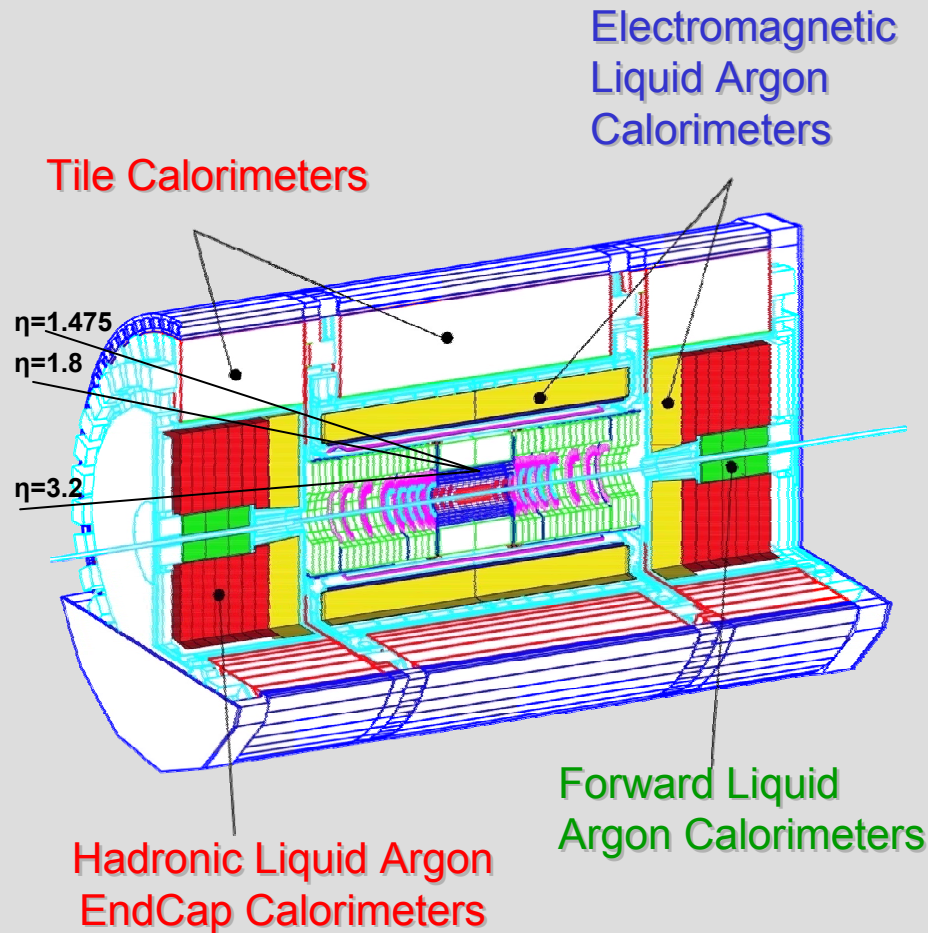


Overview of Jet measurement at ATLAS

- Calorimeters of ATLAS Jet reconstruction
 - Jet energy measurement
 - In situ energy calibration
 - Conclusions

The ATLAS calorimeter

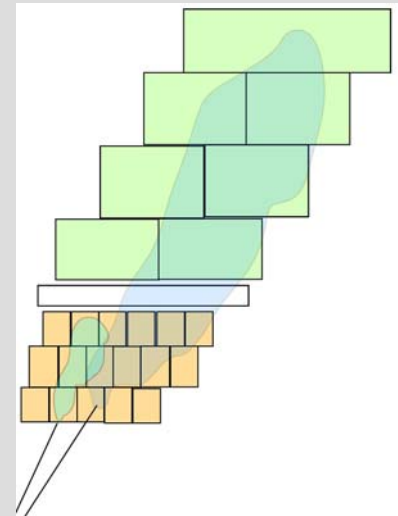


EM accordion 3 long. sections

Central Hadronic 3 long. sections

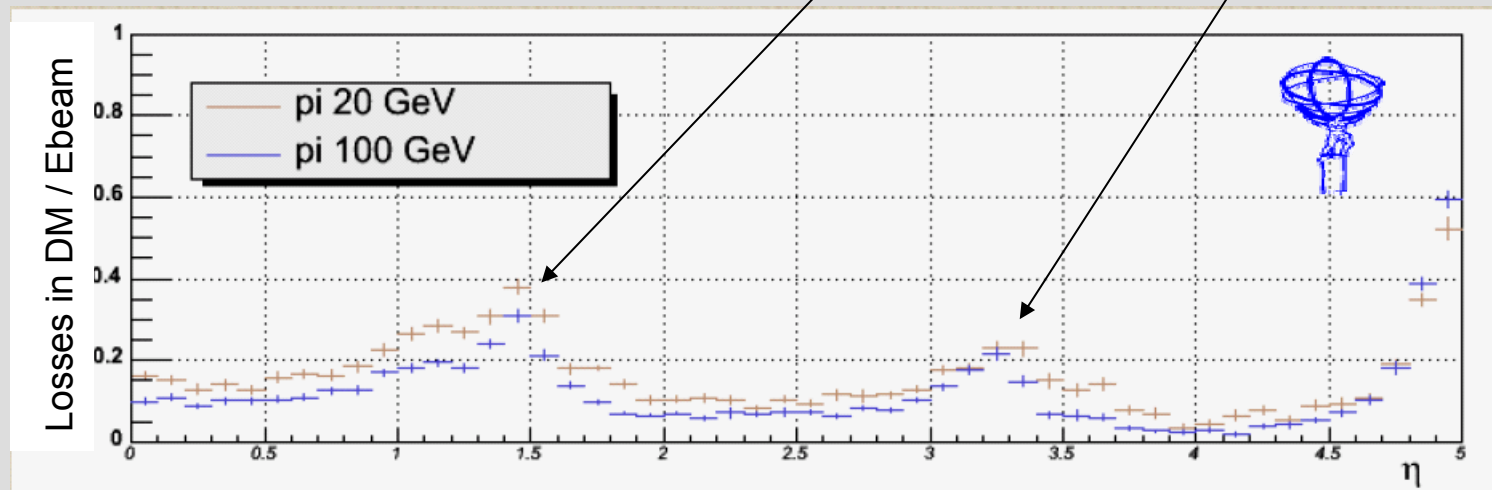
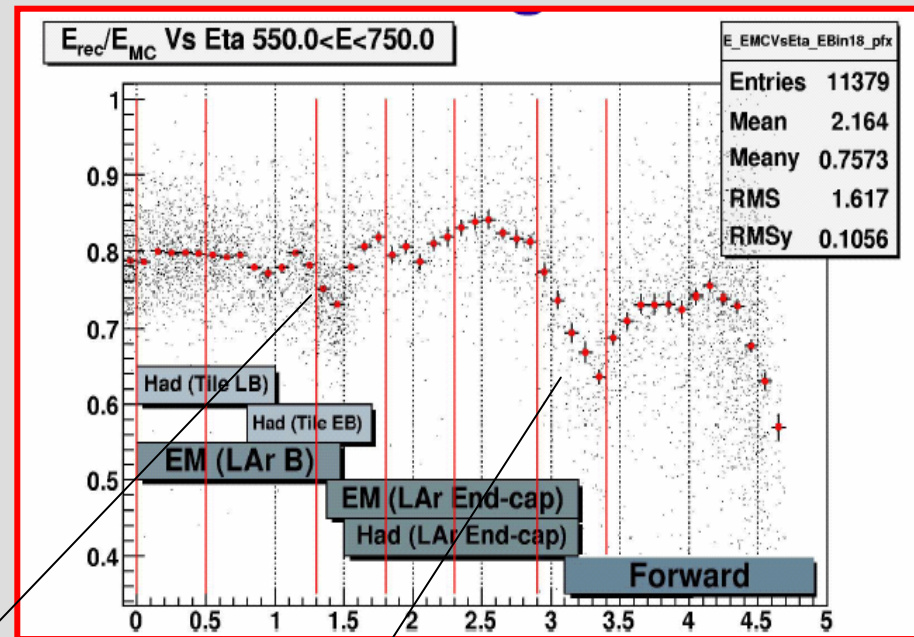
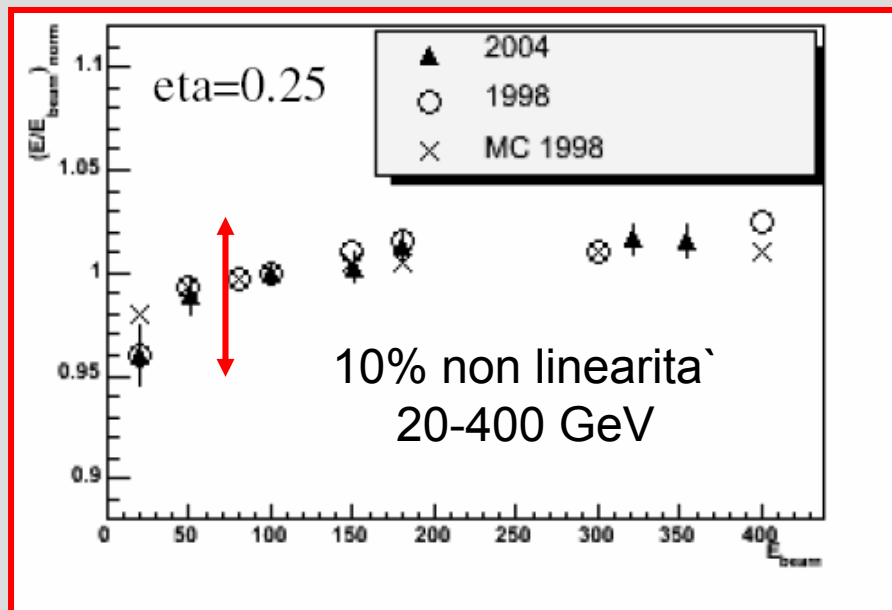
End Cap Hadronic 4 long. sections

Forward calorimeter 3 long. sections

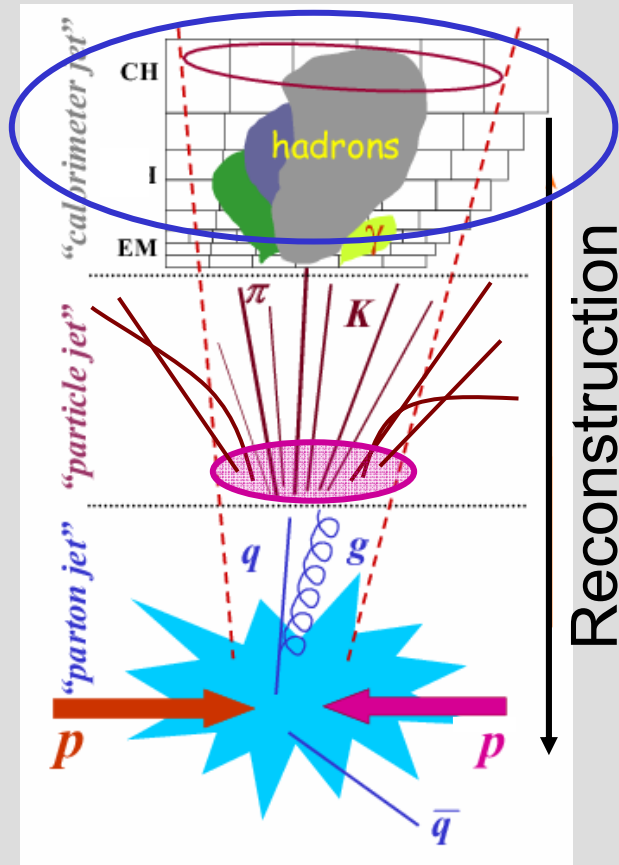


Fine
longitudinal
and lateral
segmentation.
Fundamental
for recovering
compensation.

Detector effects: e/h and cracks



Jet reconstruction phase 1: the calorimeter jet



The cell energy deposits are clusterized to obtain the base objects for jet reconstruction.

Noise (electronic noise and pile-up) suppression algorithms are applied at this stage

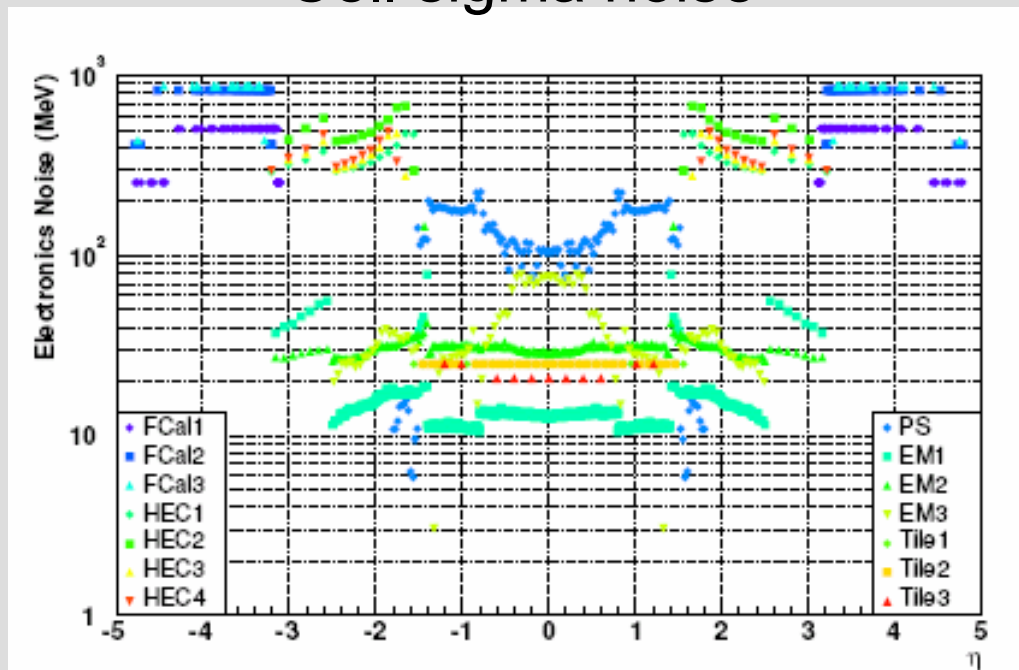
Jet reconstruction algorithms are applied and recombination scheme is used to obtain jet kinematics

Cell energy clusterization

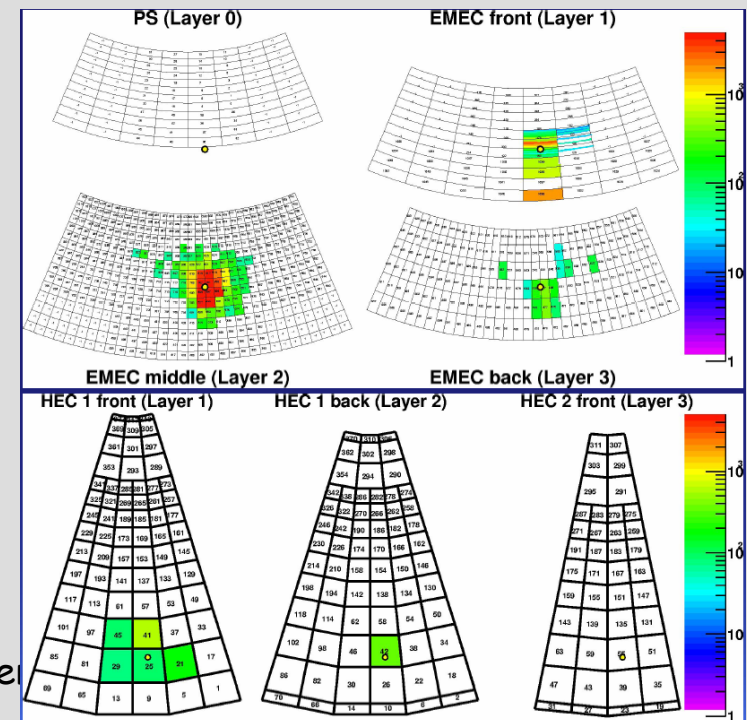
Two objects may be used as input for the jet reconstruction algorithms:

- Calorimeter Towers of dimension $\Delta\eta \times \Delta\phi = 0.1 \times 0.1$
- 3D Energy blobs – Topological Clusters: 3 levels of E/σ thresholds are used for **seed cells – neighbour cells - final expansion (4-2-0)**
- TopoClusters give a better noise suppression

Cell sigma noise



Cluster for 120 GeV pion in
EMEC and HEC (2002 Test Beam data)



Jet reconstruction algorithms

Iterative cone

$$\Delta R = \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$$

	ΔR	$E_{T,seed}$	ΔS
ATLAS	0.7/0.4	2 GeV	50%
CMS	0.5	2 GeV	-

Split and merge procedure (ΔS)
for ATLAS

CMS exclude jet algorithm input
objects as soon as they belong to
a built jet

E Recombination scheme: 4 vector
sum of object components to obtain
jet kinematics

Cone algorithm is the most
widely used for physics
analysis in ATLAS up to
recently: cone 0.4 for top
physics and cone 0.7 for all
other analysis.

Calibrations are tuned on
cone 0.7 jets.

In last year a lot of activity
has started also on other jet
algorithms.

Jet reconstruction algorithms

Midpoint Algorithm

Implementation based on CDF approach.

0 Seed $E_t > 2 \text{ GeV}$

1 Cone precluster with radius $0.5 \times \Delta R$

2 Add midpoints if preclusters i, j are separated $< 2 \times \Delta R$

3 Cone jets of radius ΔR are searched

4 Merge if $> 50\%$ of p_T of lowest jet is shared, else split

K_T algorithm

Preclusters objects with $\Delta R < 0.2$

For each object i of transverse impulse k_T calculate:

$$d = \begin{cases} d_{ii} = k_{T,i}^2 \\ d_{ij} = \min(k_{T,i}^2, k_{T,j}^2) \frac{\Delta R_{ij}^2}{D^2} \end{cases}$$

If $d_{\min} = d_{ii} \Rightarrow$ jet

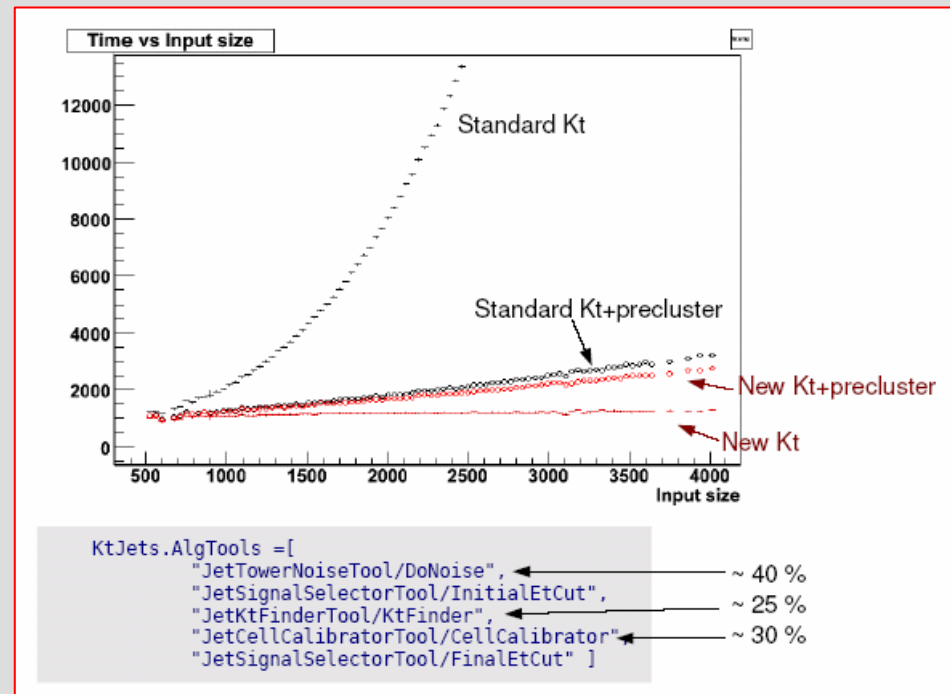
else if $d_{\min} = d_{ij} \Rightarrow ij$ (4-vector sum) in a new d_{ii}

$D = 1$ is default – more later

E Recombination scheme: 4 vector sum of object components to obtain jet kinematics

Fast K_T

- KT algorithms are typically slow since speed scales with $O(N^3)$
- It has been shown that they can be made faster by using nearest neighbour information (Cacciari, Salam hep-ph/0512210)
- FastKT has been implemented in ATLAS and it also allows to skip the preclustering phase.



Choice of jet algorithm and parameters

What is the best way to choose a jet algorithm for my analysis and the value of jet algorithm parameters ?

Some examples from what has been done in ATLAS.

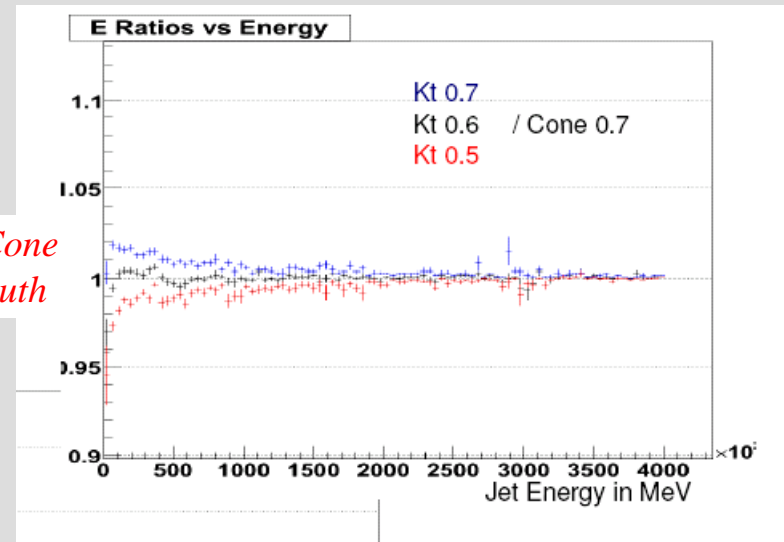
Look at MC particles jets reconstructed with Cone and Kt and compare the reconstructed energy.

With this “recepty” we obtain:

D = 0.6 for Cone 0.7

D = 0.3 for Cone 0.4

$$E_{truth}^{KT} / E_{truth}^{Cone}$$

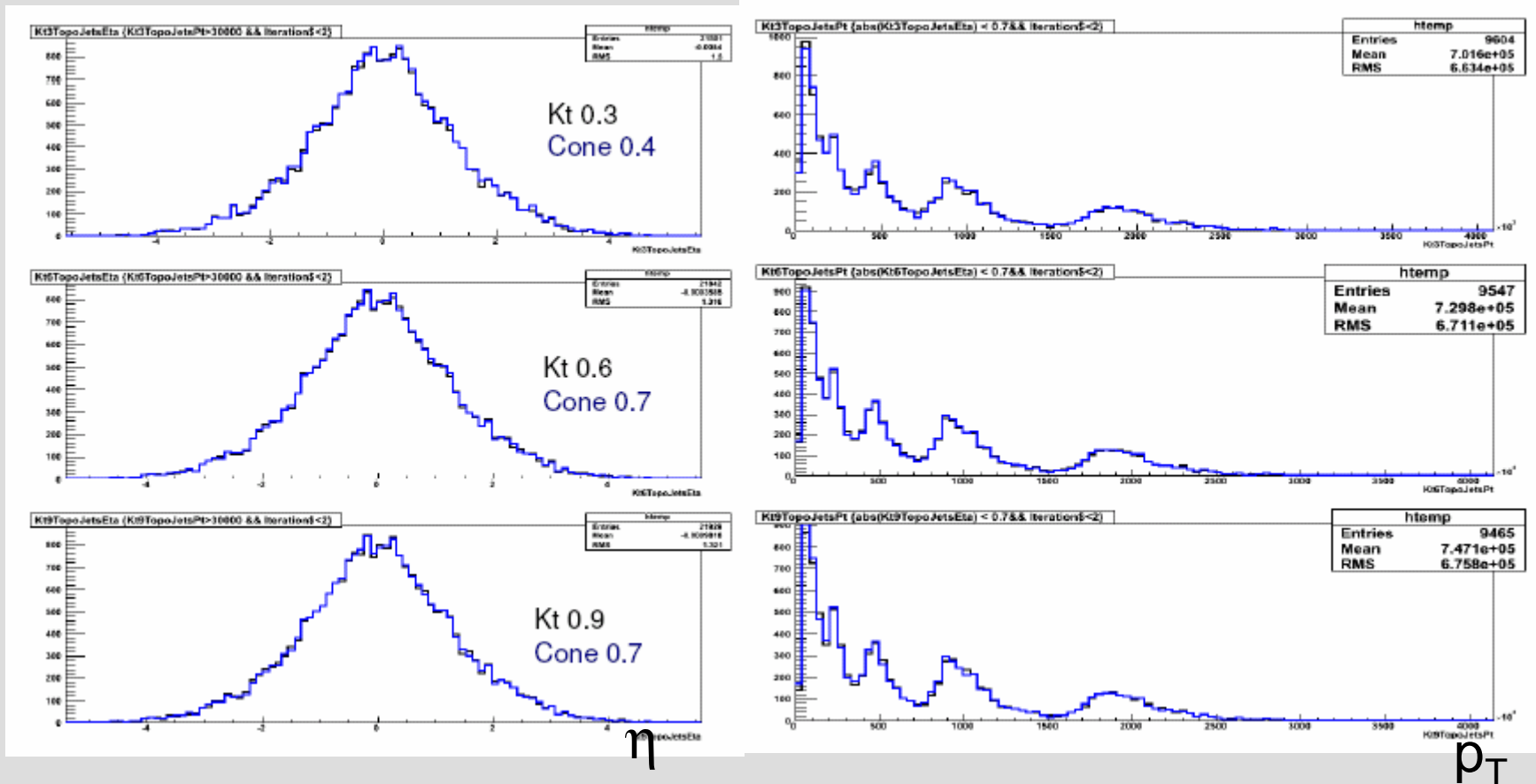


In Ellis-Soper article (PRD 48, 3160–3166 (1993)) in order to have the same inclusive one-jet cross section dependance on renormalization and factorization scale for cone KT indicates $D=1.35 \times \Delta R$

Are these two results inconsistent ?

Cone vs K_T

Looking at the 2 highest pt jet

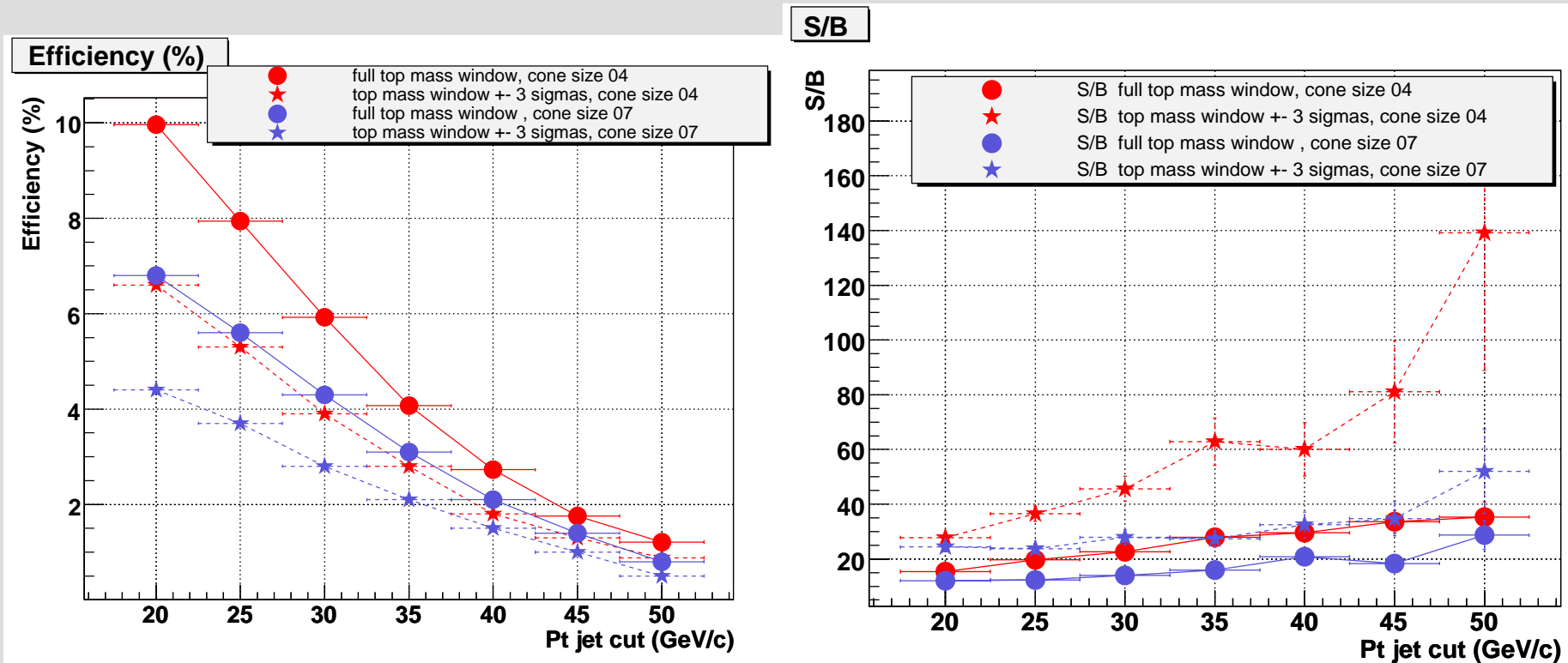


Tuning cone algorithm on top events

Studied for top mass measurement in $l\nu b j\bar{b}$ channel

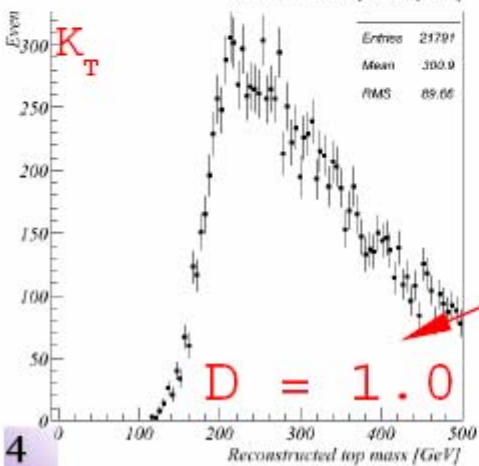
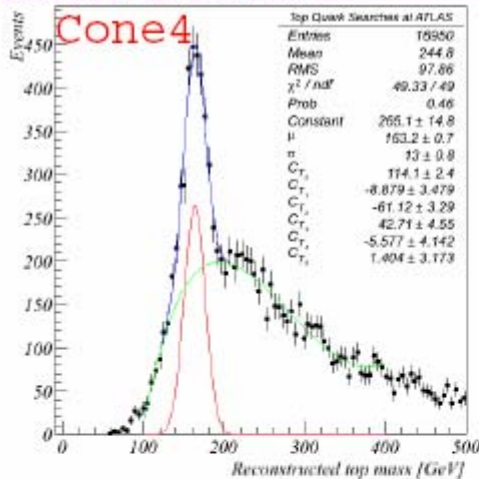
Looked at efficiency, S/B, W and top « purity » (correct jets taken)

→ Cone 0.4 (much) better than cone 0.7



Tuning KT algorithm on top data

Introduction



K_T ALGORITHM (inclusive)

Ellis & Soper, PRD48, 3160, (1993)

Start from a set of objects:

$$\{p_1, p_2, \dots, p_i, p_j, \dots, p_n\}$$

For each object pair ij :

Define d_i and d_{ij} in the ΔR scheme:

$$d_i = (p_{T_i})^2$$

$$d_{ij} = \text{Min}((p_{T_i})^2, (p_{T_j})^2) * \Delta R_{ij}^2 / \mathbf{D}^2$$

$$\Delta R_{ij}^2 = (\phi_i - \phi_j)^2 + (\eta_i - \eta_j)^2$$

if $d_{ij} < d_i$, merge objects i and j :

by several recombination schemes

$$(E, p_t, p_t^2, E_t \text{ or } E_t^2).$$

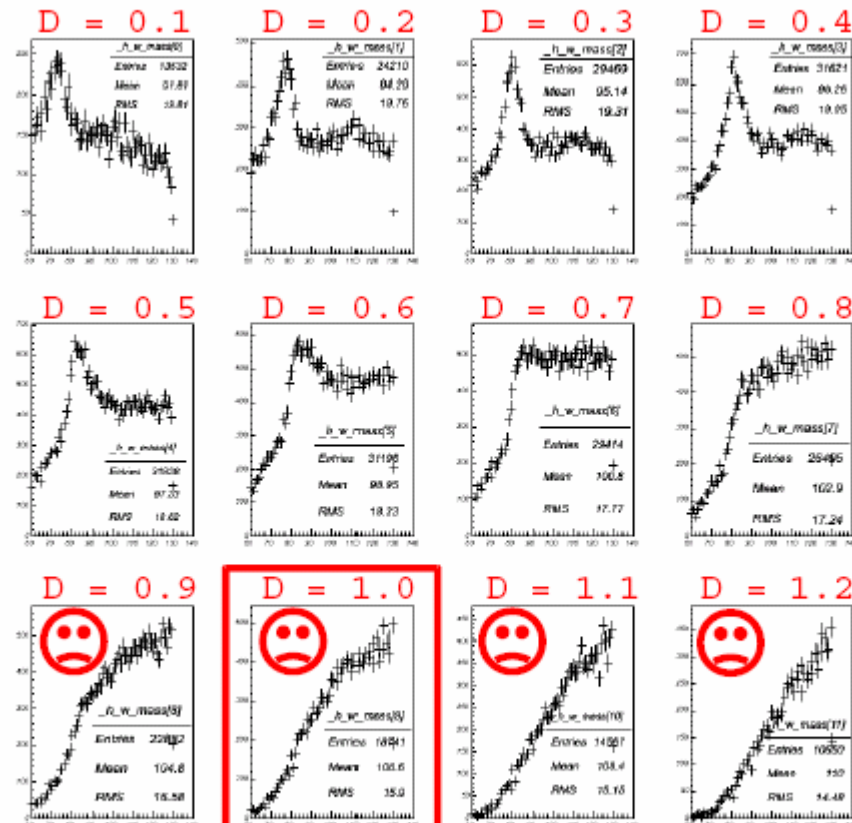
For instance, the E scheme:

$$p_{ij} = p_i + p_j$$

if $d_i < d_{ij}$, object i is a jet

Tuning the KT algorithm

W mass reconstruction: e.g. ΔR , E scheme [Default]

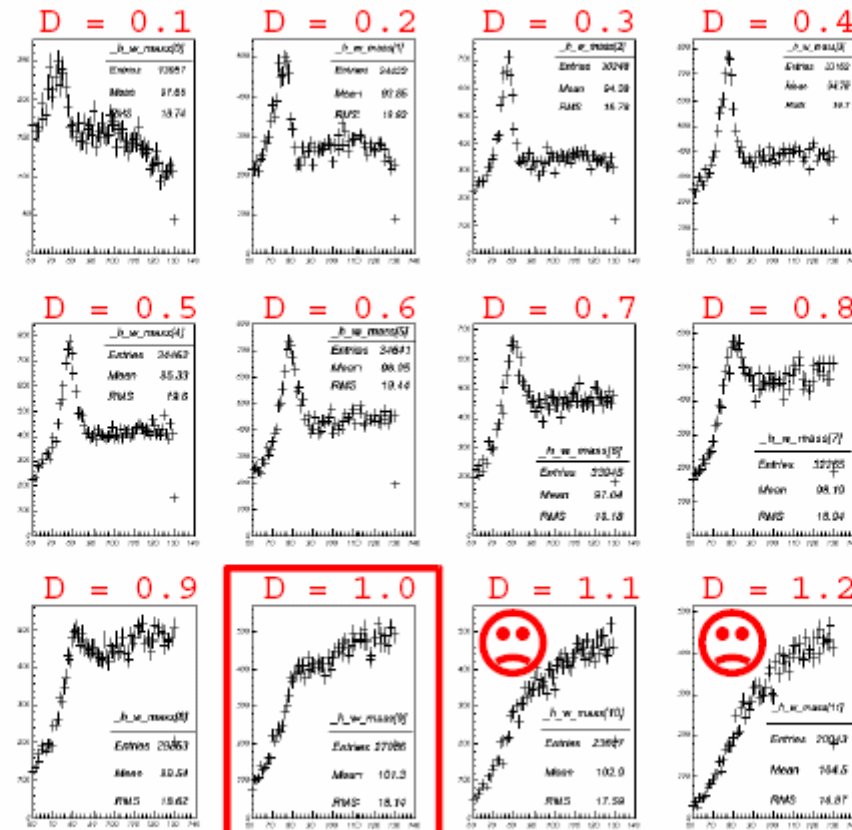


Increasing D , the reconstructed W shifts to larger mass values, but D should not be too large...

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Which recombination scheme...

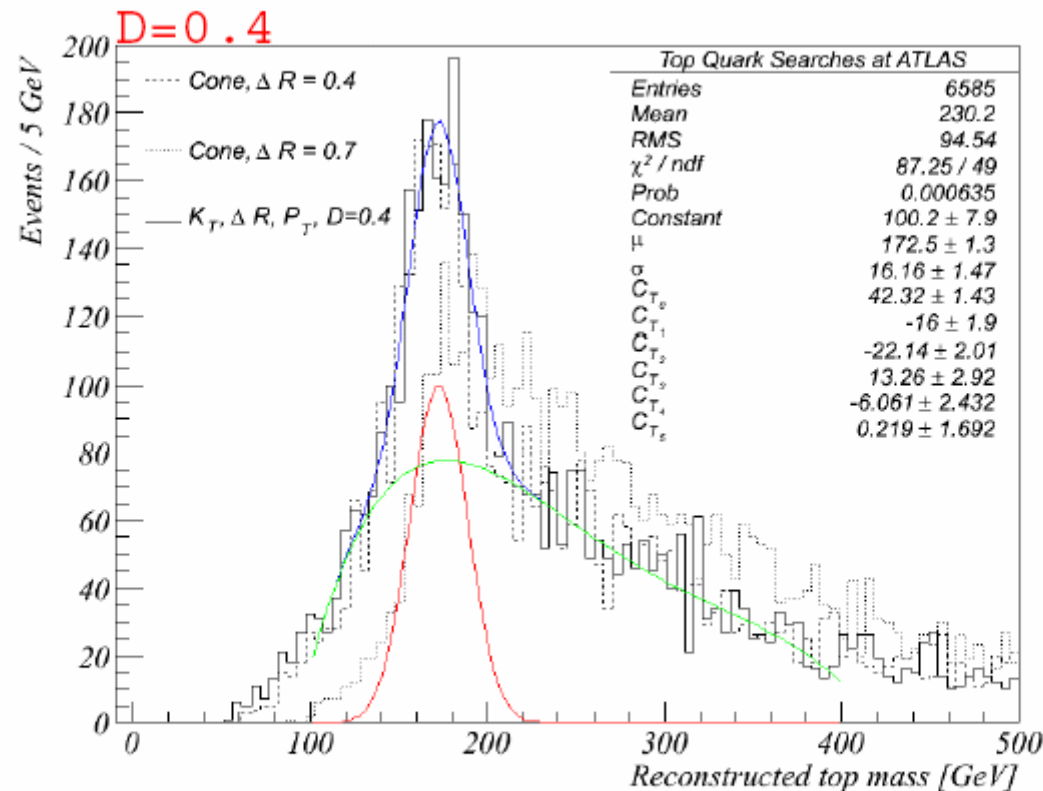
W mass reconstruction: e.g. ΔR , P_T scheme



13

W mass with tuned KT

T1-4100: full simulation using $\Delta R, P_T$ scheme



Clearly, we improve the top mass peak reconstruction with a reasonable **D** value of the K_T algorithm.

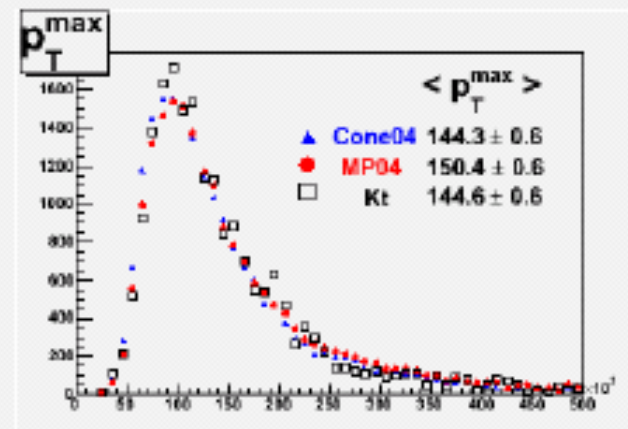
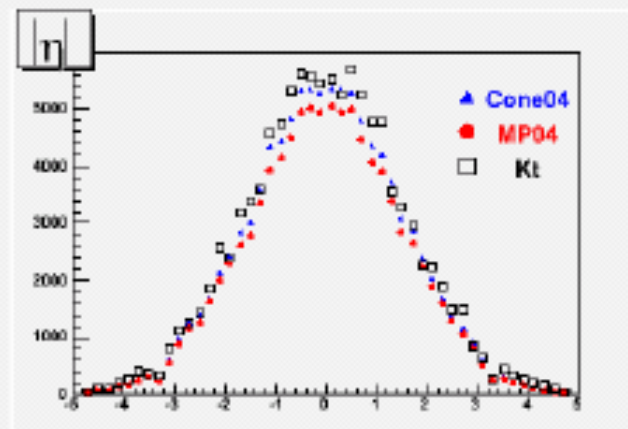
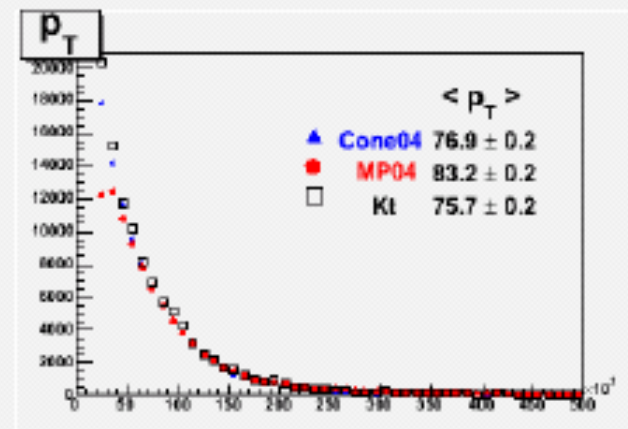
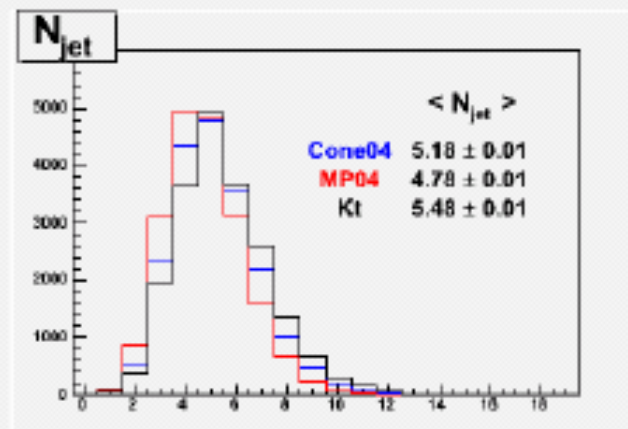
18

Midpoint first comparison

Cone 0.4

Midpoint 0.4

KT 0.45



- Excess of low- p_T jets ($20 \text{ GeV} < p_T < 50 \text{ GeV}$) for Cone and Kt
- MidPoint tends to merge low- p_T jets

Choosing and tuning

Various methods to tune the favorite algorithm

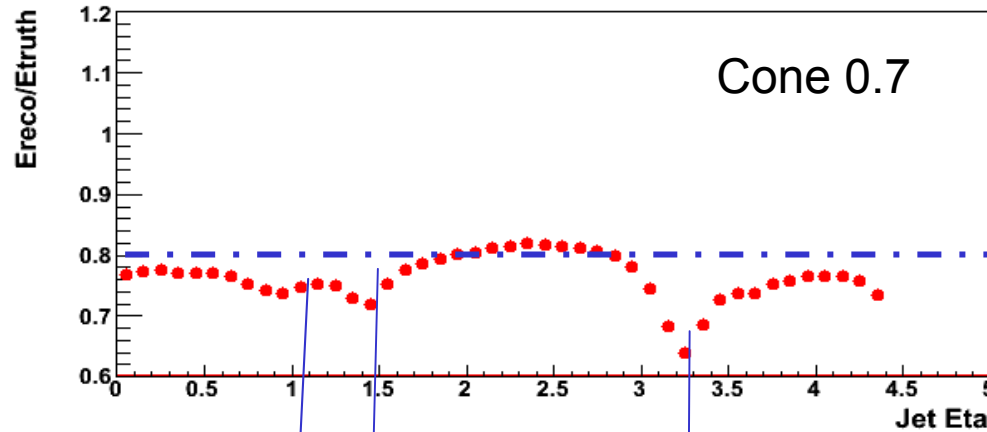
Proliferation of jet collections, we need to find a well defined path to make the choice and to define the parameter tuning

Calibration and UE, pileup subtraction will have to be understood in detail and it will not be straightforward to generalize it to any clustering algorithm

So a deciding on benchmark jet clustering would help a lot since

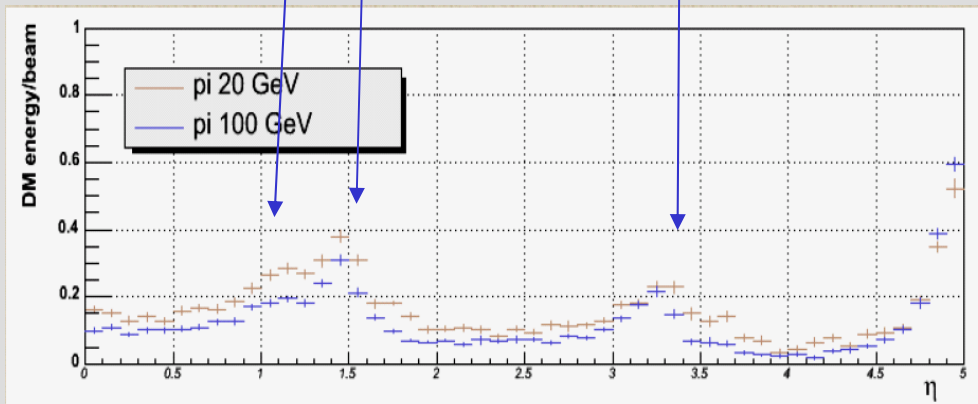
Now we go to calibration ...

Uncalibrated reconstructed Jets: detector effects at work...

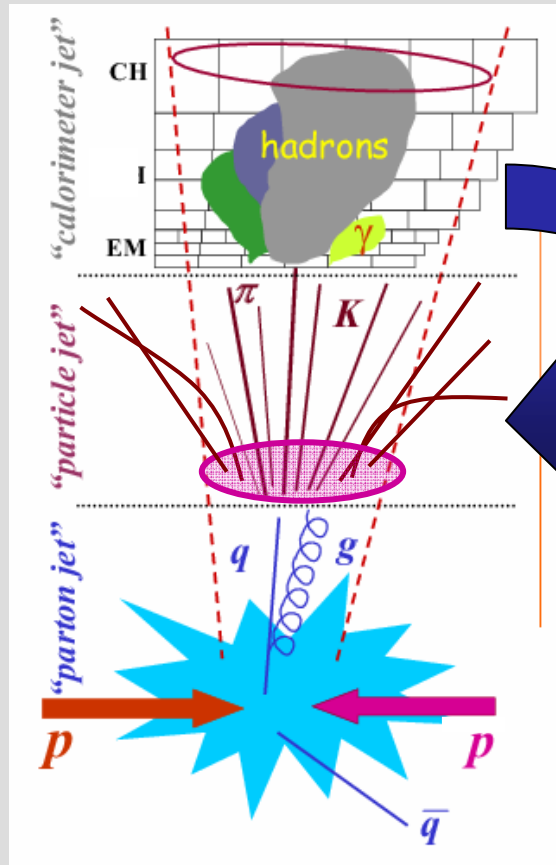


Detector effects are clearly seen on $E_{\text{reconstructed}}/E_{\text{truth}}$:

Cracks, e/h, B field
(tracks with $p_T < 350$ MeV do not reach the calorimeter)



Phase 2: calibrating to the particle jet



Phase 2 : the detector effects are corrected calibrating the reconstructed jet to the particle jet

Truth:

- Particle Jets are reconstructed applying jet algorithm to stable particles (excluding neutrinos and muons) and are matched to Calorimeter Jet (ATLAS)

- All particle falling in the angular region of calorimeter jet (only for cone) (ATLAS)

Both truths contain the particles swept off from B field:

ATLAS $p_T < 350$ MeV



Calibrating to Particle Jet

2 step procedure

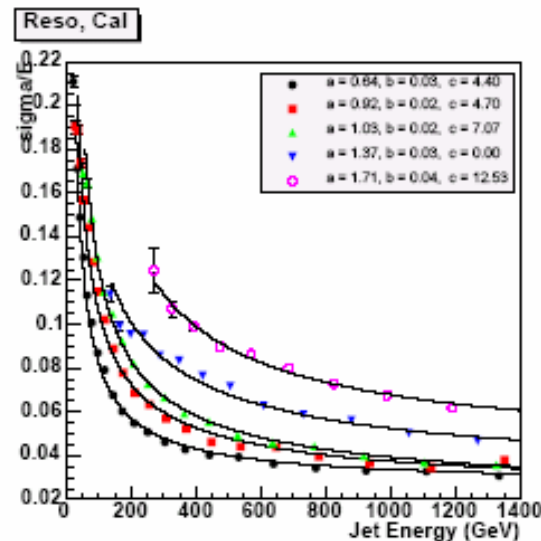
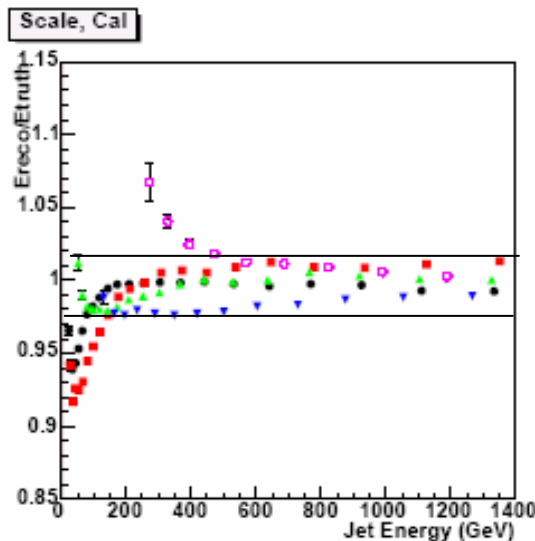
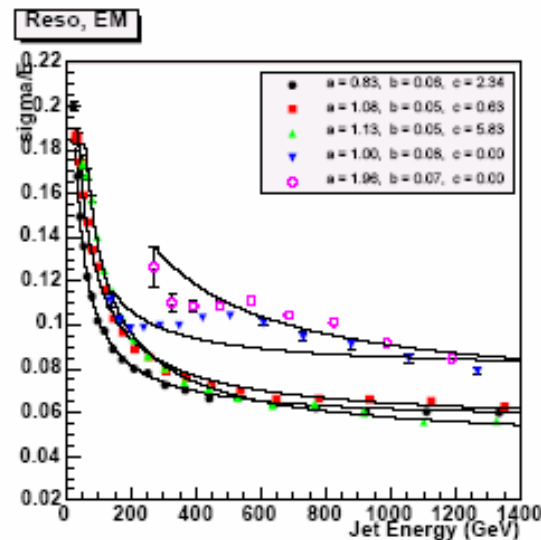
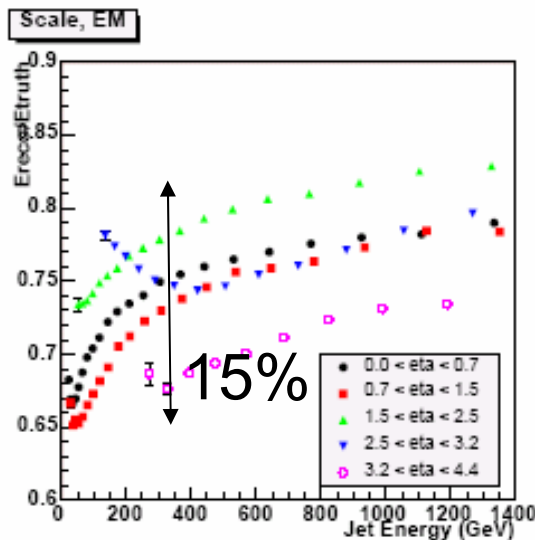
1. Calibrated energy is calculated as:

$$\begin{aligned} E_{Raw} &= \sum_s E_{cell_s} \\ E_{Rec} &= \sum_s w(E_{cell}, CellPosition) E_{cell_s} \end{aligned} \quad \left. \vphantom{\sum_s} \right\} \begin{array}{l} \text{Cell} \\ \text{weighting} \end{array}$$

the $w(E_{cell}, CellPosition)$ coefficients are obtained by minimizing the energy resolution to the MC truth with the linearity constraint. Same weights are used for different algorithms.

2. A factor $R(E_T, \eta) = E_{Trec}/E_{TMC}$ is applied to correct for residual non linearities and for algorithm effects.

From the Calorimeter jet to the Particle Jet



QCD – 0.7 cone jet built from calorimeter towers

Uncalibrated $|\eta| < 0.7$

$$\frac{\sigma(E)}{E} = \frac{0.83}{\sqrt{E(\text{GeV})}} \oplus 0.05 \oplus \frac{2.3}{E(\text{GeV})}$$

Calibrated $|\eta| < 0.7$

$$\frac{\sigma(E)}{E} = \frac{0.67}{\sqrt{E(\text{GeV})}} \oplus 0.02 \oplus \frac{4.4}{E(\text{GeV})}$$

Linearity $\pm 2\%$ $E_T > 20$ GeV

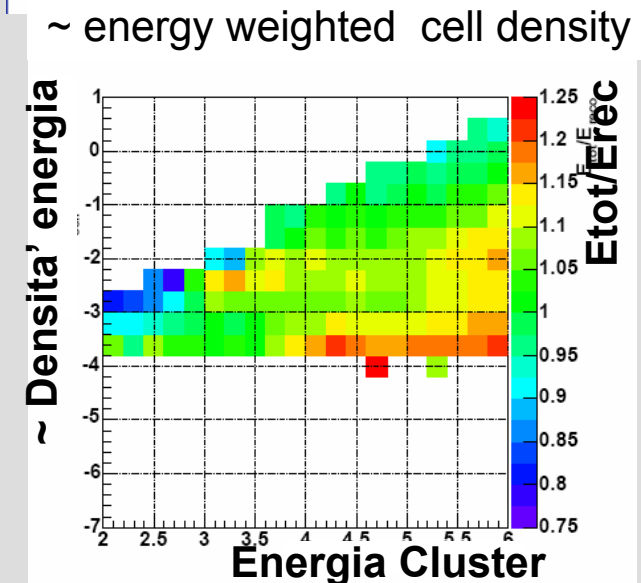
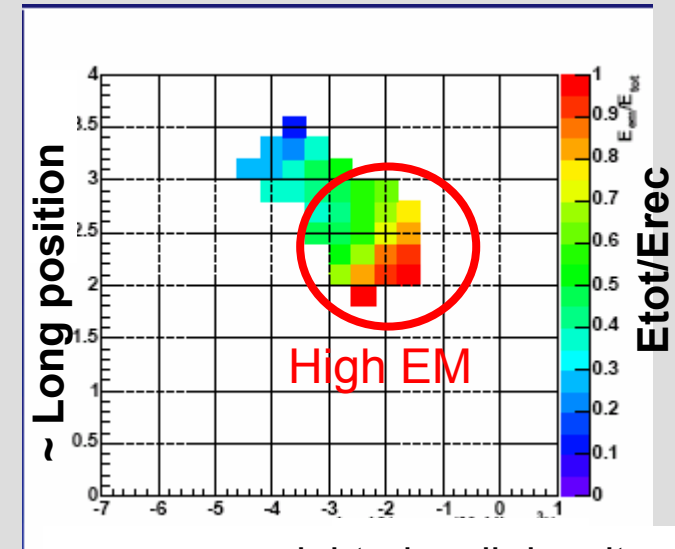
However changing the jet algorithms requires new tuning of calibration....what we are trying to do to generalize the calibration approach

Local hadron calibration

The aim is to Calibrate the TopoClusters before reconstructing the jets. The calibration is based on MC information: for each cell EM energy, Escaped energy, Invisible energy, Non EM energy.

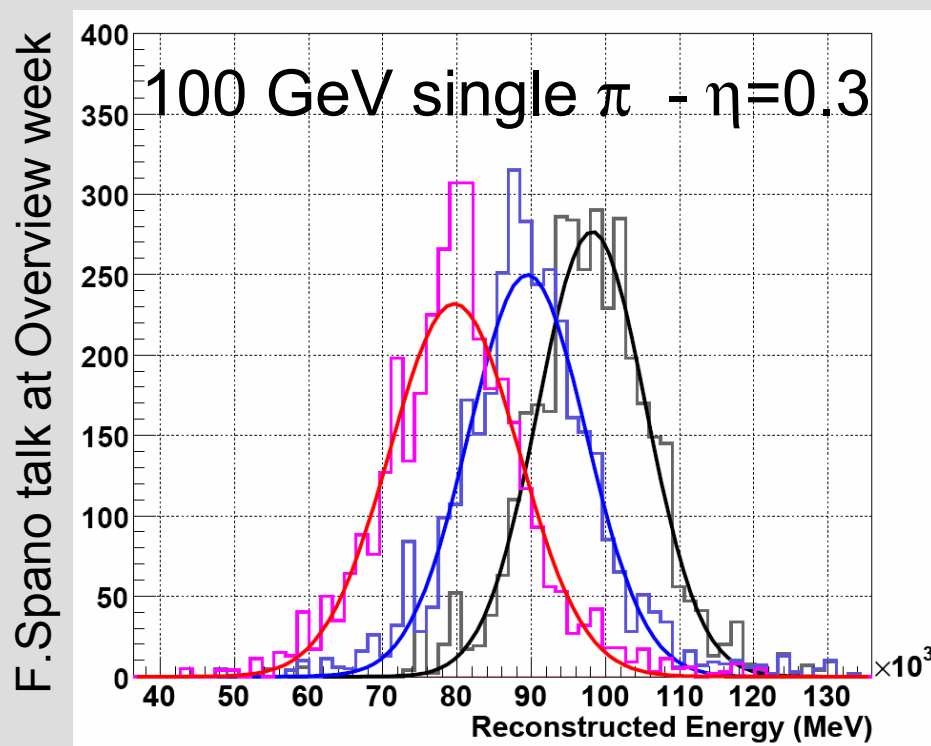
This information are used to:

- 1 – **Classify** the calorimetric deposit: EM not weighted, NonEM to be weighted
- 2 – **Calbration**: topoclusters characteristics are used to define the calibration weights
- 3 – **Dead Material Correction**



Local Hadron Calibration

Example of how the three step works on single pions: TopoCluster classification (EM or not); weighting to correct e/h; dead material correction



EM scale

Classified + Weighted

Classified +Weighted +Dead
Material corr.

Linearity is recovered and
resolution is comparable to what
we obtain with different methods

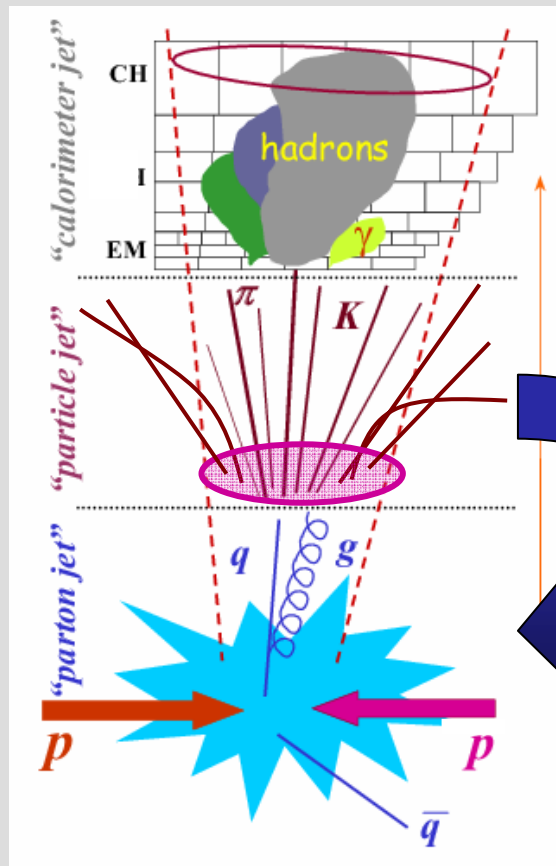
Local calibration

Local calibration would allow to:

- better understand the contribution of each factor to calibration, it relies on a deep understanding of our data
- Start jet reconstruction at a calibrated scale
- Obtain reconstructed and calibrated jets with any algorithm. Calibration would include the corrections for e/h and dead material
- Corrections for B swept tracks, out of cone ... would be added on top of this calibration

Local calibration is being developed and it still needs testing and validation before being used for physics analysis.

Phase 3: back to the parton energy



Phase 3: absolute energy measurement of parton energy. Goal precision 1%....

- correct for energy losses out of jet clustering
- correct for energy physics effect such as: underlying event, ISR, FSR

Back to the parton energy

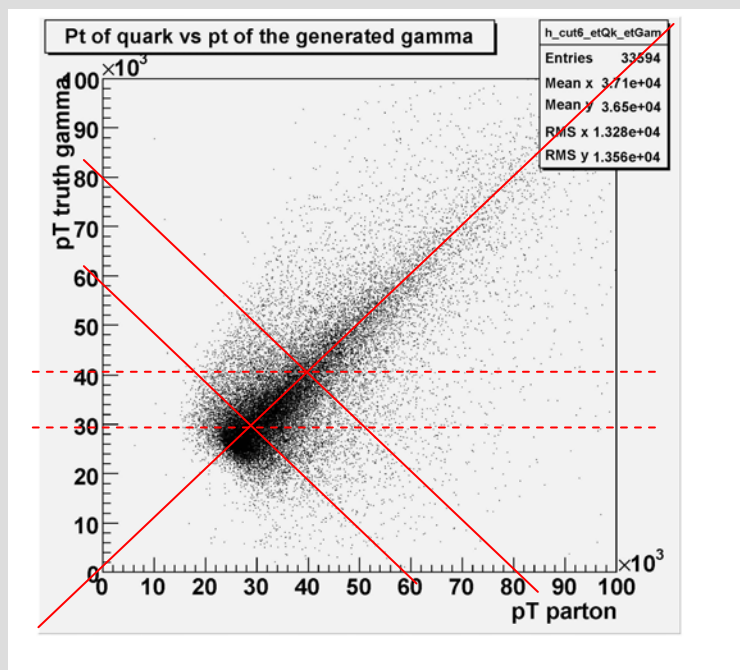
From data sample (in-situ):

- $W \rightarrow jj$: imposing W mass. Maximum Energy ≈ 200 GeV, jet overlapping. Events from $t\bar{t}$ with 1 lepton used to trigger.
- $Z(e^+e^-, \mu^+\mu^-) + \text{jet}$: p_T balance or E_{miss} projection method. Useable for light and b jets, about 5% of the total event rate. p_T range ~ 40 -400 GeV.
- $\gamma+j$: p_T balance or E_{miss} projection method. Higher statistics but high QCD background. More on next slides...

Calibration using data: Gamma + jet

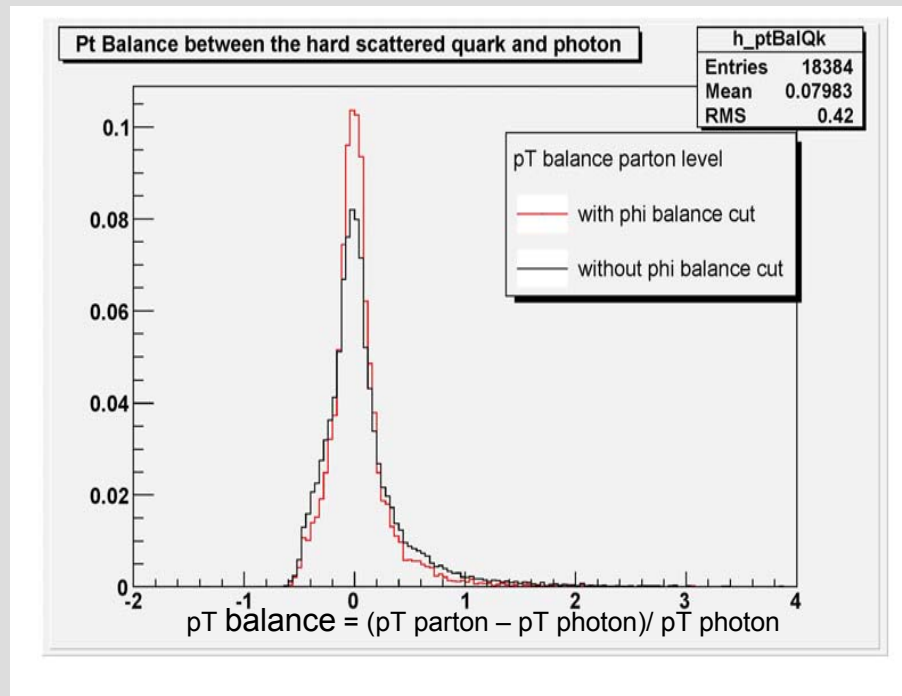
Direct photon production: $qg \rightarrow q\gamma$ (90%) $q\bar{q} \rightarrow g\gamma$ (10%)

$$p_{T, \text{Jet}} = p_{T, \text{Gamma}} \rightarrow k_{\text{jet}} = p_{T, \text{Jet}} / p_{T, \gamma}$$



Gamma selection isolation & $E_T > 30$ GeV
Select Highest p_T jet
apply phi back-to-back cut $\Delta\phi > 175^\circ$

Calibration with Gamma+jet events



$$p_T balance = \frac{p_T Jet - p_T Photon}{p_T Photon}$$

Fit peak region iterating a gaussian fit
between $\pm\sigma$ around the most probable value

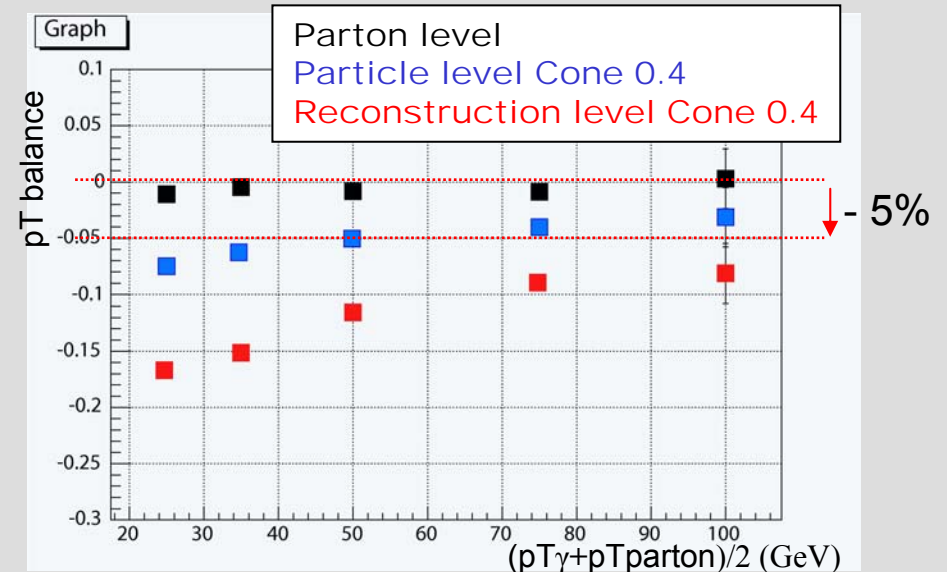
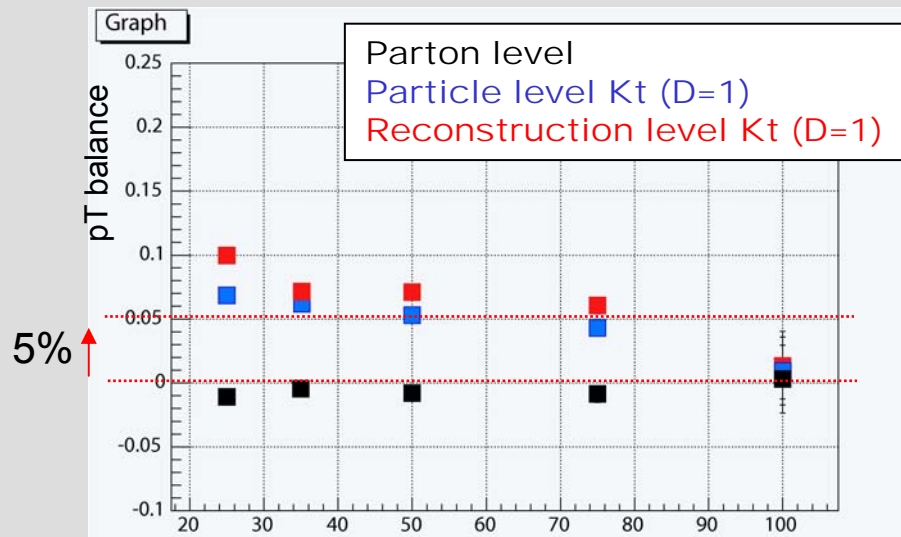
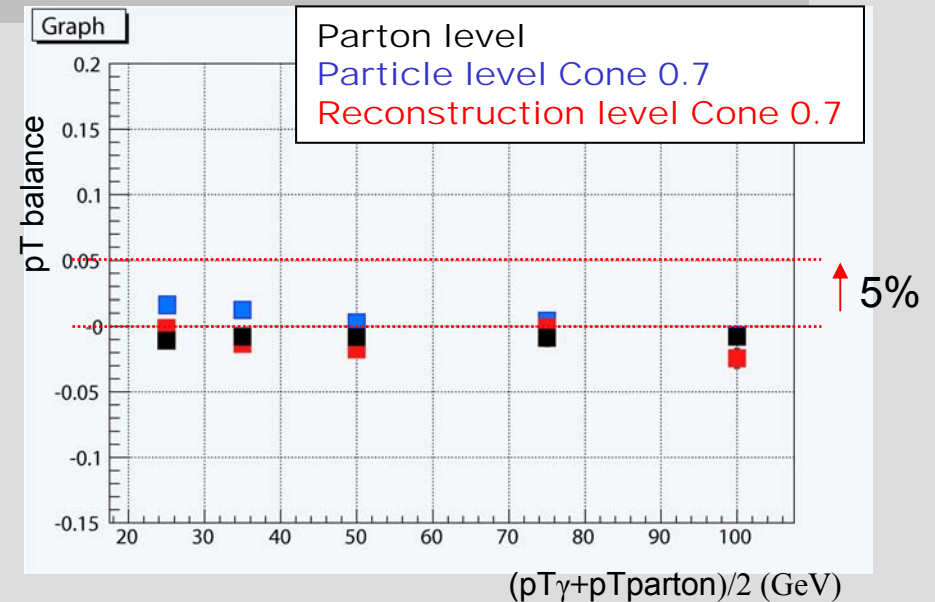
Calibration using data: Gamma + jet

Comparing balance at reconstruction, MC jet and parton level gives indication on:

1. calibration biases

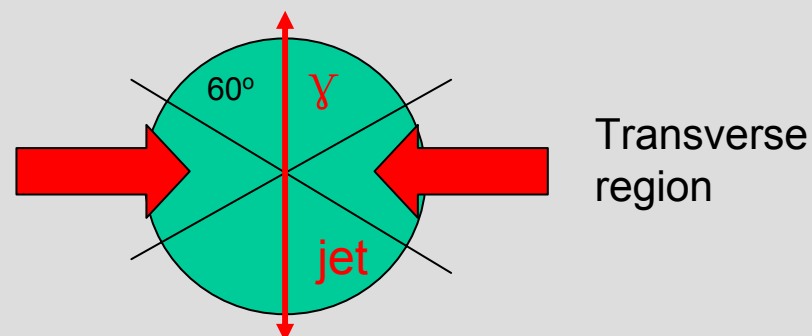
2. contribution of UE event, contribution of Out of cone energy, 3. effect of ISR contribution.

More work needed to disentangle UE from Out of cone. Work is in progress.

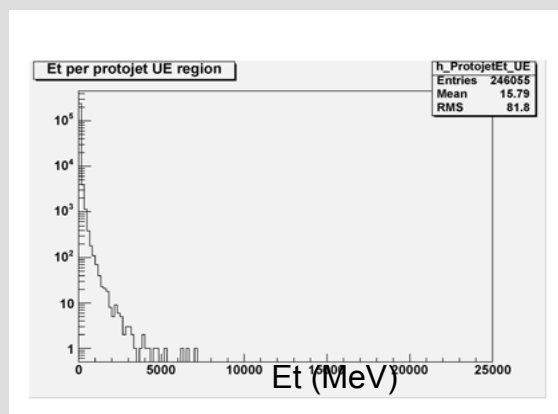


Estimating the UE event contribution

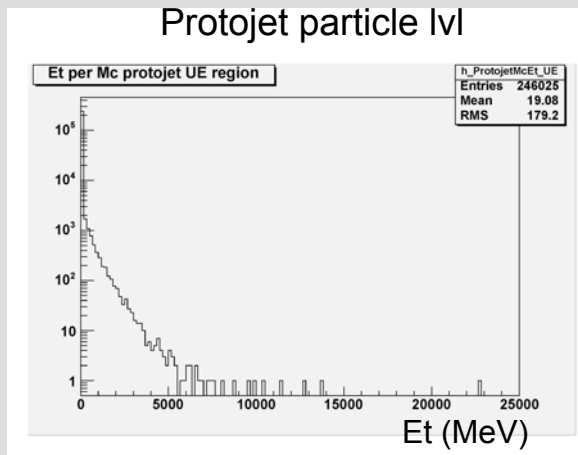
To find the mean Et for UE, we consider the transverse region of the event: avoiding 60 degrees on both sides of photon and jet



Protojet recon lvl



Protojet particle lvl



Considering the number of jet components we estimate UE contribution to reco jets:

Cone 0.4 - 2.7 GeV

Cone 0.7 – 1.2 GeV

KT (D=1) – 5.1 GeV

Very preliminar results ...
work in progress.

Mean transverse energy

Recon protojet	15.8 ± 0.2 MeV
Particle protojet	19.1 ± 0.4 MeV

A candle for jet energy: top mass

Top at detector commissioning: no b tagging, only calibration for detector effects. Realistic scenario for first data taking phase.

Search for $t\bar{t}$ → l nubbjj.

Selection cut:

$E_{\text{miss}} > 20$ GeV

1 lepton $P_T > 20$ GeV

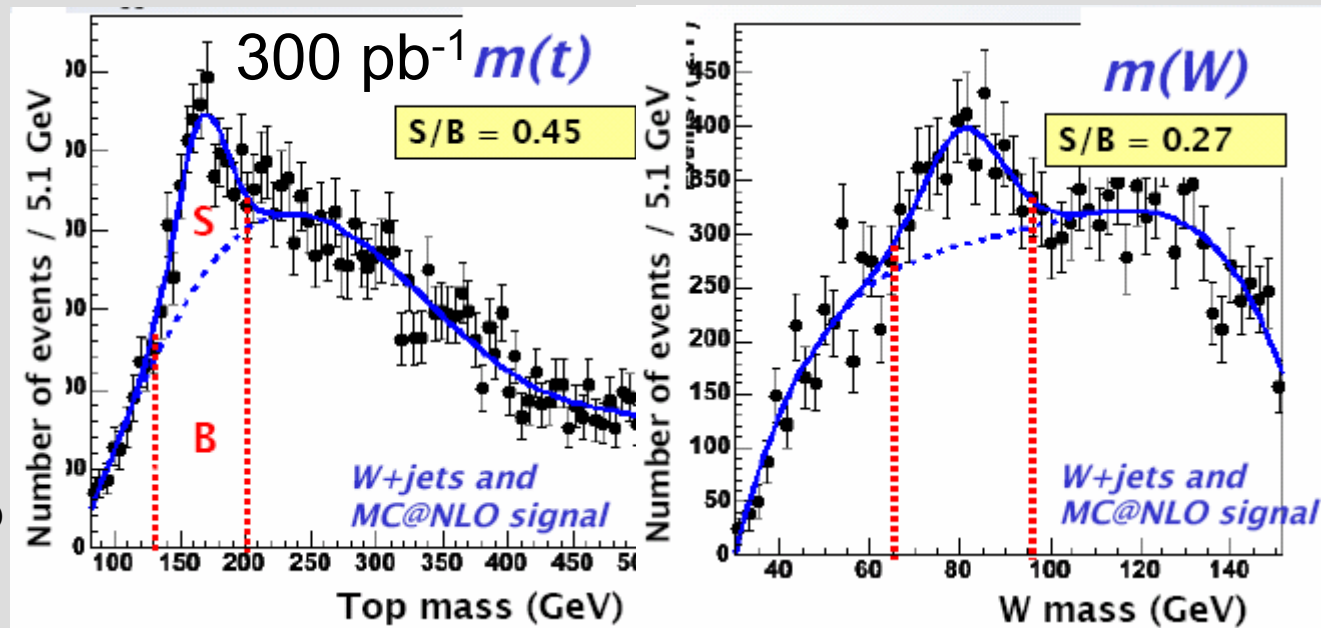
4 jet $P_T > 40$ GeV

Reconstruction:

$W \rightarrow jj$ selection: 2 highest p_T jet in CM of jjj

Third jet giving highest p_T top

W+4j background
increased by three
times to consider
W+3/5j and uncertainty



Peak and width can be used to understand
MC/data agreement and calibration
performance

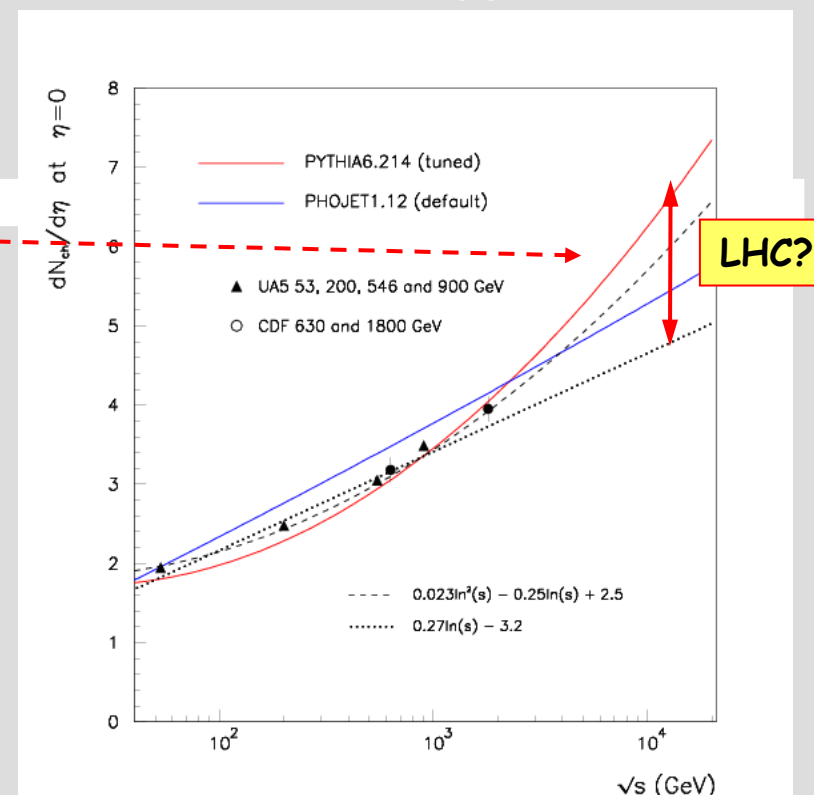
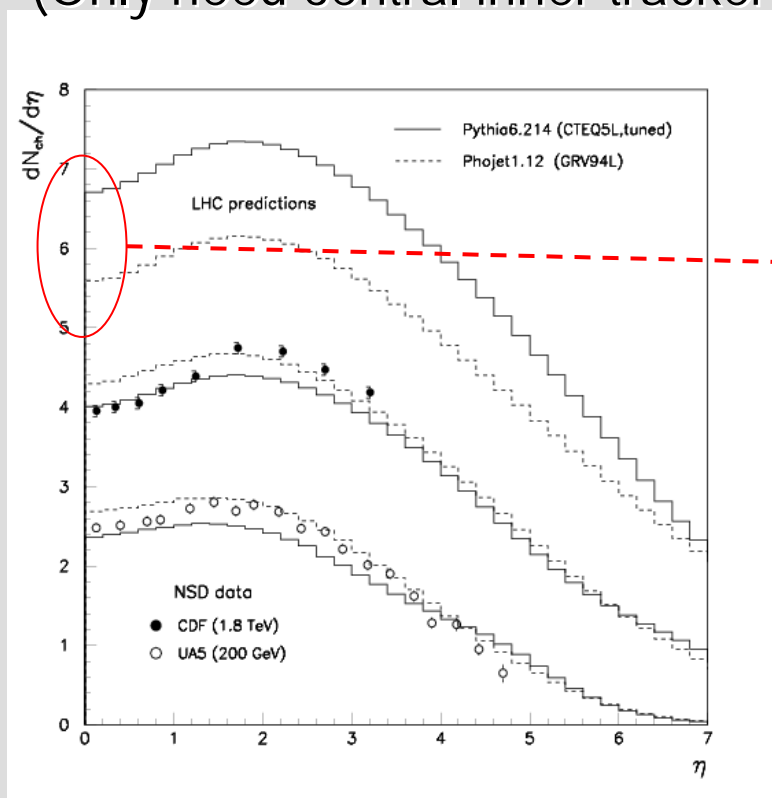
Summary

- Cone, KT and midpoint algorithms have been implemented
- They can take as input MC particles, CaloTowers, TopoClusters
- Tuning of various algorithm to various analysis needs is under going
- Calibration to correct for detector effects has been developed and possible alternatives are being studied
- Study on UE and pileup subtraction are just starting (re-starting)
- My personal opinion is that we should at the beginning concentrate on a well defined and justified jet algorithm in order to understand all the issue about jet reconstruction and energy scale and than we can move to more general scheme.
- Which is the best way to define the jet algorithm to become the benchmark has to be understood.

Back up slides

Charged particle density at $\eta = 0$

(Only need central inner tracker and a few thousand pp events)



Multiple interaction model in PHOJET predicts a $\ln(s)$ rise in energy dependence. PYTHIA suggests a rise dominated by the

$\ln^2(s)$ term.

MoLHC CERN July 20 2006

C.Roda - INFN & University of Pisa

The Underlying Event in jet physics

- The underlying event in charged jet evolution:

Phys. Rev. D, 65 092002 (2002)

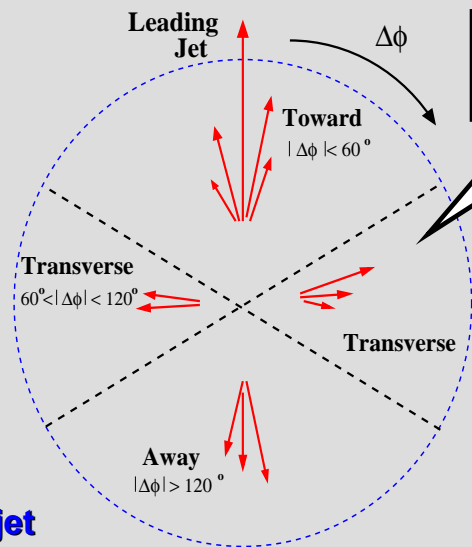
CDF analysis:

- charged particles:
 $p_t > 0.5 \text{ GeV}$ and $|\eta| < 1$

- cone jet finder:

$$R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.7$$

$$\Delta\phi = \phi - \phi_{\text{jet}}$$



UE is defined as the
Transverse Region

➤ The underlying event is defined as *everything in the collision except the hard process*.

➤ *It is not* a minimum bias event!

➤ The underlying event has *hard* (multiple “semi-hard” parton scatterings) and *soft* components (beam-beam remnants).

LHC predictions: pp collisions at $\sqrt{s} = 14$ TeV

ATL-PHYS-PUB-2005-007

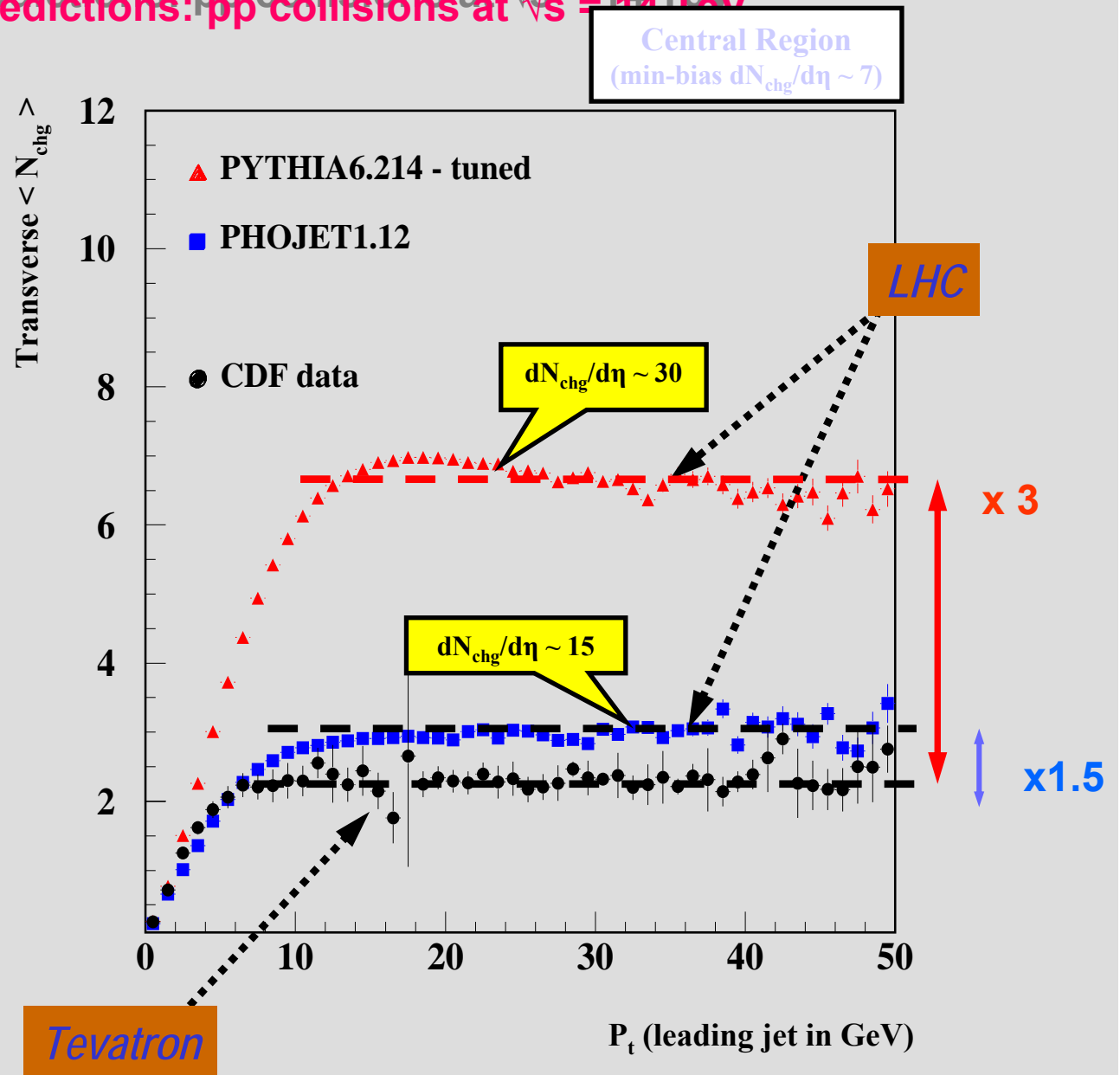
Charged particles:

$p_t > 0.5$ GeV and $|\eta| < 1$

Cone jet finder:

$$R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.7$$

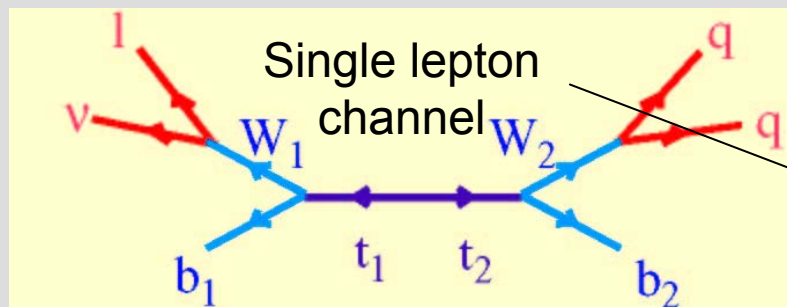
UE particles come from region transverse to the leading jet.



Systematic on m_{top}

$10^{33} \text{cm}^{-2} \text{sec}^{-1}$

Events in top mass window

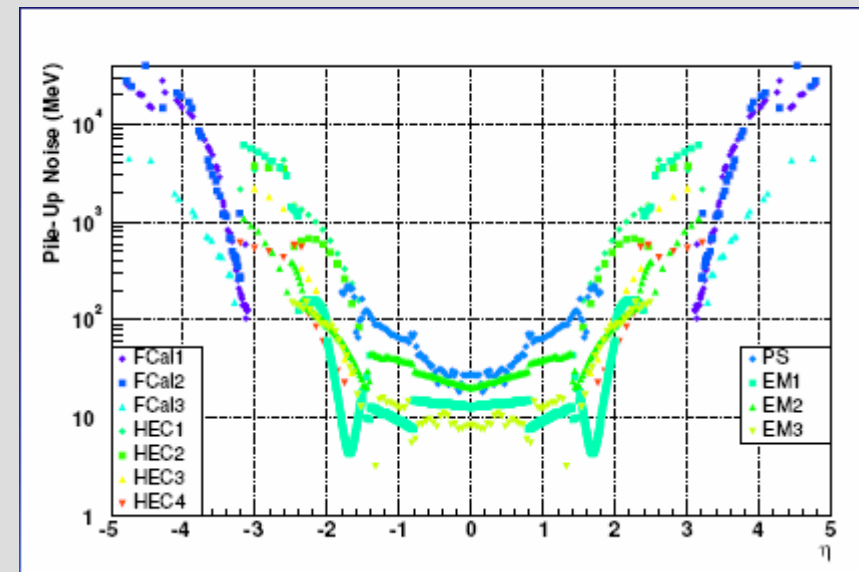
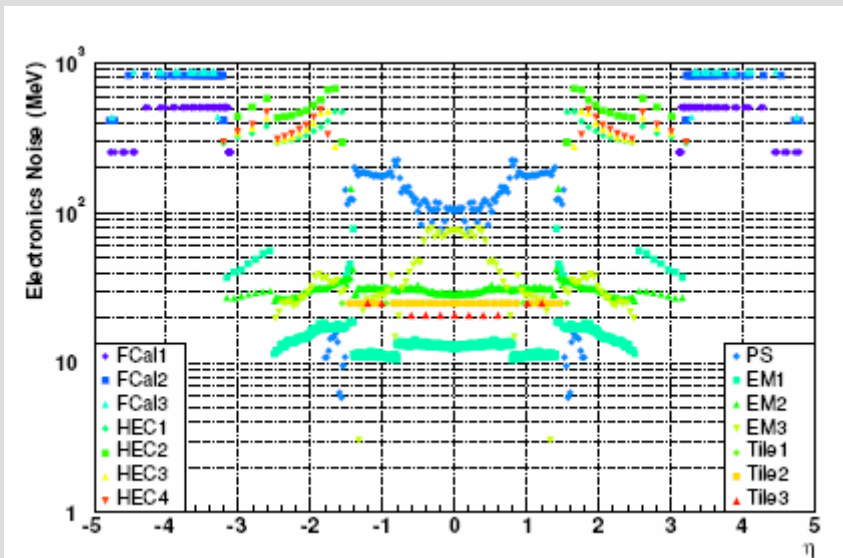


<u>Period</u>	<u>evts</u>	<u>$dM_{\text{top}}(\text{stat})$</u>
1 year	3×10^5	0.1 GeV
1 week	1.9×10^3	0.4 GeV

Systematic error on top mass from light and bjet energy scale uncertainty

jet unc.	$\Delta M_{\text{top}}(\text{jet})$	$\Delta M_{\text{top}}(\text{bjet})$
1%	0.9 GeV	0.7 GeV
5%	11 GeV	3.5 GeV

Noise level

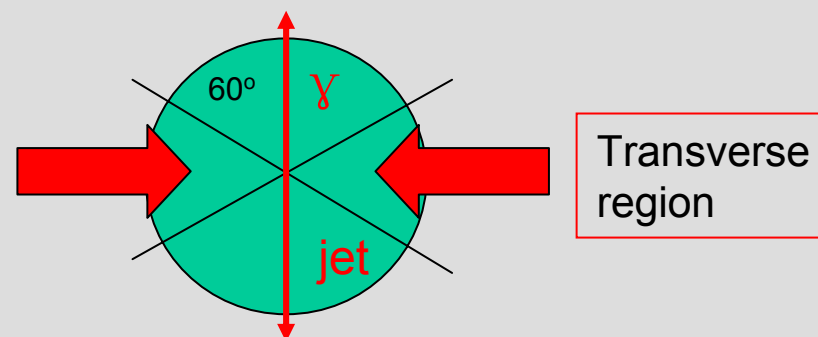




Gamma + jet and Underlying event

Try to measure the mean ET of UE from the event sample

Select the “transverse region” of the event: avoiding 60 degrees in Phi around both photon and the jet.



Mean transverse energy per $\eta \times \phi = 0.1 \times 0.1$

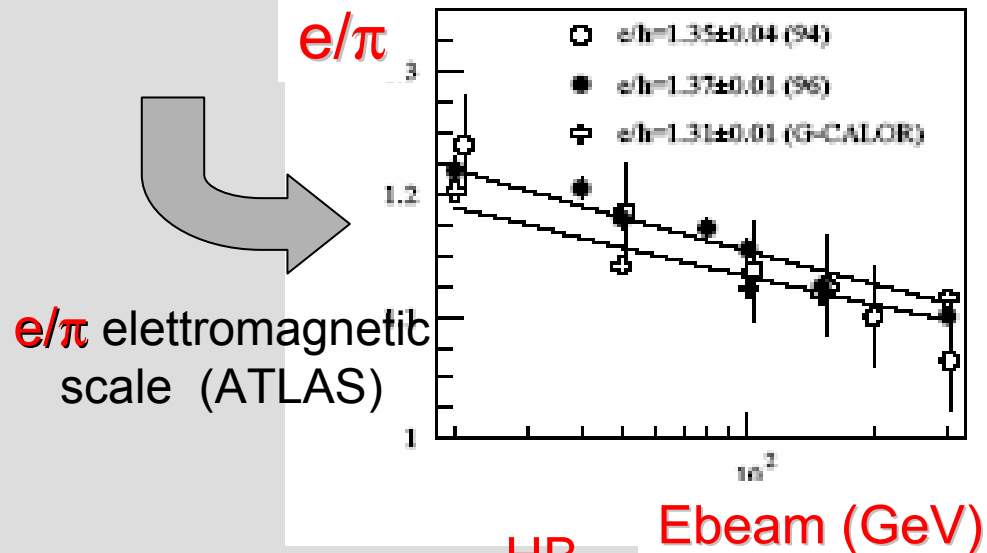
Tower (RMS of el.noise ~140 MeV)	15.1 ± 0.2 MeV
Particle protojet (Σ particles per tower)	19.1 ± 0.4 MeV

3 GeV in cone 0.7

Average UE level
~10% RMS of el.noise
(very sensitive to noise suppression)

Subtraction algorithm and biases introduced are under study.

Compensation



CMS $e/h = 1.4$
 ATLAS $e/h = 1.36/1.5$ } Had. Cal.
 $\Delta S = 15\%-12\%$ per $E\pi$ 20-300 GeV

$$E_{rec} = E_{EM} + (\alpha \times H1 + H2 + H3)$$

A high signal H1 indicates hadronic signal thus the EM scale is too low and $\alpha > 1$ corrects for $e/h < 1$. The correction may or not depend on energy.

➤ benchmark

➤ H1 [NIM-A1809(1981)429] :

$$E_{rec} = \sum W_{EM}(E_{cell}, E_{part}) E_{cell} +$$

$$\sum W_{HAD}(E_{cell}, E_{part}) E_{cell}$$

W obtained by minimizing resolution



Noise suppression performance

— Zero suppression
- - - 2 sigma symmetric
- - - CaloTopoCluster

Plots show how much negative energy is left in jets after noise cancellation/subtraction algorithm is applied in each calorimeter region.

Topological Cluster < 2 sigma symmetric < Zero suppression

