

Krzysztof Turzyński (University of Michigan)

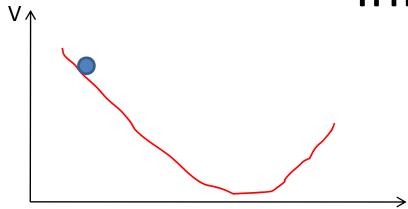
based on arXiv:0704.0212

with **Z.Lalak**¹, **D.Langlois**² and **S.Pokorski**¹

¹ Warsaw University, Poland

² IAP & University Paris 7, France

Inflation



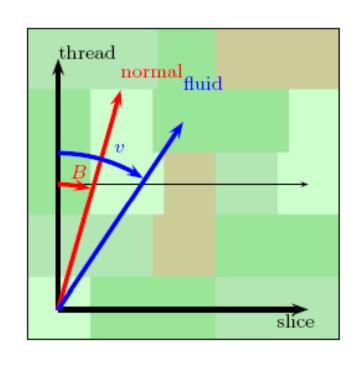
$$\ddot{\sigma} + 3H\dot{\sigma} + V_{\sigma} = 0 \qquad \sigma$$

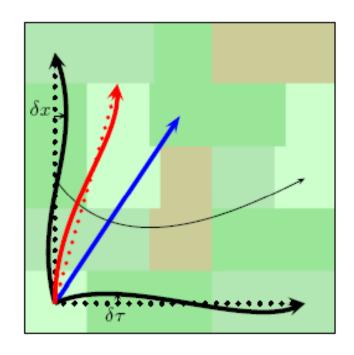
$$H = \frac{\dot{a}}{a} = \frac{1}{3M_{Pl}^{2}}\rho$$

$$\varepsilon = \frac{\dot{\sigma}^2}{2H^2 M_{Pl}^2} \qquad \eta_{\sigma\sigma} = \frac{V_{\sigma\sigma}}{3H^2}$$
small = slow-roll

- solves the flatness problem
- solves the horizon problem
- dilutes any preexisting relics
- generates primordial density perturbations through quantum fluctuations of the coupled inflaton-gravity system

Curvature perturbations

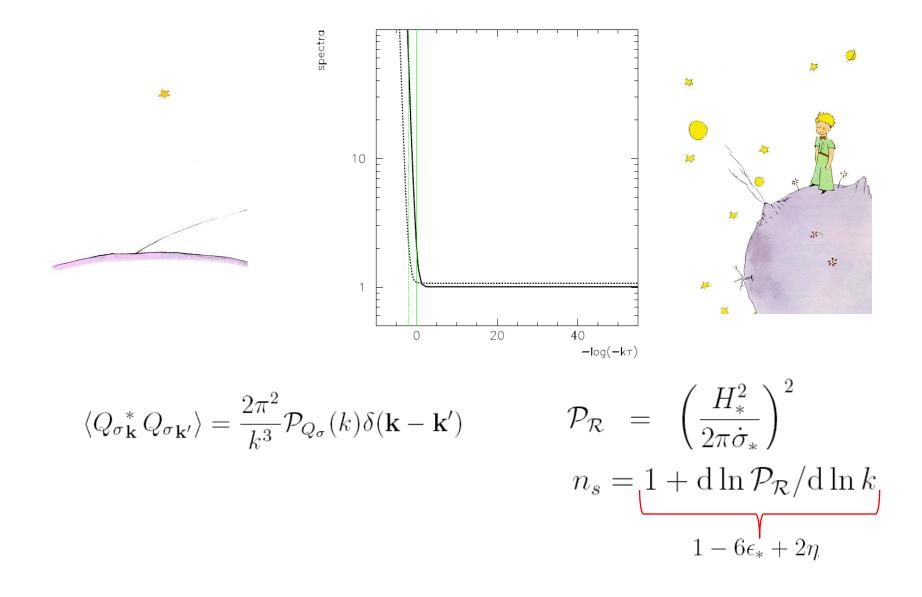




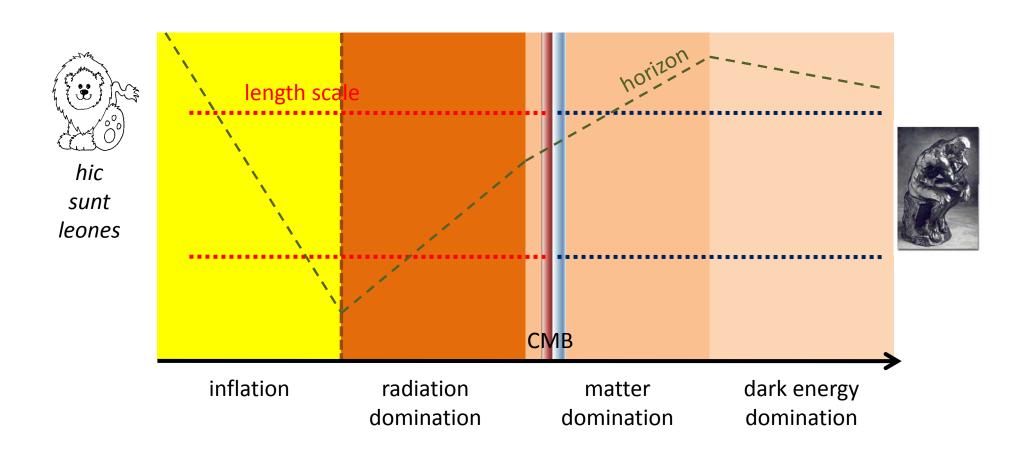
 Q_{σ} gauge invariant variable describing perturbations

$$R = \frac{H}{\dot{\phi}} Q_{\sigma}$$
 3-dim curvature of the comoving hypersurfaces

Evolution of perturbations



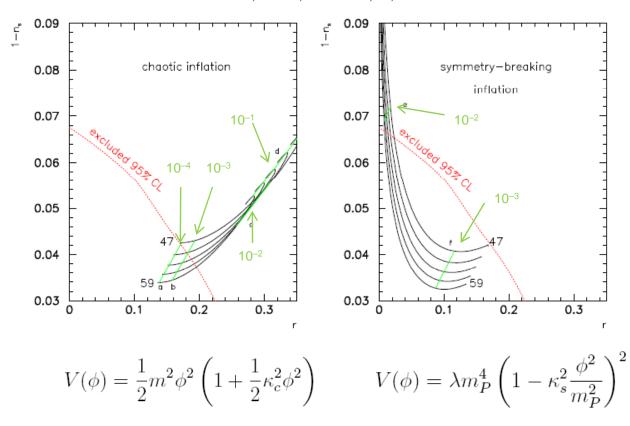
Road map





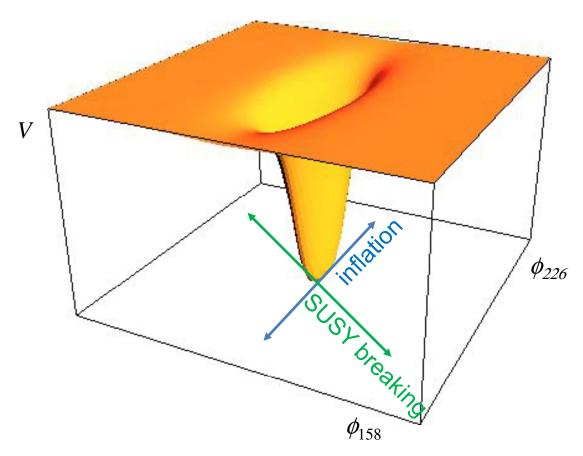
Observations and inflation

Ellis, Lalak, Pokorski, kt, 2006



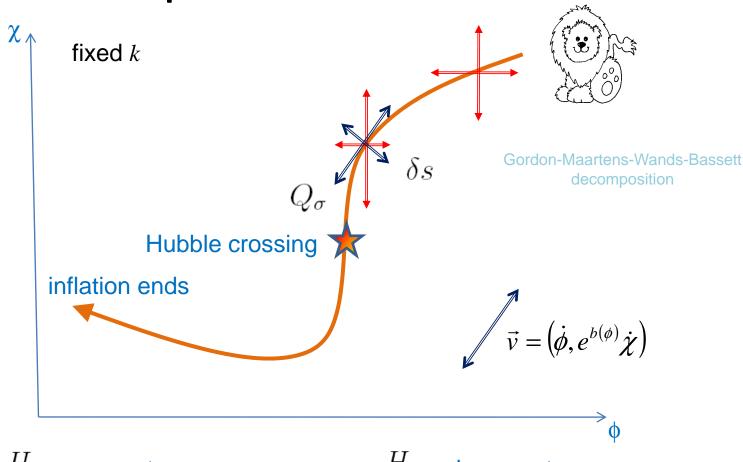


Motivation for multi-field scenarios



$$S = \int d^4x \sqrt{-g} \left[\frac{M_P^2}{2} R - \frac{1}{2} (\partial_\mu \phi) (\partial^\mu \phi) - \frac{e^{2b(\phi)}}{2} (\partial_\mu \chi) (\partial^\mu \chi) - V(\phi, \chi) \right]$$

Curvature and isocurvature perturbations



$$\mathcal{R} \equiv \frac{H}{\dot{\sigma}} Q_{\sigma} \quad \begin{array}{c} \text{curvature} \\ \text{perturbation} \end{array}$$

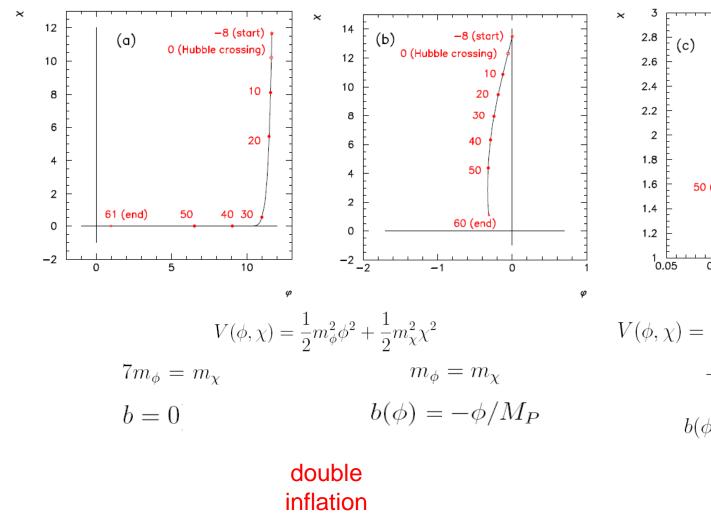
$$\mathcal{S} \equiv \frac{H}{\dot{\sigma}} \delta s \quad \begin{array}{l} \text{isocurvature} \\ \text{perturbation} \end{array}$$

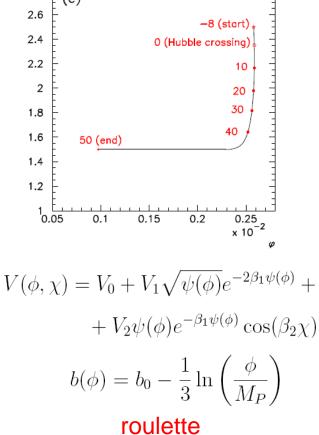
Equations

$$\ddot{Q}_{\sigma} + 3H\dot{Q}_{\sigma} + \left(\frac{k^2}{a^2} + C_{\sigma\sigma}\right)Q_{\sigma} + \frac{2V_s}{\dot{\sigma}}\dot{\delta s} + C_{\sigma s}\,\delta s = 0$$
$$\ddot{\delta s} + 3H\dot{\delta s} + \left(\frac{k^2}{a^2} + C_{ss}\right)\delta s - \frac{2V_s}{\dot{\sigma}}\dot{Q}_{\sigma} + C_{s\sigma}Q_{\sigma} = 0,$$

$$\begin{split} C_{\sigma\sigma} &= V_{\sigma\sigma} - \left(\frac{V_s}{\dot{\sigma}}\right)^2 + 2\frac{\dot{\sigma}V_{\sigma}}{M_P^2H} + \frac{3\dot{\sigma}^2}{M_P^2} - \frac{\dot{\sigma}^4}{2M_P^4H^2} - b_{\phi}\left(s_{\theta}^2c_{\theta}V_{\sigma} + (c_{\theta}^2 + 1)s_{\theta}V_s\right) \\ C_{\sigma s} &= 6H\frac{V_s}{\dot{\sigma}} + \frac{2V_{\sigma}V_s}{\dot{\sigma}^2} + 2V_{\sigma s} + \frac{\dot{\sigma}V_s}{M_P^2H} + 2b_{\phi}(s_{\theta}^3V_{\sigma} - c_{\theta}^3V_s) \\ C_{ss} &= V_{ss} - \left(\frac{V_s}{\dot{\sigma}}\right)^2 + b_{\phi}(1 + s_{\theta}^2)c_{\theta}V_{\sigma} + b_{\phi}c_{\theta}^2s_{\theta}V_s - \dot{\sigma}^2(b_{\phi\phi} + b_{\phi}^2) \\ C_{s\sigma} &= -6H\frac{V_s}{\dot{\sigma}} - \frac{2V_{\sigma}V_s}{\dot{\sigma}^2} + \frac{\dot{\sigma}V_s}{M_P^2H} \\ \cos\theta &\equiv \frac{\dot{\phi}}{\dot{\sigma}}, \quad \sin\theta \equiv \frac{\dot{\chi}e^b}{\dot{\sigma}} \quad \text{with} \quad \dot{\sigma} \equiv \sqrt{\dot{\phi}^2 + e^{2b}\dot{\chi}^2} \end{split}$$

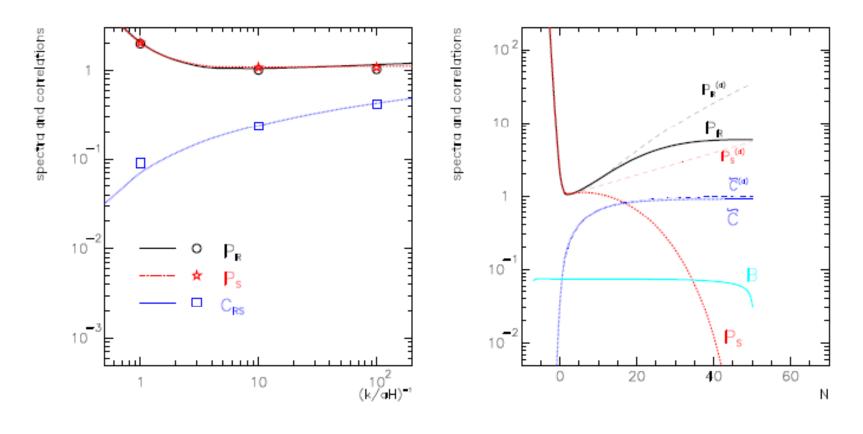
Examples





inflation Bond et al., 2006

Examples



roulette inflation

Examples

n_s	$1 - 6\epsilon_* + 2\eta_{\sigma\sigma*}$	single-field result	full result
double inflation (canonical)	0.929	0.982	0.967
double inflation (non-canonical)	0.953	0.968	0.934
roulette inflation	1.017	1.019	0.932

$$\mathcal{P}_{\mathcal{R}}^{\mathrm{sf}}(k) \simeq \frac{H^4}{4\pi^2 \dot{\sigma}^2} = \frac{H^4}{8\pi^2 \mathcal{L}_{\mathrm{kin}}}$$

Conclusions



good motivation to look for the implications of next-to-minimal models of inflation

richer dynamics; possible generation of isocurvature perturbations (after reheating)

cautionary tale: predictions may differ from the single-field case

We have a tool

(for studying the evolution of the multi-inflaton systems) and won't hesitate to use it!