Can LHC Test the See-Saw Mechanism?

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Based on J.K., A.Yu. Smirnov, arXiv:0705.3221 [hep-ph]

Outline

- Introduction
- Cancellation of Neutrino Masses and Underlying Symmetries
- Signals at Colliders
- Conclusions

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The See-Saw Mechanism

Standard Model (or MSSM) + right-handed neutrinos ν_{R}

- Singlets under all gauge groups
 → Very large Majorana masses m_R possible
- Yukawa couplings to Higgs and lepton doublets
 → Electroweak-scale Dirac masses m_D

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Mass eigenstates:

- Very light Majorana neutrinos, $m_{\nu} = -m_{\rm D} m_{\rm R}^{-1} m_{\rm D}^{T}$
- Very heavy ones with masses ∼ m_R

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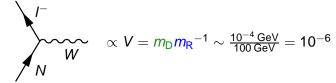
- Very light Majorana neutrinos, $m_{\nu} = -m_{\rm D} m_{\rm R}^{-1} m_{\rm D}^{T}$
- Very heavy ones with masses ∼ m_R
- Experimental limit: $m_{\nu} \lesssim 0.1 \text{ eV}$
- Common assumption: $\mathcal{O}(1)$ Yukawa couplings
 - $\Rightarrow m_{\rm R} \gtrsim 10^{14} \, {\rm GeV}$
 - ⇒ Mechanism not directly testable

Electroweak-Scale Singlets

- What if $m_{\rm R} \sim 100 \, {\rm GeV?}$ $m_{\rm D} \sim 10^{-4} \, {\rm GeV} = 100 \, {\rm keV} \sim m_{\rm e}$
 - → Not totally unreasonable
 - ⇒ RH neutrinos may be within reach of LHC and ILC

Electroweak-Scale Singlets

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 - → Not totally unreasonable
 - ⇒ RH neutrinos may be within reach of LHC and ILC
- Yukawa couplings tiny ⇒ irrelevant for colliders
- Gauge interactions via mixing, e.g.



Observation at colliders needs V ≥ 0.01

Han, Zhang, PRL **97** (2006); del Aguila, Aguilar-Saavedra, Pittau, J. Phys. Conf. Ser. **53** (2006); Bray, Lee, Pilaftsis, hep-ph/0702294

 \Rightarrow no way?



- Cancellation of Neutrino Masses and Underlying Symmetries
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- Contributions from different singlets to m_{ν} can cancel Buchmüller, Wyler, PLB **249** (1990); Pilaftsis, Z. Phys. **C55** (1992)
- 3 singlets: $m_{\nu}=0$ if and only if

•
$$m_D$$
 has rank 1, $m_D = m \begin{pmatrix} y_1 & y_2 & y_3 \\ \alpha y_1 & \alpha y_2 & \alpha y_3 \\ \beta y_1 & \beta y_2 & \beta y_3 \end{pmatrix}$

Buchmüller, Greub, NPB **363** (1991); Ingelman, Rathsman, Z. Phys. **C60** (1993); Heusch, Minkowski, NPB **416** (1994)

Less Naive Point of View

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- Size of Yukawa couplings arbitrary ⇒ large mixing allowed
- Experimental limit: V ≤ 0.1
- Cancellation at least at the level $10^{-8} \Rightarrow$ severe fine-tuning → Symmetry motivation?

Lepton Number Conservation

Most straightforward: conserved lepton number

Wyler, Wolfenstein, NPB 218 (1983); Bernabéu, Santamaria, Vidal, Mendez, Valle, PLB 187 (1987); Tommasini, Barenboim, Bernabéu, Jarlskog, NPB 444 (1995); Pilaftsis, PRL 95 (2005); Pilaftsis, Underwood, PRD 72 (2005)

$$L(\nu_{L}) = 1, L(\nu_{R}^{1}) = 1, L(\nu_{R}^{2}) = -1, L(\nu_{R}^{3}) = 0$$

$$\Rightarrow m_{R} = \begin{pmatrix} 0 & M & 0 \\ M & 0 & 0 \\ 0 & 0 & M_{3} \end{pmatrix}, m_{D} = m \begin{pmatrix} a & 0 & 0 \\ b & 0 & 0 \\ c & 0 & 0 \end{pmatrix}$$

- $\nu_{\rm R}^1, \nu_{\rm R}^2$ form a Dirac neutrino with mass M
- $\nu_{\rm p}^3$ is decoupled

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Are there symmetries realizing the cancellation without L conservation?

Cancellation Without *L* Conservation?

2 or 3 singlets with equal masses involved in cancellation \Rightarrow *L* conservation \rightsquigarrow try $M_1 \neq M_2$

- Suppose singlets $\nu_{\rm R}^1$ and $\nu_{\rm R}^2$ participate in cancellation
- Some symmetry $\rightsquigarrow m_{\nu} = 0$ at scale M_2
- Symmetry broken below M₂

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- Running of contributions from $\nu_{\rm R}^1$ and $\nu_{\rm R}^2$ different Antusch, J.K., Lindner, Ratz, PLB 538 (2002); Antusch, J.K., Lindner, Ratz, Schmidt, JHEP 03 (2005)
 - ⇒ Cancellation unstable

$$m_{
u} \sim$$
 100 keV In $\frac{M_2}{M_1}$ at M_1

Cancellation Without L Conservation?

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$$m_{
u} \sim 100 \, \mathrm{keV} \, \ln \frac{M_2}{M_1} \, \mathrm{at} \, M_1$$

- ⇒ Singlets must be degenerate
- ⇒ Lepton number must be conserved

$$m_{\mathsf{R}} = \left(\begin{array}{ccc} 0 & M & 0 \\ M & 0 & 0 \\ 0 & 0 & M_3 \end{array} \right) \; , \; m_{\mathsf{D}} = m \left(\begin{array}{ccc} a & 0 & 0 \\ b & 0 & 0 \\ c & 0 & 0 \end{array} \right)$$

$$\begin{aligned} \textit{m}_{R} &= \begin{pmatrix} \epsilon_{1}\textit{M} & \textit{M} & \epsilon_{13}\textit{M} \\ \textit{M} & \epsilon_{2}\textit{M} & \epsilon_{23}\textit{M} \\ \epsilon_{13}\textit{M} & \epsilon_{23}\textit{M} & \textit{M}_{3} \end{pmatrix} \;,\; \textit{m}_{D} = \textit{m} \begin{pmatrix} \textit{a} & \delta_{\textit{a}} & \epsilon_{\textit{a}} \\ \textit{b} & \delta_{\textit{b}} & \epsilon_{\textit{b}} \\ \textit{c} & \delta_{\textit{c}} & \epsilon_{\textit{c}} \end{pmatrix} \end{aligned}$$

 $\epsilon_2, \delta_{a,b,c} \lesssim 10^{-10}$ for max $(a,b,c) \sim 1$, $\frac{m}{M} \sim 0.1$

Perturbations Leading to Non-Zero Neutrino Masses

$$m_{\mathsf{R}} = \begin{pmatrix} \epsilon_1 M & M & \epsilon_{13} M \\ M & \epsilon_2 M & \epsilon_{23} M \\ \epsilon_{13} M & \epsilon_{23} M & M_3 \end{pmatrix} , m_{\mathsf{D}} = m \begin{pmatrix} a & \delta_a & \epsilon_a \\ b & \delta_b & \epsilon_b \\ c & \delta_c & \epsilon_c \end{pmatrix} \equiv m (v \ v_\delta \ v_\epsilon)$$

$$\epsilon_2, \delta_{a,b,c} \lesssim 10^{-10}$$
 for max $(a,b,c) \sim 1$, $\frac{m}{M} \sim 0.1$

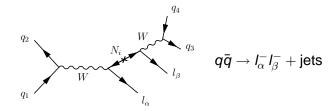
- Most general case: more parameters than observables
- Restricted cases, e.g. assuming similar size for all ϵ, δ :

$$m_{
u} pprox rac{m^2}{M} \left[\epsilon_2 \, v v^T - \left(v v_{\delta}^T + v_{\delta} v^T \right) \right]$$

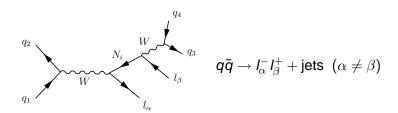
- Strong mass hierarchy
- Leading-order Yukawa couplings determined by observables
- Examples studied in leptogenesis context
 Raidal, Strumia, Turzyński, PLB 609 (2005); Pilaftsis, Underwood, PRD 72 (2005)

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Lepton Number Violation



- $m_{\nu} = 0$ due to symmetry $\Rightarrow L$ conservation \Rightarrow leading-order cross-section vanishes
- *L*-violating perturbations $\Rightarrow m_{\nu} \neq 0 \Rightarrow$ tiny
- ⇒ Unobservable without fine-tuning



- L conservation ⇒ no cancellation possible
- Strong constraints from searches for LFV decays, especially $\mu \to e \gamma \Rightarrow$ best candidate: $\mu^- \tau^+$
- ⇒ Observable in principle

Probably not at LHC

del Aguila, Aguilar-Saavedra, Pittau, hep-ph/0703261

Testing the See-Saw Mechanism

- m_{ν} small due to cancellation, not due to see-saw
- Colliders probe leading-order Yukawa couplings, not perturbations giving $m_{\nu} \neq 0$
- General case: no relation to neutrino masses and mixings

→ Decoupling of collider physics and neutrino mass generation

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→ Decoupling of collider physics and neutrino mass generation

Restricted cases:

- Strong neutrino mass hierarchy
- Leading-order Yukawas related to m_{ν}
- Correlations between LFV amplitudes
- Possible verification: measure V directly at e⁺e⁻ collider

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Conclusions

- Considered type-I see-saw scenario with light singlets
- Not considered: Right-handed neutrinos with additional interactions
- Naive expectation: Yukawa couplings tiny ⇒ unobservable
- Sizable couplings \Rightarrow cancellation needed for small m_{ν}
- Requires either fine-tuning or lepton number conservation
- Small neutrino masses due to tiny perturbations
- Colliders: lepton number violation not observable in untuned scenario
- Lepton flavour violation possibly observable
- Neutrino mass generation and collider physics decoupled in general
- Connection possible in constrained setups

