

Non-equilibrium dilepton production in hadronic transport approaches

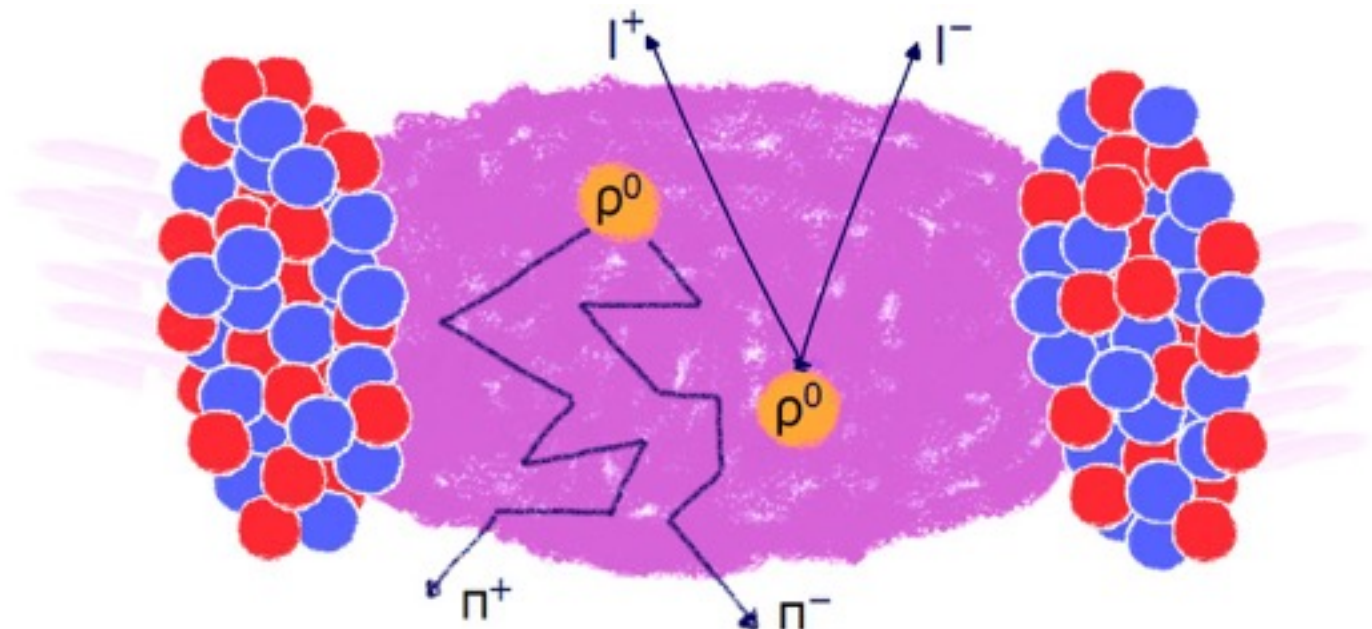
Jan Staudenmaier, G. Vujanovic, J. Weil, U. Heinz
and H. Petersen

Hot Quarks 2016, Padre Island, Texas, USA



Motivation

- Why use dileptons?
 - only electromagnetic interaction: clean probe for hadronic matter
 - extract medium properties and medium modification over whole lifetime of collision
 - > valuable to constrain dynamical models



Motivation

- no description from first principles for HIC → need for models
- understand experimental data
- combination of approaches for description of HIC
- Transport models
 - late dilute stages (e.g. important for identified particle spectra)
 - non-equilibrium
 - low beam energy collisions

<u>stage</u>	<u>model</u>
early	non-equilibrium initial state (e.g. fluctuating color fields/strings)
hot & dense	relativistic dissipative hydrodynamics with equation of state from QCD lattice calculations
dilute	hadronic transport

SMASH

- new hadronic transport approach for dilute non-equilibrium stages of HIC and low energy collisions
- SMASH = **S**imulating **M**any **A**ccelerated **S**trongly-interacting **H**adrons
- developed in group of Hannah Petersen at FIAS with modern tools (C++, Git, Redmine, Doxygen), version 0.9
- goal: standard reference for hadronic system with vacuum properties
- recent paper: *J.Weil et al, arXiv:1606.06642*

General Setup

- transport models based on the relativistic Boltzmann equation

$$p^\mu \partial_\mu f_i(x, p) + m_i F^\alpha \partial_\alpha^p f_i(x, p) = C_{\text{coll}}^i$$

- Collision Criterion: geometric criterion

$$d_{\text{trans}} < d_{\text{int}} = \sqrt{\frac{\sigma_{\text{tot}}}{\pi}}$$

- modes: nuclear collisions, infinite matter calculations, afterburner for hydrodynamic simulations
- features: Test Particle Method, Mean-Field potentials, Fermi motion, Pauli blocking

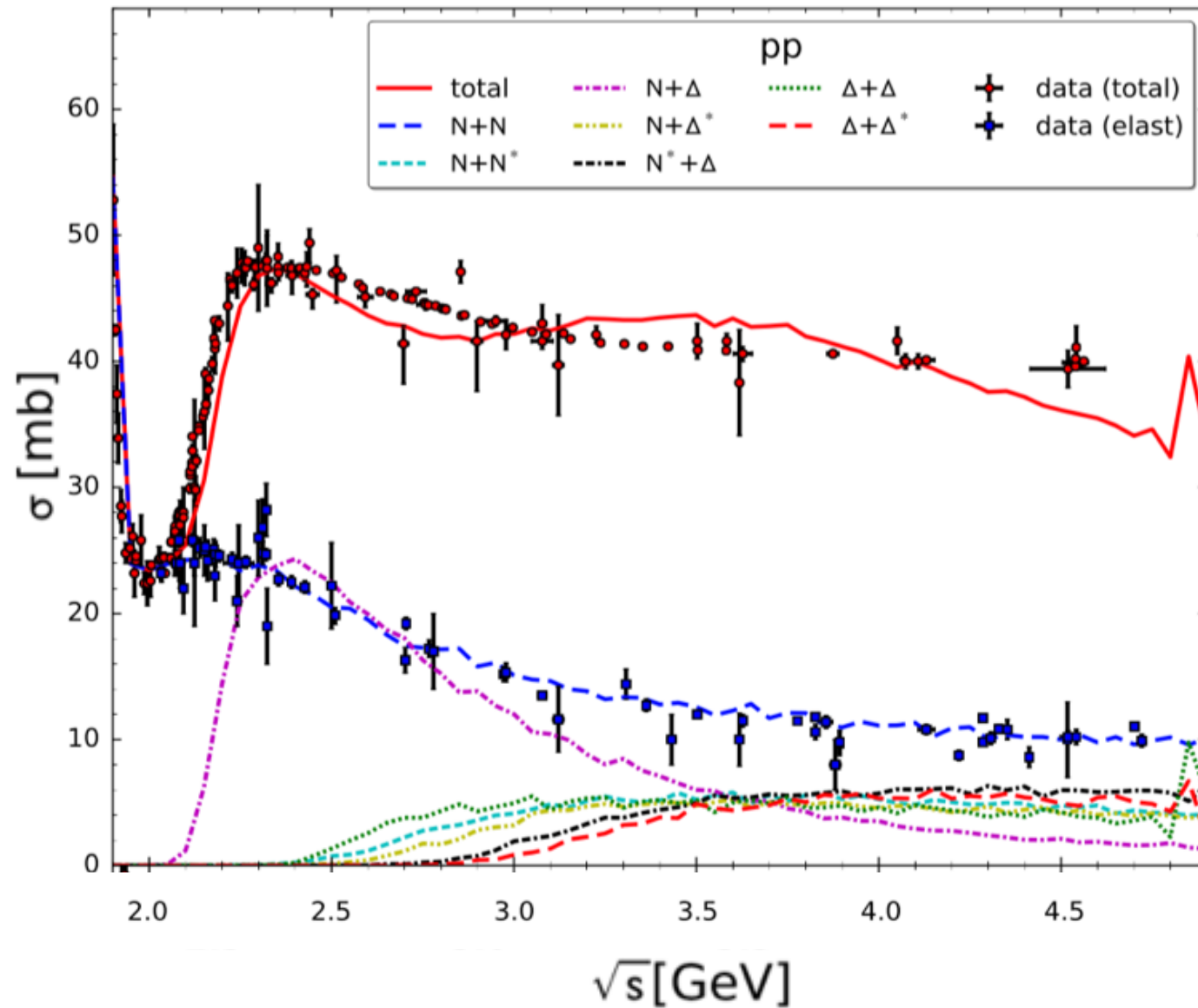
Degrees of Freedom

Mesons	Baryons
π, η	$N + 16 N^*$ states
ρ, ω, ϕ	$\Delta + 7 \Delta^*$ states
f_2	$\Lambda + 7 \Lambda^*$ states
σ	$\Sigma + 16 \Sigma^*$ states
$K, K^*(892), K^*(1410)$	Ξ, Ω

+ antiparticles

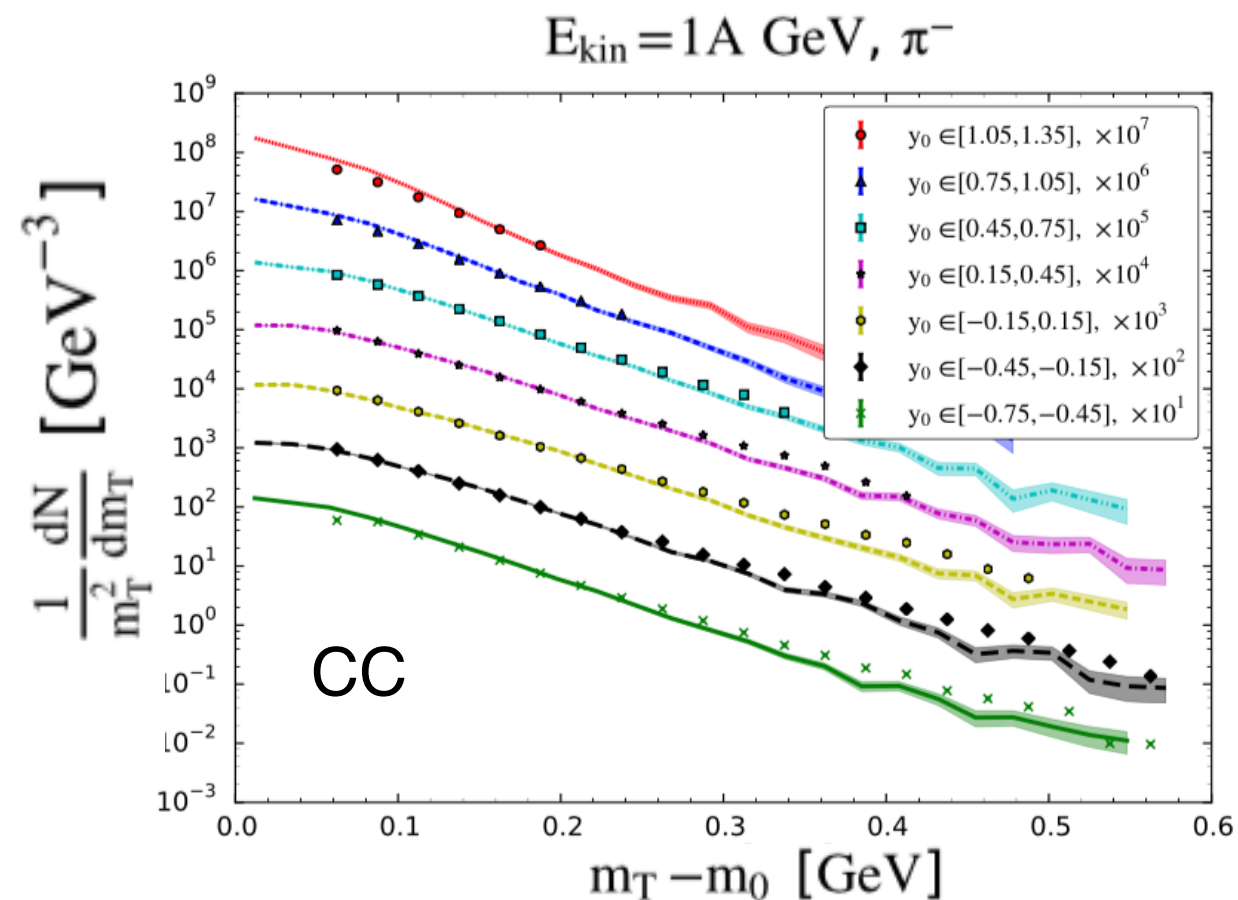
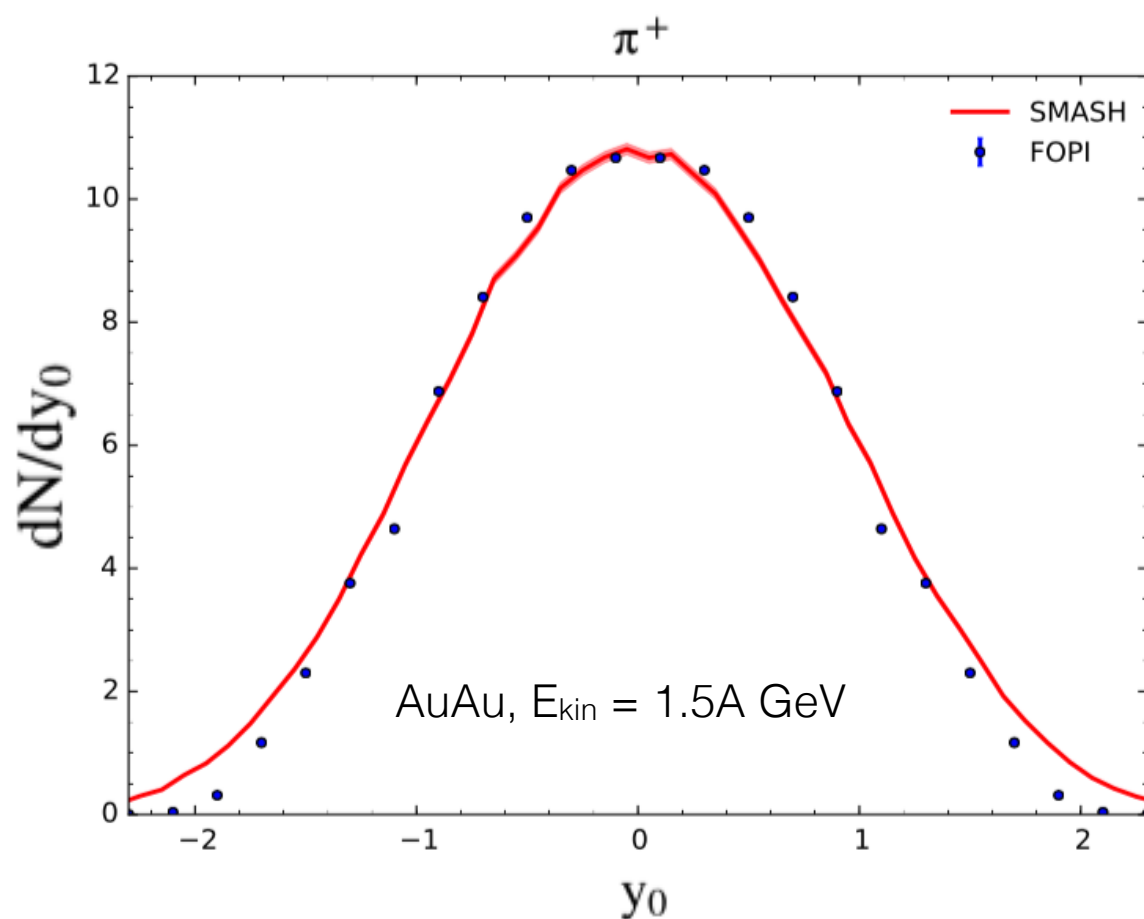
- all (well-established) hadrons from PDG up to a mass of 2 GeV
- isospin symmetry
- perturbative treatment of non-hadronic particles (photons, leptons)

Elementary Cross Sections



- total cross section for pp collisions
- parametrized elastic cross section
- many resonance contributions to inelastic cross section
- reasonable description of data up to 4 - 4.5 GeV

Results for HIC



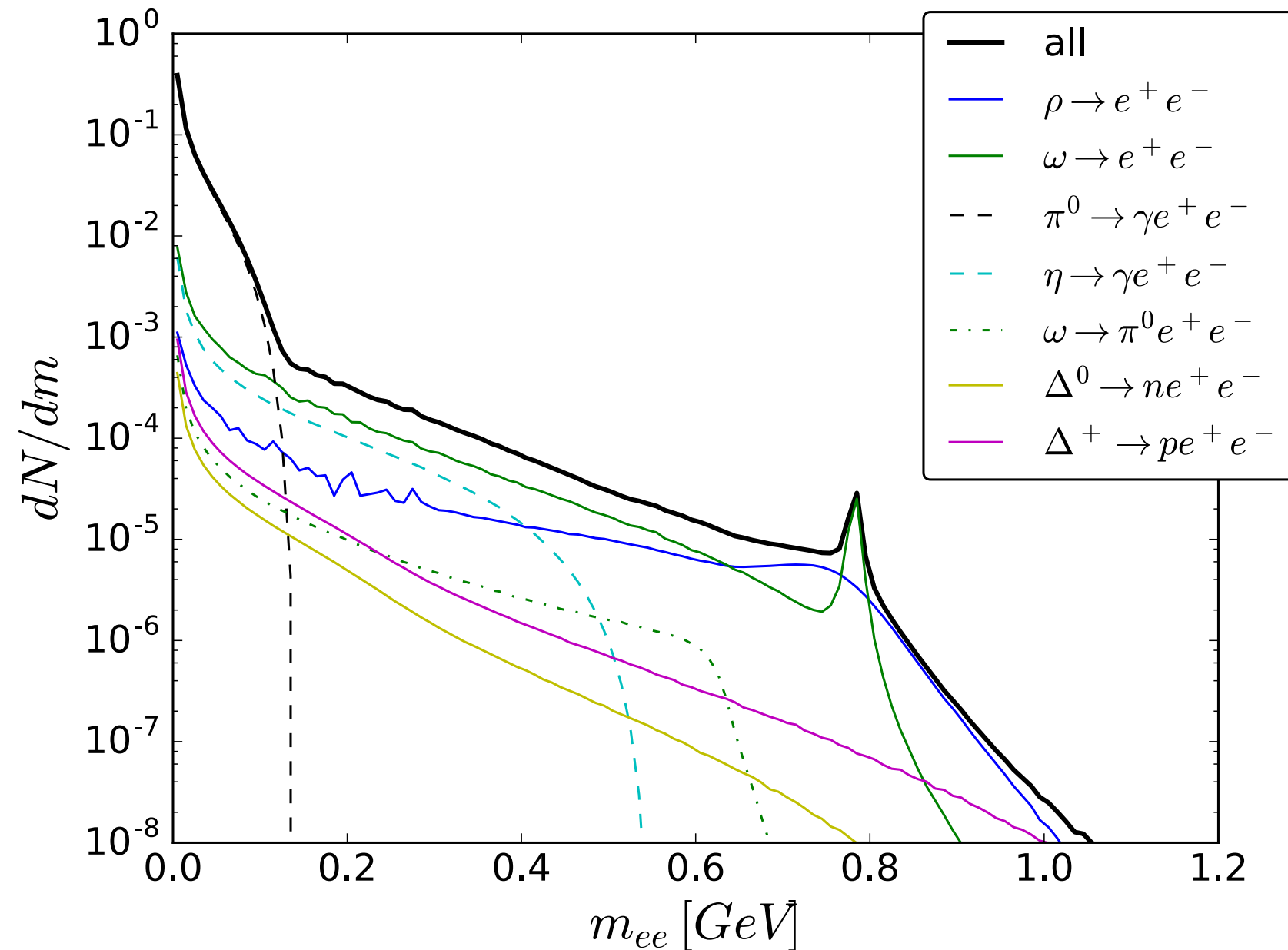
- shape of pion rapidity spectra reproduced
- transverse mass spectra compared with HADES data are in good agreement

Dileptons in SMASH

- direct and Dalitz dilepton decay channels
- rare e.m. decays \rightarrow Time-Integration-Method / *Shining* Phys.Lett. B259 (1991) 162-168
- continuously perform dilepton decays and weight them by taking their decay probability into account (better statistics)

Direct Decays
$\rho^0 \rightarrow l^+ l^-$
$\omega \rightarrow l^+ l^-$
$\phi \rightarrow l^+ l^-$
Dalitz Decays
$\pi^0 \rightarrow e^+ e^- \gamma$
$\eta \rightarrow e^+ e^- \gamma$
$\omega \rightarrow e^+ e^- \pi^0$
$\Delta^+ \rightarrow e^+ e^- p$
$\Delta^0 \rightarrow e^+ e^- n$

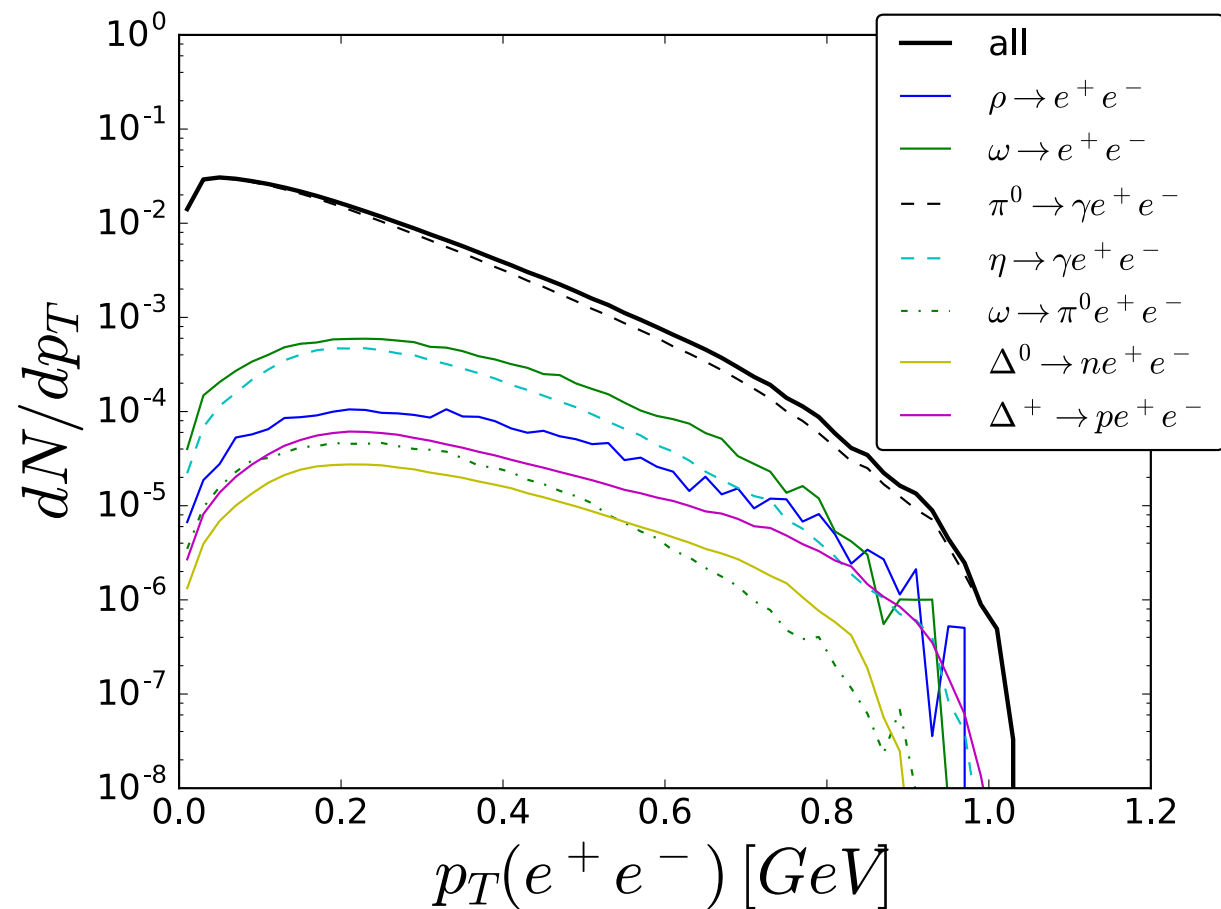
Dilepton Spectrum



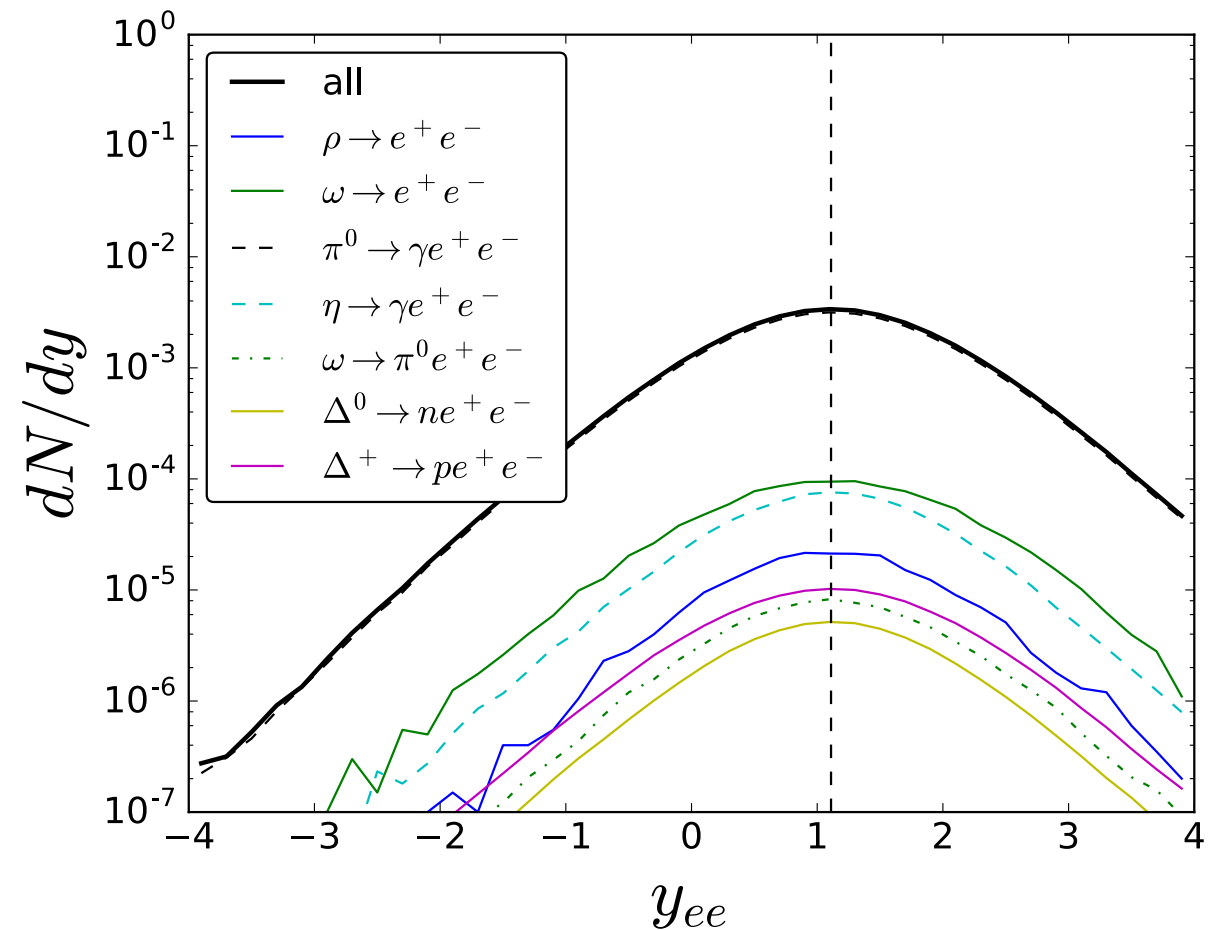
- invariant mass spectrum for pp collision with $E_{\text{kin}} = 3.5$ GeV

p_T and y Spectra

pp, 3.5 GeV



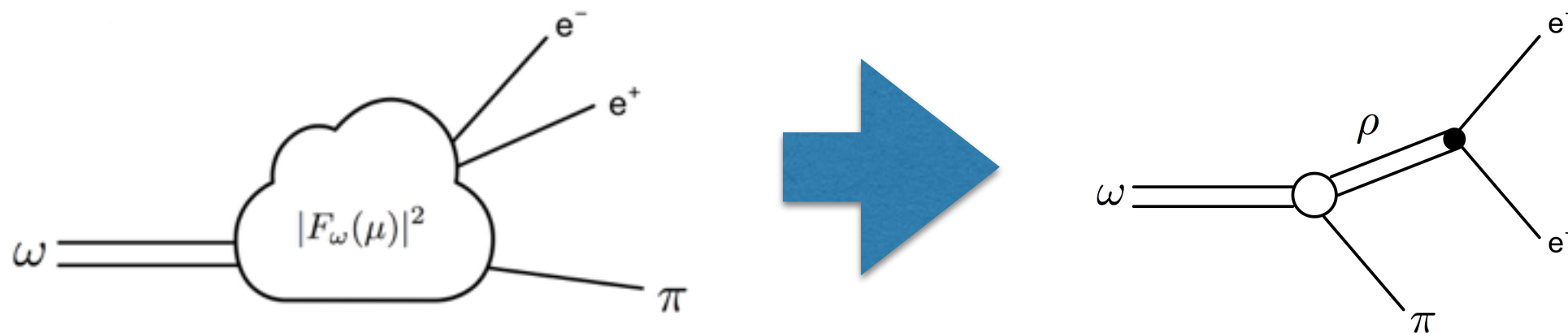
- dominant contribution from pion decays



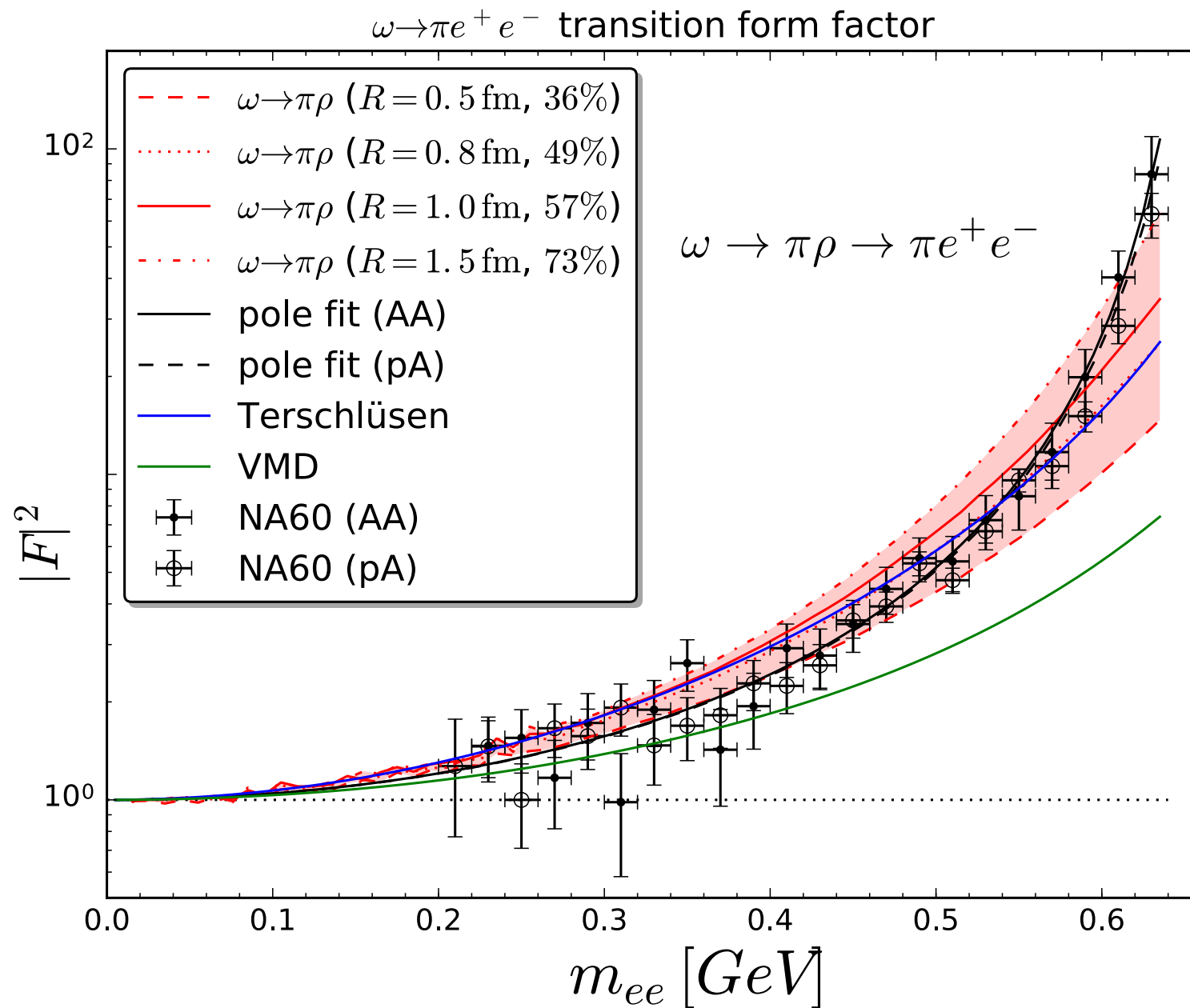
- rapidity spectrum shifted due to fixed-target setup

Form Factors

- dilepton routine was used to investigate the e.m. transition form factor of the omega dilepton Dalitz decay $\omega \rightarrow \pi^0 e^+ e^-$ *J. Weil, JS, Hannah Petersen, arXiv:1604.07028v1*
- decay described by two steps similar to hadronic ω decay into 3π : $\omega \rightarrow \rho\pi \rightarrow 3\pi$ (VMD-inspired)



Form Factors



- comparison of our ansatz with NA60 data

Phys.Lett. B677 (2009) 260-266

- variation of B.R. and cutoff-parameter R

- description of data with B.R. of 57% possible

$$B.R.(\omega \rightarrow 3\pi) = 89\%$$

- also compatible with Terschlüsen approach

Terschlüsen, Strandberg, Leupold and Eichstädt, Eur. Phys. J. A 49 116

Hybrid Approach

in coll. with G. Vujanovic and U. Heinz

Motivation

- study the behaviour of e.m. radiation in hydrodynamics relative to transport models

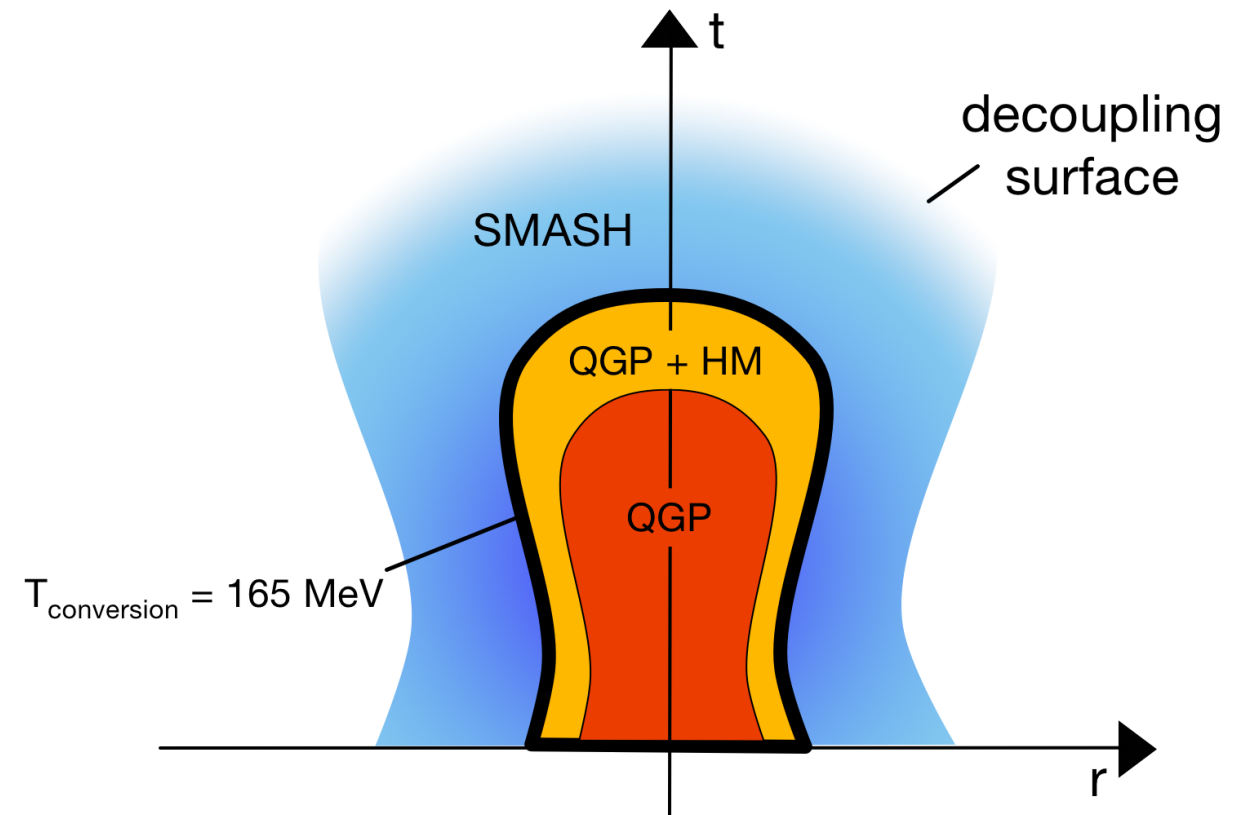
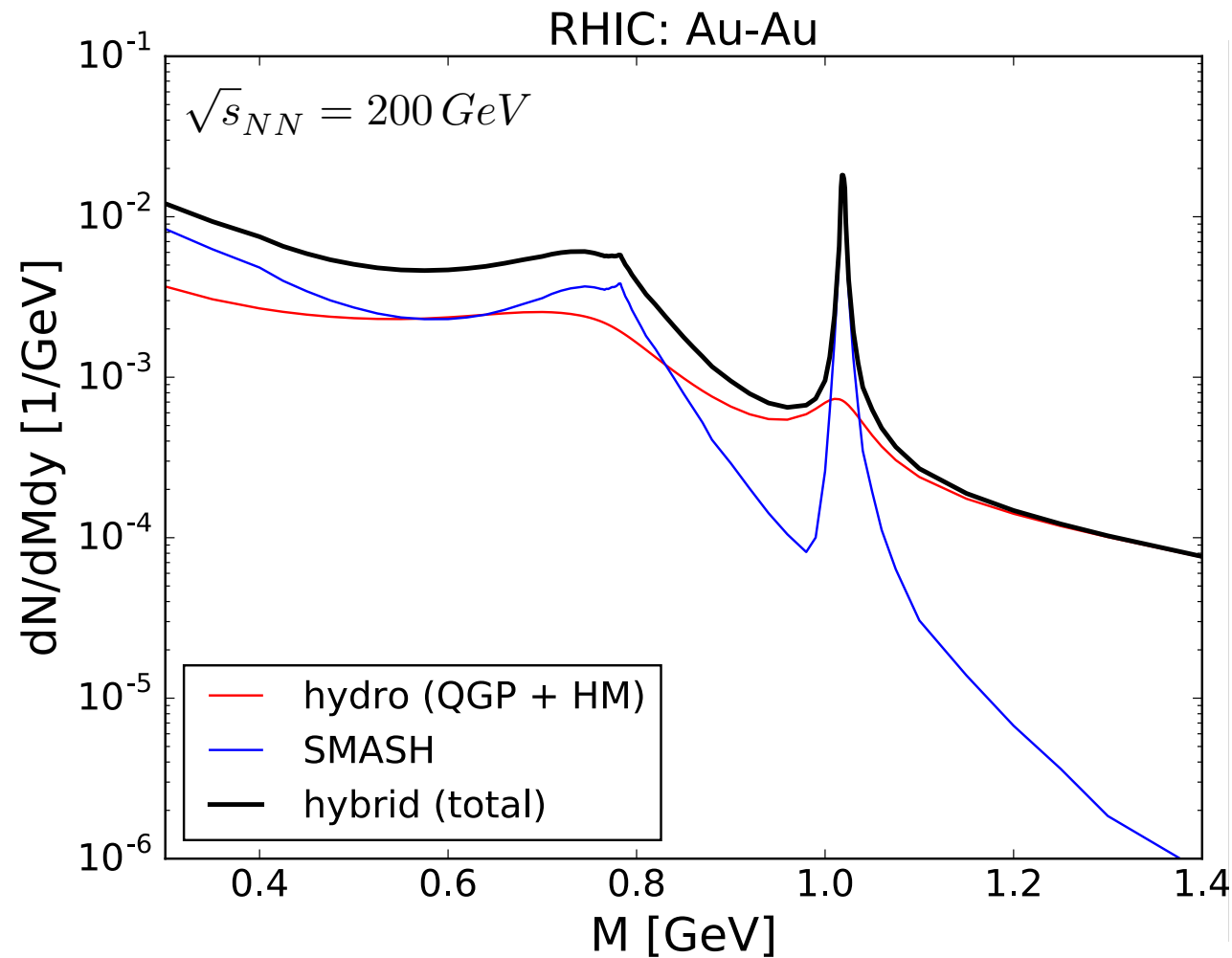
Model

- initialize (*iSS*) transport model (*SMASH*) on the chemical freeze-out hypersurface of hydrodynamical model (*MUSIC*)

arXiv:1409.8164

- combine dilepton radiation from hydro model with dilepton emission from transport model

First Result



- preliminary result
- MUSIC setup: IP-Glasma initial conditions, followed by visc. hydro (incl. Bulk and Shear viscosity), $T_{\text{conversion}} = 165 \text{ MeV}$
 - PRC 93, 044906 (2016)*
 - PRL 115, 132301 (2015)*
- thermal dilepton emission from hydro by G. Vujanovic (preliminary)
 - PRC 94, 014904 (2016)*

Summary + Outlook

Summary

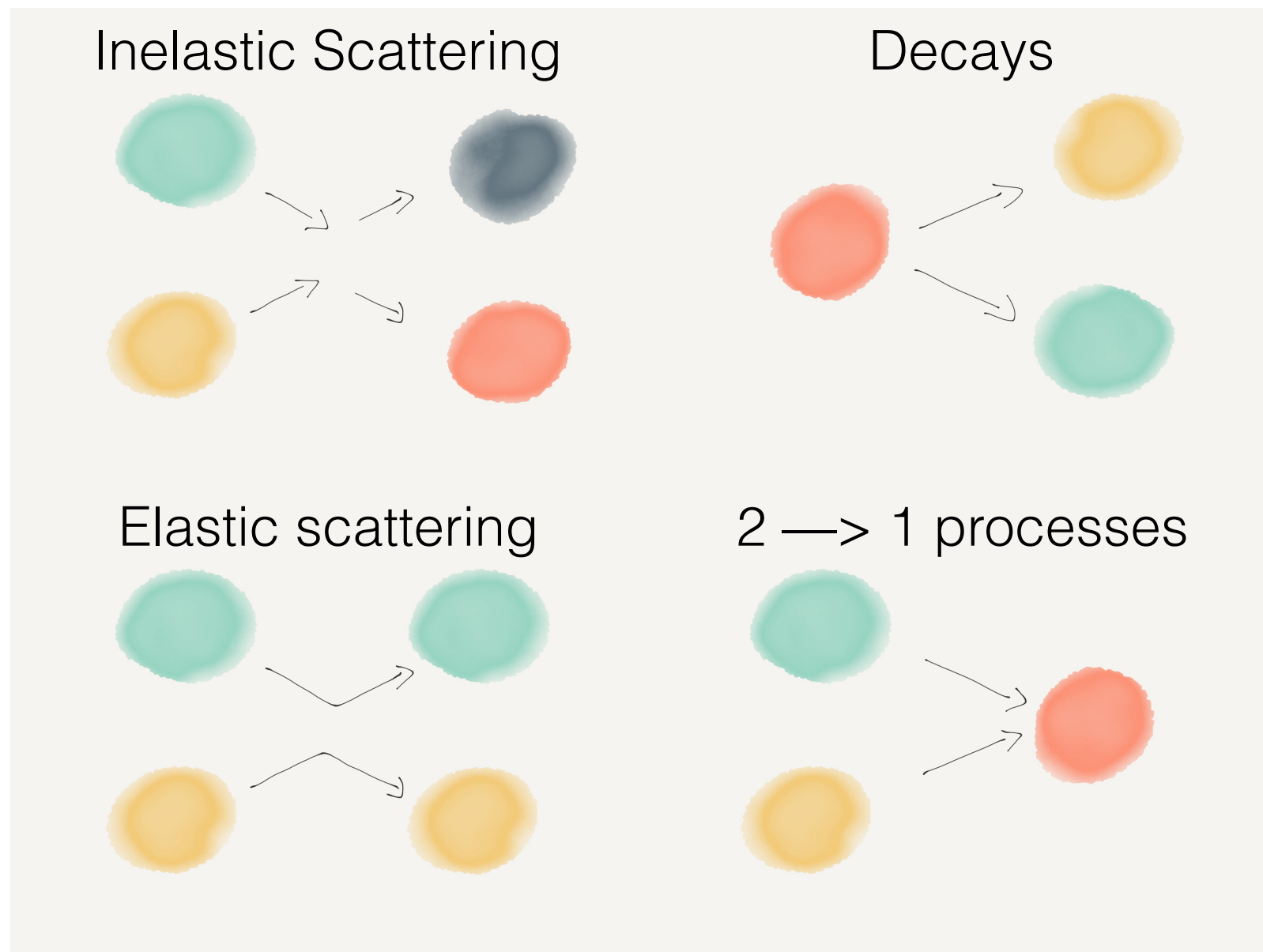
- transport models, like SMASH, are important to describe the late dilute stages of HIC
- dileptons can be used to study the properties of resonances and the e.m. transition form factor of the ω Dalitz decay
- dilepton production from the hybrid approach yields first results

Outlook

- compare dilepton production in SMASH to HADES data for low energy collisions
- conduct a systematic study with the hybrid model in order to compare with the dilepton radiation from a pure hydrodynamical approach
- SMASH: String fragmentation, Photons, ...

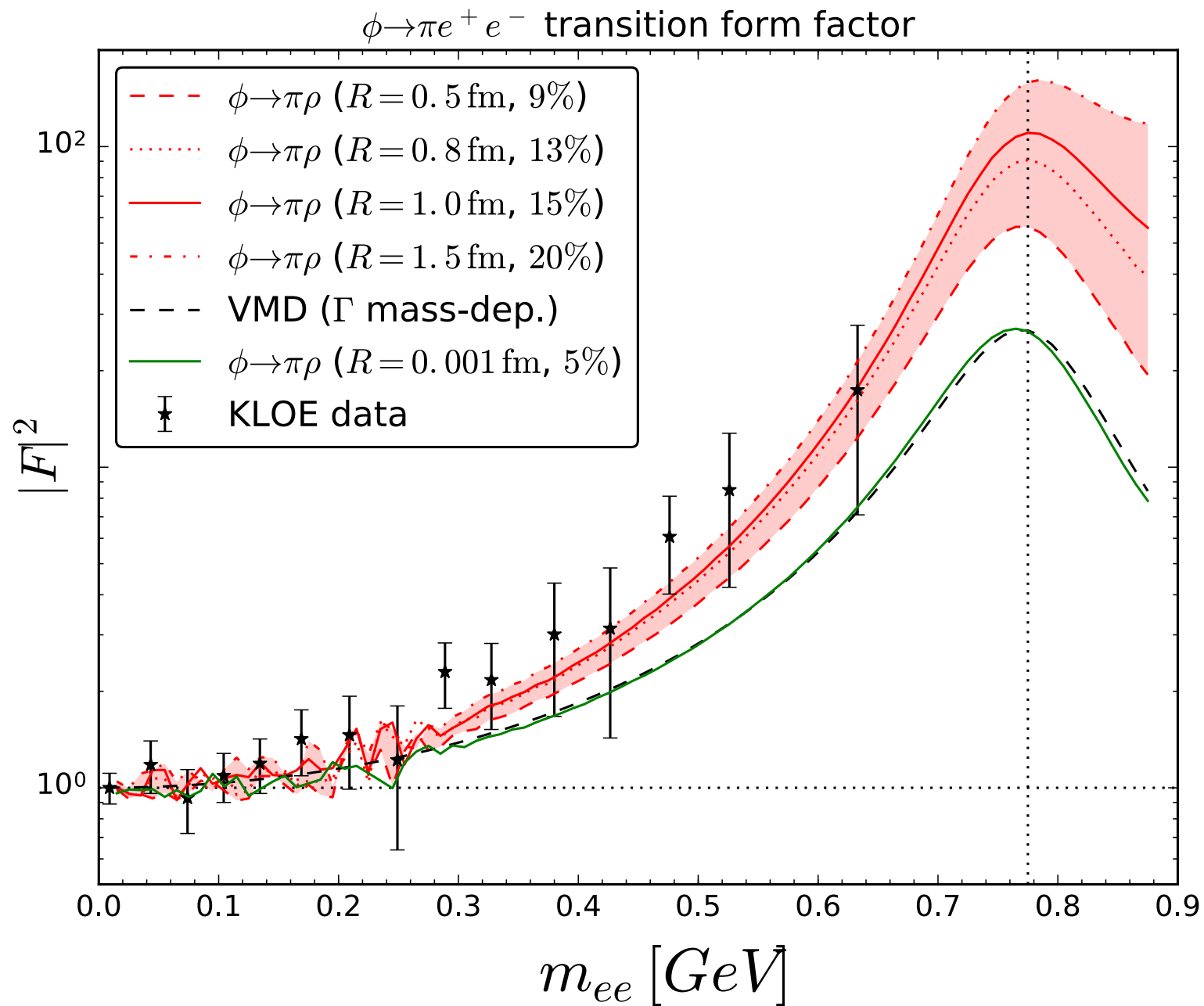
Backup Slides

Collision Term



- in few GeV energy regime decay and excitation of resonances dominate hadronic cross section
- no string fragmentation yet

Form Factors



- same ansatz for the Phi meson

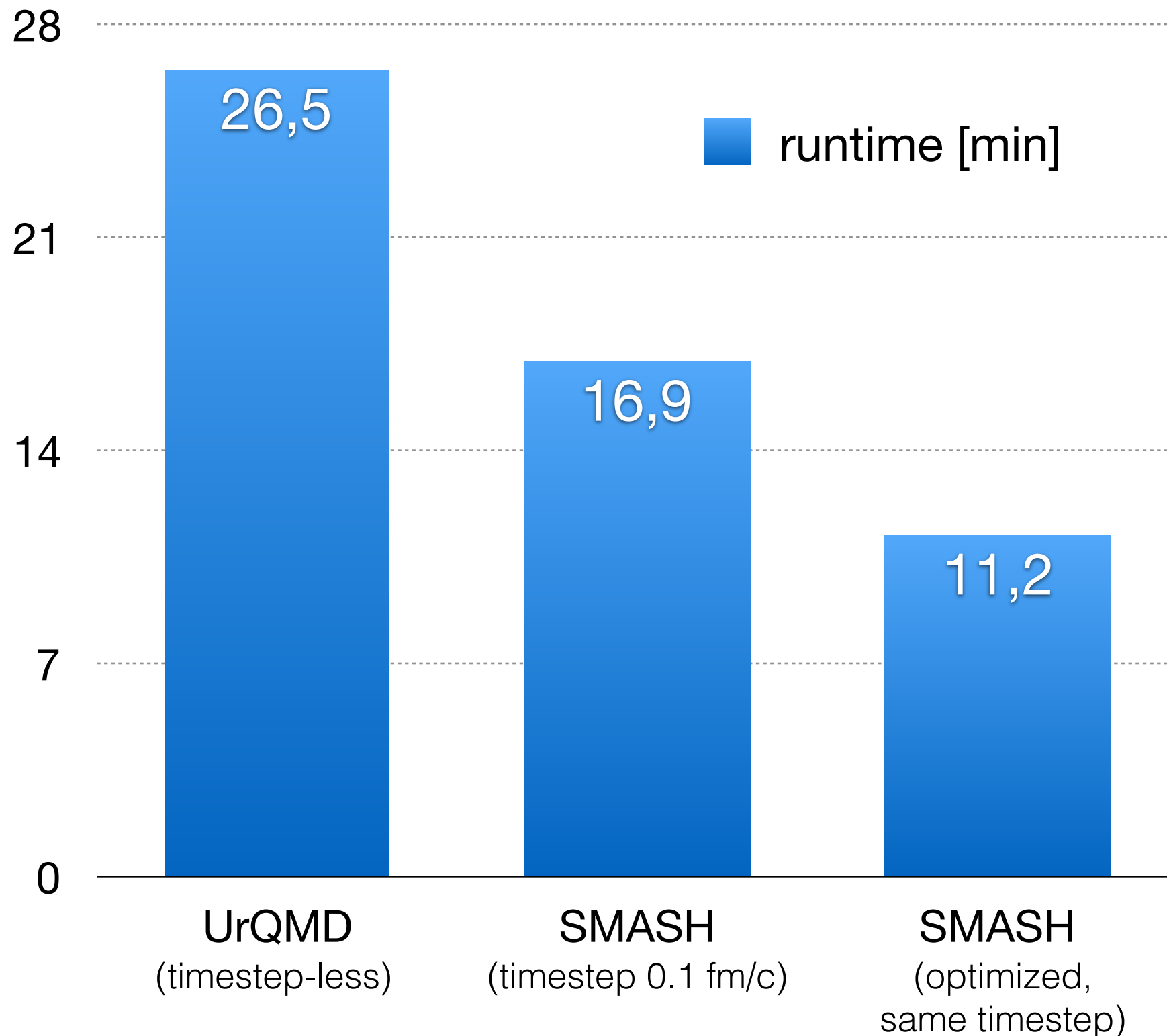
- comparison with data from KLOE

Phys. Lett. B 757 (2016) 362

- description of data suggests B.R. of 15% $B.R.(\phi \rightarrow 3\pi) = 15\%$

- again better description as „simple“ VMD

Performance compared to UrQMD

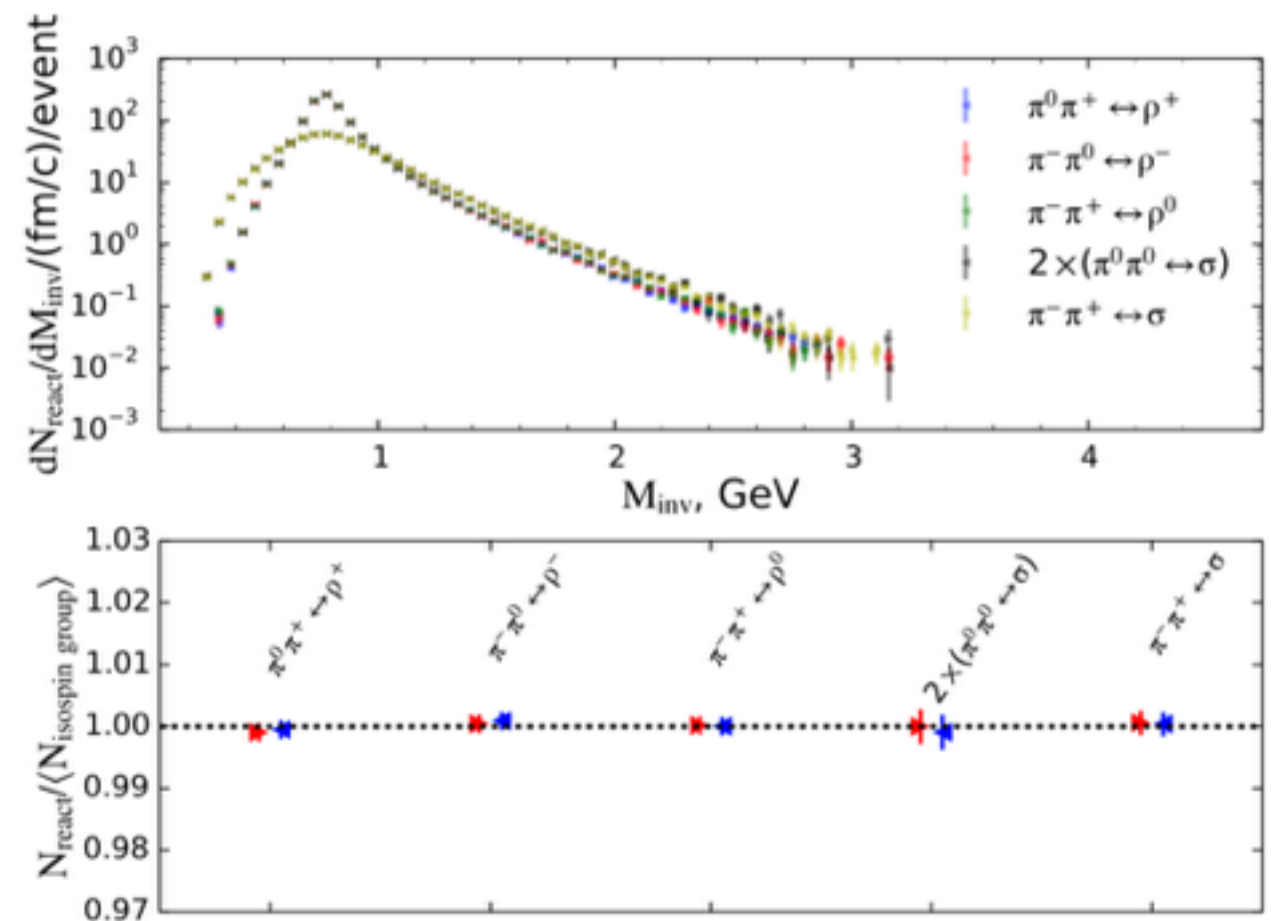
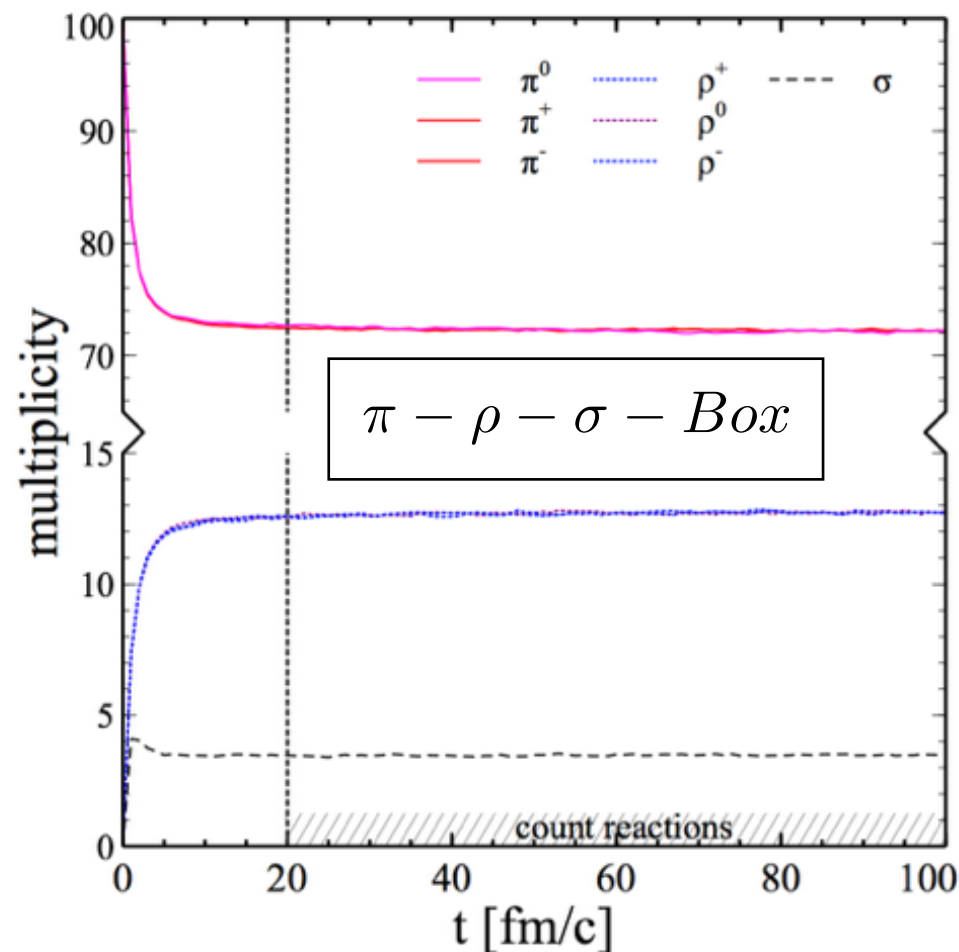


Setup	
System	central AuAu collision
Kinetic Energy	3 AGeV
N_{events}	1000
Machine	FIAS office desktop
End Time	20 fm/c
Output	final particles

—> more than 35% faster with the more expensive timestep-based evolution (more than 55% with optimizations)

Detailed Balance

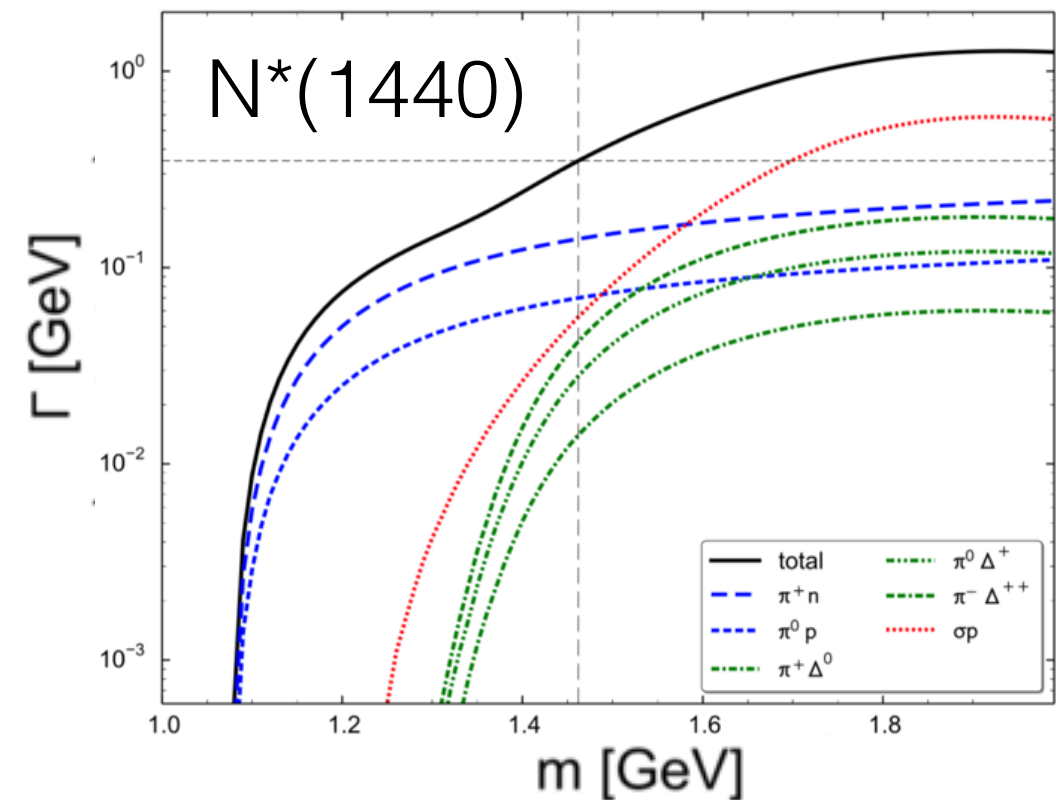
- inverse absorption cross section calculated from production cross section
- conservation of detailed balance (only $1 \leftrightarrow 2$ or $2 \leftrightarrow 2$ processes)



Resonances

- Spectral Function
 - all unstable particles („resonances“) have relativistic Breit-Wigner spectral functions
- Decay Widths
 - particles stable, if width $< 10 \text{ keV}$ (π, η, K, \dots)
 - treatment of Manley et al

$$\mathcal{A}(m) = \frac{2\mathcal{N}}{\pi} \frac{m^2 \Gamma(m)}{(m^2 - M_0^2)^2 + m^2 \Gamma(m)^2}$$



$$\Gamma_{R \rightarrow ab} = \Gamma_{R \rightarrow ab}^0 \frac{\rho_{ab}(m)}{\rho_{ab}(M_0)}$$

*D. M. Manley and E. M. Saleski,
Phys. Rev. D 45, 4002 (1992)*

Treatment of Manley

D. M. Manley and E. M. Saleski, Phys. Rev. D 45, 4002 (1992)

- scaling of on-shell decay width:

$$\Gamma_{R \rightarrow ab} = \Gamma_{R \rightarrow ab}^0 \frac{\rho_{ab}(m)}{\rho_{ab}(M_0)}$$

- definition of rho-function:

$$\rho_{ab}(m) = \int dm_a dm_b \mathcal{A}_a(m_a) \mathcal{A}_b(m_b) \times \frac{|\vec{p}_f|}{m} B_L^2(|\vec{p}_f|R) \mathcal{F}_{ab}^2(m)$$

Blatt Weisskopf functions

$$B_0^2 = 1$$

$$B_1^2(x) = x^2 / (1 + x^2)$$

...

- hadronic Form Factor:

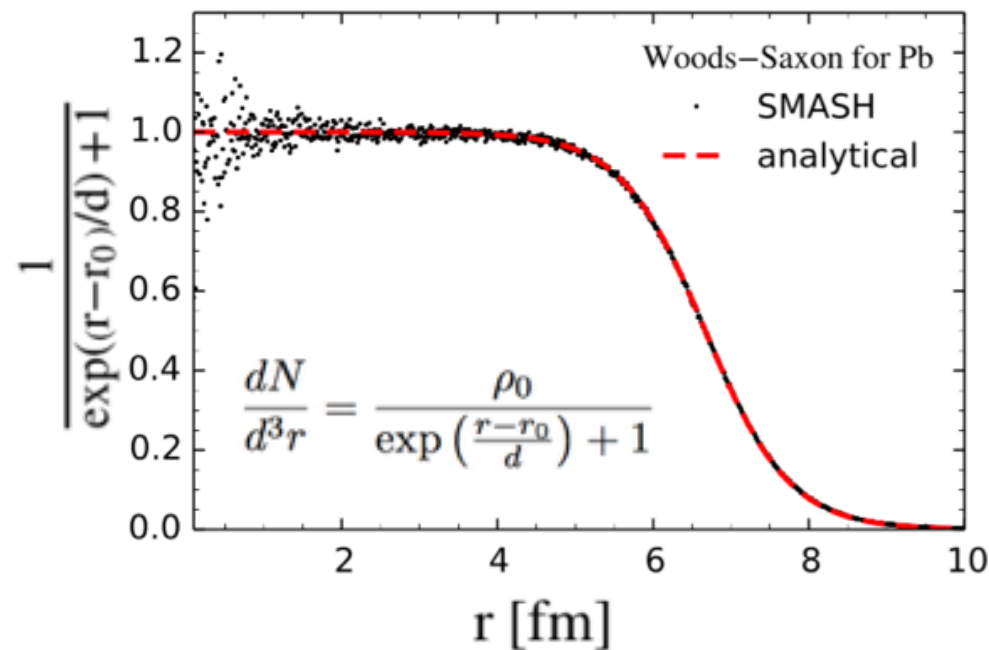
$$\mathcal{F}_{ab}(m) = \frac{\lambda^4 + 1/4(s_0 - M_0^2)^2}{\lambda^4 + (m^2 - 1/2(s_0 + M_0^2))^2}$$

decay	λ [GeV]
$\pi\rho$	0.8
unstable mesons (e.g. $\rho N, \sigma N$)	1.6
unstable baryons (e.g. $\pi\Delta$)	2.0
two unstable daughters (e.g. $\rho\rho$)	0.6

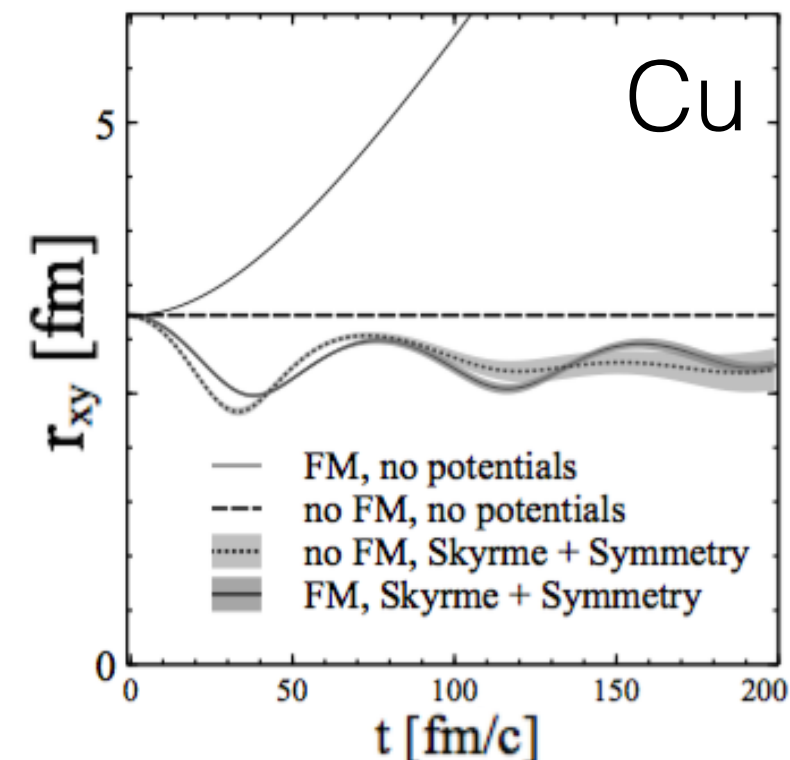
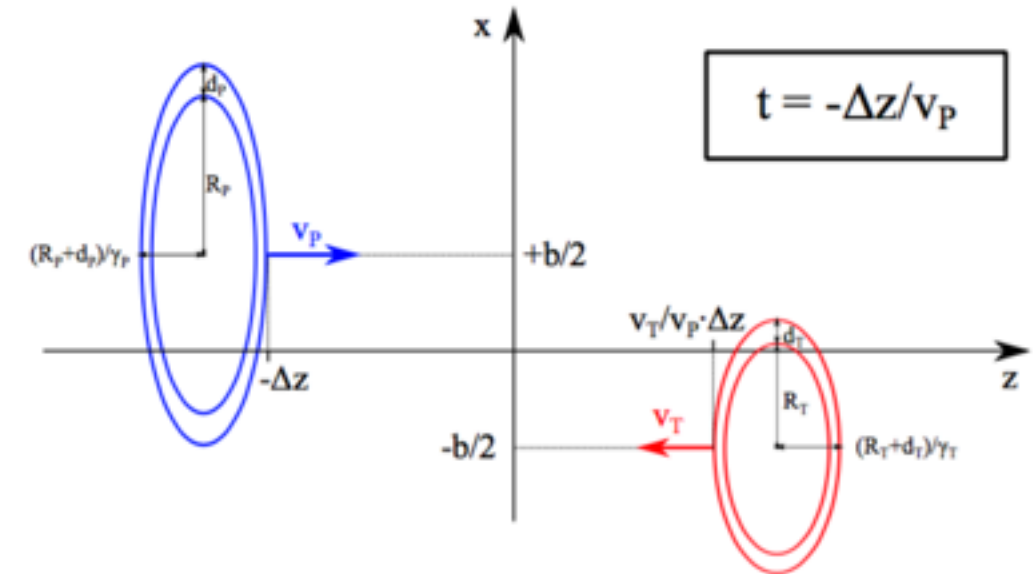
M. Post, S. Leupold, U. Mosel, Nucl. Phys. A 741, 81 (2004)

Initial Conditions

- Nuclear Collisions
 - Woods-Saxon distribution in coordinate space



- *optional*: deformed nuclei
- *optional*: Fermi Motion



Fermi Motion

- momentum distribution of nucleons in nucleus = uniform distribution with $0 \leq |\vec{p}| \leq p_F(\vec{r})$ („Fermi sphere“)
- radius depends on density of nucleons:

$$p_F(\vec{r}) = \hbar c (3\pi^2 \rho(\vec{r}))^{1/3}$$

Pauli Blocking

- probability that action is blocked

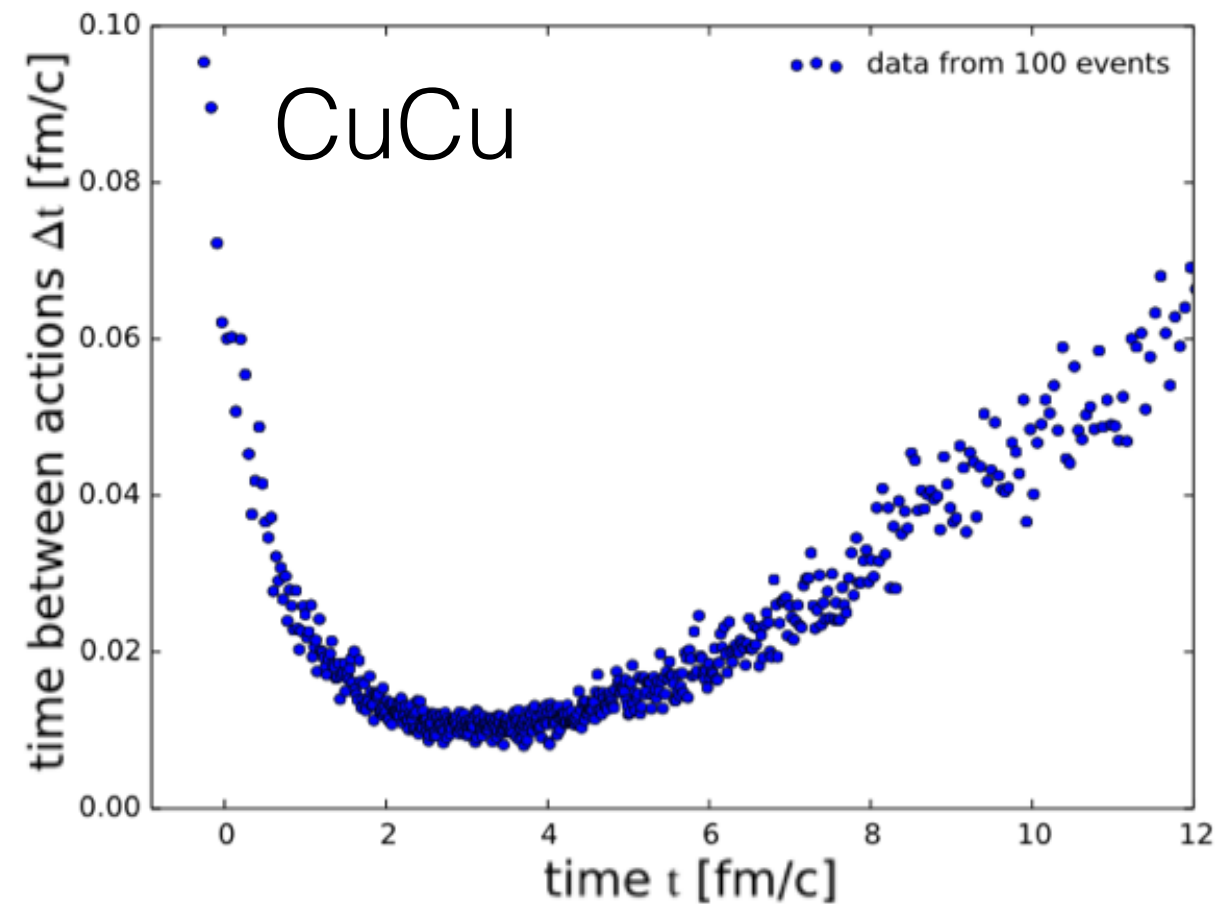
$$1 - \prod_i (1 - f_i)$$

i: final states of reaction
f_i: phase space density of species i

- „derivation“ from Boltzmann-Uehling-Uhlenbeck-equation

General Setup

- Time Steps
 - time steps either fixed or adaptive
 - one interaction per timestep per particle, afterwards propagation
 - also possible: timestep-less mode
- Mean-field potentials



$$U = a(\rho/\rho_0) + b(\rho/\rho_0)^\tau \pm 2S_{\text{pot}} \frac{\rho_{I3}}{\rho_0}$$

Xu et al, arXiv:1603.08149