# Degenerate scalar-tensor theories a geometrical approach

#### **Gianmassimo Tasinato**

Swansea University

Based on works with Javier Chagoya, Marco Crisostomi + Kazuya Koyama, David Langlois, Karim Noui

# Question

What is the most general consistent scalar-tensor theory?

- ▶ General Relativity is the only consistent 4d theory describing a massless, interacting spin two field (2 tensor dofs).
- ▶ Is there any analogue result for **covariant scalar-tensor** set-ups propagating massless spin two + spin zero fields (2+1 dofs)?

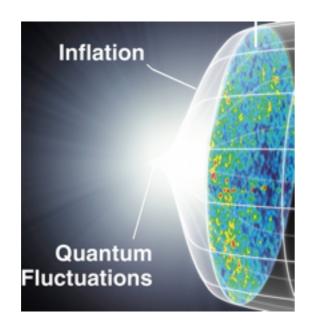
# Plan

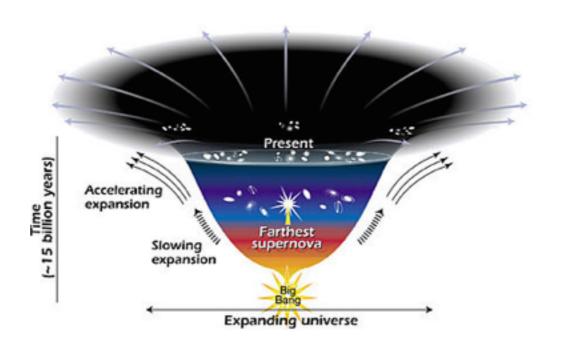
- 1. **Motivations** for considering this question
- 2. What do we know about its answer

little, but a geometrical approach can be useful (maybe)

# Why considering them?

- ► Naturally arise from **string theory**
- ► Important applications to **cosmology**





# Why considering them?

- ► Naturally arise from **string theory**
- ► Important applications to **cosmology**

Prototype 
$$S=\int d^4x \sqrt{-g} \left[ \frac{M_{Pl}^2}{2} \, R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) \right]$$
 where  $\chi$  interactions Canonical kinetic term

► Much more can be done...

#### What about derivative scalar self-interactions?

- $\blacktriangleright$  interactions involving **single derivative** of scalars. Call  $X = \partial_{\mu}\pi\partial^{\mu}\pi$ 
  - Prototype: **DBI action**  $S = \int d^4x \sqrt{-g} \sqrt{1+X}$
  - More generally: any function  $S = \int d^4x \sqrt{-g} P(X)$  is fine EOMs are second order
- interactions including **second derivatives** of scalars

[Nicolis, Rattazzi, Trincherini]

$$\mathcal{L}_2 = -\frac{1}{2} \left( \partial \pi \right)^2$$

The Galileons

$$\mathcal{L}_3 = (\partial \pi)^2 \, \Box \pi$$

$$\mathcal{L}_4 = (\partial \pi)^2 \left[ (\Box \pi)^2 - (\partial_\mu \partial_\nu \pi)^2 \right]$$

$$\mathcal{L}_{5} = (\partial \pi)^{2} \left[ (\Box \pi)^{3} + 2 \left( \partial_{\mu} \partial_{\nu} \pi \right)^{3} - 3 \Box \pi \left( \partial_{\mu} \partial_{\nu} \pi \right)^{2} \right]$$

Screening effects



Acceleration of the universe





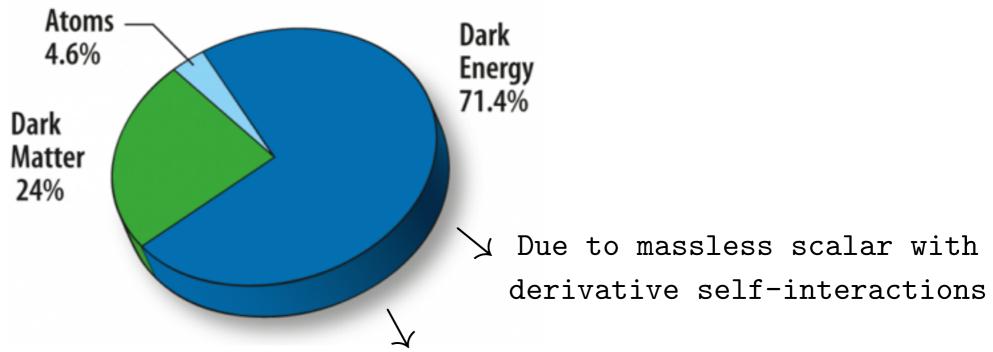
Non-renormalization theorems

The Galileons

#### Acceleration of the universe

#### self-acceleration

There exist branches of cosmological solutions that are asymptotically de Sitter, with no need of positive cosmological constant.



Regime where second derivatives  $\partial^2\pi$  are important

The Galileons

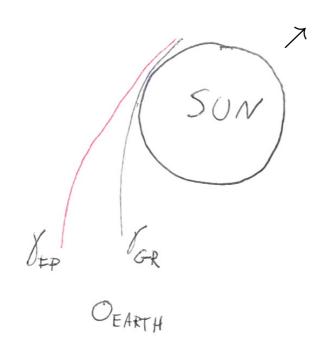
Screening effects

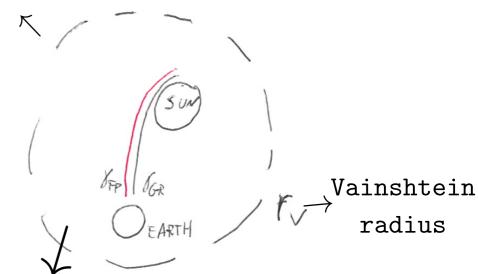
why don't we see it? Vainshtein mechanism:

non-linearities associated with self-interactions hide the effect of scalar nearby spherically symmetric sources

Evade stringent bounds on

deviations from GR





Regime where second derivatives  $\partial^2\pi$  are important

The Galileons

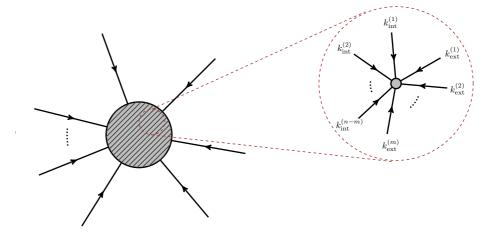
Non-renormalization theorems

[Nicolis, Rattazzi; Goon, Hinterbichler, Trodden]

Galileon interactions don't get renormalized in perturbation theory

Thanks to structure of interactions + Galilean symmetry

Safe radial interval where  $\partial^2 \pi$  is large and theory technically natural





 $\mathcal{L}_2 = -\frac{1}{2} (\partial \pi)^2$ 

The Galileons

 $\mathcal{L}_3 = (\partial \pi)^2 \, \Box \pi$ 

 $\mathcal{L}_4 = (\partial \pi)^2 \left[ (\Box \pi)^2 - (\partial_\mu \partial_\nu \pi)^2 \right]$ 

 $\mathcal{L}_{5} = (\partial \pi)^{2} \left[ (\Box \pi)^{3} + 2 \left( \partial_{\mu} \partial_{\nu} \pi \right)^{3} - 3 \Box \pi \left( \partial_{\mu} \partial_{\nu} \pi \right)^{2} \right]$ 

Screening effects



Acceleration of the universe





Good for Dark Energy!

Non-renormalization theorems

Horndenski (1974)

# Question

What's the most general scalar-tensor action leading to 2nd order equations of motion?

Horndenski (1974)

# Question

What's the most general scalar-tensor action leading to 2nd order equations of motion?

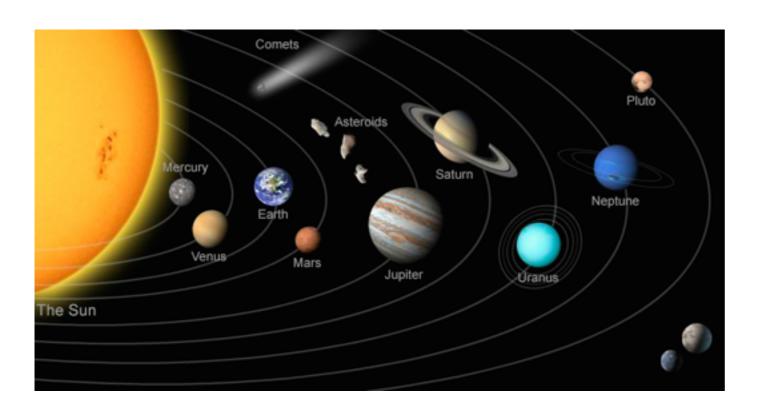
$$L_{2}^{H} \equiv G_{2}(\phi, X) , \qquad L_{3}^{H} \equiv G_{3}(\phi, X) \Box \phi ,$$

$$L_{4}^{H} \equiv G_{4}(\phi, X)^{(4)}R - 2G_{4,X}(\phi, X)(\Box \phi^{2} - \phi^{\mu\nu}\phi_{\mu\nu}) ,$$

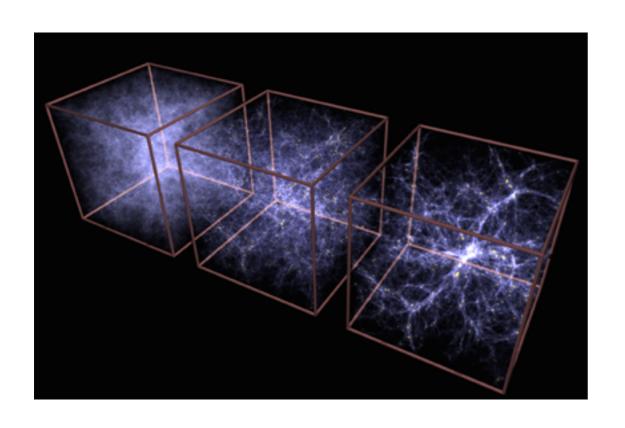
$$L_{5}^{H} \equiv G_{5}(\phi, X)^{(4)}G_{\mu\nu}\phi^{\mu\nu} + \frac{1}{3}G_{5,X}(\phi, X)(\Box \phi^{3} - 3\Box \phi \phi_{\mu\nu}\phi^{\mu\nu} + 2\phi_{\mu\nu}\phi^{\mu\sigma}\phi^{\nu}{}_{\sigma}) ,$$

where the  $G_i$  are arbitrary functions of  $\phi$ , X

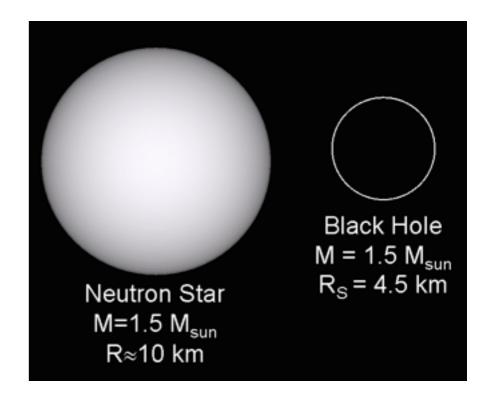
- contains up to three powers of  $\partial^2 \phi$ . Flat space  $\to$  Galileons
- lot of activity for applications to cosmology, gravity/BHs etc etc



# $\leftarrow$ ok







Horndenski (1974)

# Question

What's the most general scalar-tensor action leading to 2nd order equations of motion?



Is this really necessary?

Horndenski (1974)

# Question

What's the most general scalar-tensor action leading to 2nd order equations of motion?

Ostrogradsky theorem

[see e.g. Woodard]

Any non-degenerate theory with **EOMs of order higher than two** has **Hamiltonian unbounded from below** 

Ostrogradsky ghost

Horndenski (1974)

# Question

What's the most general scalar-tensor action leading to 2nd order equations of motion?

Ostrogradsky theorem

[see e.g. Woodard]

Any non-degenerate theory with **EOMs of order higher than two** has **Hamiltonian unbounded from below** 



Let's consider degenerate scalar-tensor theories

Relations defining conjugate momenta can't be fully inverted: velocities can't be expressed in terms of fields and their conjugate momenta

Constraint conditions exist!

In this way new consistent covariant scalar-tensor theories can be found, that propagate only up to three degrees of freedom

no Ostrogradsky ghost mode



Let's consider degenerate scalar-tensor theories

Relations defining conjugate momenta can't be fully inverted: velocities can't be expressed in terms of fields and their conjugate momenta

Beyond Horndenski [Gleyzes et al]

$$L_4^{\text{bH}} \equiv F_4(\phi, X) \epsilon^{\mu\nu\rho}{}_{\sigma} \epsilon^{\mu'\nu'\rho'\sigma} \phi_{\mu} \phi_{\mu'} \phi_{\nu\nu'} \phi_{\rho\rho'} ,$$
  

$$L_5^{\text{bH}} \equiv F_5(\phi, X) \epsilon^{\mu\nu\rho\sigma} \epsilon^{\mu'\nu'\rho'\sigma'} \phi_{\mu} \phi_{\mu'} \phi_{\nu\nu'} \phi_{\rho\rho'} \phi_{\sigma\sigma'} ,$$



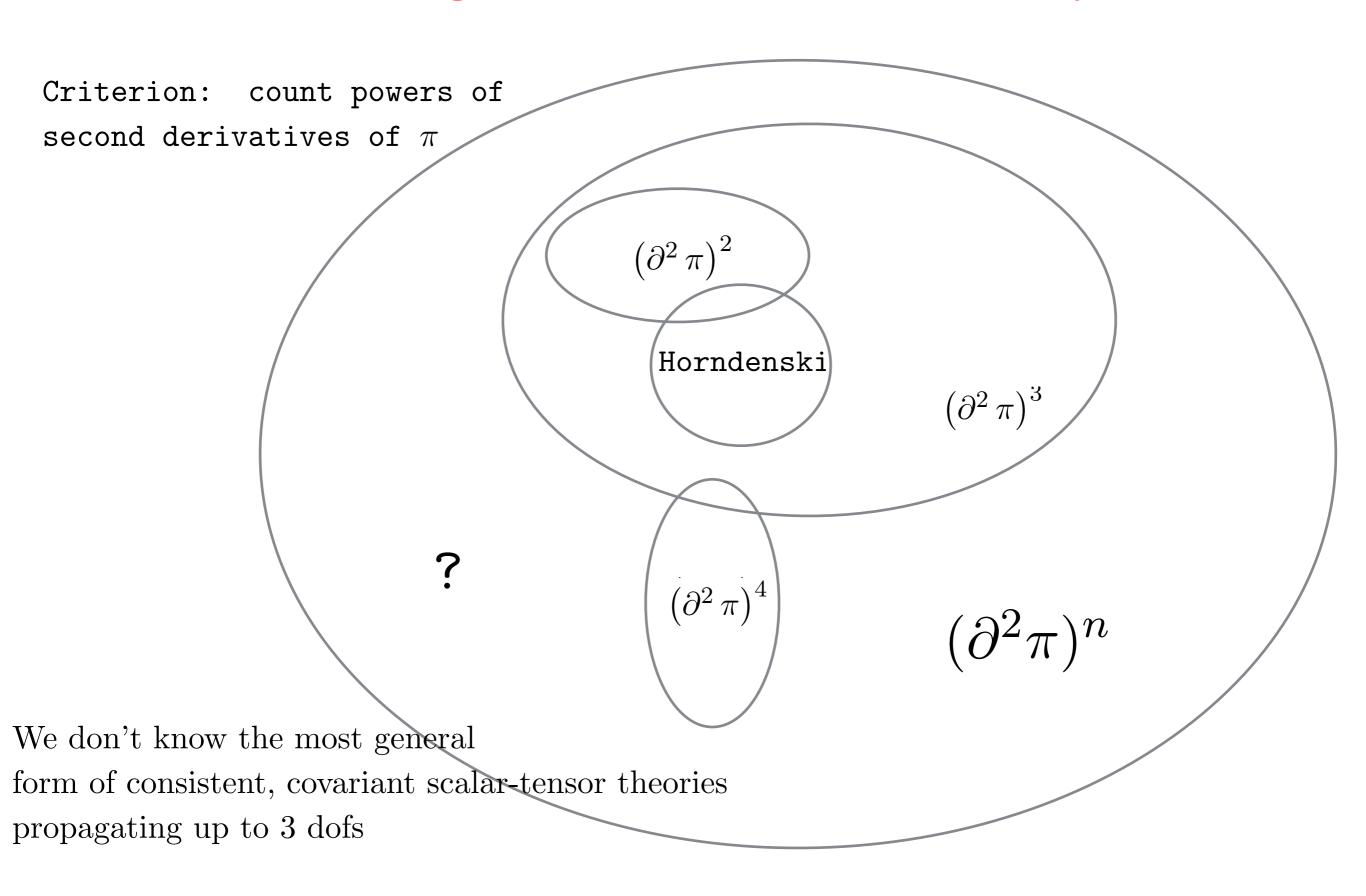
Let's consider degenerate scalar-tensor theories

Relations defining conjugate momenta can't be fully inverted: velocities can't be expressed in terms of fields and their conjugate momenta

Extended Scalar Tensor [Langlois, Noui; Crisostomi, Koyama, GT; Ben Achour et al]

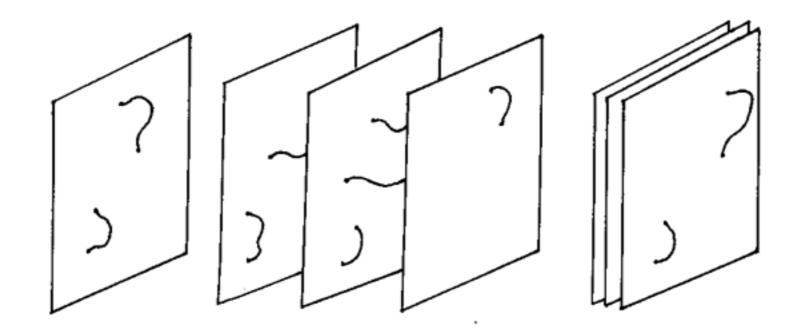
- Use Hamiltonian approach to systematically find **all consistent theories**, classifying them in terms of **number of second derivatives of scalar**
- New theories found, with interesting consequences for dark energy (growth of structure) + screening mechanisms (neutron stars etc)

# Most general scalar-tensor theory?



# Question

Do some of these theories admit a geometrical interpretation, which might reveal some symmetries (recall Galilean symmetry)



Recall lessons from Galileons: let's think to DE applications

[de Rham, Tolley]

• Probe brane in Minkowski space: brane position  $\pi$ 

$$g_{\mu\nu} = \eta_{\mu\nu} + \partial_{\mu}\pi\partial_{\nu}\pi$$

DBI action:  $S_{\lambda} = -\lambda \int d^4x \sqrt{-g} = -\lambda \int d^4x \sqrt{1 + (\partial \pi)^2}$ 

• Action is invariant under a symmetry

$$\delta_v \pi(x) = v_\mu x^\mu + \pi(x) v^\mu \partial_\mu \pi(x)$$

inherited from 5d Poincaré symmetry

#### • Generalization

Consider higher dimensional Lovelock invariants + associated Gibbons-Hawking terms

$$S_K = -M_5^3 \int d^4x \sqrt{-g}K$$
 ;  $S_R = \frac{M_4^2}{2} \int d^4x \sqrt{-g}R$  ;

$$S_{GB} = -\beta \frac{M_5^3}{m^2} \int d^4x \sqrt{-g} \, \mathcal{K}_{GB}$$
;  $\kappa_{GB} = -\frac{2}{3} K_{\mu\nu}^3 + K K_{\mu\nu}^2 - \frac{1}{3} K^3 - 2G_{\mu\nu} K^{\mu\nu}$ 

- EOMs for  $\pi$  remain second order
- Symmetry is still preserved  $\delta_v \pi(x) = v_\mu x^\mu + \pi(x) v^\mu \partial_\mu \pi(x)$

#### • Generalization

Consider higher dimensional Lovelock invariants + associated Gibbons-Hawking terms

$$\begin{split} S_{\lambda} &= -\lambda \int \mathrm{d}^4 x \sqrt{-g} = -\lambda \int \mathrm{d}^4 x \sqrt{1 + (\partial \pi)^2} \\ S_{K} &= -M_5^3 \int \mathrm{d}^4 x \sqrt{-g} K = M_5^3 \int \mathrm{d}^4 x \, \left( [\Pi] - \gamma^2 [\phi] \right) \\ S_{R} &= \frac{M_4^2}{2} \int \mathrm{d}^4 x \sqrt{-g} R = \frac{M_4^2}{2} \int \mathrm{d}^4 x \, \gamma \left( \left( [\Pi]^2 - [\Pi^2] \right) + 2 \gamma^2 \left( [\phi^2] - [\Pi] [\phi] \right) \right) \\ S_{GB} &= -\beta \frac{M_5^3}{m^2} \int \mathrm{d}^4 x \sqrt{-g} \, \mathcal{K}_{GB} \\ &= \beta \frac{M_5^3}{m^2} \int \mathrm{d}^4 x \, \gamma^2 \left( \frac{2}{3} \left( [\Pi]^3 + 2 [\Pi^3] - 3 [\Pi] [\Pi^2] \right) + 4 \gamma^2 ([\Pi] [\phi^2] - [\phi^3] \right) \\ &\quad - 2 \gamma^2 ([\Pi]^2 - [\Pi^2]) [\phi] \right), \end{split}$$

#### Generalization

Consider higher dimensional Lovelock invariants + associated Gibbons-Hawking terms

Take non-relativistic limit  $\partial\pi\ll 1$  (but still  $\partial^2\pi$  large)

$$S_{2} = S_{\lambda}^{NR} = -\frac{\lambda}{2} \int d^{4}x \, (\partial \pi)^{2}$$

$$S_{3} = S_{K}^{NR} = \frac{M_{5}^{3}}{2} \int d^{4}x \, (\partial \pi)^{2} \, \Box \pi$$

$$S_{4} = S_{R}^{NR} = \frac{M_{4}^{2}}{4} \int d^{4}x \, (\partial \pi)^{2} \, \left( (\Box \pi)^{2} - (\partial_{\mu} \partial_{\nu} \pi)^{2} \right)$$

$$S_{5} = S_{GB}^{NR} = \beta \frac{M_{5}^{3}}{3m^{2}} \int d^{4}x \, (\partial \pi)^{2} \, \left( (\Box \pi)^{3} + 2(\partial_{\mu} \partial_{\nu} \pi)^{3} - 3\Box \pi (\partial_{\mu} \partial_{\nu} \pi)^{2} \right)$$

• Symmetry becomes Galilean symmetry

$$\delta_v \pi(x) = v_\mu x^\mu + \pi(x) v^\mu \partial_\mu \pi(x) \quad \Rightarrow \quad \delta \pi = v_\mu x^\mu$$

#### Can we find an analogous geometric construction for EST?

▶ Build probe brane actions made of powers of extrinsic curvature

$$S_K = \int \sqrt{-g} K_\mu^{\ \mu}$$

$$S_{K^2} = \int \sqrt{-g} K_\mu^{\ \nu} K_\nu^{\ \mu}$$
...

▶ These actions have symmetries associated with bulk isometries

$$\delta_v \pi(x) = v_\mu x^\mu + \pi(x) v^\mu \partial_\mu \pi(x).$$

...but when coupled with gravity, their EOMs have higher order derivatives

#### Can we find an analogous geometric construction for EST?

► Take 'ultrarelativistic' limit  $\partial \pi \gg 1$ 

$$\mathcal{L}_{1} = \Lambda^{2} \sqrt{|X|} 
\mathcal{L}_{2} = \Lambda \left( [\Pi] - \frac{1}{X} [\Phi] \right) 
\mathcal{L}_{3} = \frac{1}{\sqrt{|X|}} \left( [\Pi]^{2} - [\Pi^{2}] + \frac{2}{X} \left( [\Phi^{2}] - [\Phi] [\Pi] \right) \right) 
\mathcal{L}_{4} = \frac{1}{\Lambda X} \left( [\Pi]^{3} + 2[\Pi^{3}] - 3[\Pi^{2}] [\Pi] + \frac{3}{X} \left( 2[\Pi] [\Phi^{2}] - 2[\Phi^{3}] - [\Phi] [\Pi]^{2} + [\Phi] [\Pi^{2}] \right) \right)$$

symmetry left

$$\delta\pi \,=\, \pi\,w^\mu\partial_\mu\pi$$

First time that such symmetry is noticed for these theories

#### Can we find an analogous geometric construction for EST?

► Take 'ultrarelativistic' limit  $\partial \pi \gg 1$ 

$$\mathcal{L}_{1} = \Lambda^{2} \sqrt{|X|}$$

$$\mathcal{L}_{2} = \Lambda \left( [\Pi] - \frac{1}{X} [\Phi] \right)$$

$$\mathcal{L}_{3} = \frac{1}{\sqrt{|X|}} \left( [\Pi]^{2} - [\Pi^{2}] + \frac{2}{X} \left( [\Phi^{2}] - [\Phi] [\Pi] \right) \right)$$

$$\mathcal{L}_{4} = \frac{1}{\Lambda X} \left( [\Pi]^{3} + 2[\Pi^{3}] - 3[\Pi^{2}] [\Pi] + \frac{3}{X} \left( 2[\Pi] [\Phi^{2}] - 2[\Phi^{3}] - [\Phi] [\Pi]^{2} + [\Phi] [\Pi^{2}] \right) \right)$$

▶ when coupled with gravity, one gets degenerate scalar tensor theories

▶ Good news for model building: we can model models of dark energy/inflation using EST with underlying symmetry preserving structure of action

Can we find an analogous geometric construction for EST?

#### yes

Take 'ultra-relativistic' limit of of actions built with combinations of  $K_{\mu}{}^{\nu}$ 

Does the same work with induced action made with other combinations of curvature fluctuations?  $R_{\mu\nu}R^{\mu\nu}$  etc?

? work in progress....?

# Conclusions

- ▶ The structure of **scalar-tensor** theories is richer than expected: degenerate systems
  - We still don't know the **most general** theory for massless spin 2 + spin 0 fields
  - Cosmologists are very interested to this question, for building new models of dark energy and inflation
- ► I presented
  - most recent findings about what we know of these theories
  - A geometrical approach which makes apparent previously unnoticed underlying symmetry

#### TO DO NEXT

► Continue the exploration of **geometrical approach** by studying **ultrarelativistic** limit of **brane induced actions** made with curvature invariants