

Quark Compositeness in ATLAS



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Quark compositeness in ATLAS

- To describe quark compositeness a simple isoscalar left-left four-fermion contact interaction was added to QCD (as an analogy to Fermi approximation):

$$L_{qqqq} = \frac{\eta g^2}{2\Lambda_q^2} \overline{\Psi}_q^L \gamma^\mu \Psi_q^L \overline{\Psi}_q^L \gamma^\mu \Psi_q^L,$$

$$L = L_{QCD} + L_{qqqq}$$

- Λ (TeV) – scale, η – interference sign (+1 used here, destructive), $g^2 = 4\pi$.
- Are quark composite? Tevatron Run I: not up to $\Lambda \sim 2$ TeV.
- ATLAS detector with $E_{\text{CMS}} = 14$ TeV, calorimeter $\eta_{\text{max}} = 4.9$ and ability to measure high- p_T jets is especially suitable to extend our knowledge further.

Quark compositeness

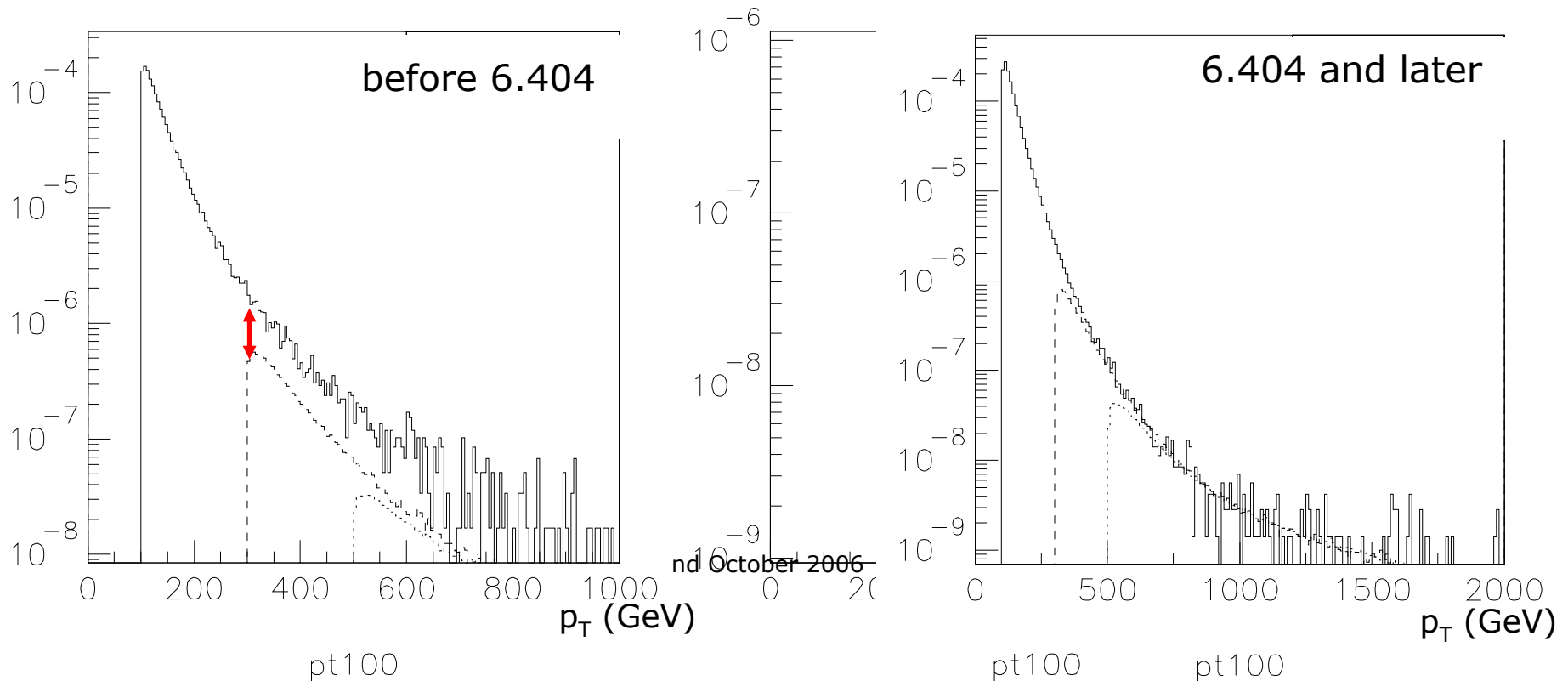
- Simulated samples:
 - $\Lambda = 3, 5, 10, 20$ and 40 TeV, QCD.
 - All quarks composite
 - Events generated in Pythia (Rome settings), fast simulation done in Athena 11.0.41 framework, using Atlfast. Analysis done on Analysis Object Data (AOD).

- What is investigated:
 - **Inclusive jet production cross-section $d\sigma/dp_T$.** For high p_T the contact term above (CT) causes excess of events above standard QCD p_T spectrum. p_T spectrum is sensitive to PDF uncertainties and systematics (e.g. calorimeter nonlinearity). This effect can mask or fake compositeness scenario.

 - **Dijet angular distribution.** Contact term causes excess of events with small pseudorapidity. Smaller dependence on systematics.

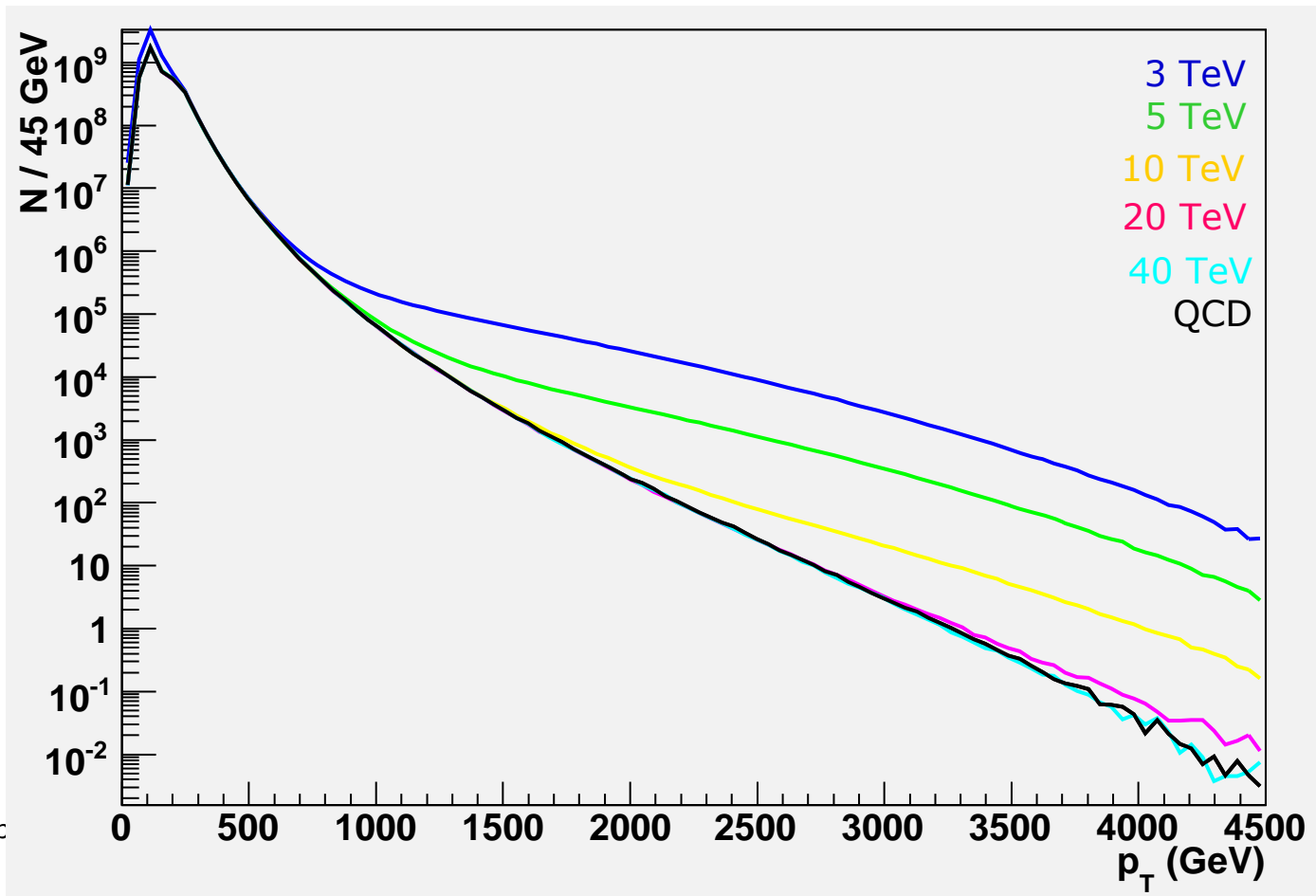
About Pythia

- Switch MSEL=51 used
 - problem with ISR encountered
 - thanks to T. Davidek for help
 - solved in 6.404 (thanks to T.Sjostrand)
 - solution compiled to Athena 11.0.41



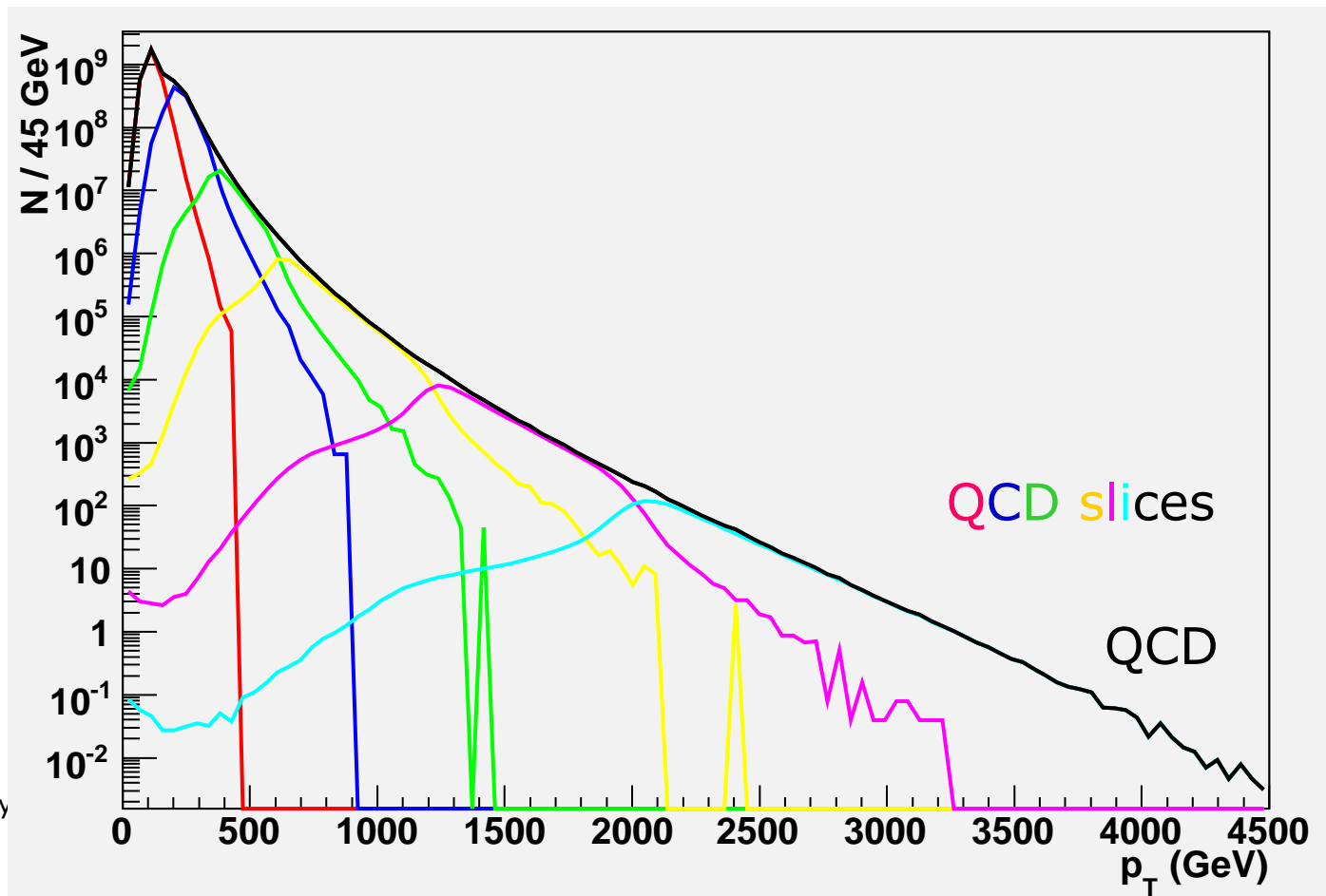
Inclusive jet production cross-section

- Inclusive leading dijet p_T spectrum for various Λ .
- Integral luminosity 20 fb^{-1} .



Inclusive jet production cross-section

- In order to cover such a large p_T range (from trigger 2j350 to TeV scale), the data had to be sewn from several p_T slices.



Inclusive jet production cross-section

- To characterize the excess of high- p_T events one can use:

$$R = \left(\frac{N(E_T > E_T^0)}{N(E_T < E_T^0)} \right)_{CT+QCD} \bigg/ \left(\frac{N(E_T > E_T^0)}{N(E_T < E_T^0)} \right)_{QCD}$$

- $E_T^0 := 1100$ GeV to optimize R_{dist} . For 10 fb^{-1} :

Λ (TeV)	R	σR	R_{dist}
3	3.39	0.02	145
5	1.81	0.01	78
10	1.05	<0.01	8.0
20	1.01	<0.01	2.2
40	1.004	0.006	0.62

To quantify the distance from QCD spectrum:

$$R_{dist} = \frac{R(\Lambda) - R(SM)}{\sigma_{R(\Lambda)}}$$

Inclusive jets – discovery limits

- Int. luminosities to achieve $R_{\text{dist}} = 3$

Λ (TeV)	3	5	10	20	40
L (fb ⁻¹)	4.3 pb ⁻¹	15 pb ⁻¹	1.4 fb ⁻¹	19 fb ⁻¹	234 fb ⁻¹

- R_{dist} values for L = 300 fb⁻¹.

Λ (TeV)	3	5	10	20	40
R_{dist}	794	427	44	12	3.4

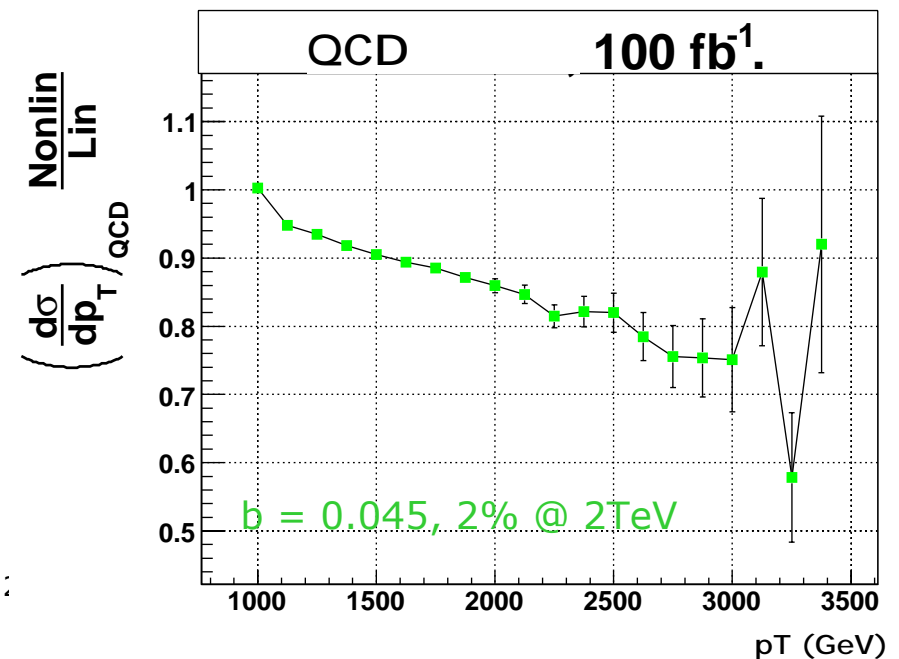
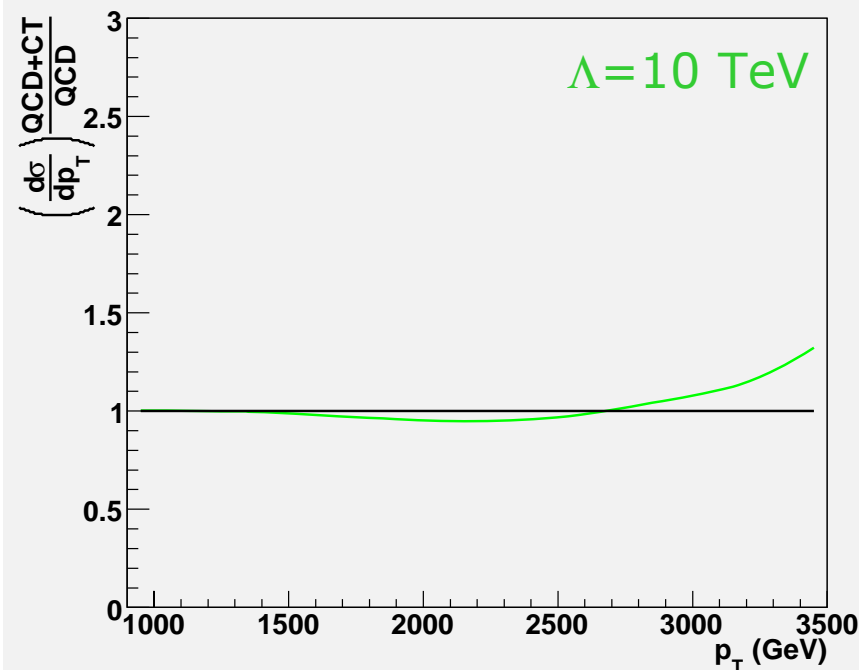
- $\Lambda = 3, 5, 10$ TeV might be ruled out or verified with first tens of pb⁻¹ of good data.
- But *no systematics* is included (PDF, nonlinearity,...)
- Therefore the required L will be larger, in case of $\Lambda = 40$ TeV the discovery is still unclear.

Inclusive jets – calorimeter nonlinearity

- What effect a calorimeter nonlinearity will make?
- To parametrize nonlinearity of ATLAS hadronic calorimeter (simple method):

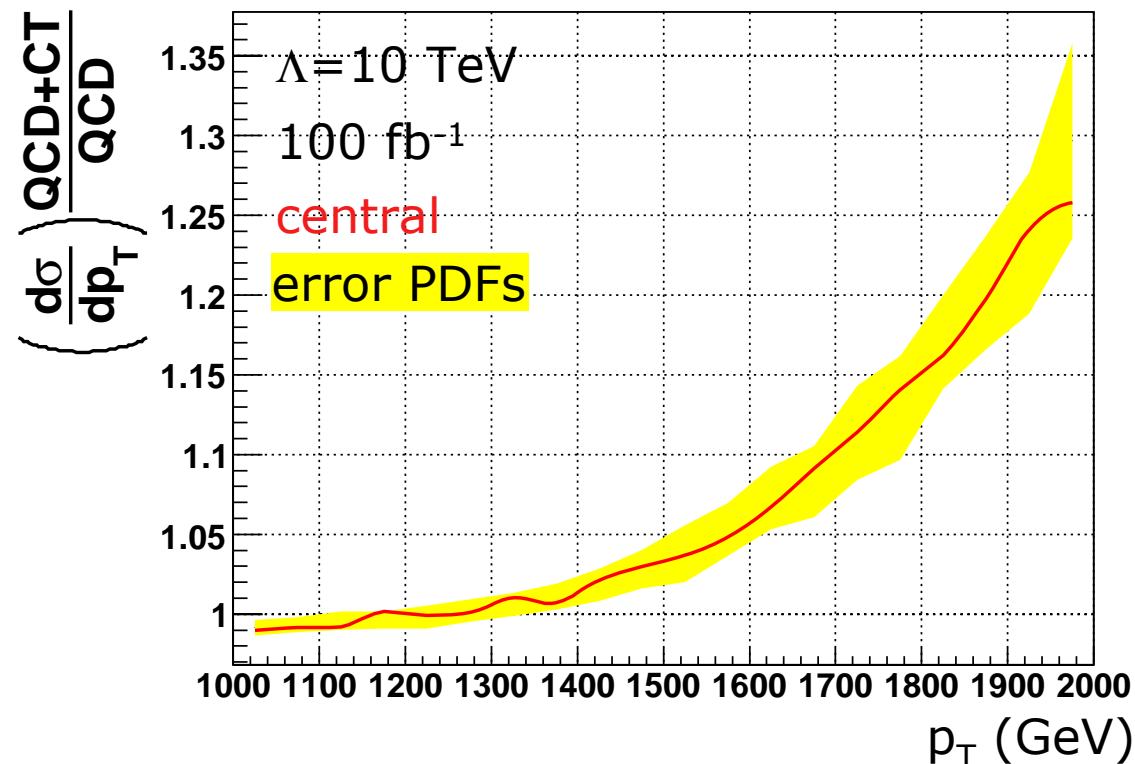
$$E_T(\text{meas.}) = E_T \frac{1}{c(1+(e/h-1)b \ln E_T)}$$

- e/h - noncompensation, b - nonlinearity (smaller values achievable by e.g. weighting method)
- c makes nonlin. and lin. spectra equal at 500 GeV.



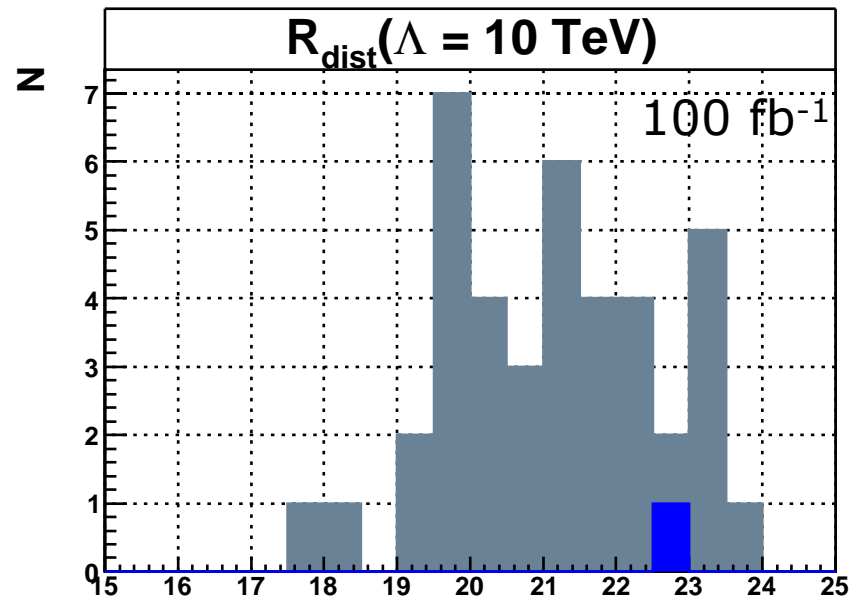
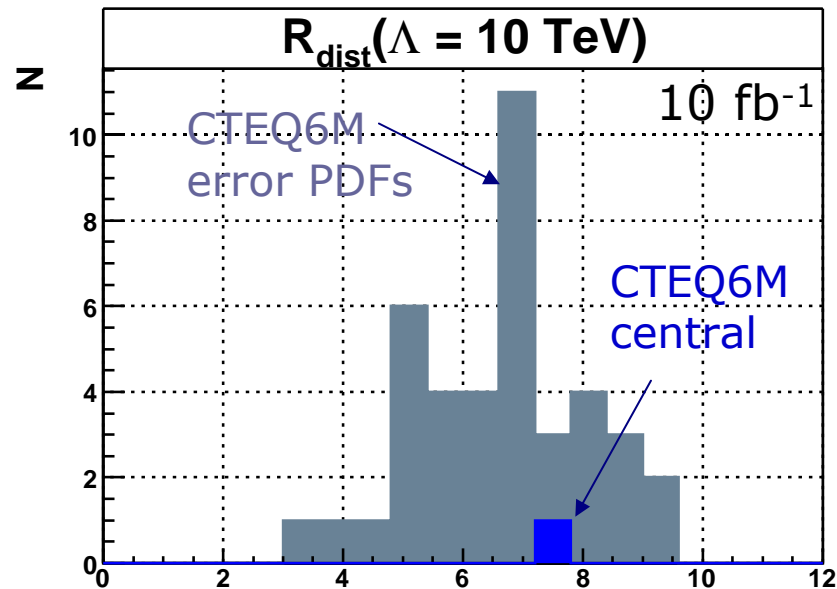
PDF uncertainty studies

- For the purpose of uncertainty calculations CTEQ6M PDFs were used. These are based on NLO calculations fitted to DIS data.
- The global fit of data is 20 parametric, thanks to that we have 40 error PDFs (+ one central value) that were used to generate the data below.
- PDF uncertainty studies done with Pythia 6.326, but repaired for ISR, $p_T > 1$ TeV.



PDF uncertainty studies

- R_{dist} can be calculated for each PDF:



- Systematic error due to PDF uncertainties in this case $\sigma_{\text{PDF}}(R_{\text{dist}}) = 1.40$. Compare it to $R_{\text{dist}}(\Lambda = 40 \text{ TeV}, 300 \text{ fb}^{-1}) = 3.40$.

Dijet angular distribution - R_χ

- We need a variable less sensitive to calo nonlinearity and more unique to compositeness (inclusive jet c.s is similar to graviton)- dijet angular distribution.

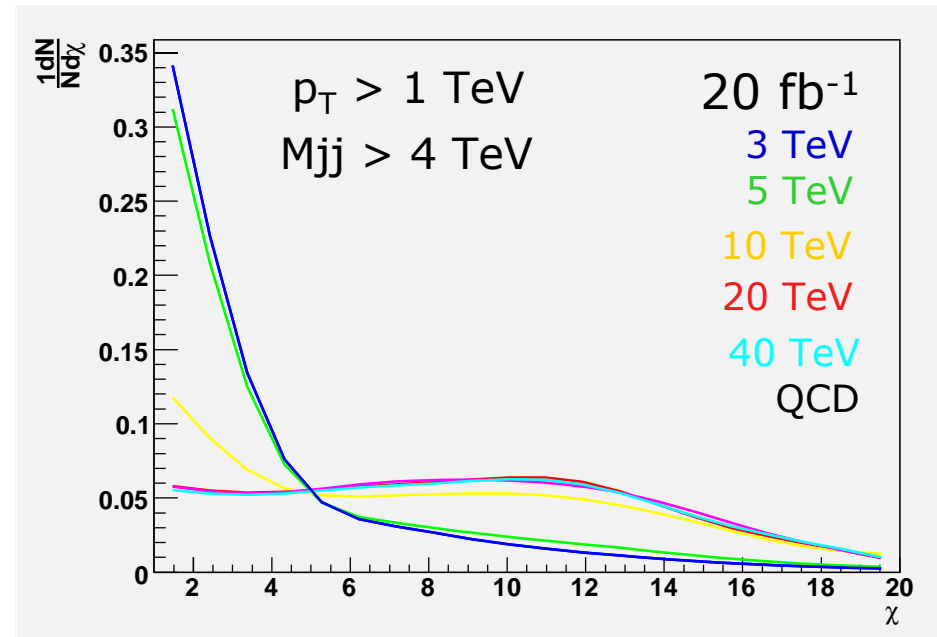
- Two leading jets with η_1, η_2 .

$$\chi = e^{|\eta_1 - \eta_2|}$$

- To characterize this distribution:

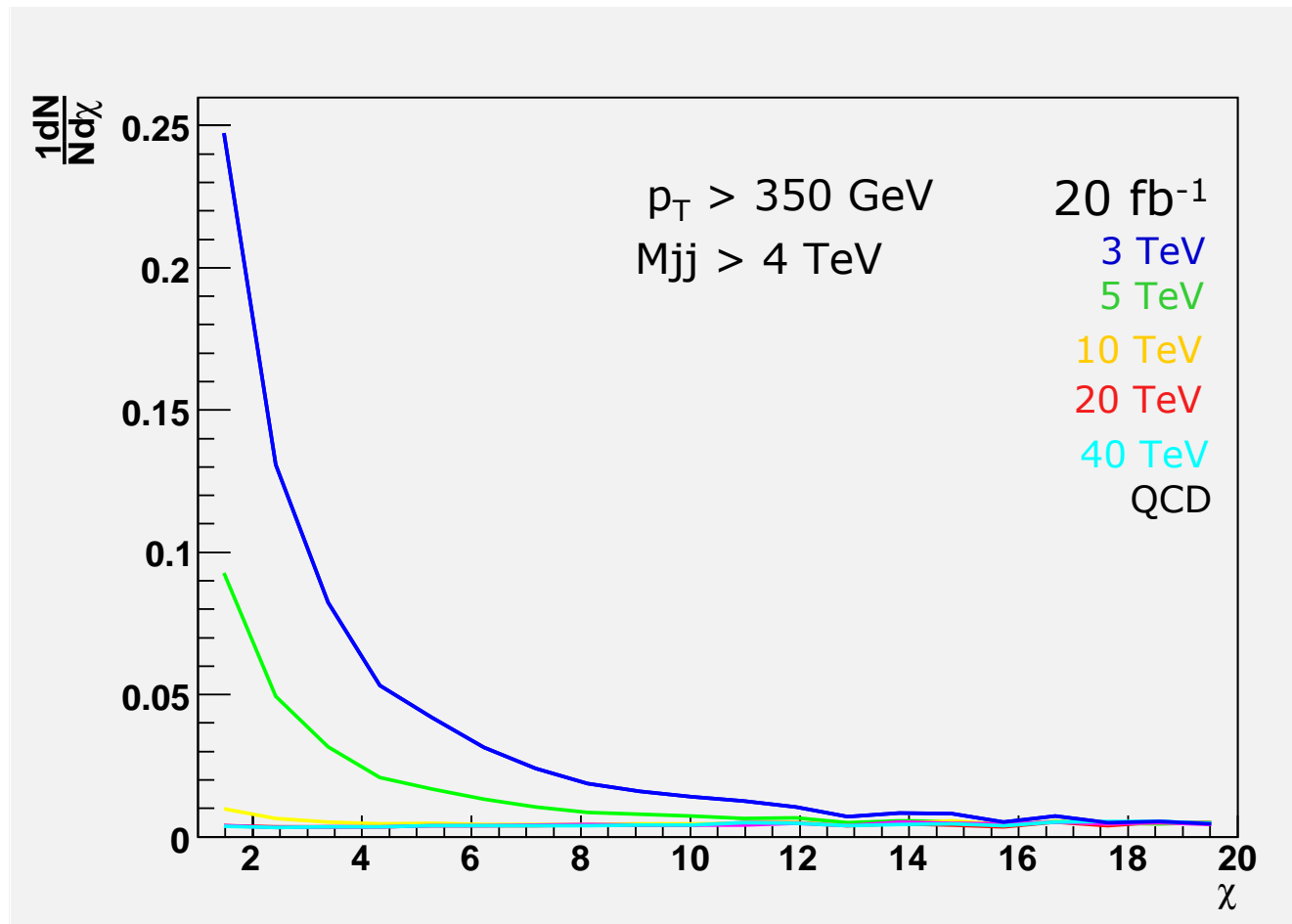
$$R_\chi = N(\chi < \chi_{cut}) / N(\chi > \chi_{cut})$$

- $\chi_{cut} = 2.8$ (to get the largest difference between Λ and SM)



Dijet angular distribution

- It is worth using again $p_T > 350$ GeV (see later).



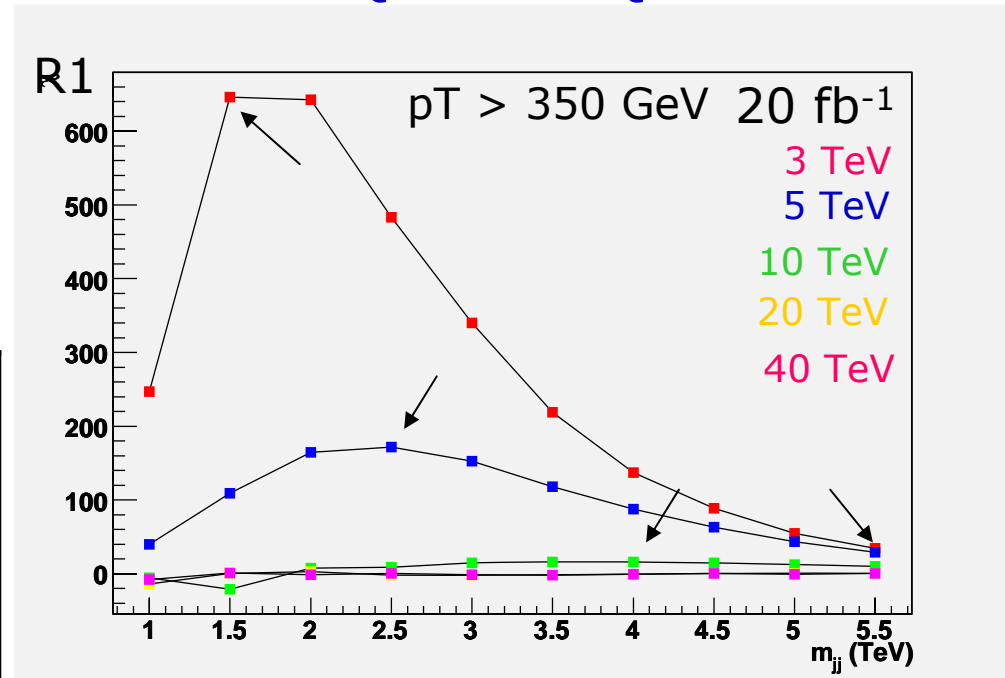
R1

- Way to quantify difference between SM QCD and QCD+CT:

$$R_1 = \frac{R_\chi(\Lambda) - R_\chi(SM)}{\sqrt{\sigma_\Lambda^2 + \sigma_{SM}^2}}$$

- For 20 fb⁻¹:

$\Lambda(\text{TeV})$	$R1_{350}$	$R1_{1000}$
3	646	100
5	172	71
10	16	10
20	2.3	0.7
40	0.65	<0.1



- Dijet invariant mass (m_{jj}) lower cut tuned also to optimum for each Λ .

R1 – discovery limits

- Int. luminosities to achieve $R1 = 3$

Λ (TeV)	3	5	10	20	40
L (fb ⁻¹)	< 1 pb ⁻¹	6 pb ⁻¹	0.7 fb ⁻¹	34 fb ⁻¹	426 fb ⁻¹

- R1 values for $L = 300$ fb⁻¹.

Λ (TeV)	3	5	10	20	40
R _{dist}	2500	665	62	8.9	2.5

- $\Lambda = 3, 5, 10$ TeV might be ruled out or verified with first tens of pb⁻¹ of good data.
- But *no systematics* is included (PDF, nonlinearity,...)
- Therefore the required L will be larger, in case of $\Lambda = 40$ TeV the discovery is unclear.

R1 – attempt without gluons

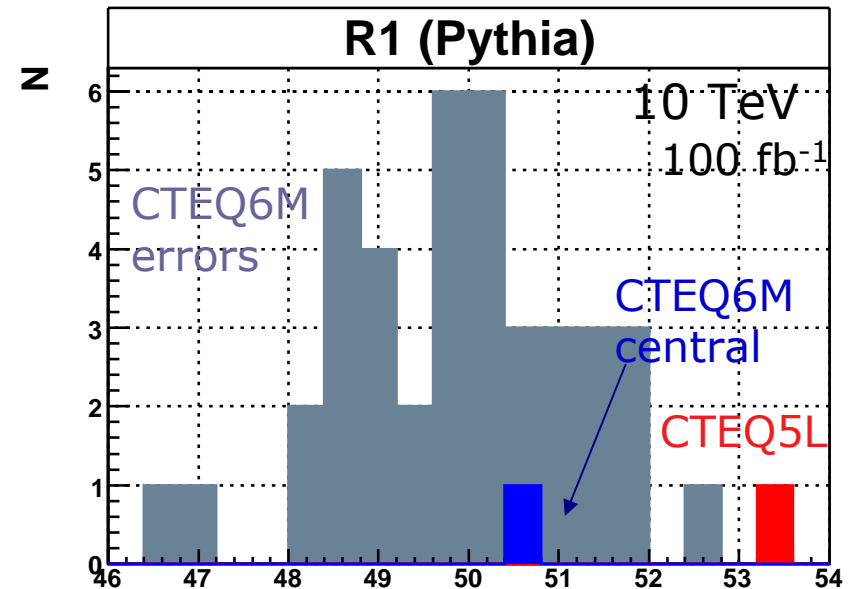
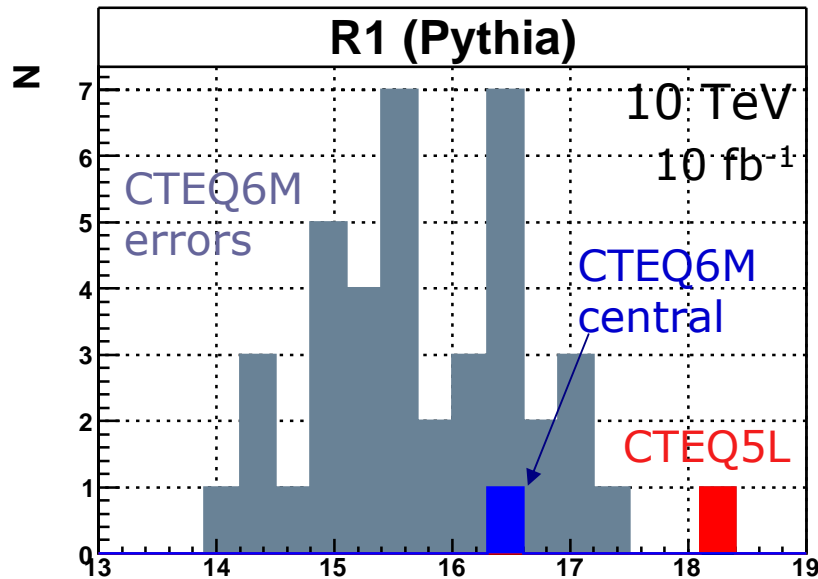
- Gluons do not enter the CT. When omitting the gluon jets, the discovery potential increases (20 fb^{-1}):

$\Lambda(\text{TeV})$	R1 no gluons	R1 w/ gluons
3	700	646
5	210	172
10	18	16
20	2.5	2.3
40	0.70	0.65

- Gluon jets have different jet shapes. The efficiency of spotting such a jet still has to be studied.

No systematic errors included.

PDF uncertainties (R1)



- R1 from Pythia (2 partons with highest p_T), $\Lambda = 10$ TeV, $m_{jj} = 4$ TeV, $p_T > 1$ TeV.
- Systematic error due to PDF uncertainties in this case $\sigma_{PDF}(R1) = 0.88$. That is comparable to $R1(\Lambda=40\text{TeV}, 30 \text{ fb}^{-1}) = 0.80$.
- Preliminary to say it is less sensitive than R_{dist} .

Conclusions

- Early limits: $\Lambda \sim 10$ TeV might be discovered or ruled out with first tens of pb^{-1} of good data.
- $\Lambda \sim 20$ TeV still should be visible with larger int. luminosity.
- Discovering $\Lambda \sim 40$ TeV requires better energy linearity than 2% @ 2TeV.
- Systematic errors still need to be understood in this study.
- To be continued ...