



Heavy Higgses

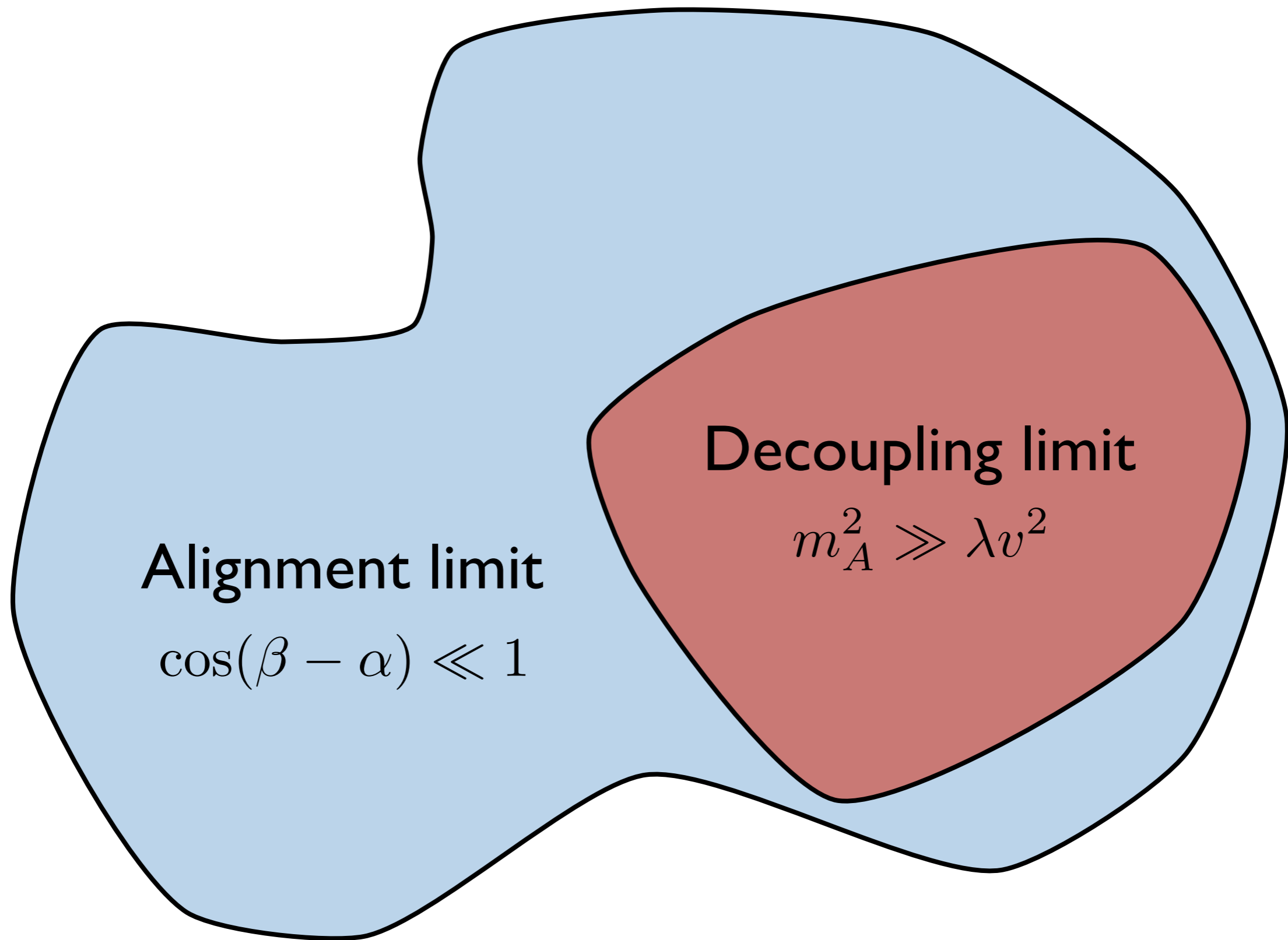
Nathaniel Craig
UCSB

FCC Higgs/EWSB WG

Including work with E.Brownson, N.Dhingra, U.Heintz,
G.Kukartsev, M.Narain, N.Parashar, J.Stupak

2HDM

- For the purposes of this talk, I'll primarily focus on extended EWSB sectors whose IR physics is described by two Higgs doublets.
- This covers a broad class of known models and allows for convenient parameterization.
- ...but many of the qualitative features are shared by other extended EWSB sectors.
- I'll keep the focus on bottom-up phenomena, generalizing beyond SUSY 2HDM.



Why heavy(ish) Higgses?

Generically, the mass scale of heavy Higgses is only constrained by the distribution of the EWSB vev, which can naturally be (reasonably) asymmetric.

E.g. SUSY:

$$\Delta \approx \sin^2(2\beta) \frac{m_H^2}{m_h^2} + \mathcal{O}(m_H^0)$$

In the limit of large $\tan\beta$ implies suppressed sensitivity;

$$\Delta(\tan\beta = 50) \leq 1 \rightarrow m_H \lesssim 3.1 \text{ TeV}$$

So multi-TeV additional Higgs states are very consistent with naturalness in this framework. Not a scale we're likely to reach at 14 TeV, but certainly could be within reach of 100 TeV.

100 TeV

Greater kinematic reach

Decoupling

Higher cross sections
[$\sigma_{ggF} 15\times\sigma(14\text{ TeV}) @ 125\text{ GeV}$]

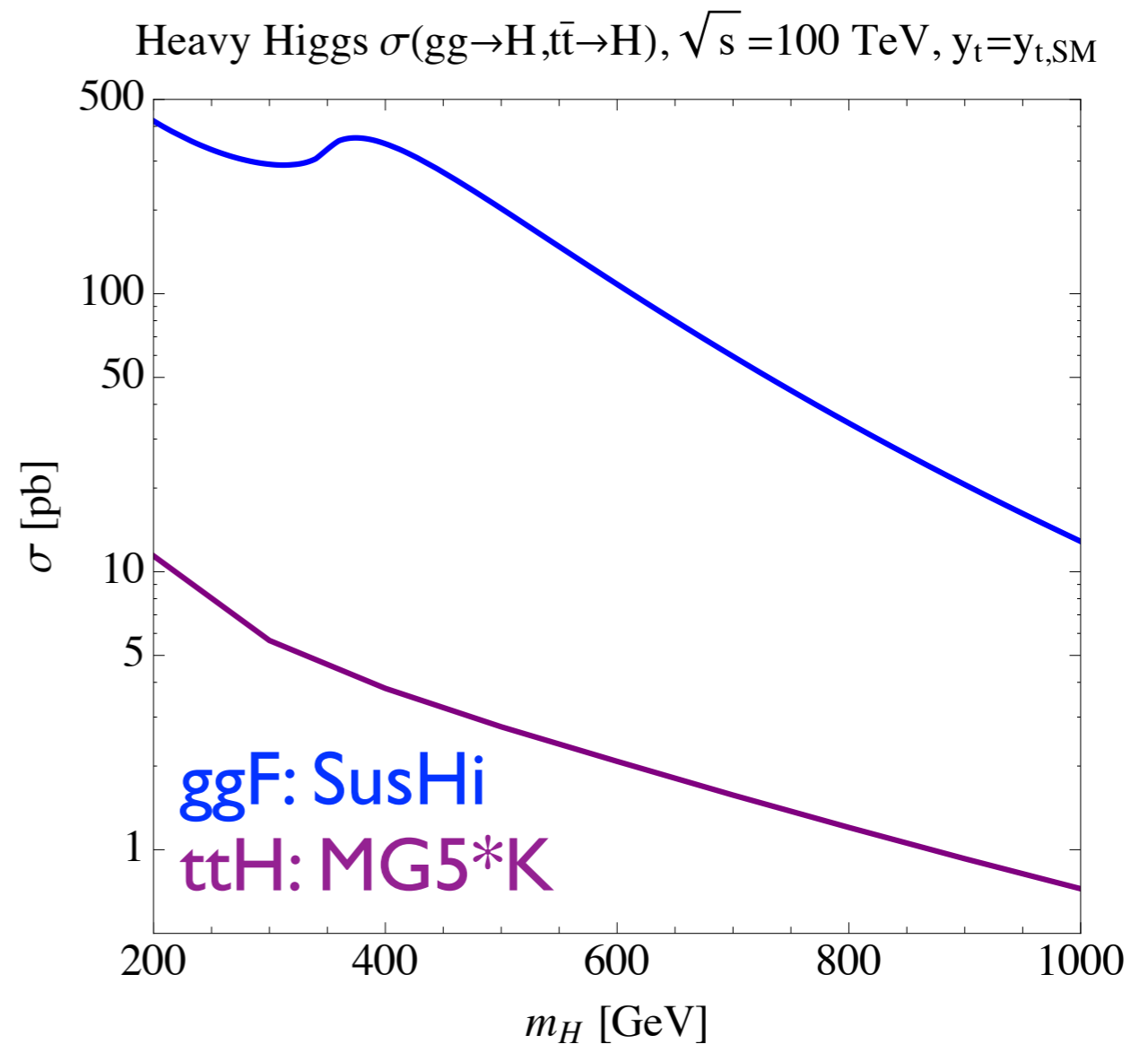
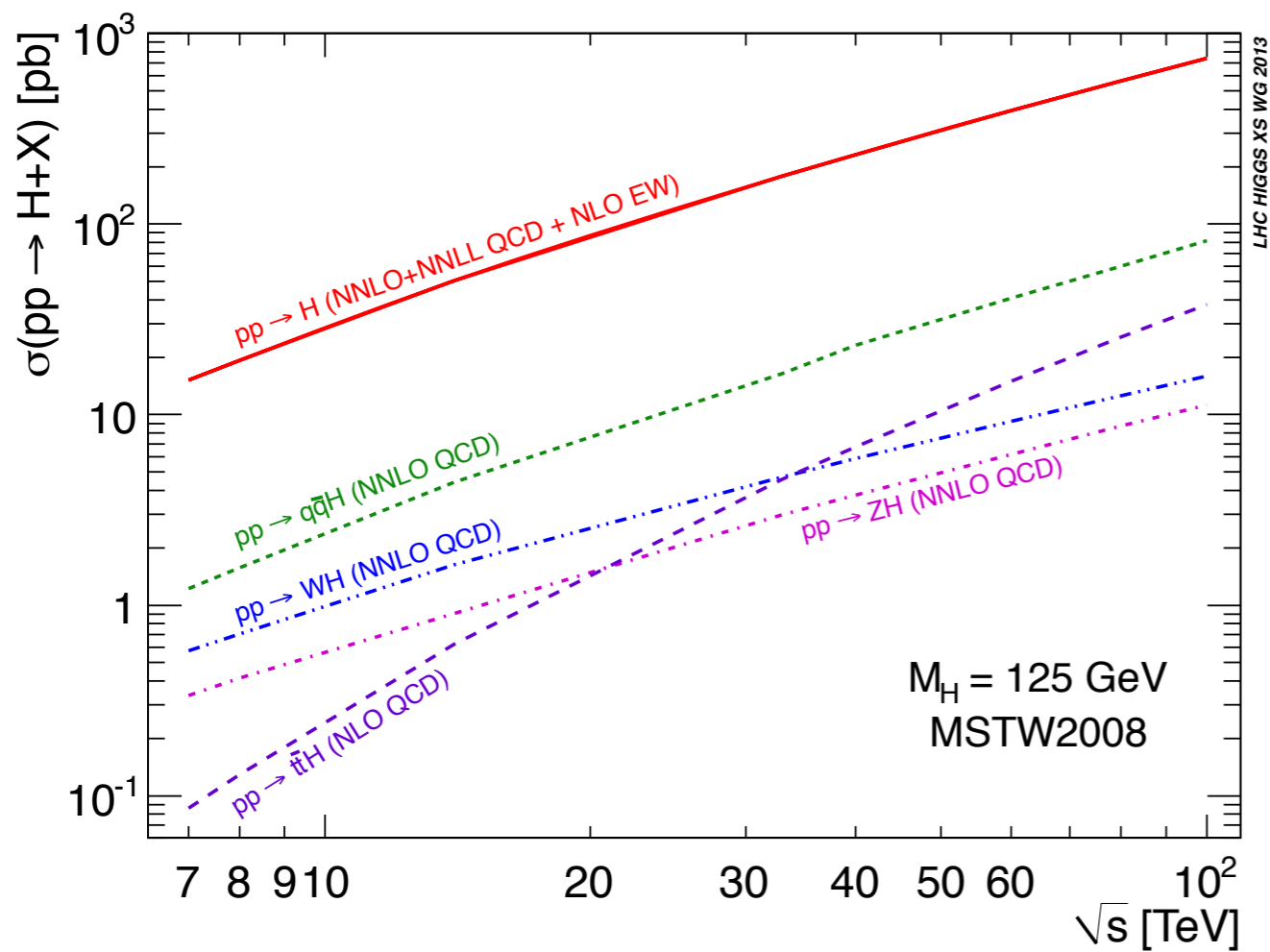
Exploit this to search for high-mass Higgses.

Alignment

Exploit this to search more efficiently for parametrically lighter states with small SM production/distinctive decay rates.

100 TeV Opportunities

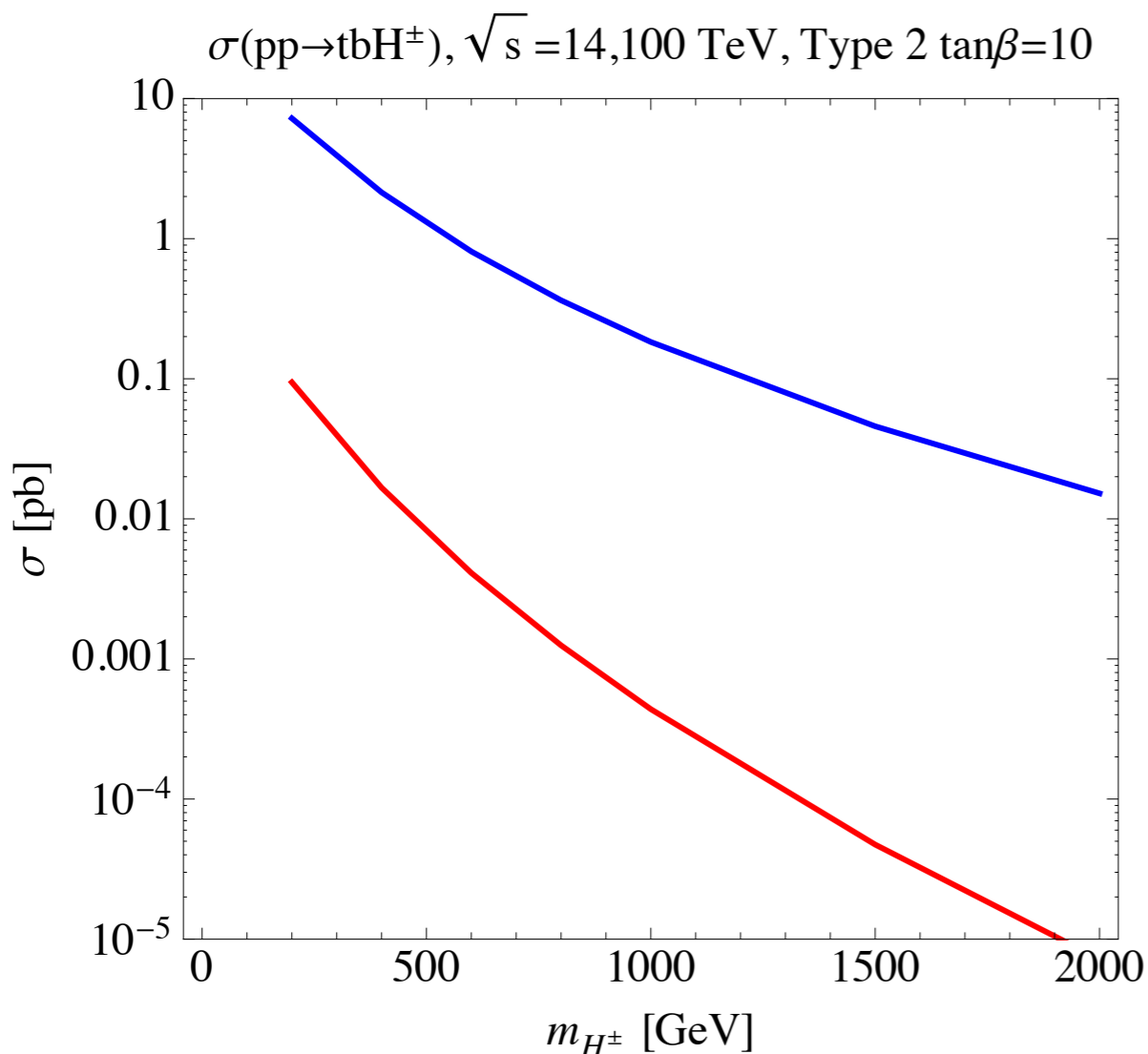
(I) ttH grows due to gluon pdfs & improved kinematics



$ttH: 61 \times \sigma(14 \text{ TeV})$; $ggF: 15 \times \sigma(14 \text{ TeV})$

100 TeV Opportunities

(2) Charged Higgs production becomes appreciable.
Also improved by kinematics, pdfs at 100 TeV



Preliminary estimate (LO MG5):

tbH^\pm : $80 \times \sigma(14 \text{ TeV})$

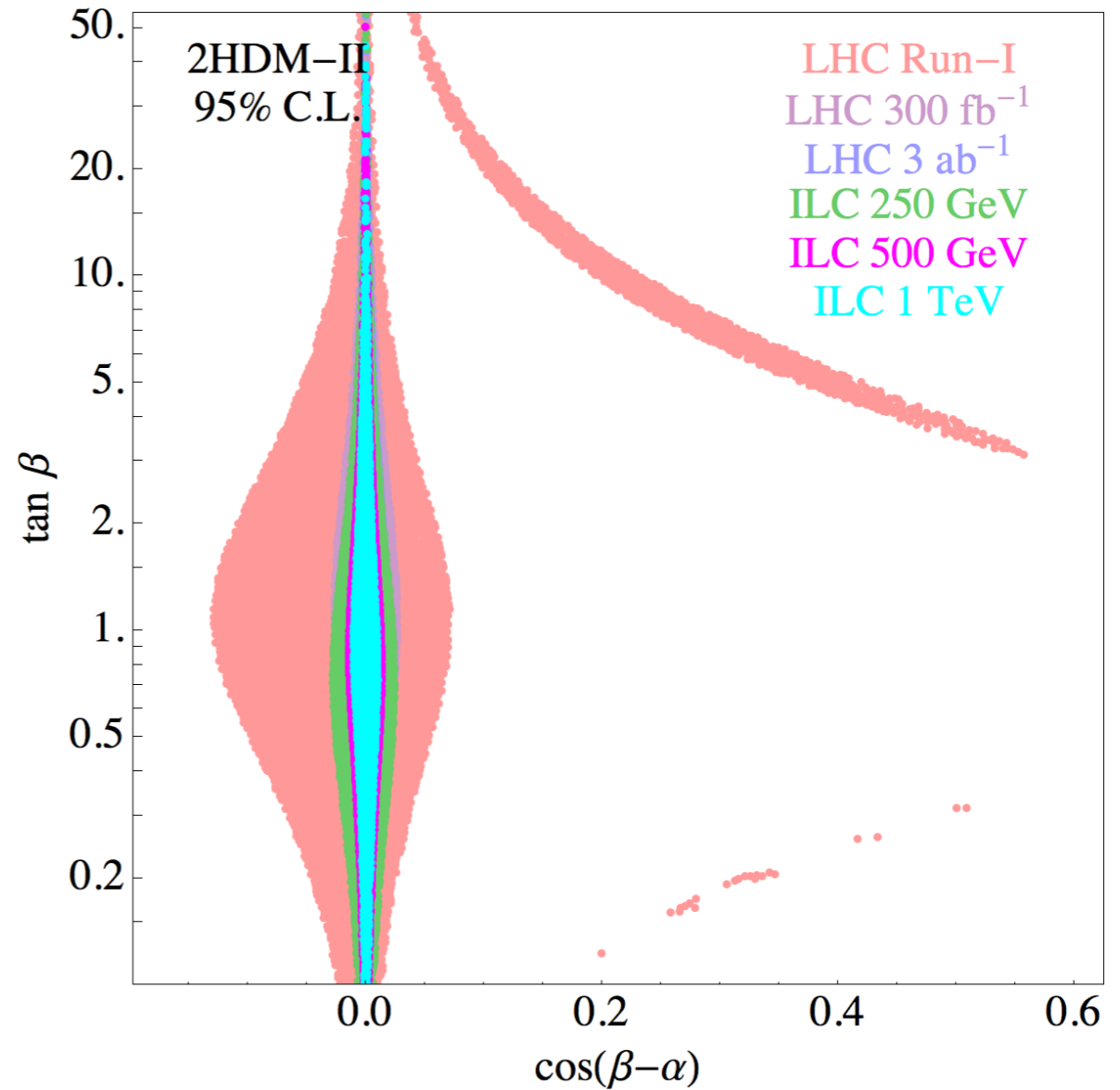
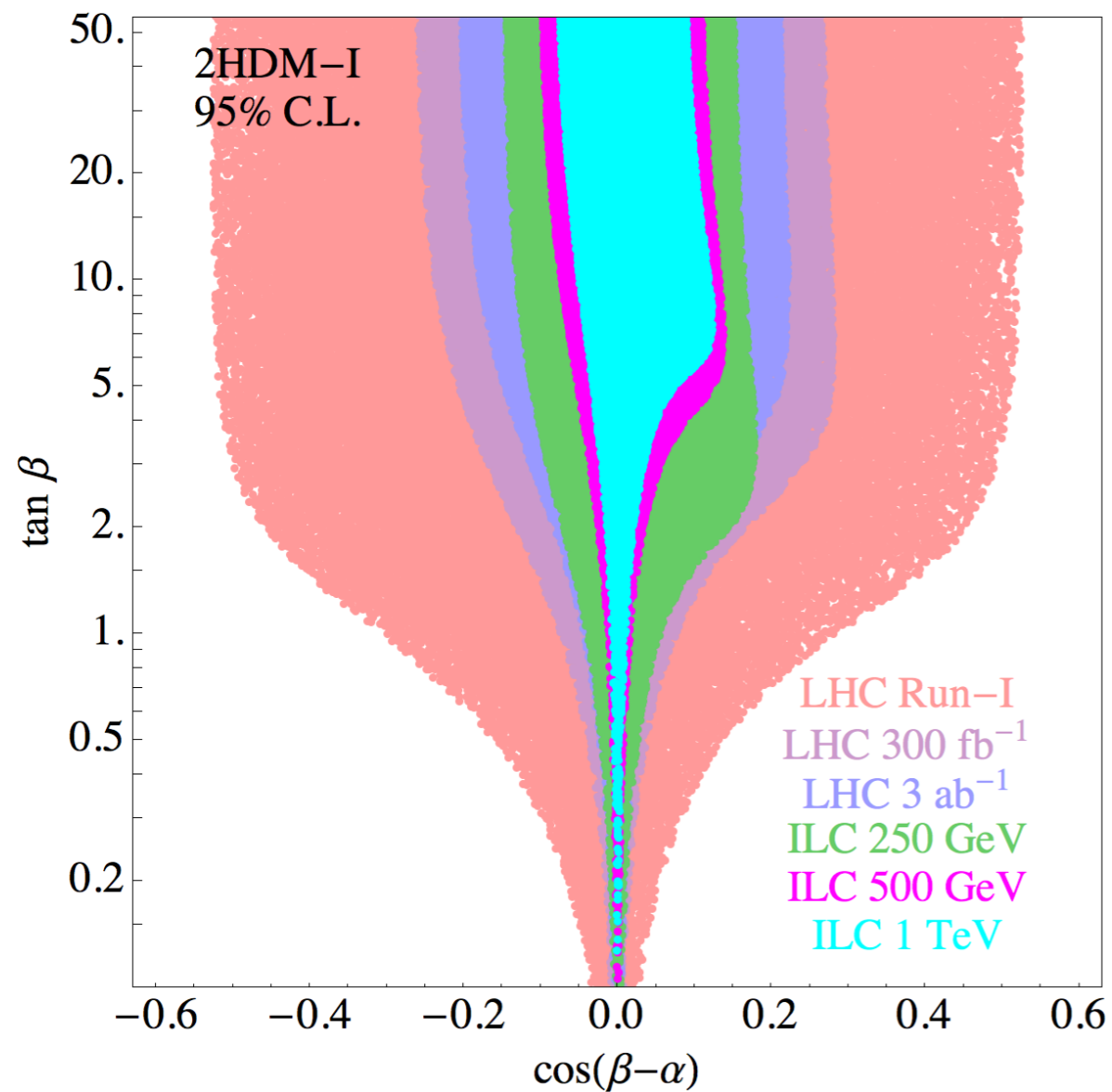
@ $\tan\beta = 10$ $m_{H^\pm} = 200$ GeV.

$\sigma > 100 \text{ fb}$ to $m_{H^\pm} \sim \text{TeV}$

This is the “worst-case” value
of $\tan\beta$; rate increases at
higher and lower $\tan\beta$

Complementarity

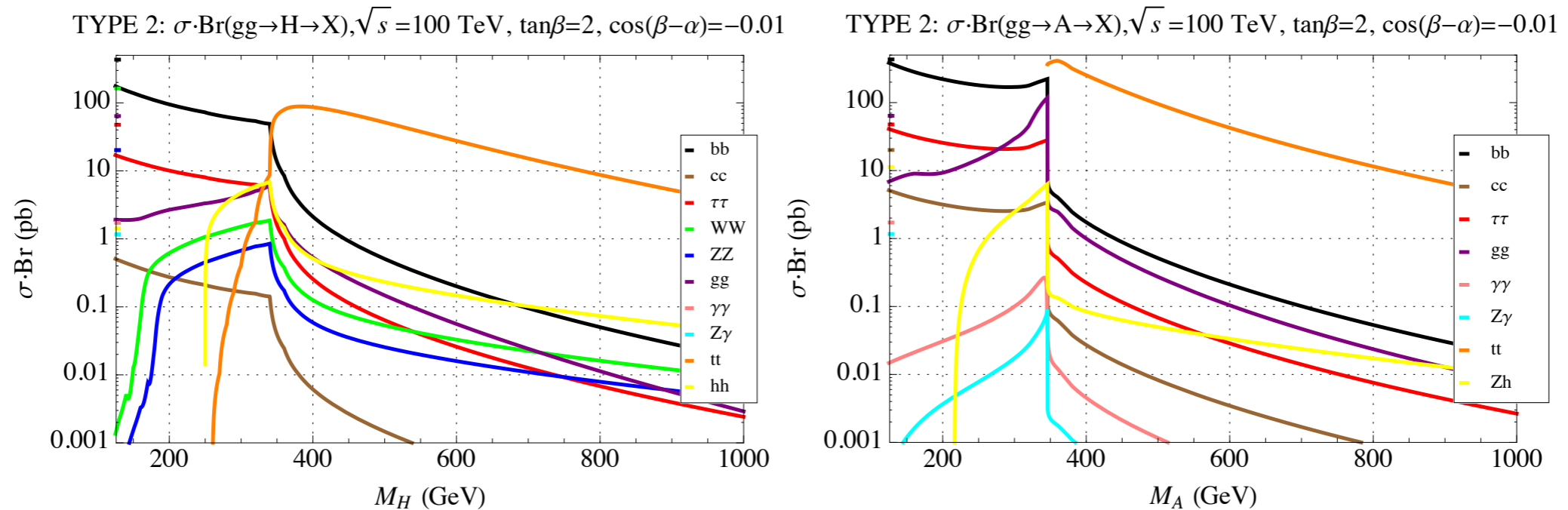
[Barger, Everett, Logan, Shaughnessy | 308.0052]



Where can we expect signs of heavy Higgses given the envelope of coupling measurements? E.g., starting the VLHC after ILC 1000-level precision?

Complementarity

If observed Higgs couplings very SM-like,
promising final for heavy (neutral) Higgses
include tt , WW , ZZ , $\tau\tau$, bb , hh , Zh .



(See backup for details)

Collider study

with E.Brownson, N.Dhingra, U.Heintz, G.Kukartsev, M.Narain, N.Parashar, J.Stupak
[1308.6334]

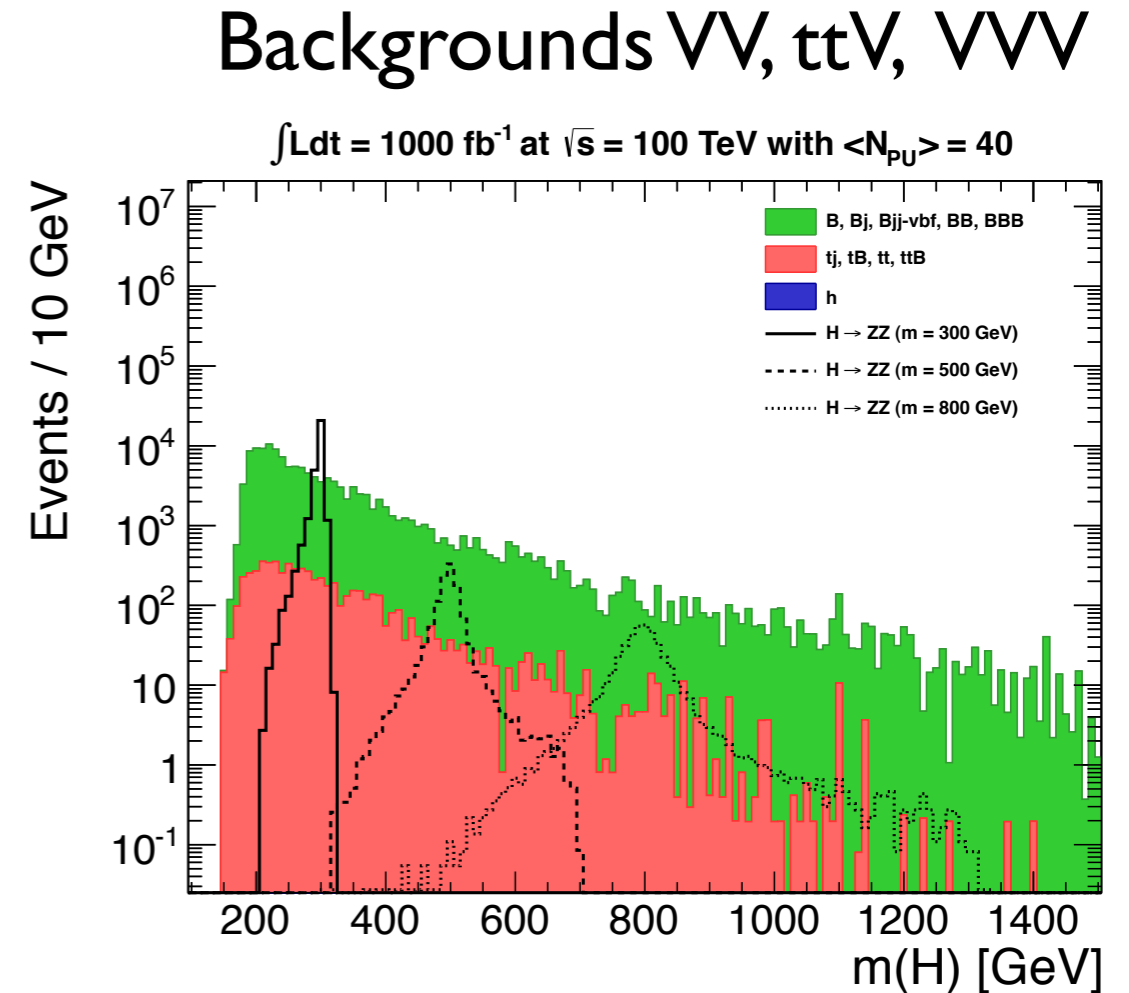
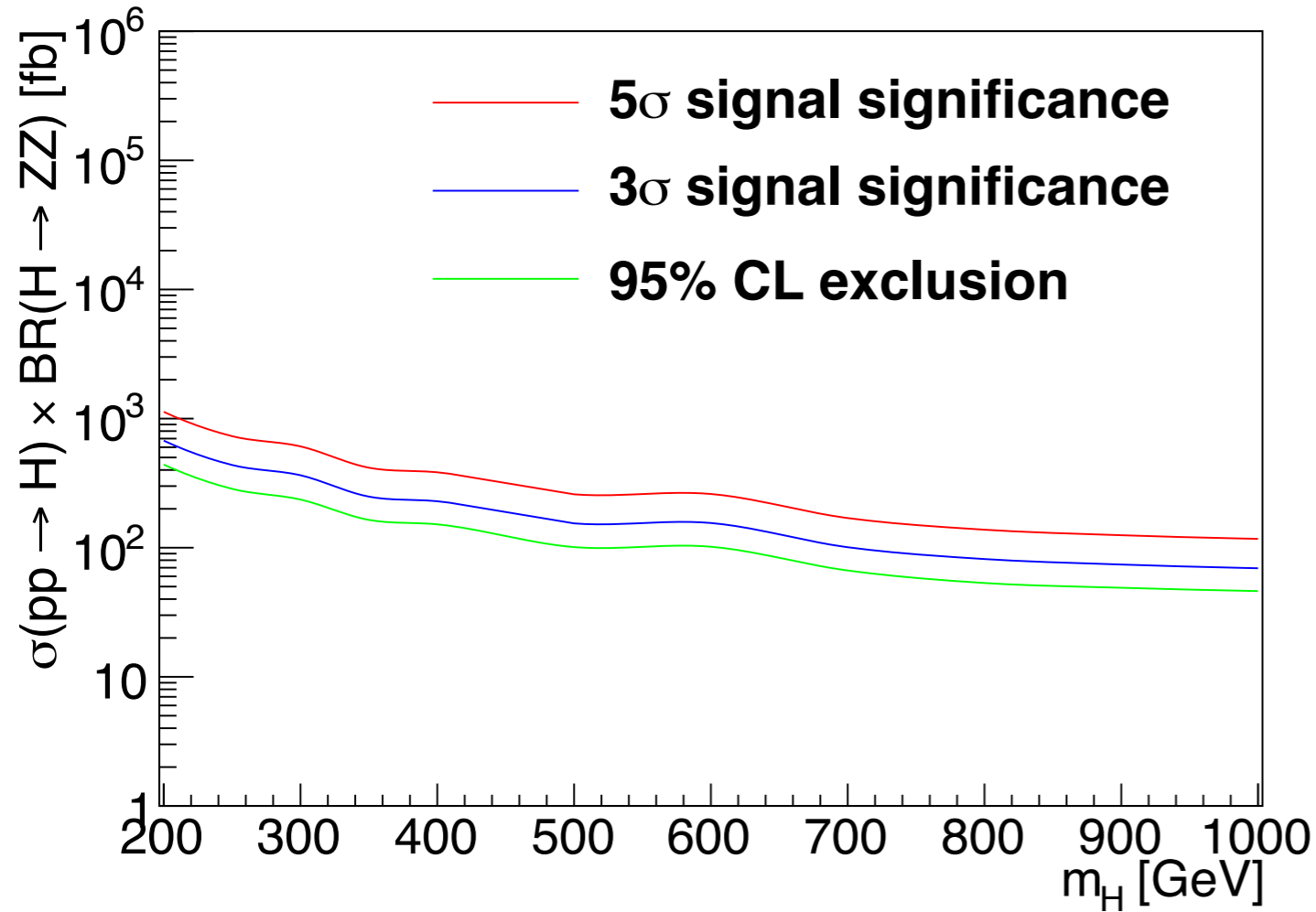
To get a sense for 100 TeV reach, investigate
two promising 2HDM modes:

$$H \rightarrow ZZ \rightarrow 4\ell \quad A \rightarrow Zh \rightarrow \ell\ell(\tau\tau+bb)$$

	LHC Run II	HL-LHC	HE-LHC	VLHC
$s^{1/2}$ [TeV]	14	14	33	100
L [fb ⁻¹]	300	3000	3000	1000
$\langle N_{PU} \rangle$	50	140	140	40

Snowmass 2013 VLHC simulation parameters
Snowmass backgrounds, pileup, etc.

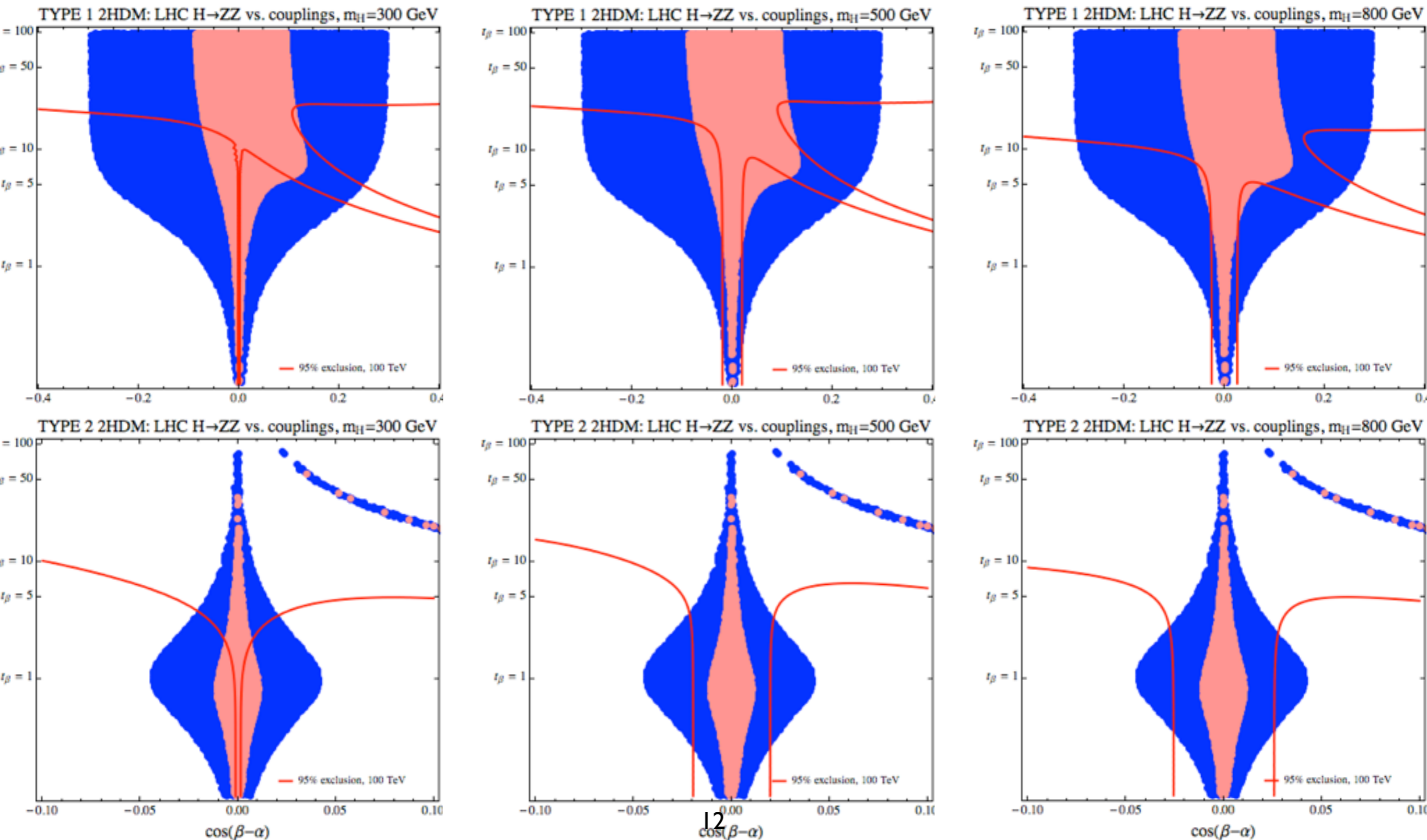
$$H \rightarrow ZZ \rightarrow 4\ell$$



With relatively crude cuts, reach into the 100's of fb @ 100 TeV

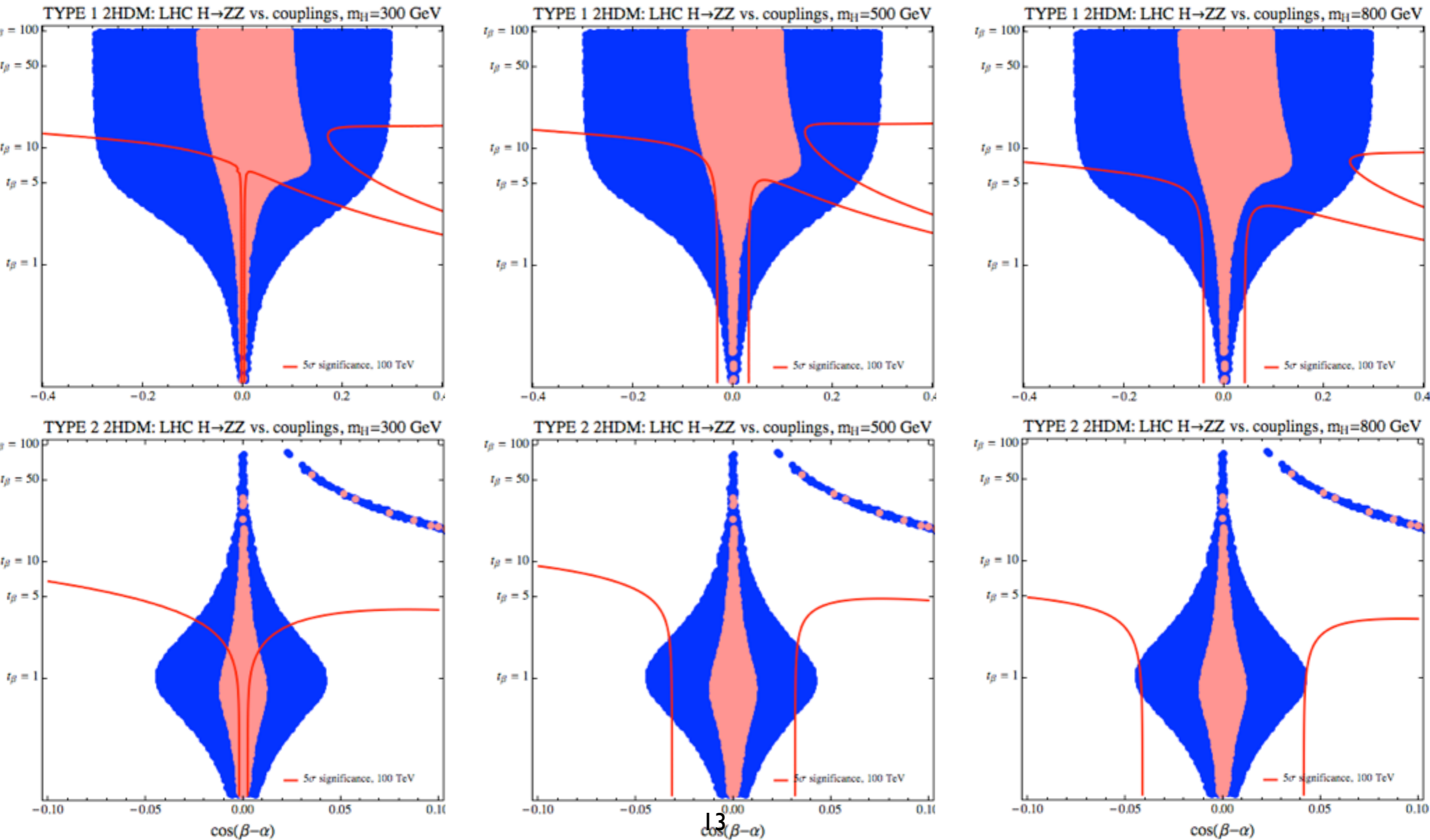
Exclusion Complementarity

(Blue: allowed by couplings after 14 TeV, 3/ab; Pink: allowed after ILC 1 TeV)

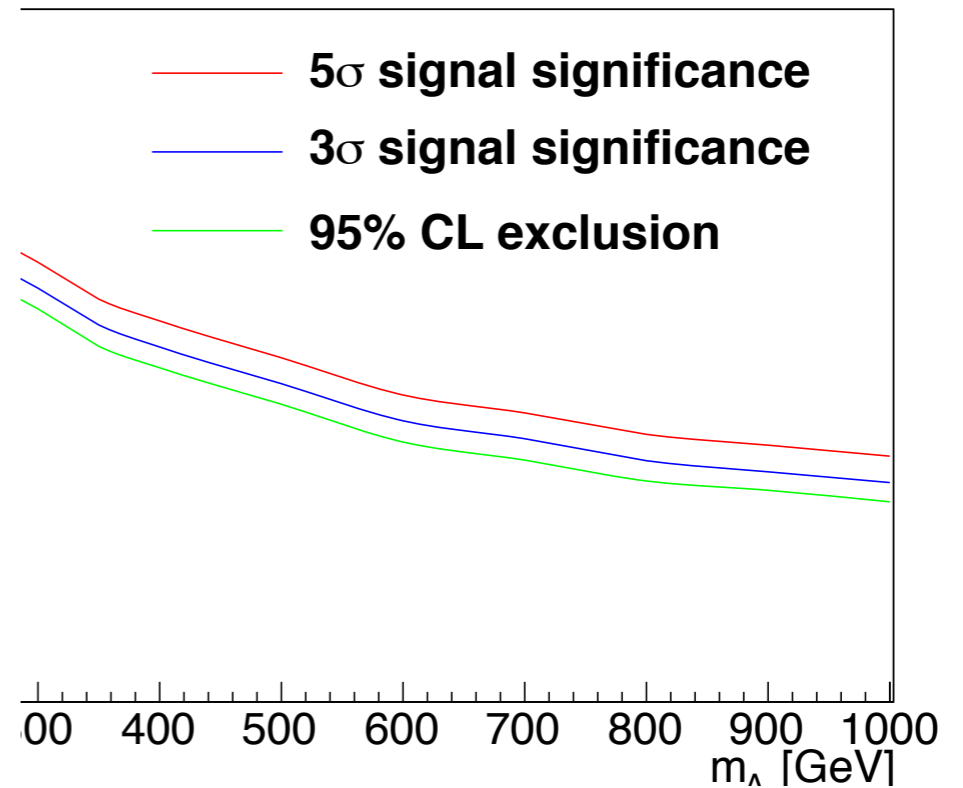
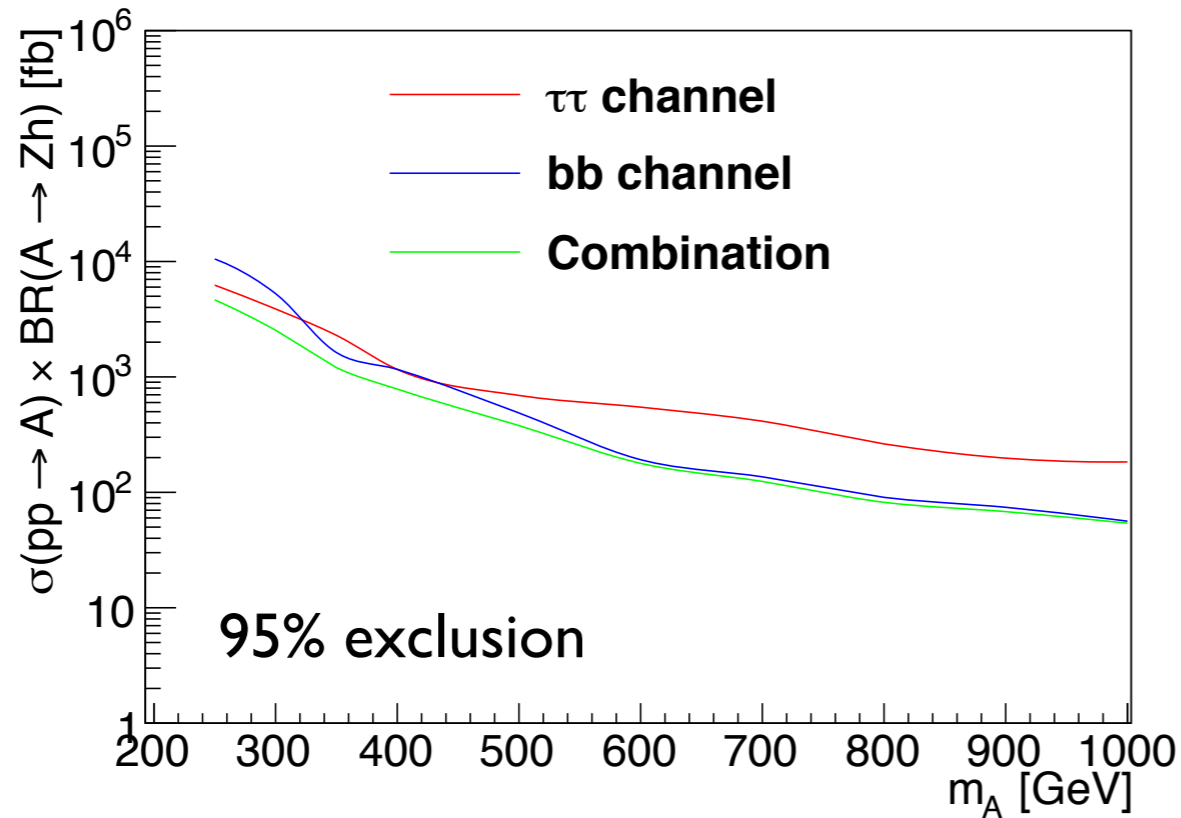


Discovery Complementarity

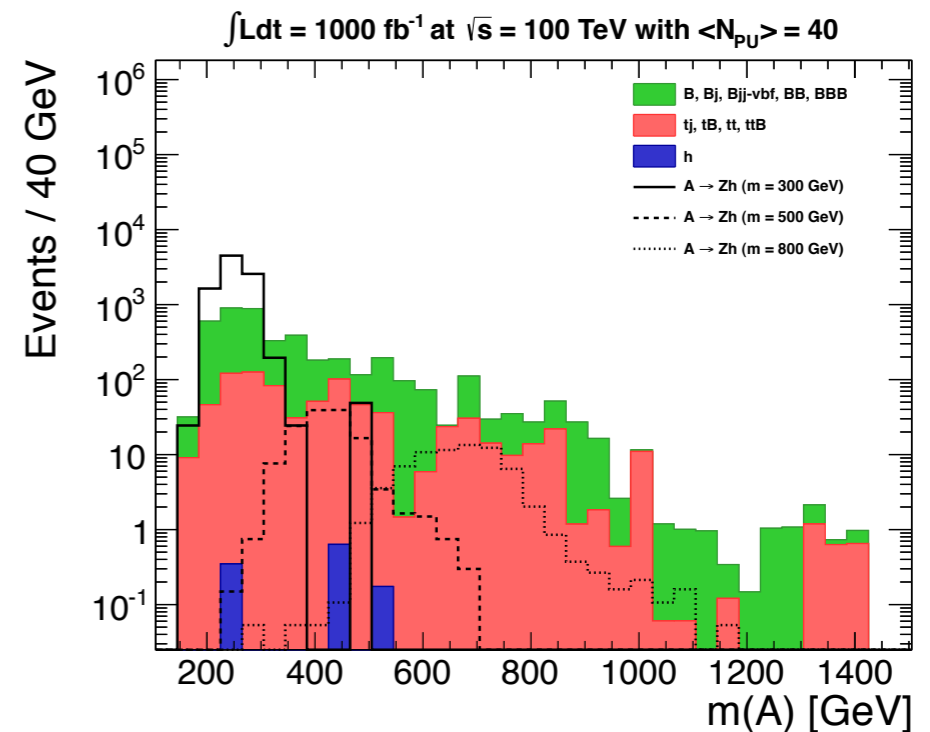
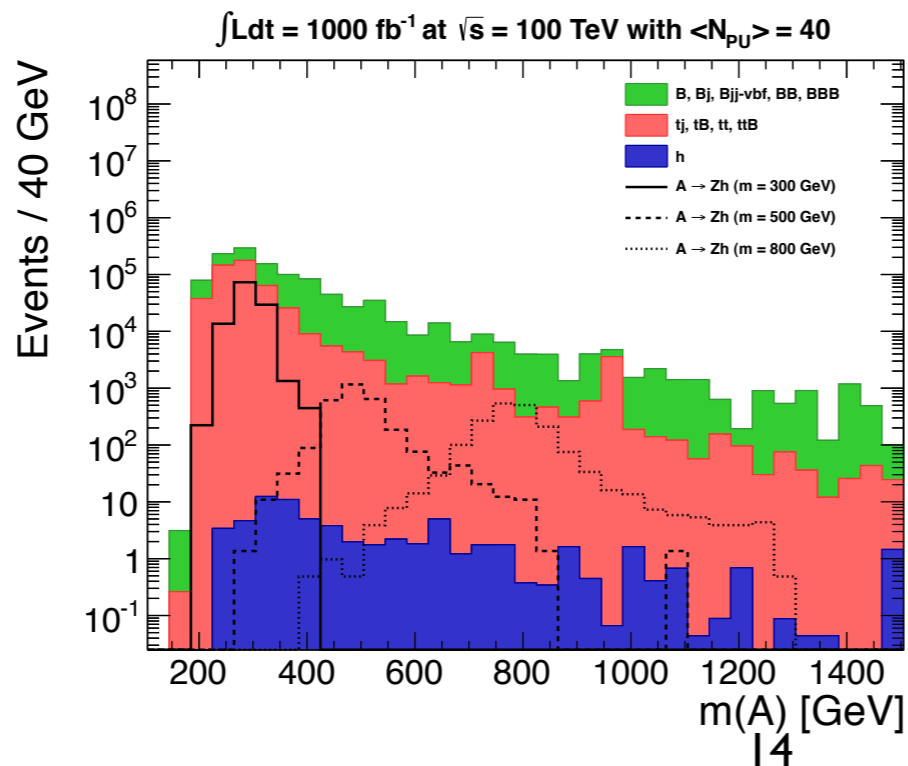
(Blue: allowed by couplings after 14 TeV, 3/ab; Pink: allowed after ILC 1 TeV)



$A \rightarrow Zh \rightarrow \ell\ell (\tau\tau + bb)$

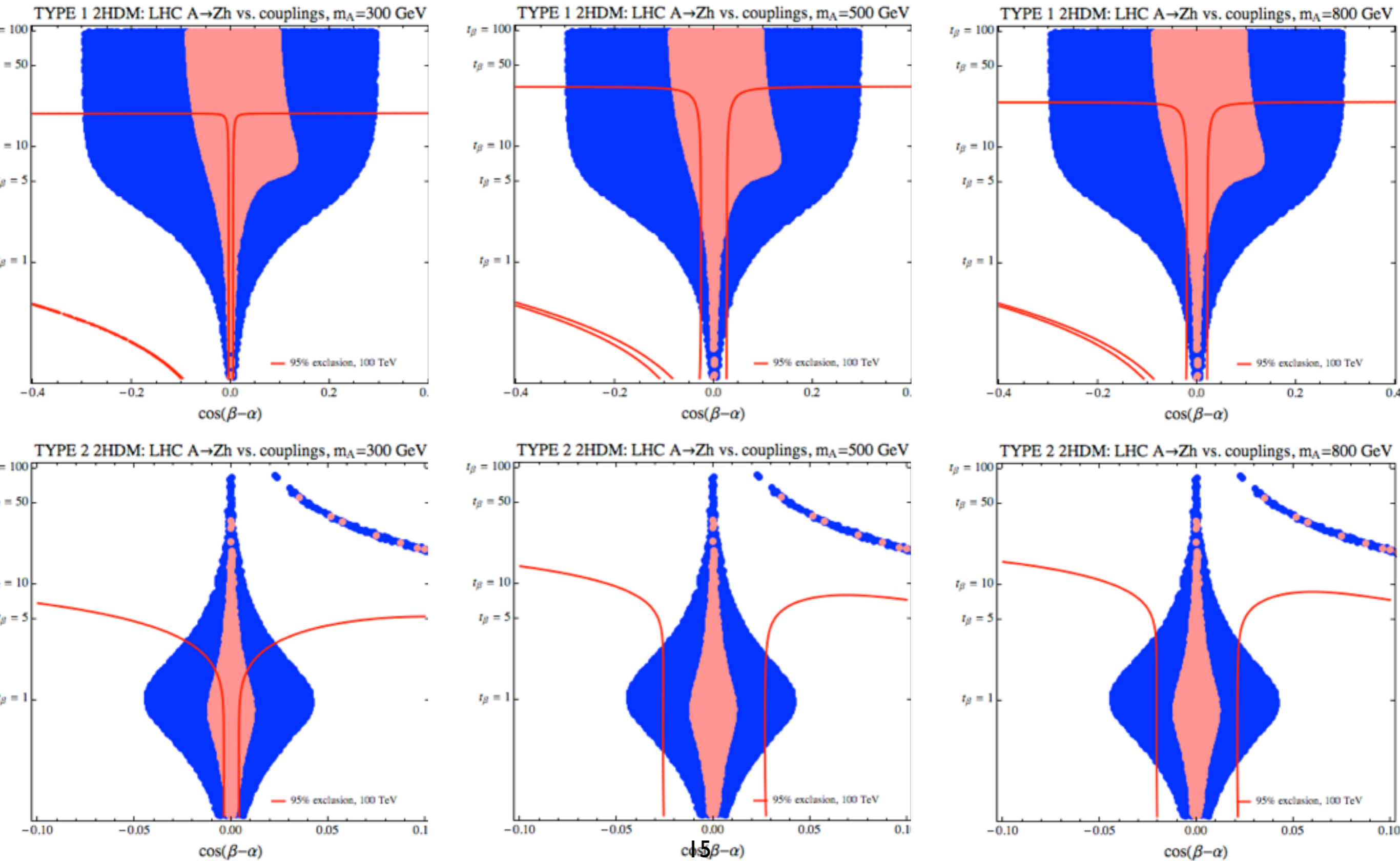


Bkgds V,
 tt, VV,
 ttV, tV,
 VVV



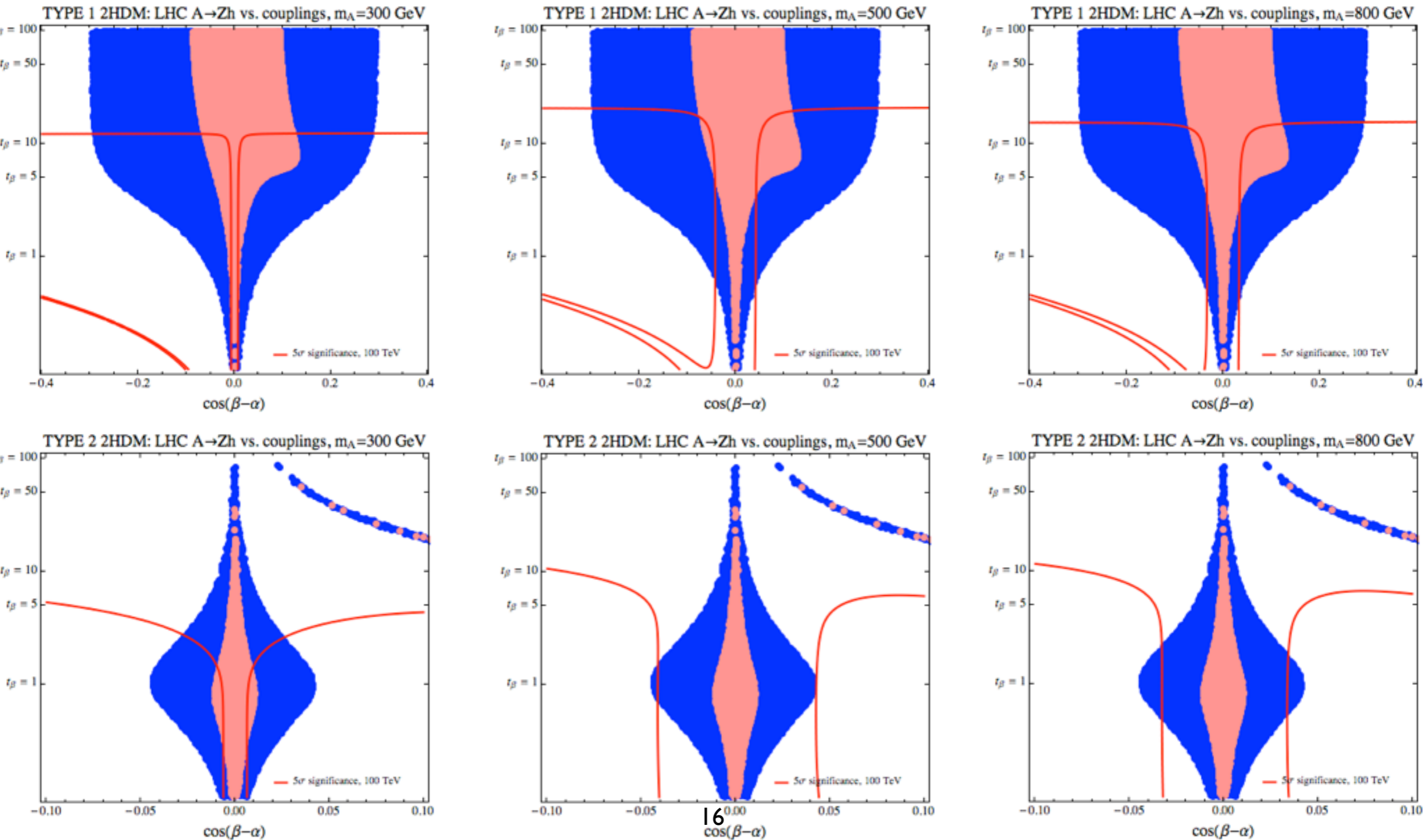
Exclusion complementarity

(Blue: allowed by couplings after 14 TeV, 3/ab; Pink: allowed after ILC 1 TeV)

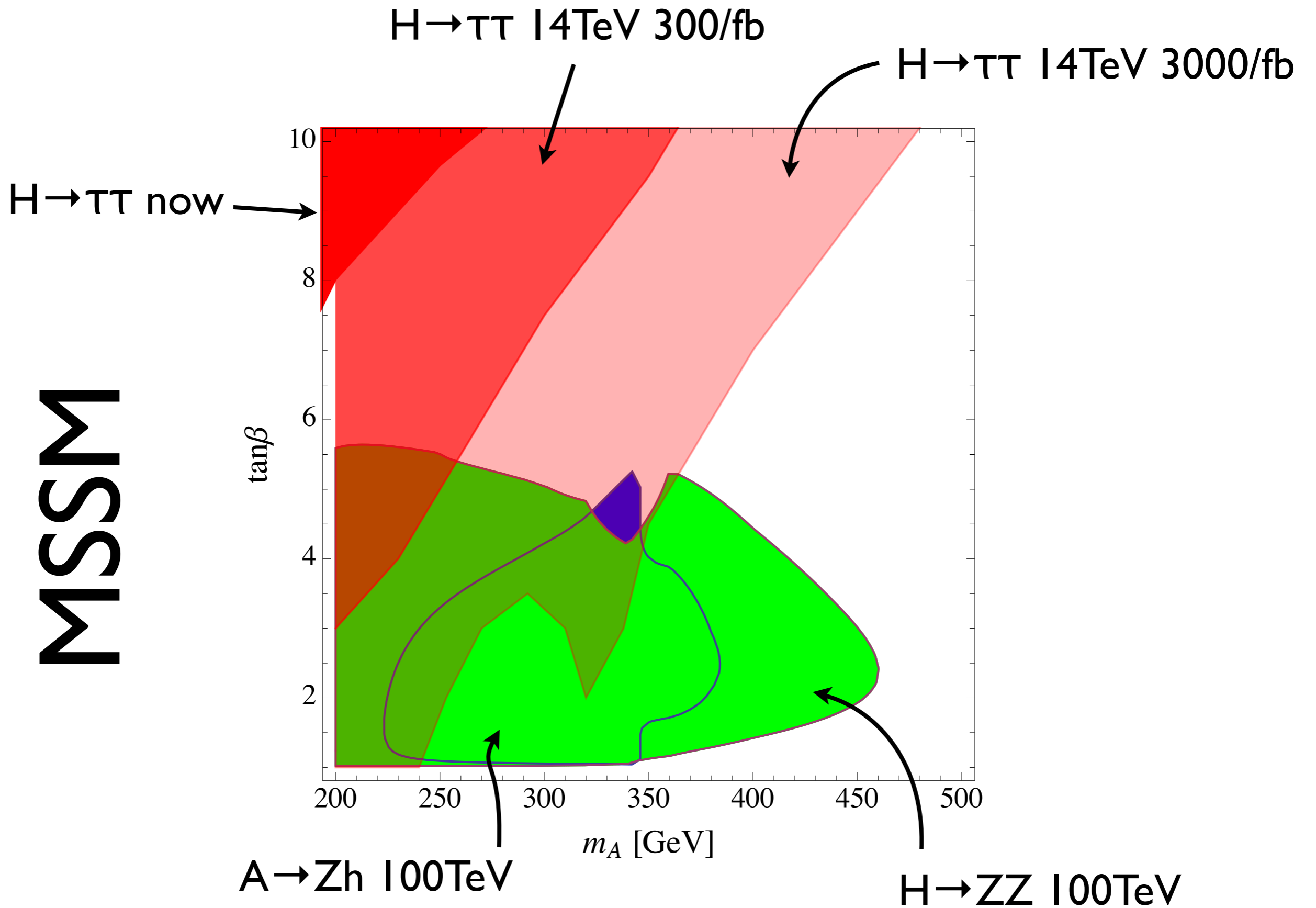


Discovery complementarity

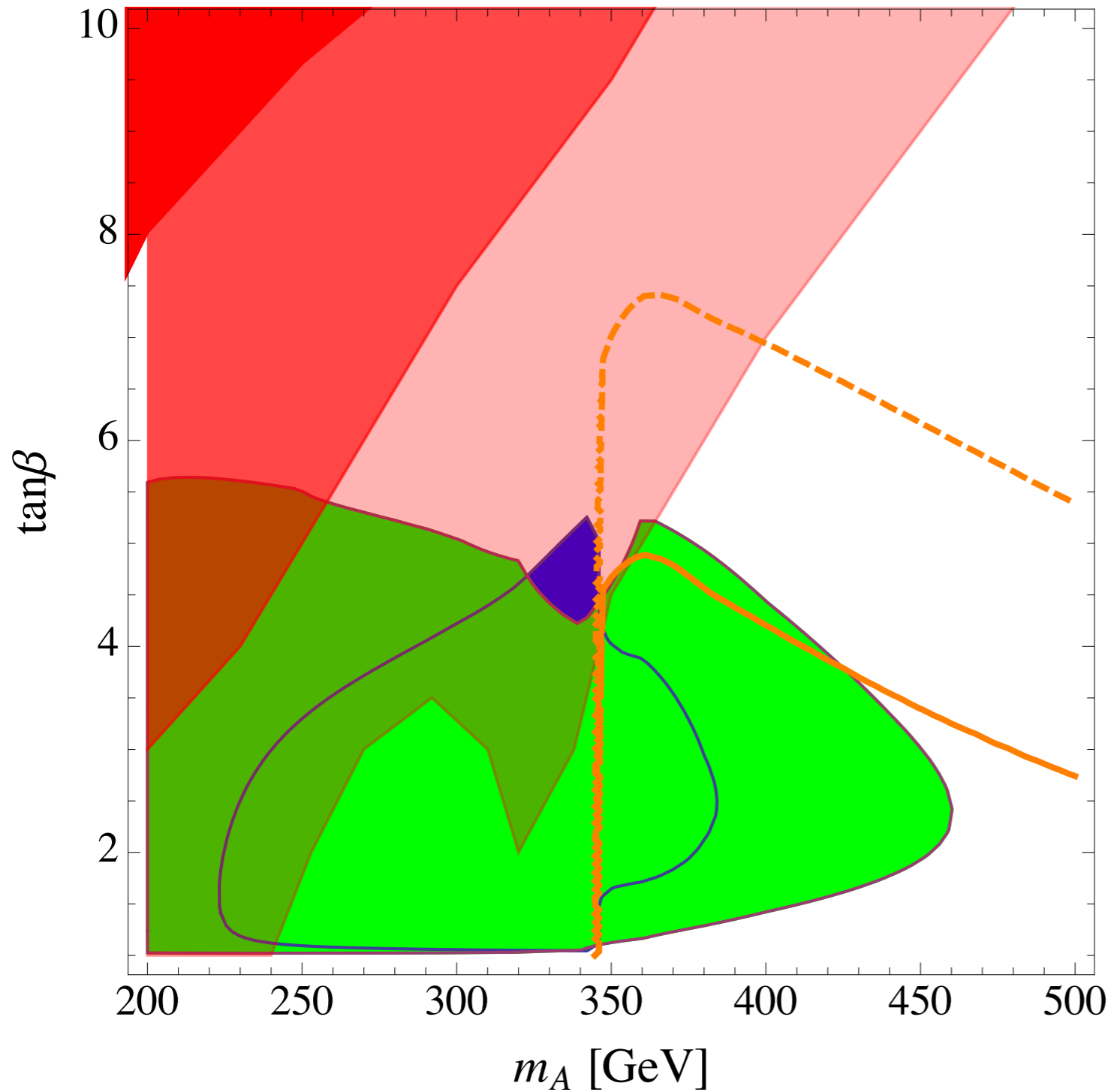
(Blue: allowed by couplings after 14 TeV, 3/ab; Pink: allowed after ILC 1 TeV)



MSSM



Likely too pessimistic, but where's the signal going?

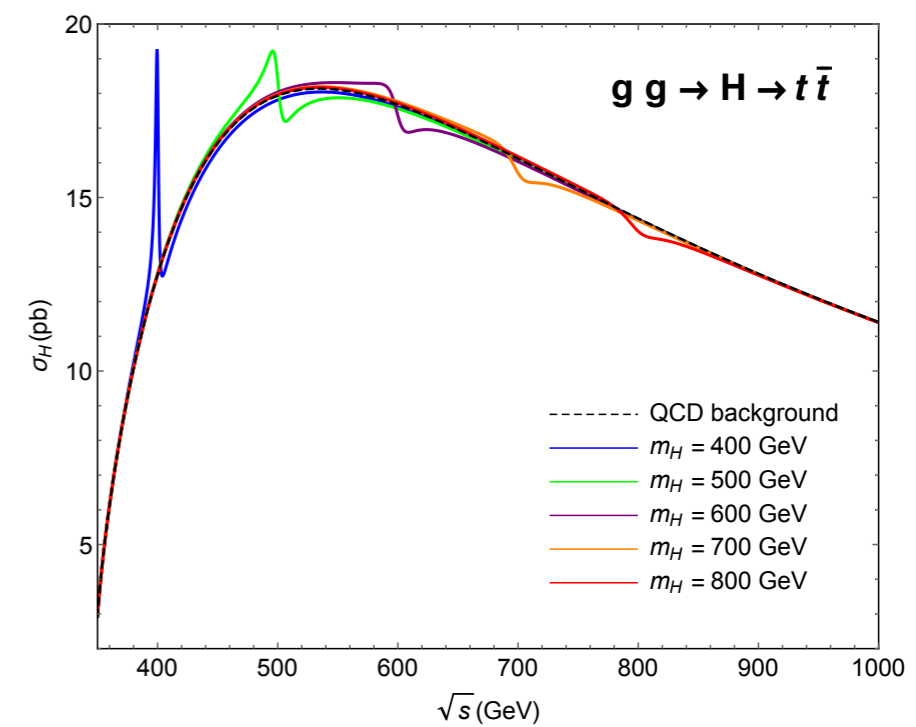


$H \rightarrow \tau\tau$ should take care of high- $\tan\beta$ region. Going to need to search for decays into top pairs. Optimizing charged Higgs search also crucial.

$\sigma \cdot \text{Br}(gg \rightarrow A \rightarrow t\bar{t}) = 10 \text{ pb}$

$\sigma \cdot \text{Br}(gg \rightarrow A \rightarrow t\bar{t}) = 50 \text{ pb}$

Challenge: interferes with SM continuum background; interference effects are dominant above 400 GeV.



Conclusions

- Strong motivation for BSM Higgses outside LHC14 reach but within 100 TeV reach in \sim few TeV range.
- 100 TeV capability to pursue both alignment and decoupling.
- New opportunities at 100 TeV from enhanced top associated production: $tt\Phi$ and tbH^\pm appreciable and provide new handles for otherwise challenging final states.
- Under-studied modes such as $\Phi \rightarrow tt$ become increasingly important at high mass given projected coupling limits.
- Excellent complementarity between coupling measurements and reach for current 100 TeV studies, demonstrates high utility of a 100 TeV BSM Higgs program.

Desiderata

Moving forward, we'll have a more complete picture from:

- Dedicated $\Phi \rightarrow \tau\tau$ study at 100 TeV.
- Dedicated $\Phi \rightarrow tt$ search at 8/14 TeV, study for 100 TeV *including interference*.
- Study & exploitation of $tt\Phi$ and tbH^\pm modes, particularly for otherwise-difficult $\Phi \rightarrow tt$ and $H^\pm \rightarrow tb$ decays.

Backup

A simplified parameter space

Physical d.o.f. are (8-3=5): h, H, A, H^\pm

After EWSB there are 9 free parameters in CP-conserving scalar potential.

Useful basis of 4 physical masses, 2 angles, 3 couplings:

$$m_h, m_H, m_A, m_{H^\pm} \qquad \tan \beta \equiv \langle \Phi_2 \rangle / \langle \Phi_1 \rangle$$

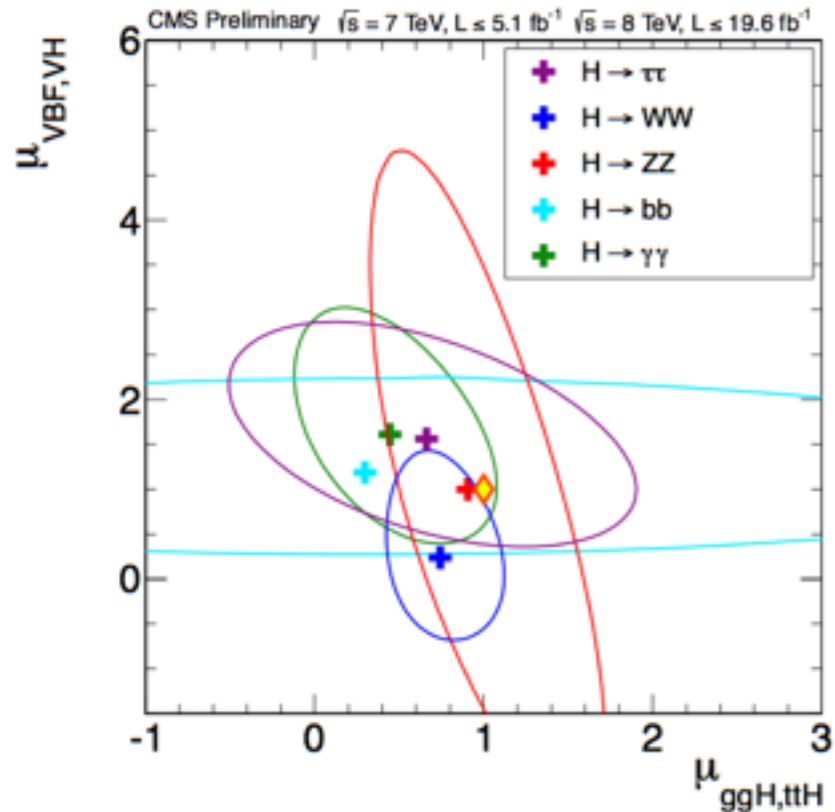
$$\alpha : \begin{pmatrix} \sqrt{2} \operatorname{Re}(\Phi_2^0) - v_2 \\ \sqrt{2} \operatorname{Re}(\Phi_1^0) - v_1 \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} h \\ H \end{pmatrix}$$

$$\lambda_5, \lambda_6, \lambda_7 \qquad (\text{only appear in trilinear couplings})$$

Couplings of scalars to fermions, vectors only depend on angles.

$$\text{Discrete symm. for flavor: } \lambda_{6,7} = 0 \qquad \text{MSSM: } \lambda_{5,6,7} = 0$$

Alignment limit

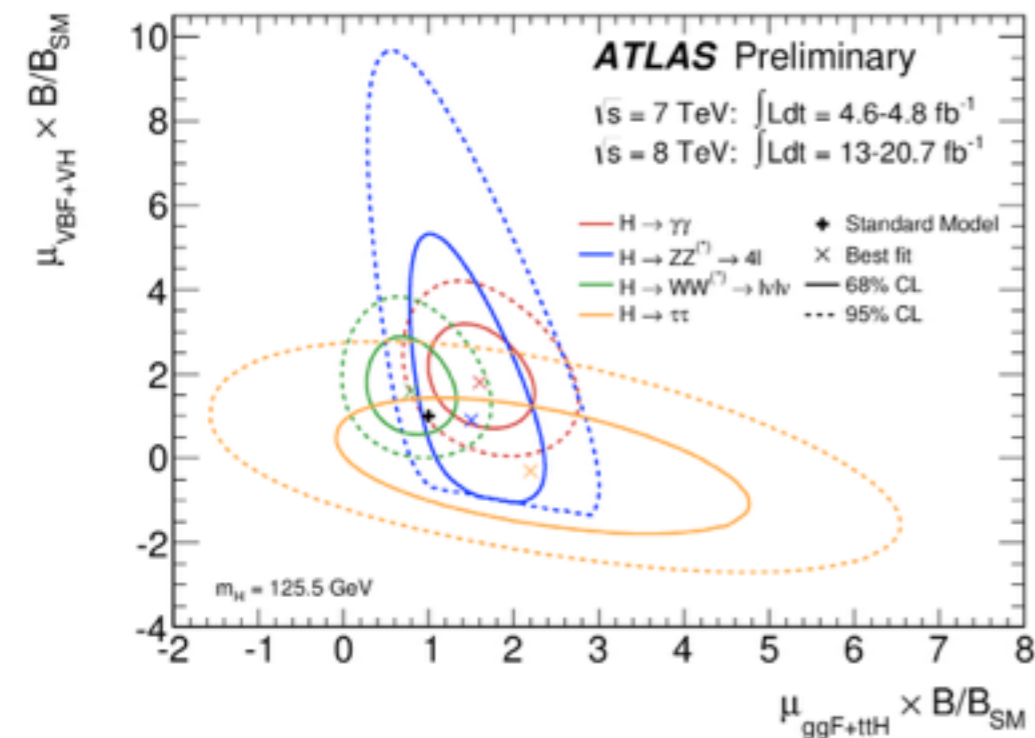


- Couplings of the observed Higgs are so far approximately SM-like
- Strongly suggests proximity to the alignment limit

$$\alpha \approx \beta - \pi/2$$

- In this limit h is the fluctuation around the vev, while remaining scalars are spectators to EWSB
- (Limit obtainable via decoupling in mass or accidentally, via dimensionless couplings)
- Useful to expand in

$$\begin{aligned} \delta &= \beta - \alpha - \pi/2 \\ &\approx -\cos(\beta - \alpha) \end{aligned}$$



Four discrete 2HDM types. All couplings to SM states fixed in terms of two angles.

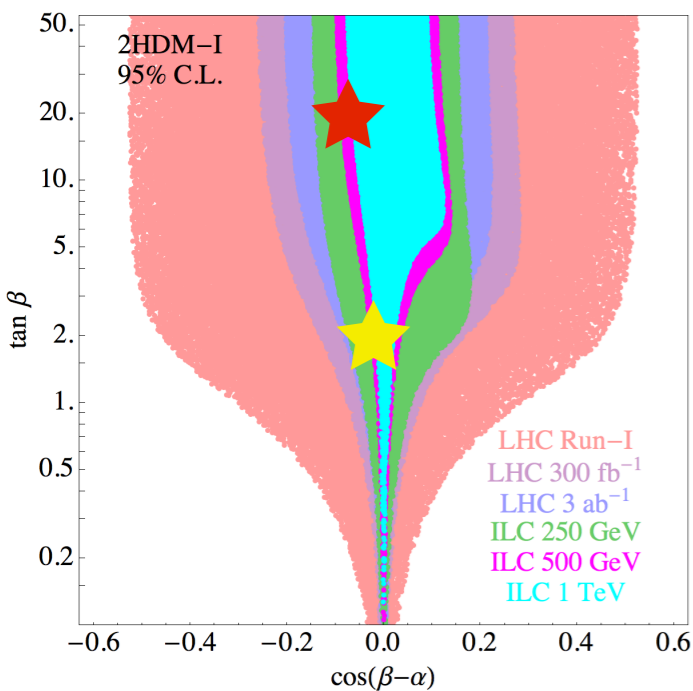
	2HDM I	2HDM II	2HDM III	2HDM IV
u	Φ_2	Φ_2	Φ_2	Φ_2
d	Φ_2	Φ_1	Φ_2	Φ_1
e	Φ_2	Φ_1	Φ_1	Φ_2

$y_{2\text{HDM}}/y_{\text{SM}}$	2HDM 1	2HDM 2
hVV	$1 - \delta^2/2$	$1 - \delta^2/2$
hQu	$1 - \delta/t_\beta$	$1 - \delta/t_\beta$
hQd	$1 - \delta/t_\beta$	$1 + \delta t_\beta$
hLe	$1 - \delta/t_\beta$	$1 + \delta t_\beta$
HVV	$-\delta$	$-\delta$
HQu	$-\delta - 1/t_\beta$	$-\delta - 1/t_\beta$
HQd	$-\delta - 1/t_\beta$	$-\delta + t_\beta$
HLe	$-\delta - 1/t_\beta$	$-\delta + t_\beta$
AVV	0	0
AQu	$1/t_\beta$	$1/t_\beta$
AQd	$-1/t_\beta$	t_β
ALe	$-1/t_\beta$	t_β

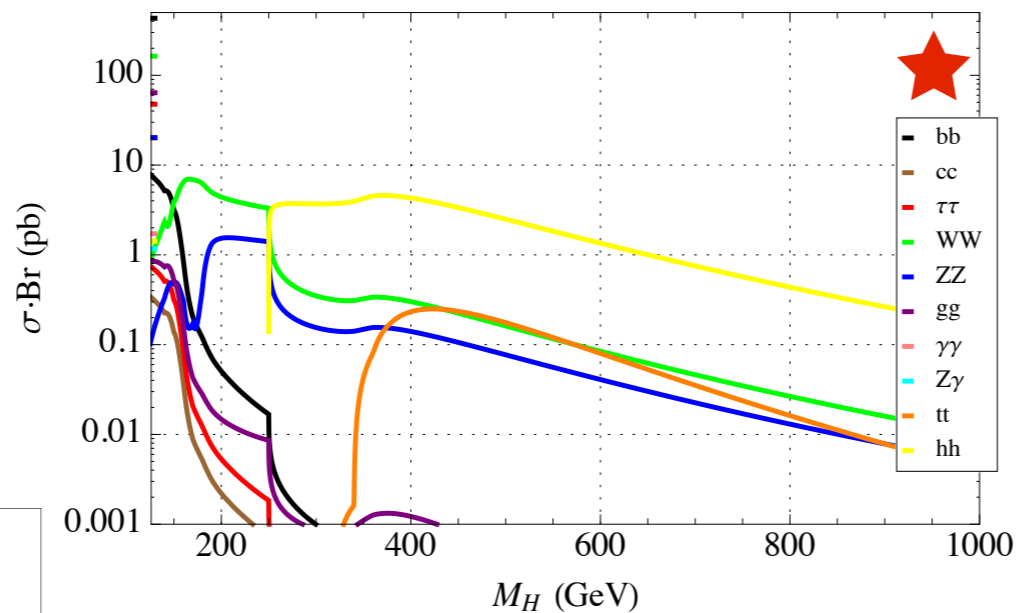
- Scalar self-couplings have additional parametric freedom.
- Gives a map between current fits to the Higgs couplings and the possible size of NP signals.
- H, A are similar d.o.f. in alignment limit; H^+ couplings analogous to A.
- Focus on the two most familiar, Types I and 2.

$$\delta = \beta - \alpha - \pi/2$$

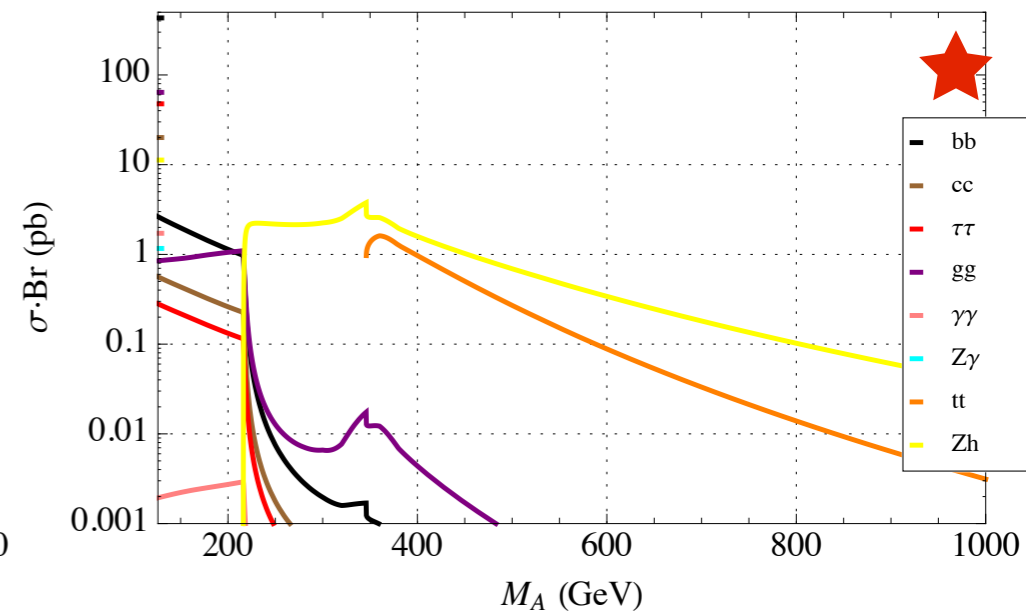
Type I



TYPE 1: $\sigma \cdot \text{Br}(gg \rightarrow H \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta = 20, \cos(\beta - \alpha) = -0.07$

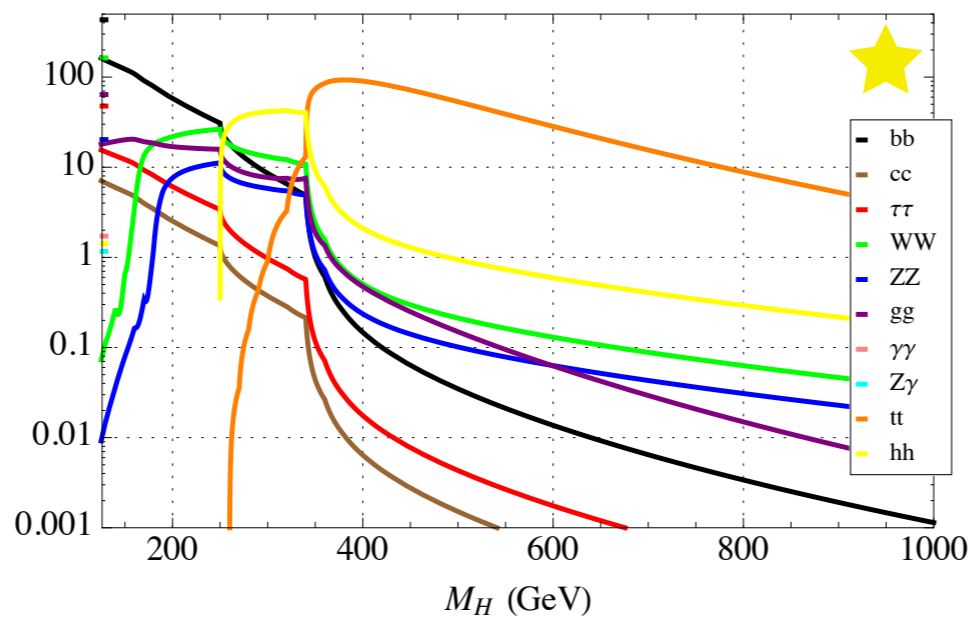


TYPE 1: $\sigma \cdot \text{Br}(gg \rightarrow A \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta = 20, \cos(\beta - \alpha) = -0.07$

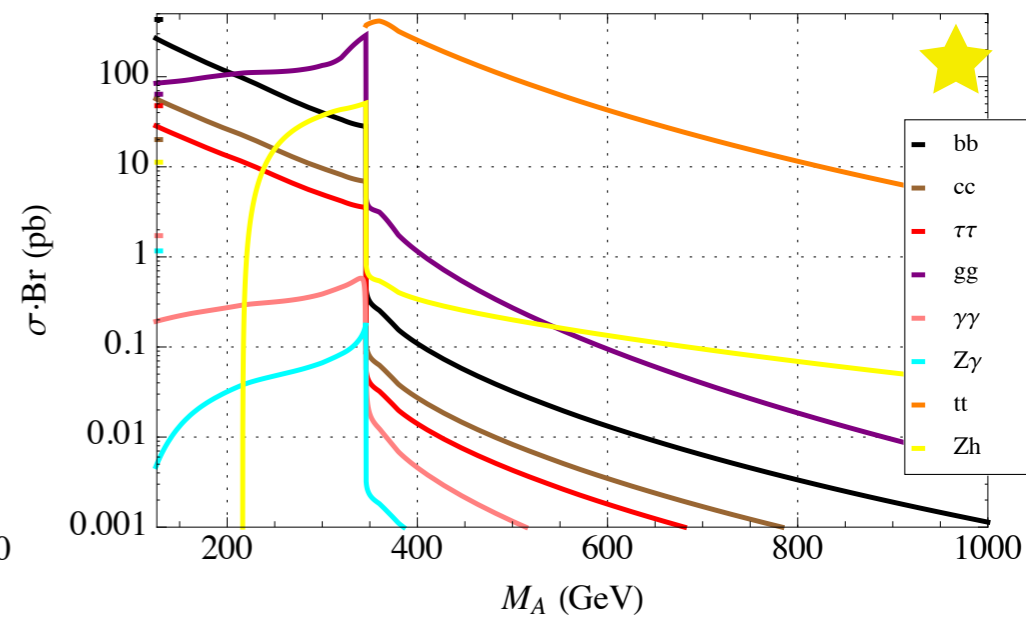


High $\tan\beta$ dominated by hh, Zh, VV, tt

TYPE 1: $\sigma \cdot \text{Br}(gg \rightarrow H \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta = 2, \cos(\beta - \alpha) = -0.02$

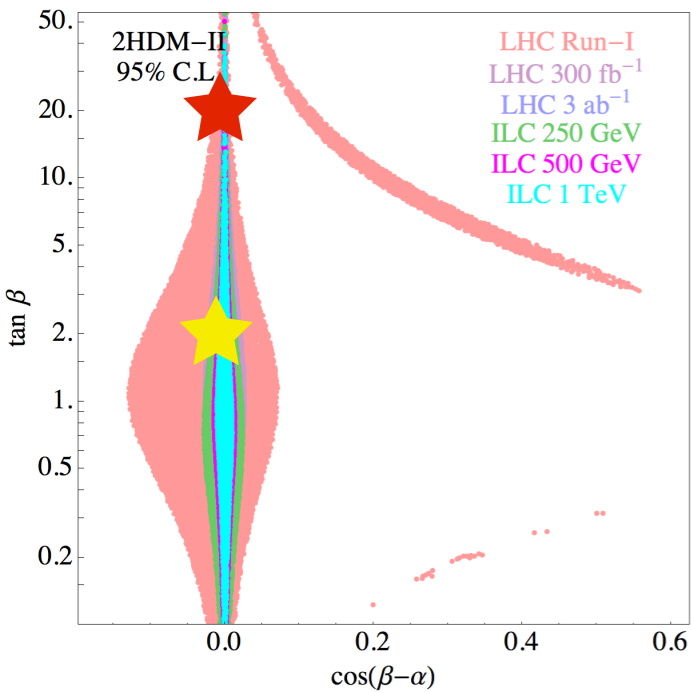


TYPE 1: $\sigma \cdot \text{Br}(gg \rightarrow A \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta = 2, \cos(\beta - \alpha) = -0.02$

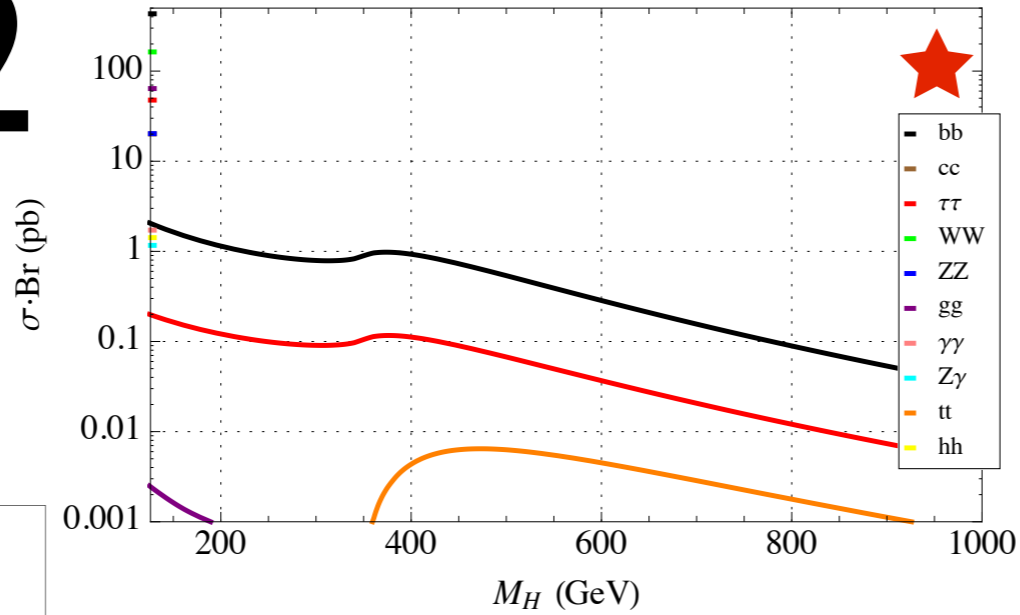


Low $\tan\beta$ dominated by tt, hh, Zh, still some distance from alignment

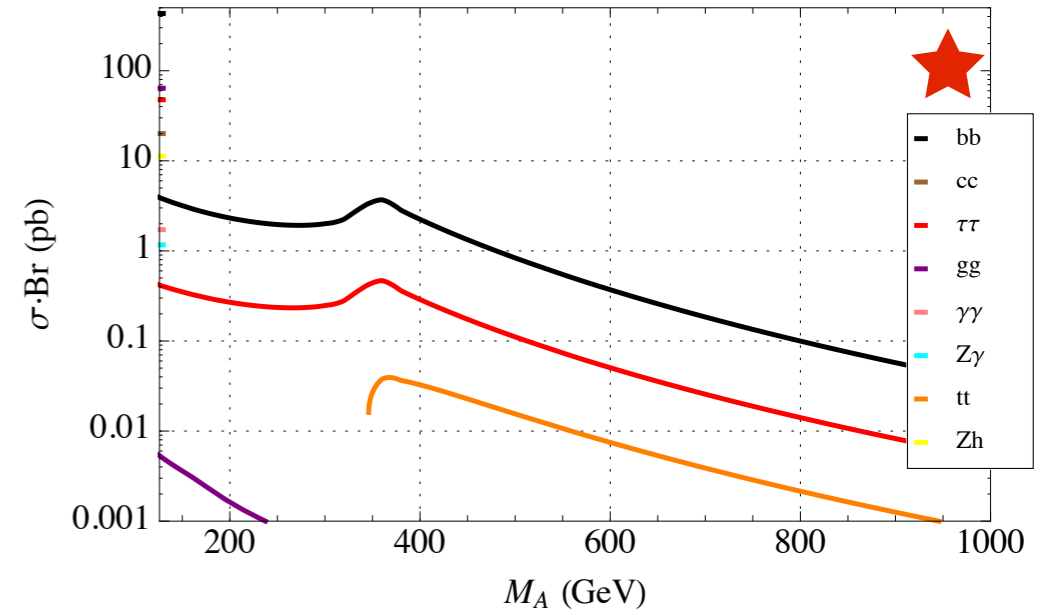
Type 2



TYPE 2: $\sigma \cdot \text{Br}(gg \rightarrow H \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta=20, \cos(\beta-\alpha)=-0.005$

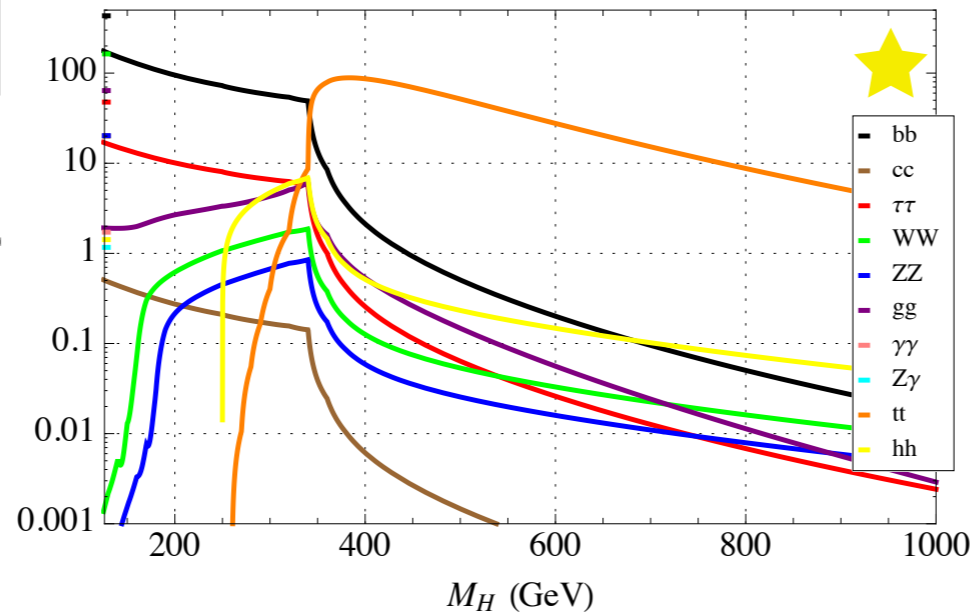


TYPE 2: $\sigma \cdot \text{Br}(gg \rightarrow A \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta=20, \cos(\beta-\alpha)=-0.005$

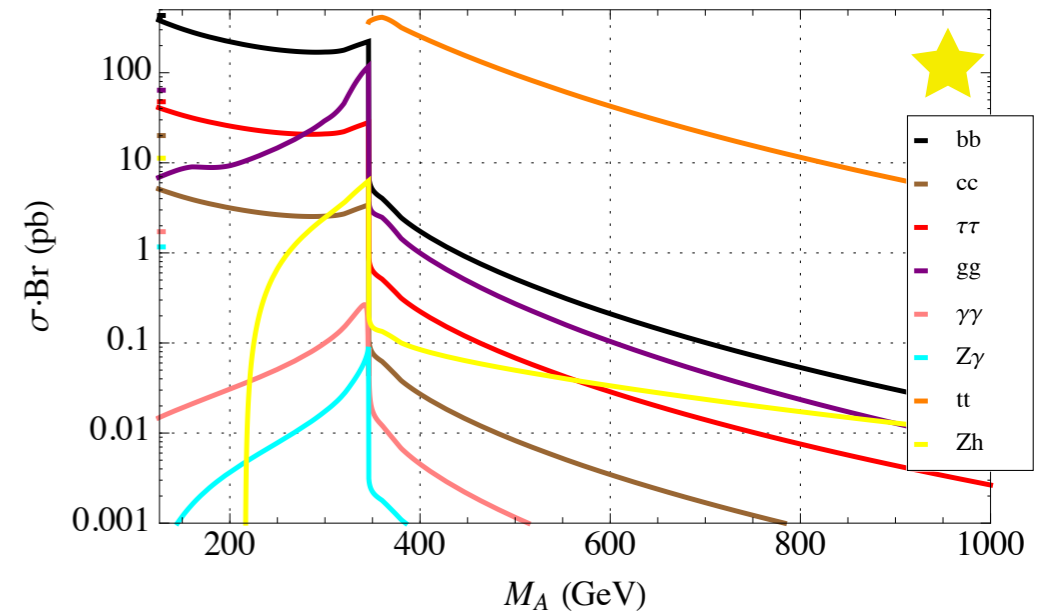


High $\tan\beta$ dominated by bb, $\tau\tau$, as expected from MSSM

TYPE 2: $\sigma \cdot \text{Br}(gg \rightarrow H \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta=2, \cos(\beta-\alpha)=-0.01$

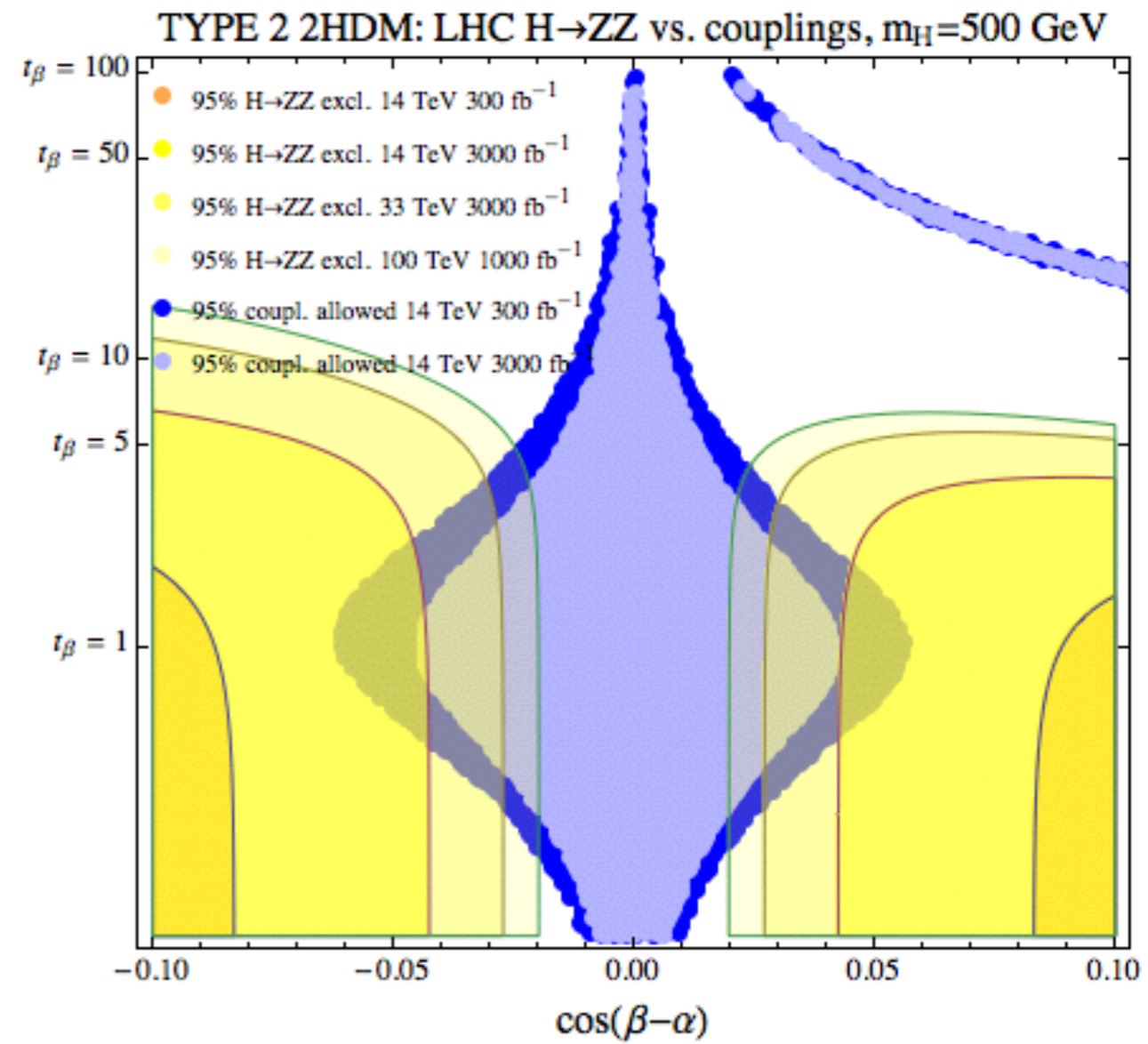
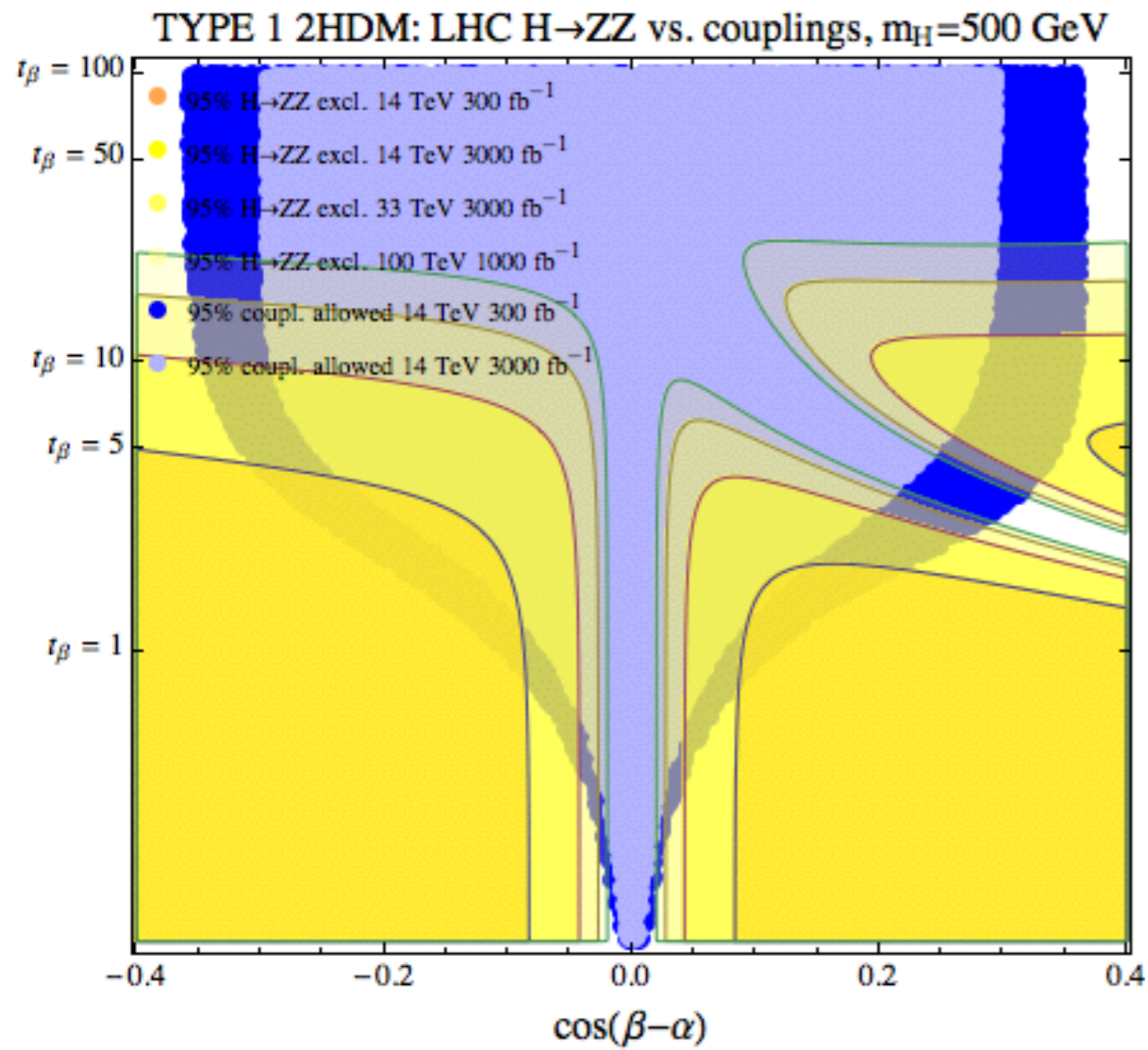


TYPE 2: $\sigma \cdot \text{Br}(gg \rightarrow A \rightarrow X), \sqrt{s} = 100 \text{ TeV}, \tan\beta=2, \cos(\beta-\alpha)=-0.01$



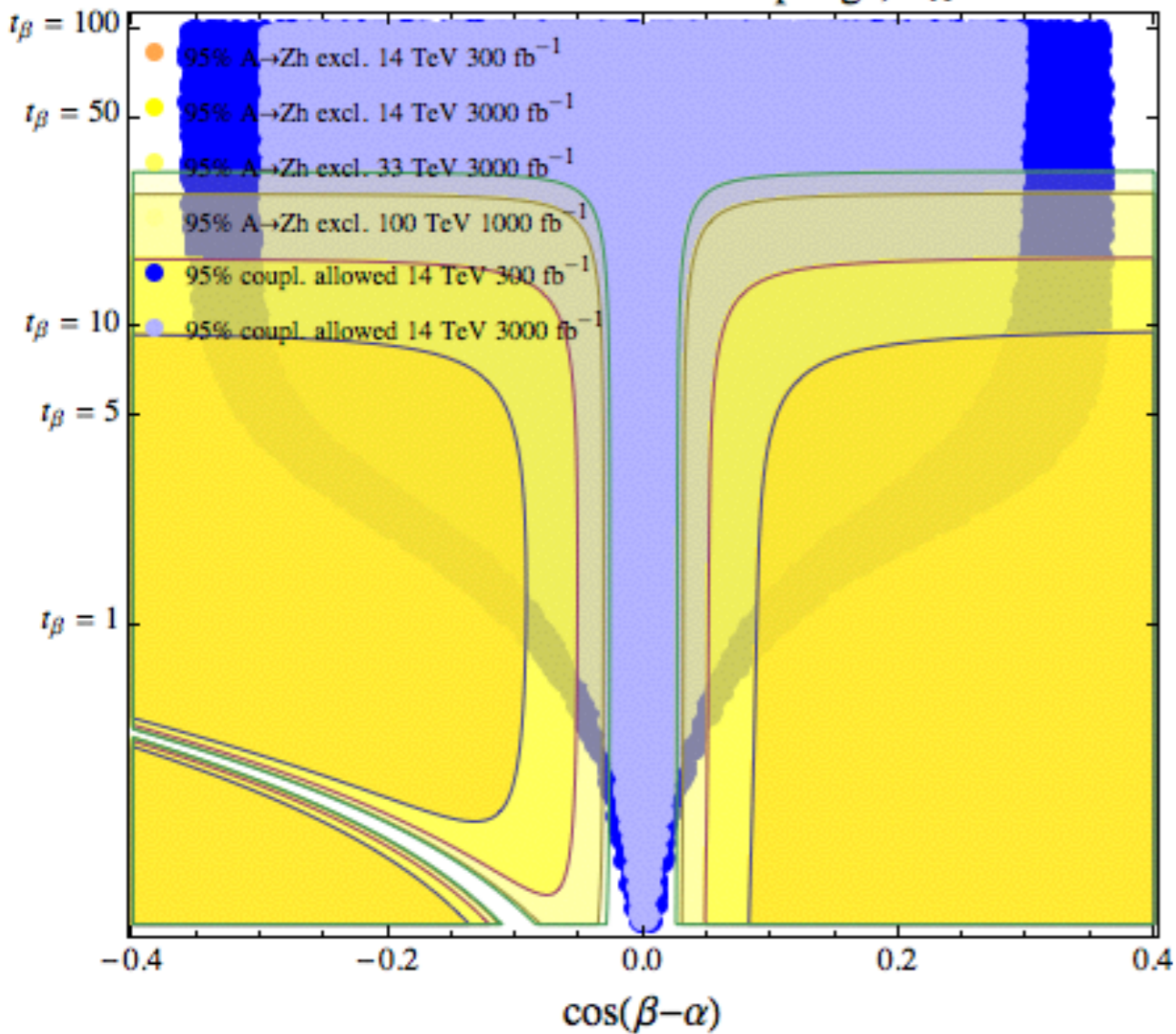
Low $\tan\beta$ dominated by tt, bb; vectors suppressed by alignment

Direct complementarity



Direct complementarity

TYPE 1 2HDM: LHC $A \rightarrow Zh$ vs. couplings, $m_A = 500$ GeV



TYPE 2 2HDM: LHC $A \rightarrow Zh$ vs. couplings, $m_A = 500$ GeV

