

Probe Seclude $U(1)$ at BES

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SLAC

Outline

- Introduction
- Invisible decays of mesons @BESII/III
 - Limit of invisible decays from BESII data;
 - Limit of invisible decays at BESIII (see Yang-Heng Zheng's talk);
- Signature of U boson at BESIII
 - Reach of U boson in direct production;
 - Reach of U boson decays in charmonium decays;
- Summary

Introduction

- Experimental signature of Super symmetry is focused on heavy particles with 1TeV or above, however,
- A possible U(1) symmetry, which requires a new gauge boson could be light and weakly interacting with SM particles:

P. Fayet PLB95, 285(1980)

B. Holdom PLB166, 196(1986)

N.Arakani-Hamed and N.Weiner
JHEP 12,104(2008)

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \epsilon_Y F^{Y,\mu\nu} F'_{\mu\nu} + \frac{1}{4} F'^{\mu\nu} F'_{\mu\nu} + m_{A'}^2 A'^{\mu} A'_{\mu},$$

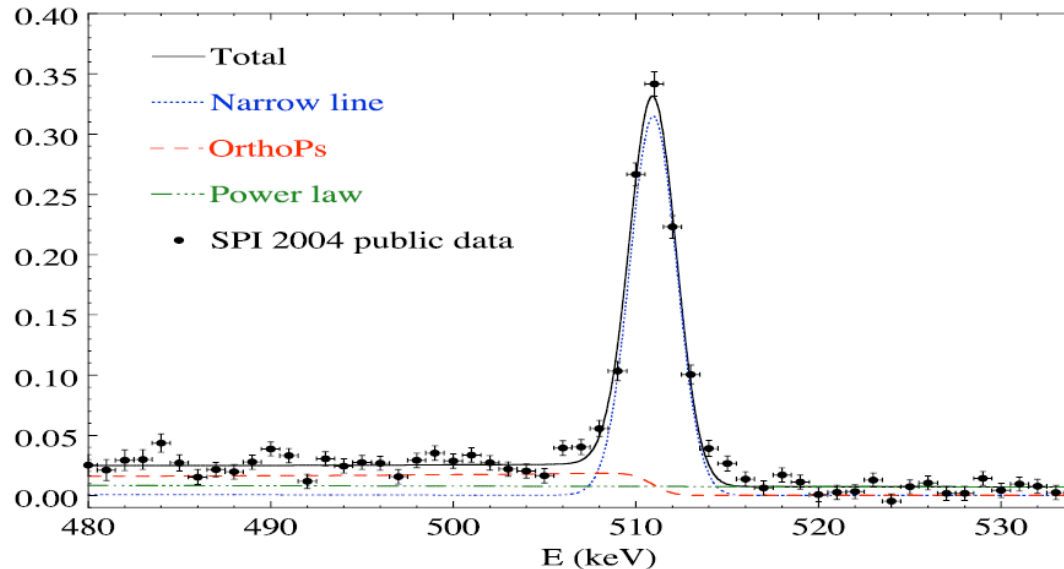
- The possible connection of the U-boson and light DM is interpreted in a view of electron /positron excesses observed by PAMELA, ATIC, FERMI.

Introduction

The SPI Spectrometer aboard the INTEGRAL satellite observed a large flux of photons with energy of 511 keV:

P. Jean et al., *Astron. Astrophys.* 407, L55(2003)

J. Knodlseder et al., *Astron. Astrophys.* 411, L457(2003).



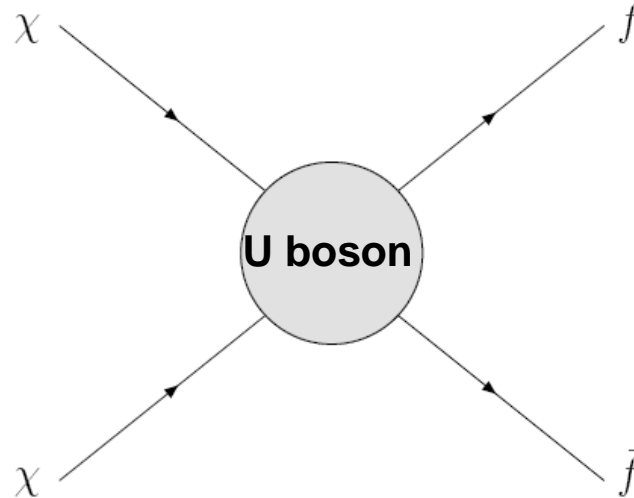
Such sharp line can only come from the annihilation of non-relativistic e^+e^- pairs .
Estimates of e^+ production by astrophysical source fall well short of required flux.

Possible interpretation: annihilation of light DM particles $\chi\chi \rightarrow e^+e^-$:

[C. Boehm, D. Hooper, J. Silk, M. Casse and J. Paul, *PRL* 92, 101301 (2004)]

Light DM candidates and light U boson

Possible massive vector mediates interaction among MeV (scalar or majorana) dark matter and electron-positron:



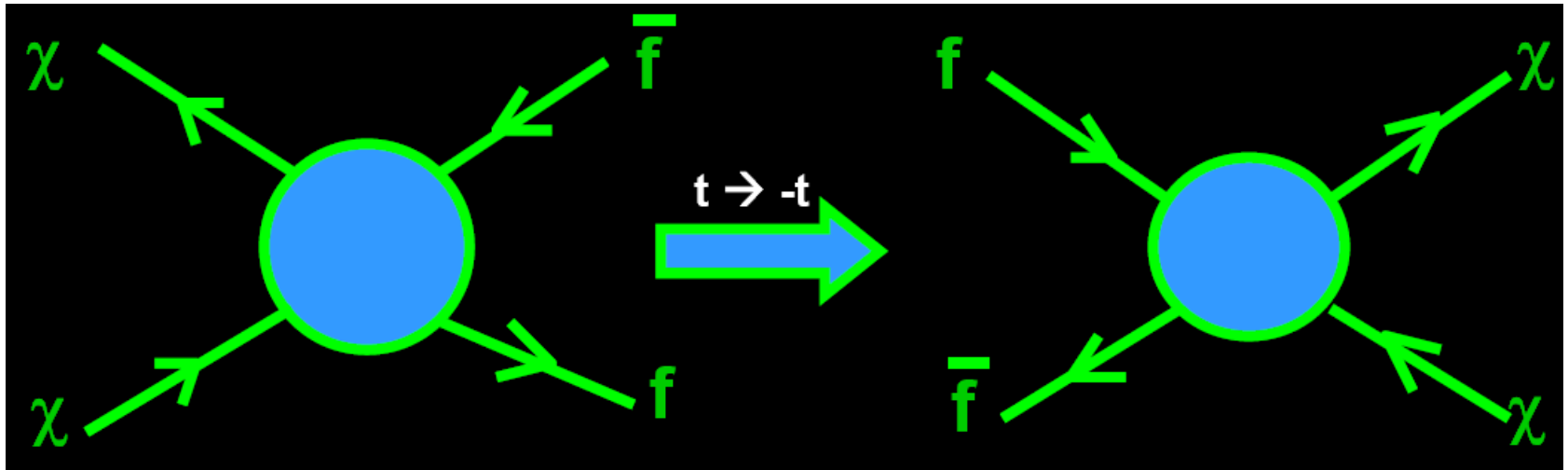
From astrophysical constraints, $m_\chi > 20\text{MeV}$ could be excluded due to over-production of MeV photons through $\chi\chi \rightarrow e^+e^-\gamma$

[J.F. Beacom, N.F. Bell and G. Bertone, PRL 94, 171301(2005)]

Hence, one has:

$$m_e \leq m_\chi \leq 20 \text{ MeV}$$

Invisible decay of mesons



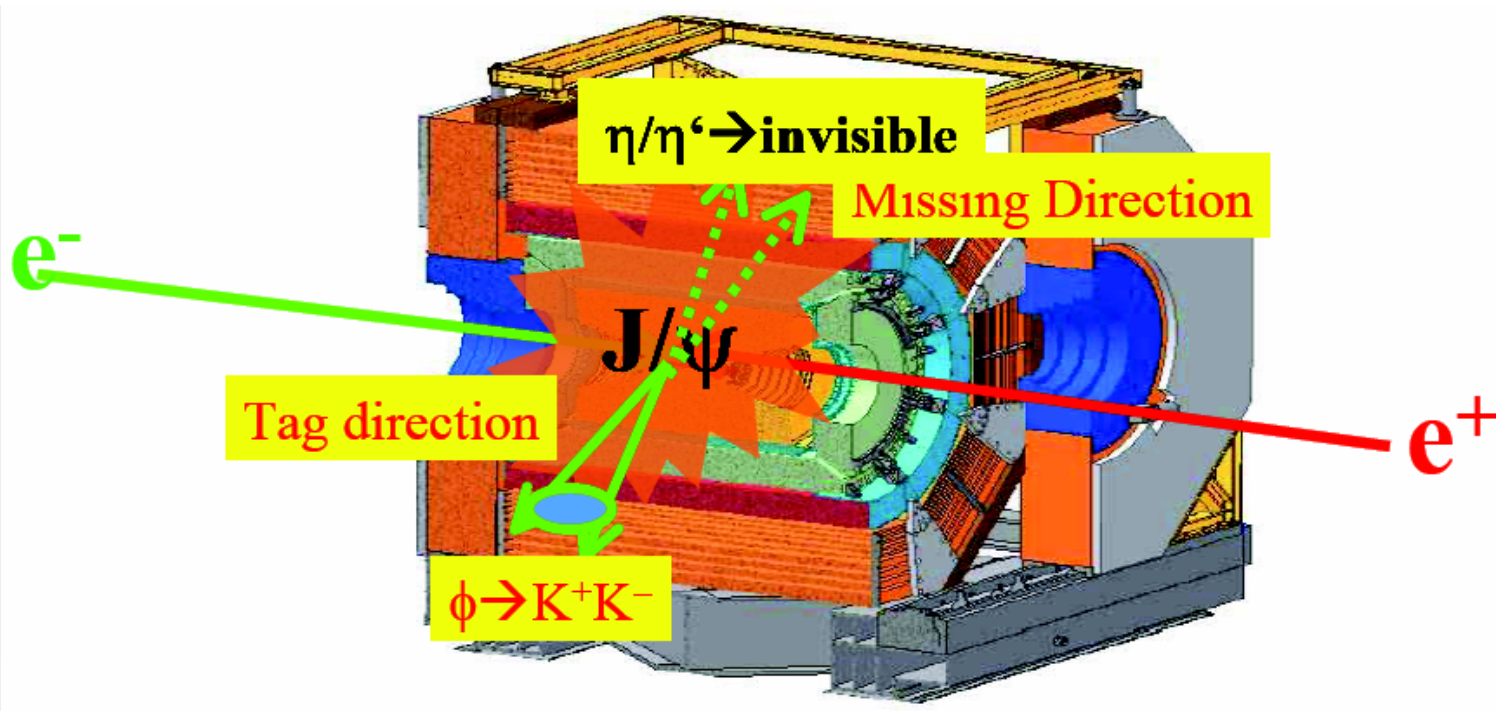
The time reversed annihilation diagram corresponds to invisible decays of quarkonia.

Measuring the invisible decays give direct sensitivity to the J^{PC} of the mediator.

B. McElrath PRD 72, 103508 (2005)

We have many $f\bar{f}$ bound states : $\pi^0, \eta, \omega, \rho, \phi, \eta_c, \chi_c, \psi(nS), Y(ns)$.

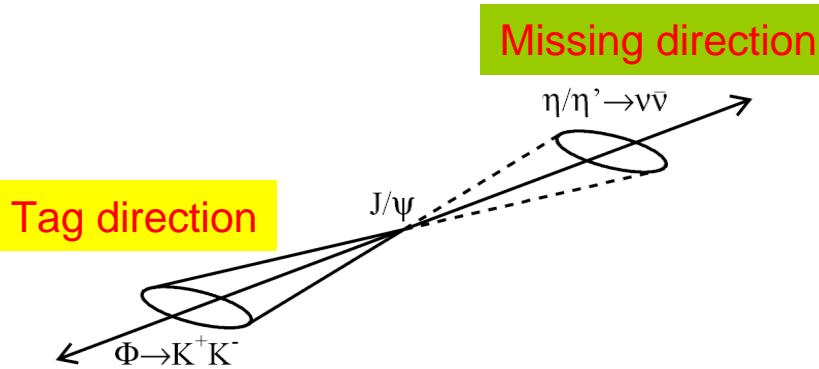
Invisible decays of η/η' at BES



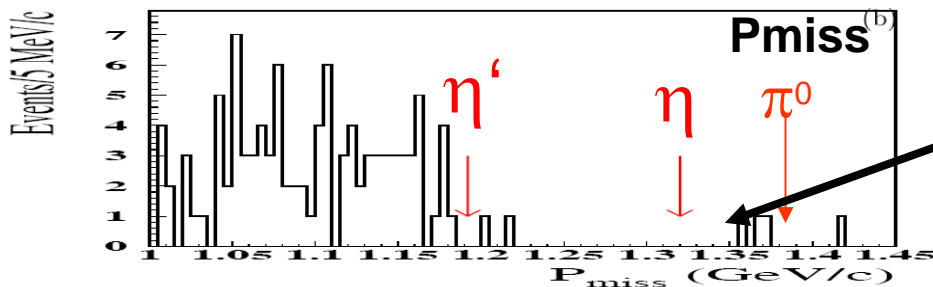
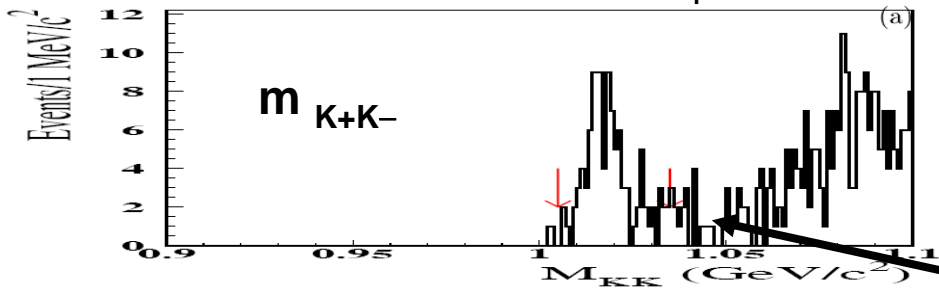
- o $J/\psi \rightarrow \phi \eta/\eta(958)$, where $\eta/\eta(958) \rightarrow \text{Invisible}$, both hadrons are boosted.
- o Reconstructing $\phi \rightarrow K^+K^-$, looking at missing mass or missing momentum of ϕ .
- o No any hit information are required outside of the core of ϕ decays.

Reconstruction of $J/\psi \rightarrow \phi(K^+K^-) + \text{nothing}$

- Two good charged Kaons in MDC, $|\cos(\theta)| < 0.8$; track from IP; PID as kaons;
- No clusters in BSC in the region of outside cones of 30° around the charged kaons.
- For the missing momentum: $|\cos(\theta_{\text{miss}})| < 0.7$;
- $1.005 < m(K^+K^-) < 1.035 \text{ GeV}$ (3σ);



58 million J/ψ @BESII



Cleanly tagging "nothing"
Background free at BES!

Background estimations in $\eta/\eta' \rightarrow$ invisible

- Class I: $J/\psi \rightarrow \phi \eta(\eta')$, $\phi \rightarrow K^+K^-$, $\eta(\eta') \rightarrow$ visible final states, these expected backgrounds are negligible small.
- Class II: $J/\psi \rightarrow \phi + \text{non-}\eta(\eta')$, or non- $\phi\eta$
 - (1) $J/\psi \rightarrow \phi K_L K_L$, $\phi f_0(980)$ [$f_0(980) \rightarrow K_L K_L$];
 - (2) $J/\psi \rightarrow K^{*0} K_L$, $K^{*0} \rightarrow K\pi$

The expected backgrounds

$$N(\eta) = 3.0 \pm 0.2$$

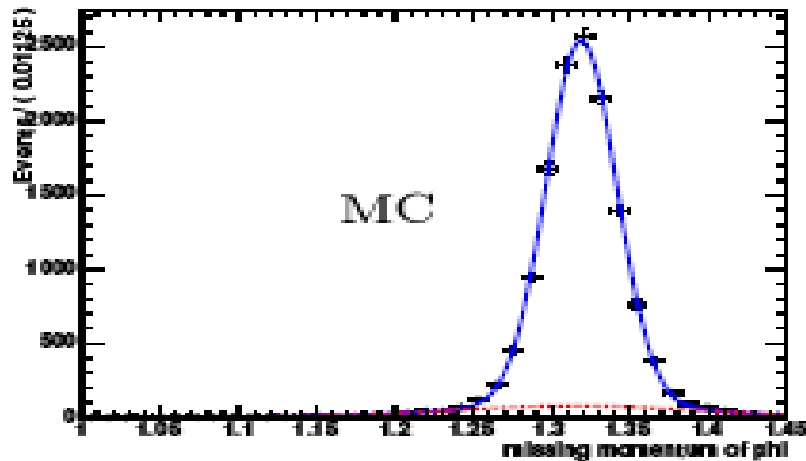
$$N(\eta') = 90 \pm 64$$

Looking for $\eta/\eta' \rightarrow$ Invisible decays in $J/\psi \rightarrow \phi\eta/\eta'$

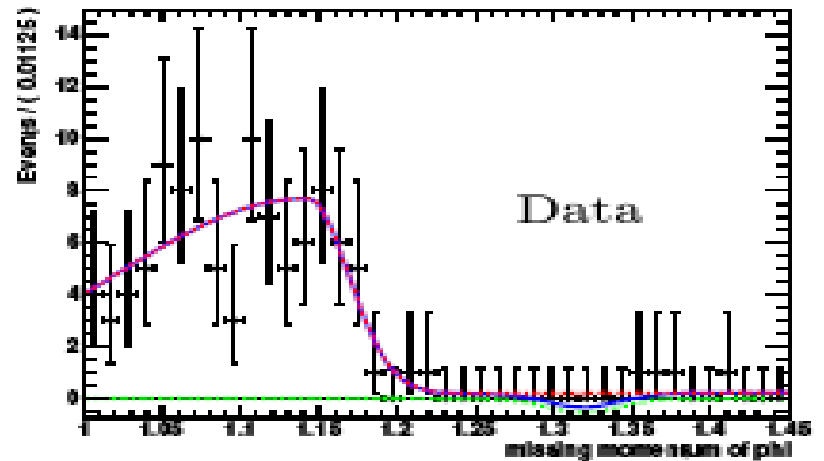
An unbinned extended Maximum Likelihood Fit:

$$\mathcal{L}(N_{sig}^{\eta}, N_{sig}^{\eta'}, N_{bkgd}) = \frac{e^{-(N_{sig}^{\eta} + N_{sig}^{\eta'} + N_{bkgd})}}{N!} \times \prod_{i=1}^N [N_{sig}^{\eta} \mathcal{F}_{sig}^{\eta}(P_{miss}^i) + N_{sig}^{\eta'} \mathcal{F}_{sig}^{\eta'}(P_{miss}^i) + N_{bkgd} \mathcal{F}_{bkgd}(P_{miss}^i)],$$

A RooPlot of "missing momentum of phi"



A RooPlot of "missing momentum of phi"



$$N[J/\psi \rightarrow \phi\eta (\eta \rightarrow \text{invisible})] = -2.8 \pm 1.4$$

$$N[J/\psi \rightarrow \phi\eta' (\eta' \rightarrow \text{invisible})] = 2.2 \pm 3.4$$

Looking for $\eta/\eta' \rightarrow$ Invisible Decays in $J/\psi \rightarrow \phi\eta/\eta'$

$$\frac{\mathcal{B}(\eta \rightarrow \text{invisible})}{\mathcal{B}(\eta \rightarrow \gamma\gamma)} < \frac{n_{UL}^\eta / \epsilon_\eta}{n_{\gamma\gamma}^\eta} \cdot \frac{1}{(1 - \sigma_\eta)}$$

n_{UL}^η : #observed signal @90%CL

n_{UL}^γ : #observed $\eta \rightarrow \gamma\gamma$

ϵ_η : Efficiency of signal mode

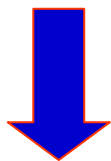
σ_η : 1 sigma uncertainty

$$\frac{BR(\eta \rightarrow \text{invisible})}{BR(\eta \rightarrow \gamma\gamma)} < 1.65 \times 10^{-3} \text{ @ 90\% C.L.}$$

$$\frac{BR(\eta' \rightarrow \text{invisible})}{BR(\eta' \rightarrow \gamma\gamma)} < 6.69 \times 10^{-2} \text{ @ 90\% C.L.}$$

Many uncertainty related to tag side and detector noise can cancel in the ratio.

BES, Phys.Rev.Lett.97:202002,2006



$$Br(\eta \rightarrow \chi\chi) < 6.4 \times 10^{-4} \text{ @90\%}$$

$$Br(\eta' \rightarrow \chi\chi) < 1.4 \times 10^{-3} \text{ @90\%}$$

Theory:

B. McElrath PRD 72, 103508 (2005)

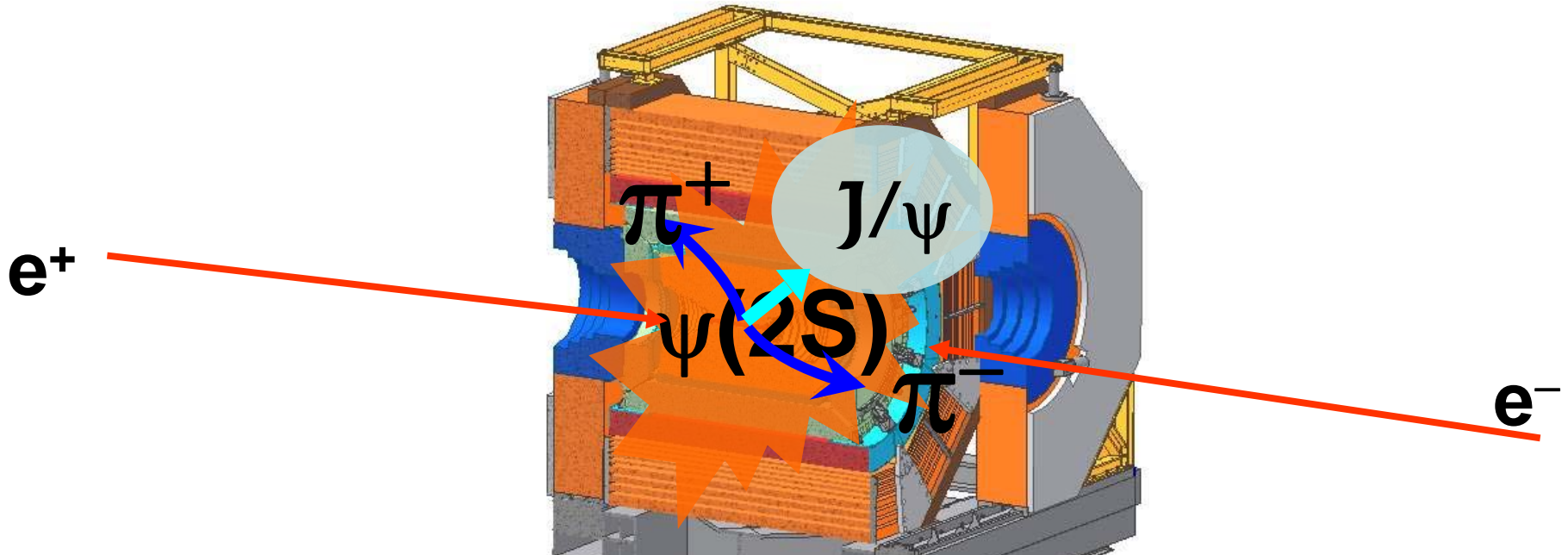
$$Br(\eta \rightarrow \chi\chi) \sim 7.4 \times 10^{-5}$$

$$Br(\eta' \rightarrow \chi\chi) \sim 8.1 \times 10^{-7}$$



$$\psi(2S) \rightarrow \pi^+ \pi^- \text{ J}/\psi \text{ invisible}$$

Invisible means nothing besides two charged pions are seen in tracking and calorimetry systems at BESII.



The charged-track trigger@BESII:

- (1) at least one hit in the 48 barrel TOF counter array
- (2) at least one track in the tracking system
- (3) at least 100 MeV of energy deposit in the BSC

The trigger is sensitive to two soft pions in this decays

$$\psi(2S) \rightarrow \pi^+ \pi^- J/\psi^{\text{invisible}}$$

- Two good charged pions:
from IP; $|\cos(\theta)| < 0.8$; each pion: $p_\pi < 0.4 \text{ GeV}/c$; $m(\pi\pi) > 0.35 \text{ GeV}$
- Requirement in Vertex Chamber(VC):
Events with additional tracks from IP penetrate at least 6 layers in inner vertex chamber, should be rejected.
- Requirement in Barrel Shower Counter:
The number of SC cluster, with $E > 10 \text{ MeV}$, and no matched with either charged tracks in MDC, is required to be zero.

The $\psi \rightarrow \pi^+ \pi^- J/\psi (\rightarrow \mu\mu)$ are also reconstructed to make the ratio, so that many systematic errors related to the soft pions cancel:

$$\frac{\mathcal{B}(J/\psi \rightarrow \text{invisible})}{\mathcal{B}(J/\psi \rightarrow \mu^+ \mu^-)}$$

$\psi(2S) \rightarrow \pi^+ \pi^- \psi(1S)$ invisible backgrounds

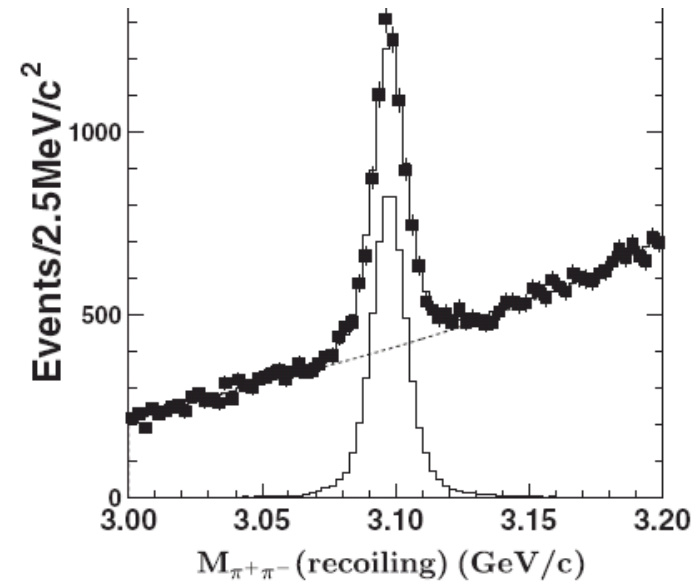
- **Two-photon**
 - 2 prong $\pi\pi, ee, \mu\mu \dots$
 - 3 body $\pi^+\pi^-\pi^0 \dots$
- $J/\psi \rightarrow \mu\mu, ee \dots$ (outside of acceptance)
 - Same distribution as signal (peaking bkg)
 - Huge contribution as signal for BESII

TABLE I. Expected number of events (N_{bg}) and efficiencies for peaking backgrounds.

Background channel [$\psi(2S) \rightarrow \pi^+ \pi^- J/\psi, J/\psi \rightarrow$]	Efficiency (%)	Expected N_{bg}
$\mu^+ \mu^-$	0.964	2543 ± 254
$e^+ e^-$	0.907	2393 ± 240
$n\bar{n}$	10.46	1011 ± 85
$p\bar{p}$	0.434	42 ± 13
$n\bar{n}\pi^0$	0.486	29 ± 10
Total		6018 ± 514

$\psi(2S) \rightarrow \pi^+ \pi^- J/\psi$ invisible

Huge standard model backgrounds as expected:



By using 14 million $\psi(2S)$ decays, 90% upper limit:

$$\frac{\mathcal{B}(J/\psi \rightarrow \text{invisible})}{\mathcal{B}(J/\psi \rightarrow \mu^+ \mu^-)} < \frac{N_{UL}^{J/\psi} / \epsilon_{\text{invisible}}}{N_{\mu^+ \mu^-}^{J/\psi} / \epsilon_{\mu^+ \mu^-}^{J/\psi}} = 1.2 \times 10^{-2}$$

BES, Phys. Rev. Lett. 100: 192001 (2008)

Within SM, the invisible quarkonium decays are predicted by:

hep-ph/9806487, L.N Chang, O. Lebedev and J. N. Ng

$$\begin{aligned} \frac{\Gamma(\Upsilon \rightarrow \nu \bar{\nu})}{\Gamma(\Upsilon \rightarrow e^+ e^-)} &= \frac{27G^2 M_\Upsilon^4}{64\pi^2 \alpha^2} \left(-1 + \frac{4}{3} \sin^2 \theta_W\right)^2 \\ &= 4.14 \times 10^{-4}, \end{aligned}$$

$$\begin{aligned} \frac{\Gamma(J/\Psi \rightarrow \nu \bar{\nu})}{\Gamma(J/\Psi \rightarrow e^+ e^-)} &= \frac{27G^2 M_{J/\Psi}^4}{256\pi^2 \alpha^2} \left(1 - \frac{8}{3} \sin^2 \theta_W\right)^2, \\ &= 4.54 \times 10^{-7}, \end{aligned}$$

BESIII milestones

Peak Lumi. @ Nov. 2008:
 $1.2 \times 10^{32} \text{cm}^{-2}\text{s}^{-1}$
Peak Lumi. @ May 2009:
 $3.2 \times 10^{32} \text{cm}^{-2}\text{s}^{-1}$

July 18, 2008: First e^+e^- collision event in BESIII

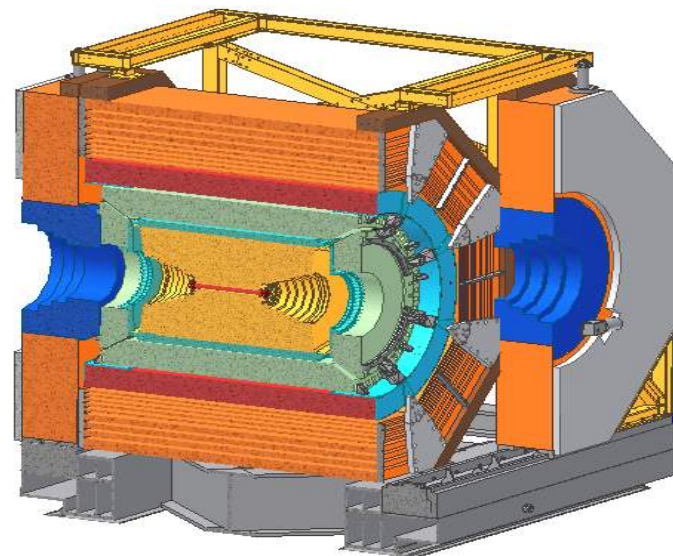
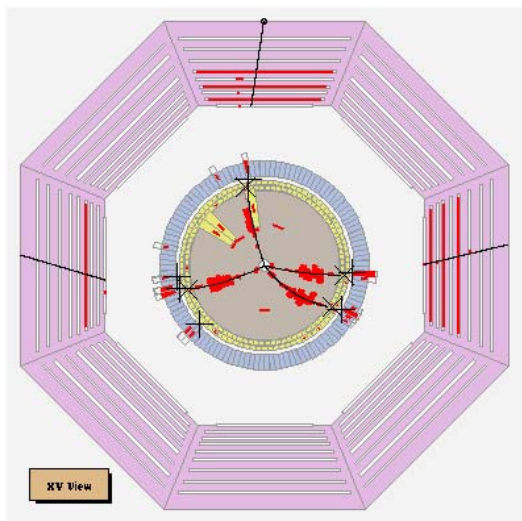
Nov. 2008: $\sim 14\text{M}$ $\psi(2\text{S})$ events collected

April 14, 2009 $\sim 100\text{M}$ $\psi(2\text{S})$ events collected

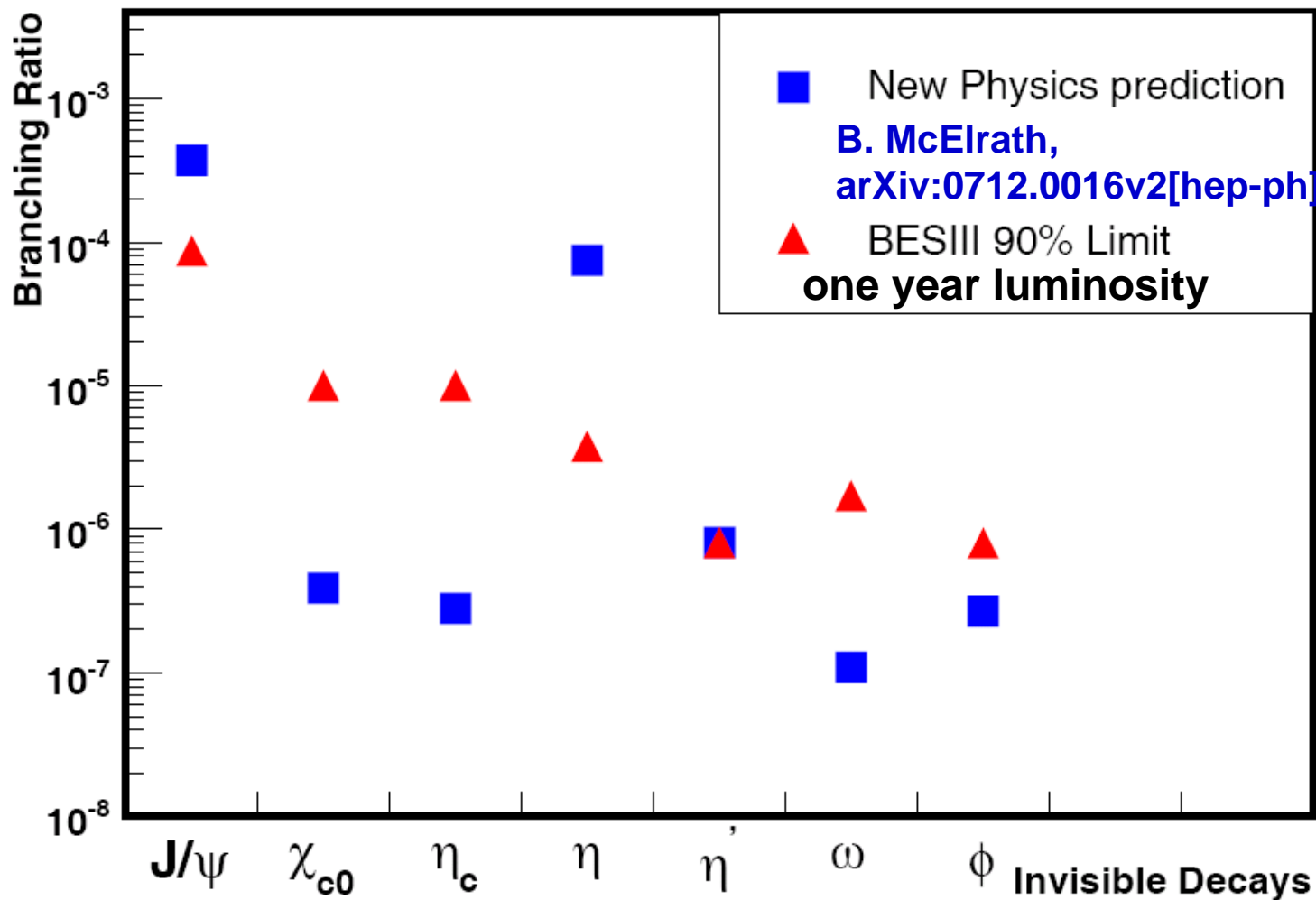
May 30, 2009 42pb^{-1} at continuum collected

July 28, 2009 $\sim 230\text{M}$ J/ψ events collected

See Yang-Heng Zheng's slides



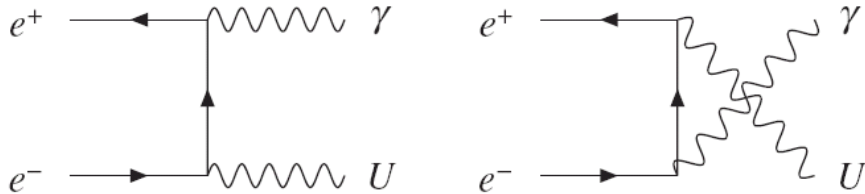
Sensitivities of invisible decays of mesons@BESIII



See Yang-Heng's talk for detail

Signature of light U boson at BEPCII/BESIII

Direct production of spin-1 boson:



P. Fayet, PRD75, 115017(2007);
 R. Essig, P.Schuster and
 N. Toro 0903.3941[he-ph];
 S.H.Zhu 0904.4644[hep-ph].

$$\frac{d\sigma_{\gamma A'}}{d\cos\theta} = \frac{2\pi\epsilon^2\alpha^2}{E_{\text{cm}}^2} \left(1 - \frac{m_{A'}^2}{E_{\text{cm}}^2}\right) \frac{1 + \cos^2\theta + \frac{4m_{A'}^2/E_{\text{cm}}^2}{(1 - m_{A'}^2/E_{\text{cm}}^2)^2}}{(1 + \cos\theta)(1 - \cos\theta)},$$

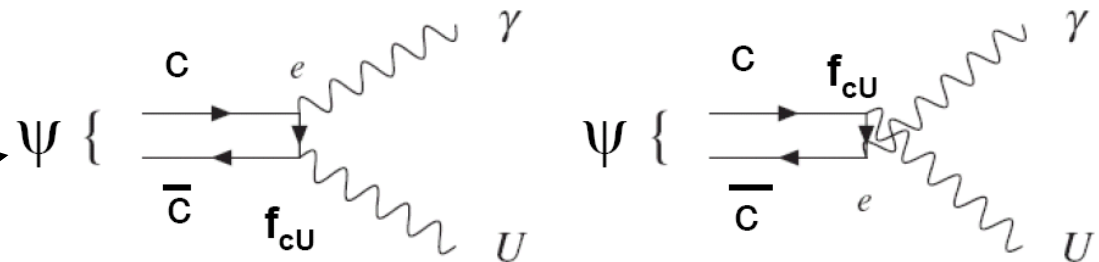
Low energy colliders better!

Spin-1 boson in ψ decay:

(1) $\psi \rightarrow e^+e^-U (U \rightarrow e^+e^-, \mu^+\mu^-)$

(2) $\psi \rightarrow \gamma U (U \rightarrow e^+e^-, \mu^+\mu^-)$

S. H. Zhu PRD75, 115004(2007)



The mass of U-boson : $U \rightarrow e^+e^-, \mu^+\mu^-, \pi^+\pi^-, \nu\nu, \chi\chi$

U-boson in direct production $e^+e^- \rightarrow U\gamma \rightarrow l^+l^-\gamma$

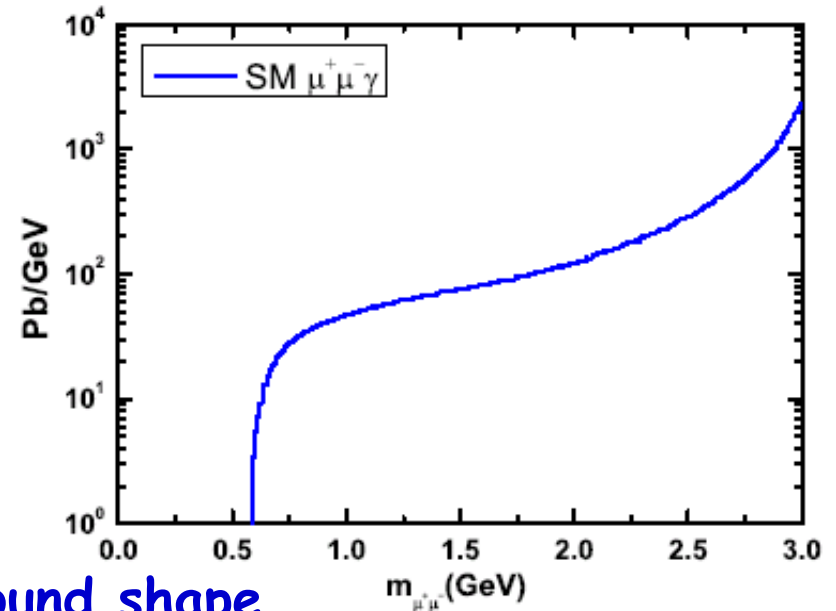
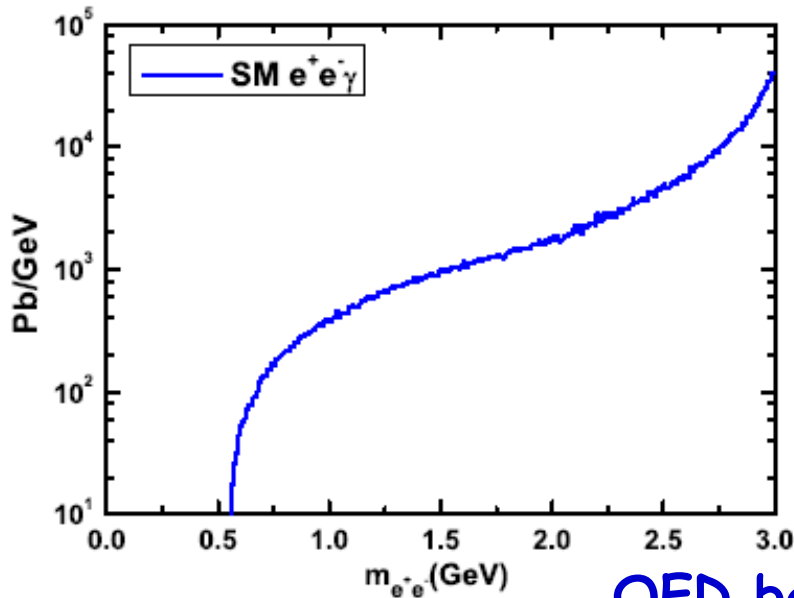
The major standard model processes: $e^+e^- \rightarrow \gamma e^+e^-$ $e^+e^- \rightarrow \gamma \mu^+\mu^-$

$$\cos(\theta_i) < 0.9,$$

θ_i : polar angle of final states

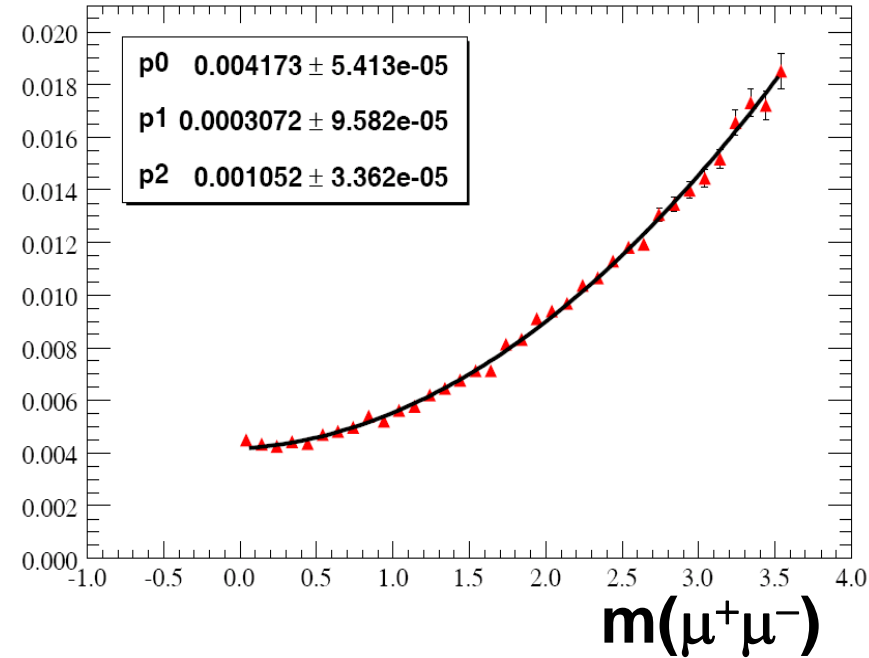
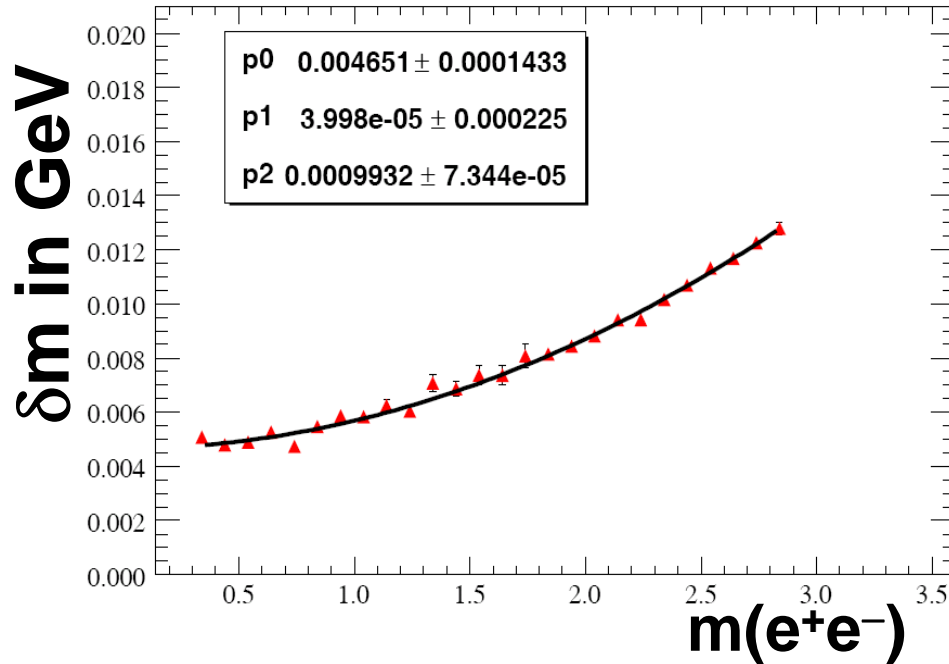
$$\cos(\theta_{l\gamma}) < 0.94,$$

$\theta_{l\gamma}$: angle between lepton and photon



QED background shape

Resolution function of $U \rightarrow l^+ l^-$



**Resolution as a function
of U boson mass:**

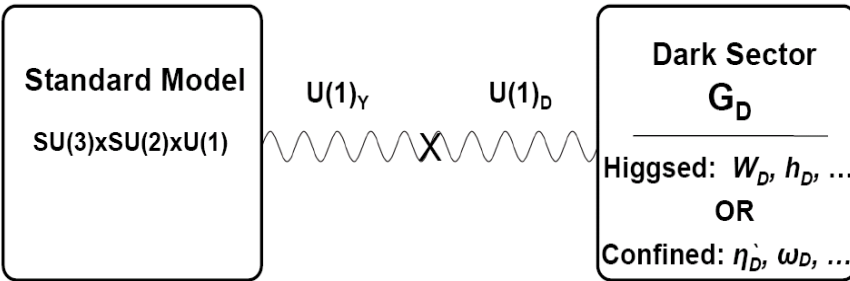
$$\delta m(\mu^+ \mu^-) = \left(4.6 + 0.04 \left(\frac{m_U}{1.0 \text{ GeV}} \right) + 1.0 \left(\frac{m_U}{1.0 \text{ GeV}} \right)^2 \right) \text{ MeV}$$

$$\delta m(e^+ e^-) = \left(4.1 + 0.3 \left(\frac{m_U}{1.0 \text{ GeV}} \right) + 1.0 \left(\frac{m_U}{1.0 \text{ GeV}} \right)^2 \right) \text{ MeV}$$

Reach of U boson decay in visible channel at BESIII

Following the study of M. Reece and L.T.Wang [0904.1743v2hep-ph]
The reach of U boson search at BESIII is estimated.

The model: (R.Essig et al, 0903.3941)



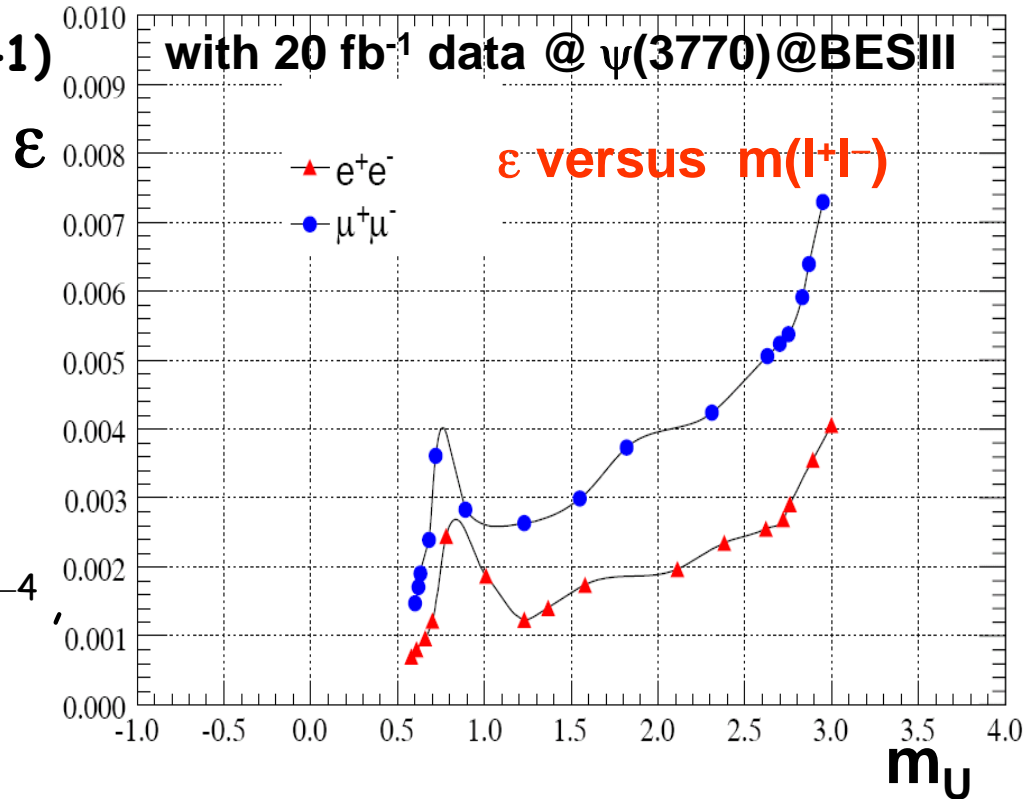
$$\mathcal{L}_{\text{kin-mix}} = -2\epsilon F_d^{\mu\nu} F_{\mu\nu}$$

We consider the case $\epsilon \sim 10^{-3} \text{ -- } 10^{-4}$,

$$\epsilon \sim 10^{-4} \left(\frac{m_U}{1 \text{ GeV}} \right)^2$$

The signal cross section:

$$\sigma_s(e^+e^- \rightarrow U\gamma) \propto \epsilon^2 \times \sigma_0(e^+e^- \rightarrow \gamma\gamma)$$



The significance of the measurement:

$$\frac{S}{\sqrt{B}} \sim \sqrt{\sigma_0 \mathcal{L}} \frac{\epsilon^2}{\sqrt{\alpha/\pi}} \sqrt{\frac{m_U}{\delta m}} \times \text{BR}(U \rightarrow l^+l^-)$$

Spin-1 boson in J/ψ decays

$Br(J/\psi \rightarrow ee) \sim 5\%$

$Br(J/\psi \rightarrow ee \gamma) \sim 0.9\%$ with $E(\gamma) > 100$ MeV

We measure the ratio of J/ψ decays:

$$R = \frac{Br(\psi \rightarrow e^+ e^- U (\rightarrow e^+ e^- \text{ or } \mu^+ \mu^-))}{Br(\psi \rightarrow e^+ e^-)}$$

With 10^{10} J/ψ decays, one can obtain ϵ

$R < 5.0 \times 10^{-6}$ for $U \rightarrow ee$

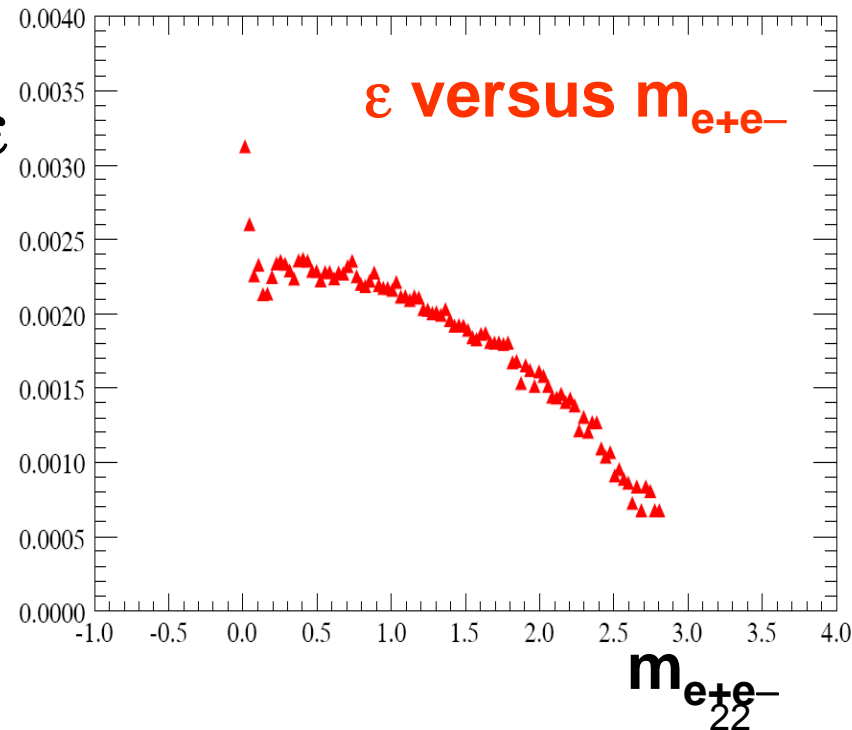
$R < 3.0 \times 10^{-6}$ for $U \rightarrow \mu\mu$

$\epsilon \sim (0.5 - 2.5) \times 10^{-3}$

2.3×10^8 J/ψ collected
in one and half month:

2009-9-25 $\epsilon \sim (0.15 - 7.5) \times 10^{-3}$

Hai-Bo Li



Spin-1 boson in $\psi(2S)$ radiative decays

We can look for U boson with mass less than 170 MeV in

$$R = \frac{Br(\psi(2S) \rightarrow U\chi_{c1,2}, U \rightarrow \text{anything})}{Br(\psi(2S) \rightarrow \gamma\chi_{c1,2})}, \chi_{c1,2} \rightarrow \gamma J/\psi (\rightarrow l^+l^-)$$

$Br(\chi_{c1,2} \rightarrow \gamma J/\psi, J/\psi \rightarrow l^+l^-)$ is large (about $0.36 \times 0.12 = 4\%$ for χ_{c1} .)

Here: anything: ee , $\nu\nu$ or $\chi\chi$ since below $\mu\mu$ threshold.

Advantage: many systematic uncertainties due to reconstruction of $\chi_{c1,2}$ cancel in the ratio R.

Technique: only reconstruct $\chi_{c1,2} \rightarrow \gamma J/\psi (\rightarrow ee \text{ or } \mu\mu)$, by looking at the missing-mass distribution, to search for narrow U boson decaying to anything.

With $3 \times 10^9 \psi(2S)$ per year, the R can be measured to 1×10^{-5} for $10 \text{ MeV} < m_U < 170 \text{ MeV}$, and the $Br(\psi(2S) \rightarrow U\chi_{c1}) \sim 1 \times 10^{-6}$, hence, ϵ can be reached at level of 10^{-3} .

$1 \times 10^8 \psi(2S)$ decays in tape, the $\epsilon < 2.5 \times 10^{-3}$ for U boson 10-170 MeV

Summary

- Invisible decays of mesons give opportunity for probing new physics beyond SM,
- The rates of invisible decay can reach 10^{-4} – 10^{-6} for light mesons at BESIII (see Yang-Heng Zheng's talk).
- The reach of searching for U boson at BESIII has been estimated in direct production and ψ decays:
The reach of mixing parameter ε is less than 10^{-3} ($20 \text{ fb}^{-1} @ \psi(3770)$).

Thank you!