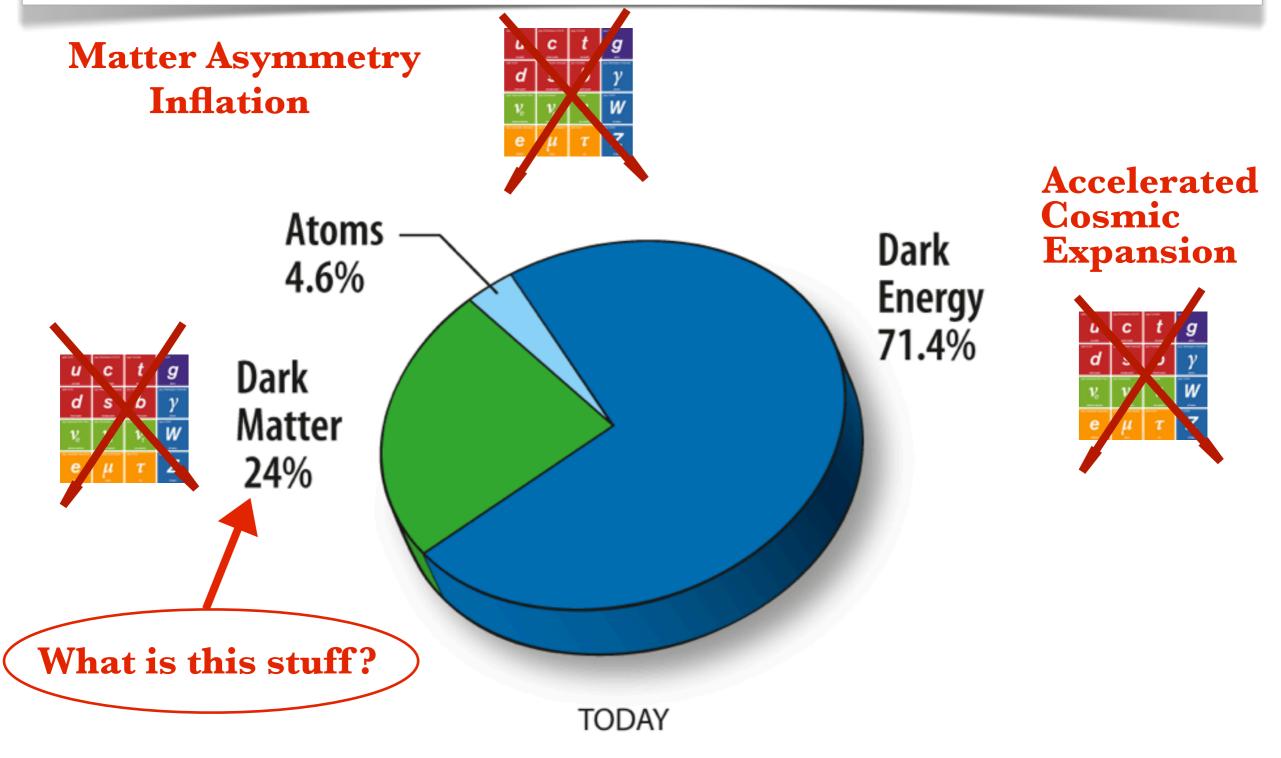
# Light Thermal DM & Higgs Portal Mediators

# Gordan Krnjaic Fermilab

Phys.Rev. D94 (2016) 1512.04119

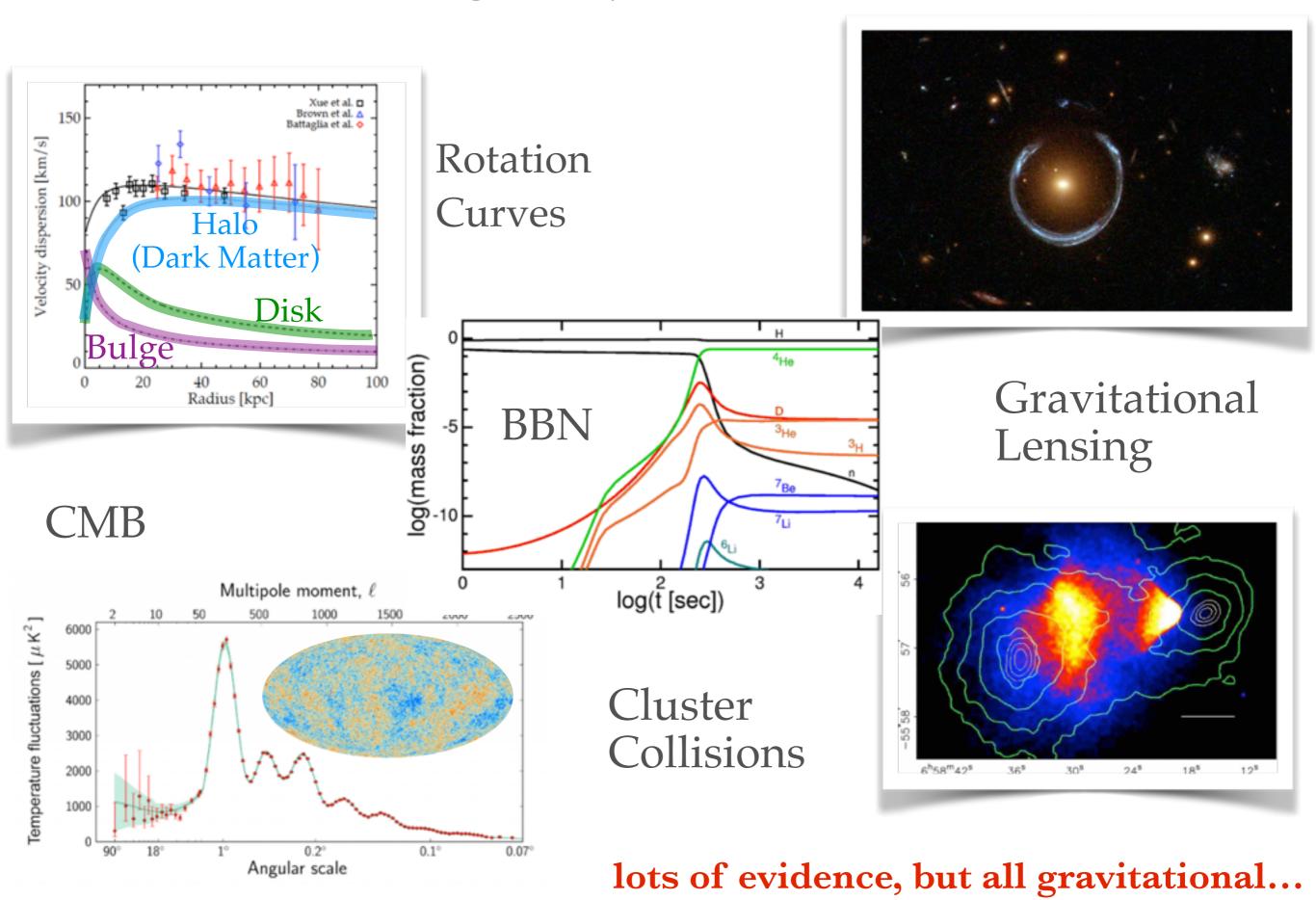
CERN-EPFL-Korea Theory institute
"New Physics at the Intensity Frontier"

# Zeroth Order Outstanding Problems



All other problems generically require extremely high scales .... DM can be much lighter

# **Obligatory DM Slide**



# Historical Analogy Understanding the Weak Force

Discovery of radioactivity

(1890s)

Fermi Scale identified

$$G_F \sim \frac{1}{(100 \text{ GeV})^2}$$
 (1930s)

Non-Abelian Gauge Theory

(1950s)

Higgs Mechanism

(1960s)

W/Z bosons

(1970s)

Higgs discovered

(2010s)

Each step required revolutionary theoretical/experimental leaps

 $t \sim 100 \text{ years}$ 

# How long will we wait for DM?

Discovery of missing mass (1930s)

Rotation curves (1970s)

CMB power spectrum (1990s)

Relevant scale? > 2017

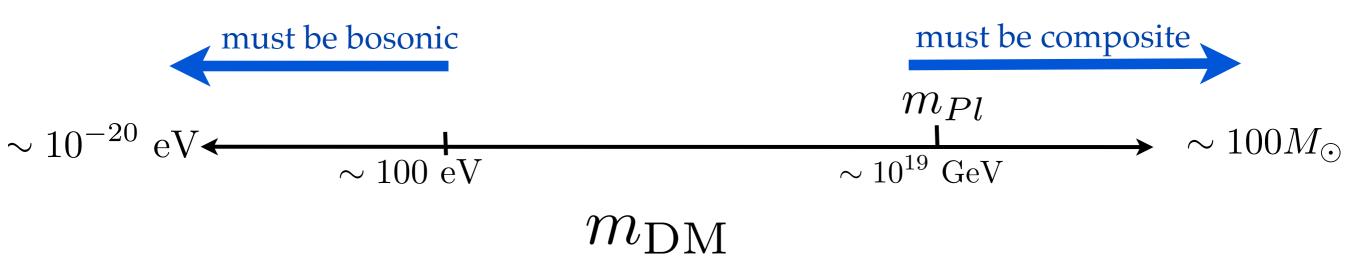
Non-gravitational interactions not guaranteed No clear target of opportunity

Discovery time frame? t > 80 yrs

# DM Prognosis?

### Bad news: DM-SM interactions are not obligatory

If nature is unkind, we may never know the right scale



**Good news:** most *discoverable* DM candidates are in thermal equilibrium with us in the early universe

# Why is this good news?

# Thermal Equilibrium Advantage #1: Minimum Annihilation Rate

Equilibrium, achieved easily with a tiny DM/SM coupling

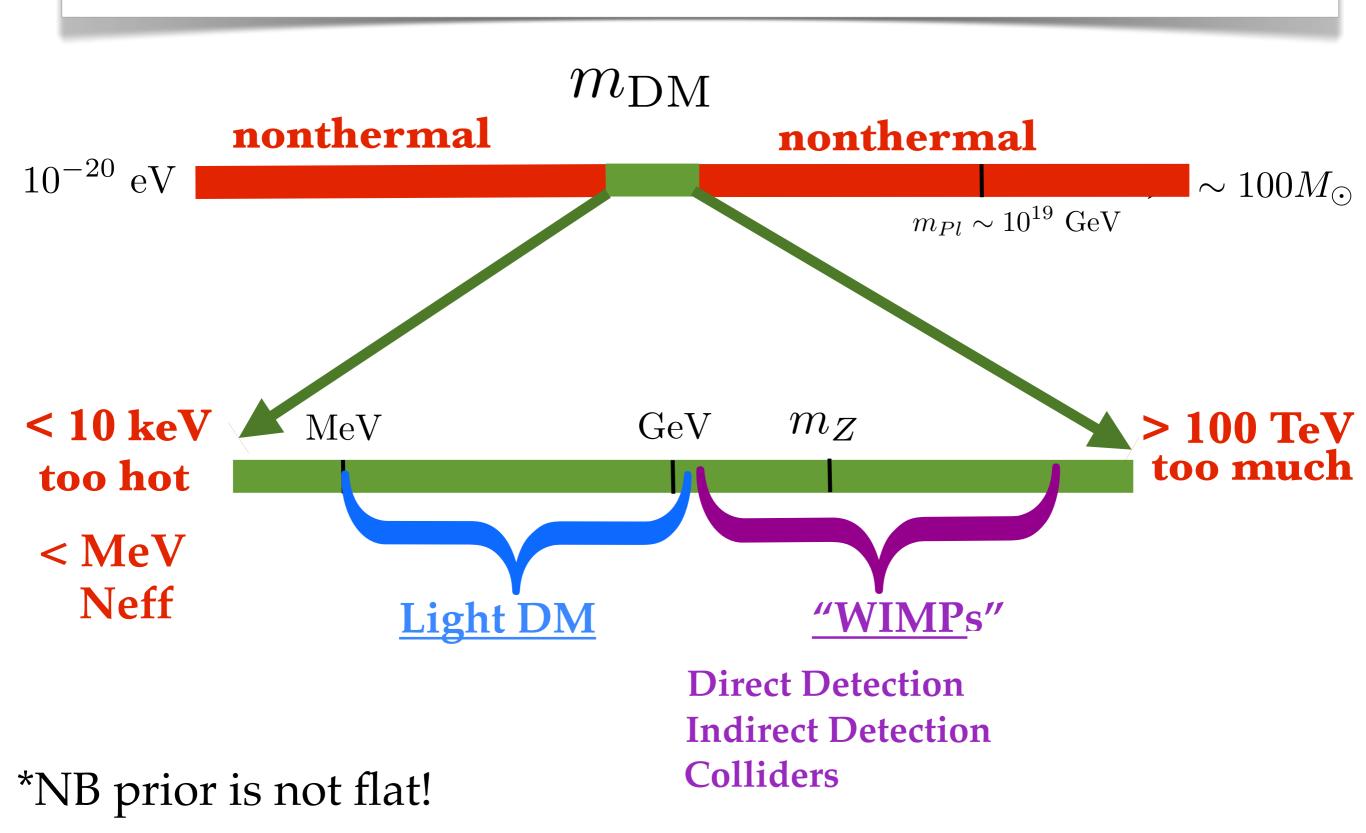
$$n_{\rm DM} = \int \frac{d^3p}{(2\pi)^3} \frac{g_i}{e^{E/T} \pm 1} \sim T^3$$

#### Generically overproduces DM

Requires much larger annihilation cross section to deplete

$$\sigma v \ge \sigma v_{\rm relic} \sim 3 \times 10^{-26} {\rm cm}^3 {\rm s}^{-1}$$

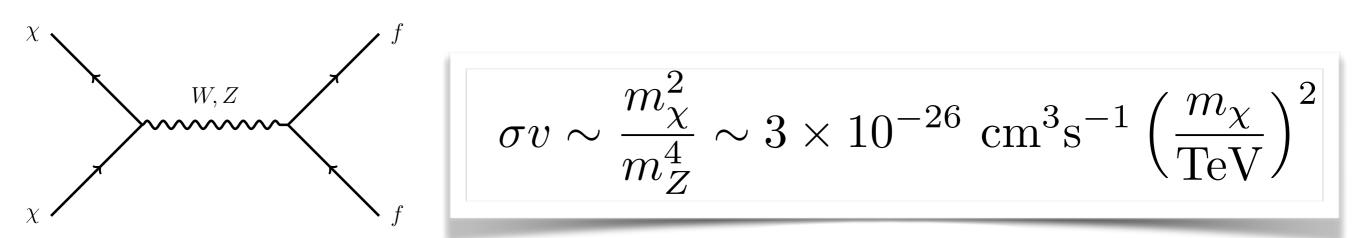
# Thermal Equilibrium Advantage #2: Narrows Mass Range



# Heavy vs. Light #1

LDM needs new forces

Heavy DM can achieve right abundance w/ SM weak force



For LDM, annihilation via SM forces is too weak so equilibrium is lost too soon

$$m_{\chi} \sim \text{GeV} \implies \sigma v \ll 3 \times 10^{-26} \text{ cm}^3/\text{s}$$

DM overproduced unless there are light new force carriers

# **Avoiding LDM Overproduction**

Choose light mediator

#### Must be SM singlet, options limited by SM gauge invariance

Vector Portal mix w/ photon (or gauge B-L etc.)

$$\epsilon F_{\mu\nu}F'_{\mu\nu}$$

Neutrino Portal mix w/ RHN

$$H^{\dagger}LN$$

Higgs Portal mix w/ SM Higgs

$$(H^{\dagger}H)\phi$$

# **Avoiding LDM Overproduction**

Choose light mediator

### Must be SM singlet, options limited by SM gauge invariance

Vector Portal mix w/ photon (or gauge B-L etc.)



Lots of attention here (wont discuss here)

Neutrino Portal mix w/ RHN



Hard to make thermal (decays,  $\nu$  masses etc.)

Higgs Portal mix w/ SM Higgs

$$(H^{\dagger}H)\phi$$

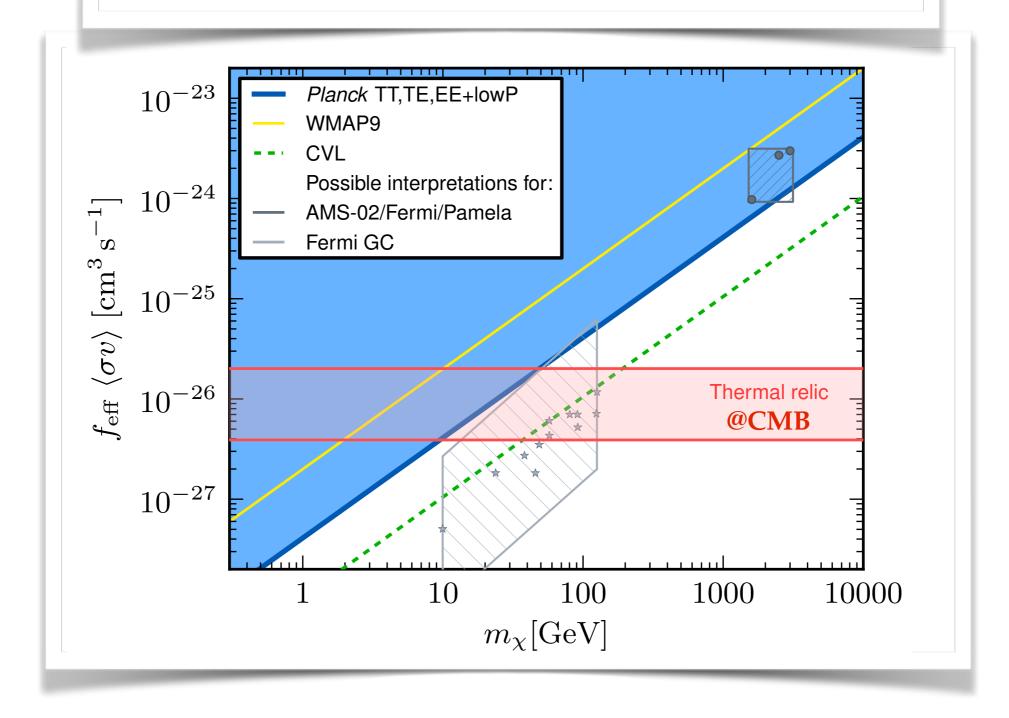
 $\sin \theta$   $- \times - -$ 

SM coupling  $\propto$  fermion mass

# Heavy vs. Light # 2

CMB rules out LDM < 10 GeV?

Planck 1303.5076



Safe models:

cross section is smaller @ CMB or annihilation stops @ CMB

# **Option 1: Smaller Cross Section**

Velocity/Temperature Dependence

$$\sigma v \propto v^2$$

Rate large at freeze-out w/  $v \sim 0.1 c$ 

$$\langle \sigma v \rangle \big|_{T=m_{\chi}} = 3 \times 10^{-26} \text{ cm}^3/\text{s} \implies \Omega_{\chi} = \Omega_{\text{DM}}$$

#### Velocity redshifted at late times

$$\langle \sigma v \rangle \big|_{T=eV} \ll 3 \times 10^{-26} \text{ cm}^3/\text{s} \implies \text{CMB safe}$$

Choose DM/mediator combination to get v-dependence

# **Option 2: Annihilation Stops Later**

Case Study: Asymmetric DM

#### Annihilation @ $T \sim m$ reduces antiparticle fraction

$$n_{\chi} \neq n_{\bar{\chi}} \propto \exp(-\langle \sigma v \rangle)$$

Counterintuitive: larger cross section is safer!

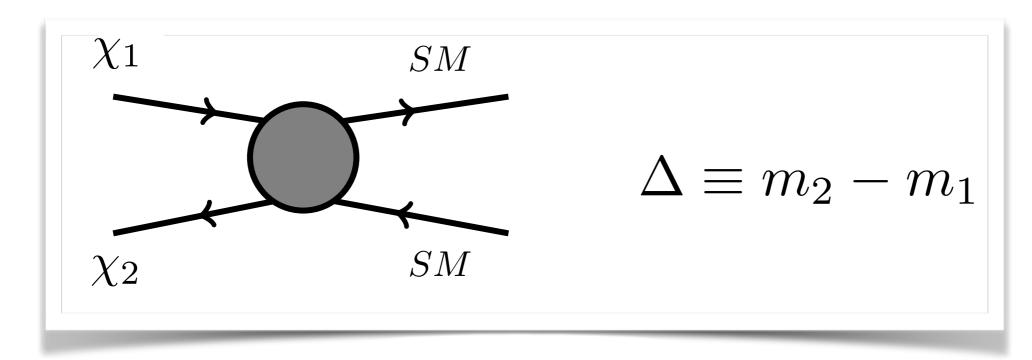
$$\frac{f_{\text{eff.}}\langle \sigma v \rangle e^{-\langle \sigma v \rangle}}{m_{\chi}} \ll 2 \times 10^{-28} \text{ cm}^3 \text{ s}^{-1} \text{ GeV}^{-1}$$

Easily satisfies CMB bound with  $\langle \sigma v \rangle > 3 \times 10^{-26} {\rm cm}^3 {\rm \ s}^{-1}$  as required for asymmetric DM

# **Option 2: Annihilation Stops Later**

Case Study: Inelastic couplings

#### Two-level co-annihilating system



As universe cools, heavier state is Boltzmann suppressed

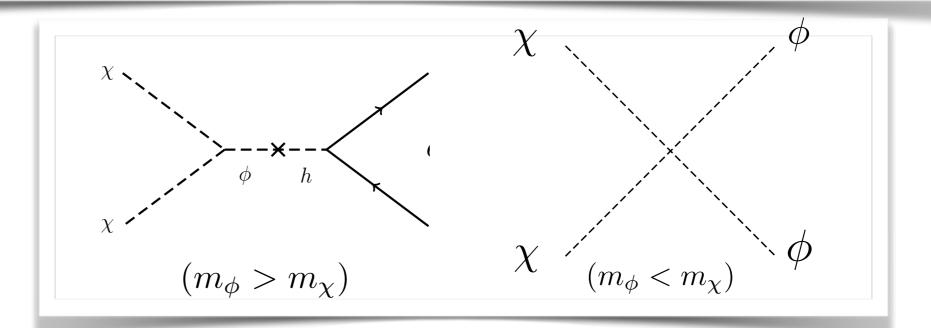
$$n_{\chi_2} \propto e^{-\Delta/T}$$

annihilation shuts off at late times

Generic (e.g if dark there are dark Dirac & Majorana masses)

### Choose CMB safe DM for for scalar mediator

**Scalar DM** s-wave annihilation

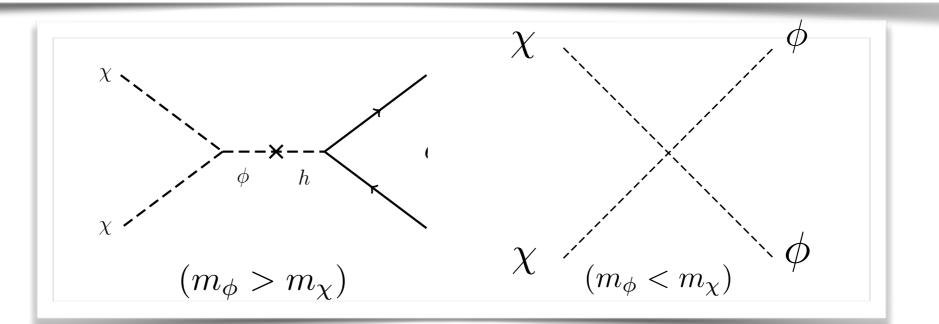


Need particle asymmetry and/or inelasticity for CMB safety

#### Choose CMB safe DM for for scalar mediator

#### Scalar DM

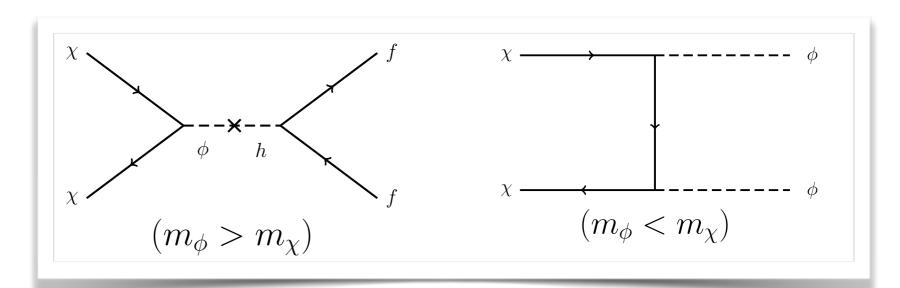
s-wave annihilation



Need particle asymmetry and/or inelasticity for CMB safety

### Fermion DM

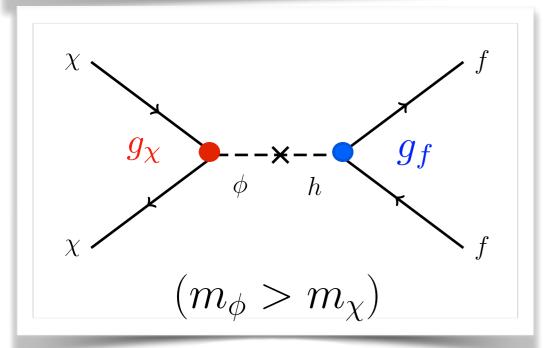
p-wave annihilation  $g_\chi \phi \bar{\chi} \chi$ 



Can also include  $\phi \bar{\chi} \gamma^5 \chi$  must be small (adds s-wave terms)

Both CMB safe  $\sigma v \propto v^2$ 

# Thermal Target: Direct Annihilation to SM



$$g_f = g_e \left(\frac{m_f}{m_e}\right)$$

Coupling scales with SM fermion mass

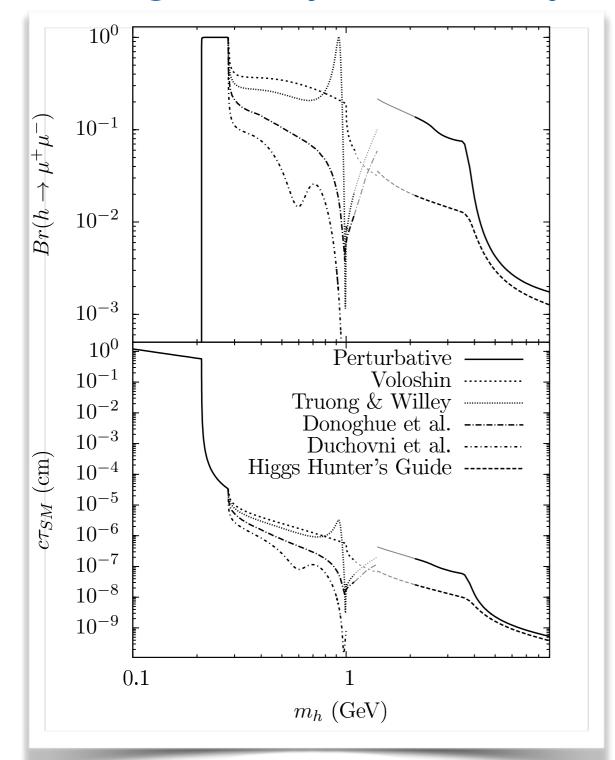
$$\sigma v = \sum_{f} (\sigma v)_f \propto g_{\chi}^2 g_e^2 \left(\frac{m_{\chi}}{m_{\phi}}\right)^4 \frac{1}{m_{\chi}^2}$$

Define dimensionless target 
$$y\equiv g_\chi^2g_e^2\left(\frac{m_\chi}{m_\phi}\right)^4$$

Normalized to electron coupling because rate exists for every mass point

## Thermal Target: Direct Annihilation to SM

#### Large theory uncertainty in SM coupling near QCD scale

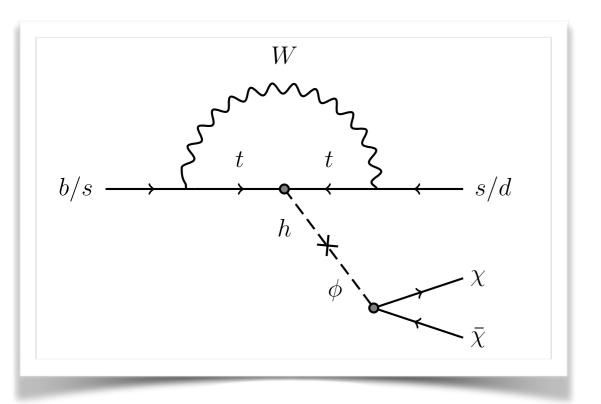


Estimate from different numerical extractions from light Higgs literature

$$g_f(s) \simeq \sin \theta \sqrt{\frac{8\pi}{m_h} \Gamma(h \to SM)} \Big|_{m_h = \sqrt{s}}$$

Clarke, Foot, Volkas 1310.8042

## Thermal Target: Rare B/K Decays



2/3 body decays to LDM/mediator

$$B^+ \to K^+ \phi, \ K^+ \chi \chi$$
  
 $K^+ \to \pi^+ \phi, \ \pi^+ \chi \chi$ 

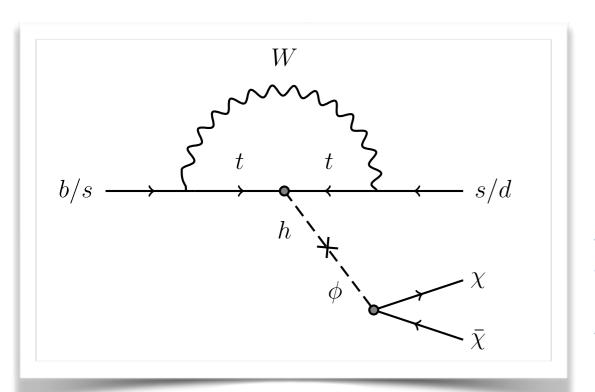
constrained by

$$B^+ \to K^+ \nu \bar{\nu} \qquad K^+ \to \pi^+ \nu \bar{\nu}$$

Annihilation rate set by \*small\* yukawas FCNC constraints set by \*top\* yukawa

Bird, Jackson, Kowalewski, Pospelov arXiv: 0401195

# How to compare meson decays w/target?



On shell decays to mediator independent of DM

Need to assume DM mass/coupling for thermal comparison

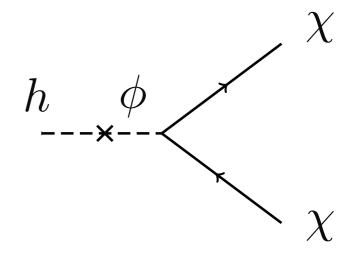
$$\Delta \text{Br}(B^+ \to K^+ \cancel{E}) \propto g_f^2 = y \times \frac{1}{g_\chi^2} \left(\frac{m_\phi}{m_\chi}\right)^4$$

Conservative worst-case "y" reach  $g_{\chi}$ ,  $m_{\phi}/m_{\chi} \to \mathcal{O}(1)$  (choose smallest mass ratio still consistent with direct annihilation)

# Thermal Target: Other Constraints

#### Higgs Invisible Width

Possibly compensate with additional *h* production, but can't avoid interference with 4*l* final state

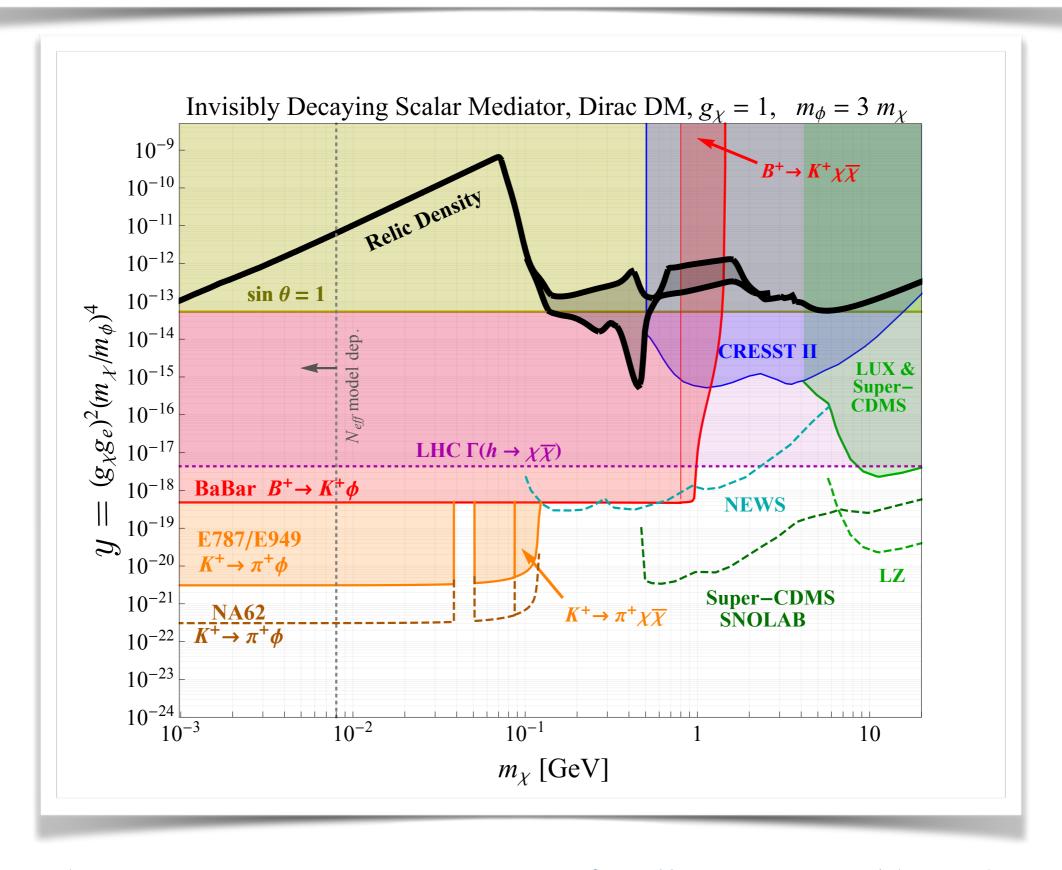


#### **Low Mass Direct Detection ~ GeV**

Assuming elastic scattering invariant comparison with thermal target

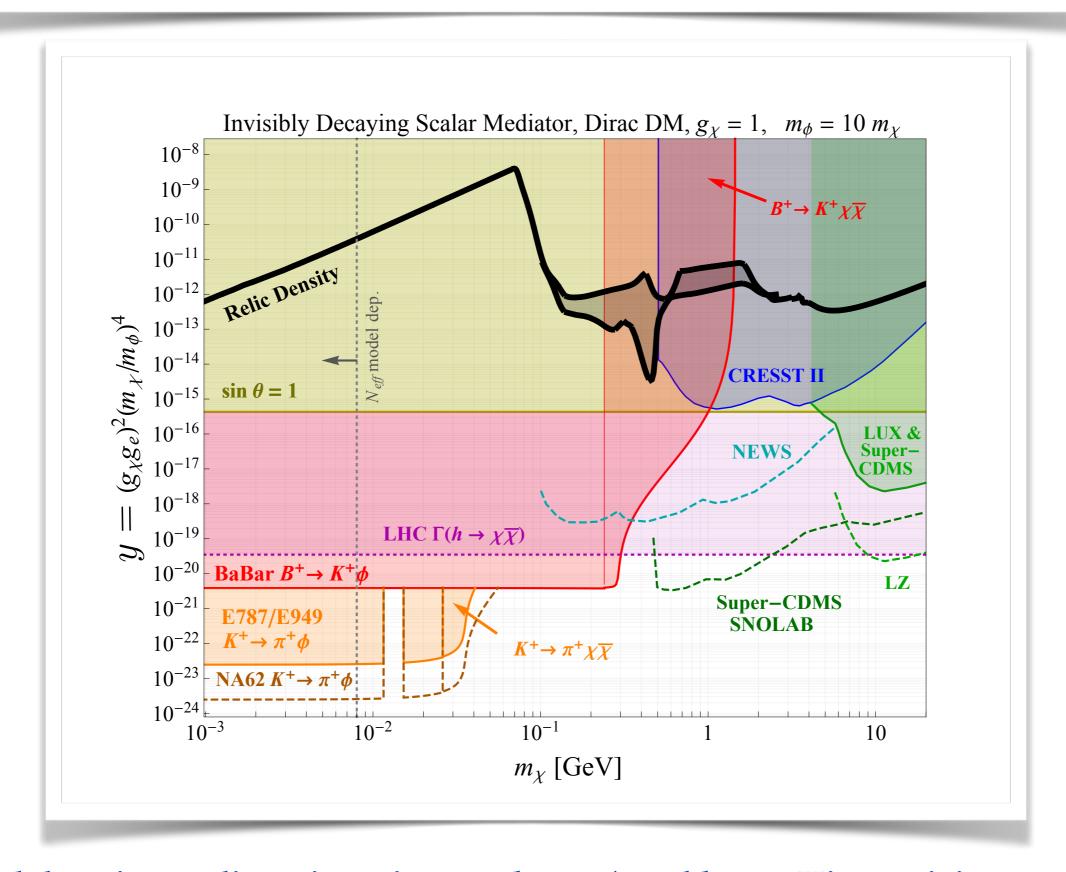
$$X = \frac{\chi}{h}$$
 $N = N$ 

### Direct Annihilation: Ruled Out



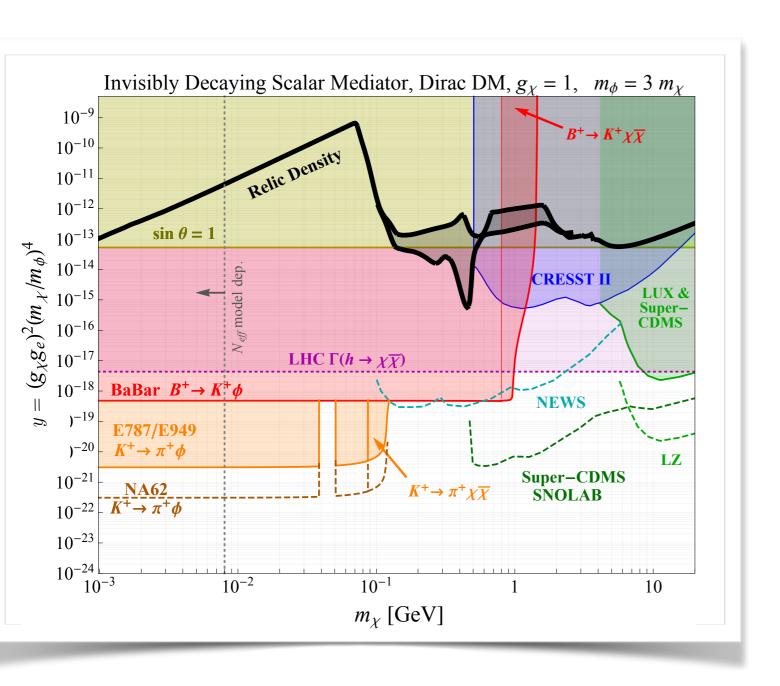
This is the most conservative prescription for all experimental bounds

### Direct Annihilation: Ruled Out



Much heavier mediator is easier to rule out (need larger Higgs mixing angle)

# DM Candidate Variations? (for direct annihilation)



#### Asymmetric DM?

No, annihilation rate bigger

#### Scalar symmetric DM?

No, death by CMB

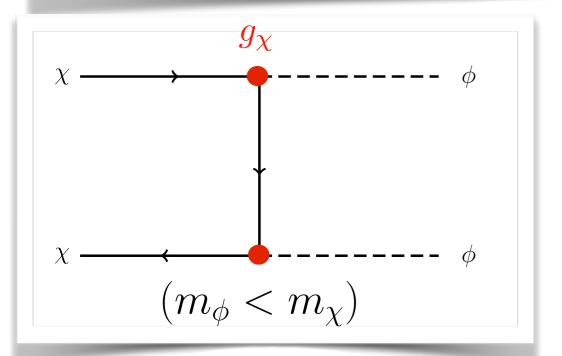
#### **Inelastic couplings?**

No, colliders/mesons don't care

# Motivates vector mediators for direct annihilation

Izaguirre, GK, Schuster, Toro PRL 115 (2015) 1505.00011

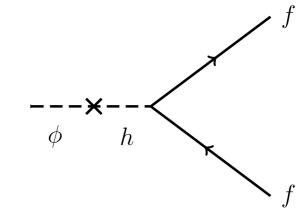
## Secluded Annihilation to Mediators: Thermal Target?



#### Annihilation rate independent of SM

$$\sigma v(\chi \chi \to \phi \phi) = \frac{3g_{\chi}^4 v^2}{128\pi m_{\chi}^2}$$

Mediator decays visibly to SM final states through Higgs portal mixing



Can still produce/observe mediator, but no direct target

So long as annihilation is p-wave DM doesn't matter for bounds

# Next best thing? Minimum mixing for thermal production

#### Assuming Higgs-mediator mixing \*alone\* produces thermal DM

$$\Gamma_{\text{SM}\to\chi\bar{\chi}} = \sum_{f} n_f(T) \langle \sigma | v | \rangle_{f\bar{f}\to\chi\bar{\chi}}$$

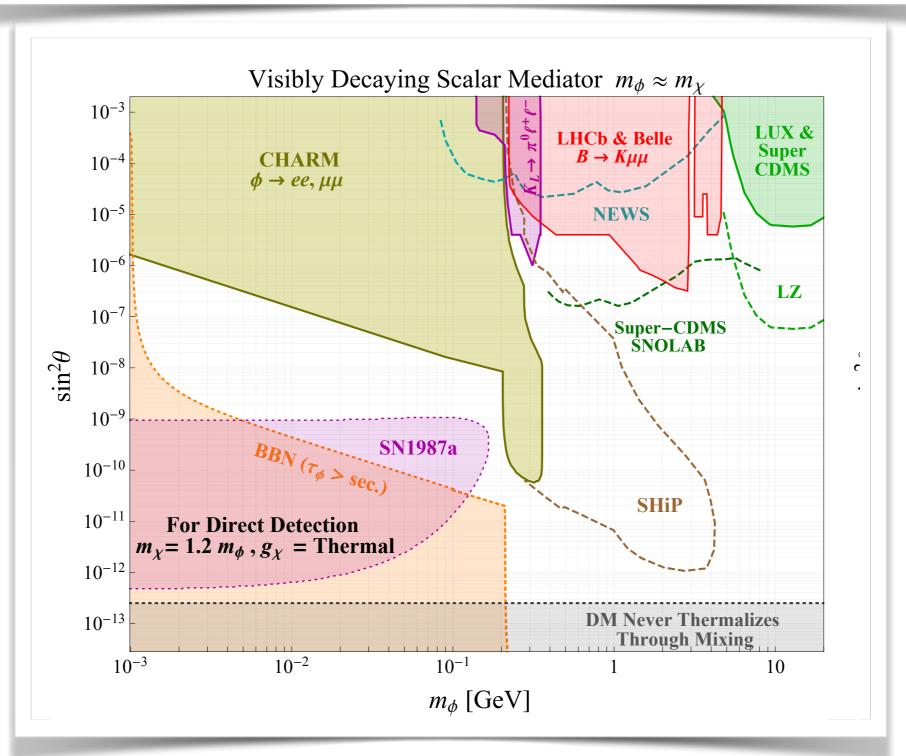
Requiring DM production rate > Hubble in early universe

$$\sin^2 \theta \gtrsim \frac{53\pi^3 \sqrt{g_*(m_t)} m_t}{g_\chi^2 \zeta(3) m_{Pl}} \approx 2 \times 10^{-13} ,$$

Reaching this sensitivity for  $g_{\chi} \sim \mathcal{O}(1)$ 

would rule out thermal DM for the secluded annihilation scenario

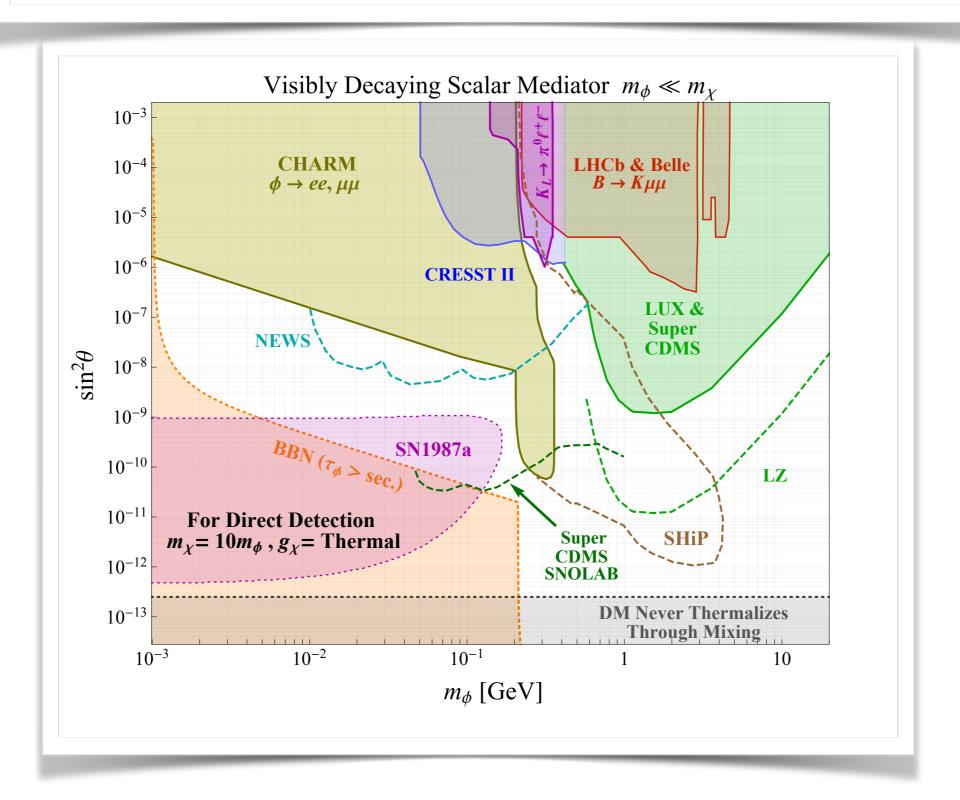
# Mediator mixing bounds (Comparable DM/mediator mass)



beam dump and meson bounds Clarke, Foot, Volkas 1310.8042

Thermal DM can live anywhere unshaded

# Mediator mixing bounds (hierarchical DM/mediator mass)



Larger hierarchy is no longer Light DM (more WIMP-like)

#### **Conclusions**

#### Thermal DM important organizing principle for discovery effort

Viable over MeV-TeV range, need new BSM forces for MeV-GeV Finite, comprehensive list of mediator options for light new forces

#### Higgs portal thermal target already covered for direct annihilation

Direct annihilation scenario ruled out by rare meson and Higgs bounds Independent of DM candidate variations (fermion/scalar/asymmetric/inelastic)

#### "Secluded" thermal annihilation to mediators still viable/testable

Future direct detection (LZ, NEWS, Super-CDMS SNOLAB, Xenon1T...)

Current/proposed hadronic production searches (SHiP, NA62...)

#### Would be interesting to see if other experiments are also sensitive

SeaQuest, DUNE, MiniBooNE, future colliders (FCC,ILC...)?