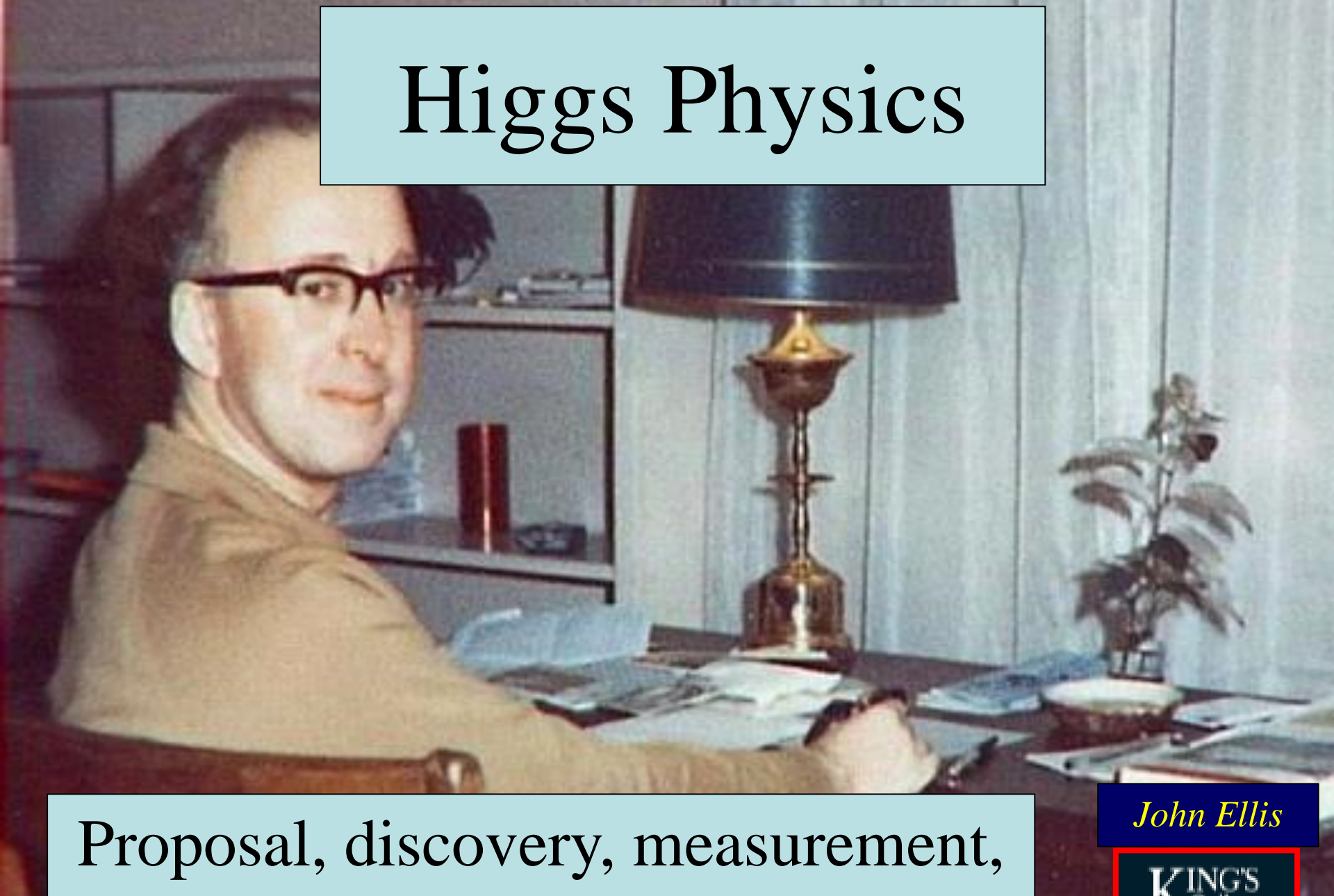


Higgs Physics



Proposal, discovery, measurement,
interpretation, what next?

John Ellis

KING'S
College
LONDON

The (G)AEBHGHKMP'tH Mechanism

BROKEN SYMMETRY AND THE MASS OF GAUGE VECTOR MESONS*

F. Englert and R. Brout

Faculté des Sciences, Université Libre de Bruxelles, Bruxelles, Belgium

(Received 26 June 1964)

BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS

P. W. HIGGS

Tait Institute of Mathematical Physics, University of Edinburgh, Scotland

Received 27 July 1964

VOLUME 13, NUMBER 16

PHYSICAL REVIEW LETTERS

BROKEN SYMMETRIES AND THE MASSES OF GAUGE VECTOR MESONS

Peter W. Higgs

Tait Institute of Mathematical Physics, University of Edinburgh,

(Received 31 August 1964)

The only one
who mentioned a
massive scalar boson

GLOBAL CONSERVATION LAWS AND MASSLESS PARTICLES*

G. S. Guralnik,[†] C. R. Hagen,[‡] and T. W. B. Kibble
Department of Physics, Imperial College, London, England
(Received 12 October 1964)

SPONTANEOUS BREAKDOWN OF STRONG INTERACTION SYMMETRY AND THE
ABSENCE OF MASSLESS PARTICLES

A. A. MIGDAL and A. M. YAKOVLEV

Submitted to JETP editor November 30, 1965; resubmitted February 16, 1966

J. Exp. Theor. Phys. (USSR) 51: 195-198 (1966)

The occurrence of massless particles in the presence of spontaneous symmetry breakdown is discussed. By summing all Feynman diagrams, one obtains for the difference of the mass

The Englert-Brout-Higgs Mechanism

- **Vacuum expectation value of scalar field**
- Englert & Brout: June 26th 1964
- First Higgs paper: July 27th 1964
- Pointed out loophole in argument of Gilbert if gauge theory described in Coulomb gauge
- Accepted by Physics Letters
- Second Higgs paper with explicit example sent on July 31st 1964 to Physics Letters, rejected!
- Revised version (Aug. 31st 1964) accepted by PRL

But the Higgs Boson

Englert &
Brout

Higgs

Guralnik, Hagen & Kibble



$$\partial^\mu \{ \partial_\mu (\Delta\varphi_1) - e\varphi_0 A_\mu \} = 0, \quad (2a)$$

$$\{ \partial^2 - 4\varphi_0^2 V''(\varphi_0^2) \} (\Delta\varphi_2) = 0, \quad (2b)$$

$$\partial_\nu F^{\mu\nu} = e\varphi_0 \{ \partial^\mu (\Delta\varphi_1) - e\varphi_0 A_\mu \}. \quad (2c)$$

It is worth noting that an essential feature of the type of theory which has been described in this note is the prediction of incomplete multiplets of scalar and vector bosons.⁸ It is to be

Also
Goldstone in
global case

We consider, as our example, a theory which was partially solved by Englert and Brout,⁵ and bears some resemblance to the classical theory of Higgs.⁶ Our starting point is the ordinary electrodynamics of massless spin-zero particles, characterized by the Lagrangian

$$\mathcal{L} = -\frac{1}{2} F^{\mu\nu} (\partial_\mu A_\nu - \partial_\nu A_\mu) + \frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \varphi^\mu \partial_\mu \varphi + \frac{1}{2} \varphi^\mu \varphi_\mu + ie_0 \varphi^\mu q \varphi A_\mu,$$

With no loss of generality, we can take $\eta_2 = 0$, and find

$$(-\partial^2 + \eta_1^2) A_k^T = 0, \\ -\partial^2 \varphi_2 = 0, \\ (-\partial^2 + \eta_1^2) A_k^T = 0,$$

where the superscript T denotes the transverse part. The two degrees of freedom of A_k^T combine with φ_1 to form the three components of a

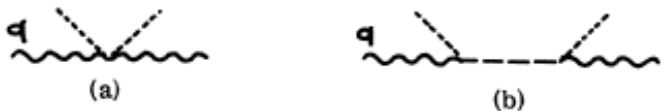


FIG. 1. Broken-symmetry diagram leading to a mass for the gauge field. Short-dashed line, $\langle\varphi_1\rangle$; long-dashed line, φ_2 propagator; wavy line, A_μ propagator. (a) $\rightarrow (2\pi)^4 i e^2 g_{\mu\nu} \langle\varphi_1\rangle^2$, (b) $\rightarrow -(2\pi)^4 i e^2 (q_\mu q_\nu / q^2) \times \langle\varphi_1\rangle^2$.

Summary of the Standard Model

- Particles and $SU(3) \times SU(2) \times U(1)$ quantum numbers:

L_L E_R	$\begin{pmatrix} \nu_e \\ e^- \end{pmatrix}_L, \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}_L, \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix}_L$ e_R^-, μ_R^-, τ_R^-	$(1,2,-1)$ $(1,1,-2)$
Q_L U_R D_R	$\begin{pmatrix} u \\ d \end{pmatrix}_L, \begin{pmatrix} c \\ s \end{pmatrix}_L, \begin{pmatrix} t \\ b \end{pmatrix}_L$ u_R, c_R, t_R d_R, s_R, b_R	$(3,2,+1/3)$ $(3,1,+4/3)$ $(3,1,-2/3)$

- Lagrangian: $\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F^{a\ \mu\nu}$ gauge interactions
 $+ i\bar{\psi} \not{D}\psi + h.c.$ matter fermions
 $+ \psi_i y_{ij} \psi_j \phi + h.c.$ Yukawa interactions
 $+ |D_\mu \phi|^2 - V(\phi)$ Higgs potential

Untested
before 2012

The Standard Model Lagrangian

$$\mathcal{L}_{SM} = \mathcal{L}_m + \mathcal{L}_g + \mathcal{L}_h + \mathcal{L}_y \quad ,$$

$$\mathcal{L}_m = \bar{Q}_L i \gamma^\mu D_\mu^L Q_L + \bar{q}_R i \gamma^\mu D_\mu^R q_R + \bar{L}_L i \gamma^\mu D_\mu^L L_L + \bar{l}_R i \gamma^\mu D_\mu^R l_R$$

$$\mathcal{L}_G = -\frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{4} W_{\mu\nu}^a W^{a\mu\nu} \quad \checkmark \quad \text{Experiment: accuracy} < \%$$

$$\mathcal{L}_H = (D_\mu^L \phi)^\dagger (D^{L\mu} \phi) - V(\phi)$$

$$\mathcal{L}_Y = y_d \bar{Q}_L \phi q_R^d + y_u \bar{Q}_L \phi^c q_R^u + y_L \bar{L}_L \phi l_R +$$

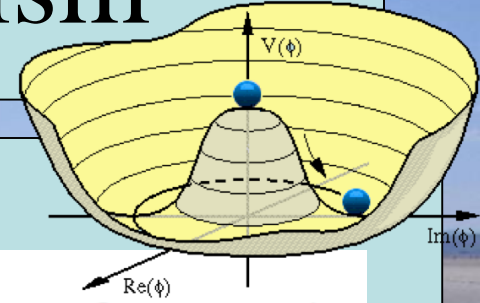
No direct evidence
until July 4, 2012

$$D_\mu^L = \partial_\mu - ig W_\mu^a T^a - iY g' B_\mu \quad , \quad D_\mu^R = \partial_\mu - iY g' B_\mu$$

$$V(\phi) = -\mu^2 \phi^2 + \lambda \phi^4 \quad .$$

1964

Abelian Higgs Mechanism



- Lagrangian

$$\mathcal{L} = (D_\mu \phi)^\dagger (D^\mu \phi) - V(|\phi|) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}, \quad D_\mu = \partial_\mu - ieA_\mu$$

- Gauge transformation $\phi'(x) = e^{i\alpha(x)} \phi(x) = e^{i\alpha(x)} e^{i\theta(x)} \eta(x)$

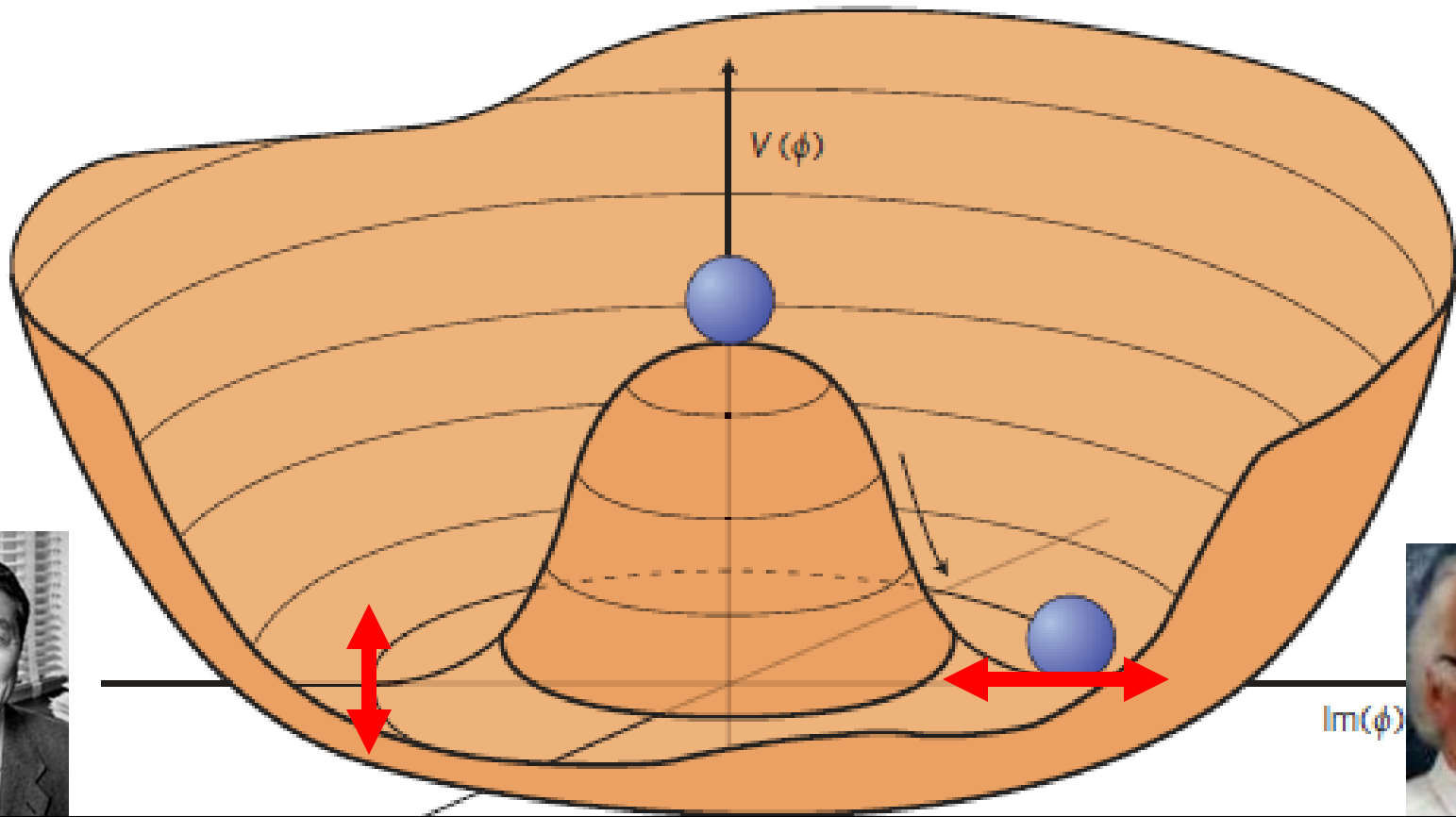
$$A'_\mu(x) = A_\mu(x) + \frac{1}{e} \partial_\mu \alpha(x)$$

- Choose $\alpha(x) = -\theta(x)$: $\phi'(x) = \eta(x)$

- Rewrite Lagrangian: $\mathcal{L} = |(\partial - ieA'_\mu)\eta|^2 - V(\eta) - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu}$

$$\begin{aligned} \mathcal{L} &= |(\partial_\mu - ieA'_\mu)(v + \frac{1}{\sqrt{2}}H)|^2 - \frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} - V \\ &= \underbrace{-\frac{1}{4} F'_{\mu\nu} F'^{\mu\nu} + v^2 e^2 A'_\mu A'^\mu}_{\text{massive } A\text{-field, } m_A \sim ev} + \underbrace{\frac{1}{2} [(\partial_\mu H)^2 - m_H^2 H^2]}_{\text{neutral scalar, } m_H \neq 0} + \dots \end{aligned}$$

Nambu **EB, H, GHK** and Higgs



Spontaneous symmetry breaking: massless Nambu-Goldstone boson **'eaten'** by massless gauge boson

Accompanied by massive particle

1967

Masses for SM Gauge Bosons

- Kinetic terms for SU(2) and U(1) gauge bosons:

$$\mathcal{L} = -\frac{1}{4} G_{\mu\nu}^i G^{i\mu\nu} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

where $G_{\mu\nu}^i \equiv \partial_\mu W_\nu^i - \partial_\nu W_\mu^i + ig\epsilon_{ijk} W_\mu^j W_\nu^k$ $F_{\mu\nu} \equiv \partial_\mu W_\nu - \partial_\nu W_\mu$

- Kinetic term for Higgs field:

$$\mathcal{L}_\phi = -|D_\mu \phi|^2 \quad D_\mu \equiv \partial_\mu - i g \sigma_i W_\mu^i - i g' Y B_\mu$$

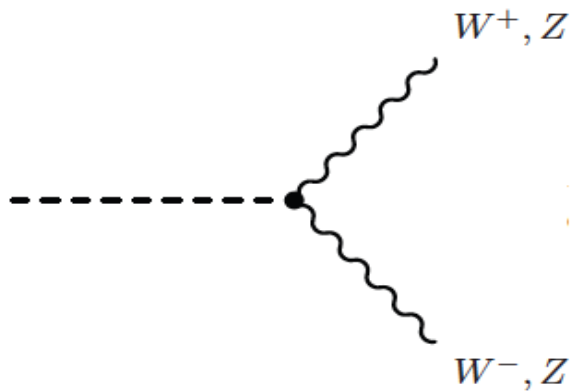
- Expanding around vacuum: $\phi = \langle 0|\phi|0 \rangle + \hat{\phi}$

$$\mathcal{L}_\phi \ni -\frac{g^2 v^2}{2} W_\mu^+ W^{\mu-} - \frac{g'^2 v^2}{2} B_\mu B^\mu + g g' v^2 B_\mu W^{\mu 3} - g^2 \frac{v^2}{2} W_\mu^3 W^{\mu 3}$$

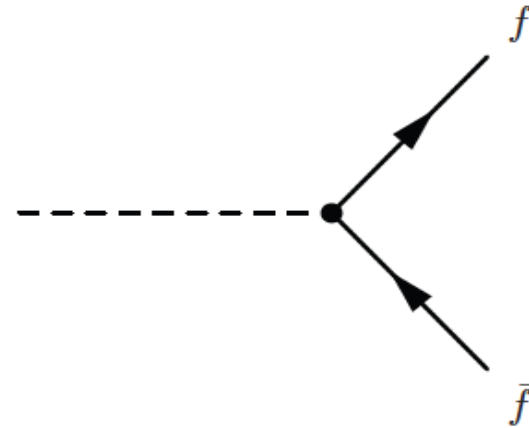
- Boson masses:

$$m_{W^\pm} = \frac{gv}{2} \quad Z_\mu = \frac{gW_\mu^3 - g'B_\mu}{\sqrt{g^2 + g'^2}} : m_Z = \frac{1}{2} \sqrt{g^2 + g'^2} v ; \quad A_\mu = \frac{g'W_\mu^3 + gB_\mu}{\sqrt{g^2 + g'^2}} : m_A = 0$$

Higgs Boson Couplings



$$g_2 M_W, \quad g_2 \frac{M_Z}{c_W}$$



$$\frac{m_f}{v} = \frac{g_2 m_f}{2M_W}$$

$$\Gamma(H \rightarrow f\bar{f}) = N_c \frac{G_F M_H}{4\pi\sqrt{2}} m_f^2, \quad N_C = 3 (1) \text{ for quarks (leptons)}$$

$$\Gamma(H \rightarrow VV) = \frac{G_F M_H^3}{8\pi\sqrt{2}} F(r) \left(\frac{1}{2} \right)_Z, \quad r = \frac{M_V}{M_H}$$

1975

A Phenomenological Profile of the Higgs Boson

- First attempt at systematic survey

A PHENOMENOLOGICAL PROFILE OF THE HIGGS BOSON

John ELLIS, Mary K. GAILLARD * and D.V. NANOPOULOS **
CERN, Geneva

Received 7 November 1975

A discussion is given of the production, decay and observability of the scalar Higgs boson H expected in gauge theories of the weak and electromagnetic interactions such as the Weinberg-Salam model. After reviewing previous experimental limits on the mass of

We should perhaps finish with an apology and a caution. We apologize to experimentalists for having no idea what is the mass of the Higgs boson, unlike the case with charm [3,4] and for not being sure of its couplings to other particles, except that they are probably all very small. For these reasons we do not want to encourage big experimental searches for the Higgs boson, but we do feel that people performing experiments vulnerable to the Higgs boson should know how it may turn up.

1984

A Preview of the Higgs Boson @ LHC

- Prepared for LHC Lausanne workshop 1984

DEUTSCHES ELEKTRONEN-SYNCHROTRON DESY

DESY 84-071
August 1984
CERN-TH-3943/84

NEW PARTICLES AND THEIR EXPERIMENTAL SIGNATURES

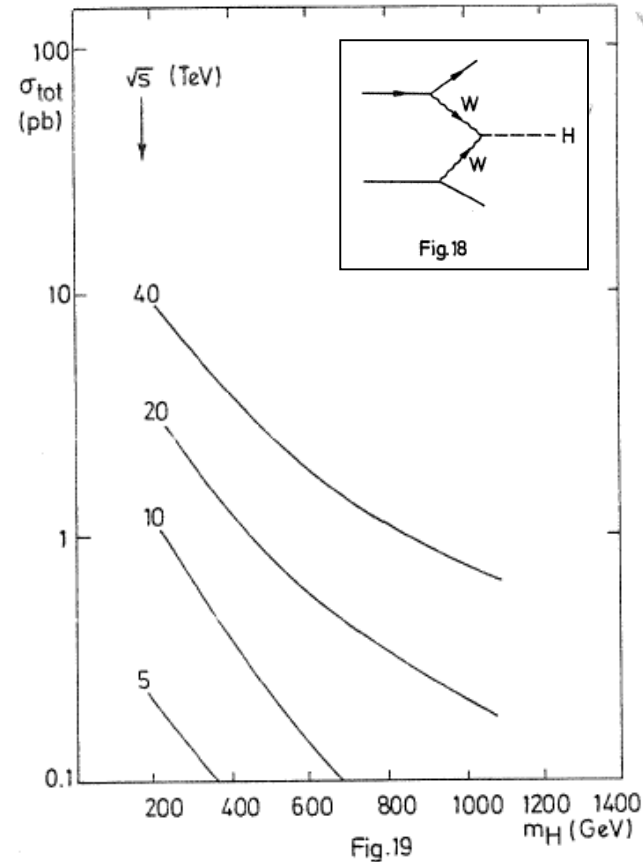
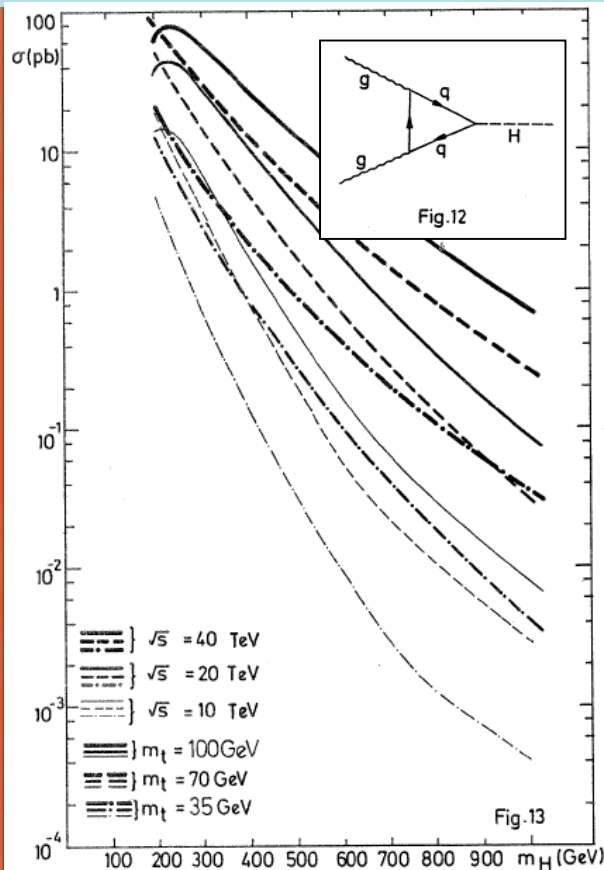
by

J. Ellis and G. Gelmini
CERN, Geneva

H. Kowalski
Deutsches Elektronen-Synchrotron DESY, Hamburg

ISSN 0418-9833

NOTKESTRASSE 85 · 2 HAMBURG 52



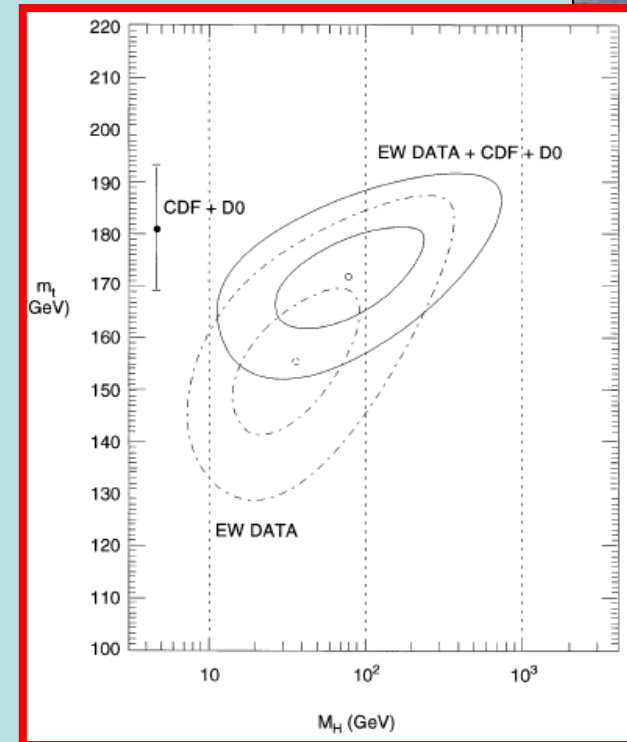
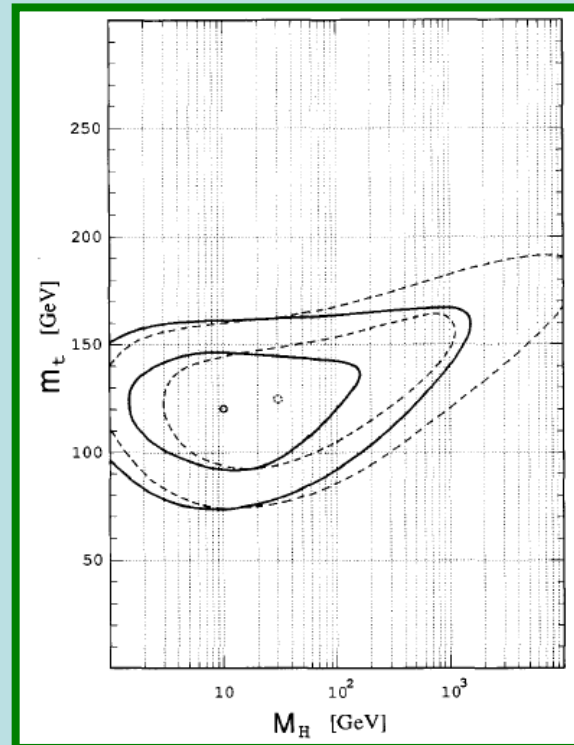
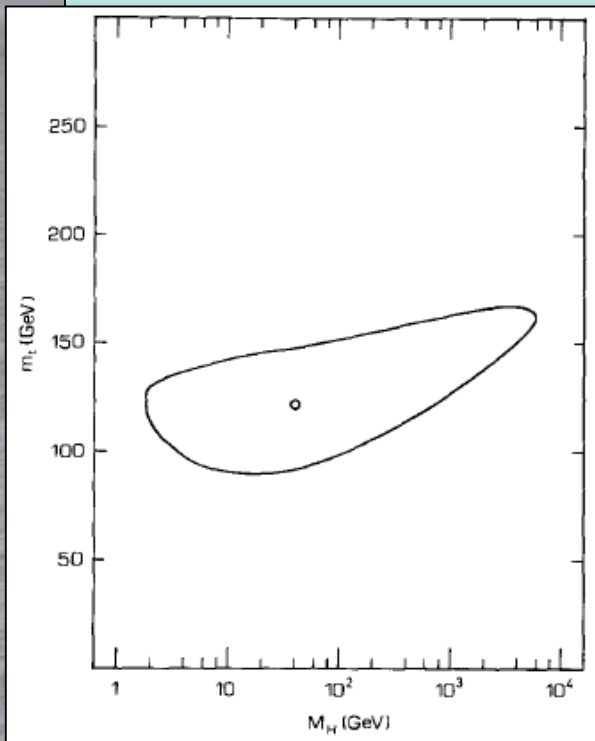
Status of the Standard Model

- Perfect agreement with all *confirmed* accelerator data
- Consistency with precision electroweak data (LEP et al) *only if there is a Higgs boson*
- Agreement seems to require *a relatively light Higgs boson* weighing $< \sim 180 \text{ GeV}$
- Raises many unanswered questions:
mass? flavour? unification?

1990

Estimating the Mass of the Higgs Boson

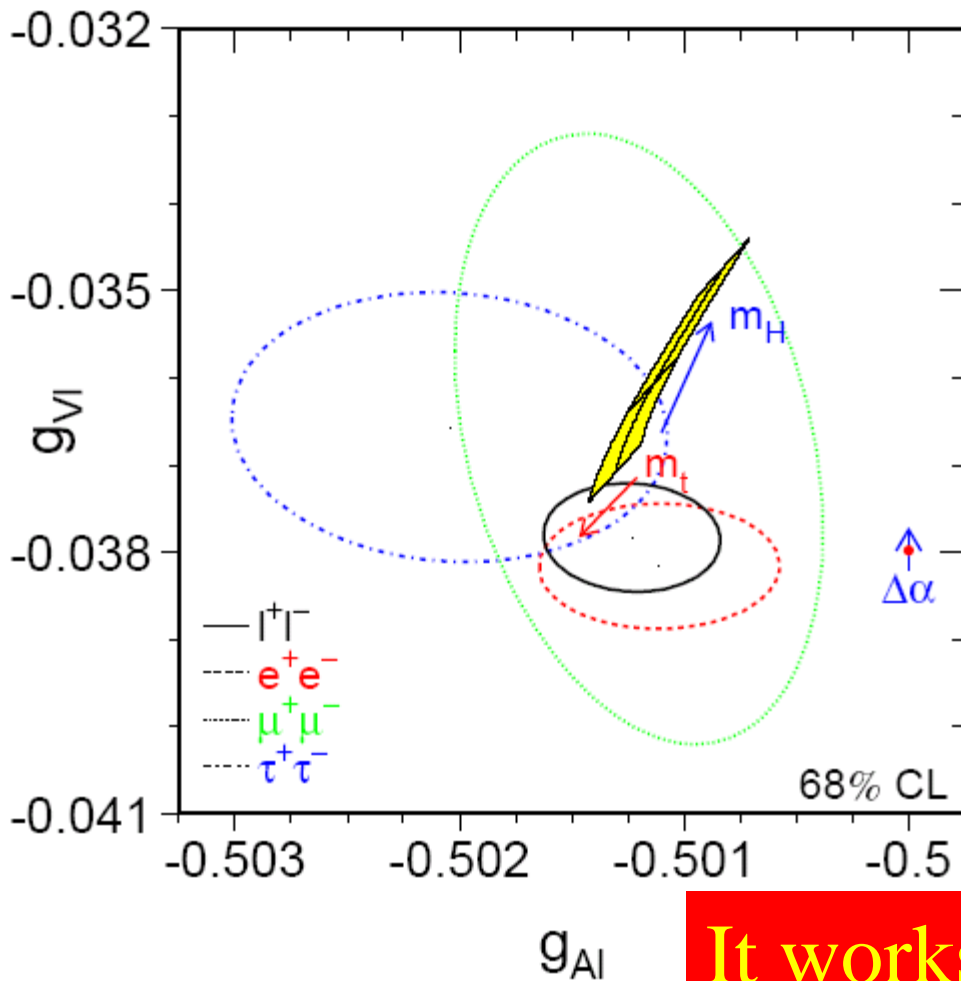
- First attempts in 1990, 1991



- Easier after the discovery of the top quark

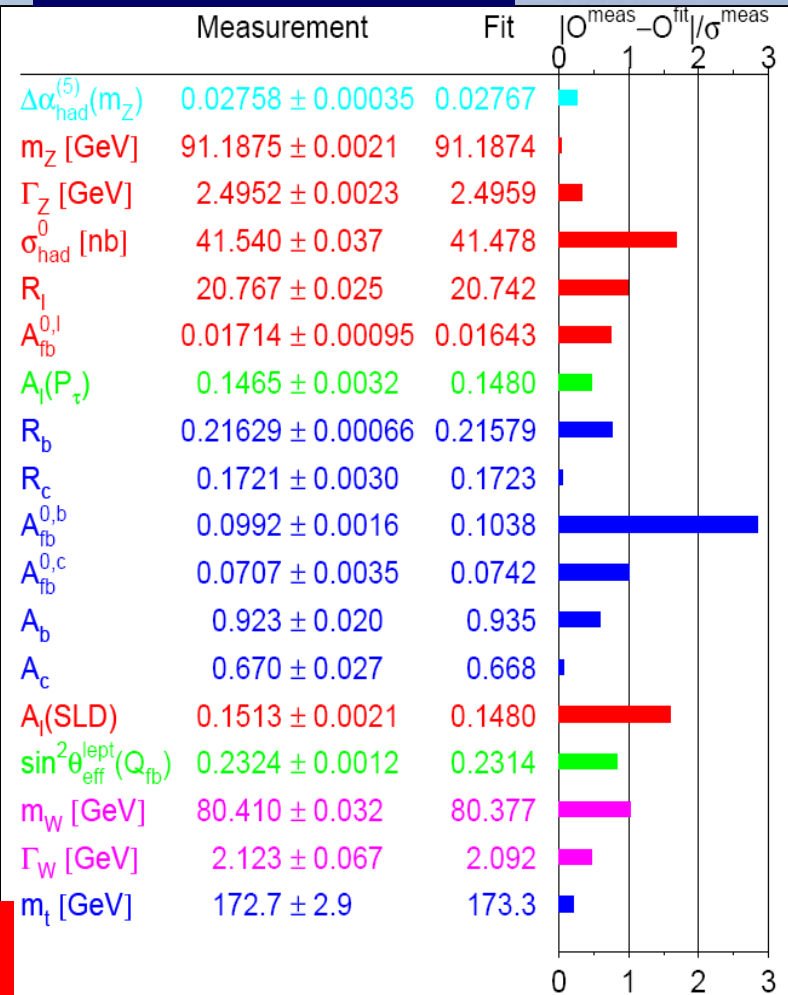
Precision Tests of the Standard Model

Lepton couplings



It works!

Pulls in global fit



Constraints on Higgs Mass

- Electroweak observables sensitive via quantum loop corrections:

$$m_W^2 \sin^2 \theta_W = m_Z^2 \cos^2 \theta_W \sin^2 \theta_W = \frac{\pi\alpha}{\sqrt{2}G_F}(1 + \Delta r)$$

- Sensitivity to top, Higgs masses:

$$\frac{3G_F}{8\pi^2\sqrt{2}}m_t^2$$

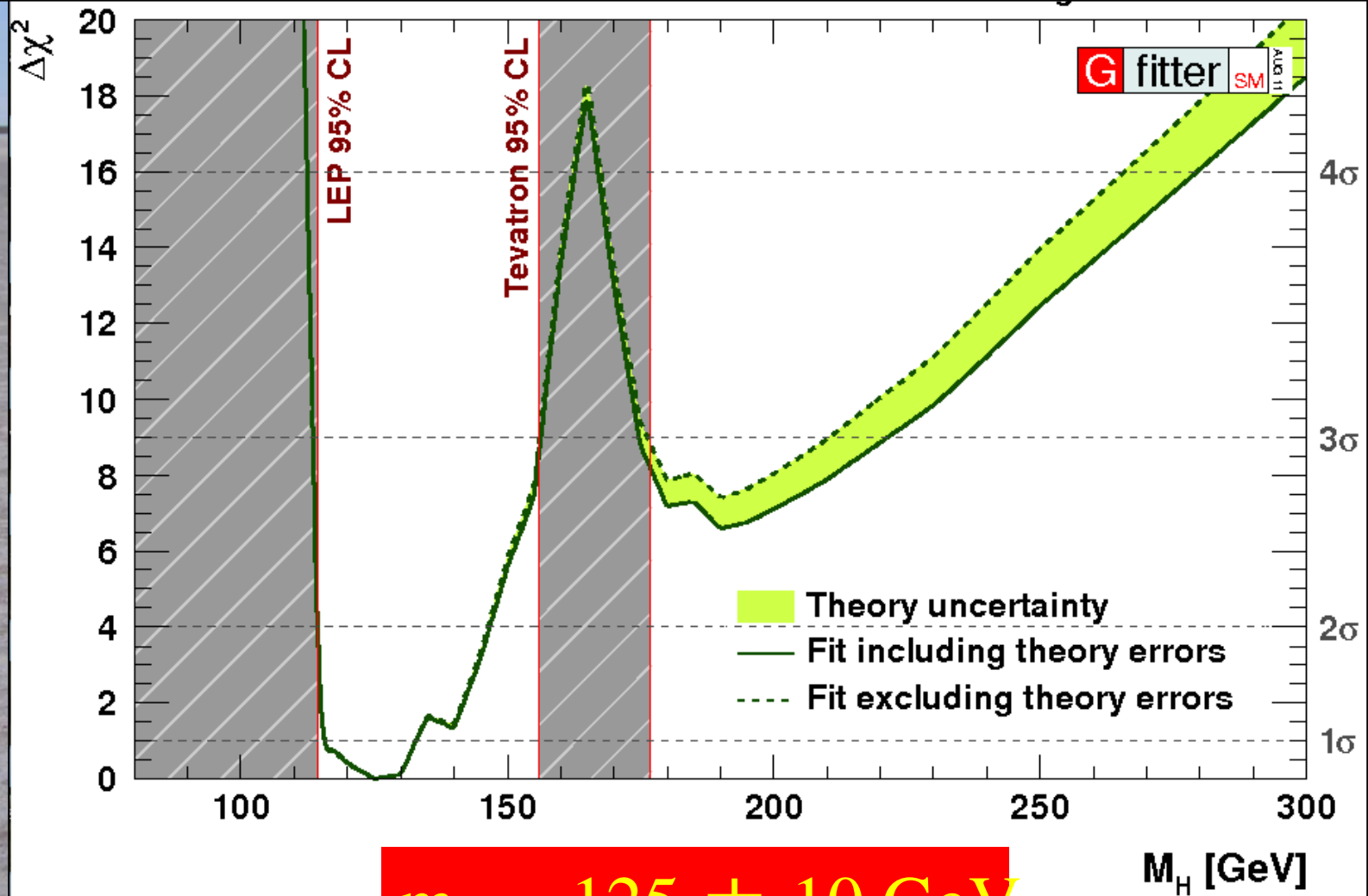
$$\frac{\sqrt{2}G_F}{16\pi^2}m_W^2\left(\frac{11}{3}\ln\frac{M_H^2}{m_Z^2} + \dots\right), M_H \gg m_W$$

- Preferred Higgs mass: **$m_H \sim 100 \pm 30 \text{ GeV}$**
- Compare with lower limit from direct search at LEP:

$$\mathbf{m_H > 114 \text{ GeV}}$$

and exclusion around **(160, 170 GeV)** at TeVatron

2011: Combining Information from Previous Direct Searches and Indirect Data

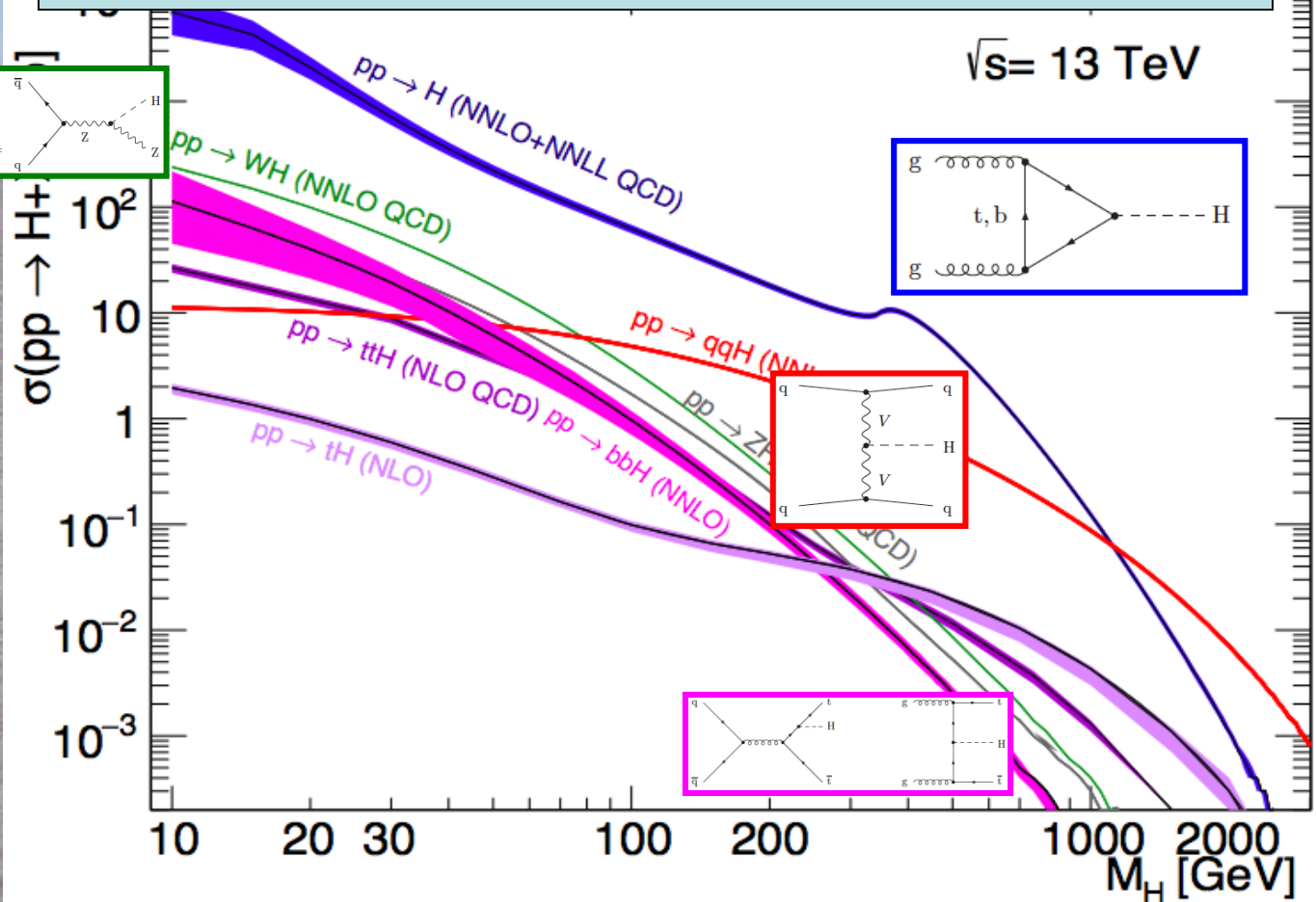
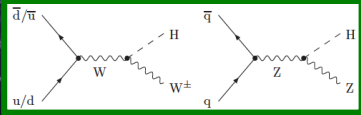
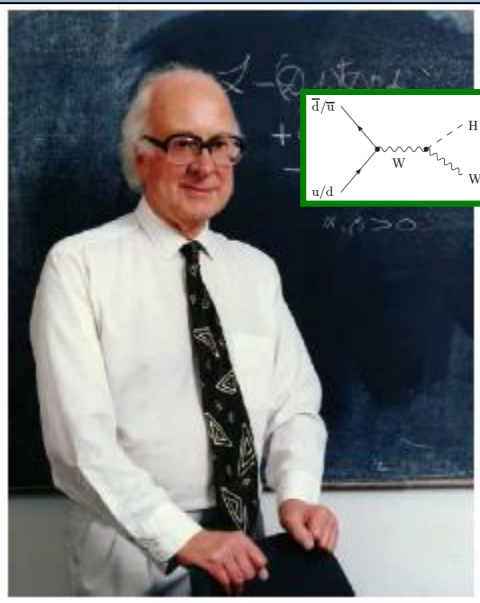


$m_H = 125 \pm 10 \text{ GeV}$

Gfitter collaboration

A la recherche
du
Higgs perdu ...

Higgs Production at the LHC



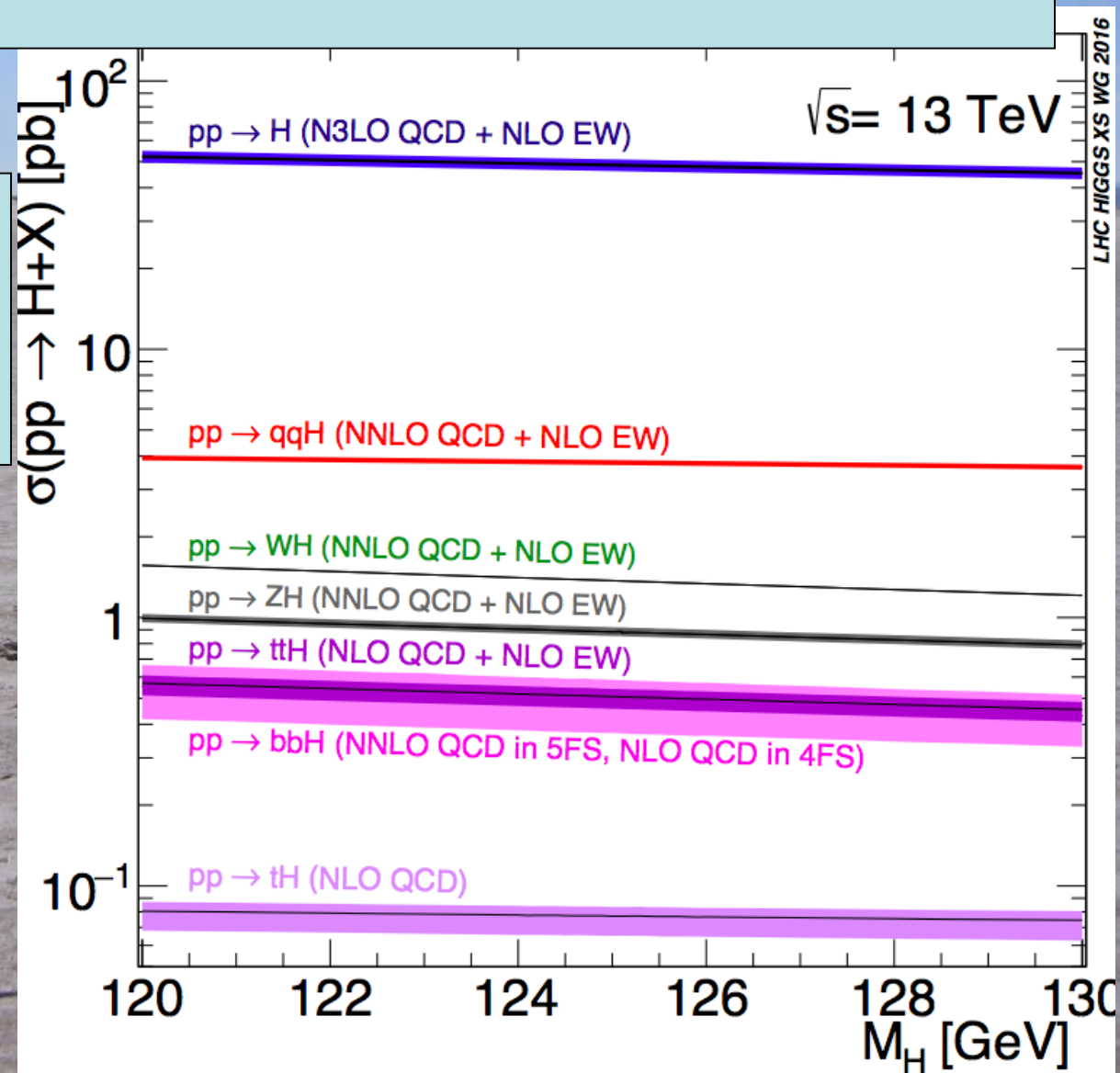
LHC Higgs Cross-Section
Working Group
(LHXS WG)

Many production modes measurable if $M_h \sim 125 \text{ GeV}$

Higgs Production at the LHC

Cross sections for
Higgs mass near
125 GeV

LHC Higgs Cross-Section
Working Group
(LHXS WG)



gg → Higgs Production at the LHC

- Calculated at Next-to-Next-to-Next-to-Leading Order (N³LO) in limit of heavy t quark:

$$\sigma = 48.58 \text{ pb}^{+2.22 \text{ pb} (+4.56\%)}_{-3.27 \text{ pb} (-6.72\%)} \text{ (theory)} \pm 1.56 \text{ pb} (3.20\%) \text{ (PDF}+\alpha_s)$$

- Contributions at different orders:

48.58 pb =	16.00 pb	(+32.9%)	(LO, rEFT)
	+ 20.84 pb	(+42.9%)	(NLO, rEFT)
	- 2.05 pb	(-4.2%)	((t, b, c), exact NLO)
	+ 9.56 pb	(+19.7%)	(NNLO, rEFT)
	+ 0.34 pb	(+0.7%)	(NNLO, 1/m _t)
	+ 2.40 pb	(+4.9%)	(EW, QCD-EW)
	+ 1.49 pb	(+3.1%)	(N ³ LO, rEFT)

- Breakdown of theoretical uncertainties:

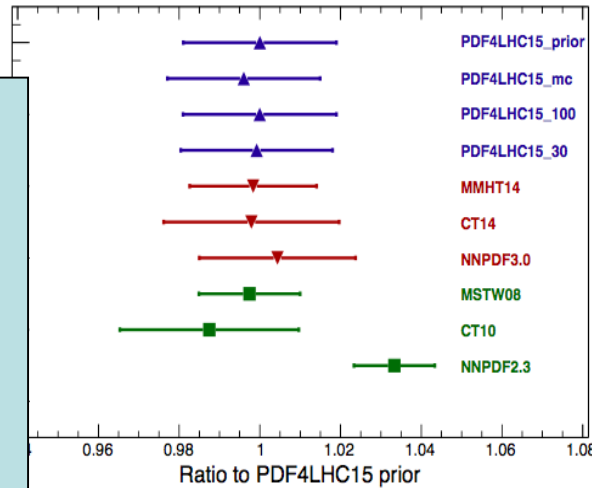
$\delta(\text{scale})$	$\delta(\text{trunc})$	$\delta(\text{PDF-TH})$	$\delta(\text{EW})$	$\delta(t, b, c)$	$\delta(1/m_t)$
+0.10 pb -1.15 pb	±0.18 pb	±0.56 pb	±0.49 pb	±0.40 pb	±0.49 pb
+0.21% -2.37%	±0.37%	±1.16%	±1%	±0.83%	±1%

Dependences on Parton Distributions

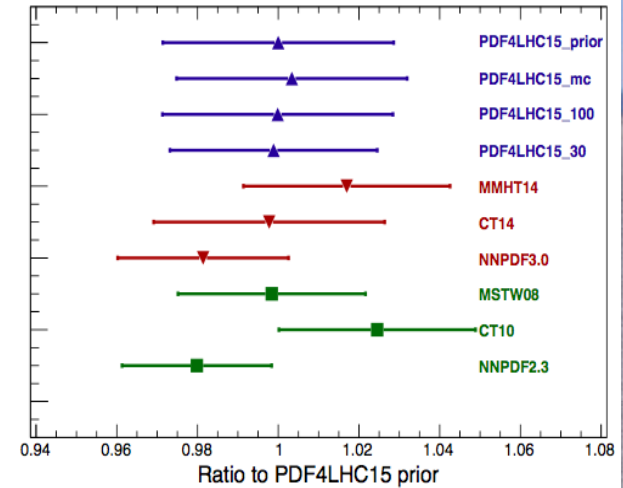
Uncertainties
now $\sim 2\%$,
Similar
uncertainties
in $gg \rightarrow H, ttH$
from α_s

LHC Higgs Cross-Section
Working Group
(LHXS WG)

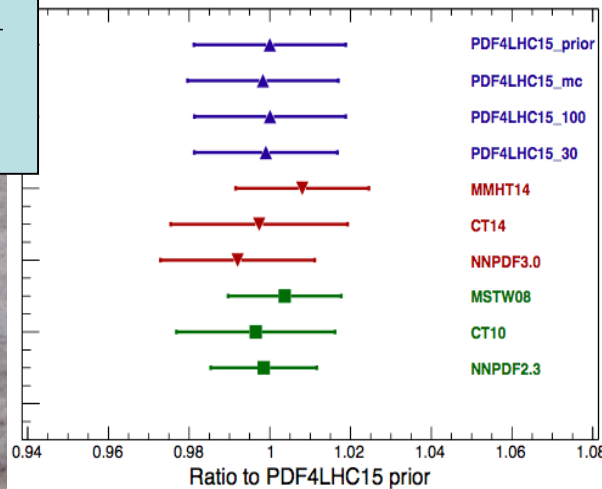
Gluon-Fusion Higgs production, LHC 13 TeV



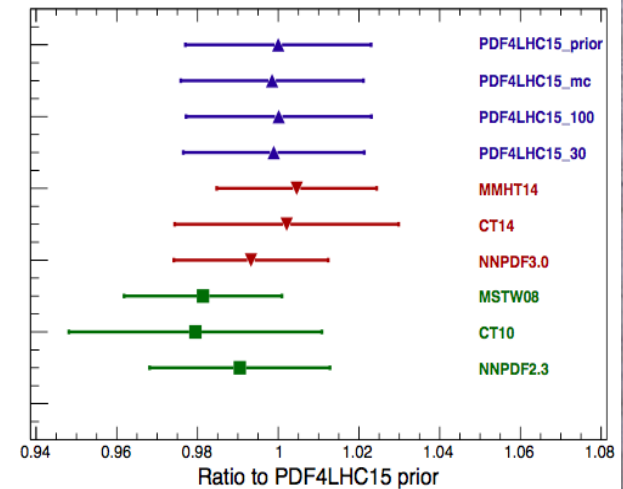
Higgs+ $t\bar{t}$ production, LHC 13 TeV



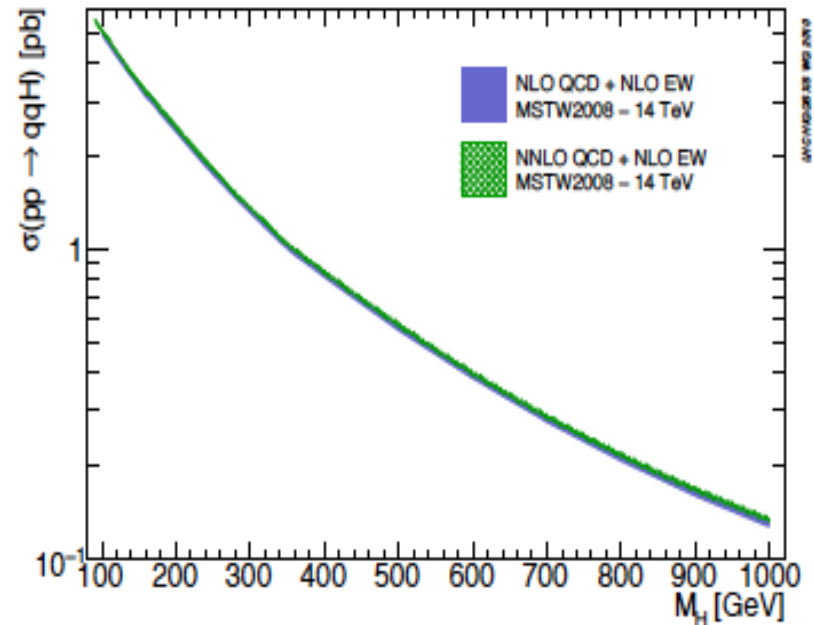
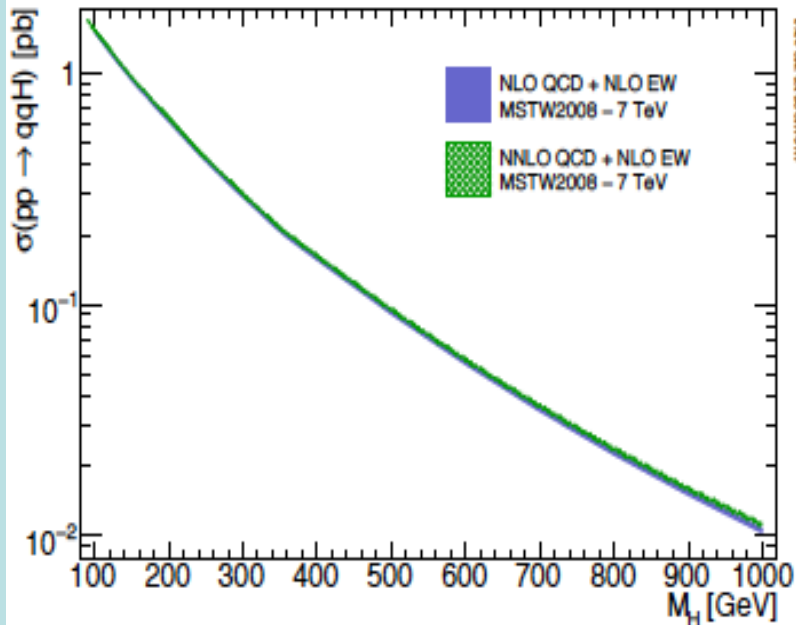
W+Higgs production in a, LHC 13 TeV



Vector-Boson Fusion Higgs production, LHC 13 TeV

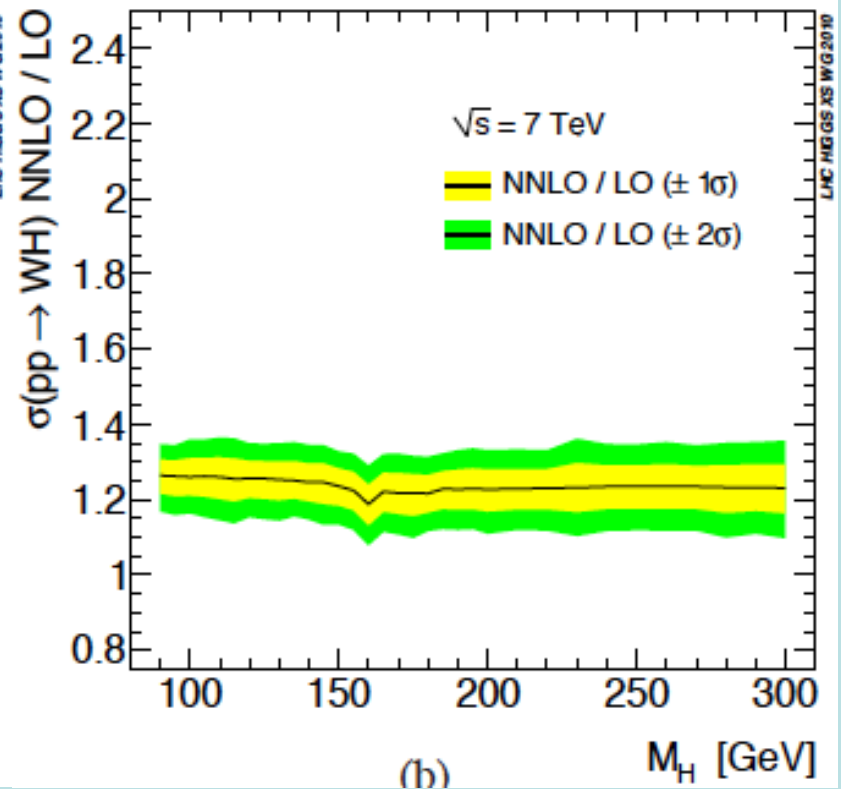
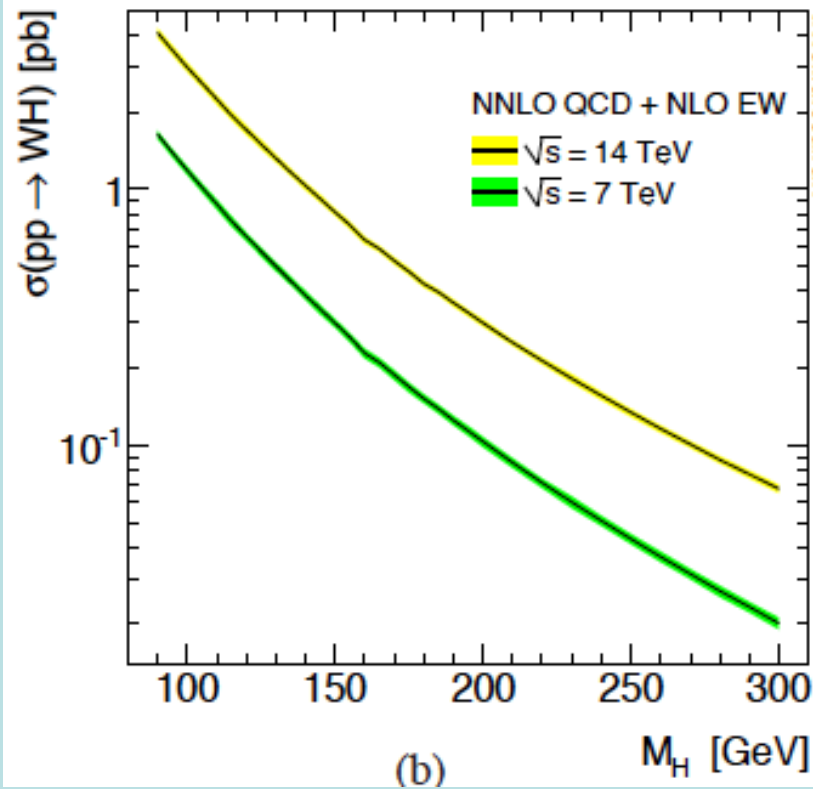


VBF Higgs Production at the LHC



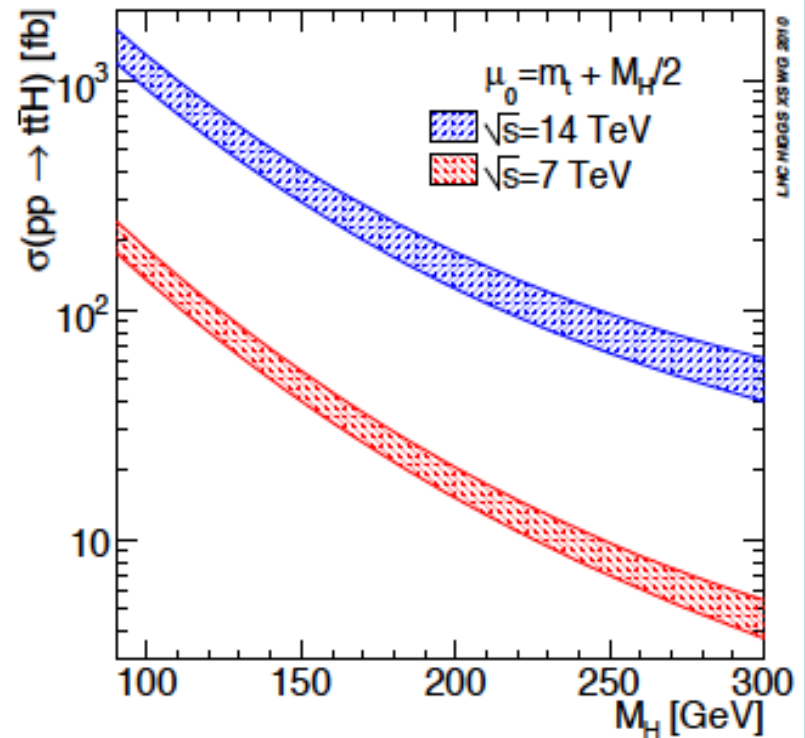
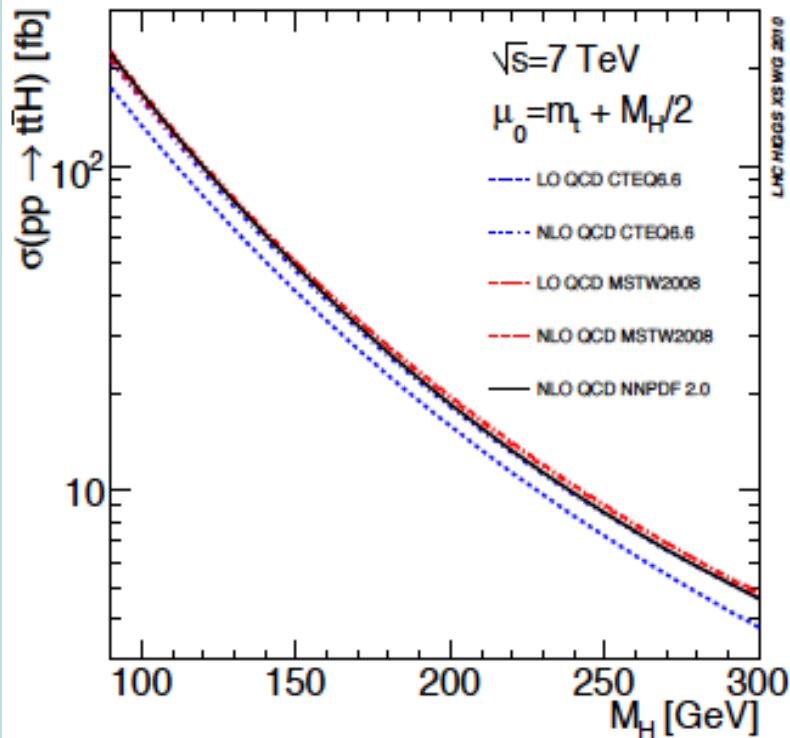
- Calculated at NNLO including electroweak corrections at NLO
- Good convergence of perturbation expansion
- Small uncertainties in quark parton dist'ns

Associated V+H Production at the LHC



- Calculated at NNLO including electroweak corrections at NLO
- Good convergence of perturbation expansion

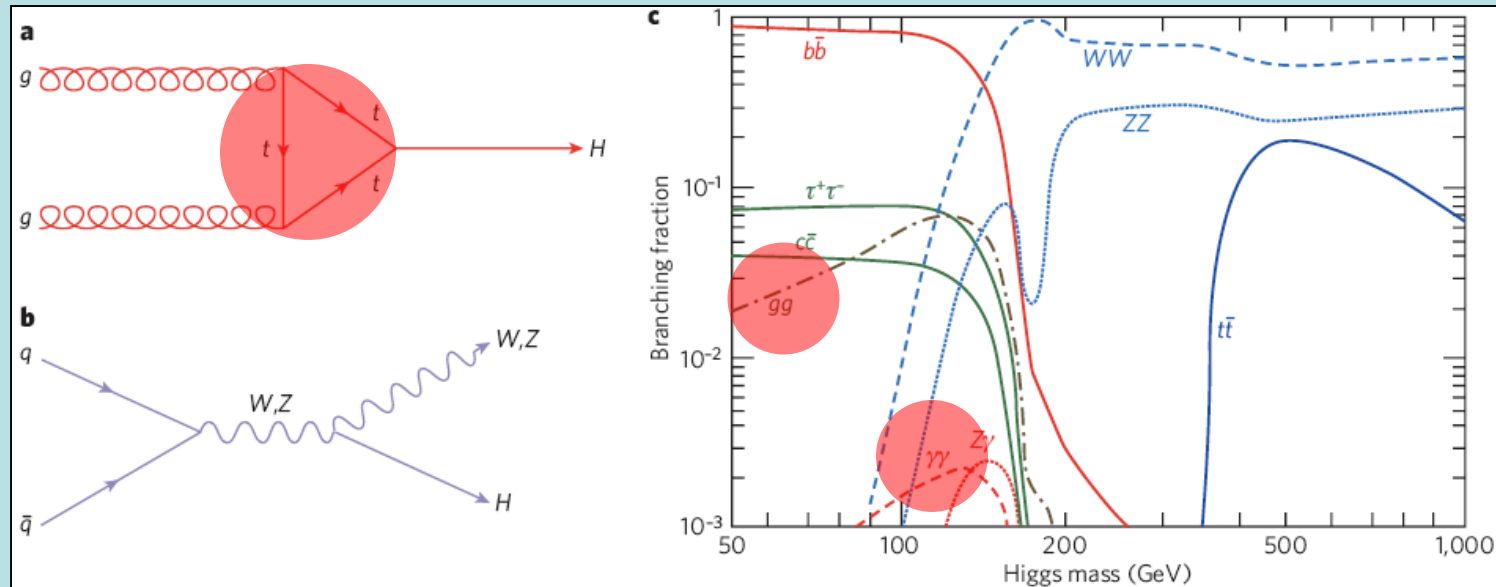
Associated $t\bar{t}+H$ Production



- Calculated at NLO: uncertainties due to
 - Perturbation expansion
 - Choice of parton distributions

Higgs Decay Branching Ratios

- Couplings proportional to masses (?)



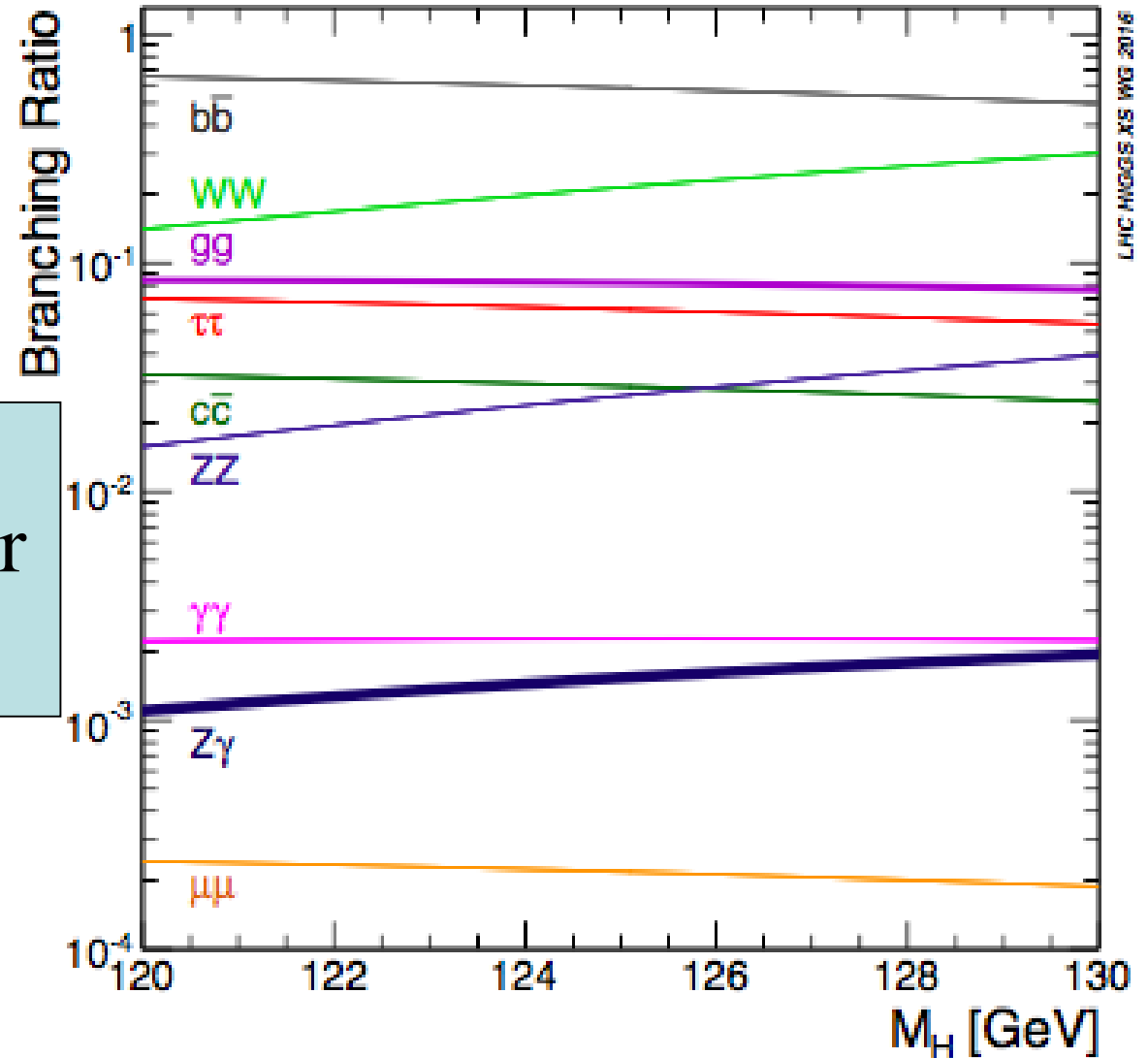
- Important couplings through loops:
 - gluon + gluon \rightarrow Higgs \rightarrow $\gamma\gamma$

Many decay modes measurable if $M_h \sim 125$ GeV

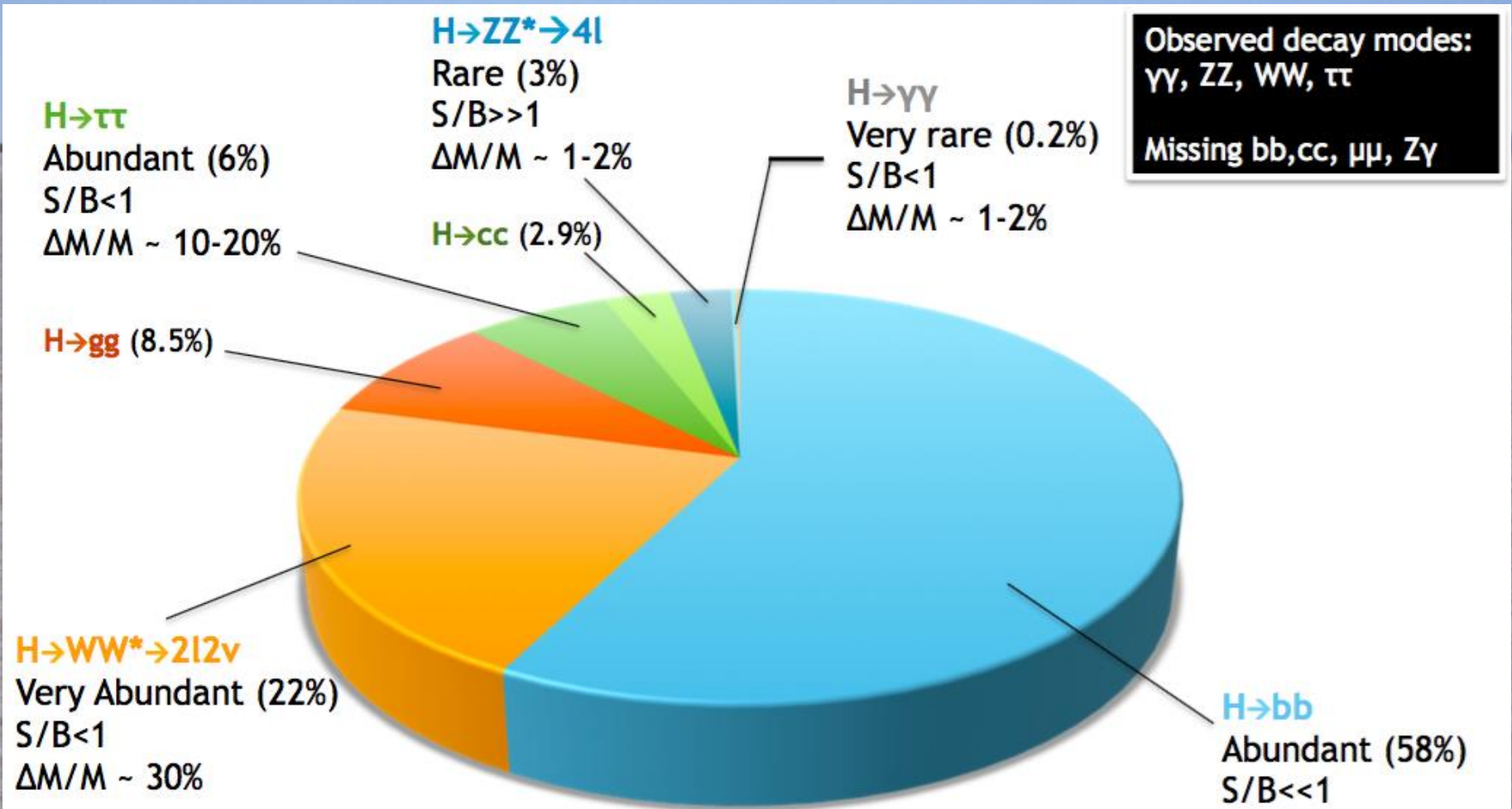
Higgs Decay Branching Ratios

Dominant decay branching ratios for $m_H \sim 125$ GeV

LHC Higgs Cross-Section
Working Group
(LHXS WG)



What we Expect



What do we know?

The Stakes in the Higgs Search

- How is gauge symmetry broken?
- Is there any elementary scalar field?
- **Likely portal to new physics**
- Would have caused phase transition in the Universe when it was about 10^{-12} seconds old
- May have generated then the matter in the Universe: **electroweak baryogenesis**
- A related **inflaton** might have expanded the Universe when it was about 10^{-35} seconds old
- Contributes to today's **dark energy: 10^{60} too much!**

Higgs champagne in Singapore



The wall is over
Champagne bottle, 2012
展览日期
2012年

In 1964, Robert Brout, François Englert, Gerard 't Hooft, Carl Richard Hagen, Tom Kibble and Peter Higgs put forward a bold new idea to explain how fundamental particles acquire mass. Higgs suggested that if they were right, a new particle should exist.

1964年，Robert Brout, François Englert, Gerard 't Hooft, Carl Richard Hagen, Tom Kibble 和 Peter Higgs 提出了一个大胆的新的理论来解释基本粒子如何获得质量。Higgs 提出，如果他们是对的，应该存在一种新的粒子。

On the evening of 3 July 2012, Higgs shared this bottle of champagne with his friends, theoretical physicist John Ellis and former CERN Director-General Chris Llewellyn-Smith. The next day CERN announced the discovery of the particle he had proposed almost five decades earlier, the Higgs boson.

2012年7月3日晚上，Higgs 和他的朋友们——理论物理学家 John Ellis 和前任 CERN 总干事 Chris Llewellyn-Smith 分享了这个香槟瓶。第二天，CERN 宣布发现了 Higgs 粒子，这是他在五十年前提出的粒子。



Peter Higgs at the University of North Carolina, 1965.
展览日期
2012年

In the early 20th century, the electron was the only subatomic particle known. When Ernest Rutherford and his team at Manchester discovered that atoms had another particle at their centre – later named the 'proton' – a new theory was needed to describe the atom's inner structure.

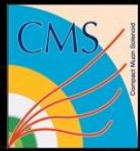
在 20 世纪初，人们只知道一种亚原子粒子，即电子。Ernest Rutherford 和他的团队在曼彻斯特发现原子中心还有另一种粒子——后来被称为“质子”。这就需要一种新的理论来描述原子的内部结构。



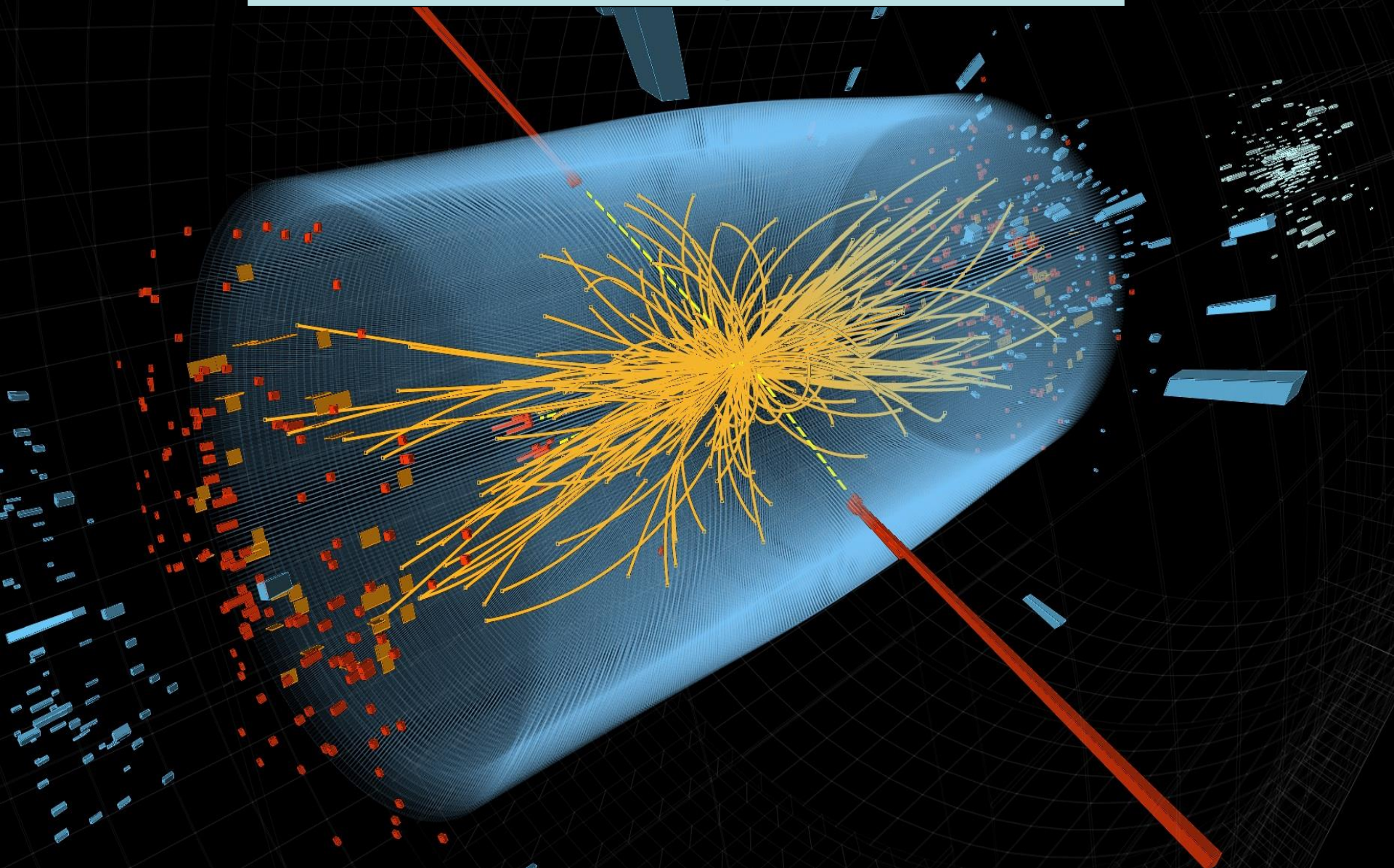
Niels Bohr
Bohr's atomic model, 1913
展览日期
2012年

These models, showing rings of electrons orbiting a nucleus, might be able to be proven by Rutherford in 1911 and revised by Niels Bohr in 1913. This planetary structure became the classic image of the atom in popular culture.

这些模型，展示了电子环围绕原子核，可能在 1911 年由 Rutherford 证明，并在 1913 年由 Niels Bohr 修正。这种行星状结构成为了原子在流行文化中的经典形象。



Interesting Events



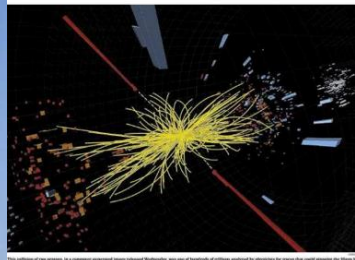
Higgsdependence Day!



July 4th 2012 The Higgs discovery

Discovery upends world of physics

CERN reports finding particle that could solve mysteries large and small
BY MICHAEL COOPER
PARIS — The discovery of a new particle, which scientists believe is the Higgs boson, has upended the world of physics. The particle, which was discovered at the Large Hadron Collider (LHC) at CERN, is the last missing piece of the Standard Model of particle physics. It is the particle that gives other particles their mass. The discovery is a triumph for the LHC and for the scientists who have worked so hard to build it and to understand its results.



The Economist
A giant leap for science
Finding the Higgs boson

Discovery upends world of physics
CERN reports finding particle that could solve mysteries large and small

В ТЕАТРЫ БУДУТ ПУСКАТЬ ПО МОБИЛЬНЫМ ТЕЛЕФОНАМ
MK
ПОСЛЕДНИЙ КИРПИЧ В СТЕНУ МИРОЗДАНИЯ
«КРЕМЛЕВСКИЕ САМОЛЕТЫ ПРИШЛОСЬ МЕНЯТЬ НА ПЕРЕГРABE»
METPO CПУCТAT HA BOДY

又粒子発見か
ミチ子検出年内に結論
日米欧2チーム
Milhares de moradores de bairros sociais em risco de perderem RSI
A mudança está a passar despercebida, mas deve afectar milhares de beneficiários de RSI que vivem em habitação social, agora, morar numa casa construída a uma forma de rendimento Portugal 12

Le boson de Higgs, particule manquante pour expliquer l'Univers, vient d'être découvert
Les physiciens du CERN de Genève ont prouvé son existence à 99,9999%
7,2 milliards de plus dès 2012
Algérie: L'INDÉPENDANCE
Une fête sans panache
La souffrance, mais pas de haine
Ces livres qui explorent l'histoire

AD ALGEMEEN DAGBLAD
Zieke Kaj en zijn moeder toch samen in de VS
EINDELIJK BELIJK NA 48 JAAR
Große Mehrheit im Europäischen Parlament

Frankfurter Allgemeine
Masse macht's
Große Mehrheit im Europäischen Parlament

Physicists Find Elusive Particle Seen as Key to Universe
ROMNEY NOW SAYS HEALTH MANDATE BY OBAMA IS A TAX
Subterfuge at Tankers as Embargo Tightens
Move Algeria, Hen With Conservative Nod
SBIIT BENEWS CRITICISM
Yodpattaya discovery has scientists giddy

The Gazette
EL PAIS
A solas con la prueba del VIH
De Villota pierde el ojo derecho
Pistorius estará en los Juegos
fallada la partícula clave para a comprensión del universo
La Audiencia Nacional imputa a toda la cúpula de Bankia

CHINADAILY
THURSDAY, July 5, 2012
UNDER FIVE FRANK COUNTRIES
ADJUSTING ESSEL PRICE BETTER THAN TAKING CAR DONAZA FOR AS 6
Big bang moment: Scientists may have found 'God particle'

THE TIMES OF INDIA
UNDER FIVE FRANK COUNTRIES
ADJUSTING ESSEL PRICE BETTER THAN TAKING CAR DONAZA FOR AS 6
Big bang moment: Scientists may have found 'God particle'

THE HINDU
Elusive particle found, looks like Higgs boson
CERN physicists hail evidence of game-changing discovery of subatomic particle

CORRIERE DELLA SERA
La particella che può svelare i segreti dell'universo
CERN: il bosone di Higgs, la particella mancante per spiegare l'Universo, viene scoperto

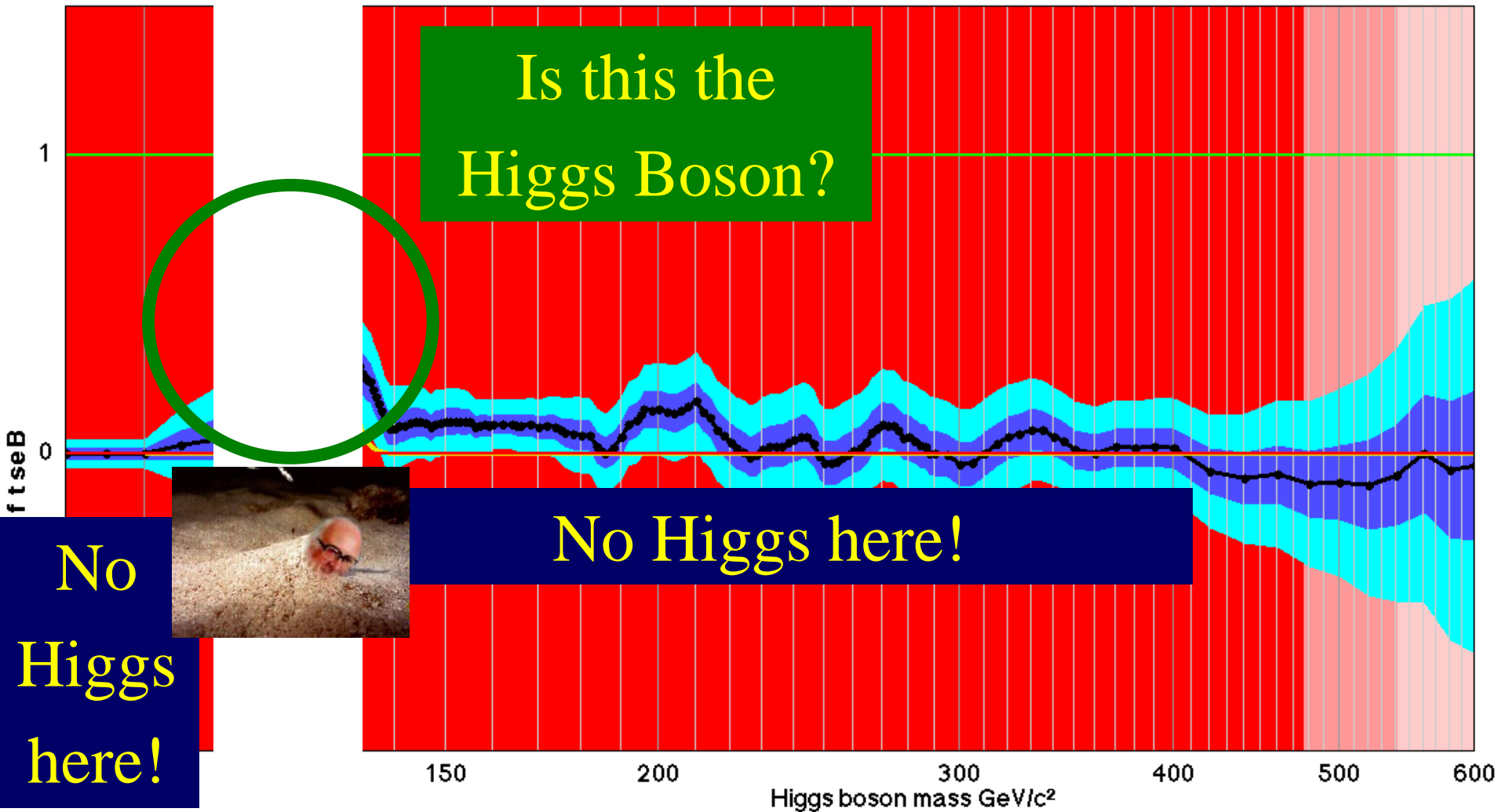
gazeta
WYBORCZA.PL
Ukraincy bije się o Białą Rosyjską
Czastke Higgsa fizycy najpierw wymyślili, potem szukali 40 lat
BOSKA MASA

বিশ্বজ্ঞানের 'সিঁধুর' দর্শন
সত্যেন্দ্রনাথকে বিনস প্রণাম
পেয়েছি, যা খুঁজছিলাম

Unofficial Combination of Higgs Data

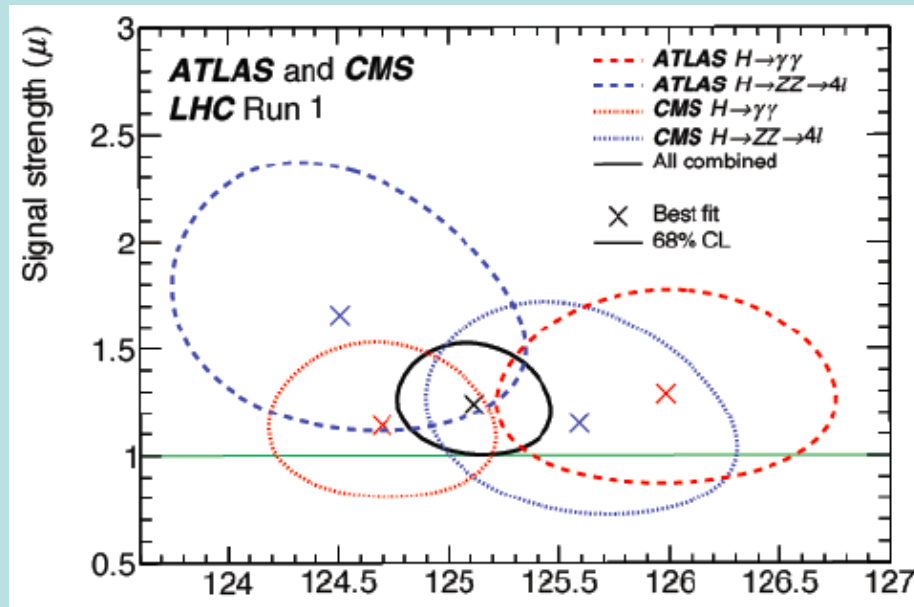
1/fb - 10/fb

06/03/2013



Higgs Mass Measurements

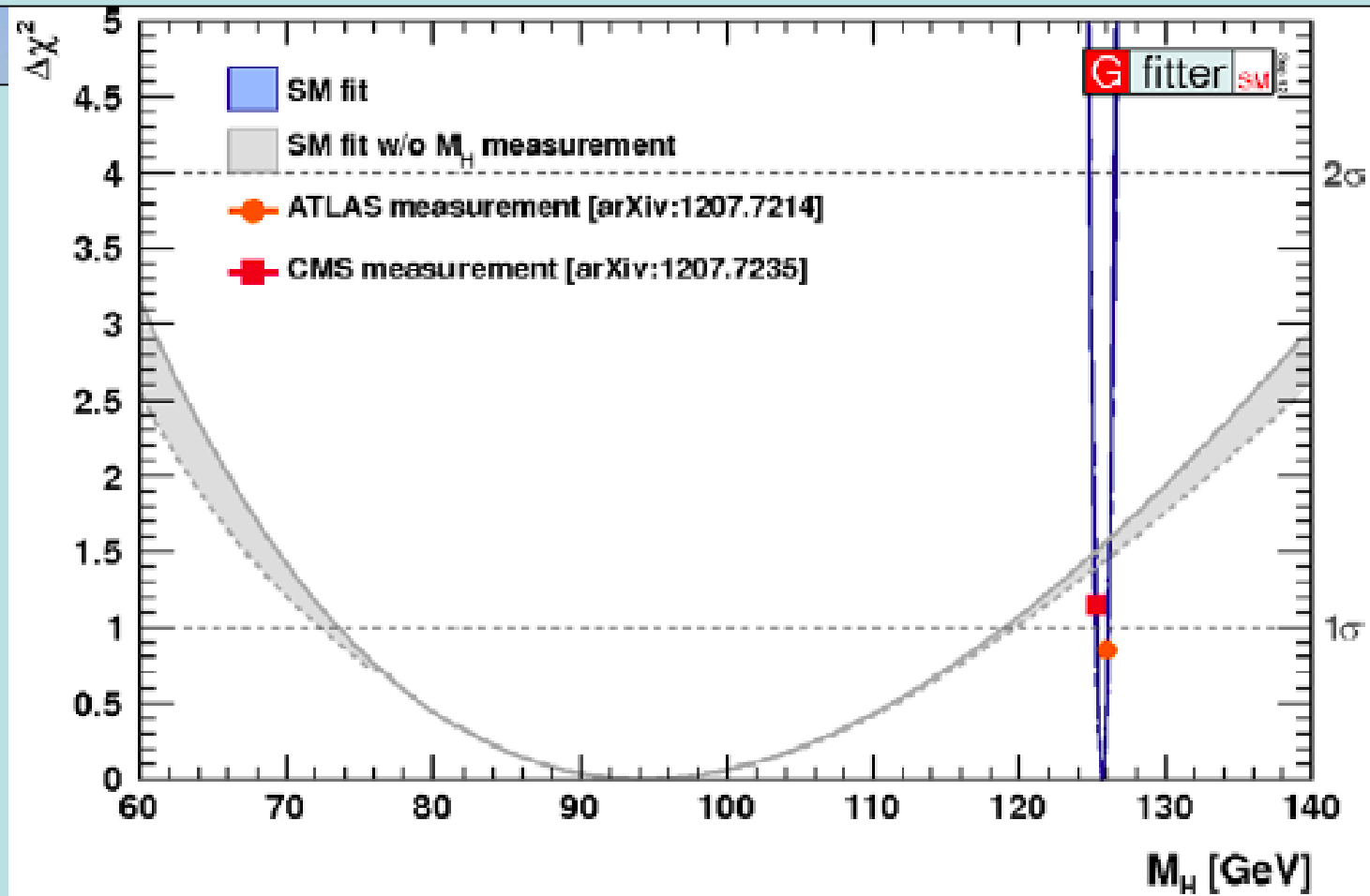
- ATLAS + CMS ZZ^* and $\gamma\gamma$ final states



125.09 ± 0.21 (stat) ± 0.11 (syst)

- Statistical uncertainties dominate
- Allows precision tests
- **Crucial for stability of electroweak vacuum**

Comparison with Electroweak Fit



Quite consistent: $\Delta\chi^2 \sim 1.5$

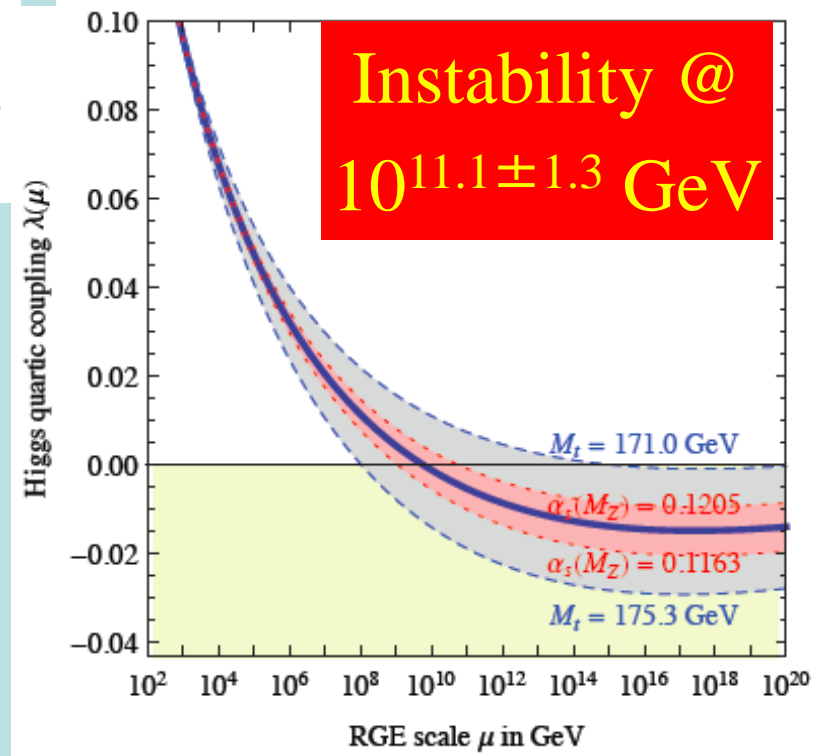
Theoretical Constraints on Higgs Mass

- Large $M_h \rightarrow$ large self-coupling \rightarrow blow up at

$$\lambda(Q) = \lambda(v) - \frac{3m_t^4}{2\pi^2 v^4} \log \frac{Q}{v}$$

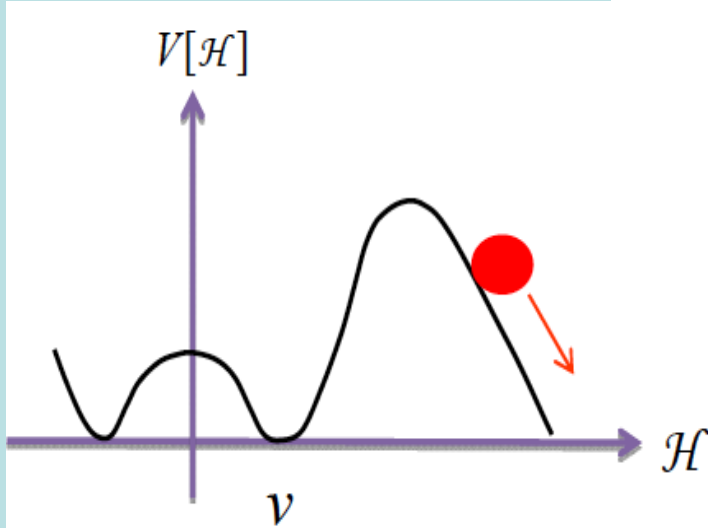
- Small: renormalization due to t quark drives quartic coupling < 0 at some scale Λ
 \rightarrow vacuum unstable

- Vacuum could be stabilized by **Supersymmetry**



Vacuum Instability in the Standard Model

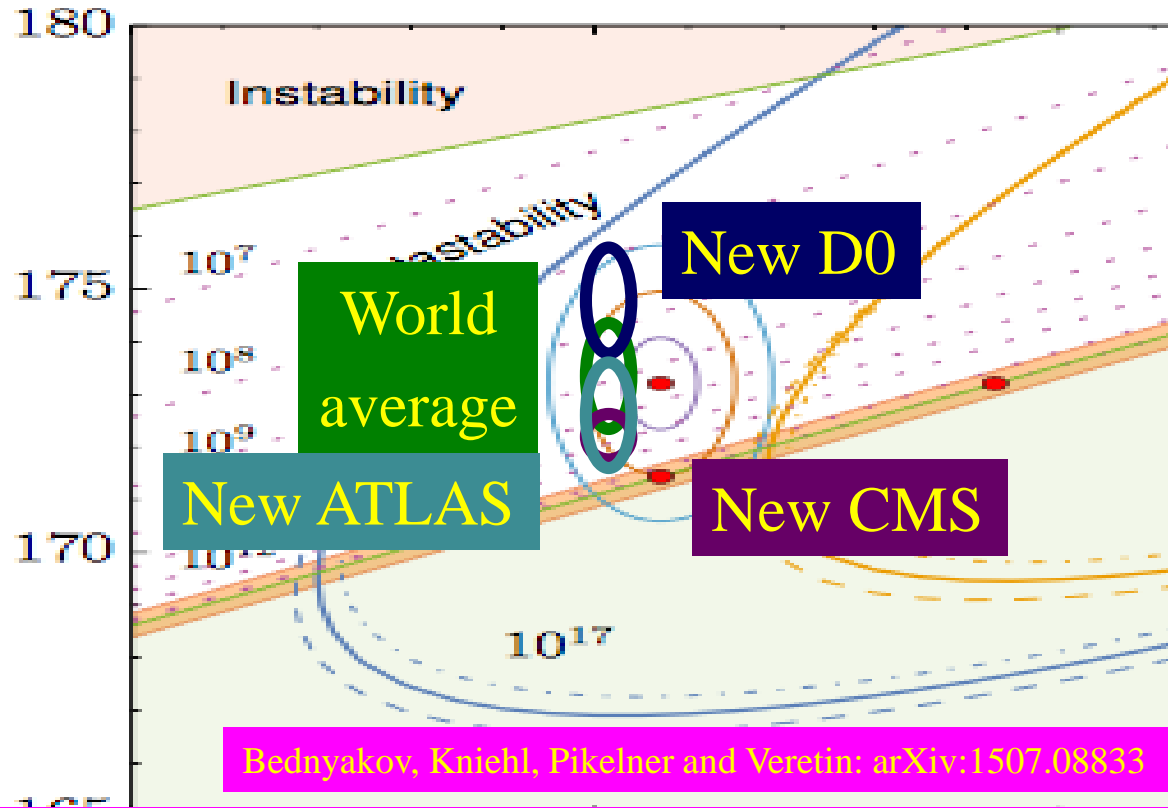
- Very sensitive to



- Instability scale

$$\log_{10} \frac{\Lambda_I}{\text{GeV}} = 11.3 + 1.0 \left(\frac{M_h}{\text{GeV}} - 125.66 \right) - 1.2 \left(\frac{M_t}{\text{GeV}} - 173.10 \right) + 0.4 \frac{\alpha_3(M_Z) - 0.1184}{0.0007}$$

$$m_t = 173.3 \pm 1.0 \text{ GeV} \rightarrow \log_{10}(\Lambda/\text{GeV}) = 11.1 \pm$$



Bednyakov, Kniehl, Pikelner and Veretin: arXiv:1507.08833

Buttazzo, Degrandi, Giardino, Giudice, Sala, Salvio & Strumia, arXiv:1307.3536

Hard QCD: the Top Mass

- Basic parameter of SM; **stability of EW vacuum?**

- World average:

$$m_t = 173.34 \pm 0.76 \text{ GeV}$$

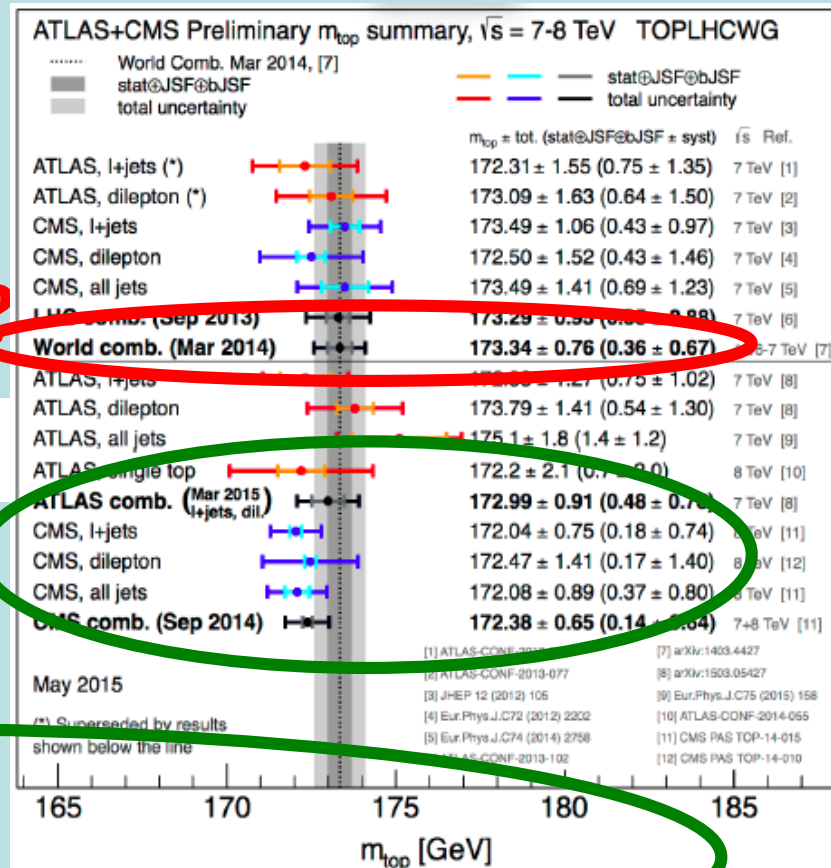
- **Running \rightarrow pole mass OK?**

$$+ 7.557 + 1.617 + 0.501 + 0.195)$$

- Monte Carlo mass \checkmark ?
- New measurements:

$$\text{ATLAS: } 172.99 \pm 0.91 \text{ GeV}$$

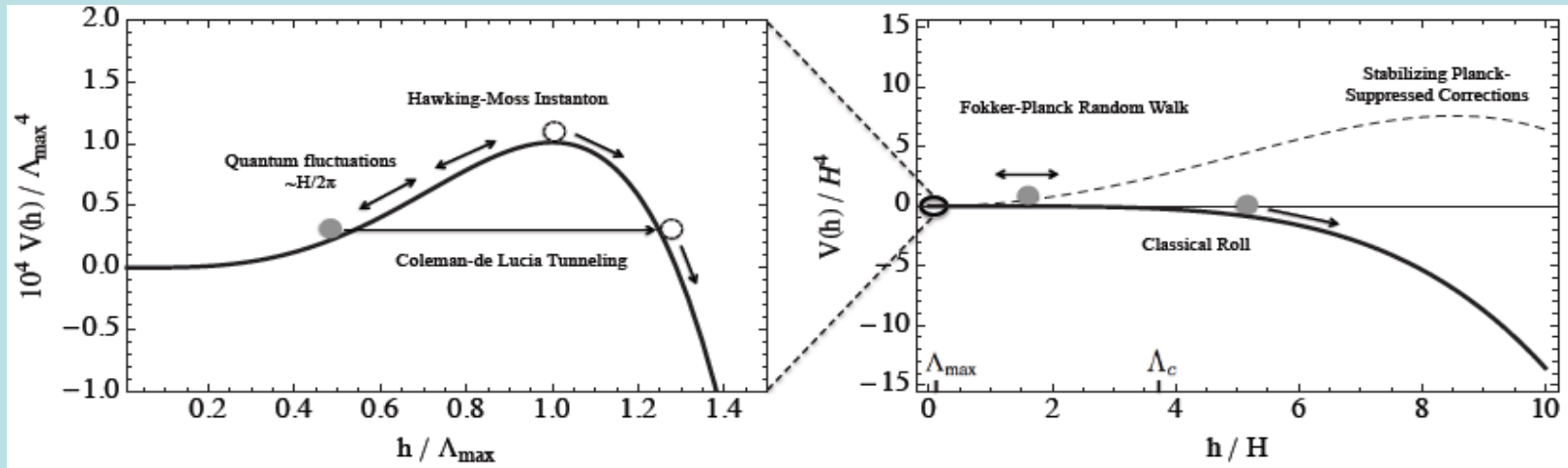
$$\text{CMS: } 172.38 \pm 0.65, \text{ D0: } 174.98 \pm 0.58 \pm 0.49 \text{ GeV}$$



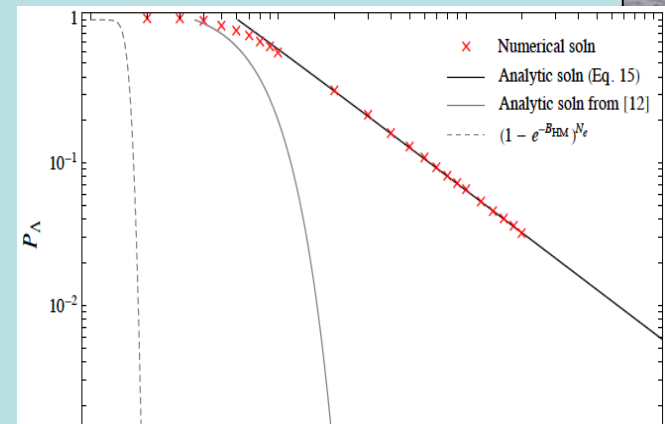
Instability during Inflation?

Hook, Kearns, Shakya & Zurek: arXiv:1404.5953

- Do inflation fluctuations drive us over the hill?



- Then Fokker-Planck evolution
- Do AdS regions eat us?
 - Disaster if so
 - If not, OK if more inflation



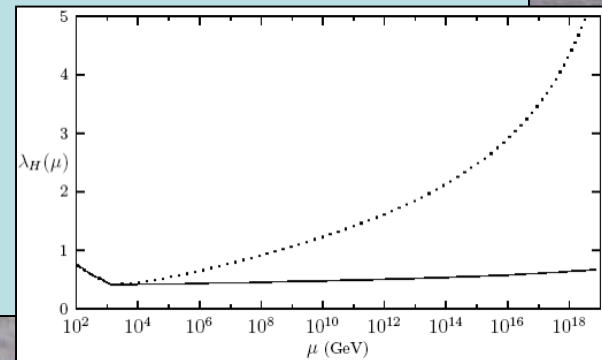
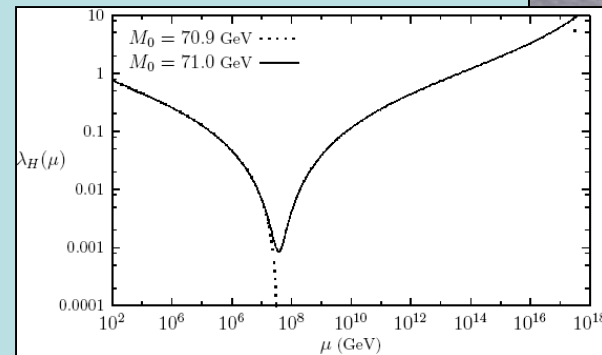
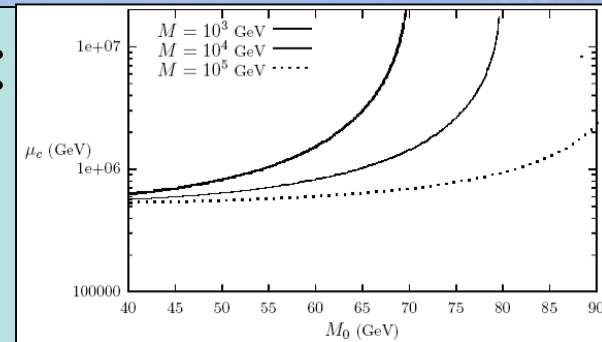
Stabilize vacuum with some physics beyond the SM?

How to Stabilize a Light Higgs Boson?

- Top quark destabilizes potential:
introduce stop-like scalar:

$$\mathcal{L} \supset M^2 |\phi|^2 + \frac{M_0}{v^2} |H|^2 |\phi|^2$$

- Can delay collapse of potential:
- But new coupling must be fine-tuned to avoid blow-up:
- Stabilize with new fermions:
 - just like Higgsinos
- Very like **Supersymmetry!**



‘God Particle’ no big Deal?

- Peter Higgs as quoted in the London Times:
- *“A discovery widely acclaimed as the most important scientific advance in a generation has been “overhyped”, the British scientist behind it has said.”*

Without Higgs ...

... there would be no atoms

- Electrons would escape at the speed of light

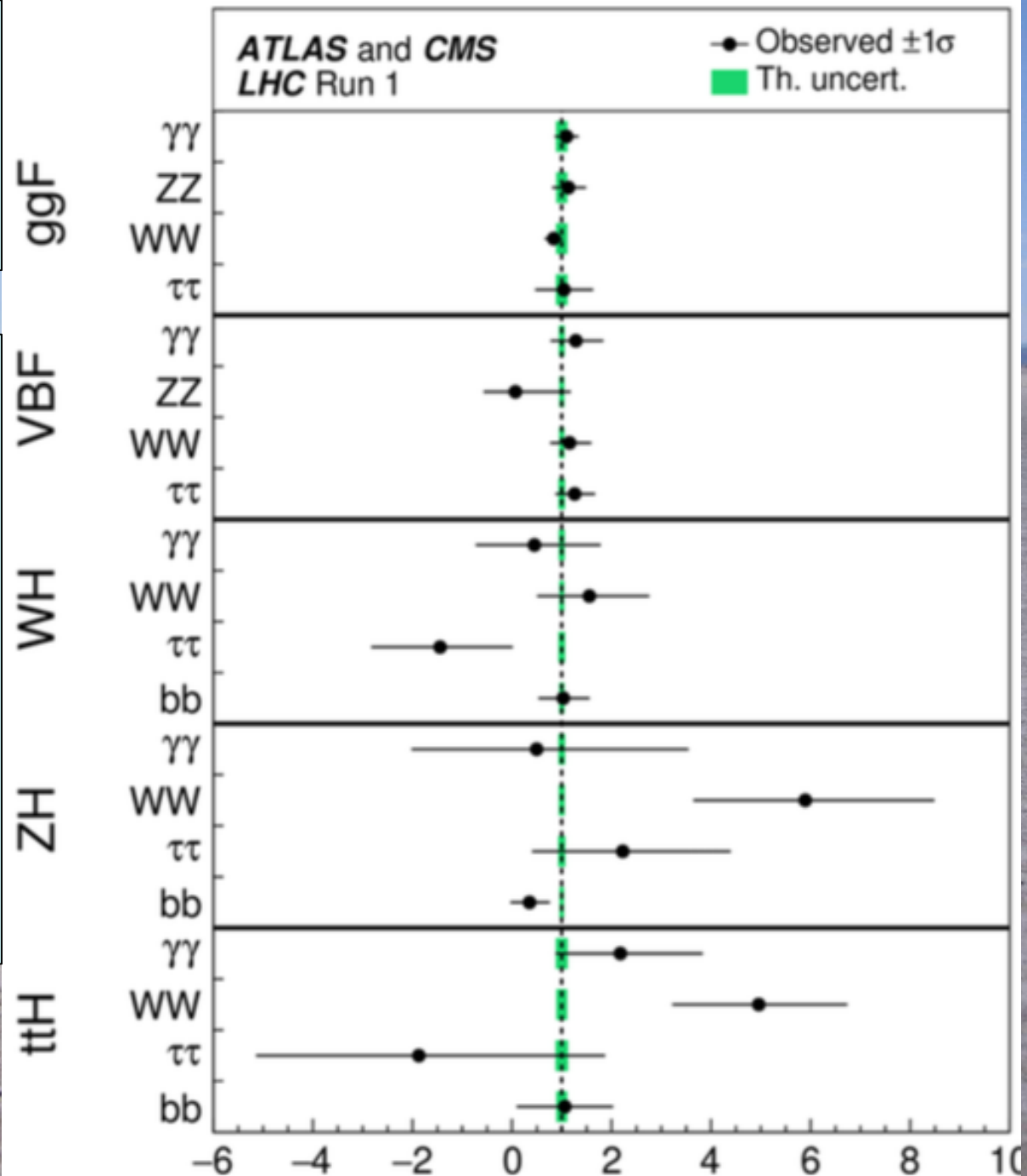
... weak interactions would not be weak

- Life would be impossible: there would be no nuclei, everything would be radioactive

The discovery of the Higgs Boson is a big deal

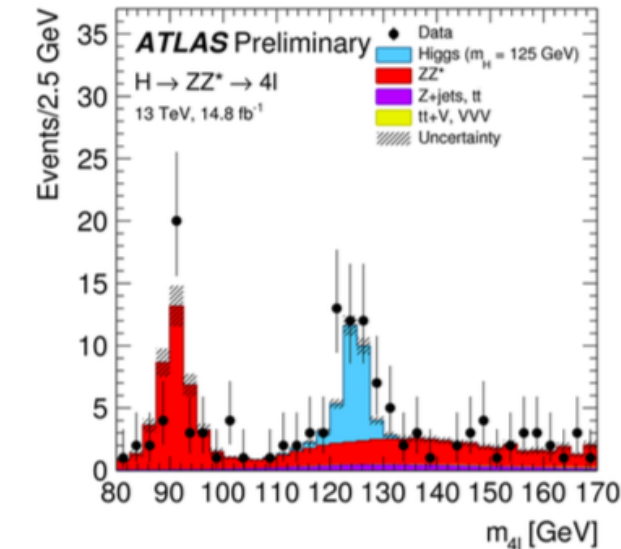
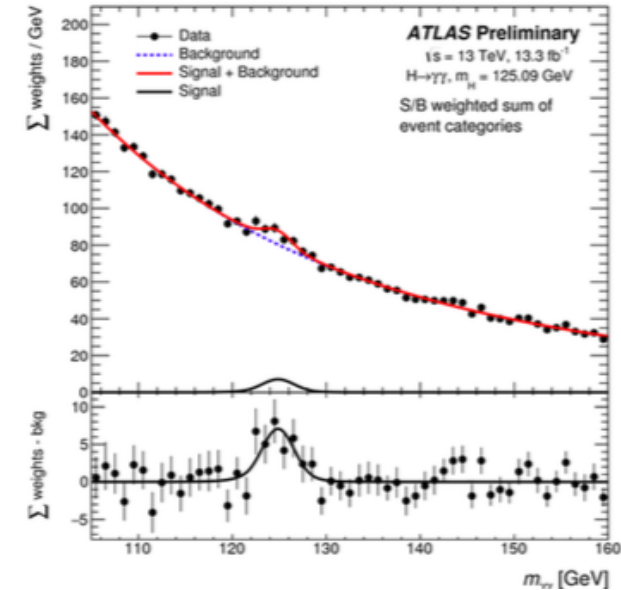
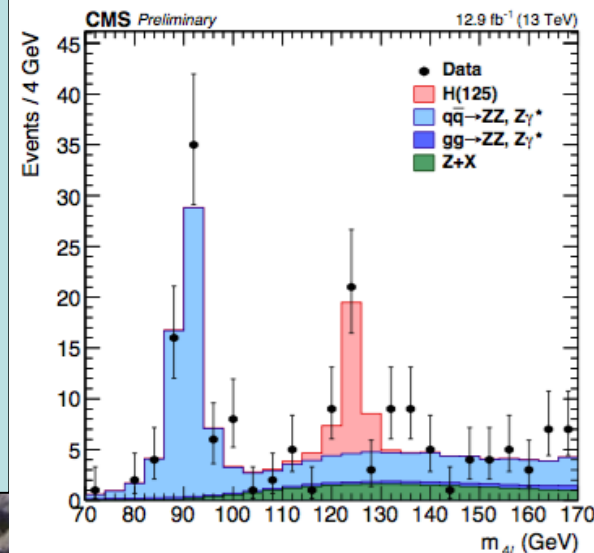
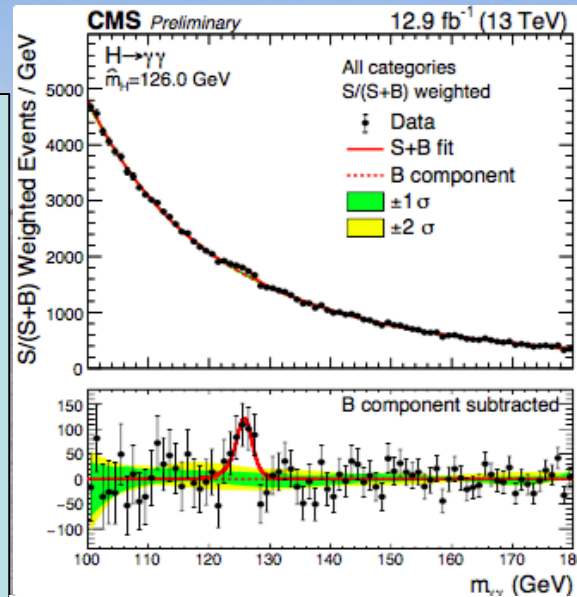
Measurements in Run 1

- Open questions:
 - $H \rightarrow bb$?
 - 2.6σ @ LHC
 - 2.8σ @ FNAL
 - $H \rightarrow \mu\mu$?
 - ttH production?
 - tH production?



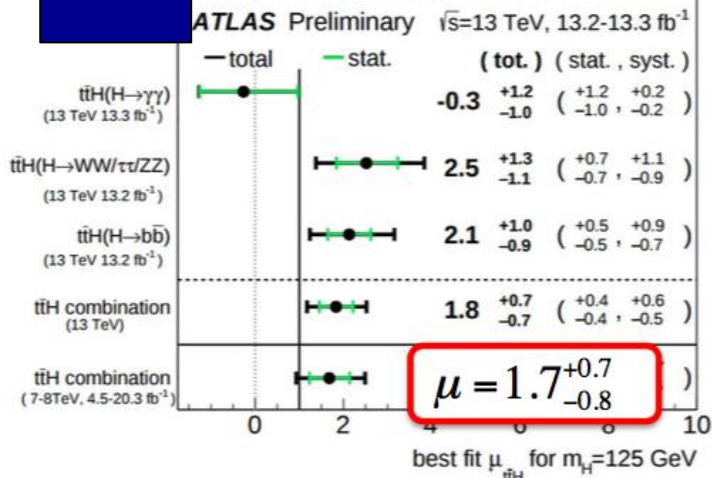
Measurements in Run 2

- Measurements in $\gamma\gamma$ and ZZ^* final states
- 10σ significance
- Cross-section agrees with theory
- Searching for ttH , $H \rightarrow \mu\mu$ and tH

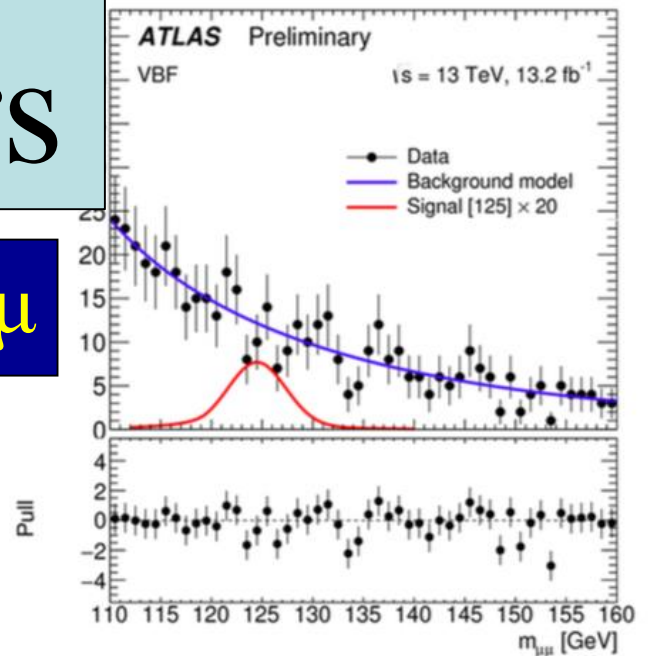
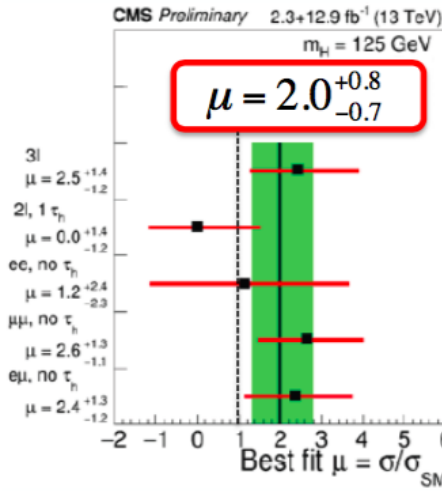


The Next Frontiers

ttH



$H \rightarrow \mu\mu$



ATLAS	Upper limit x SM (expected)
Run 1	7.1 (7.2)
Run 2	4.4 (5.5)
Combined Run 1 and Run 2	3.5 (4.5)

tH

CMS Upper limit x SM (expected)

SM 113.7 (98.6)

Opposite sign of t Yukawa coupling 6.0 (6.4)

