



# Structure and properties of oxygen-containing thin films and bulk $\text{MgB}_2$

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## Abstract

## Motivation

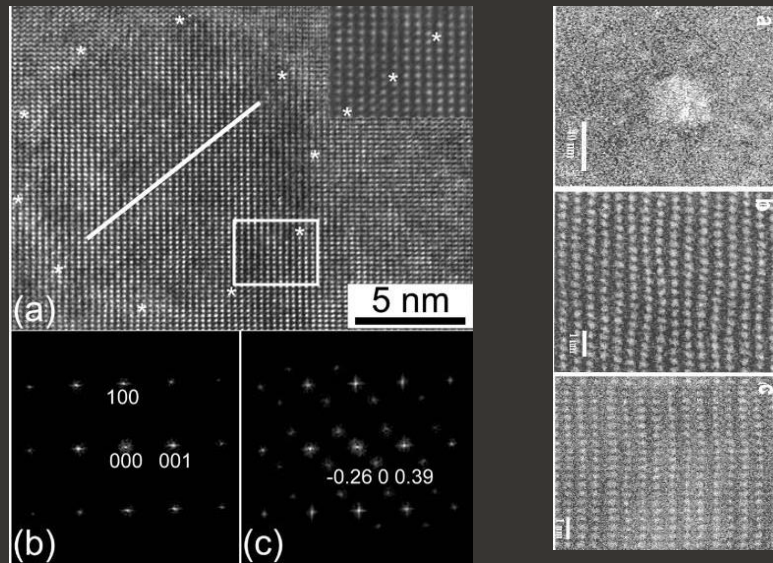
Despite  $\text{MgB}_2$  is oxygen-free superconductor the high affinity of Mg toward oxygen makes practically impossible to synthesize  $\text{MgB}_2$ -based materials free of oxygen.

The structural Auger spectroscopy study of  $\text{MgB}_2$  thin (140 nm) oxygen-containing polycrystalline films obtained by magnetron sputtering with  $j_c(0-1\text{T}, 20\text{K})=7.8-2.7 \text{ MA/cm}^2$

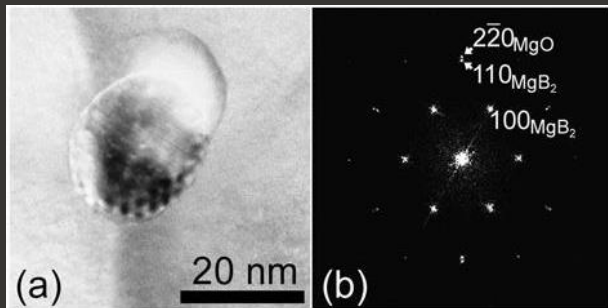
and **99% dense**  $\text{MgB}_2$  bulks synthesized at 2 GPa having  $j_c(0-1 \text{ T}, 20 \text{ K})=0.9 - 0.7 \text{ MA/cm}^2$  or **0.4-0.3 MA/cm<sup>2</sup>** (depending on initial boron),  
up to **80% connectivity**,  
containing up to **94% of shielding fraction**

*allows to conclude that  $j_c$  of  $\text{MgB}_2$  to high extend depends on distribution of admixture oxygen.*

# $\text{Mg}(\text{B},\text{O})_2$ precipitation in $\text{MgB}_2$ (X. Z. Liao et al. J. Appl. Phys. 93, 6208 (2003))



- (a) An HREM image along [010] direction.  
(b) Fourier transformation of the matrix.  
(c) Fourier transformation of the precipitate showing satellite spots with a vector of  $[-0.26, 0, 0.39]$



Longer exposure to trace amount of O at high temperatures also results in the transformation of the precipitates from the hexagonal  $\text{Mg}(\text{B},\text{O})_2$  phase with compositional modulations to the face-centered cubic MgO phase.

$\text{MgB}_2$  grains to form coherently ordered  $\text{MgB}_{2-x}\text{O}_x$  precipitates from about 5 up to 100 nm in size and that such precipitates can act as pinning centers

Ordered replacement of B atoms by O and have the same basic structure as the  $\text{MgB}_2$  matrix but with composition modulations.

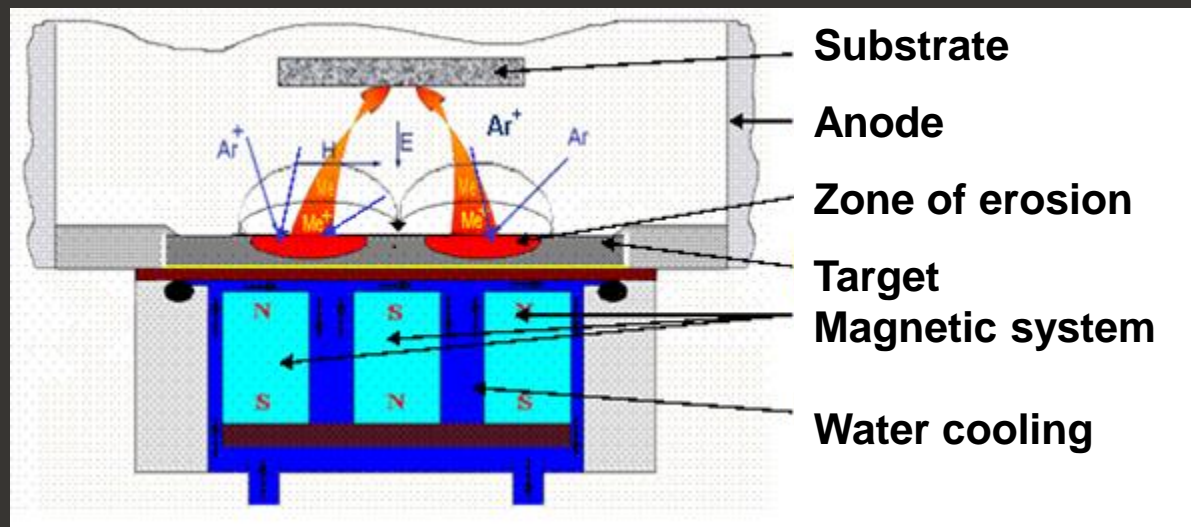
No difference in the lattice parameters between the precipitates and the matrix can be detected in conventional electron diffraction patterns.

However, extra satellite diffraction spots were revealed by HREM study in some directions implying a structural modulation caused by the precipitates.

The periodicity of O atom ordering depends on the concentration of O atoms in the precipitate and primarily occurred in the (010) plane.

# Methods of preparation

## Magnetron sputtering of $\text{MgB}_2$ thin film



The thin  $\text{MgB}_2$  films were deposited by DC magnetron sputtering under 1 Pa Ar pressure on  $8 \times 8 \times 0.2 \text{ mm}^3$  sized substrates from **leucosapphire in (0001) orientation** having room temperature, which subsequently were annealed at 600-650 °C for 5 min in Ar under 10 Pa. For deposition we used  $\text{MgB}_2$  target synthesized by hot pressing at 30 MPa, 800 °C, 1 h from Mg turning (Technical Specification of Ukraine 48-10-93-88) and commercially available precursor B powders HC Starck (grain size 4  $\mu\text{m}$ , 1.5 wt% O, 0.47 wt% C, 0.40 wt% N, 0.37 wt% H).



# Methods of preparation

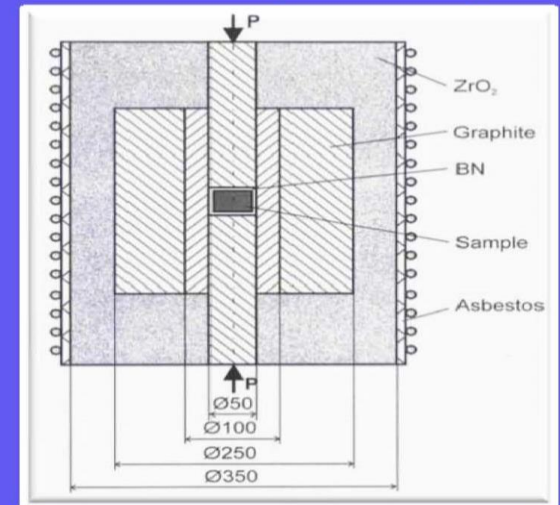
High pressure (2 GPa) and hot pressure (30 MPa) synthesis of bulk  $\text{MgB}_2$



Recessed-anvil-type high-pressure apparatus (HPA)



Hot – P = Hot-pressing  
30 MPa (300 atm)

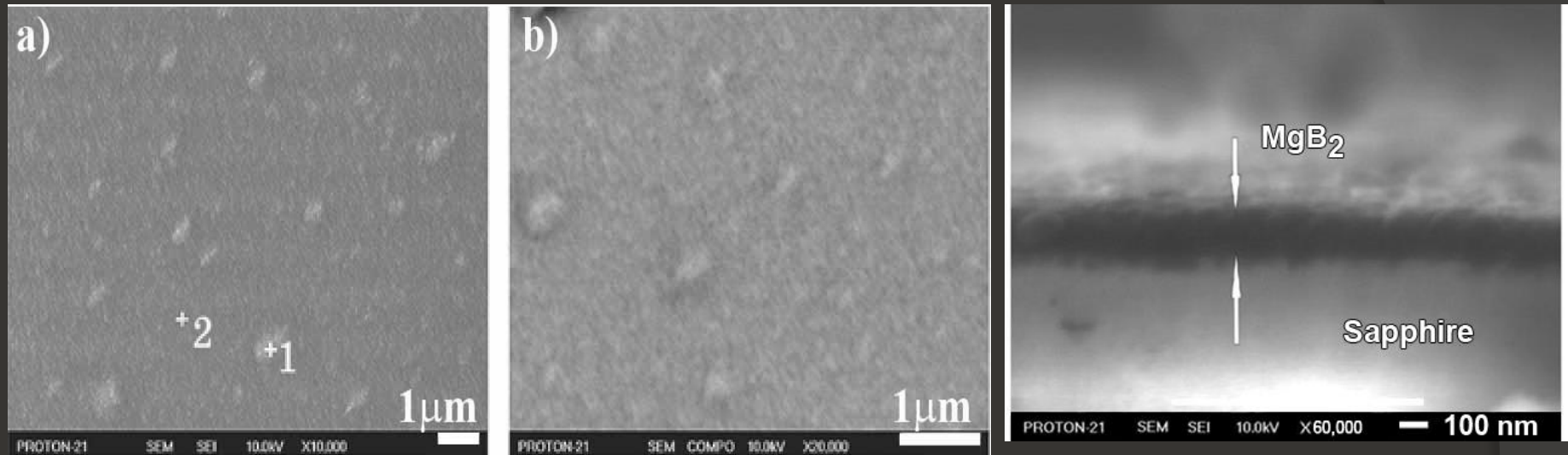


Possibility to press blocks up to  
200 mm in diameter



ASEA (Sweeden) 140 MN effort

# Structure of $\text{MgB}_2$ thin (140 nm) film deposited by magnetron sputtering (**Type 1**)



Structure of  $\text{MgB}_2$  film after etching in Ar of oxidized surface layer in Auger spectrometer JAMP-9500F

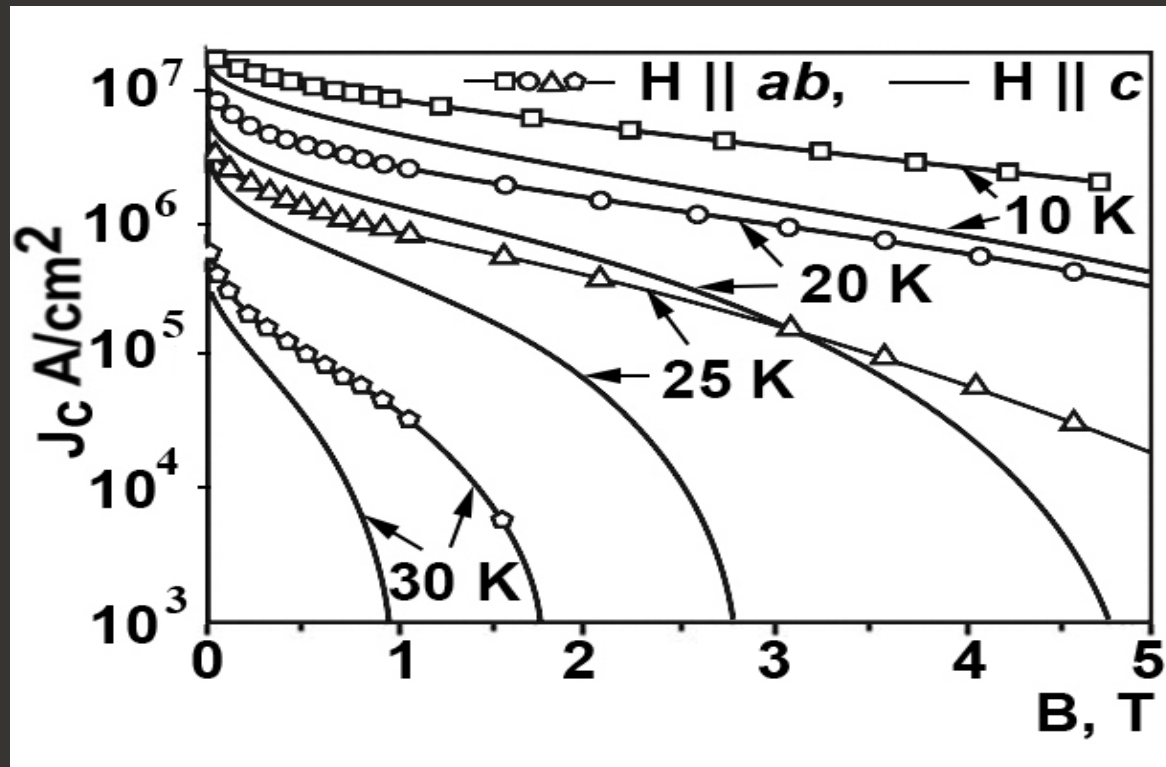
- (a) – SEI (secondary electron image showing the relief of the film) and
- (b) – COMPO (compositional electron image – the brightness increase with increase in Z )

## Composition (Fig.a, at.%)

in point 1: 47.9 % B, 7.1 % C, **31.6 % O**, 13.4 % Mg.  
and in point 2: 58.3 % B, 10.8 % C, **16.6 % O**, 14.3 % Mg.

The structure of film **Type 1** which demonstrates the highest  $J_c$  contains 16.6-31.6 at% of oxygen and nanostructural inhomogenities which can be pinning centers influencing  $J_c$  are much smaller and more homogeneously distributed as compare to bulk  $\text{MgB}_2$

# Critical current density, $j_c$ , vs. magnetic field, $B$ , of thin $\text{MgB}_2$ film (Type 1)

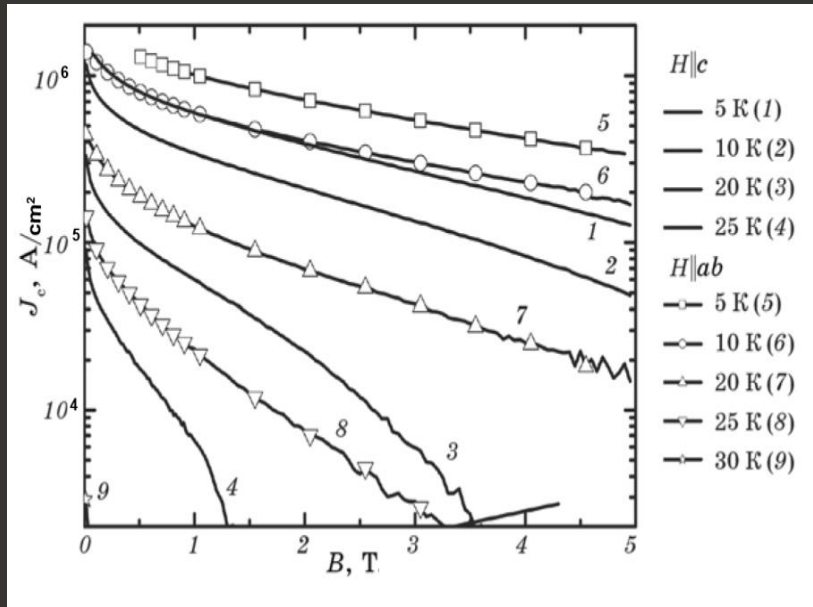


field parallel ( $H \parallel ab$ ) and perpendicular ( $H \parallel c$ ) to the substrate surface

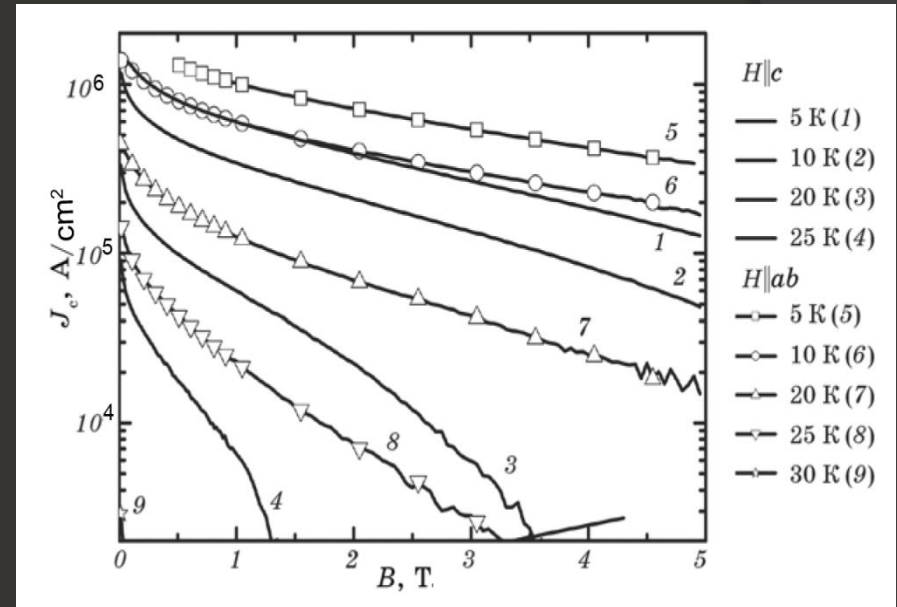
$$j_c(0-1\text{T}, 20\text{K})=7.8-2.7 \text{ MA/cm}^2$$

The thin  $\text{MgB}_2$  films were deposited by DC magnetron sputtering under 1 Pa Ar pressure on  $8 \times 8 \times 0.2 \text{ mm}^3$  sized substrates from leucosapphire in (0001) orientation having room temperature, which subsequently were annealed at 600-650 °C for 5 min in Ar under 10 Pa (Type 1)

# Critical current density, $j_c$ , vs. magnetic field, $B$ , of thin $\text{MgB}_2$ films after repeated annealing



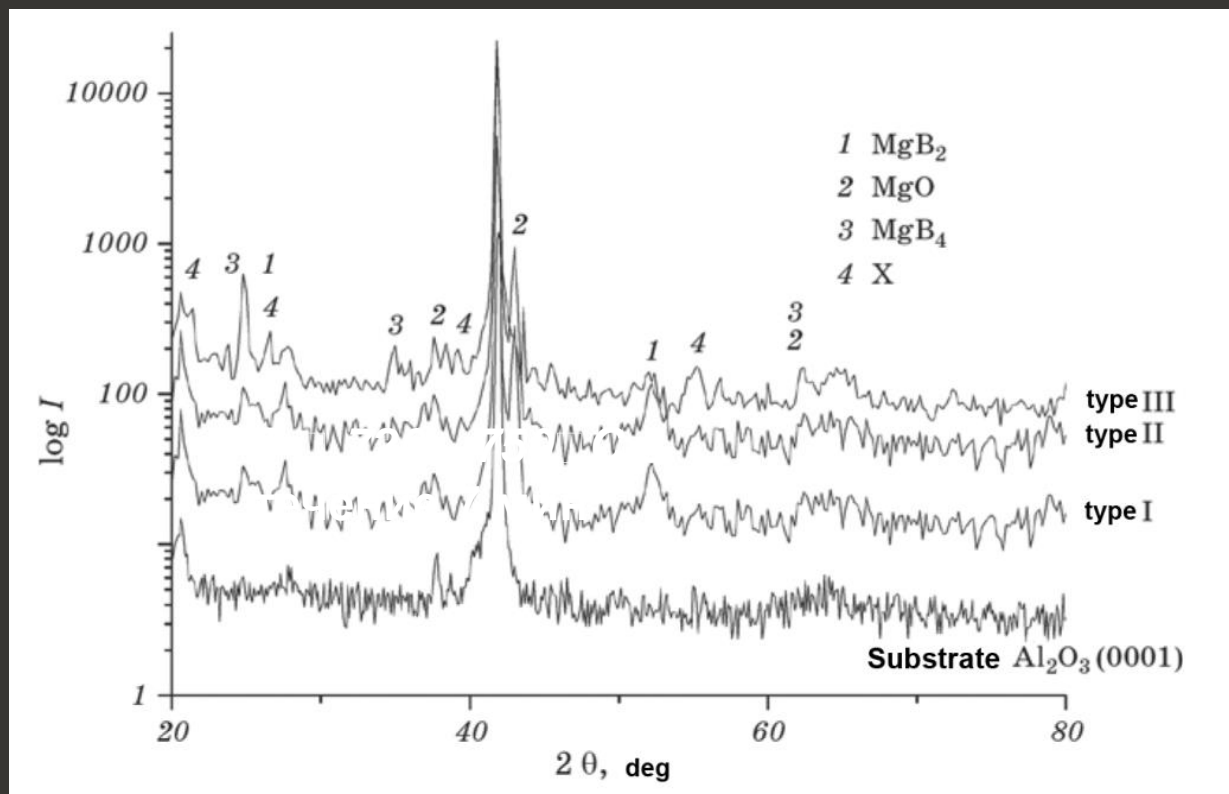
**Type 2** repeatedly annealed at 720—750°C for 7 min in Mg vapor at  $P = 100 \text{ Pa}$



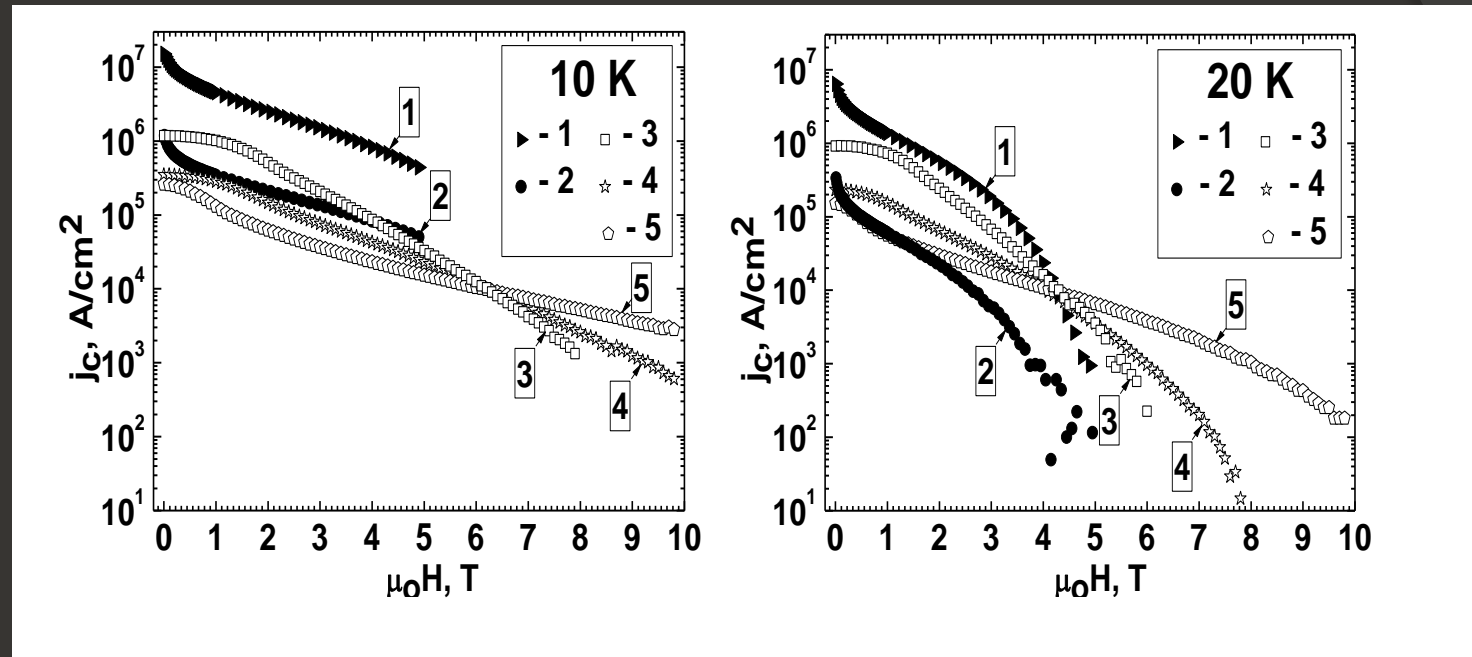
**Type 3** repeatedly annealed at 720—750°C for 7 min vacuumed under  $P = 0.1 \text{ Pa}$ .



# X-ray diffraction patterns of three types thin films



X-ray of films deposited on substrates from leucosapphire in (0001) orientation having room temperature and then heated up to 600-650 °C for 5 min under 10 Pa of Ar (**Type 1**) and films which after this y were repitedly annealed at 720—750 °C for 7 min in Mg vapor at  $P = 100$  Pa (**Type 2**) and in 0.1 Pa vacuum (**Type 3**)



Critical current density,  $J_c$ , vs. magnetic field,  $\mu_0 H$ , for thin films and bulk  $\text{MgB}_2$ :

1 – film as deposited in Ar from  $\text{MgB}_2$  target (Type 1),

2 – the same film after repeated annealing in vacuum (Type 3);

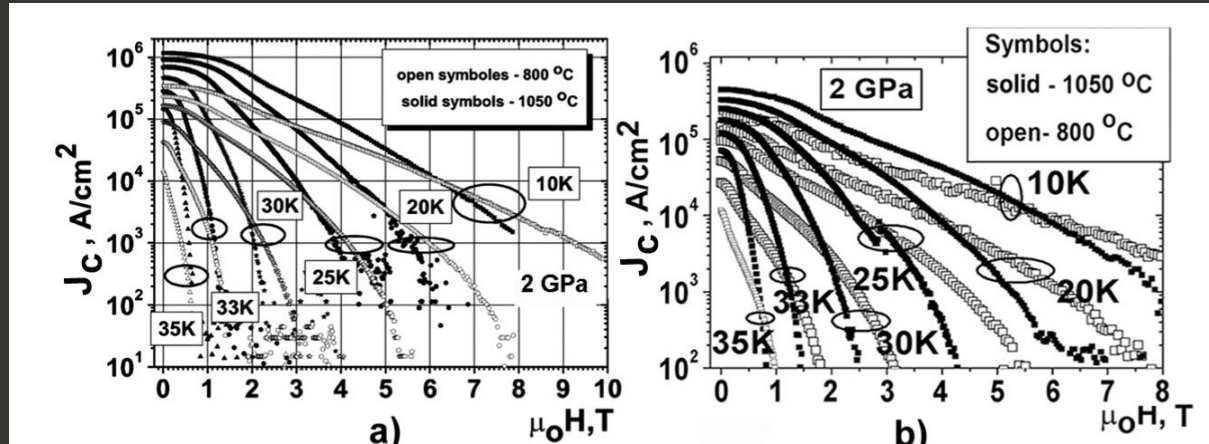
and

3-5 - bulk  $\text{MgB}_2$  high-pressure (2 GPa, 1 h) synthesized at 1050, 800 and 600 °C, respectively;

# Introduction

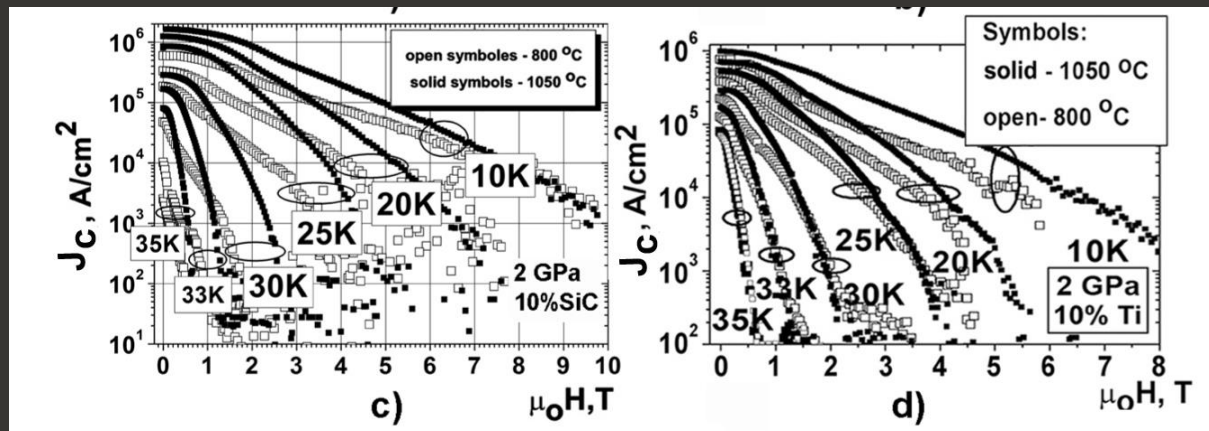
Critical current density,  $j_c$ , vs. magnetic field,  $\mu_0 H$ , of MgB<sub>2</sub> materials synthesized at **2 GPa** for 1 h from Mg+2B at **800** and **1050 °C**;

From B (1)  
with grain  
size  $<5 \mu\text{m}$ ,  
**0.66 wt% O**,  
0.31 wt% C,  
0.48 wt% N,  
0.32 wt % H



$$j_c(0-1 \text{ T}, 20 \text{ K}) = \mathbf{0.9 - 0.7 \text{ MA/cm}^2} \quad j_c(0-1 \text{ T}, 20 \text{ K}) = \mathbf{0.4-0.3 \text{ MA/cm}^2}$$

From B(1)  
with  
**10 wt% SiC**



$$j_c(0-1 \text{ T}, 20 \text{ K}) = \mathbf{1.2 - 1.0 \text{ MA/cm}^2} \quad j_c(0-1 \text{ T}, 20 \text{ K}) = \mathbf{0.7-0.5 \text{ MA/cm}^2}$$

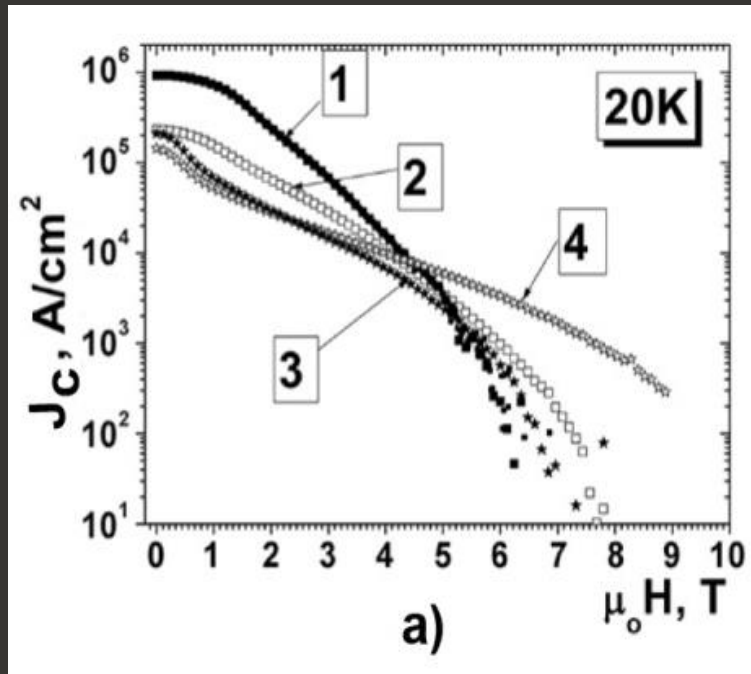
# Motivation

From B (2)  
with grain  
size  $4 \mu\text{m}$ ,  
**1.5 wt% O**,  
0.47 wt% C,  
0.40 wt% N,  
0.37 wt% H

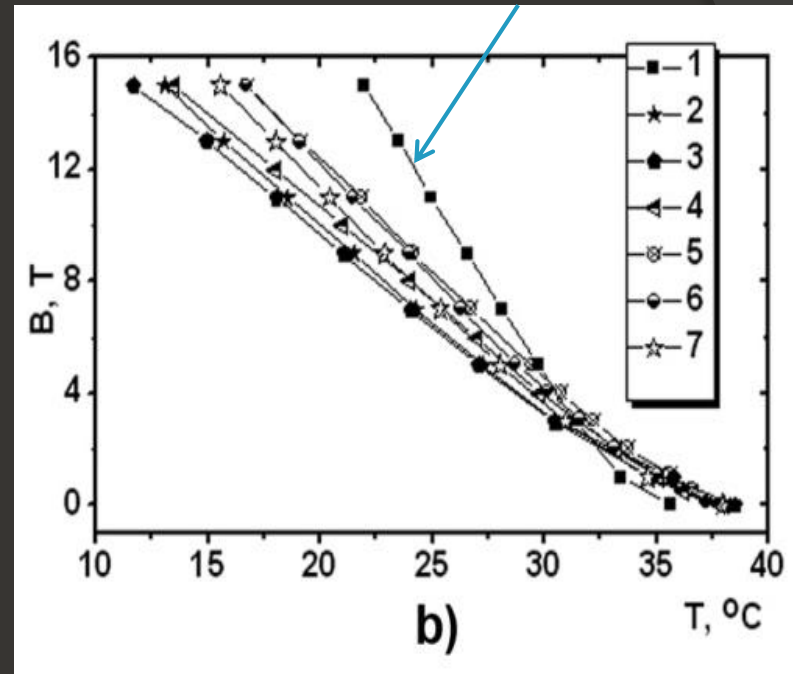
From B(2)  
with  
**10 wt% Ti**

# Introduction

$B_{c2}$  (22 K)=15 T (No 1) and extrapolated  $B_{c2}(0\text{ K})=42.1\text{ T}$ , as well, as irreversibility fields:  $B_{irr}(18.5\text{ K})=15\text{ T}$  and  $B_{irr}(0\text{ K})=32.5\text{ T}$



a)



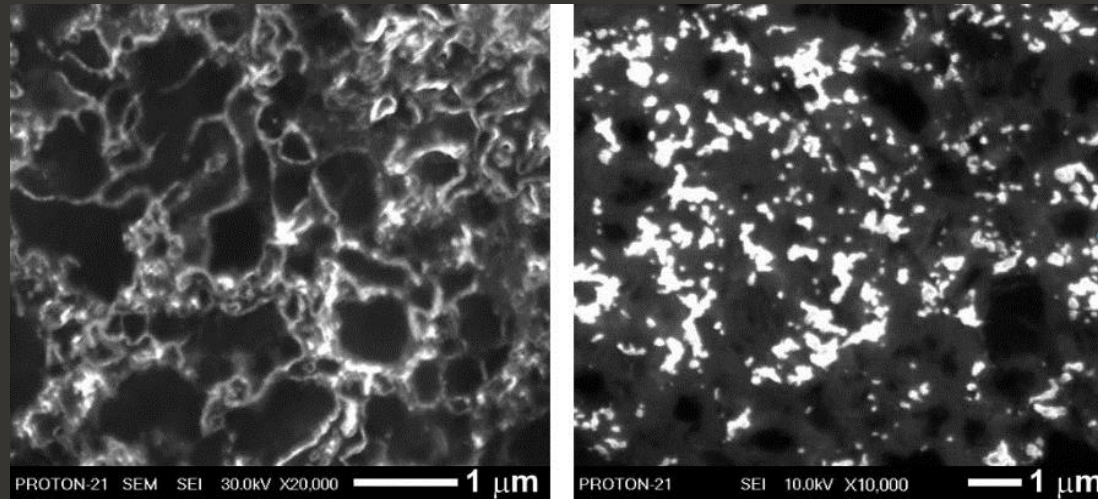
b)

(a) Comparison of  $J_c$  at 20 K in bulk  $\text{MgB}_2$  prepared under various temperatures at 2 GPa:  
 1, 2: samples made of B(I) at 1050 °C, ( $k=0.43$ ) and 800 °C, ( $k=0.37$ ), respectively;  
 3, 4 of B(II) at 800 °C and 600 °C, ( $k=0.32$ ), respectively.

(b)  $B_{c2}$  of bulk  $\text{MgB}_2$  prepared from Mg:2B :  
 1:B(II) at 2 GPa, 600 °C, 1 h, (HP);  
 2:B(III) at 30 MPa, 800 °C, 2h, (HotP);  
 3: B(III) at 50 MPa, 600-1050 °C, 1 h, (SPS);  
 4: B(III) at 2 GPa, 900 °C, 1 h (HP);  
 5:B(V)+10 wt.% Zr at 2 GPa, 800 °C, 1 h (HP);  
 6:B(V)+10 wt.% Ti at 2 GPa, 800 °C, 1 h (HP);  
 7:B(I)+10 wt.% SiC at 2 GPa, 1050 °C, 1 h (HP).



## Structural transformations in bulk $\text{MgB}_2$ with manufacturing temperature increase from 800 to 1050 °C (under 2 GPa)



2 GPa, 800 °C, 1 h

2 GPa, 1050 °C, 1 h

99% dense  $\text{MgB}_2$

$j_c(0-1 \text{ T}, 20 \text{ K}) =$   
0.4-0.3 MA/cm<sup>2</sup>

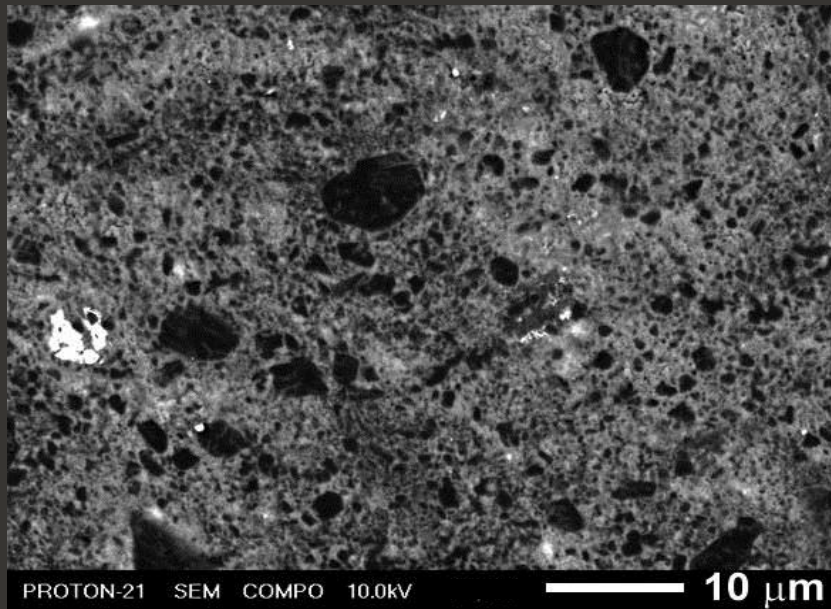
up to 80% connectivity

up to 94% of shielding  
fraction

As Auger analysis shown

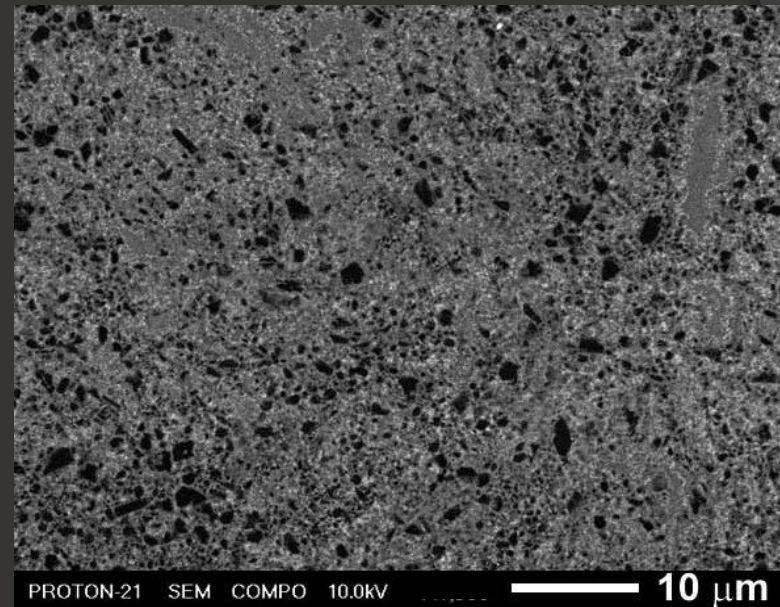
- material matrix (for the sample synthesized at 2GPa 1050 °C) showed near  $\text{MgB}_{2.2-1.7}\text{O}_{0.4-0.6}$  stoichiometry as a result of multiple analyzing after multiple etching steps (providing depth profile)
- bright Mg-B-O nanolayers (10-20 nm) had near  $\text{MgB}_{1.2-2.7}\text{O}_{1.8-2.5}$  stoichiometry
- separate bright Mg-B-O inclusions had near  $\text{MgB}_{0.6-0.8}\text{O}_{0.8-0.9}$  stoichiometry
- higher magnesium borides  $\text{MgB}_x$  which appear the most black have near  $\text{MgB}_{11-13}\text{O}_{0.2-0.3}$  stoichiometry

***The size and amount of higher magnesium borides  
(black inclusions with near  $\text{MgB}_{11-13}\text{O}_{0.2-0.3}$   
stoichiometry according to Auger study) decreasing  
with manufacturing temperature increase***



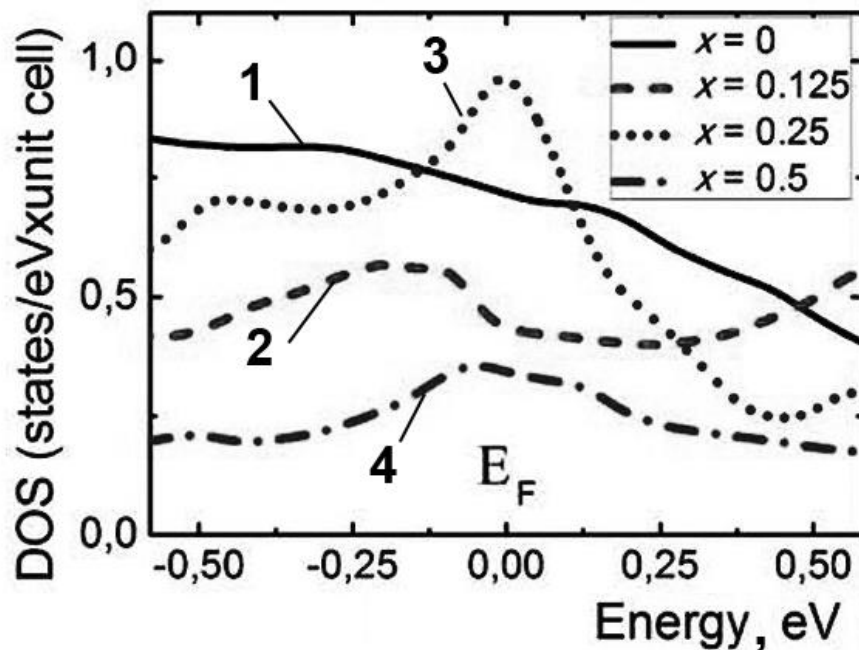
**Mg:2B, 2 GPa, 800 °C, 1 h**

The Berkovich nanohardness and Young modulus of such inclusions as estimated under the 10-60 mn-load were  $32.2 \pm 1.7$  and  $385 \pm 14$  GPa, respectively, what were even higher than that for sapphire:  $31.1 \pm 2.0$  and  $416 \pm 22$  GPa.



**Mg:2B, 2 GPa, 1050 °C, 1 h**

## Density of states (DOS)



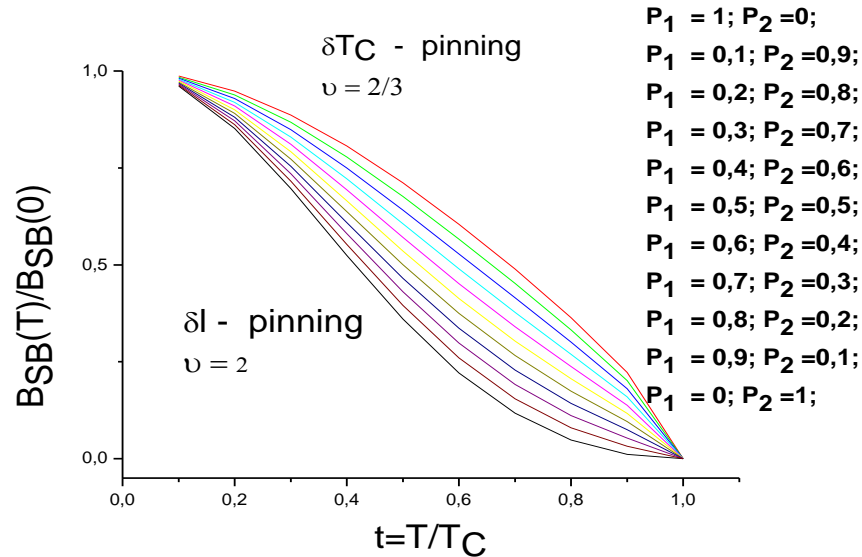
Calculated density of states (DOS)  
for  $\text{MgB}_2$  (1)  
and  $\text{Mg}(\text{B}_{1-x}\text{O}_x)_2$ :  $\text{MgB}_{1.75}\text{O}_{0.25}$  (2),  
 $\text{MgB}_{1.5}\text{O}_{0.5}$  (3)  
 $\text{MgBO}$  (4).

We have applied the density functional theory [1] based on a full-potential linearized augmented plane wave method with the generalized gradient correction to exchange-correlation potential [2] (the program package WIEN2k [3]). In the nearest vicinity of the Fermi energy  $E_F$  corresponding dependencies are not constant and their changes with  $x$  are non-monotonous. But, in any case, all studied compounds are conductors with a non-zero density of states at the Fermi level. Large density of delocalized states proves **the metallic type of the charge transport in them.**

- [1] Parr R G and Yang W 1989 Density-Functional Theory of Atoms and Molecules, *Oxford Univ. Press.*
- [2] Perdew J P, Burke S and Ernzerhof M Phys. 1996 *Rev. Lett.* 77 3865
- [3] Blaha P, Schwarz K, Madsen G K H, Kvasnicka D and Luitz J 2001 WIEN2K, An augmented plane wave + local orbitals program for calculating crystal properties. *Techn. Univ. Wien.*

The highest DOS was observed in the case of  $\text{MgB}_{1.5}\text{O}_{0.5}$  (which was even higher than that for  $\text{MgB}_2$ ) and according to Auger estimation the superconducting matrix demonstrated near  $\text{MgB}_{2.2-1.7}\text{O}_{0.4-0.6}$  stoichiometry.

# Estimation of the type of pinning



**Dependences  $B_{SB}(T)$**

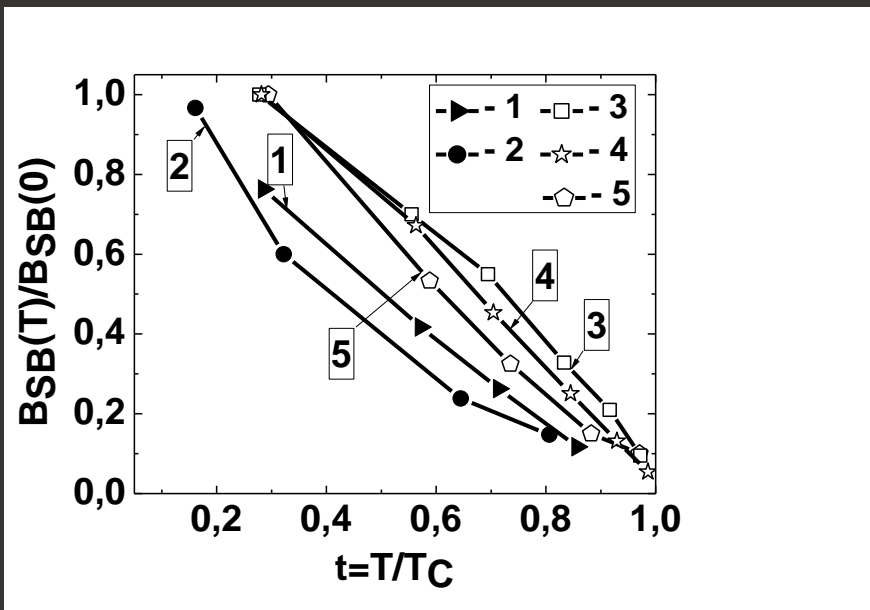
$$B_{SB}(T) = P_1 B_{SB}^{T_C} + P_2 B_{SB}^l$$

$$P_1 - \delta T_C; P_2 - \delta l$$

The work [G.Blatter, et.al, *Reviews of Modern Physics* 1994, v. 66, No. 4, p.p. 1125 – 1380] considers the models of Abrikosov vortices behavior in HTS and the following types of their pinning:  $\delta T_C$  pinning (in the areas of fluctuation of critical temperature,  $T_C$ ) and  $\delta l$  pinning (in the areas of fluctuation of free path length,  $l$ ), their combinations. The partial input of  $\delta T_C$  and  $\delta l$  into collective pinning was estimated based on the determination of temperature dependence  $B_{SB}(T)$  from  $J_C(B)$  temperature dependences using equation  $B_{SB}(T) = P_1 B_{SB}^{T_C} + P_2 B_{SB}^l$ , where  $B_{SB}$  is the field in which replacement of single vortex pinning for pinning of small vortex bundle takes place.



# Temperature dependences of transitions between single vortex pinning and small vortex bundle pinning



The estimated contributions in pinning were:

1-  $P_1=0.5$ ,  $P_2=0.5$ ;

2 -  $P_1=0.2$ ;  $P_2=0.8$ ;

3 -  $P_1=0.7$ ,  $P_2=0.3$ ;

4 -  $P_1=0.6$ ,  $P_2=0.4$ ;

5 -  $P_1=0.4$ ;  $P_2=0.6$

( $P_1$  -  $\delta-T_c$ ;  $P_2$  -  $\delta-I$ )

1 – film as deposited in Ar from MgB<sub>2</sub> target (Type 1),

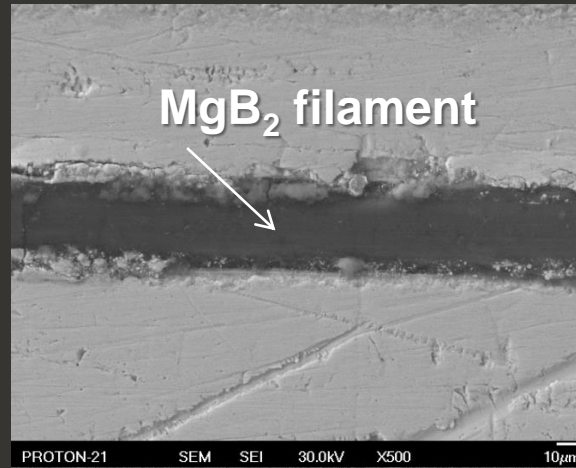
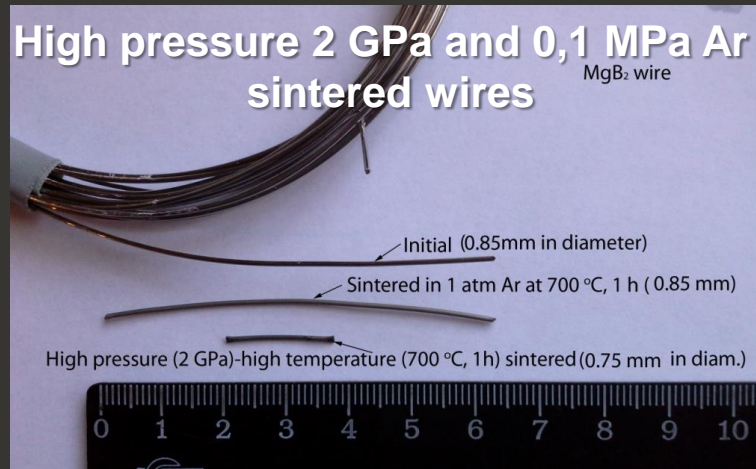
2 – the same film after repeated annealing in vacuum (Type 3);

and

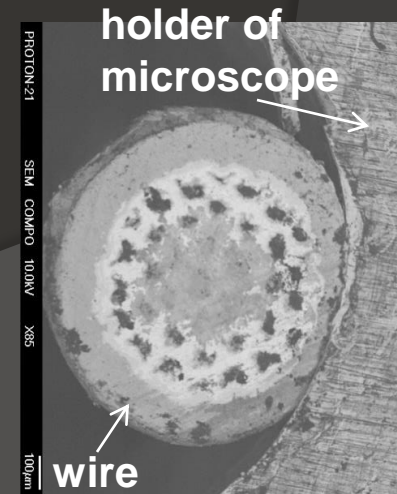
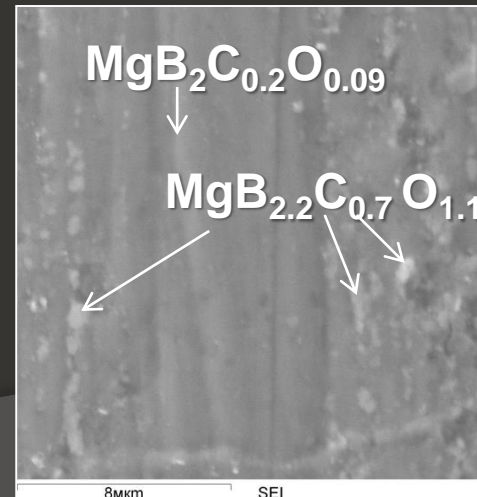
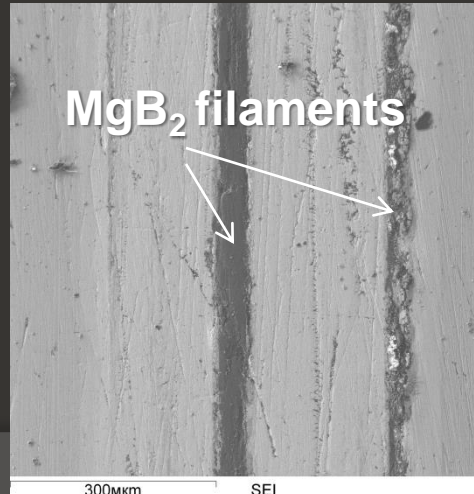
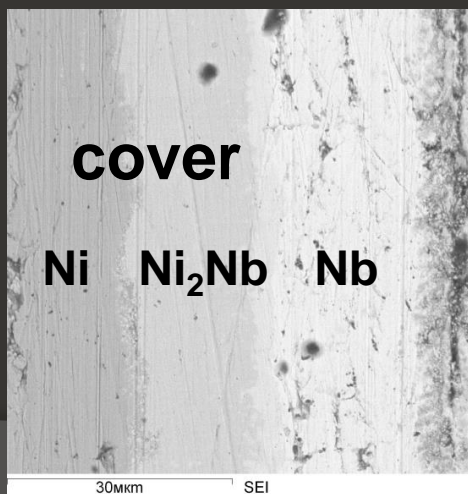
3-5 - bulk MgB<sub>2</sub> high-pressure (2 GPa, 1 h) synthesized at 1050, 800 and 600 °C, respectively;

*It has been shown that for attaining maximal  $J_c$  in low and medium magnetic fields in bulk MgB<sub>2</sub> the most effective is  $\delta-T_c$  pinning or combined  $\delta-T_c$  (dominant) with  $\delta-I$ . For attaining high  $J_c$  in high magnetic fields the  $\delta-I$  pinning should be prevailed in bulk MgB<sub>2</sub>. For nanocrystalline oxygen-containing films the increase of  $J_c$  in low and high magnetic fields correlates with increase of  $\delta-T_c$  pinning contribution. Films with combined type of pinning with equal partial contributions of  $\delta-T_c$  and  $\delta-I$  ( $P_1=P_2=0.5$ ) demonstrate the  $J_c(0\text{ T})=1.8 \times 10^7\text{ A/cm}^2$  and  $J_c(5\text{ T})=2 \times 10^6\text{ A/cm}^2$  at 10 K.*

# High-pressure (2 GPa, 800 °C, 1 h) treated wire from $\text{MgB}_2$



## Panorama of $\text{MgB}_2$ filament after high pressure-high temperature treatment

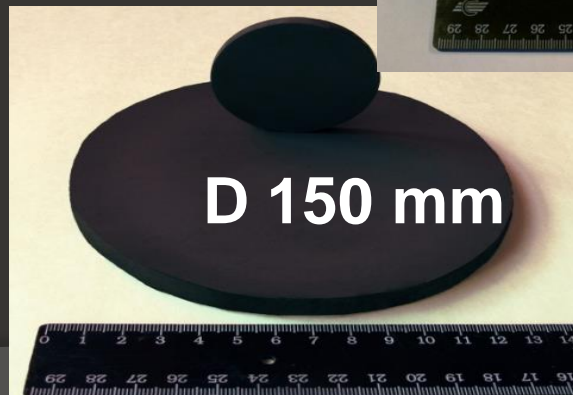
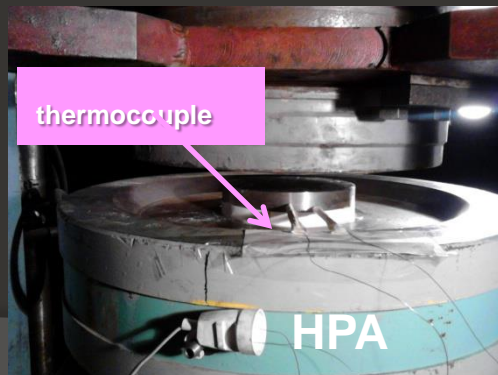
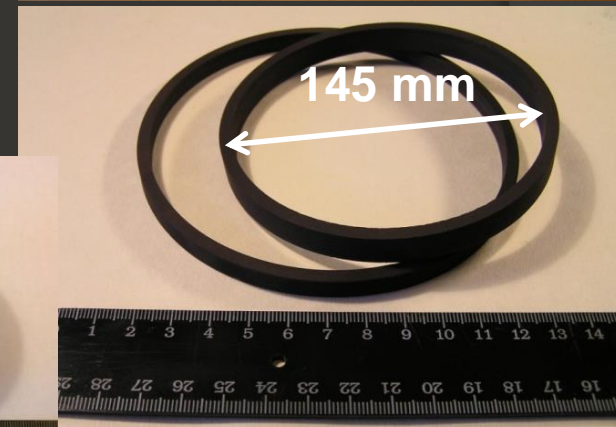
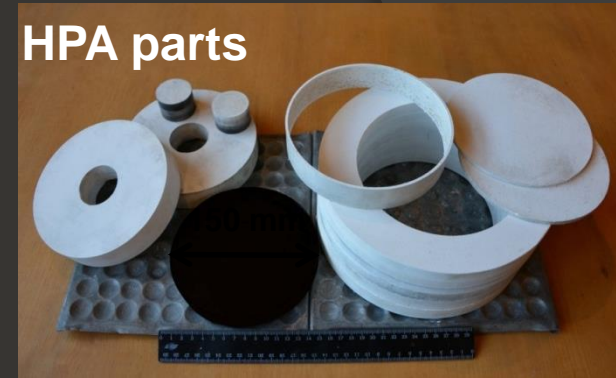




HPA cylinder-piston(  $V_p = 6300 \text{ cm}^3$ )  $P = 1\text{-}2 \text{ GPa}$ ,  
 $T = 600\text{-}1100 \text{ }^\circ\text{C}$ , Press ASEA 140 MN effort



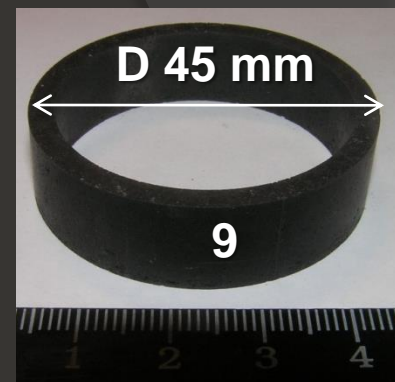
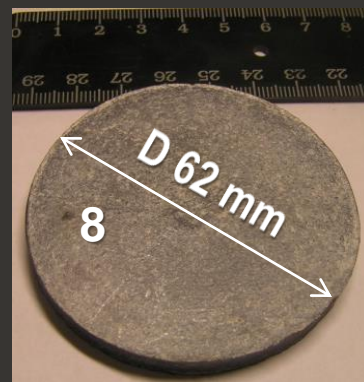
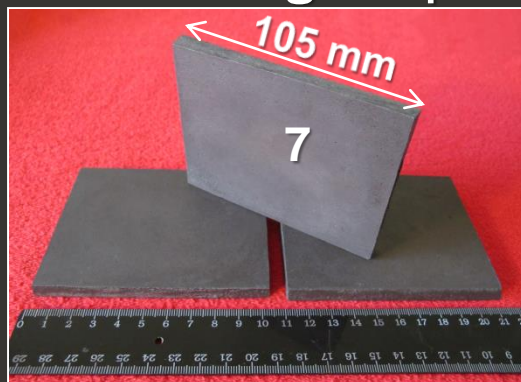
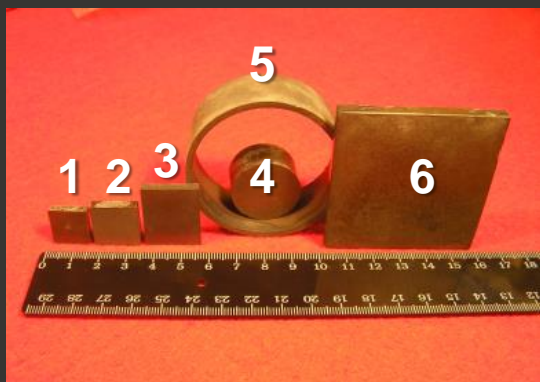
HPA parts



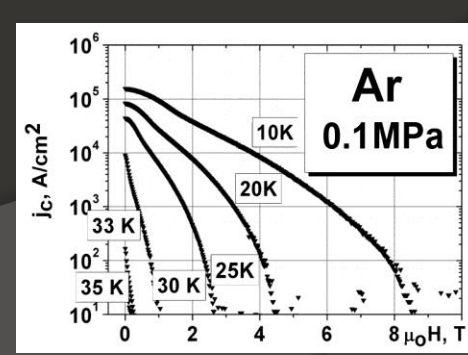
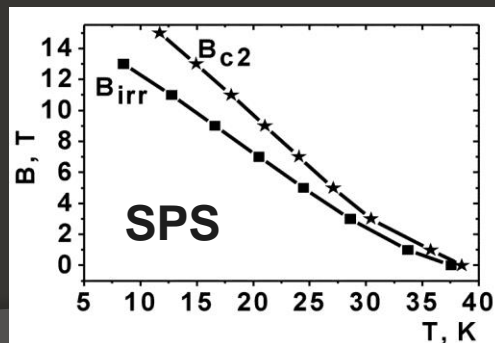
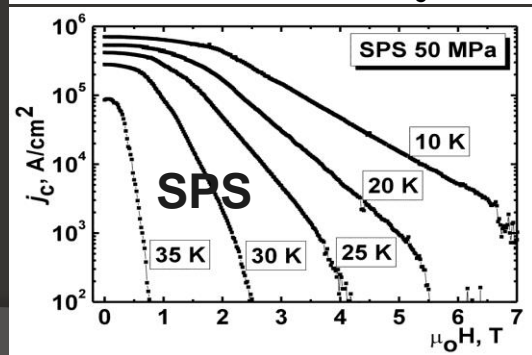
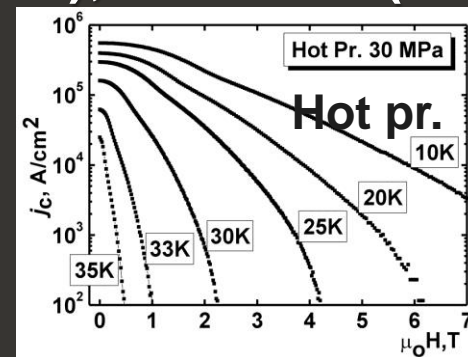
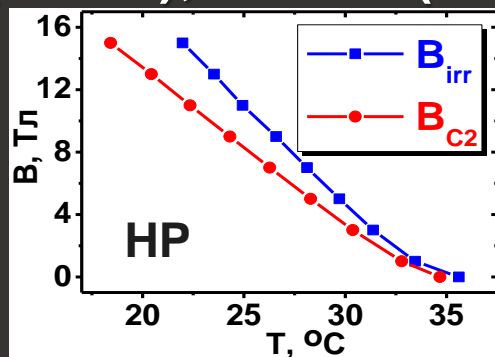
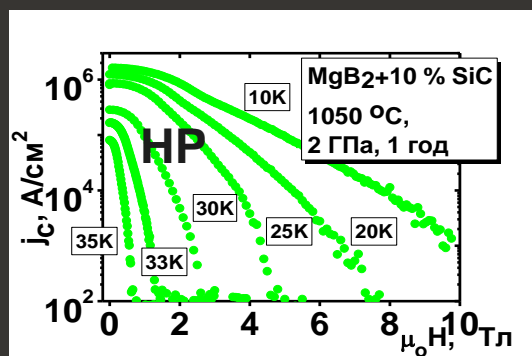
Items from  $\text{MgB}_2$

# MgB<sub>2</sub>-based superconductors

working temperature 20-30 K



1-4, 8, 9 - High pressed (2 GPa); 5 – HIP (0.1 GPa); 7- Hot Pr. (30 MPa)





## Conclusions

- The comparison of nanostructures of  $\text{MgB}_2$  oxygen-containing thin film (with thickness of 140 nm) and highly dense oxygen-containing bulk materials demonstrating high SC performance gave us reasons to conclude that **critical current density,  $j_c$ , of  $\text{MgB}_2$  –based materials to high extend depends on distribution of admixture oxygen** what can be regulated by manufacturing conditions and by synthesis temperature, in particular.
- The higher  $j_c$  of  $\text{MgB}_2$  thin films as compared to bulk was explained by the **higher density of finer oxygen-enriched Mg-B-O inhomogeneities** in the film structure.
- The results of calculations of the electron density of states in  $\text{MgB}_2$ ,  $\text{MgB}_{1.75}\text{O}_{0.25}$ ,  $\text{MgB}_{1.5}\text{O}_{0.5}$  and  $\text{MgBO}$  showed that **all the compounds are conductors with metal-like behavior** with maximum for  $\text{MgB}_{1.5}\text{O}_{0.5}$  (what it is even higher than for oxygen-free  $\text{MgB}_2$ ).
- The stoichiometry of the matrix of bulk  $\text{MgB}_2$  materials having high SC performance was about  **$\text{MgB}_{2.2-1.7}\text{O}_{0.4-0.6}$**  (according to Auger quantitative analysis), i.e. very near to the composition with calculated highest DOS:  **$\text{MgB}_{1.5}\text{O}_{0.5}$** . The introduction of high amount of oxygen into  $\text{MgB}_2$  structure cannot dramatically reduce material's  $T_c$  and allows obtaining high  $j_c$ .
- Films with **combined type of pinning** with equal partial contributions of  $\delta-T_c$  and  $\delta-I$  ( **$P_1 = P_2 = 0.5$** ) demonstrate the  $J_c(10\text{K}, 0\text{ T}) = 1.8 \times 10^7\text{ A/cm}^2$  and  $j_c(10\text{K}, 5\text{ T}) = 2 \times 10^6\text{ A/cm}^2$  at 10 K.



# Thank you for your kind attention!





# Estimation of MgB<sub>2</sub> density

**Table 2.** Characteristics (critical current density- $j_c$ , maximal pinning force- $F_{p(max)}$ , amount of MgB<sub>2</sub>, MgO and MgB<sub>4</sub> estimated from Rietveld refinements, and density- $\rho$ ) of MgB<sub>2</sub>-based materials prepared under different  $P$ - $T$  conditions from Mg and B taken in MgB<sub>2</sub> stoichiometry or from MgB<sub>2</sub> powder using different methods: under high quasihydro-static pressure –HP, by hot pressing HotP, by spark plasma sintering – SPS, by pressureless sintering under 0.1 MPa of Ar – PL.

No	Preparation	Type of B	$P$ , MPa	$T$ , °C	$j_c$ , MA/cm <sup>2</sup> , at 0-1 T, at 20 K	$F_{p(max)}/10^9(N/m^3)$ at 20 K	MgB <sub>2</sub> /MgO/MgB <sub>4</sub> , wt%	Density, $\rho$
1.	in-situ, HP	I	2000	1050	0.9 – 0.7	7.6	94/6/0	99%
2.	in-situ, HP	I	2000	800	0.2 – 0.15	1.6	91/5.5/0	98%
3.	in-situ, HP	III	2000	1050	0.4 – 0.3	2.3	87/13/0	99%
4.	in-situ, HP	III	2000	1000	0.36-0.23	2.4	86/14/0	99%
5.	in-situ, HP	III	2000	950	0.3 – 0.2	2.1	-	98%
6.	in-situ, HP	III	2000	900	0.12 – 0.06	1.1	71/13/0	98%
7.	in-situ, HP	III	2000	800	0.12 – 0.07	0.8	73/12/0	97%
8.	in-situ, HP	II	2000	600	0.14 – 0.05	0.6	64/30/0	83%
9.	in-situ, SPS	III	50	1050	0.5 – 0.4	4.6	83/4.5/12.5	94%
10.	in-situ, SPS	III	50	800	0.4 – 0.36	2.7	-	74%
11.	ex-situ, SPS	-	50	1050	0.4 – 0.3	3.3	83/6.5/10.5	96%
12.	in-situ, HotP	III	30	1050	0.08 – 0.016	0.2	46/8.5/45.5	99%
13.	in-situ, HotP	III	30	800	0.3-0.2	0.6	-	72%
14.	in-situ, PL	IV	0.1	800	0.08-0.03	1.9	90/10/0	55%
15.	in-situ, cold dens. 2 GPa, PL	III	0.1	600	0.26-0.13	1.3	94.5/5.5/0	65%

# Estimation of grain sizes in MgB<sub>2</sub>

The average crystallite sizes were calculated from line broadening in X-ray diffraction pattern by the standard program in accordance with the following:

$$\text{Crystallite size} = \frac{K \cdot \lambda}{W_{\text{size}} \cdot \cos \theta} \quad \text{with} \quad W_{\text{size}} = W_b - W_s$$

where  $W_{\text{size}}$  is the broadening caused by small crystallites;  $W_b$  is the broadened profile width;  $W_s$  is the standard profile width (0.08 °);  $K$  is the form factor;  $\lambda$  is the wavelength.

Table. Critical current density,  $j_c$ , vs. relative average grain size of crystallites of high-pressure sintered from MgB<sub>2</sub> and synthesized from Mg and B taken in 1:2 ratio materials

High pressed at 2 GPa for 1 h at T <sub>s</sub> , °C	average crystal sizes	lattice parameters		$j_c$ , kA/cm <sup>2</sup> at 10 K		$j_c$ , kA/cm <sup>2</sup> at 20 K	
		$a$ (nm)	$c$ (nm)	0 T	1 T	0 T	1 T
From MgB <sub>2</sub> (10 μm and 0.8 % of O )							
700	19.7 nm	0.30805	0.35188	-	-	-	-
800	18.8 nm	0.30822	0.35212	-	-	-	-
900	18.5 nm	0.30820	0.35208	56	14	36	8
1000	25.0 nm	0.30797	0.35200	28	8	19	5
From Mg chips and B (4 μm, 1.5 % O) mixed and milled in 1:2 ratio							
800	15.0 nm	0.30747	0.35188	245	142	138	79
900	21.0 nm	0.30819	0.35174	205	136	128	61
1000	37.0 nm	0.30808	0.35192	485	364	360	237



# Estimation of connectivity, $A_F$

$$A_F = \frac{9 \mu\Omega \cdot \text{cm}}{\rho_{300} - \rho_{40}}$$

$$9 \mu\Omega \cdot \text{cm}$$

perfect connectivity for  $\text{MgB}_2$

resistivity

$$\rho = \frac{V \cdot \sigma}{I \cdot \Delta l}$$

$$\rho_{300} = \frac{V_{300} \cdot \sigma}{I \cdot \Delta l} \quad \rho_{40} = \frac{V_{40} \cdot \sigma}{I \cdot \Delta l}$$

$\Delta l$  - distance between voltage contacts

$\sigma = a \cdot b$  - cross-section area of the sample,  $a \cdot b \cdot c$  - sample dimensions

$V_{300}$  - arithmetic average of all  $V$  data at room temperature

$V_{40}$  - voltage at 40 K

$I = 100 \text{ mA}$  - current

$\rho$  - by four probe method

# Estimation of pinning mechanism

M. Eisterer “ Calculation of the volume pinning force in  $\text{MgB}_2$  superconductors” (PHYSICAL REVIEW B 77, 144524 2008)

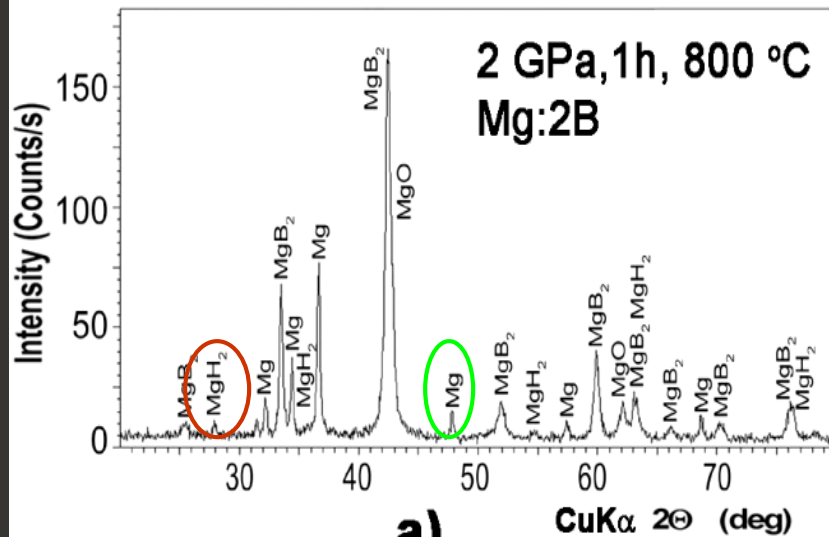
**Identification of the dominant pinning mechanism from the peak position of pinning volume force:**

$$k = B_{\text{peak}} / B_n ;$$

$$B_{\text{peak}} = F_{p(\text{max})}, \quad B_n = F_{p(\text{max})} / 2$$

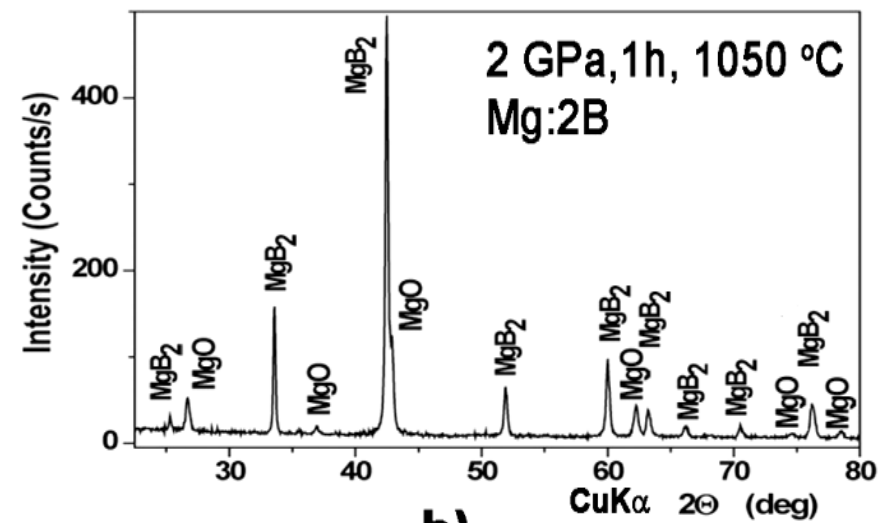
The dominant pinning mechanism was determined from the volume pinning force  $j_c B$  according to a scaling procedure proposed by M.Eisterer , which eliminates the influence of the intrinsic anisotropy of  $\text{MgB}_2$  on the field dependence of the volume pinning force. The field  $B_{\text{peak}}$ , where the maximum of the volume pinning force occurs, was normalized by the field  $B_n$ , at which the volume pinning force drops to half its maximum (on the high field side), instead of the upper critical field,  $B_{c2}$ , in the original approach for isotropic superconductors . For typical values for the anisotropy (4) and the percolation threshold (0.25) in bulk  $\text{MgB}_2$  samples, the position of the peak,  $B_{\text{peak}}/B_n$ , is expected to be at 0.34 and 0.47 for grain boundary and point pinning, respectively. Changes in the anisotropy and the percolation threshold only slightly change the peak position, in particular in the realistic range of these parameters. Another advantage of this scaling approach is that  $B_n$  is easily available and no separate experiment for determining  $B_{c2}$  or any extrapolation to estimate the irreversibility (or Kramer) field is necessary to define the scaling field.

**X-ray diffraction patterns of the materials prepared from Mg:2B under 2 GPa for 1 h without additions at 800 °C (a) and 1050 °C (b) and with 10 wt. % Ti at 800 °C (c) and 1050 °C (d).**

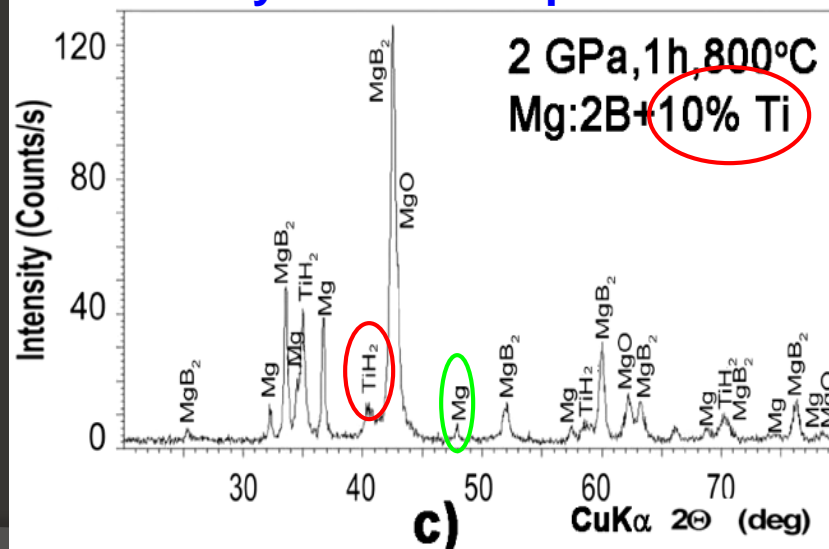


**a)**

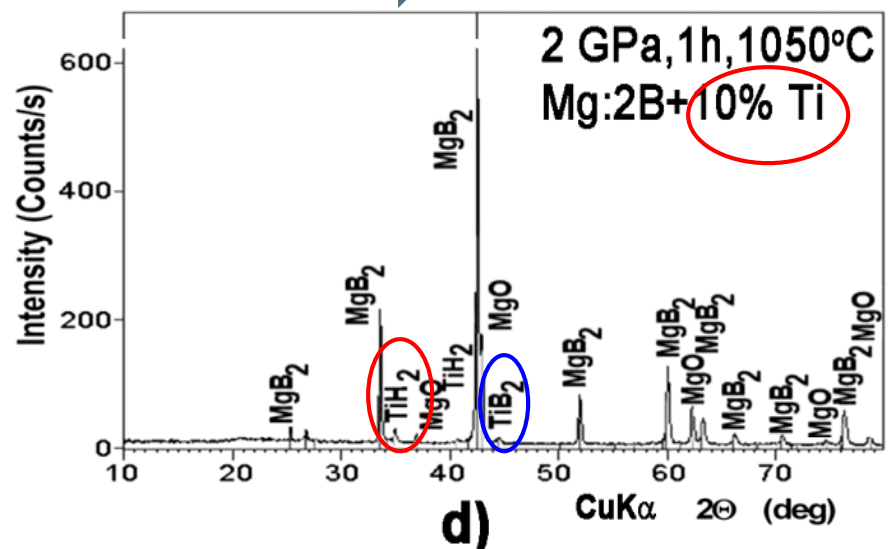
**Synthesis temperature increase**



**b)**



**c)**



**d)**

# Estimation of connectivity, $A_F$ , amount of shielding fraction, $S$ , and type of pinning, $k$ .

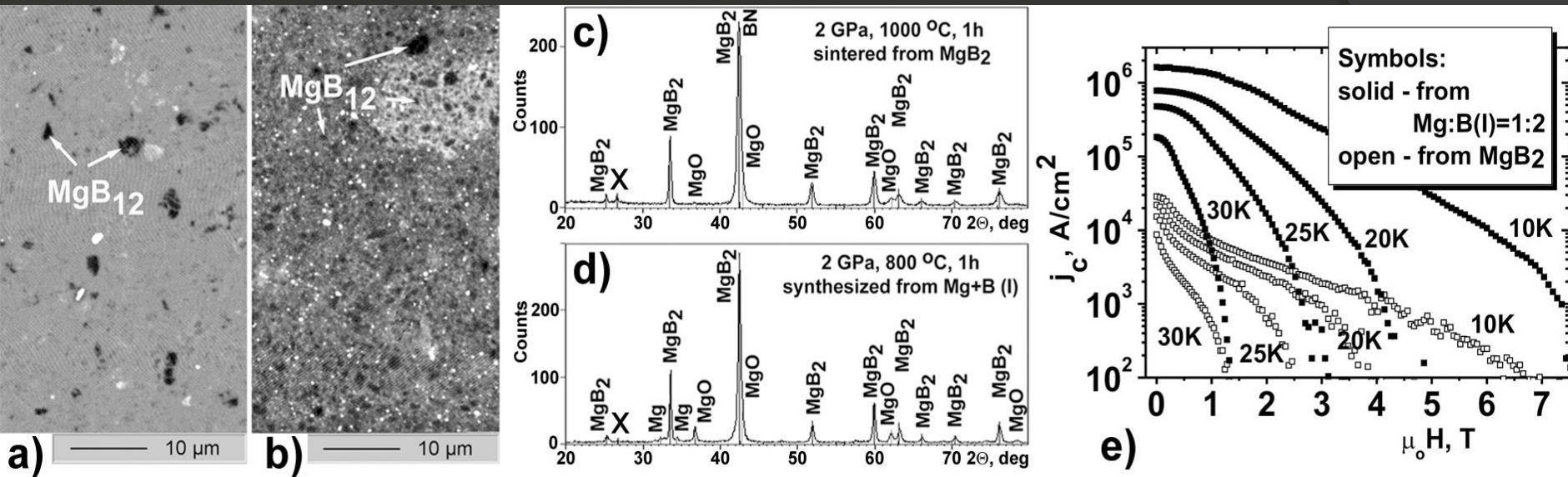
Connectivity,  $A_F$ ,  $k$ -ratio:  $k = B_{\text{peak}}/B_n$  and amount of shielding fraction,  $S$ , of  $\text{MgB}_2$  materials prepared under different conditions. The numbering in tables 1 and 2 is the same.

No	Preparation	Type of B	$P$ , MPa	$T$ , °C	$k = B_{\text{peak}}/B_n$	$A_F$ , %	$S$ , %
1.	in-situ, HP	I	2000	1050	0.43	-	-
2.	in-situ, HP	I	2000	800	0.37	-	-
3.	in-situ, HP	III	2000	1050	-	79	94
5.	in-situ, HP	III	2000	950	-	59	81
7.	in-situ, HP	III	2000	800	-	57	91-100
8.	in-situ, HP	II	2000	600	0.32	18	90
9.	in-situ, SPS	III	50	1050	-	98	91
11.	ex-situ, SPS	-	30	1050	-	80	100
12.	in-situ, HotP	III	30	1050	-	32	-
13.	in-situ, HotP	III	30	800	-	73	-
15.	in-situ, cold dens. 2 GPa, PL	III	0.1	600	-	-	75
16.	in-situ, cold dens. 2 GPa, PL	III	0.1	800	-	-	84

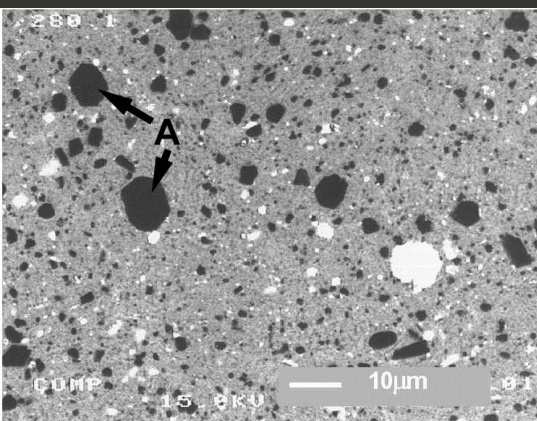
**With synthesis temperature increase the shift from grain boundary pinning towards point pinning occurs (about what  $k$ -ration increase witnessed)**



# CONTRADICTIONS BETWEEN X-RAY DATA AND SEM OBSERVATIONS

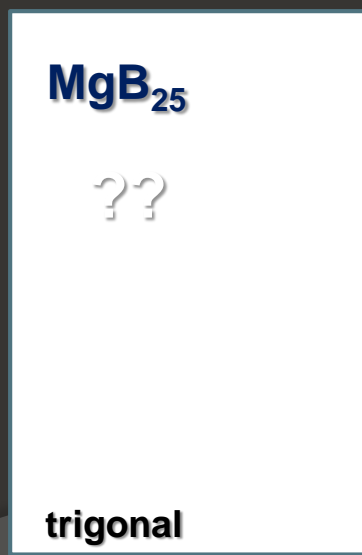
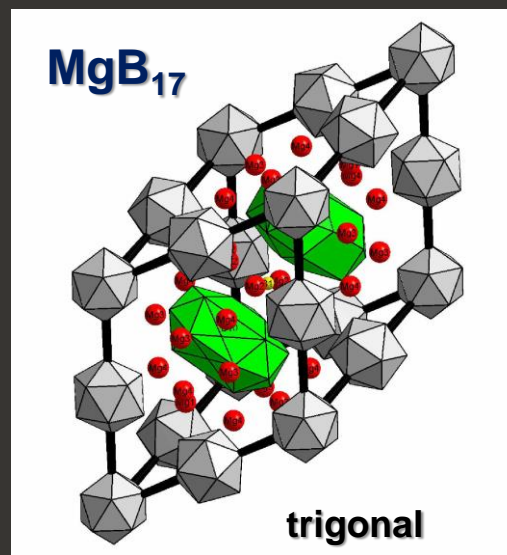
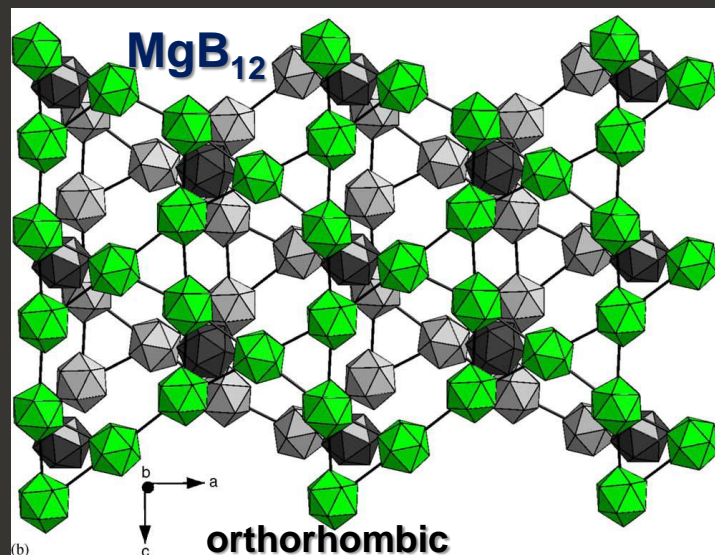
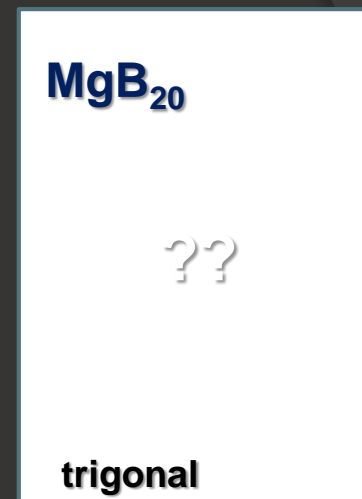
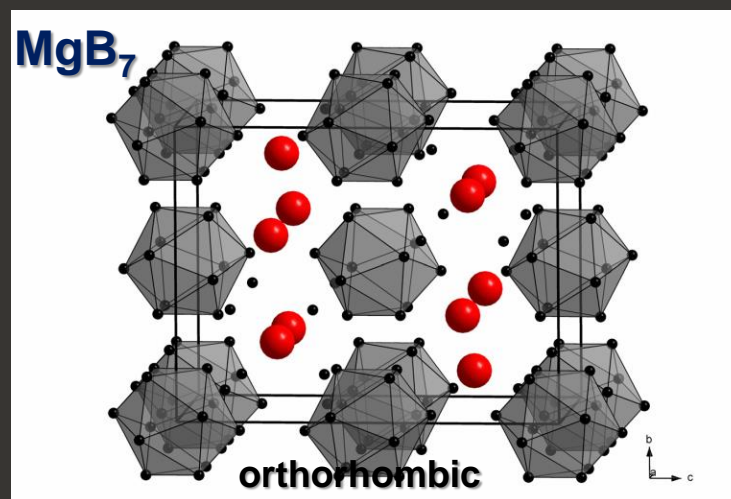
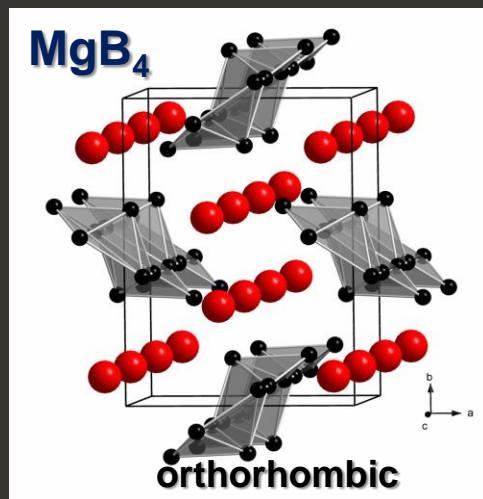


- (a), (b) structures of the samples obtained by SEM in COMPOsitional contrast:  
 (a) sintered from  $\text{MgB}_2$  at 2 GPa, 1000 °C, 1 h;  
 (b) synthesized from Mg and B taken into 1:2 ratio at 2 GPa, 800 °C, 1 h;  
 (c), (d) –X-ray patterns of the samples shown in Figs. 5a, b;  
 (e) dependences of critical current density,  $j_c$ , on magnetic fields,  $\mu_0 H$ , at different temperatures of the samples shown in Fig. e: open symbols – sintered from  $\text{MgB}_2$  material and solid symbols - synthesized from Mg and B taken in 1:2 ratio



**Structure of  $\text{MgB}_2$ -based material synthesized from amorphous boron and magnesium chips taken in  $\text{Mg}:\text{B}=1:2$  ratio in contact with Ta foil at 2GPa, 800°C, 1h (BEI SEM image) shows that inclusions of black phase with near  $\text{MgB}_{12}$  stoichiometry may crystallize in hexagonal habit or have near hexagonal cross-section (see, for example, inclusion marked "A")**

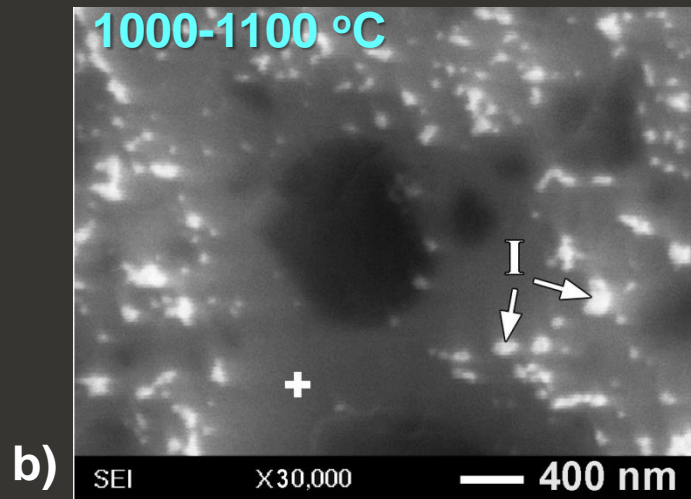
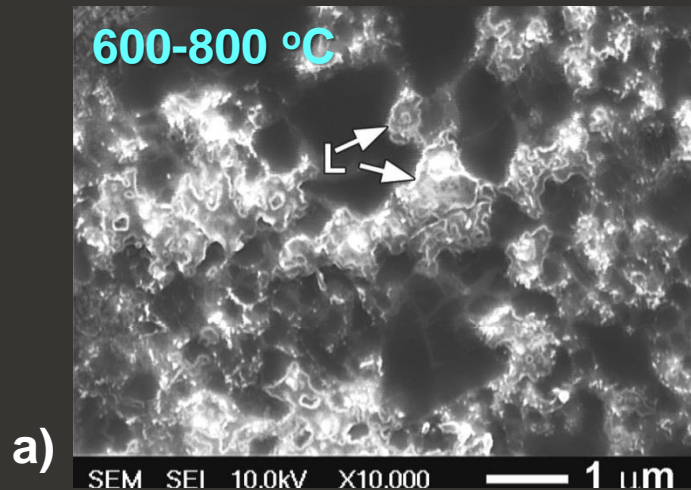
Each of the following non-superconducting higher magnesium borides in the Mg-B system can crystallize in the  $\text{MgB}_2$  matrix and can influence pinning



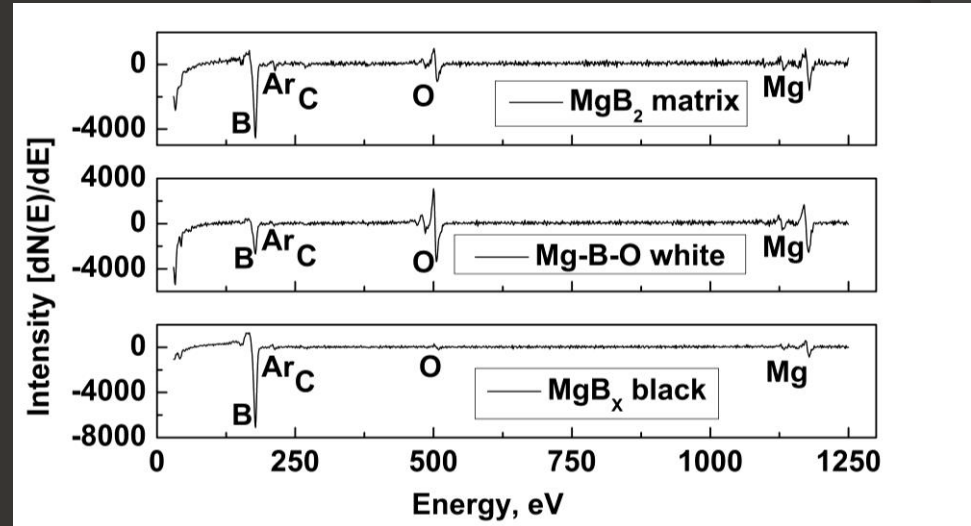
The method of preparation (pressure-temperature-time conditions) can affect the stoichiometry of the  $\text{MgB}_x$  inclusions in  $\text{MgB}_2$ . (2 GPa –  $\text{MgB}_{12}$ )

# Typical structures of $\text{MgB}_2$ prepared at 0.1 MPa – 2 GPa at low temperatures (a) - 600-800 °C and high (b) -1000-1100 °C and their compositions

0.1 MPa – 2 GPa



“L” -  $\text{MgB}_{0.9-3.5}\text{O}_{1.6-2}$  white layers, thickness (15-20 nm)



Auger spectrum

“P” -  $\text{MgB}_{0.5-0.8}\text{O}_{0.8-0.9}$  white inclusions

$\text{MgB}_{11-13}\text{O}_{0.2-0.5}$  black inclusions

$\text{MgB}_{2.2-1.7}\text{O}_{0.3-0.6}$  gray matrix.

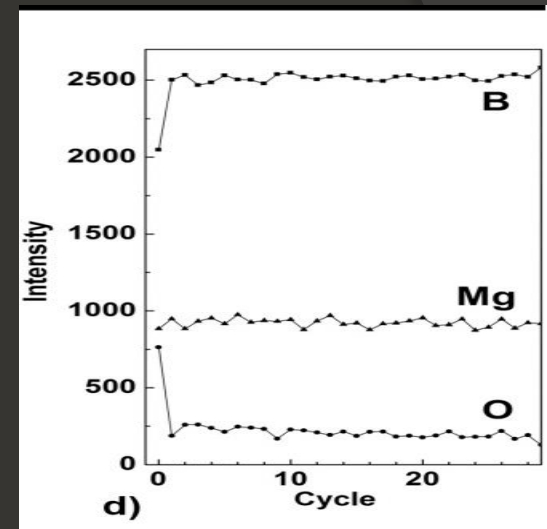
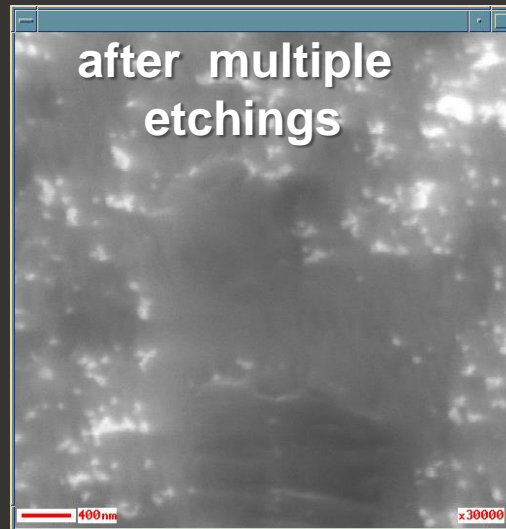
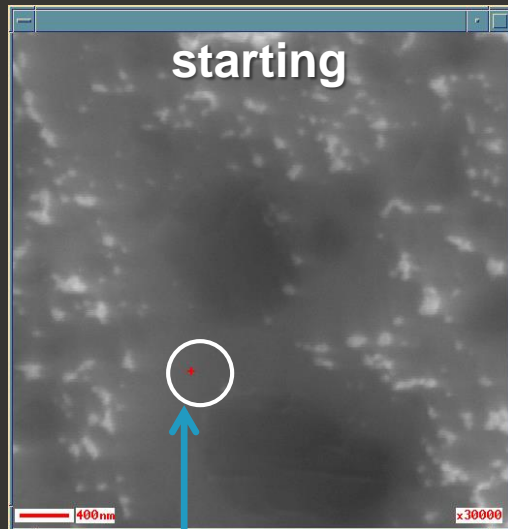
Auger study by JAMP-9500F

Zone of excitation being 10 nm in diameter and two lattice parameters in depth etching in Ar allows to exclude analyzing of oxidized layer on the surface!



# Auger study by JAMP-9500F

zone of excitation being 10 nm in diameter and two lattice parameters in depth  
Mg(I):2B(III), 1200 °C, 2 GPa, 1h



## Place of analysis

$\text{MgB}_{0.5-0.8}\text{O}_{0.8-0.9}$  white inclusions  
 $\text{MgB}_{11-13}\text{O}_{0.2-0.5}$  black inclusions  
 $\text{MgB}_{2.2-1.7}\text{O}_{0.3-0.6}$  gray matrix.

Spectrums  
of gray matrix

Depth profile  
of gray matrix

Comparing these results with that of X-ray structure analysis we can conclude that **boron incorporated into the  $\text{MgB}_2$  structure.**

Repeated etching in Ar (in the JAMP-9500F chamber during study) and quantitative Auger analysis showed practically the same oxygen concentration in  $\text{MgB}_2$  matrix of the structure after 30 cycles, which points to the **oxygen incorporation into the  $\text{MgB}_2$  lattice.**

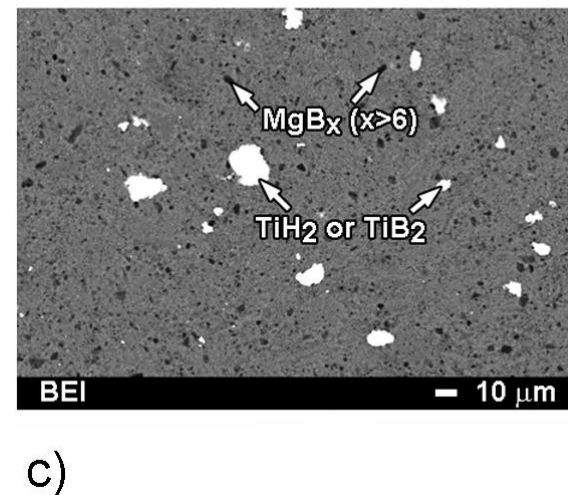
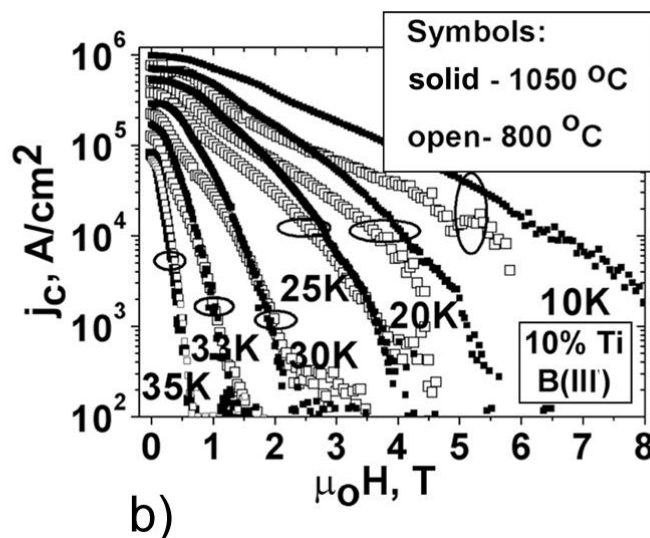
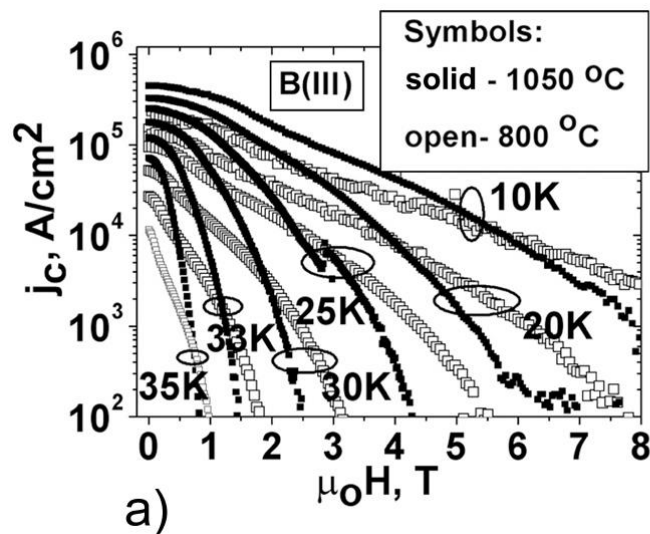


# Increase of $j_c$ in low and medium magnetic fields with increase of manufacturing temperature and effect of Ti addition

2 GPa, 1 h

from Mg:2B

from Mg:2B + 10% Ti



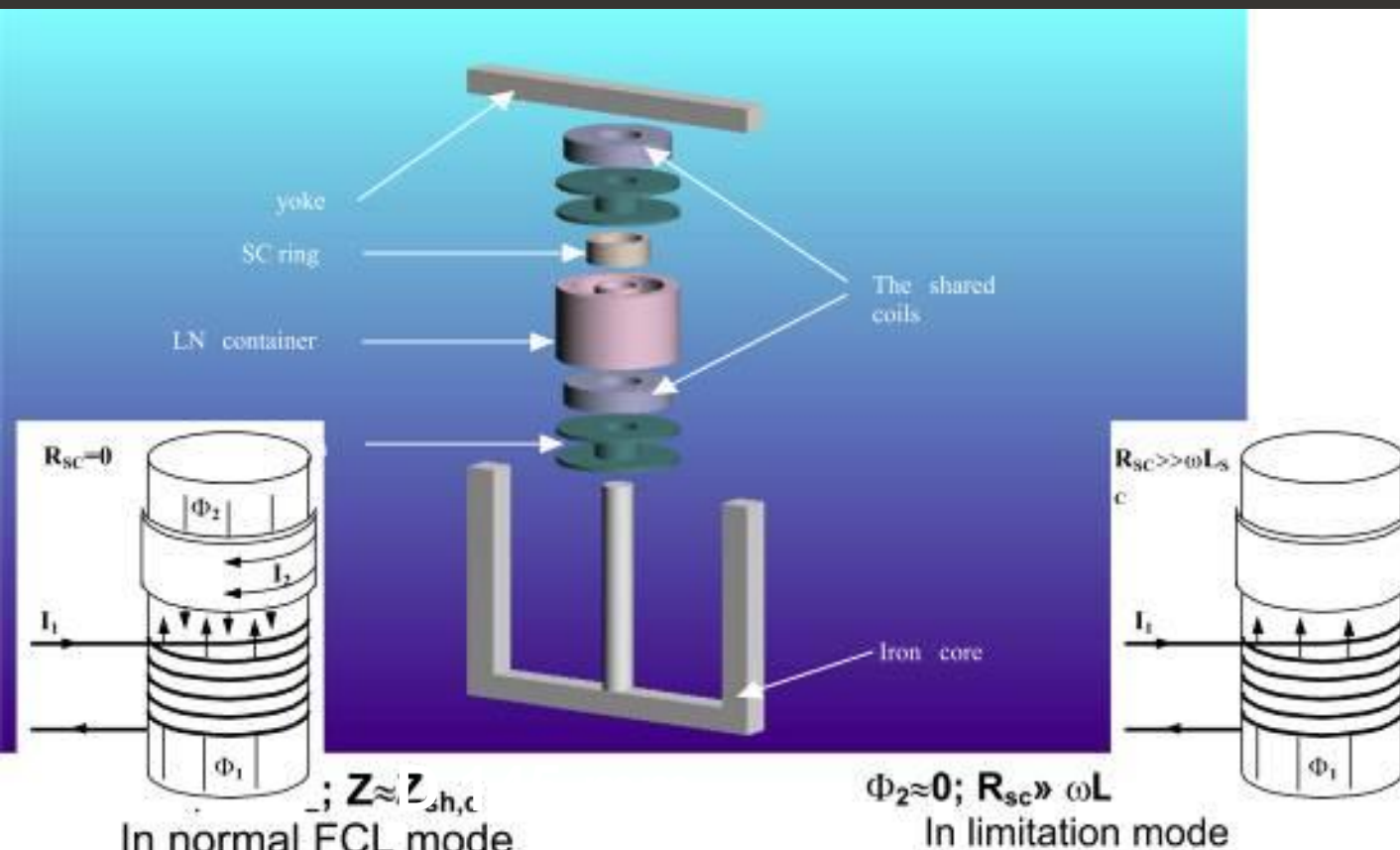
Boron amorphous: grain size 4  $\mu\text{m}$ , 1.5 wt% O, 0.47 wt% C, 0.40 wt% N, 0.37 wt% H

➤ The critical current density of  $\text{MgB}_2$  – based materials usually increases at low and medium magnetic fields with increasing manufacturing temperature, but lower manufacturing temperatures can result in a  $j_c$  increase at high magnetic fields.

➤ A shift from grain boundary pinning to point pinning was observed

➤ A positive effect in enhancing the critical currents of  $\text{MgB}_2$  is played, by dopants, such as Ti, for example.

# MgB<sub>2</sub> for fault current limiter



Rings cut from MgB<sub>2</sub> blocks

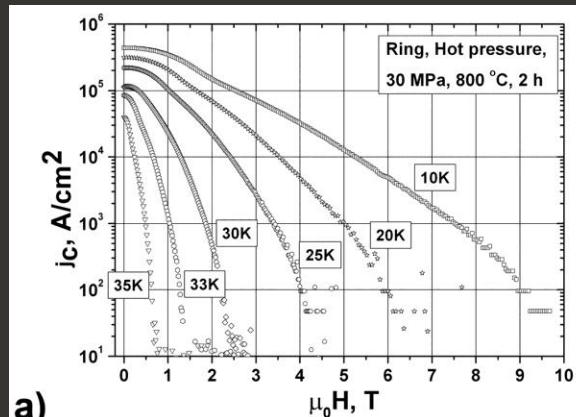


**Transformer- type fault current limiter (Budapest University of Technology and Economics)**

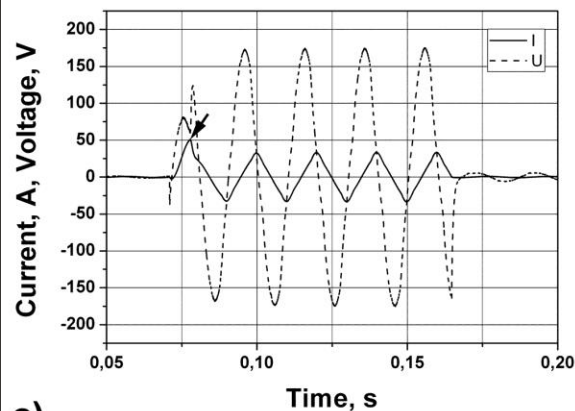
At fault event an increase in the circuit current causes a growth of the current in a superconductor higher than the critical value. The SC-or passes into the resistive state introducing a required for limitation impedance in the circuit ( 2–4 ms and 1 ms ).

# Smart application of HP and hot-pressed $\text{MgB}_2$ -based materials

Characteristics of the  $\text{MgB}_2$ -based material synthesized at 30 MPa, 800 °C, 2 h from amorphous boron (III) (H.C. Starck, 1.5 % O, 4  $\mu\text{m}$  grains) and magnesium (I) taken in the  $\text{MgB}_2$  stoichiometry and of the ring cut from this material:

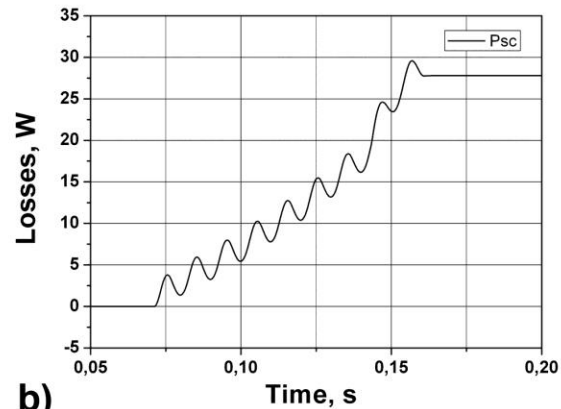


a)

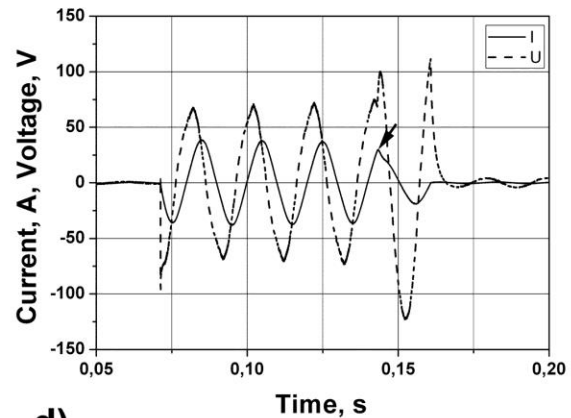


c)

voltage of source 170 V



b)



d)

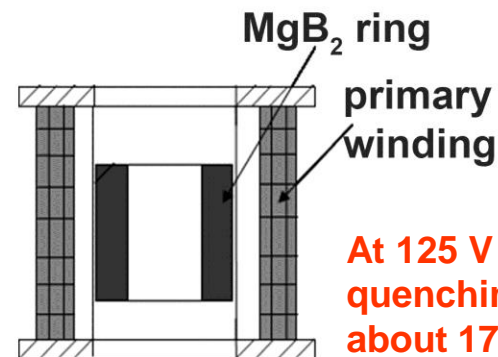
voltage of source 125 V



e)

Quenching current 24000 A.

$j_c = 63200 \text{ A/cm}^2$ .



f)

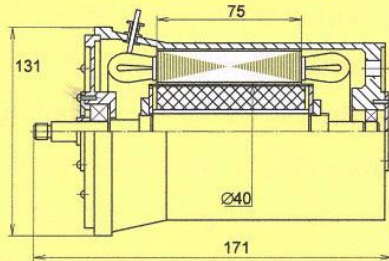
At 125 V losses till quenching are about 17 J. Power of the losses is about 200 W.

Typical oscilloscope traces of the current (solid lines) and voltage drop across the primary normal-metal winding (dashed lines) at 4.2 K. Arrows show the quenching current. Voltage of source was 170 V (c) and 125 V (d); e) ring: outer diameter – 45 mm, height - 11.6 mm, wall thickness – 3.3 mm.



## Experimental HTS Motor for 20 K Temperature

Cross-section of HTS Motor

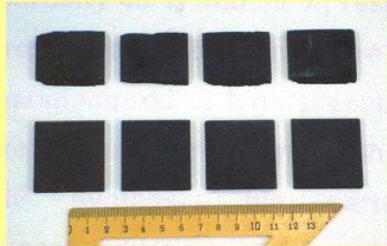


General View of HTS Motor



## Compound HTS-Ferromagnetic Rotor With YBCO Bulk Elements

General View of YBCO Samples

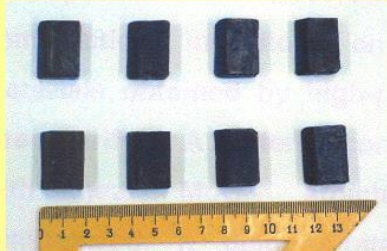


Rotor With YBCO Bulks



## Compound HTS-Ferromagnetic Rotor With MgB<sub>2</sub> Bulk Elements

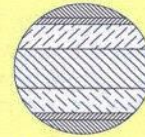
General View of MgB<sub>2</sub> Bulk Samples



Rotor With MgB<sub>2</sub> Samples



## Test results of reluctance motor with MgB<sub>2</sub> blocks



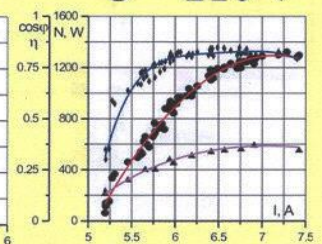
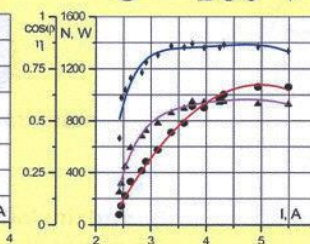
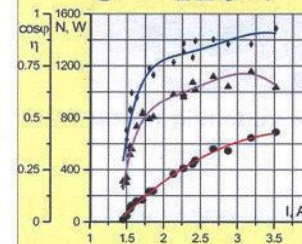
Temperature:  $T = 15 \dots 20 \text{ K}$

Phase voltage: Phase voltage: Phase voltage:

$U = 120 \text{ V}$

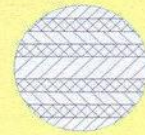
$U = 160 \text{ V}$

$U = 210 \text{ V}$



◆ - efficiency, ● - output power, ▲ - power factor

## Test results of reluctance motor with YBCO blocks



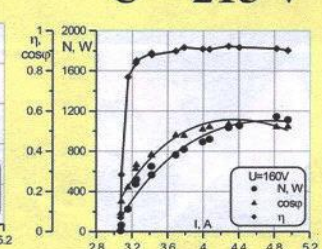
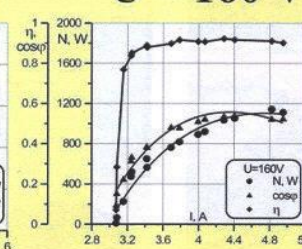
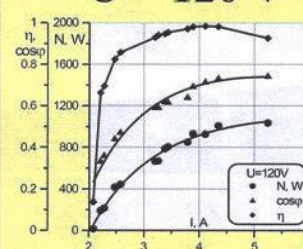
Temperature:  $T = 15 \dots 20 \text{ K}$

Phase voltage: Phase voltage: Phase voltage:

$U = 120 \text{ V}$

$U = 160 \text{ V}$

$U = 215 \text{ V}$



◆ - efficiency, ● - output power, ▲ - power factor

# Conclusions

- The distribution of nanostructural inhomogeneities, such as areas with a high concentration of B and impurity O, plays a key role in the variation of  $j_c$  in  $\text{MgB}_2$  materials.
- Auger and SEM studies show that with increasing manufacturing temperature oxygen enriched 15-20 nm thick nano-layers transform into distinct dispersed Mg-B-O inclusions.
- In parallel, the Ti addition results in a further increase in  $j_c$ , because Ti fosters the localization (or segregation) of oxygen and of higher magnesium borides and facilitates the formation of a homogeneous  $\text{MgB}_2$  matrix with low oxygen content, but with Mg-B-O and  $\text{MgB}_x$  pinning centers.
- At low synthesis temperature Ti can absorb hydrogen forming titanium hydrides, thus preventing the formation of  $\text{MgH}_2$  and promoting the material densification.
- The positive effect of Ti addition results from the high ability of Ti to absorb H, O, and Mg.

# Experimental

2 GPa - HP, 30 MPa- HotP, 50 MPa – SPS, 0.1 MPa Ar - PL

## Initial materials

- (I) *in-situ* from Mg:B=1:2 ( $\text{MgB}_2$ ), from Mg:B=1:7 and 1:12 mixing and milling in a high speed activator for 3-5 min
- ✓ Boron (B) amorphous

### *Characteristics of initial boron powders.*

Type of B	Grain size	O, wt%	C, wt%	N, wt%	H, wt%
I	<5 $\mu\text{m}$	0.66	0.31	0.48	0.32
II	<1 $\mu\text{m}$	-	3.5	1.02	0.87
III	4 $\mu\text{m}$	1.5	0.47	0.40	0.37

- ✓ Magnesium (Mg)
  - Mg(I) turning (Technical Specifications of Ukraine 48-10-93-88)
  - Mg(II) powder 325 mesh (HyperTec, USA)

(II) *ex-situ* from  $\text{MgB}_2$  powder

- ✓ Magnesium diboride (98% purity)