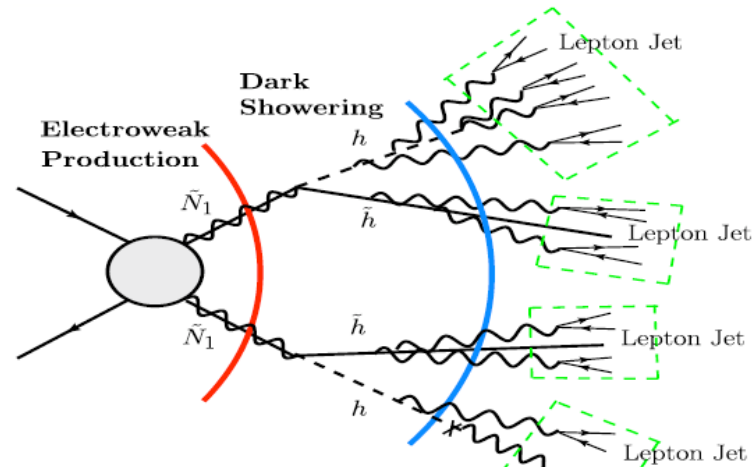


Hunting for Lepton Jets

Itay Yavin

New York University

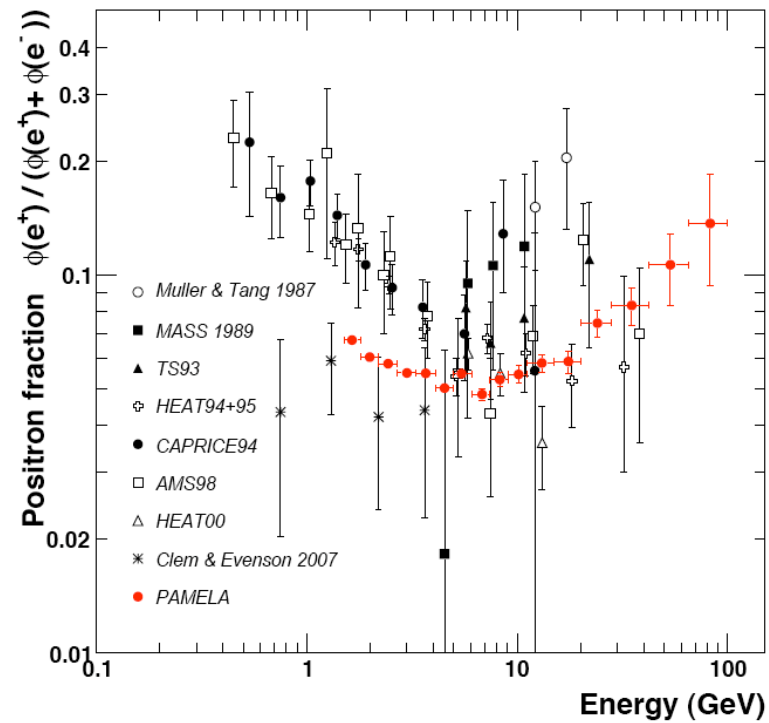


M. Baumgart, C. Cheung, J. T. Ruderman, L. T. Wang and I. Y. 0901.0283 [hep-ph]

C. Cheung, J. T. Ruderman, L. T. Wang and I. Y. 0909.0290[hep-ph]

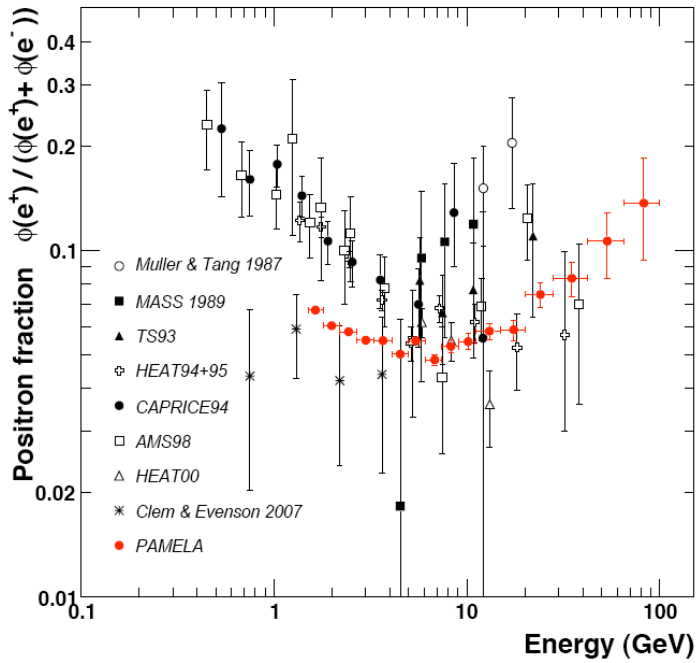
Part I

Introduction

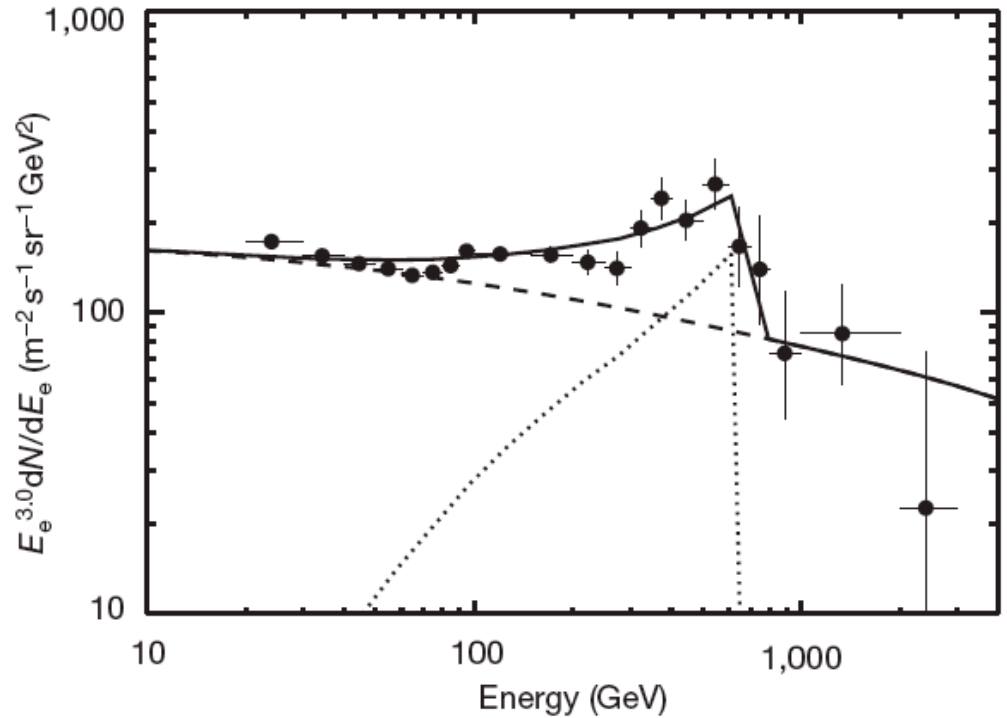
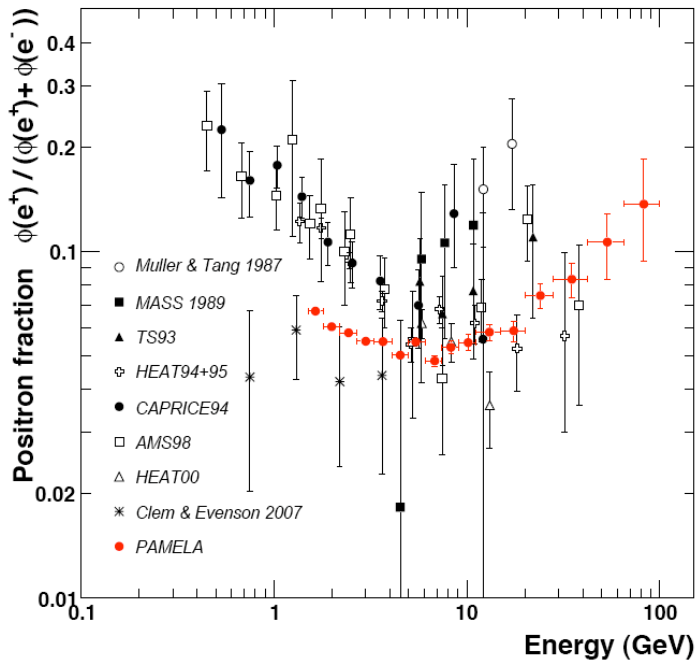


Motivation 1

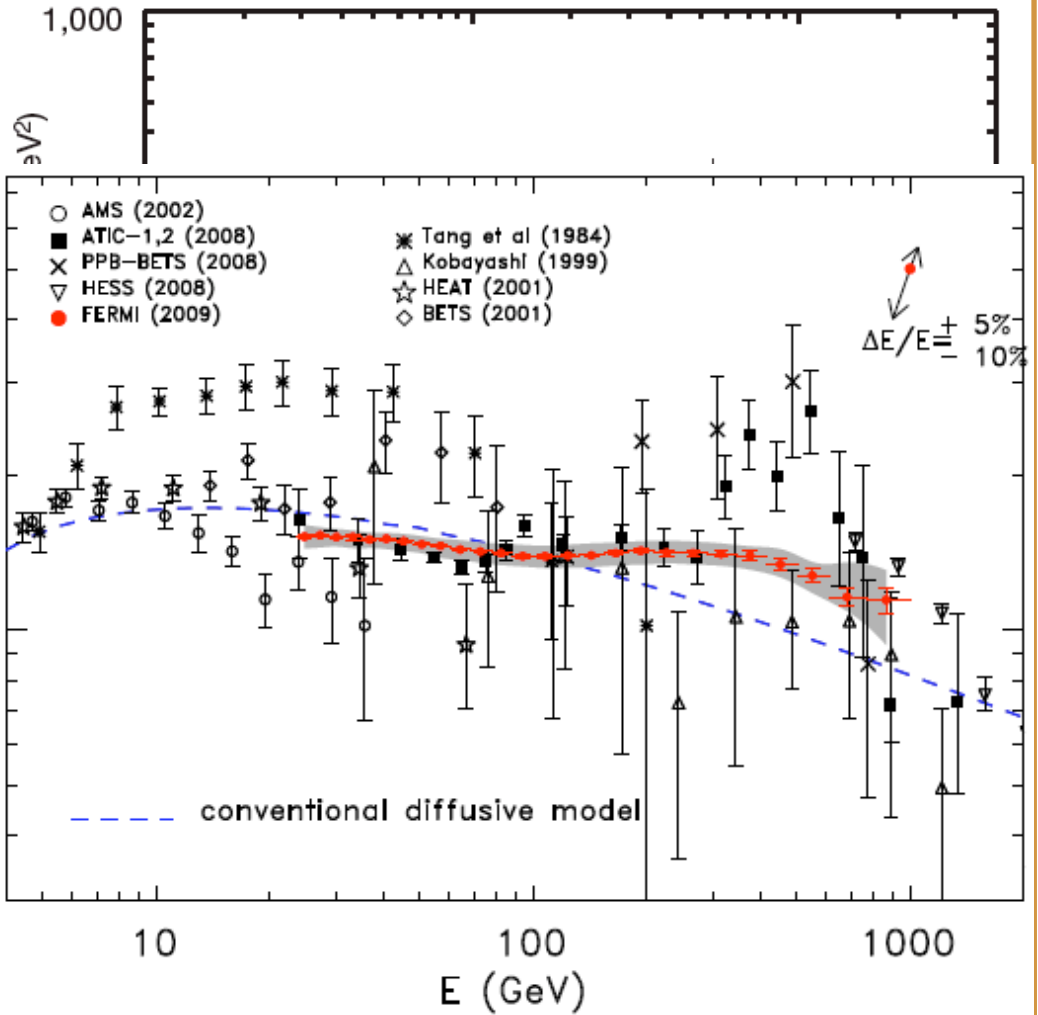
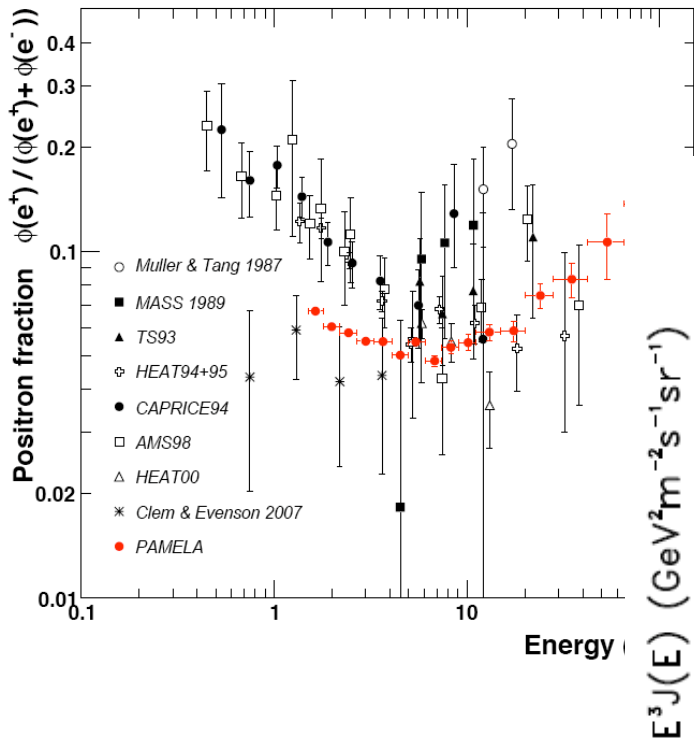
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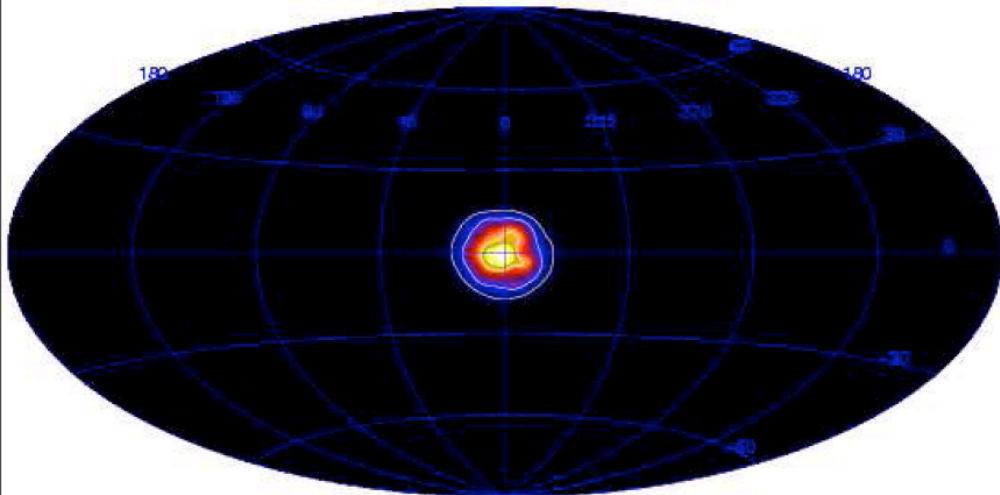
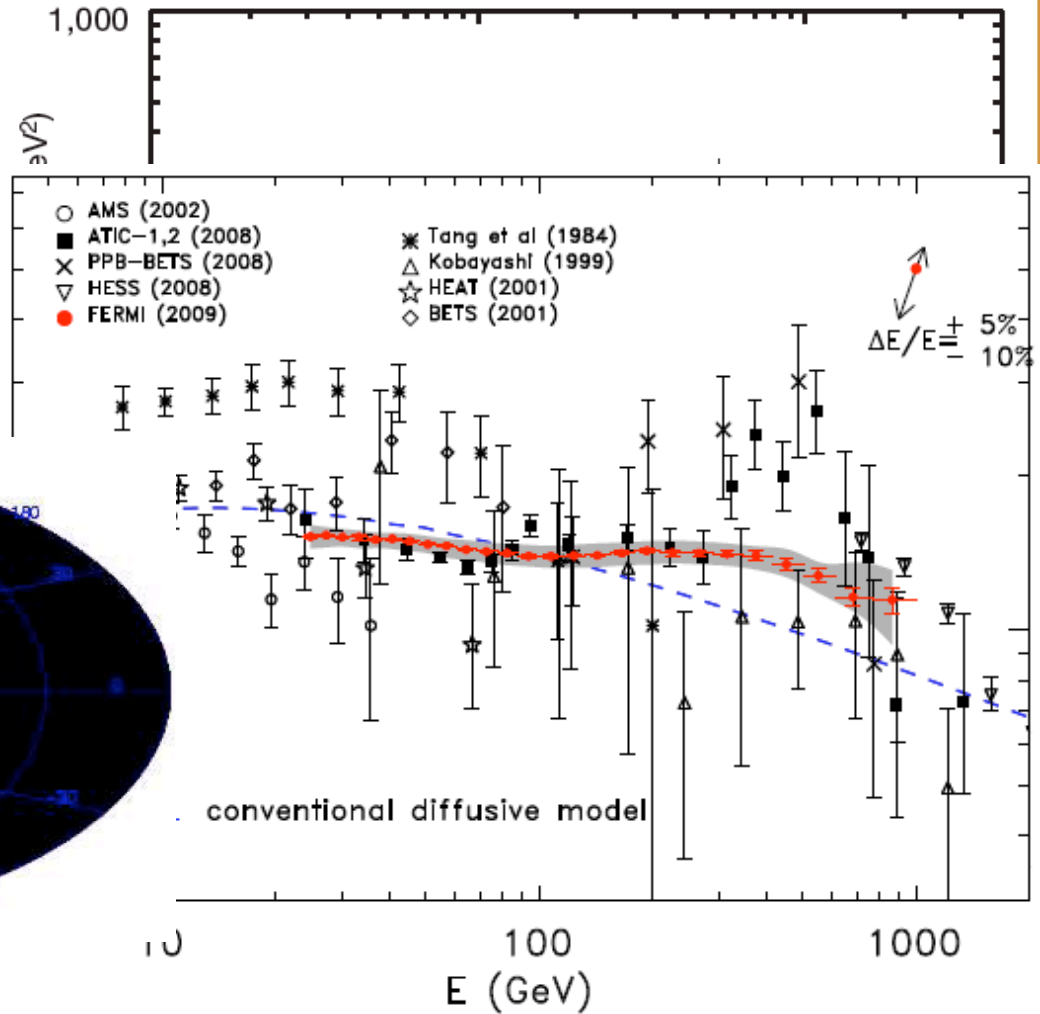
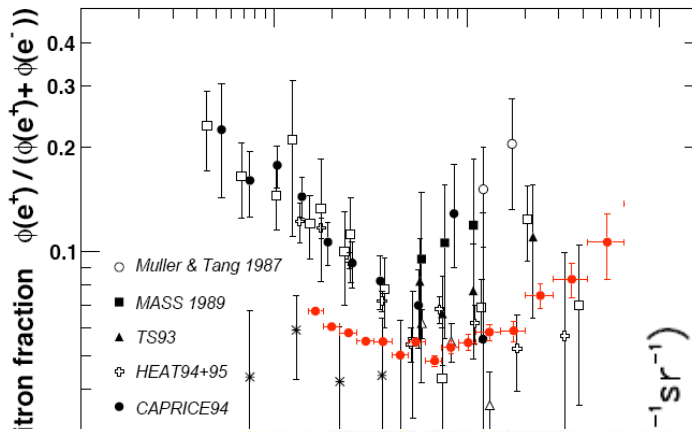
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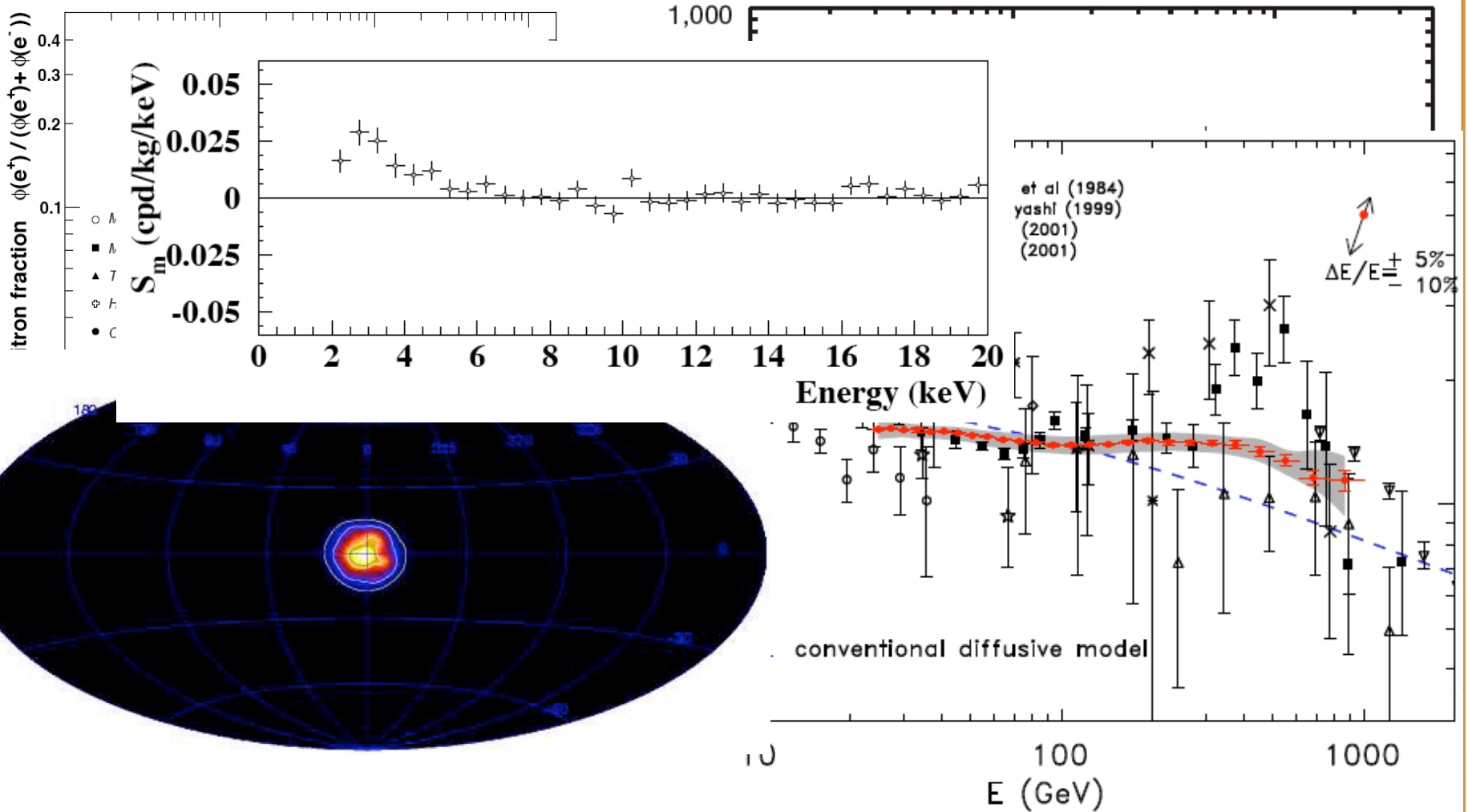
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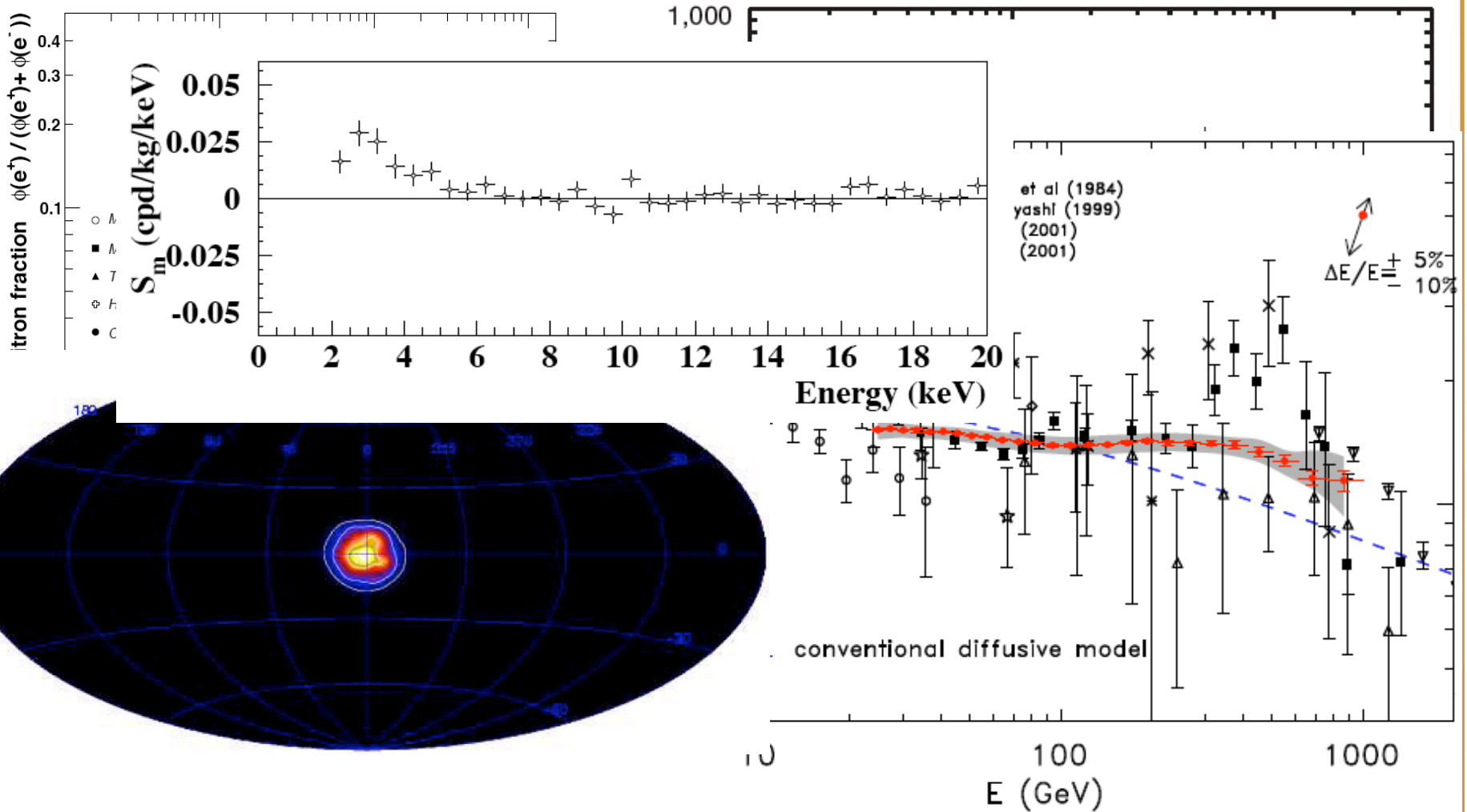
Motivation 1



Motivation 1



Motivation 1



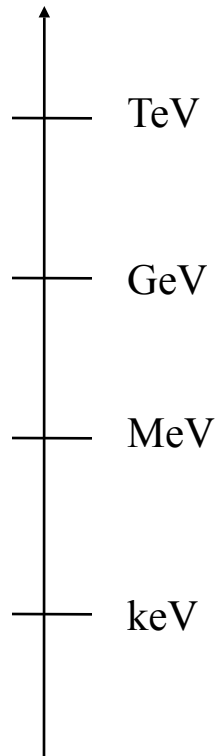
A Theory of DM - Arkani-Hamed, Finkbeiner, Slatyer, and Weiner (0810.0713)

Motivation 2

Aside from the recent astrophysical observations, there can be another motivation for looking for such objects.

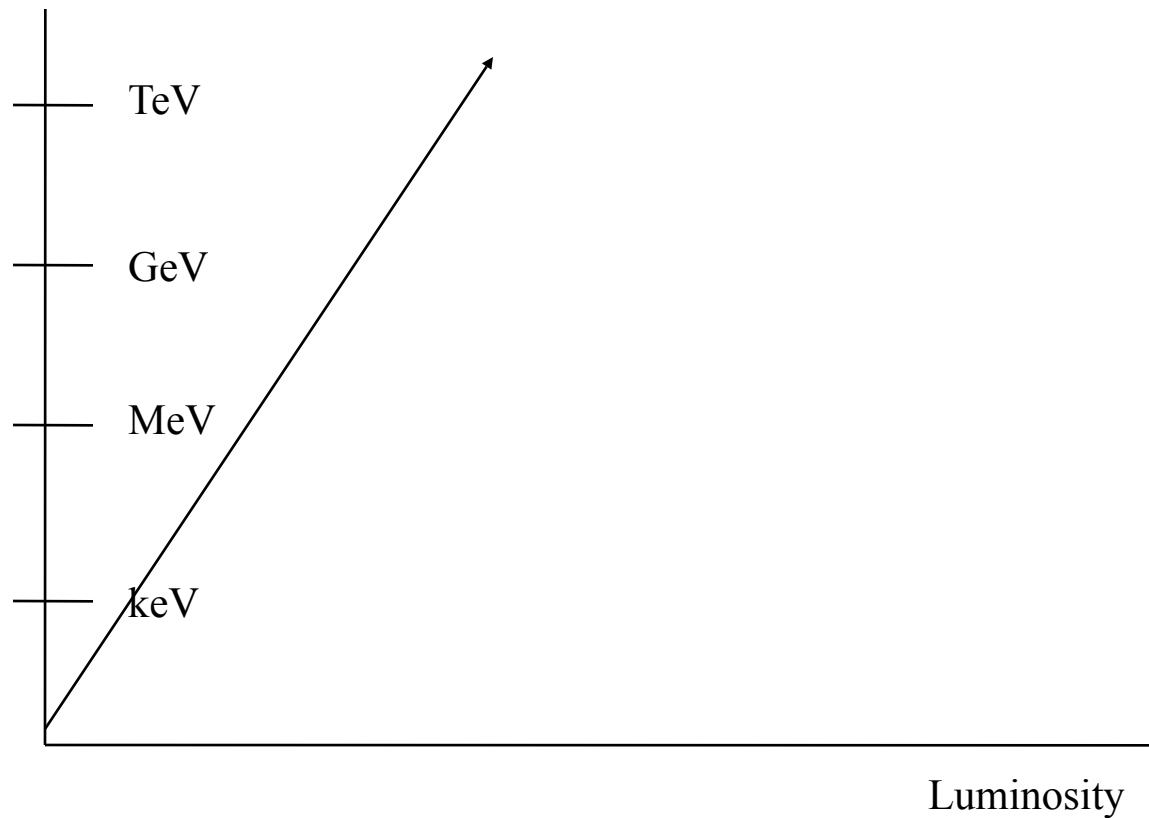
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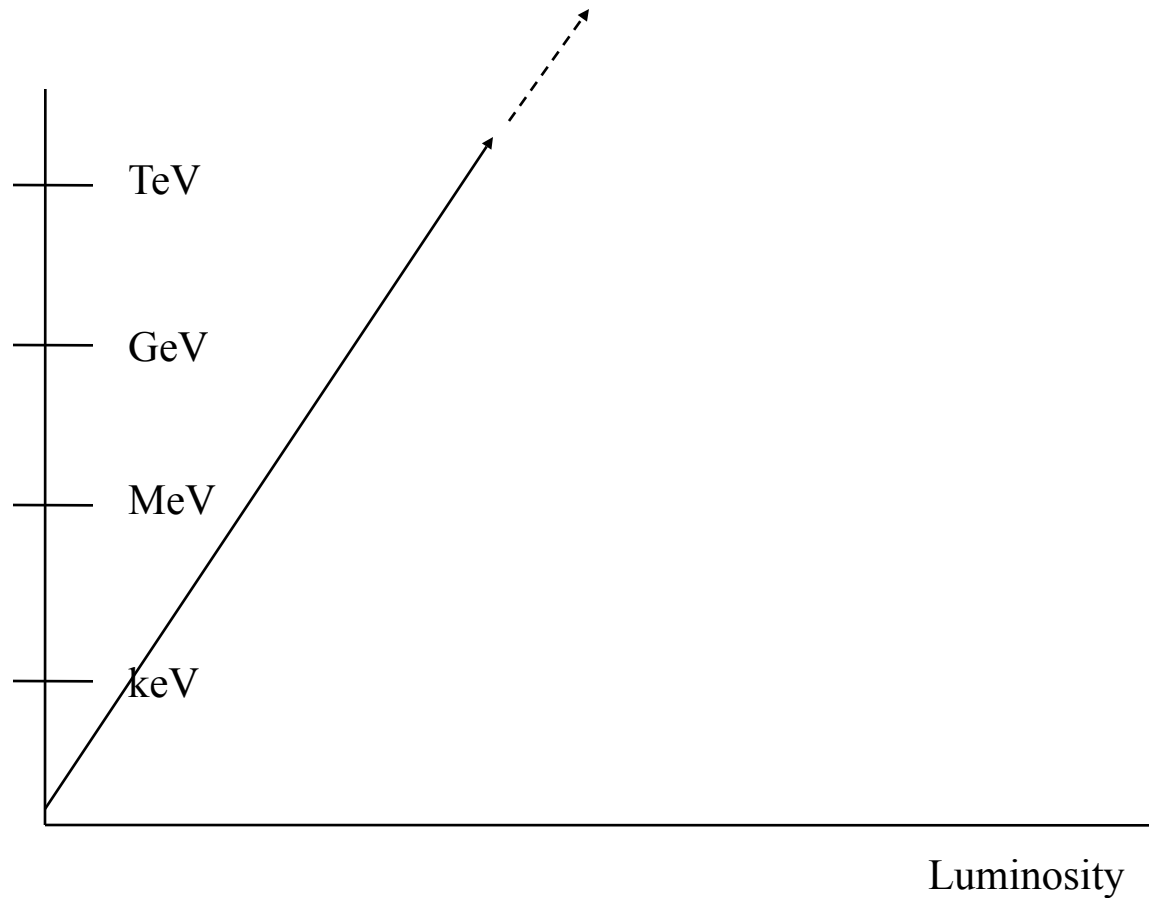
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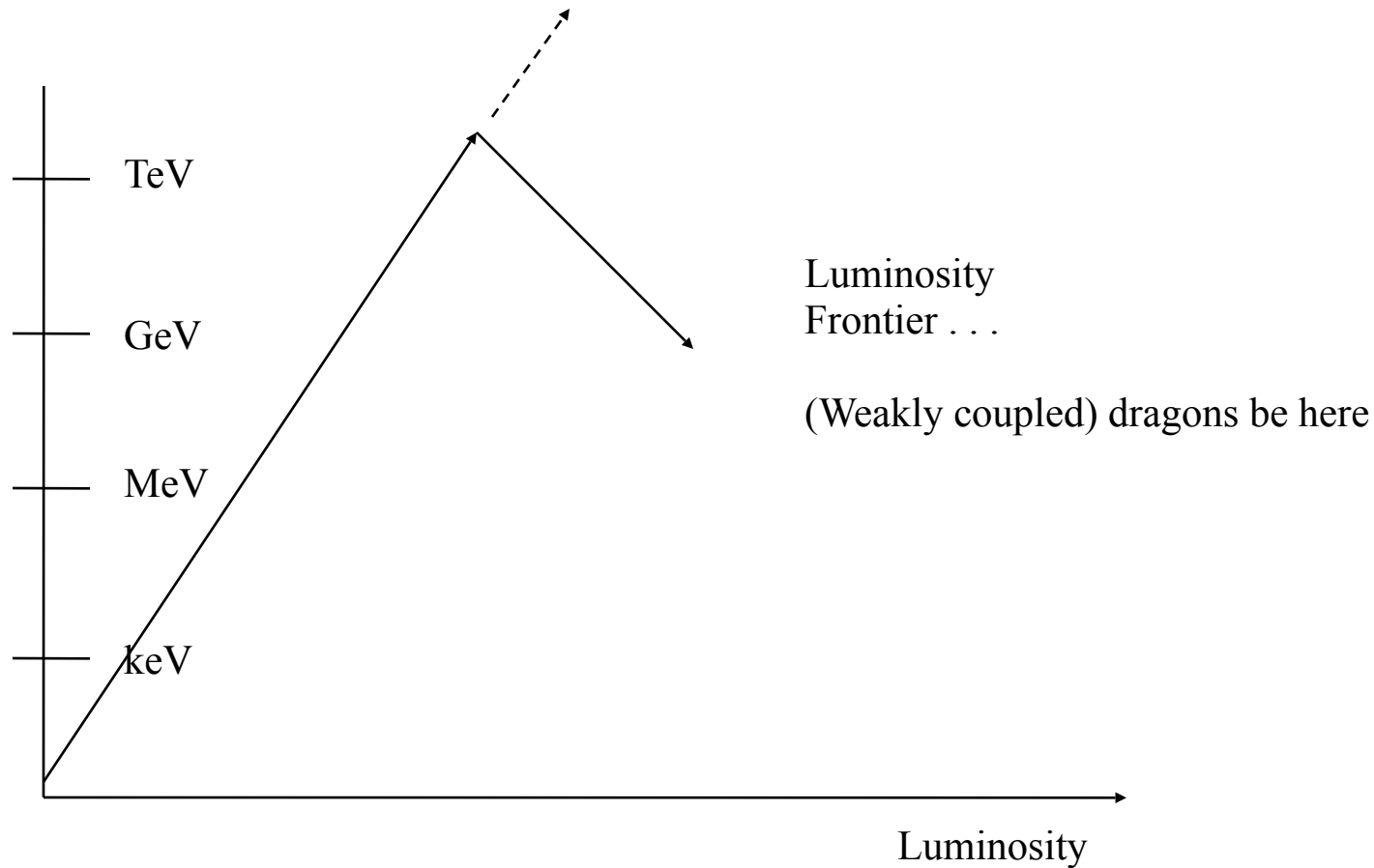
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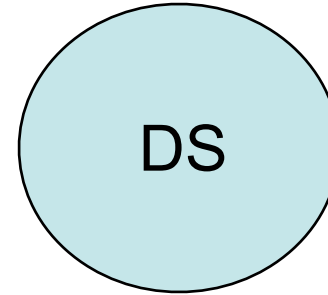
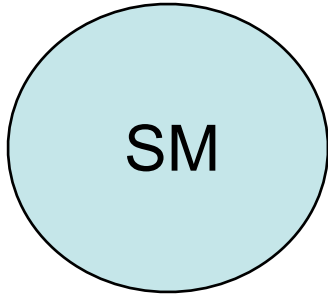


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A Tale of Two Sectors



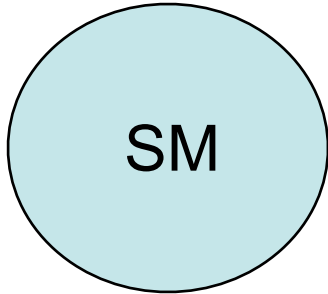
$$\mathcal{L}_{\text{SM}} = \dots$$

<http://pdg.lbl.gov>

Secluded DM - Pospelov, Ritz, and Voloshin (0711.4866)

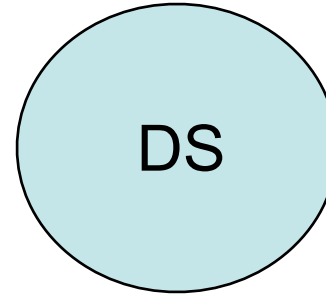
A Theory of DM - Arkani-Hamed, Finkbeiner, Slatyer, and Weiner (0810.0713)

A Tale of Two Sectors



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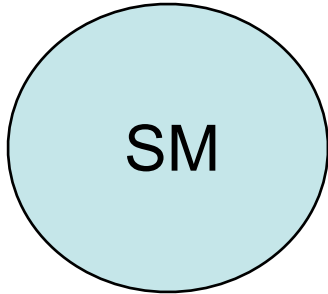


$$\mathcal{L}_{\text{DS}} \supset i\bar{\chi}\gamma^\mu D_\mu\chi + M\bar{\chi}\chi$$

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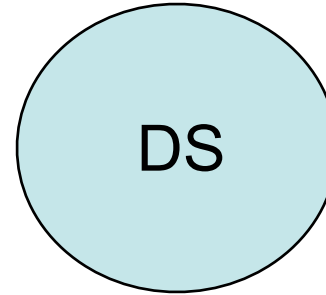
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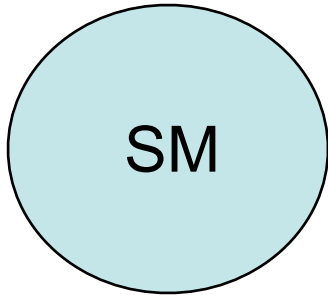
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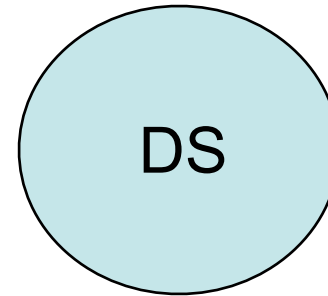
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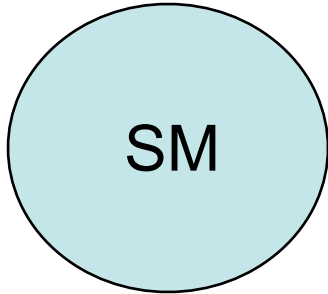


$$\begin{aligned} \mathcal{L}_{\text{DS}} \supset & i\bar{\chi}\gamma^\mu D_\mu\chi + M\bar{\chi}\chi \\ & -\frac{1}{4}f_{\mu\nu}f^{\mu\nu} + \frac{1}{2}m^2 b_\mu b^\mu \\ & +|D_\mu h_i|^2 - V(h_i) \end{aligned}$$

Secluded DM - Pospelov, Ritz, and Voloshin (0711.4866)

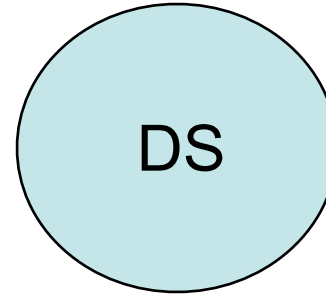
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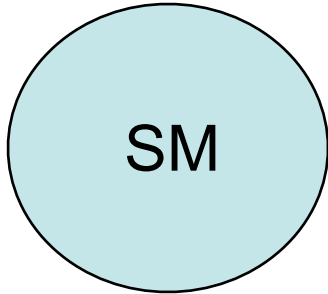
$$M \sim \text{TeV}$$

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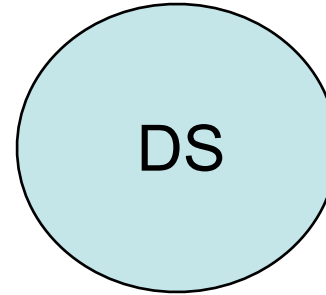
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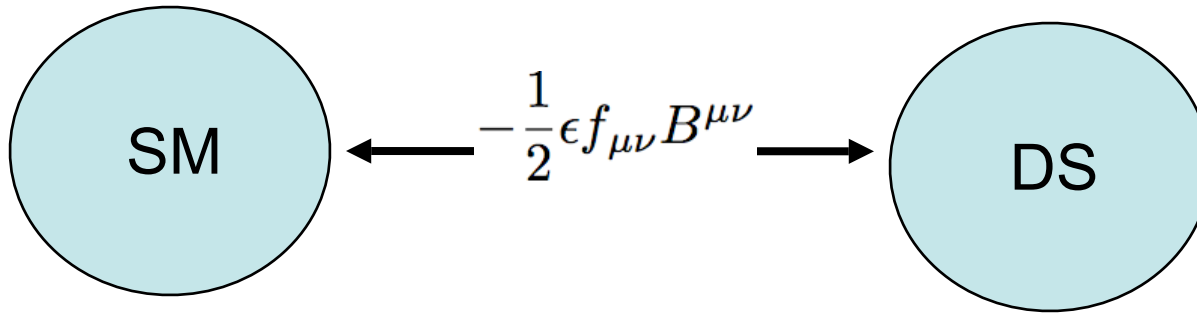
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A Theory of DM - Arkani-Hamed, Finkbeiner, Slatyer, and Weiner (0810.0713)

Dark Spectrum

$$\text{TeV} \frac{\quad}{M\bar{\chi}\chi}$$

Dark Spectrum

$$\text{TeV} \quad \overline{M\bar{\chi}\chi}$$

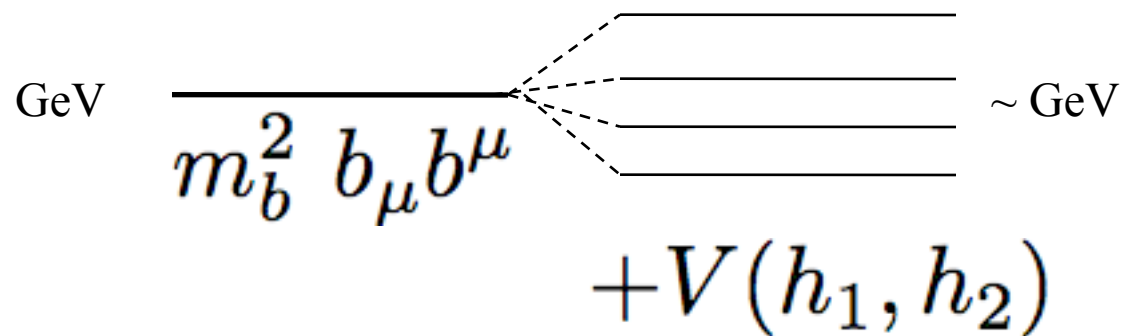
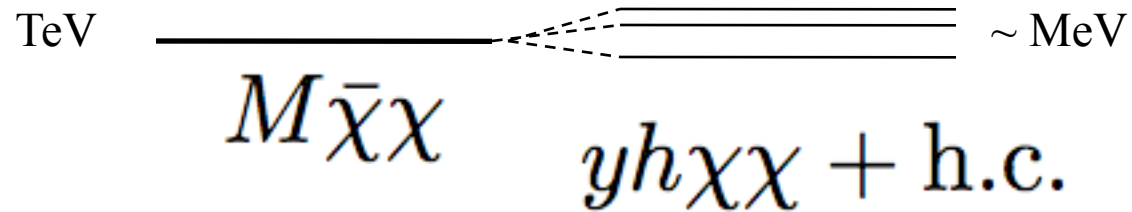
$$\text{GeV} \quad \overline{m_b^2 b_\mu b^\mu}$$

Dark Spectrum

TeV $\overline{\hspace{10em}}$ $M\bar{\chi}\chi$ $\begin{array}{l} \text{---} \\ \text{---} \\ \text{---} \end{array}$ $\sim \text{MeV}$
 $yh\chi\chi + \text{h.c.}$

GeV $\overline{\hspace{10em}}$
 $m_b^2 b_\mu b^\mu$

Dark Spectrum



Dark Spectrum

TeV $\overline{\hspace{1.5cm}}$ $M\bar{\chi}\chi$ $\begin{array}{l} \text{---} \\ \text{---} \\ \text{---} \end{array}$ $\sim \text{MeV}$

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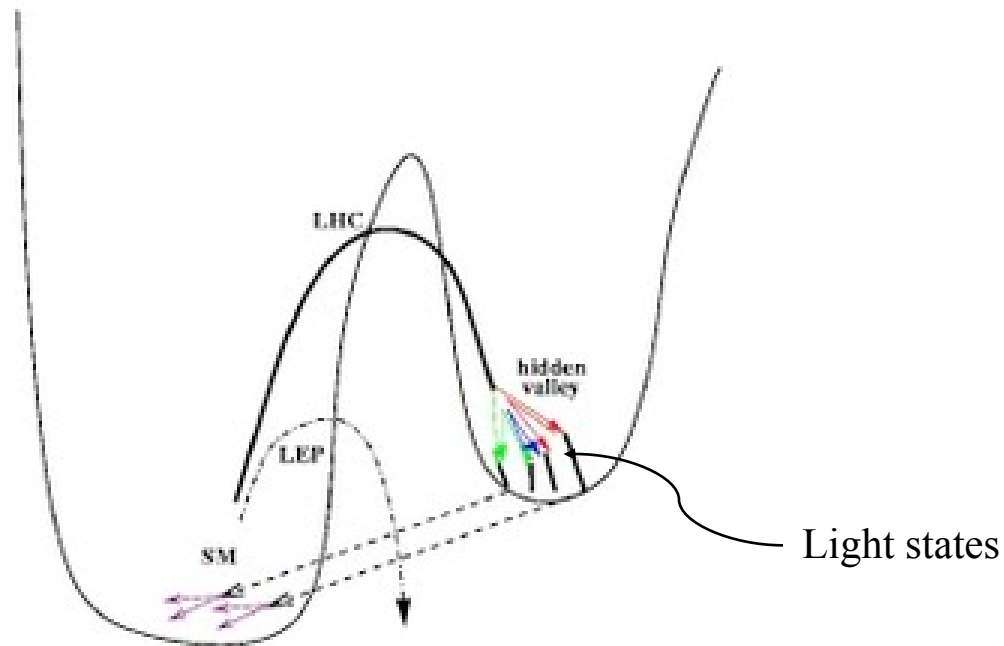
Can produce at colliders!!!

GeV $\overline{\hspace{1.5cm}}$ $m_b^2 b_\mu b^\mu$ $\begin{array}{l} \text{---} \\ \text{---} \\ \text{---} \end{array}$ $\sim \text{GeV}$

$+V(h_1, h_2)$

Hidden Valleys

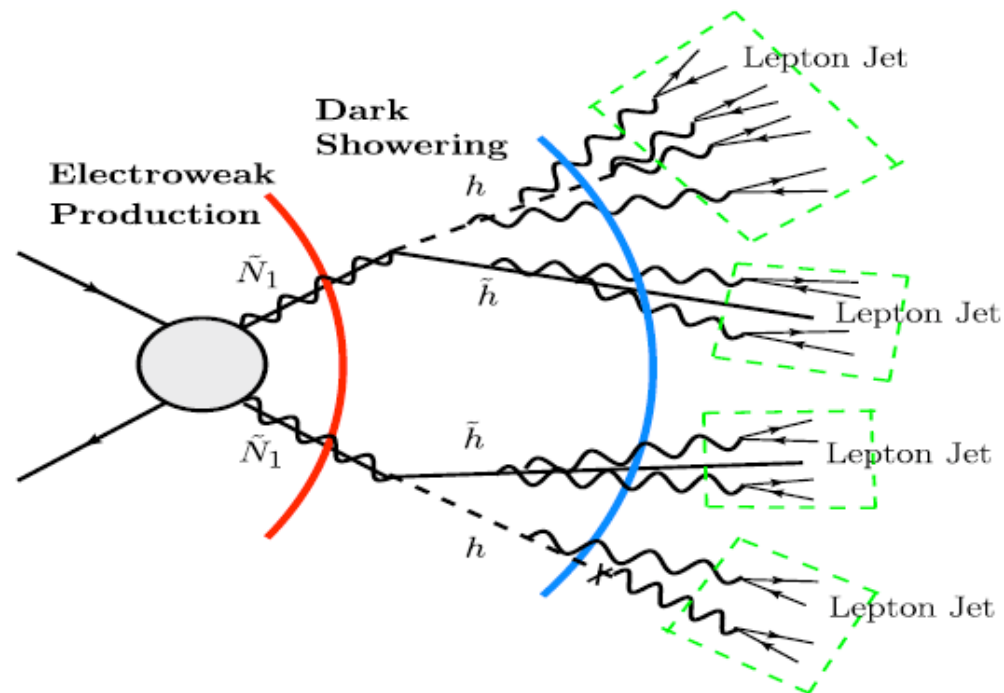
Strassler and Zurek's proposal of hidden valleys share some of the phenomenology and lepton jet searches can in principle be sensitive to these type of models as well,



* Taken (without permission) from Strassler's talk.

Part II

Production and Evolution of Dark States



Coupling to the Standard Model

Coupling to the Standard Model

In general the dark gauge-boson can mix with both the photon and the Z^0 ,

$$\begin{aligned}\mathcal{L}_{\text{gauge mix}} &= -\frac{1}{2}\epsilon_1 b_{\mu\nu} A^{\mu\nu} - \frac{1}{2}\epsilon_2 b_{\mu\nu} Z^{\mu\nu} \\ &= -\frac{1}{2}\epsilon'_1 b_{\mu\nu} B^{\mu\nu} - \frac{1}{2}\epsilon'_2 b_{\mu\nu} W_3^{\mu\nu}\end{aligned}$$

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If **supersymmetry** is only softly broken in the dark sector, then there is also an important mixing of the electroweak gauginos with the dark gaugino:

$$\mathcal{L}_{\text{gaugino mix}} = -2i\epsilon'_1 \tilde{b}^\dagger \bar{\sigma}^\mu \partial_\mu \tilde{B} - 2i\epsilon'_2 \tilde{b}^\dagger \bar{\sigma}^\mu \partial_\mu \tilde{W}_3 + \text{h.c.}$$

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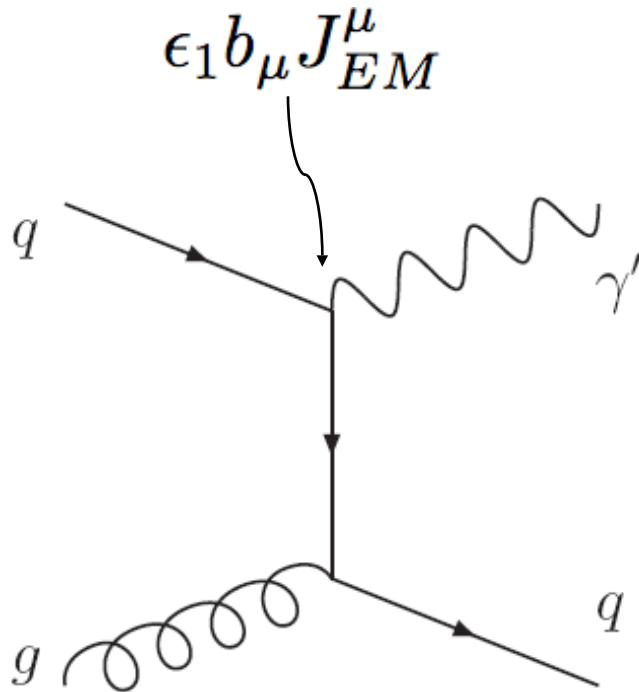
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All in all we have the following couplings (after diagonalization and etc.), which act as a portal to the dark sector

$$\mathcal{L}_{\text{portal}} = \epsilon_1 b_\mu J_{\text{EM}}^\mu + \epsilon_2 Z_\mu J_b^\mu + \epsilon'_1 \tilde{B} \tilde{J}_{\tilde{b}} + \epsilon'_2 \tilde{W}_3 \tilde{J}_{\tilde{b}}$$

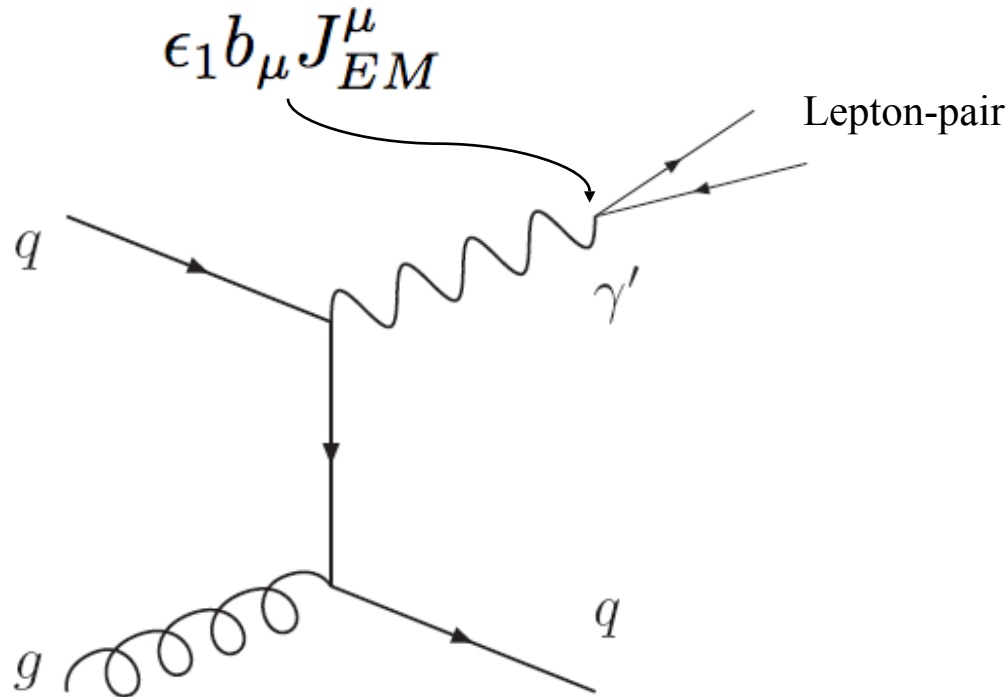
Prompt (Dark) Photon Production

The coupling to the electromagnetic current allows for the production of the dark photon in a similar manner to a prompt photon ([Baumgart et al. 0901.0283](#)) :



Prompt (Dark) Photon Production

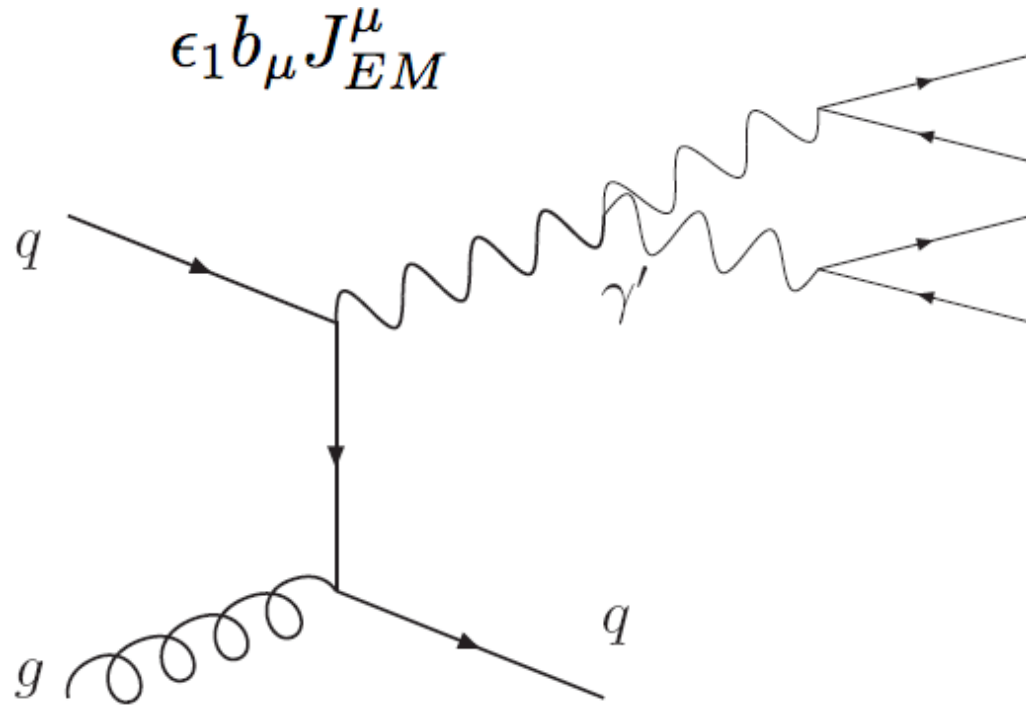
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Too much background!!!

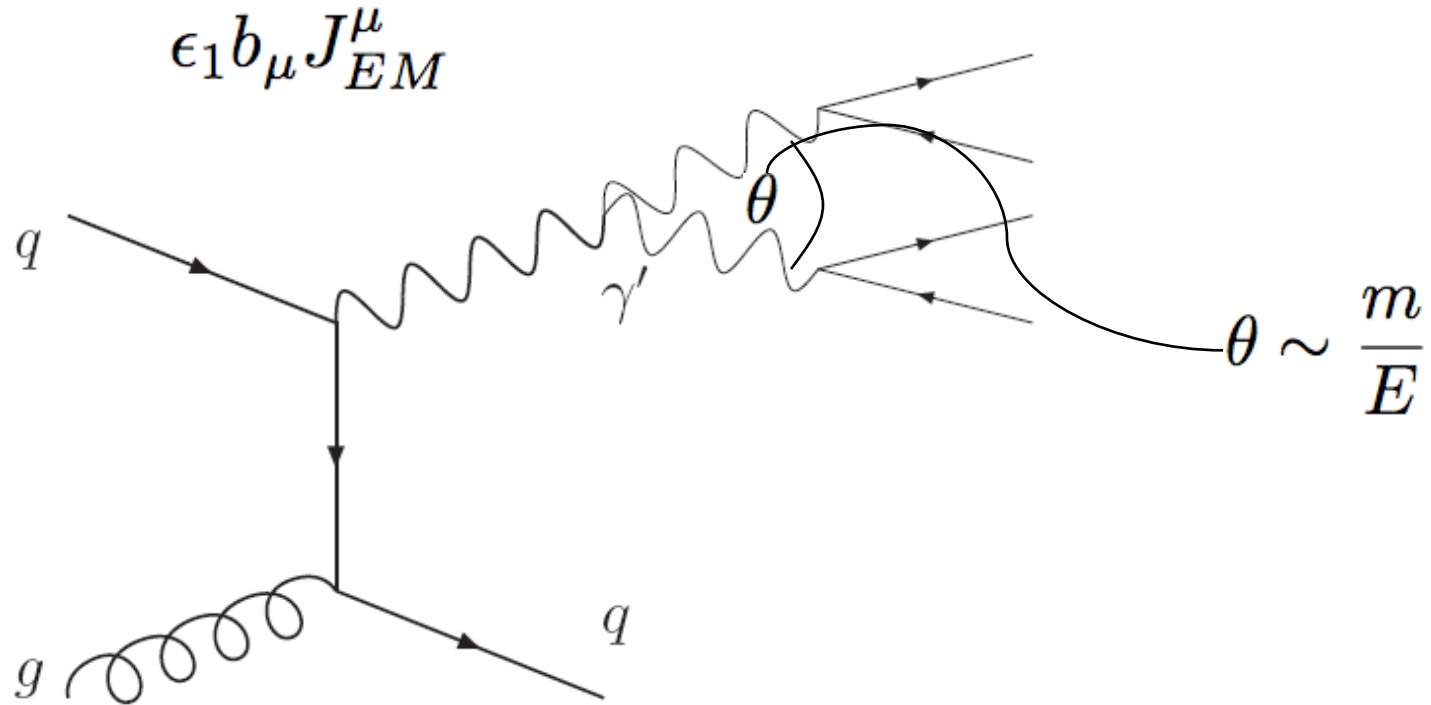
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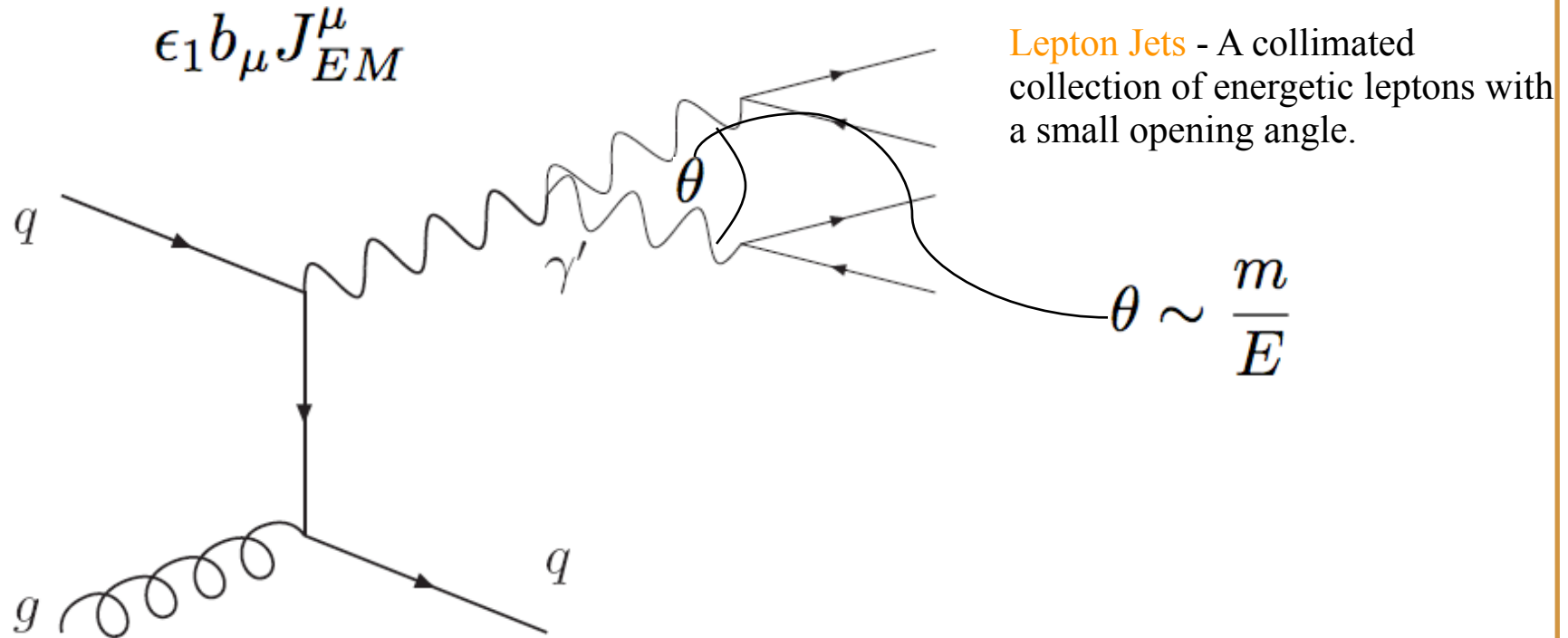
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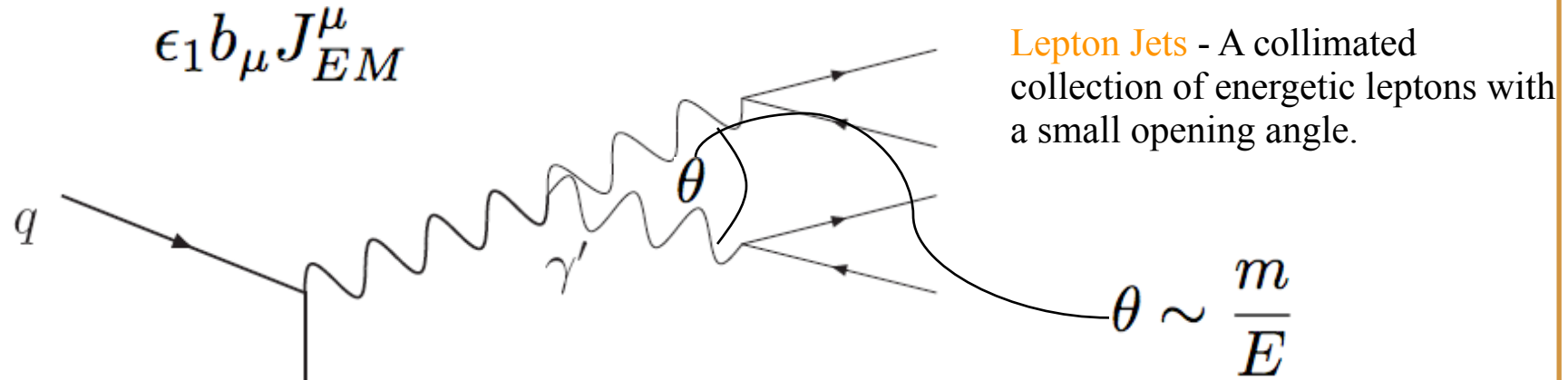
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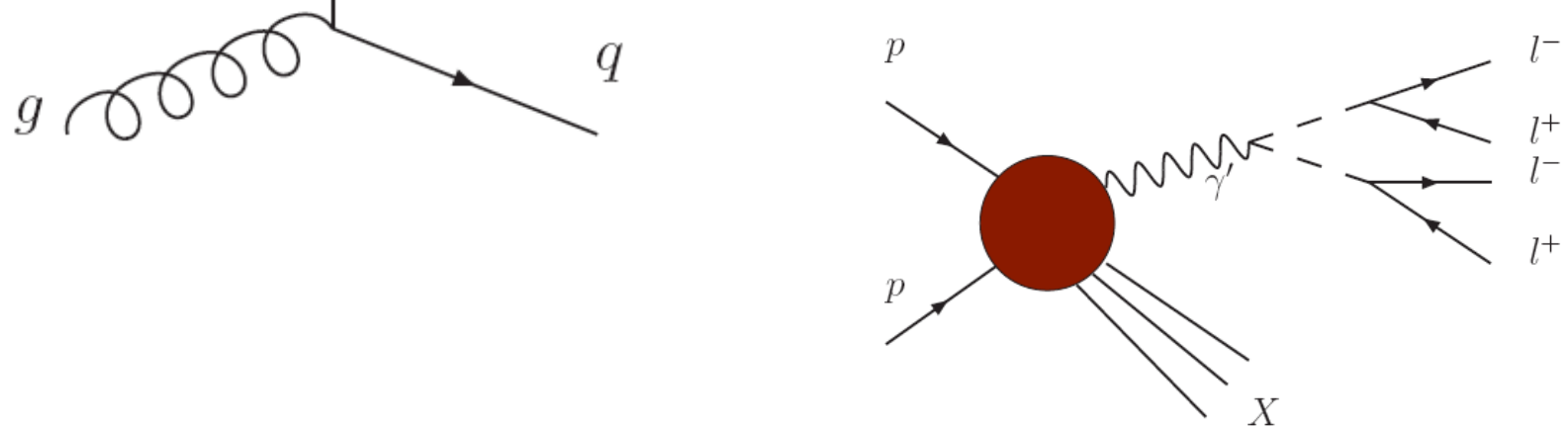


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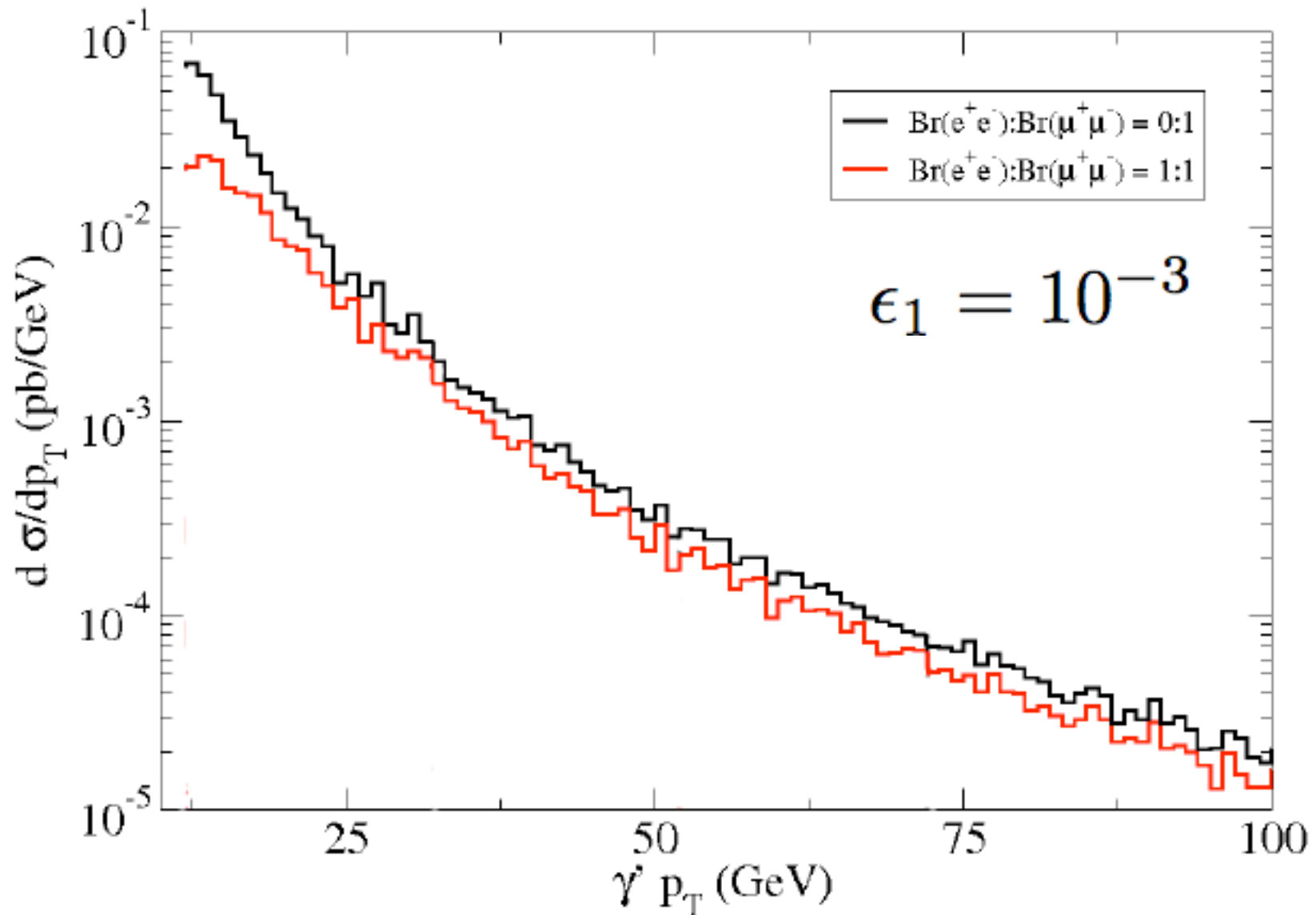
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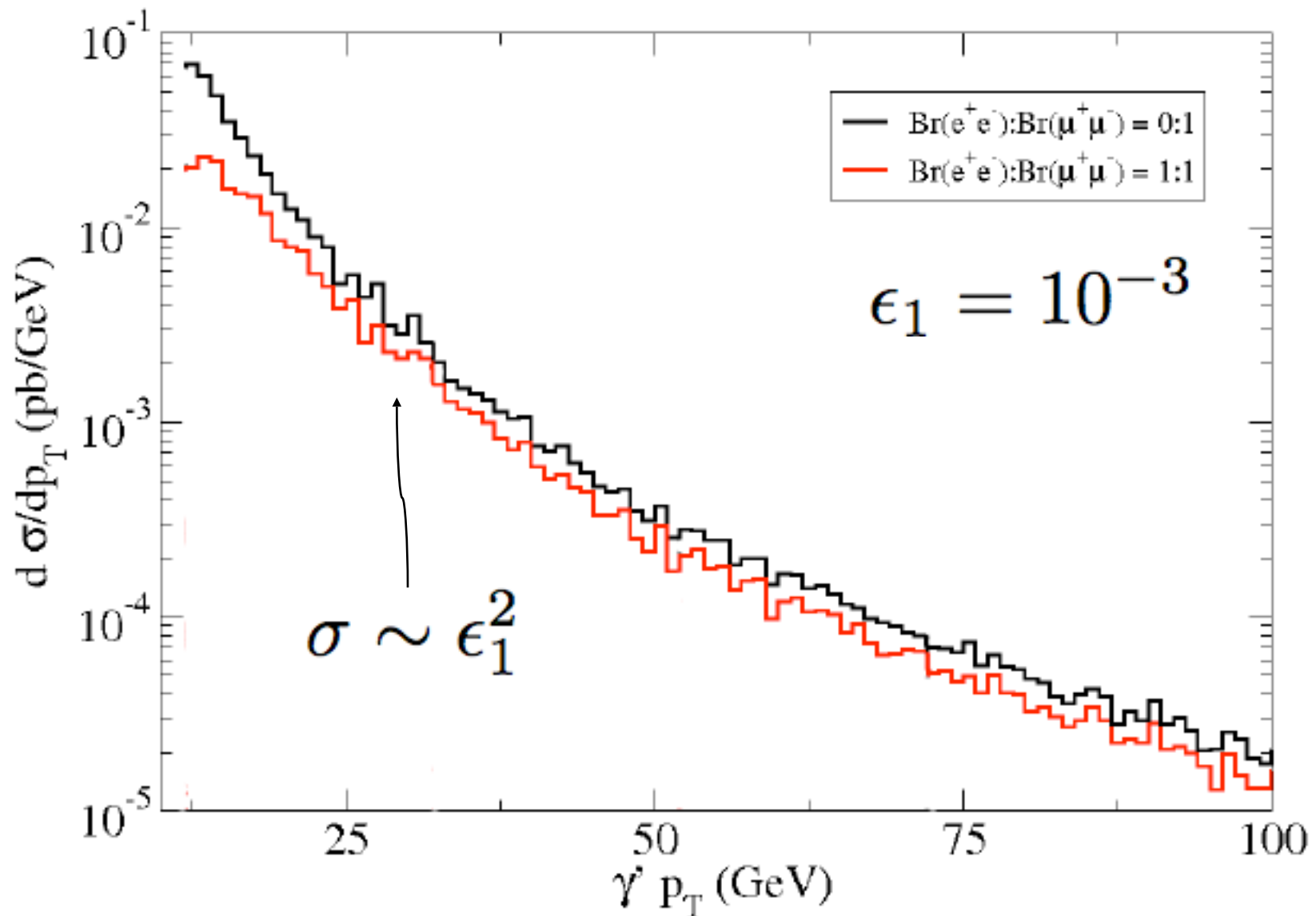
Lepton-Jet recoiling against a QCD jet:



Cross-Sections at the LHC

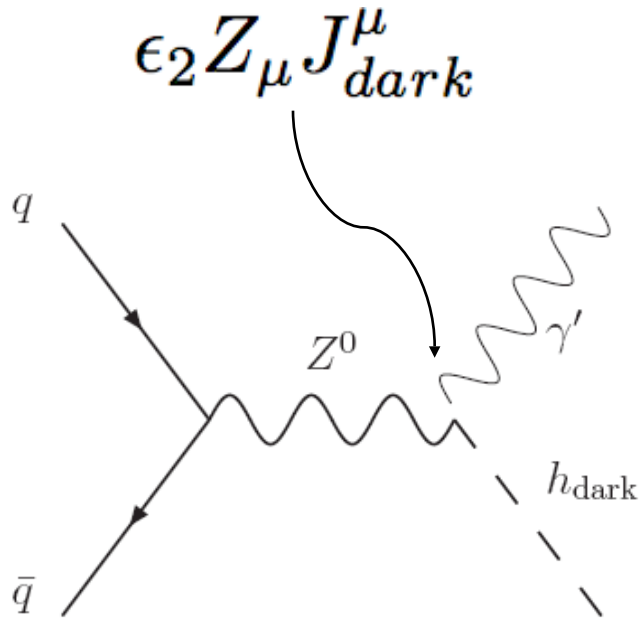


Cross-Sections at the LHC



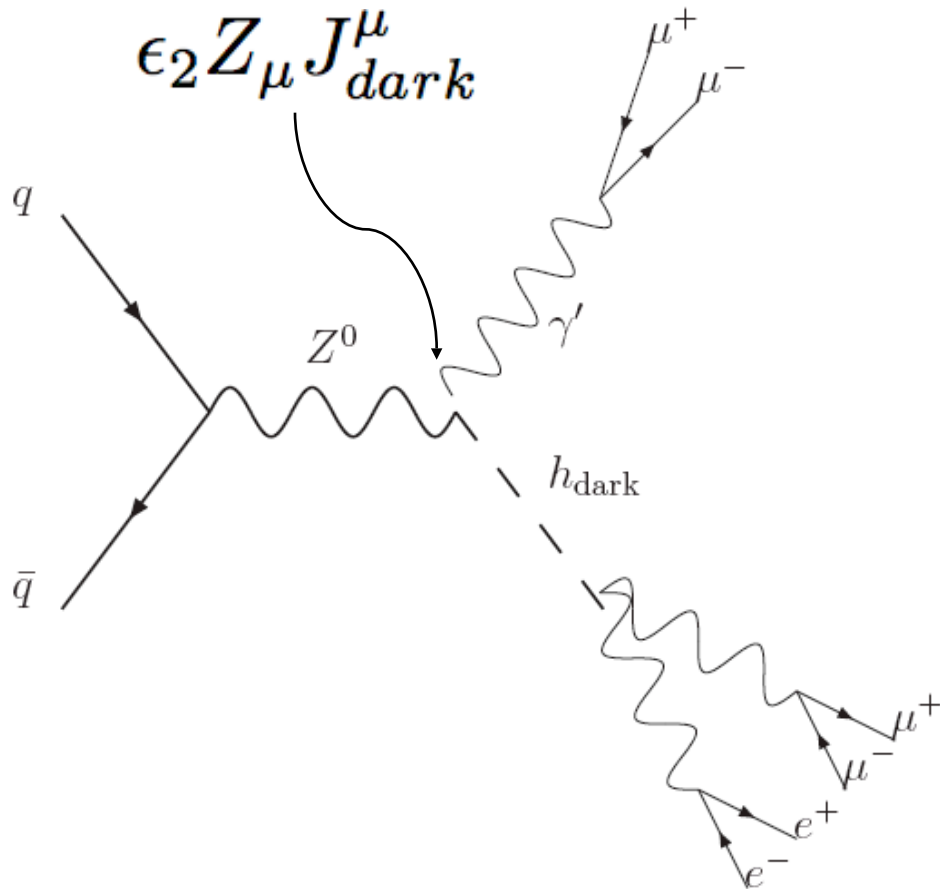
Rare Z^0 Decay

The neutral vector-boson couples directly to the dark current (Baumgart et al. and Cheung et al.). Therefore, the dark higgses and can be directly produced:



Rare Z^0 Decay

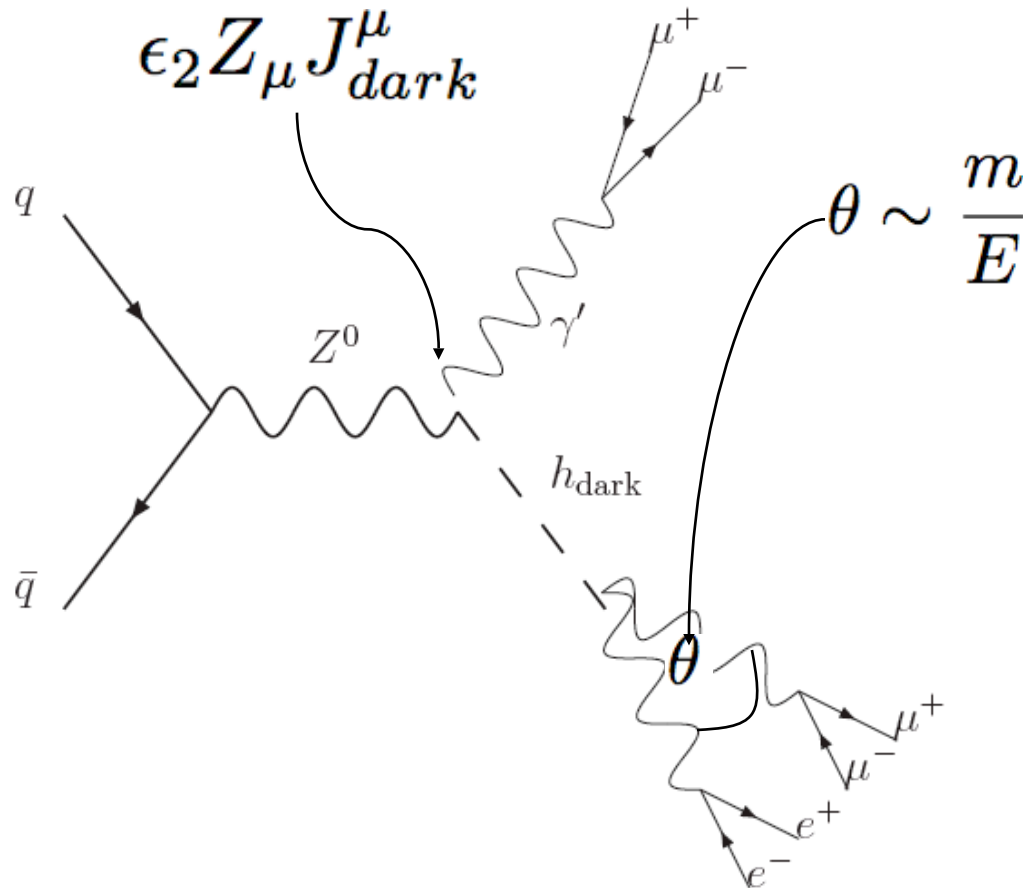
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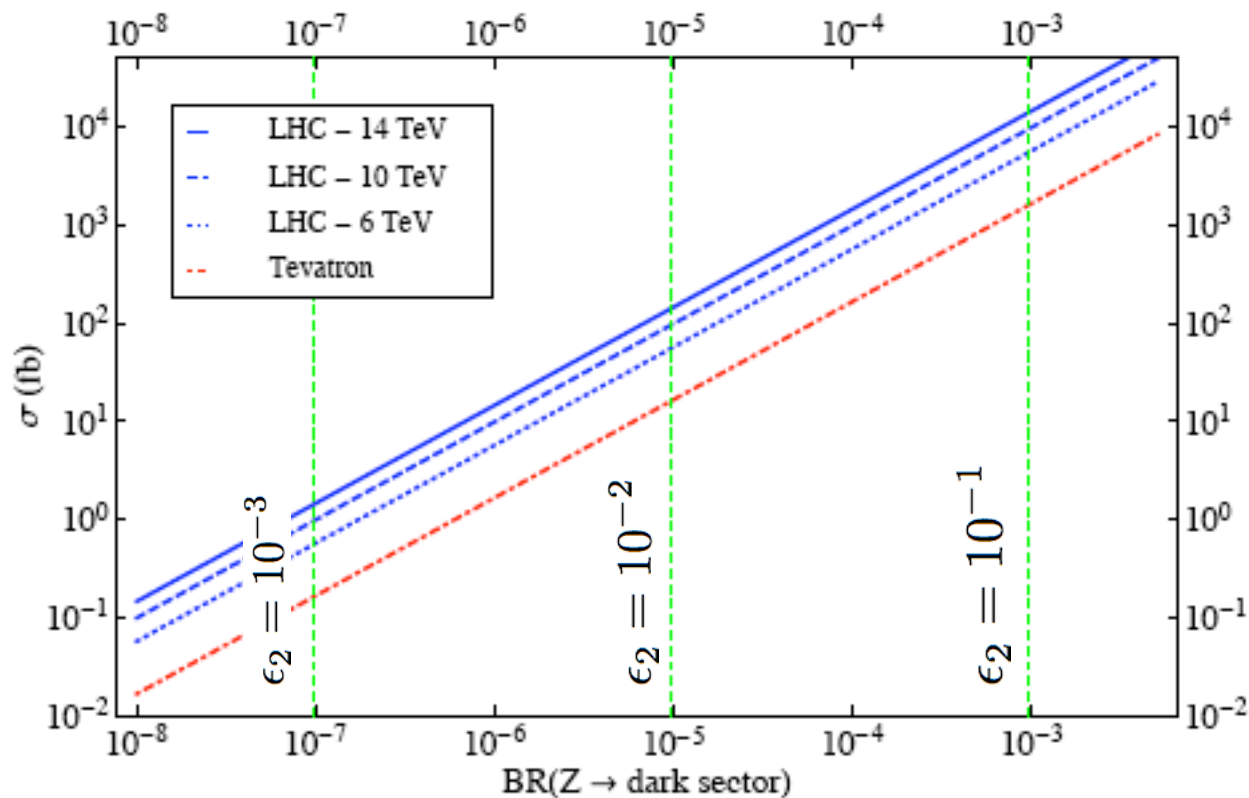
Lepton Jets - A collimated collection of energetic leptons with a small opening angle.



Rare Z^0 Decays - Reach

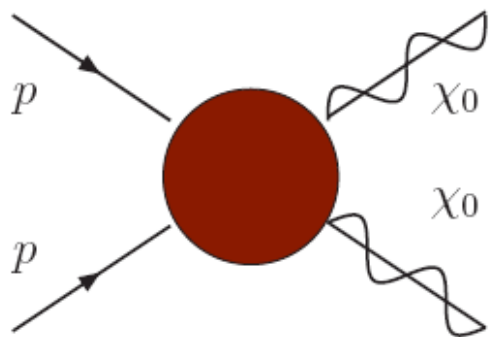
At LEP : $\text{BR}(Z \rightarrow f\bar{f}) = \frac{\epsilon_2^2 g_{\text{dark}}^2 M_{Z^0}}{12\pi \Gamma_{Z^0}} \rightarrow \mathcal{O}(100)$ events for $\epsilon_2 = 10^{-2}$

At Tevatron and LHC :



Neutralino Decay

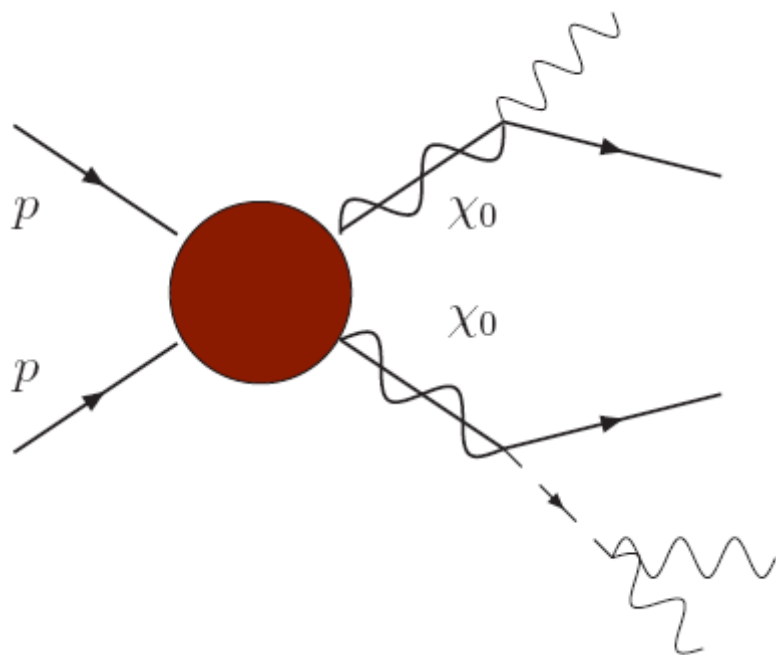
The bottom of the SUSY cascade is no longer stable ([Arkani-Hamed and Weiner](#)). It will decay into the dark sector. A clean channel is electroweak-ino production ([Cheung et al.](#))



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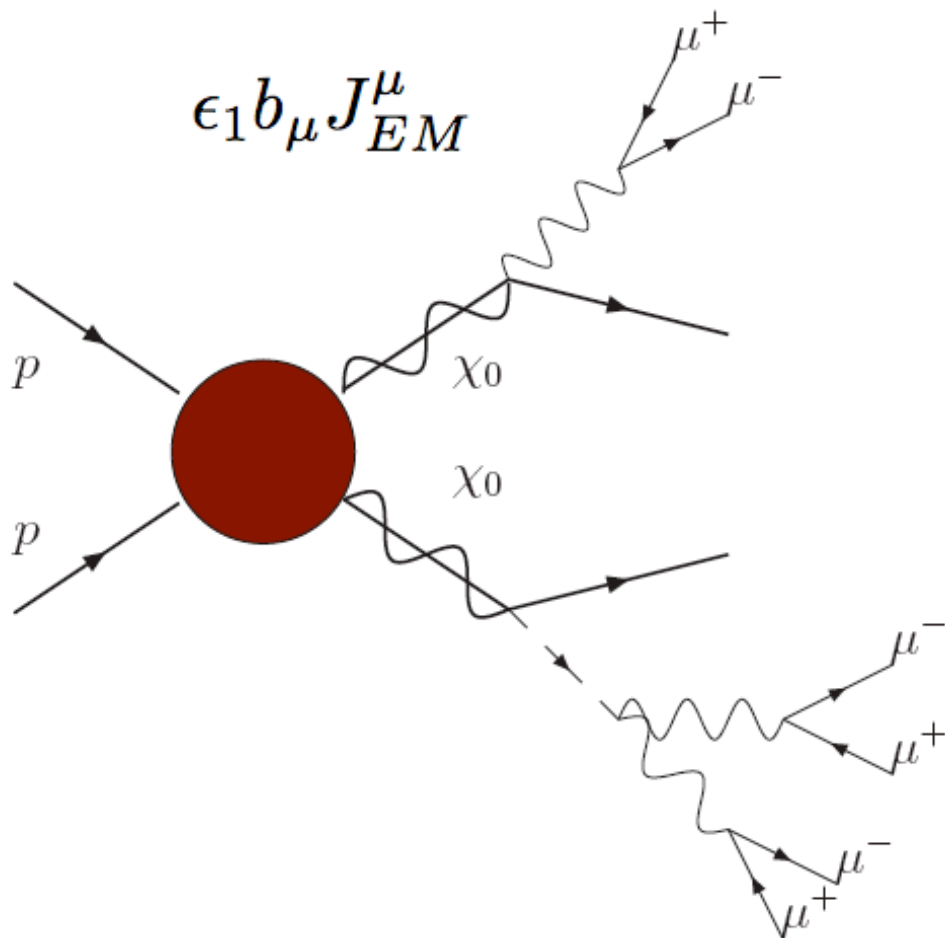
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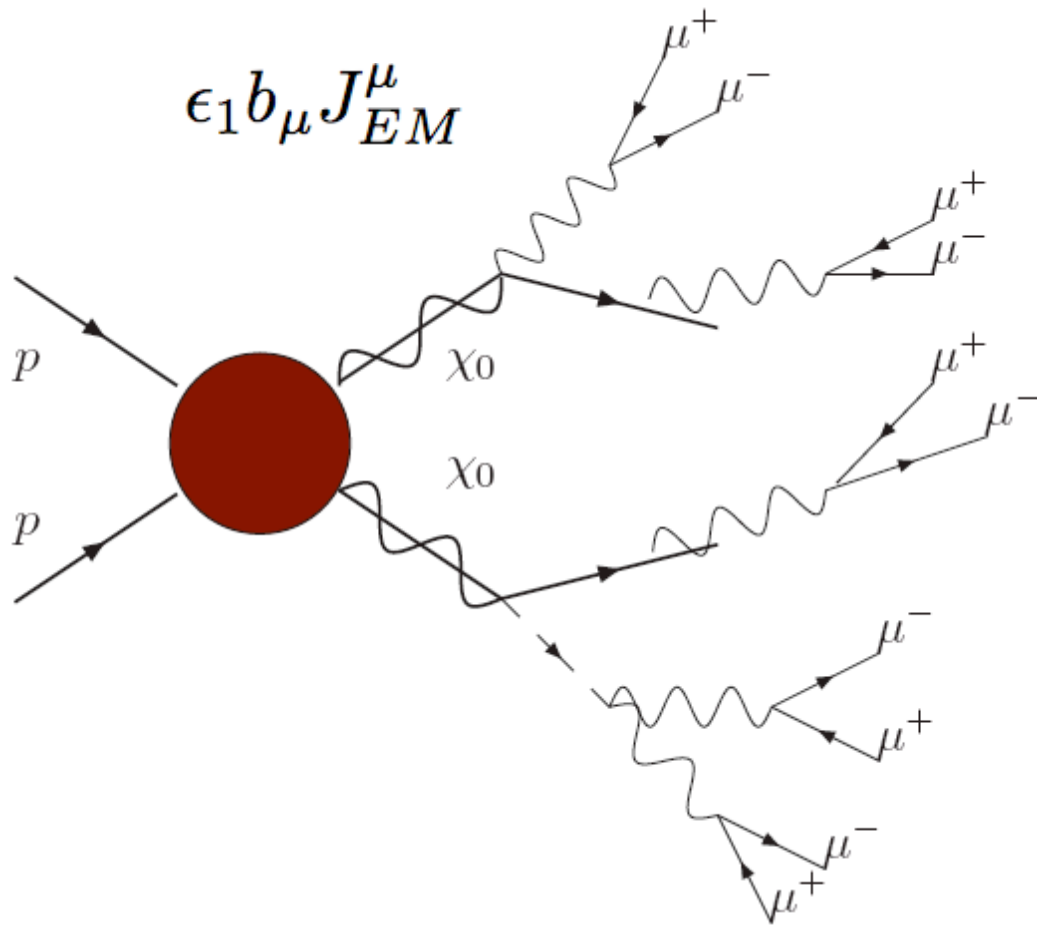
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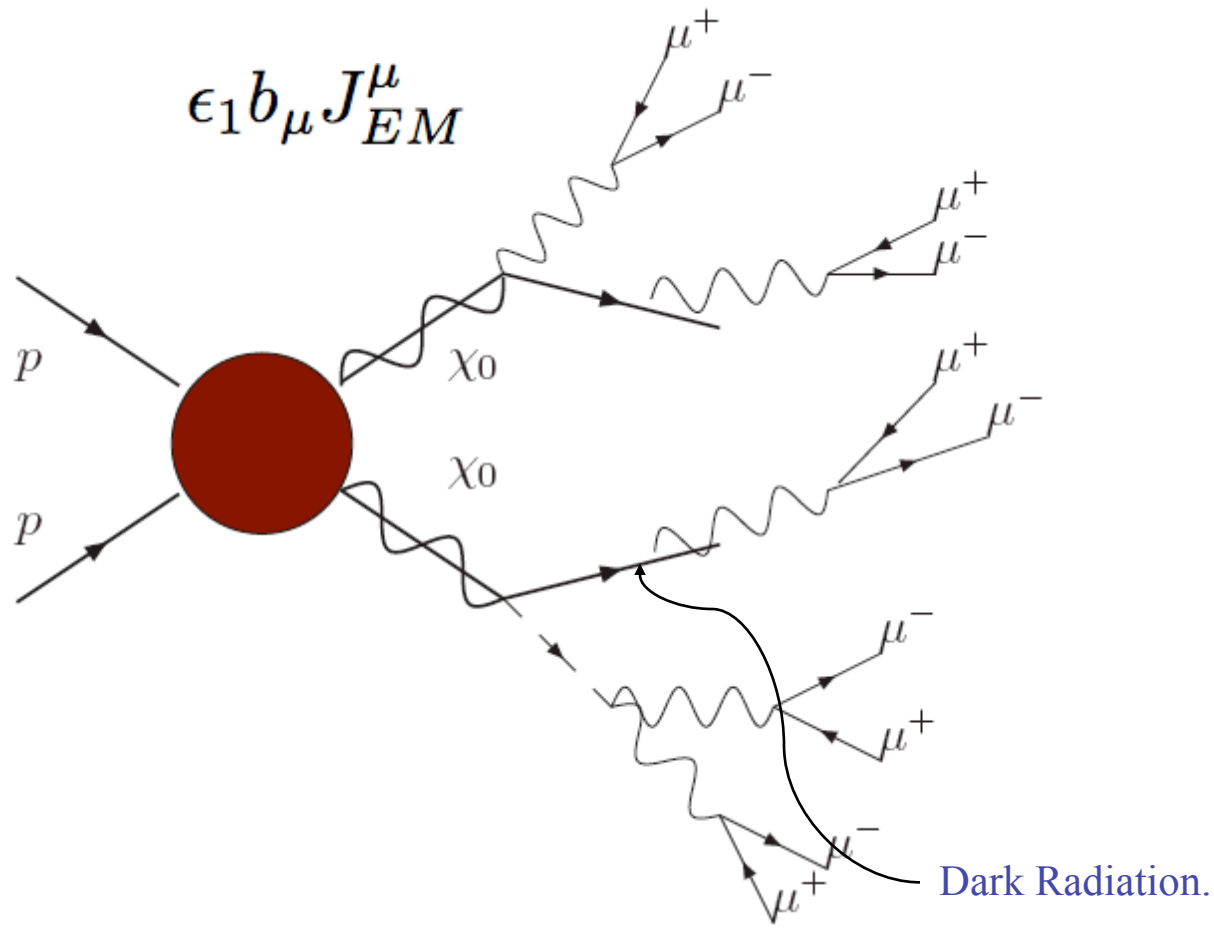
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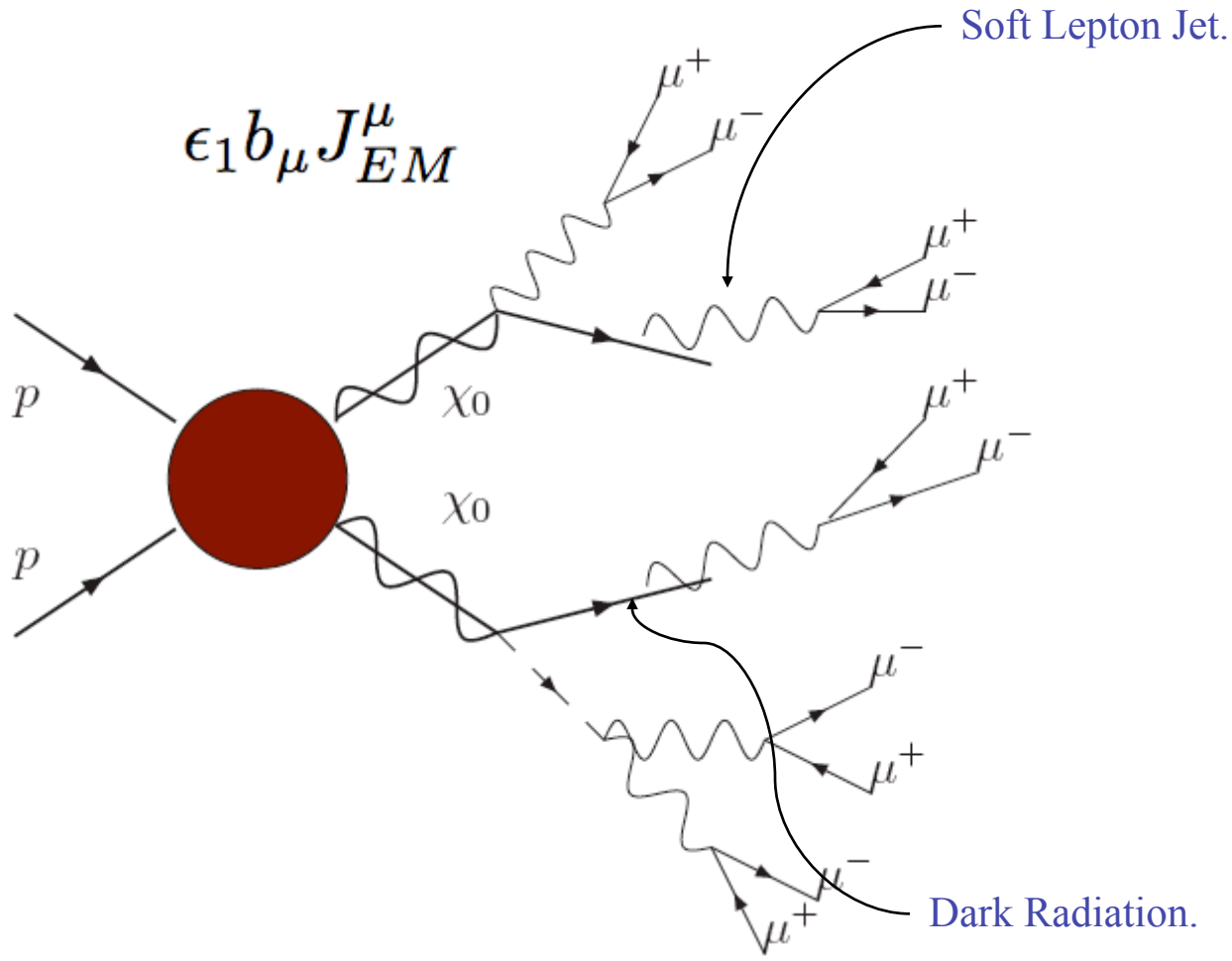
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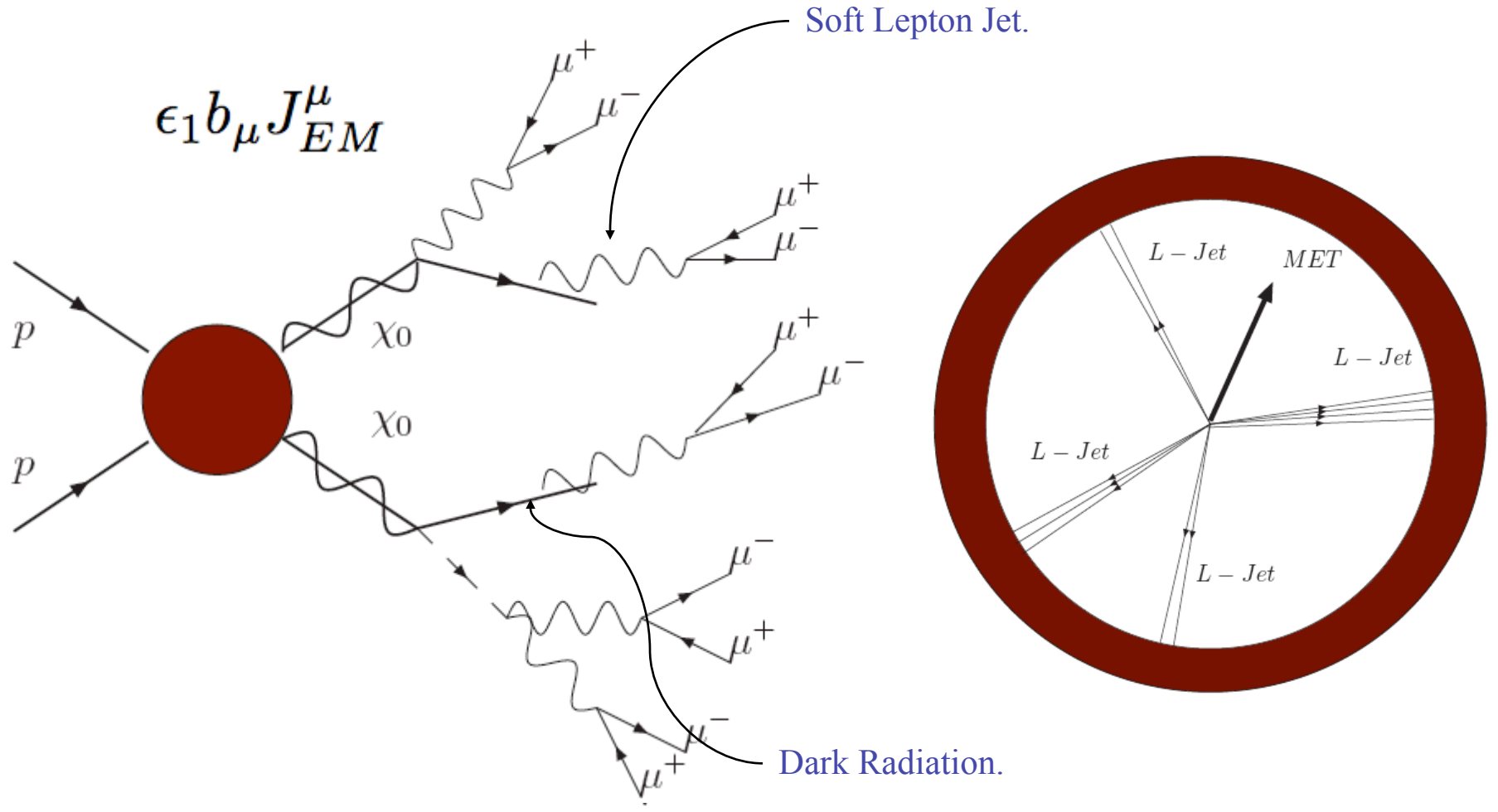
Neutralino Decay

The bottom of the SUSY cascade is no longer stable ([Arkani-Hamed and Weiner](#)). It will decay into the dark sector. A clean channel is electroweak-ino production ([Cheung et al.](#))

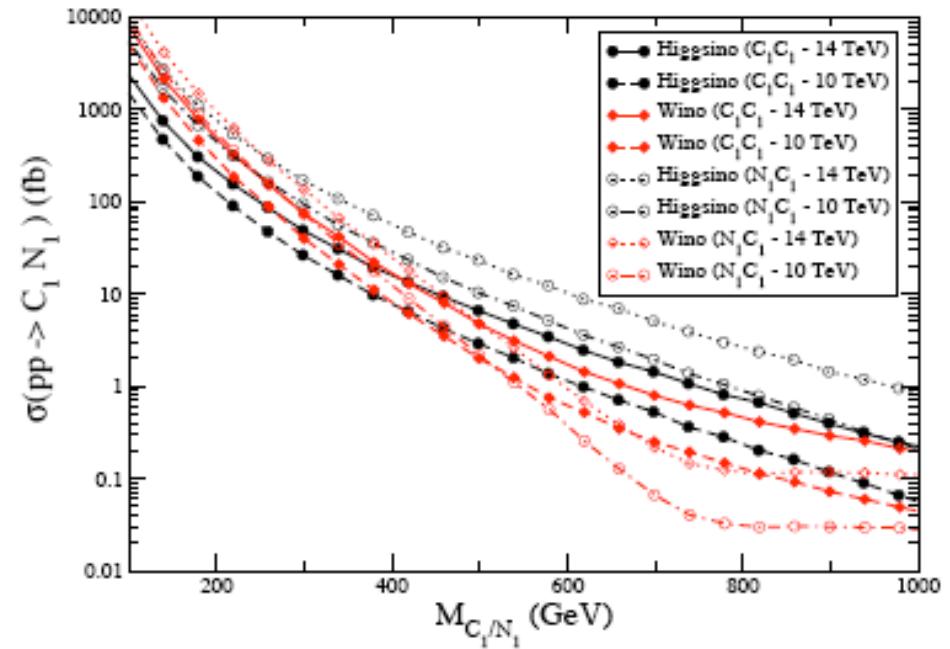
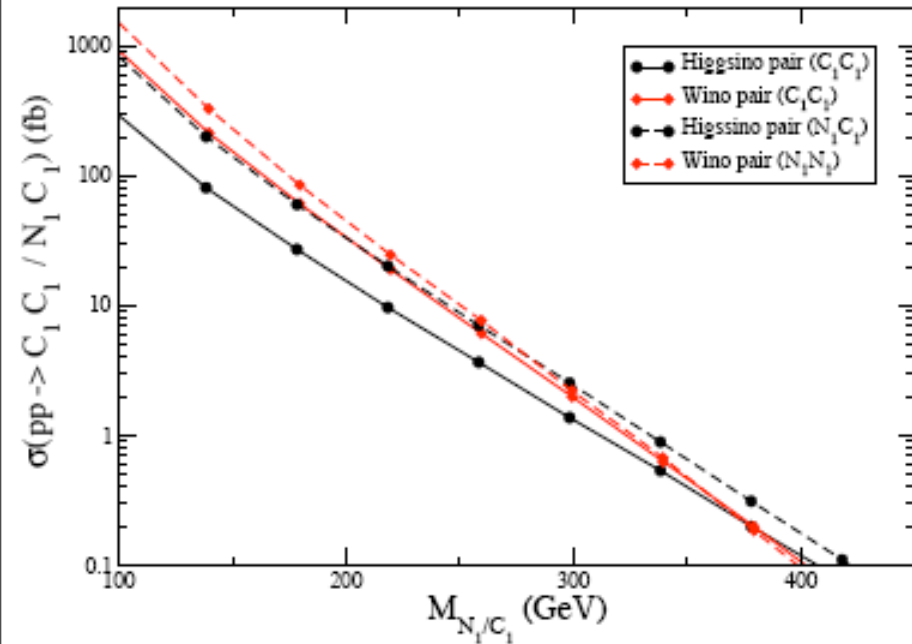


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LHC/Tevatron Reach



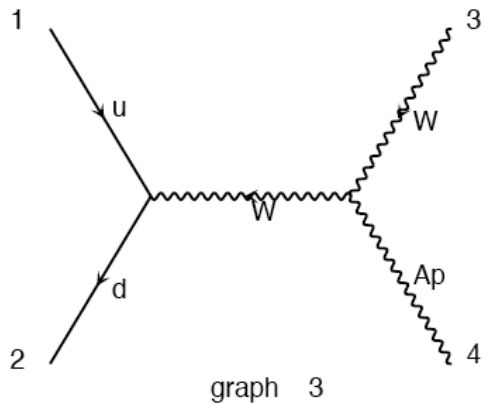
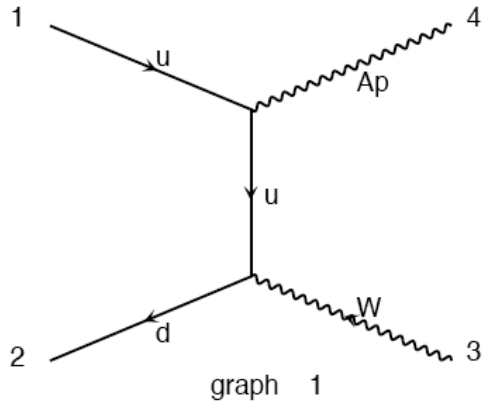
* This is for a squark mass of 750 GeV.

These are large cross-sections.

Some of the parameter space can already be excluded by Tevatron searches...

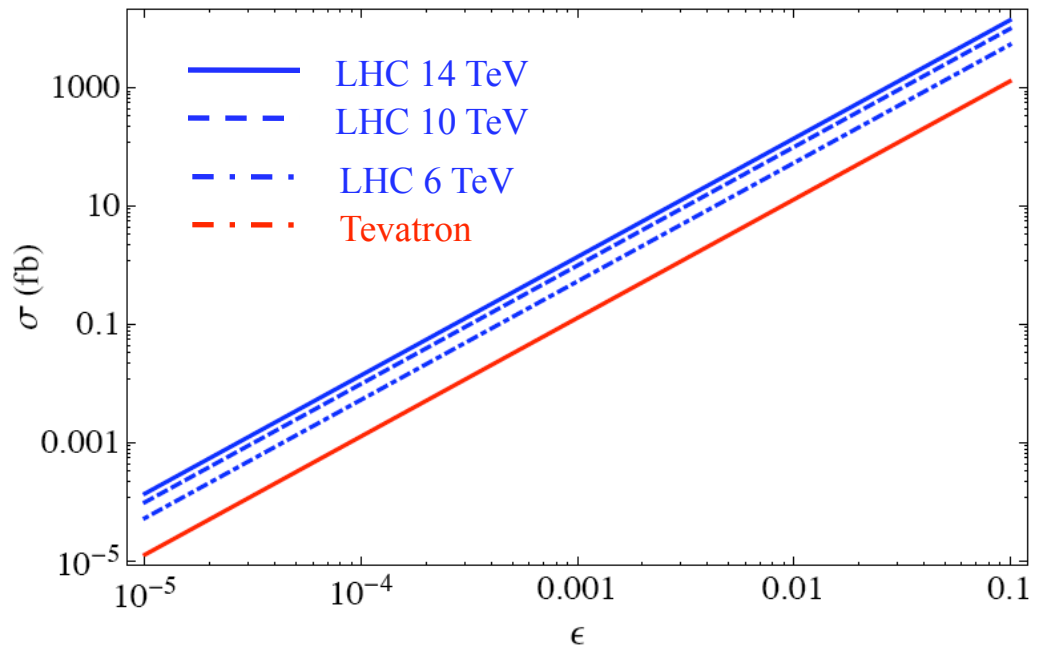
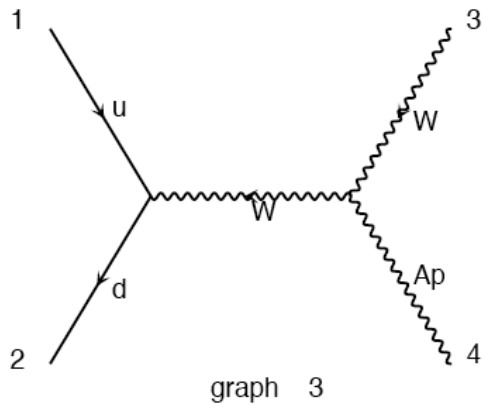
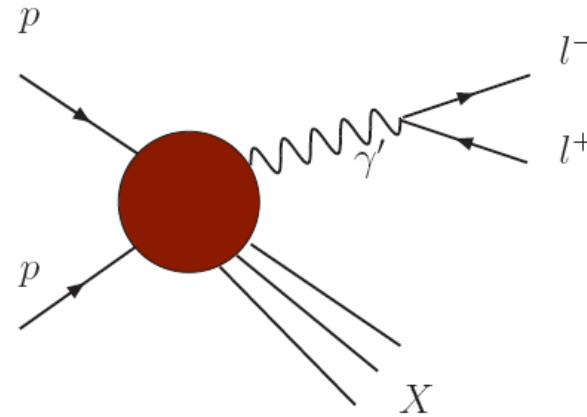
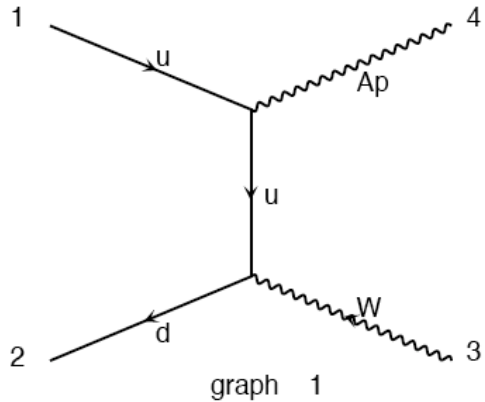
Dark Photon + W (work in progress. . .)

Similar to the prompt dark photon production we can consider the associate production of a dark photon together with a W boson. You lose on the cross-section, but you gain from the W mass peak.



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Lepton Jets

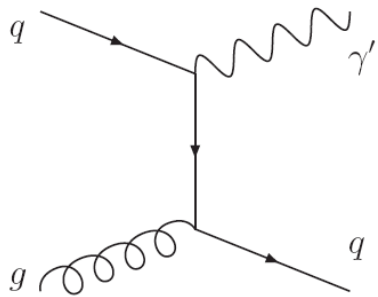
Origin of a Species

Production:

Origin of a Species

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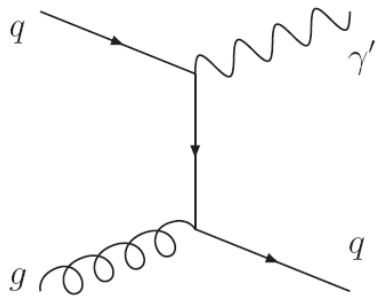
Prompt Dark Photon



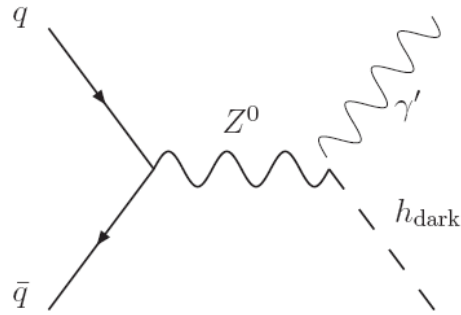
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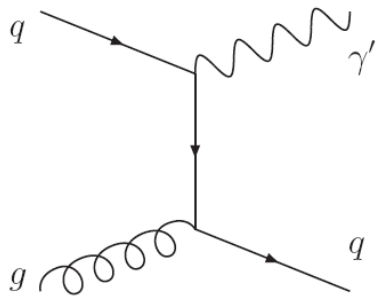
Rare Z^0 decays



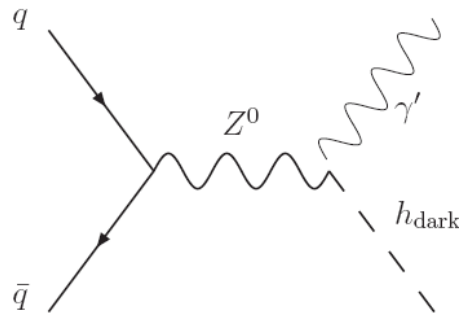
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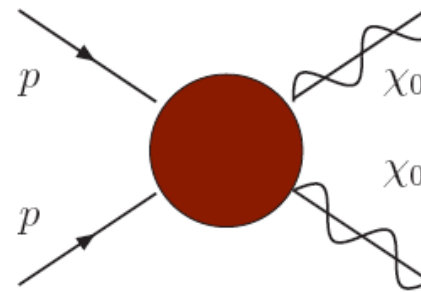
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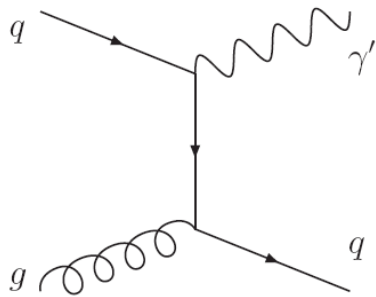
Gaugino Pair Prod



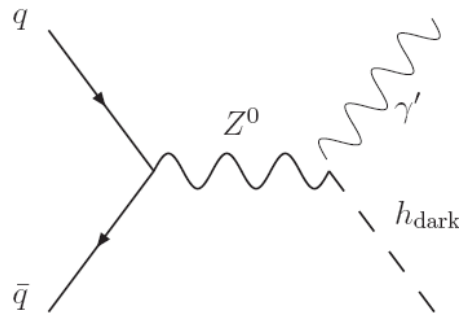
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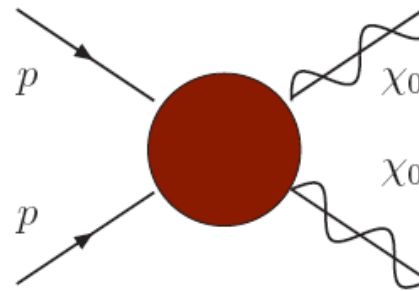
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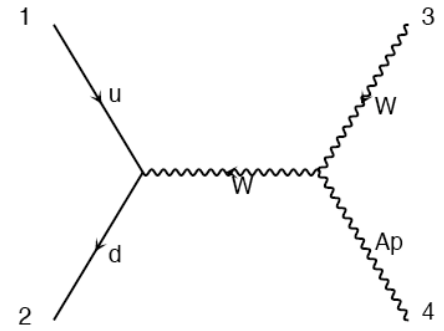
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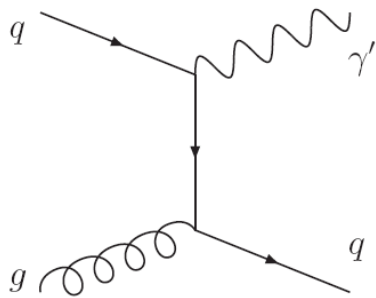
Dark Photon+W



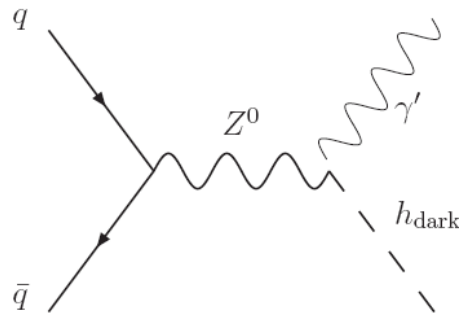
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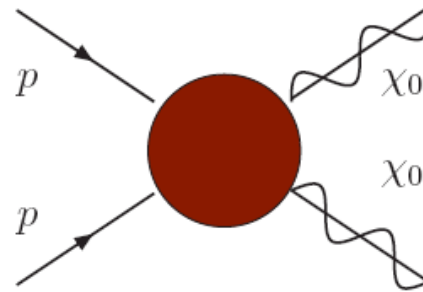
Prompt Dark Photon



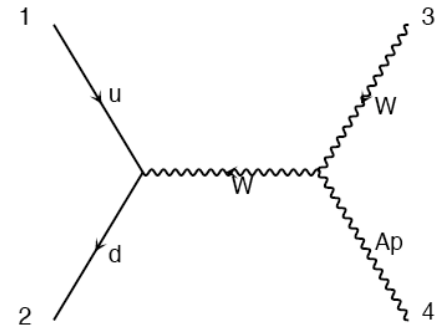
Rare Z^0 decays



Gaugino Pair Prod



Dark Photon+W



Evolution:

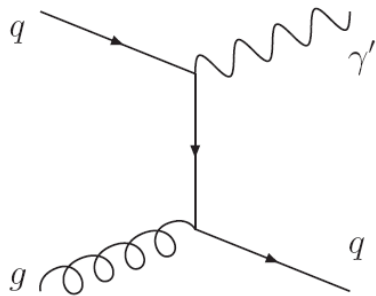
Dark Radiation



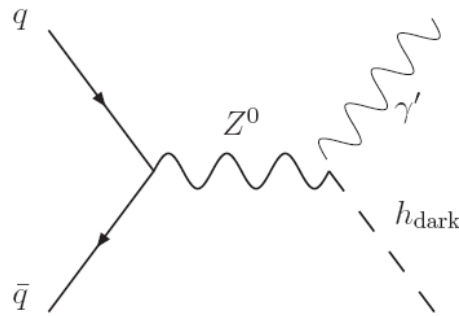
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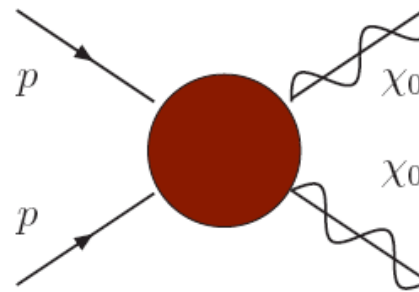
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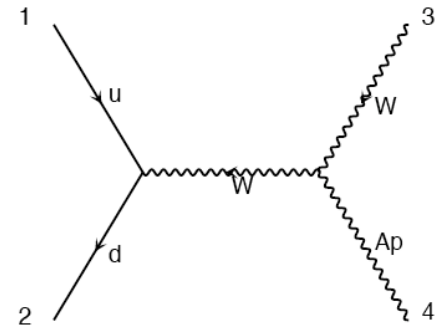
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Gaugino Pair Prod



Dark Photon+W

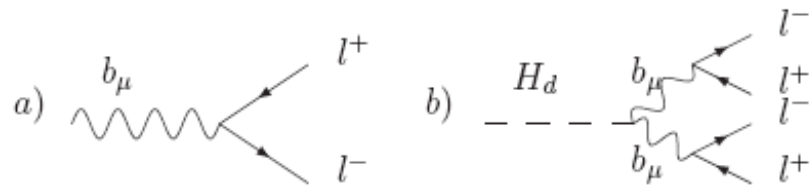


Evolution:

Dark Radiation



Dark Cascades and Lepton Jets



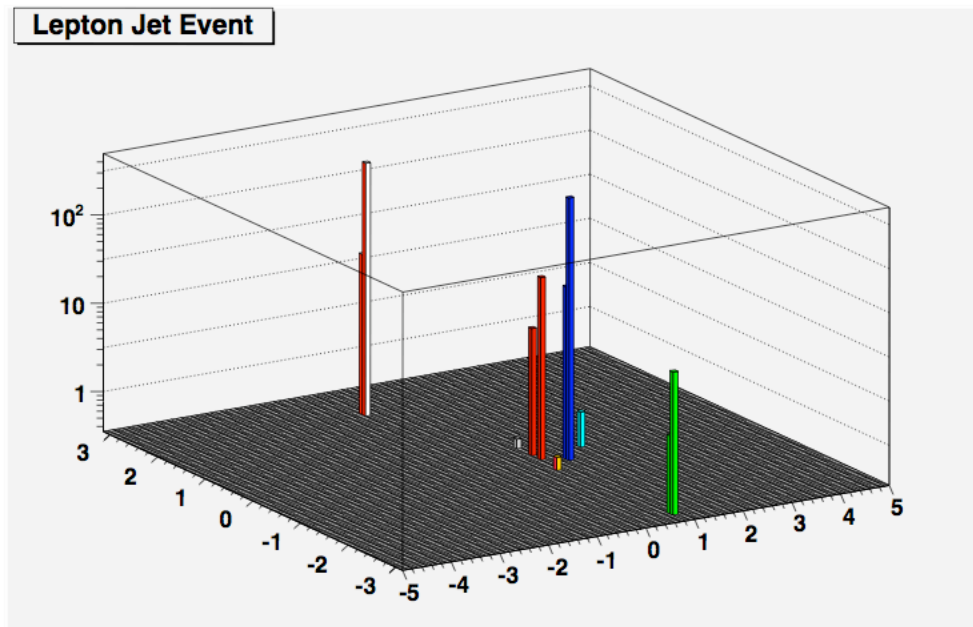
Lepton Jets

SLAC

Itay Yavin

Part III

Lepton Jets



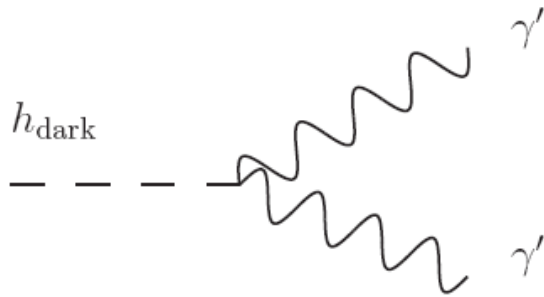
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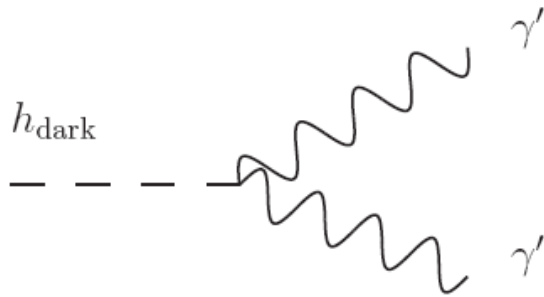
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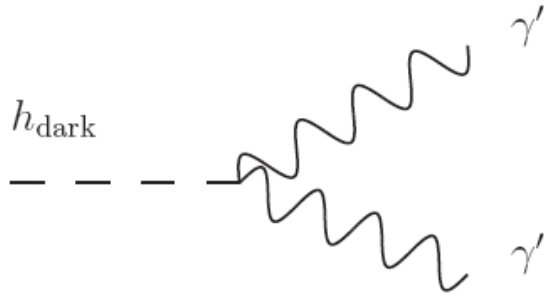


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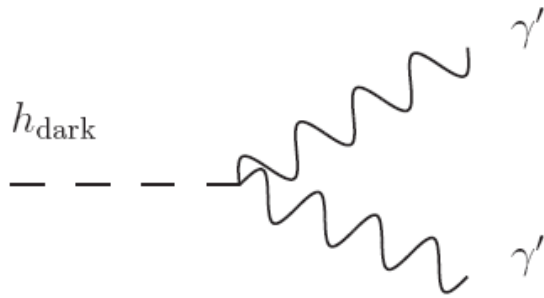
Prompt decay
(Many leptons)

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dark photon!

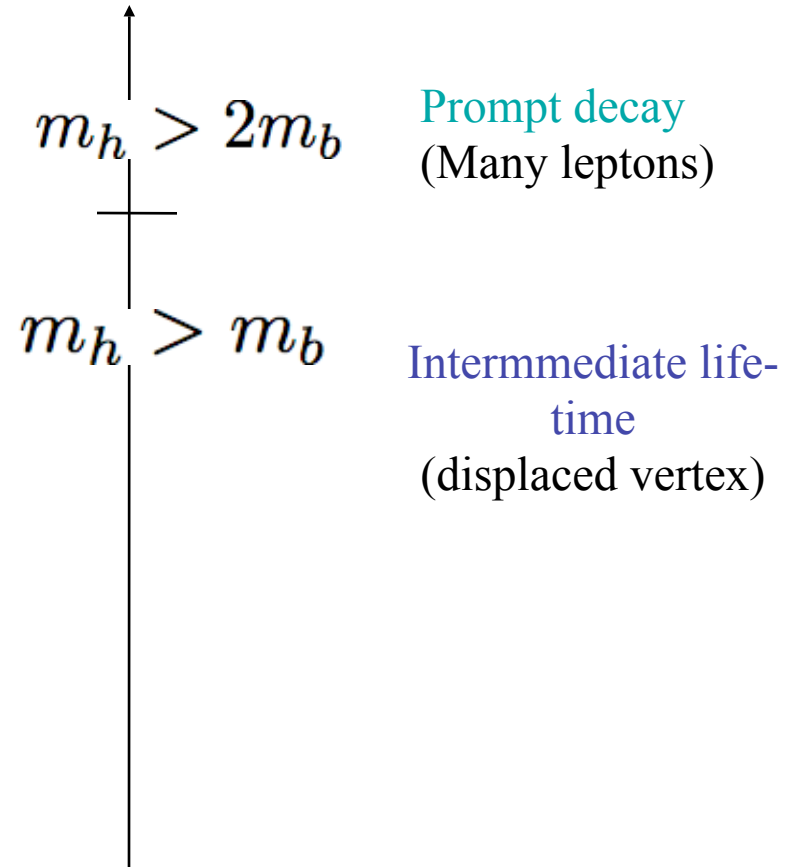
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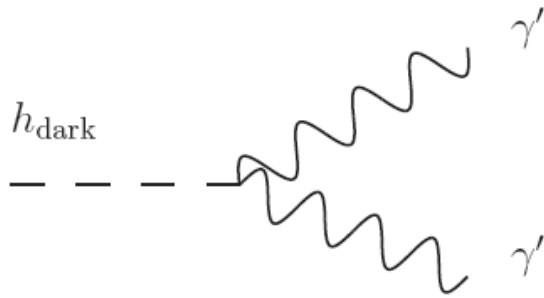
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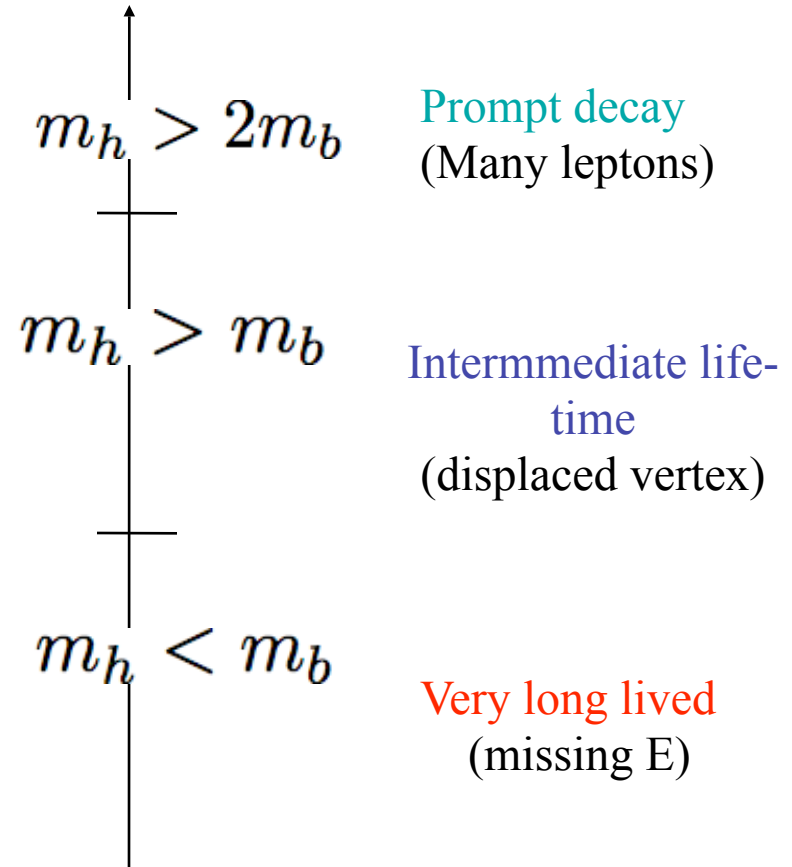
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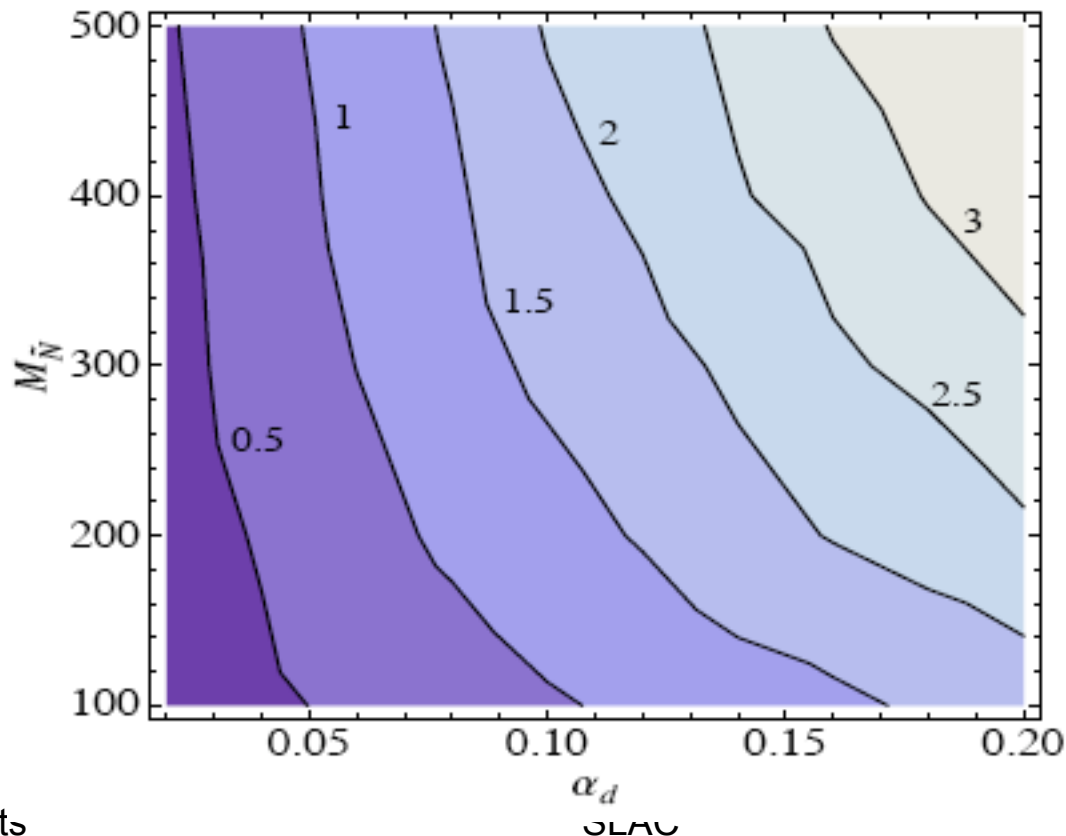
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Dark Radiation

Since the dark state are extremely boosted, they will radiate dark gauge-bosons,

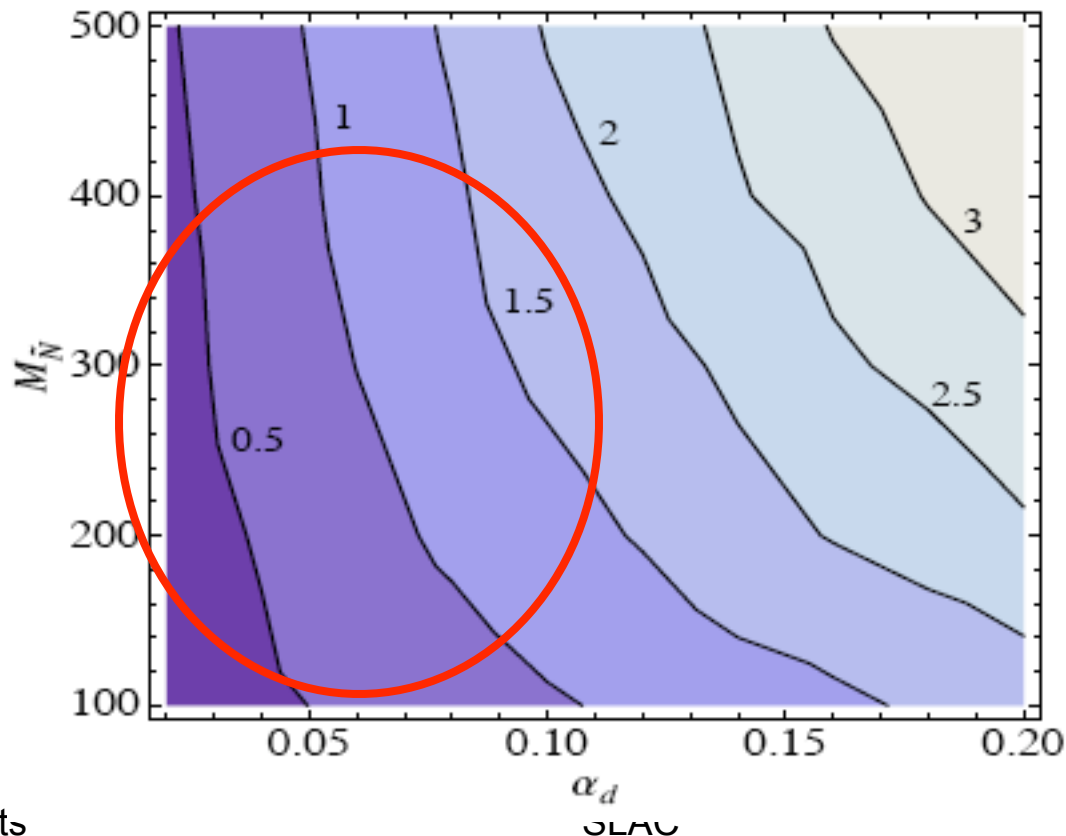
$$N_{\gamma'} \sim \frac{\alpha_d}{2\pi} \log \left(\frac{M_{\text{EW}}^2}{M_{\text{dark}}^2} \right)^2 \simeq 1.4 \left(\frac{\alpha_d}{0.1} \right)$$



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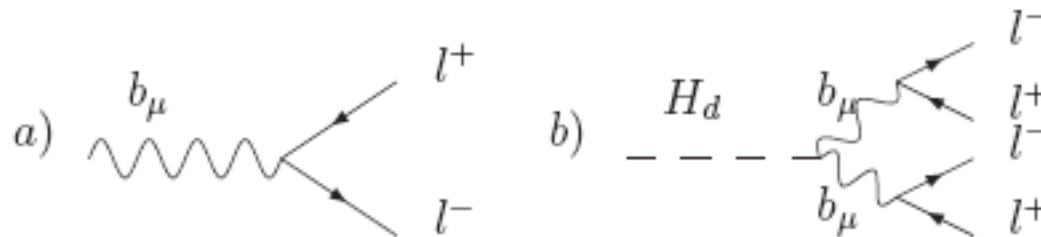


Cascades in the Dark

After showering finishes, the dark higgses will cascade down to the standard model. If we consider a simple model with 2 dark higgses, then there are several possibilities:

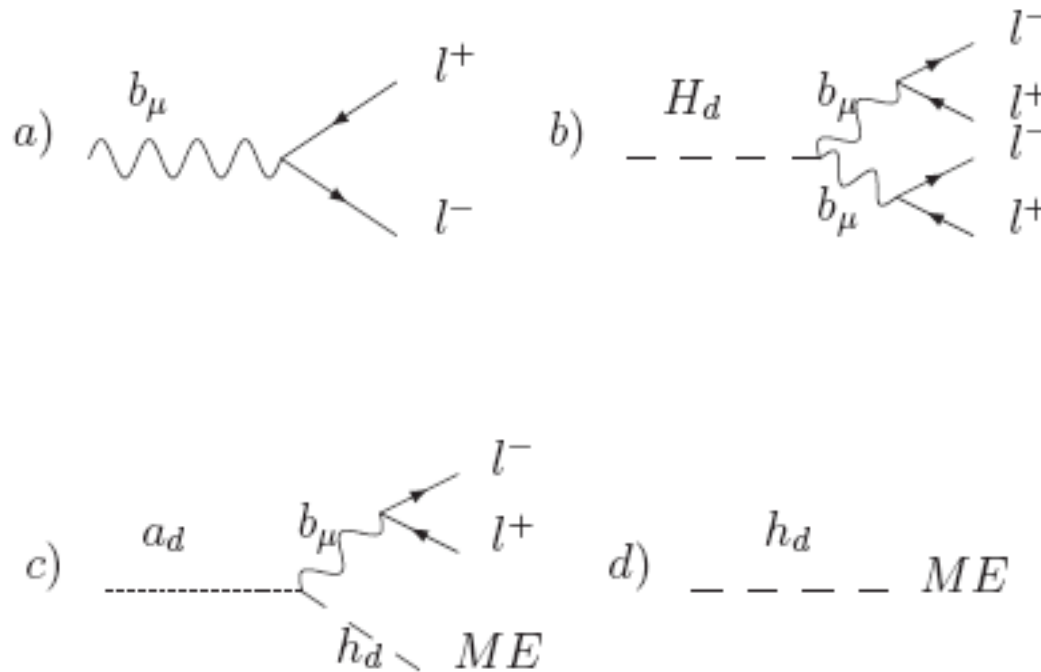
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Lepton Jets

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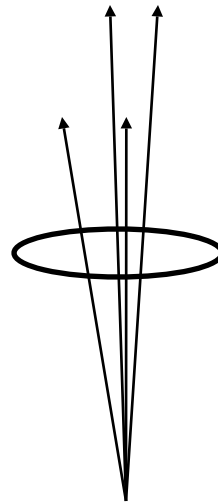
Two or more leptons with $p_T > 10$ GeV inside a cone of $\Delta R < 0.1$ with hadronic/leptonic isolation cut of $p_T < 3$ GeV in an annulus of $0.1 < \Delta R < 0.4$ around the leptons.

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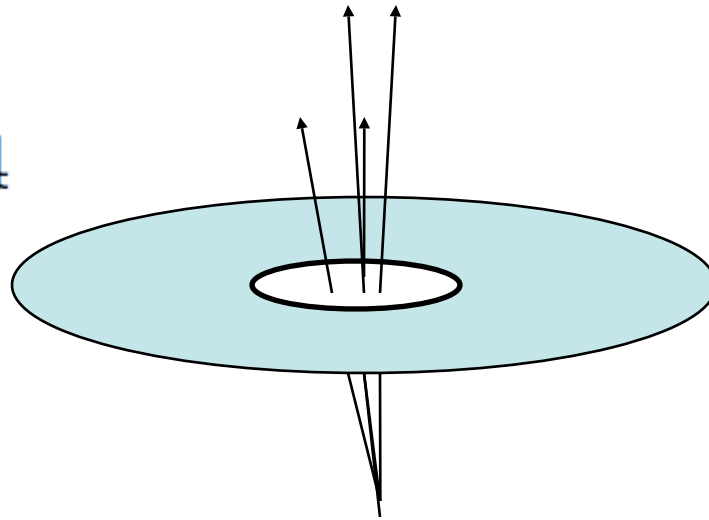


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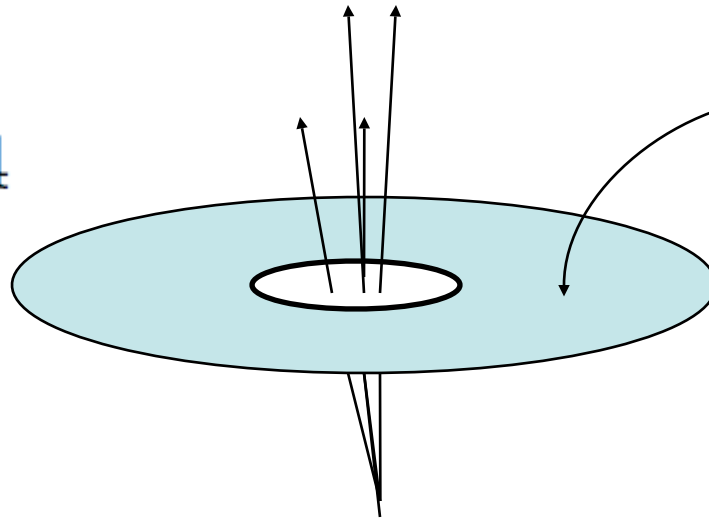


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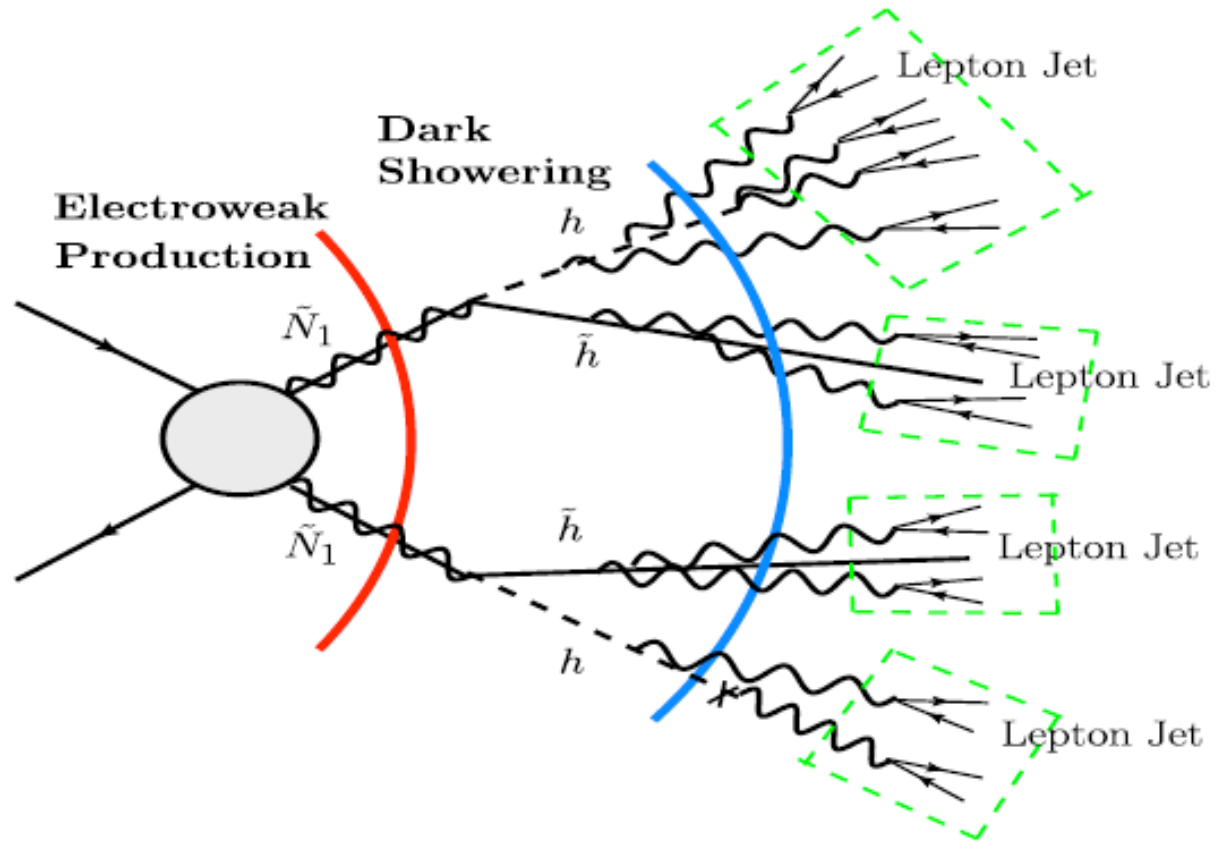
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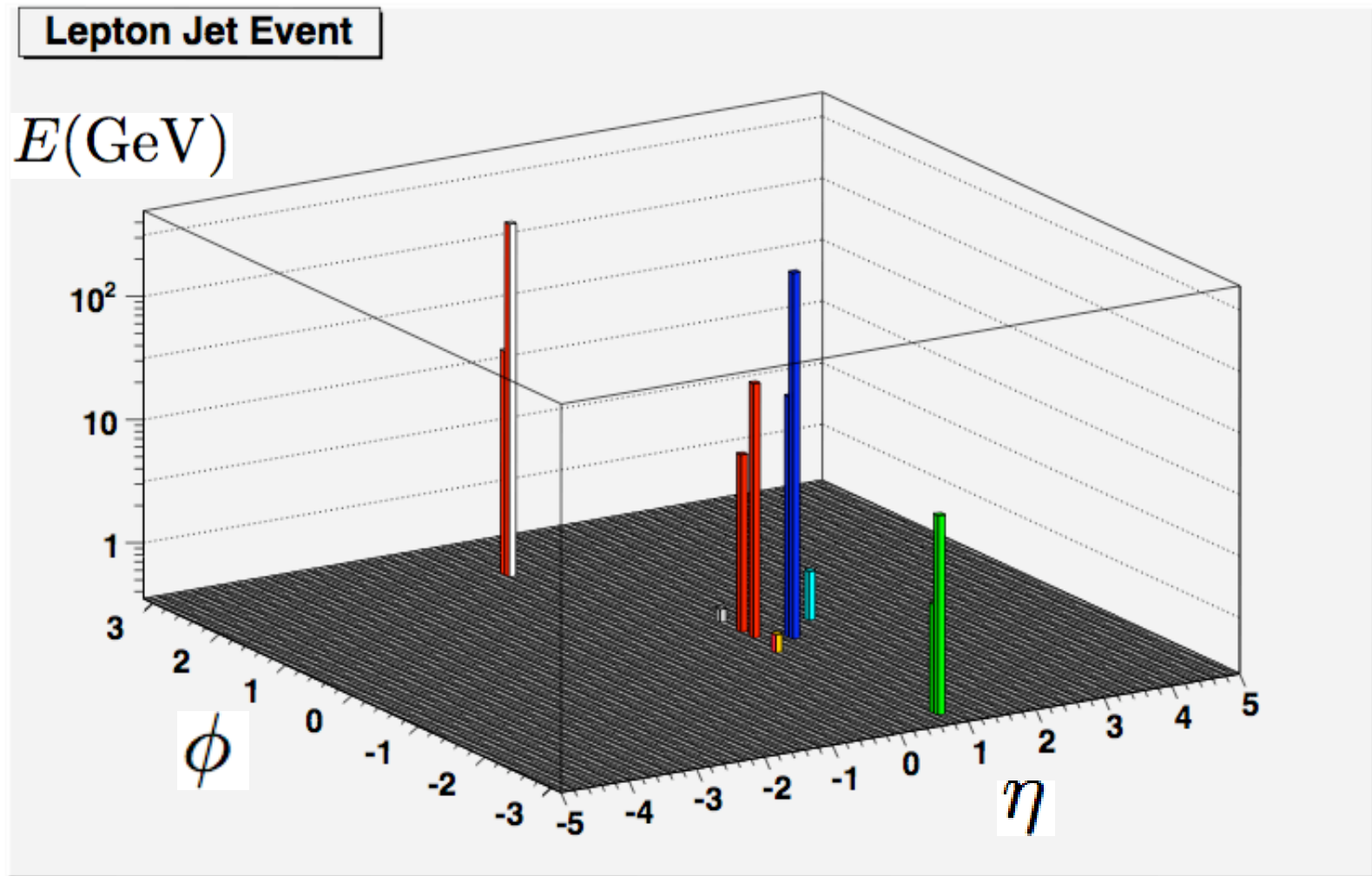
$$\sum p_T < 3 \text{ GeV}$$

Full Evolution



Lego Plots

For a 500 GeV LSP pair production, the event looks like:



Experimental Discovery

By defining lepton jets as a searchable object one can look for:

- 1) Lepton-jets + ME
- 2) Lepton-jets + QCD-jets
- 3) Lepton-jets + isolated leptons

While a resonance structure is probably present, since we don't know the mass, it may not very useful to implement mass-window cuts and etc.

Experimental Efforts

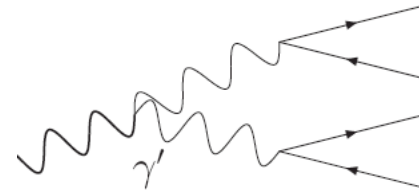
Several experimental groups are working on (designing) searches for lepton-jets

- 1) A. Haas and Y. Gershtein for D0 - Phys. Rev. Lett. 103, 061801 (2009), arXiv:0905:3381
- 2) B. Demirkoz and R. Moore for ATLAS - designing proper triggers for lepton jets.
- 3) K. Cranmer and the NYU group - lepton jet gun.
- 4) H. Lubatti and the Washington group - triggering on long lived neutral particles.
- 5) V. Halyo for CMS - searches for lepton jets.
- 6) Searches at BaBar - See all the local experts.

Sources of Lepton Jets

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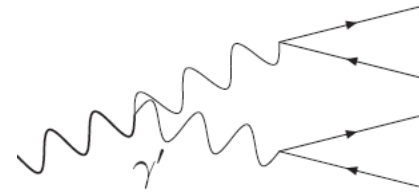
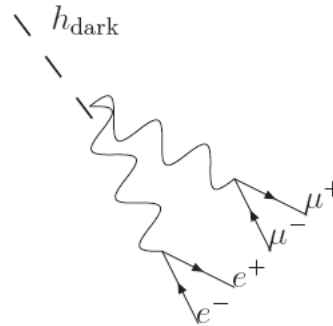
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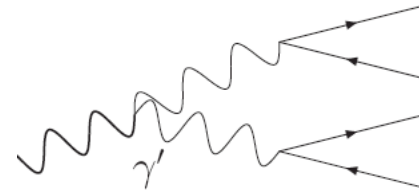
- 1) A non-abelian structure in the dark sector



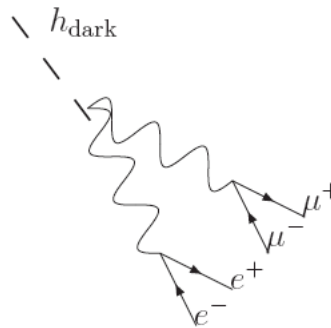
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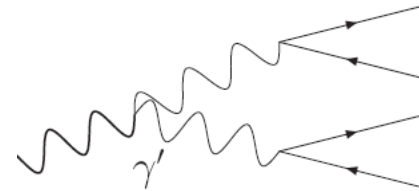
- 2) Dark higgs(es) decay



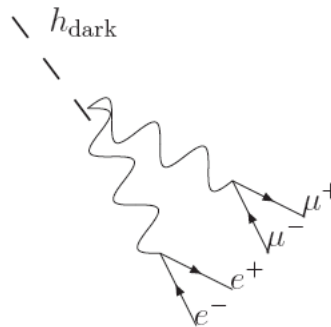
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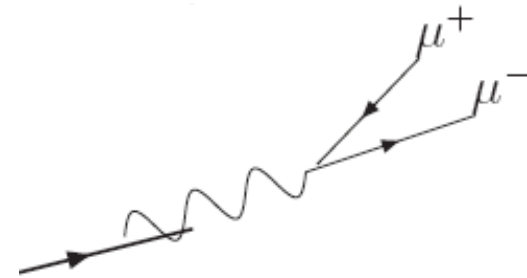
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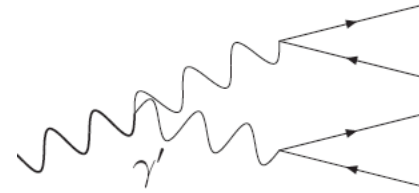
- 3) Dark radiation



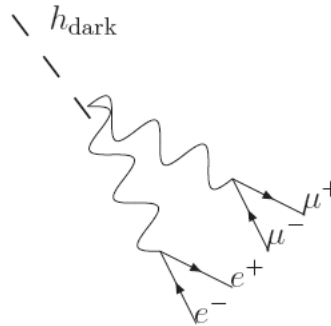
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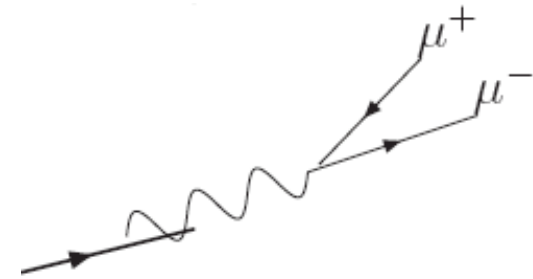
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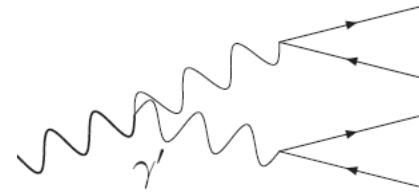


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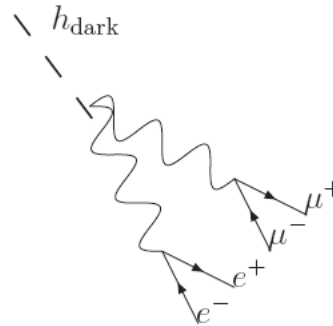
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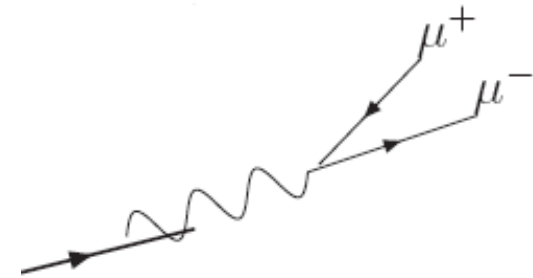
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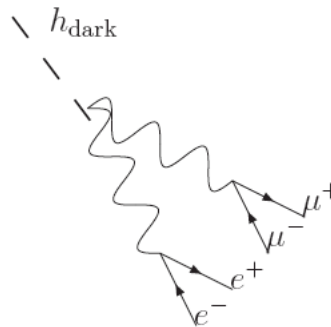
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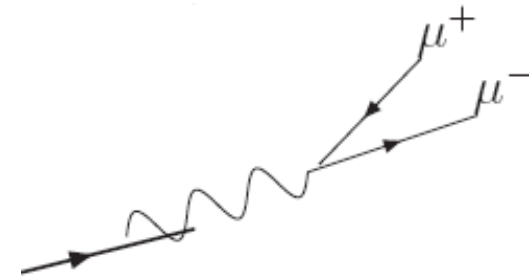
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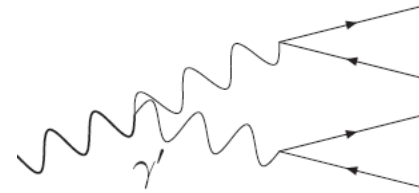
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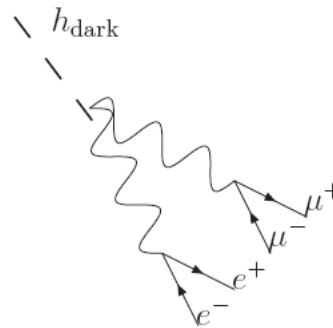
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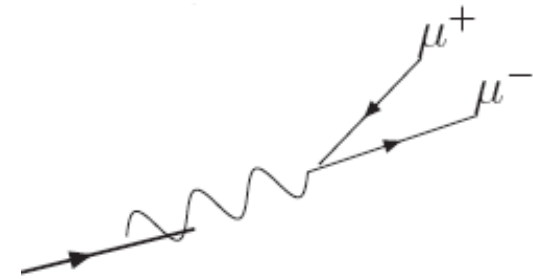
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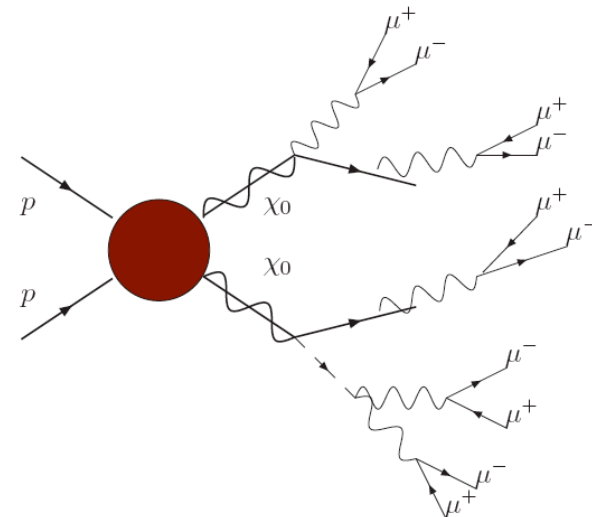


- 3) Dark radiation



There are different ways of producing dark states:

- 1) Prompt dark photon
- 2) Rare Z decays
- 3) Susy cascades



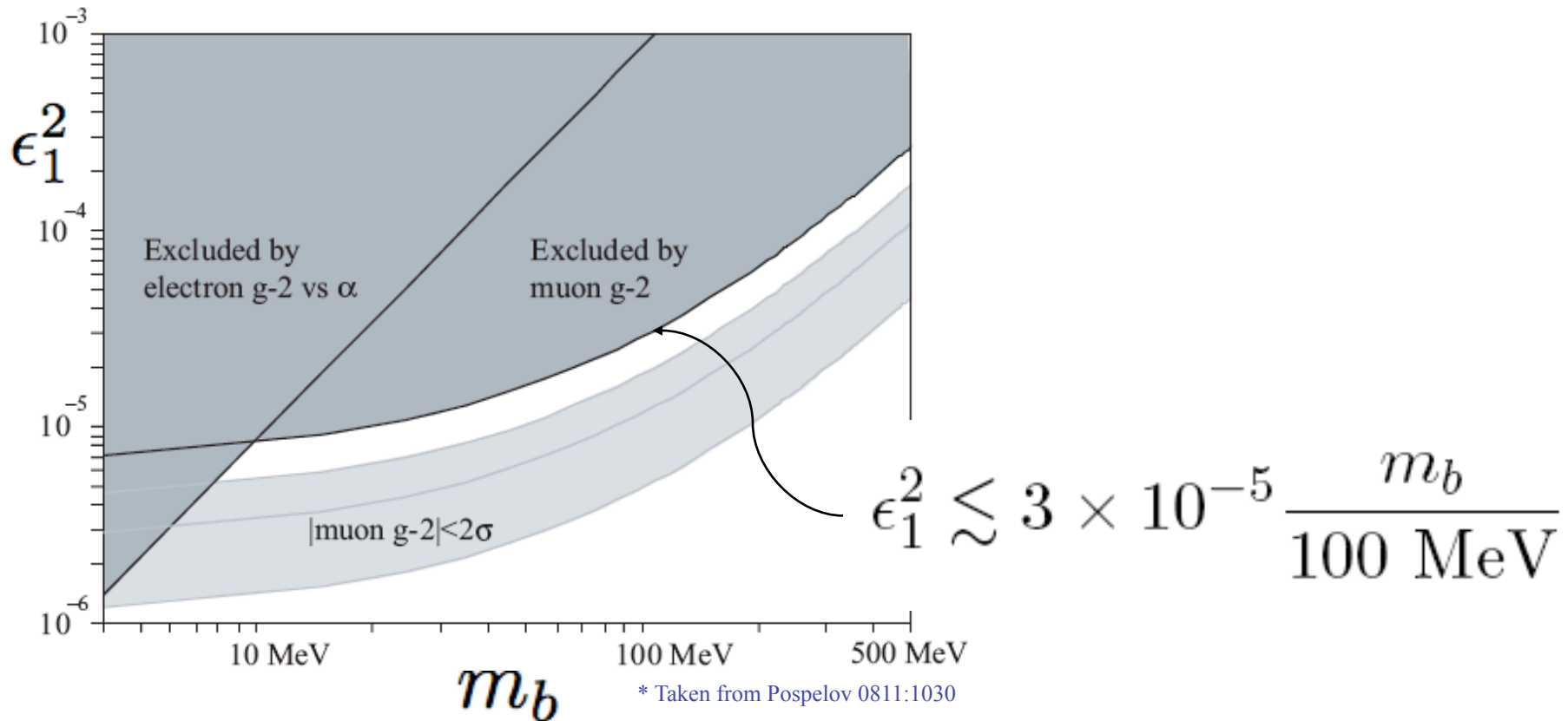
Thank You

Future Directions

- Searches at LEP and Tevatron.
- Searches at BaBar/Belle (see [Essig, Schuster, and Toro](#)) .
- Tune and modify triggers (see [Demirkoz](#) and [Moore](#)).
- Lepton-Jet observables?
- Other scenarios with similar signatures? (see [Strassler](#) and [Zurek](#)).

Limits on Kinetic Mixing

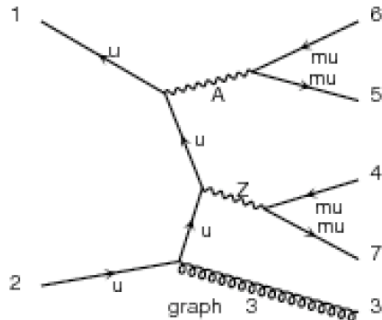
The kinetic mixing with the photon is bounded by low energy experiments, in particular the muonic g-2 ratio (Pospelov 0811:1030):



Notice that this measurement does not bound ϵ_2

Standard Model Background

The SM can give 2 muon pairs recoiling against a jet and that is an irreducible background. Simulation with Madgraph suggest that this is not going to be a serious obstacle:



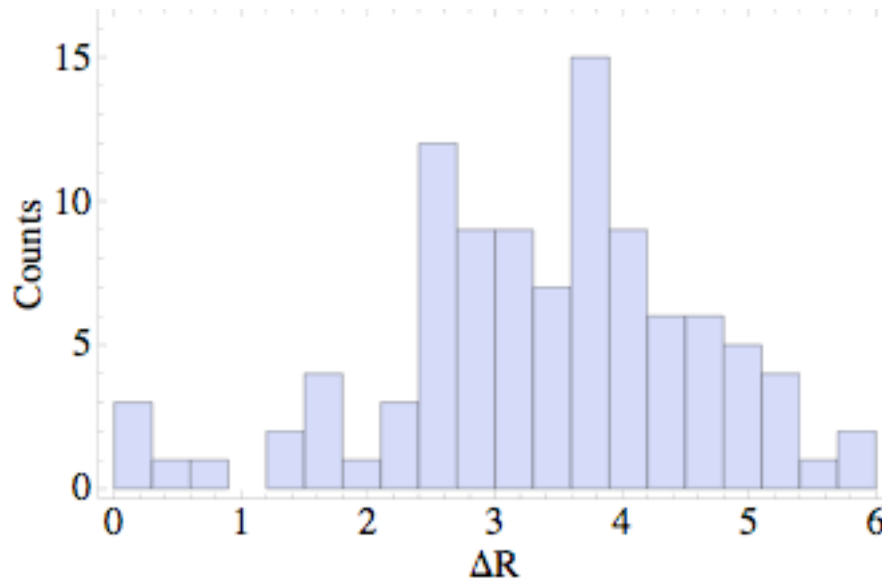
... and many more ...

$$\sigma = 37 \text{ fb}$$

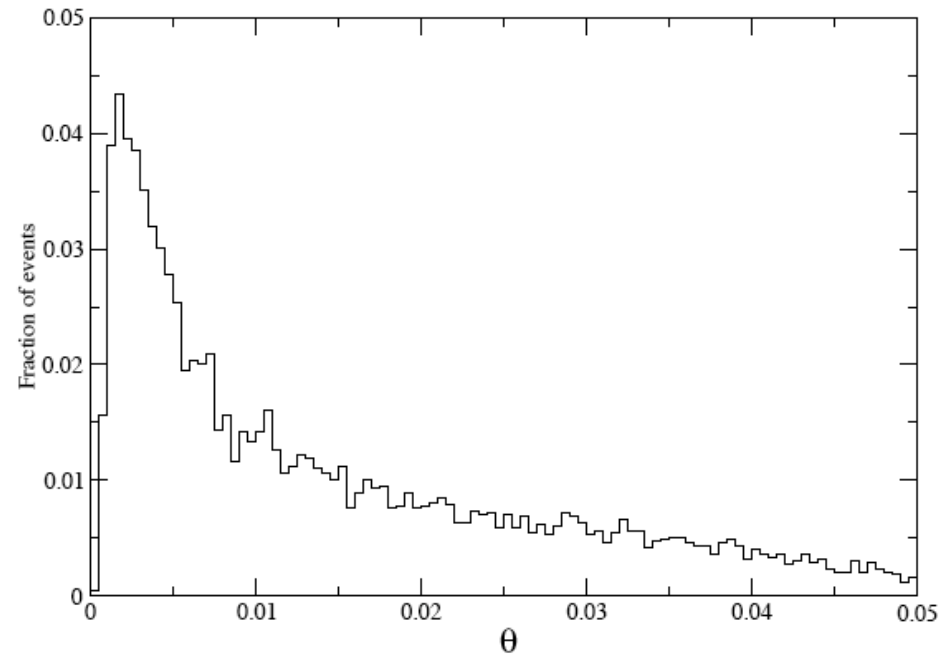
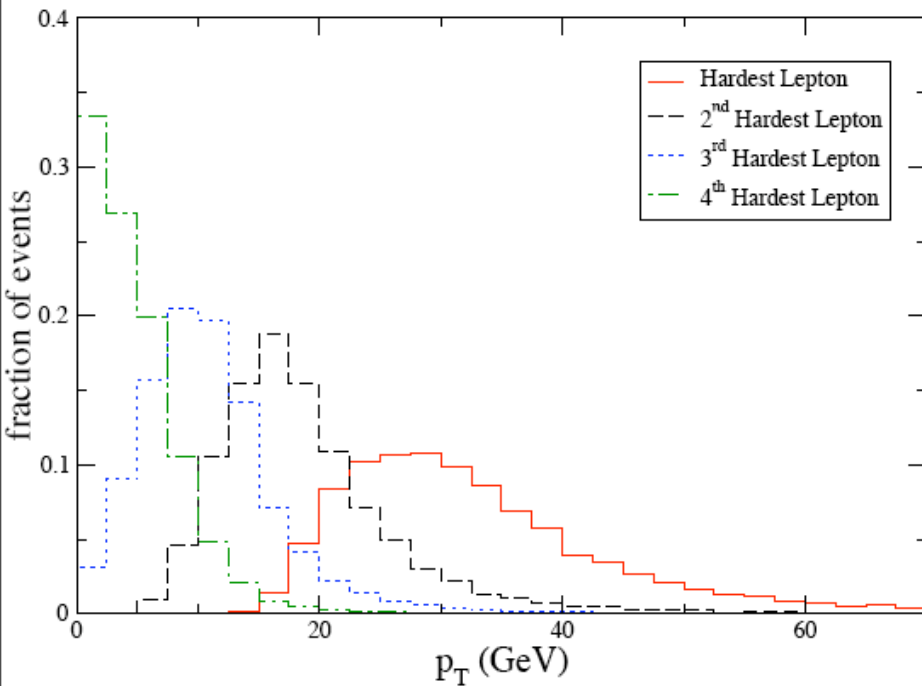
This can be reduced to $< 1 \text{ fb}$ by requiring:

$$\Delta R < 0.1$$
$$p_T > 3 \text{ GeV}$$

For all 4 muons.



Detector Performance



Bilge Demirköz and **Roger Moore** investigated ATLAS performance using the prompt dark photon production as a benchmark.

Bilge Demirköz also implemented new triggers to help improve the efficiency associated with such events.

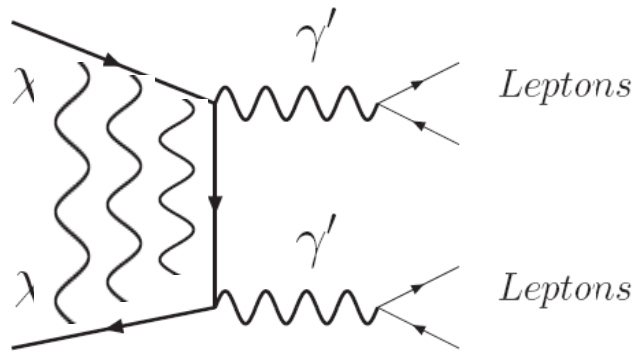
Lepton Jet Efficiency

Lepton Jet Efficiencies						
	1 Lepton-Jet			2 Lepton-Jet		
$\text{Br}_{b \rightarrow \pi\pi}$	1/7	1/3	3/5	1/7	1/3	3/5
α_d						
0	0.49 (0.49)	0.47 (0.47)	0.31 (0.31)	0.28 (0.28)	0.14 (0.15)	0.05 (0.05)
0.01	0.47 (0.47)	0.44 (0.45)	0.31 (0.32)	0.3 (0.31)	0.16 (0.16)	0.04 (0.04)
0.03	0.43 (0.41)	0.47 (0.48)	0.3 (0.3)	0.27 (0.3)	0.14 (0.16)	0.04 (0.05)
0.1	0.43 (0.39)	0.41 (0.44)	0.29 (0.32)	0.23 (0.3)	0.13 (0.18)	0.05 (0.07)
0.3	0.38 (0.32)	0.34 (0.36)	0.25 (0.34)	0.16 (0.3)	0.11 (0.22)	0.05 (0.09)

Table 1: Clean lepton jet efficiencies for different values of the dark gauge-coupling and $\text{Br}(b \rightarrow \pi^+ \pi^-)$. The neutralino mass was set to $\tilde{M} = 300$ GeV. For $\alpha_d = 0$ dark radiation was switched off. The number of lepton jets increases with α_d as radiation becomes more likely. The requirement for "clean" lepton jets, as described in the text, results in a decrease in efficiency with the growth of the branching ratio into pion. In brackets are efficiencies for the case where only hadronic isolation is required in the $0.1 < \Delta R < 0.4$ annulus.

Resolution of PAMELA

So dark matter annihilates to dark photons first. The dark photons then decay into leptons.



- 1) **Protons** are kinematically disallowed.
- 2) The leptons are **direct products** of the annihilations.
- 3) **Sommerfeld enhancement** of the cross-section due to light particle exchange.

