

Perspective for DM searches @ FCCee

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What we expect DM to look like

A. Strumia, talk @ FCCee workshop

<https://agenda.infn.it/getFile.py/access?contribId=19&sessionId=7&resId=0&materialId=slides&confId=8830>

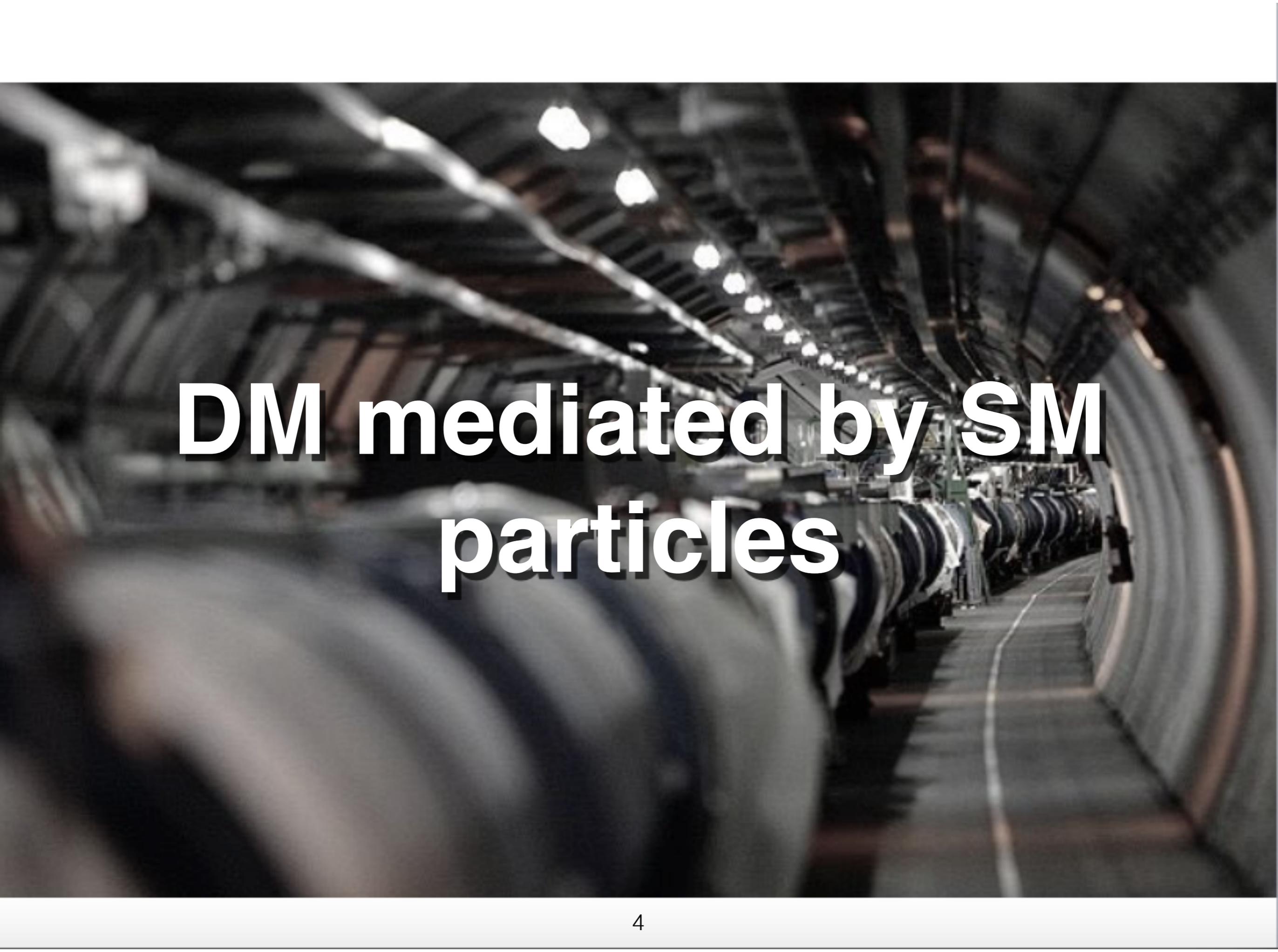
- DM particles should be in some Z_2 symmetry making them stable (e.g. R-parity in SUSY)
 - Should be produced in pairs
- DM is not detectable with a typical collider detector
 - Main signature @FCCee is missing momentum/mass
 - some other object could be radiates (**monophoton, monojet, mono....**)
- Most likely, DM is part of a new sector, with other particles offering an easier path to discovery

Subject of this talk

Not covered today

How could DM be produced

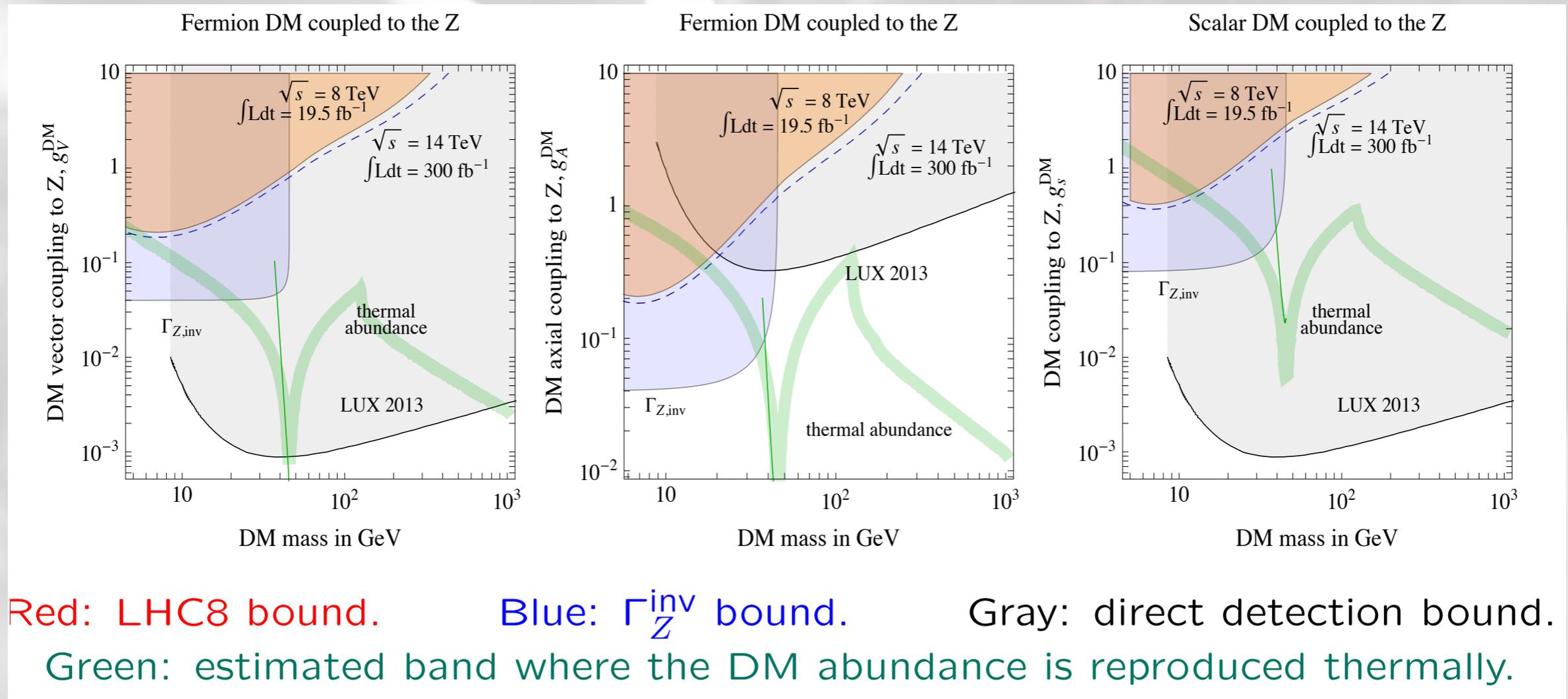
- SM mediator [De Simone et al. http://arxiv.org/abs/1402.6287](http://arxiv.org/abs/1402.6287)
 - DM couples to SM particles (Z, H, ...)
 - Widths to invisible for Z and H play major role
- New heavy particles coupling both to SM and DM
 - off-shell mediator connect colliding ee to DM pair (s-channel or t-channel)
 - Effective field theory used in this case (and justified more than in other scenarios, e.g. LHC)
 - Nature of the mediator determines the effective operator that rules the interaction

A long, dimly lit tunnel with a series of lights on the ceiling, receding into the distance. The perspective is from the center of the tunnel, looking down its length. The walls and floor are dark, and the lights create a strong sense of depth and perspective.

DM mediated by SM particles

The Z portal option

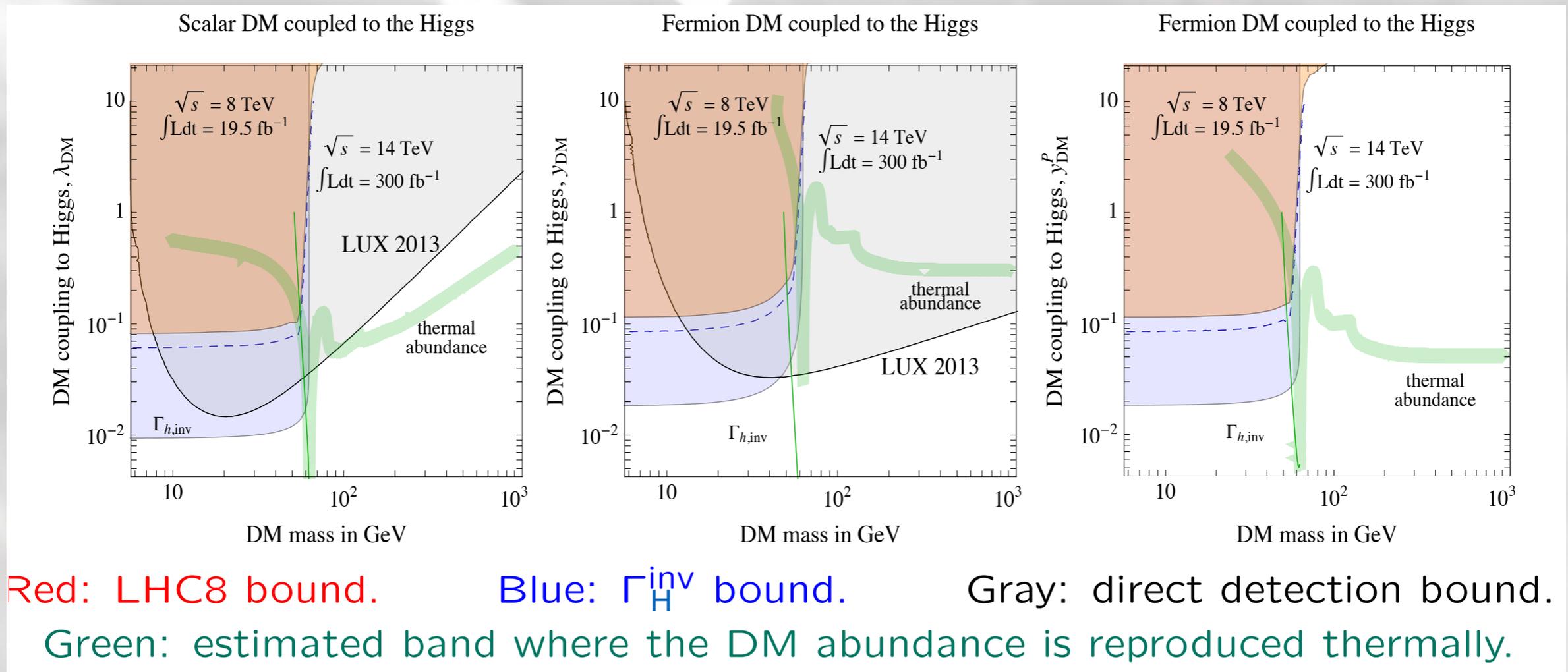
De Simone et al. <http://arxiv.org/abs/1402.6287>



- Direct detection excludes large part of the coupling vs mass plane
- LHC searches complementary (can probe masses below the recoil threshold of direct detection)
- Bound from $\Gamma(Z \rightarrow inv)$ from LEP-SLD stronger than LHC where it applies ($m_{DM} < 45$ GeV)

The H portal option

De Simone et al. <http://arxiv.org/abs/1402.6287>



- Similar picture for Higgs portal
- LHC limit from monojet and $H \rightarrow inv$
- Some scenario only probed at colliders

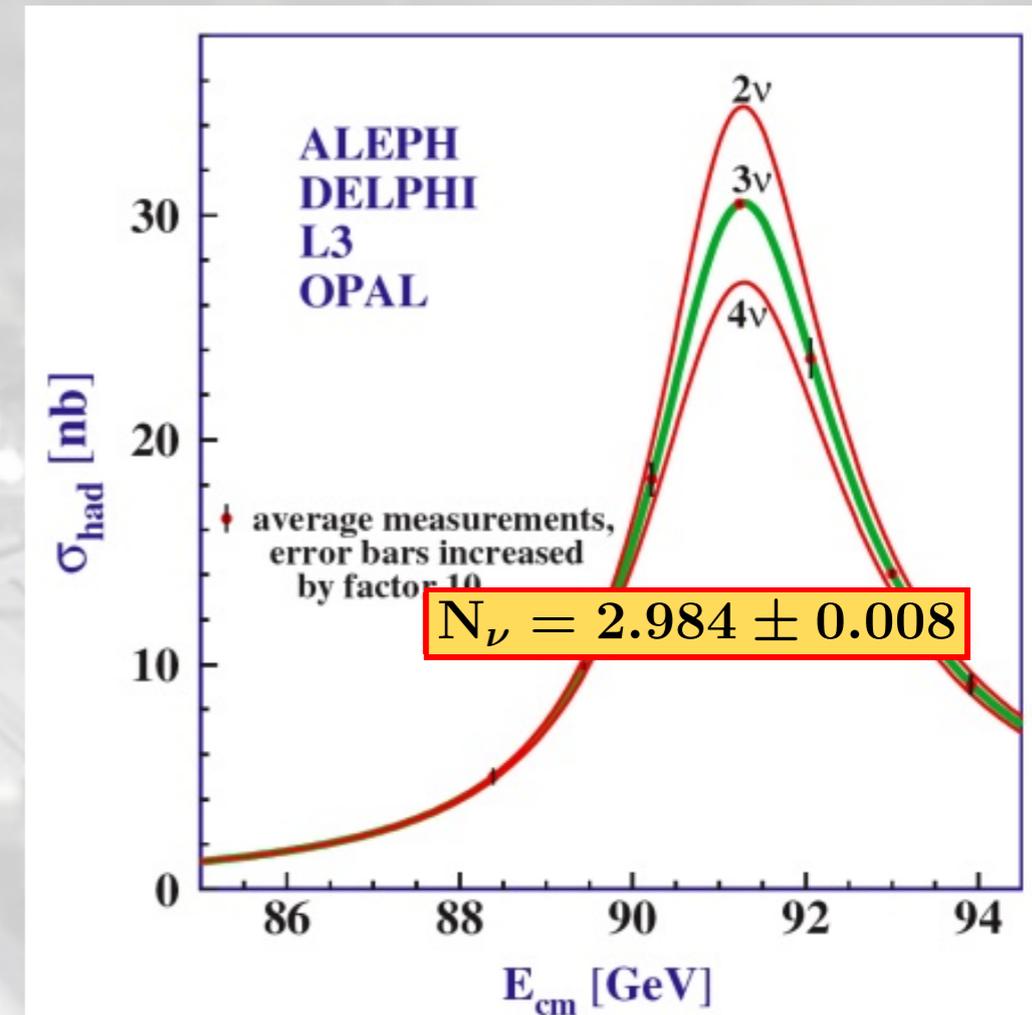
What can be done @FCCee?

- $\Gamma(Z \rightarrow \text{inv})$ measured @LEP from Z lineshape
 - need to know shad precisely
 - need to control luminosity

Expected statistical precision on δN_ν

Mode	$N_{inv}(240)$	$N_{inv}(350)$	δN_ν	Efficiency
Z(l)Z(inv)	$1.9 \cdot 10^5$	$0.4 \cdot 10^5$	0.013	CMS LEP3 [arXiv:1208.1662]
Z(bb)Z(inv)	$2.0 \cdot 10^5$	$0.4 \cdot 10^5$	0.012	21%, CMS LEP3 [arXiv:1208.1662]
H(bb)Z(inv)	$0.6 \cdot 10^5$	$0.1 \cdot 10^5$	0.022	21%, same as for Z(bb)Z(inv)
Total	$4.5 \cdot 10^5$	$0.9 \cdot 10^5$	0.008	
$\gamma Z(\text{inv})$	$3.6 \cdot 10^7$ (240+350)		0.001	100% eff., no bkg.

- This measurement will be systematic dominated @FCCee
- The alternative would be a monophoton search, for which systematics could cancel in a ratio
- In general, monophoton search is the DM probe @FCCee. More in next slides

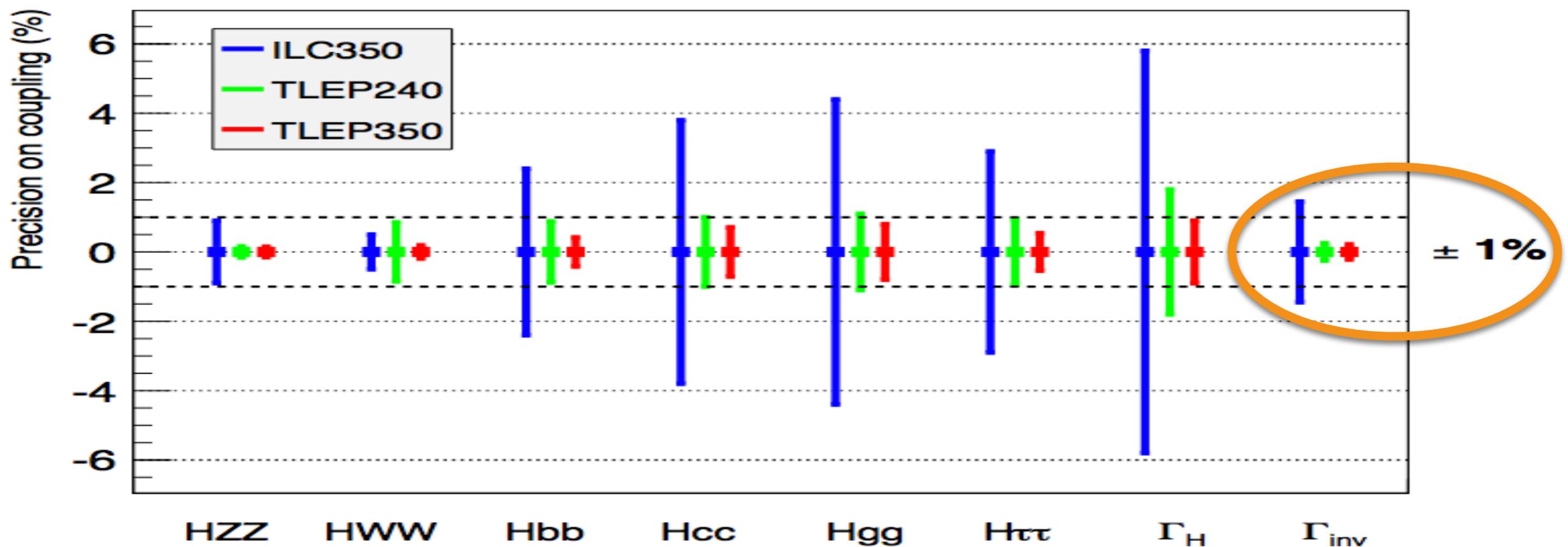
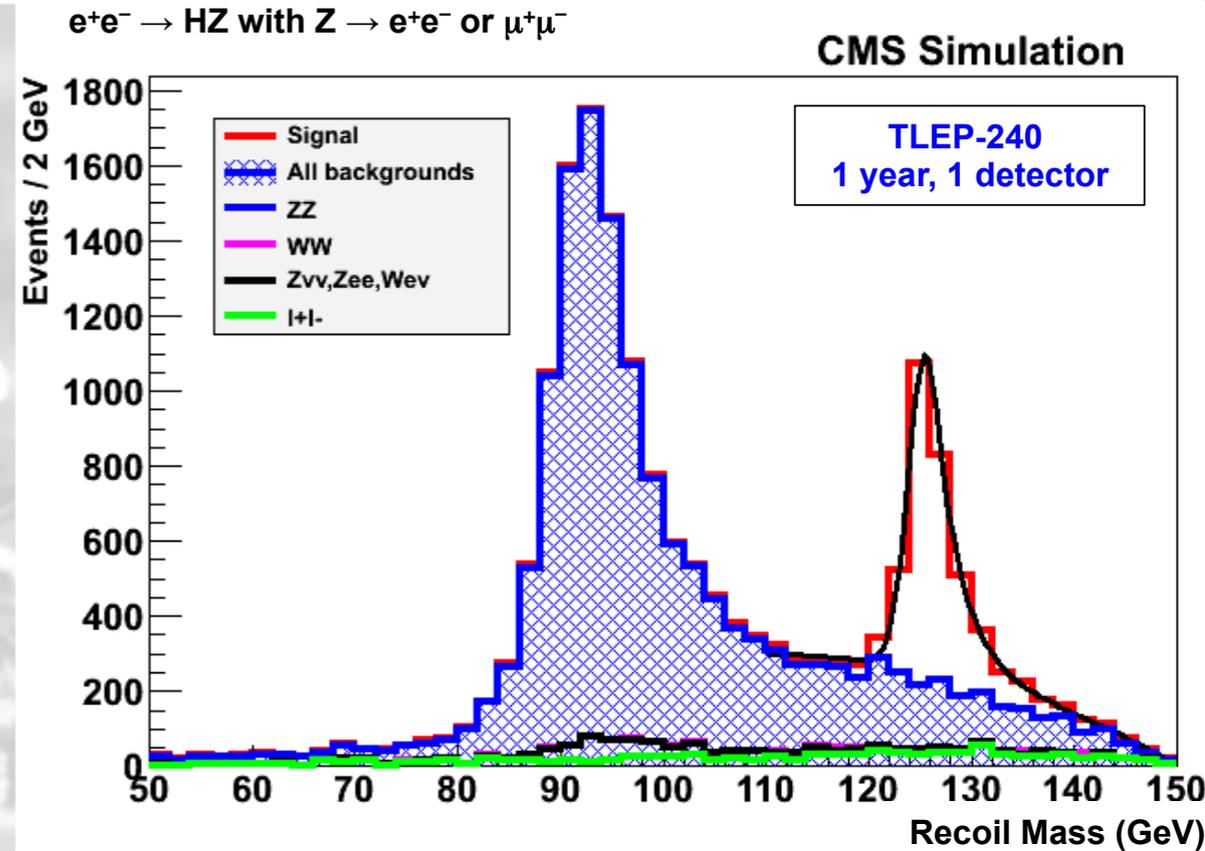


$$\delta N_\nu = 10.5 \frac{\delta \sigma_{had}}{\sigma_{had}} \oplus 3.0 \frac{\delta n_{lept}}{n_{lept}} \oplus 7.5 \frac{\delta L}{L}$$

$$N_\nu = \frac{N(\gamma Z_{inv})}{N(\gamma Z_{ee,\mu\mu})} / \left(\frac{\Gamma_{\nu_l}}{\Gamma_l} \right)_{SM}$$

What can be done @ FCCee?

- Higgs boson production can be tagged from the Z
- One can then veto known Higgs final states
 - mono-Z search
- **Can reach sensitivity $\sim 1\%$**



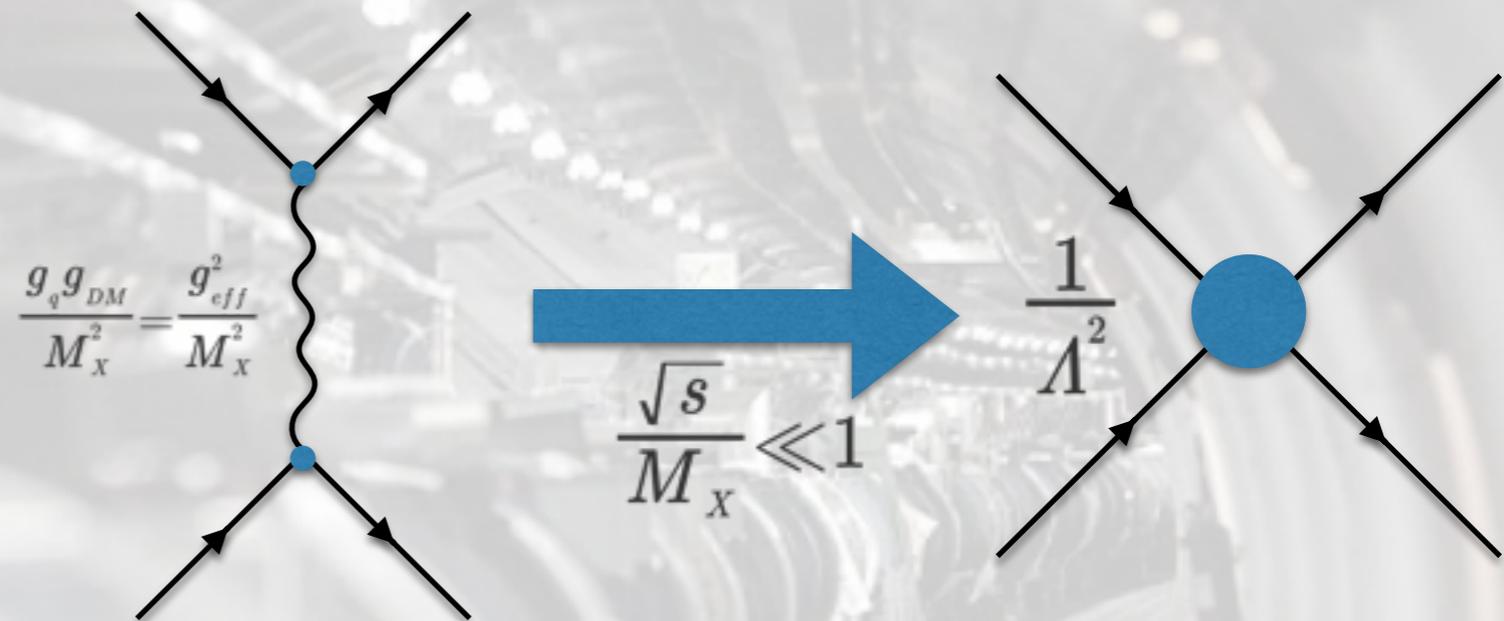


DM mediated by New Heavy Resonances

Effective Field Theory

- DM produced through a heavy mediator X
- In the low-energy limit, one expands the propagator (OPE) and describes the interaction as a local higher-order operator

- Similar to Fermi's description of β decays



- Rather than a dimension-full operator (GF for Fermi), one introduces a new-physics scale Λ , where the UV completion of the theory (the W for Fermi) become relevant

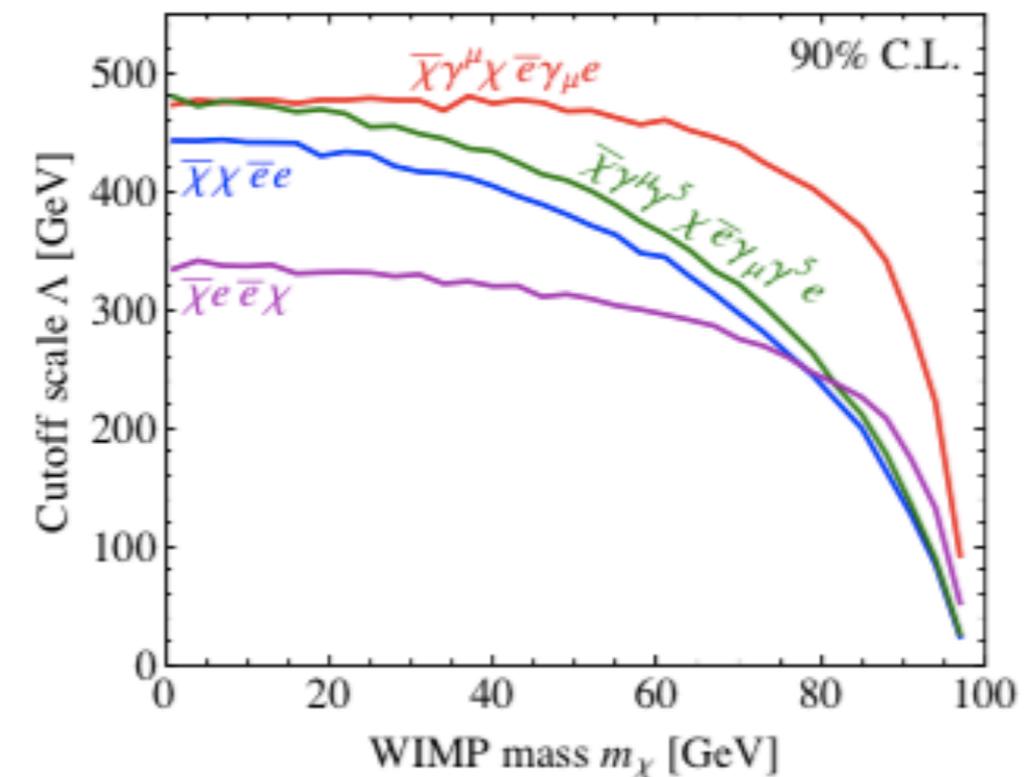
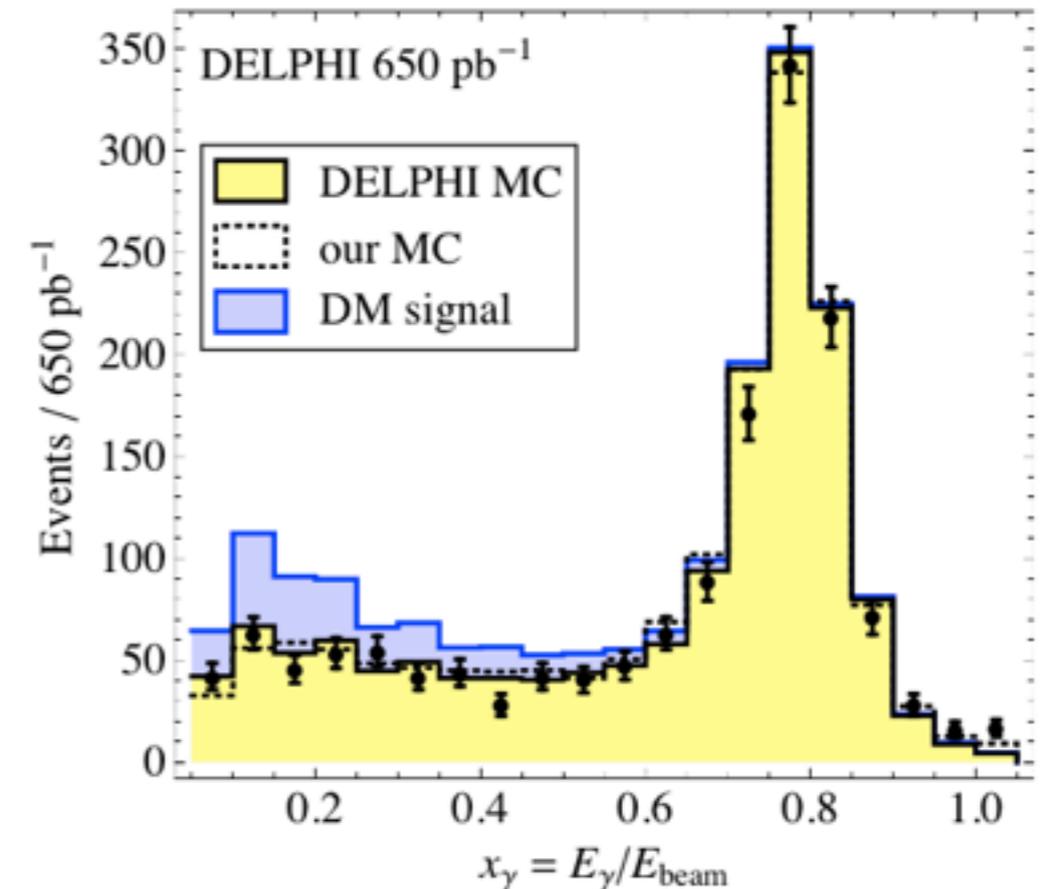
- Different choices of the nature of X result in different effective operators

$$\begin{aligned} \mathcal{O}_V &= \frac{(\bar{\chi}\gamma_\mu\chi)(\bar{\ell}\gamma^\mu\ell)}{\Lambda^2}, && \text{(vector, } s\text{-channel)} \\ \mathcal{O}_S &= \frac{(\bar{\chi}\chi)(\bar{\ell}\ell)}{\Lambda^2}, && \text{(scalar, } s\text{-channel)} \\ \mathcal{O}_A &= \frac{(\bar{\chi}\gamma_\mu\gamma_5\chi)(\bar{\ell}\gamma^\mu\gamma_5\ell)}{\Lambda^2}, && \text{(axial vector, } s\text{-channel)} \\ \mathcal{O}_t &= \frac{(\bar{\chi}\ell)(\bar{\ell}\chi)}{\Lambda^2}, && \text{(scalar, } t\text{-channel)} \end{aligned}$$

DELPHI Monojet Search Recast

P. Fox et al. <http://arxiv.org/abs/1103.0240>

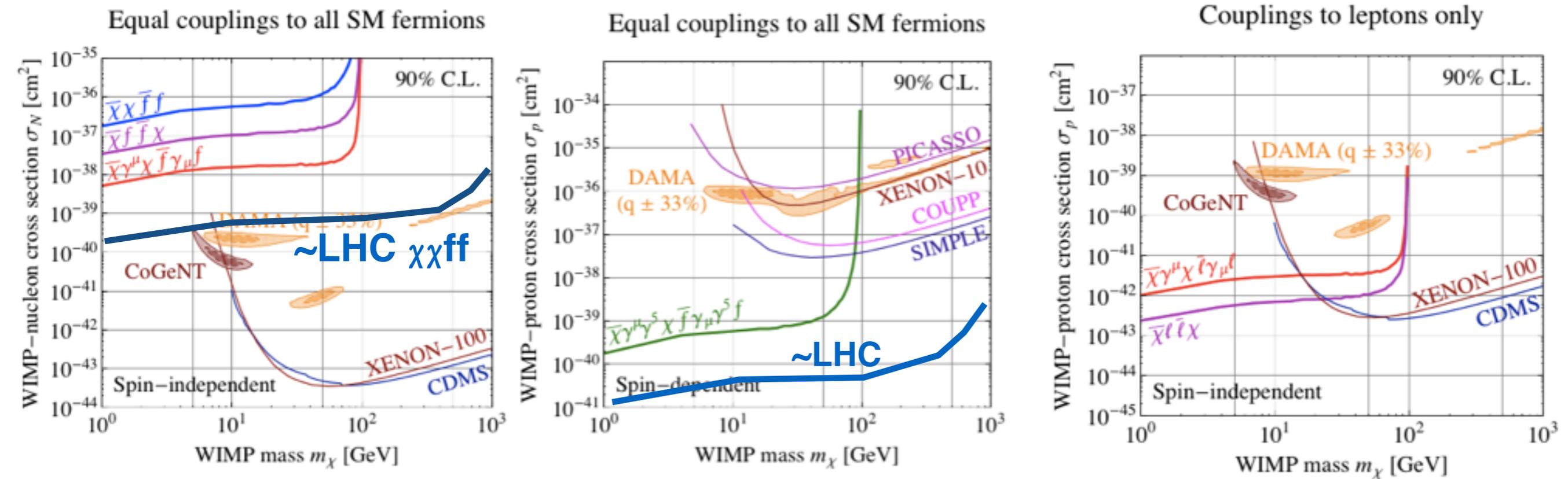
- Search for monophoton from LEP experiments
 - interpreted as SUSY searches
- Recast in EFT limit for different operators
 - reproduce DELPHI selection and plot
 - simulated DM sample for different EFT
 - Interpreted result as limit on Λ vs m_{DM}



Limit on nucleon-DM scattering

P. Fox et al. <http://arxiv.org/abs/1103.0240>

- When the result was presented, bound was relevant at small values of m_{DM}
- Limit less stringent than the LHC
- Need to push to exotic scenarios (DM coupling only to leptons) to cover unexplored territory

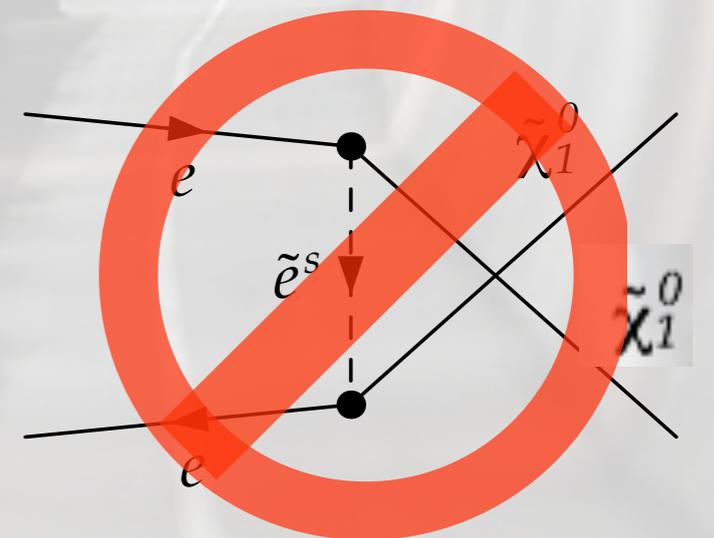
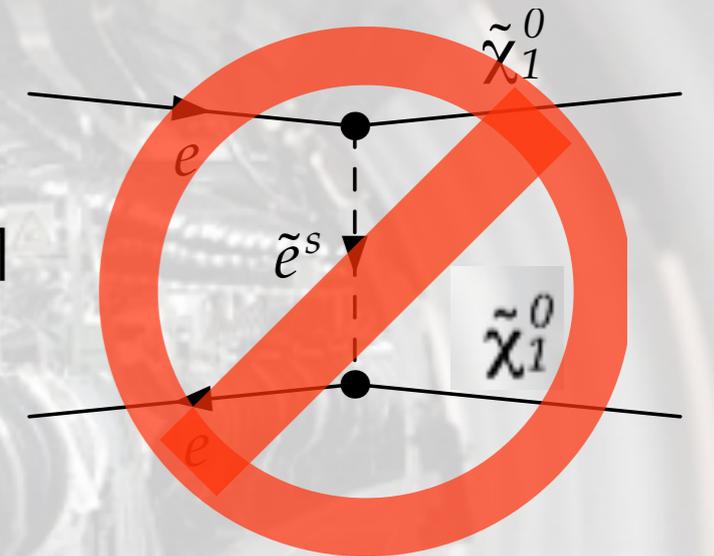
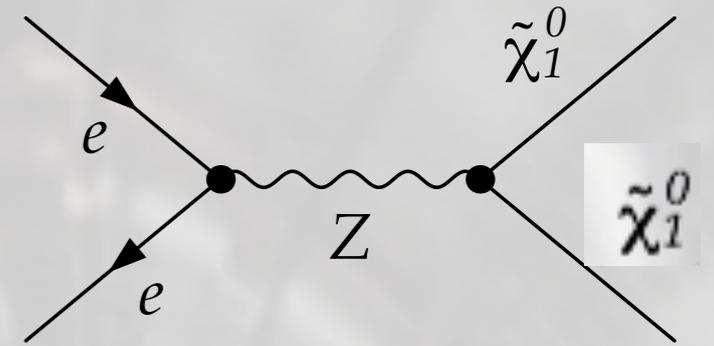




A MonoPhoton Search @ FCCee

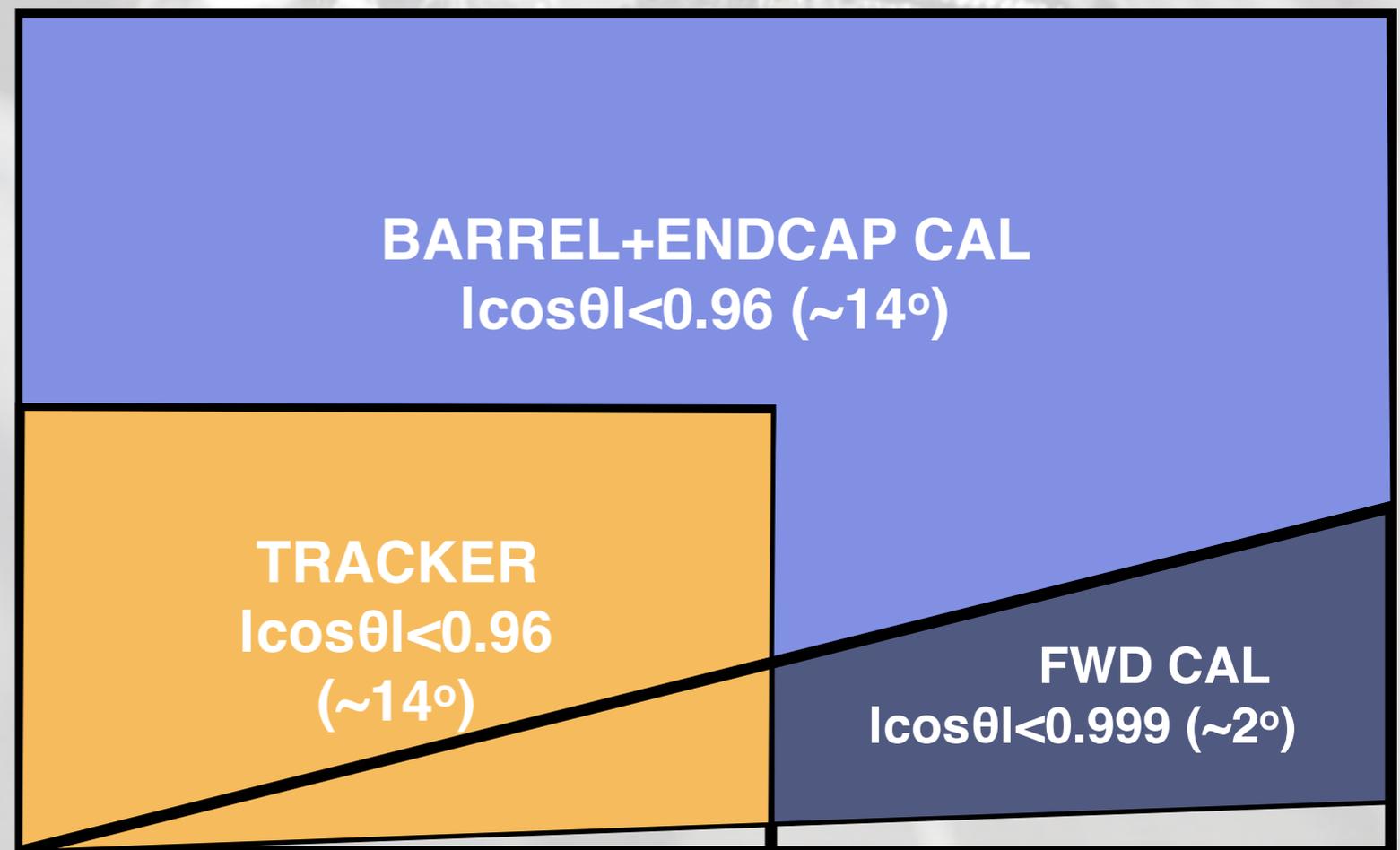
Samples

- Generated events with pythia8 at different \sqrt{s}
 - WW(160 GeV), HZ(240 GeV), and tt (350 GeV)
- No detector effects at this stage
- Only applied acceptance cut (see next slide)
- Simulated relevant SM processes and a prototype of DM signal (neutralino pair production) to derive plausible signal shapes
- For signal, used $\chi^0\chi^0$ production with other SUSY particles decoupled (essentially a Z-portal model)
 - To be extended to the models in the future (t-channel, heavy mediator, etc)
 - No kinematic discrimination vs $\nu\nu\gamma$ for $2m_\chi < m_Z$



Analysis Selection

- Used the ALEPH analysis as a reference, simplifying the detector geometry
- Required one photon within acceptance and $p_T > 0.00357\sqrt{s}$ ($\theta > 2^\circ$)
- Veto on tracks within acceptance ($\theta > 14^\circ$)
- Require less than 1 GeV of additional energy within the acceptance ($\theta > 14^\circ$)
- Require no energy in the fwd part of the detector ($2^\circ < \theta < 14^\circ$)



Selection Efficiency

160 GeV	$m_{\text{DM}} = 15 \text{ GeV}$	$m_{\text{DM}} = 35 \text{ GeV}$	$m_{\text{DM}} = 55 \text{ GeV}$
>1 γ and no track	16%	13%	2%
1 γ	15%	12%	2%
$E_{\text{tot}} < 1 \text{ GeV}$	13%	11%	2%
No Efwd	11%	9%	2%
240 GeV	$m_{\text{DM}} = 15 \text{ GeV}$	$m_{\text{DM}} = 35 \text{ GeV}$	$m_{\text{DM}} = 55 \text{ GeV}$
>1 γ and no track	15%	12%	3%
1 γ	15%	12%	3%
$E_{\text{tot}} < 1 \text{ GeV}$	12%	10%	3%
No Efwd	10%	8%	2%
350	$m_{\text{DM}} = 15 \text{ GeV}$	$m_{\text{DM}} = 35 \text{ GeV}$	$m_{\text{DM}} = 55 \text{ GeV}$
>1 γ and no track	15%	12%	4%
1 γ	14%	12%	4%
$E_{\text{tot}} < 1 \text{ GeV}$	12%	10%	3%
No Efwd	10%	8%	3%

Used higgsino-like DM as a benchmark (too low χ sec for Bino or Wino)
Efficiency computed wrt all events, i.e. not all events have a photon

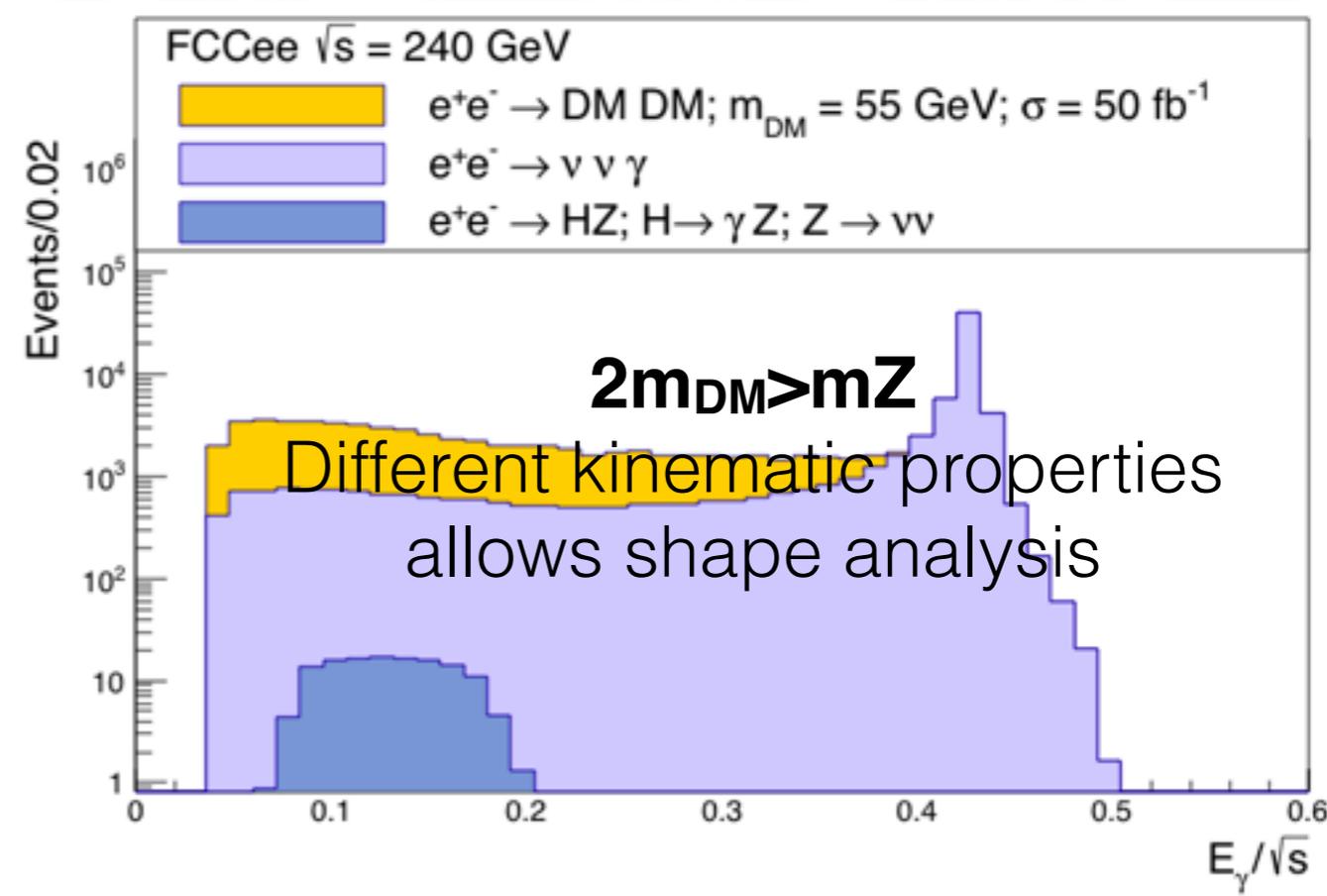
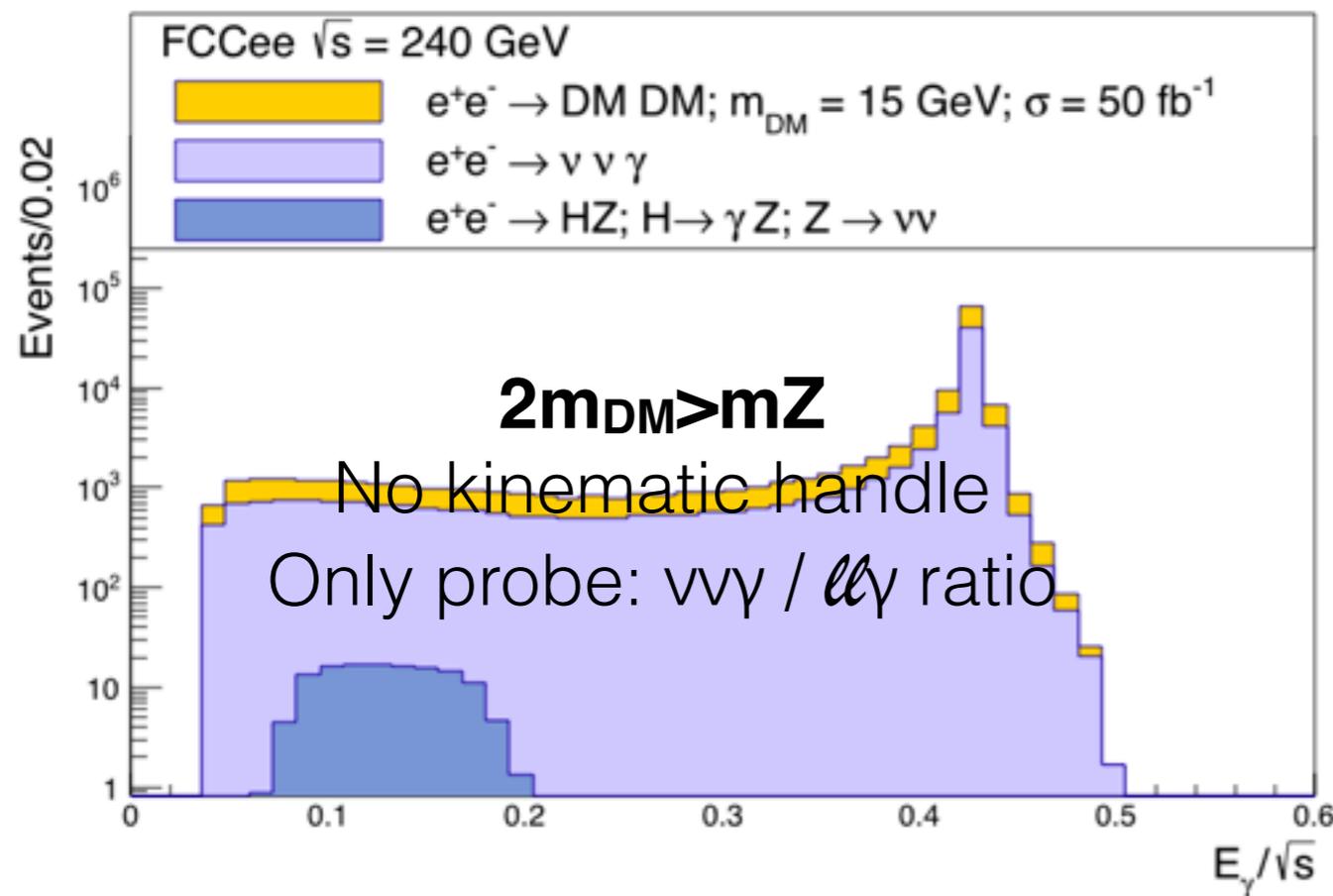
Background Efficiency

vvy	160 GeV	240 GeV	350 GeV
>1 γ and no track	39%	32%	28%
1 γ	38%	32%	28%
Etot<1 GeV	33%	26%	23%
No Efwd	27%	22%	19%
H($\gamma\gamma$)Z(vv)	160 GeV	240 GeV	350 GeV
>1 γ and no track	92%	94%	94%
1 γ	4%	7%	8%
Etot<1 GeV	0.03%	0.1%	0.08%
No Efwd	0.03%	0.1%	0.06%
H(γZ)Z Z\rightarrowvv	160 GeV	240 GeV	350 GeV
>1 γ and no track	93.0%	91.3%	84.1%
1 γ	92.8%	90.8%	81.4%
Etot<1 GeV	87%	79.4%	66.5%
No Efwd	71%	64.3%	54.1%

Other backgrounds ($qq\gamma$, $\gamma\gamma$, $ll\gamma$) have $<10^{-6}$ efficiency (neglected for now)
 Very loose selection on E_γ : expected large bkg to be reduced
 (due to shape separation)

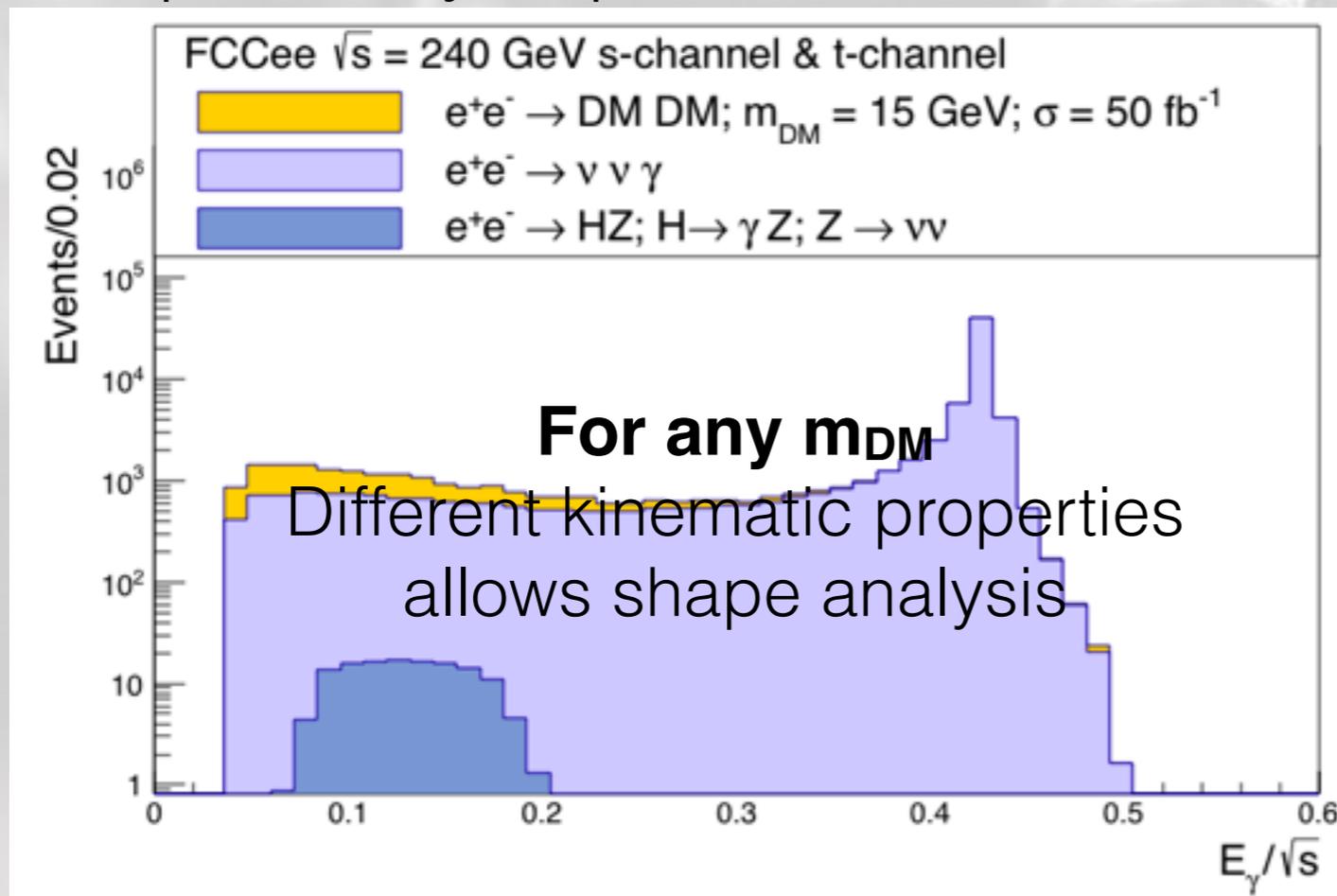
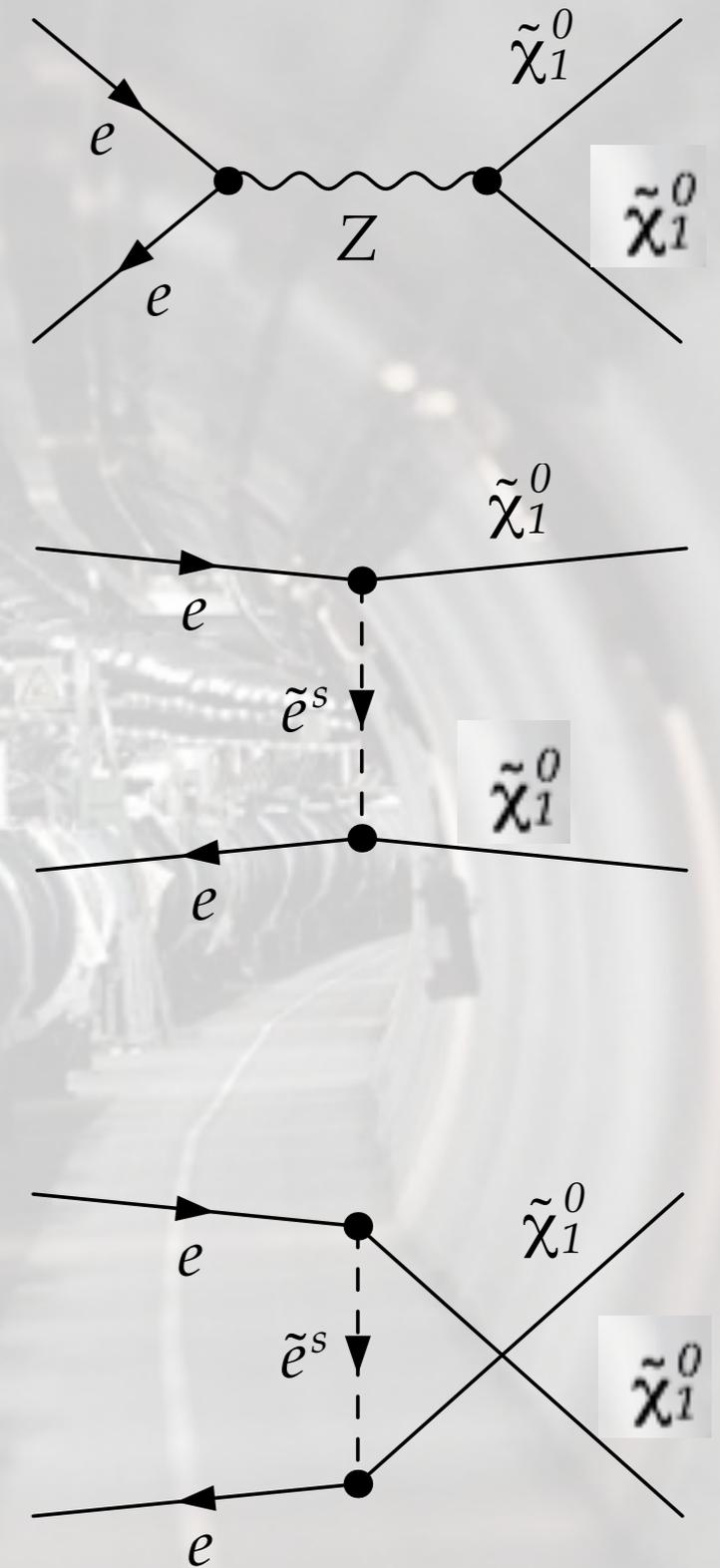
Expected Background

Ebeam	Run time [year]	# IPs	$\int L \cdot dt$	$N_{\nu\nu\gamma}$	$N_{H(\gamma Z)Z Z \rightarrow \nu\nu}$
160 GeV	1	4	42 ab ⁻¹	8.9E+05	-
240 GeV	5	4	11 ab ⁻¹	7.3E+04	134
350 GeV	5	4	3 ab ⁻¹	8.3E+03	22



Adding the t-channel

- Allow for selections (eL and eR) at 500 GeV
 - No impact on Higgsino-like DM
 - Large xsec increase for Wino-like DM
 - t-channel provides for extra kinematic separation
 - Small efficiency at small masses
 - Shape barely depends on DM mass



Benchmark Needed

- At this point, any statement on DM sensitivity would be premature
 - strong dependence on the model
 - strong dependence on the rest of NP spectrum (heavy mediator, heavy (s)leptons, etc)
- We need to establish a minimal set of simplified models to work with
 - Follow the work happening between ATLAS and CMS
 - Work on obvious minimal models (Z and H portals) requiring no extra BSM particle
- This work could then be pushed beyond DM itself (SUSY simplified models, etc)

Summary

- Dark Matter could be produced at colliders
 - through some new heavy mediator which also couples to SM particles
 - coupling directly to SM particles (e.g. Z or H portals)
- Existing bounds from LEP put interesting bounds already
 - Measurement of $\Gamma(Z \rightarrow \text{invisible})$
 - Monophoton search
- Improved sensitivity @ FCCee, thanks to large target $\int \text{Lumi}$
 - $\Gamma(Z \rightarrow \text{invisible}) / \Gamma(Z \rightarrow \ell\ell)$ @ Z pole
 - $\sigma(\nu\nu\gamma) / \sigma(\ell\ell\gamma)$ above Z pole with monophoton (Z portal)
 - Study photon spectrum for monophoton events (other DM models)
 - Add extra sensitivity with $H \rightarrow \text{invisible}$ search
- Quantitative conclusions are model dependent: need to establish benchmark simplified models