# Local SU(5) Unification from the Heterotic String 

## Christoph Lüdeling <br> ITP, Universität Heidelberg

W. Buchmüller, CL, J. Schmidt, arXiv:0707.1651
(1) Introduction
(2) The Model
(3) Anomaly Cancellation
(4) Local GUT
(5) Outlook

## Introduction

- GUT: Attractive features:
- $\mathrm{SU}(3) \times \mathrm{SU}(2) \times \mathrm{U}(1) \subset \mathrm{SU}(5), \mathrm{SO}_{10} \ldots$, gauge couplings unify
- Unification matter into larger multiplets
- Large Higgs representations required
- Doublet-triplet-splitting
- Yukawa couplings do not unify
- Drawbacks can be addressed in higher-dimensional orbifold GUTs - Nice nocsihility. Heterotic Strins.
- $E_{8} \times E_{8}$ gauge symmetry included
- Simple orbifold compactifications with realistic four-dimensional matter content and gauge group possible
- UV completion


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[Kobayashi, Raby, Zhang; Buchmüller, Hamaguchi, Lebedev, Ratz; Kim, Kim, Kyae; Förste, Nilles, Vaudrevange, Wingerter, Ramos-Sanchez,...]


## Heterotic Orbifold Compactification

- Choose a torus with discrete isometry ("twist") with fixed points
- Mod out by this isometry, fixed points become singularities
- Fixing boundary conditions at fixed points requires embedding the twist into gauge group and choosing Wilson lines
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- Twisted sectors: States localised at fixed points


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## The Model: Geometry

[Buchmüller, Hamaguchi, Lebedev, Ratz]

- Torus: $\mathrm{G}_{2} \times \mathrm{SU}(3) \times \mathrm{SO}(4)$ root lattice, $\mathbb{Z}_{6-I I}=\mathbb{Z}_{3} \times \mathbb{Z}_{2}$ twist:
[Kobayashi,Raby,Zhang]

- Obtain effective 6D Theory on $T^{2} / \mathbb{Z}_{2}$ orbifold
- Internal zero modes and twisted states show up as bulk states, twisted states are localised at orbifold fixed points


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## String Scale



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## Anomalies

- Orbifold have bulk and brane anomalies
- Anomaly cancellation by Green-Schwarz mechanism requires factorisation of anomaly polynomials, $I_{8}=X_{4} Y_{4}$ and $I_{6}^{f}=X_{4}^{f} Y_{2}$
- $\mathcal{O}(500)$ conditions but guaranteed by string theorv (and modular invariance conditions on twist vectors and Wilson lines): Check of spectrum
- Anomalous U(1)'s induce localised FI terms

- These lead to localisation of bulk fields, break the $U(1)$ and need to be cancelled to obtain SUSY vacuum


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\begin{aligned}
& \xi_{0}=148\left(\frac{g M_{P}^{2}}{384 \pi^{2}}\right) \delta^{(2)}\left(z-z_{0}\right) \\
& \xi_{1}=80\left(\frac{g M_{P}^{2}}{384 \pi^{2}}\right) \delta^{(2)}\left(z-z_{1}\right)
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## Local SU(5) GUT

- Local GUT: At fixed points, boundary conditions break bulk gauge group to smaller groups,

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\text { in our case: } \quad \mathrm{SU}(6) \longrightarrow\left\{\begin{array}{c}
\mathrm{SU}(5) \\
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\end{array}\right.
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- Localised fields come in complete multiplets of local GUT group
- Due to other branes, bulk fields form split multiplets
- Due to higher symmetry, decoupling of exotics much more transparent that in four-dimensional limit


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## Projection



- On branes, SUSY is broken to $\mathcal{N}=1$
- Bulk Matter: Hypermultiplets, split as $H=\left(H_{L}, H_{R}\right)$ into chiral multiplet
- Bulk vector multiplets split as $V=(A, \phi)$ into vector and chiral multiplets


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- Bulk vector multiplets split as $V=(A, \phi)$ into vector and chiral multiplets
- Only one $\mathcal{N}=1$ multiplet survives projection


## Decoupling

- Several pairs of $\mathbf{5}+\overline{\mathbf{5}}$ and most exotics decoupled easily

| Bulk: | 5 | 51 | $\overline{5}_{0}^{\square}$ | $\overline{5}$ | $\overline{5}$ | $\overline{5}_{2}$ | $5{ }_{0}^{c}$ | $5{ }_{2}^{c}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{SU}(3) \times \mathrm{SU}(2)$ | $(1,2)$ | $(1,2)$ | $(3,1)$ | (1, 2) | $(1,2)$ | $(\overline{3}, 1)$ | $(\overline{3}, 1)$ | $(1,2)$ |
| $\mathrm{U}(1)_{B-L}$ | 0 | 0 | $-\frac{2}{3}$ | 0 | 0 | $-\frac{1}{3}$ | $\frac{2}{3}$ | -1 |
| MSSM content | $H_{u}$ |  |  |  | $H_{d}$ | $d_{3}$ |  | 13 |

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- Remaining 5's and $\overline{\mathbf{5}}$ 's:

| Bulk: | $\mathbf{5}$ | $\mathbf{5}_{1}$ | $\overline{5}_{0}^{c}$ | $\overline{\mathbf{5}}$ | $\overline{\mathbf{5}}_{1}$ | $\overline{\mathbf{5}}_{2}$ | $5_{0}^{c}$ | $\mathbf{5}_{2}^{c}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{SU}(3) \times \mathrm{SU}(2)$ | $(1, \mathbf{2})$ | $(1, \mathbf{2})$ | $(\mathbf{3}, 1)$ | $(1, \mathbf{2})$ | $(1, \mathbf{2})$ | $(\overline{\mathbf{3}}, 1)$ | $(\overline{\mathbf{3}}, 1)$ | $(1, \mathbf{2})$ |
| $\mathrm{U}(1)_{B-L}$ | 0 | 0 | $-\frac{2}{3}$ | 0 | 0 | $-\frac{1}{3}$ | $\frac{2}{3}$ | -1 |
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$$
\begin{gathered}
2 \times(\overline{\mathbf{5}}+\mathbf{1 0}) \text { generations on the branes } \\
2 \times(\overline{\mathbf{5}}+\mathbf{1 0}) \text { generations in the bulk } \\
\mathbf{5}+\overline{\mathbf{5}} \text { Higgses in the bulk }
\end{gathered}
$$

## Split Multiplets

- Bulk generations:

$$
\begin{array}{ll}
\overline{\mathbf{5}}_{(3)}=(\overline{\mathbf{3}}, 1)+(1, \mathbf{2}) & \mathbf{1 0}_{(3)}=(\mathbf{3}, \mathbf{2})+(\overline{\mathbf{3}}, 1)+(1,1) \\
\overline{\mathbf{5}}_{(4)}=(\overline{\mathbf{3}}, 1)+(1, \mathbf{2}) & \mathbf{1 0}_{(4)}=(\mathbf{3}, 2)+(\overline{\mathbf{3}}, 1)+(1,1)
\end{array}
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$$

$$
\mathbf{1 0}_{(3)}=(3,2)+(\overline{\mathbf{3}}, 1)+(1,1)
$$

$$
\mathbf{1 0}_{(4)}=(3,2)+(3,1)+(1,2)
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$$
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\overline{\mathbf{5}}_{(4)}=(\overline{\mathbf{3}}, 1)+(1,2) & \mathbf{1 0}_{(4)}=(\mathbf{3}, 2)+(\overline{3}, 1)+(1,1)
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One generation remains, avoiding $\operatorname{SU}(5)$ mass relations

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\mathbf{1 0}_{(4)}=(\mathbf{3}, 2)+(\overline{3}, \mathbf{2})+\left(1, \frac{1}{2}\right)
\end{array}
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One generation remains, avoiding $\operatorname{SU}(5)$ mass relations

- Higgses:

$$
\begin{aligned}
& \mathbf{5}_{u}=(\mathbf{3}, 1)+(1, \mathbf{2}) \\
& \overline{\mathbf{5}}_{d}=(\overline{\mathbf{3}}, 1)+(1, \mathbf{2})
\end{aligned}
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## Split Multiplets

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$$
\begin{aligned}
& \overline{5}_{(3)}=\left(\overline{3}{ }^{2}\right)+(1,2) \\
& 10_{(3)}=(3,2)+(\overline{3}, 1)+(1,1) \\
& \overline{\mathbf{5}}_{(4)}=(\overline{3}, 1)+(1,2) \\
& 10_{(4)}=(3,2)+(3,1)+(1,2)
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Orbifold projection solves doublet-triplet-splitting

## Yukawa Couplings

$$
\begin{gathered}
W=C_{(i j)}^{(u)} \mathbf{5}_{u} \mathbf{1 0}_{(i)} \mathbf{1 0 _ { ( j ) }}+C_{(i j)}^{(d)} \mathbf{5}_{d} \overline{\mathbf{5}}_{(i)} \mathbf{1 0}_{(j)} \\
C_{(i j)}^{(u)}=\left(\begin{array}{cccc}
a_{1} & 0 & a_{2} & a_{3} \\
0 & a_{1} & a_{2} & a_{3} \\
a_{2} & a_{2} & 0 & g \\
a_{3} & a_{3} & g & a_{4}
\end{array}\right), \quad C_{i j}^{(d)}=\left(\begin{array}{cccc}
0 & 0 & b_{1} & b_{2} \\
0 & 0 & b_{1} & b_{2} \\
b_{3} & b_{3} & b_{4} & 0 \\
b_{5} & b_{5} & b_{6} & b_{5}^{2}
\end{array}\right)
\end{gathered}
$$

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\end{array}\right), \quad C_{i j}^{(d)}=\left(\begin{array}{cccc}
0 & 0 & b_{1} & b_{2} \\
0 & 0 & b_{1} & b_{2} \\
b_{3} & b_{3} & b_{4} & 0 \\
b_{5} & b_{5} & b_{6} & b_{5}^{2}
\end{array}\right) \\
& a_{1}=\left\langle Y_{0}^{c} \bar{Y}_{0}^{c} S_{1} S_{3}\right\rangle, \quad a_{2}=\left\langle\left(\bar{Y}_{0}^{c} S_{1}\right)^{2} S_{5}\right\rangle, \quad a_{3}=\left\langle Y_{0}^{c} \bar{Y}_{0}^{c} S_{1} S_{3} S_{5}\right\rangle, \\
& a_{4}=\left\langle Y_{0}^{c} \bar{Y}_{0}^{c} S_{1} S_{3}\left(S_{5}\right)^{2}\right\rangle, \\
& b_{1}=\left\langle Y_{0} \bar{Y}_{1}\left(S_{5}\right)^{3}\left(S_{7}\right)^{2}\right\rangle, \quad b_{2}=\left\langle X_{1}^{c} \bar{Y}_{2}^{c} U_{1}^{c} S_{7}\right\rangle, \quad b_{3}=\left\langle X_{1}^{c} \bar{Y}_{1} S_{3}\left(S_{5} S_{7}\right)^{2}\right\rangle, \\
& b_{4}=\left\langle\left(X_{1}^{c}\right)^{2} \bar{Y}_{1} U_{1}^{c} S_{4} S_{7}\right\rangle, \quad b_{5}=\left\langle S_{5}\right\rangle, \quad b_{6}=\left\langle\left(X_{1}^{c}\right)^{2} Y_{1} S_{1} S_{7}\right\rangle
\end{aligned}
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## Yukawa Couplings

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\begin{gathered}
W=C_{(i j)}^{(u)} \mathbf{5}_{u} \mathbf{1 0}_{(i)} \mathbf{1 0}_{(j)}+C_{(i j)}^{(d)} \mathbf{5}_{d} \overline{\mathbf{5}}_{(i)} \mathbf{1 0}_{(j)} \\
C_{(i j)}^{(u)}=\left(\begin{array}{cccc}
a_{1} & 0 & a_{2} & a_{3} \\
0 & a_{1} & a_{2} & a_{3} \\
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0 & 0 & b_{1} & b_{2} \\
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b_{3} & b_{3} & b_{4} & 0 \\
b_{5} & b_{5} & b_{6} & b_{5}^{2}
\end{array}\right) \\
W=Y_{i j}^{u} h_{u} u_{i}^{c} q_{j}+Y_{i j}^{d} h_{d} d_{i}^{c} q_{j}+Y_{i j}^{\prime} h_{d} l_{i} e_{j}^{c} \\
Y_{i j}^{u}=\left(\begin{array}{ccc}
a_{1} & 0 & a_{3} \\
0 & a_{1} & a_{3} \\
a_{2} & a_{2} & g
\end{array}\right), \quad Y_{i j}^{d}=\left(\begin{array}{ccc}
0 & 0 & b_{2} \\
0 & 0 & b_{2} \\
b_{5} & b_{5} & b_{7}
\end{array}\right), \quad Y_{i j}^{\prime}=\left(\begin{array}{ccc}
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## Outlook

- Constructed local 6D GUT from the heterotic string
- Doublet-triplet splitting achieved easily, SU(5) mass relations avoided dur to split bulk multiplets
- More symmetry in $6 D \rightsquigarrow$ simple decoupling of unwanted states
- Supersymmetric vacuum: four-dimensional D-term vanishes
- Open Questions:


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- Phenomenology needs to be improved (CKM mixing, $R$-parity)
- Profiles of bulk fields due to localised FI terms
- Rlownum/resolution of sinmularitios, senoralisation to k3 internal space


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- Stabilisation of moduli, in particular, size of two-dimensional torus
$\qquad$


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- Phenomenology needs to be improved (CKM mixing, $R$-parity)
- Stabilisation of moduli, in particular, size of two-dimensional torus
- Profiles of bulk fields due to localised FI terms


## Outlook

- Constructed local 6D GUT from the heterotic string
- Doublet-triplet splitting achieved easily, SU(5) mass relations avoided dur to split bulk multiplets
- More symmetry in 6D simple decoupling of unwanted states
- Supersymmetric vacuum: four-dimensional $D$-term vanishes
- Open Questions:
- Phenomenology needs to be improved (CKM mixing, $R$-parity)
- Stabilisation of moduli, in particular, size of two-dimensional torus
- Profiles of bulk fields due to localised FI terms
- Blowup/resolution of singularities, generalisation to K3 internal space

