

New constraints from isospin asymmetry in $b \rightarrow s\gamma$ transitions

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Karlsruhe - July 31, 2007

Outline

Introduction

Theoretical framework

Effective Hamiltonian

Supersymmetric contributions

Isospin Asymmetry

Experimental limits and data

New constraints from Isospin Asymmetry

Conclusion

Motivations

- ▶ $b \rightarrow s\gamma$ transition: very sensitive to new physics
 - ▶ forbidden at the tree level in SM and can only be induced via loop diagrams,
 - ▶ SM contributions are vanishingly small,
- ▶ branching ratios have been extensively used to constrain SUSY parameter space
- ▶ Study another observable: isospin asymmetry
 - ▶ already measured by BELLE and BABAR

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Effective Hamiltonian

The idea of $B \rightarrow X_s \gamma$ decay begins with introducing an effective Hamiltonian:

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_{i=1}^8 C_i(\mu) O_i(\mu)$$

$$\left\{ \begin{array}{ll} O_1 = (\bar{s}_L \gamma_\mu T^a c_L)(\bar{c}_L \gamma^\mu T^a b_L) & O_2 = (\bar{s}_L \gamma_\mu c_L)(\bar{c}_L \gamma^\mu b_L) \\ O_3 = (\bar{s}_L \gamma_\mu b_L) \sum_q (\bar{q} \gamma^\mu q) & O_4 = (\bar{s}_L \gamma_\mu T^a b_L) \sum_q (\bar{q} \gamma^\mu T^a q) \\ O_5 = (\bar{s}_L \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} b_L) \sum_q (\bar{q} \gamma^{\mu_1} \gamma^{\mu_2} \gamma^{\mu_3} q) & \\ O_6 = (\bar{s}_L \gamma_{\mu_1} \gamma_{\mu_2} \gamma_{\mu_3} T^a b_L) \sum_q (\bar{q} \gamma^{\mu_1} \gamma^{\mu_2} \gamma^{\mu_3} T^a q) & \\ O_7 = \frac{e}{16\pi^2} m_b (\bar{s}_L \sigma^{\mu\nu} b_R) F_{\mu\nu} & O_8 = \frac{g}{16\pi^2} m_b (\bar{s}_L \sigma^{\mu\nu} T^a b_R) G_{\mu\nu}^a \end{array} \right.$$

Wilson Coefficients

$$C_i^{\text{eff}}(\mu) = C_i^{(0)\text{eff}}(\mu) + \frac{\alpha_s(\mu)}{4\pi} C_i^{(1)\text{eff}}(\mu) + \dots$$

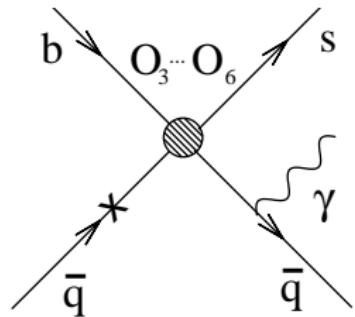
The effective coefficients evolve according to their RGE:

$$\mu \frac{d}{d\mu} C_i^{\text{eff}}(\mu) = C_j^{\text{eff}}(\mu) \gamma_{ji}^{\text{eff}}(\mu)$$

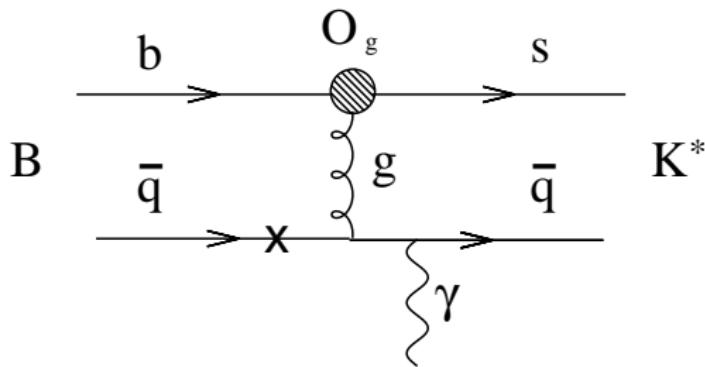
driven by the anomalous dimension matrix $\hat{\gamma}^{\text{eff}}(\mu)$:

$$\hat{\gamma}^{\text{eff}}(\mu) = \frac{\alpha_s(\mu)}{4\pi} \hat{\gamma}^{(0)\text{eff}} + \frac{\alpha_s^2(\mu)}{(4\pi)^2} \hat{\gamma}^{(1)\text{eff}} + \dots$$

Contribution to Isospin Asymmetry



QCD penguin operators

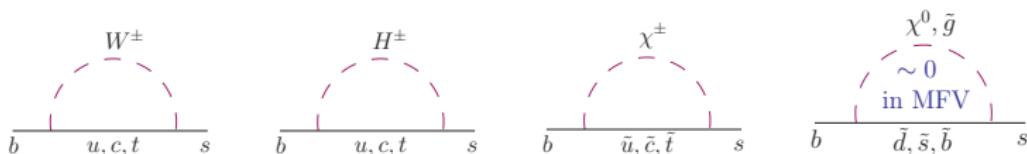


Electro- and chromo-magnetic
operators

Supersymmetric contributions

MSSM with minimal flavor violation (MFV)

↪ no more flavor/CP violation than in SM



Calculation of the coefficients at $\mu = M_W$:

$$C_i(\mu) = C_i^{W^\pm}(\mu) + C_i^{H^\pm}(\mu) + C_i^{\chi^\pm}(\mu)$$

Gómez et al. Phys. Rev. D74, 015015 (2006)

Degrassi et al. JHEP 12, 009 (2000)

Ciuchini et al. Nucl. Phys. B 534, 3 (1998)

Ciuchini et al. Nucl. Phys. B 527, 21 (1998)

Isospin Asymmetry

$$\Delta_{0-} \equiv \frac{\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0}\gamma) - \Gamma(B^- \rightarrow K^{*-}\gamma)}{\Gamma(\bar{B}^0 \rightarrow \bar{K}^{*0}\gamma) + \Gamma(B^- \rightarrow K^{*-}\gamma)}$$

$$\Delta_{0-} = \text{Re}(b_d - b_u).$$

$$b_q = \frac{12\pi^2 f_B Q_q}{m_b T_1^{B \rightarrow K^*} a_7^c} \left(\frac{f_{K^*}^\perp}{m_b} K_1 + \frac{f_{K^*} m_{K^*}}{6\lambda_B m_B} K_2 \right)$$

$$a_7^c = C_7 + \frac{\alpha_s(\mu) C_F}{4\pi} \left(C_1(\mu) G_1(s_p) + C_8(\mu) G_8 \right) + \frac{\alpha_s(\mu_h) C_F}{4\pi} \left(C_1(\mu_h) H_1(s_p) + C_8(\mu_h) H_8 \right)$$

In the Standard Model: $\Delta_{0-} \simeq 8\%$

Kagan and Neubert, Phys. Lett. B 539, 227 (2002)

Bosch and Buchalla, Nucl. Phys. B 621, 459 (2002)

Experimental data

BABAR

$$\Delta_{0-} = -0.006 \pm 0.058(\text{stat}) \pm 0.009(\text{syst}) \pm 0.024(R^{+/0})$$

Aubert et al. (BABAR Collaboration) Phys. Rev. D72 (2005)

BELLE

$$\Delta_{0+} = +0.012 \pm 0.044(\text{stat}) \pm 0.026(\text{syst})$$

Nakao et al. (BELLE Collaboration) Phys. Rev. D69 (2004)

Allowed Region: $-0.047 < \Delta_{0-} < 0.093$

Experimental limits

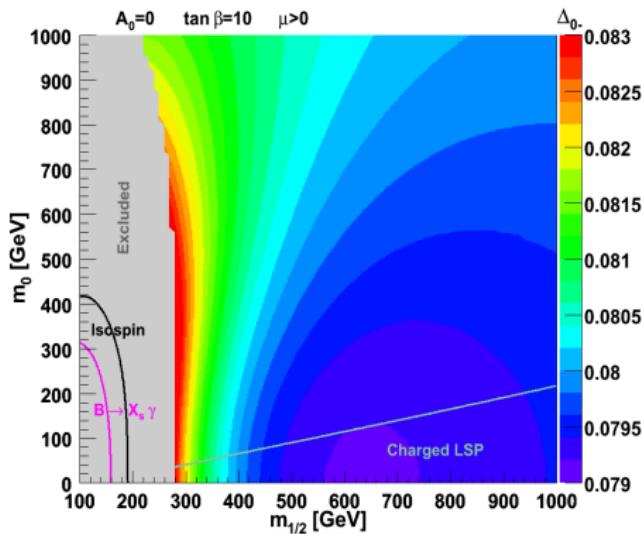
Lower bounds on sparticle masses in GeV:

Particle	h^0	χ_1^0	\tilde{l}_R	$\tilde{\nu}_{e,\mu}$	χ_1^\pm	\tilde{t}_1	\tilde{g}	\tilde{b}_1	$\tilde{\tau}_1$	\tilde{q}_R
Lower bound	111	46	88	43.7	67.7	92.6	195	89	81.9	250

Yao et al. J. Phys. G33 (2006)

Results: mSUGRA

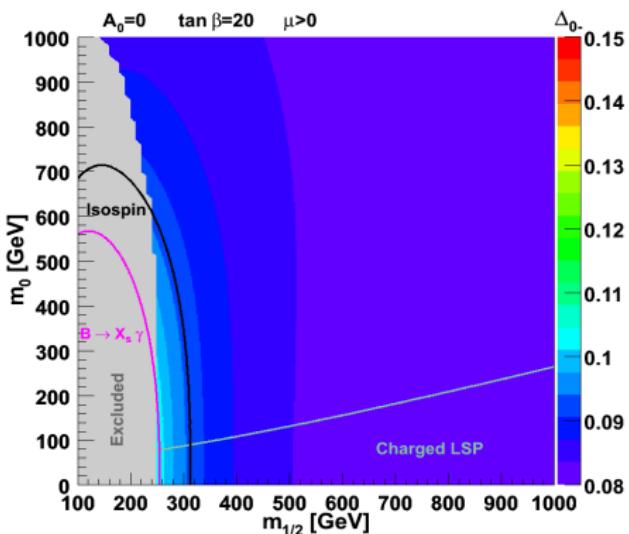
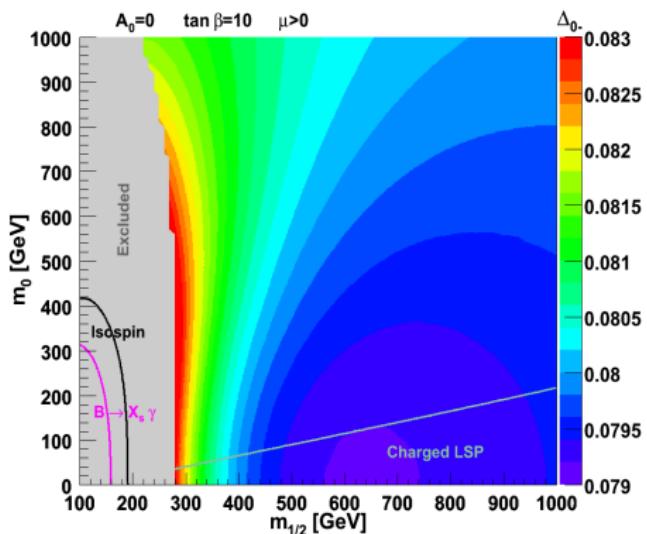
Masses and couplings calculated with IsaJet



Ahmady & Mahmoudi, Phys. Rev. D75 (2007)

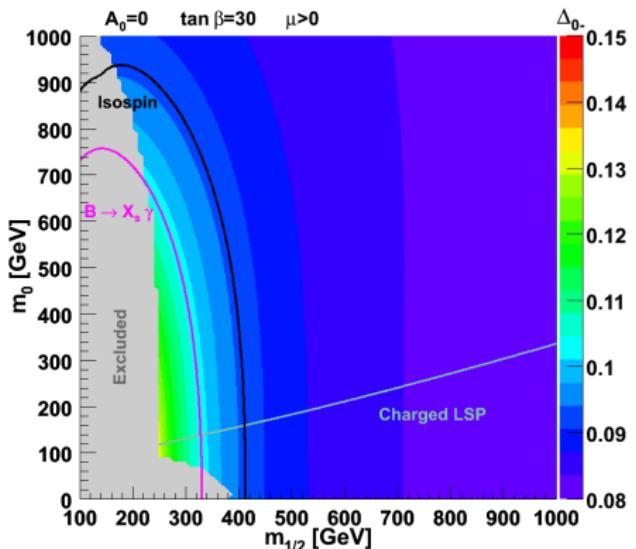
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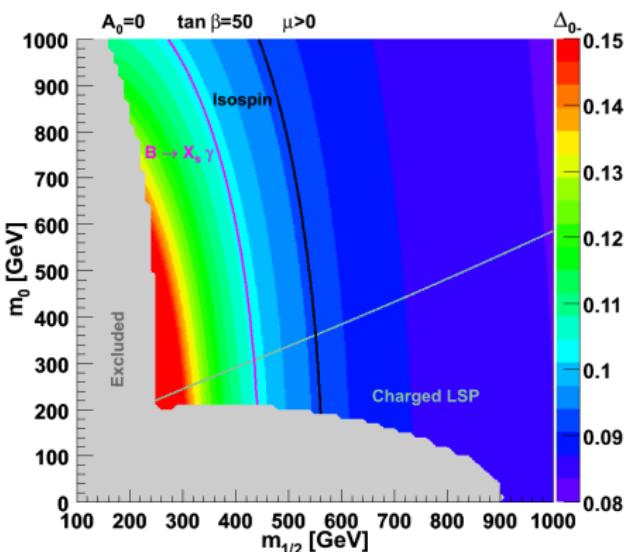
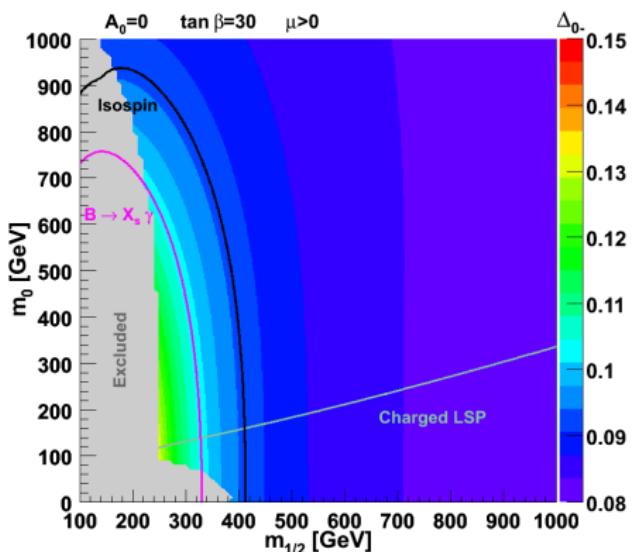
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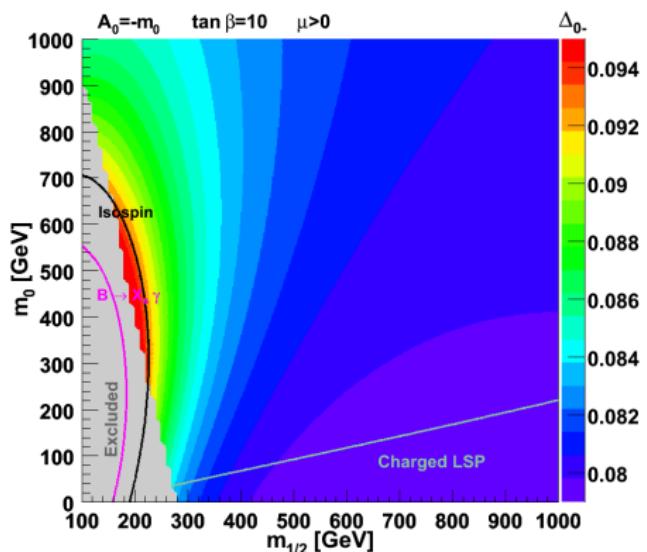
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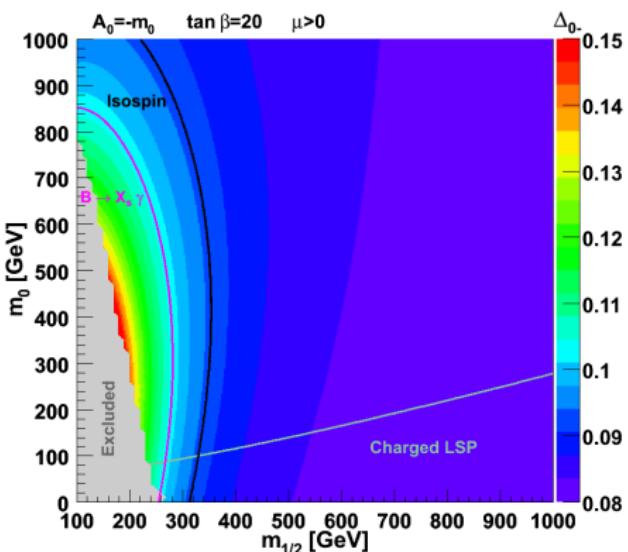
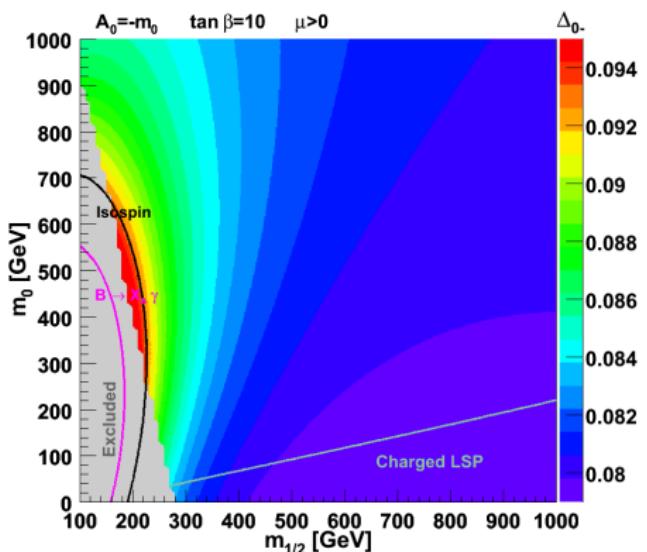
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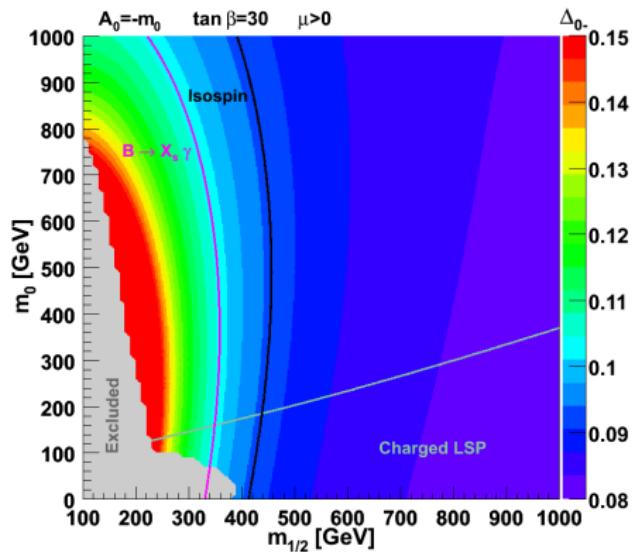
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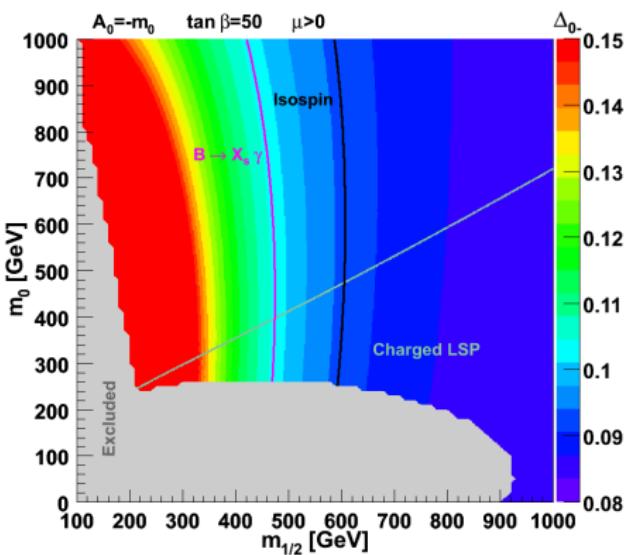
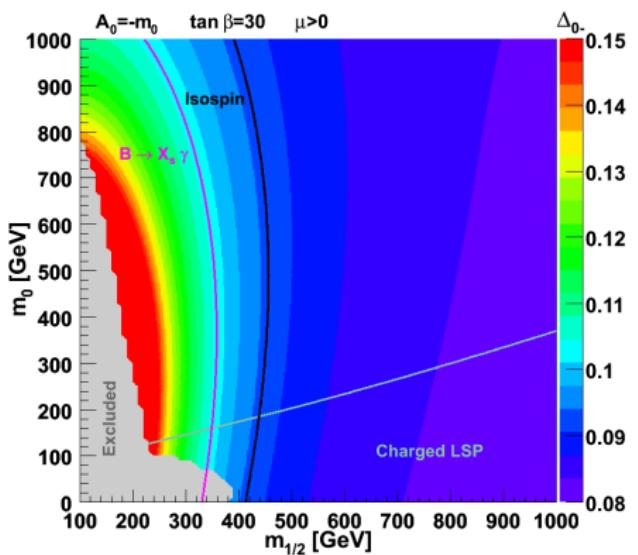
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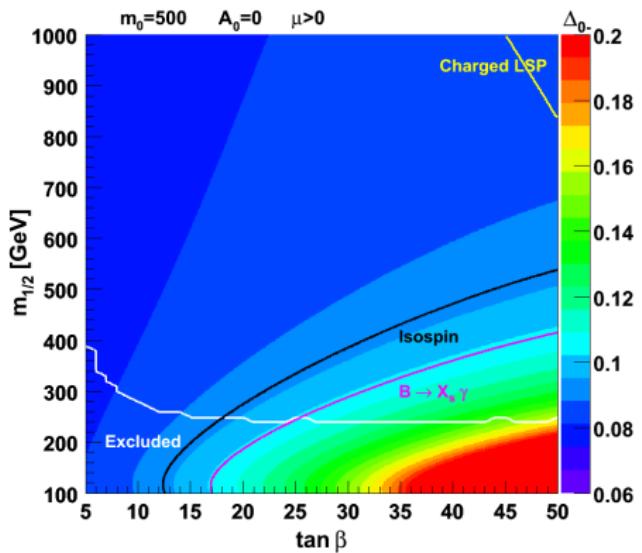


Ahmady & Mahmoudi, Phys. Rev. D75 (2007)

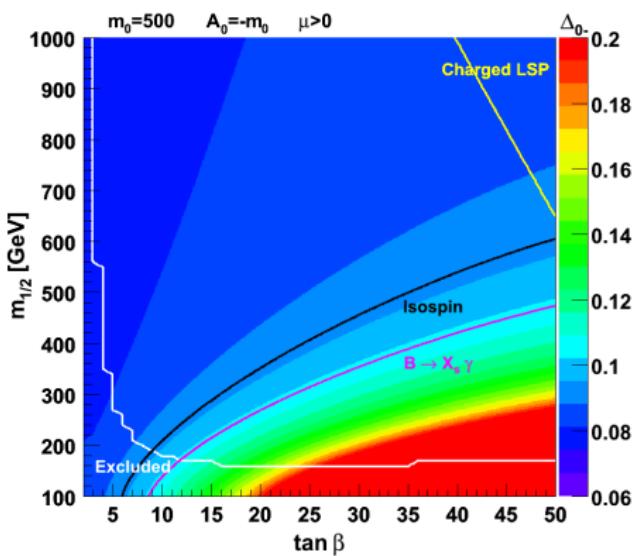
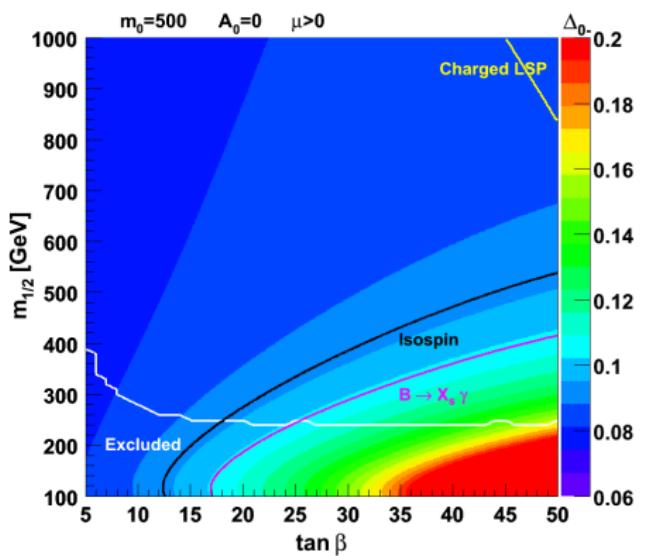
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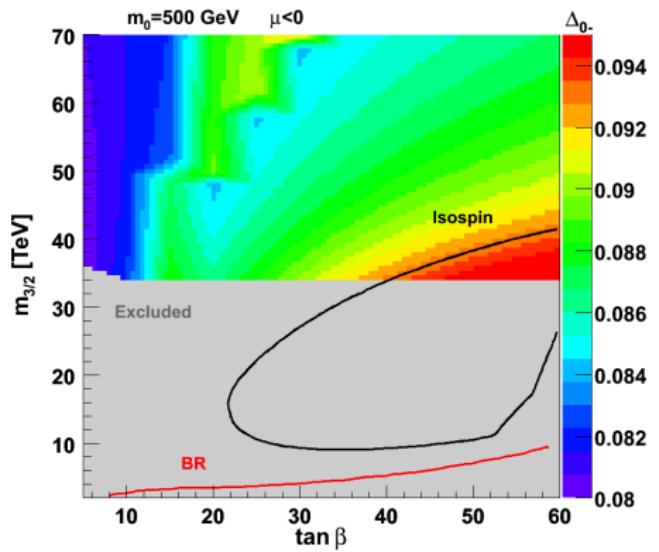
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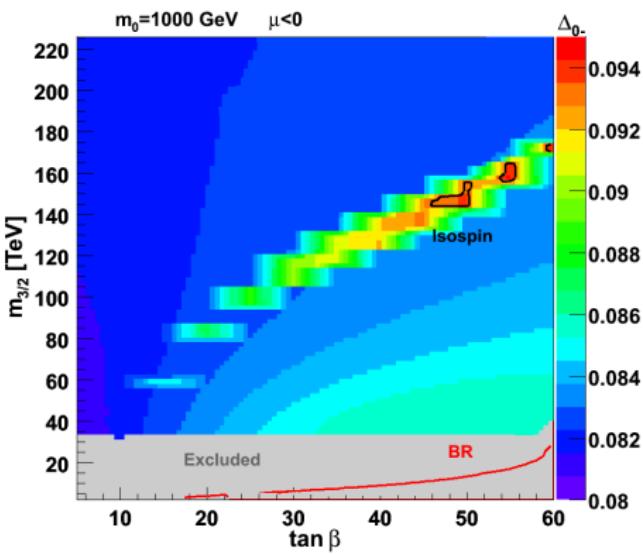
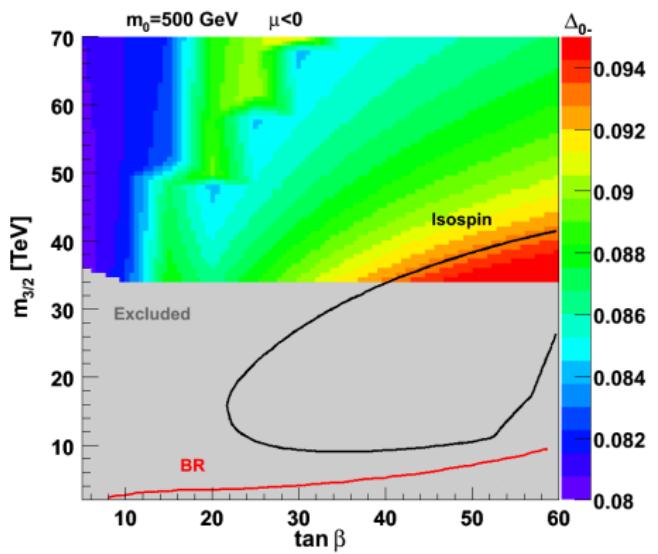
Results: mSUGRA



Results: AMSB

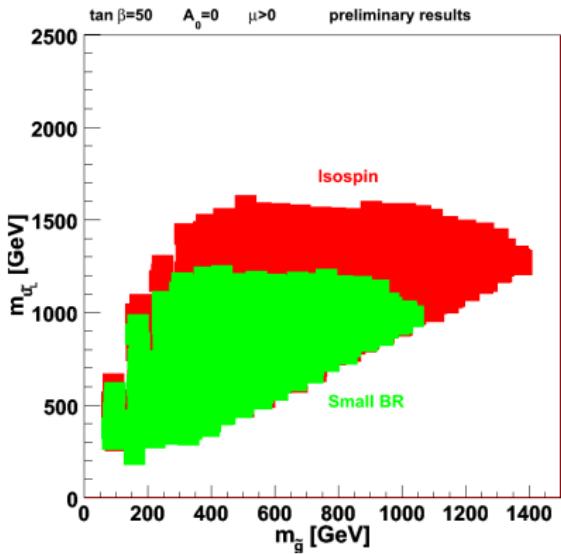


Results: AMSB



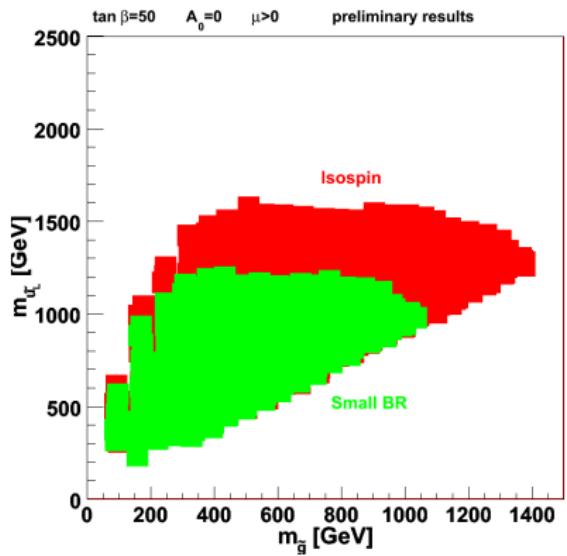
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mSUGRA

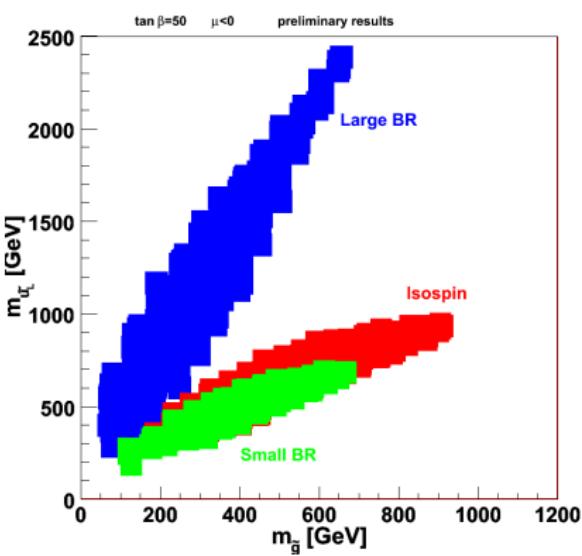


Results:

mSUGRA



AMSB



Conclusion

- ▶ We can obtain new constraints (new contours) using the Isospin asymmetry
- ▶ Very tight constraints on the mSUGRA parameter space, better than the inclusive branching ratio
- ▶ Can be applied to other models
- ▶ Isospin asymmetry seems to be an important observable in the precision test of the SM and in constraining new physics parameters

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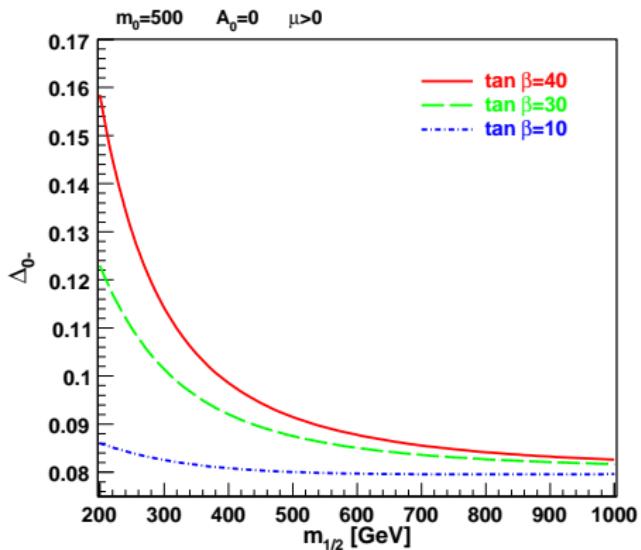
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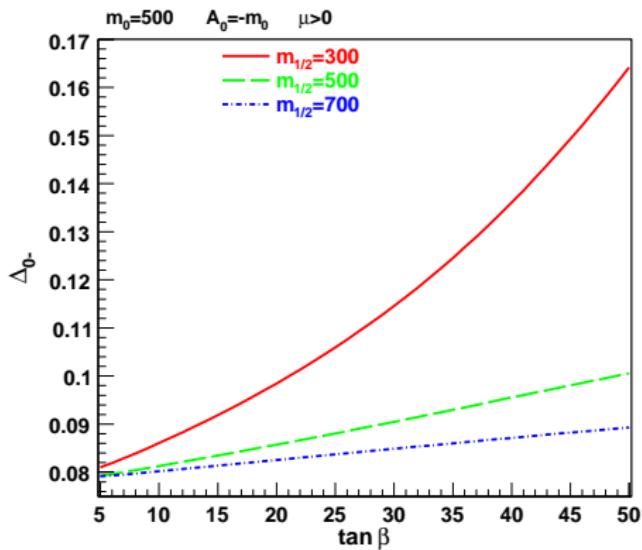
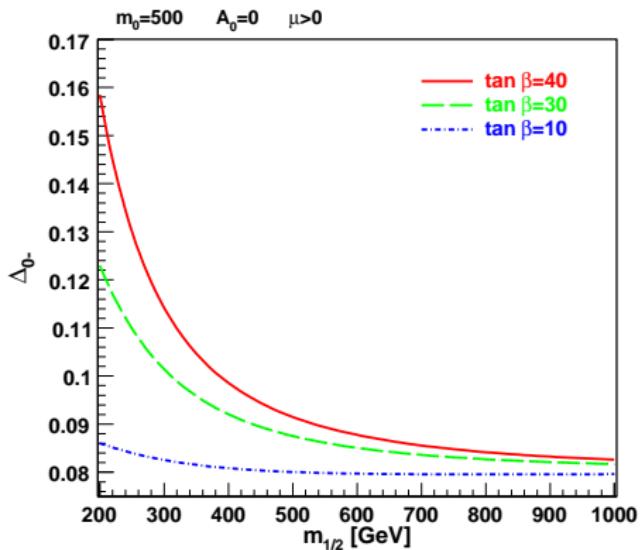
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Backup

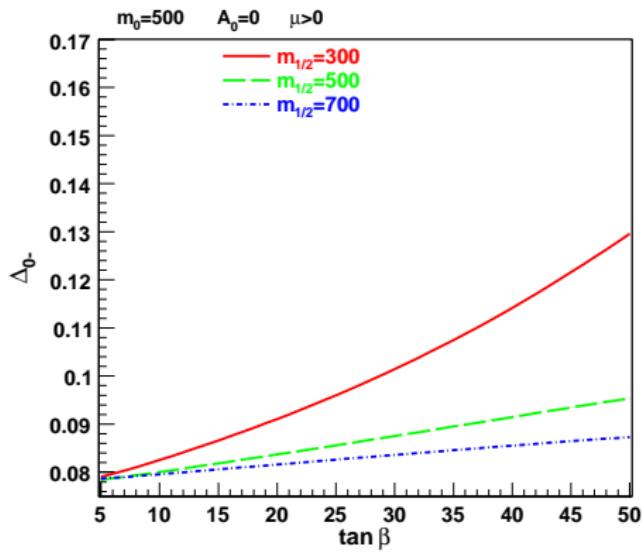
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