

An aerial photograph of a cyclotron's two semi-circular dees, which are painted a light teal color. The dees are connected to a central vertical axis. The surrounding area is filled with complex machinery, including various pipes, cables, and structural supports, all set within a large industrial facility.

# **Injection and Extraction in Cyclotrons**

**CERN Accelerator School – Specialised Course  
Erice, March 12, 2017**

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Paul Scherrer Institut

# Outline

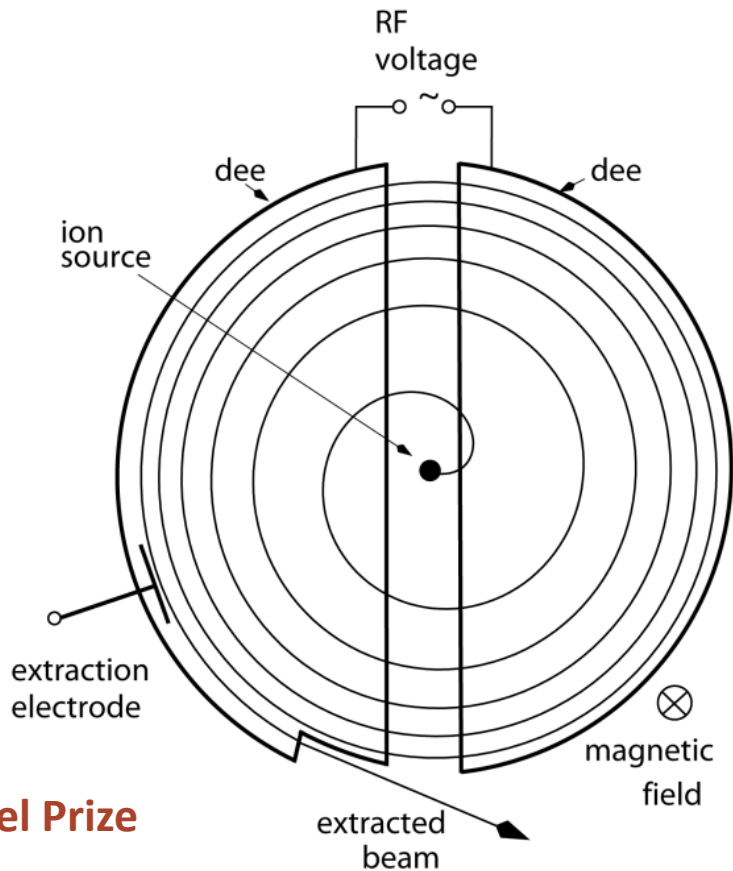
- **Cyclotron Basics**  
scaling and isochronicity, focusing, turn separation, classical cyclotrons and derived types
- **Injection for Cyclotrons**  
internal source, electrostatic deflectors, horizontal injection, optics matching, bunching
- **Extraction for Cyclotrons**  
electrostatic septum, stepwidth calculation, charge exchange extraction



# The Classical Cyclotron

two capacitive electrodes  
„Dees“, two gaps per turn  
internal ion source  
homogenous B field  
**constant revolution time**  
(for low energy,  $\gamma \approx 1$ )

$$\omega_c = \frac{eB_z}{\gamma m}$$



**invented 1930, Lawrence, Nobel Prize**

**powerful concept:**

- ➔ **simplicity, compactness**
- ➔ **continuous injection/extraction**
- ➔ **multiple usage of accelerating voltage**



# wide spectrum of cyclotrons ...

compact and cost optimized for series production  
e.g. medical nuclide production  
→ Internal source, extraction or internal target



CYCLONE 30 (IBA) : H- 15 à 30 MeV

huge and complex for variable research purposes, e.g. R.I.B. production or high intensity  
→ External source, injection



RIKEN s.c. Ring Cyclotron- „as big as a house“



# cyclotron basics: isochronicity and scalings

continuous acceleration → revolution time should stay constant, though  $E_k$ ,  $R$  vary

magnetic rigidity:

$$BR = \frac{1}{e} p = \beta\gamma \frac{m_0 c}{e}$$

orbit radius from isochronicity:

$$R = \frac{c}{\omega_c} \beta = R_\infty \beta$$

deduced scaling of  $B$ :

$$R \propto \beta; BR \propto \beta\gamma \longrightarrow B(R) \propto \gamma(R)$$

**thus, to keep the isochronous condition,  $B$  must be raised in proportion to  $\gamma(R)$ ; this contradicts the focusing requirements!**

$$\begin{aligned} \text{field index } k &= \frac{R}{B} \frac{dB}{dR} \\ &= \frac{\beta}{\gamma} \frac{d\gamma}{d\beta} \\ &= \gamma^2 - 1 \end{aligned}$$



# cyclotron basics: stepwidth (nonrelativistic, B const)

relation between  
energy and radius

$$qRB_z = \sqrt{2mE_k}$$
$$\frac{dR}{R} = \frac{1}{2} \frac{dE_k}{E_k}$$

use:  $\Delta E_k = \text{const}; B_z = \text{const}; E_k \propto R^2$

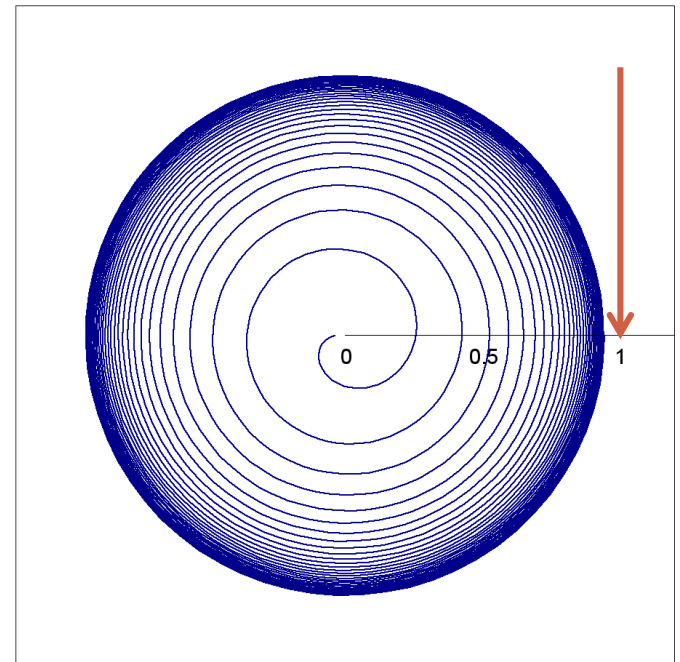
thus:

$$\Delta R \propto \frac{R}{E_k} \propto \frac{1}{R}$$

radius increment per turn  
decreases with increasing radius  
→ **extraction becomes more and more difficult at higher energies**

“cyclotron  
language”

$$R_\infty = R/\beta$$



# focusing in a cyclotron

centrifugal force  $mv^2/r$



Lorentz force  $qv \times B$



$$m\ddot{r} = mr\dot{\theta}^2 - qr\dot{\theta}B_z$$

focusing: consider small deviations  $x$  from beam orbit  $R$  ( $r = R+x$ ):

$$\ddot{x} + \frac{q}{m}vB_z(R+x) - \frac{v^2}{R+x} = 0,$$

$$\ddot{x} + \omega_c^2(1+k)x = 0.$$

using :

$$\omega_c = qB_z/m = v/R$$

$$r\dot{\theta} \approx v$$

$$k = \frac{R}{B} \frac{dB}{dR}$$

**thus in radial plane:**

$$\omega_r = \omega_c \sqrt{1+k} = \omega_c \nu_r$$

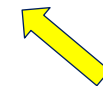
$$\nu_r = \sqrt{1+k}$$

$$\approx \gamma$$

using isochronicity condition

**in vertical plane:**

$$\nu_z = \sqrt{-k}$$

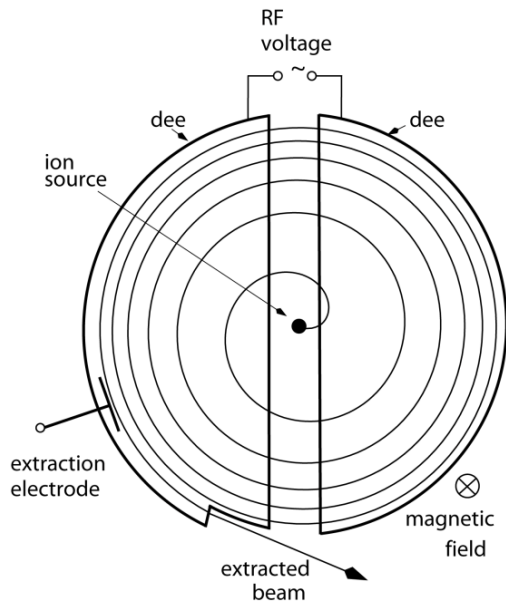


$k < 0$  to obtain vertical focus.



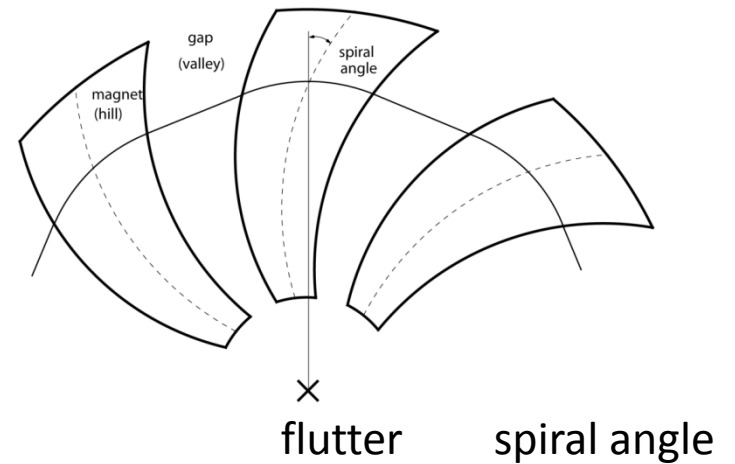
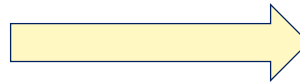
# Classical vs Isochronous Cyclotron

classical cyclotron



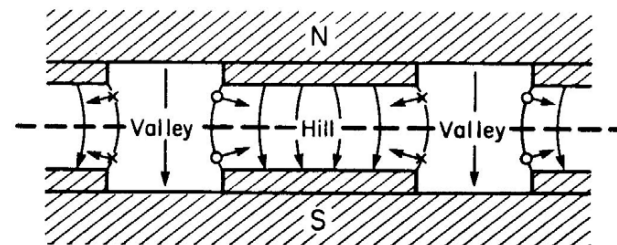
- insufficient vertical focusing
- limited energy reach

Sector/AVF cyclotron



$$v_z^2 = -\frac{R}{B_z} \frac{dB_z}{dR} + F^2(1 + 2 \tan^2 \delta)$$

$$F^2 = \frac{\overline{B_z^2} - \overline{B_z}^2}{\overline{B_z}^2}$$

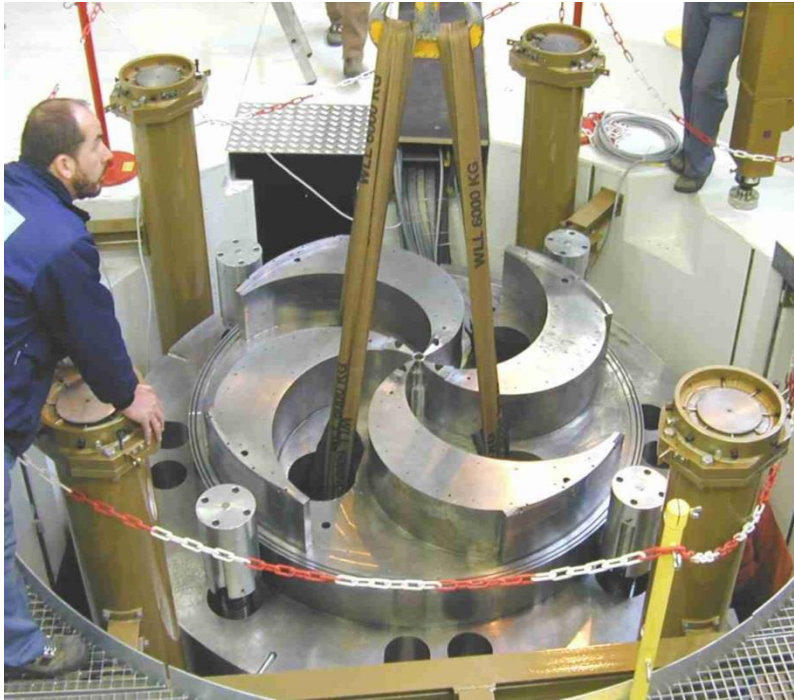


[illustration of focusing at edges]



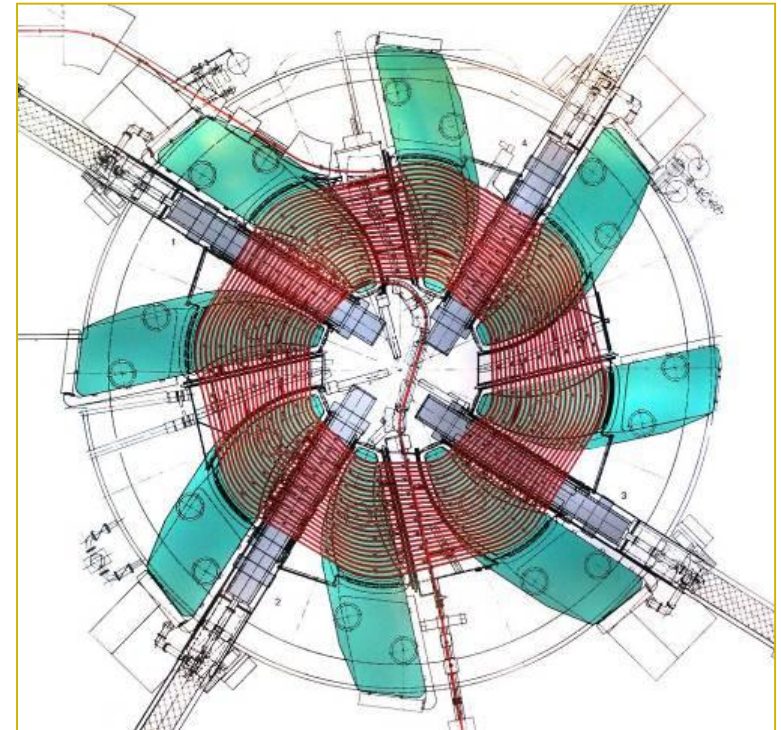


# Azimuthally Varying Field vs. Separated Sector Cyclotrons



PSI/Varian comet: 250MeV sc. medical cyclotron

- **AVF = single pole with shaping**
- often **spiral poles** used
- **internal source** possible
- **D-type RF electrodes**, rel. low energy gain
- **compact**, cost effective
- depicted Varian cyclotron: 80% extraction efficiency; **not suited for high power**

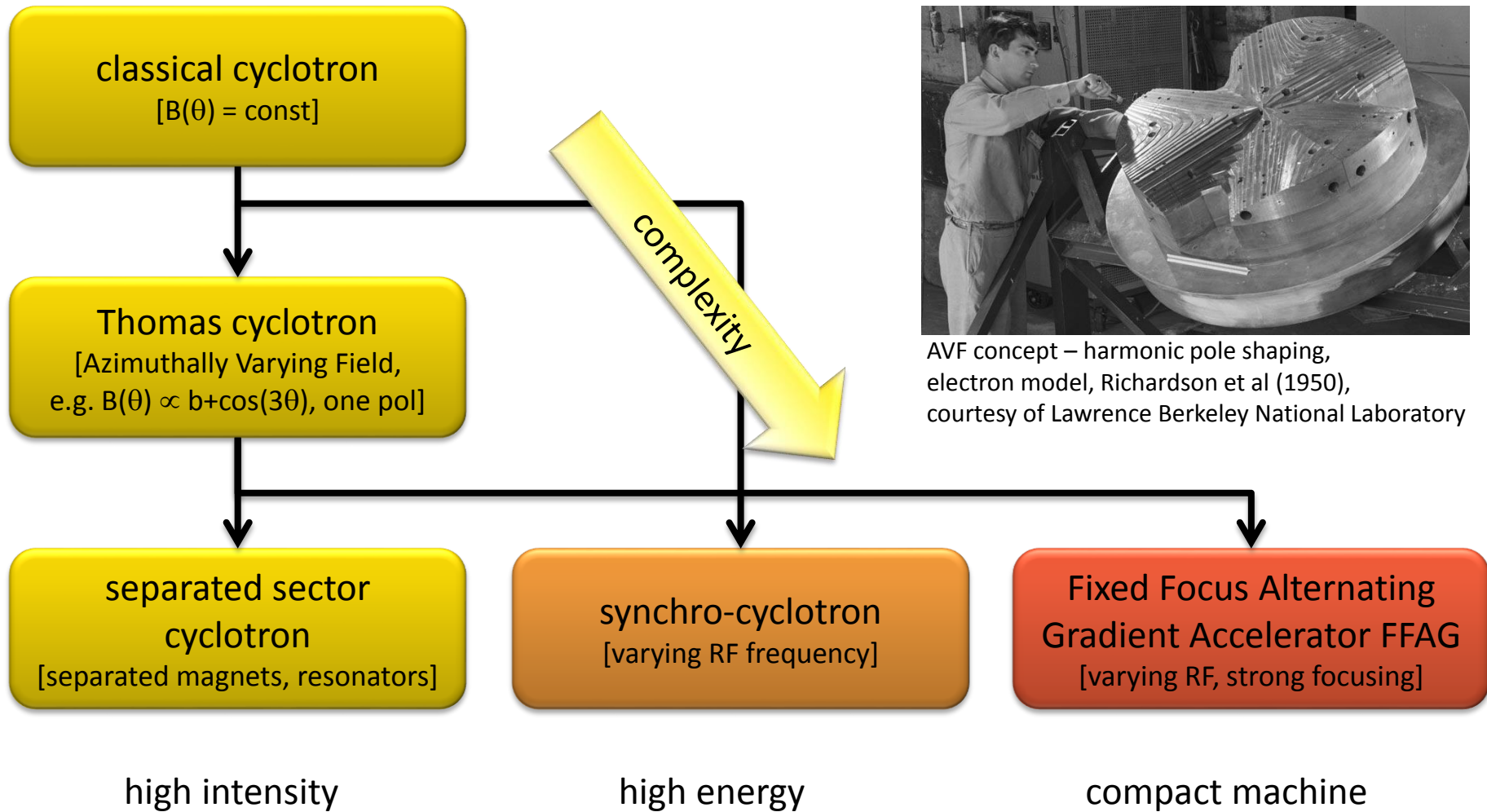


PSI Ring cyclotron

- **modular layout**, larger cyclotrons possible, sector magnets, box resonators, stronger focusing, injection/extraction in straight sections
- **external injection** required, i.e. pre-accelerator
- **box-resonators** (high voltage gain)
- high **extraction efficiency** possible:  
e.g. PSI: 99.98% =  $(1 - 2 \cdot 10^{-4})$



# classification of cyclotron like accelerators



AVF concept – harmonic pole shaping, electron model, Richardson et al (1950), courtesy of Lawrence Berkeley National Laboratory



## next: **injection for cyclotrons**

- internal source, axial injection, horizontal injection
- electrostatic inflector, electrostatic deflectors
- transverse matching, bunching
- space charge



# Injection – Overview

## Injection Techniques

- **internal source**
- axial injection
  - mirror inflector
  - **spiral inflector**
  - hyperbolic inflector
- radial injection
  - **electrostatic septum**
  - stripping injection

## Aspects to be considered

- overall central region design
- radial centering
- matching of beam optics
- vertical centering
- bunching / long. capture
- minimize overall losses for high intensity application



# Internal Ion Source

## Example: Cold Cathode, Penning Ionisation Gauge (PIG)

cylindrical „chimney“ with slit as extraction aperture for protons

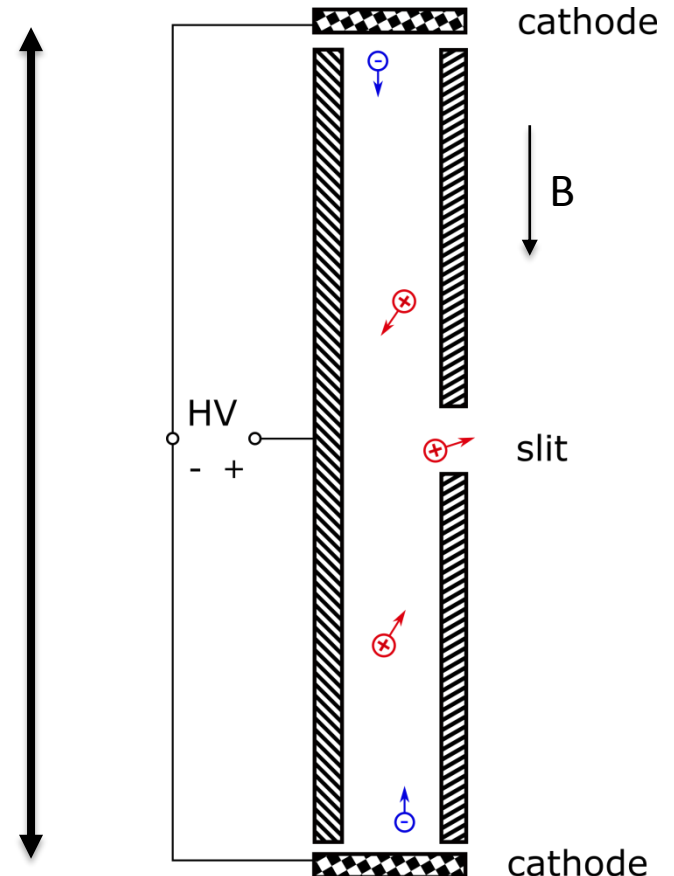
### advantage:

- simple concept
- no heating required

### critical:

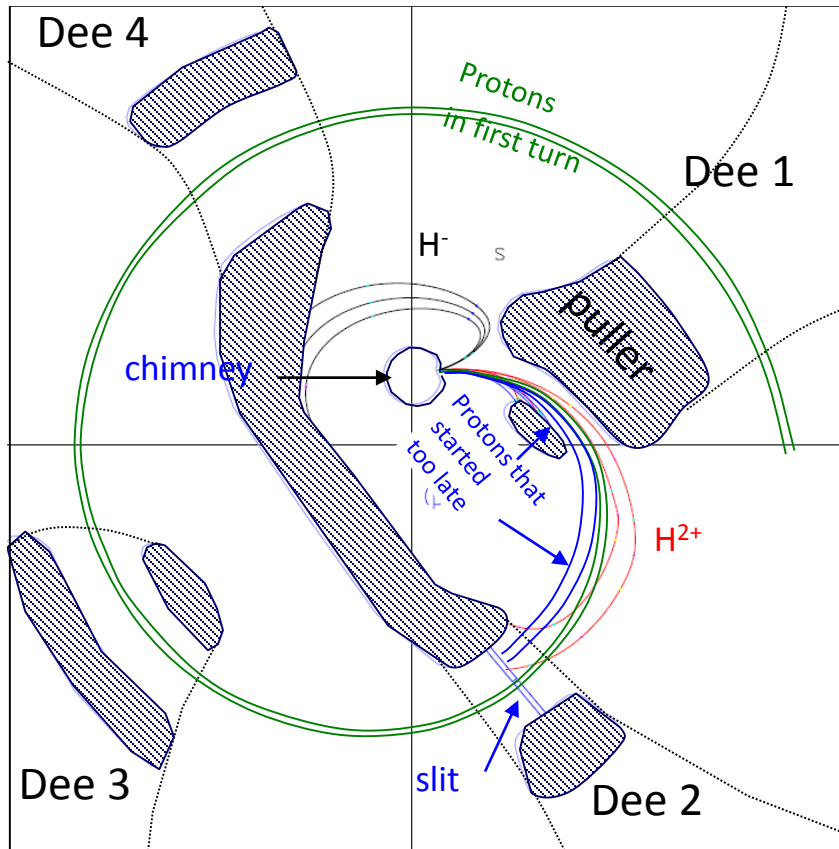
- reproducibility of captured current (geometry related sensitivity)
- current stability on short (ms) timescale

$O(10\text{cm})$



# internal ion source

→ example COMET (Accel/Varian)

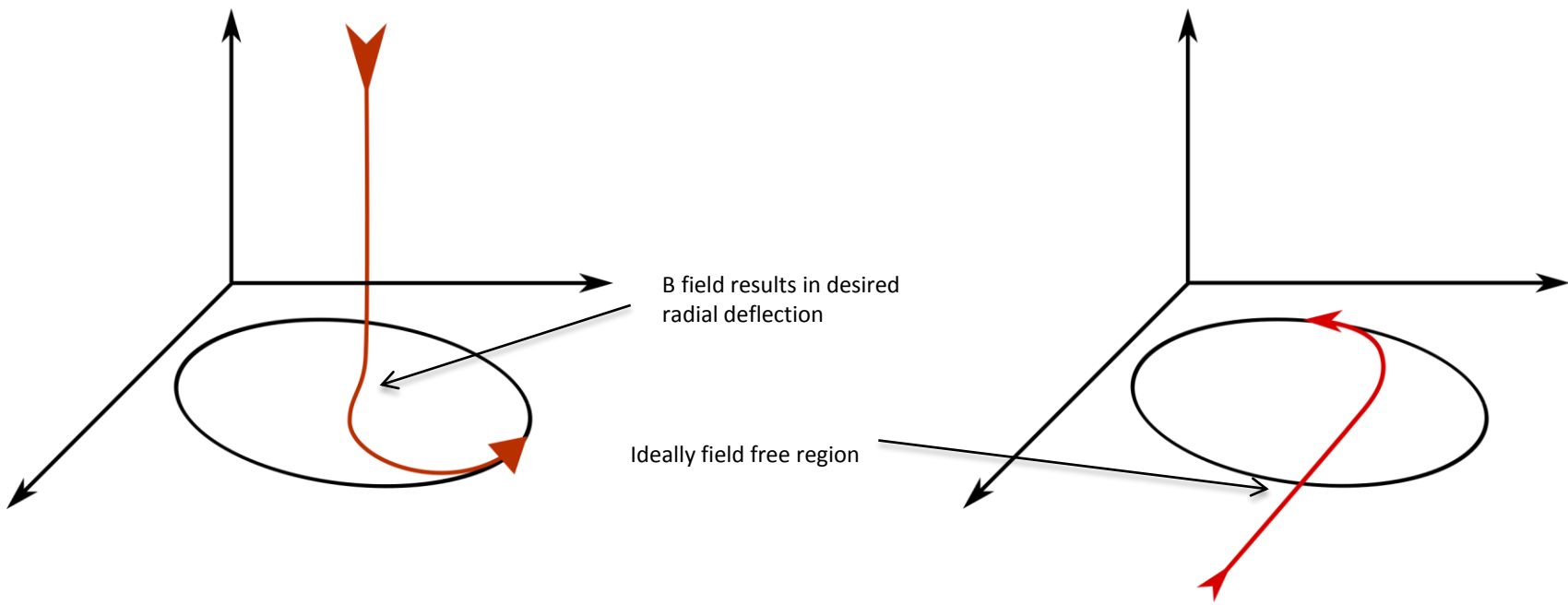


- Hydrogen is injected and ionized through chimney
- first acceleration by puller, connected to one Dee (80kV)

chimney  
= ion source  
deflector  
electrode  
for intensity  
regulation



# external source: axial vs. horizontal injection



axial: suited for compact cyclotron with field covering entire plane

horizontal: suited for sector cyclotron with gaps between magnets



# Beam Deflection by Electric Field

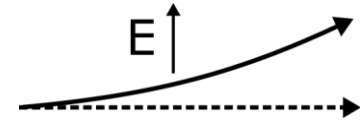
momentum change:  $\Delta p_{\perp} = \int F_{\perp} dt = \int \frac{F_{\perp}}{\beta c} ds, F = qE$

resulting angle:  $\theta = \frac{\Delta p_{\perp}}{p} = \frac{qlE}{\gamma\beta^2 E_0}$

bending radius:  $\rho = \frac{l}{\theta}$

electric rigidity:  $E\rho = \frac{\gamma^2 - 1}{\gamma} \frac{E_0}{q} = \frac{\gamma + 1}{\gamma} \frac{E_k}{q}$

low energy at source:  $E\rho \approx 2U_{\text{acc}}$



- $E_0$  rest energy
- $E$  field strength deflector
- $E_k$  kinetic energy
- $\rho$  bending radius
- $q$  charge
- $l$  length
- $U_{\text{acc}}$  acc.voltage (source)

## Bending radius in B and E:

comparison electric and magnetic force on protons

$$\vec{F}_E = e \cdot \vec{E}, \quad \vec{F}_B = e \cdot \vec{v} \times \vec{B}$$

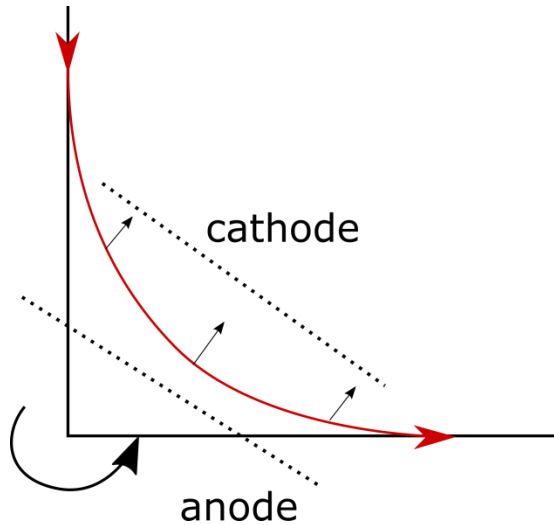
table: bending radius, varying  $E_k$

$E_k$	$B = 1\text{T}$	$E = 10\text{MV/m}$
60 keV	35 mm	12 mm
1 MeV	140 mm	200 mm
1 GeV	5.6 m	150 m

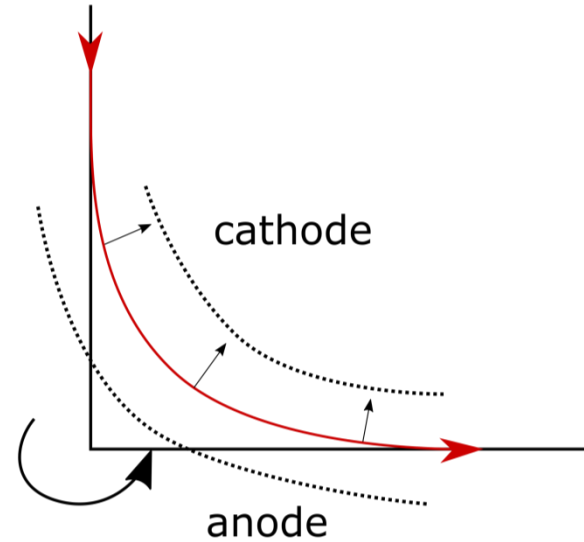




# electrostatic inflectors



mirror inflector: particle energy is variable, simple design



spiral inflector: force always perpendicular to velocity vector, no energy change

velocity vector rotates around vertical axis due to action of magnetic field;  
other solutions exist, e.g. hyperbolic inflector or even magnetostatic inflector



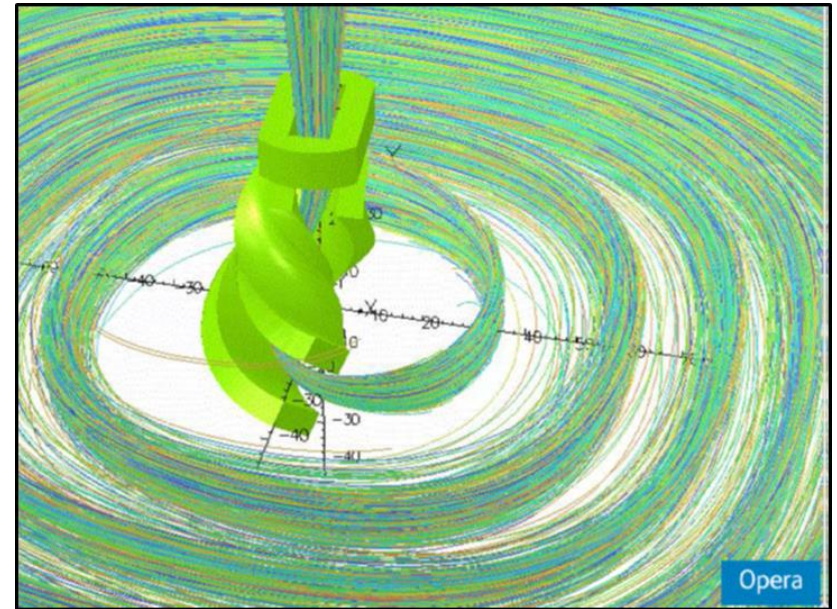
# injection schemes – spiral inflector

- an electrostatic component, basically a capacitor
- E-field arranged perpendicular to orbit, particles move on equipotential surfaces

simulation of orbits injected through a spiral inflector



[inflector IBA Cyclone 30 cyclotron]



[courtesy: W.Kleeven (IBA)]

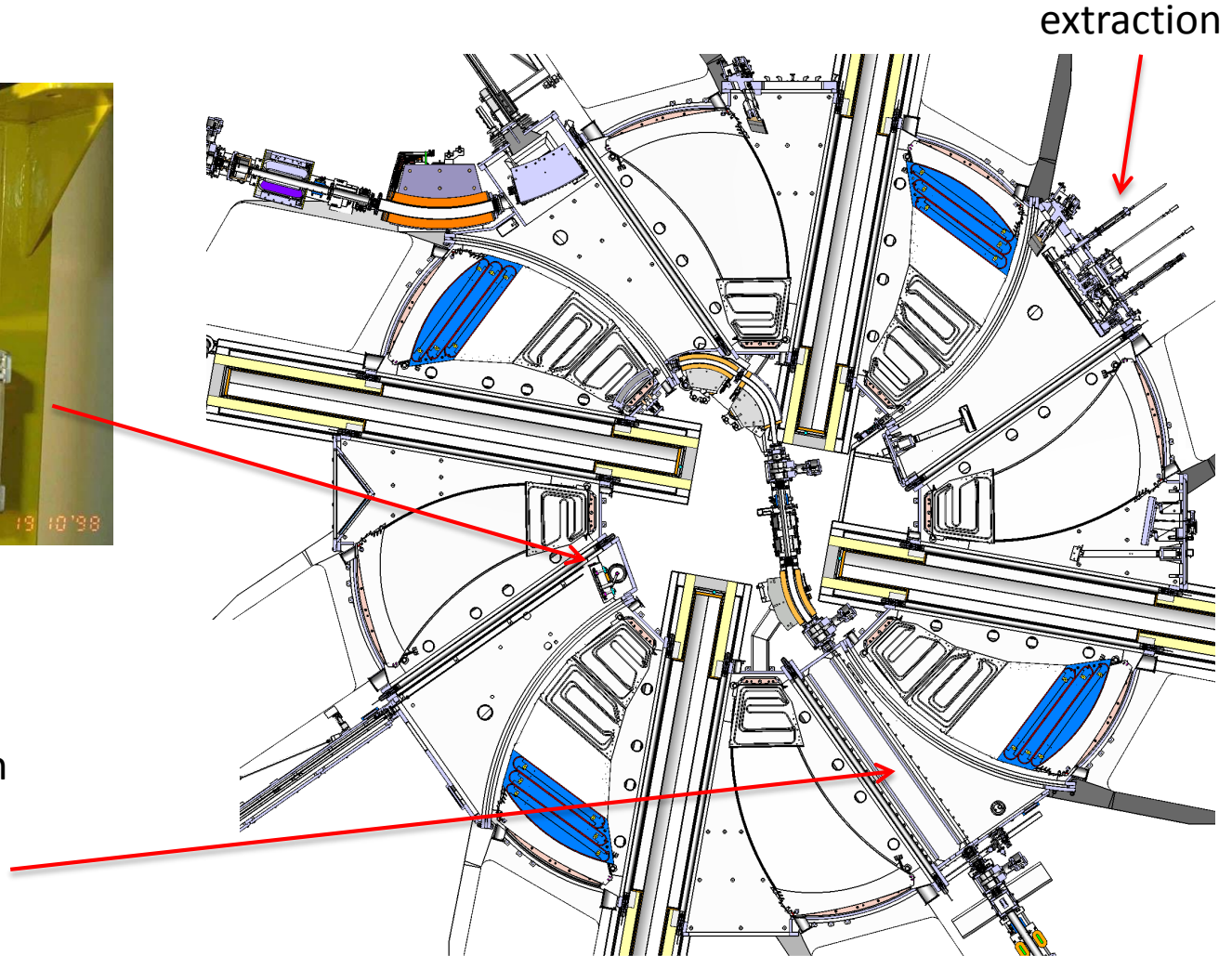


# Horizontal Injection – Example PSI Ring Cyclotron

Injection element



**Injection path (72MeV)** in region of low field, passing along 3rd-harmonic (150MHz) resonator



# Bunching for Cyclotrons

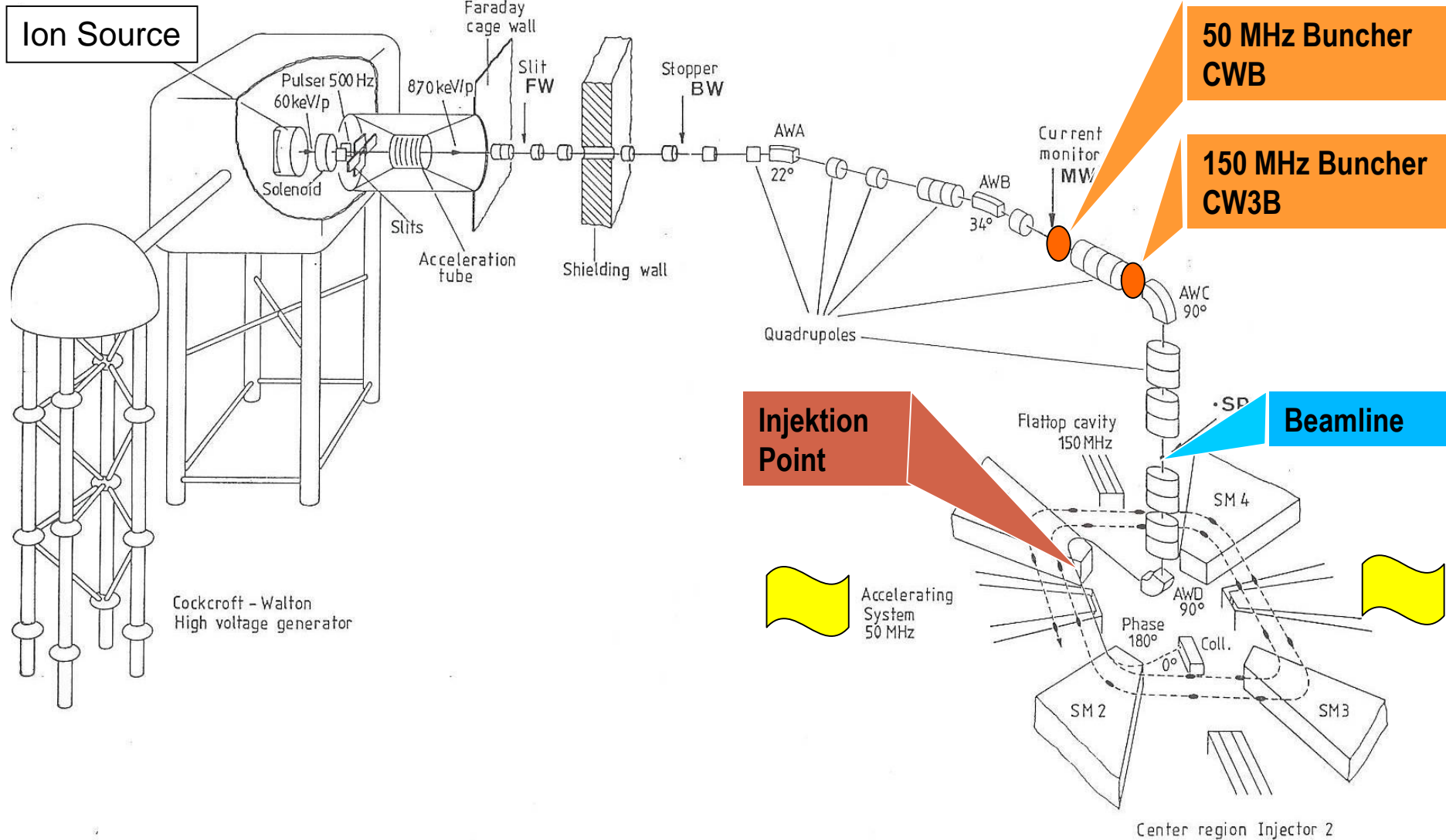
Ion sources deliver DC beam; for acceleration in an RF field the beam must be bunched; unbunched beam should be removed at low energy ( $\leq 5\text{MeV}$ ) to avoid uncontrolled losses and activation

schemes applied in practice:

	bunching in cyclotron	external buncher cavities	comment
internal source	x		lowest cost and complication
external source	x	x	higher intensity, variety of ions
DC pre-accelerator Cockcroft-Walton		x	low $\Delta E$ , costly
Radio Freq. Quadrupole (RFQ)		x	compact, costly



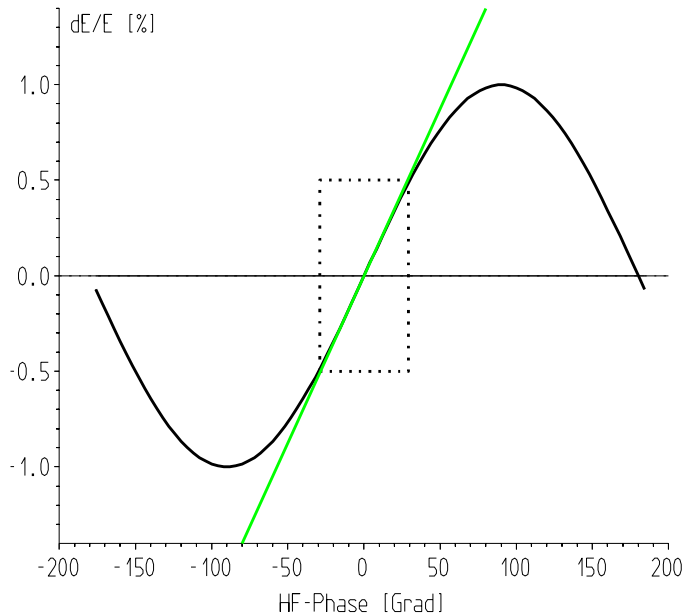
# Sketch of 870 keV Injektion Beam Line



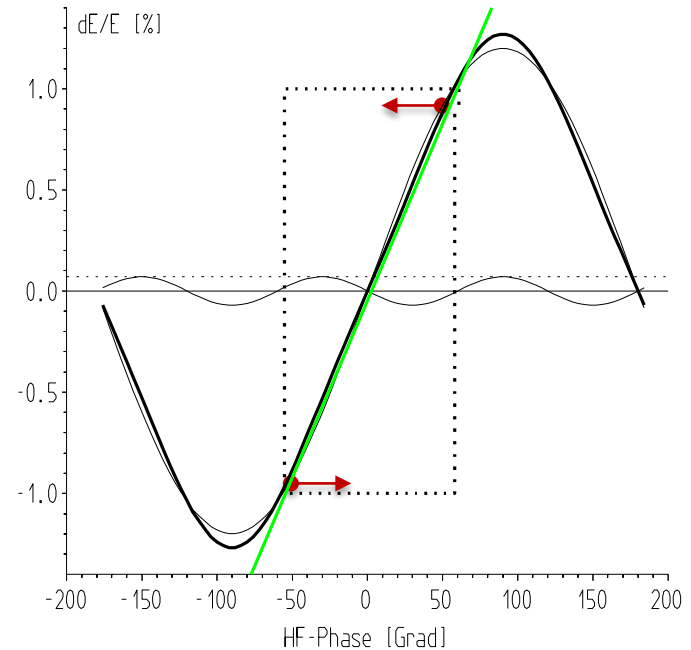
# 50 MHz and 50/150 MHz Harmonic Oscillation

→ by utilizing a harmonic buncher ( $3\omega$ ), a larger fraction of a DC beam can be captured in the cyclotron

only 50MHz buncher

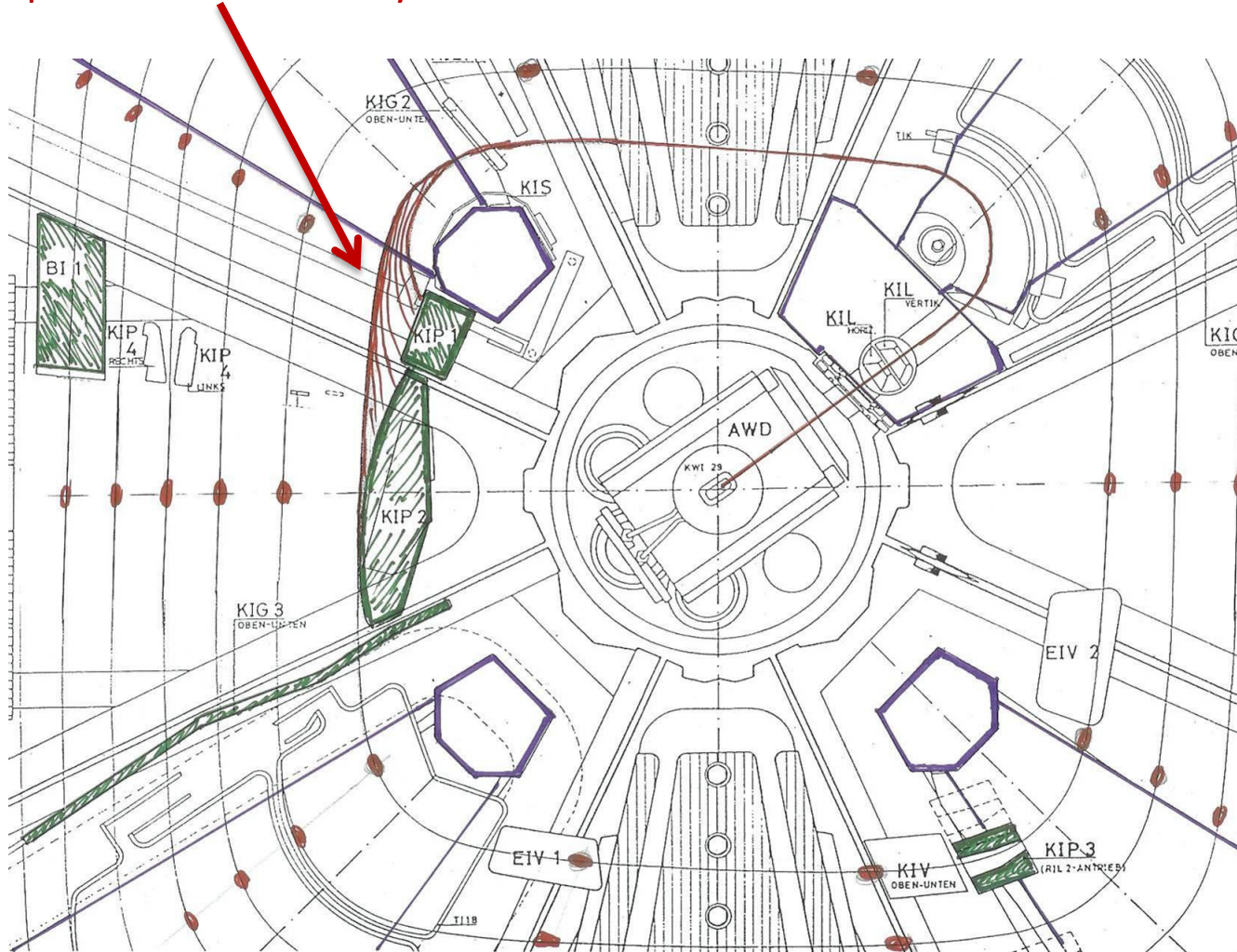


additional 150MHz buncher



# Center Region of PSI Injector 2

collimation of low energy  
protons and intensity control



**0.86 → 72MeV**  
**max 2.5mA, 180kW**



# PSI Injector 2 and Injection Beamline



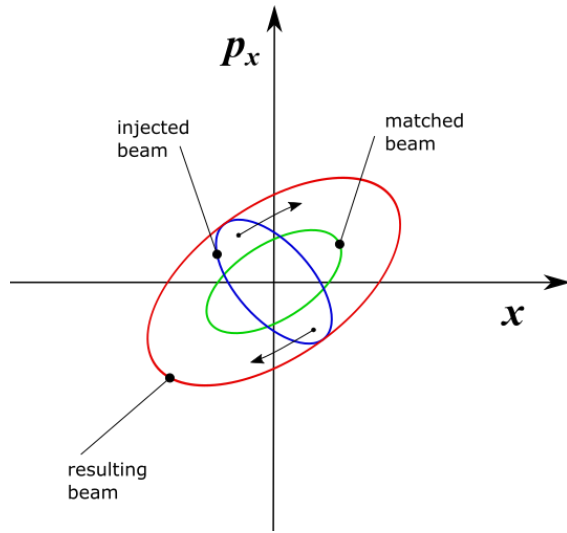
To Beamdump BX2





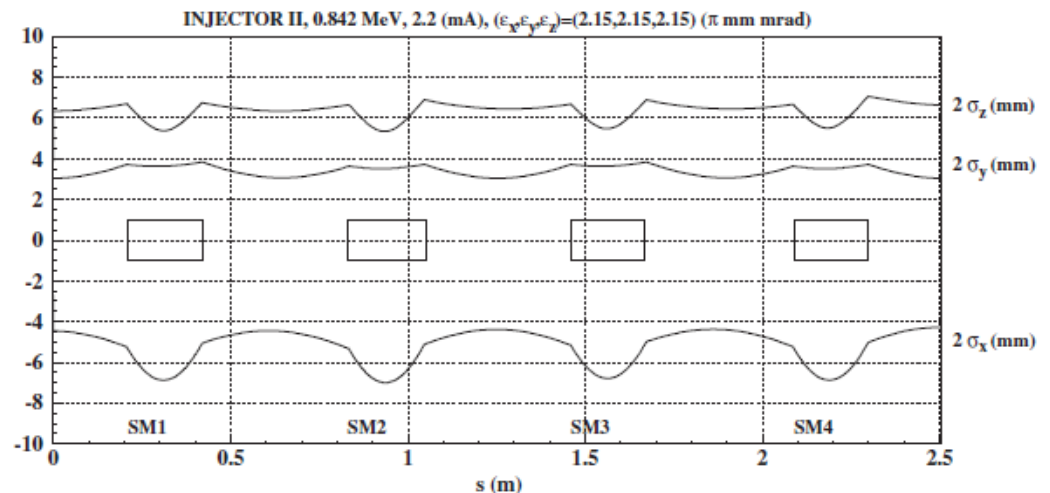
# Transverse Matching

- Similar to a synchrotron the envelope function  $\beta$  varies around the circumference; the beam at injection must be matched to avoid blow up and sub-optimal beam distributions



nonetheless of the short «storage» time of a beam in a cyclotron, the distribution starts to filament, if not properly matched

example: beam sizes around the circumference for Inj II cyclotron, PSI [Ch.Baumgarten, [7]]



# transverse space charge

especially at low energy space charge effects are critical for the injection of high intensity beams

vertical force from space charge:  $F_y = \frac{n_v e^2}{\epsilon_0 \gamma^2} \cdot y$ ,  $n_v = \frac{N}{(2\pi)^{\frac{3}{2}} \sigma_y D_f R \Delta R}$   
[constant charge density,  $D_f = I_{\text{avg}}/I_{\text{peak}}$ ]

thus, eqn. of motion:  $\ddot{y} + \left( \omega_c^2 \nu_{y0}^2 - \frac{n_v e^2}{\epsilon_0 m_0 \gamma^3} \right) y = 0$

→ tune shift results in **intensity limit** (see [6])!

tune shift from forces:  $\Delta \nu_y \approx -n_v \frac{2\pi r_p R^2}{\beta^2 \gamma^3 \nu_{y0}}$





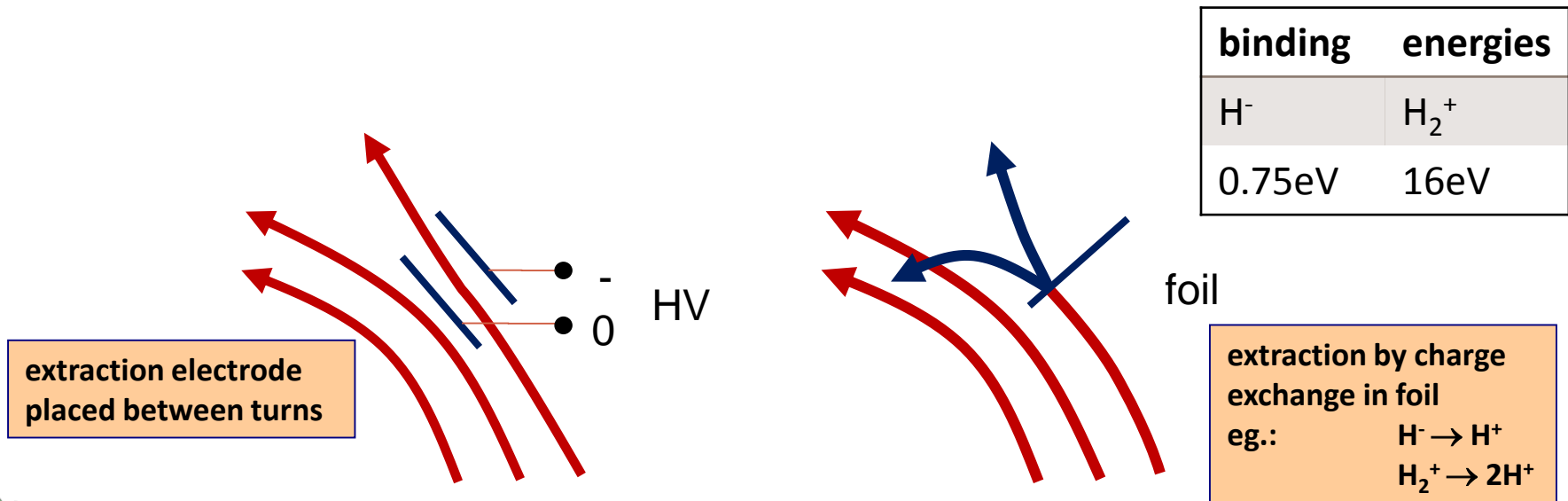
next: **extraction for cyclotrons**

- review of schemes: internal targets, electrostatic deflectors, stripping
- maximizing extraction efficiency: stepwidth, coherent oscillations, avoid tails



# electrostatic septum and charge exchange extraction

- simplest solution: use beam without extraction → **internal target**; use some mechanism to exchange target
- **electrostatic deflectors** with thin electrodes, deflecting element should affect just one turn, not neighboured turn → critical, cause of losses
- alternative: **charge exchange by stripping foil**; accelerate  $H^-$  or  $H_2^+$  to extract protons (problem: significant probability for unwanted loss of electron; Lorentz dissociation: B-field low, scattering: vacuum  $10^{-8}$ mbar)



# derivation of relativistic turn separation in a cyclotron

starting point: bending strength

→ compute total log.differential

→ use field index  $k = R/B \cdot dB/dR$

$$BR = \sqrt{\gamma^2 - 1} \frac{m_0 c}{e}$$

$$\frac{dB}{B} + \frac{dR}{R} = \frac{\gamma d\gamma}{\gamma^2 - 1}$$

$$\frac{dR}{d\gamma} = \frac{\gamma R}{\gamma^2 - 1} \frac{1}{1 + k}$$

slide 5

radius change per turn

$$\frac{dR}{dn_t} = \frac{dR}{d\gamma} \frac{d\gamma}{dn_t} \quad [U_t = \text{energy gain per turn}]$$

$$= \frac{U_t}{m_0 c^2} \frac{\gamma R}{(\gamma^2 - 1)(1 + k)}$$

} isochronicity not conserved (last turns)

$$= \frac{U_t}{m_0 c^2} \frac{R}{(\gamma^2 - 1)\gamma}$$

} isochronicity conserved (general scaling)



# discussion: scaling of turn separation

for clean extraction a large stepwidth (turn separation) is of utmost importance; in the PSI Ring most efforts were directed towards maximizing the turn separation

general scaling at extraction:

$$\Delta R(R_{\text{extr}}) = \frac{U_t}{m_0 c^2} \frac{R_{\text{extr}}}{(\gamma^2 - 1)\gamma}$$

desirable:

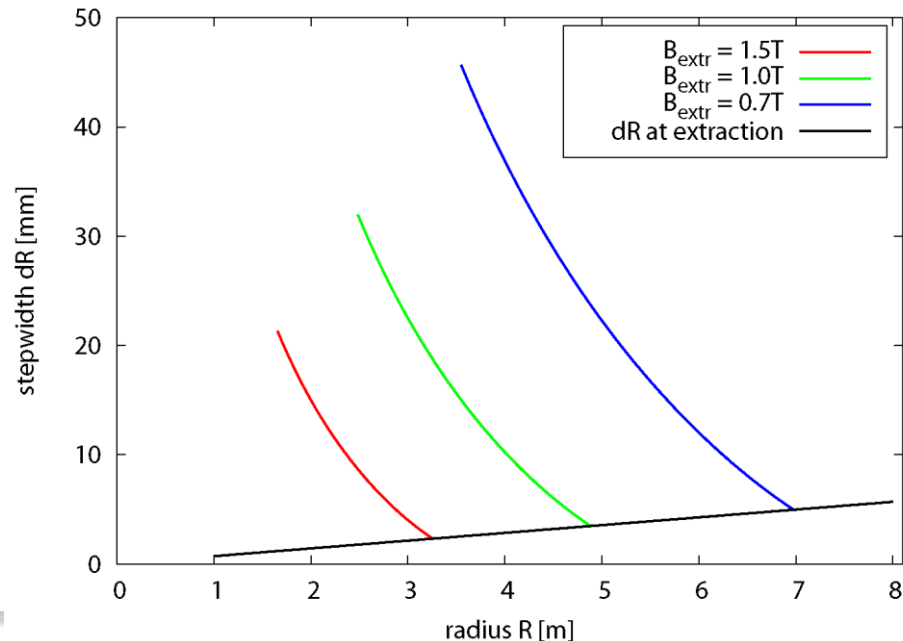
- limited energy (< 1GeV)
- large radius  $R_{\text{extr}}$
- high energy gain  $U_t$

scaling during acceleration:

$$\frac{dR}{dn_t} \approx \frac{U_t}{m_0 c^2} \frac{R}{\beta^2} \rightarrow \Delta R(R) \propto \frac{1}{R}$$

illustration:

**stepwidth vs. radius** in cyclotrons of different sizes; 100MeV inj  $\rightarrow$  800MeV extr



# methods to enhance turn separation

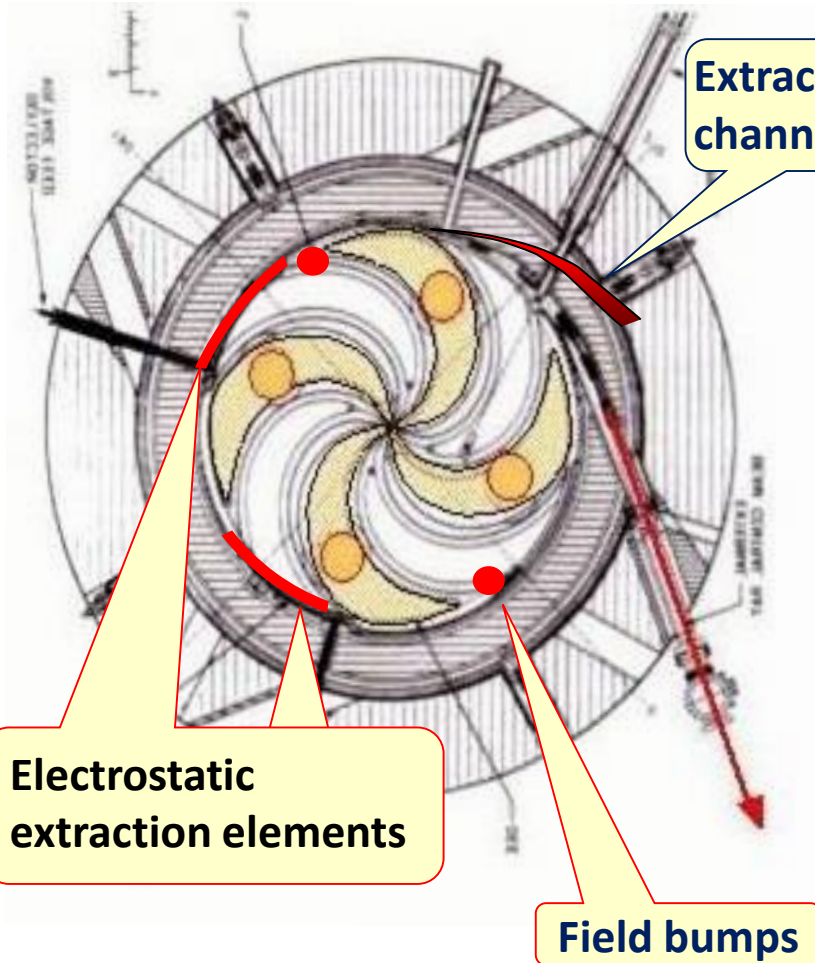
several techniques were invented to „artificially“ increase turn separation beyond the magnitude achieved by simple acceleration

<b>„brute force“</b>	resonant orbit distortion is excited by harmonic coils beyond a certain radius
<b>precessional extraction</b>	resonant excitation at $\nu_r=1$ plus steep $\nu_r$ slope in fringe field
<b>regenerative extraction</b>	using coherent excitation at half integer resonance by gradient bump

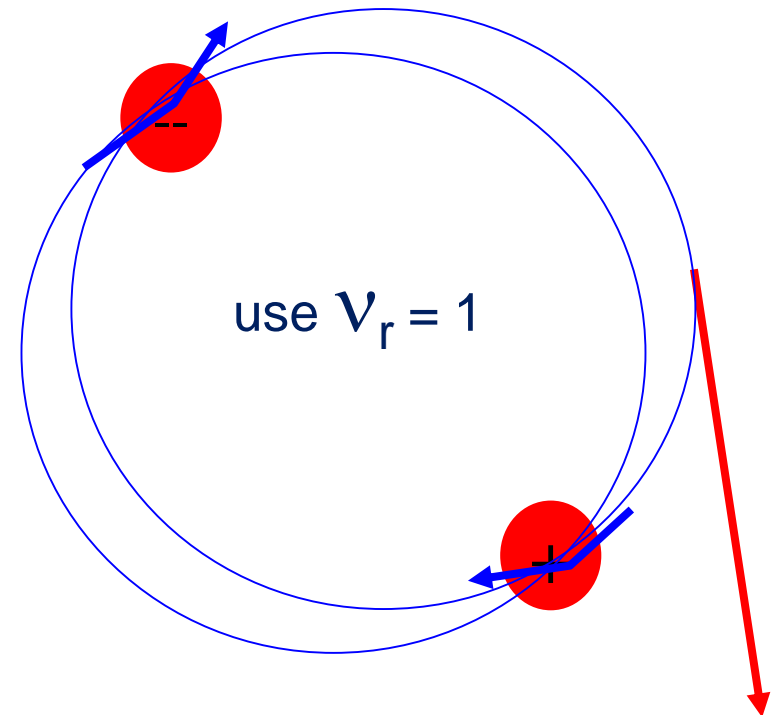
taken from Kleeven [1]



# Resonant Extraction (Varian/Accel cyclotron)



extraction efficiency:  
up to 80%



[M.Schippers, PSI]





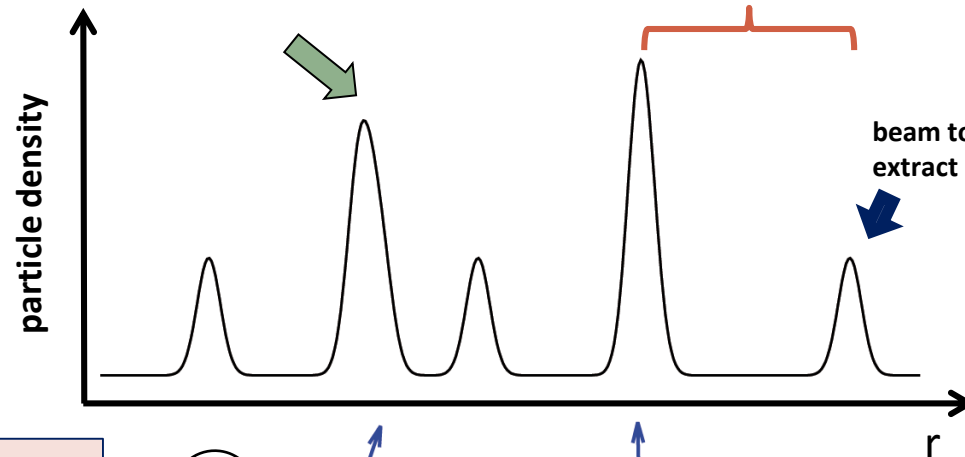
# extraction with coherent oscillations (PSI)

betatron oscillations around the “closed orbit” can be used to increase the radial stepwidth by a factor 3 !

**without orbit oscillations:** stepwidth from  $E_k$ -gain



**with orbit oscillations:** extraction gap; up to 3 x stepwidth possible for  $\nu_r = 1.5\pi$  (phase advance)



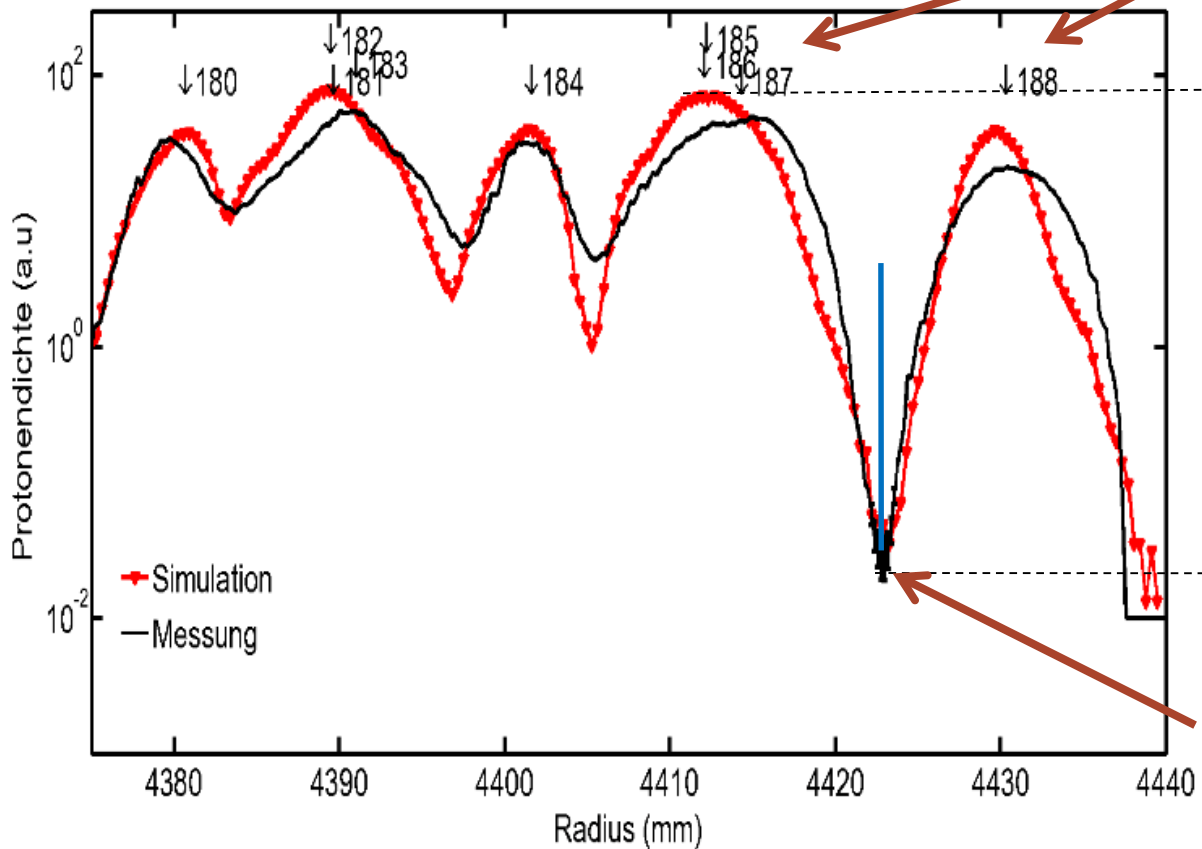
phase vector of orbit oscillations ( $r, r'$ )

$\nu_r$  decreases from 1.75 to 1.5



# extraction profile measured at PSI Ring Cyclotron

red: tracking simulation [OPAL]  
black: measurement



turn numbers  
from simulation

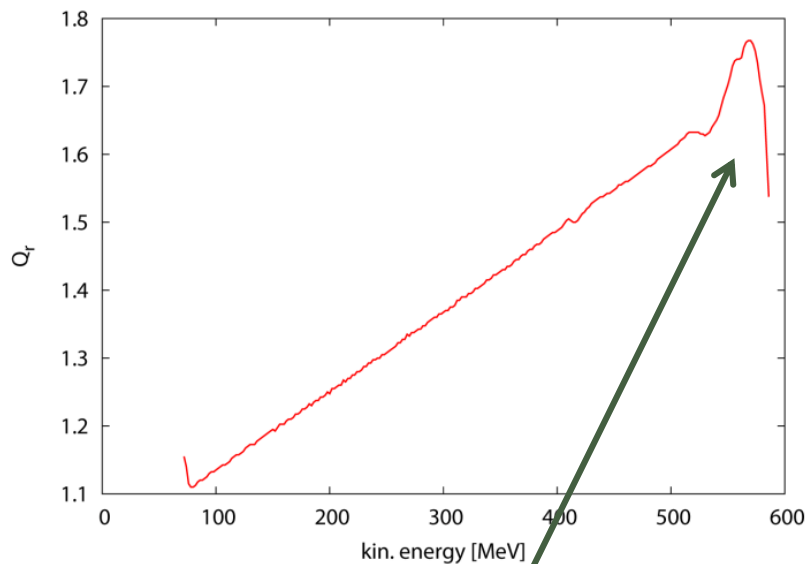
dynamic range:  
factor 2.000 in  
particle density

position of extraction septum  
 $d=50\mu\text{m}$

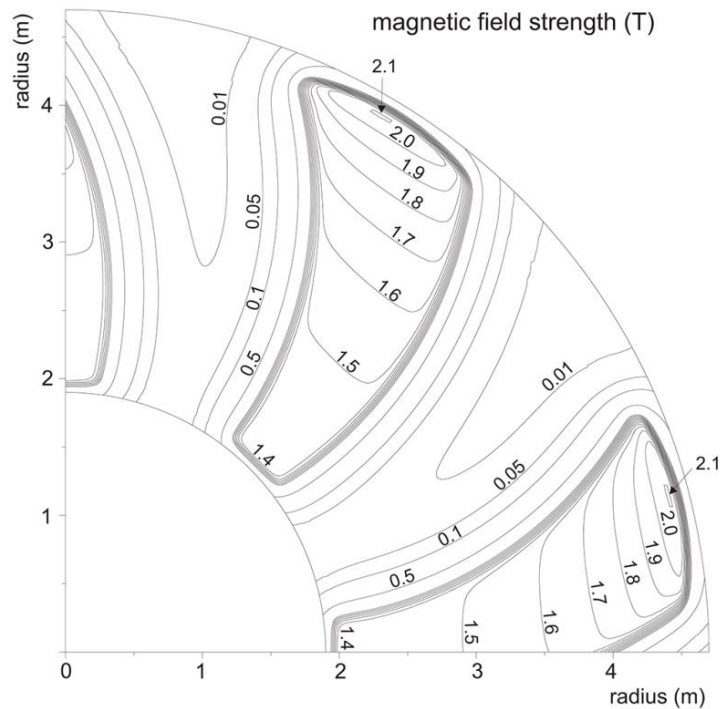
[Y.Bi et al]



# vertical tune in Ring cyclotron supports extraction



**radial tune vs. energy (PSI Ring)**  
typically  $\nu_r \approx \gamma$  during acceleration;  
but decrease in outer fringe field



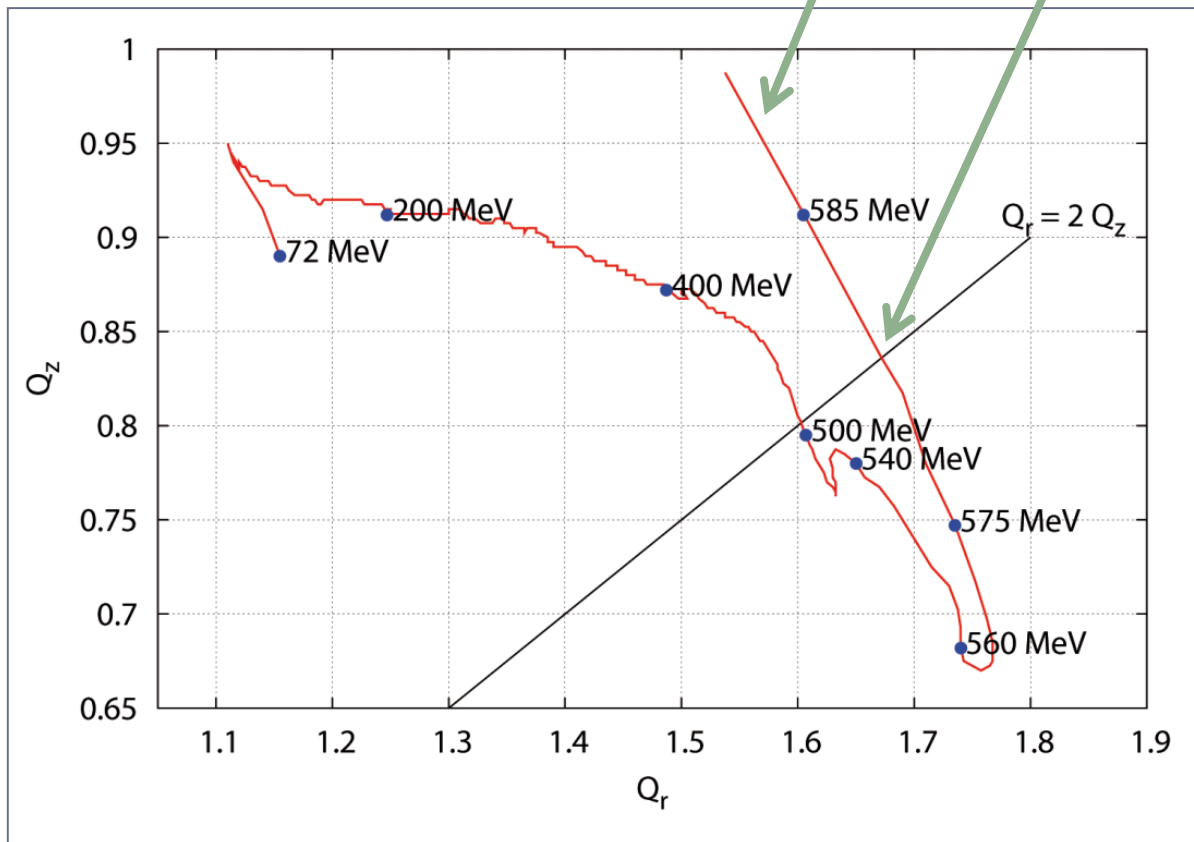
field map showing increase and steep decline of field with radius



# PSI Ring Cyclotron – tune diagram

coupling resonance – pass quickly!

$Q_r$  decreases towards extraction  
– enhance turn separation

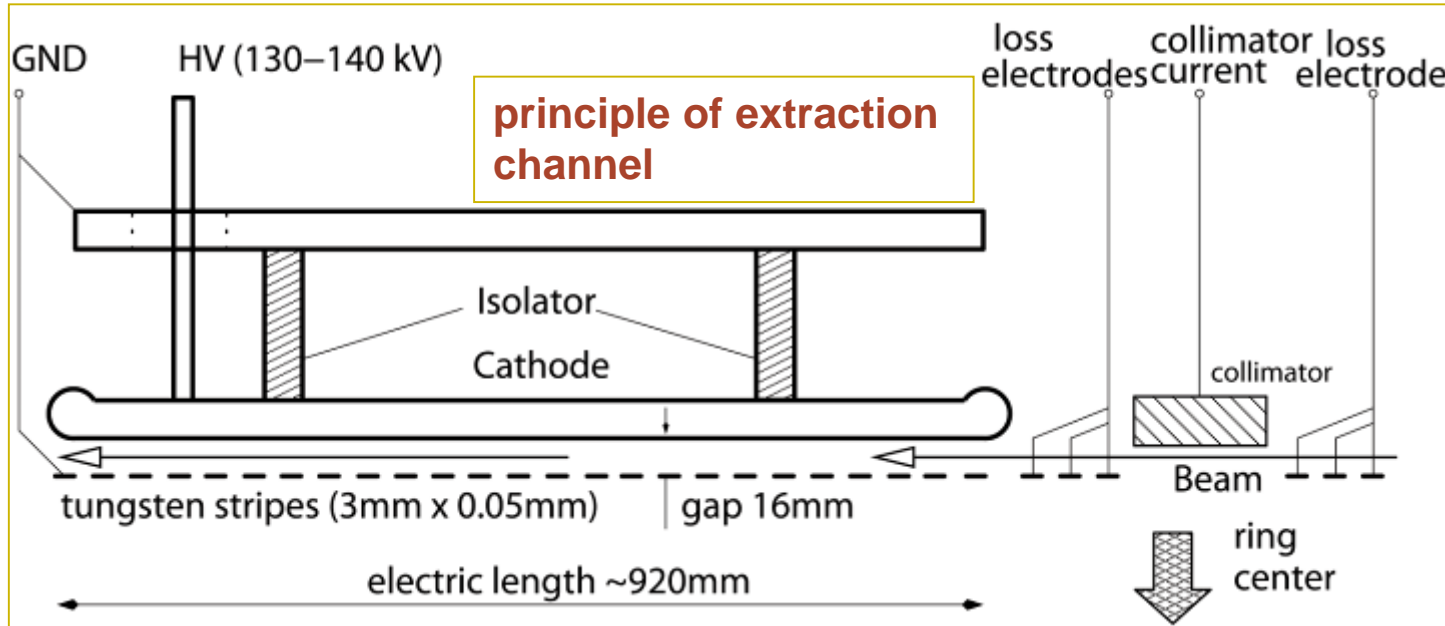


## comments:

- running on the coupling resonance would transfer the large radial betatron amplitude into vertical oscillations, which must be avoided
- special care has to be taken with fine-tuning the bending field in the extraction region



# injection/extraction with electrostatic elements



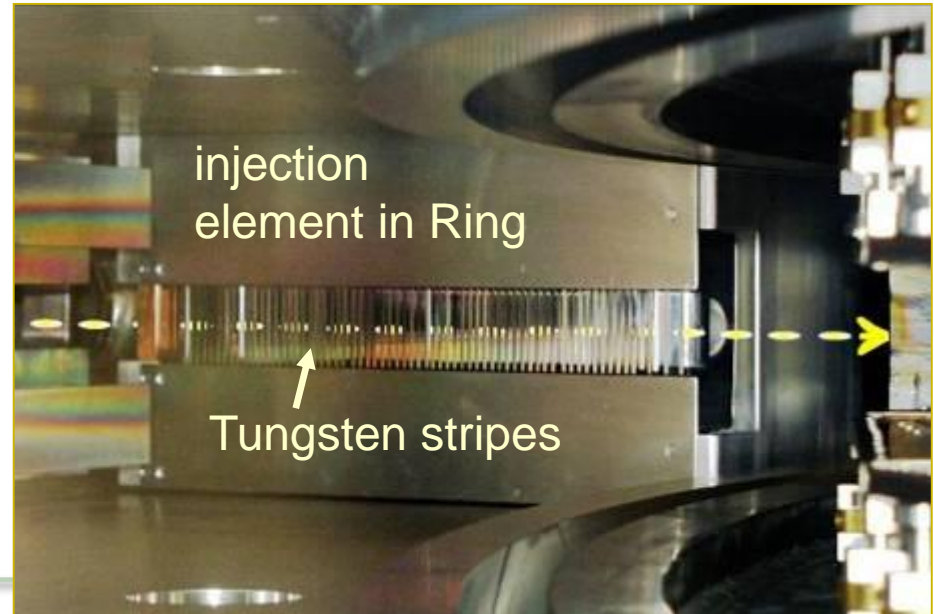
**parameters  
extraction chan.:**

$E_k = 590 \text{ MeV}$   
 $E = 8.8 \text{ MV/m}$   
 $\theta = 8.2 \text{ mrad}$   
 $\rho = 115 \text{ m}$   
 $U = 144 \text{ kV}$

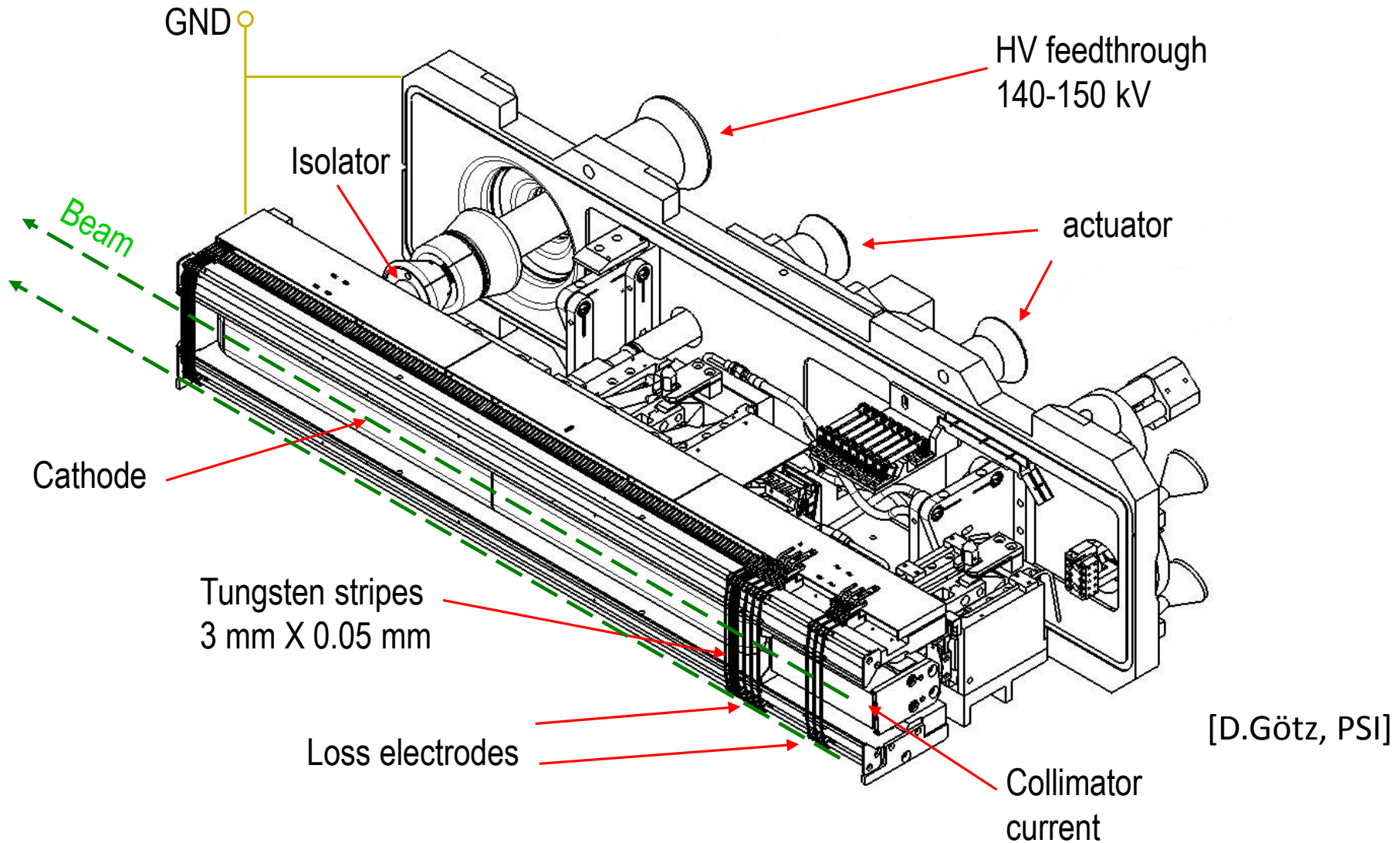
**major loss  
mechanism is  
scattering in 50 $\mu\text{m}$   
electrode!**

**electrostatic rigidity:**

$$E\rho = \frac{\gamma + 1}{\gamma} \frac{E_k}{q}$$



# Electrostatic Elements for High Energy/High Intensity



# longitudinal space charge (tails at extraction)

## sector model:

- accumulated energy spread transforms into transverse tails
- consider rotating uniform sectors of charge (overlapping turns)
- test particle “sees” only fraction of sector due to shielding of vacuum chamber with gap height  $2w$

two factors are proportional to the number of turns:

- 1) the charge density in the sector
- 2) the time span the force acts

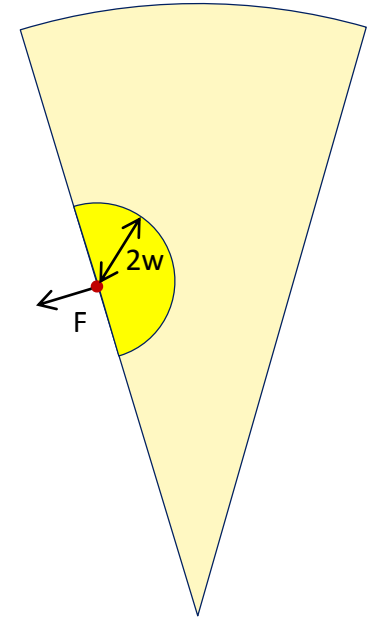
$$\Delta U_{sc} \approx 2.800\Omega \cdot eI_p \cdot \frac{n_{\max}^2}{\beta_{\max}}$$

derivation see [4]: Joho 1981

in addition:

- 3) the inverse of turn separation at extraction:  $\frac{1}{\Delta R_{\text{extr}}} \propto n_{\max}$

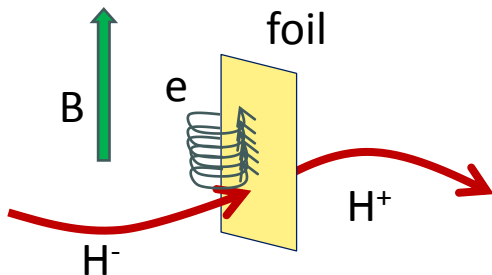
→ the attainable current at constant losses scales as  $n_{\max}^{-3}$



# extraction foil

- thin foil, for example carbon, removes the electron(s) with high probability
- new charge state of ion brings it on a new trajectory → separation from circulating beam
- lifetime of foil is critical due to heating, fatigue effects, radiation damage
- conversion efficiencies, e.g. generation of neutrals, must be considered carefully

electrons removed from the ions spiral in the magnetic field and may deposit energy in the foil



## How much power is carried by the electrons?

→ velocity and thus  $\gamma$  are equal for  $p$  and  $e$

$$E_k = (\gamma - 1)E_0$$

$$\rightarrow E_k^e = \frac{E_0^e}{E_0^p} E_k^p = 5.4 \cdot 10^{-4} E_k^p$$

## Bending radius of electrons?

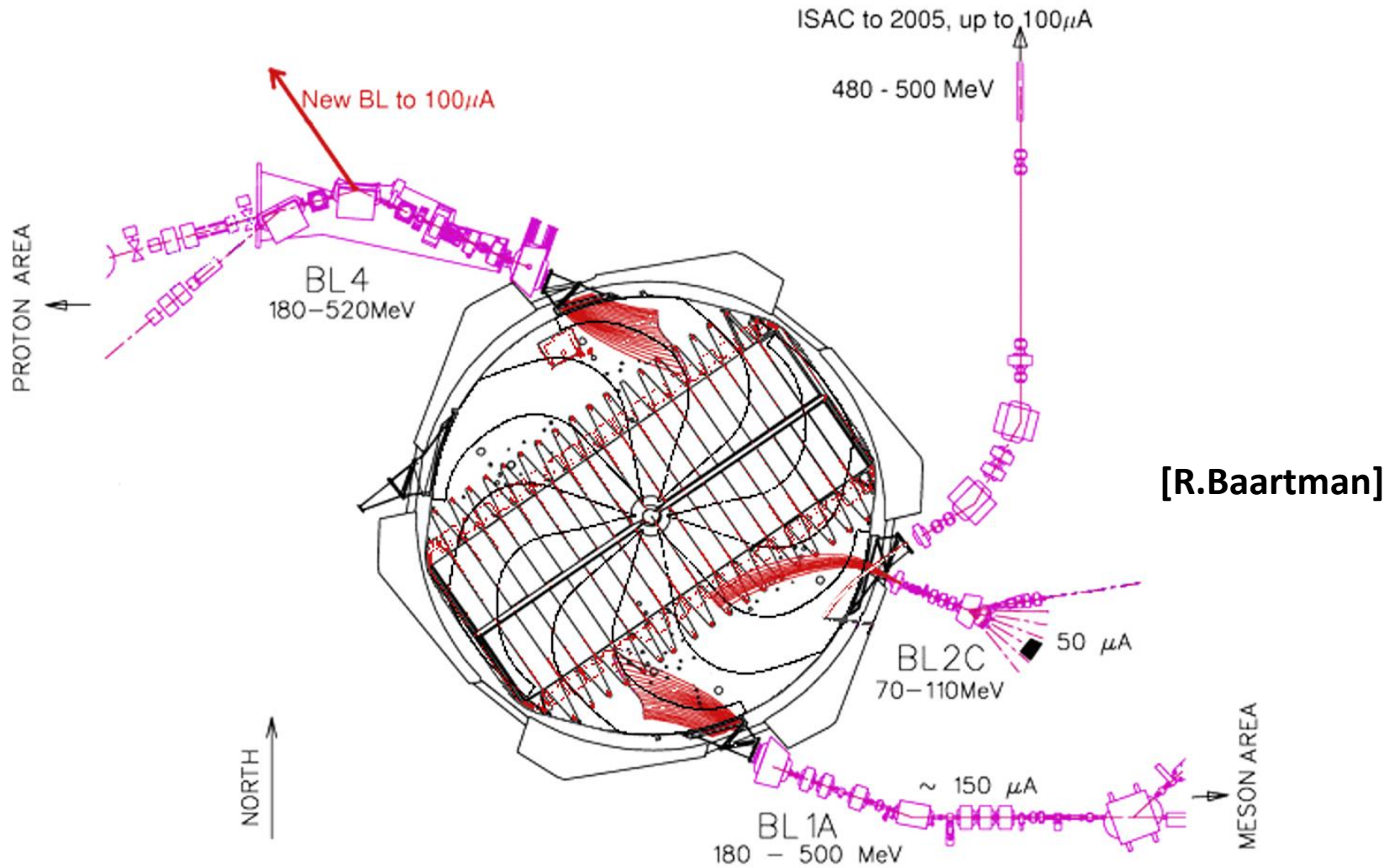
$$\rho^e = \frac{E_0^e}{E_0^p} \rho^p$$

→ typically mm

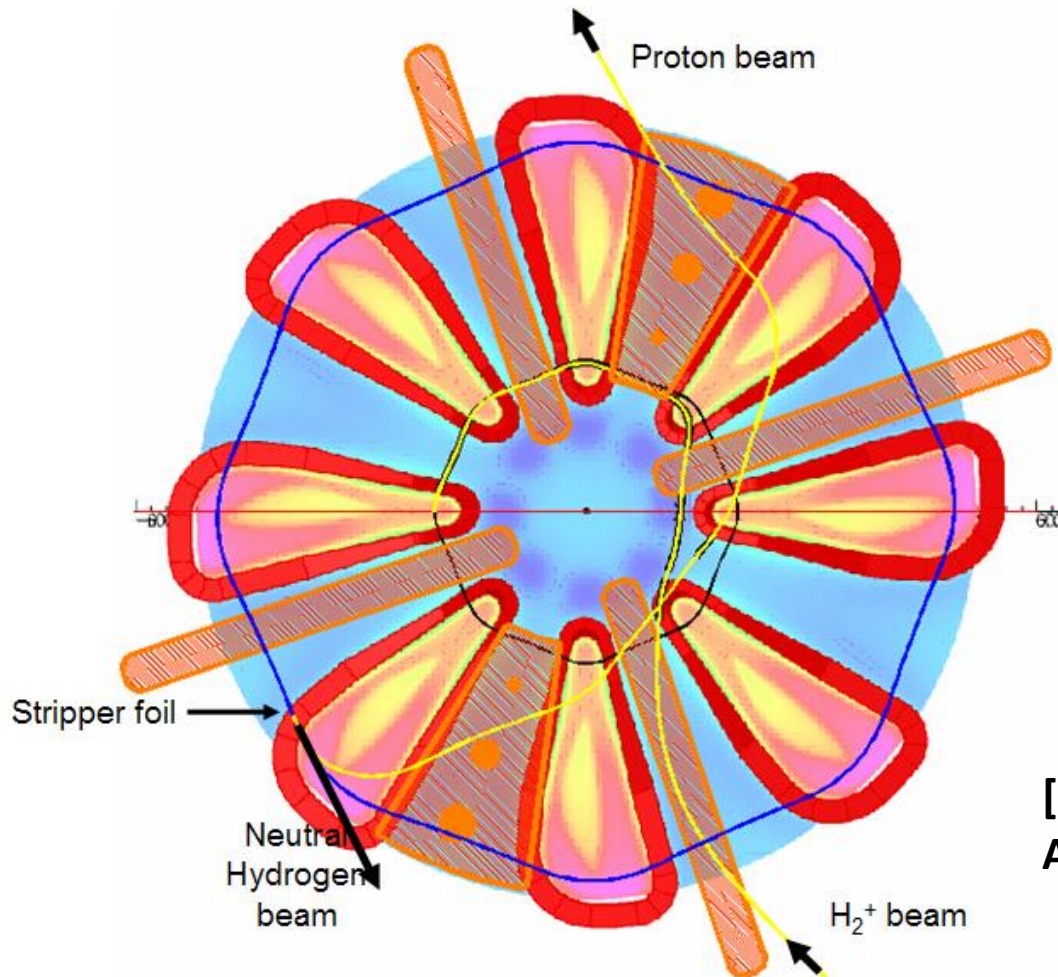




# example: multiple H<sup>-</sup> stripping extraction at TRIUMF



# example: $H_2^+$ stripping extraction in proposed Daedalus cyclotron [neutrino source]

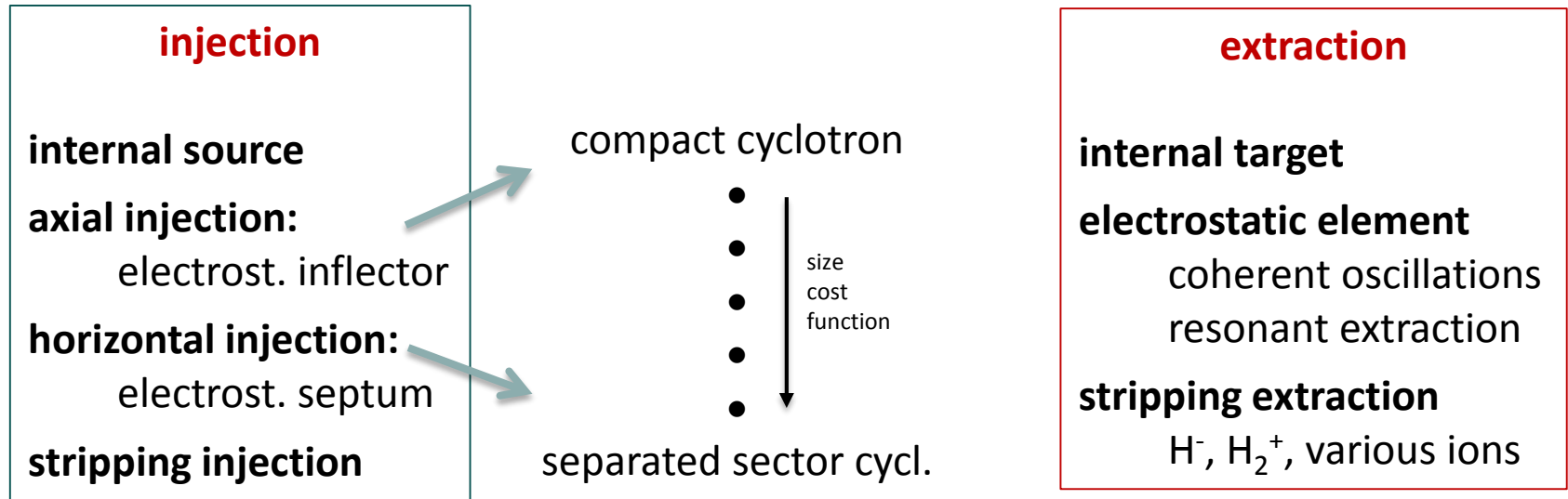


- purpose: pulsed high power beam for neutrino production, goals:
- 800MeV kin. energy
  - 5MW avg. beam power

[L.Calabretta,  
A.Calanna et al]



# Summary: Injection & Extraction for Cyclotrons



## beam physics aspects:

central region design, beam centering, transverse matching, bunching, beam blowup/tails & loss minimization & activation, space charge



# literature w.r.t. cyclotron injection/extraction

[1]	comprehensive review of inj./extr. concepts	W.Kleeven (IBA), Injection and Extraction for Cyclotrons <a href="https://cds.cern.ch/record/1005057/files/p271.pdf">https://cds.cern.ch/record/1005057/files/p271.pdf</a>
[2]	many examples and calculations for compact machines	P.Heikkinen (Jyväskylä), Injection and Extraction for Cyclotrons <a href="http://www.iaea.org/inis/collection/NCLCollectionStore/Public/26/001/26001643.pdf">http://www.iaea.org/inis/collection/NCLCollectionStore/Public/26/001/26001643.pdf</a>
[3]	calculations and matching on spiral inflectors	W.Kleeven & R.Baartman, 2x paper on spiral inflectors, Particle Accelerators 41 (1993), pages 41 and 55
[4]	extraction for very high intensity	W.Joho, High Intensity Problems in Cyclotrons, Proc. 5th intl. Conf. on Cyclotrons and their Applications, Caen, 337-347 (1981) <a href="http://accelconf.web.cern.ch/AccelConf/c81/papers/ei-03.pdf">http://accelconf.web.cern.ch/AccelConf/c81/papers/ei-03.pdf</a>
[5]	OPAL simulations; extraction profile	Y.Bi, A. Adelman, et al. Phys. Rev. STAB Vol. 14, 054402 (2011) <a href="http://journals.aps.org/prab/pdf/10.1103/PhysRevSTAB.14.054402">http://journals.aps.org/prab/pdf/10.1103/PhysRevSTAB.14.054402</a>
[6]	Intensity limitation	R.Baartman, Space Charge limit in separate Turn Cyclotrons, <i>Intl. Cycl. Conf. (2013)</i> <a href="http://accelconf.web.cern.ch/AccelConf/CYCLOTRONS2013/papers/we2pb01.pdf">http://accelconf.web.cern.ch/AccelConf/CYCLOTRONS2013/papers/we2pb01.pdf</a>
[7]	formation of round bunches and matching approach	Ch.Baumgarten, transverse-longitudinal coupling by space charge in cyclotrons <a href="http://journals.aps.org/prab/abstract/10.1103/PhysRevSTAB.14.114201">http://journals.aps.org/prab/abstract/10.1103/PhysRevSTAB.14.114201</a>



Thank you for your attention !

