



a $\mu \rightarrow 3e$ search at 10^{-16} ?

ETHZ, November 17 2008

andries van der schaar, Zürich

1 Why do it?¹

- Flavor puzzle: what is the origin of the patterns of quark and lepton masses and mixings?
- Most models predict observable rates for LFV decays involving charged leptons which may even be required to solve the flavor puzzle.
- Rare μ decays give best constraints in most scenarios and their relative strengths would help in the interpretation.
- Experimental prospects differ considerably between the three modes of interest:
- **Previous talk:** searches for $\mu - e$ conversion are limited by the available beam intensities and large improvements in sensitivity may still be achieved.
- **Next talk:** searches for $\mu^+ \rightarrow e^+ \gamma$ are limited by accidental $e^+ \gamma$ coincidences and muon beam intensities have to be reduced now already.
- **This talk:** what about $\mu \rightarrow 3e$ ² (20 year old upper limit 10^{-12}) ?

¹"Flavor physics of leptons and dipole moments", report of Working Group 3, CERN Workshop "Flavor in the era of the LHC", November 2005 - March 2007.

<http://www.springerlink.com/content/784n166817gp2gj8/>

²Discussed at a PSI workshop four months ago



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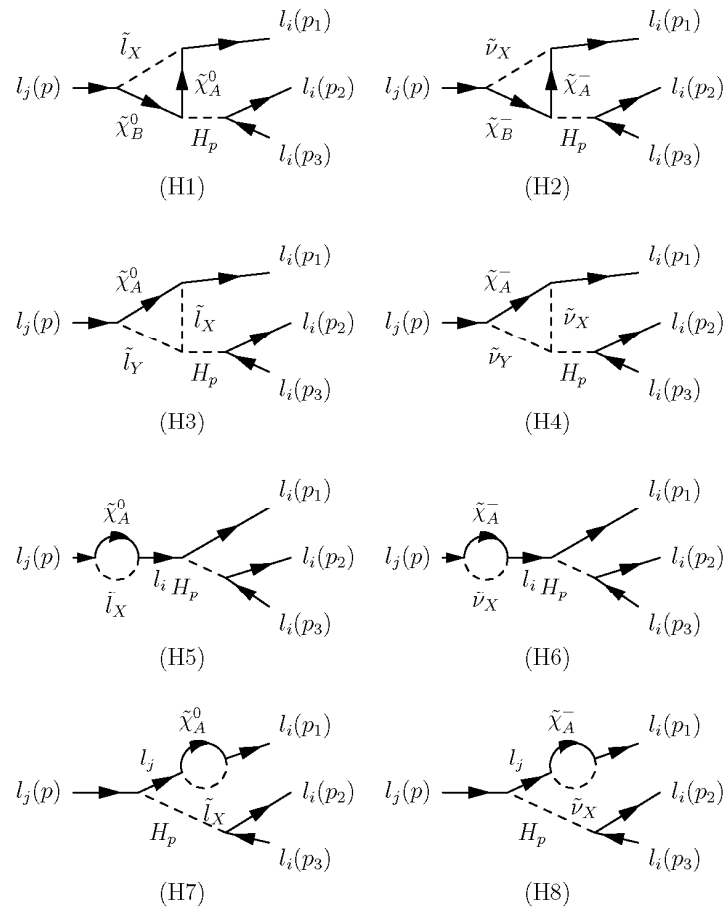


FIG. 4: Higgs-penguin diagrams contributing to the $l_j^- \rightarrow l_i^- l_i^- l_i^+$ decay. Here

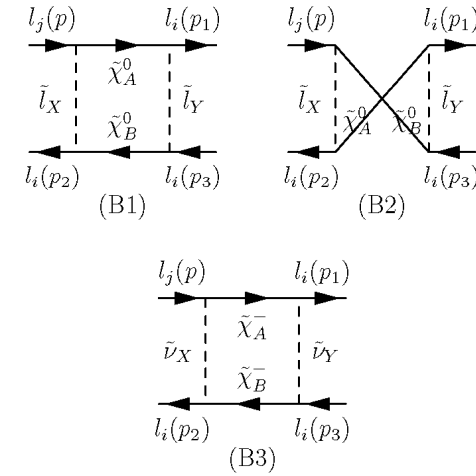
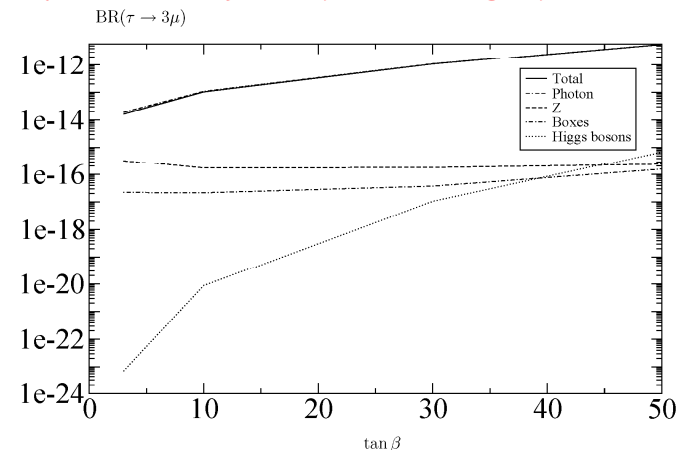


FIG. 3: Box-type diagrams contributing to the $l_j^- \rightarrow l_i^- l_i^- l_i^+$ decay

unfortunately, the photonic graphs dominate





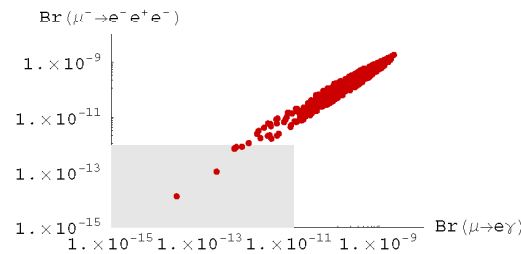
1.2 A recent example: the Littlest Higgs Model

An alternative to SUSY recently developed by Arkani-Hamed et al.

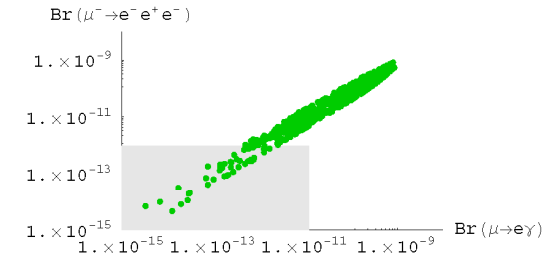
A (The ?) minimal extension of the SM "weakly coupled to new physics" at the TeV scale:

- below 1 TeV nothing changes and around 1 TeV a handful of additional particles are predicted.

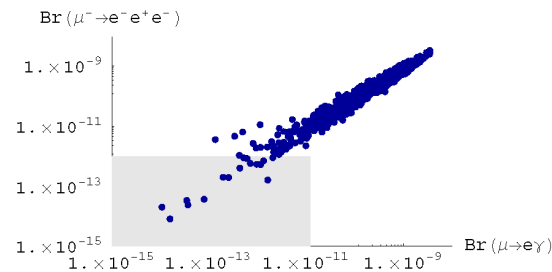
Buras et al., 2007



Scenario A



Scenario B



Scenario C

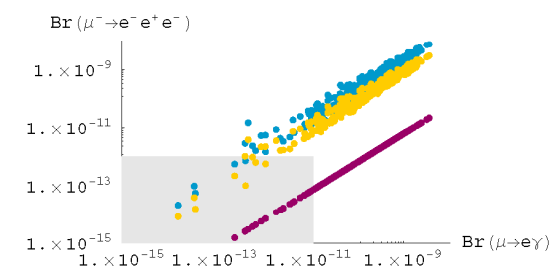


Figure 9: Correlation between $\mu \rightarrow e\gamma$ and $\mu^- \rightarrow e^-e^+e^-$ in the scenarios of Section 12.2. In the right plot of Scenario C we show the contributions to $\mu^- \rightarrow e^-e^+e^-$ from $\bar{D}'_{odd}{}^{\mu e}$ (purple, lowermost), $\bar{Z}'_{odd}{}^{\mu e}$ (orange, middle) and $\bar{Y}_{e,odd}{}^{\mu e}$ (light-blue, uppermost) separately. The shaded area represents the experimental constraints.



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decay	$f = 1000 \text{ GeV}$	$f = 500 \text{ GeV}$	exp. upper bound
$\mu \rightarrow e\gamma$	$1.2 \cdot 10^{-11}$ ($1 \cdot 10^{-11}$)	$1.2 \cdot 10^{-11}$ ($1 \cdot 10^{-11}$)	$1.2 \cdot 10^{-11}$ [17]
$\mu \rightarrow e e^+ e^-$	$1.0 \cdot 10^{-12}$ ($1 \cdot 10^{-12}$)	$1.0 \cdot 10^{-12}$ ($1 \cdot 10^{-12}$)	$1.0 \cdot 10^{-12}$ [42]
$\mu \text{Ti} \rightarrow e \text{Ti}$	$2 \cdot 10^{-10}$ ($5 \cdot 10^{-12}$)	$4 \cdot 10^{-11}$ ($5 \cdot 10^{-12}$)	$4.3 \cdot 10^{-12}$ [29]
$\tau \rightarrow e\gamma$	$8 \cdot 10^{-10}$ ($7 \cdot 10^{-10}$)	$1 \cdot 10^{-8}$ ($1 \cdot 10^{-8}$)	$9.4 \cdot 10^{-8}$ [33]
$\tau \rightarrow \mu\gamma$	$8 \cdot 10^{-10}$ ($8 \cdot 10^{-10}$)	$2 \cdot 10^{-8}$ ($1 \cdot 10^{-8}$)	$1.6 \cdot 10^{-8}$ [33]
$\tau^- \rightarrow e^- e^+ e^-$	$7 \cdot 10^{-10}$ ($6 \cdot 10^{-10}$)	$2 \cdot 10^{-8}$ ($2 \cdot 10^{-8}$)	$2.0 \cdot 10^{-7}$ [71]
$\tau^- \rightarrow \mu^- \mu^+ \mu^-$	$7 \cdot 10^{-10}$ ($6 \cdot 10^{-10}$)	$3 \cdot 10^{-8}$ ($3 \cdot 10^{-8}$)	$1.9 \cdot 10^{-7}$ [71]
$\tau^- \rightarrow e^- \mu^+ \mu^-$	$5 \cdot 10^{-10}$ ($5 \cdot 10^{-10}$)	$2 \cdot 10^{-8}$ ($2 \cdot 10^{-8}$)	$2.0 \cdot 10^{-7}$ [72]
$\tau^- \rightarrow \mu^- e^+ e^-$	$5 \cdot 10^{-10}$ ($5 \cdot 10^{-10}$)	$2 \cdot 10^{-8}$ ($2 \cdot 10^{-8}$)	$1.9 \cdot 10^{-7}$ [72]
$\tau^- \rightarrow \mu^- e^+ \mu^-$	$5 \cdot 10^{-14}$ ($3 \cdot 10^{-14}$)	$2 \cdot 10^{-14}$ ($2 \cdot 10^{-14}$)	$1.3 \cdot 10^{-7}$ [71]
$\tau^- \rightarrow e^- \mu^+ e^-$	$5 \cdot 10^{-14}$ ($3 \cdot 10^{-14}$)	$2 \cdot 10^{-14}$ ($2 \cdot 10^{-14}$)	$1.1 \cdot 10^{-7}$ [71]
$\tau \rightarrow \mu\pi$	$2 \cdot 10^{-9}$ ($2 \cdot 10^{-9}$)	$5.8 \cdot 10^{-8}$ ($5.8 \cdot 10^{-8}$)	$5.8 \cdot 10^{-8}$ [33]
$\tau \rightarrow e\pi$	$2 \cdot 10^{-9}$ ($2 \cdot 10^{-9}$)	$4.4 \cdot 10^{-8}$ ($4.4 \cdot 10^{-8}$)	$4.4 \cdot 10^{-8}$ [33]
$\tau \rightarrow \mu\eta$	$6 \cdot 10^{-10}$ ($6 \cdot 10^{-10}$)	$2 \cdot 10^{-8}$ ($2 \cdot 10^{-8}$)	$5.1 \cdot 10^{-8}$ [33]
$\tau \rightarrow e\eta$	$6 \cdot 10^{-10}$ ($6 \cdot 10^{-10}$)	$2 \cdot 10^{-8}$ ($2 \cdot 10^{-8}$)	$4.5 \cdot 10^{-8}$ [33]
$\tau \rightarrow \mu\eta'$	$7 \cdot 10^{-10}$ ($7 \cdot 10^{-10}$)	$3 \cdot 10^{-8}$ ($3 \cdot 10^{-8}$)	$5.3 \cdot 10^{-8}$ [33]
$\tau \rightarrow e\eta'$	$7 \cdot 10^{-10}$ ($7 \cdot 10^{-10}$)	$3 \cdot 10^{-8}$ ($3 \cdot 10^{-8}$)	$9.0 \cdot 10^{-8}$ [33]
$K_L \rightarrow \mu e$	$4 \cdot 10^{-13}$ ($2 \cdot 10^{-13}$)	$3 \cdot 10^{-14}$ ($3 \cdot 10^{-14}$)	$4.7 \cdot 10^{-12}$ [50]
$K_L \rightarrow \pi^0 \mu e$	$4 \cdot 10^{-15}$ ($2 \cdot 10^{-15}$)	$5 \cdot 10^{-16}$ ($5 \cdot 10^{-16}$)	$6.2 \cdot 10^{-9}$ [73]
$B_d \rightarrow \mu e$	$5 \cdot 10^{-16}$ ($2 \cdot 10^{-16}$)	$9 \cdot 10^{-17}$ ($9 \cdot 10^{-17}$)	$1.7 \cdot 10^{-7}$ [74]
$B_s \rightarrow \mu e$	$5 \cdot 10^{-15}$ ($2 \cdot 10^{-15}$)	$9 \cdot 10^{-16}$ ($9 \cdot 10^{-16}$)	$6.1 \cdot 10^{-6}$ [75]
$B_d \rightarrow \tau e$	$3 \cdot 10^{-11}$ ($2 \cdot 10^{-11}$)	$2 \cdot 10^{-10}$ ($2 \cdot 10^{-10}$)	$1.1 \cdot 10^{-4}$ [76]
$B_s \rightarrow \tau e$	$2 \cdot 10^{-10}$ ($2 \cdot 10^{-10}$)	$2 \cdot 10^{-9}$ ($2 \cdot 10^{-9}$)	—
$B_d \rightarrow \tau \mu$	$3 \cdot 10^{-11}$ ($3 \cdot 10^{-11}$)	$3 \cdot 10^{-10}$ ($3 \cdot 10^{-10}$)	$3.8 \cdot 10^{-5}$ [76]
$B_s \rightarrow \tau \mu$	$2 \cdot 10^{-10}$ ($2 \cdot 10^{-10}$)	$3 \cdot 10^{-9}$ ($3 \cdot 10^{-9}$)	—

Table 2: Upper bounds on LFV decay branching ratios in the LHT model, for two different values of the scale f , after imposing the constraints on $\mu \rightarrow e\gamma$ and $\mu^- \rightarrow e^- e^+ e^-$. The numbers given in brackets are obtained after imposing the additional constraint $R(\mu \text{Ti} \rightarrow e \text{Ti}) < 5 \cdot 10^{-12}$. For $f = 500 \text{ GeV}$, also the bounds on $\tau \rightarrow \mu\pi, e\pi$ have been included. The current experimental upper bounds are also given.

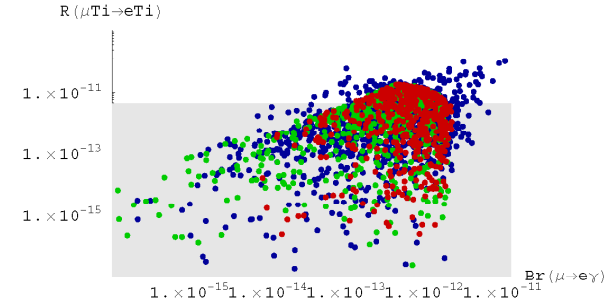


Figure 10: $\mu - e$ conversion rate in ${}^{48}_{22}\text{Ti}$ as a function of $\text{Br}(\mu \rightarrow e\gamma)$, after imposing the existing constraints on $\mu \rightarrow e\gamma$ and $\mu^- \rightarrow e^- e^+ e^-$. The shaded area represents the current experimental upper bound on $R(\mu \text{Ti} \rightarrow e \text{Ti})$.

ratio	LHT	MSSM (dipole)	MSSM (Higgs)
$\frac{\text{Br}(\mu^- \rightarrow e^- e^+ e^-)}{\text{Br}(\mu \rightarrow e\gamma)}$	0.4...2.5	$\sim 6 \cdot 10^{-3}$	$\sim 6 \cdot 10^{-3}$
$\frac{\text{Br}(\tau^- \rightarrow e^- e^+ e^-)}{\text{Br}(\tau \rightarrow e\gamma)}$	0.4...2.3	$\sim 1 \cdot 10^{-2}$	$\sim 1 \cdot 10^{-2}$
$\frac{\text{Br}(\tau^- \rightarrow \mu^- \mu^+ \mu^-)}{\text{Br}(\tau \rightarrow \mu\gamma)}$	0.4...2.3	$\sim 2 \cdot 10^{-3}$	0.06...0.1
$\frac{\text{Br}(\tau^- \rightarrow e^- \mu^+ \mu^-)}{\text{Br}(\tau \rightarrow e\gamma)}$	0.3...1.6	$\sim 2 \cdot 10^{-3}$	0.02...0.04
$\frac{\text{Br}(\tau^- \rightarrow \mu^- e^+ e^-)}{\text{Br}(\tau \rightarrow \mu\gamma)}$	0.3...1.6	$\sim 1 \cdot 10^{-2}$	$\sim 1 \cdot 10^{-2}$
$\frac{\text{Br}(\tau^- \rightarrow e^- e^+ e^-)}{\text{Br}(\tau^- \rightarrow e^- \mu^+ \mu^-)}$	1.3...1.7	~ 5	0.3...0.5
$\frac{\text{Br}(\tau^- \rightarrow \mu^- \mu^+ \mu^-)}{\text{Br}(\tau^- \rightarrow \mu^- e^+ e^-)}$	1.2...1.6	~ 0.2	5...10
$\frac{R(\mu \text{Ti} \rightarrow e \text{Ti})}{\text{Br}(\mu \rightarrow e\gamma)}$	$10^{-2} \dots 10^2$	$\sim 5 \cdot 10^{-3}$	0.08...0.15

Table 3: Comparison of various ratios of branching ratios in the LHT model and in the MSSM without and with significant Higgs contributions.



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Programm $\mu \rightarrow 3e$ Workshop (<https://midas.psi.ch/eLogs/MEEE/>)

Welcome

Stefan Ritt

PSI

purpose of the exercise

Motivation

Andries van der Schaaf

UZH

$\mu \rightarrow 3e$ v.s. $\mu \rightarrow e\gamma$ v.s. $\mu - e$ conversion

SINDRUM I

Willi Bertl

PSI

the best result since 20 years

Design criteria for a new $\mu \rightarrow 3e$ experiment

Andries van der Schaaf

UZH

limitations to the sensitivity

Ideas for a new $\mu \rightarrow e^+e^+e^-$ experiment

Roland Horisberger

PSI

a large radial TPC with fine-grained readout

Π E5 beam line

Peter-Raymond Kettle

PSI

the MEG experience

Active targets IKAR & MAYA

Oleg Kiselev

PSI

alternatives mainly for heavy fragments

MuCAP TPC

Malte Hildebrandt

PSI

a TPC based on hydrogen

Geiger mode APDs

Dieter Renker

PSI

from strips to pads for the plastic scintillator



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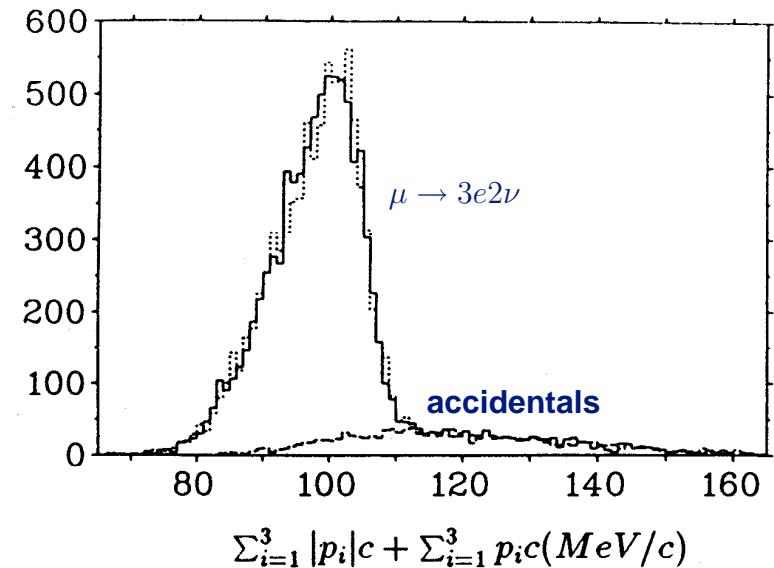
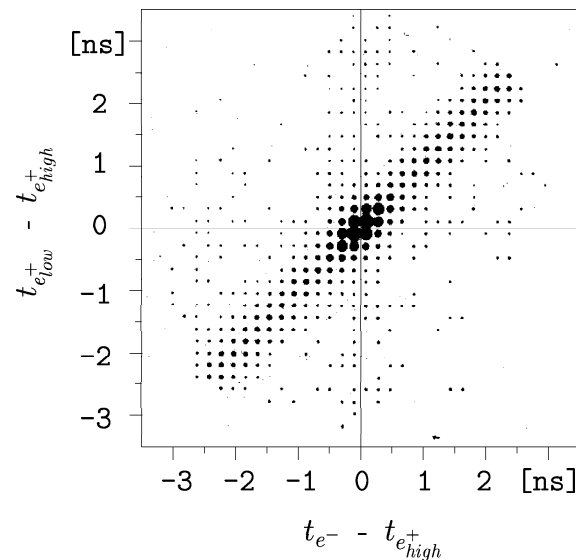
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2 signature $\mu \rightarrow 3e$ at rest

- total energy, total momentum (\rightarrow coplanarity).
- Phase space distribution gives additional information if observed.
- In a constant B field the acceptance is defined by the p_t threshold.

3 background

SINDRUM I



Accidental background

involves low invariant mass e^+e^- pairs produced by photons or by Bhabha scattering.



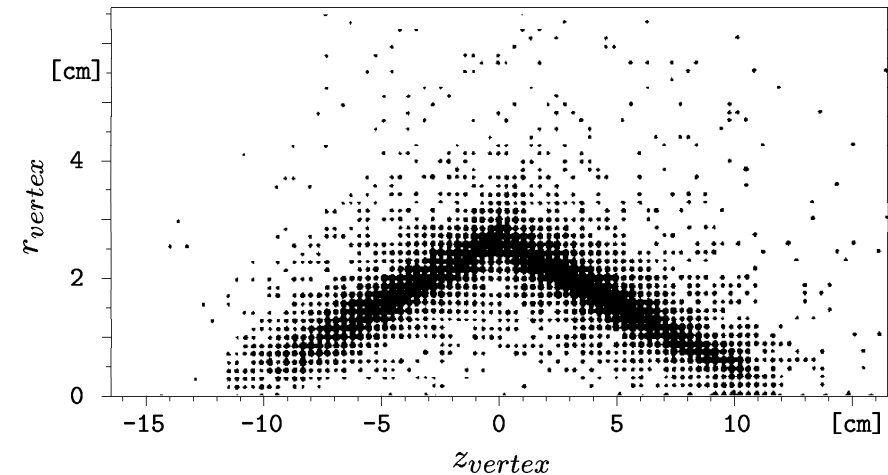
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Suppressing accidental background:

- The three trajectories meet in a common vertex.
- The common vertex has to be in a muon-stop region. For this reason SINDRUM I used a relatively large surface target.
- An active target could lead to a dramatic suppression since one would know the interaction point of γ conversions and Bhabha scatterings. ³



³Peter Kammel is gratefully acknowledged for bringing this up



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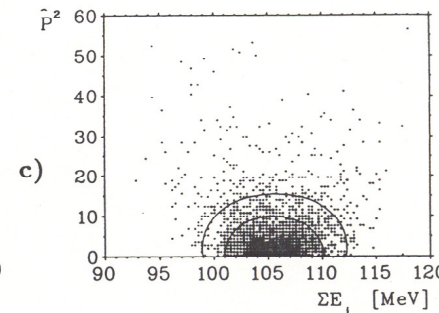
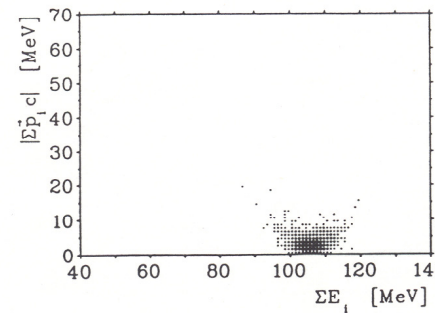
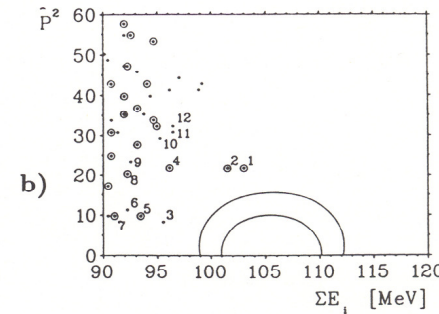
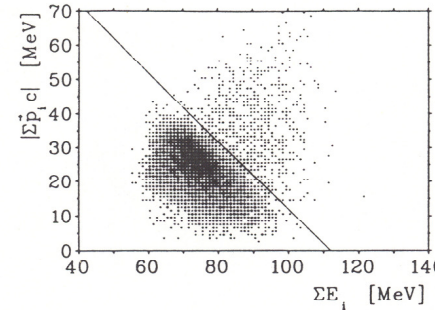
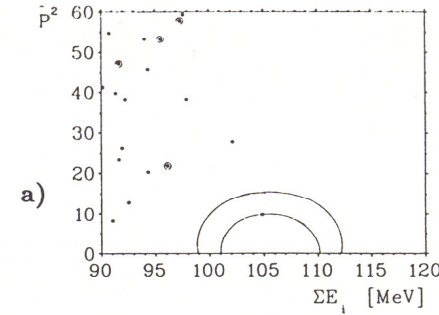
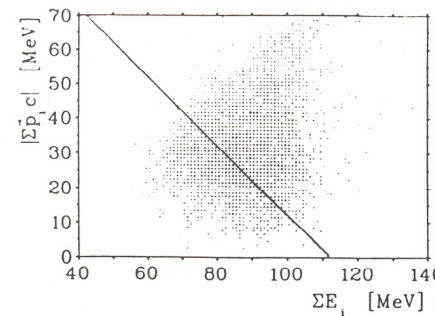
3.1 How to reach a single-event sensitivity of $O(10^{-16})$?

- Measure 100 instead of 10 weeks.
- Raise stop rate from 5×10^6 to $10^9/s$.
- Lower threshold on p_t to gain in acceptance.

χ^2 is a test of the $e^+e^+e^-$ correlation based on time and vertex variables

$$\hat{P}^2 \equiv \left(\frac{P_{\parallel}}{\sigma_{P_{\parallel}}} \right)^2 + \left(\frac{P_{\perp}}{\sigma_{P_{\perp}}} \right)^2$$

\parallel and \perp are defined w.r.t. the decay plane.





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3.2 What about background ?

reducing accidental background by improving detector resolutions

assumption ^a	gain factor	background
SINDRUM I	1	40000
$\Delta t \times 0.25$	4	10000
vertex $\times 0.5$	4	2500
energy $\times 0.5$	2	1250
momentum $\times 0.5$	4	300
target size $\times 2$	2	150
target mass/area $\times 0.5$	2	75

^afor example by linear scaling the detector by factor 2

So one would need an additional factor 100.

A vertex detector would do the job if it would stand the rate.



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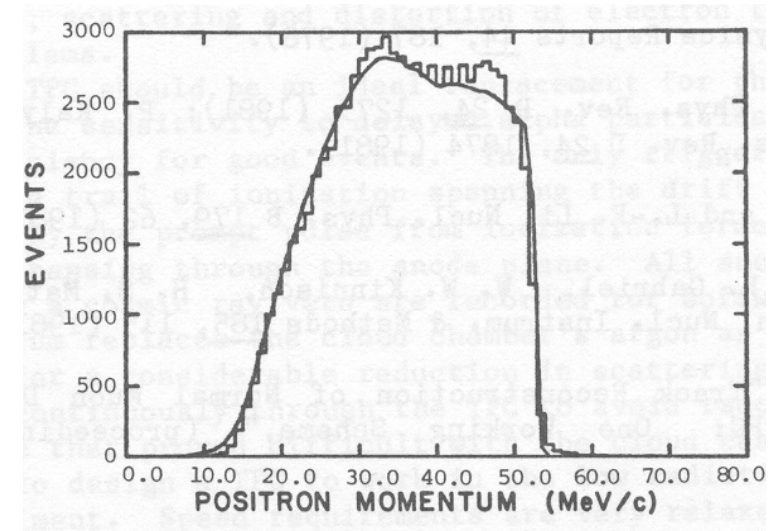
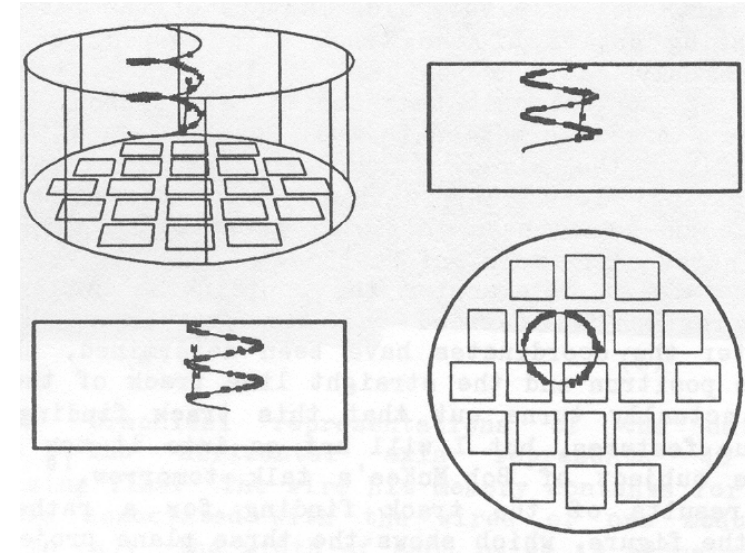
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3.3 1985: LAMPF TPC

The Time Projection Chamber

AIP Conference Proceedings 108, ed. J.A. Macdonald
contributions by W.W. Kinnison and R.J. McKee

- six authors!
- diameter 122 cm, length 55 cm
- Both the incoming surface muon and its decay positron are observed.
- momentum resolution 1%





$a \mu \rightarrow 3e$ search at 10^{-16} ?

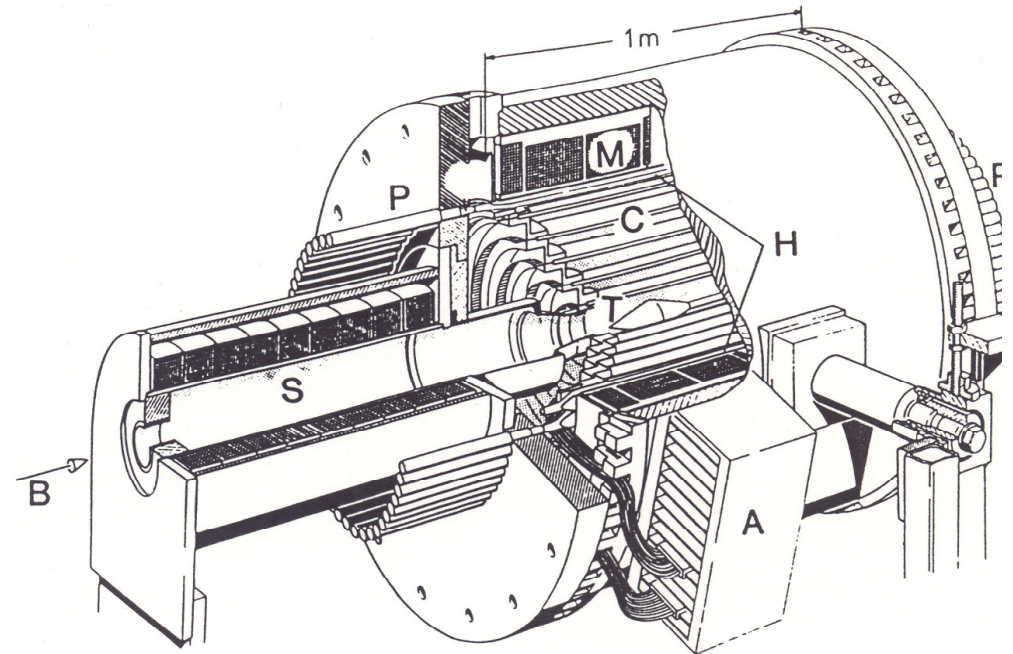
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3.4 Detector issues

SINDRUM I

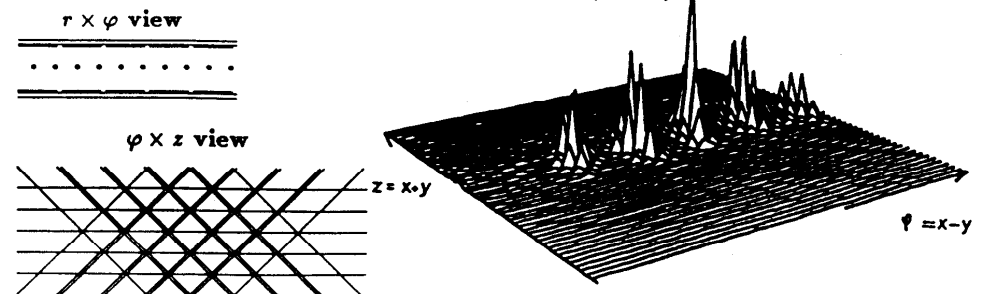
- B* beam
- S* focussing solenoid
- T* hollow target
- C* MWPC tracking
- H* plastic hodoscope



Events triggered with an ultra-thin scintillator.

- Cathodes image the avalanches at the anodes.*
- Phi resolution given by number of anode wires.*
- z resolution 0.2 mm.*

MWPC structure (total mass 30 mg/cm²)





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Could one stand the rate?

- **extra tracks, combinatorial background**

SINDRUM I saw about 0.1 extra track per event at the 50 - 70 ns gating time. If the detector would be twice faster there would be 10 random tracks. No problem with sufficient granularity (at least 500 anode wires and cathode strips per plane).

SINDRUM I v.s. MEGA

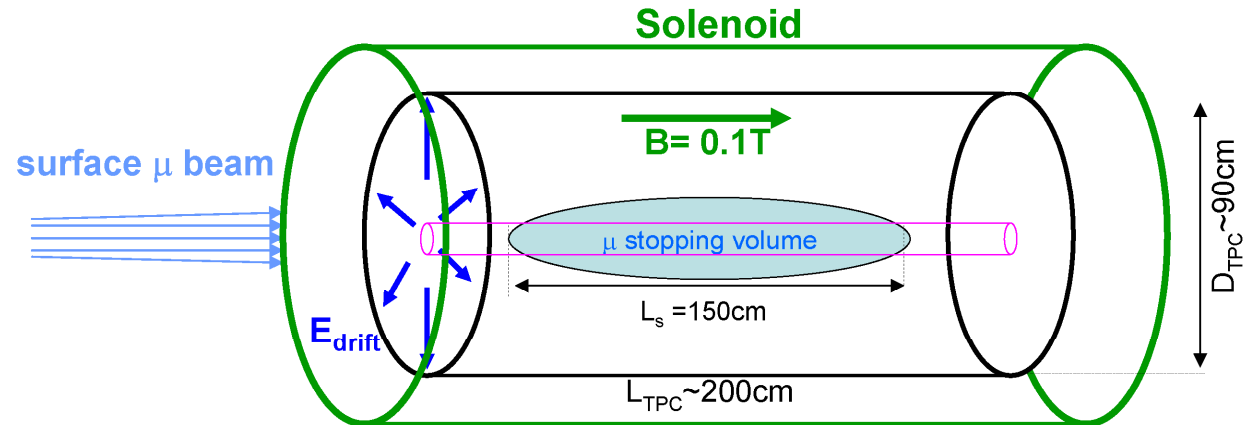
	SINDRUM I	MEGA	
self-supporting	yes	no	
thickness	10^{-3}	0.3×10^{-3}	rad. length
wire spacing	2	1.3	mm
gas	Ar-C ₂ H ₆ (50-50)	CF ₄ -C ₄ H ₁₀ (80-20)	
gate width	60	30	ns
turns/helix	≈ 1	≈ 5	
peak stop rate	5×10^6	2.5×10^8	1/s
rate per anode	10^5	10^7	1/s
max. fluence	3×10^2	4×10^4	1/mm ² .s

Conclusion: it could work



4 A radial TPC ?

(Roland Horisberger)

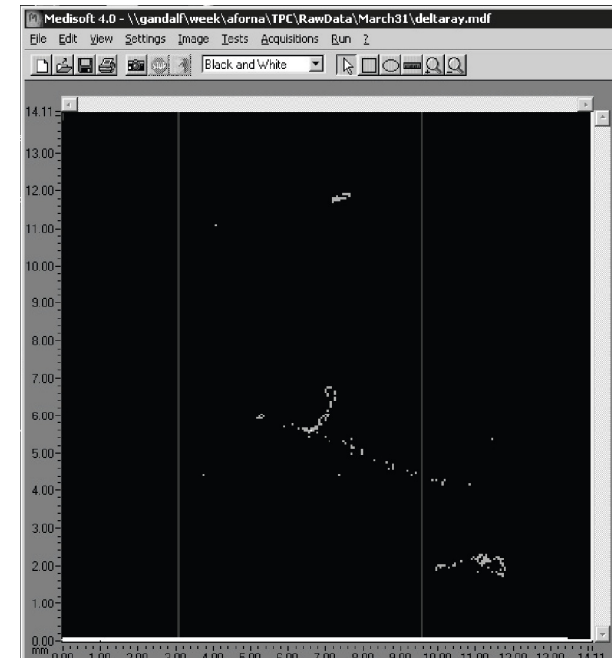


Micro-pattern readout schemes as studied by LCTPC and CERN RD51 (5 years starting now, Geneva is in) would:

- match the intrinsic precision offered by TPC's,
- stand high particle fluxes by suppression of ion back-flow,
- allow curved structures for radial drift field.

delta electron imaged by LCTPC prototype

$14 \times 14 \text{ mm}^2$





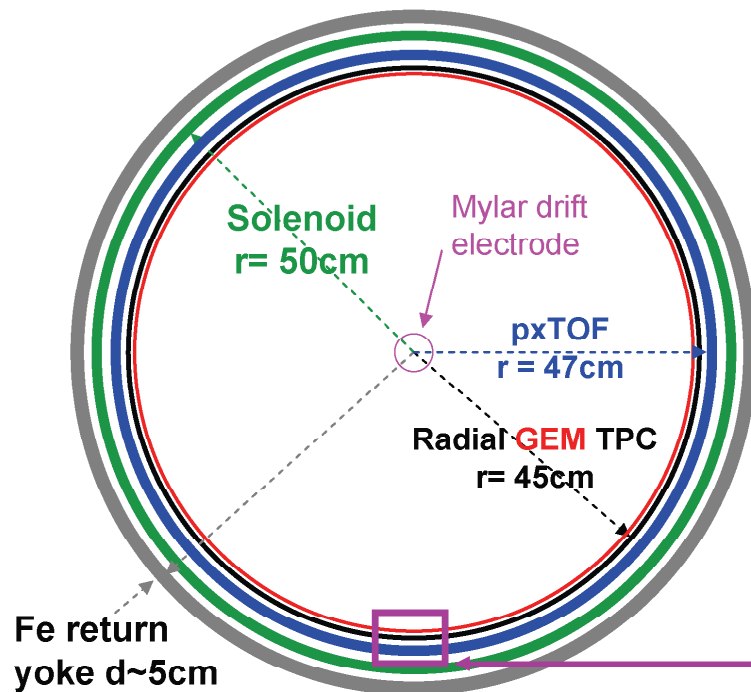
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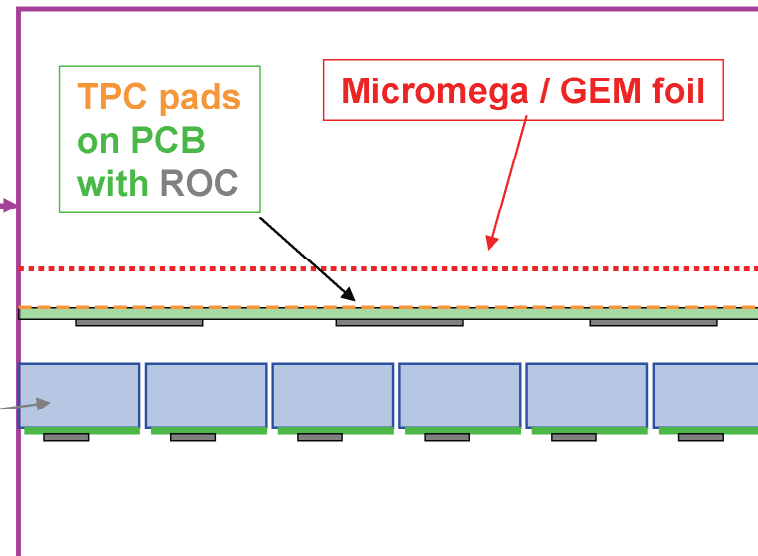


Cross-section of $\mu \rightarrow 3e$ Experiment



Inner surface of solenoid ($\sim 6\text{m}^2$) is equipped:

- 1) **GEM or Micromega TPC Pad readout**
0.75mm x 0,75mm pad size, gain $\sim 3\text{-}4\text{K}$
resolution $\sim 200\mu$ in x,y & t (ILC exp.)
few mm double track resolution
- 2) **Scint. Pixel TOF system**
1cm x 1cm time resolution $\sim 50\text{psec}$



Scintillator based pixel TOF with Si-PM & TDC ROC



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4.1 Open issues

- What is the highest beam intensity that PSI can deliver in 5 years?
Proton current 2→3 mA, optimized target geometry.
- How harmful is loss of central region? One would like to see the $e^+e^+e^-$ vertex.
- A TPC is a slow device. Can events with 10^4 additional muon tracks and decay positrons be analyzed?
- Can triggering be solved? Would a second plastic layer help to trigger fast on charge?
- Would a hybrid scheme (much smaller and faster gated TPC for vertex only combined with 25 ns tracker) solve some of the above?
- Budget? Comparable to MEG?
- Sufficient interest to form an international collaboration?
No problem if MEG finds a signal!