Local SU(5) Unification from the Heterotic String

Christoph Lüdeling ITP, Universität Heidelberg

W. Buchmüller, CL, J. Schmidt, arXiv:0707.1651

- Introduction
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- 3 Anomaly Cancellation
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- Outlook

- GUT: Attractive features:
 - $SU(3) \times SU(2) \times U(1) \subset SU(5), SO_{10} \dots$, gauge couplings unify
 - Unification matter into larger multiplets
- Drawbacks in 4d GUTS
 - Large Higgs representations required
 - Doublet-triplet-splitting
 - Yukawa couplings do not unify
- Drawbacks can be addressed in higher-dimensional orbifold GUTs
- Nice possibility: Heterotic String
 - E₈ × E₈ gauge symmetry included
 - Simple orbifold compactifications with realistic four-dimensional matter content and gauge group possible
 - UV completion

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[Kobayashi, Raby, Zhang; Buchmüller, Hamaguchi, Lebedev, Ratz; Kim, Kim, Kyae; Förste, Nilles, Vaudrevange, Wingerter, Ramos-Sanchez,...]

- Choose a torus with discrete isometry ("twist") with fixed points
- Mod out by this isometry, fixed points become singularities
- Fixing boundary conditions at fixed points requires embedding the twist into gauge group and choosing Wilson lines
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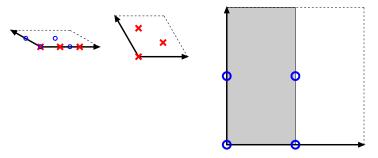
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The Model: Geometry

[Buchmüller, Hamaguchi, Lebedev, Ratz]

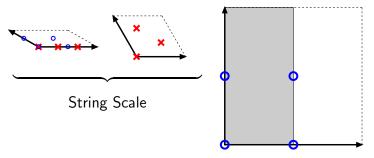
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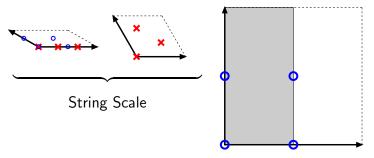
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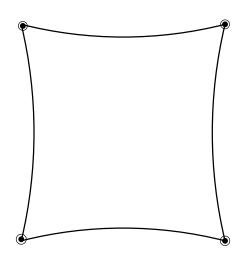
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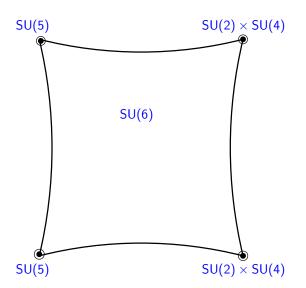
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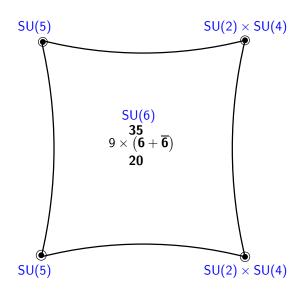
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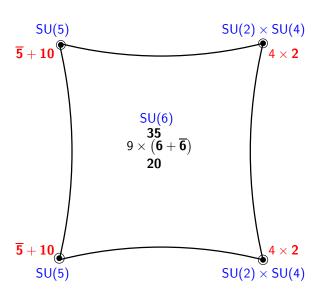


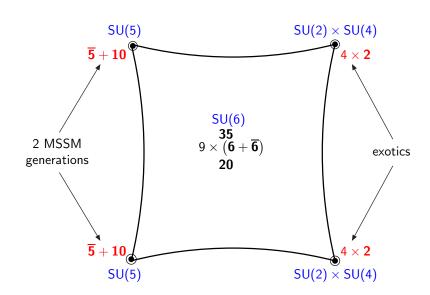
- Obtain effective 6D Theory on T^2/\mathbb{Z}_2 orbifold
- Internal zero modes and \mathbb{Z}_3 twisted states show up as bulk states, \mathbb{Z}_2 twisted states are localised at orbifold fixed points











Orbifold have bulk and brane anomalies

- Anomaly cancellation by Green–Schwarz mechanism requires factorisation of anomaly polynomials, $I_8 = X_4 Y_4$ and $I_6^f = X_4^f Y_2$
- $\mathcal{O}(500)$ conditions, but guaranteed by string theory (and modular invariance conditions on twist vectors and Wilson lines): Check of spectrum
- Anomalous U(1)'s induce localised FI terms

$$\xi_0 = 148 \left(\frac{gM_{\rm P}^2}{384\pi^2} \right) \delta^{(2)}(z - z_0)$$

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• These lead to localisation of bulk fields, break the U(1) and need to be cancelled to obtain SUSY vacuum [Lee, Nilles, Zucker]

in our case:
$$SU(6) \longrightarrow \left\{ \begin{array}{c} SU(5) \\ SU(2) \times SU(4) \end{array} \right.$$

- In zero mode spectrum, only the intersection of local groups survives, which is $G_{SM} = SU(3) \times SU(2) \times U(1)$
- Localised fields come in complete multiplets of local GUT group
- Due to other branes, bulk fields form split multiplets
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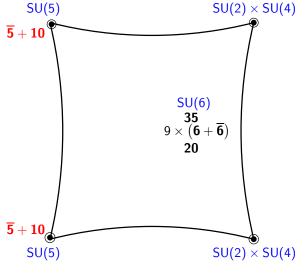
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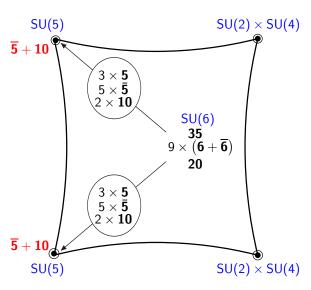
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Projection



- ullet On branes, SUSY is broken to $\mathcal{N}=1$
- Bulk Matter: Hypermultiplets, split as H = (H_L, H_R) into chiral multiplet
- Bulk vector multiplets split as $V = (A, \phi)$ into vector and chiral multiplets
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Decoupling

- \bullet Several pairs of ${\bf 5}+\bar{\bf 5}$ and most exotics decoupled easily
- Remaining $\mathbf{5}$'s and $\mathbf{\bar{5}}$'s:

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| Bulk: | 5 | 5 ₁ | 5 ^c ₀ | 5 | $\bar{5}_{1}$ | $\bar{5}_2$ | 5 ^c ₀ | 5 ^c ₂ |
|----------------------|----------------|-----------------------|------------------------------------|---------------|----------------|----------------|------------------------------------|------------------------------------|
| Zero modes: | | | | | | | | |
| $SU(3) \times SU(2)$ | (1, 2) | (1, 2) | (3,1) | (1, 2) | (1, 2) | $(\bar{3},1)$ | $(\bar{3},1)$ | (1, 2) |
| $U(1)_{B-L}$ | 0 | 0 | $-\frac{2}{3}$ | 0 | 0 | $-\frac{1}{3}$ | <u>2</u> 3 | -1 |
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 $2\times\left(\boldsymbol{\bar{5}}+\boldsymbol{10}\right)$ generations on the branes $2\times\left(\boldsymbol{\bar{5}}+\boldsymbol{10}\right)$ generations in the bulk $\boldsymbol{5}+\boldsymbol{\bar{5}}$ Higgses in the bulk

• Bulk generations:

$$\begin{split} \mathbf{\bar{5}}_{(3)} &= \left(\mathbf{\bar{3}}, 1\right) + (1, \mathbf{2}) \\ \mathbf{\bar{5}}_{(4)} &= \left(\mathbf{\bar{3}}, 1\right) + (1, \mathbf{2}) \end{split} \qquad \quad \mathbf{10}_{(3)} &= \left(\mathbf{3}, \mathbf{2}\right) + \left(\mathbf{\bar{3}}, 1\right) + (1, 1) \\ \mathbf{10}_{(4)} &= \left(\mathbf{3}, \mathbf{2}\right) + \left(\mathbf{\bar{3}}, 1\right) + (1, 1) \end{split}$$

Higgses

$$\mathbf{5}_u = (\mathbf{3}, \mathbf{1}) + (\mathbf{1}, \mathbf{2})$$

$$ar{\mathbf{5}}_d = ig(ar{\mathbf{3}},1ig) + ig(1,\mathbf{2}ig)$$

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One generation remains, avoiding SU(5) mass relations

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Orbifold projection solves doublet-triplet-splitting

Yukawa Couplings

$$W = C_{(ij)}^{(u)} \mathbf{5}_{u} \mathbf{10}_{(i)} \mathbf{10}_{(j)} + C_{(ij)}^{(d)} \mathbf{5}_{d} \mathbf{\bar{5}}_{(i)} \mathbf{10}_{(j)}$$

$$C_{(ij)}^{(u)} = \begin{pmatrix} a_1 & 0 & a_2 & a_3 \\ 0 & a_1 & a_2 & a_3 \\ a_2 & a_2 & 0 & \mathbf{g} \\ a_3 & a_3 & \mathbf{g} & a_4 \end{pmatrix}, \qquad C_{ij}^{(d)} = \begin{pmatrix} 0 & 0 & b_1 & b_2 \\ 0 & 0 & b_1 & b_2 \\ b_3 & b_3 & b_4 & 0 \\ b_5 & b_5 & b_6 & b_5^2 \end{pmatrix}$$

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$$\begin{split} &a_1 = \langle Y_0^c \bar{Y}_0^c S_1 S_3 \rangle, & a_2 = \langle \left(\bar{Y}_0^c S_1 \right)^2 S_5 \rangle, & a_3 = \langle Y_0^c \bar{Y}_0^c S_1 S_3 S_5 \rangle, \\ &a_4 = \langle Y_0^c \bar{Y}_0^c S_1 S_3 \left(S_5 \right)^2 \rangle, & \\ &b_1 = \langle Y_0 \bar{Y}_1 \left(S_5 \right)^3 \left(S_7 \right)^2 \rangle, & b_2 = \langle X_1^c \bar{Y}_2^c U_1^c S_7 \rangle, & b_3 = \langle X_1^c \bar{Y}_1 S_3 \left(S_5 S_7 \right)^2 \rangle, \\ &b_4 = \langle \left(X_1^c \right)^2 \bar{Y}_1 U_1^c S_4 S_7 \rangle, & b_5 = \langle S_5 \rangle, & b_6 = \langle \left(X_1^c \right)^2 Y_1 S_1 S_7 \rangle \end{split}$$

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$$W = Y^u_{ij} h_u u^c_i q_j + Y^d_{ij} h_d d^c_i q_j + Y^I_{ij} h_d l_i e^c_j$$

$$Y^{\textit{u}}_{ij} = \begin{pmatrix} a_1 & 0 & a_3 \\ 0 & a_1 & a_3 \\ a_2 & a_2 & \textit{g} \end{pmatrix}, \quad \ \ \, \begin{array}{c} \textbf{Y}^{\textit{d}}_{ij} = \begin{pmatrix} 0 & 0 & b_2 \\ 0 & 0 & b_2 \\ b_5 & b_5 & b_7 \end{pmatrix}, \quad \ \, \textbf{Y}^{\textit{I}}_{ij} = \begin{pmatrix} 0 & 0 & b_1 \\ 0 & 0 & b_1 \\ b_3 & b_3 & b_4 \end{pmatrix}$$



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