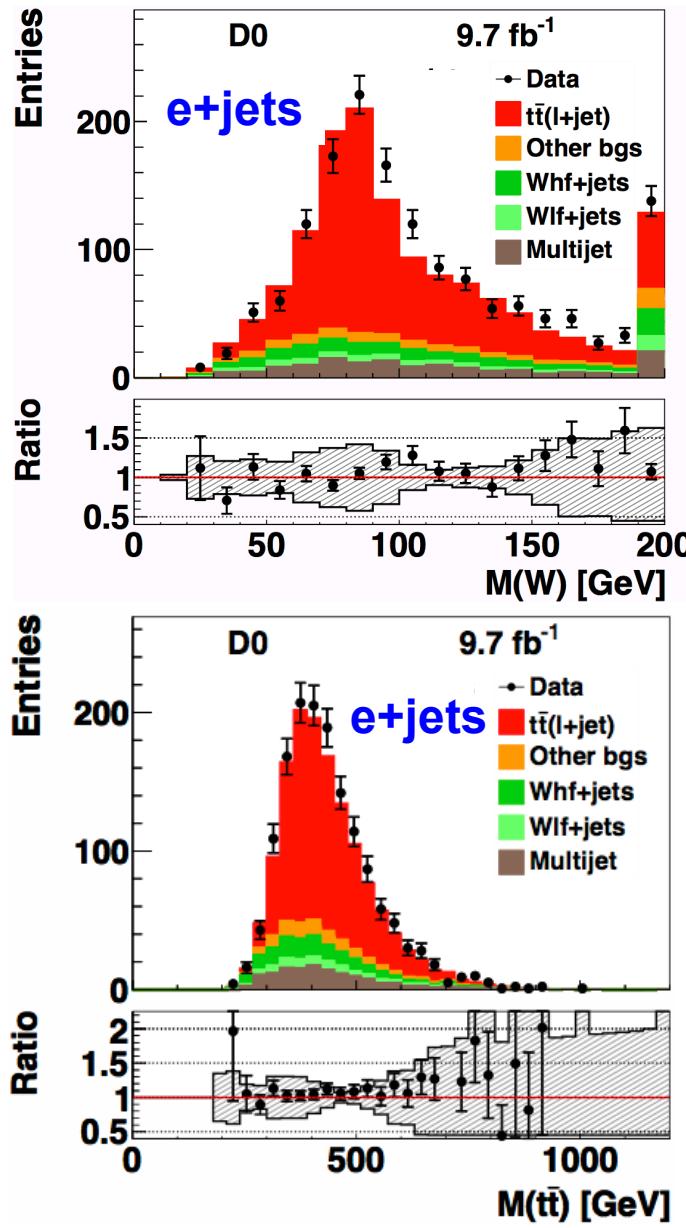
# Data/MC comparisons for best fit $m_t$ , $k_{JES}$ , and $f$



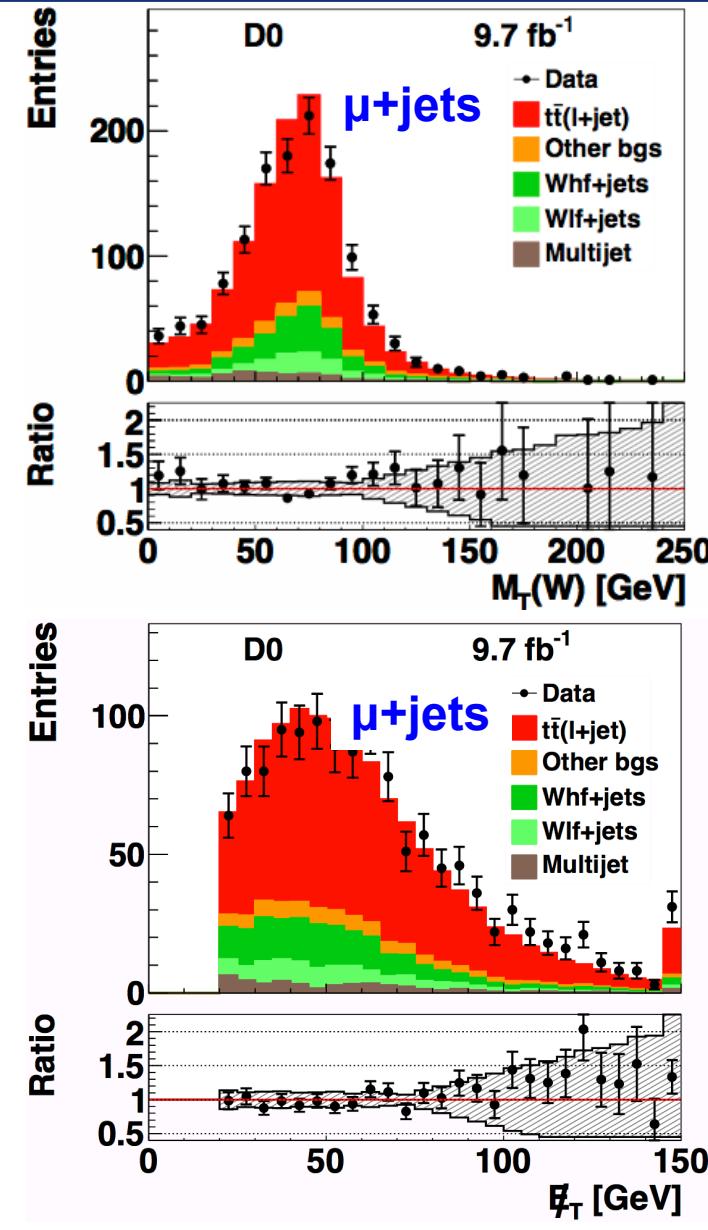
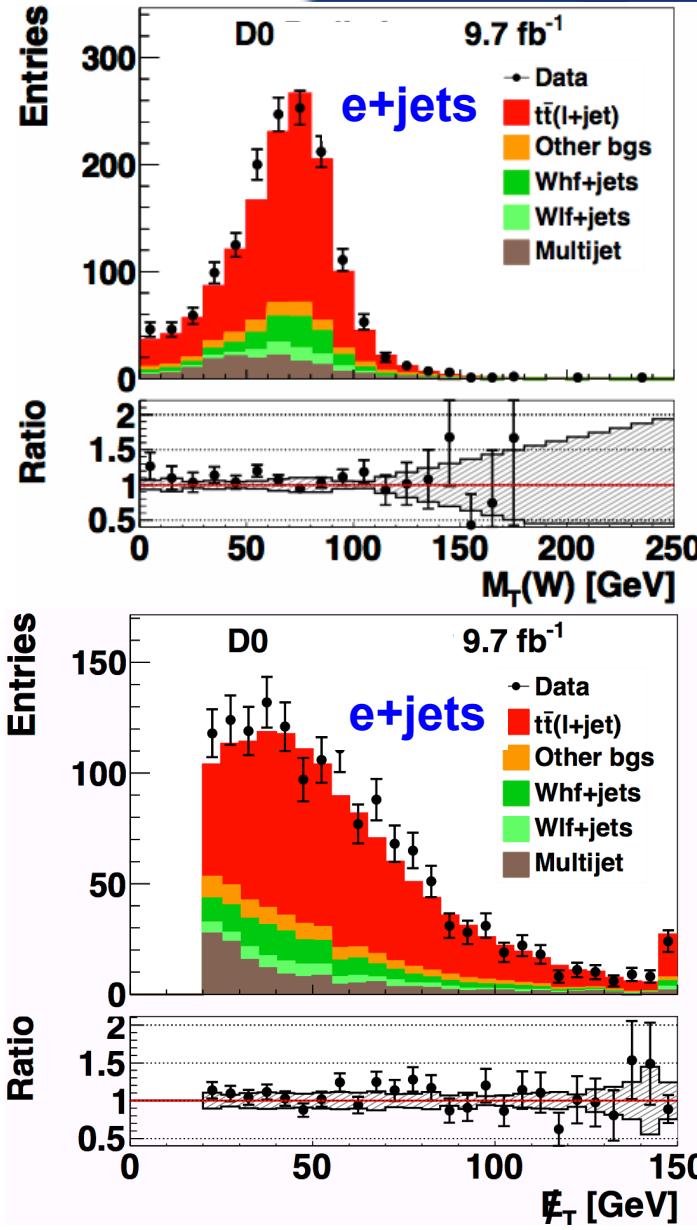
- In the following, showing “post-fit” data/MC comparison plots
  - Use cross sections fitted with the ME technique:
    - 7.8 pb for e+jets
    - 7.6 pb for  $\mu$ +jets
  - Use  $t\bar{t}$  simulations with  $m_t = 175 \text{ GeV}$ ,  $k_{\text{JES}}=1.025$



More data/MC comparisons in backup



More data/MC comparisons in backup



More data/MC comparisons in backup

A close-up, shallow depth-of-field photograph showing a hand in a white glove placing a small, round, yellow cookie dough onto a dark baking sheet. Other cookies are visible in the background.

# Systematic uncertainties



- **Procedures:**
  - Construct pseudo-experiments identically to the default calibration:
    - including dominant W+jets background
    - Method related uncertainties also include next-to-dominant MJ background
  - For four signal modeling uncertainties:
    - Compare different models for  $m_t=172.5 \text{ GeV}$ ,  $k_{\text{JES}}=1$
    - **Stat. uncert.  $O(0.05 \text{ GeV})$**
  - For all other uncertainties, re-derive calibration
    - **Stat. uncert.  $O(0.01 \text{ GeV})$**
- **Discuss refinements in the following:**
  1. Uncertainty due to limited size of MC samples
  2. New calibration of the detector
  3. Refined treatment of signal modeling uncertainties



This measurement

Source of uncertainty	Effect on $m_t$ (GeV)
<i>Signal and background modeling:</i>	
Higher order corrections*	0.15
Initial/final state radiation*	0.09
Hadronization & UE*	0.26
Color reconnection*	0.10
Multiple $p\bar{p}$ interactions	0.06
Heavy flavor scale factor	0.06
$b$ -jet modeling	0.09
PDF uncertainty	0.11
<i>Detector modeling:</i>	
Residual jet energy scale	0.21
Data-MC jet response difference	0.16
$b$ -tagging	0.10
Trigger	0.01
Lepton momentum scale	0.01
Jet energy resolution	0.07
Jet ID efficiency	0.01
<i>Method:</i>	
Modeling of multijet events	0.04
Signal fraction	0.08
MC calibration	0.07
<i>Total systematic uncertainty</i>	
<i>Total statistical uncertainty</i>	
<i>Total uncertainty</i>	

Source	Uncertainty (GeV)
<i>Modeling of production:</i>	
Higher-order effects	$\pm 0.25$
ISR/FSR	$\pm 0.26$
Hadronization and UE	$\pm 0.58$
Color reconnection	$\pm 0.28$
Multiple $p\bar{p}$ interactions	$\pm 0.07$
<i>Modeling of background</i>	$\pm 0.16$
$W+jets$ heavy-flavor scale factor	$\pm 0.07$
<i>Modeling of <math>b</math> jets</i>	$\pm 0.09$
Choice of PDF	$\pm 0.24$
<i>Modeling of detector:</i>	
Residual jet energy scale	$\pm 0.21$
Data-MC jet response difference	$\pm 0.28$
$b$ -tagging efficiency	$\pm 0.08$
Trigger efficiency	$\pm 0.01$
Lepton momentum scale	$\pm 0.17$
Jet energy resolution	$\pm 0.32$
Jet ID efficiency	$\pm 0.26$
<i>Method:</i>	
Multijet contamination	$\pm 0.14$
Signal fraction	$\pm 0.10$
MC calibration	$\pm 0.20$
<i>Total</i>	$\pm 1.02$

Note: for a given source of uncertainty, we cite:  
**max{ statistical uncertainty, |face value of systematic| }**



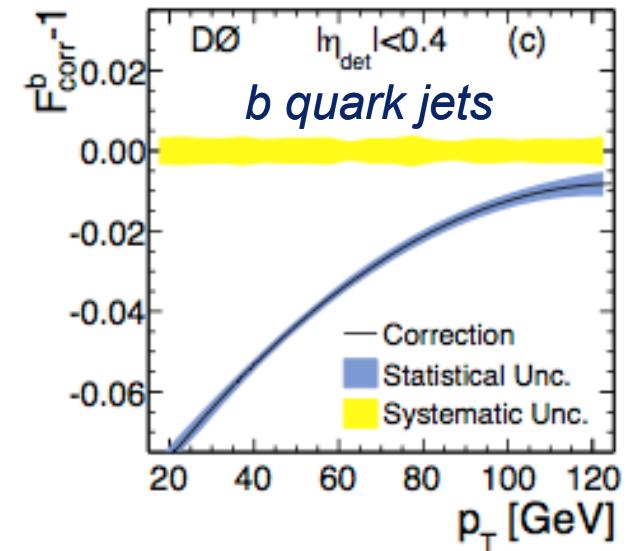
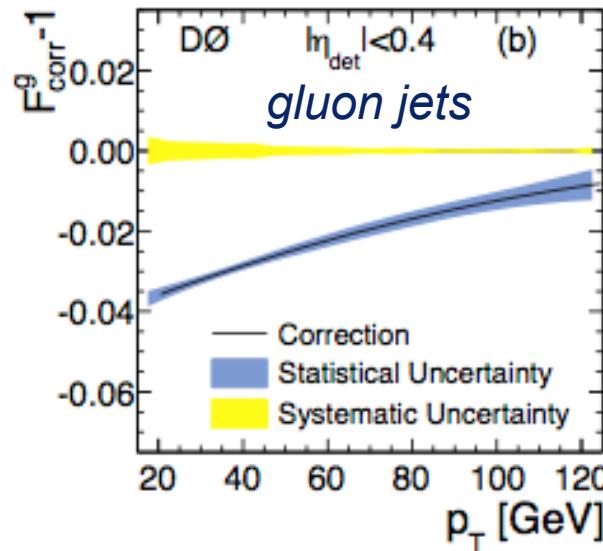
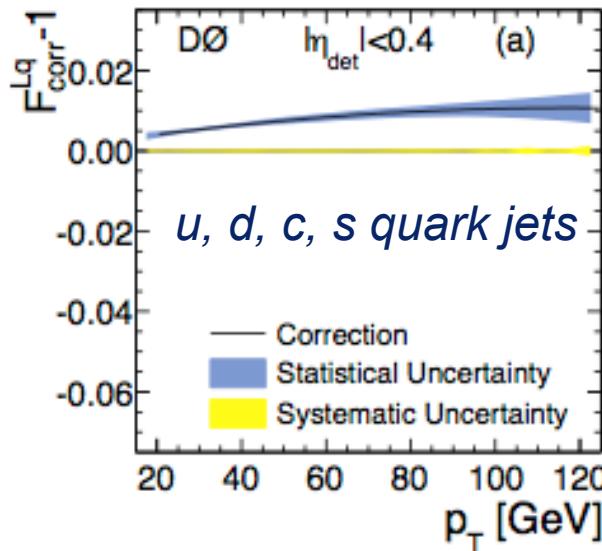
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<i>Total statistical uncertainty</i>	
<i>Total uncertainty</i>	

Source	Uncertainty (GeV)
<i>Modeling of production:</i>	
Modeling of signal:	
Higher-order effects	±0.25
ISR/FSR	±0.26
Hadronization and UE	±0.58
Color reconnection	±0.28
Multiple $p\bar{p}$ interactions	±0.07
Modeling of background	±0.16
$W+jets$ heavy-flavor scale factor	±0.07
Modeling of $b$ jets	±0.09
Choice of PDF	±0.24
<i>Modeling of detector:</i>	
Residual jet energy scale	±0.21
Data-MC jet response difference	±0.28
$b$ -tagging efficiency	±0.08
Trigger efficiency	±0.01
Lepton momentum scale	±0.17
Jet energy resolution	±0.32
Jet ID efficiency	±0.26
<i>Method:</i>	
Multijet contamination	±0.14
Signal fraction	±0.10
MC calibration	±0.20
<b>Total</b>	<b>±1.02</b>

Also here we profit from the increased size of calibration samples

- Use new JES calibration including flavour-dependent response correction as integral part [1]:



- → Uncertainty from flavor-dependent response:
  - 0.16 GeV (was 0.28 GeV)
- This uncertainty accounts for JES difference between light quark jets and b quark jets

[1] in the previous measurement (D $\bar{\nu}$  Coll., PRD 84, 032004 (2011)), the flavour-dependent response correction was used as an a posteriori correction.



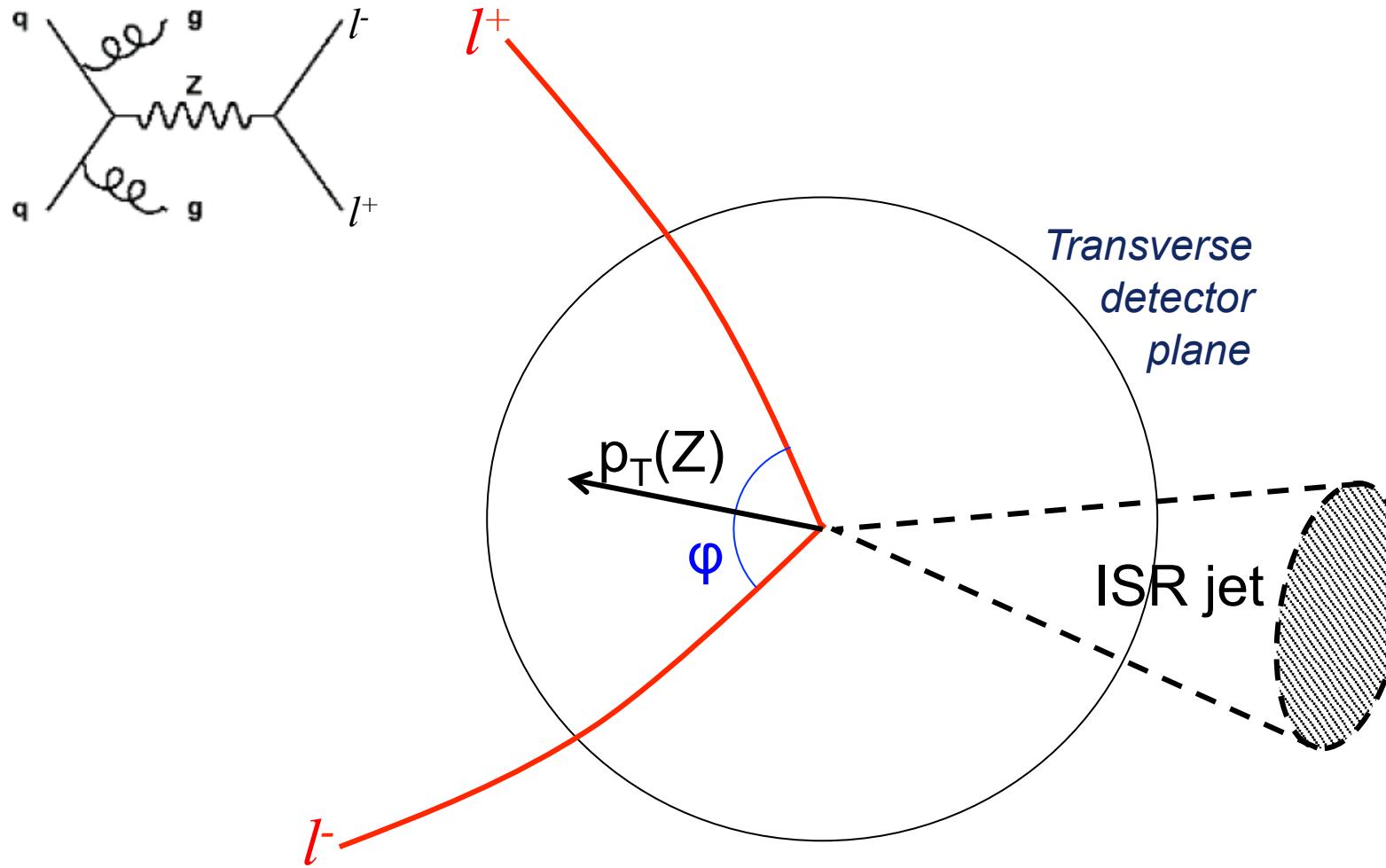
This measurement

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<i>Signal and background modeling:</i>	
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<i>Modeling of <math>b</math> jets</i>	
Choice of PDF	$\pm 0.24$
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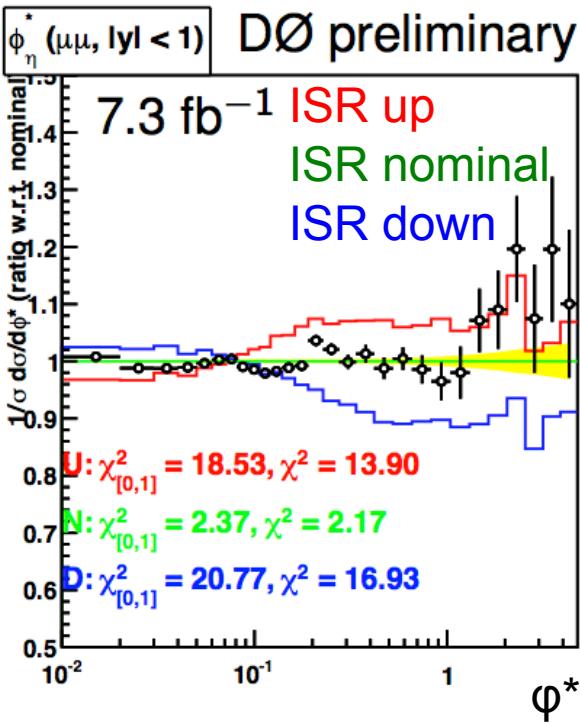
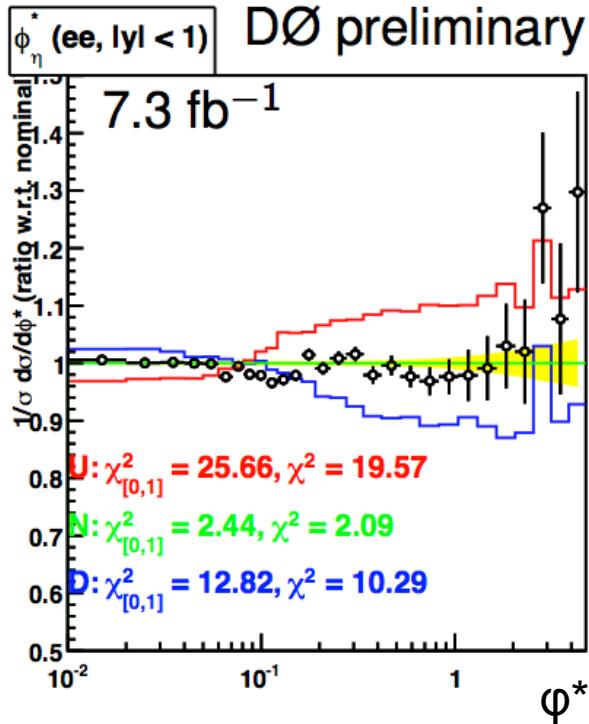
Also here we profit from the increased size of calibration samples

- Constrain ISR/FSR by studying Drell-Yan events
  - Measurement of  $p_T(Z)$  using  $\varphi^*$  variable [1]



[1]  $D\emptyset$  Coll., PRL 106, 122001 (2011)

- Constrain ISR/FSR by studying Drell-Yan events
  - Measurement of  $p_T(Z)$  using  $\varphi^*$  variable [1]
  - Vary ISR/FSR via CKKW renormalization scale in **alpgen** (**ktfac**), as suggested in [2]
    - **ktfac** variations by  $\pm 1.5$  cover excursions of MC from data



Also tune in other kinematic regions:  

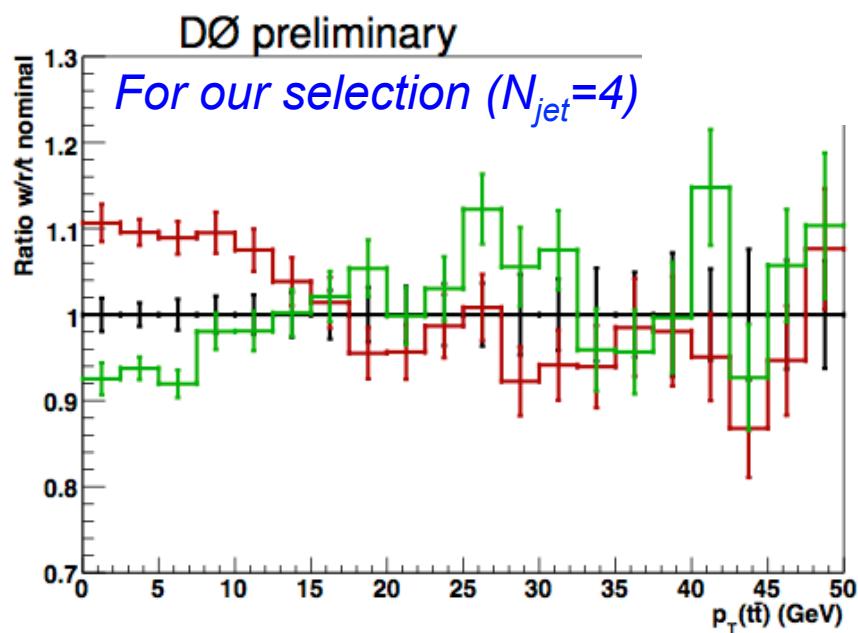
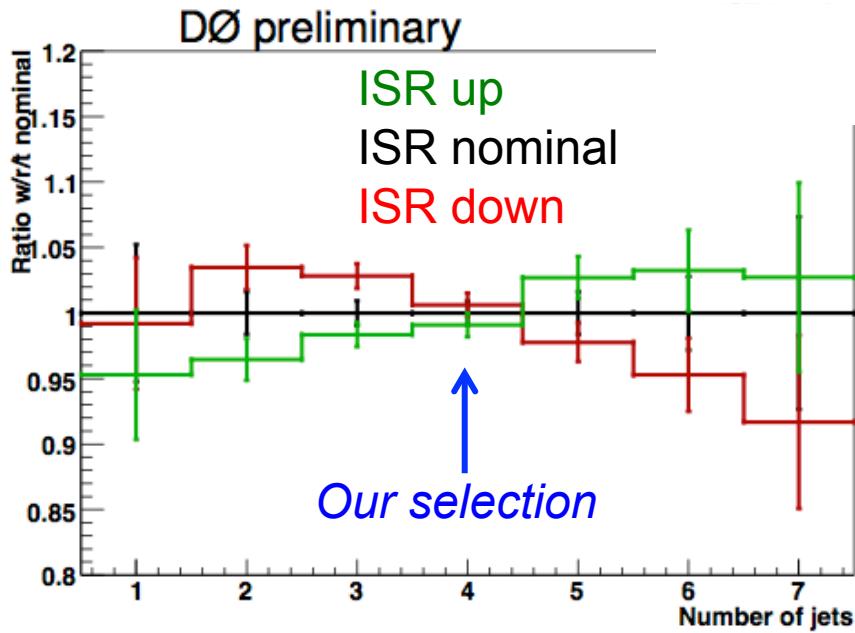
- $1 < |y| < 2$
- $|y| > 2$   
(cf. backup)

[1] DØ Coll., PRL 106, 122001 (2011)

[2] M. Mangano, P. Skands et al, EPJ C72 2078 (2012)

- Constrain ISR/FSR by studying Drell-Yan events
  - Measurement of  $p_T(Z)$  using  $\varphi^*$  variable [1]
- Vary ISR/FSR via CKKW renormalization scale in alpgen (ktfac), as suggested in [2]
  - ktfac variations by  $\pm 1.5$  cover excursions of MC from data

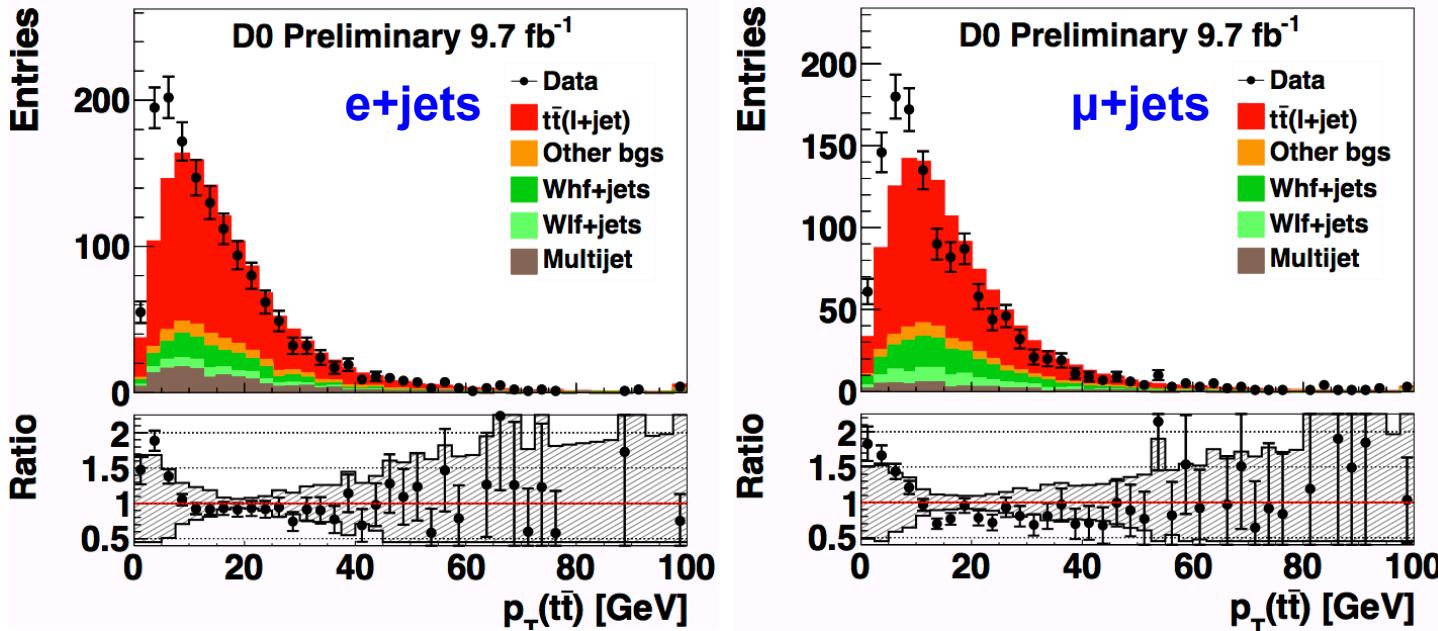
## The effect of ISR/FSR variations in top-antitop events



[1] DØ Coll., PRL 106, 122001 (2011)

[2] M. Mangano, P. Skands et al, EPJ C72 2078 (2012)

- Constrain ISR/FSR by studying Drell-Yan events
  - Measurement of  $p_T(Z)$  using  $\varphi^*$  variable [1]
- Vary ISR/FSR via CKKW renormalization scale in alpgen (ktfac), as suggested in [2]
  - ktfac variations by  $\pm 1.5$  cover excursions of MC from data
- In addition: reweight  $t\bar{t}$  simulations in  $p_T(t\bar{t})$  to data

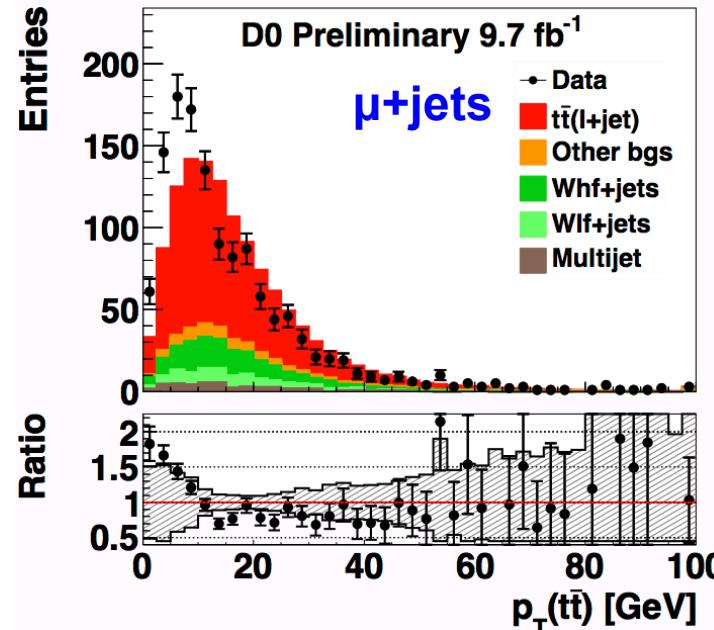
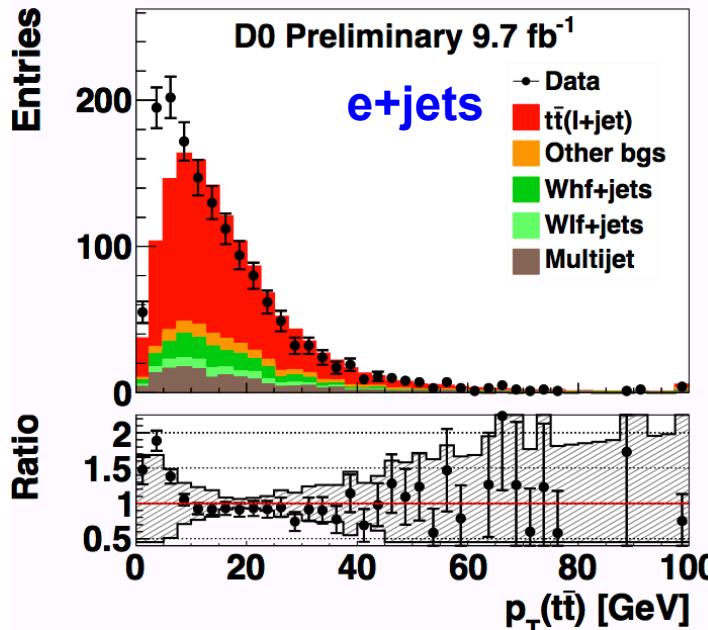


- Effect may be related to ISR/FSR mismodelling

[1] *DØ Coll., PRL 106, 122001 (2011)*

[2] *M. Mangano, P. Skands et al, EPJ C72 2078 (2012)*

- Constrain ISR/FSR by studying Drell-Yan events
  - Measurement of  $p_T(Z)$  using  $\varphi^*$  variable [1]
  - Vary ISR/FSR via CKKW renormalization scale in alpgen ( $k_{tfac}$ ), as suggested in [2]
    - $k_{tfac}$  variations by  $\pm 1.5$  cover excursions of MC
- In addition: reweight  $t\bar{t}$  simulations in  $p_T(t\bar{t})$  to data



- Effect may be related to ISR/FSR mismodelling

0.06 GeV

0.07 GeV



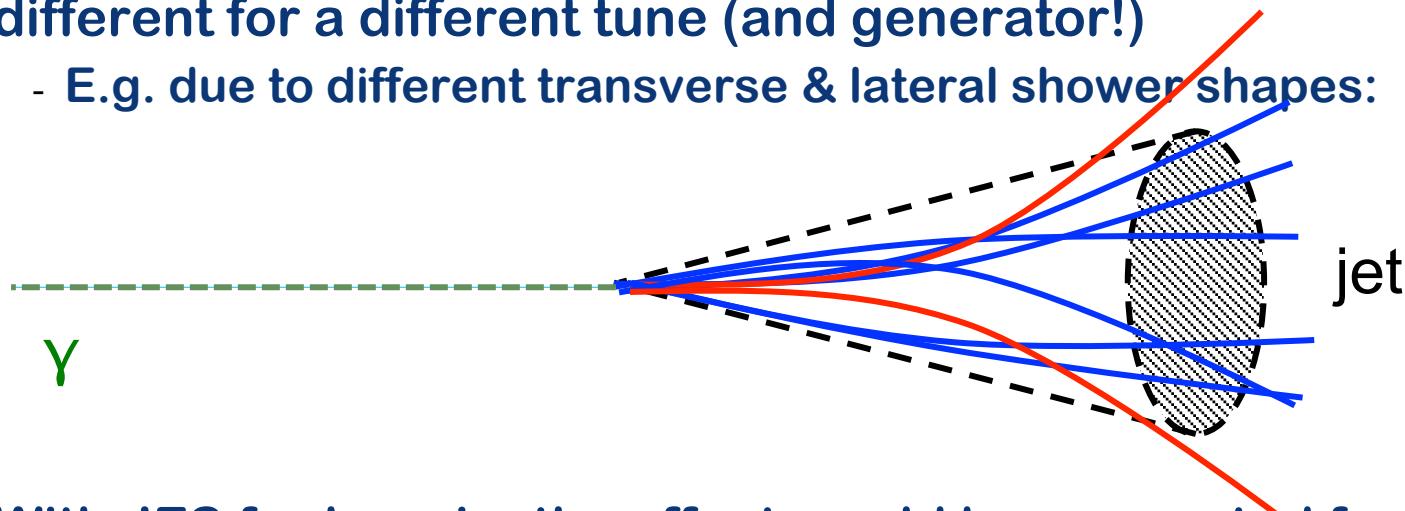
This measurement

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<i>Signal and background modeling:</i>	
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Initial/final state radiation*	0.09
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Signal fraction	0.08
MC calibration	0.07
<i>Total systematic uncertainty</i>	
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<i>Total uncertainty</i>	

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Signal fraction	$\pm 0.10$
MC calibration	$\pm 0.20$
<i>Total</i>	$\pm 1.02$

Also here we profit from the increased size of calibration samples

- Compare **alpgen+herwig** vs **alpgen+pythia** (default)
- Combination of two effects:
  1. The **actual effect** we are interested in
  2. Component from **different JES** (differential in  $p_T, \eta$ )
    - Strictly, our **JES** is valid only for **pythia** with D0 Tune A
    - One can argue from first principles that JES will be different for a different tune (and generator!)
      - E.g. due to different transverse & lateral shower shapes:



- With JES for herwig, the effect would be accounted for
  - Factorize uncertainties to avoid double-counting:
    - We have an uncertainty for dependence of JES on  $p_T, \eta, \dots$



- Factor out the component from different JES
  - Evaluate using the momenta of particle level jets matched to detector level jets with  $\Delta R = 0.25 = R_{\text{cone}}/2$
  - Apply default selection at detector level
    - → minimize bias from acceptance etc.
- We also factor out the effect of different  $p_T(t\bar{t})$  in:
  - Default (alpgen+pythia)
  - Alternative model (alpgen+herwig)
  - Achieved by reweighting default simulation in  $p_T(t\bar{t})$  to match the alternative model
  - This effect is already taken into account in ISR/FSR uncertainty
- → Hadronization and underlying event uncertainty:
  - 0.26 GeV (was: 0.58 GeV)



This measurement

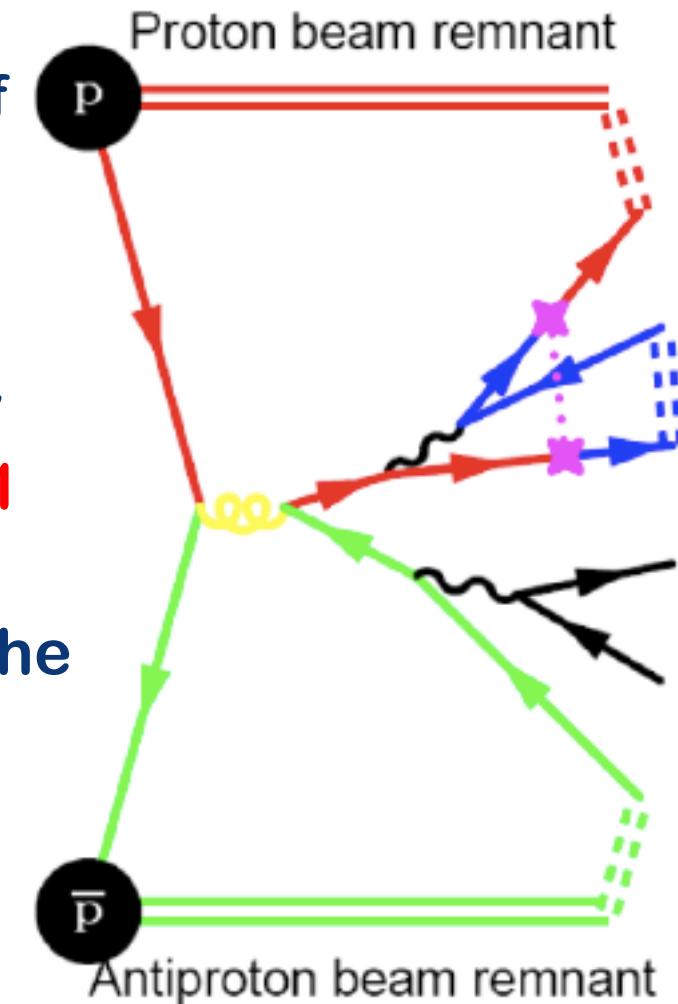
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Heavy flavor scale factor	0.06
$b$ -jet modeling	0.09
PDF uncertainty	0.11
<i>Detector modeling:</i>	
Residual jet energy scale	0.21
Data-MC jet response difference	0.16
$b$ -tagging	0.10
Trigger	0.01
Lepton momentum scale	0.01
Jet energy resolution	0.07
Jet ID efficiency	0.01
<i>Method:</i>	
Modeling of multijet events	0.04
Signal fraction	0.08
MC calibration	0.07
<i>Total systematic uncertainty</i>	
<i>Total statistical uncertainty</i>	
<i>Total uncertainty</i>	

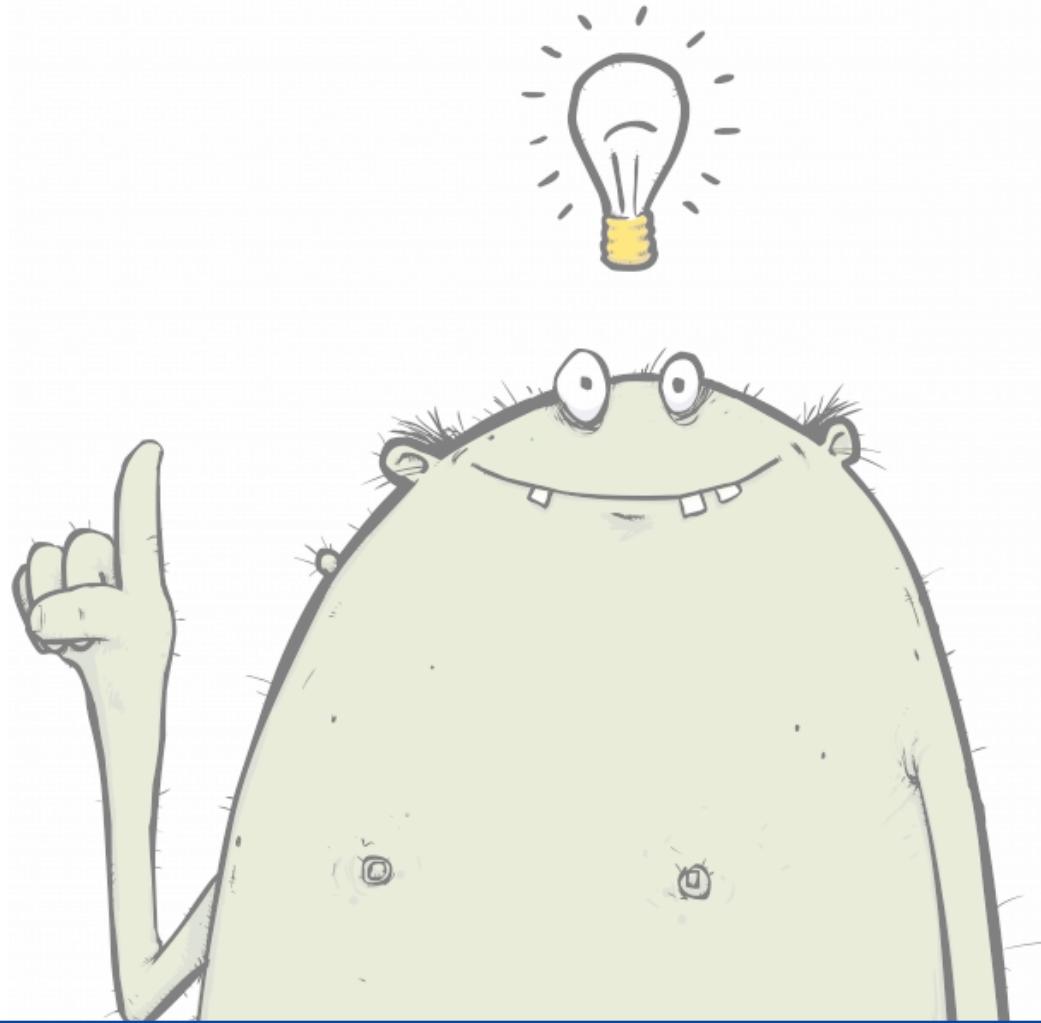
Source	Uncertainty (GeV)
<i>Modeling of production:</i>	
Higher-order effects	$\pm 0.25$
<i>Modeling of signal:</i>	
ISR/FSR	$\pm 0.26$
Hadronization and UE	$\pm 0.58$
Color reconnection	$\pm 0.28$
Multiple $p\bar{p}$ interactions	$\pm 0.07$
<i>Modeling of background</i>	
$W+jets$ heavy-flavor scale factor	$\pm 0.07$
<i>Modeling of <math>b</math> jets</i>	
Choice of PDF	$\pm 0.24$
<i>Modeling of detector:</i>	
Residual jet energy scale	$\pm 0.21$
Data-MC jet response difference	$\pm 0.28$
$b$ -tagging efficiency	$\pm 0.08$
Trigger efficiency	$\pm 0.01$
Lepton momentum scale	$\pm 0.17$
Jet energy resolution	$\pm 0.32$
Jet ID efficiency	$\pm 0.26$
<i>Method:</i>	
Multijet contamination	$\pm 0.14$
Signal fraction	$\pm 0.10$
MC calibration	$\pm 0.20$
<i>Total</i>	$\pm 1.02$

Also here we profit from the increased size of calibration samples



- Use new color reconnection model:
  - Parametrises colour string survival probability in terms of the rapidity difference of beginning and end of color string
  - Old crude model:
    - ad-hoc breaking up of color connections with some probability
- Compare pythia with Perugia 2011 vs Perugia 2011NOCR tunes
- Use identical hard ME events for the comparison
- → Uncertainty from color reconnection:
  - 0.10 GeV (was: 0.28 GeV)





**(Intermediate) conclusions**



## Results of three years of hard work and countless studies...

This measurement

Source of uncertainty	Effect on $m_t$ (GeV)
<i>Signal and background modeling:</i>	
Higher order corrections*	0.15
Initial/final state radiation*	0.09
Hadronization & UE*	0.26
Color reconnection*	0.10
Multiple $p\bar{p}$ interactions	0.06
Heavy flavor scale factor	0.06
$b$ -jet modeling	0.09
PDF uncertainty	0.11
<i>Detector modeling:</i>	
Residual jet energy scale	0.21
Data-MC jet response difference	0.16
$b$ -tagging	0.10
Trigger	0.01
Lepton momentum scale	0.01
Jet energy resolution	0.07
Jet ID efficiency	0.01
<i>Method:</i>	
Modeling of multijet events	0.04
Signal fraction	0.08
MC calibration	0.07
<i>Total systematic uncertainty</i>	0.49
<i>Total statistical uncertainty</i>	
<i>Total uncertainty</i>	

Source	Uncertainty (GeV)
<i>Modeling of production:</i>	
Modeling of signal:	
Higher-order effects	±0.25
ISR/FSR	±0.26
Hadronization and UE	±0.58
Color reconnection	±0.28
Multiple $p\bar{p}$ interactions	±0.07
Modeling of background	±0.16
$W+jets$ heavy-flavor scale factor	±0.07
Modeling of $b$ jets	±0.09
Choice of PDF	±0.24
<i>Modeling of detector:</i>	
Residual jet energy scale	±0.21
Data-MC jet response difference	±0.28
$b$ -tagging efficiency	±0.08
Trigger efficiency	±0.01
Lepton momentum scale	±0.17
Jet energy resolution	±0.32
Jet ID efficiency	±0.26
<i>Method:</i>	
Multijet contamination	±0.14
Signal fraction	±0.10
MC calibration	±0.20
<b>Total</b>	±1.02

1.02 GeV

0.49 GeV



- Lots of work went into **refining** the evaluation of **systematic uncertainties** for this measurement:
  - 1) **Acceleration of the ME technique** by  $O(100)$ :
    - Eliminate the statistical component from systematics
  - 2) For “free”: the great work done by the **JES group!!!**
  - 3) **Refine the evaluation of uncertainties**
    - include new, modern models + factorise sources



- Lots of work went into refining the evaluation of systematic uncertainties for this measurement:
  - 1) Acceleration of the ME technique by  $O(100)$ :
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  - 3) Refine the evaluation of uncertainties
    - include new, modern models + factorise sources

$$m_t = 174.98 \pm 0.76 \text{ GeV} \quad 0.43\%$$

$$m_t = 174.98 \pm 0.58 \text{ (stat + JES)} \pm 0.49 \text{ (syst) GeV}$$

0.44% - Most precise @ LHC:  $172.04 \pm 0.19 \text{ (stat.+JSF)} \pm 0.75 \text{ (syst.) GeV}$ .

0.64% - 2<sup>nd</sup> most precise @ Tevatron:  $172.9 \pm 0.5 \text{ (stat)} \pm 1.0 \text{ (syst) GeV}$

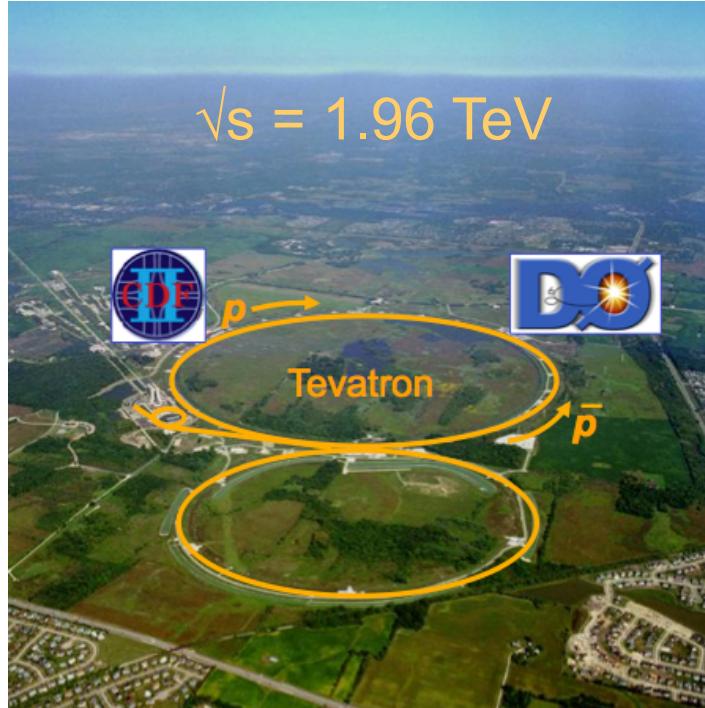
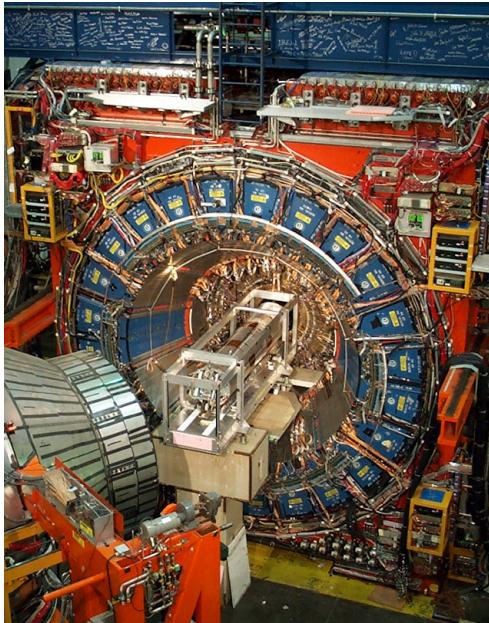


# 2014 Tevatron average of $m_t$





# Tevatron average: input measurements



- **Run I:**
  - PRD 63, 032003 (2001) [CDF I+jets,  $0.1 \text{ fb}^{-1}$ ]
  - Nature 429, 638 (2004) [DØ I+jets,  $0.1 \text{ fb}^{-1}$ ]
  - PRL 82, 271 (1999) [CDF dilep.,  $0.1 \text{ fb}^{-1}$ ]
  - PRL 80, 2063 (1998) [DØ dilep.,  $0.1 \text{ fb}^{-1}$ ]
  - PRL 79, 1992 (1997) [CDF all-jets,  $0.1 \text{ fb}^{-1}$ ]
- **Run II:**
  - PRL 109, 152003 (2012) [CDF I+jets,  $8.7 \text{ fb}^{-1}$ ]
  - PRL 113, 032002 (2014) [DØ I+jets,  $9.7 \text{ fb}^{-1}$ ]
  - CDF-Conf 11072 (2014) [CDF dilep.,  $9.1 \text{ fb}^{-1}$ ]
  - PRD 86, 051103 (2012) [DØ dilepton  $5.4 \text{ fb}^{-1}$ ]
  - CDF-Conf 11084 (2014) [CDF all-jets,  $9.3 \text{ fb}^{-1}$ ]
  - PRD 88, 011101 (2013) [CDF  $E_T$ +jets,  $8.7 \text{ fb}^{-1}$ ]
  - PRD 81 032002 (2009) [CDF  $L_{xy}$ ,  $1.9 \text{ fb}^{-1}$ ]

arXiv:1407.2682 [hep-ex]



	Run I published					Run II published					Run II prel.	
	CDF			DØ		CDF			DØ		CDF	
	$\ell+$ jets	$\ell\ell$	all-jets	$\ell+$ jets	$\ell\ell$	Lxy	Met	$\ell+$ jets	$\ell\ell$	all-jets	$\ell\ell$	all-jets
$\int \mathcal{L} dt$	0.1	0.1	0.1	0.1	0.1	8.7	1.9	8.7	9.7	5.4	9.1	9.3
Result	176.1	167.4	186.0	180.1	168.4	172.85	166.90	173.93	174.98	174.00	170.80	175.07
<i>In situ</i> light-jet calibration (iJES)	n/a	n/a	n/a	n/a	n/a	0.49	n/a	1.05	0.41	0.55	n/a	0.97
Response to $b/q/g$ jets (aJES)	n/a	n/a	n/a	0.0	0.0	0.09	0.00	0.10	0.16	0.40	0.18	0.02
Model for $b$ jets (bJES)	0.6	0.8	0.6	0.7	0.7	0.16	0.00	0.17	0.09	0.20	0.28	0.20
Out-of-cone correction (cJES)	2.7	2.6	3.0	2.0	2.0	0.21	0.36	0.18	n/a	n/a	1.65	0.37
Light-jet response (1) (rJES)	3.4	2.7	4.0	n/a	n/a	0.48	0.24	0.40	n/a	n/a	1.72	0.42
Light-jet response (2) (dJES)	0.7	0.6	0.3	2.5	1.1	0.07	0.06	0.04	0.21	0.56	0.46	0.09
Lepton modeling (LepPt)	n/e	n/e	n/e	n/e	n/e	0.03	0.00	n/a	0.01	0.35	0.36	n/a
Signal modeling (Signal)	2.6	2.9	2.0	1.1	1.8	0.61	0.90	0.63	0.35	0.86	0.96	0.53
Jet modeling (DetMod)	0.0	0.0	0.0	0.0	0.0	0.00	0.00	0.00	0.07	0.50	0.00	0.00
$b$ -tag modeling (b-tag)	0.4	0.0	0.0	0.0	0.0	0.03	0.00	0.03	0.10	0.00	0.05	0.04
Background from theory (BGMC)	1.3	0.3	0.0	1.0	1.1	0.12	0.80	0.00	0.06	0.00	0.30	0.00
Background based on data (BGData)	0.0	0.0	1.7	0.0	0.0	0.16	0.20	0.15	0.09	0.20	0.33	0.15
Calibration method (Method)	0.0	0.7	0.6	0.6	1.1	0.05	2.50	0.21	0.07	0.51	0.19	0.87
Offset (UN/MI)	n/a	n/a	n/a	1.3	1.3	n/a	n/a	n/a	n/a	n/a	n/a	n/a
Multiple interactions model (MHI)	n/e	n/e	n/e	n/e	n/e	0.07	0.00	0.18	0.06	0.00	0.30	0.22
Systematic uncertainty (syst)	5.3	4.9	5.7	3.9	3.6	0.98	2.82	1.36	0.63	1.49	2.69	1.55
Statistical uncertainty (stat)	5.1	10.3	10.0	3.6	12.3	0.52	9.00	1.26	0.41	2.36	1.83	1.19
Total uncertainty	7.3	11.4	11.5	5.3	12.8	1.12	9.43	1.85	0.76	2.80	3.26	1.95



- The combination is performed using BLUE:
  - **BLUE: Best Linear Unbiased Estimate**
    - Obtain **best unbiased estimate** via **linear** combination of input measurements by minimising the total uncertainty on the combined result
  - Assume **uncertainties are Gaussian**
  - Use **categories** of uncertainties that are **uncorrelated**
  - For a given category of uncertainty:
    - take into account **correlations** between **individual measurements**

arXiv:1407.2682 [hep-ex]



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  - **BLUE**: Best Linear Unbiased Estimate
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  - Use **categories** of uncertainties that are **uncorrelated**
  - For a given category of uncertainty:
    - take into account **correlations** between **individual measurements**

This is the tricky bit,  
discuss in the following  
+ detailed cross-checks

arXiv:1407.2682 [hep-ex]



- **Jet energy scale (JES):**
  - iJES → statistical component (cf. next page)
  - aJES → response variations versus jet flavour
  - bJES → modelling of b quark jets
  - cJES → out-of-cone correction
  - dJES → response variations versus  $\eta$  and  $p_T$
  - rJES → absolute response calibration (CDF)
- **Theory and signal modelling:**
  - Signal MC generator + associated PDF, hadronisation model, underlying event + colour reconnection
- **Detector modelling**
  - Detector resolution effects, reconstruction efficiencies, b-quark jet identification (b-tagging), ...

arXiv:1407.2682 [hep-ex]



# Correlations between uncertainties

Category	$\rho_{\text{all}}$	$\rho_{\text{CDF}}$	$\rho_{\text{D}\emptyset}$	$\rho_{\text{RunXCDF}}$	$\rho_{\text{RunXD}\emptyset}$	$\rho_{\text{channel}}$	$\rho_{\text{chnIData}}$
Stat	0	0	0	0	0	0	0
iJES	0	0	0	0	1	0	0
aJES	0	0	0	1	1	0	0
bJES	1	1	1	1	1	1	1
cJES	1	1	1	1	1	1	1
rJES	0	1	1	1	1	0	0
dJES	0	0	0	1	1	0	0
LepPt	0	0	0	1	1	0	0
Signal	1	1	1	1	1	1	1
DetMod	0	1	1	1	1	0	0
B-tag	0	0	0	1	1	0	0
BGMC	0	0	0	0	0	1	0
BGData	0	0	0	0	0	0	1
Method	0	0	0	0	0	0	0
Offset	0	1	1	1	1	0	0
MHI (pile-up)	0	0	0	1	1	0	0

$\rho_{\text{CDF(D}\emptyset)}$

→  $\rho$  within an experiment

$\rho_{\text{channel}}$  →  $\rho$  within a channel

$\rho_{\text{RunXCDF(D}\emptyset)}$

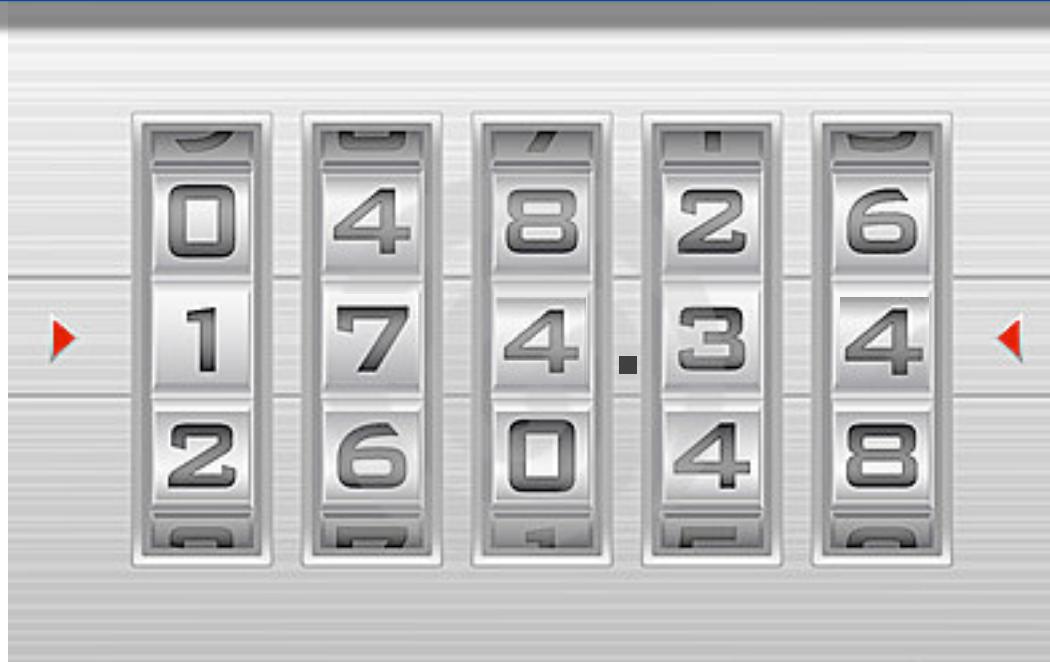
→  $\rho$  within an experiment and within an epoch

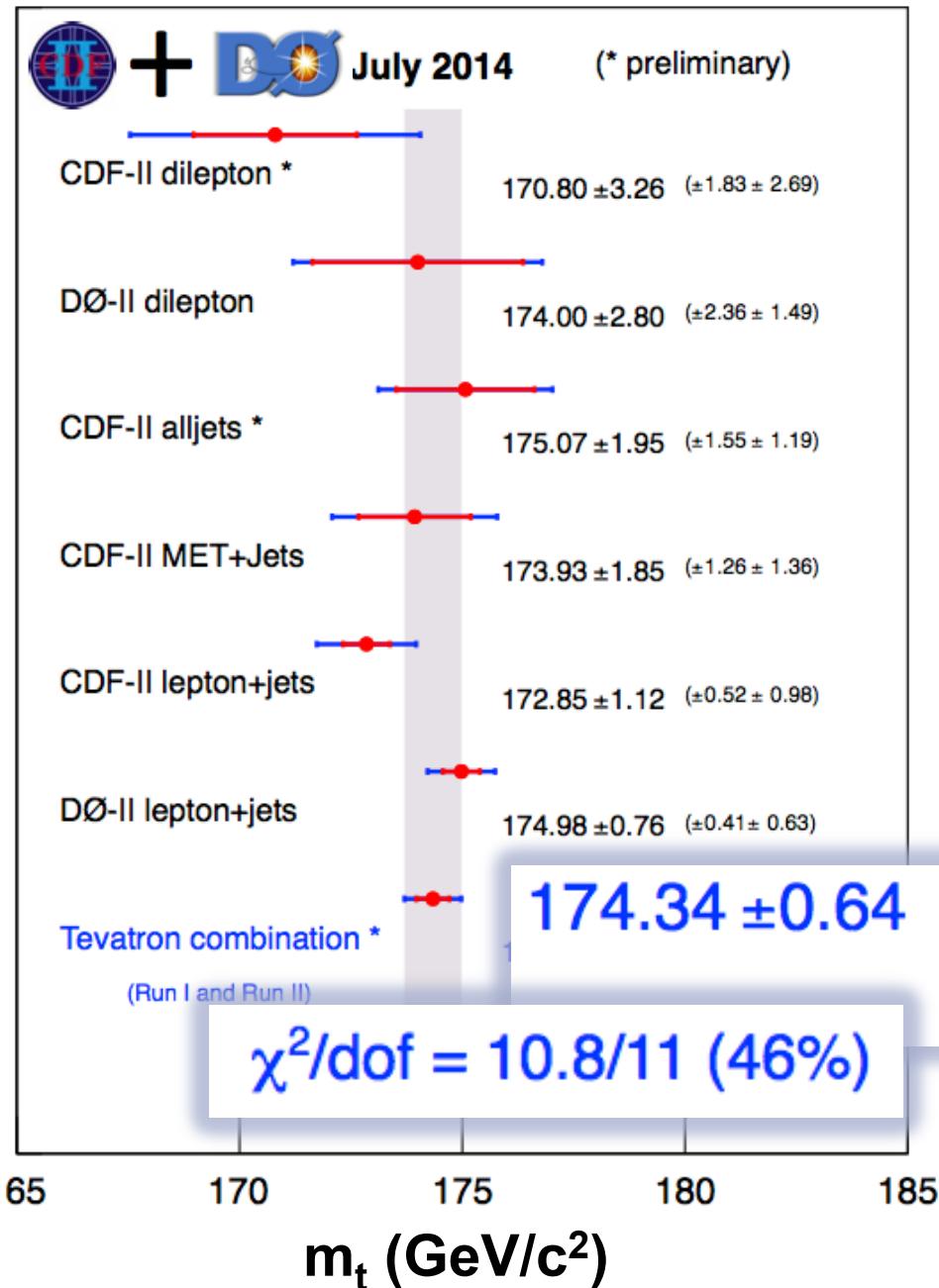
$\rho_{\text{chnIData}}$

→  $\rho$  within a channel, within an experiment, within an epoch



# Combined result



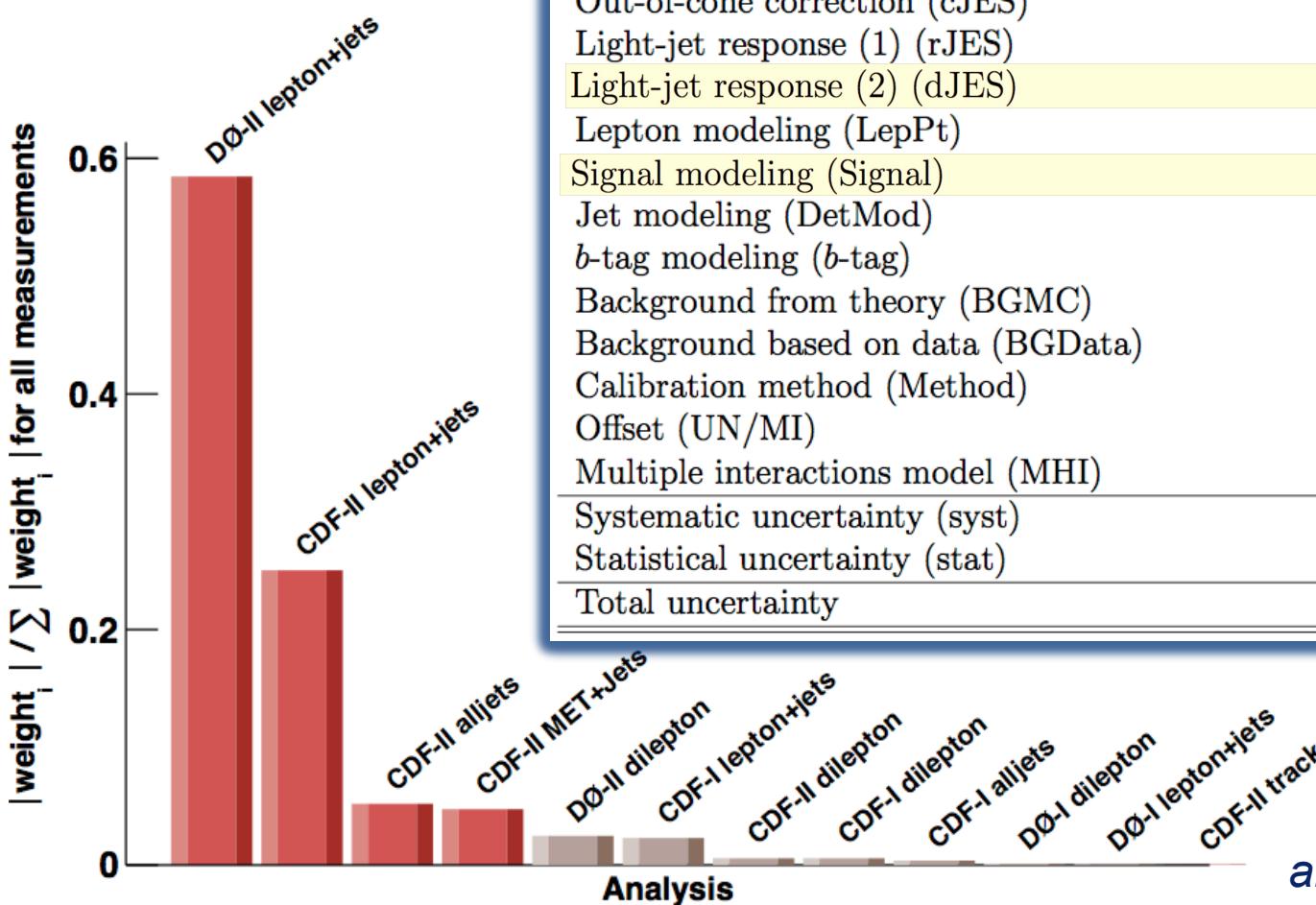
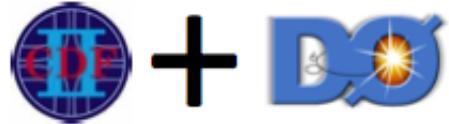


Expect further improvements as new measurements enter the game

Most precise determination of  $m_t$  comes from the Tevatron (yet)

**0.37%**  
**total uncertainty**

arXiv:1407.2682 [hep-ex]

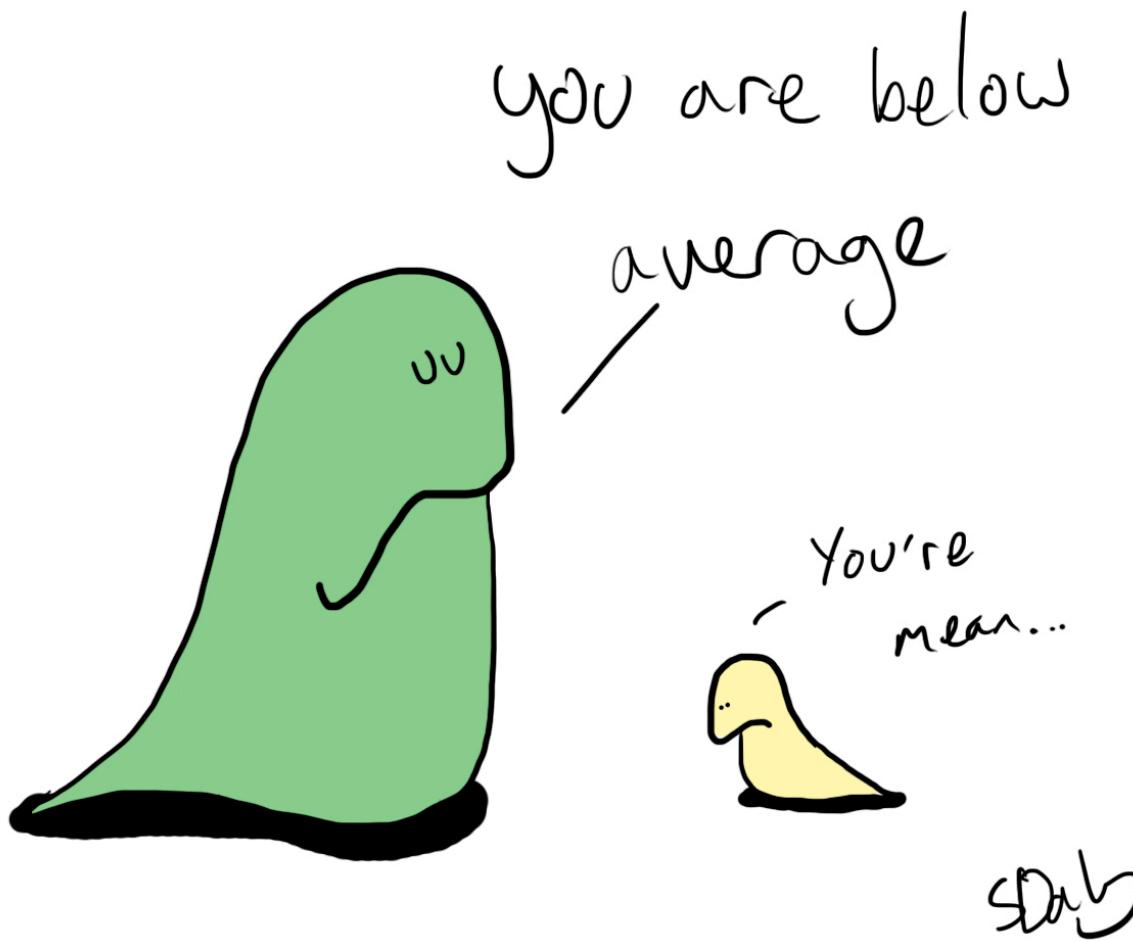


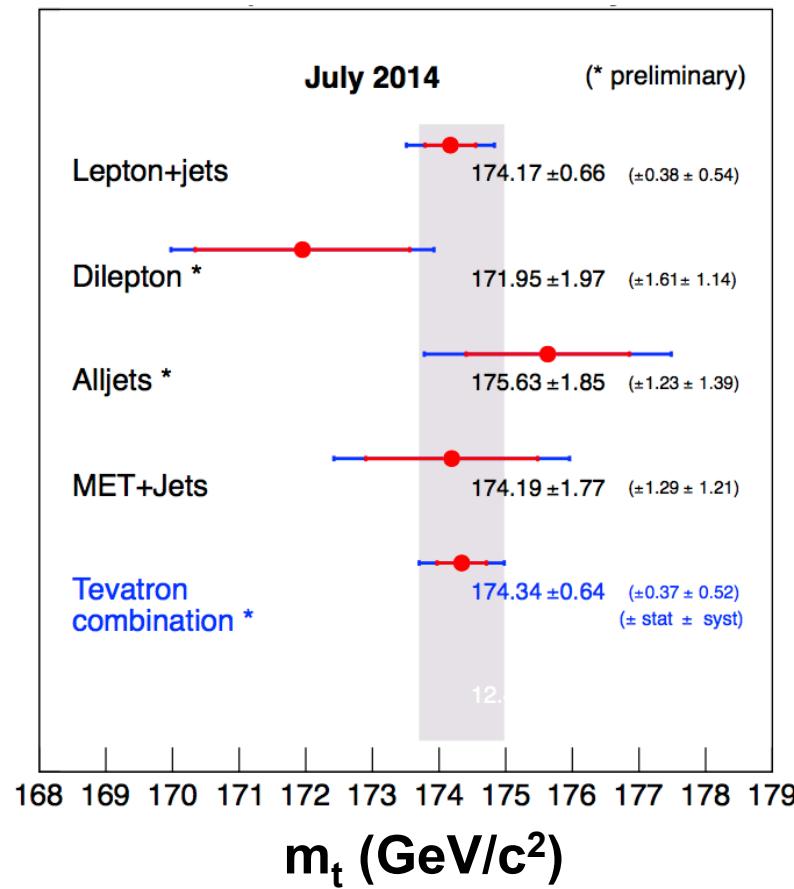
	Tevatron combined values ( $\text{GeV}/c^2$ )
$M_t$	174.34
In situ light-jet calibration (iJES)	0.31
Response to $b/q/g$ jets (aJES)	0.10
Model for $b$ jets (bJES)	0.10
Out-of-cone correction (cJES)	0.02
Light-jet response (1) (rJES)	0.05
Light-jet response (2) (dJES)	0.13
Lepton modeling (LepPt)	0.07
Signal modeling (Signal)	0.34
Jet modeling (DetMod)	0.03
$b$ -tag modeling ( $b$ -tag)	0.07
Background from theory (BGMC)	0.04
Background based on data (BGData)	0.08
Calibration method (Method)	0.07
Offset (UN/MI)	0.00
Multiple interactions model (MHI)	0.06
Systematic uncertainty (syst)	0.52
Statistical uncertainty (stat)	0.37
Total uncertainty	0.64

arXiv:1407.2682 [hep-ex]



# Cross-checks





Parameter	Value ( $\text{GeV}/c^2$ )	Correlations			
		$M_t^{\text{all-jets}}$	$M_t^{\ell+\text{jets}}$	$M_t^{\ell\ell}$	$M_t^{\text{MET}}$
$M_t^{\text{all-jets}}$	$175.63 \pm 1.85$	1.00			
$M_t^{\ell+\text{jets}}$	$174.17 \pm 0.66$	0.21	1.00		
$M_t^{\ell\ell}$	$171.95 \pm 1.97$	0.21	0.41	1.00	
$M_t^{\text{MET}}$	$174.19 \pm 1.77$	0.11	0.23	0.18	1.00

Probability of the  $\chi^2$ -statistic  
of pair-wise differences:  
22%, 42%, 99%, 13%, 35%, 55%

	Run I published					Run II published					Run II prel.	
	CDF			DØ		CDF			DØ		CDF	all-jets
	$\ell+\text{jets}$	$\ell\ell$	all-jets	$\ell+\text{jets}$	$\ell\ell$	$\ell+\text{jets}$	$L_{XY}$	MEt	$\ell+\text{jets}$	$\ell\ell$		
Pull	0.24	-0.61	+1.01	+1.09	-0.46	-1.64	-0.791	-0.24	+1.60	-0.13	-1.11	0.39
Weight [%]	-2.6	-0.7	-0.4	-0.1	-0.14	+28.8	+0.1	+5.5	+67.2	-2.9	-0.66	+6.0

arXiv:1407.2682 [hep-ex]



**Anything  
to be  
learned?**





## Disclaimer:

This is just my **personal view** what could be useful in the context of  $m_t$  measurements in the next decade

This refers to “conventional” measurements of  $m_t$  (but may also apply to e.g.  $m_t$  extraction from  $\sigma_{tt}$ )

- 1) It is crucial to understand the **origin of systematic uncertainties**:
  - Factorise effects **into different categories**, if applicable
  - The total uncertainty is a quadratic sum of all effects
    - Double-counting counts!



- 2) Finite statistics of MC samples results in a sizable statistical component of systematic uncertainties
  - Increase size of MC samples
    - Typical statistical component @ LHC: 0.1 GeV
    - The bottle neck can be:
      - Generation and simulation of MC events
      - Their analysis with advanced method like the ME technique
- 3) Ironically, the dominant limitation in precision comes from the soft part of the tt event
  - We need to better understand and constrain the hadronisation model
  - OR we fit the b-quark JES in-situ `a la ATLAS
    - No free lunch: increased statistical uncertainty
      - Not an issue with Run II data (hopefully!)
      - Clearly an issue with MC samples



- 4) We need to improve our understanding of the signal modeling part (hard+soft)
  - Include new generators as they become available, e.g.:
    - Sherpa, aMC@NLO, herwig++, pythia8, etc
  - Reject models which are in tension with datasets used for tuning of generators or parton shower simulation
    - E.g. fHerwig?
- 5) We are eagerly awaiting MC generators which can simulate the full  $t\bar{t}$  decay at NLO
  - Then the finite width of t propagator is accounted for
    - → well-defined concept of  $m_t$ !



- 6) The ME technique may be less sensitive to systematic uncertainties
  - Evaluates the impact of an uncertainty in context of a concrete model for  $t\bar{t}$  production described by the ME
  - Canonical example:
    - Reduced sensitivity of LO ME to  $t\bar{t}$  events with initial/final state radiation due to lower  $P_{\text{sig}}$
- 7) Measure  $m_t$  in various regions of phase space
  - Check for biases
    - Reject models
  - Evaluate systematic uncertainties from data



GAME OVER

BONUS MATERIAL



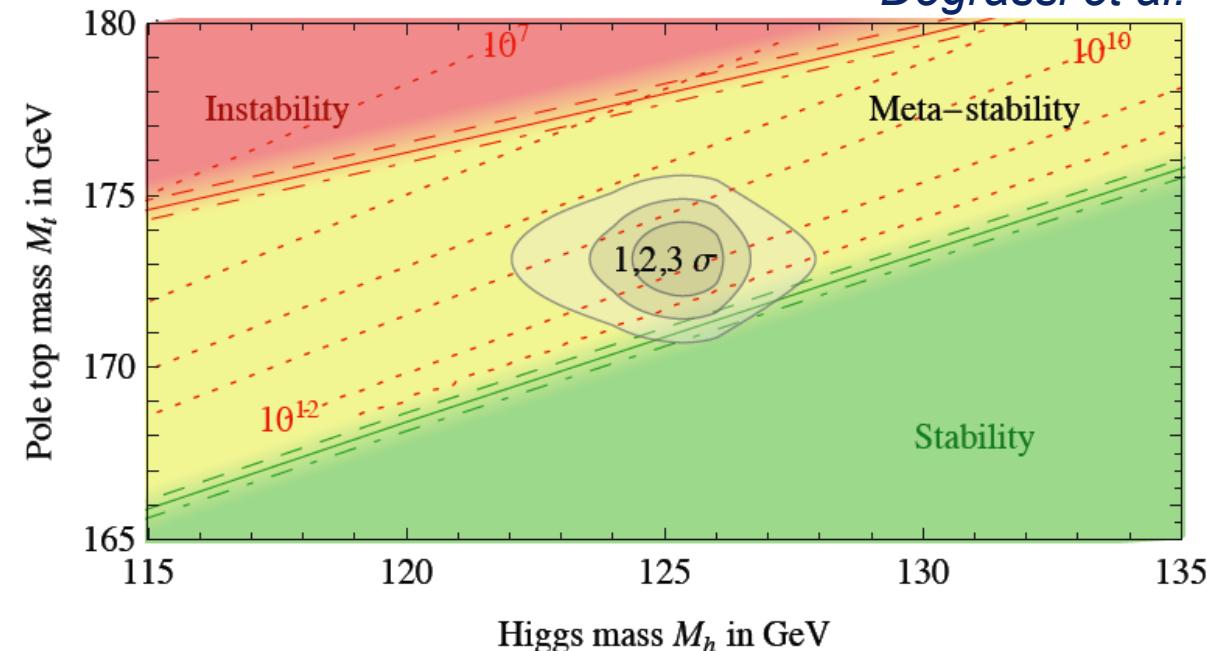
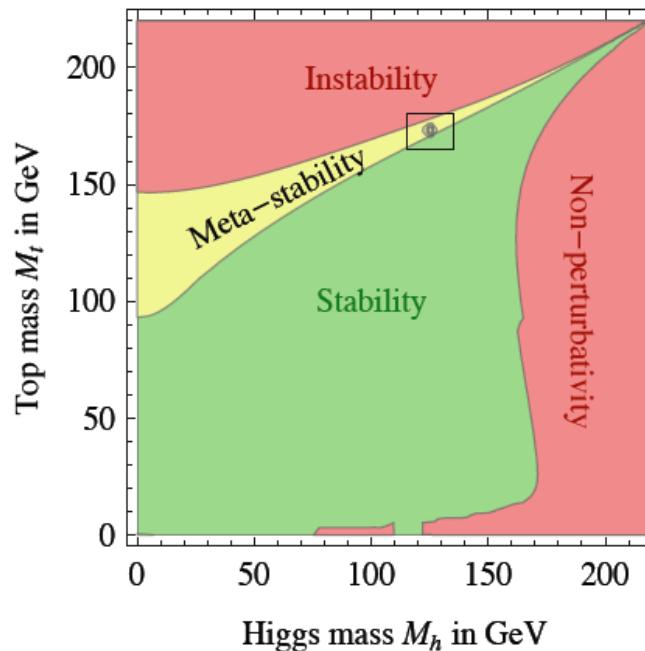
- Current World average:

$$m_t = 173.34 \pm 0.76 \text{ GeV}$$

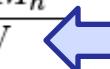
*arXiv:1403.4427 [hep-ex]*

- Assuming no statistical correlation between this result and the combination
  - Taking full uncertainty for the Tevatron average
  - Taking statistical uncertainty only for this measurement
- Consistency at 1.71 SD level (p-value of 3.1%)

- The phase diagram of the Universe:


*Degrasse et al.*

Type of error	Estimate of the error	Impact on $M_h$
$M_t$	experimental uncertainty in $M_t$	$\pm 1.4$ GeV
$\alpha_s$	experimental uncertainty in $\alpha_s$	$\pm 0.5$ GeV
<b>Experiment</b>	<b>Total combined in quadrature</b>	<b><math>\pm 1.5</math> GeV</b>
$\lambda$	scale variation in $\lambda$	$\pm 0.7$ GeV
$y_t$	$\mathcal{O}(\Lambda_{\text{QCD}})$ correction to $M_t$	$\pm 0.6$ GeV
$y_t$	QCD threshold at 4 loops	$\pm 0.3$ GeV
RGE	EW at 3 loops + QCD at 4 loops	$\pm 0.2$ GeV
<b>Theory</b>	<b>Total combined in quadrature</b>	<b><math>\pm 1.0</math> GeV</b>





- Apply dedicated corrections for:

- **u, d, c, s** quark jets
- **b** quark jets
- **gluon** jets

- The correction is given by:

$$F_{\text{corr}} = \frac{1}{\langle F \rangle_{\gamma+\text{jet}}} \cdot \frac{\sum_i E_i \cdot R_i^{\text{data}}}{\sum_i E_i \cdot R_i^{\text{MC}}}$$

- $F_{\text{corr}}$  preserves default JES by construction
- Derive single particle responses  $R_i$  in data/MC for:
  - $\gamma, e^\pm, \mu^\pm, \pi^\pm, K^\pm, K_0^S, K_0^L, p^\pm, n$  and  $\Lambda$
  - (Keep in mind that DØ corrects jet energies to particle level in data and in MC)

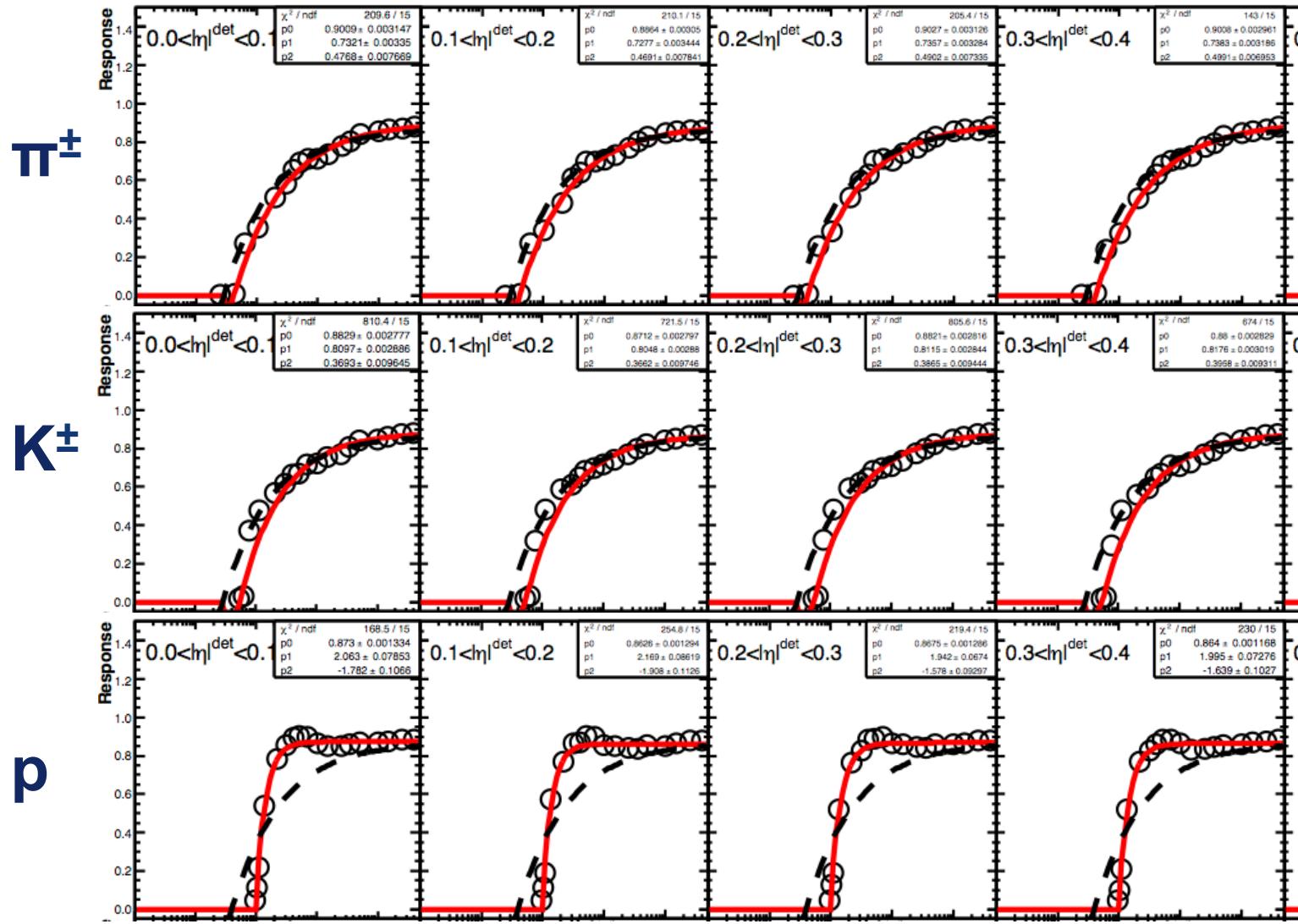
DØ Coll, Section 14 in arXiv:1312.6873 [hep-ex], submitted to NIM



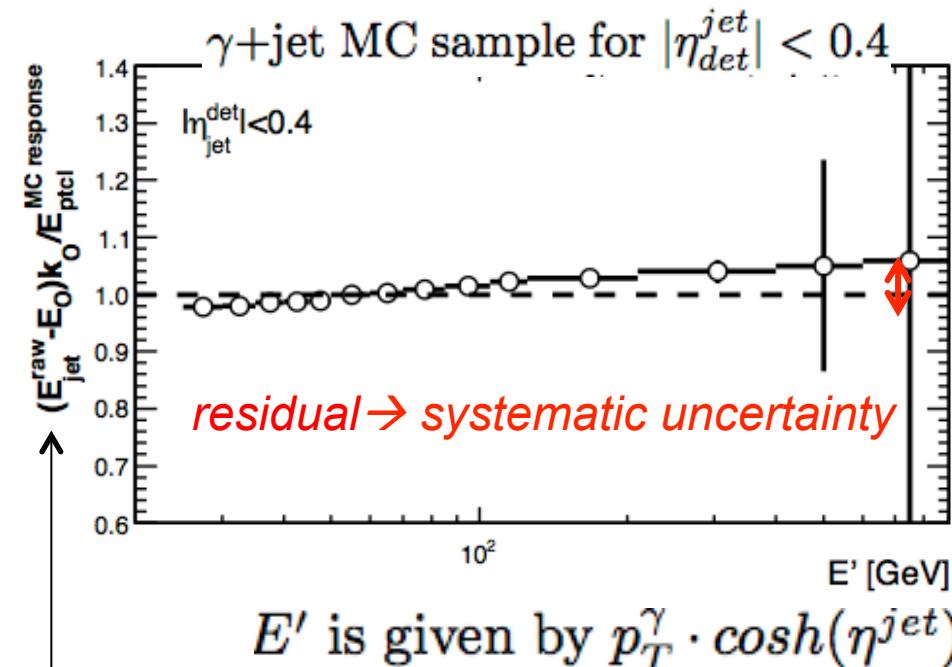
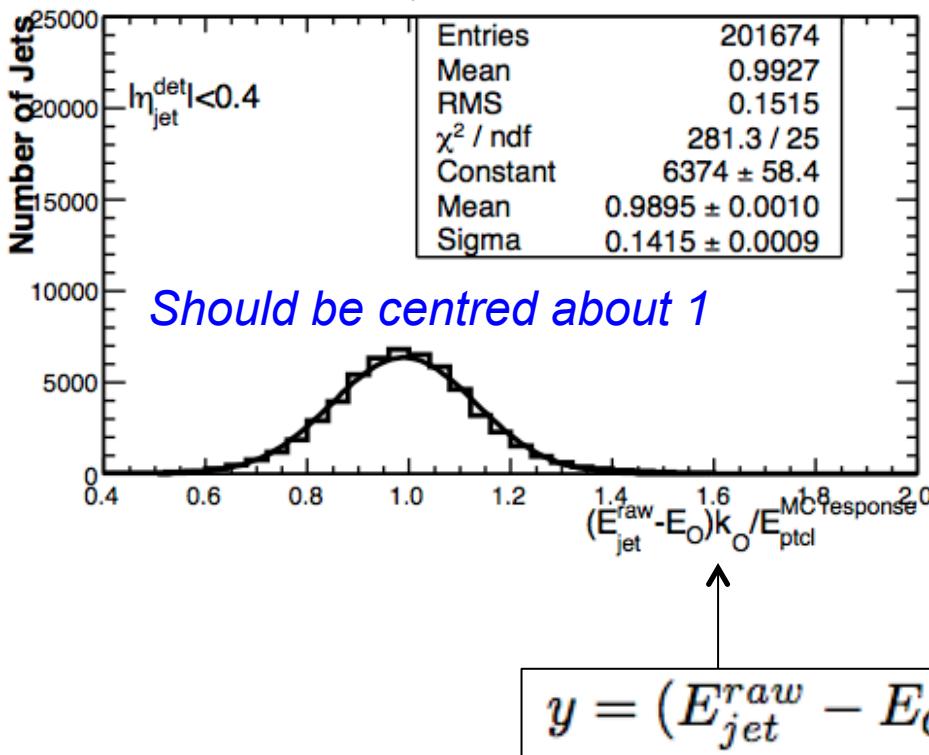
- Derive single particle responses  $R_i$  in MC:
  - Use single particle MC samples for each of
    - $\gamma, e^\pm, \mu^\pm, \pi^\pm, K^\pm, K_0^S, K_0^L, p^\pm, n$  and  $\Lambda$
  - Using:
    - Zero energy noise suppression off for default
    - (Noise suppression on → systematic uncertainty)
  - Fit with appropriate function:
    - $e, \mu, \gamma$  (not shown):
      - Calibrated separately and have one function each
    - For all hadrons:
      - Response function is (but different fit parameters!):
        - $R_h^{MC} = p_h^0 \cdot \left[ 1 - p_h^1 \cdot (E/0.75)^{p_h^2 - 1} \right]$  if  $p_T > m_h$ ; 0 if  $p_T < m_h$

DØ Coll, Section 14 in arXiv:1312.6873 [hep-ex], submitted to NIM

- Few example fits of MC response for three particles:



- Closure test: check that  $\sum_i E_i \cdot R_i^{MC}$  describes the raw offset-corrected energy  $(E_{jet}^{raw} - E_O) \cdot k_O$  correctly
  - Example:



$E_{jet}^{raw}$  is the raw jet energy

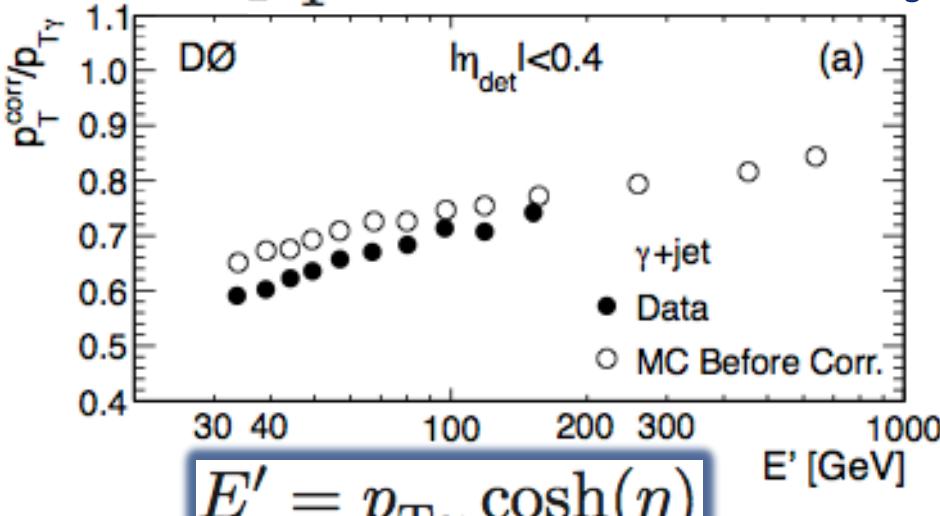
$E_O$  is the offset correction for noise and pile-up (in- and out-of-time)

$k_O$  is the correction for noise suppression bias & only needed to perform closure test

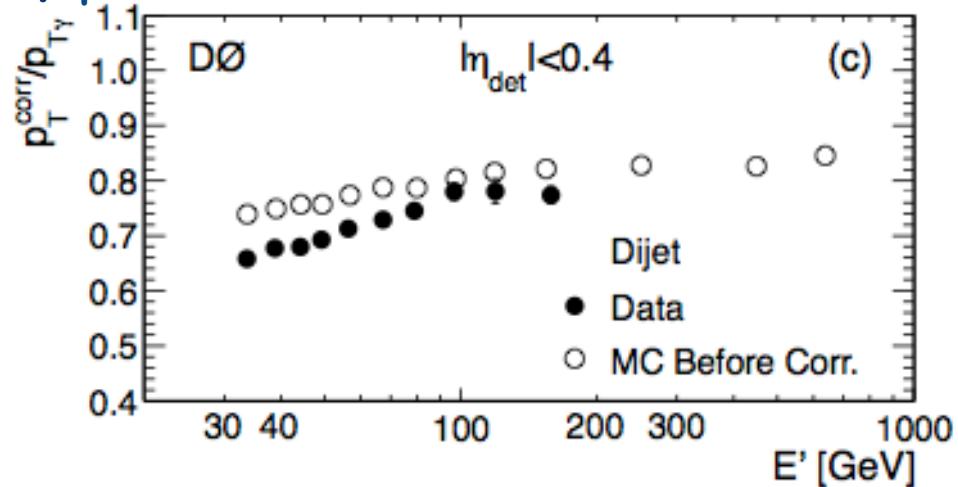


- Deriving single particle responses in data:
    - For  $e, \mu, \gamma$ :
      - Assume perfect modelling of detector response by MC
        - Systematic uncertainty estimated on this assumption
    - For hadrons  $\pi^\pm, K^\pm, K_0^S, K_0^L, p^\pm, n$  and  $\Lambda$ :
      - Basic shapes in  $E, \eta$  (i.e. per-hadron fit parameters) from MC
      - Fit **unique** (not per-bin or hadron) parameters  $A, B, C$ :
- $R_h^{data} = C \cdot p_h^0 \cdot \left[ 1 - A \cdot p_h^1 \cdot (E/0.75)^{p_h^2 + B - 1} \right]$  if  $p_T > m_h$ ; 0 if  $p_T < m_h$
- Identical to  $R_h^{MC}$  for  $A = C = 1$  and  $B = 0$ .
  - Find an optimal set of  $A, B, C$  to tune MC jet responses such that the ratios  $p_{T, corr}^{jet}/p_T^\gamma$  are consistent in data and MC
    - Use the particle composition of the jet from MC as a function of the jet energy and  $\eta$
    - Measure  $p_{T, corr}^{jet}/p_T^\gamma$  in data and MC samples enriched with isolated photons (“ $\gamma$ +jet”) and inverted photon isolation (“dijet”)
    - Here,  $p_{T, corr}^{jet}$  is reconstructed jet  $p_T$  with offset correction

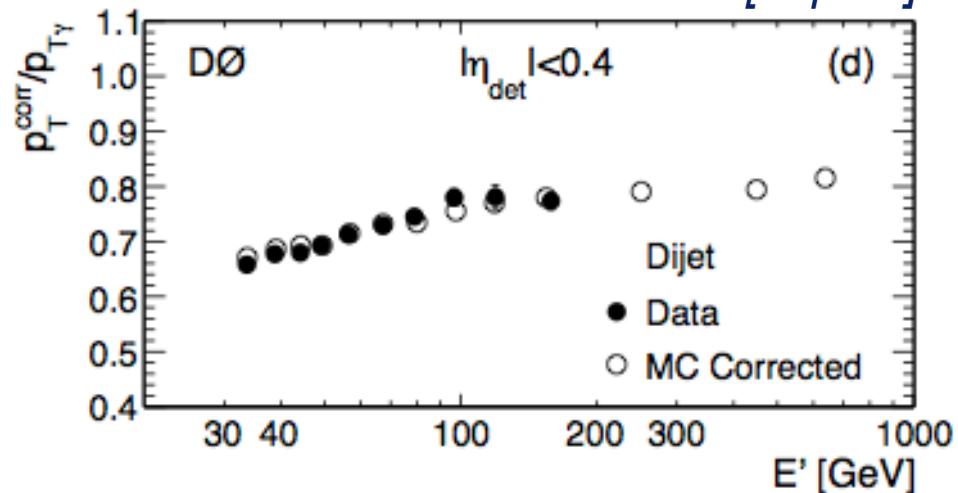
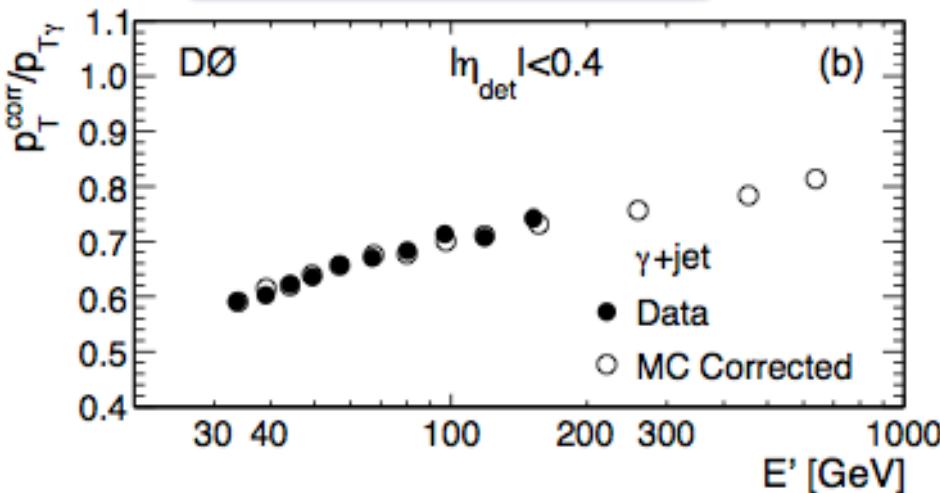
- Closure test of the flavour-dependent response:  
-  $p_T^{\text{corr}}$  is reconstructed jet  $p_T$  with offset correction



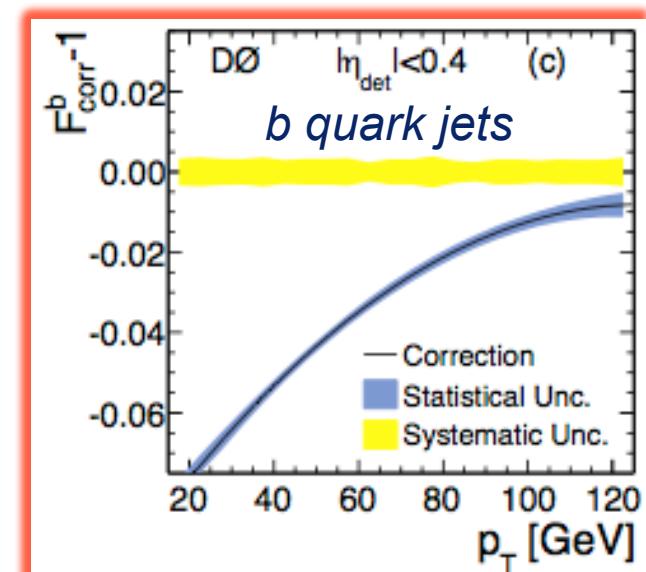
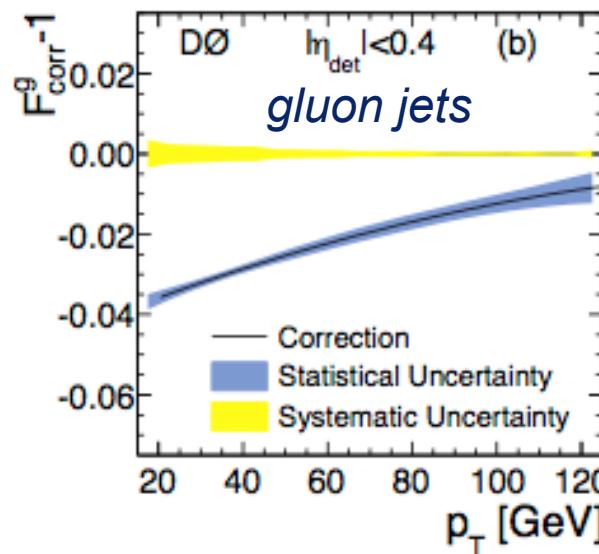
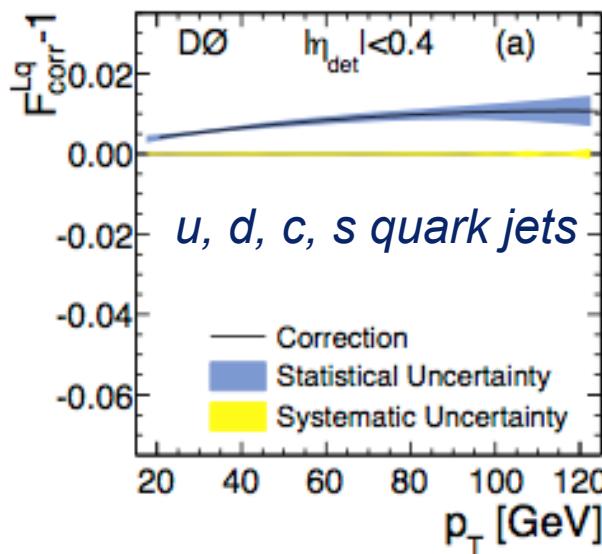
$$E' = p_{T\gamma} \cosh(\eta)$$



arXiv:1312.6873 [hep-ex]



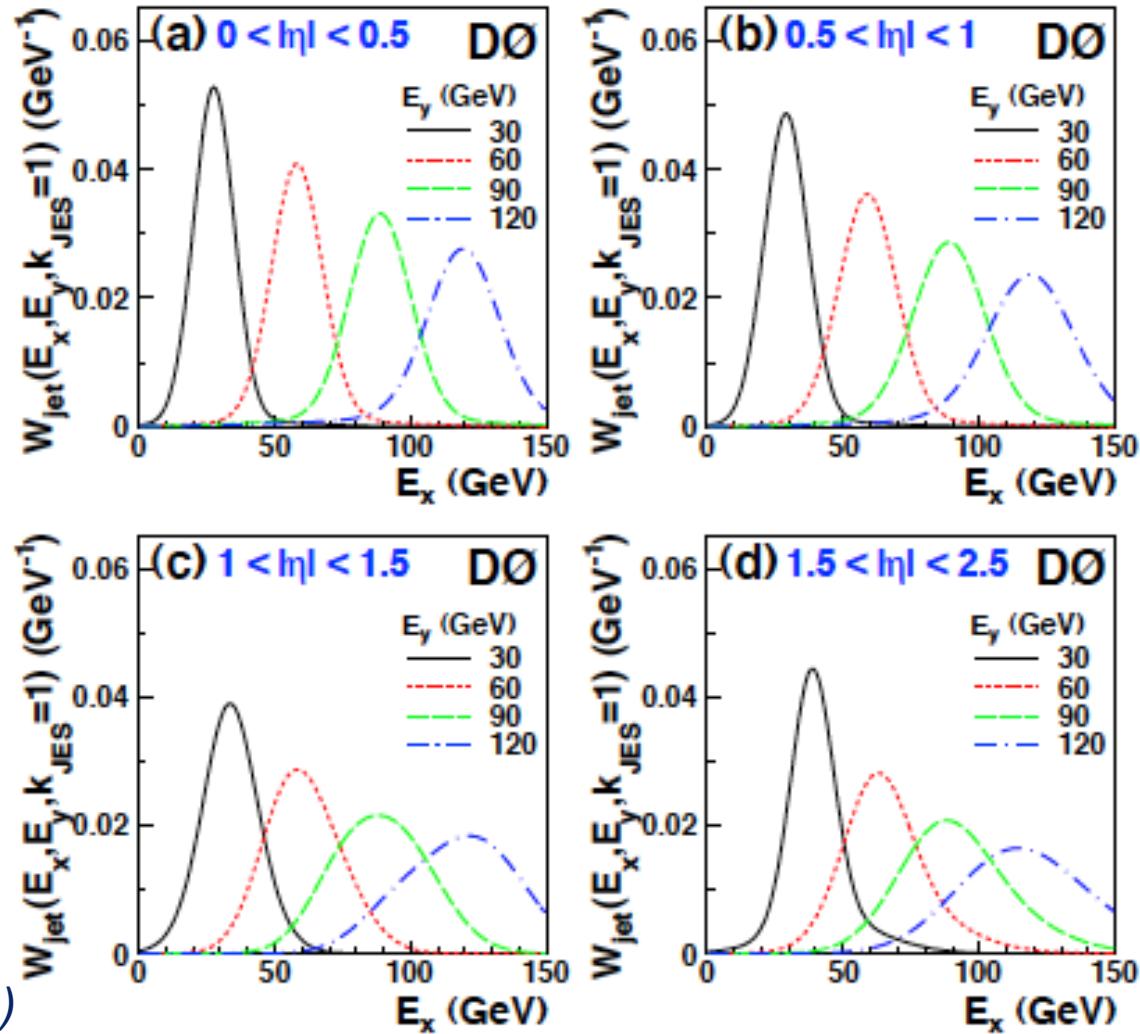
- The final flavour-dependent correction:



- The correction accounts for the difference in JES for b quark jets and light quark jets:
  - Substantial reduction of one of the dominant systematic uncertainties!

D $\emptyset$  Coll, Section 14 in arXiv:1312.6873 [hep-ex], submitted to NIM

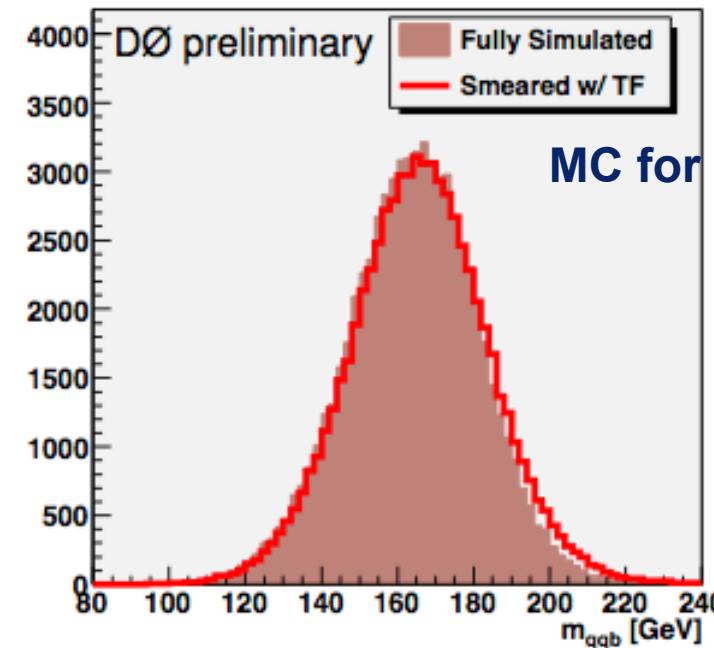
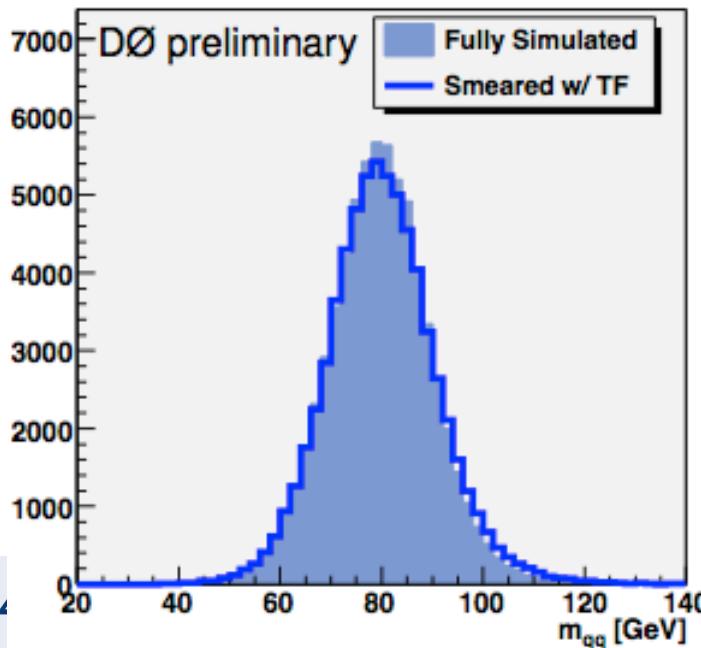
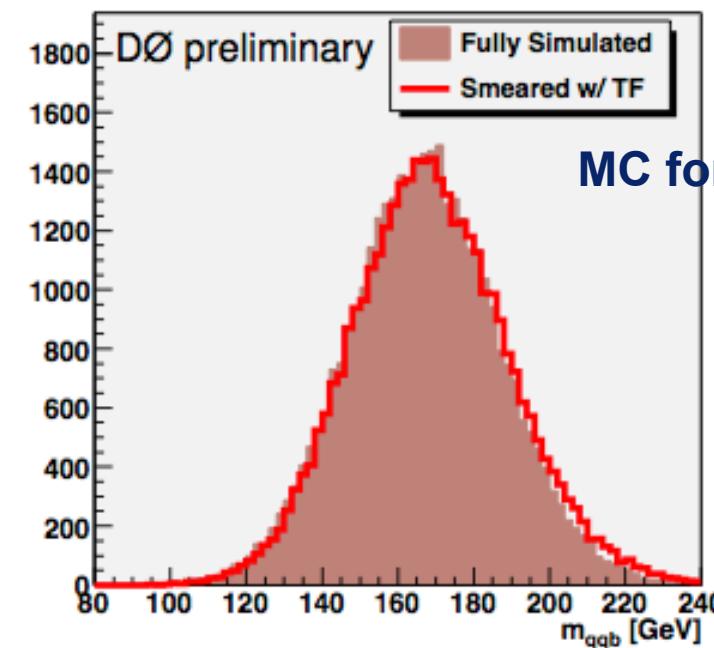
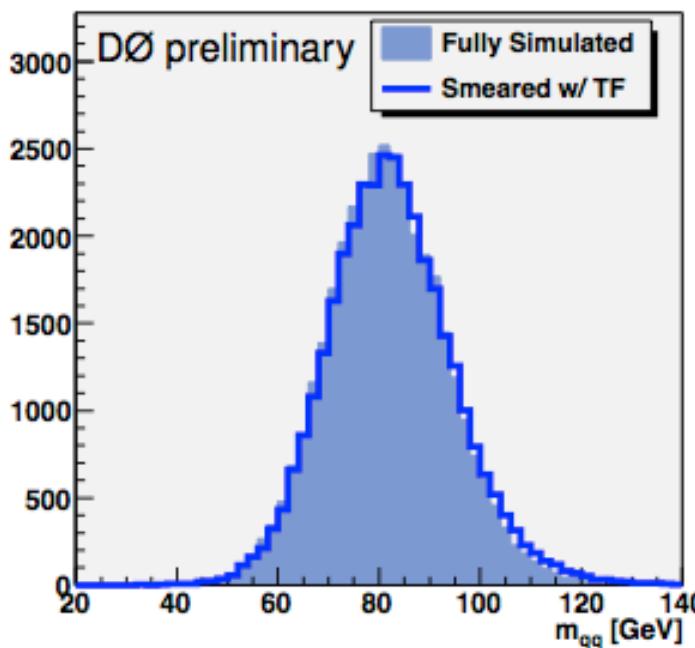
- The Transfer Functions  $W(x, y; JES)$  relate parton-level quantities to reconstruction-level ones
- Some typical examples for light quark jets from [1]



[1] DØ Coll, PRD 84, 032004 (2011)

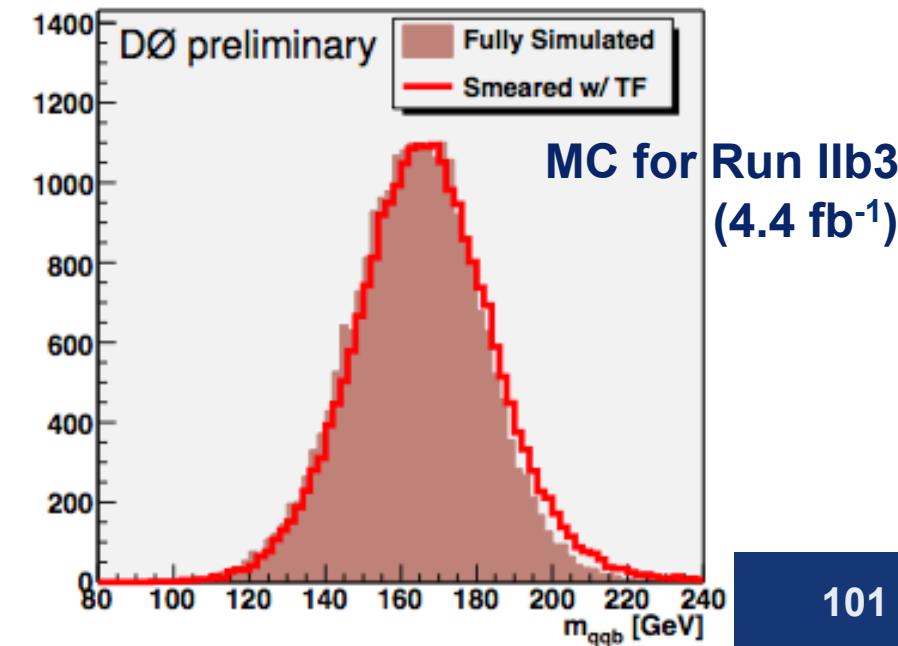
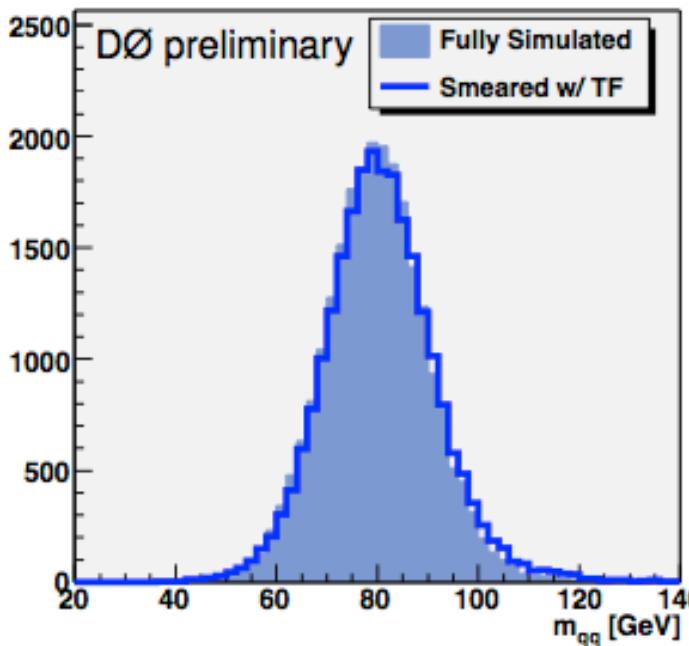
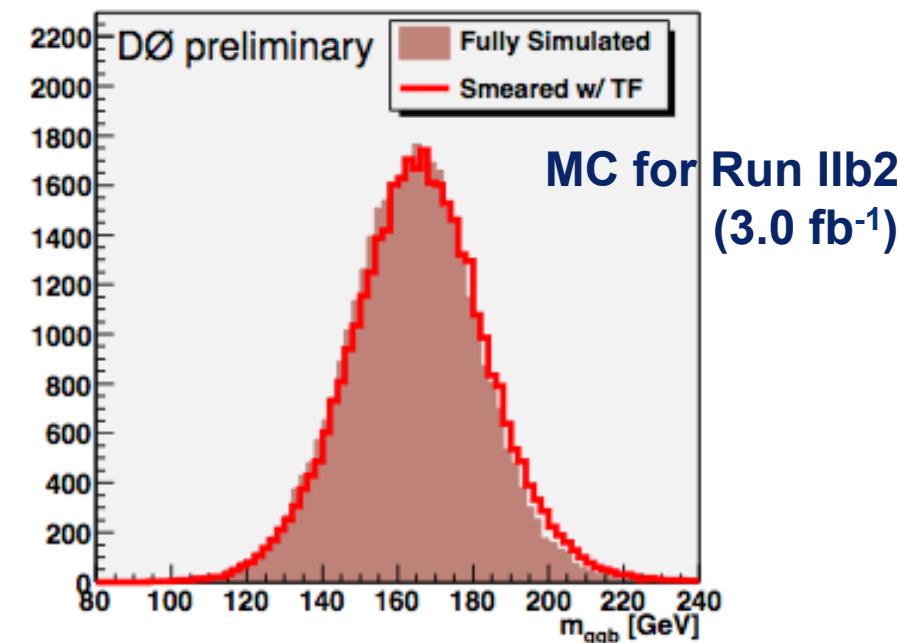
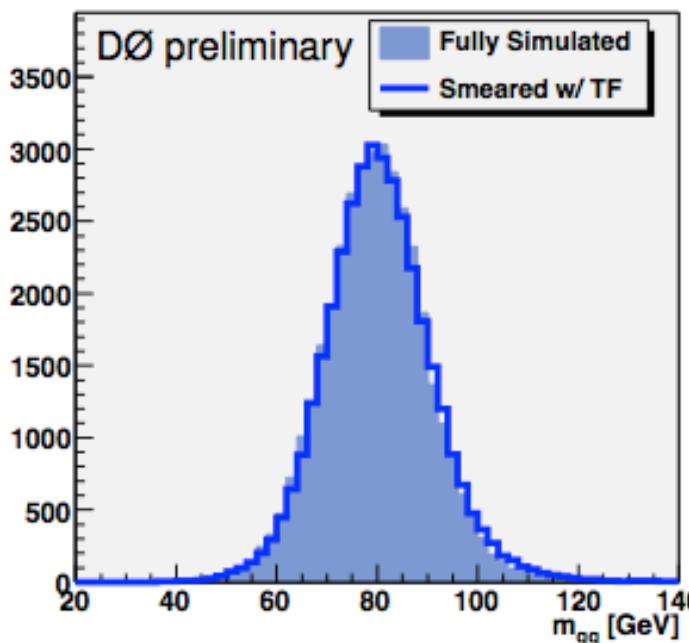


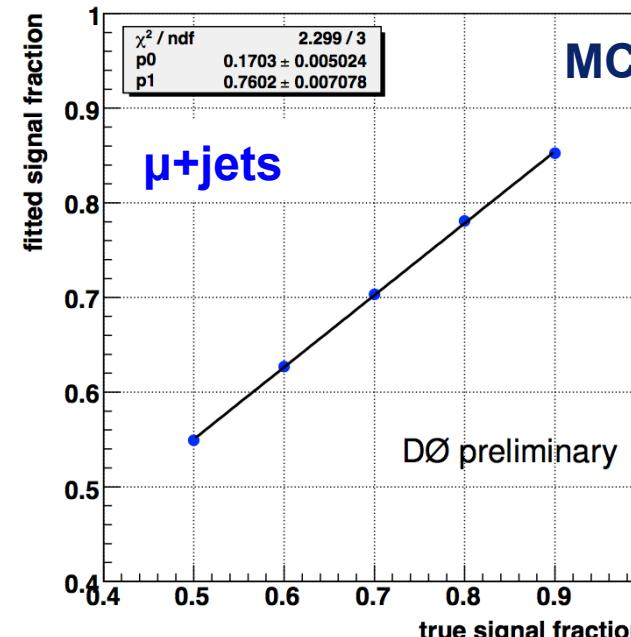
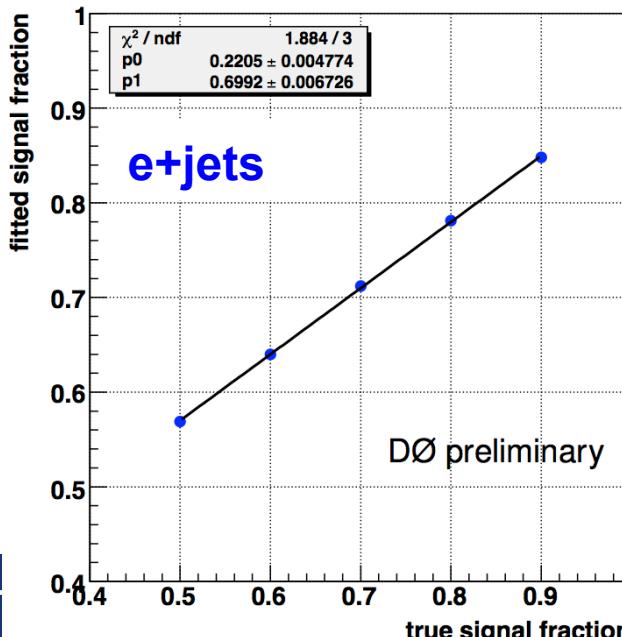
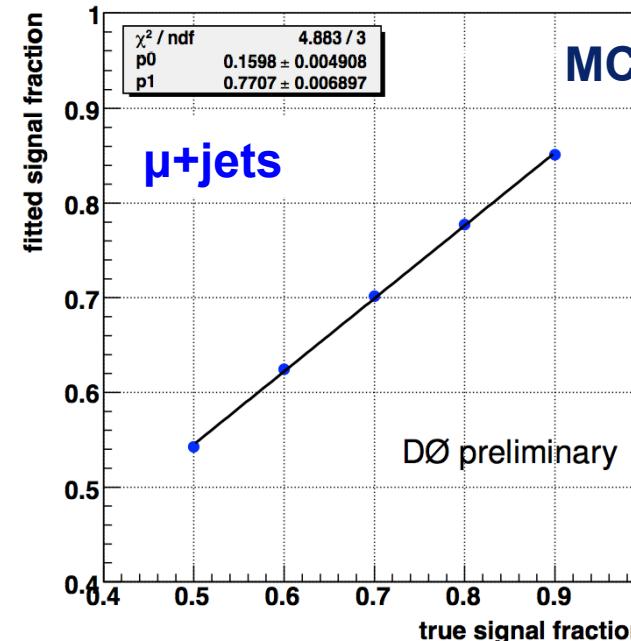
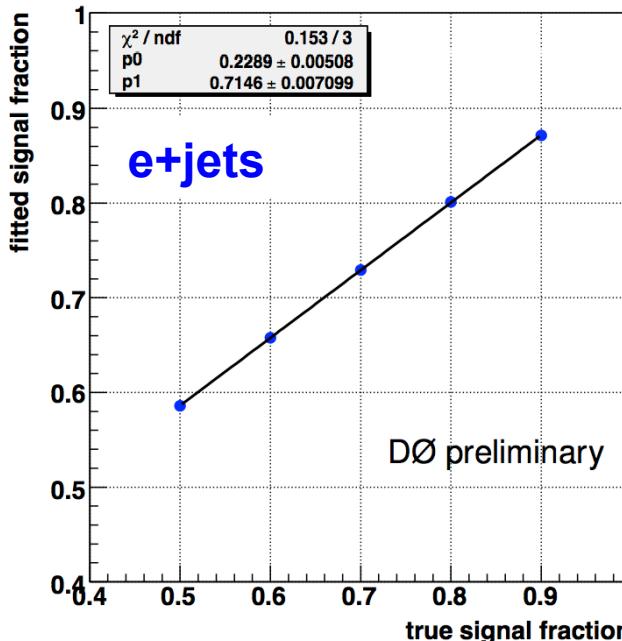
# Transfer function validation





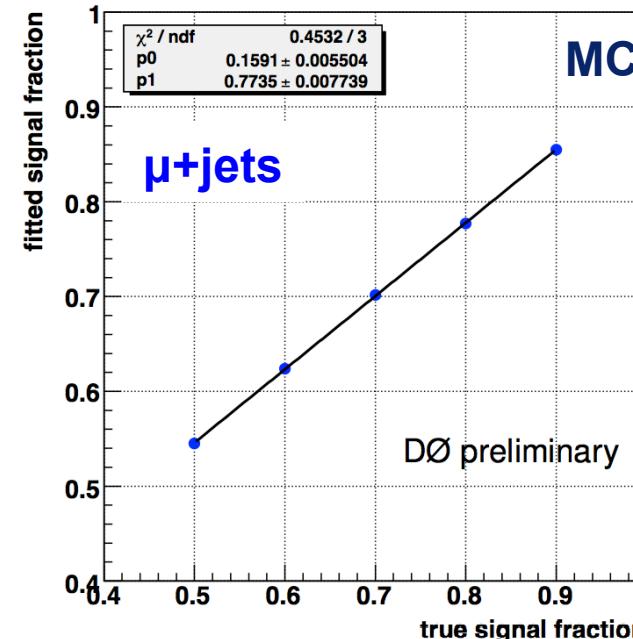
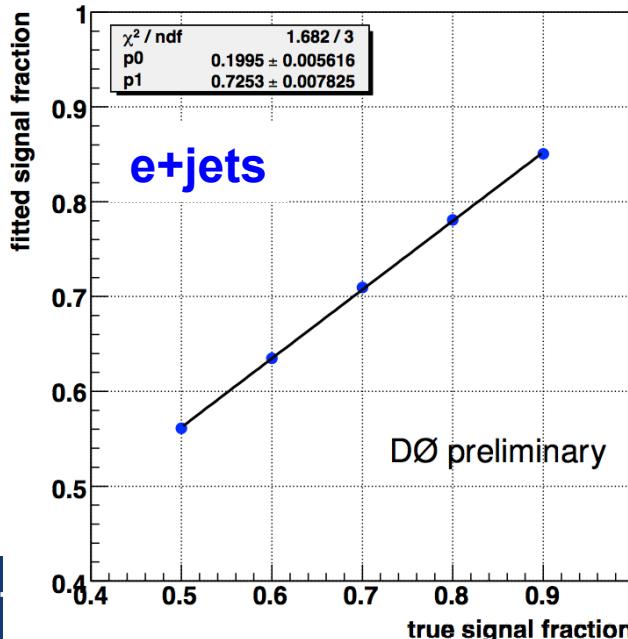
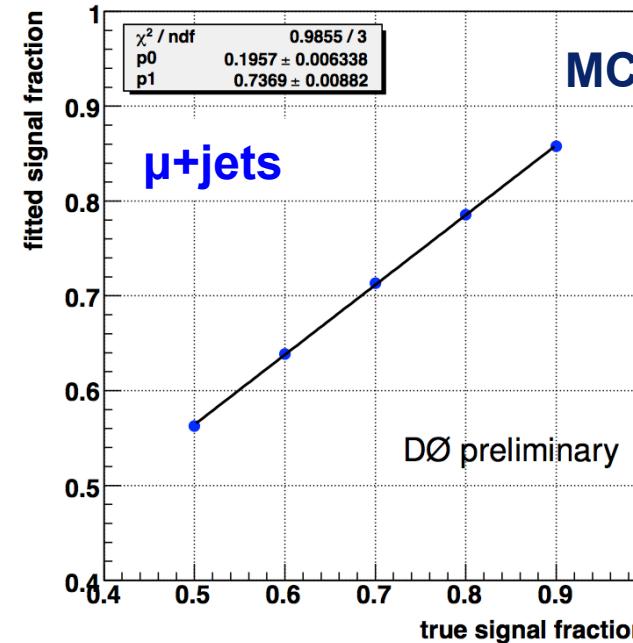
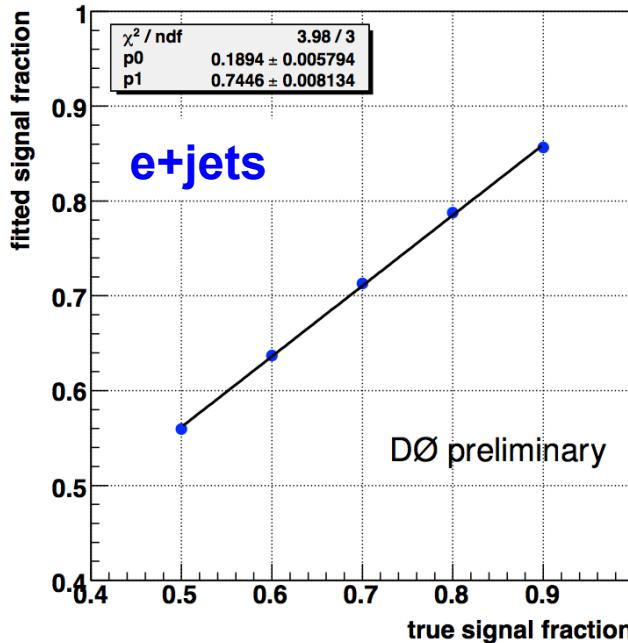
# Transfer function validation





MC for Run Ila  
(1.1  $\text{fb}^{-1}$ )

MC for Run IIb1  
(1.3  $\text{fb}^{-1}$ )

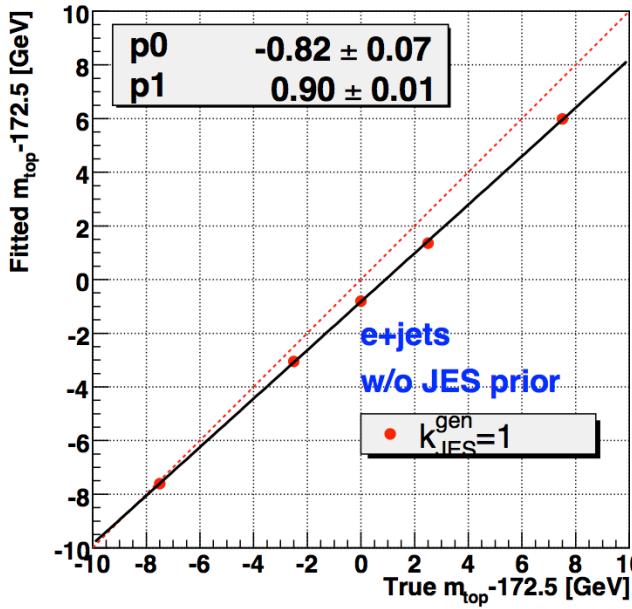


MC for Run IIb2  
(3.0  $\text{fb}^{-1}$ )

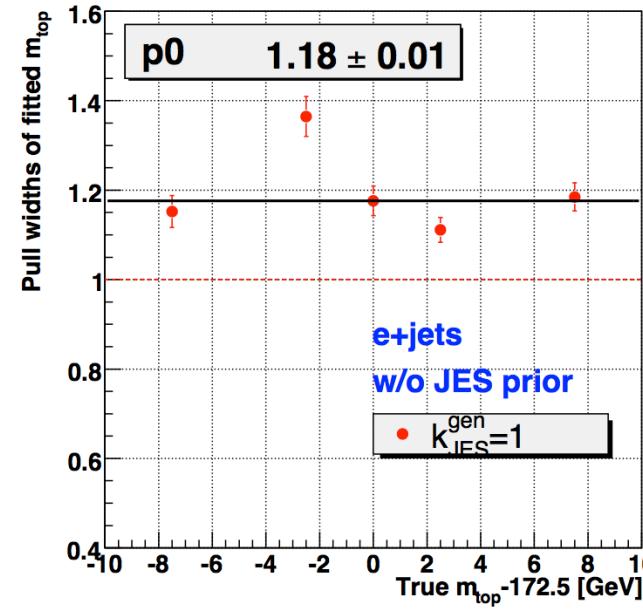
MC for Run IIb3  
(4.4  $\text{fb}^{-1}$ )



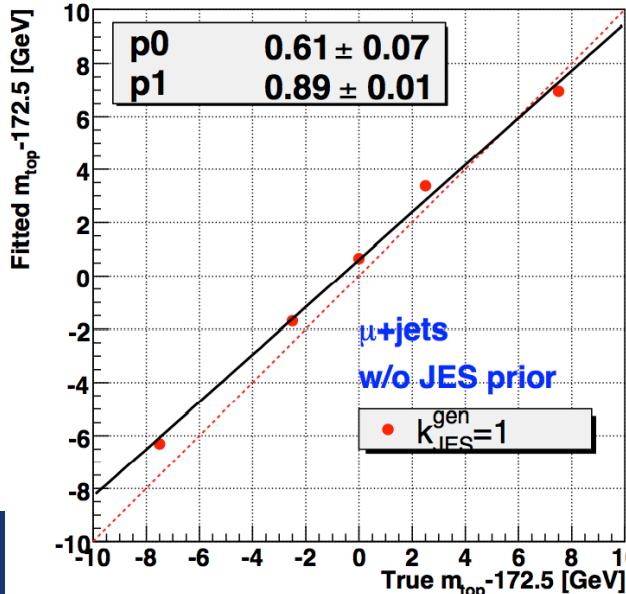
DØ Run Ila Preliminary



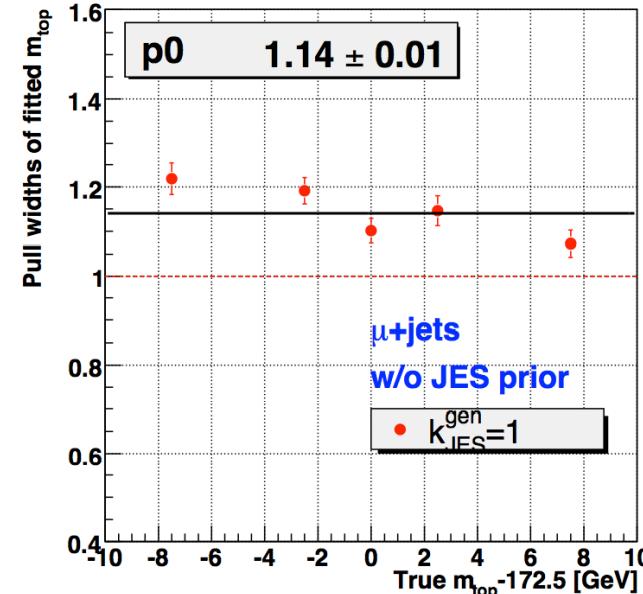
DØ Run Ila Preliminary



DØ Run Ila Preliminary

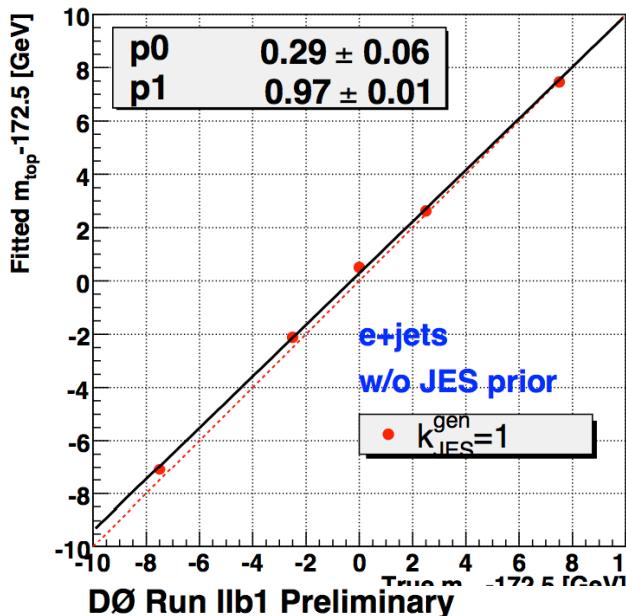


DØ Run Ila Preliminary

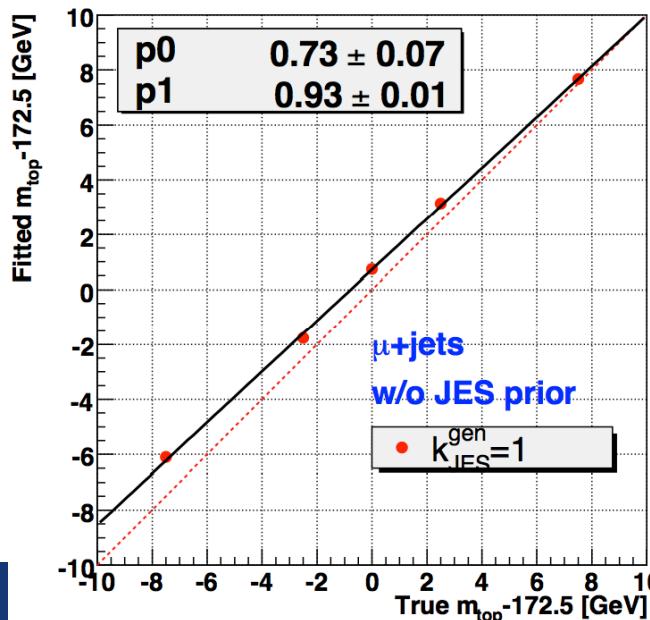




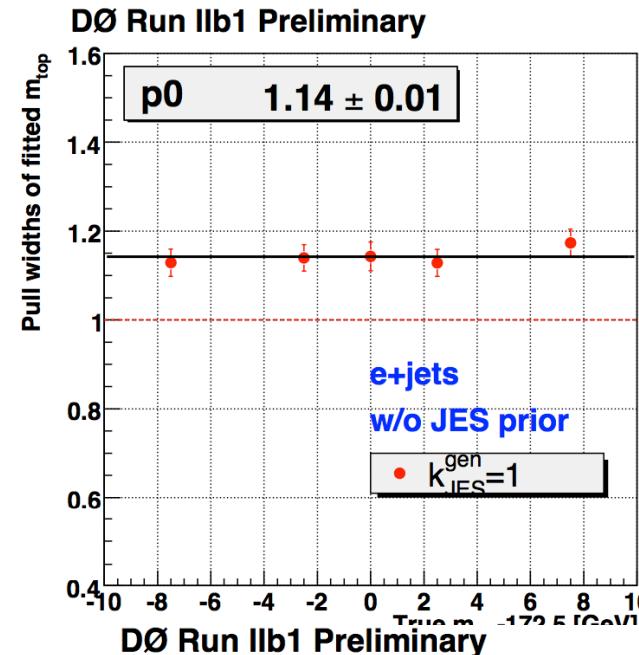
DØ Run IIb1 Preliminary



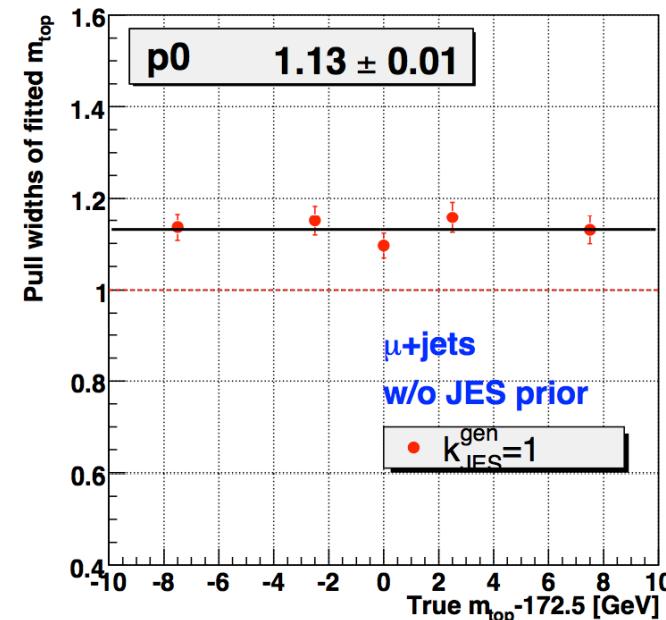
DØ Run IIb1 Preliminary

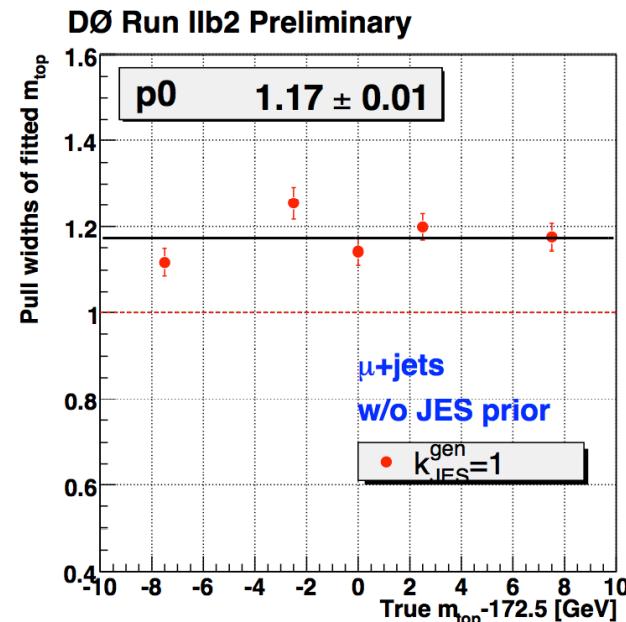
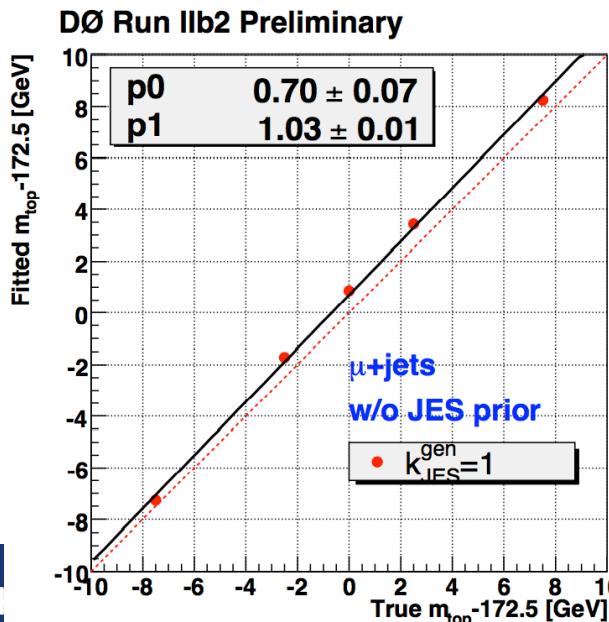
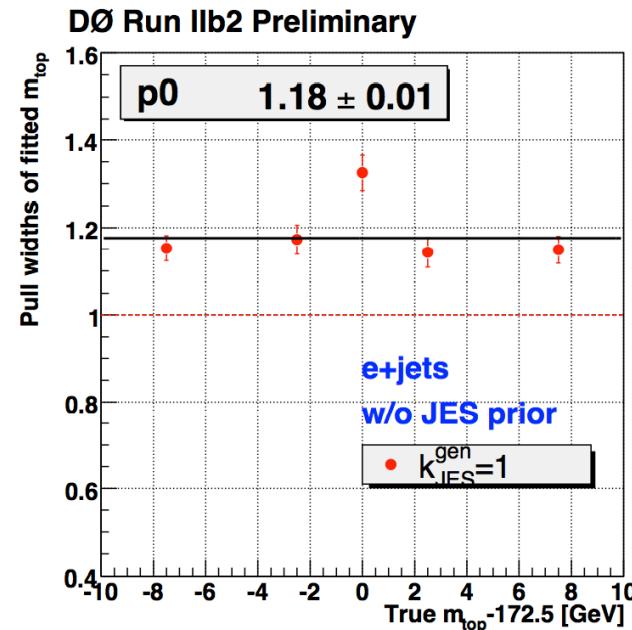
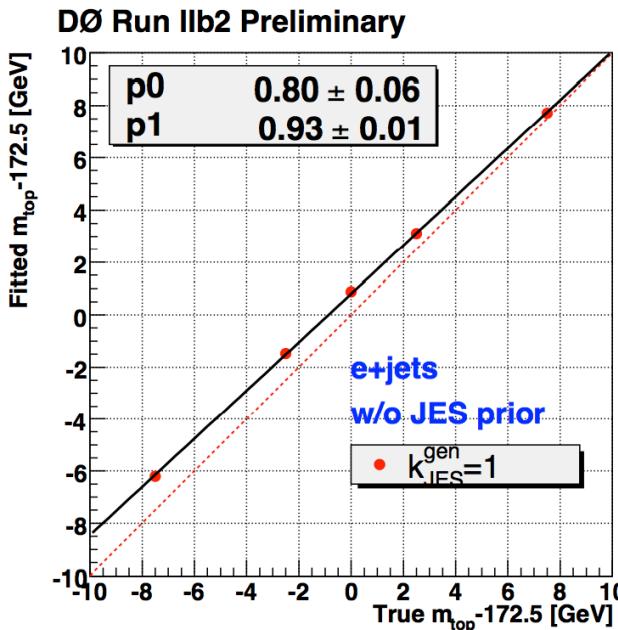


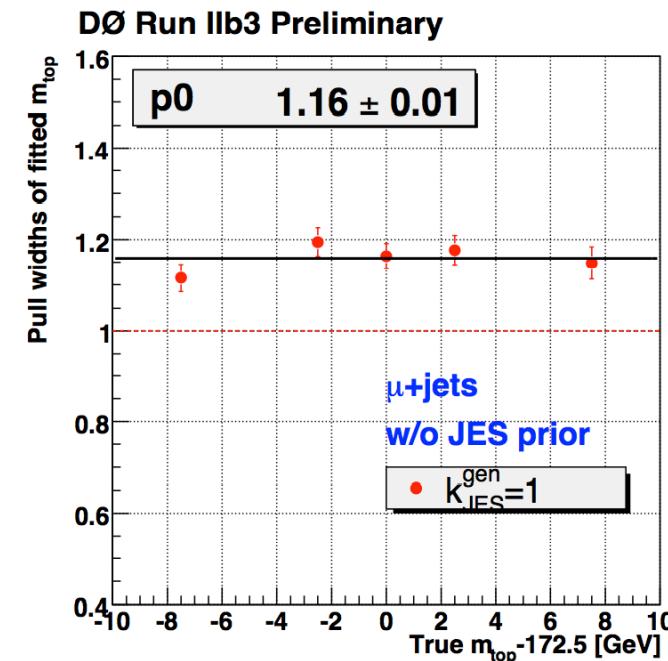
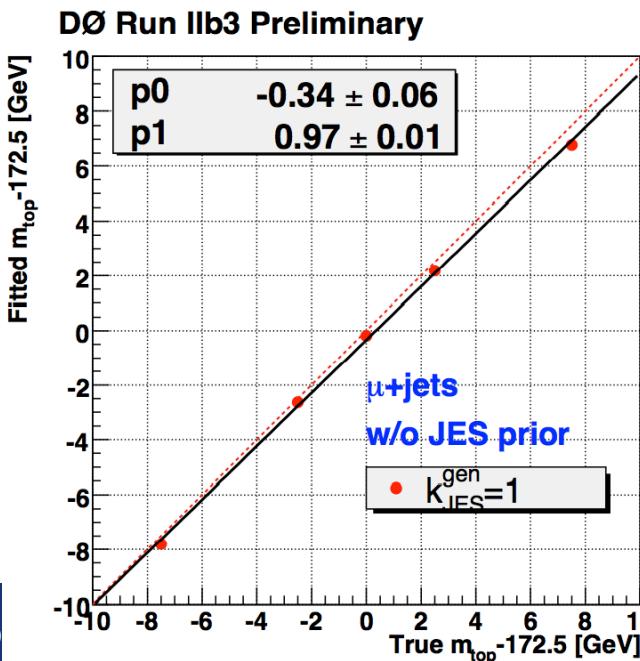
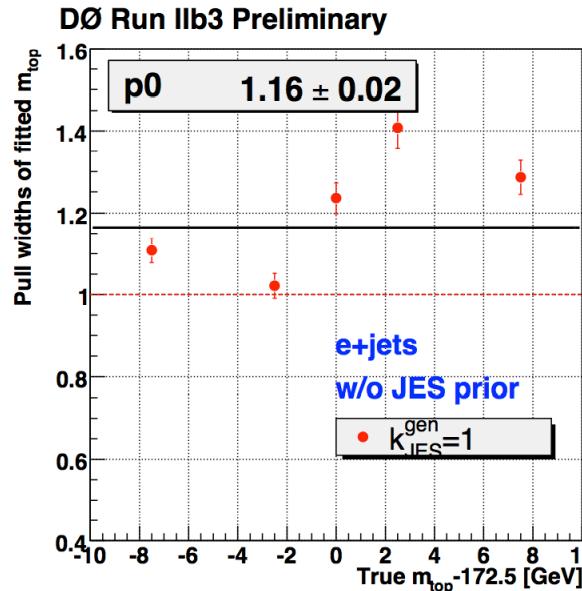
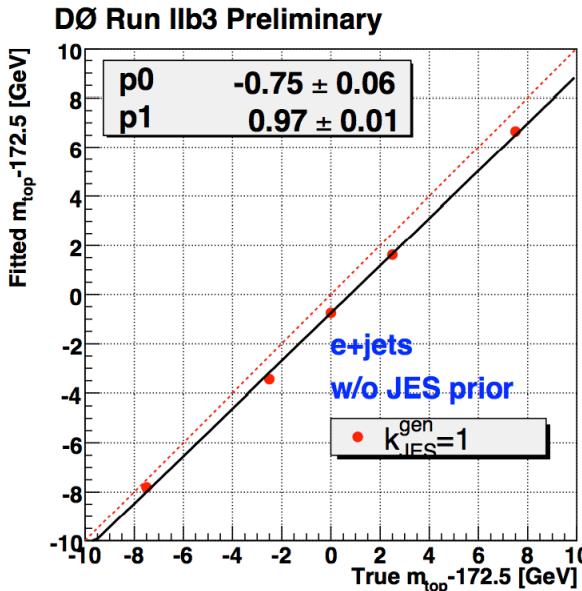
DØ Run IIb1 Preliminary

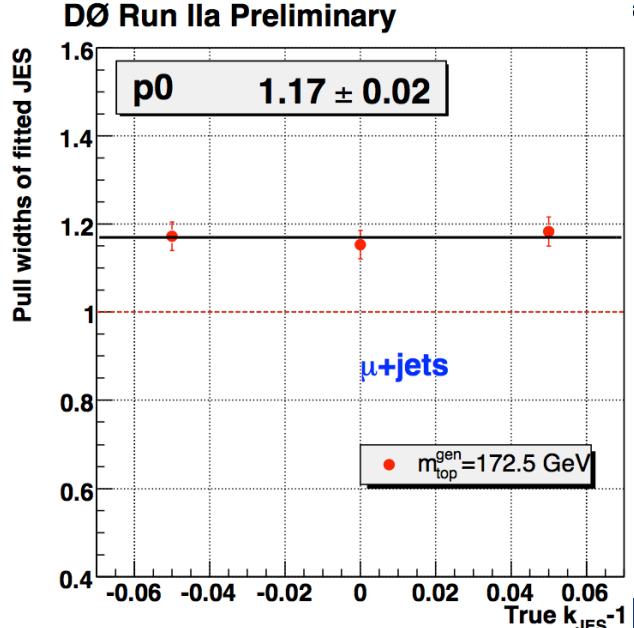
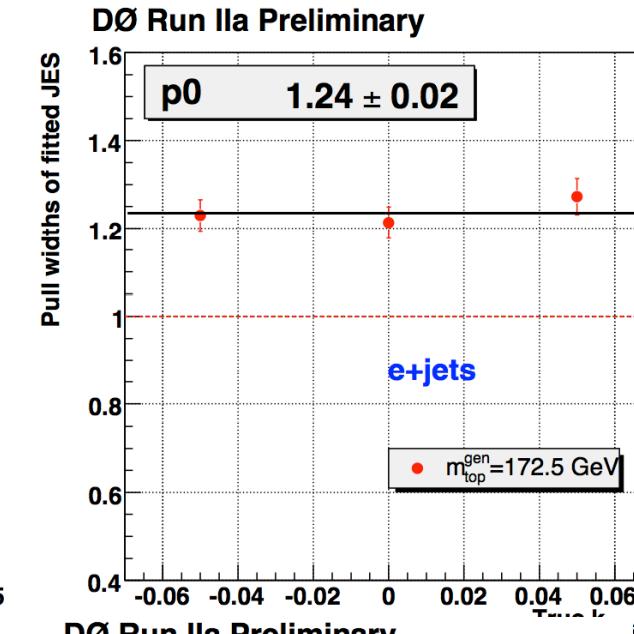
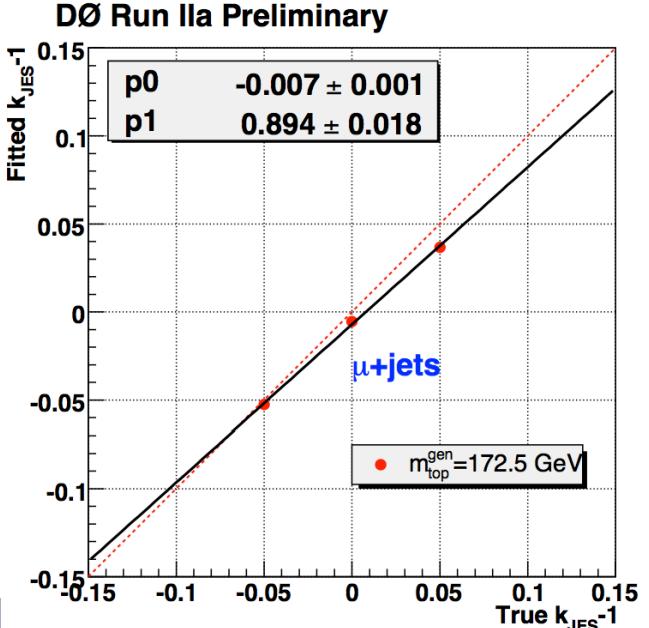
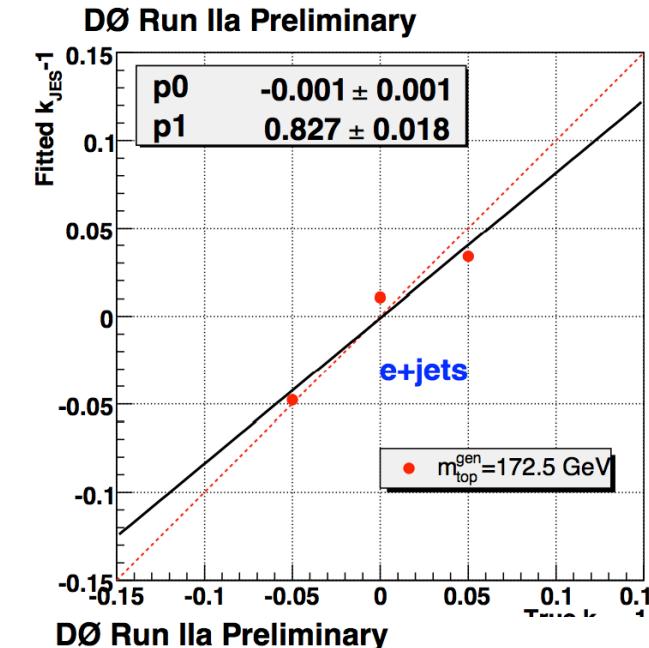


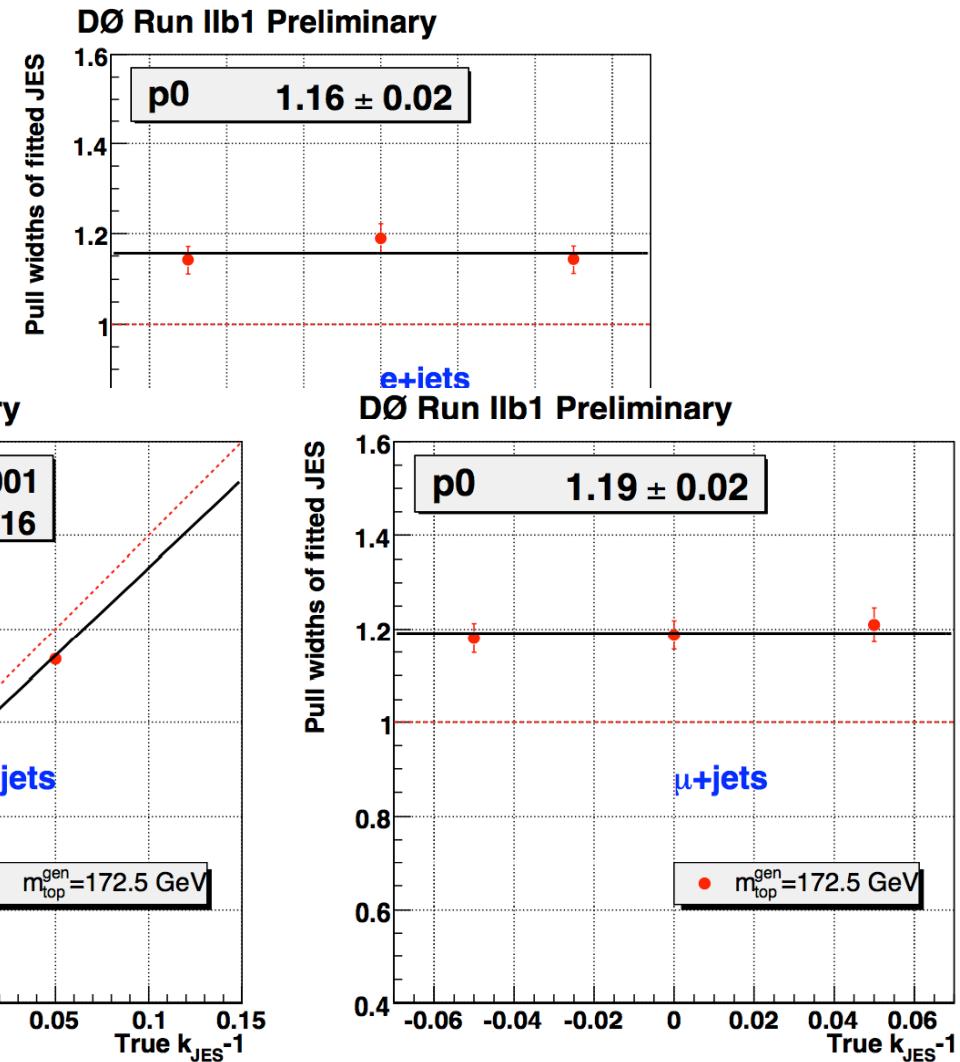
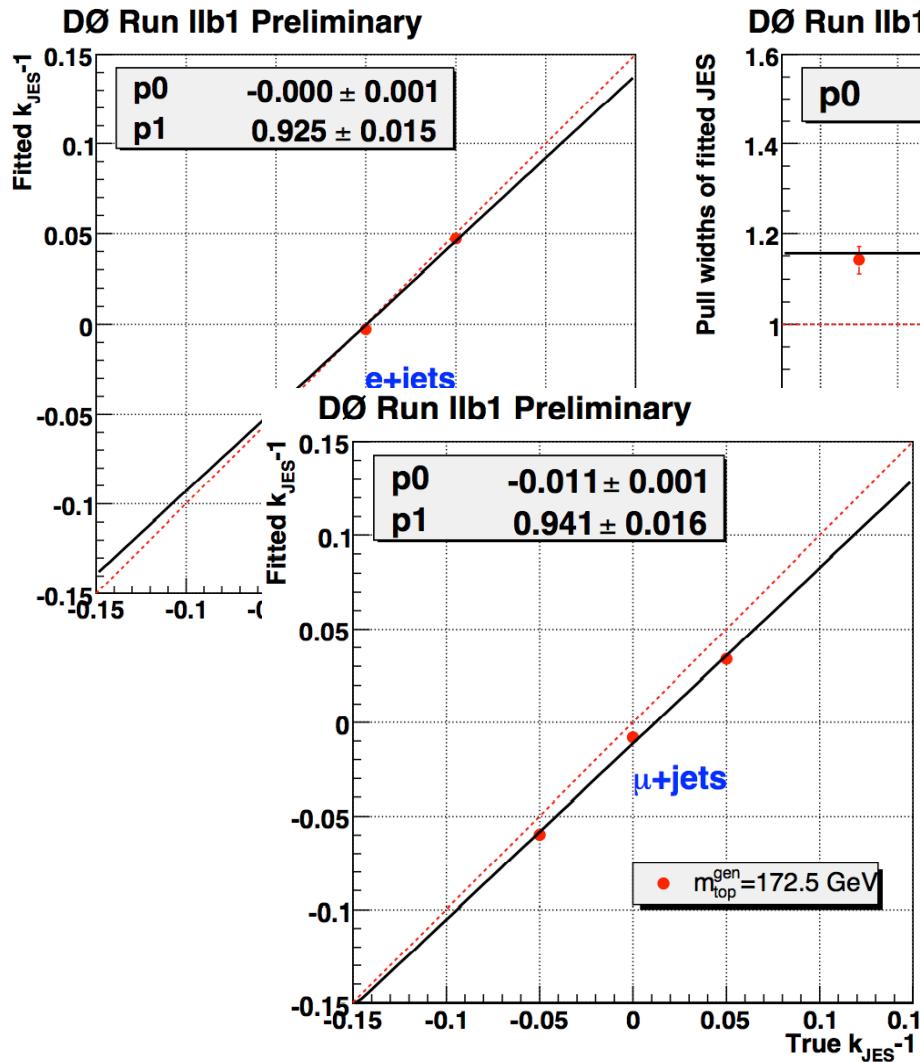
DØ Run IIb1 Preliminary





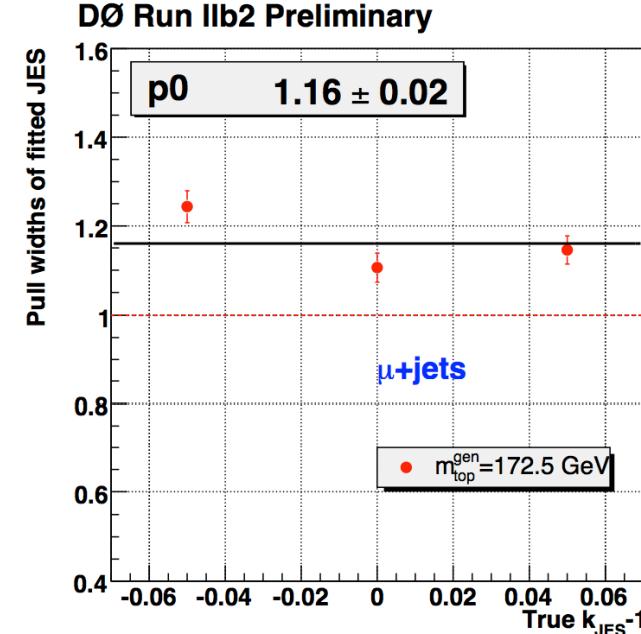
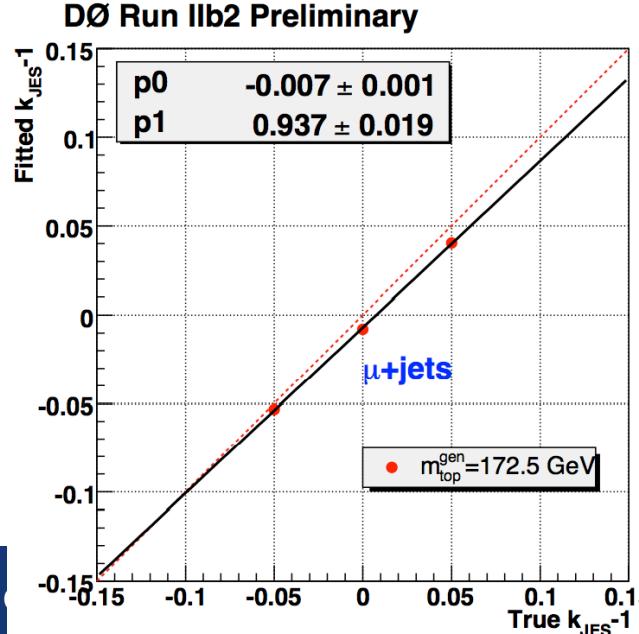
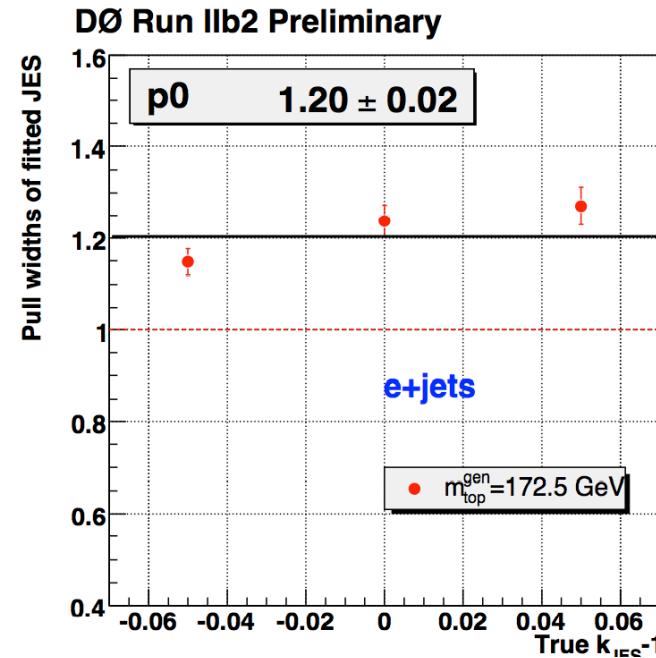
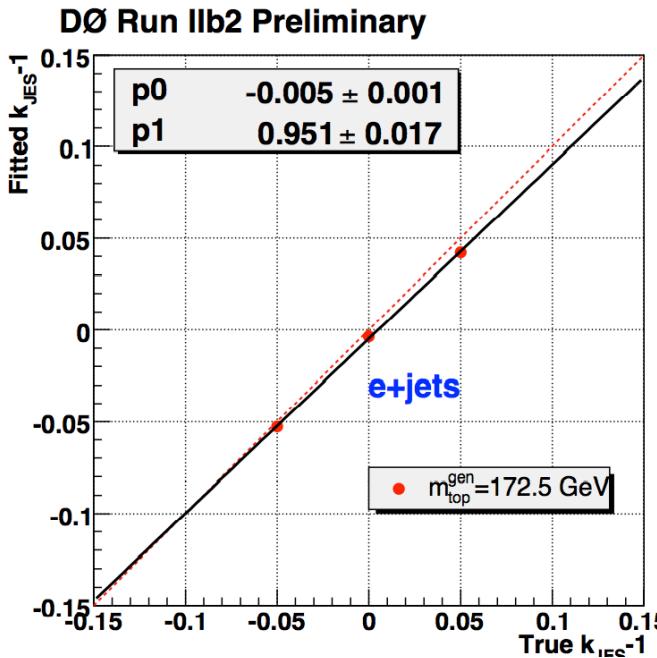


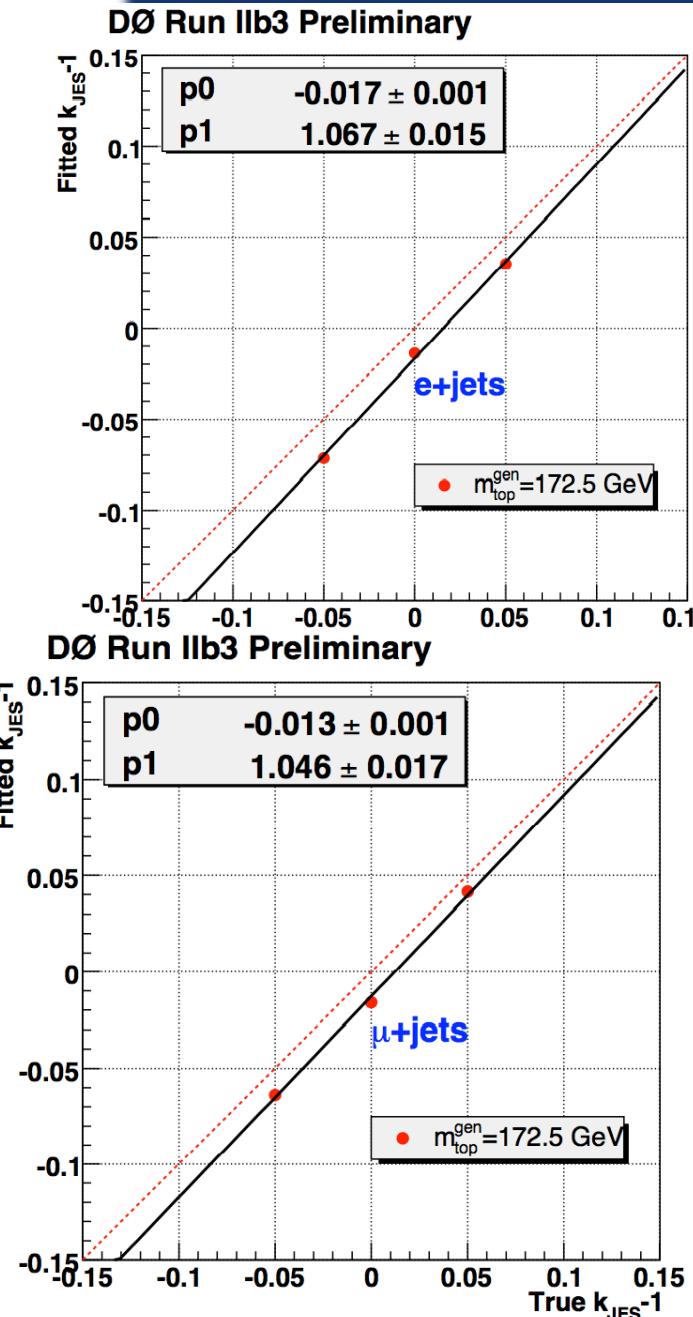






# Method response in $k_{\text{JES}}$





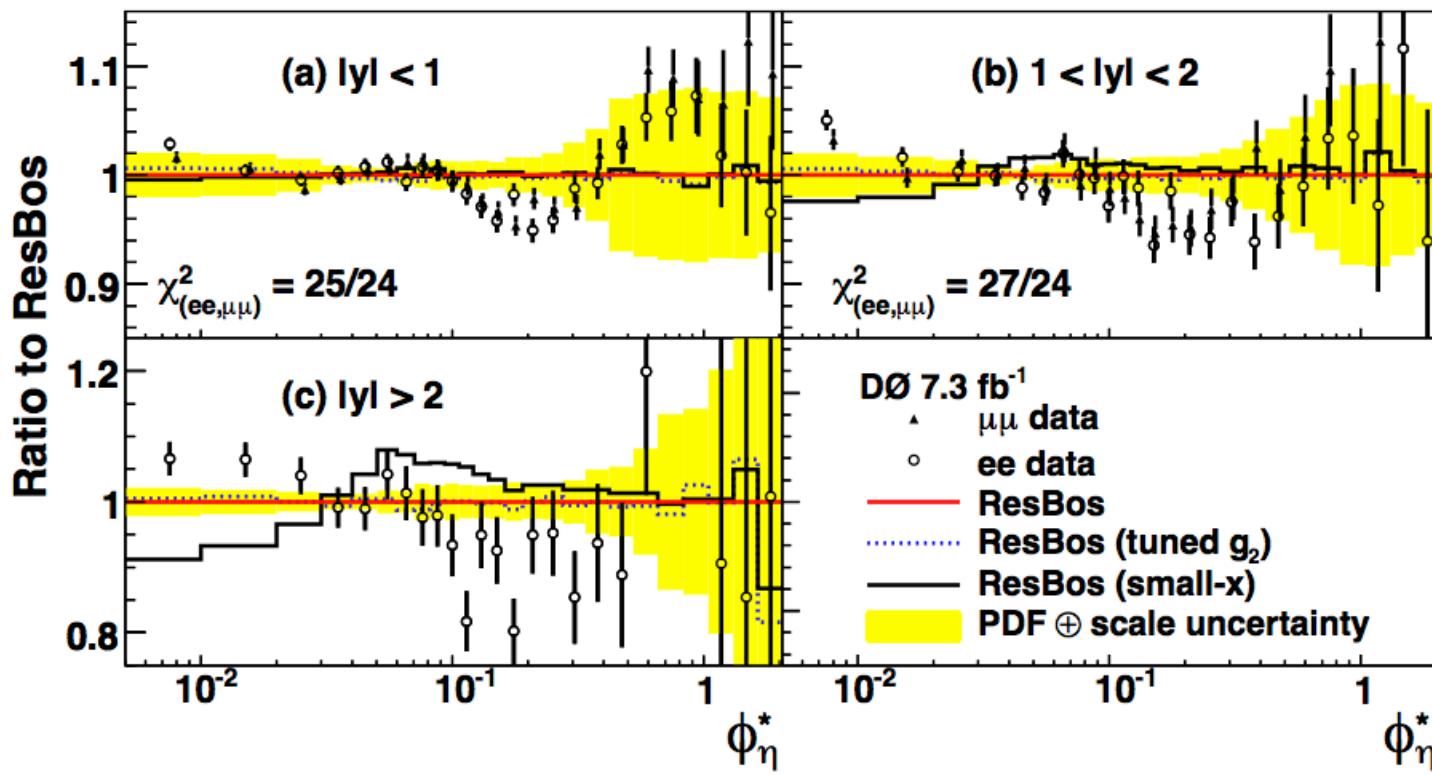
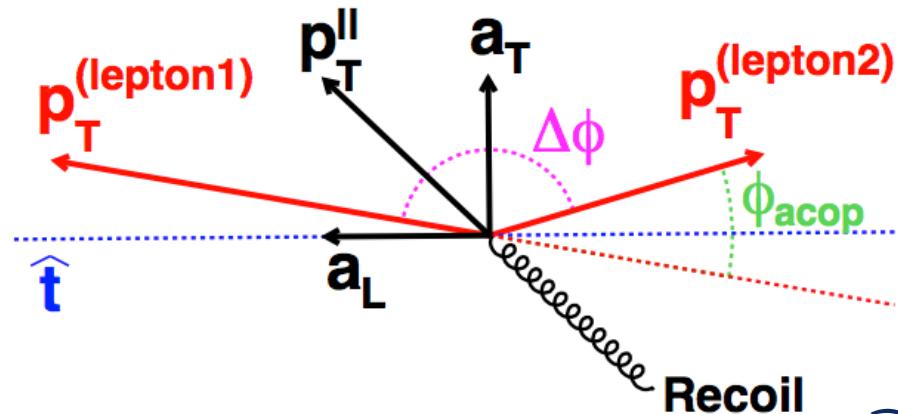


- Even though we perform an in-situ calibration of JES, this is only an overall calibration
  - $k_{\text{JES}}$  cannot account for any effects differential in  $(p_T, \eta)$
  - Study various parametrisations:
    - Vary jet energies according to upper error corridor on the JES, differentially in  $(p_T, \eta)$  using a parametrisation
    - Same for lower error corridor using a parametrisation
    - Vary jet energies according to upper error on JES jet-by-jet, i.e. w/o parametrisation
    - Assuming a linear increase in JES which is 0 for  $E=0$  and increases such as to touch the upper error corridor
      - In reality, only one parametrisation is correct
        - → take envelope
    - → Uncertainty from residual JES variations in  $(pT, \eta)$ :
      - 0.21 GeV (was: 0.21 GeV)

- Based on  $7.3 \text{ fb}^{-1}$  of data

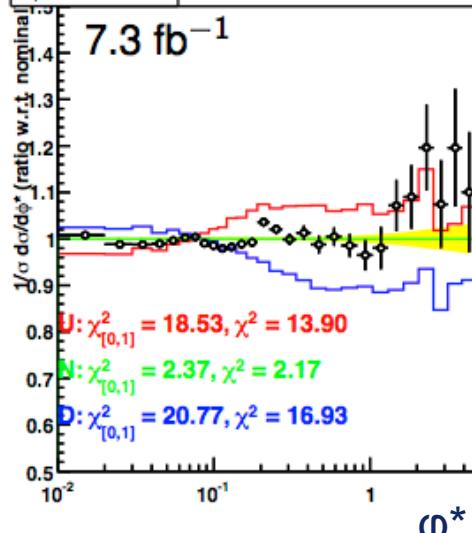
- Observable:

$$\phi_\eta^* = \tan(\phi_{\text{acop}}/2) \sin(\theta_\eta^*)$$

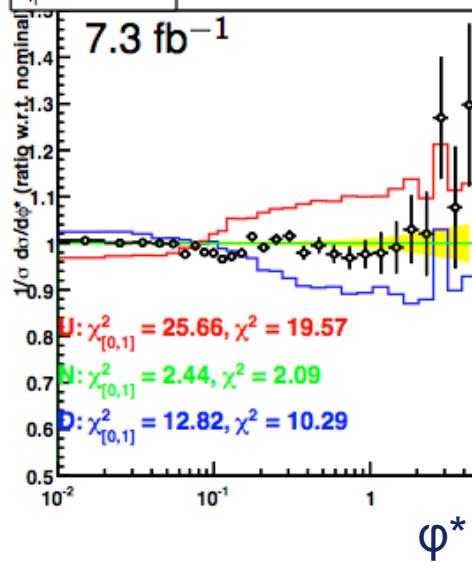




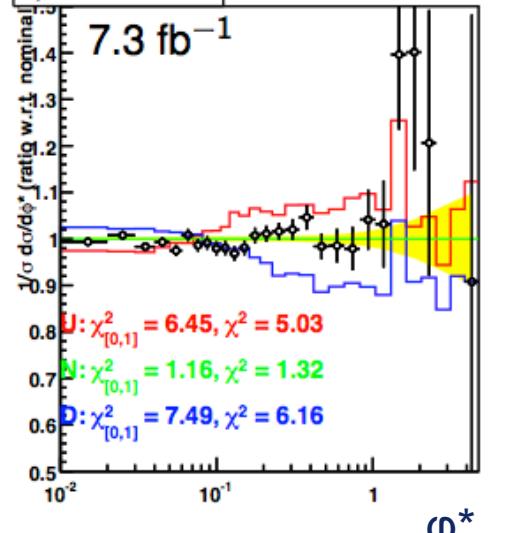
$\phi_\eta$  ( $\mu\mu$ ,  $|y| < 1$ ) DØ preliminary



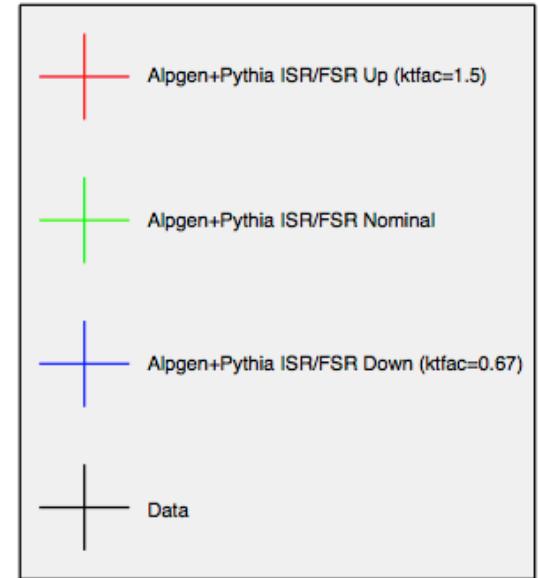
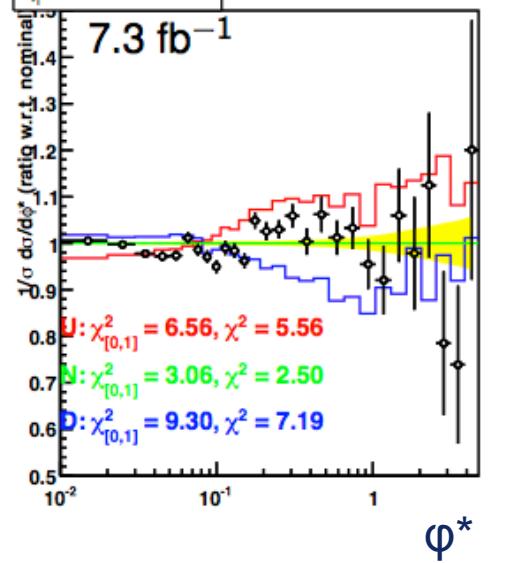
$\phi_\eta$  (ee,  $|y| < 1$ ) DØ preliminary



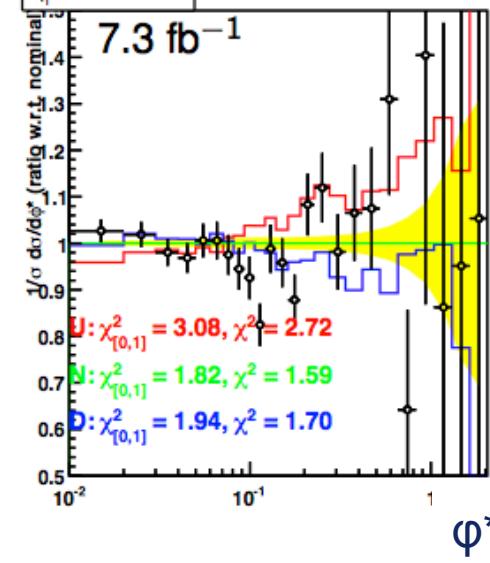
$\phi_\eta$  ( $\mu\mu$ ,  $1 < |y| < 2$ ) DØ preliminary



$\phi_\eta$  (ee,  $1 < |y| < 2$ ) DØ preliminary



$\phi_\eta$  (ee,  $|y| > 2$ ) DØ preliminary





- For reference:

$$m'_t = \frac{(m_t - 172.5 \text{ GeV}) - p_0^{m_t}}{p_1^{m_t}} + 172.5 \text{ GeV}$$

$$k'_{\text{JES}} = \frac{(k_{\text{JES}} - 1) - p_0^{k_{\text{JES}}}}{p_1^{k_{\text{JES}}}} + 1.$$

$$\sigma'(m'_t) = \sigma(m'_t) \times w_{\text{pull}}(m_t),$$

$$\sigma'(k'_{\text{JES}}) = \sigma(k'_{\text{JES}}) \times w_{\text{pull}}(k_{\text{JES}})$$

$$pull = \frac{\langle x \rangle - \bar{x}}{\sigma}$$



# Top quark mass: Tevatron combination

	Run I published					Run II published					Run II prel.	
	CDF			DØ		CDF			DØ		CDF	
	$\ell+$ jets	$\ell\ell$	all-jets	$\ell+$ jets	$\ell\ell$	$\ell+$ jets	$L_{XY}$	MET	$\ell+$ jets	$\ell\ell$	$\ell\ell$	all-jets
CDF-I $\ell+$ jets	1.00	0.29	0.32	0.26	0.11	0.49	0.07	0.26	0.19	0.12	0.54	0.27
CDF-I $\ell\ell$	0.29	1.00	0.19	0.15	0.08	0.29	0.04	0.16	0.12	0.08	0.32	0.17
CDF-I all-jets	0.32	0.19	1.00	0.14	0.07	0.30	0.04	0.16	0.08	0.06	0.37	0.18
DØ-I $\ell+$ jets	0.26	0.15	0.14	1.00	0.16	0.22	0.05	0.12	0.13	0.07	0.26	0.14
DØ-I $\ell\ell$	0.11	0.08	0.07	0.16	1.00	0.11	0.02	0.07	0.07	0.05	0.13	0.07
CDF-II $\ell+$ jets	0.49	0.29	0.30	0.22	0.11	1.00	0.08	0.32	0.28	0.18	0.52	0.30
CDF-II $L_{XY}$	0.07	0.04	0.04	0.05	0.02	0.08	1.00	0.04	0.05	0.03	0.06	0.04
CDF-II MET	0.26	0.16	0.16	0.12	0.07	0.32	0.04	1.00	0.17	0.11	0.29	0.18
DØ-II $\ell+$ jets	0.19	0.12	0.08	0.13	0.07	0.28	0.05	0.17	1.00	0.36	0.15	0.14
DØ-II $\ell\ell$	0.12	0.08	0.06	0.07	0.05	0.18	0.03	0.11	0.36	1.00	0.10	0.09
CDF-II $\ell\ell$	0.54	0.32	0.37	0.26	0.13	0.52	0.06	0.29	0.15	0.10	1.00	0.32
CDF-II all-jets	0.27	0.17	0.18	0.14	0.07	0.30	0.04	0.18	0.14	0.09	0.32	1.00

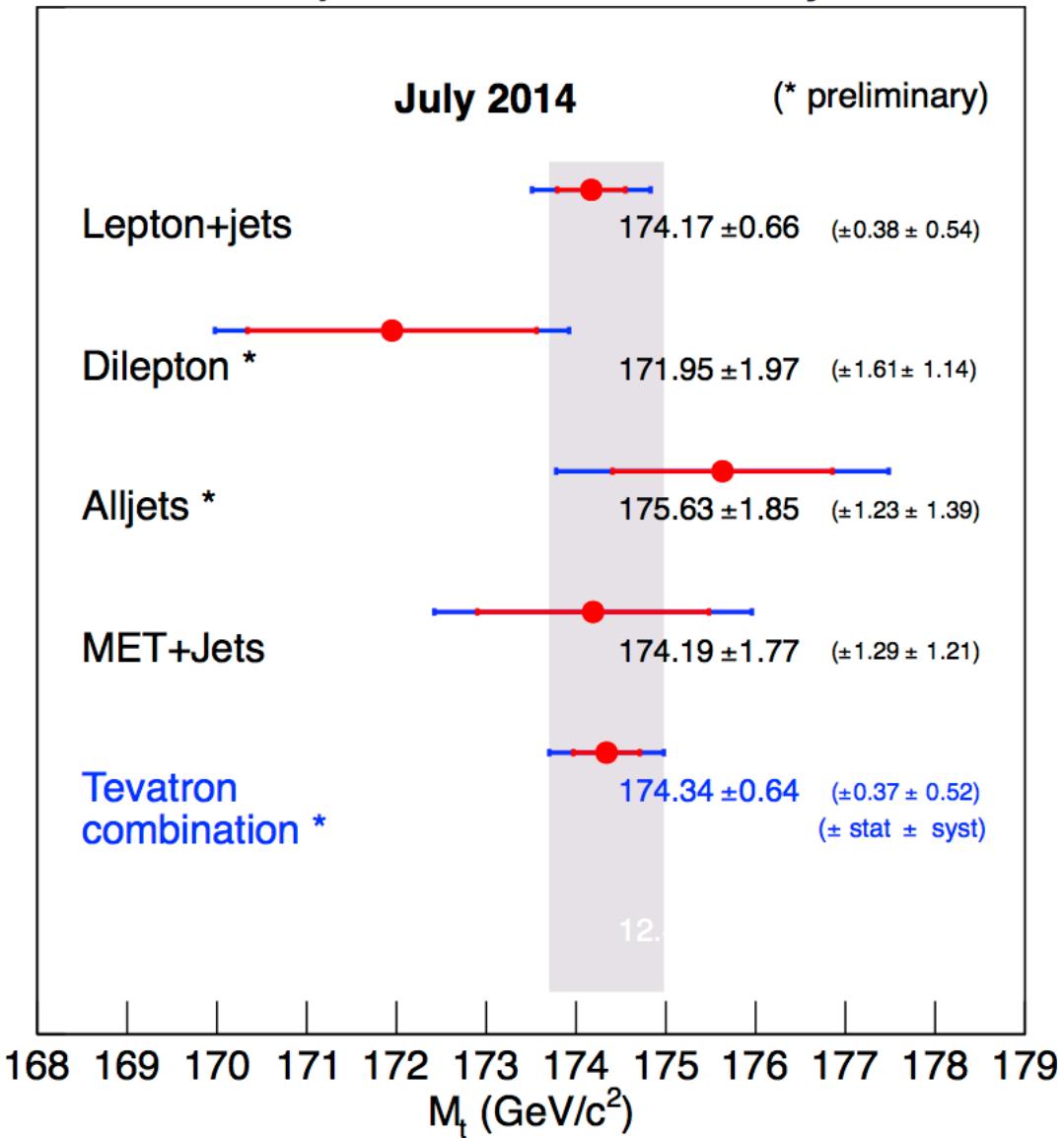
	Run I published					Run II published					Run II prel.	
	CDF			DØ		CDF			DØ		CDF	
	$\ell+$ jets	$\ell\ell$	all-jets	$\ell+$ jets	$\ell\ell$	$\ell+$ jets	$L_{XY}$	MET	$\ell+$ jets	$\ell\ell$	$\ell\ell$	all-jets
Pull	0.24	-0.61	+1.01	+1.09	-0.46	-1.64	-0.791	-0.24	+1.60	-0.13	-1.11	0.39
Weight [%]	-2.6	-0.7	-0.4	-0.1	-0.14	+28.8	+0.1	+5.5	+67.2	-2.9	-0.66	+6.0

Parameter	Value (GeV/c <sup>2</sup> )	Correlations			
		$M_t^{\text{all-jets}}$	$M_t^{\ell+\text{jets}}$	$M_t^{\ell\ell}$	$M_t^{\text{MET}}$
$M_t^{\text{all-jets}}$	$175.63 \pm 1.85$	1.00			
$M_t^{\ell+\text{jets}}$	$174.17 \pm 0.66$	0.21	1.00		
$M_t^{\ell\ell}$	$171.95 \pm 1.97$	0.21	0.41	1.00	
$M_t^{\text{MET}}$	$174.19 \pm 1.77$	0.11	0.23	0.18	1.00

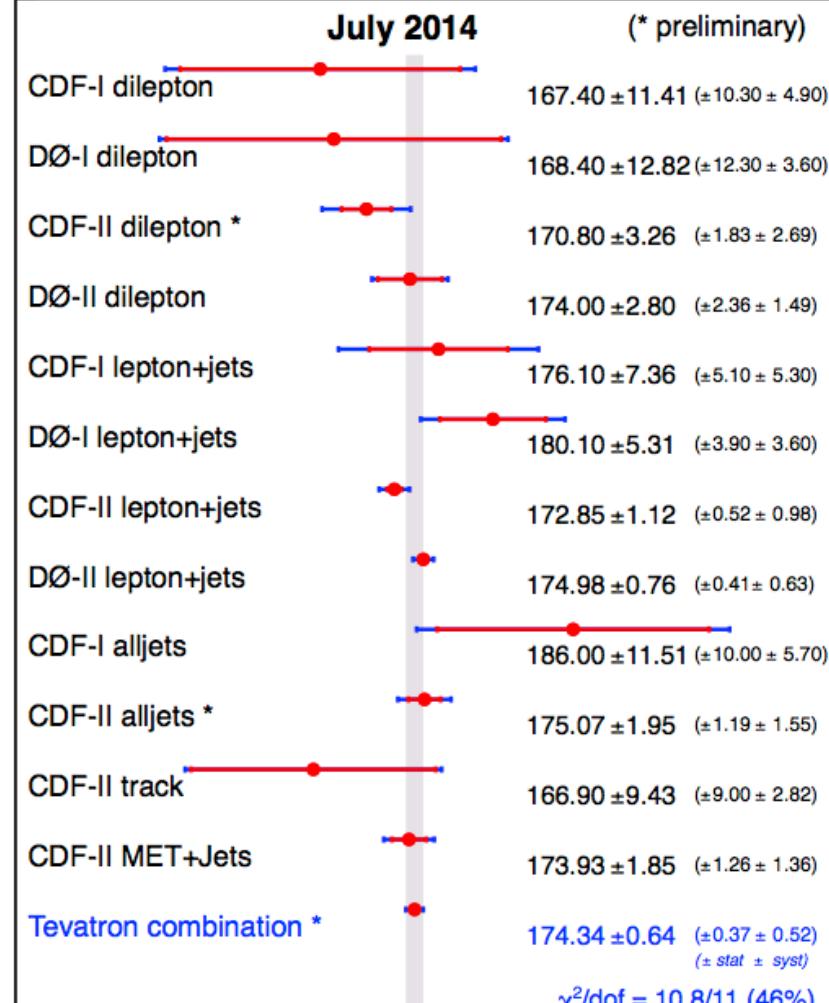


## Top quark mass: Tevatron combination

## Mass of the Top Quark in Different Decay Channels



## Mass of the Top Quark

 $\chi^2/\text{dof} = 10.8/11$  (46%)

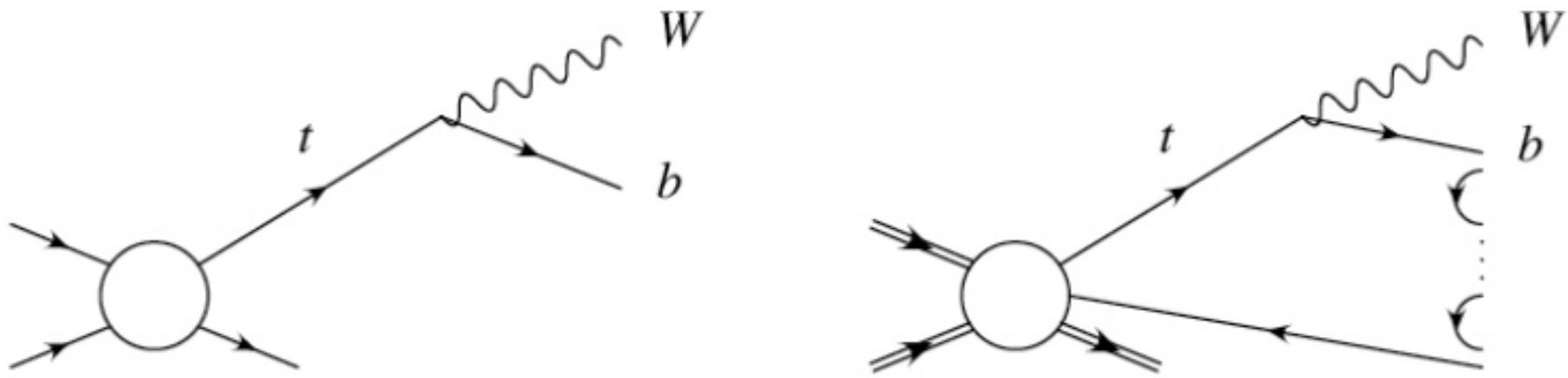


- Input: correlations between uncertainty categories:

	$\rho_{\text{EXP}}$				$\rho_{\text{LHC}}$	$\rho_{\text{TEV}}$	$\rho_{\text{COL}}$	
	$\rho_{\text{CDF}}$	$\rho_{\text{D0}}$	$\rho_{\text{ATL}}$	$\rho_{\text{CMS}}$			$\rho_{\text{ATL-TEV}}$	$\rho_{\text{CMS-TEV}}$
Stat	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
iJES	0.0	1.0	0.0	0.0	0.0	0.0	0.0	0.0
stdJES	1.0	1.0	1.0	1.0	0.0	0.0	0.0	0.0
flavourJES	1.0	1.0	1.0	1.0	0.0	0.0	0.0	0.0
bJES	1.0	1.0	1.0	1.0	0.5	1.0	1.0	0.5
MC	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0
Rad	1.0	1.0	1.0	1.0	1.0	1.0	0.5	0.5
CR	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0
PDF	1.0	1.0	1.0	1.0	1.0	1.0	0.5	0.5
DetMod	1.0	1.0	1.0	1.0	0.0	0.0	0.0	0.0
<i>b</i> -tag	1.0	1.0	1.0	1.0	0.0	0.0	0.0	0.0
LepPt	1.0	1.0	1.0	1.0	0.0	0.0	0.0	0.0
BGMC <sup>†</sup>	1.0	1.0	1.0	1.0	1.0	1.0	1.0	1.0
BGData	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
Meth	0.0	0.0	0.0	0.0	0.0	0.0	0.0	0.0
MHI	1.0	1.0	1.0	1.0	1.0	0.0	0.0	0.0

- $\rho_{\text{CDF}}, \rho_{\text{D0}}, \rho_{\text{ATL}}, \rho_{\text{CMS}}$ : correlations within an experiment
- $\rho_{\text{LHC}}, \rho_{\text{TEV}}$ : correlations within the collider (LHC/Tevatron)
- $\rho_{\text{ATL-TEV}}, \rho_{\text{CMS-TEV}}$ : correlations between ATLAS or CMS and Tevatron

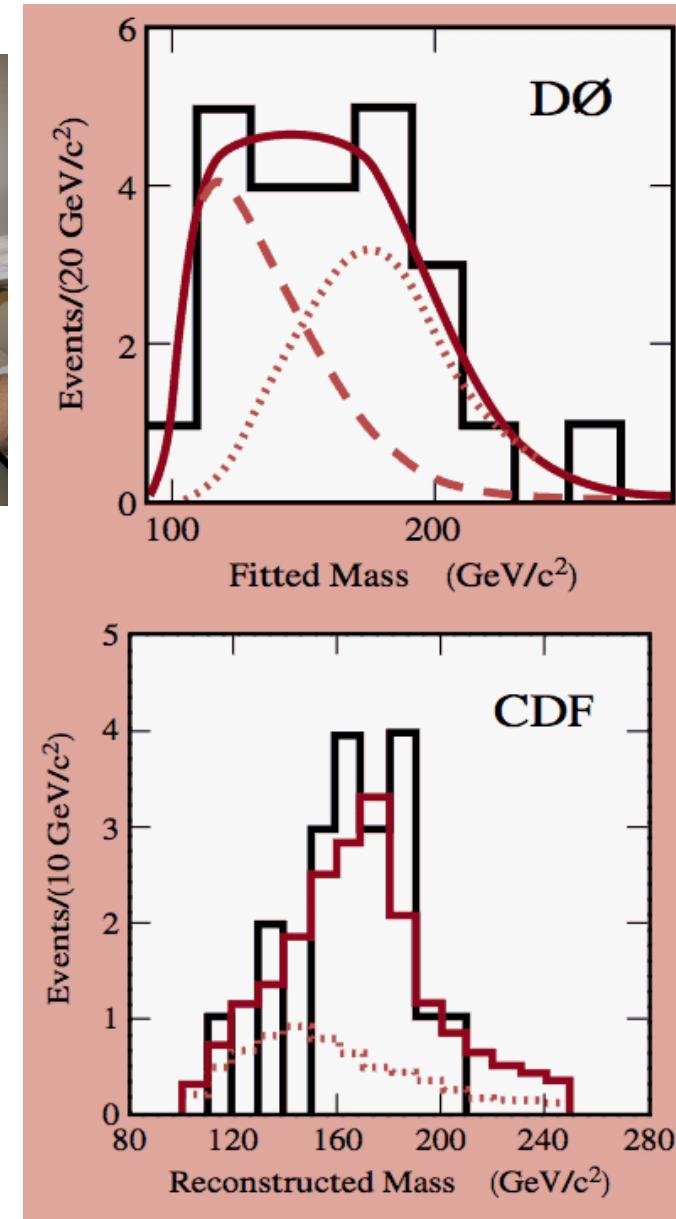
- (I only want to refresh our memory here)
- The top **mass** is **not an observable** per se and has to be inferred from its effect on kinematic observables
- The mass **cannot be well-defined at LO**
- The **pole mass** corresponds to our physical intuition of a stable particle
  - $m_{\text{top}}$  is the “pole” in the top quark propagator
    - Although this is not fully correct (hadronisation effects)
  - The pole mass can never be determined with **precision** better than  $\Lambda_{\text{QCD}}$ :





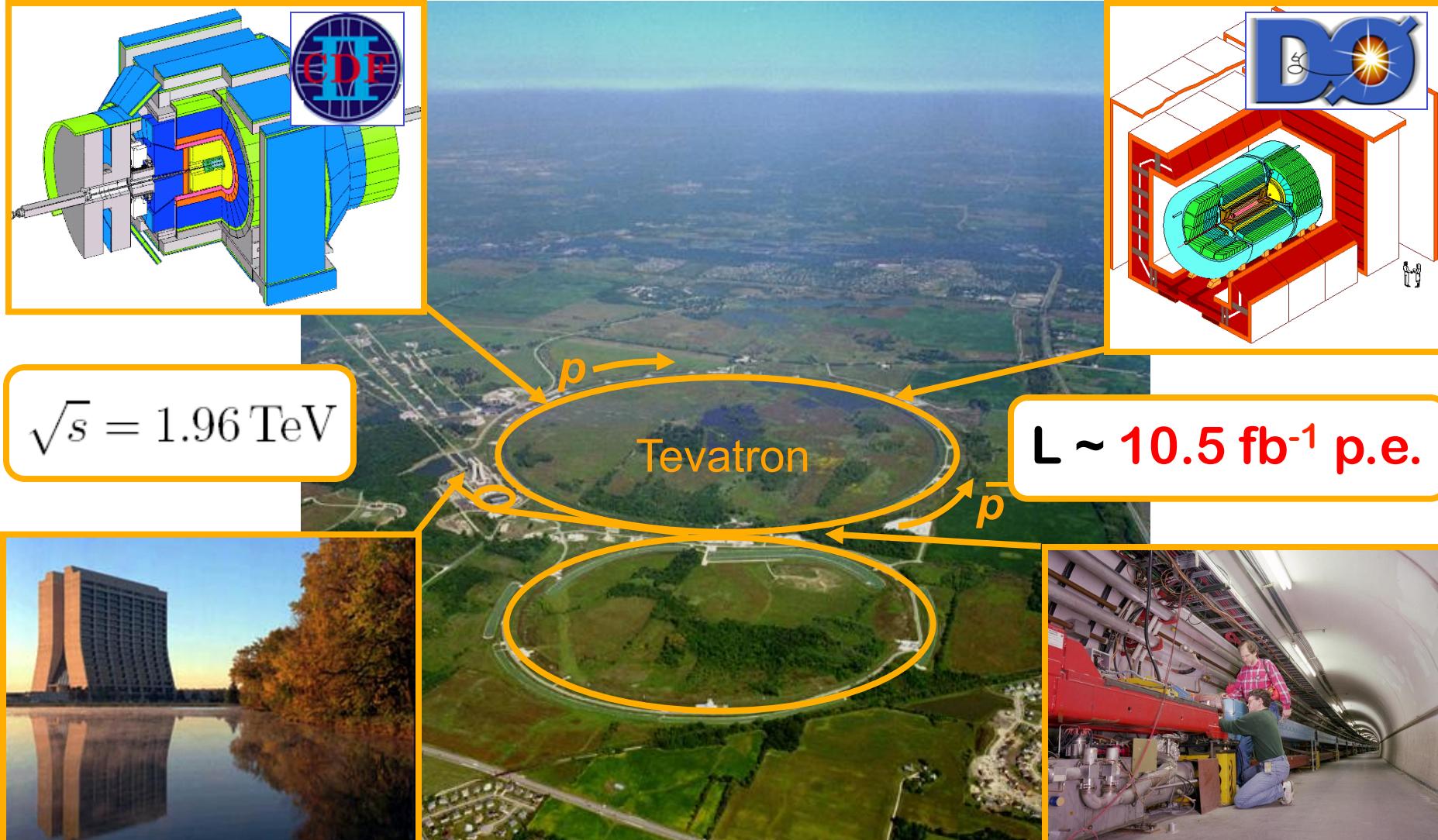
- Other popular mass definition schemes:
  - e.g. modified minimal subtraction scheme ( $\overline{\text{MS}}$ ), also referred to as running mass  $m_{\text{top}}(\mu_r)$ 
    - The  $\mu_r$  dependence can be used to absorb logarithmic corrections through resummation (in specific cases)
      - better behaviour of perturbative predictions
    - The  $\overline{\text{MS}}$  mass can be translated into the pole mass at any fixed order of perturbation theory
- What we typically measure at hadron colliders, is:
  - Neither the  $\overline{\text{MS}}$  mass, nor the pole mass  $\rightarrow m^{\text{MC}}$
  - “Close” to the pole mass
    - “Close” not quantified yet
- True also for NLO generators like e.g. powheg
  - finite width effects of top propagator are not simulated, but generated via reweighting

- 24 Feb. 1995:
  - Simultaneous PRL submission by CDF and DØ
- CDF ( $67 \text{ pb}^{-1}$ ):
  - $\sigma = 6.8^{+3.6}_{-2.4} \text{ pb}$ ,
  - observed 19 events, expected 6.9 bkg
    - bkg-only hypothesis rejected at  $4.8\sigma$
  - $m_{\text{top}} = 176 \pm 13 \text{ GeV}$
- DØ ( $50 \text{ pb}^{-1}$ ):
  - $\sigma = 6.4 \pm 2.2 \text{ pb}$ ,
  - observed 17 events, expected 3.8 bkg
    - → bkg-only hypothesis rejected at  $4.6\sigma$
  - $m_{\text{top}} = 199 \pm 30 \text{ GeV}$



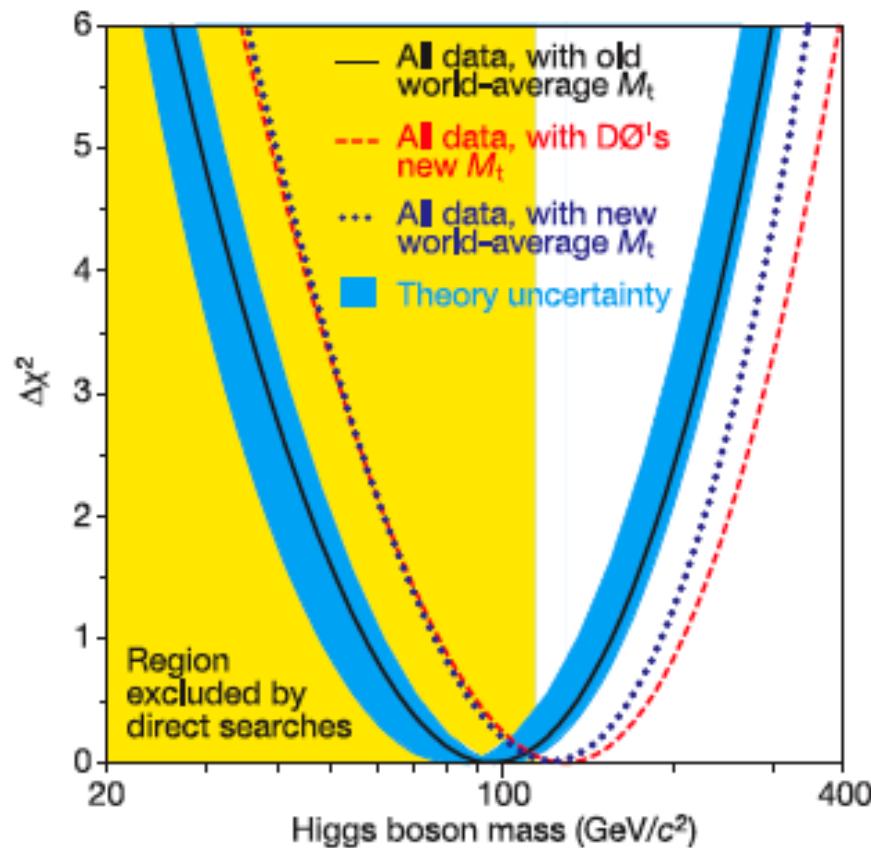


# More about the birth place...





- The first (published) measurement in HEP using the MEM:



$\text{N}_2\text{H}^+$  obser-  
far-ultraviolet  
gen chemistry  
lar gas. □

citic molecular clouds.

interstellar clouds  
222–243 (1995).  
implications for  
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$\text{N}_2$  in diffuse clouds.

T. C. V. S. The  
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variations in the

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00–1012 (2003).

149–L52 (1997).

m. *Astrophys. J.* 538,

Explorer satellite.

## A precision measurement of the mass of the top quark

D0 Collaboration\*

\*A list of authors and their affiliations appear at the end of the paper

The standard model of particle physics contains parameters—such as particle masses—whose origins are still unknown and which cannot be predicted, but whose values are constrained through their interactions. In particular, the masses of the top quark ( $M_t$ ) and  $W$  boson ( $M_W$ )<sup>1</sup> constrain the mass of the long-hypothesized, but thus far not observed, Higgs boson. A precise measurement of  $M_t$  can therefore indicate where to look for the Higgs, and indeed whether the hypothesis of a standard model Higgs is consistent with experimental data. As top quarks are produced in pairs and decay in only about  $10^{-24}$  s into various final states, reconstructing their masses from their decay products is very challenging. Here we report a technique that extracts more information from each top-quark event and yields a greatly improved precision (of  $\pm 5.3 \text{ GeV}/c^2$ ) when compared to previous measurements<sup>2</sup>. When our new result is combined with our published measurement in a complementary decay mode<sup>3</sup> and with the only other measurements available<sup>2</sup>, the new world average for  $M_t$  becomes<sup>4</sup>  $178.0 \pm 4.3 \text{ GeV}/c^2$ . As a

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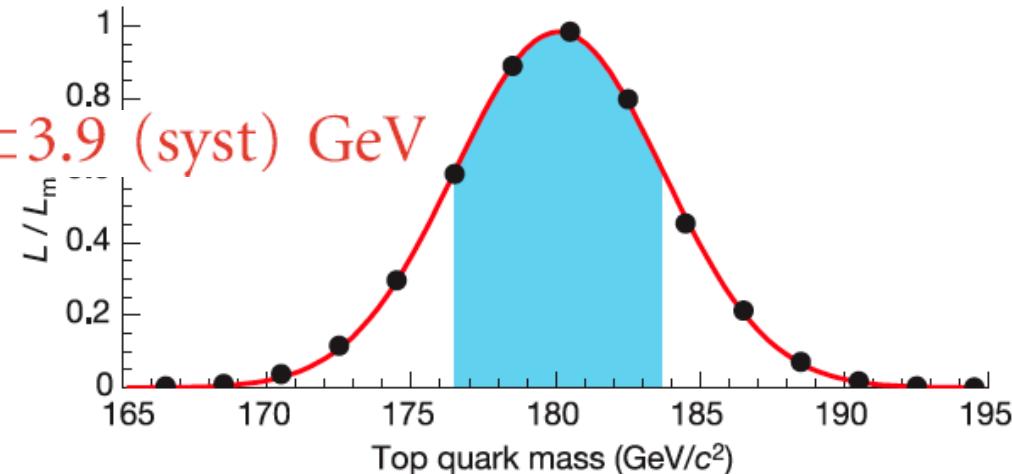
NATURE | VOL 429 | 10 JUNE 2004 | www.nature.com/nature

**letters to nature**

The experiment- top quark in our previous publication, and correspond to an |

- The final result:

- $M_t = 180.1 \pm 3.6 \text{ (stat)} \pm 3.9 \text{ (syst) GeV}$ 
  - Using 125 pb<sup>-1</sup> of p-pbar collisions @ 1.8 TeV, 71 events



- Previous result:

- $M_t = 173.3 \pm 5.6 \text{ (stat)} \pm 5.5 \text{ (syst) GeV}$ 
  - same dataset, 91 candidates

- Much higher statistical sensitivity:

- Corresponding to 2.4x more data with old method!
- Systematic uncertainties are also smaller

- Already this analysis

- Was using jet-parton transfer functions
  - Looked at 12 possible jet-parton assignments (4 jets)
  - Used numerical integration in 5 variables