

Hadron Level Cross-Sections for W + Jets Production at CDF



Ben Cooper



MICHIGAN STATE
UNIVERSITY

MC4LHC

CERN

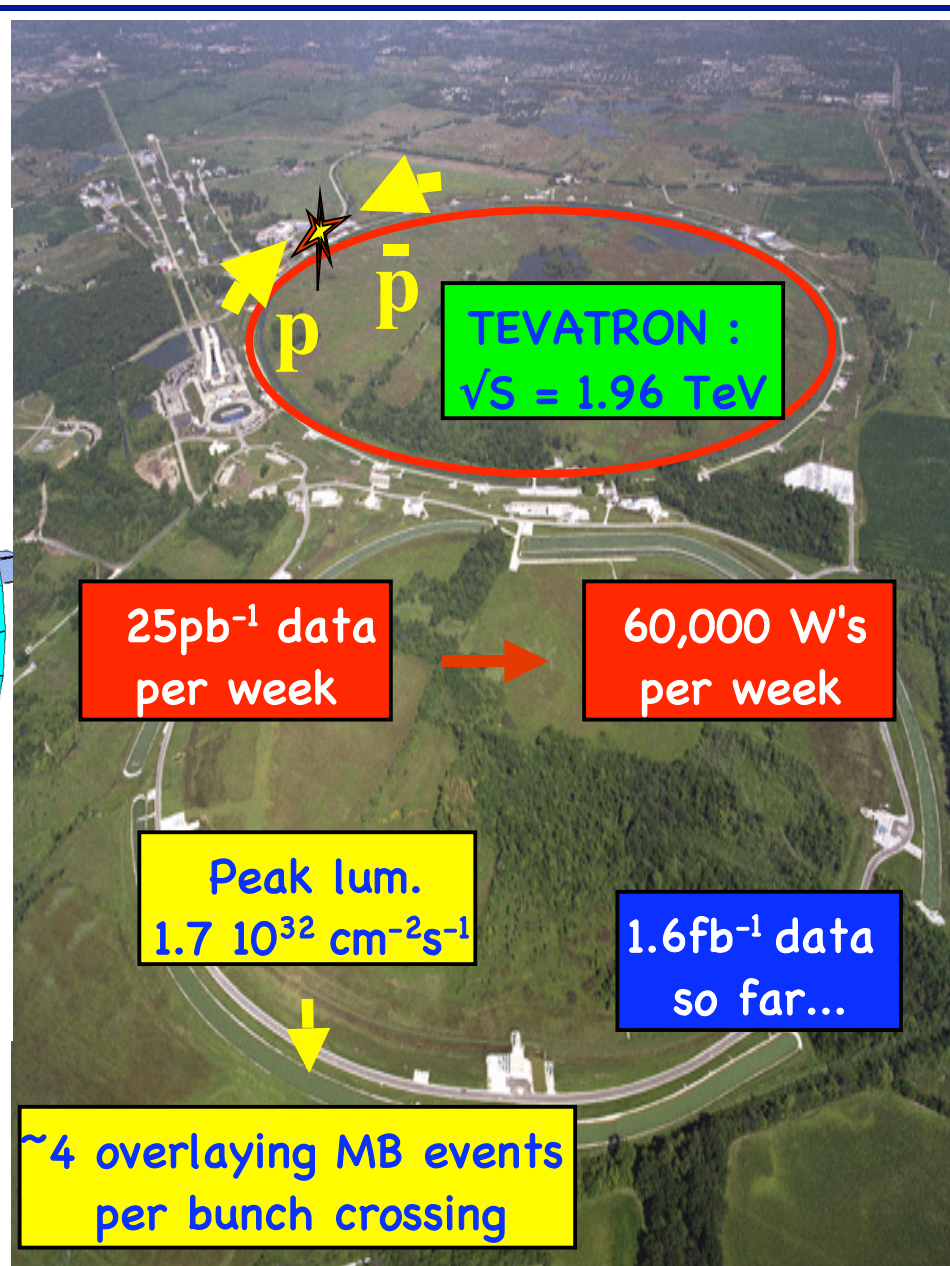
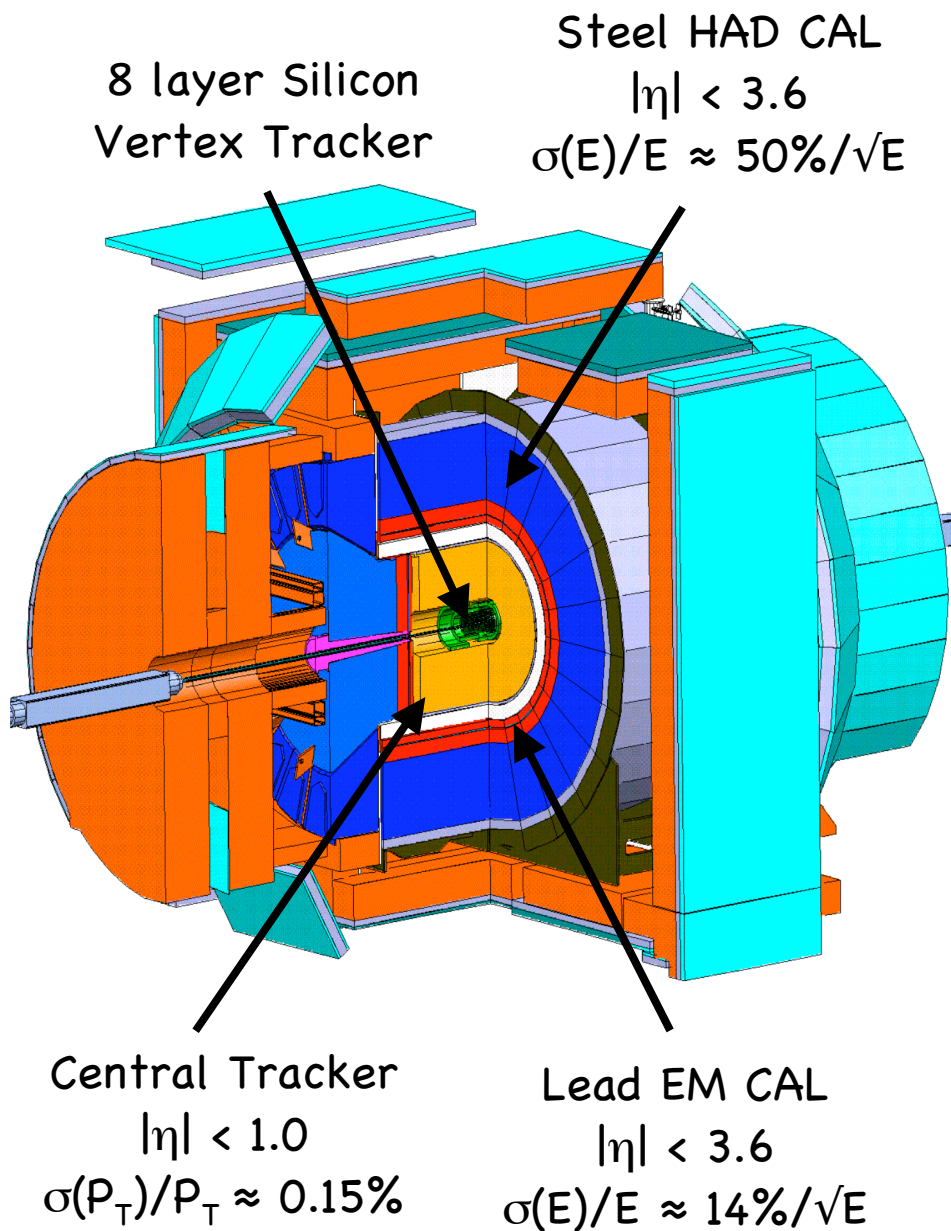
July 2006

OUTLINE

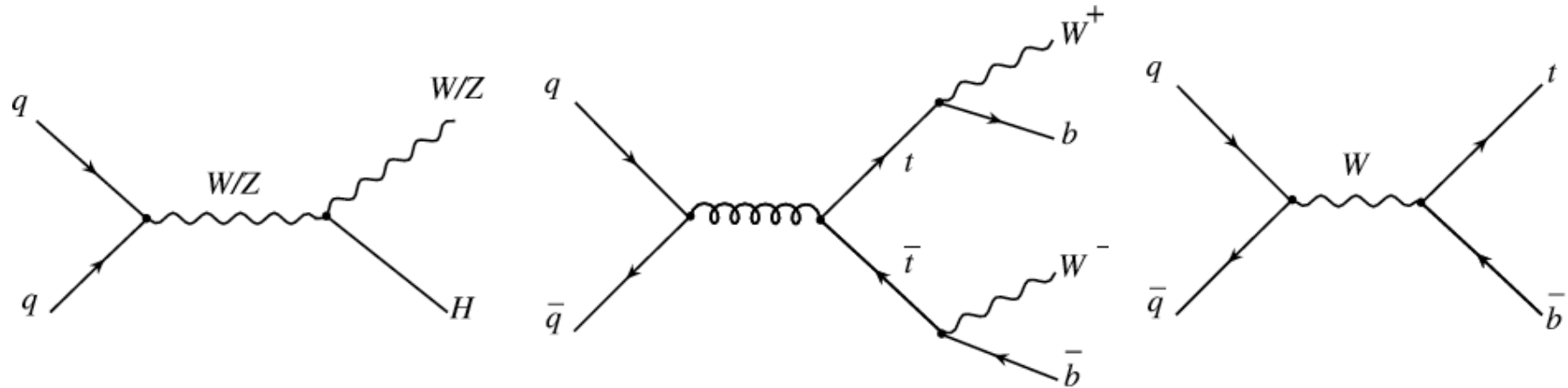
- Motivations for Study
- Monte Carlo Issues
- Measurement definition
- Correcting to hadron level
- Background estimation
- Results
- Conclusions & plans



CDF and the Tevatron



Motivation: Understanding a vital background



- Boson + jet is the final state for a number of important high p_T physics processes:
 - Top pair & single top production.
 - Higgs boson searches.
 - Searches for super-symmetric particles.
- All these signals are overwhelmed by large QCD production of boson + jets.
- It is crucial to have a good understanding of the boson + jets process.



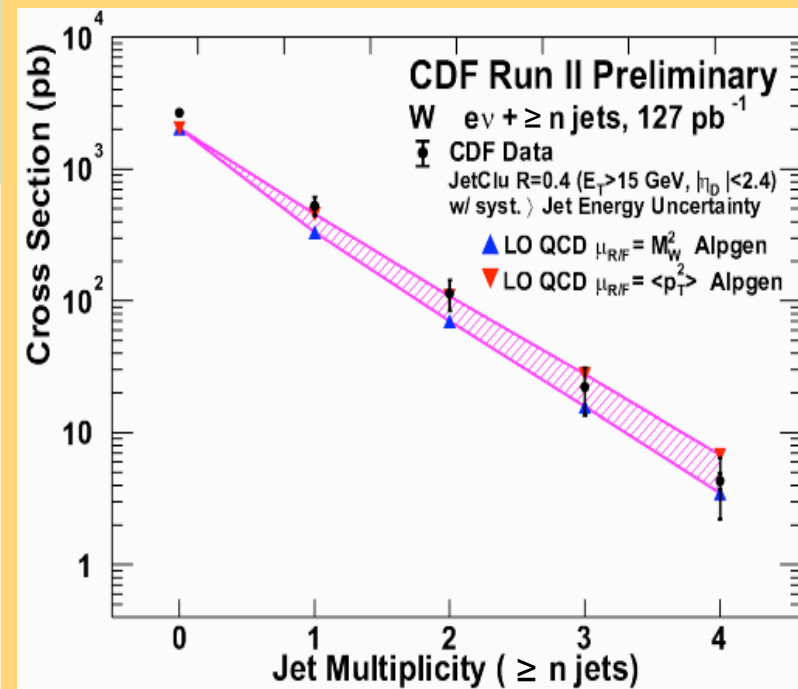
Motivation: Test of pQCD Predictions

- Testing ground for pQCD in multijet environment
- The presence of a W/Z boson:
 - Ensures high Q^2 – pQCD
 - Large BR into leptons – easy to detect experimentally
- Key sample to test LO and NLO pQCD calculations
 - Pythia, Herwig: parton shower & hadronization, limited ME
 - AlpGen : W + n parton ME, interface to Pythia/Herwig for PS, MLM ME-PS matching scheme
 - Sherpa : W + n parton ME, APACIC showering, CKKW ME-PS matching scheme
 - MCFM: NLO ME W + 1 or 2 partons
 - MC@NLO: W+X (NLO ME + herwig shower)
- Study the underlying event in an alternative topology than inclusive multijets.



W + n Jets LO Predictions

- W + n parton LO ME calculation + parton shower + hadronisation:
 - W + $\geq n$ jets Cross-section
 - Jet kinematics for $\geq n$ jets
- Issues:
 - Dependence on Q^2 scale
 - Dependence on parton cuts
 - Phase space overlap when combine n parton samples
- Advances:
 - ME-PS matching – CKKW and MLM prescriptions.
 - NLO predictions.



LO W + $\geq n$ parton cross-section vs data: good within large Q^2 uncertainties.



Definition of our measurement

- Aim for a definition as much as possible independent of theoretical predictions and detector effects.

$$\frac{d\hat{\sigma}_{W \rightarrow e\nu}}{d(\text{jet})}$$

- Restrict W decay to analysis acceptance.
- $P_T^{\text{ele}} > 20 \text{ GeV}$, $P_T^{\nu} > 30 \text{ GeV}$
- $WM_T > 30 \text{ GeV}/c^2$, $|\eta^{\text{ele}}| < 1.1$
- Reduces theoretical dependence of measurement, without comprising usefulness.

$$d(\text{jet})$$

- Jets: JETCLU cone 0.4, $E_T > 15 \text{ GeV}$, $|\eta| < 2.0$.
- Jet energies corrected to hadron level and for multiple interactions – underlying event remains.
- Differential w.r.t. 1st, 2nd, 3rd and 4th jet E_T , 1st-2nd jet invariant mass and ΔR .

- This is not an EWK measurement – W is a clean signal for high Q^2 events within which we can examine jet kinematics.



Making the Measurement

Identify $W \rightarrow e\nu$ candidate events from high E_T electron and large missing E_T .
Reconstruct jets.

In each bin of the jet kinematic variable calculate:

$$\sigma = \frac{N^{cand} - N^{bkgd}}{A \cdot \epsilon_{ID} \cdot L}$$

Backgrounds:

- QCD multijet
- Top pair
- $Z \rightarrow ee$, $W \rightarrow \tau\nu$
- WW
- Multiple interactions

Dataset luminosity
 $320 \pm 20 \text{ pb}^{-1}$

Acceptance and efficiency both estimated using detector simulated LO $W + \text{Jets}$ Monte Carlo.

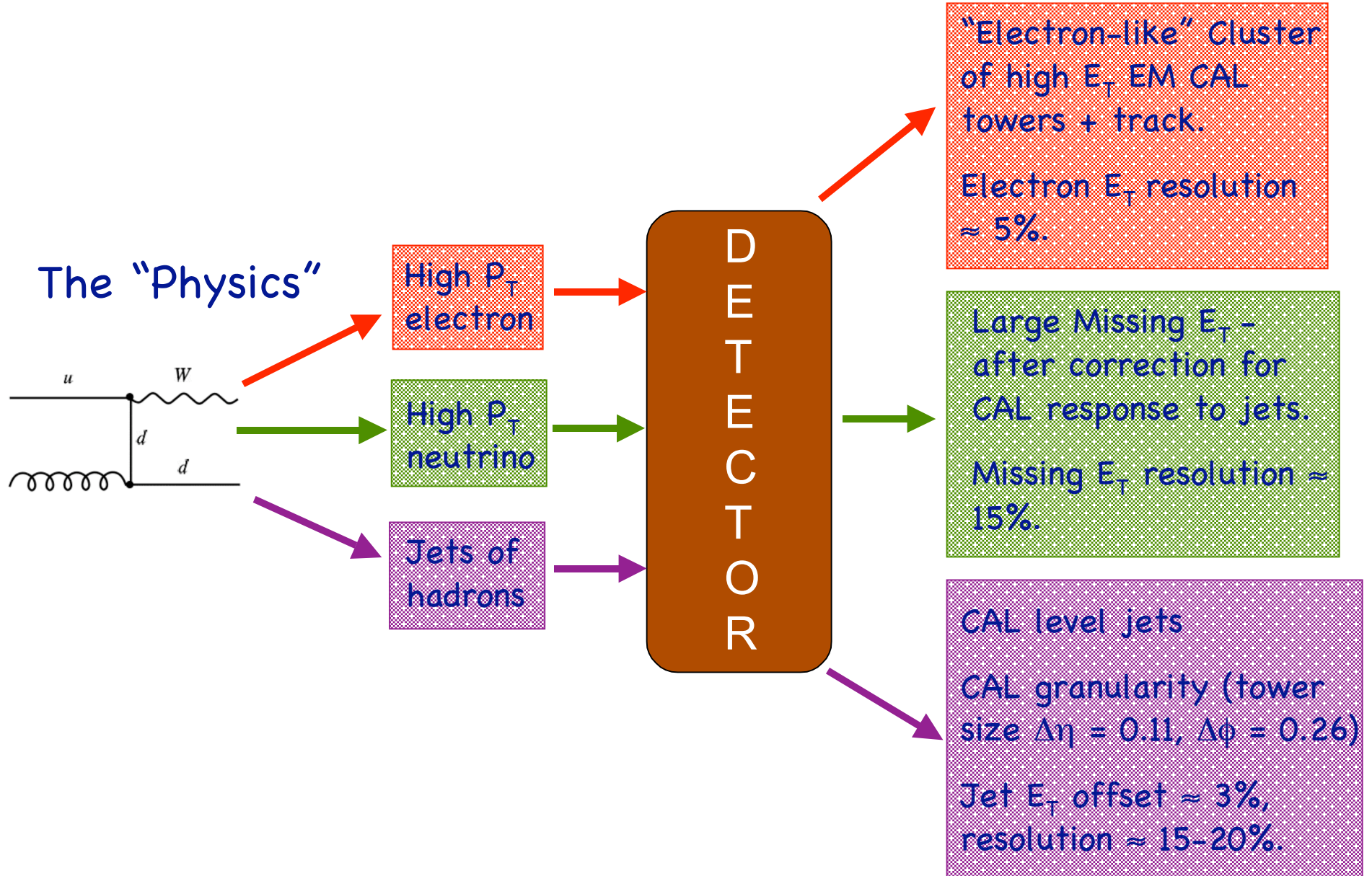
- Theoretical dependence enters the measurement via background and acceptance estimation - covered by systematics.
- Detector dependence removed by jet energy scale corrections and acceptance.



Making a CDF-independent Cross-Section Measurement



The Detector Effects



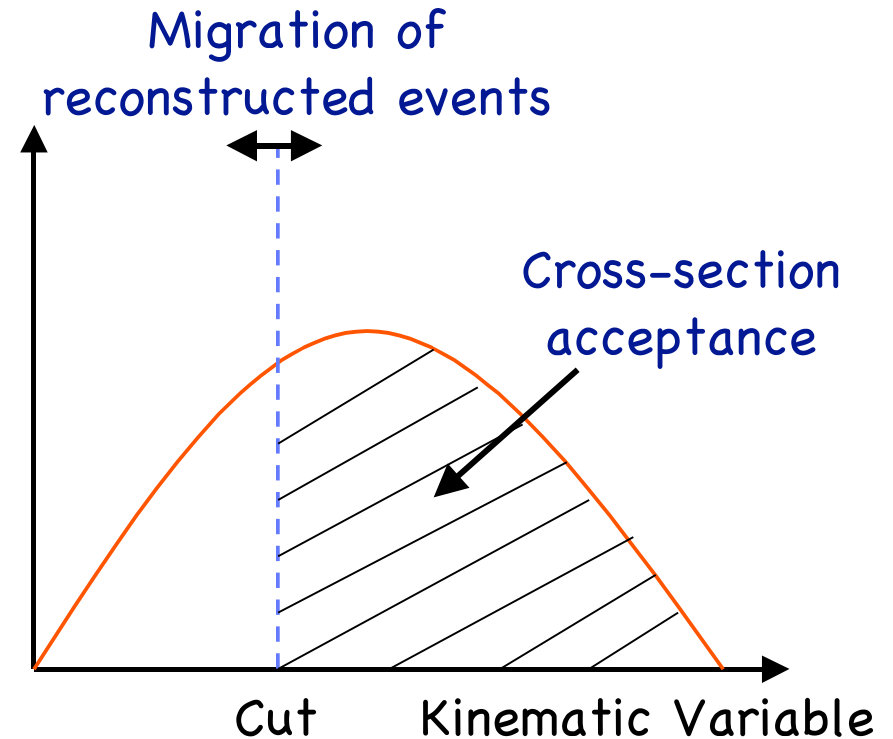
Making a detector independent measurement

- We rely largely on detector simulated Monte Carlo to remove detector effects.
- Affects on W selection accounted for using Monte Carlo derived acceptance and electron ID efficiency corrections.
- Jets corrected back to hadron-level energy:
 - Using generic CDF jet energy scale corrections: on average correction of energies back to hadron-level.
 - On average correction for energy from multiple p-pbar interactions (but no correction for UE here).
 - Additional unfolding of detector resolution effects on jet E_T spectrum.



Acceptance and Efficiency Correction

- W Cross-section phase space same as analysis acceptance.
- Acceptance factor reduced to accounting for detector resolution and local shape around cut.
- Reduces theoretical dependence of measurement.



- Use W MC for acceptance and electron ID efficiency:
 - Systematic on ID efficiency by comparing MC and Z data
 - Estimation of acceptance systematic by comparing different MC models } 5%

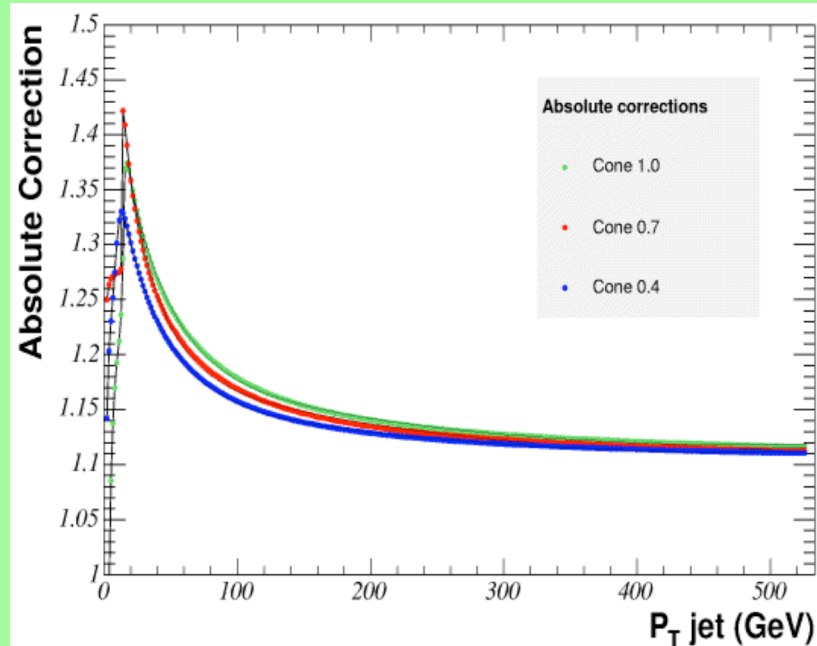
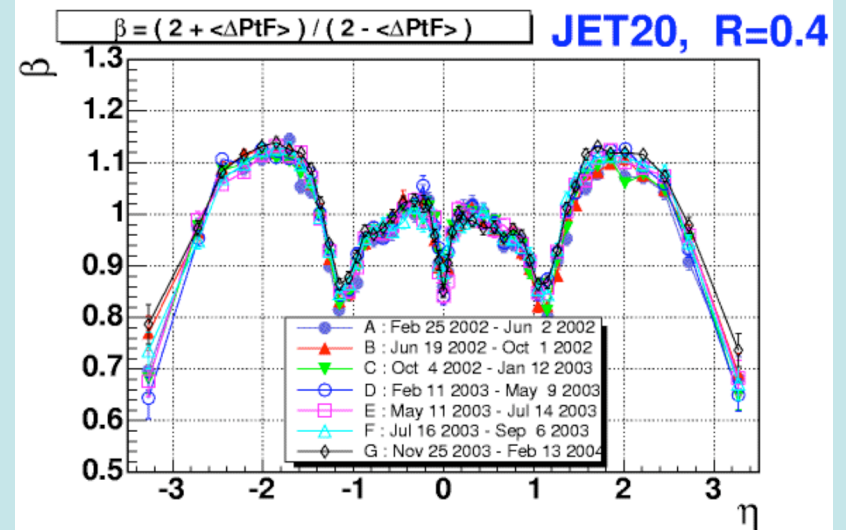
$$A \cdot \epsilon_{ID} \approx 0.6 \pm 0.03$$

Largely independent of jet kinematics



CDF Jet Energy Scale Corrections

- RELATIVE corrections from “dijet balancing” in data.
- Purpose: Equalise response over all η to that in central CAL.
- Identify QCD $2 \rightarrow 2$ processes - one central “trigger” and one “probe” jet.
- Probe E_T should on average equal trigger E_T .

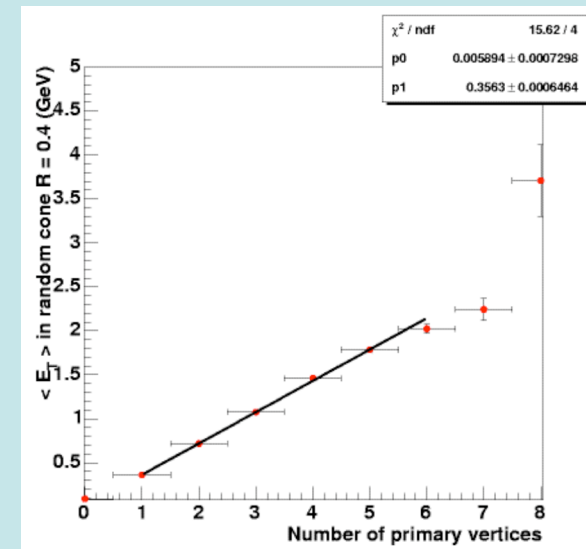


- ABSOLUTE corrections from detector simulated dijet MC.
- Purpose: Correct jets on average from CAL to hadron energy.
- In MC jets clustered at both hadron level and CAL level and “matched”.
- Determine most probably HAD E_T for each CAL E_T - defines correction.
- Systematic 2-3%.



Multiple Interactions Correction

- At $1 \times 10^{32} \text{cm}^{-2}\text{s}^{-1}$ most probably no. of interactions is 2 per bunch crossing.
- Additional “soft” interaction deposits energy in jet cone.
- Using minimum bias data – parameterise average energy in random cone as function of number of additional reconstructed vertices.



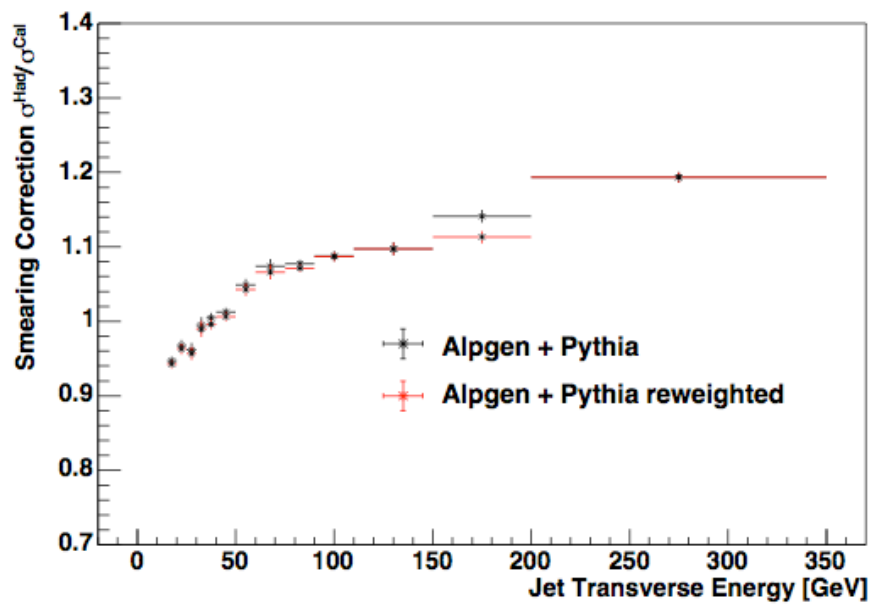
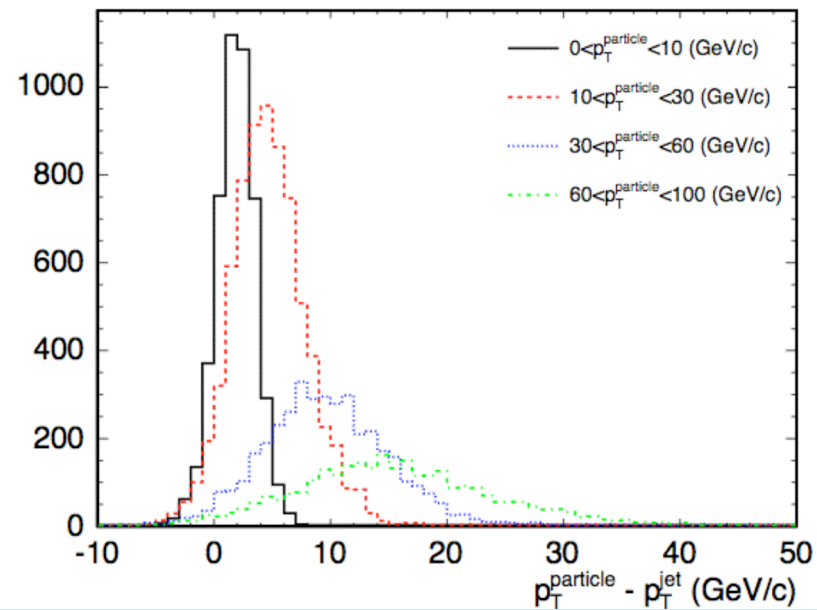
- In addition – rarely MI can result in an extra jet.
- “Promotion bkgd” – causes migration of signal events between jet multiplicity bins.
- Using MB data get extra jet rate per additional vertex.
- W candidate N_{vtx} distbn – normalise MB jet E_T spectra.

njet \geq	1	2	3	4
bkgd _{frac} %	2.4	3.6	3.8	3.3



Analysis Specific Unsmearing Correction

- Absolute corrections only correct on average for offset of CAL energy response.
- What about impact of non-zero CAL resolution $\sigma \approx 15\text{--}20\%$.
- Resolution alters shape of measured E_T spectrum – necessitates “unsmearing” correction.



- Use detector simulated W+Jets MC.
- Form jet E_T spectra at HAD level and at (JES corrected) CAL level.
- Define unsmearing factor in each bin as $N_{\text{HAD}}/N_{\text{CAL}}$.
- To remove dependance on MC HAD spectrum – use iterative procedure, reweighting until CAL spectrum equal to that observed in data.

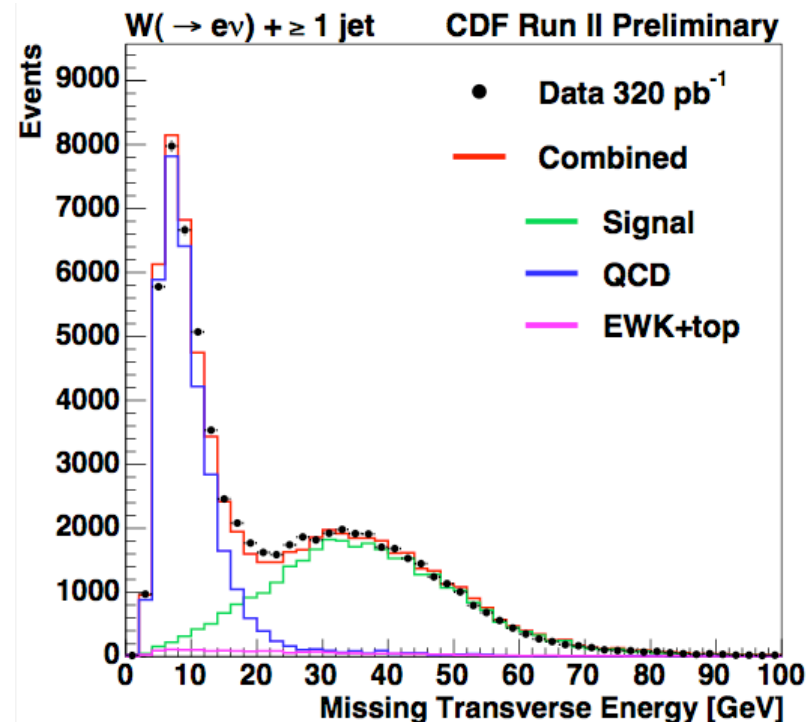
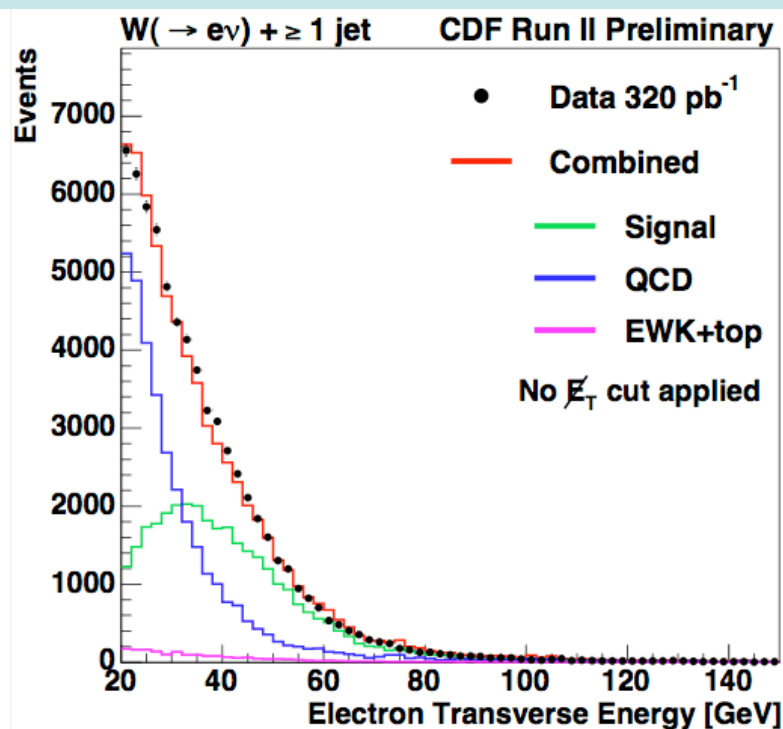


Other Elements of the Measurement



Background Estimation

- QCD modelled by fake-electron sample formed from dataset.
- MC for other bkgds and signal
- Background normalisation from fit to data missing E_T distribution.
- Excellent agreement in other kinematic variables.

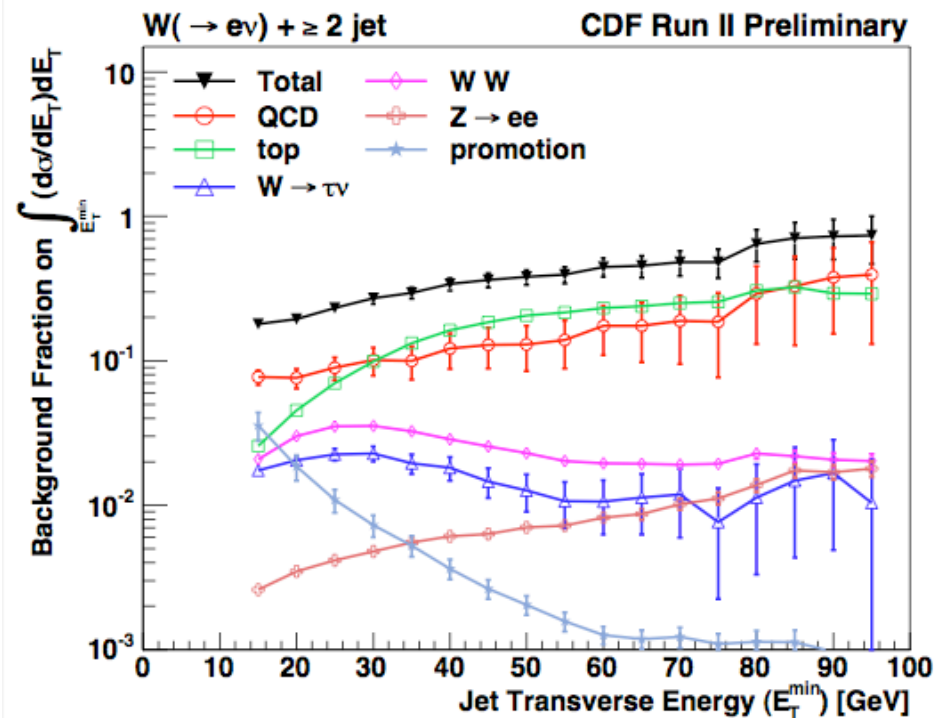
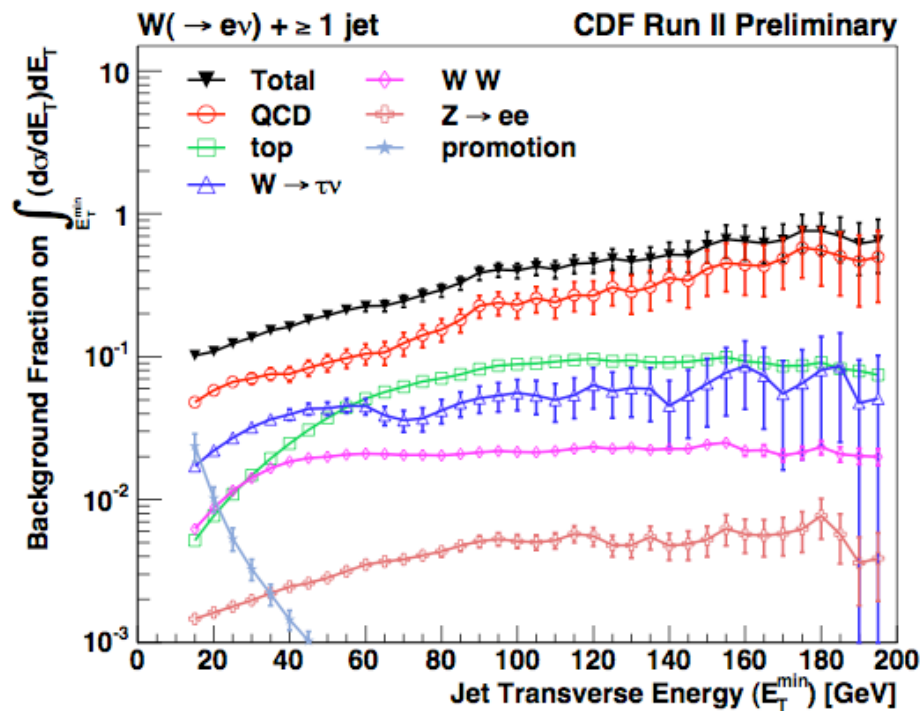


Background Systematics:

- Fake-electron statistics (dominant)
- Fake-electron QCD model (5-20%)
- Top cross-section (10-20%)
- MC model dependence (5%)



Background Fractions



- QCD is a substantial background contribution, dominating at low E_T .
- But in high E_T region Top pair production is dominant.



The Cross-Section Results

(MC comparisons by no means exhaustive)

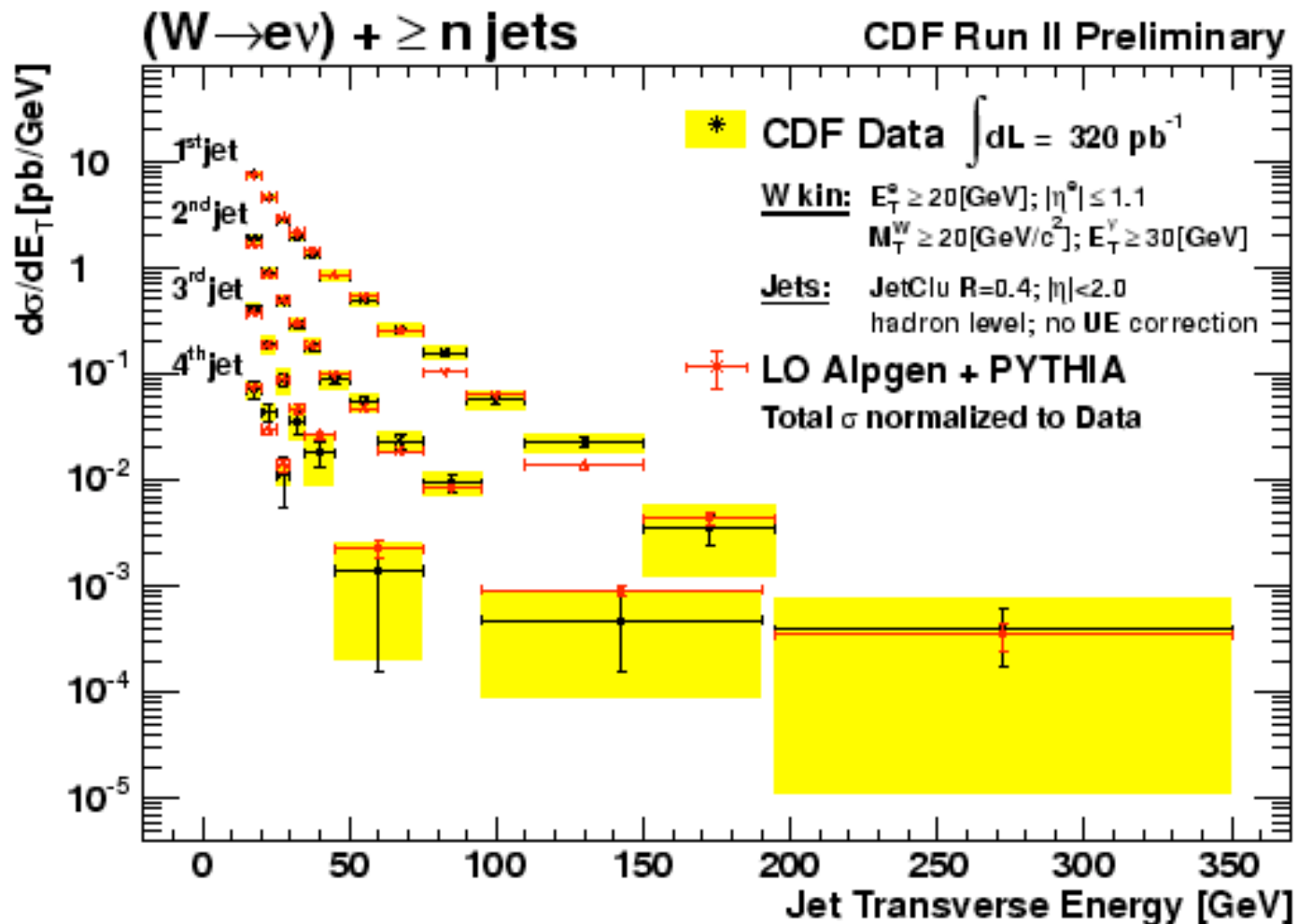




CDF W+jets Differential Cross-Section

$$\frac{d\hat{\sigma}_{W \rightarrow e\nu}}{dE_T}$$

For 1st, 2nd, 3rd
and 4th Jet E_T



MC has been normalized to inclusive data cross section in each jet sample!

LO W + n parton prediction reproduces shape of $d\sigma/dE_T$ reasonably well.

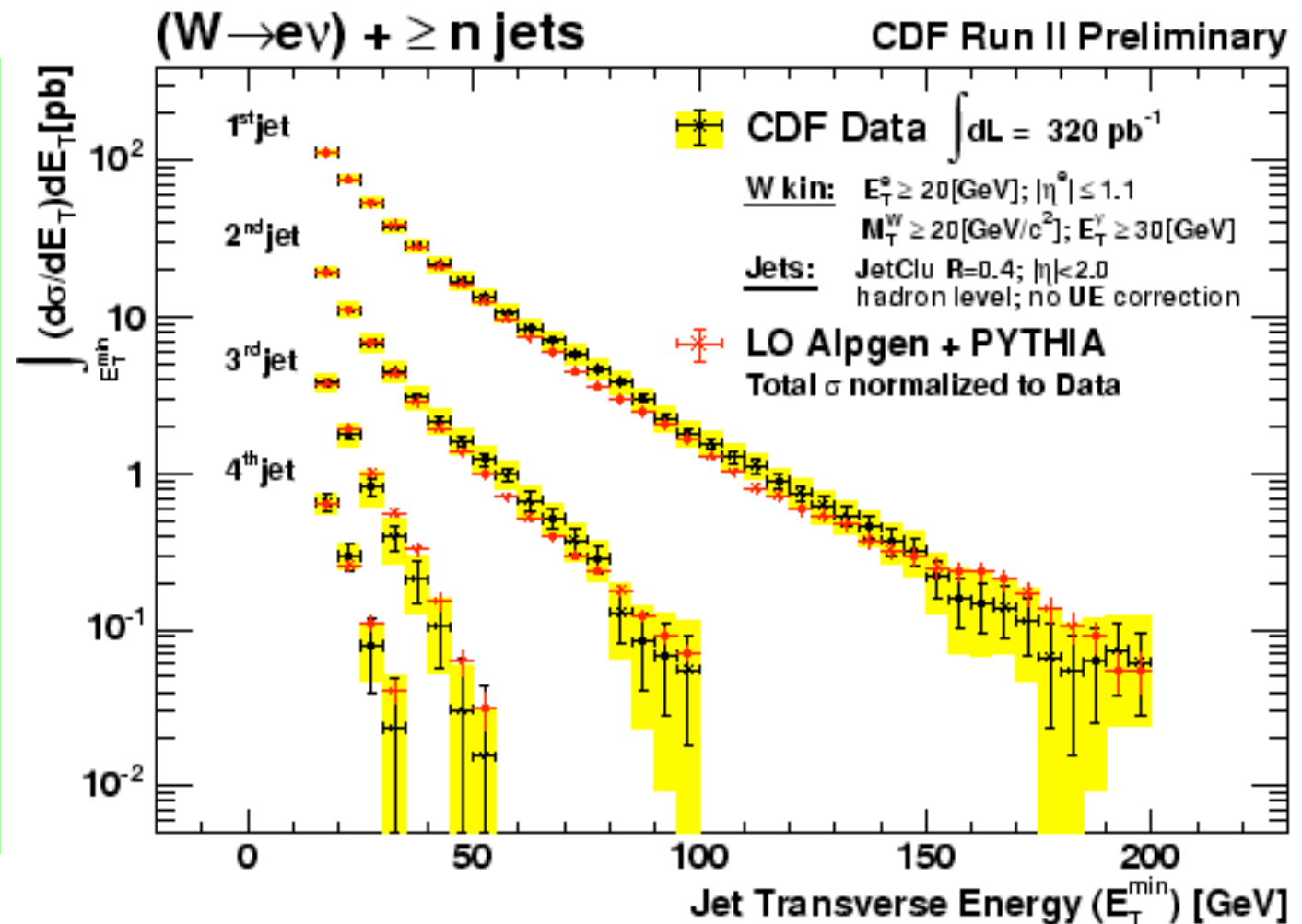


CDF W+jets Integrated Cross-Section

$$\int_{E_T^{\min}} \frac{d\hat{\sigma}_{W \rightarrow ev}}{dE_T} dE_T$$

For 1st, 2nd, 3rd
and 4th Jet E_T

Essentially a
measurement of
 $\sigma(W + \geq n \text{ jets})$ for
different jet E_T
thresholds.



MC has been normalized to inclusive data cross section in each jet sample!

First bin MC & data is in perfect agreement by construction.



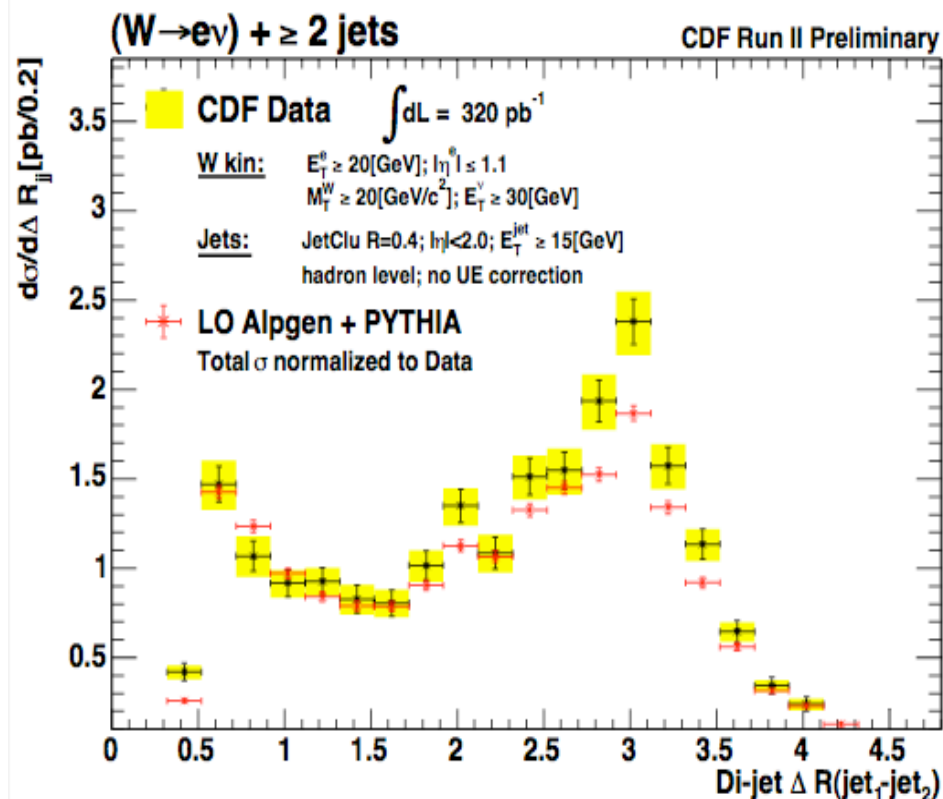
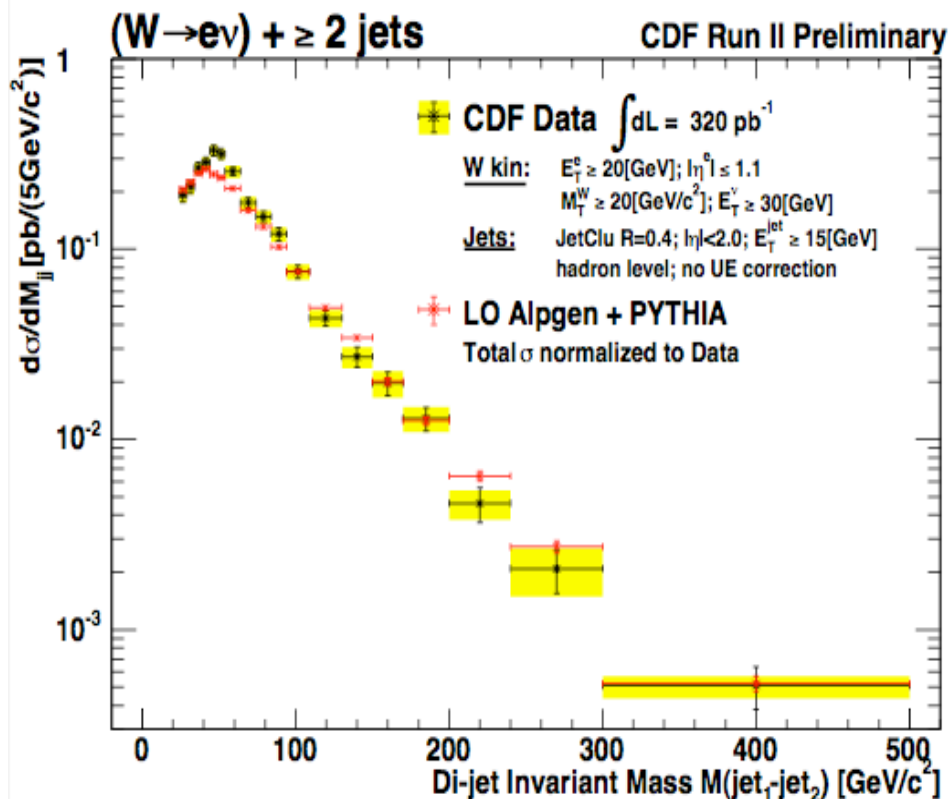
CDF W+jets Differential Cross-Section

$$\frac{d\hat{\sigma}_{W \rightarrow e\nu}}{dM_{jj}}$$

Differential σ w.r.t 1st-2nd jet invariant mass in the $W + \geq 2$ jet sample

$$\frac{d\hat{\sigma}_{W \rightarrow e\nu}}{d\Delta R_{jj}}$$

Differential σ w.r.t 1st-2nd jet ΔR in the $W + \geq 2$ jet sample



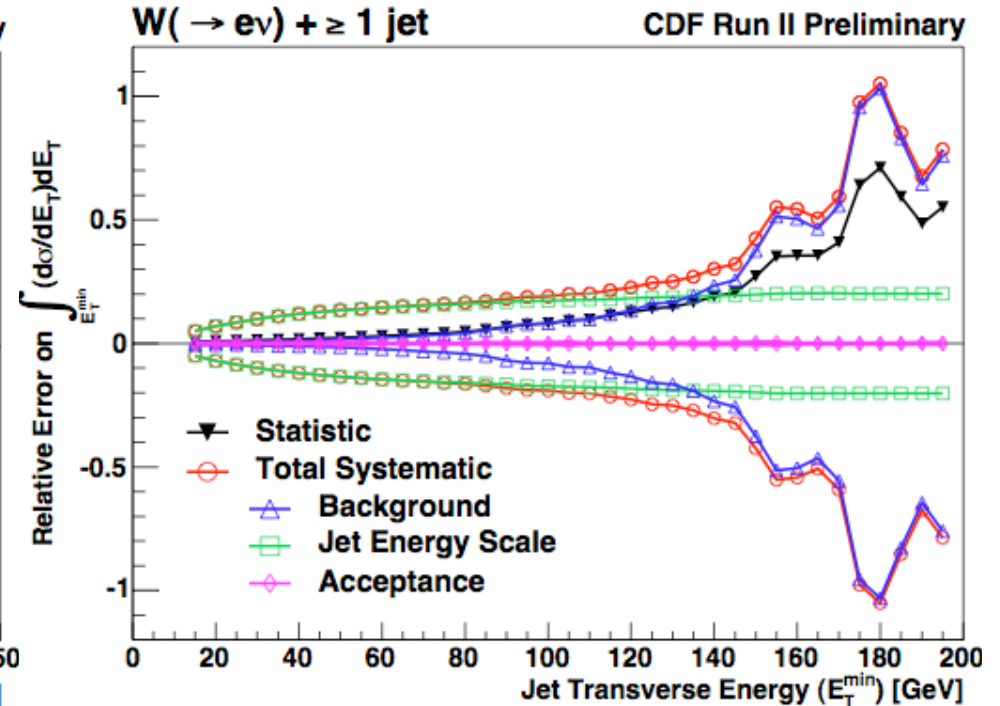
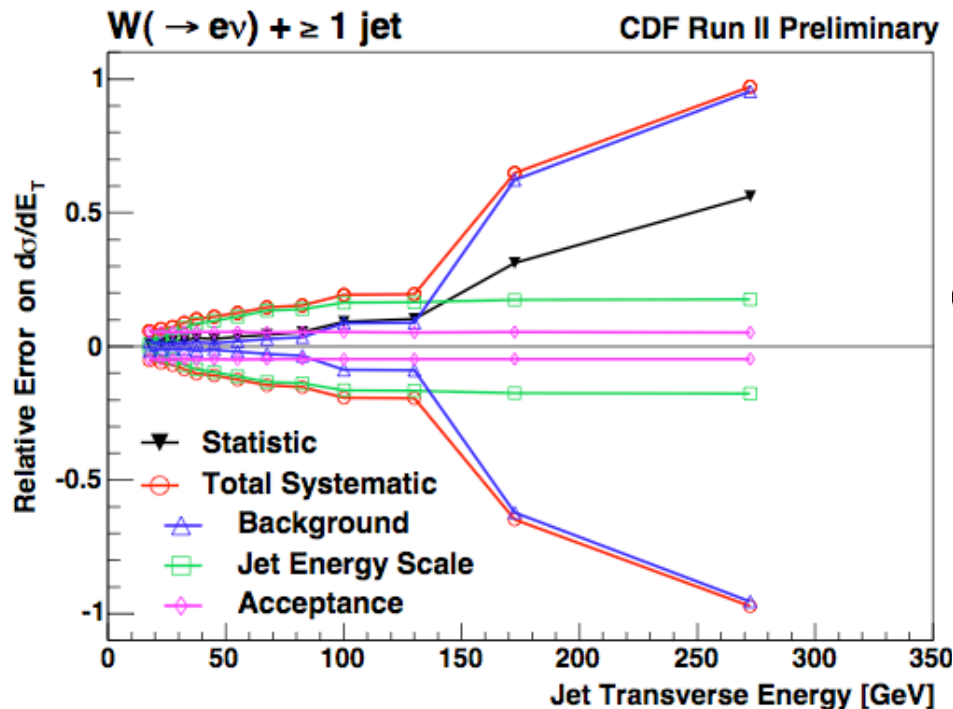
MC has been normalized to measured W+2 jet inclusive cross section!



Errors breakdown

Representative of the behavior of errors in the measurements

Relative error on leading jet $d\sigma/dE_T$ Error on leading jet $\int (d\sigma/dE_T) dE_T$



- Large statistical uncertainty at large E_T .
- Systematic dominated by jet energy scale at low E_T , and by the (QCD) background subtraction at high E_T .



Work in progress and Plans

- Extend the measurement to use muons and to 1fb^{-1} :
 - Larger E_T range, more sensitive to the tail of the cross-section.
 - Better control on data driven QCD background subtraction.
 - Move to the preferred midpoint algorithm – don't expect big changes.
- Make extensive comparisons to theory, both shape and rate predictions:
 - LO ME-PS matching prescriptions – CKKW and MLM
 - NLO predictions: MCFM (parton level), MC@NLO (hadron level)
- Measure the Z + Jets cross-section:
 - Reduced statistics but backgrounds greatly reduced also.
 - Z + Jet events provide an alternative and cleaner environment for UE studies than multijet events.



Conclusions

- A new measurement of differential $\sigma(W + \text{jets})$ w.r.t jet kinematics, more suitable for theoretical comparisons:
 - Hadron level measurement: jet detector effects removed.
 - Differential measurement: background, acceptance and efficiency impact on shape accounted for.
 - Restricted W decay cross-section definition: reduced theoretical dependence.
- Any theorist can overlay their hadron-level predictions without need for CDF detector simulation.
- The systematic on many high p_T measurements receives a substantial contribution from boson + jet knowledge.
- Crucial to have a robust simulation of boson + jets to explore for new physics at Tevatron & LHC.