

Single microwave photon detectors for ADMX

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ADMX-G2, an Axion Dark Matter Radio

Tests QCD axion models which would solve the strong-CP problem

Coherently scatter classical axion dark matter wave (10²⁰ axions/m³ coherent state) on DC magnetic field.

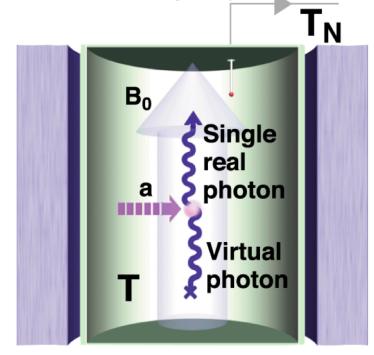
Resonantly enhance scattering cross section using tunable high-Q microwave cavity as the "antenna."

Use quantum-limited amplifier with "irreducible" single photon readout noise.

At each radio tuning, look for tiny excess power by averaging away noise over >10⁵ measured power spectra.

Confirm/exclude the predicted axion-photon coupling at that frequency.

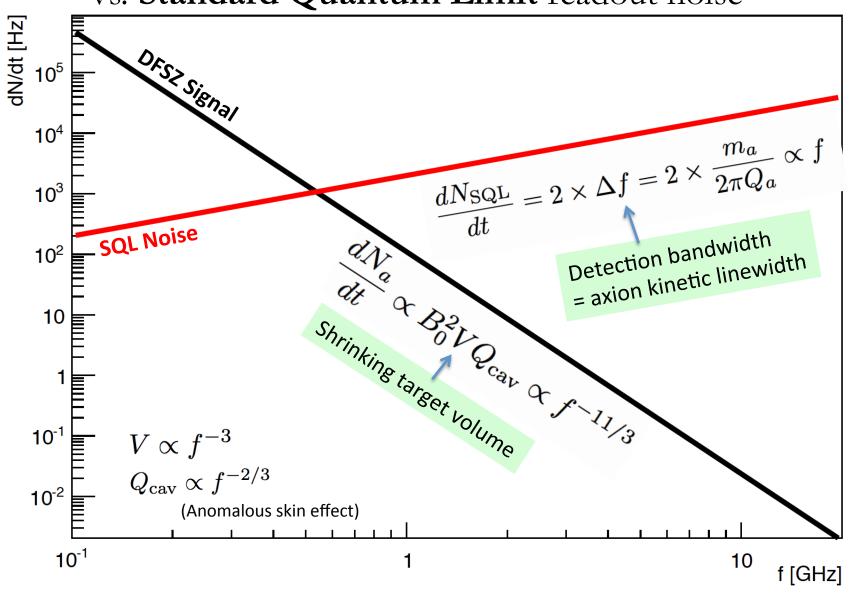
Aaron S. Chou (FNAL), ICHEP 2016



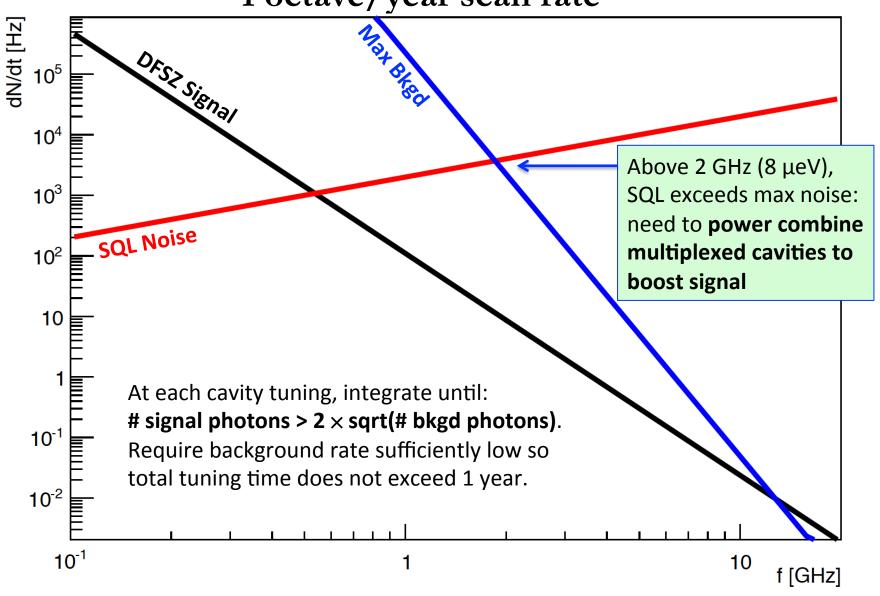


ADMX copper cavity

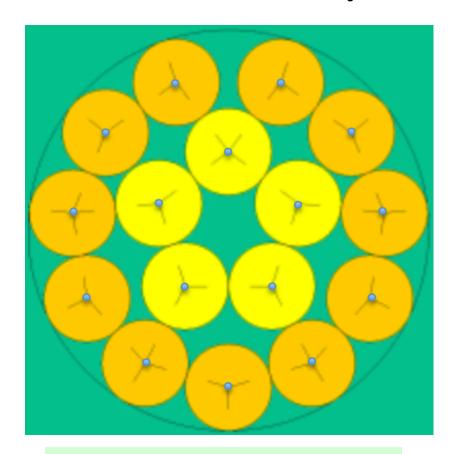
DFSZ signal photon rate for single volume= λ^3 cavity vs. **Standard Quantum Limit** readout noise



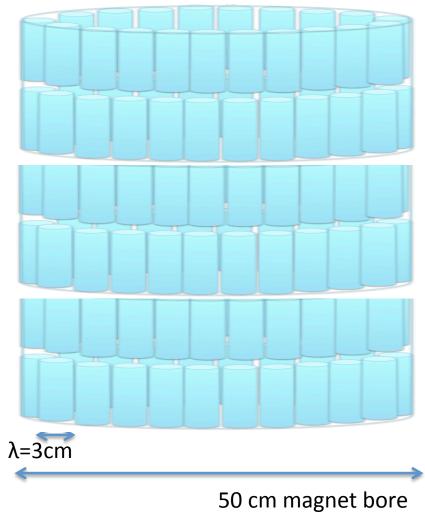
Maximum tolerable background photon rate for 1 octave/year scan rate



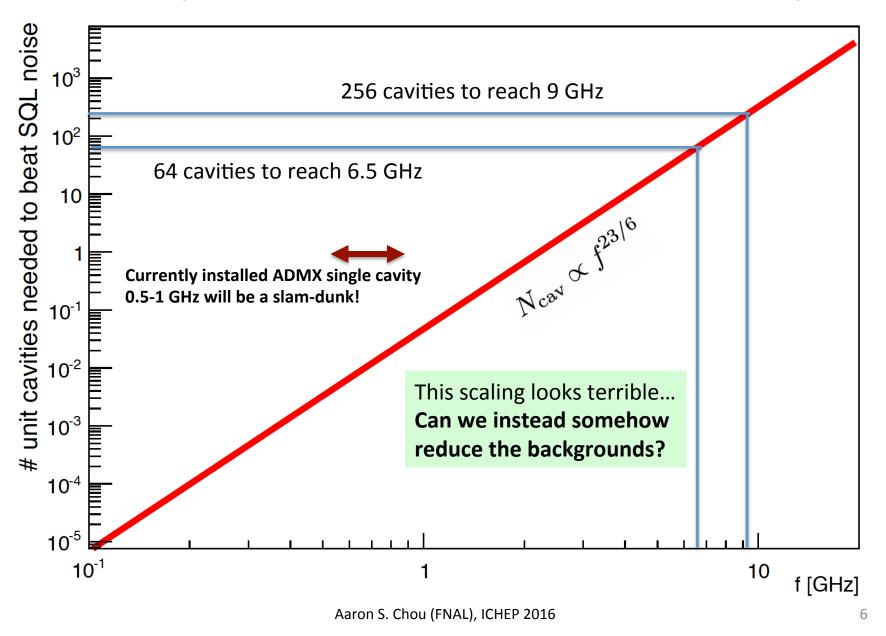
Swiss watch problem: Many resonant elements must be simultaneously tuned to the same frequency



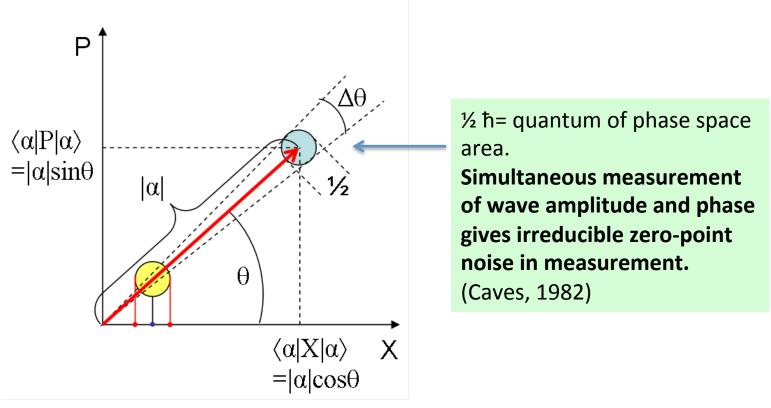
Cost and complexity scale at least linearly with N_{cav}



How many cavities are needed to cover 1 octave/year?



Quantum-limited amplifiers suffer from zero-point readout noise – the Standard Quantum Limit (SQL)



Thermal noise = kT of energy per resolved mode

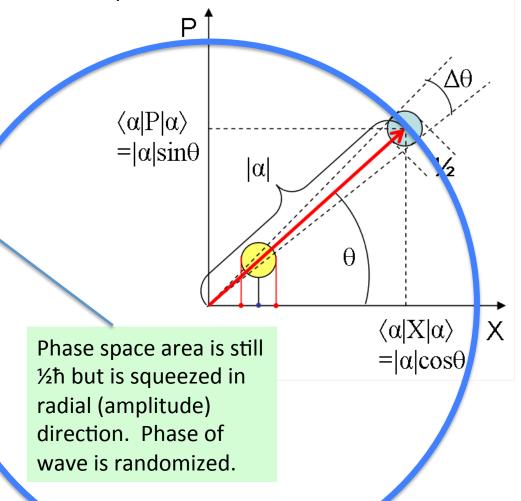
→ Quantum noise = 1 photon per resolved mode in the T=0 limit.

Noise photon rate exceeds signal rate in high frequency dark matter axion searches.

Need new sensor technology....

Quantum non-demolition (QND) single photon detection can do much better

Number operator commutes with the Hamiltonian → all backreaction is put into the phase. Measure exact photon number. Noise = shot noise, thermal backgrounds, read noise.



Demonstrated with Rydberg atoms, (Haroche/Wineland Nobel Prize 2012)

Implementation using solid state artificial atom qubits, (D.Schuster et.al, 2007)

Proposed for axion search: (Lamoreaux, et.al, 2013, Zheng, et.al, 2016)

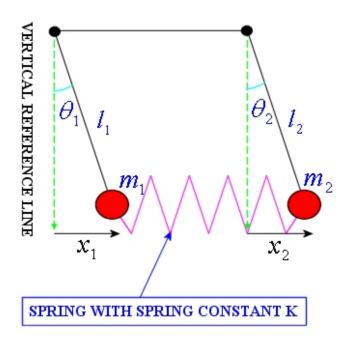
Note: different from Carrack technique of non-QND single photon counting using Rydberg atoms (Ogawa, Matsuki, Yamamoto, 1996).

What does a QND single microwave photon detector look like?



Coupled oscillators

Energy stored in the spring



$$H = \begin{pmatrix} \omega_1 & g \\ g & \omega_2 \end{pmatrix}$$

Mixing angle to diagonalize:

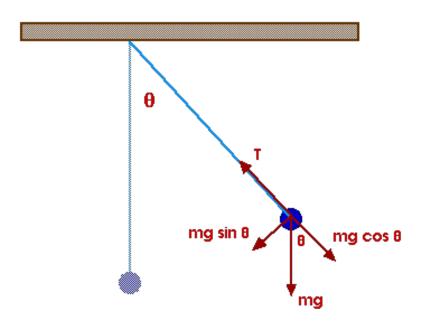
$$\tan 2\theta = 2g/(\omega_1 - \omega_2)$$

Normal mode frequencies for small g

$$\tilde{\omega}_1 = \omega_1 + \frac{2g^2}{\omega_1 - \omega_2}$$

$$\tilde{\omega}_2 = \omega_2 - \frac{2g^2}{\omega_1 - \omega_2}$$

Suppose one oscillator has non-linear restoring force



For example:

Resonant frequency of a realworld pendulum increases with oscillation amplitude

$$\omega_2 = \omega_2(A_2)$$

Then the instantaneous resonant frequency of **linear** oscillator 1 depends weakly on the amplitude or occupation number of **nonlinear** oscillator 2

$$\tilde{\omega}_1 = \omega_1 + \frac{2g^2}{\omega_1 - \omega_2(A_2)}$$

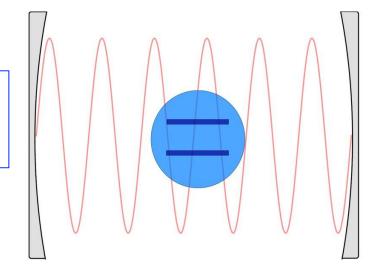
Measuring the frequency of oscillator 1 performs a QND measurement on the number of quanta stored in oscillator 2 (and vice-versa)

Cavity QED:

Use 2-level atom to measure cavity photon population

Linear cavity

Bosonic oscillator, Number operator = $a^{\dagger}a$



2-level "atom"

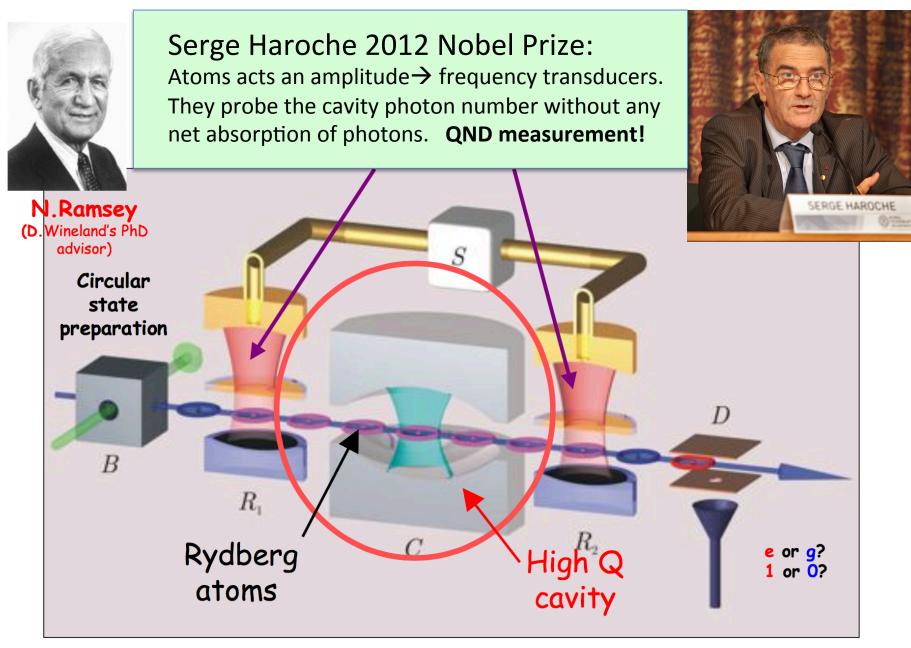
Fermionic oscillator, Number operator = σ_z

The 1^{st} order non-linearity in (number operator)² in the undiagonalized Hamiltonian is:

$$Hpprox \hbar\omega_{
m r}\left(a^{\dagger}a+1/2
ight)+rac{\hbar}{2}\left(\omega_{
m a}+rac{2g^2}{\Delta}a^{\dagger}a+rac{g^2}{\Delta}
ight)\sigma_{
m z}$$
 Δ = $\omega_{
m r}$ - $\omega_{
m r}$

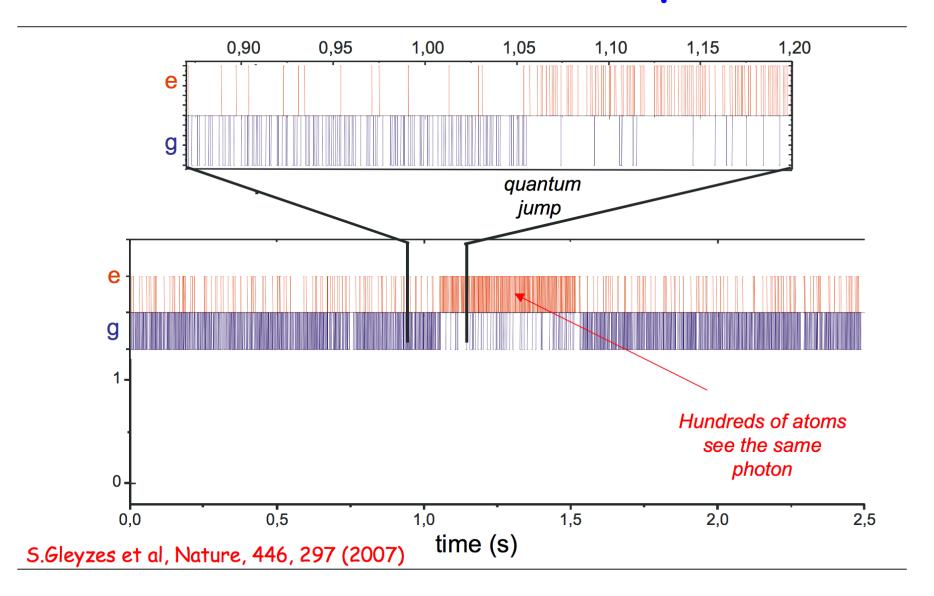
The atom frequency depends on the cavity resonator's occupation number!

This product of number operators commutes with H and allows QND measurement.



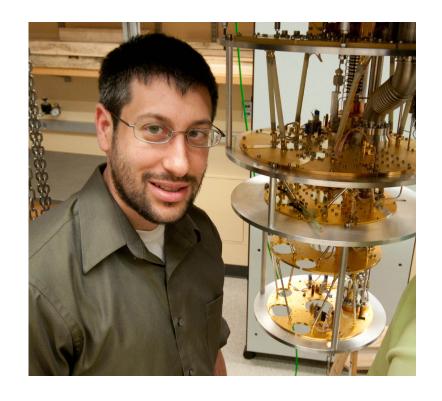
An atomic clock delayed by photons trapped inside

Birth, life and death of a photon

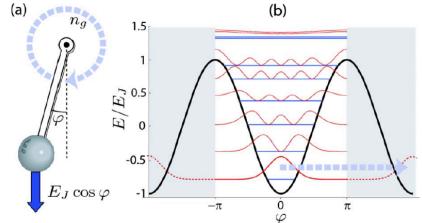


Nonlinear circuit oscillators have non-degenerate energy level spacings and hence behave just like 2-level atoms

Slides from Dave Schuster (U.Chicago)



Transmon qubit based on the Cooper pair box J.Koch, e.al, Phys.Rev.A76, 042319 (2007)



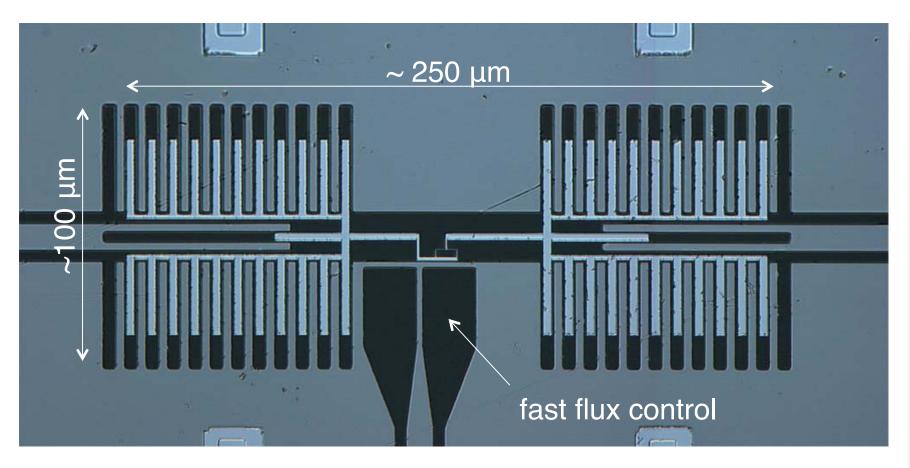
QND detectors developed for high fidelity quantum computing qubit readout.

B.R. Johnson, et.al, Nature Physics 6, 663-667 (2010)

The sarantapede

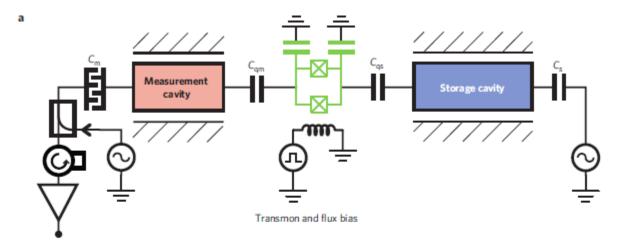


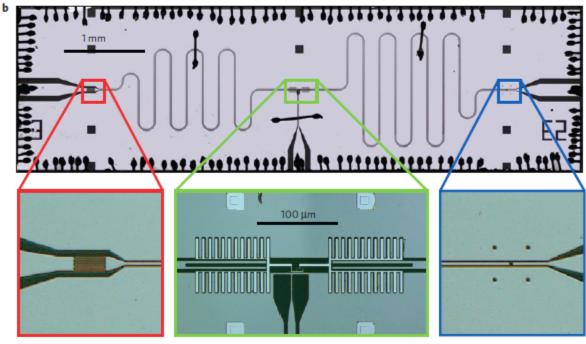
An end-coupled "transmon" qubit with ~40 legs



QND Detector = qubit + fast cavity





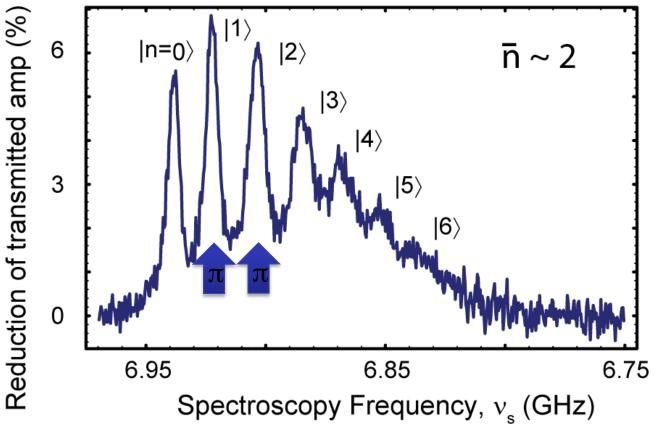


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Sensing photon number with a qubit



$$H \approx \hbar \omega_r a^{\dagger} a + \frac{\hbar}{2} (\omega_a' + 2\chi a^{\dagger} a) \sigma_z$$



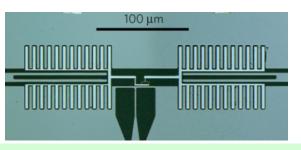
 Qubit transition frequency depends on photon number in cavity

Theory: J. Gambetta, A. Blais, ..., S. Girvin, and R. J. Schoelkopf, *PRA 94 123602 (2005)*Aaron S. Chou (FNAL), ICHEP 2016

Experiment: D. I. Schuster, ..., S. M. Girvin, R. J. Schoelkopf, Nature (London) **445** 515 (2007)

Prototype for 10 GHz axion QND detector





Superconducting qubit in field-free bucking coil region acts as an amplitude \rightarrow frequency transducer for QND measurements.

Qubit frequency shifts by 10 MHz per photon deposited in axion cavity. Successful "spin-flip" of qubit confirms presence of cavity photon.



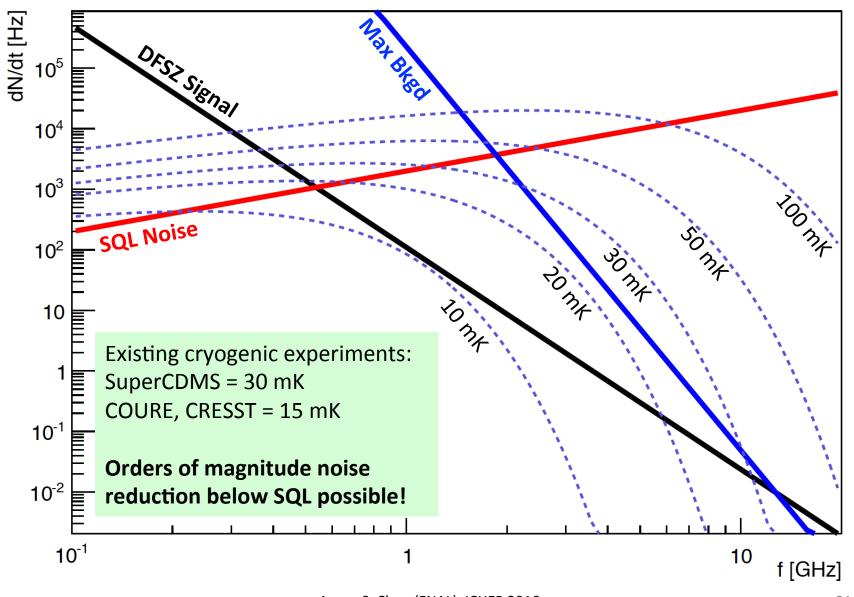
Axion scattering cavity dipped into high B-field region

Akash Dixit, Aaron Chou, David Schuster, R&D in progress

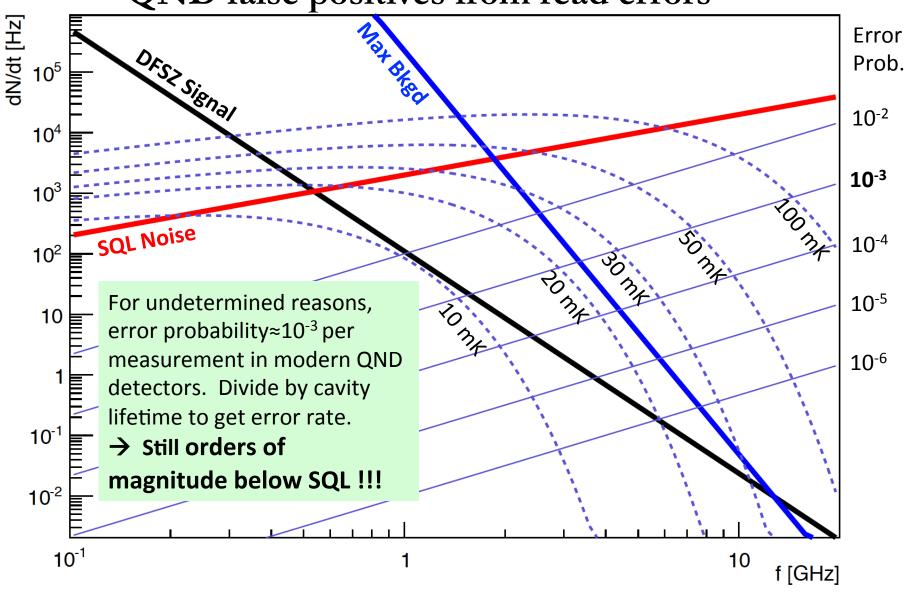
What are the improvements one can realistically expect from qubit-based QND readout for high frequency axion detectors?



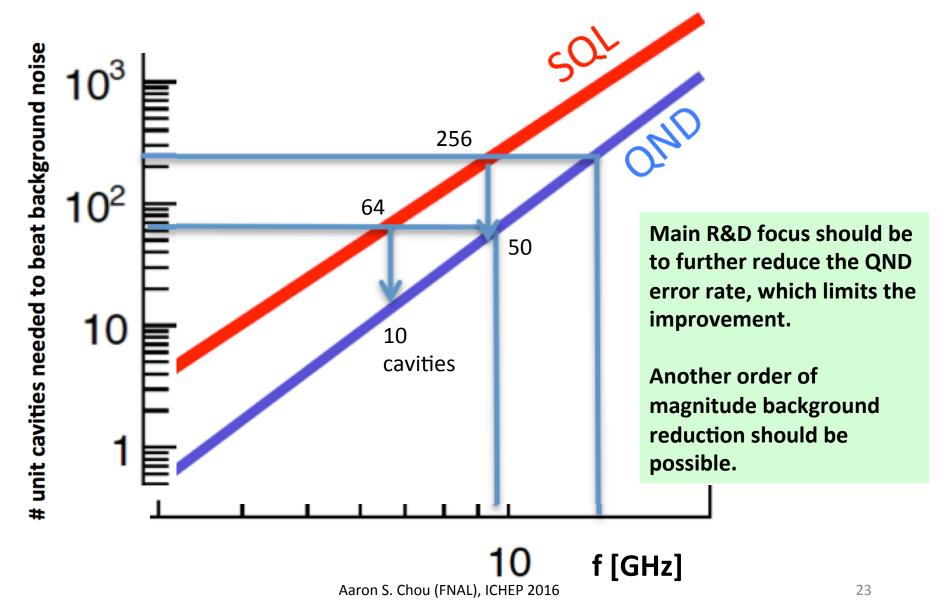
Thermal photon emission rates are negligible



Another background "photon" source: QND false positives from read errors



QND reduces required N_{cav} by factor of 6, or alternatively extends frequency range for fixed N_{cav}



Reducing noise is cheaper than increasing signal



A fun summer at Fermilab!

