

Right-handed currents and heavy neutrinos

Proposal for the Report on Physics at 100 TeV

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In collaboration with:

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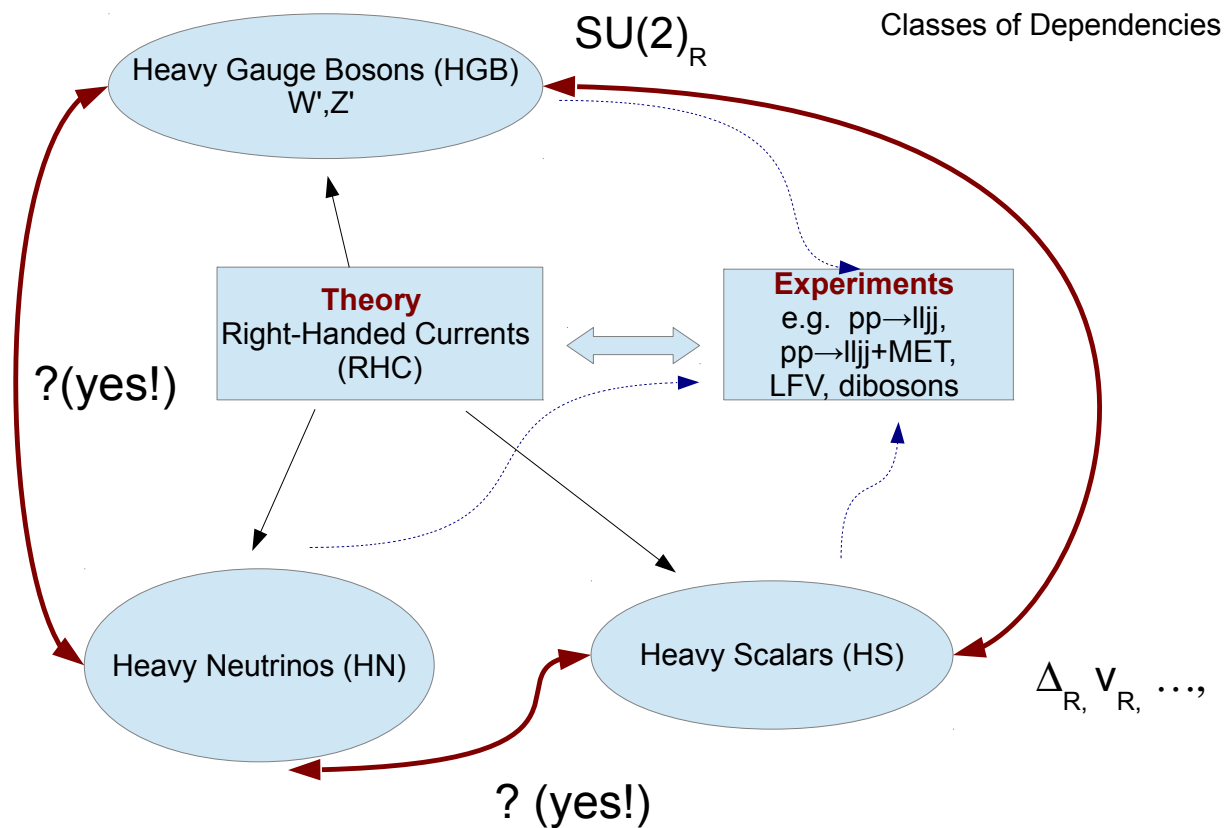
T. Jeliński, M. Kordiaczyńska (Silesia U.)

R. Szafron (Alberta U.)

CERN, 5.10.2015, Vidyo transmission

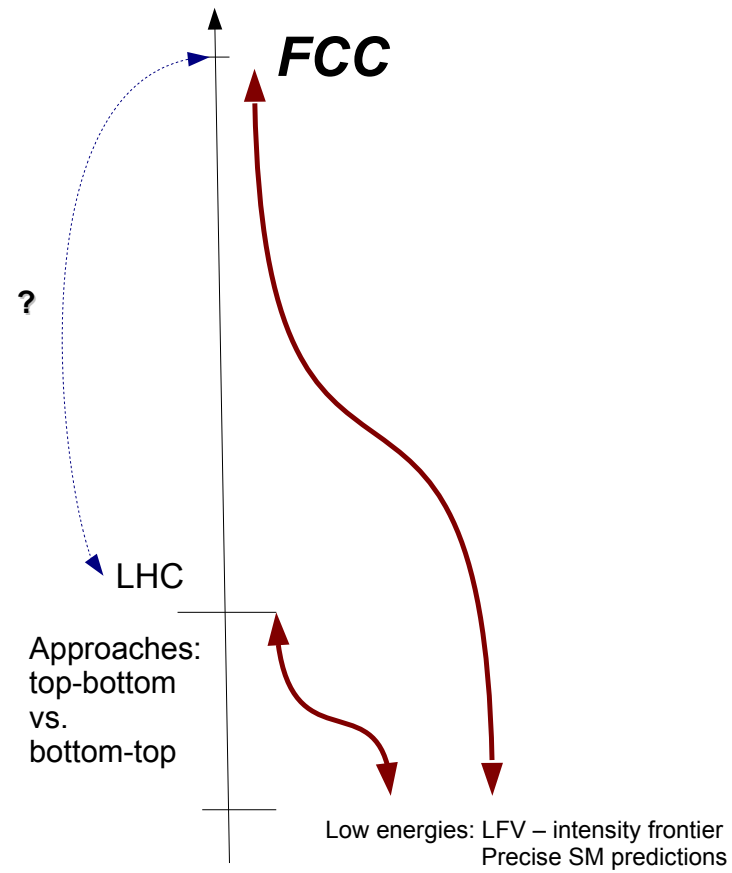
FCC-hh BSM group informal meeting, [indico](#)

RHC includes plenty of connected issues



We focus on non-supersymmetric theories

From the lowest to the high(est) energies



This talk:
 $pp \rightarrow lljj$

LHC-1 excess data

A few deviations from the SM predictions in invariant mass distributions near 2 TeV:

(i) a **3.4 σ excess at ~ 2 TeV** in the ATLAS search interpreted as a W' boson decaying into $WZ \rightarrow jj$, The mass range with significance above 2σ is ~ 1.9 – 2.1 TeV; the global significance is 2.5σ - see [pdf](#).

(ii) A CMS search for jj resonances without distinguishing between the W - and Z -tagged jets, a 1.4σ excess at ~ 1.9 TeV - see [pdf](#)

(iii)

a **2.8 σ excess in the 1.8 – 2.2 TeV** bin in the CMS search for a W' and a heavy “right-handed” neutrino, N_R , through the $W' \rightarrow N_R e \rightarrow eejj$ process - see [pdf](#)

(iv) a **2.2 σ excess in the 1.8 – 1.9 TeV** bin in the CMS search for $W' \rightarrow Wh^0$, where the SM Higgs boson, h^0 , is highly boosted and decays into $b\bar{b}$, while $W \rightarrow \ell\nu$ - see [pdf](#)

(v) a $\sim 2\sigma$ excess at ~ 1.8 TeV in the CMS dijet resonance search - see [pdf](#). The ATLAS search in the same channel has yielded only a 1σ excess at 1.8 TeV - see [pdf](#)

Theory vs experiment

Three examples:

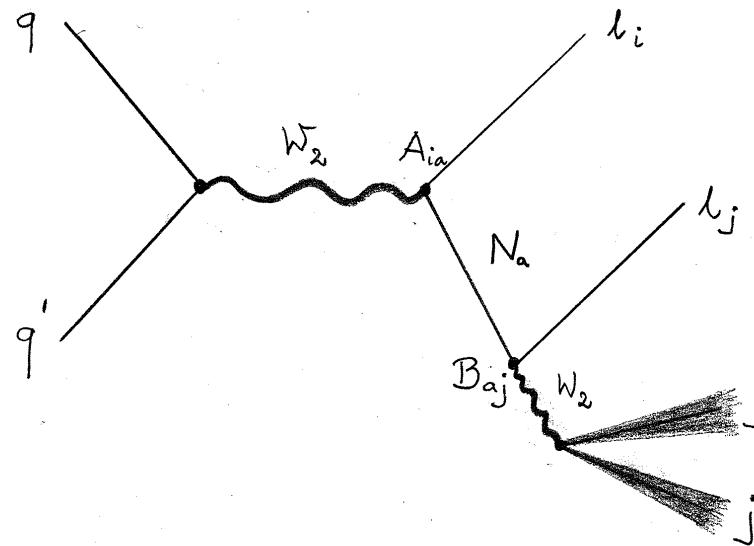
- (i) “Symmetry Restored in Dibosons at the LHC?” [Brehmer et al., arXiv:1507.00013]
- (ii) “Reconciling the 2 TeV Excesses at the LHC in a Linear Seesaw Left-Right Model” [Deppisch et al., arXiv:1508.05940]
- (iii) “Unified explanation of the $eejj$, diboson and dijet resonances at the LHC” [Bhupal Dev and Mohapatra, arXiv:1508.02277]

In (ii) linear see-saw is favorable, in (iii) inverse see-saw is favorable over type I see-saw

What is surprising about CMS data on $pp \rightarrow lljj$?

1. Ratio of opposite-sign (OS) $pp \rightarrow e^\pm e^\mp jj$ to the same-sign (SS) $pp \rightarrow e^\pm e^\pm jj$ leptons:

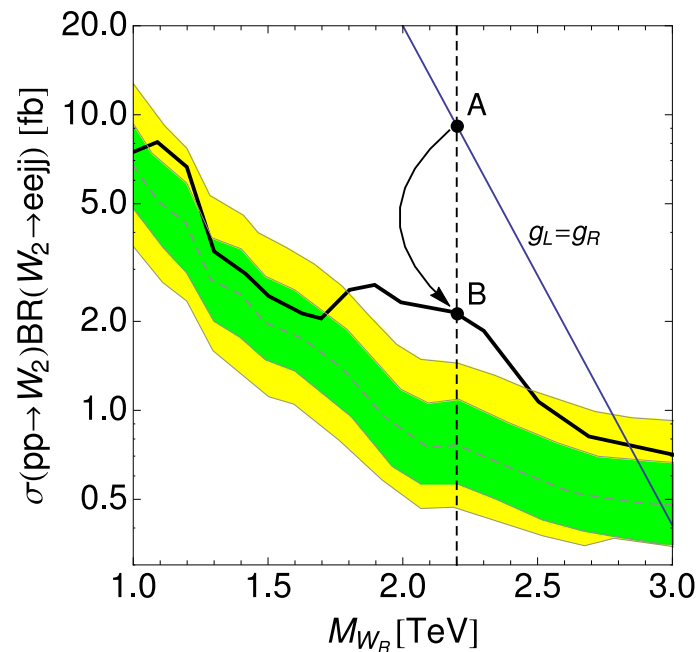
$$r_{CMS} = \frac{N_{SS}}{N_{OS}} = \frac{1}{13},$$



Almost no lepton number violation! Dirac/Pseudo-Dirac heavy neutrinos

In addition:

2. No excess in the $\mu\mu$ channel
3. Overall excess in $eejj$ production (point B), interpreted by CMS with $g_L = g_R$ (point A)



ATLAS case!

Only SS leptons
considered -
assumption of
Majorana neutrinos

Keung, Senjanovic
1983 paper ...

How to explain CMS excess data with/without RHC?

- ❖ supersymmetric models,...
 - ❖ RHC, $A \rightarrow B$: $g_R < g_L$
 - ❖ RHC, $A \rightarrow B$: take into account interferences among neutrino states
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Our point of view:

‘Heavy neutrinos and the $pp \rightarrow lljj$ CMS data,’ JG and T. Jeliński,
PLB **748** (2015) 125, e-Print: [arXiv:1504.05568](https://arxiv.org/abs/1504.05568)



All the above facts can be **reconciled** with data if RHC with **heavy neutrinos with CP phases and mixings** are taken into account (CMS takes in analysis degenerate heavy neutrinos, without mixings and CP phases),

$$\begin{aligned}\mathcal{L} \quad \supset \quad & \frac{g}{\sqrt{2}} \sum_{a=1}^3 \bar{\nu}_a \gamma^\mu P_L (U_{PMNS})_{aj} l_j W_{1\mu}^+ + \text{h.c.} \\ & + \frac{\tilde{g}}{\sqrt{2}} \sum_{a=1}^3 \bar{N}_a \gamma^\mu P_R (K_R)_{aj} l_j W_{2\mu}^+ + \text{h.c.}\end{aligned}$$

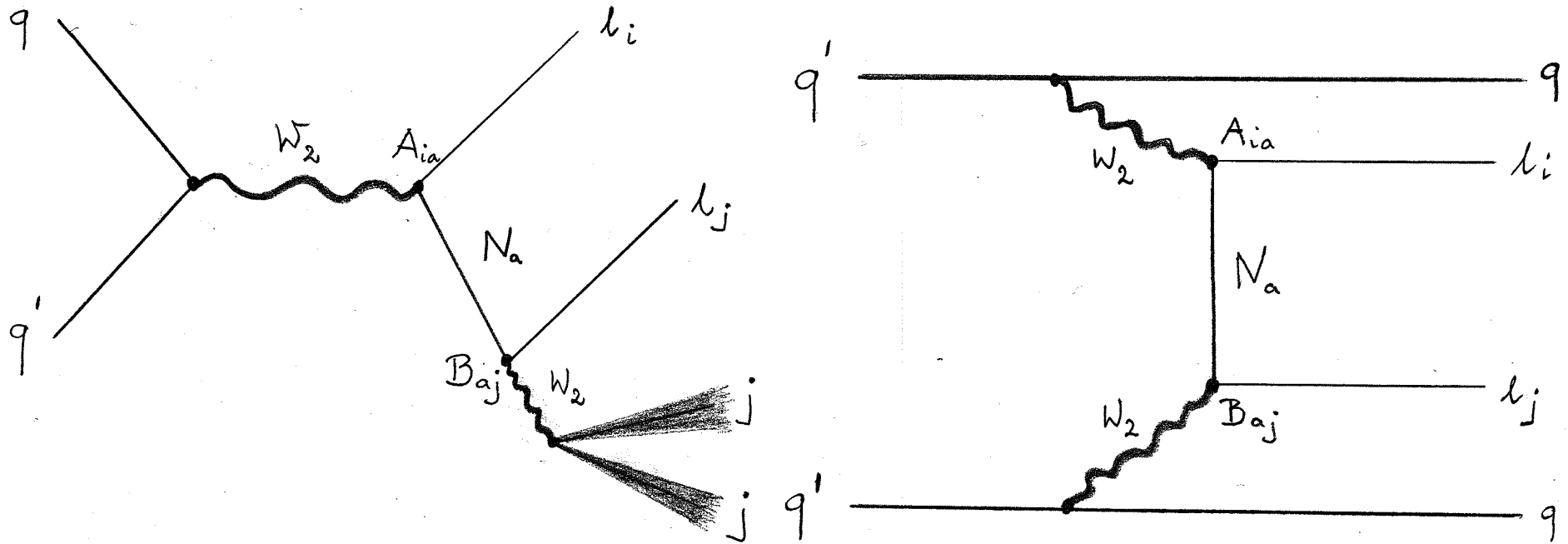
e.g.

$$M_{N_{1,3}} = 0.925 \text{ TeV}, \quad M_{N_2} = 10 \text{ TeV}, \quad K_R = \begin{pmatrix} \cos \theta_{13} & 0 & \sin \theta_{13} \\ 0 & 1 & 0 \\ -e^{i\phi_3} \sin \theta_{13} & 0 & e^{i\phi_3} \cos \theta_{13} \end{pmatrix}.$$

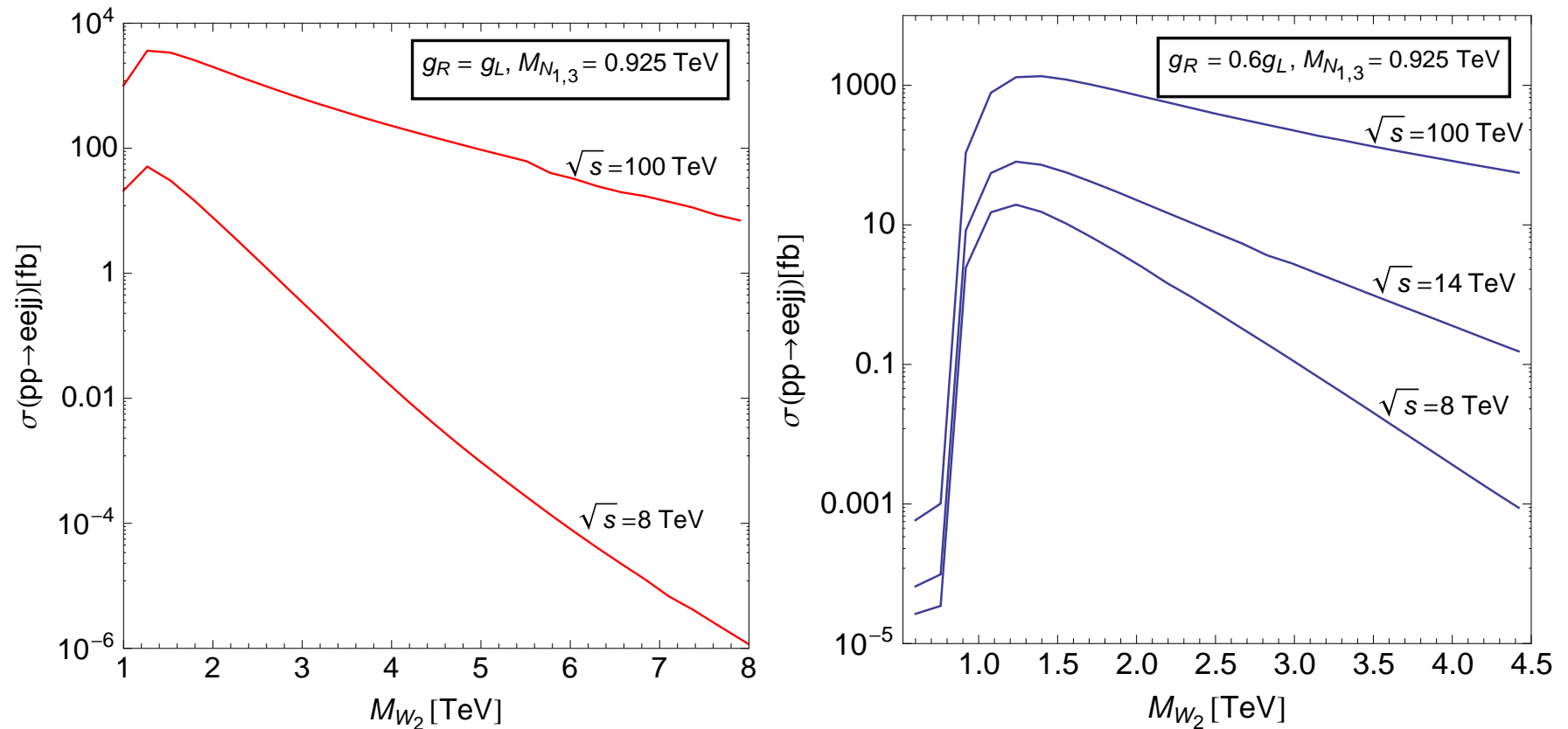
$$\mathcal{L}_{\mathcal{RHC}} = \overbrace{\left(\cos \theta_{13} \bar{N}_1 - e^{i\phi_3} \sin \theta_{13} \bar{N}_3 \right)}^{N_{eff}} \frac{\tilde{g}}{\sqrt{2}} \gamma^\mu P_R (K_R)_{aj} l_j W_{2\mu}^+$$

There is no a problem with $\mu \rightarrow e\gamma$ here (zero or very small K_{R12} elements).

Also problem with neutrinoless doubly beta decay can be avoided (destructive interferences)

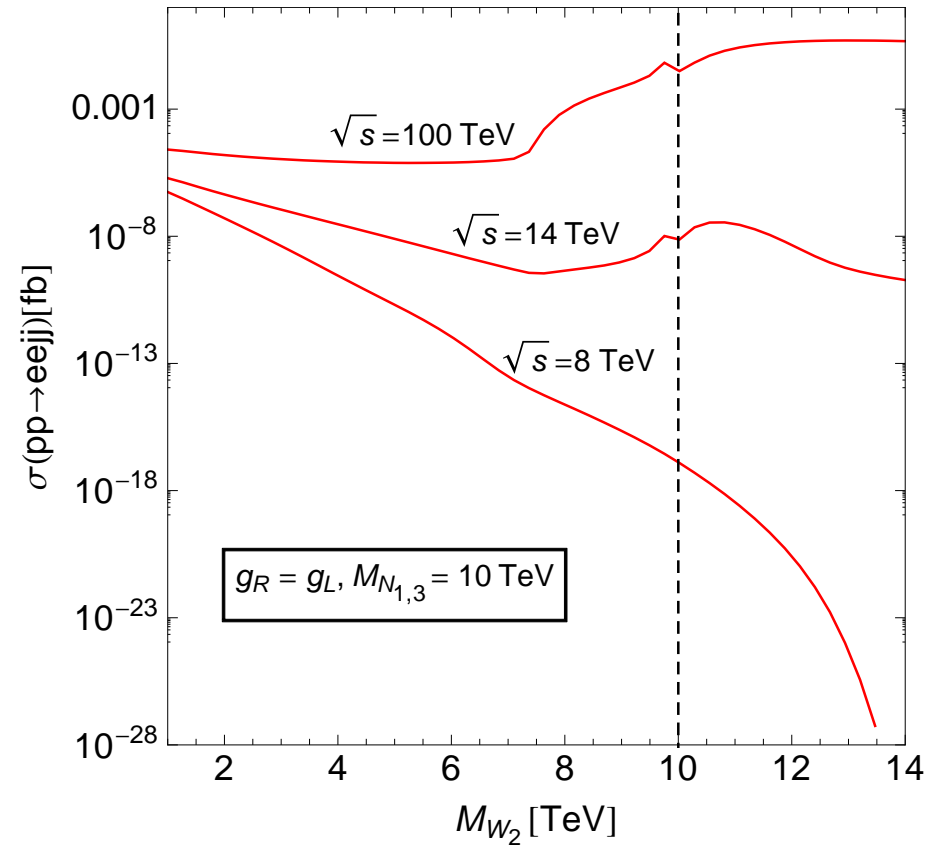
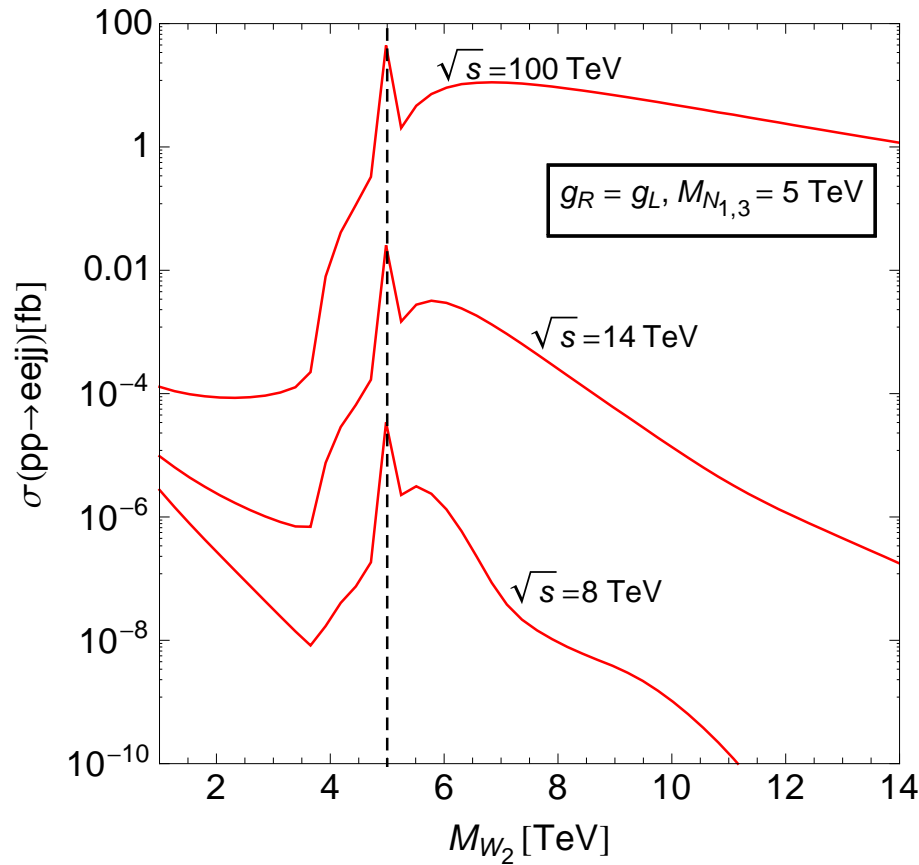


Preliminary plots, proposal for BSM section of the Report on Physics at 100 TeV



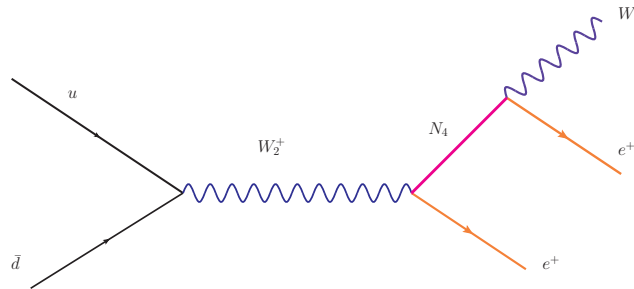
1. What about 33 TeV option?
2. We plan to tweak predictions next weeks: Cuts, Significance, other flavours,...

Preliminary plots, proposal for BSM section of the Report on Physics at 100 TeV



Remark, models consistency, radiative corrections

"Left-Right Symmetry at LHC and Precise 1-Loop Low Energy Data", J. Chakraborty et al, JHEP 1207 (2012) 038



Muon decay constrain parameter space of a model

$$\frac{G_F}{\sqrt{2}} = \frac{e^2}{8(1 - M_W^2/M_Z^2)M_W^2}(1 + \Delta r[M_N, M_H, M_{W_2}]).$$

⇒ calculate Δr in LR

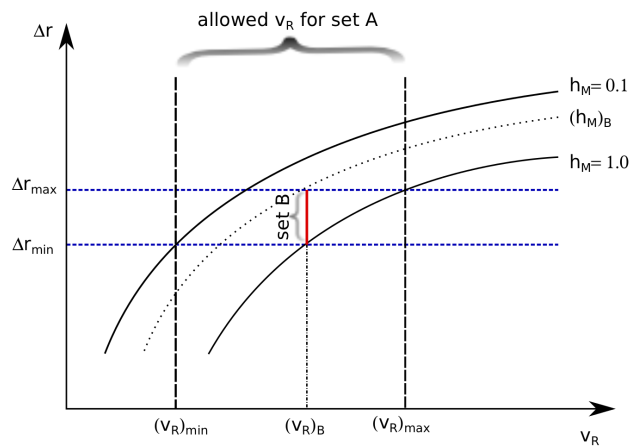
⇒ do the matching with SM

⇒ compare with data

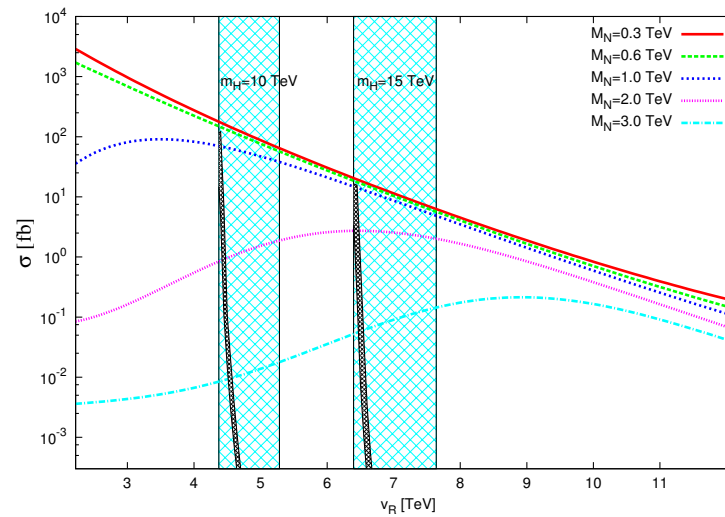
M. Czakon, J. Gluza, M. Zralek, Nucl. Phys. **B573** (2000) 57 and

M. Czakon, J. Gluza, J. Hejczyk, Nucl. Phys. **B642** (2002) 157-172.

Corrections narrow parameter space, $\sqrt{s} = 14$ TeV



$$pp \rightarrow eeW_2$$



Summary

$$pp \rightarrow lljj$$

- ❖ We are ready for preparing and delivering updated analysis for the "smoking-gun" $pp \rightarrow lljj$ process at 100 TeV, taking into account
 1. Knowledge on heavy neutrino sector in various scenarios for neutrino masses and mixings (e.g. see-saw, inverse see-saw)
 2. Low energy constraints
 3. Various RHC scenarios (e.g. $g_R \neq g_L$)
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