

# SUSY tools and Predictions

by

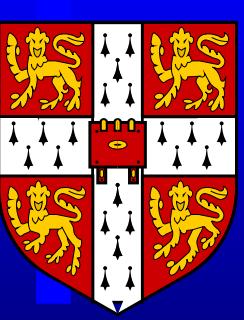
Ben Allanach (University of Cambridge)



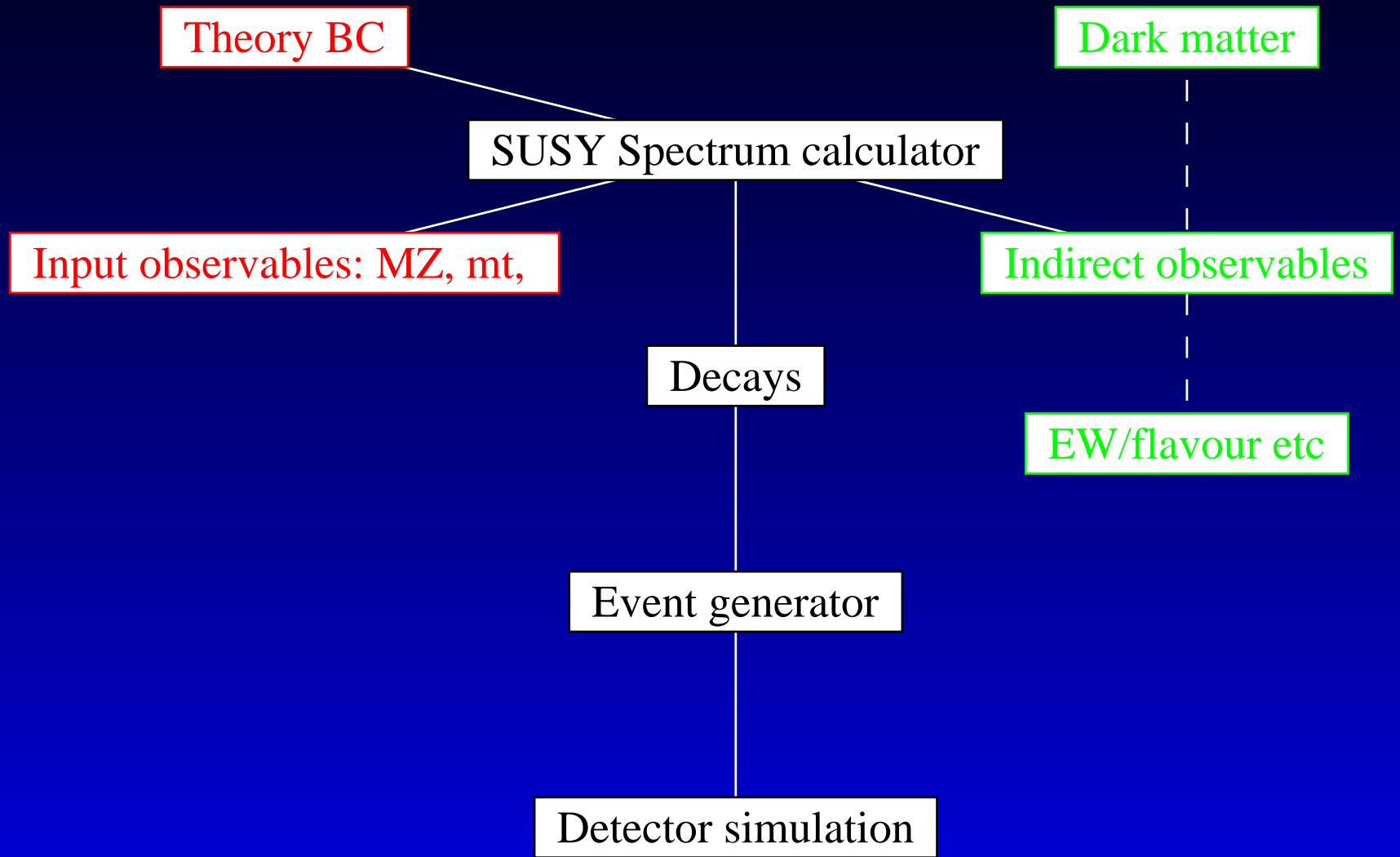
## Talk outline

- Bestiary of **public** codes only: supposedly **impartial**
- Apology: no author lists
- Predictions for the LHC: **partial**

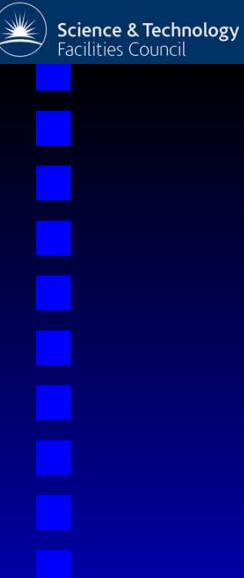


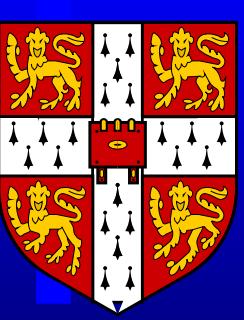


# MSSM Tools



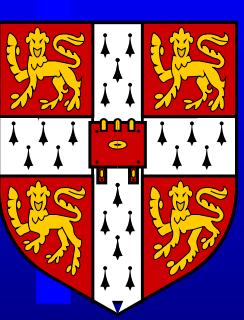
SLHA: Skands *et al*, JHEP 0407 (2004) 036, SLHA2 on its way  
(NMSSM, RPV, FV, CPV)





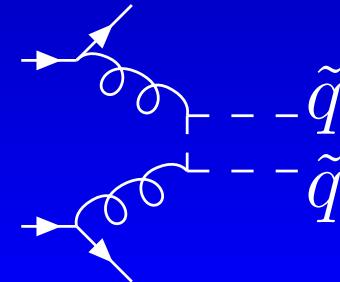
# Spectrum and decays

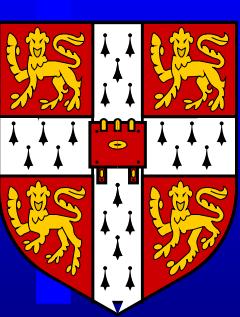
- **ISASUSY** decouples particles at the mass thresholds but misses some finite terms in the matching: re-sums log splittings.
- **SOFTSUSY**, **sPHENO**, **SUSPECT** all catch the finite terms but do the splittings to leading log in RPC-MSSM.
- **CPsuperH**, **FeynHiggs** do Higgs mass spectrum and decays with of CP violating MSSM
- **NMSPEC** does the **CNMSSM** spectrum, **NMHDECAY** gives the decays widths etc
- **PYTHIA**, **ISASUSY**, **sPHENO** and **SusyHIT** do decays of Higgs and SUSY particles in MSSM.



# Matrix Element Generators

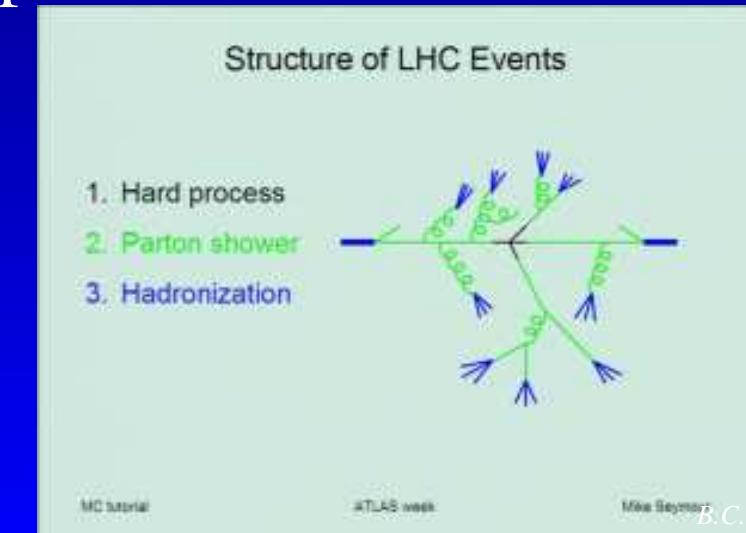
- Additional hard jets *cannot* be modelled reliably using the parton shower - you need to simulate the matrix element.
- **SMADGRAPH**, **compHEP**, **calcHEP**, **GRACE** do SUSY and more general models at tree level. 2 to 4 possible. **BRIDGE** can be used to remember spin information in the decays.
- **WHIZARD**, **SUSYGEN** - polarisation included for  $e^+e^-$
- **PROSPINO** does NLO-QCD sparticle production

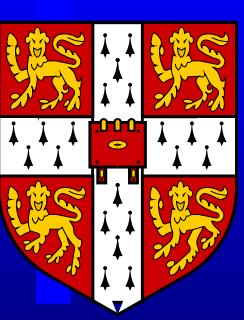




# Event Generation

- Can pass matrix-element generated events to event generators with the (original) *Les Houches Accord*
- **PYTHIA** used extensively. Includes RPV. phase-space decays. **ISAJET** too.
- **HERWIG** maintains spin info down cascade decays. RPV too.
- **SHERPA** matches up ME with more standard event generation.
- Shift toward C++

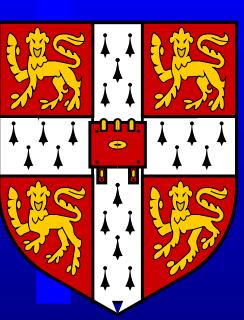




# SUSY Prediction of $\Omega h^2$

- Assume relic in thermal equilibrium with  $n_{eq} \propto (MT)^{3/2} \exp(-M/T)$ .
- Freeze-out with  $T_f \sim M_f/25$  once **interaction rate < expansion rate** ( $t_{eq}$  critical)
- **microMEGAs** uses **calcHEP** to automatically calculate relevant Feynman diagrams for some given model Lagrangian: *flexible*.
- **darkSUSY** has MSSM annihilation channels hard-coded. Much work on (in)-*direct* detection possibilities.





# Implementation

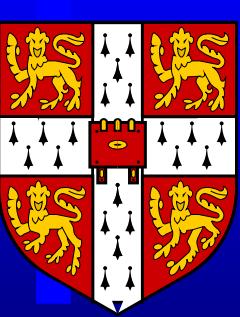
Input parameters are:  $m_0$ ,  $A_0$ ,  $M_{1/2}$ ,  $\tan \beta$ ,

- $m_t = 171.4 \pm 2.9$ ,  $m_b(m_b) = 4.24 \pm 0.11$  GeV,
- $\alpha_s(M_Z)^{\overline{MS}} = 0.1176 \pm 0.002$ ,
- $\alpha^{-1}(M_Z)^{\overline{MS}} = 127.918 \pm 0.018$

For the likelihood, we also use

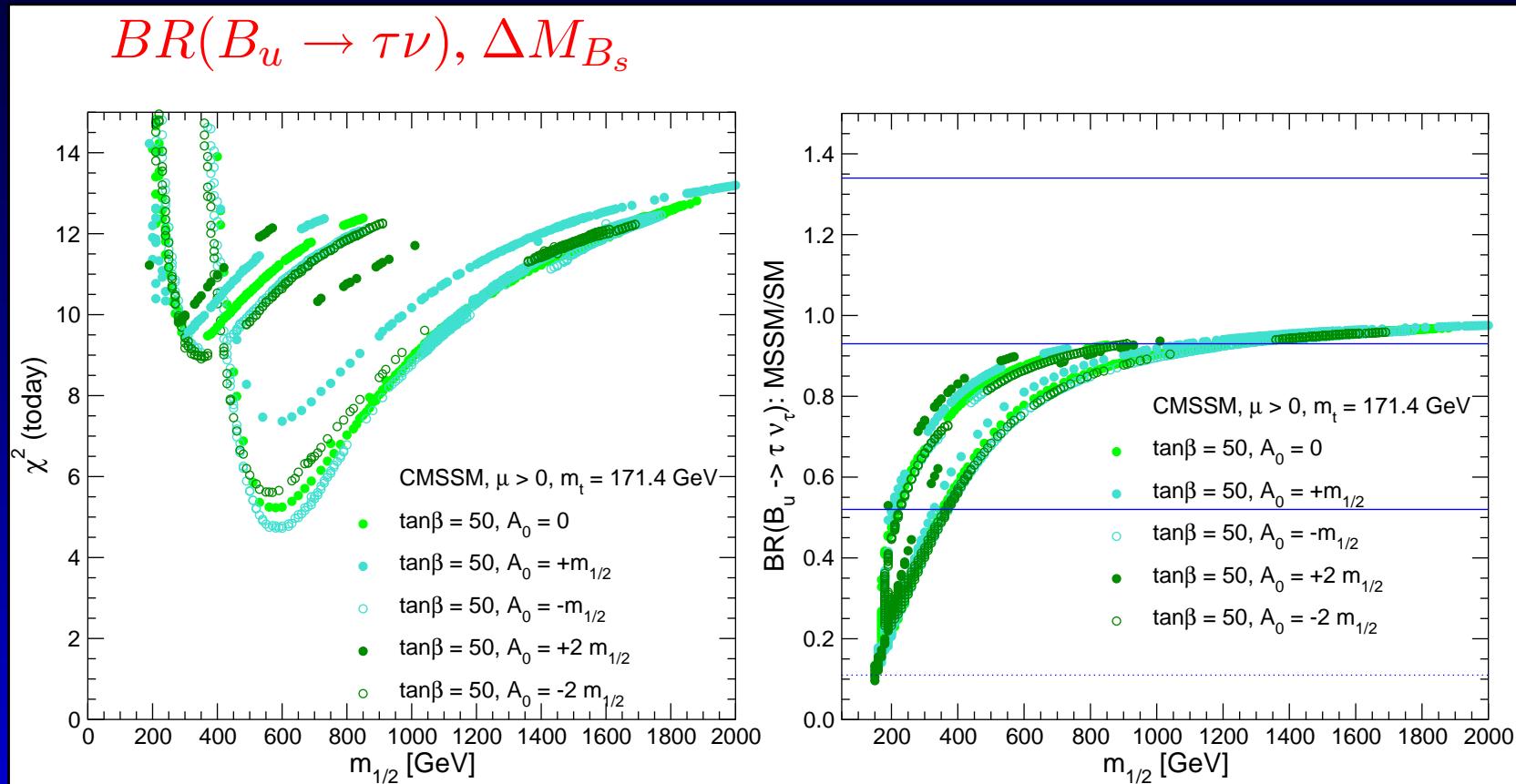
- $\Omega_{DM} h^2 = 0.104^{+0.0073}_{-0.0128}$  micrOMEGAs, darkSUSY
- $\delta(g - 2)_\mu / 2 = (22 \pm 10) \times 10^{-10}$  Stöckinger *et al*
- $BR[b \rightarrow s\gamma] = (3.55 \pm 0.38) \times 10^{-5}$  Misiak *et al*
- $\sin^2 \theta_w^l(\text{eff}) = 0.23153 \pm 0.000175$
- $M_W = 80.392 \pm 0.031$  GeV W Hollik, A Weber *et al*

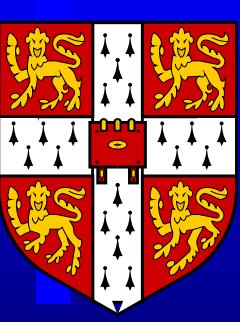




# *b* Observables

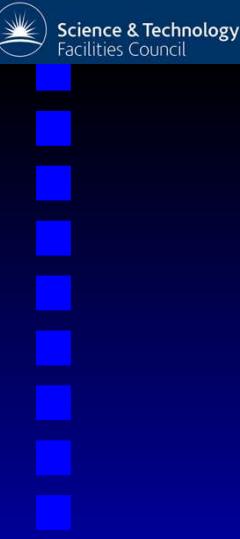
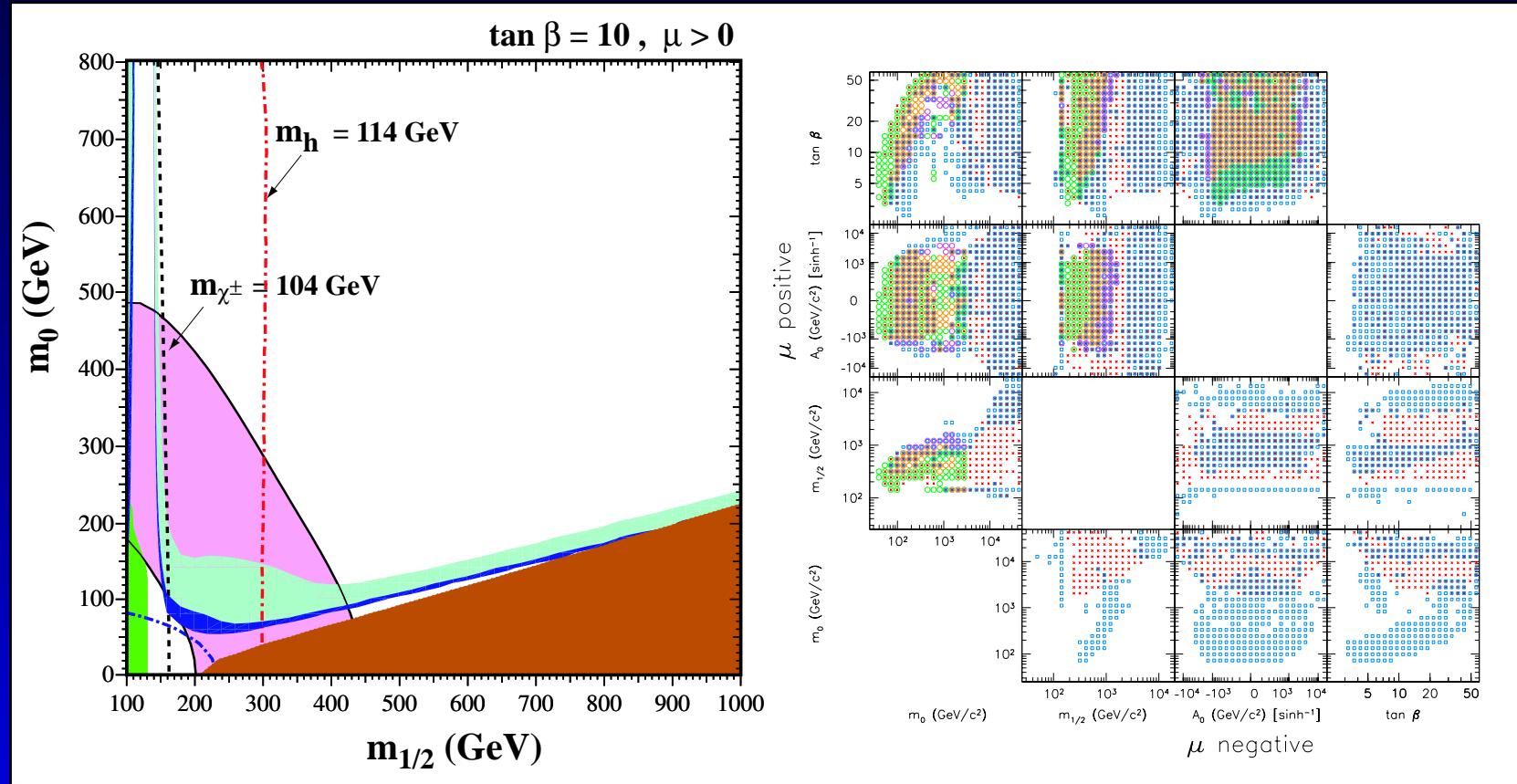
CMSSM: Ellis, Heinemeyer, Olive, Weber, Weiglein, arXiv:0706.0652

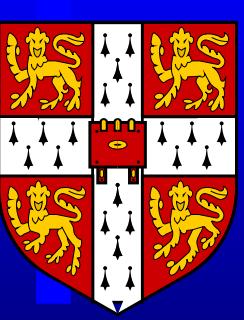




# Constraints on SUSY Models

CMSSM well-studied in literature: eg Ellis, Olive *et al* PLB565 (2003) 176; Roszkowski *et al* JHEP 0108 (2001) 024; Baltz, Gondolo, JHEP 0410 (2004) 052; ...





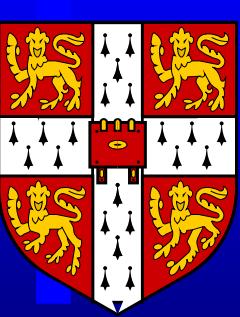
# Fit Development

- Typically done 2d scans with  $2\sigma$  exclusion regions, but in general we have  $\alpha(M_Z)$ ,  $\alpha_s(M_Z)$ ,  $m_t$ ,  $m_b$ ,  $m_0$ ,  $M_{1/2}$ ,  $A_0$ ,  $\tan \beta$  to vary
- Effective 3d type scan done <sup>a</sup> which parameterises a 2d surface of central  $\Omega h^2$
- 4d scan <sup>b</sup> used the impressive *Markov Chain Monte Carlo technique* like in cosmology.
- Combine likelihoods from all of the different measurements.

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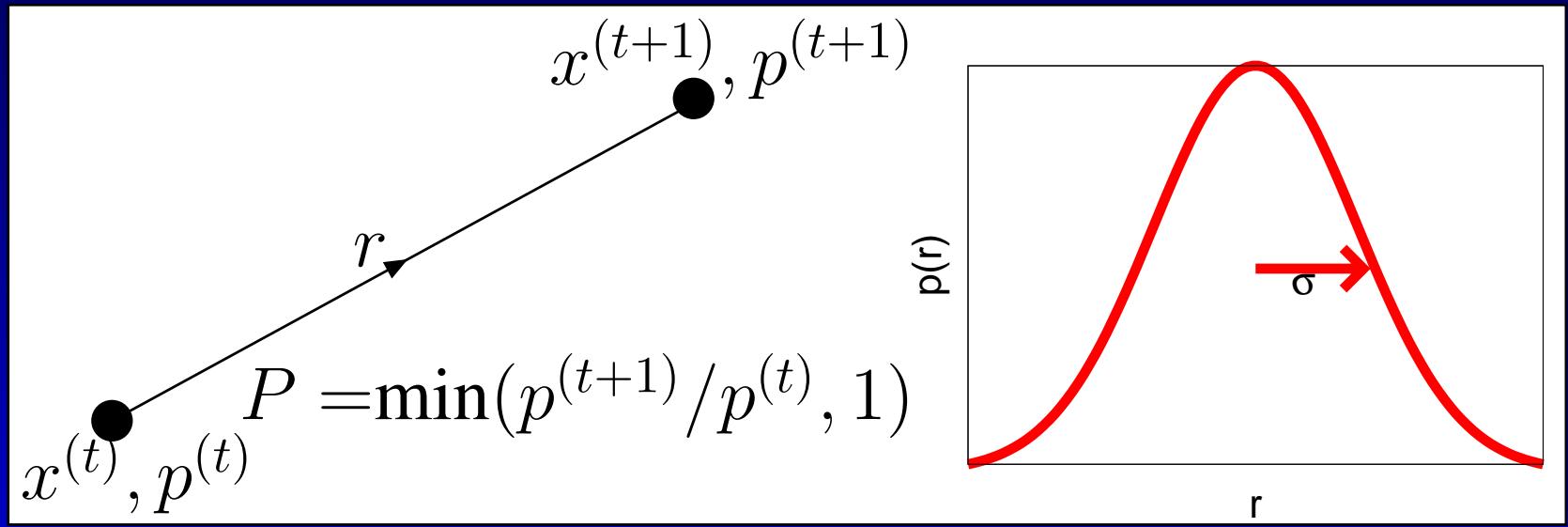
<sup>a</sup>Ellis *et al*, arXiv:0706.0652

<sup>b</sup>Baltz, Gondolo, JHEP 0410 (2004) 052

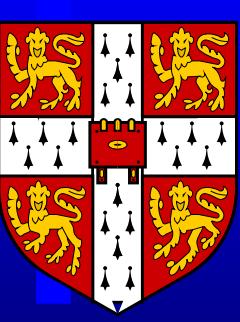


# Markov-Chain Monte Carlo

Metropolis-Hastings Markov chain sampling consists of list of parameter points  $x^{(t)}$  and associated posterior probabilities  $p^{(t)}$ .

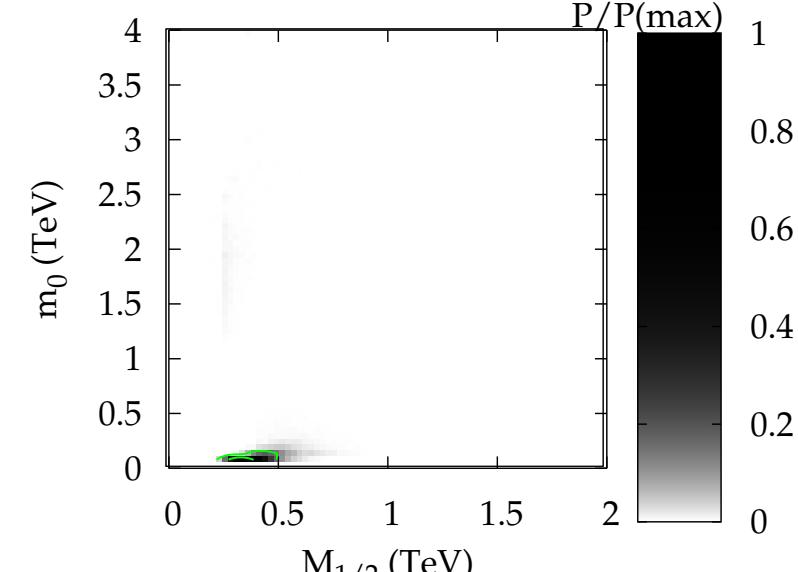
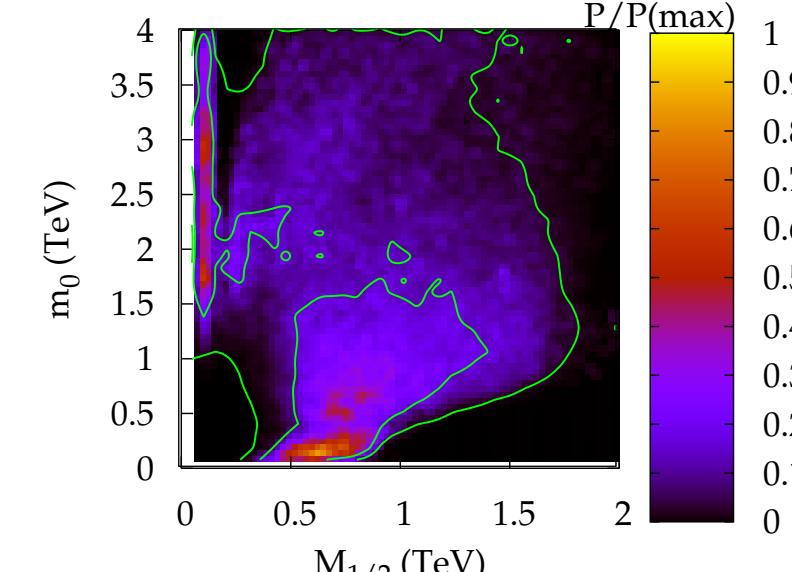


Final density of  $x$  points  $\propto p$ . Required number of points relatively *insensitive* to number of dimensions.

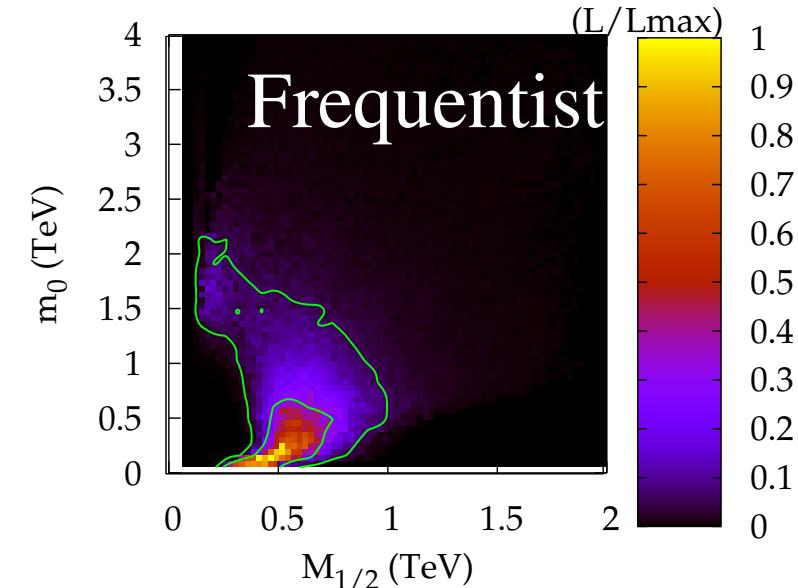
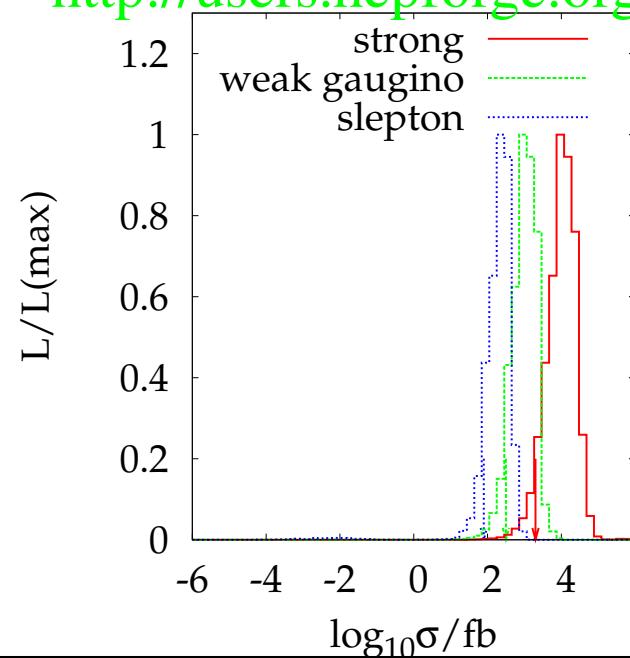


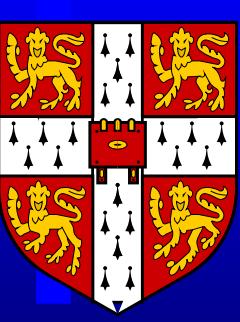
# Killer Inference for Susy METeorology

BCA, Cranmer, Weber, Lester, arXiv:0705.0487



<http://users.hepforge.org/~allanach/benchmarks/kismet.html>



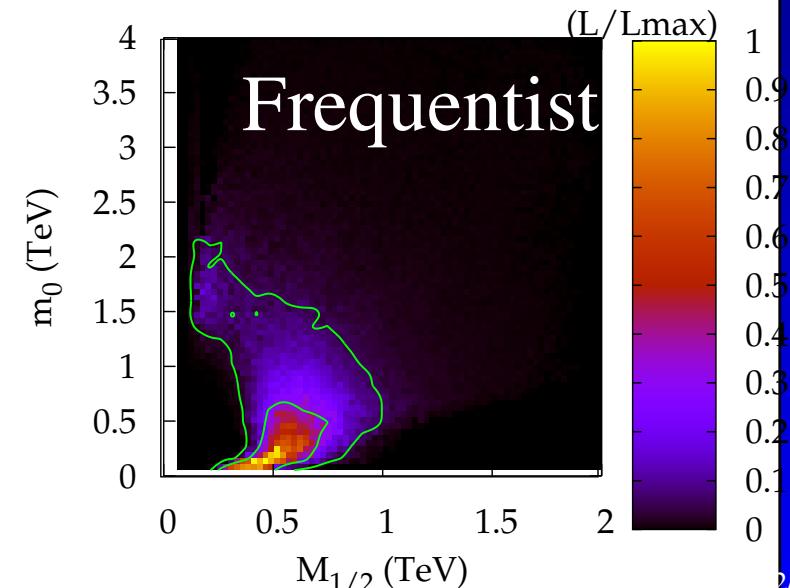
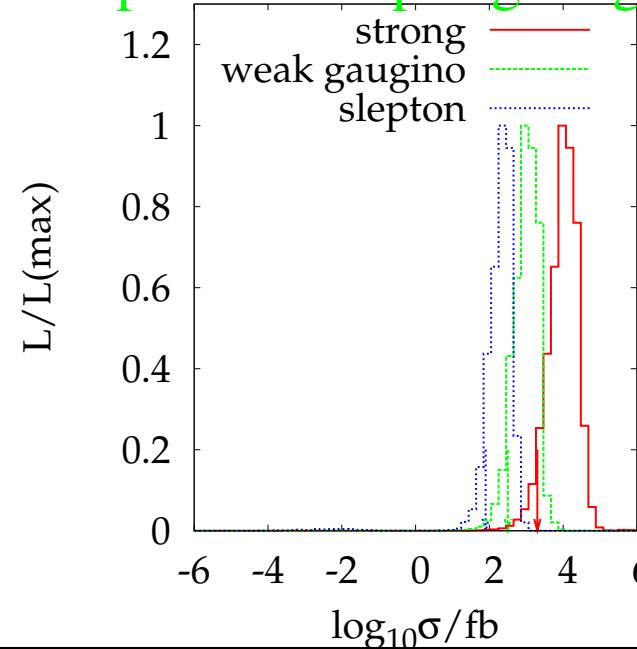


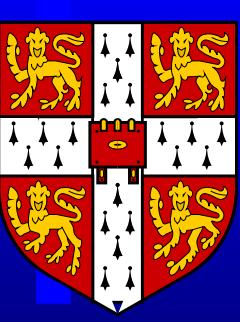
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# Higgs Meteorology

BCA, Cranmer, Lester, Weber arXiv:0705:0487

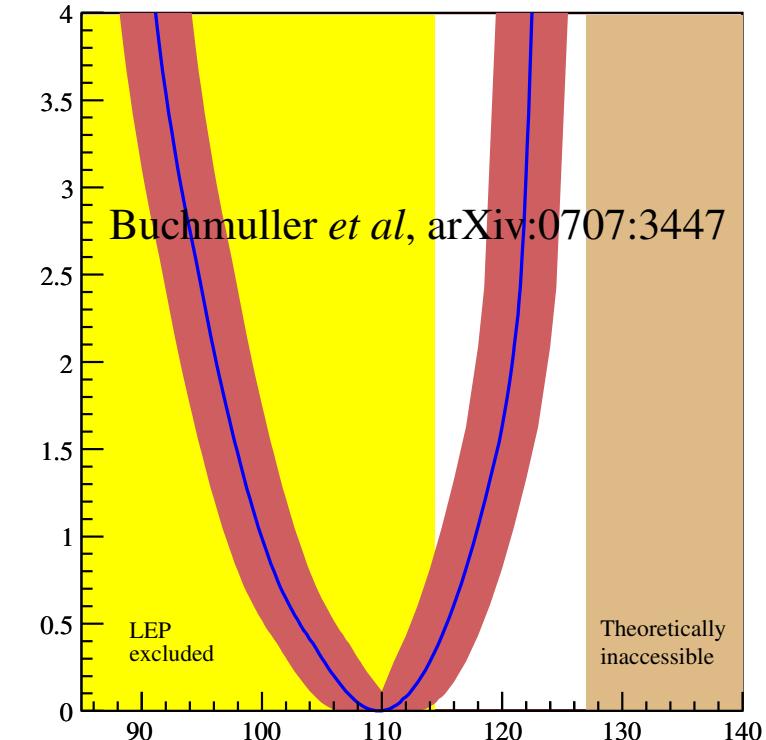
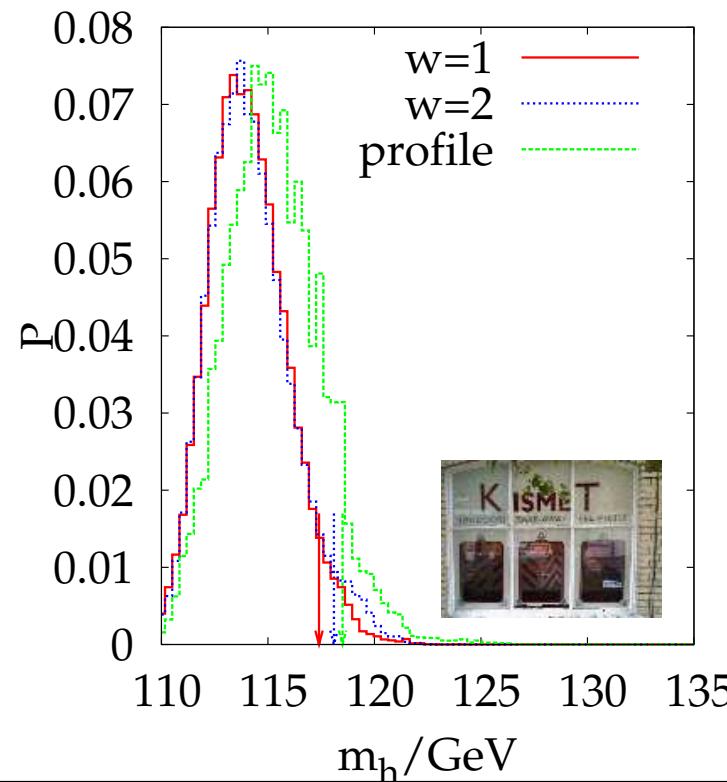
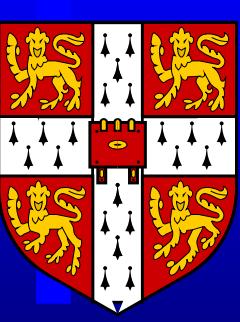
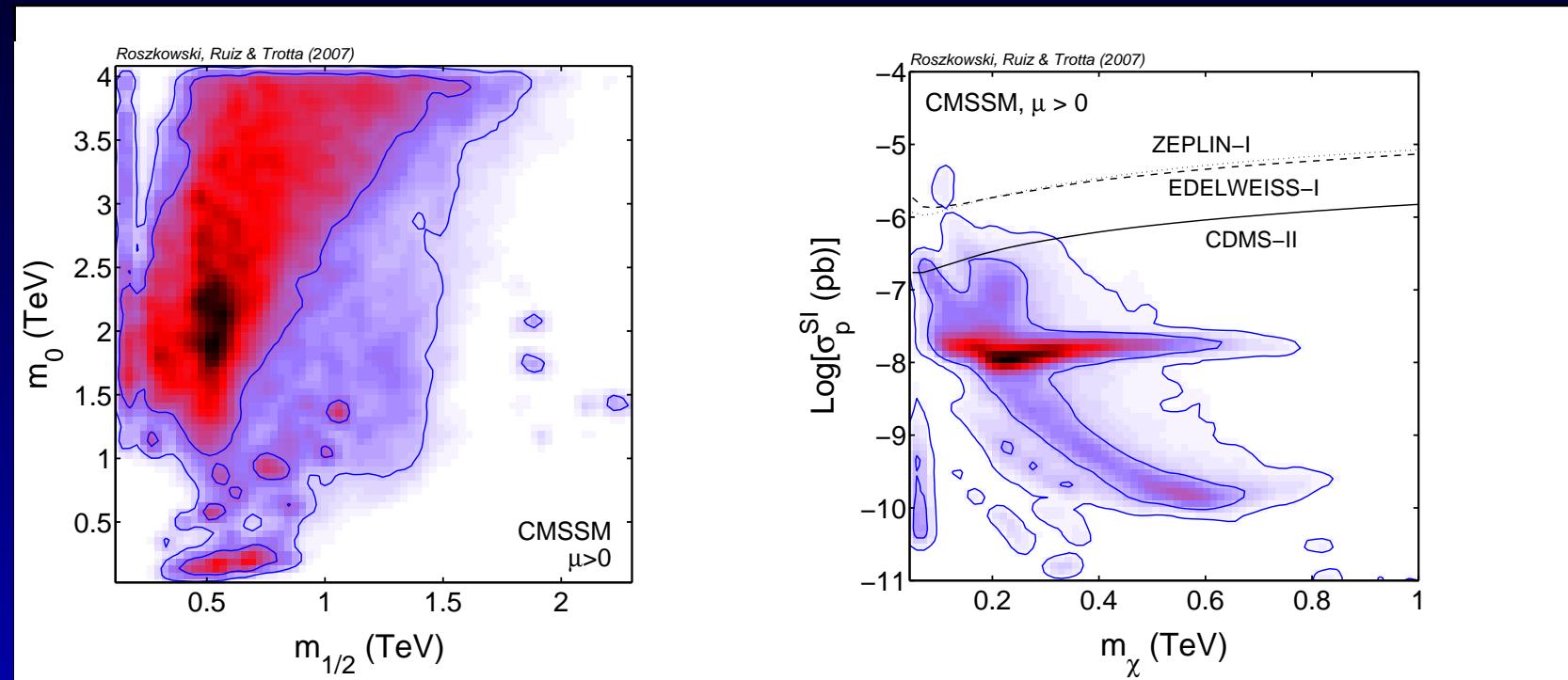


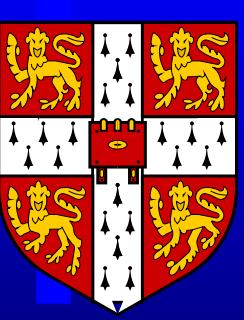
Figure 0: Including (LHS) or *not* including (RHS) the LEP2 direct Higgs mass constraints on the CMSSM.



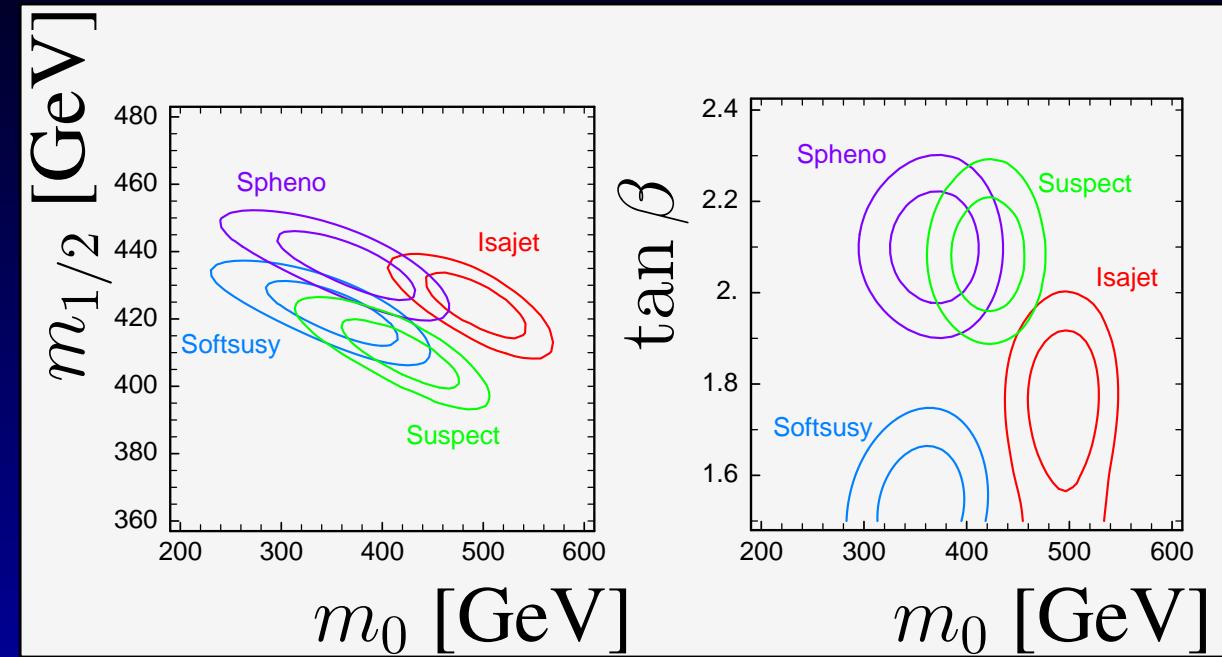
# Other literature



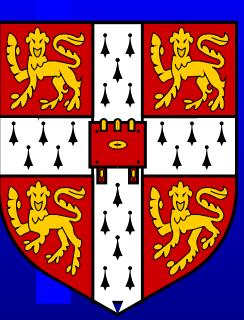
R. R. de Austri, R. Trotta and L. Roszkowski,  
arXiv:0705.2012, including some NNLO  $b \rightarrow s\gamma$   
pieces. **susyBayes**



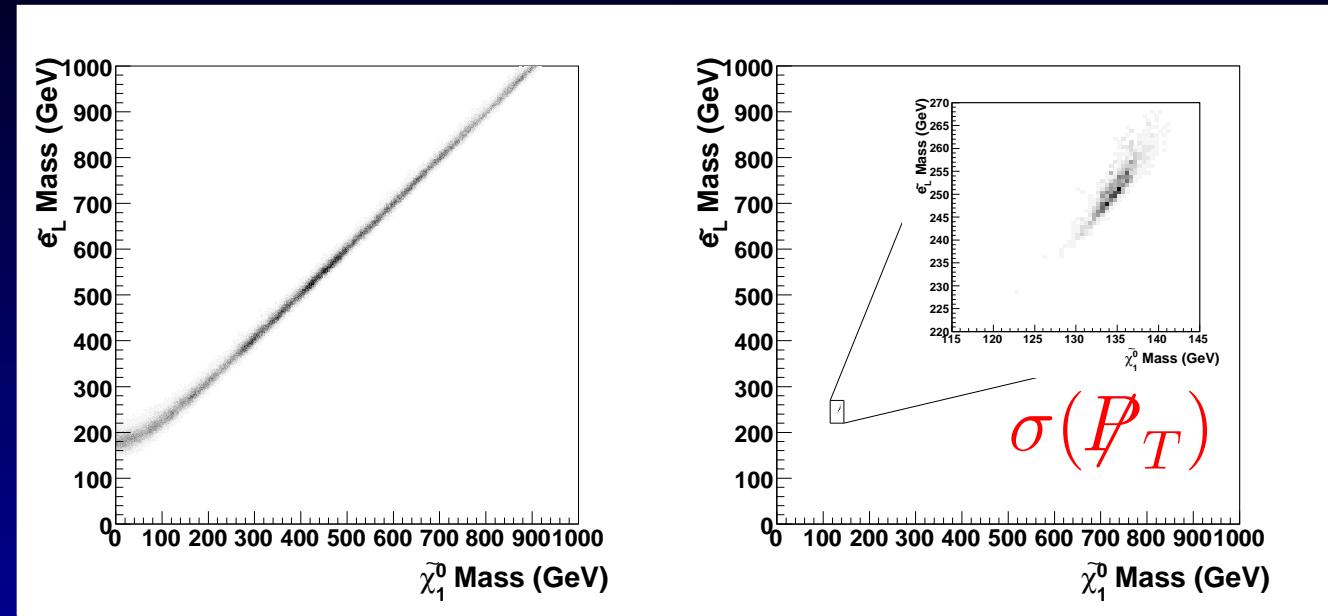
# Fitting to SUSY Breaking Model



- Experimenters pick a SUSY breaking point
- They derive observables and errors after detector simulation
- We fit this “data” with our codes

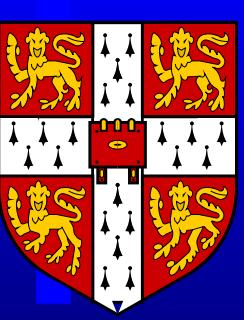


# Fits to future collider data



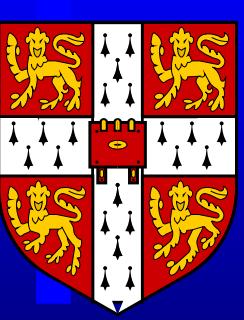
Lester, Parker, White, JHEP 0601 (2006) 080

- Assume edge measurements from some SUSY point: what constraints exist on the phenomenological MSSM?
- SFITTER/FITTINO: see talks by **Rauch**, **Bechtle** in SUSY pheno 1.

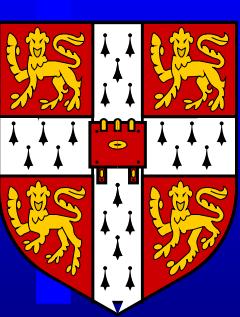


# Summary

- There is now a bewildering multitude of codes for calculating SUSY related observables.
- There has been some organisation and consolidation between them, notably in the form of *Les Houches Accords*.
- SUSY fitting in the **multi-dimensional régime**, currently. Could easily still be in this situation after early LHC data.
- *Markov Chain Monte Carlos* are a very useful tool for exploring such a régime.
- Current dependence on priors should **not** be a surprise: probably only eliminated after ILC data.



# Supplementary Material



# Likelihood and Posterior

*Q:* What's the chance of observing someone to be pregnant, given that they are female?

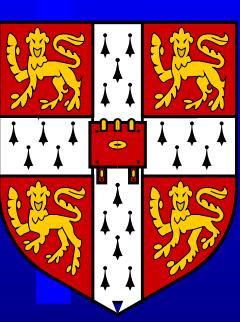


Likelihood

$$p(\text{pregnant} \mid \text{female, human}) = 0.01$$

Posterior

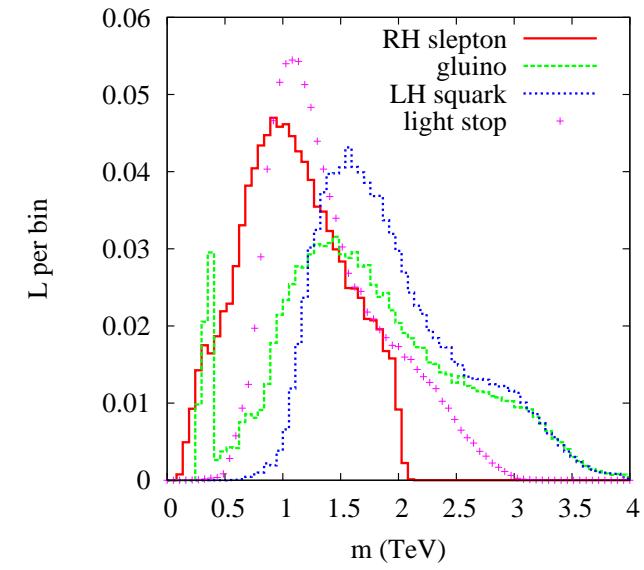
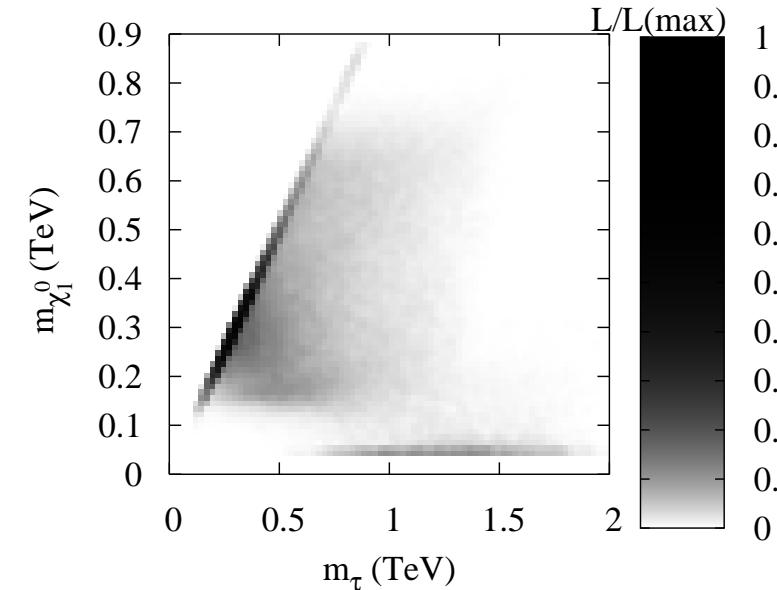
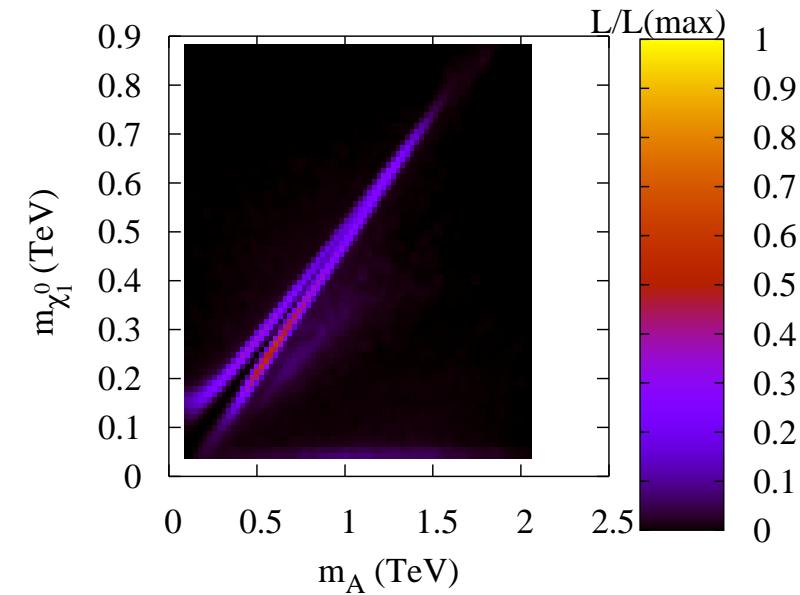
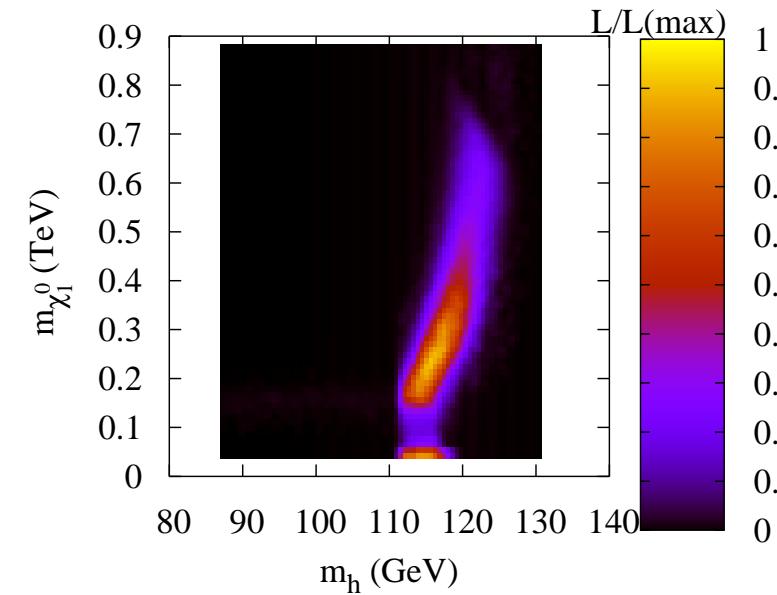
$$p(\text{female} \mid \text{pregnant, human}) = 1.00$$

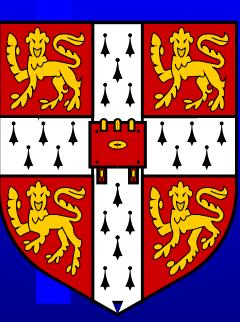


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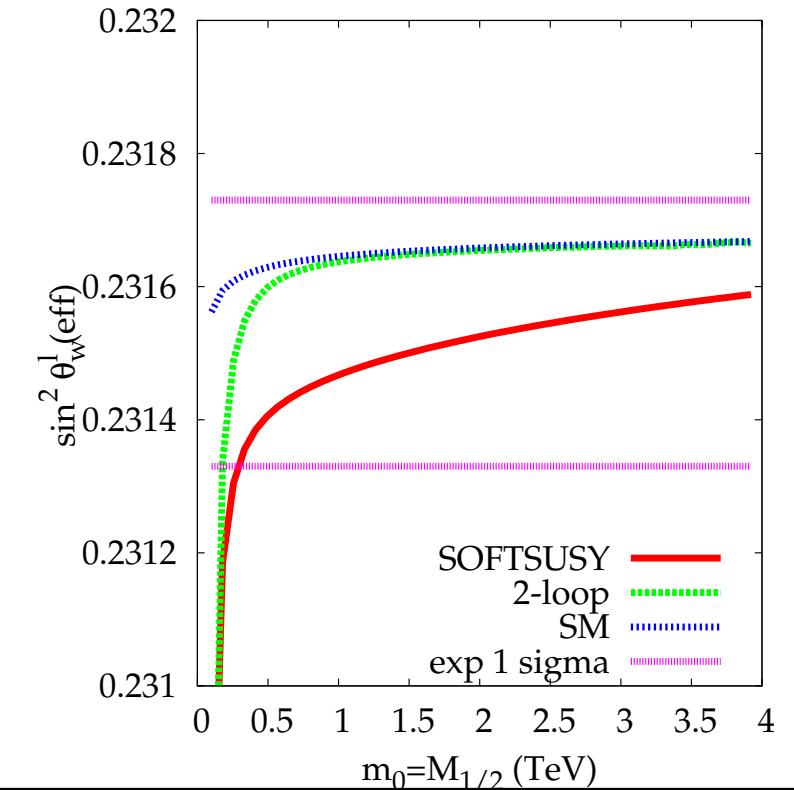
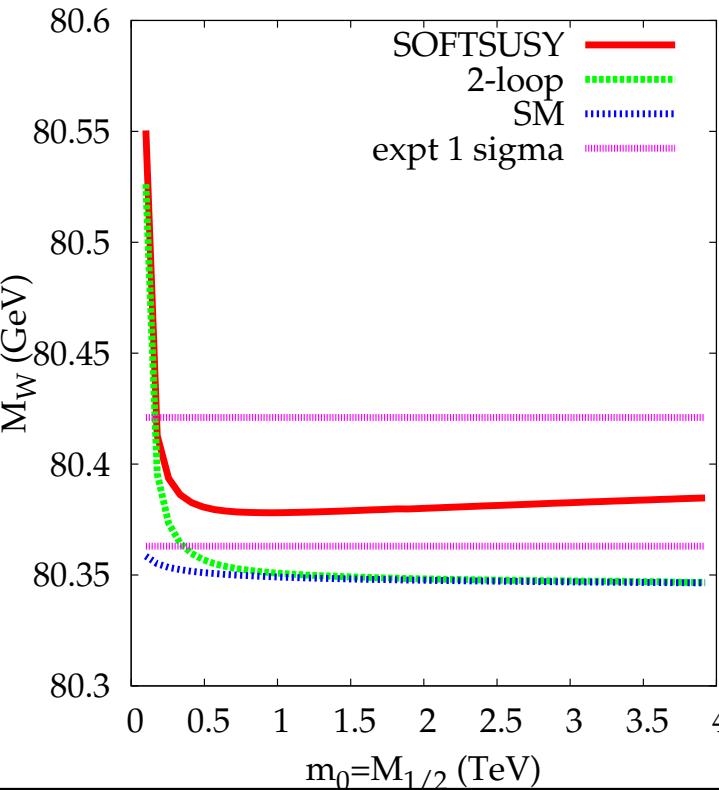
Supersymmetry  
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# Sanity Check





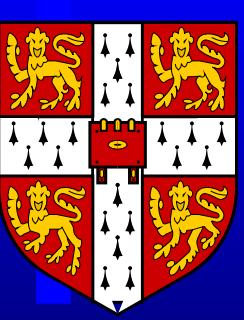
# Electroweak Observables



They prefer light SUSY  
prox.

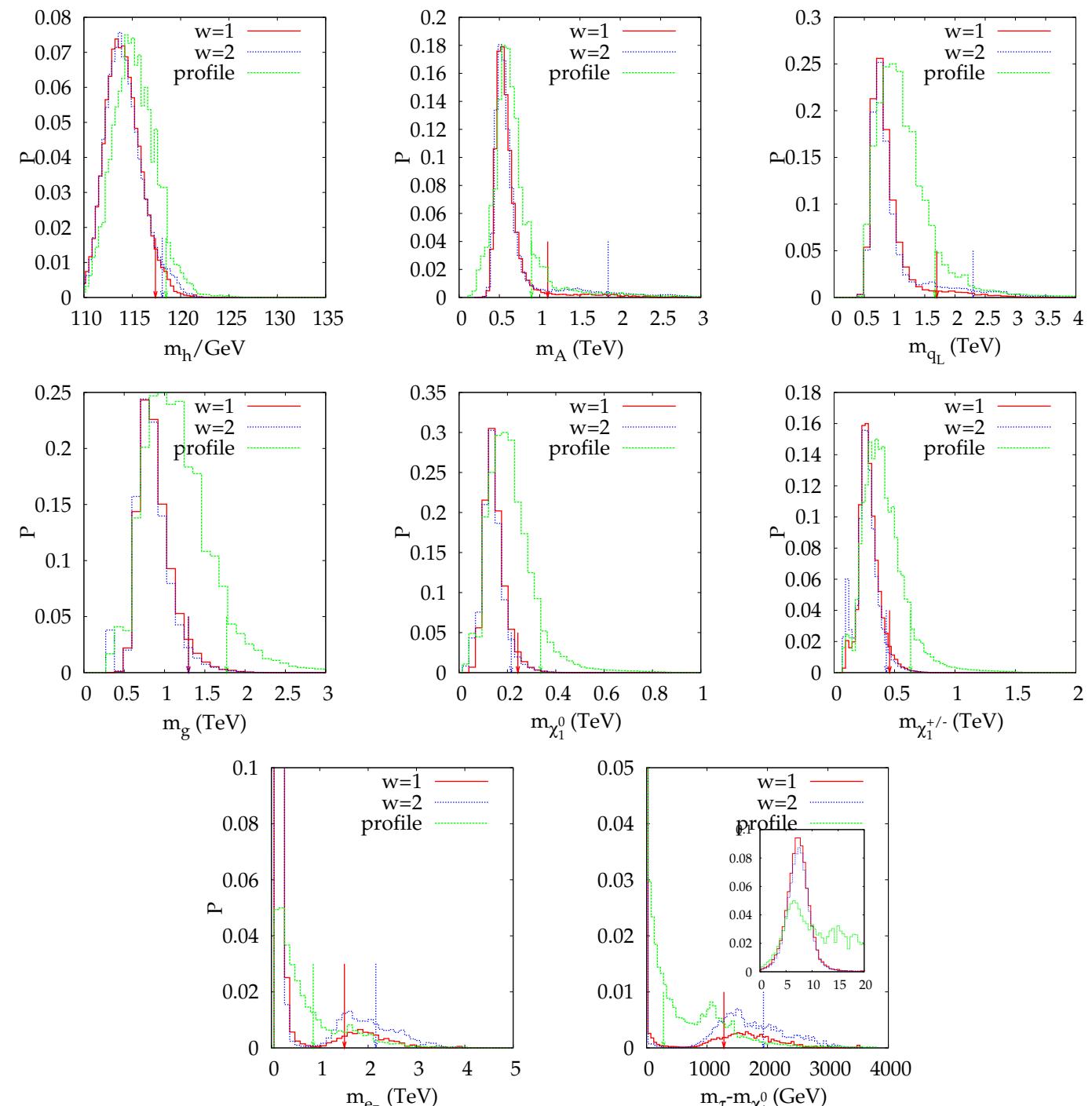
. Be careful of 1-loop ap-

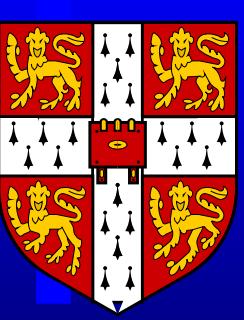
Ellis *et al*, hep-ph/0411216; hep-ph/0602220.



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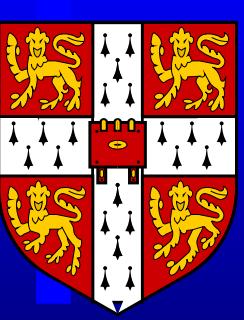
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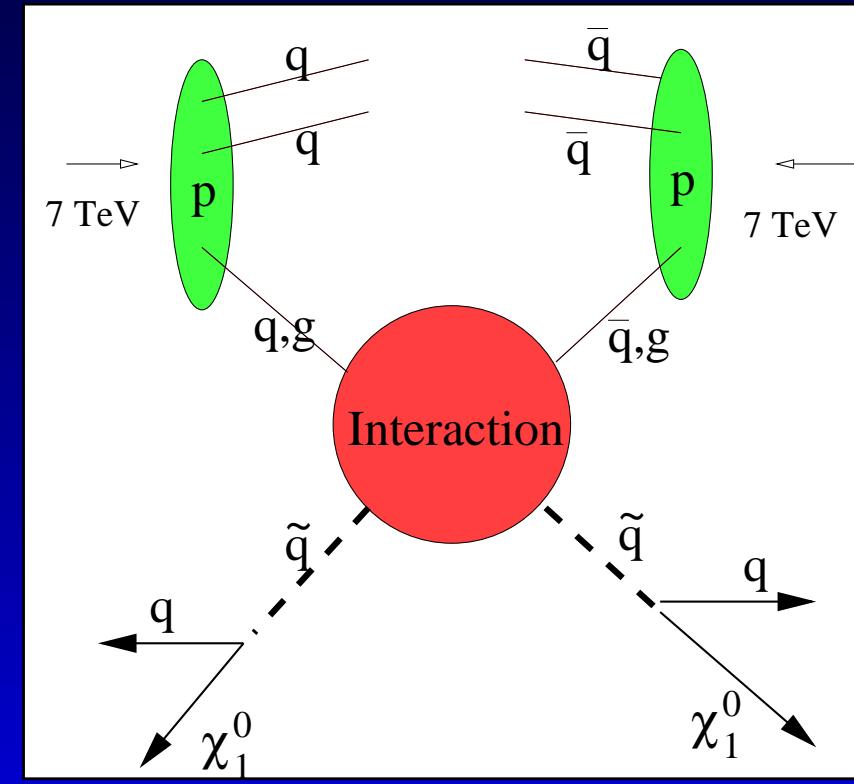
# Caveats

- Implicitly assumed that LSP constitutes *all* of dark matter
- Assumed radiation domination in post-inflation era. No clear evidence between freeze-out+BBN that this is the case ( $t_{eq}$  changes).
- Examples of non-standard cosmology that would change the prediction:
  - Extra degrees of freedom
  - Low reheating temperature
  - Extra dimensional models
  - Anisotropic cosmologies
  - Non-thermal production of neutralinos (late decays?)



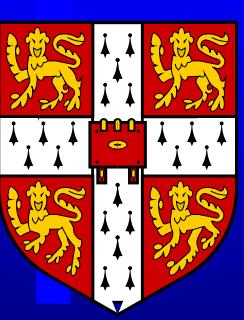
# Collider SUSY Dark Matter Production

Strong sparticle production and decay to dark matter particles.



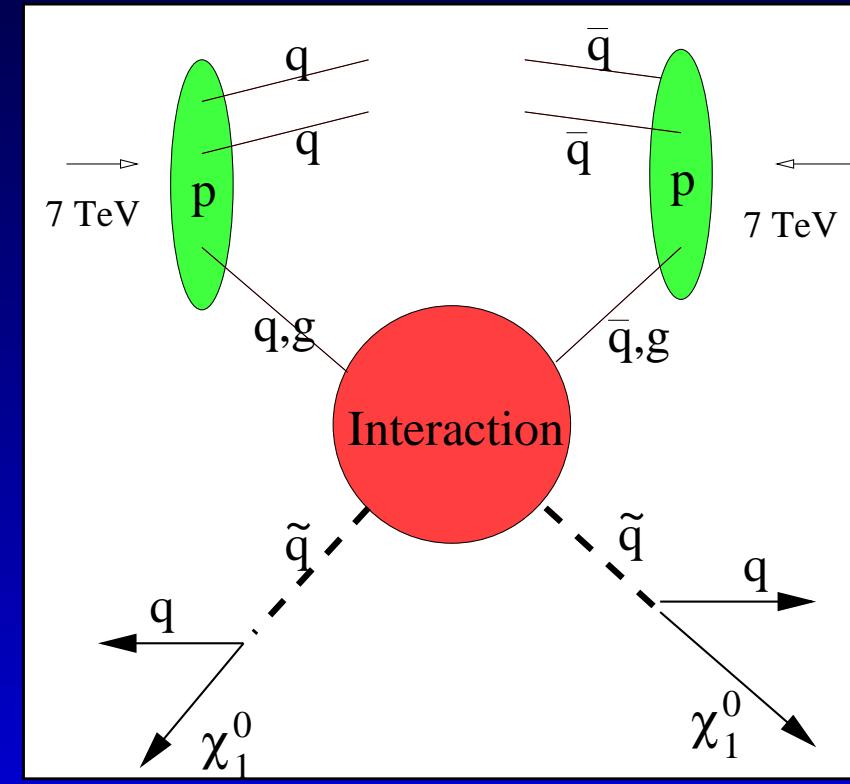
*Q:* Can we measure enough to predict  $\sigma$ ?





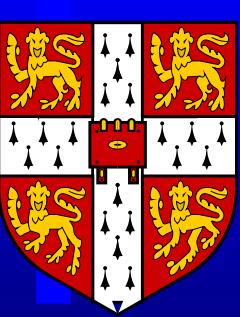
# Collider SUSY Dark Matter Production

Strong sparticle production and decay to dark matter particles.



*Any dark matter candidate that couples to hadrons can  
be produced at the LHC*





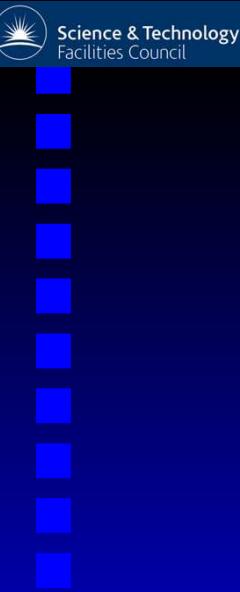
# Collider Check

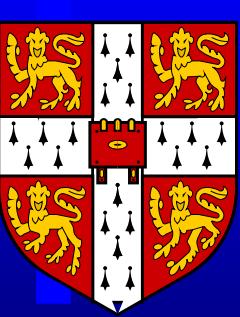
Need corroboration with *direct detection*.  
If we can pin particle physics down, a comparison  
between the predicted relic density and that observed  
is a test of the cosmological assumptions used in the  
prediction.

Thus, if it doesn't fit, you change the cosmology until  
it does.

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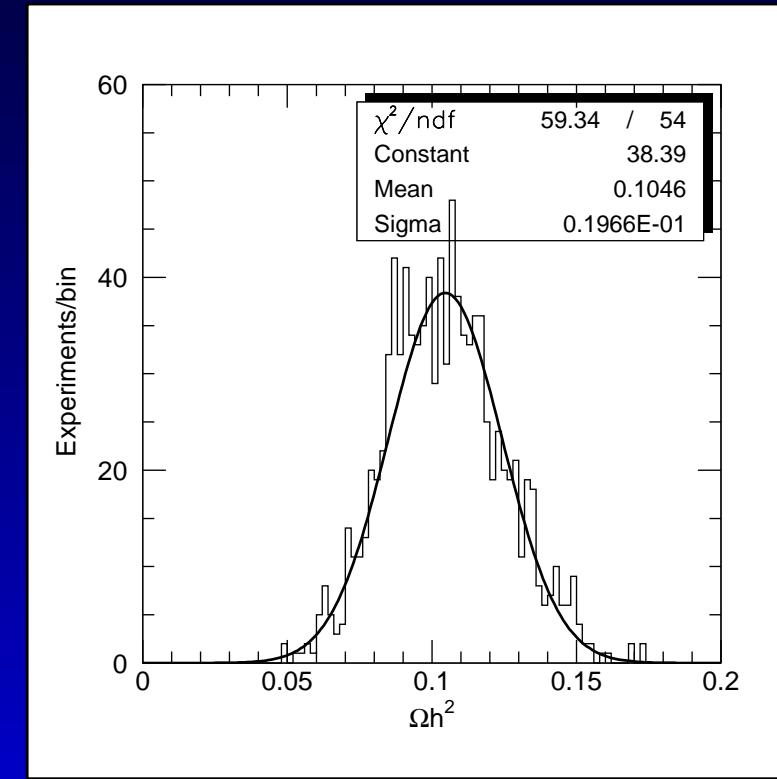
BCA, G. Belanger, F. Boudjema, A. Pukhov, JHEP 0412 (2004)  
020.; M. Nojiri, D. Tovey, JHEP 0603 (2006) 063





# Predicting $\Omega h^2$

Not much left that's allowed but edge measurements  
allow reasonable  $\Omega h^2$  error for  $300 \text{ fb}^{-1}$ .

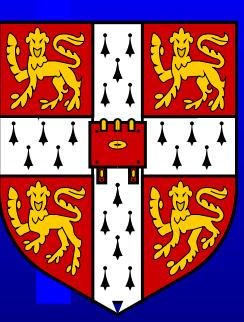


Q: What about other bits of parameter space?

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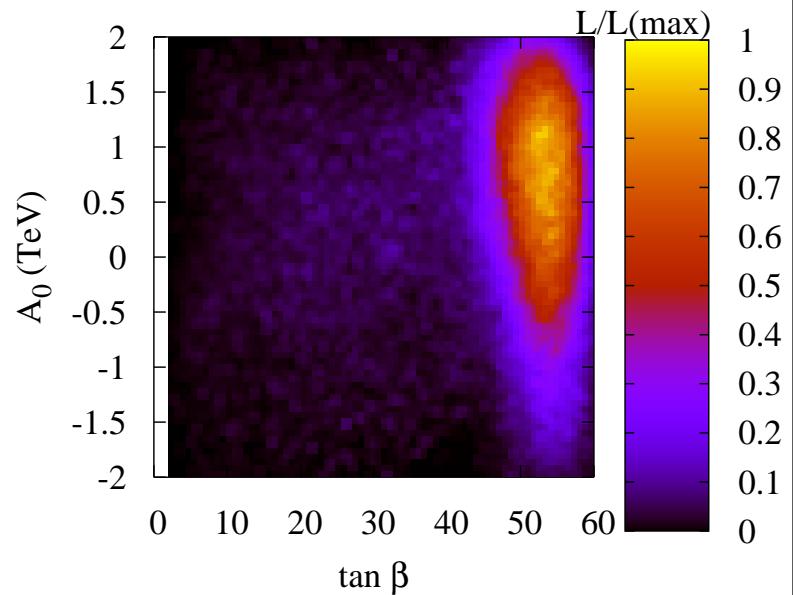
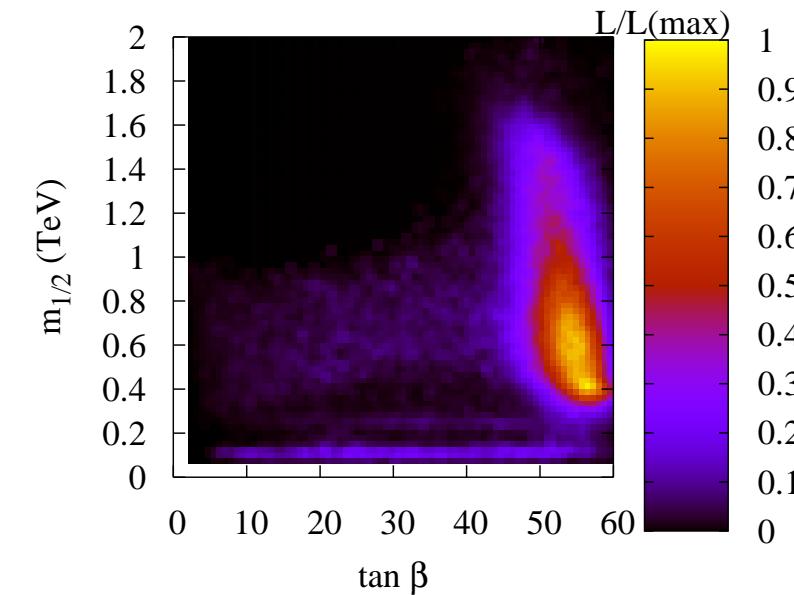
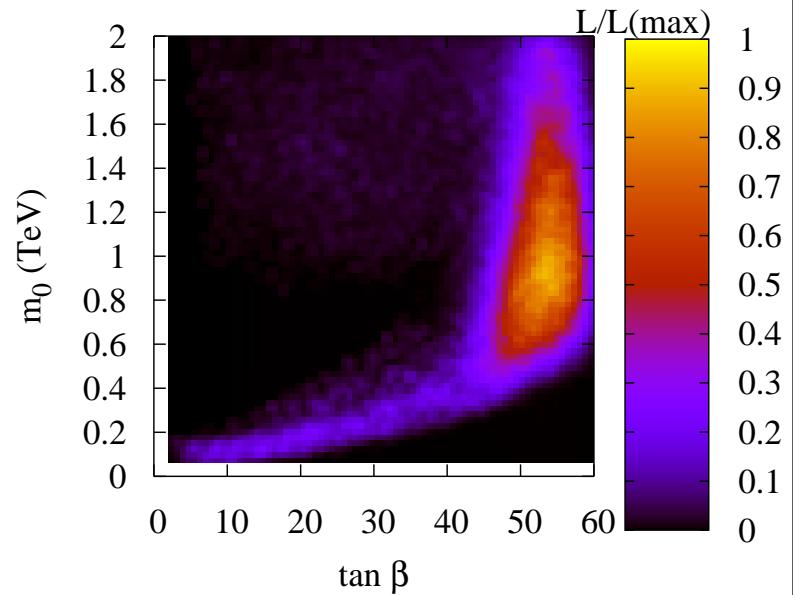
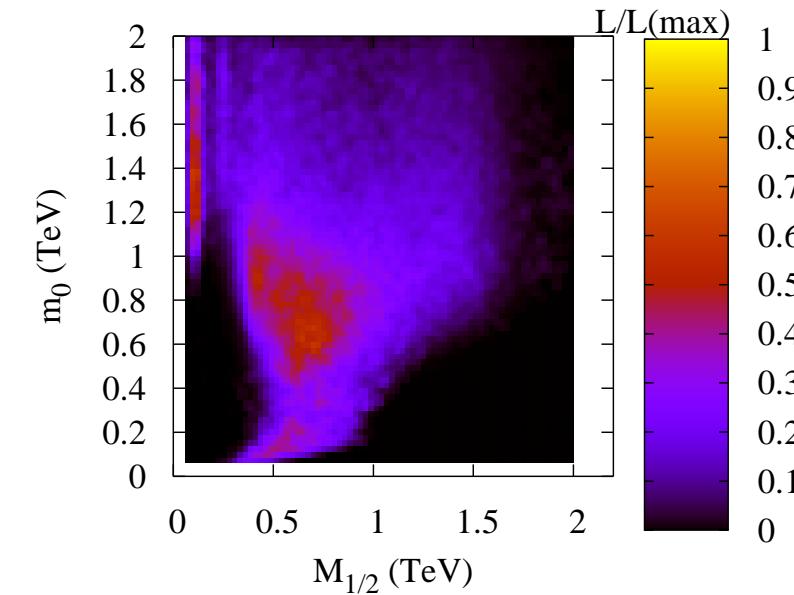
M Nojiri, G Polesello, D Tovey, JHEP 0603 (2006) 063,

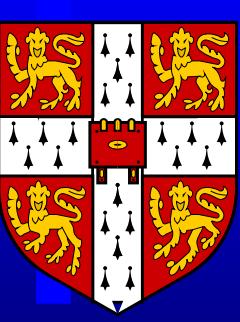
[hep-ph/0512204](https://arxiv.org/abs/hep-ph/0512204).



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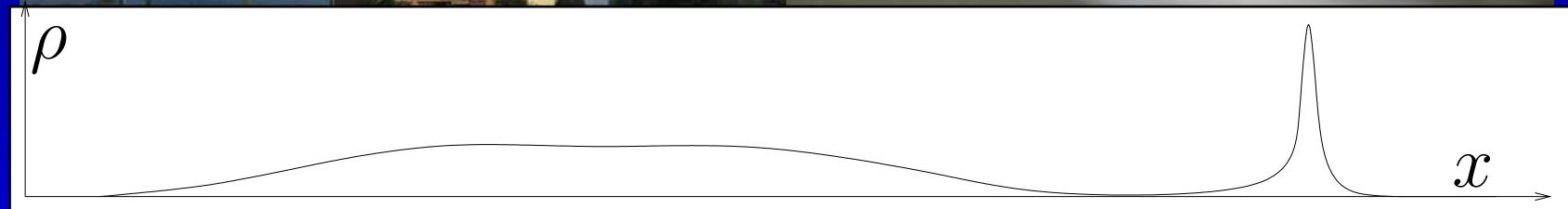
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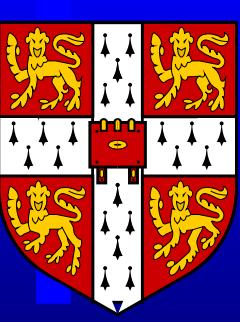




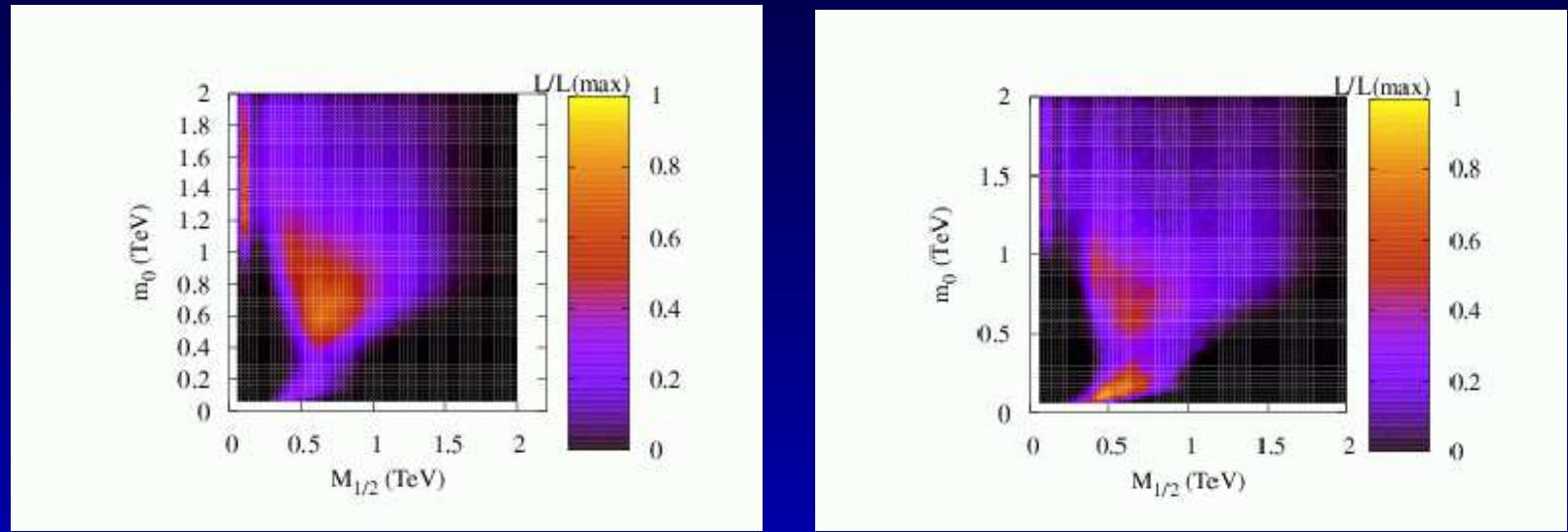
# Volume Effects

*Can't rely on a good  $\chi^2$  in non-Gaussian situation*

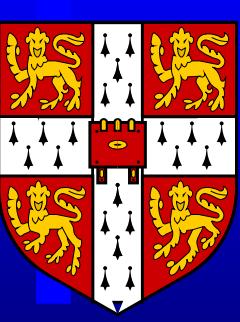




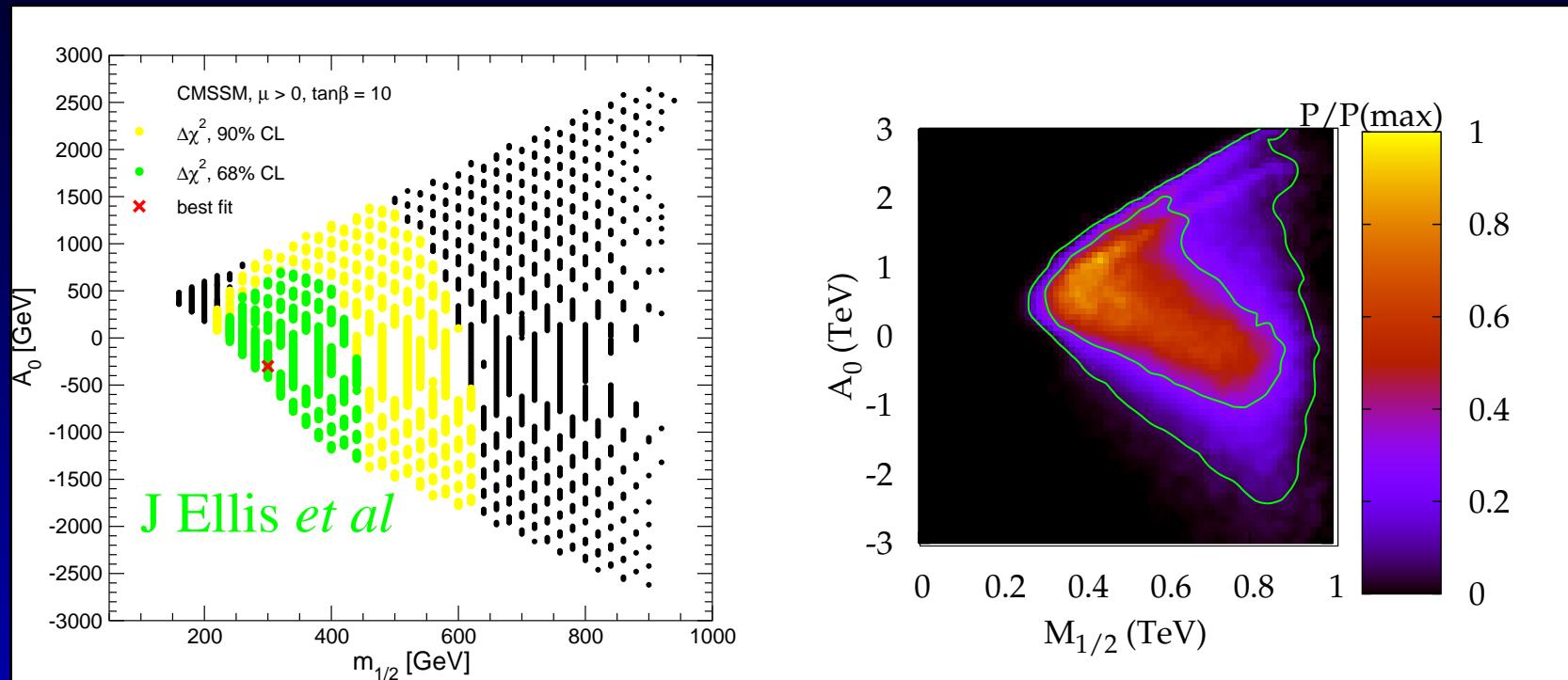
# Comparison



- LHS: allowing non thermal- $\chi_1^0$  contribution
- RHS: only  $\chi_1^0$  dark matter
- (*flat priors*)

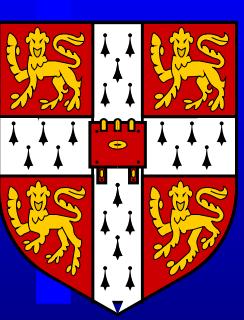


# Comparison



- Fix  $\tan\beta = 10$  and all SM inputs
- Restrict  $m_0, M_{1/2} < 1$  TeV.
- *Same* fits!





# Priors

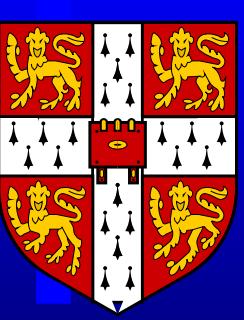
We have assumed a flat prior in  $\tan \beta$ , implies a measure:

$$p(m_0|\text{data}) = \int dM_{1/2} dA_0 d\tan \beta ds p(m_0, M_{1/2}, A_0, \tan \beta, s|\text{data}).$$

$$\begin{aligned} \mu B &= \frac{\sin 2\beta}{2} (\bar{m}_{H_1}^2 + \bar{m}_{H_2}^2 + 2\mu^2), \\ \mu^2 &= \frac{\bar{m}_{H_1}^2 - \bar{m}_{H_2}^2 \tan^2 \beta}{\tan^2 \beta - 1} - \frac{M_Z^2}{2}. \end{aligned}$$

Change variables:  $\int d\mu dB \rightarrow \int dM_Z d\tan \beta |J|$





# EWSB prior

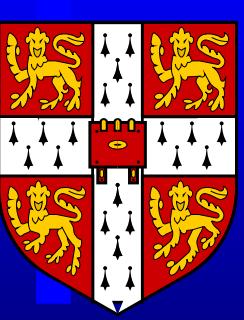
$$\begin{aligned} & p(\text{all data} | m_0, M_{1/2}, A_0, \mu, B, s) \\ & \approx p(\text{data} | m_0, M_{1/2}, A_0, \mu, B, s) \times \\ & \quad p(M_Z | m_0, M_{1/2}, A_0, \mu, B, s). \\ & \approx p(\text{data} | m_0, M_{1/2}, A_0, \mu, B, s) \times \delta(M_Z - M_Z^{cen}) \end{aligned}$$

Change variables

$$\int d\mu dB \delta(M_Z - M_Z^{cen}) \rightarrow \int d\tan\beta |J|:$$

$$J = \frac{B}{\mu \tan\beta} \frac{\tan^2\beta + 1}{\tan^2\beta - 1}$$





# Same order prior

We wish to encode the idea that “SUSY breaking terms should be of the same order of magnitude”

$$p(m_0|M_S) = \frac{1}{\sqrt{2\pi w^2} m_0} \exp\left(-\frac{1}{2w^2} \log^2\left(\frac{m_0}{M_S}\right)\right),$$

$$p(A_0|M_S) = \frac{1}{\sqrt{2\pi e^{2w}} M_S} \exp\left(-\frac{1}{2e^{2w}} \frac{A_0^2}{M_S^2}\right),$$

We don't know SUSY breaking scale  $M_S$ :

$$p(m_0, M_{1/2}, A_0, \mu, B) =$$

$$\int_0^\infty dM_S p(m_0, M_{1/2}, A_0, \mu, B|M_S) p(M_S)$$

