Study of charged and neutral Higgs boson decays at CLIC

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Introduction

- In the Standard Model, only 1 doublet of Higgs scalars is responsible for the electroweak symmetry breaking. As a result, there is only one neutral Higgs boson h⁰.
- Other theoretical models, in particular with Supersymmetry (MSSM), predict the existence of 2 complex Higgs doublets \rightarrow 5 physical states: H^+ , H^- , h^0 , H^0 and A^0 .
- At tree level, the MSSM Higgs sector is fully determined by two independent parameters only: m_A and tan β .
- By comparing the signal rates of H[±] → tb and H[±] → τν, or of H⁰/A⁰ → tt and H⁰/A⁰ → bb, one can derive tan β, whether or not the charged and neutral Higgs bosons also decay into non-SM particles.

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Monte-Carlo simulation studies

- The charged and neutral Higgs boson decay widths and branching ratios are computed with HDECAY.
- Signal events are generated with PYTHIA and the CLIC beam-beam effects (beamstrahlung, ISR, γγ → hadrons) are included.
- The physics background events are generated with MadEvent/MadGraph. A home-made subroutine was written to include the CLIC beam-beam effects and PYTHIA is used for the fragmentation of the quarks.
- Fast detector simulation and event reconstruction with SIMDET, 70% tagging efficiency for *b* and τ jets.

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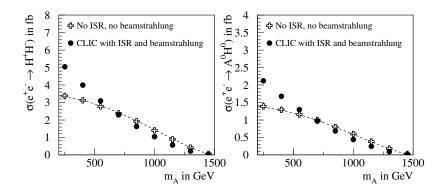
CLIC beam parameters at 3 TeV

Center-of-mass energy	3	TeV
Main linac RF frequency	30	GHz
Accelerating gradient	150	MV/m
Linac and site lengths	28/33.2	km
Linac repetition rate	150	Hz
No. of bunches per pulse	220	
No. of particles per bunch	2.56	10 ⁹
Bunch spacing	0.267	ns
Primary beam power	20.4	MW
Total site AC power	418	MW
Wall plug to main beam efficiency	12.5	%
Horizontal emittance $(\beta\gamma)\epsilon_x$	0.660	mm.mrad
Vertical emittance $(\beta\gamma)\epsilon_y$	0.001	mm.mrad
Horizontal beam size σ_x	60	nm
Vertical beam size σ_y	0.7	nm
Bunch length σ_z	30.8	μ m
Peak luminosity	6.5	10 ³⁴ cm ⁻² s ⁻¹
Luminosity within 1% of Ecm	3.3	10 ³⁴ cm ⁻² s ⁻¹
Photons per e^+/e^-	1.1	
Beamstrahlung loss	16.0	%
Coherent pairs per bunch crossing	5	10 ⁷
$\gamma\gamma ightarrow$ hadrons per bunch crossing	0.73	

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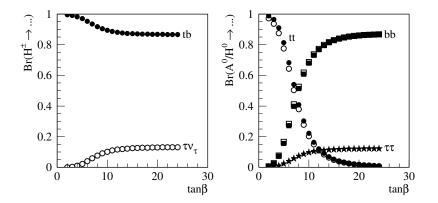
Charged and neutral Higgs boson pair production



In the following, we use an integrated luminosity of 3000 fb^{-1} .

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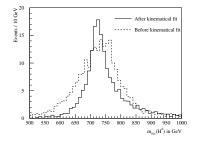
Charged and neutral Higgs boson decays



In the following, we assume that the charged and neutral Higgs bosons only decay into SM particles.

 $e^+e^- \rightarrow H^+H^- \rightarrow tbtb$

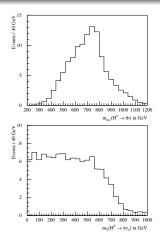
- Events with no isolated lepton, at least 8 jets including 4 b-jets,
- Assignment of the non-b jets to 2 W bosons, reconstruction of top quarks and of the charged Higgs bosons,
- Mass constrained kinematical fit to improve the reconstruction.



A cut on m_{bb} is then applied to further reduce the $e^+e^- \rightarrow tbtb$ background.

 $e^+e^- \rightarrow H^+H^- \rightarrow tb\tau\nu$

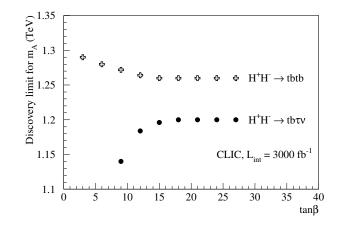
- Events with no isolated lepton, at least 5 jets including 2 b-jets and 1 τ-jet,
- Assignment of 2 non-b jets to a W boson, reconstruction of the top quark and of H[±] → tb,
- Transverse mass reconstruction for $H^{\pm} \rightarrow \tau \nu$.



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Cuts on the missing P_T , on the transverse angle between the charged Higgs boson candidates and on the transverse mass are applied to further reduce the $e^+e^- \rightarrow tb\tau\nu$ background.

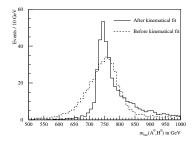
Charged Higgs boson discovery potential at CLIC



For a discovery, one requires $S \ge 10$ and $S/\sqrt{B} \ge 5$.

This decay chain has the largest branching ratio at large tan β .

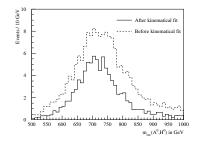
- Events with no isolated lepton and 4 b-jets,
- Assignment of two bb pairs to the neutral Higgs bosons,
- Mass constrained kinematical fit to improve the reconstruction.



$e^+e^- \rightarrow A^0 H^0 \rightarrow tttt$

This decay chain has the largest branching ratio at small $\tan \beta$.

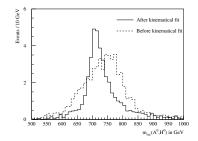
- Events with no isolated lepton, at least 12 jets, including 4 b-jets,
- Assignment of 8 non-b jets to 4 W bosons, reconstruction of 4 top quarks and assignment of *tt* pairs to the neutral Higgs bosons,
- Mass constrained kinematical fit to improve the reconstruction...
 Poor convergence efficiency, due to the complex event topology.



$e^+e^- ightarrow A^0 H^0 ightarrow ttbb$

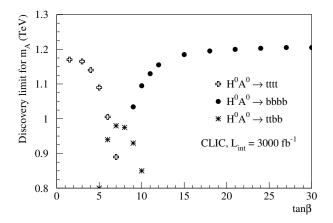
This decay chain has a significant branching ratio in the intermediate tan β region (around 7).

- Events with no isolated lepton, at least 8 jets including 4 b-jets,
- Assignment of the non-b jets to 2 W bosons, reconstruction of top quarks and of the neutral Higgs bosons (*tt* and *bb*),
- Mass constrained kinematical fit to improve the reconstruction.



A cut on $|\Delta m_{tb}|$ can be applied to reduce the contribution of $e^+e^- \rightarrow H^+H^- \rightarrow tbtb$ events.

Neutral Higgs boson discovery potential at CLIC



The discovery limit is set by the *bbbb* and *tttt* channels, except in the intermediate $\tan \beta$ region, where the *tbtb* cascade decay can also be observed.

$\tan \beta$ determination with charged Higgs bosons (1)

$$\frac{\Gamma(H^{\pm} \to tb)}{\Gamma(H^{\pm} \to \tau\nu)} \simeq \frac{3\Delta_{\text{QCD}}}{m_{\tau}^{2}} \times \left[\bar{m}_{t}^{2}(m_{H^{\pm}})\cot^{4}\beta + \bar{m}_{b}^{2}(m_{H^{\pm}})\right]$$
$$R = \frac{N_{tbtb}}{N_{tb\tau\nu}} = \frac{\epsilon_{tbtb}}{2\epsilon_{tb\tau\nu}} \times \frac{\Gamma(H^{\pm} \to tb)}{\Gamma(H^{\pm} \to \tau\nu)}$$

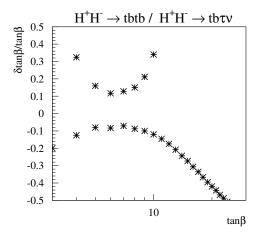
- One can determine tan β from the ratio *R* between the signal rates for $H^+H^- \rightarrow tbtb$ and $H^+H^- \rightarrow tb\tau\nu$.
- The result does not depend on possible other H^{\pm} decays.
- The (statistical) error on $\tan \beta$ is directly derived from:

$$\frac{\Delta R}{R} = \sqrt{\left[\frac{\Delta(\sigma \times Br)}{\sigma \times Br}\right]_{tbtb}^{2} + \left[\frac{\Delta(\sigma \times Br)}{\sigma \times Br}\right]_{tb\tau\nu}^{2}}$$

$\tan \beta$ determination with charged Higgs bosons (2)

The statistical error on tan β is smallest in the 4-10 region.

- At low tan β , the signal rate for $H^+H^- \rightarrow tb\tau\nu$ is very small.
- At large $\tan \beta$, the ratio *R* is constant.



$\tan \beta$ determination with neutral Higgs bosons (1)

 $\frac{\Gamma(H^0/A^0 \to tt)}{\Gamma(H^0/A^0 \to bb)}$ varies with $\tan \beta$ similarly to $\frac{\Gamma(H^{\pm} \to tb)}{\Gamma(H^{\pm} \to \tau\nu)}$.

• One can determine $\tan \beta$ from three ratios:

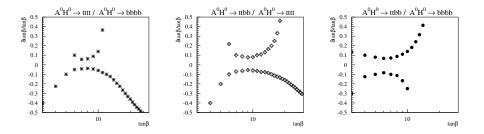
- R_{bbbb}^{tttt} between $H^0A^0 \rightarrow tttt$ and $H^0A^0 \rightarrow bbbb$,
- R_{bbbb}^{ttbb} between $H^0 A^0 \rightarrow ttbb$ and $H^0 A^0 \rightarrow bbbb$,
- R_{ttbb}^{tttt} between $H^0A^0 \rightarrow tttt$ and $H^0A^0 \rightarrow ttbb$.
- The results do not depend on other H^0/A^0 decays.
- The (statistical) error on $\tan \beta$ is directly derived from:

$$\frac{\Delta R_2^1}{R_2^1} = \sqrt{\left[\frac{\Delta(\sigma \times Br)}{\sigma \times Br}\right]_{\text{Signal 1}}^2 + \left[\frac{\Delta(\sigma \times Br)}{\sigma \times Br}\right]_{\text{Signal 2}}^2}$$

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$\tan \beta$ determination with neutral Higgs bosons (2)

The statistical errors on tan β are smallest in the 4-10 region.



Conclusion and outlooks

- New simulation studies of the charged and neutral Higgs boson decays show that CLIC will be sensitive to these new particles over the whole tan β spectrum, for masses beyond 1 TeV.
- At CLIC, tan β can be measured with a good accuracy in the intermediate region (not accessible at LHC) through a comparison of the signal rates for various H[±] and H⁰/A⁰ decays.
- A study of H⁰A⁰ → ttττ/bbττ and a combined analysis of the charged and neutral Higgs sectors will be performed next...

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