

Cross Section Measurement of $pp \rightarrow Z(e^+e^-) + X$ with First Data from the ATLAS Experiment

Motivation • Cross Section Determination • Prospects

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Why studying the cross section for the Z ?

Clean and fully reconstructed leptonic final states of the Z decay provides “standard candle” for detector calibration

From the physics point of view

- ▶ quite early sufficient statistics for $d\sigma/dp_T$ and $d\sigma/d\eta$ measurements
- ▶ $d\sigma/dp_T$ spectrum provides constraints on QCD, $d\sigma/d\eta$ is direct probe of the PDFs
- ▶ advantage of $Z \rightarrow ee$ over $Z \rightarrow \mu\mu$ is larger η acceptance

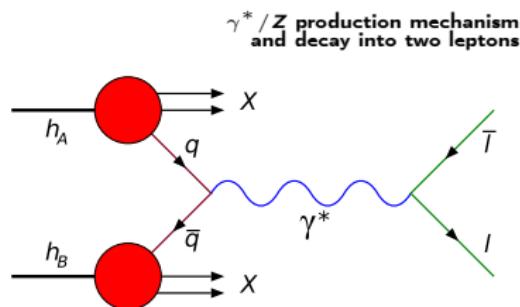
“Master formula” for calculation:

$$\sigma = \frac{N-B}{A \cdot \varepsilon \cdot \mathcal{L}}$$

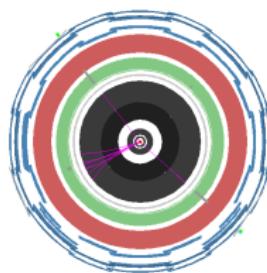
Ingredients:

- | | |
|--|-----------------------|
| ▶ N | counted events |
| ▶ B | background events |
| ▶ A | acceptance |
| ▶ \mathcal{L} | integrated luminosity |
| ▶ $\varepsilon = \varepsilon_{\text{trigger}} \cdot \varepsilon_{\text{reco}} \cdot \varepsilon_{\text{ID}}$ | efficiencies |
| ▶ $\varepsilon_{\text{trigger}}$ | trigger efficiency |
| ▶ $\varepsilon_{\text{reco}}$ | reconstruction |
| ▶ ε_{ID} | identification |

Simulations used for this Analysis



ATLAS event display



- ▶ center-of-mass energy 14 TeV
- ▶ simulated with PYTHIA 6.411
- ▶ PDFs CTEQ6L

$Z \rightarrow e^+e^-$ signal simulation:

- ▶ limit on hard process $\sqrt{s} > 60$ GeV
- ▶ $\sigma(pp \rightarrow Z/\gamma^*(e^+e^-) + X)$

2.015 nb (NNLO)

- ▶ full detector simulation GEANT4

background (QCD) samples:

- ▶ p_T^{jet} > 15 GeV
- ▶ $\sigma(pp \rightarrow \text{Dijet})$

$2.333 \cdot 10^6 \text{ nb}$

Presentation is based on an integrated luminosity of 50 pb^{-1} , if not stated otherwise.

How to find signal events

$$\sigma = \frac{N-B}{A \cdot \varepsilon \cdot \mathcal{L}}$$

Event selection, $L = 10^{31} \text{ cm}^{-2}\text{s}^{-1}$

- 1 trigger threshold

$$p_T > 10 \text{ GeV}$$

- 2 two electromagnetic clusters

$$E_T > 15 \text{ GeV}$$

$$|\eta| < 2.4$$

$$80 \text{ GeV} < M_{\text{cl,cl}} < 100 \text{ GeV}$$

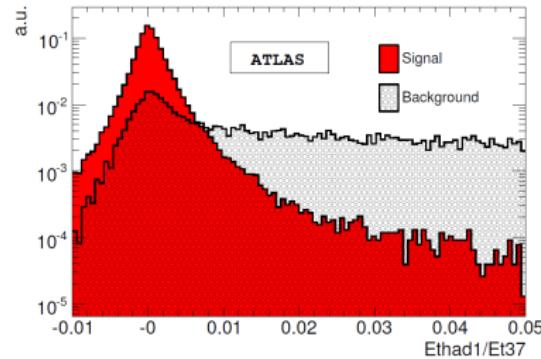
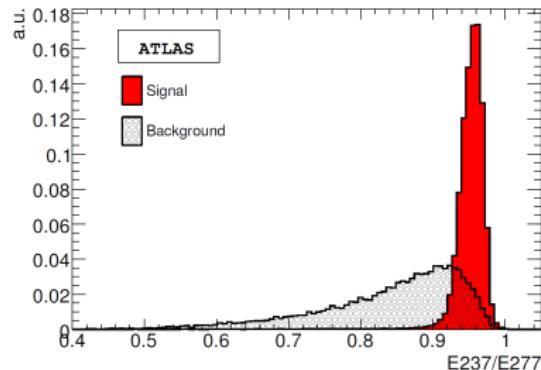
- 3 "loose" (i.e. calorimeter only) electron identification criteria

hadronic leakage

shower shapes

- 4 isolation in ΔR cone

$$\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} < 0.45$$



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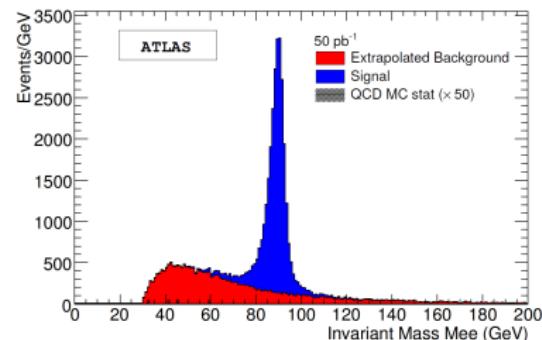
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Expected number of $Z \rightarrow e^+e^-$ -signal and background events ($\times 10^4$) for $\mathcal{L} = 50 \text{ pb}^{-1}$ after the successive cuts

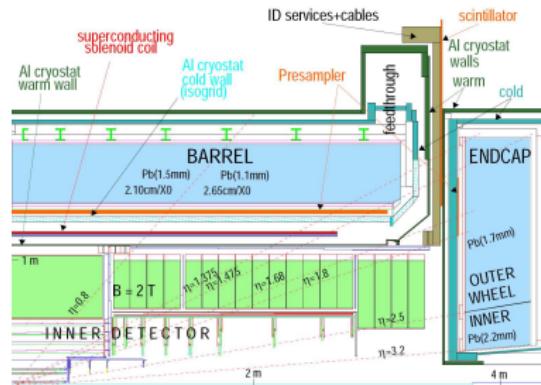
1	6.70 ± 0.01	3110 ± 40
2	2.76 ± 0.01	11.1 ± 0.8
3	2.64 ± 0.01	0.8 ± 0.2
4	2.48 ± 0.01	0.2 ± 0.1

Acceptance studies

The acceptance is defined as

- fraction of the signal that passes kinematic and geometric cuts

$$\mathcal{A} = \frac{N_{\text{after kinematical and geometrical cuts}}}{N_{\text{before cuts}}}$$



Uncertainties w.r.t. physics

- limited knowledge of underlying physics
non-perturbative mechanisms

PDFs

initial state radiation

intrinsic k_T of the incoming partons

Uncertainties $\delta\mathcal{A}$ were estimated by switching on and off the different effects, and a variety of generators have been used, namely

PYTHIA

Herwig

MC@NLO

- combined result: $\delta\mathcal{A}/\mathcal{A} = 2.3\%$

Intermezzo: The “tag and probe” method

$$\sigma = \frac{N-B}{A \cdot \varepsilon \cdot \mathcal{L}}$$

Data-driven approach to get efficiencies

- ▶ motivation for data-driven technique is that one don't want to rely on Monte Carlo for efficiency studies
- ▶ as real data won't contain “true information”, methods need to be approved

Common use of the “tag and probe” method

- ▶ for $Z \rightarrow ee$, use hard cuts to find one “good” lepton as reference (“tag”)
- ▶ look for a cluster, such that

$$m_{\min} < M_{\text{tag,cluster}} < m_{\max}$$

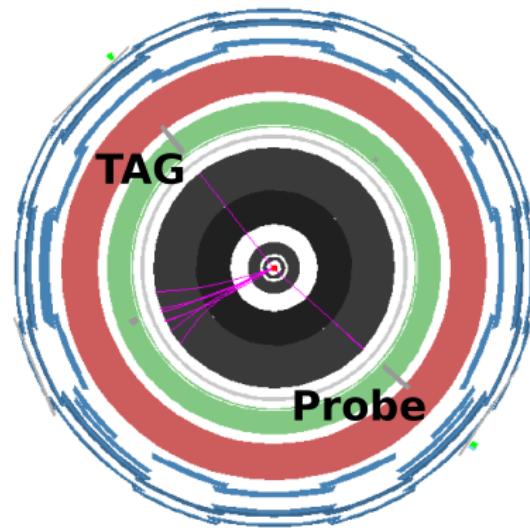
$$\Rightarrow N_1$$

- ▶ use this second cluster as “probe” to get a sub-sample when additional cuts are passed

$$\Rightarrow N_2$$

- ▶ “tag and probe” efficiency is then defined as

$$\varepsilon_{\text{tag and probe}} = \frac{N_2}{N_1}$$



Reconstruction and identification efficiencies

$$\sigma = \frac{N-B}{A \cdot \varepsilon \cdot \mathcal{L}}$$

Technique for efficiency determination

- ▶ based on seed-clusters in the EM calorimeter, matched to tracks
- ▶ conditions on the "tag"

reconstructed electron

"tight" isolation cuts

outside $1.37 < |\eta| < 1.52$

pass trigger threshold

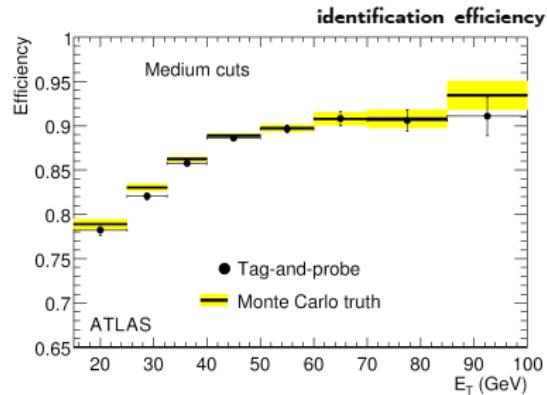
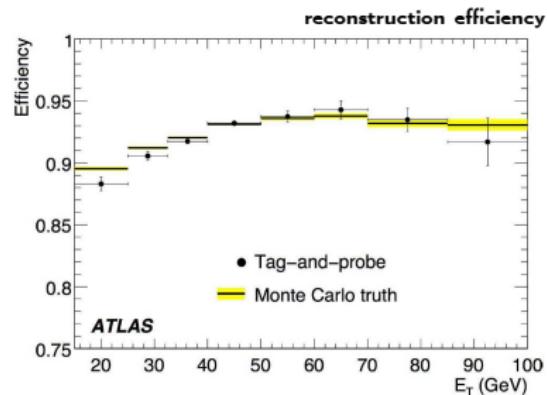
$E_T > 15$ GeV

- ▶ conditions on the "probe"

$\Delta\Phi_{\text{tag, probe}} > 3/4\pi$

$E_T > 15$ GeV

$80 \text{ GeV} < M_{\text{tag, probe}} < 100 \text{ GeV}$



Tag and probe method shows good agreement with MC truth

Trigger efficiencies

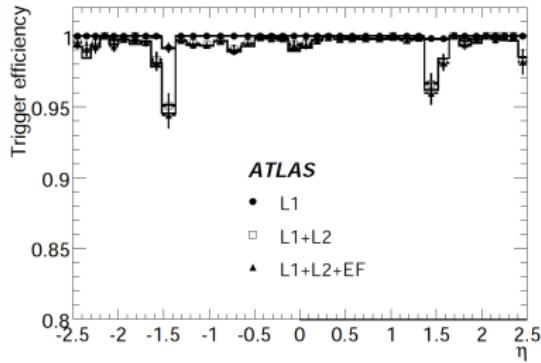
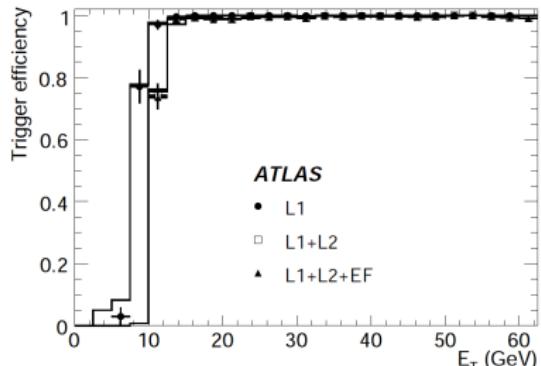
$$\sigma = \frac{N-B}{A \cdot \varepsilon \cdot \mathcal{L}}$$

Early running at ATLAS

- ▶ vital for understanding of
 - detector performance
 - offline reconstruction
 - trigger performance
- ▶ low luminosity run allows
 - loose selections
 - low trigger thresholds

For early data, 10 GeV trigger threshold

- ▶ shown here: single electron trigger efficiency vs. E_T and $|\eta|$
- ▶ agreement better than 1% between tag and probe method and MC truth
- ▶ high efficiency



Luminosity measurement at ATLAS

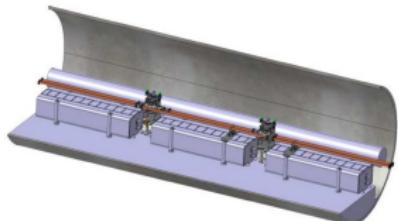
$$\sigma = \frac{N-B}{A \cdot \varepsilon \cdot \mathcal{L}}$$

Luminosity parameters at LHC-startup

- ▶ instantaneous luminosity

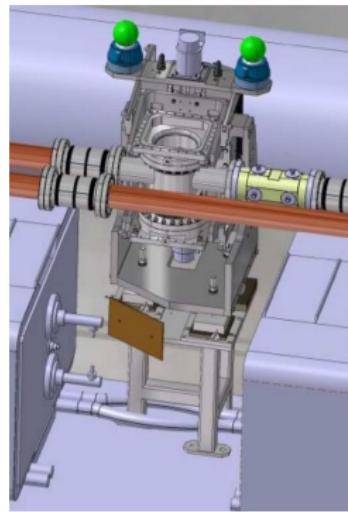
$$L = 10^{31} \text{ cm}^{-2}\text{s}^{-1}, \frac{\delta L}{L} \approx 10\%$$

- ▶ L "measurement" based on machine parameters



Measurement of the luminosity via "Roman pots"

- ▶ installed on both sides of the IP
- ▶ measuring elastic scattered protons
- ▶ $\mathcal{O}(1 \text{ mm})$ distance from beam
- ▶ measurement done via dedicated runs, i.e. transfer to normal run conditions necessary



Aim is to

- ▶ access "Coulomb Interference Region"
- ▶ reach $\frac{\delta L}{L} \approx 2-3\%$

Results for the total cross section

With this toy-study for 50 pb^{-1} at a startup-luminosity of $10^{31} \text{ cm}^{-2}\text{s}^{-1}$, the following results were obtained:

► N	$(2.71 \pm 0.02) \times 10^4$
► B	$(0.23 \pm 0.04) \times 10^4$
► $\mathcal{A} \times \varepsilon$	0.246
► $\delta\mathcal{A}/\mathcal{A}$	0.023
► $\delta\varepsilon/\varepsilon$	0.030
⇒	$\sigma = (2016 \pm 16_{\text{stat}} \pm 83_{\text{syst}}) \text{ pb}$

Agreement with nominal total cross section ($\sigma = 2015 \text{ pb}$), result is dominated by systematic uncertainties

(Uncertainty on luminosity not included in these results!)

$$\sigma = \frac{N-B}{\mathcal{A} \cdot \varepsilon \cdot \mathcal{L}}$$

Extrapolation to $\mathcal{L} \rightarrow \mathcal{O}(1 \text{ fb}^{-1})$

- statistical uncertainty for N becomes negligible
- efficiency uncertainty decreases for part which is determined from measurement
- trigger threshold will be increased with “tight” selection criteria
- uncertainty on the acceptance will not scale with statistics and therefore will be dominating the result

$$\Rightarrow \sigma = (2016 \pm 4_{\text{stat}} \pm 49_{\text{syst}}) \text{ pb}$$

Differential Cross Section

$$\sigma_\alpha = \frac{S_\alpha}{\mathcal{L}} \frac{n_\alpha - b_\alpha}{\varepsilon_\alpha \cdot \mathcal{A}_\alpha}$$

Differential cross sections useful to put constraints on uncertainties of the total cross section measurement, at the cost of higher statistics needed

- ▶ S_α in formula for σ_α represents the detector smearing correction
- ▶ signal selection

12 GeV trigger threshold

$p_T > 20$ GeV

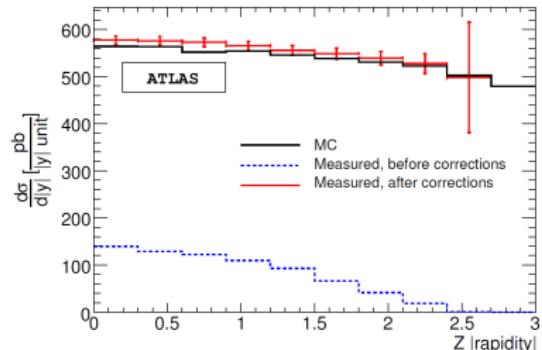
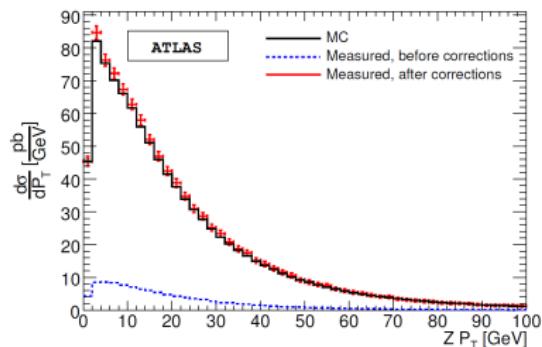
$|\eta| < 2.5$

- ▶ extraction of $d\sigma_Z/dp_T$ and $d\sigma_Z/dy$ in different regions α

50 p_T bins

9 y bins

- ▶ results show good agreement between reconstructed and true cross section



Summary and Outlook

Summary

- ▶ Motivation for this study was to test the tools for measuring the (differential) cross section and the expected uncertainty for this “standard candle”-process
- ▶ With the method presented, cross section measurement for $pp \rightarrow Z/\gamma^*(e^+ e^-) + X$ possible with an uncertainty (including the one on luminosity) of
 - ▶ $\leq 10\%$ for the first 50 pb^{-1}
 - ▶ $\leq 5\%$ extrapolated to $\mathcal{L} \approx 1 \text{ fb}^{-1}$
- ▶ Measurement of the differential cross section also feasible within statistics intended for the first LHC-run

Outlook

- ▶ Ultimately, the uncertainty on this measurement is driven by strong interaction effects, affecting the y and p_T distributions of the W and the Z
- ▶ In the end, the lower bound is fixed by the precision of luminosity and acceptance
- ▶ Without further improvement of these two parameters, measurement of the total cross section will be limited to $\delta\sigma/\sigma \gtrapprox 2\%$

waiting for data