# Cosmology of Gravitino LSP Scenario with Right-Handed Sneutrino NLSP

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#### 1. Introduction

# Supersymmetric model

- strongly motivated as a beyond standard model
- analyzed in detail from cosmological point of view recently

Supersymmetric model with right-handed (s)neutrinos where neutrino masses are *purely Dirac type* 

#### Right-handed sneutrino $\tilde{\nu}_R$

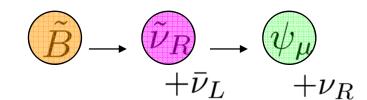
- Never thermalized ( : small neutrino Yukawa coupling)
- Can be relatively light among superparticles
   (: no EW scale corrections for its mass)
- --- Possibility to affect cosmology [T.Asaka,K.I.,T.Moroi(2005)]
- Gravitino-LSP scenario with right-handed sneutrino NLSP changes usual gravitino-LSP scenario

# Gravitino-LSP scenario with right-handed sneutrino NLSP

- bino as NNLSP case  $(m_{\tilde{B}} > m_{\tilde{
u}_R} > m_{3/2})$  -

#### Bino:

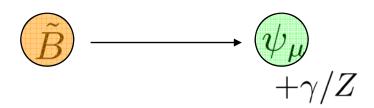
- Freezes out, and decays
- Decay mode  $\,\tilde{B} \to \tilde{\nu}_R \bar{\nu}_L \,$  competes with  $\,\tilde{B} \to \psi_\mu \gamma/Z$  , or dominates in total decay



### Cf.) Without right-handed sneutrino

- $\bullet \tilde{B} \to \psi_{\mu} \gamma/Z$  dominates in total decay to emit more hadrons, which may spoil success of Big-Bang Nucleosynthesis (BBN)
- $\bullet$  Gravitino mass is strictly constrained as  $m_{3/2} \lesssim 0.1~{\rm GeV}$  [ J.L.Feng, S.Su, F.Takayama (2005) ]





# 2. Model

# MSSM with right-handed (s)neutrinos where neutrino masses are *purely Dirac type*

# Superpotential

$$\begin{split} W &= W_{\rm MSSM} + y_{\nu} \hat{H}_{u} \hat{L} \hat{\nu}_{R}^{c} \\ & \longrightarrow y_{\nu} \sin \beta = 3.0 \times 10^{-13} \times \left(\frac{m_{\nu}^{2}}{2.8 \times 10^{-3}~{\rm eV^{2}}}\right)^{1/2} \\ & \left(\begin{array}{c} y_{\nu} : {\rm neutrino~Yukawa~coupling} \\ m_{\nu} : {\rm neutrino~mass} \end{array}\right) \end{split}$$

# Soft SUSY breaking terms

$$\mathcal{L}_{\text{soft}} = -M_{\tilde{L}}^2 \tilde{L}^{\dagger} \tilde{L} - m_{\tilde{\nu}_R}^2 \tilde{\nu}_R^* \tilde{\nu}_R + (A_{\nu} H_u \tilde{L} \tilde{\nu}_R^c + \text{h.c.}) + \cdots$$

# Assumptions:

- Three right-handed sneutrino masses are degenerate
- $A_{
  u}$  is parametrized as

$$A_{\nu} = a_{\nu} y_{\nu} M_{\tilde{L}} \quad \text{where} \quad a_{\nu} \sim \mathcal{O}(1)$$
 
$$\longrightarrow \tilde{\nu}_{L} - \tilde{\nu}_{R} \quad \text{mixing term} \qquad - \tilde{\nu}_{L} - \otimes - \tilde{\nu}_{R} - - \otimes - \tilde{\nu}_{R} - - \otimes - \tilde{\nu}_{R}$$
 
$$\sim m_{\nu} M_{\tilde{L}}$$

# Mass spectrum

$$m_{\tilde{B}} > m_{\tilde{\nu}_R} > m_{3/2}$$

→ gravitino as LSP, right-handed sneutrino as NLSP, and bino as NNLSP(MSSM-LSP)

# 3. Cosmological constraints

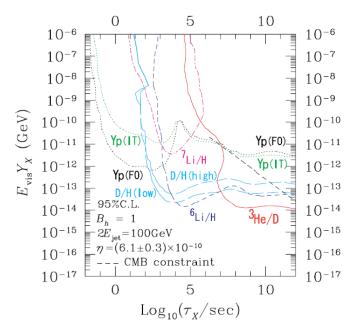
#### BBN constraints [M.Kawasaki, K.Kohri, T.Moroi (2005)]

Upper bound of  $Y_X E_{vis}$  as a function of  $T_X$ 

 $Y_X$ : yield variable of particle X

 $E_{
m vis}$  : mean energy of hadrons emitted

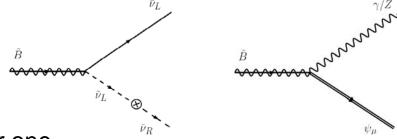
: lifetime of particle X



# Bino decay

 After freezing out, bino decays to right-handed sneutrino or gravitino :

$$\tilde{B} \to \tilde{\nu}_R \bar{\nu}_L$$
, or  $\tilde{B} \to \psi_\mu \gamma/Z$ 



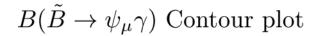
The former mode competes with the latter one or dominates in total decay

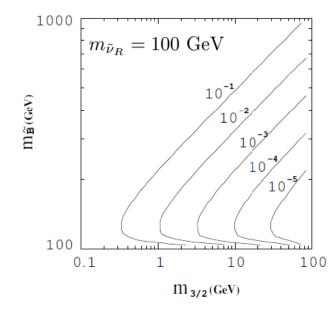
• Bino is long-lived :  $\tau_{\tilde{B}} \sim 10^2 \ {\rm sec}$ 

Cf.) Without  $\,\tilde{\nu}_R\,$ ,  $\,\tilde{B} \to \psi_\mu \gamma/Z\,$  dominates in total decay to emit many hadrons and  $m_{3/2}$  is strictly constrained from BBN

Although bino is long-lived, less hadrons are emitted in its decay compared to no  $\tilde{\nu}_R$  case

Thus constraints from BBN is relaxed





# No hadrons are produced in two-body decays In order to calculate $B_{\mathrm{had}}$ and $E_{\mathrm{vis}}$ , we consider

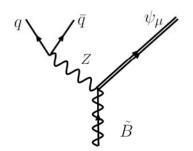
• Three- or four- body decays to produce hadrons :

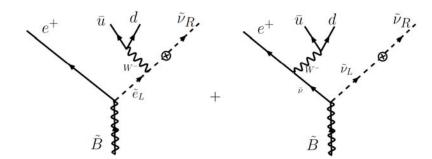
$$\tilde{B} \to \psi_{\mu} q \bar{q}$$

$$\tilde{B} \to \tilde{\nu}_R e_L^+ q \bar{q}'$$

$$\tilde{B} \to \tilde{\nu}_R \bar{\nu}_L q \bar{q}$$

(and CP conjugate final states for last two)





# Results (from BBN constraints)

# Constraints from BBN scenario is drastically relaxed

$$\rightarrow m_{3/2} < 100 \text{ GeV}$$

For a new allowed region,

ullet Upper bound of  $\,m_{ ilde{B}}\,$  (given by BBN constraints)

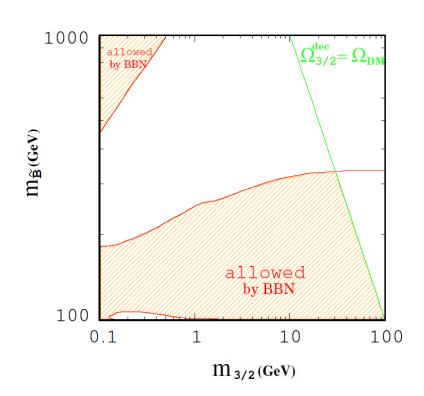
Mainly comes from four-body decay modes,

$$\tilde{B} \to \tilde{\nu}_R e_L^+ q \bar{q}', \tilde{B} \to \tilde{\nu}_R \bar{\nu}_L q \bar{q}$$

Large  $m_{\tilde{B}}$  region is excluded because

 $B_{
m had}, E_{
m vis}$  are enhanced

•  $\Omega_{3/2}^{
m dec} < \Omega_{
m DM}$  (not to overclose universe)



$$m_{\tilde{\nu}_R} = 100 \text{ GeV}$$

#### Constraints from structure formation

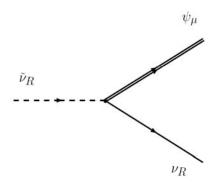
# Right-handed sneutrino decay

• Right-handed sneutrino is long-lived :

$$\tau_{\tilde{\nu}_R} \sim 10^{5-8} \text{ sec}$$

• Emitted gravitino is relativistic :

$$\lambda_{\rm FS} \sim 6 \; {\rm Mpc}$$



Emitted gravitino acts as warm dark matter

On the other hand, gravitino is also produced by thermal scattering processes and acts as cold dark matter

We consider constraints on WDM+CDM scenario from structure formation

#### WDM+CDM scenario

$$ho_{3/2}=
ho_{3/2}^{
m dec}+
ho_{3/2}^{
m th}$$
  $ho_{3/2}^{
m dec}$  ,  $ho_{3/2}^{
m th}$  ; energy density of

 $\left[ \begin{array}{ccc} \rho_{3/2}^{\rm dec}, \ \rho_{3/2}^{\rm th} \ : \ {\rm energy\ density\ of\ gravitino\ produced\ by\ decays,\ or} \\ & {\rm by\ thermal\ scattering\ respectively} \end{array} \right]$ 

A step-like decrease in power spectrum where

$$k \sim 2\pi/\lambda_{\rm FS}$$

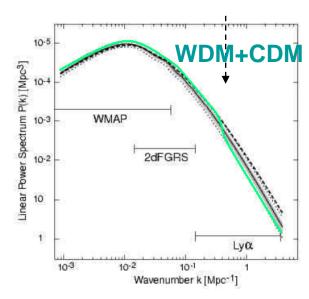
( \*.\* Only power spectrum of WDM component dumps )

In a condition that power spectrum is in 95% C.L. region of observation data,

$$f \lesssim 0.4$$
  $(f \equiv \rho_{3/2}^{
m dec}/\rho_{3/2})$   $\longrightarrow \Omega_{3/2}^{
m dec} \lesssim 0.4\Omega_{
m DM}$ 

$$\rightarrow \Omega_{3/2}^{\rm dec} \lesssim 0.4\Omega_{\rm DM}$$

$$\sim 2\pi/\lambda_{\rm FS}$$



[ D.N.Spergel (2003) ]

#### Results

(including constraints from structure formation)

Parameter region allowed by BBN is cut a little bit as

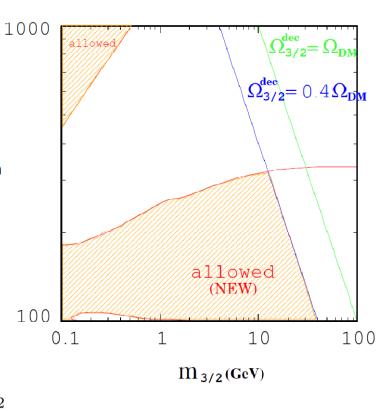
$$m_{3/2} \lesssim 40 \text{ GeV}$$

For a new allowed region,

•  $\Omega_{3/2}^{
m dec} < 0.4 \Omega_{
m DM}$  (from structure formation)

$$\left( \begin{array}{c} \Omega_{3/2}^{\rm dec} = \frac{m_{3/2}}{m_{\tilde{B}}} \Omega_{\tilde{B}} \ \ \text{with} \ \ \Omega_{\tilde{B}} h^2 \sim 0.1 \left[ \frac{m_{\tilde{B}}}{100 \ {\rm GeV}} \right]^2 \\ \Omega_{\rm DM} h^2 \simeq 0.105 \ \ [$$
 D. N. Spergel (2006) ]

$$\rightarrow m_{\tilde{R}} m_{3/2} < 0.4 \times 10^4 \text{ GeV}^2$$



$$m_{\tilde{\nu}_R} = 100 \text{ GeV}$$

### 4. Conclusion

#### Model:

MSSM with right-handed (s)neutrinos where neutrino masses are *purely Dirac type* 

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- bino as NNLSP case  $(m_{\tilde{B}} > m_{\tilde{\nu}_R} > m_{3/2})$  -

# Cosmological constraints

- Bino decay mode  $\tilde{B} \to \tilde{\nu}_R \bar{\nu}_L$  competes  $\tilde{B} \to \psi_\mu \gamma/Z$  or dominates in total decay Less hadrons are emitted than in no right-handed sneutrino case
- Constraints from BBN is drastically relaxed as

$$m_{3/2} \lesssim 40~{
m GeV}$$
 (Cf.)  $m_{3/2} \lesssim 0.1~{
m GeV}$  without  $ilde{
u}_R$ 

ullet Upper bound of  $\,m_{3/2}\,$  is given by constraints from structure formation