# Local SUSY-breaking minima in $N_f=N_c$ SQCD?

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Andrey Katz

Work done with Yael Shadmi and Tomer Volansky

Technion-Israel Institute of Technology

### **Outline**

- 1. ISS review and  $N_f = N_c$  conjecture
- 2. Another deformation, saddle point
- 3. Phenomenological consequences
- 4. Conclusions

#### ISS review

#### Itriligator, Seiberg, Shih, 2006

- framework  $SU(N_c)$  SQCD;  $N_c+1\leq N_f<\frac{3}{2}N_c$ ,  $W_{\rm tree}=(m_Q)_{ij}\bar{Q}_iQ_j$  and  ${\rm rank}[m_Q]>N_c$
- 6 SUSY-breaking local minimum near the origin
- magnetic dual: SUSY is broken by rank condition
- possesses SUSY vacuum

#### Demand $m_Q \ll \Lambda$ to get

- calculablity
- SUSY vacuum far from the orgin ⇒ metastability

## $N_f = N_c$ conjecture

### **ISS** approach:

- 6 take  $N_f = N_c + 1$ , local SUSY-breaking minimum exists
- 6 consider the limit  $(m_Q)_{N_f,N_f} o \infty$
- 6 conjecture:  $N_f = N_c$  has a similar vacuum,

## $N_f = N_c$ conjecture

#### **ISS** approach:

- 6 take  $N_f = N_c + 1$ , local SUSY-breaking minimum exists
- 6 consider the limit  $(m_Q)_{N_f,N_f} \to \infty$
- 6 conjecture:  $N_f = N_c$  has a similar vacuum,

## but $N_f = N_c$ is very different

- 6  $N_F > N_c$ :  $\frac{m_Q}{\Lambda} \ll 1 \Rightarrow$  non-calculable terms (Kähler) are under control
- 6  $N_f = N_c$ :  $m_Q/\Lambda$  small, but Kähler is not under control

# Why is $N_f = N_c$ so different? Kähler potential

#### Kähler metric

$$g_{MM^{\dagger}}^{-1} \sim \frac{\operatorname{Tr}M^{\dagger}M}{\Lambda^{2}} + \frac{\operatorname{Tr}M\operatorname{Tr}M^{\dagger}}{\Lambda^{2}} + \frac{(B_{+} + B_{+}^{\dagger})^{2}}{\Lambda^{2}} + \dots$$

#### Potential:

$$V = g_{MM^{\dagger}}^{-1} |F_M|^2$$
$$F_M = m_Q \Lambda.$$

$$V \sim m_Q^2 (\text{Tr}MM^{\dagger} + \text{Tr}M\text{Tr}M^{\dagger} + (B_+ + B_+^{\dagger})^2)$$

Contribution stays finite at  $N_f = N_c$ 

# Limit $N_f = N_c$ : calculabe versus uncalculable

Consider  $N_f = N_c + 1$  with one very heavy pair of quarks

- $(m_Q)_{N_c+1}$  the heaviest mass
- $\hat{\Lambda}$  confining scale of  $N_f=N_c+1$

 $(m_Q)_{N_c+1} \to \infty$  we approach  $N_f = N_c$  limit

- 6 calculable, tree  $m^2 \sim \frac{m_Q \hat{\Lambda}}{m_{N_c+1}} \to 0$
- 6 pseudo-moduli  $m^2 \sim \frac{1}{16\pi^2} \frac{m_Q \hat{\Lambda}}{m_{N_C+1}} \to 0$
- ${\rm ^6}$  uncalculable  $m_{\rm uncalc}^2 \sim m_Q^2$  finite
- ullet we do not even know the sign of  $m^2$

#### Another deformation - ITIY

Itriligator, Thomas, 1996; Izawa, Yanagida, 1996 Try another deformation

- 6 does the extremum survive?
- 6 is it still a minimum?
- 6 where is calculability lost?

Add singlets. Under  $SU(N_f)_L \times SU(N_f)_R$ :

- 6  $S_{ij}$   $(\bar{\mathbf{N}}, \mathbf{N})$
- 6 T (1,1),  $\bar{T}$  (1,1)

## Low-energy Superpotential

$$W = \mathcal{A}(\det M - B\bar{B} - \Lambda^{2N}) \longrightarrow +$$

$$\lambda \text{Tr}(SM) + \kappa (TB + \bar{T}\bar{B})$$

$$m_Q \text{Tr} M$$
  $+$   $\frac{m_S}{2} S^2 + \frac{m_T}{2} (T^2 + \bar{T}^2)$ 

moduli space deformation

\*ISS mass-term

coupling

 $singlet\ masses$  -  $m_Q$  has no effect w/o them

## ISS $N_f = N_c$ limit

#### Decoupling limit:

$$\frac{\lambda^2}{m_S} \to 0; \quad \frac{\kappa^2}{m_T} \to 0$$

#### SUSY-breaking solution should:

- 6  $M \to 0$ ,  $B_+ \to 0$ ,  $B_- \to \Lambda$  at the decoupling limit
- $_{6}$   $F_{M} \propto m_{Q}$

#### SUSY solution:

- 6 decoupling limit finite distance from the origin
- sufficiently far from SUSY-breaking solution

# Non-SUSY solution and decoupling

- look for solution along baryonic branch
- $^{6}$  take  $(m_Q)_{ij} \propto \delta_{ij}$

Non-SUSY solution near the origin:

$$M \sim \left(\frac{\lambda^2}{m_S} \frac{m_T}{\kappa^2}\right)^{\frac{1}{N-2}}$$

**Decoupling:** 

$$M \to 0 \qquad \Longrightarrow \qquad \frac{\lambda^2/m_S}{\kappa^2/m_T} \to 0$$

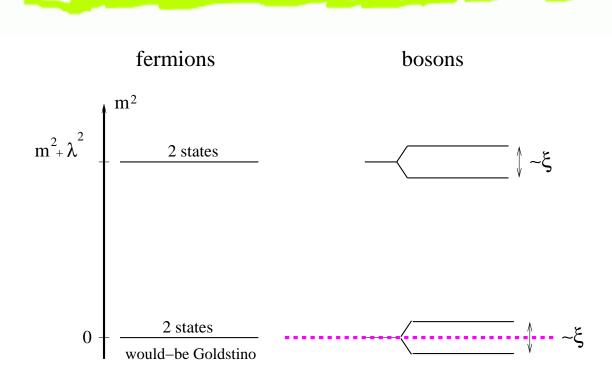
 $S_{ij}$  decouples faster than  $T, \bar{T}$ 

## Spectrum at SUSY-breaking point

- ${f 6}$  spectrum for B and T supersymmetric
- 6 ISS: mesino is Goldstino
- expect: Goldstino reduces to mesino in decoupling limit
- ullet before: admixture of M and S
- 6 one SUSY-breaking parameter:  $\xi$

$$\xi \propto \frac{\lambda^2}{m_S} m_Q^*$$

### Saddle point

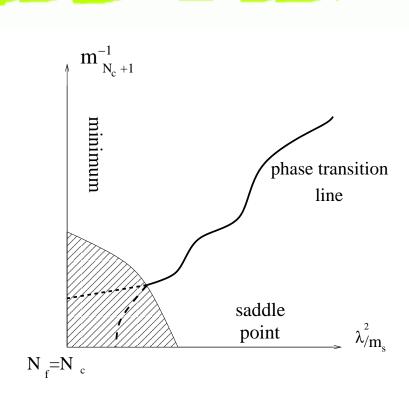


Always one state below zero - instability.

## Range of validity

- 6 wanted: calculable contributions ≫ uncalculable
- 6 the lowest state  $m^2 \sim \pm |\xi|$
- 6 demand  $\frac{(\lambda\Lambda)^2 m_Q}{m_S} \gg m_Q^2$  calculable uncalculable
- one can choose e.g.  $m_Q \ll (\lambda \Lambda) \ll m_S \lesssim \Lambda$

## $N_f = N_c$ conjecture overview



- ISS points undergo phase transition
- shaded region: ISS point is governed by non-calculable contributions from Kähler

## Phenomenological consequences - Pentagon

#### Banks, 2006, "Pentagon" model

- 6  $N_f=N_c=5$  with diagonal ISS mass  $\Delta W=m_Q{
  m Tr}M$
- use ISS-conjectered minimum
- 6 flavor symmetry  $SU(5)_{
  m diag}$
- embed the SM into the flavor symmetry
- 6 μ-problem: need the singlet  $S: \Delta W = SH_uH_d$
- $^{\rm 6}~SU(3)\times SU(2)\times U(1)$  unbroken S couples to the quarks through  $Y_{ij}$
- ${f 6}$  messengers off diagonal components of M

## Spectrum of Pentagon

Metastability  $m_Q \ll \Lambda_5$ 

Consider first  $\lambda \ll 1$  to avoid destabilization

ISS minimum? Answer in  $\Delta K$ 

let's believe ISS conjecture

Statement: weakly coupled messengers - STr[mess] > 0

Poppitz and Trivedi, 1997:

large **negative** contributions to squarks  $m^2$ 

# Wrong-sign contributions to squark masses

- 6 small  $\lambda$  tachyonic squarks,  $SU(3)_C$  is broken
- 6  $m^2[\text{squark}] \propto \log(\Lambda_5/(m_F))$  $\sim \lambda^{\#} \sqrt{m_Q \Lambda_5}$
- $\delta$  large back to ITIY-like, no stable minimum
- 6 minimum may exist in intermediate range not ISS-conjectured minimum!
- $\lambda$  is large or small Pentagon is ruled out. Intermediate  $\lambda$  we do not know. Unlikely to have viable minimum.

#### **Conclusions**

- there is no clear indication that the meta-stable SUSY-breaking vacuum exists in  $N_f=N_c$  SQCD
- o no information can be gained by deforming the theory
- minimum of one deformation saddle point in another
- 6 coupled singlets the instability may be generic
- Pentagon the coupling to singlet can not be too large or too small
- of if the conjectured minimum of Pentagon exists:
  - it's uncalculable
  - it's not ISS minimum